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Time-Space Position of Warm Dense Matter in Laser Plasma Interaction Process[a]

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Abstract: Laser plasma interaction experiments have been performed using an fs Titanium Sapphire laser. Plasmas have been generated from planar PMMA targets using single laser pulses with 3.3 mJ pulse energy, 50 fs pulse duration at 800 nm wavelength. Electron density distributions of the plasmas in different delay times have been characterized by means of Nomarski Interferometry. Experimental data were cautiously compared with relevant 1D numerical simulation. Finally these results provide a first experience of searching for the time-space position of the so-called warm dense plasma in an ultra fast laser target interaction process. These experiments aim to prepare near solid-density

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plasmas for Thomson scattering experiments using the short wavelength free-electron laser FLASH, DESY Hamburg.

I. Introduction

Ultra short-pulse high-intensity laser interact with planar PMMA target had been suggested as a methodology to generate the so-called warm dense matter, which could provide a possible target for the proposed short wavelength Thomson scattering experiments what should be done at the FLASH vacuum ultraviolet free electron laser (VUV-FEL) facility at DESY in Hamburg [1]. The goal of these experiments is to verify a consistent many-particle theory for Thomson scattering in warm dense matter [2]. Plasma densities of $10^{21} \sim 10^{22}$ cm^{-3} and plasma temperature of several eV will be required in these experiments for a detailed analysis of the shape of the expected scattering signal according to a previous work [3]. In this paper, we report our experimental work to search for the time-space position of these dense plasmas in an ultra-fast- laser interacting with planar PMMA target process. Numerical simulations were firstly calculated by using MED 103 [4] which it is a one dimensional Lagrangian hydrodynamic code. According to that results a series of experiments were designed and then performed on a Titanium Sapphire fs laser facility in our laboratory. Electron density lower than 2×10^{20} cm^{-3} of the laser-induced plasmas were measured by a modified time resolved Nomarski interferometer [5] with frequency-doubled probe light. The time-space position of the warm dense matters then was determined by the measured gradient tendency of the plasmas along the axis perpendicular to the target.

II. Numerical simulation of laser PMMA target interaction

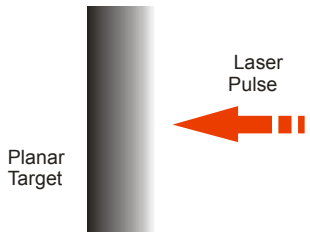
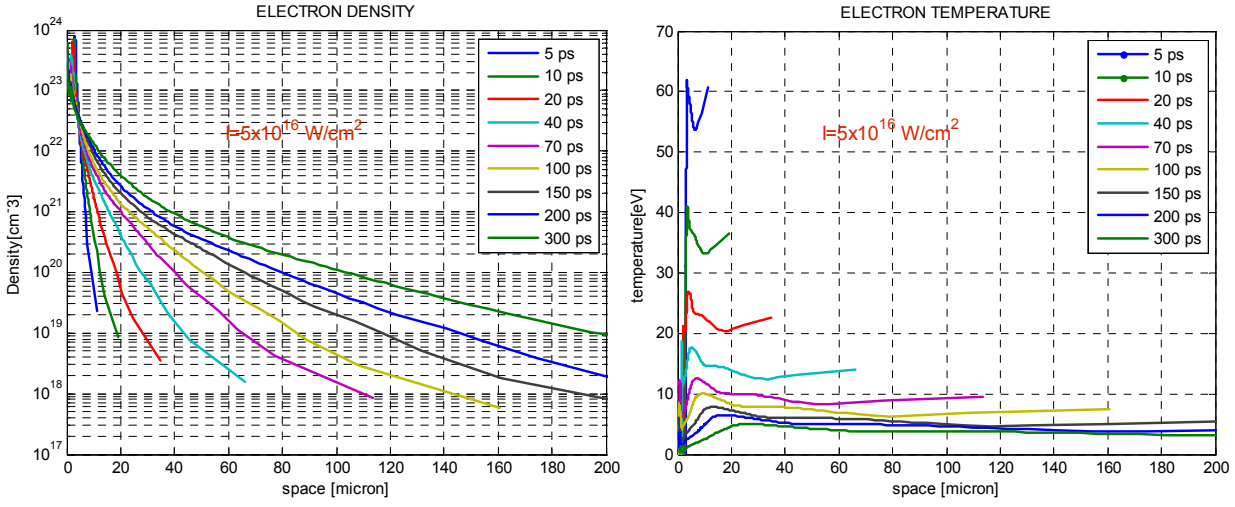


