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Lidar (DIAL)

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Submitted to: SPIE
AeroSense Conference
Orlando, Fl
April 1999

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Huygens-Fresnel wave-optics simulation of atmospheric optical turbulence and reflective speckle in CO₂ differential absorption lidar (DIAL)

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ABSTRACT

The measurement sensitivity of CO₂ differential absorption lidar (DIAL) can be affected by a number of different processes. We have previously developed a Huygens-Fresnel wave optics propagation code to simulate the effects of two of these processes: effects caused by beam propagation through atmospheric optical turbulence and effects caused by reflective speckle. Atmospheric optical turbulence affects the beam distribution of energy and phase on target. These effects include beam spreading, beam wander and scintillation which can result in increased shot-to-shot signal noise. In addition, reflective speckle alone has been shown to have a major impact on the sensitivity of CO₂ DIAL. However, in real DIAL systems it is a combination of these phenomena, the interaction of atmospheric optical turbulence and reflective speckle, that influences the results. The performance of our modified code with respect to experimental measurements affected by atmospheric optical turbulence and reflective speckle is examined. The results of computer simulations are directly compared with lidar measurements. The limitations of our model are also discussed. In addition, studies have been performed to determine the importance of key parameters in the simulation. The results of these studies and their impact on the overall results will be presented.

Keywords: atmospheric turbulence, laser speckle, beam propagation

1. INTRODUCTION

The measurement sensitivity of CO₂ differential absorption lidar (DIAL) can be affected by a number of different processes. Two of these processes are atmospheric optical turbulence and reflective speckle. Atmospheric optical turbulence affects the spatial distribution of energy and phase. Measurable effects include beam spreading, beam wander and scintillation which can result in increased shot-to-shot signal noise. In addition, reflective speckle alone has been shown to have a major impact on the sensitivity of CO₂ DIAL.

The geometry of a hard target reflection scheme is shown in Figure 1. As the laser beam propagates toward the target, index of refraction fluctuations in the atmosphere cause phase distortions in the transverse electric field distribution. Once the laser beam reaches the target, its spatial intensity distribution has been altered compared to what would be observed in propagating through a vacuum. At the target, light is scattered back toward the transmitter. This light passes through essentially the same turbulent atmosphere that altered the outgoing beam (since the atmosphere is "frozen" during the transit time of the pulse for our typical lidar geometries). The return signal will be reduced by any chemical absorbers in the path in accordance with Beer's law. Absorption will also occur from normal atmospheric constituents.

We have developed a Huygens-Fresnel wave optics propagation code to simulate the effects of reflective speckle and atmospheric optical turbulence. Previously, we compared the ability of our model to predict these separate effects with a combination of theory and experimental observations.^{1,2} However, in real DIAL systems it is a combination of these phenomena, the interaction of atmospheric optical turbulence and reflective speckle, which influences the results.³ We

present preliminary results of the comparison of our combined effects simulation with experimental measurements over a finite aperture. We have begun advanced studies to determine the nature of the reflective speckle-atmospheric turbulence interaction and provide some of our findings thus far.

2. MODEL

The model employs the Fresnel-Kirchoff theorem with the Fresnel approximation and assumes paraxial, on-axis propagation.^{4,5} The atmospheric optical turbulence effects are approximated by a series of phase screens over several propagation steps.^{6,7} This model is applied to a lidar geometry in which the beam propagates from the transmitter/receiver through an optically turbulent atmosphere to a diffuse hard target. To simulate reflection from a diffuse hard target, we randomize the phase. After scattering from the target, the portion of the beam that reflects back to our receiver propagates along the same optically turbulent path.

Our Huygens-Fresnel wave optics simulation uses an $N \times N$ array of complex numbers to represent the electric field in a plane perpendicular to the propagation axis. The initial electric field, a Gaussian TEM₀₀ spatial intensity and phase distribution with the properties of our experimental transmitter beam, is used as the input for the simulation. The simulation propagates this initial electric field by dividing the path from lidar system to the target into equal sized steps and applying a phase screen to simulate turbulence effects at each step. The expression for the electric field after a step of distance Δz is determined from⁷

$$E(\hat{\rho}, \Delta z) = IFT \left[\exp \left(i \cdot \pi \cdot \lambda \cdot \Delta z \cdot |\hat{f}|^2 \right) \cdot FT [E(\hat{\rho}, 0) \cdot \exp \{ i \cdot \theta(\hat{\rho}) \}] \right], \quad (1)$$

