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LAMINATED BEAMS - DEFLECTION
AND STRESS AS A FUNCTION OF
EPOXY SHEAR MODULUS

BY

J. BIALEK

PLASMA PHYSICS
LABORATORY



PRINCETON UNIVERSITY
PRINCETON, NEW JERSEY

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J. Bialek
Princeton Plasma Physics Laboratory
Princeton University, Princeton, N. J. 08540

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J. Bialek

Princeton Plasma Physics Laboratory
Princeton University, Princeton, N. J. 08540

ABSTRACT

The large toroidal field coil deflections observed during the PLT power test are due to the poor shear behavior of the insulation material used between layers of copper. Standard techniques for analyzing such laminated structures do not account for this effect. This paper presents an analysis of laminated beams that corrects this deficiency. The analysis explicitly models the mechanical behavior of each layer in a laminated beam and hence avoids the pitfalls involved in any averaging technique. In particular, the shear modulus of the epoxy in a laminated beam (consisting of alternate layers of metal and epoxy) may span the entire range of values from zero to classical. Solution of the governing differential equations defines the stress, strain, and deflection for any point within a laminated beam. The paper summarizes these governing equations and also includes a parametric study of a simple laminated beam.

INTRODUCTION

The toroidal field coils in the Princeton Large Torus (PLT) Tokamak are subjected to a non uniform self generated magnetic field. This results in a force distribution that causes each toroidal field coil to bend in its own plane. The deflections observed on the PLT toroidal field coils^{1,2} are not in agreement with calculations based on standard composite beam theory. In fact, the observed deflections are three to eight times larger than composite beam theory predicts. This paper presents a solution that models deflection and stress in straight laminated beams. The beams analyzed have mechanical properties similar to the PLT toroidal field coils.

The laminated beams that are considered in this paper are made of alternate layers of copper and epoxy. The unique mechanical behavior of these beams is due to the poor shear behavior of the epoxy layers. Experimental measurements³ of effective shear modulus for an epoxy layer give values for the shear modulus thirty (30) times smaller than expected. This of course effects the stiffness and stress distribution within a laminated beam. The goal of this work was to develop a model that describes the behavior of a laminated beam as the shear modulus of the epoxy varies from classical values ($G_{\text{epoxy}} = E_{\text{epoxy}} / 2(1+\nu)$) to zero.

The differential equations governing deflection and stress in each layer of a straight laminated beam are presented. These equations treat every layer of the beam explicitly and do not depend on any averaging of the beam's cross sectional properties.

Hence this analysis provides an explicit statement of deflection and stress throughout a laminated beam and avoids the pitfalls of all averaging procedures. The solution of these equations for the particular problem of a simply supported uniformly loaded beam is also presented. Results are summarized for a parametric study of this problem and finally, the stiffness and stress distribution in a straight laminated beam having the same cross section as the PLT toroidal field coil is presented.

ANALYSIS

The displacement energy method is used to analyze this problem. First the total potential energy of the laminated beam is expressed in terms of the displacement of each layer of metal. Then, variational calculus is applied to the total potential energy expression to obtain the equations governing the displacement of every layer in the beam. These equations are the equations of static equilibrium and the stresses in the beam are obtained by taking derivatives of the displacement solution. The governing equations form a set of coupled ordinary differential equations and are solved by using the Ritz procedure. The small computer program written to obtain numerical results is listed in Appendix 1.

The total potential energy of a laminated beam (π^*) is given by

$$\pi^* = E - V \quad (1a)$$

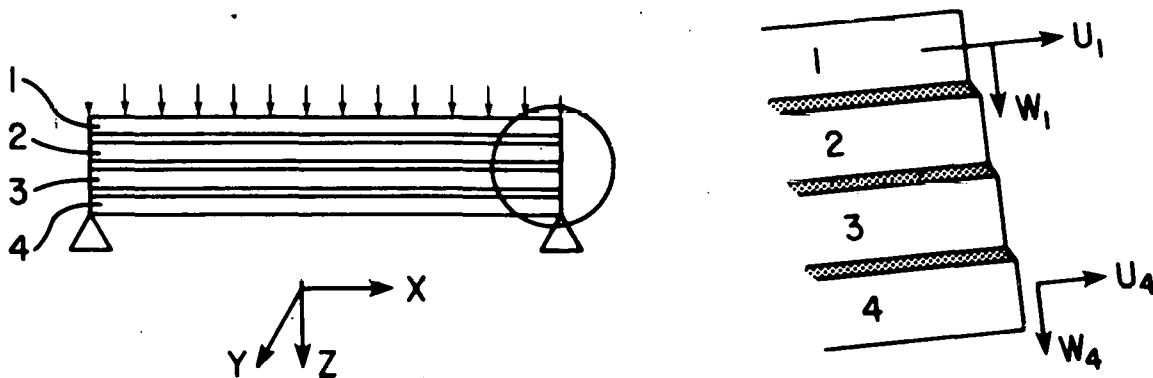
$$E = \int_{\text{volume}} \left[\frac{E_m \epsilon_{xx}^2}{2} + \frac{G_m \gamma_{xz}^2}{2} + \frac{E_e \epsilon_{zz}^2}{2} + \frac{G_e \gamma_{xz}^2}{2} \right] dV \quad (1b)$$

$$V = \int_0^l q(x)w(x)dx \quad (1c)$$

Where

- E = Strain energy in the laminated beam
- V = The potential energy of the applied load intensity "q".
- E_m = Youngs modulus in the metal
- G_m = Shear modulus in the metal
- E_e = Youngs modulus in the epoxy
- G_e = Shear modulus in the epoxy
- $\epsilon_{xx})_m$ = xx component of strain in the metal
- $\epsilon_{zz})_e$ = zz component of strain in the epoxy
- $\gamma_{xz})_m$ = Shear angle in the metal
- $\gamma_{xz})_e$ = Shear angle in the epoxy

Figure 1. describes the coordinate system used in the problem.