Figure 1 schematic for the calculation

Figure 1 shows the sketch for illustrating the laser target interaction in our calculation. An ultra short pulse laser irradiates perpendicular on a planar PMMA target and ablative plasmas are generated and evolve via time. This evolution can be calculated by using one-dimension hydrodynamic realistic code MED 103 [4] when providing all the parameters of the laser pulse what they are the wavelength, the energy, the duration and the pulse shape, the target parameters including the material, the density and the geometric configuration. The output of the code is a pack of time-space dependent plasma parameters including plasma density, plasma temperature, ionization state, pressure, and plasma velocity. The time resolved plasma density distribution and the plasma temperature distribution are of the most important interests here.



a b

Figure 2. 1-D Numerical simulation result for laser PMMA planar target interaction. a) electron density profile; b) electron temperature profile along the axis. Laser parameter: 5×10^{16} W/cm^2

Figure 2 shows the calculation results when choosing laser parameters as 800 nm wavelength, 50 fs pulse duration, 5×10^{16} W/cm^2 intensity and Gaussian pulse shape. Very high density plasma (>7×10^{23} cm^{-3}) with high temperature (>60 eV) was firstly generated on the target surface, and then expanded into and out of the target. The whole space that the plasmas occupied became lager and lager while the time proceeding. As a result, the electron density and the temperature decreased accompany with such a process. At a time point 300 ps after the laser shot, one can find that the expected warm dense matter with desired density and temperature emerged in a region bordered by the target and a plane 35 microns distant from the target.

Similar simulations for less intense laser were also calculated. Table I.lists the region where the dense plasmas with densities of $10^{21} \sim 10^{22}$ cm^{-3} existed and the relevant temperatures within these region at the point time point 300 ps after the laser shots. These data indicate that short-pulse laser with pulse duration of 100 fs, intensity less than 10^{14} W/cm^2 can nearly generate the desired warm dense matter, lasers with intensities of 10^{15} W/cm^2 and 10^{16} W/cm^2 can but the region where the warm dense matter exist were too close to the target. This means that it is impractical to generate the warm dense matter by laser beam with less intensity than 10^{16} W/cm^2. Higher laser intensity than 5×10^{16} W/^{cm2} might be better, but it would beyond the capability of our Titanium Sapphire laser system.

Table I Comparison of the Space position and temperature of the plasmas at the density of 10^{21}~10^{22} cm^{-3} existed in laser PMMA target interactions for different laser intensity. The time point taken here is 300 ps after the laser pules.

Laser intensity I (W/cm2)	Pulse duration τ (fs)	Distance to the target to limit the Region of the plasma at the density of	Electron density of the warm dense matter T_e (eV)

		10^{21}~10^{22} cm^{-3} Δz (μm)	
10^{13}	100	⩽1	<0.1
10^{14}	100	<4	<0.7
10^{15}	100	<7	<2
10^{16}	100	<16	<7
5×10^{16}	50	<35	<6

III. Electron density measurement with optical laser interferometry

Optical laser interferometry [5-11] had been verified a convenient, efficient and accurate technique to measure the electron density of plasma generated by short pulse laser. But the capability is limited by two factors: the first one is the so-called critical density and the second is due to the deflection of the probe light in plasma with gradient density [12, 13]. Plasma density higher than 1.7×10^{20} cm^{-3} had not been reached by using this way up to present. Further more, the measurement should be not very accurate if any deflection exists ineligible. The warm dense matter with electron density of 10^{21}~10^{22} cm^{-3} then cannot be measured.

However, the gradient tendency of electron density profile (see Figure 2-a) provides an index to search the position of higher density plasma from the lower one. Herein this work, we will measure the profiles of the electron density of the lower density (<2×10^{20} cm^{-3}) plasma generated by the laser pulse, and then determine the possible position of the warm dense plasma with desire density.