where $E(\hat{\rho}, 0)$ is the electric field at the initial portion of the step ($z = 0$) with the transverse coordinate given by $\hat{\rho}$. FT is the discrete two-dimensional Fourier transform, $\exp \left(i \cdot \pi \cdot \lambda \cdot \Delta z \cdot |\hat{f}|^2 \right)$ is the Fresnel propagator in the spatial frequency domain \hat{f} , λ is the laser wavelength and IFT is the discrete two-dimensional inverse Fourier transform. The phase screen, $\theta(\hat{\rho})$, is^{6,7}

$$\theta(\hat{\rho}) = 0.0984 \cdot k_o \cdot \sqrt{C_n^2(z) \cdot \Delta z} \cdot (N \cdot \delta x)^{5/6} FT \left[\left(\sqrt{n_x^2 + n_y^2} \right)^{1/6} \cdot \Theta_o(n_x, n_y) \right], \quad (2)$$

where $k_o = 2\pi/\lambda$, $C_n^2(z)$ is the path dependent index of refraction structure constant which characterizes the level of turbulence⁸, Δz is again the step propagation distance, N is the number of pixels along one dimension of the transverse array, $\delta x = \sqrt{\lambda L/N}$ is the optimized pixel width for a transmitter-to-hard target distance of L , n_x and n_y denote integer pixel coordinates within the two-dimensional transverse array, $\Theta_o(n_x, n_y)$ represents an $N \times N$ array of complex unit-variance Gaussian random numbers and FT again implies the two-dimensional discrete Fourier transform. The argument of the FT operation is an array in the spatial frequency domain produced by taking a Gaussian random number distribution and applying the $\left(\sqrt{n_x^2 + n_y^2} \right)^{-1/6}$ factor to impose properties of the Kolmogorov spectrum, which describes the spatial frequency distribution of index of refraction fluctuations, in the transverse plane.⁹ As seen in Equation (2), the magnitude of the turbulence induced phase is dependent on the strength of the turbulence, the length of the propagation step and the lidar wavelength.

The number of propagation steps has a direct impact on computation time. To keep this computation time at a minimum, one goal is to keep the number of propagation steps as low as possible. However, our model is limited in that there are constraints on propagation step size. The assumptions used to approximate a propagation step dictate that the step be within the near field propagation distance. For our lidar geometry, this means that the step can be no longer than the Rayleigh range of our laser transmitter.

Another constraint is that phase effects over this step must not be dominated by amplitude effects. Martin and Flatté¹¹ found that for the phase screen approach to be valid, the normalized point irradiance, σ_I^2 , defined below, for a single propagation step must be less than 1/10 of the total normalized point irradiance variance for the total propagation distance, L ,

$$\sigma_I^2(\Delta z) < 0.1 \sigma_I^2(L). \quad (3)$$

In addition, they found that this variance must be less than 0.1 for one step

$$\sigma_I^2(\Delta z) < 0.1. \quad (4)$$

For spherical wave propagation, assuming weak turbulence, the RMS noise or scintillation at an on-axis point detector is¹¹

$$\sigma_I = \sqrt{\exp(4\sigma_\chi^2) - 1}, \quad (5)$$

where $\sigma_I = \sigma_I / \bar{I}$ is the normalized standard deviation of irradiance. This is the square root of the normalized point irradiance variance discussed earlier. The value σ_χ^2 is the spherical wave log amplitude variance for a point detector. For a path of length L with uniform turbulence along the entire path (i.e. constant turbulence level - C_n^2) the spherical wave log amplitude variance is⁸

$$\sigma_\chi^2 = 0.124 C_n^2 k_o^{7/6} L^{1/6}, \quad (6)$$

again with $k_o = 2\pi/\lambda$.

3. SIMULATION OF INDIVIDUAL EFFECTS

We have shown that this model works well predicting separately the effects of atmospheric optical turbulence and reflective speckle.^{1,2} The simulation of long-term turbulent beam spreading was found to be in agreement with both experimental data and analytical predictions. Simulation values for point detector scintillation due to atmospheric optical turbulence showed agreement with analytical predictions. This last comparison is provided in Figure 2 as an example of our previous work.