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Fig. 1. Coordinate system for laminated beam.

The Youngs moduli and shear moduli in Equation 1 are constants for the beam. The first and second terms in Equation 1 represent the axial bending and shearing behavior in each layer of metal. The third term models the lateral compression or tension in the epoxy bond and serves to transmit force in the z direction from layer to layer within the beam. The fourth term in Equation 1 represents the shear behavior in the epoxy layers. G_e in this term should be considered an effective shear modulus and independent of E_e . This fourth term transmits shear between adjacent layers of metal. As G_e becomes smaller, greater shear deflection in the epoxy bond is allowed, less shear exists between layers of metal, and the laminated beam becomes more flexible.

The strain terms in Equation 1 are defined in the usual manner. Within each metal layer the following relations are used:

$$\left. \begin{aligned} \epsilon_{xx})_i &= \frac{\partial U_x)_i}{\partial x} & (a) \\ \gamma_{xz})_i &= \frac{\partial U_z)_i}{\partial x} + \frac{\partial U_x)_i}{\partial z} & (b) \end{aligned} \right\} \quad (2)$$

where:

- i = index referring to layer number in laminated beam $i = 1, 2, \dots, n$ (see Fig. 1.)
- $\epsilon_{xx})_i$ = xx strain component in i^{th} metal layer
- $\gamma_{xz})_i$ = shear angle in i^{th} metal layer $\gamma_{xz})_i = 2\epsilon_{xz})_i$
- $U_x)_i$ = x component of deflection field i^{th} metal layer
- $U_z)_i$ = z component of deflection field i^{th} metal layer

The deflection field used in Equation 2a and 2b must be defined with care. If the deflection field is not simplified a system of partial differential equations will result. If the deflection field is too simple the equations of equilibrium will not be satisfied.⁴ After considerable study, a deflection field that represents a typical Timoshenko beam was selected and found to be adequate. This deflection field maintains the strength of materials assumption that plane sections remain plane but this is restricted to hold only within each layer of metal. The total cross section of the beam will deform from a plane to a series of ridges and ramps. The shear deflection (γ_i) in the deflection field is also worthy of note. Although this term is a very small part of the deflection in each layer of the beam, it couples the equations together. If this term was not included, the equations would not satisfy static equilibrium.⁴ Equation 3 summarizes the deflection field.

$$\left. \begin{aligned} U_x)_i &= u_i(x) - z \left(\frac{dw_i(x)}{dx} - \gamma_i(x) \right) & (a) \\ U_y)_i &\equiv 0 & (b) \\ U_z)_i &= w_i(x) & (c) \end{aligned} \right\} \quad (3)$$

where:

z = the value of the z coordinate within the i^{th} layer
measured from a coordinate system with the $z = 0$ value
located at the geometric center of the i^{th} metal layer

$u_i(x)$ = the x deflection of the center line of the i^{th} metal
layer

$w_i(x)$ = the z deflection of the center line of the i^{th} metal
layer

$\gamma_i(x)$ = the shear deflection of the center line of the i^{th}
metal layer

Hence the strains in the metal are:

$$\left. \begin{aligned} \epsilon_{xx})_i &= \frac{\partial U_x}{\partial x} = \frac{du_i}{dx} - z \left(\frac{d^2 w_i}{dx^2} - \frac{d\gamma_i}{dx} \right) & (a) \\ \gamma_{xz})_i &= \frac{\partial U_z}{\partial x} + \frac{\partial U_x}{\partial z} = \frac{dw_i}{dx} - \left(\frac{dw_i}{dx} - \gamma_i \right) = \gamma_i & (b) \end{aligned} \right\} \quad (4)$$

The $3n$ functions u_i , w_i , and γ_i ($i = 1, 2, \dots, n$) are the unknowns of this problem. The first two terms in Equation 1 are now explicitly defined. The remaining two terms in Equation 1 will be defined in terms of the deflection field of the metal layers. By doing this, the layers of the beam are connected to each other, no new unknowns are introduced, the strain energy in the epoxy is described, and compatibility of the deflection fields is insured.