IV. Interferometry Experiments

The pump-probe interferometry experiments were arranged as shown in Figure 3. A single laser pulse with wavelength of 800 nm, pulse duration 50 fs and pulse energy of 3.3 mJ is delivered to an unbalance beam splitter. 99% of the energy of the beam is reflected to a mirror M5 and then focused by a parabola to the planar PMMA target. The diameter is about 10~15 microns as we measured primarily before a real experiment. Hence the laser intensity reached to the target was confidently near 5×10^{16} W/cm^2. Another 1% part of the beam was frequency doubled by a BBO, and then via a time delay line which was composed of two mirrors M1 and M2, and finally was reflected and sent to the plasma by the two Mirrors M3 and M4. Lens1, Lens 2 and an optical CCD made up a good imaging system to take a microscopy of the small laser-induced plasma with a magnification 10. The Wollaston prism and an additional polarizer

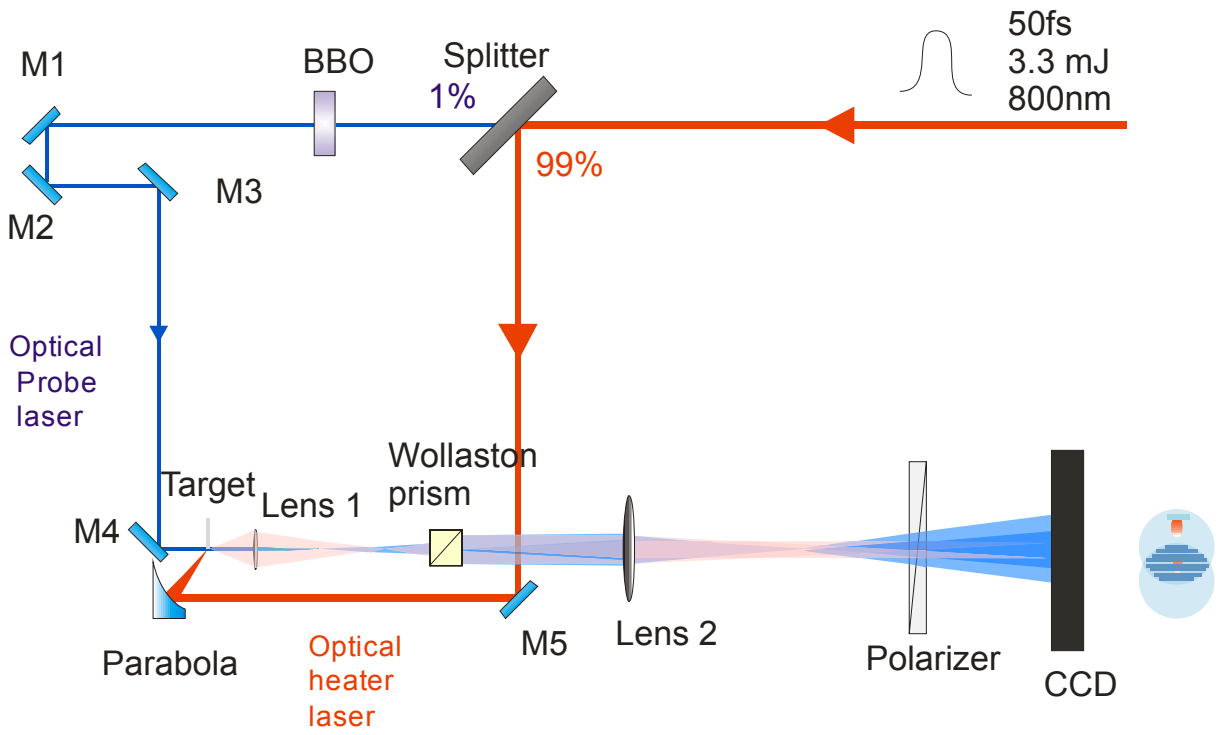


Figure 3. Schematic of the laser plasma interferometry experiments.