We also considered separately the reflective speckle effects in the absence of atmospheric optical turbulence.^{1,2} A surface that is rough on the scale of the laser wavelength scatters the coherent lidar pulse, which produces a complex interference pattern.¹² This pattern is granular in appearance and is commonly referred to as a speckle pattern. Simulated speckle coherence "sizes" were found to be in excellent agreement with those predicted by theory. The intensity probability distributions predicted by our simulation for circular receiver apertures of varying radii agreed with those observed in experiment and expected from theory. Characterized by the parameter M , which is interpreted as the number of speckle inside the receiver for an average pulse, we compared these probability distributions to geometrical predictions from the ratio of the receiver aperture area to the estimated speckle correlation area. We also compared the M value from these probability distributions to the signal-to-noise ratio obtained from the simulation. These comparisons are presented in Figure 3 and Figure 4 as examples of the resulting excellent agreement between the simulation and theoretical predictions. The simulated intensity probability distributions were consistent with those measured experimentally.¹³

4. COMBINED EFFECTS SIMULATION AND COMPARISON WITH EXPERIMENT

We conducted experiments during June and July 1998 in the Nevada desert under conditions of diurnally varying levels of atmospheric turbulence (C_n^2) at ranges of ~1340 m and ~2160 m. The target at ~1340 m was a rotating drum with a diffuse surface. This rotating drum was specifically designed to provide independent speckle realizations as the drum turns at ~ 2 revolutions per minute. The target at ~2160 m had a diffuse surface fixed to plywood. Our lidar consisted of a CO₂ laser with an effective pulse rate per line of ~113 Hz. The receiver configuration was annular with an inner diameter of ~ 4.5" and an outer diameter of ~ 12". The propagation path was horizontal over flat, featureless desert terrain. We concurrently measured the turbulence level with an incoherent near infrared scintillometer propagating over a path that was approximately parallel azimuthally to our lidar beam but on a slant path at a different height.¹⁴ In determining effective turbulence levels for the experiment, we took these differences in the paths into account.⁸

The simulation employed a 1024 x 1024 array and five propagation steps. A total of 100 realizations were run for each turbulence level. The signal was integrated over an annulus that was the same size and configuration as our experimental receiver. Independent speckle realizations were modeled in addition to independent turbulence realizations. The turbulence level (C_n^2) was simulated as uniform over the propagation path.

The model predictions for the combined effects on single-shot RMS noise and comparison to experimental results appear in Figure 5 and Figure 6 for the two ranges mentioned above. The two laser lines we used for this experimental comparison were chosen because of their negligible atmospheric absorption under normal operating conditions. The model, which neglects atmospheric absorption, accurately predicts the level of single-shot RMS noise for our annular aperture. It also correctly predicts the trend of increasing noise with increasing C_n^2 .

5. ADVANCED STUDIES

We have begun advanced studies to learn more about the reflective speckle-atmospheric turbulence interaction. In practice, there are lidar geometries in which the turbulence is not uniform along the entire propagation path. An example of these studies, in which we investigate the impact of cases of uniform/non-uniform turbulence on the speckle correlation diameter, is shown in Figure 7 and Figure 8. The simulation consisted of 100 realizations of atmospheric turbulence and reflective speckle for each specific case, which will be discussed below, using a 512 x 512 grid. The 2000 m propagation path was divided into five 400 m steps. The normalized autocovariance of the speckle intensity at the receiver was determined for each realization and a Gaussian curve was fit to each to determine the speckle correlation diameter.¹ The mean and standard deviation of the 100 Gaussian curve fits was retrieved for each case and is depicted in the figures. The beam divergence input to the simulation was $\sim 200 \mu\text{rad}$ resulting in a relatively large speckle correlation diameter ($\sim 6.61 \pm 0.01 \text{ cm}$) for the zero turbulence case.

It should be pointed out that the Gaussian fit of the normalized autocovariance of the speckle intensity strictly applies to the case of a Gaussian TEM₀₀ beam scattered at the diffuse target.^{12,13} In the case of turbulence along the outgoing portion of the propagation path, the beam on target will not be perfectly Gaussian. However, for the weak levels of turbulence we examine, the deviation from the vacuum case is minimal and provides a semi-quantitative insight into the relative effects of turbulence for the cases we explore.

In Figure 7 we consider two cases in comparison with the vacuum case (Zero Turbulence). In one case, turbulence is present along the entire propagation path (All Phase Screens Round Trip). Here, the speckle correlation diameter is smaller than the vacuum case. We also consider the weighting effect of turbulence contained in only one propagation step for both the outgoing and return legs of the path (Phase Screen Weight Round Trip). The effect here is greatest near the center of the propagation path where the smallest speckle correlation diameters occur. For the portion of the path nearest the target, the effect is negligible since the target will provide randomization that will far outweigh the effects of the phase screen. The phase screen for the step shown in Figure 7 and Figure 8 has a phase standard deviation of $\sim 0.2\pi$ radians with maximum phase variations of $\sim 0.6\pi$ radians. The target phase randomization is uniformly distributed between 0 and 2π radians.