The lateral compressive or tensile strain in the epoxy layer i is defined by

$$\epsilon_{zz}^{(i)} = \frac{w_{i+1} - w_i}{t_i} \quad (5)$$

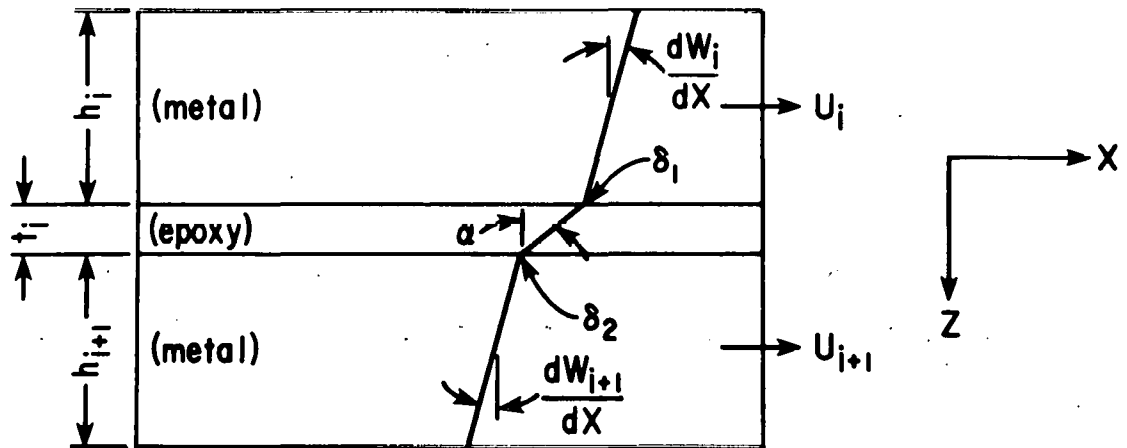
where

$\epsilon_{zz}^{(i)}$ = the zz strain component in the i^{th} epoxy layer
(between metal laminates i and $i+1$)

t_i = the thickness of the i^{th} epoxy layer ($i = 1, 2, \dots, n-1$)

The shear strain in the i^{th} epoxy layer (between metal layers i and $i+1$) is defined by the following argument.

Figure 2 describes the geometry of a typical epoxy layer in its deformed state.



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Fig. 2. Shear in epoxy bond.

The angle α is defined by

$$\alpha = \frac{\delta_1 - \delta_2}{t} \quad (6)$$

where

$$\left. \begin{aligned} \delta_1 &= u_i - \frac{h_i}{2} \left(\frac{dw_i}{dx} - \gamma_i \right) \\ \text{and} \\ \delta_2 &= u_{i+1} + \frac{h_{i+1}}{2} \left(\frac{dw_{i+1}}{dx} - \gamma_{i+1} \right) \end{aligned} \right\} \quad (7)$$

The angle α includes both the shear angle $\gamma_{xz}^*)_i$ and the average slope ($\bar{\alpha}$) of the epoxy layer

$$\text{i.e.} \quad \alpha = \gamma_{xz}^*)_i + \bar{\alpha} \quad (8)$$

The average slope of the i^{th} epoxy layer is simply a weighted average of the slopes for the i^{th} and $(i+1)^{\text{st}}$ metal layers.

$$\bar{\alpha} = \left[h_i \left(\frac{dw_i}{dx} - \gamma_i \right) + h_{i+1} \left(\frac{dw_{i+1}}{dx} - \gamma_{i+1} \right) \right] \frac{1}{(h_i + h_{i+1})} \quad (9)$$

where h_i is the thickness of the i^{th} metal layer.

Substitution of Eq. (9) into Eq. (8) yields

$$\gamma_{xz}^*)_i = \frac{1}{t_i} \left[u_i - u_{i+1} - K_i \left(\frac{dw_i}{dx} - \gamma_i \right) - L_i \left(\frac{dw_{i+1}}{dx} - \gamma_{i+1} \right) \right] \quad (10)$$

where

$$\left. \begin{aligned} K_i &= \frac{h_i}{2} + \frac{h_i t_i}{(h_i + h_{i+1})} \\ L_i &= \frac{h_{i+1}}{2} + \frac{h_{i+1} t_i}{(h_i + h_{i+1})} \end{aligned} \right\} \begin{array}{l} \text{(a)} \\ \text{(b)} \end{array} \quad \begin{array}{l} i = 1, 2, \dots, n-1 \end{array} \quad (11)$$

Then the total potential energy of the laminated beam π^* is:

$$\begin{aligned}
 \pi^* = & \int_0^l \sum_{i=1}^n \frac{Eb}{2} \left[\left(\frac{du_i}{dx} \right)^2 h_i + \frac{h_i^3}{12} \left(\left(\frac{d^2 w_i}{dx^2} \right)^2 + \left(\frac{d\gamma_i}{dx} \right)^2 - 2 \frac{d\gamma_i}{dx} \frac{d^2 w_i}{dx^2} \right) \right] dx \\
 & + \int_0^l \sum_{i=1}^n \frac{Gb_i}{2} \gamma_i^2 dx + \int_0^l \sum_{i=1}^{n-1} \frac{E b}{2 t_i} (w_{i+1}^2 + w_i^2 - 2 w_i w_{i+1}) dx \\
 & + \int_0^l \sum_{i=1}^{n-1} \frac{G_e}{2} \frac{b}{t_i} \left[u_i^2 + u_{i+1}^2 - 2 u_i u_{i+1} + K_i^2 \left(\left(\frac{dw_i}{dx} \right)^2 + \gamma_i^2 - 2 \frac{dw_i}{dx} \gamma_i \right) \right. \\
 & \quad \left. + L_i^2 \left(\left(\frac{dw_{i+1}}{dx} \right)^2 + \gamma_{i+1}^2 - 2 \frac{dw_{i+1}}{dx} \gamma_{i+1} \right) \right. \\
 & \quad - 2 u_i K_i \frac{dw_i}{dx} + 2 u_i K_i \gamma_i - 2 u_i L_i \frac{dw_{i+1}}{dx} + 2 u_i L_i \gamma_{i+1} \\
 & \quad + 2 u_{i+1} K_i \frac{dw_i}{dx} - 2 u_{i+1} K_i \gamma_i + 2 u_{i+1} L_i \frac{dw_{i+1}}{dx} - 2 u_{i+1} L_i \gamma_{i+1} \\
 & \quad \left. + 2 K_i L_i \frac{dw_i}{dx} \frac{dw_{i+1}}{dx} - 2 K_i L_i \frac{dw_i}{dx} \gamma_{i+1} - 2 K_i L_i \gamma_i \frac{dw_{i+1}}{dx} + 2 K_i L_i \gamma_i \gamma_{i+1} \right] dx \\
 & - \sum_{i=1}^n \int_0^l q_i w_i dx
 \end{aligned} \tag{12}$$

where:

b = the uniform width of the laminated beam

ℓ = the length of the beam

q_i = the lateral load (lbs/in) on the i^{th} layer

Taking variational derivatives of this expression with respect to w_i , u_i , and γ_i gives the following differential equations and boundary conditions. These are the equations that describe the behavior of a laminated beam. In the following six expressions those terms labeled "(a)" are the natural boundary conditions that are defined by the variational procedure. The vertical line $\left|_0^\ell$ indicates evaluation at the limits of integration. Those terms labeled "(b)" are the governing differential equations.

$$Ebh_i \frac{du_i}{dx} \delta u_i \Big|_0^\ell = 0 \quad (13a)$$

$$- Ebh_i \frac{d^2 u_i}{dx^2}$$

$$+ (1-\delta_{1i}) \frac{G_e b}{t_{i-1}} \left(u_i - u_{i-1} + K_{i-1} \frac{dw_{i-1}}{dx} - K_{i-1} \gamma_{i-1} + L_{i-1} \frac{dw_i}{dx} - L_{i-1} \gamma_i \right)$$

$$+ (1-\delta_{ni}) \frac{G_e b}{t_i} \left(u_i - u_{i+1} - K_i \frac{dw_i}{dx} + K_i \gamma_i - L_i \frac{dw_{i+1}}{dx} + L_i \gamma_{i+1} \right) = 0 \quad (13b)$$

$$\begin{aligned}
& \frac{Ebh_i^3}{12} \left(\frac{d^2 w_i}{dx^2} - \frac{d\gamma_i}{dx} \right) \frac{d\delta w_i}{dx} \Big|_0^\ell - \frac{Ebh_i^3}{12} \left(\frac{d^3 w_i}{dx^3} - \frac{d^2 \gamma_i}{dx^2} \right) \delta w_i \Big|_0^\ell \\
& + (1-\delta_{il}) \frac{G_{eb}}{t_{i-1}} \left(L_{i-1}^2 \frac{dw_i}{dx} - L_{i-1}^2 \gamma_i - L_{i-1} u_{i-1} \right. \\
& \left. + L_{i-1} u_i + K_{i-1} L_{i-1} \frac{dw_{i-1}}{dx} - K_{i-1} L_{i-1} \gamma_{i-1} \right) \delta w_i \Big|_0^\ell \\
& + (1-\delta_{in}) \frac{G_{eb}}{t_i} \left(K_i^2 \frac{dw_i}{dx} - K_i^2 \gamma_i - K_i u_i \right. \\
& \left. + K_i u_{i+1} + K_i L_i \frac{dw_{i+1}}{dx} - K_i L_i \gamma_{i+1} \right) \delta w_i \Big|_0^\ell = 0 \quad (14a)
\end{aligned}$$

$$\begin{aligned}
& \frac{Ebh_i^3}{12} \left(\frac{d^4 w_i}{dx^4} - \frac{d^3 \gamma_i}{dx^3} \right) + \frac{E_e}{t_{i-1}} (w_i - w_{i-1}) (1-\delta_{il}) + \frac{E_e}{t_i} (w_i - w_{i+1}) (1-\delta_{in}) \\
& + (1-\delta_{il}) \frac{G_{eb}}{t_{i-1}} \left(- L_{i-1}^2 \frac{d^2 w_i}{dx^2} + L_{i-1}^2 \frac{d\gamma_i}{dx} + L_{i-1} \frac{du_{i-1}}{dx} \right. \\
& \left. - L_{i-1} \frac{du_i}{dx} - K_{i-1} L_{i-1} \frac{d^2 w_{i-1}}{dx^2} + K_{i-1} L_{i-1} \frac{d\gamma_{i-1}}{dx} \right) \\
& + (1-\delta_{in}) \frac{G_{eb}}{t_i} \left(- K_i^2 \frac{d^2 w_i}{dx^2} + K_i^2 \frac{d\gamma_i}{dx} + K_i \frac{du_i}{dx} \right. \\
& \left. - K_i \frac{du_{i+1}}{dx} - K_i L_i \frac{d^2 w_{i+1}}{dx^2} + K_i L_i \frac{d\gamma_{i+1}}{dx} \right)
\end{aligned}$$