which plugged in the light road made the imaging system become a modified Nomarski interferometer. The probe laser including a portion which was distorted by the plasma was split by the Wollaston prism into two beams with perpendicular polarization directions and formed two images on the CCD. The 45º polarizer rotated the tow polarized beams into a same polarization direction again and made them interfered each other on the CCD. Notice that the interferometer used here was a little different from the former used one [5]. The modified interferometer here can collect deflected light with larger deflection angle, and hence, can reach higher plasma density. According to Fermat' principle, stronger deflection means steeper density gradient. Further, steeper gradient always is connected with higher density in a laser-induced plasma (please look at Figure 2-a). This modification permits us to reach higher plasma density but unfortunately can leads to some how measurement error [14].

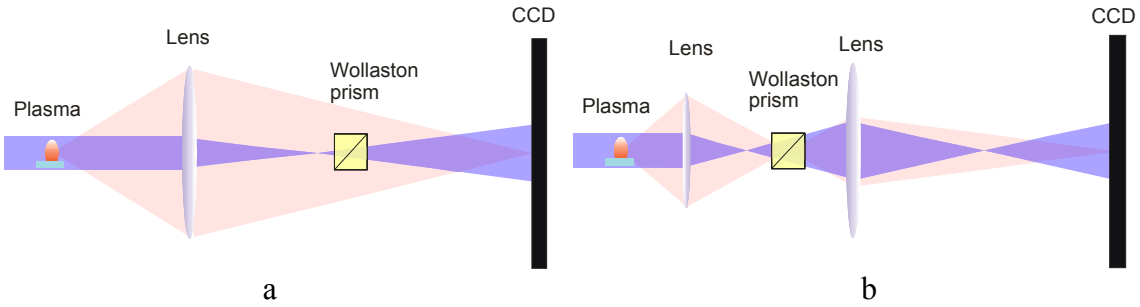


Figure 4. Modification of the interferometer: a). former used Nomarski interferometer in other place; b). presently used interferometer here in this work. The latter one has a lager acceptance solid angle and then can collect stronger deflection light.

V. Results and analysis

Pump-probe experiments had been repeated for lots of times in different time delay (see Figure 2-a) aimed to get a time resolved measurement. We could not get any plasma information when took time delays less than 40 ps. The six interfergram obtained with longer delay time than 40 ps were listed in Figure 5. All the images in Figure 5 had been reconstructed and listed in

Figure 6. These reconstructions were performed by specially developed software for interfergram analysis from Pisa, Italy [15].