We examine in Figure 8 the weighting effect for the cases of turbulence contained in one step on: (1) the outgoing path (Phase Screen Weight Out) and (2) the return path (Phase Screen Weight Back). For the outgoing path case (Phase Screen Weight Out), the return path had zero turbulence. Likewise, for the return path case (Phase Screen Weight Back), the outgoing path had zero turbulence. In the case of the outgoing path (Phase Screen Weight Out), this effect is minimal and result in speckle correlation diameters very close to the value found in the vacuum case (Zero Turbulence). By contrast, the weighting effect of one step on the return path (Phase Screen Weight Back) is similar to our round trip weighting results (All Phase Screens Round Trip) of Figure 7.

We also found that uniform turbulence along the entire outgoing path with a return propagation through vacuum resulted in a speckle correlation diameter that is very close to our vacuum value. On the other hand, turbulence along the entire return path after an outgoing propagation in vacuum resulted in a value very close to the full turbulence case (All Phase Screens Round Trip) depicted in Figure 7. Hence, turbulence on the return path has the greater effect.

6. CONCLUSIONS

For the lidar geometry of our experiment, the single-shot RMS noise is 40-50% larger under the higher turbulence conditions. The impact of this trend for lidar operations is significant. Even if multi-shot averaging is used to improve the lidar measurement, the initial noise level will be markedly higher for conditions of increased turbulence.

Our model accurately predicted the level of RMS noise for our finite aperture. It also predicted the trend of increasing RMS noise with increasing turbulence level (C_n^2) for our lidar geometry.

These preliminary results provide experimental verification for our modeling of the combined effects of atmospheric optical turbulence and reflective speckle. The results also emphasize, for this lidar geometry, the impact of increased turbulence levels on lidar operations and provide motivation for further study.

The study of the relative effects of turbulence on speckle correlation diameter provides some insight into the reflective speckle-atmospheric turbulence interaction. One finding is that the turbulence on the return path from the target has a much greater impact on reducing the size of relatively large speckle than turbulence on the outgoing path. We also found that the turbulence near the center of the return path has the greatest effect. The effect of turbulence near the target had almost no impact on reducing the relatively large speckle correlation diameter.

Our advanced studies have thus far involved very weak turbulence regimes. Future work will include further quantification of the model's limitations so that stronger turbulence values may be used.

ACKNOWLEDGEMENTS

This work was part of a large lidar project with many important team members and we wish to thank all of them. In particular we thank L. John Jolin for his contributions to this research. The authors would also like to acknowledge the cooperation of the Nevada Test Site Spill Test Facility, the DOE Remote Sensing Laboratory, Nellis Air Force Base and Air Force Research Laboratory. We are very grateful to the project leader, John F. Schultz, who directed this research. This research was fully supported by the U.S. Department of Energy under contract W-7405-ENG-36.