$$- q_i = 0 \quad (14b)$$

$$\frac{Ebh_i^3}{12} \left(\frac{d\gamma_i}{dx} - \frac{d^2w_i}{dx^2} \right) \delta\gamma_i \Big|_0^l = 0 \quad (15a)$$

$$\begin{aligned} & \frac{Ebh_i^3}{12} \left(\frac{d^3w_i}{dx^3} - \frac{d^2\gamma_i}{dx^2} \right) + Gbh_i\gamma_i \\ & + (1-\delta_{i1}) \frac{G_e b}{t_{i-1}} \left(-L_{i-1}^2 \frac{d\gamma_i}{dx} + L_{i-1}^2 \frac{d^2w_i}{dx^2} - L_{i-1} \frac{du_{i-1}}{dx} \right. \\ & \left. + L_{i-1} \frac{du_i}{dx} + K_{i-1}L_{i-1} \frac{d^2w_{i-1}}{dx^2} - K_{i-1}L_{i-1} \frac{d\gamma_{i-1}}{dx} \right) \\ & + (1-\delta_{in}) \frac{G_e b}{t_i} \left(-K_i^2 \frac{d\gamma_i}{dx} + K_i^2 \frac{d^2w_i}{dx^2} - K_i \frac{du_i}{dx} \right. \\ & \left. + K_i \frac{du_{i+1}}{dx} + K_iL_i \frac{d^2w_{i+1}}{dx^2} - K_iL_i \frac{d\gamma_{i+1}}{dx} \right) = 0 \end{aligned} \quad (15b)$$

where δ_{ij} is the kronecker delta ($\delta_{ij} = 1$ for $i=j$, $\delta_{ij}=0$ otherwise)

Equations 13a, 14a, and 15a are boundary conditions that must always be satisfied. These equations may be satisfied in many different ways. The physical meaning of the solution is dependent on how the boundary conditions are satisfied. This is similar to the situation in a homogeneous beam problem where the choice of end condition determines whether a beam is simply supported, built in, cantilevered, etc.

Equations 13b, 14b, and 15b are solved by the Ritz technique. The particular physical problem considered in this paper is a uniformly loaded simply supported laminated beam. The approximating functions appropriate for this problem are:

$$\left. \begin{aligned} w_i(x) &= \sum_{j \text{ odd}}^n a_{ij} \sin \frac{j\pi x}{\ell} & (a) \\ u_i(x) &= \sum_{j \text{ odd}}^n b_{ij} \cos \frac{j\pi x}{\ell} & (b) \\ \gamma_i(x) &= \sum_{j \text{ odd}}^n c_{ij} \cos \frac{j\pi x}{\ell} & (c) \end{aligned} \right\} \quad (16)$$

These functions satisfy 13a, 14a, and 15a. However the constants a_{ij} , b_{ij} and c_{ij} are not known and need to be determined. This is accomplished by substituting the Equations 16a, 16b, and 16c into Equation 12 and solving for values of the a_{ij} , b_{ij} , and c_{ij} that will minimize π^* . After evaluating the integrals in Equation 12, the total potential energy of the laminated beam is expressed as a quadratic form in the a_{ij} , b_{ij} , and c_{ij} . Taking derivatives with respect to a_{ij} , b_{ij} , and c_{ij} yields a set of linear simultaneous algebraic equations in the a_{ij} , b_{ij} , and c_{ij} . After much algebraic manipulation the governing algebraic equations may be written as follows:

$$\begin{aligned}
& \frac{Ebh_i^3}{12} \left(\frac{j\pi}{\ell} \right)^4 a_{ij} - \frac{Ebh_i^3}{12} \left(\frac{j\pi}{\ell} \right)^3 c_{ij} \\
& + (1-\delta_{il}) \frac{E_e b}{t_{i-1}} (a_{ij} - a_{i-1 j}) + (1-\delta_{in}) \frac{E_e b}{t_i} (a_{ij} - a_{i+1 j}) \\
& + (1-\delta_{il}) \frac{G_e b}{t_{i-1}} \left(L_{i-1}^2 \left(\frac{j\pi}{\ell} \right)^2 a_{ij} - L_{i-1}^2 \left(\frac{j\pi}{\ell} \right) c_{ij} - L_{i-1} \left(\frac{j\pi}{\ell} \right) b_{i-1 j} \right. \\
& \quad \left. + L_{i-1} \left(\frac{j\pi}{\ell} \right) b_{ij} \right. \\
& \quad \left. K_{i-1} L_{i-1} \left(\frac{j\pi}{\ell} \right)^2 a_{i-1 j} - K_{i-1} L_{i-1} \left(\frac{j\pi}{\ell} \right) c_{ij} \right) \\
& + (1-\delta_{in}) \frac{G_e b}{t_i} \left(K_i^2 \left(\frac{j\pi}{\ell} \right)^2 a_{ij} - K_i^2 \left(\frac{j\pi}{\ell} \right) c_{ij} - K_i \left(\frac{j\pi}{\ell} \right) b_{ij} + K_i \left(\frac{j\pi}{\ell} \right) b_{i+1 j} \right. \\
& \quad \left. + K_i L_i \left(\frac{j\pi}{\ell} \right)^2 a_{i+1 j} - K_i L_i \left(\frac{j\pi}{\ell} \right) c_{ij} \right) = q_i \frac{4}{j\pi} \quad (j \text{ odd})
\end{aligned}$$