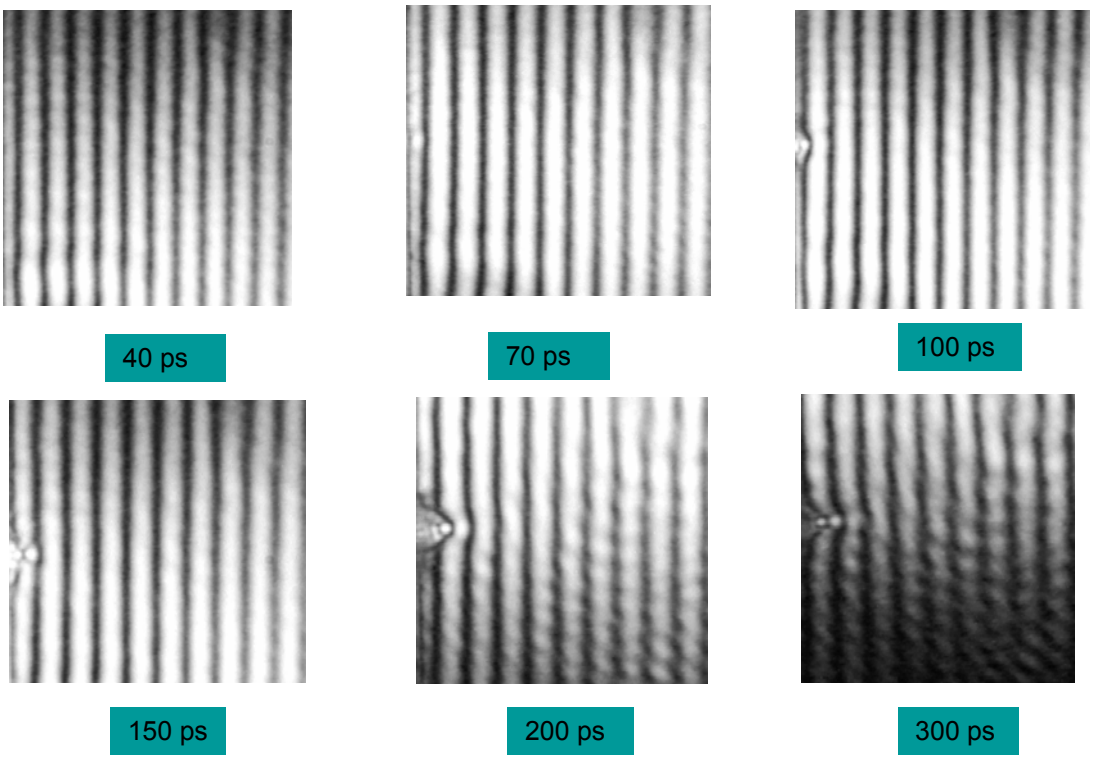


Figure 5. Interfergram obtained in different experiments for different time delay.

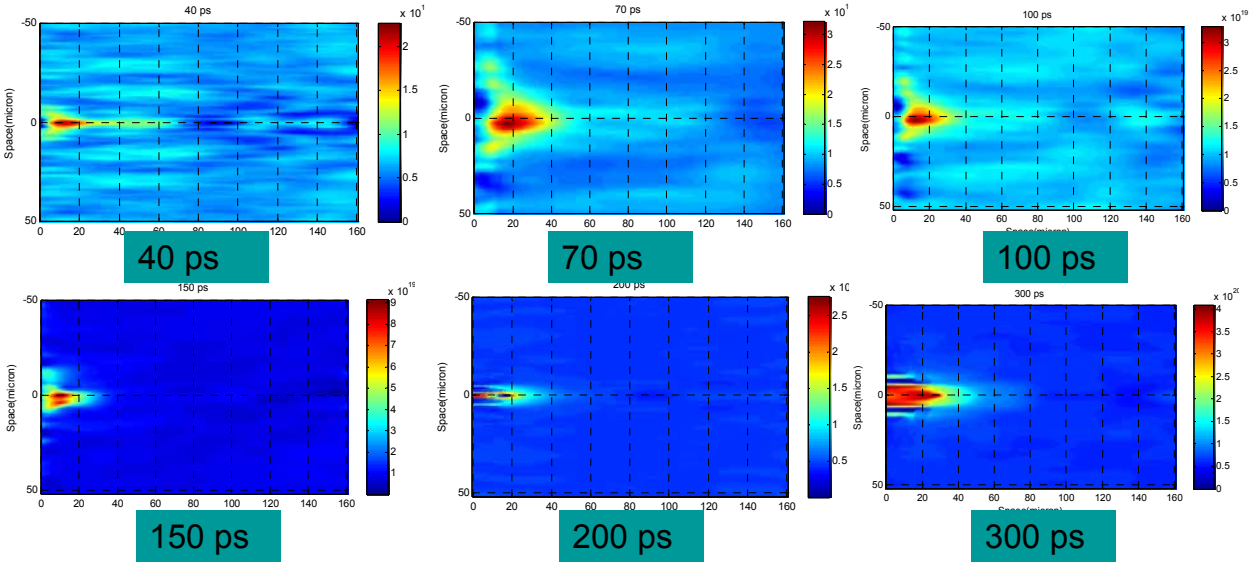


Figure 6. analysis of the interfergram in Figure 5 .