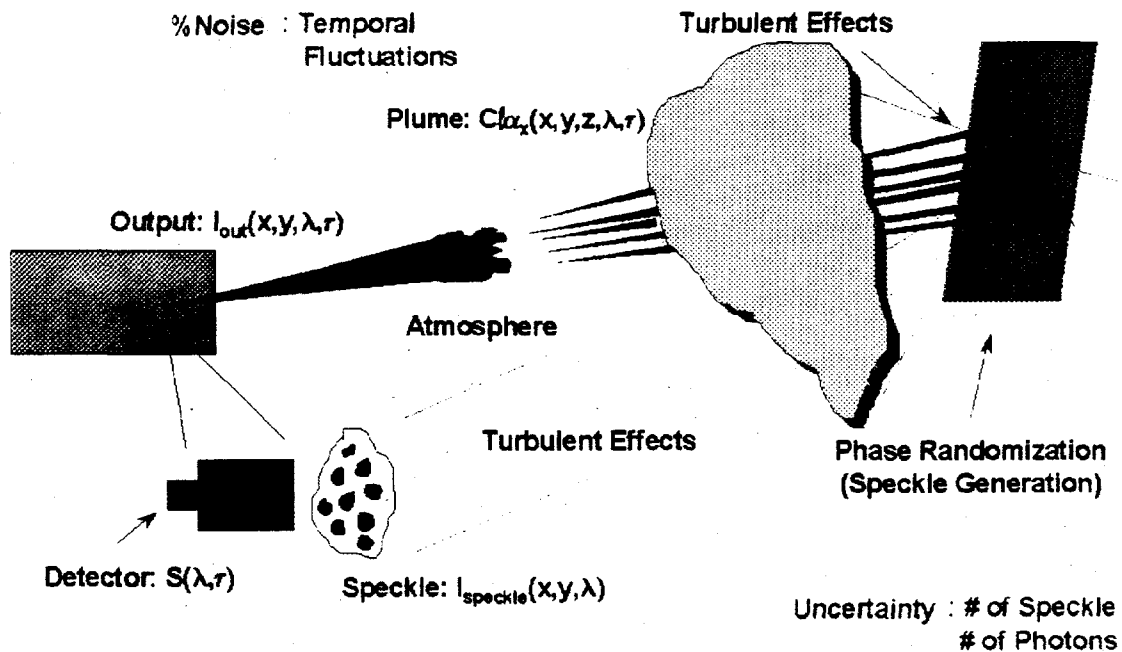


Figure 1. Hard target reflection scheme of lidar highlighting effects of the atmosphere and the target on the return signal.

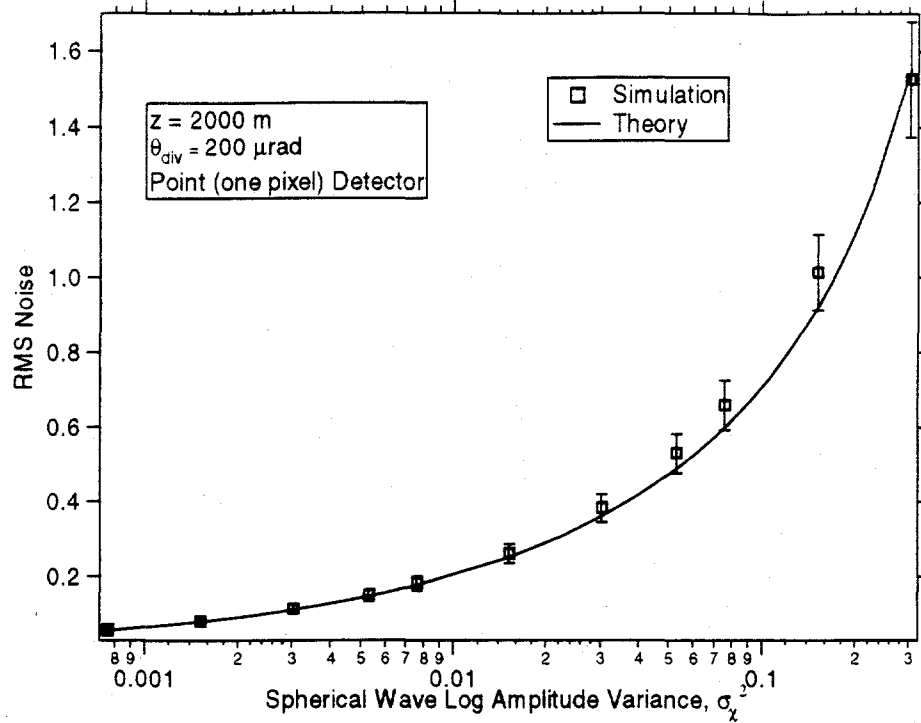


Figure 2. Scintillation of a lidar beam for a point (one pixel ~ 0.0046 m x 0.0046 m square) receiver after a 2000 m round trip propagation assuming no reflective speckle generation. The beam divergence is 200 μ rad and the laser wavelength is 10.6 μ m.