for $i = 1, 2, \dots, n$ (17)

$$\begin{aligned}
& Ebh_i \left(\frac{j\pi}{\ell} \right)^2 b_{ij} \\
& + (1-\delta_{il}) \frac{G_e b}{t_{i-1}} \left(b_{ij} - b_{i-1 j} + K_{i-1} \left(\frac{j\pi}{\ell} \right) a_{i-1 j} - K_{i-1} c_{i-1 j} + L_{i-1} \left(\frac{j\pi}{\ell} \right) a_{ij} \right. \\
& \quad \left. - L_{i-1} c_{ij} \right) + (1-\delta_{in}) \frac{G_e b}{t_i} \left(b_{ij} - b_{i+1 j} - K_i \left(\frac{j\pi}{\ell} \right) a_{ij} + K_i c_{ij} - L_i \left(\frac{j\pi}{\ell} \right) a_{i+1 j} \right. \\
& \quad \left. + L_i c_{i+1 j} \right) = 0 \quad \text{for } i = 1, 2, \dots, n \quad (18)
\end{aligned}$$

$$\begin{aligned}
 & \frac{Eb_i^3}{12} \left(\frac{j\pi}{\ell} \right)^2 c_{ij} - \frac{Eb_i^3}{12} \left(\frac{j\pi}{\ell} \right)^3 a_{ij} + Gb_i c_{ij} \\
 & (1-\delta_{il}) \frac{G_e b}{t_{i-1}} \left(L_{i-1}^2 c_{ij} - L_{i-1}^2 \left(\frac{j\pi}{\ell} \right) a_{ij} + L_{i-1} b_{i-1 j} - L_{i-1} b_{ij} \right. \\
 & \quad \left. - K_{i-1} L_{i-1} \left(\frac{j\pi}{\ell} \right) a_{i-1 j} + K_{i-1} L_{i-1} c_{i-1 j} \right) \\
 & (1-\delta_{in}) \frac{G_e b}{t_i} \left(K_i^2 c_{ij} - K_i^2 \left(\frac{j\pi}{\ell} \right) a_{ij} + K_i b_{ij} - K_i b_{i+1 j} \right. \\
 & \quad \left. - K_i L_i \left(\frac{j\pi}{\ell} \right) a_{i+1 j} + K_i L_i c_{i+1 j} \right) = 0
 \end{aligned}$$

for $i = 1, 2, \dots, n$ (19)

After close investigation of Equations 17, 18, and 19 it is apparent that the index j is decoupled from the $j-1$ and the $j+1$ terms. This means that the solution procedure for finding a_{ij} , b_{ij} , and c_{ij} ($i=1, \dots, n$; $j = 1, 3, 5, \dots, N$) is a series of small problems (a $3n \times 3n$ set of algebraic equations). In other words, the solution may be approximated to any degree of accuracy desired by solving the Equations 17, 18, and 19 repeatedly and making the value of N in Equations (16) as large as desired. Experience to date indicates that four or five terms ($j = 1, 3, 5, 7$) is adequate.

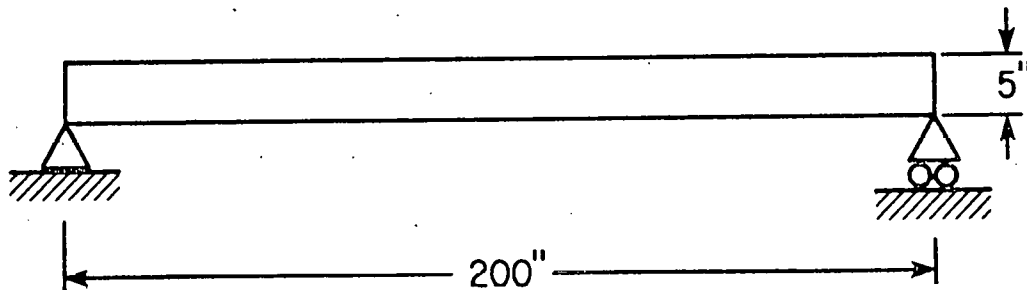
PARAMETRIC STUDIES

Numerical results obtained by solving equations 17, 18, and 19 give deflections and stresses that are in good agreement with expectations. When the shear modulus of the epoxy is set equal to zero the laminated beam acts as if the layers were bending independently. When the shear modulus is set equal to the classical $E_e/2(1+\nu)$ value the results agree with standard composite beam theory. In all of the following parametric studies the following items are kept constant.

$$\begin{aligned} E \quad (\text{metal}) &= 16. \times 10^6 \text{ psi} \\ G \quad (\text{metal}) &= 6. \times 10^6 \text{ psi} \\ E_e \quad (\text{epoxy}) &= 1.0 \times 10^6 \text{ psi} \end{aligned}$$

The value of epoxy shear modulus is treated as a parameter in the following study. However, particular attention should be directed at the results for $G_e = 10000. \text{ psi}$. Experimental work³ at Princeton indicates that the PLT toroidal field coils have an epoxy shear modulus near this value.