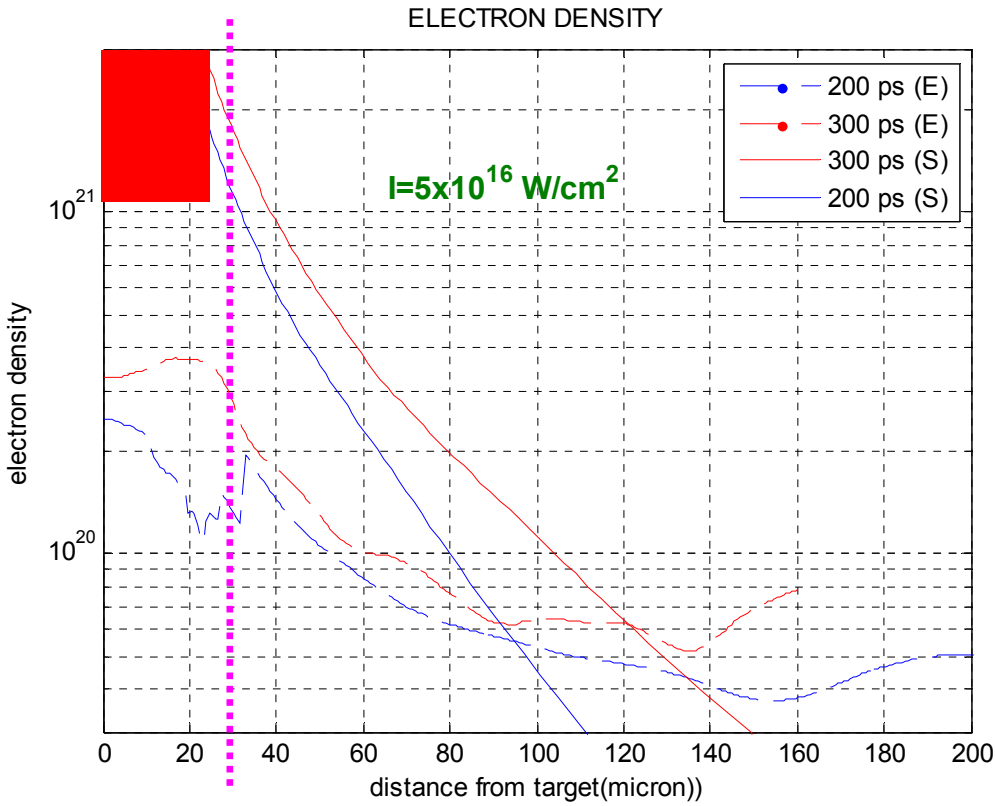


Figure 7. Electron density profiles of the laser-induced along the axis perpendicular to the target. The solid lines give the numerical results while the dashed lines denote the experimental results. The red and the blue lines response to 300 ps and 200 ps time delay respectively.

Figure 7 shows the electron density profiles along the axis perpendicular the target. The experimental curves can be divided into two parts in space. The first part is closed to the target and the distance to the target is not more 30 microns, the second is farer and the distance is great than 30 microns. The first part cannot reflect the truth density profile of the plasma at that region because the gradient there is too steep. In the second part, the curves varied slowly and flatly up to an electron density near to 2×10^{20} cm^{-3}. Judging by the gradient tendency of these curves on can expect that the dense matter should exist at a region where the distance to the target is less than 25 microns.

Agreement between the experiments and the simulations could be found from Figure 7, but obviously there exist some difference. Several possible factors could be the reason which

result it: first of all, the laser parameters we take in the simulation might not so exactly fulfill the practical case; secondly, there could be some errors when we try to find the edges of the target from the experimental data; thirdly and most importantly, according to an previously introduction [14], when deflection of the probe light exists, relevant error similar as the difference in Figure 7 could not be avoided.

VI. Conclusion

Interaction of ultra-short pulse laser and flat PMMA target has been investigated numerically and experimentally. Numerical results show that, the warm dense matter can be found in a region of $z<35$ μm, at the time 300 ps after the laser beam pulse with parameters of 50 fs in duration, $5x10^{16}$ cm^{-3} in intensity and 800 nm in wavelength. Experimental results show that, the dense matter can be expected in a region of $z<25$ μm, at the time 300 ps after the laser beam pulse (50fs, 3.3J, 10~15μm in focus diameter). According to such results, we think that laser induced PMMA plasma reaches warm dense matter conditions to be diagnosed by Thomson diagnostics on the FLASH FEL.

VII. Acknowledgement

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