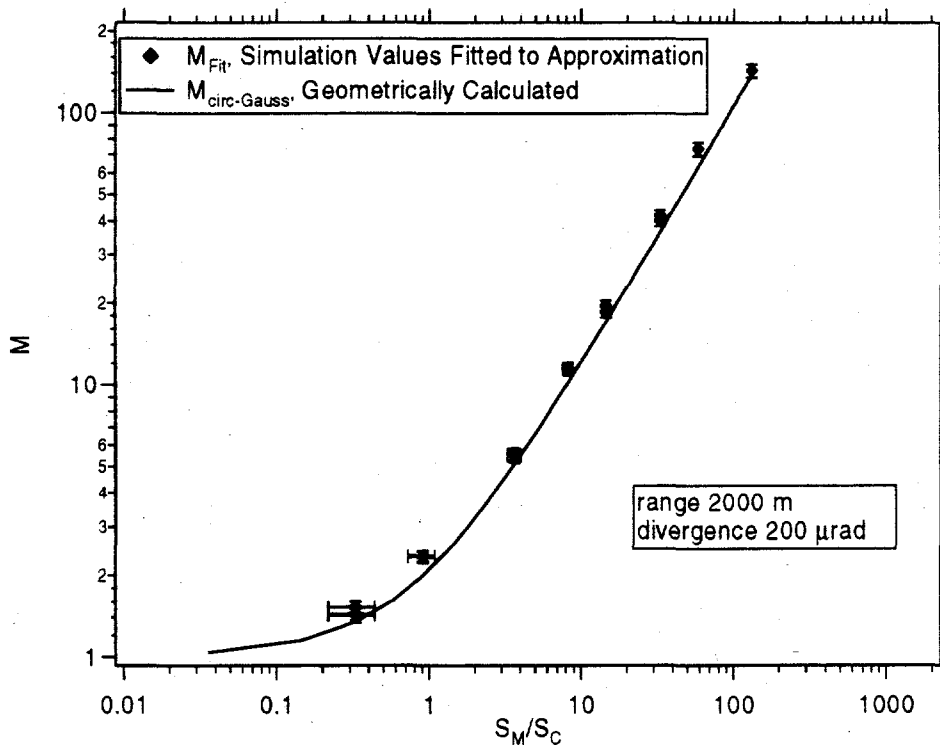


Figure 3. M values compared with receiver area/speckle "area" ratio for 1000 pulse intensity probability distributions and a 2000 m one way propagation with independent speckle realizations. Five propagation steps were used for each leg of the round trip with a 512 x 512 grid, a beam divergence of 200 μ rad and a laser wavelength of 10.6 μ m.

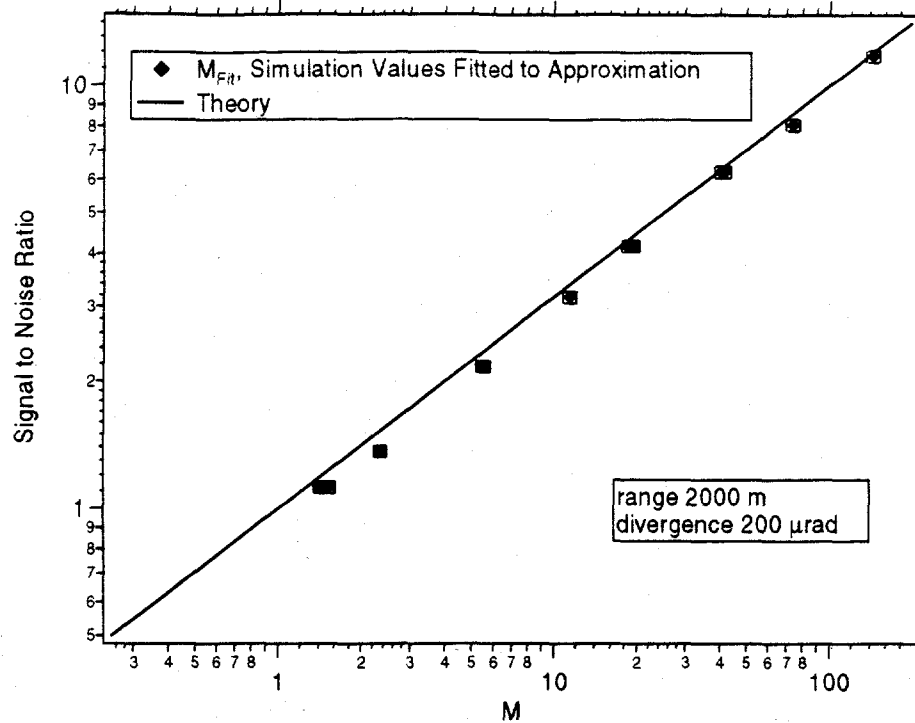


Figure 4. Signal-to-noise ratio versus M for 1000 pulse intensity probability distributions and a 2000 m one way propagation with independent speckle realizations. Five propagation steps were used for each leg of the round trip with a 512 x 512 grid, a beam divergence of 200 μrad and a laser wavelength of 10.6 μm .