The parametric studies are based on a beam with dimensions and end conditions shown in Figure 3.



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Fig. 3.

The net loading intensity in all cases is 1.0 lb/in. The parameters that were varied are the epoxy shear modulus (G_e), the fraction of epoxy in the total cross section (R), and the number of layers in the beam. In all cases, the total height of the beam has remained a constant 5 inches. This means that as the number of layers in the beam changes, the height of each metal layer and each epoxy layer changes. (All metal layers in a calculation have the same thickness.) Table 1 contains a summary of the cross sections of all cases considered.

Figures 4, 5, and 6 illustrate the relationships among epoxy shear modulus, fraction of epoxy, and number of laminates. In all cases the figure labeled B presents the details of the lower part of the figure labeled A.

Figures 7, 8, and 9 illustrate dependence on shear modulus.

Figure 10 illustrates that maximum deflection of a laminated beam still varies linearly with l^4 .

Figures 11 and 12 illustrate strain development across a laminated beam as shear modulus varies. As G_e approaches larger values the strains approach classical results.

Table 2, 3, and 4 contain the numerical data presented in Figures 4 to 12. Those readers interested in maximum strains in epoxy and copper can find the results summarized in Tables 5, 6, and 7. These last tables present the maximum fiber strains in the outermost surface of copper and the maximum shear strain in the epoxy layer closest to the neutral axis.

The results presented so far have been for the beam illustrated in Figure 3. This beam was assumed to have a regular pattern of laminates. The equations developed in this paper can also model

a straight beam made of laminates of different thicknesses. Figures 13 and 14 present the results for a straight laminated beam having the same cross section as the PLT toroidal field coil. This cross section was calculated to be 2.32 times more flexible in bending than a beam of solid copper having the same total build.

CONCLUSIONS

The equations presented here give a description of deflection and stress in straight laminated beams. The results obtained to date indicate that

1. The greater the fraction of epoxy in a cross section the more flexible that cross section becomes.
2. The lower the epoxy shear modulus the more flexible the cross section becomes.
3. The stresses in typical laminated beams are rather insensitive to the epoxy shear modulus. For values of the epoxy shear modulus greater than 10,000 psi the laminated beam is very similar to a solid beam of metal. However, once the shear modulus of the epoxy falls below some particular value the laminated beam becomes radically weaker. This result was not intuitively expected.

ACKNOWLEDGMENT

This work was supported by U. S. Energy Research and Development Administration Contract E(11-1)-3073.

```

0001      IMPLICIT REAL*8(A-H,O-Z)
0002      REAL*8 ZZZ(14400)
0003      REAL*8 A(120,120),ANS(120),AK(40),AL(40),H(40),T(40),Q(40)
0004      REAL*4 AIJ(40,10),BIJ(40,10),CIJ(40,10),TITLE(20)
0005      REAL*8 W(40),U(40),GAMMA(40),DGAMDX(40),ST1(40),ST2(40),GAMMAB(40)
0006      REAL*8 DUDX(40),DWDX(40),D2UDX2(40),D2WDX2(40)
0007      REAL*8 MOM(40),V(40),VB(40),Z(40)
0008      EQUIVALENCE (A(1,1),ZZZ(1)),(ZZZ(1),W(1)),(ZZZ(101),U(1)),
1      (ZZZ(201),GAMMA(1)),(ZZZ(301),DGAMDX(1)),(ZZZ(401),ST1(1)),
2      (ZZZ(501),ST2(1)),(ZZZ(601),GAMMAB(1)),(ZZZ(701),DUDX(1)),
3      (ZZZ(801),DWDX(1)),(ZZZ(901),D2UDX2(1)),(ZZZ(1001),D2WDX2(1)),
4      (ZZZ(1101),MOM(1)),(ZZZ(1201),V(1)),(ZZZ(1301),VB(1))
0009      PIE = 3.141592654

C
0010      1 READ(5,100,END=1000) TITLE
0011      READ(5,102) EM,GM,EX,GX
0012      READ(5,102) ALEN,B,CENTZ
0013      READ(5,101) N,M,IJK,KEY1,KEY2,IBUG
0014      N3 = 3*N
0015      IF (KEY1.NE.0) GO TO 3
C      TEST IF THE BEAM HAS A UNIFORM CROSS SECTION
C      IF KEY1.EQ.0 CROSS SECTION IS UNIFORM
C      IF KEY1.NE.0 CROSS SECTION VARIES
C      CODE FOR UNIFORM CROSS SECTION
0016      READ(5,102) HA,TA
0017      DO 2 I = 1,N
0018      H(I) = HA
0019      T(I) = TA
0020      T(N) = 0.
0021      GO TO 5
C      CODE FOR A VARIABLE CROSS SECTION
0022      3 DO 4 I = 1,N
0023      READ(5,102) H(I),T(I)
0024      4 CONTINUE
C      READ IN THE LOADING - IF KEY2=0 WE HAVE THE SAME LOAD ON ALL LAYERS
0025      5 IF (KEY2.NE.0) GO TO 50
0026      READ(5,102) QJ
0027      DO 48 I = 1,N
0028      Q(I) = QJ
0029      QT = N*QJ
0030      GO TO 70
0031      50 READ(5,101) NLOAD
0032      DO 6 I = 1,N
0033      Q(I) = 0.
0034      QT = 0.
0035      DO 7 I = 1,NLOAD
0036      READ(5,103) J,QJ
0037      QT = QT + QJ
0038      7 Q(J) = QJ
C
C      DEFINE THE AK AND AL ARRAYS
0039      70 NL = N-1
0040      DO 8 I = 1,NL
0041      IP = I+1
0042      AK(I) = H(I)/2. + H(I)*T(I)/(H(I)+H(IP))
0043      8 AL(I) = H(IP)/2. + H(IP)*T(I)/(H(I)+H(IP))