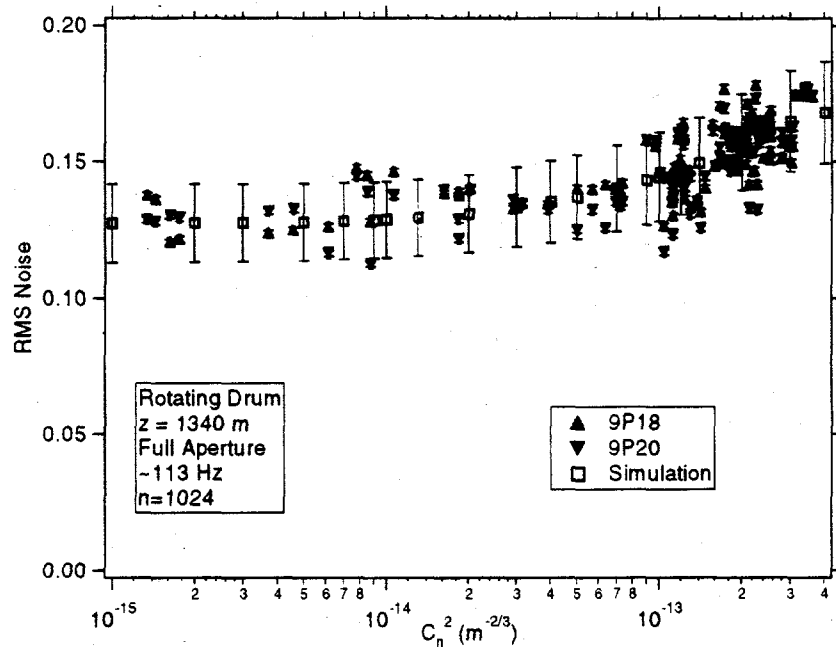


Figure 5. Comparison of simulation with experiment for a target at ~1340 m. The beam divergence was approximated as ~340 μrad . The receiver is annular with an area of ~0.06 m^2 . The simulation grid was 1024 x 1024. Five propagation steps were used for each leg of the round trip path. The propagation path was assumed horizontal with a uniform turbulence level along the path.

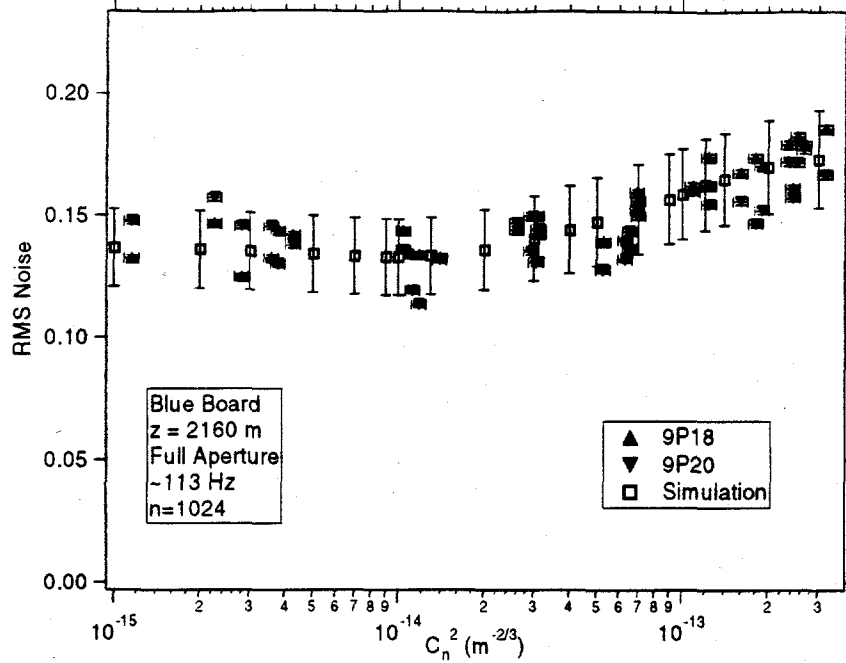


Figure 6. Comparison of simulation with experiment for a target at ~ 2160 m. The beam divergence was approximated as ~ 340 μ rad. The receiver is annular with an area of ~ 0.06 m^2 . The simulation grid was 1024×1024 . Five propagation steps were used for each leg of the round trip path. The propagation path was assumed horizontal with a uniform turbulence level along the path.

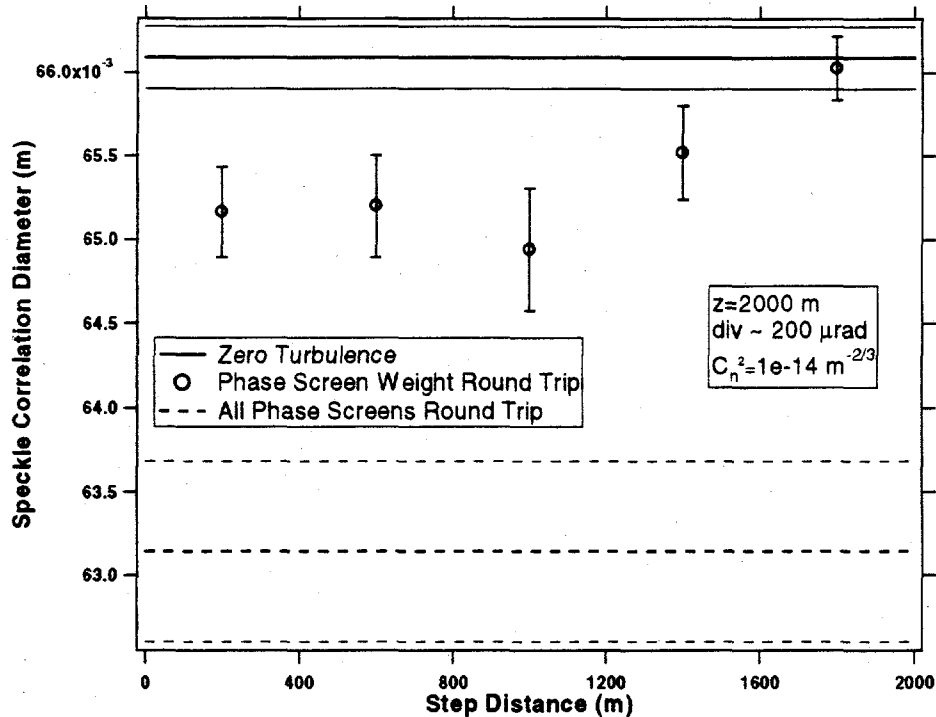


Figure 7. A comparison of the weighting effect of turbulence along the propagation path on speckle correlation diameter. The solid lines depict the mean and standard deviation for the vacuum case (Zero Turbulence). The dashed lines (All Phase Screens Round Trip) indicate the mean and standard deviation for the case in which there is turbulence in both the outgoing and return portion of the entire propagation path. The circles (Phase Screen Weight Round Trip) indicate the case in which turbulence exists on the outgoing and return portions of only one propagation step (400 m).

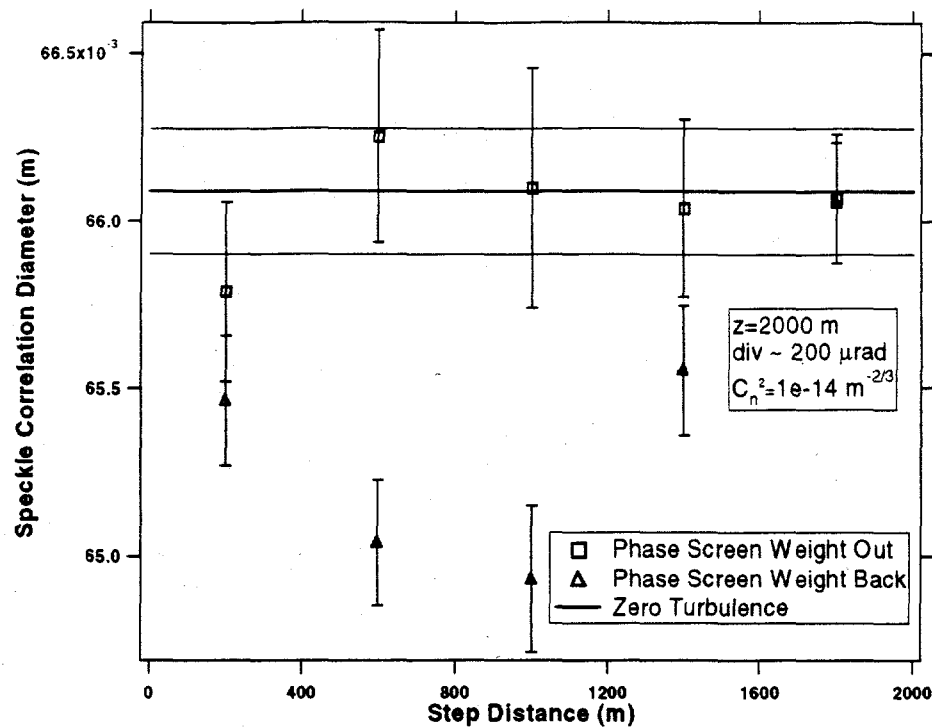


Figure 8. The effect of outgoing/return weighting of turbulence along the propagation path on the speckle correlation diameter. The solid lines depict the mean and standard deviation for the vacuum case (Zero Turbulence) as in Figure 7. The squares (Phase Screen Weight Out) indicate the case in which turbulence exists only on the outgoing portion of one propagation step (400 m). The triangles (Phase Screen Weight Back) indicate the case in which turbulence exists only on the return portion of one propagation step (400 m).

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