```

```

C
C -----
C          SET UP ALL EQUATIONS
C
0044      PRINT 110,  TITLE
0045      PRINT 111,  N,M,IJK
0046      PRINT 112,  ALEN,B,CENTZ
0047      PRINT 113,  EM,GM,EX,GX
0048      PRINT 114
0049      DO 88  I = 1,N
0050      PRINT 115,  LQ(I)  ,H(I),T(I)
0051  88    CONTINUE
0052      PRINT 116,  QT
0053      DO 20  JJ = 1,M
C
0054      DO 9   I = 1,120
0055      ANS(I) = 0.
0056      DO 9   J = 1,120
0057  9      A(I,J) = 0.
C
0058      AJ = (JJ-1)*2 + 1
0059      PJ1 = AJ*PIE/ALEN
0060      PJ2 = PJ1*PJ1
0061      PJ3 = PJ2*PJ1
0062      PJ4 = PJ3*PJ1
C
0063      DO 15  I = 1,N
0064      IL = I-1
0065      IP = I+1
0066      I1 = (I-1)*3 + 1
0067      I2 = I1 + 1
0068      I3 = I1 + 2
0069      IF ( I .EQ. N ) GO TO 10
C      SOME TERMS THAT APPEAR FREQUENTLY IN THE EQUATIONS ARE
0070      CA1 = EM*B*H(I)*H(I)*H(I)/12.
0071      CA2 = EM*B*H(I)
0072      EA1 = EX*B/T(I)
0073      GA1 = GX*B/T(I)
0074      IF ( I .GT. 1 ) GO TO 10
C
C          SET UP THE EQUATIONS FOR THE FIRST LAYER
C
0075      A(1,1) = CA1*PJ4 + EA1 + GA1*AK(1)*AK(1)*PJ2
0076      A(1,2) = -GA1*AK(1)*PJ1
0077      A(1,3) = -CA1*PJ3 - GA1*AK(1)*AK(1)*PJ1
C
0078      A(1,4) = -EA1 + GA1*AK(1)*AL(1)*PJ2
0079      A(1,5) =  GA1*AK(1)*PJ1
0080      A(1,6) = -GA1*AK(1)*AL(1)*PJ1
0081      ANS(1) = 4.49(1)/(AJ*PIE)
C
C
C
0082      A(2,1) = -GA1*AK(1)*PJ1
0083      A(2,2) = CA2*PJ2 + GA1
0084      A(2,3) = GA1*AK(1)

```

```

      C
0085      A(2,4) = -GA1*AL(1)*PJ1
0086      A(2,5) = -GA1
0087      A(2,6) = GA1*AL(1)
0088      ANS(2) = 0.
      C
      C
      C
0089      A(3,1) = -CA1*PJ3 - GA1*AK(1)*AK(1)*PJ1
0090      A(3,2) = GA1*AK(1)
0091      A(3,3) = CA1*PJ2 + GM*B*H(1) + GA1*AK(1)*AK(1)
      C
0092      A(3,4) = -GA1*AK(1)*AL(1)*PJ1
0093      A(3,5) = -GA1*AK(1)
0094      A(3,6) = GA1*AK(1)*AL(1)
0095      ANS(3) = 0.
0096      GO TO 15
      C
10      SOME MORE USEFULL CONSTANTS ARE
0097      EA2 = E*B/T(IL)
0098      GA2 = G*B/T(IL)
0099      IF ( I.EQ. N ) GO TO 14
      C
      C
      C
      SET UP THE EQUATIONS FOR THE MIDDLE LAYERS
      C
0100      A(I1,I1-3) = -EA2 + GA2*AK(IL)*AL(IL)*PJ2
0101      A(I1,I1-2) = -GA2*AL(IL)*PJ1
0102      A(I1,I1-1) = -GA2*AK(IL)*AL(IL)*PJ1
      C
0103      A(I1,I1) = CA1*PJ4 + EA2 + EA1 + GA2*AL(IL)*AL(IL)*PJ2 +
1      GA1*AK(IL)*AK(IL)*PJ2
0104      A(I1,I1+1) = GA2*AL(IL)*PJ1 - GA1*AK(IL)*PJ1
0105      A(I1,I1+2) = -CA1*PJ3 - GA2*AL(IL)*AL(IL)*PJ1
0106      A(I1,I1+2) = -CA1*PJ3 - GA2*AL(IL)*AL(IL)*PJ1 - GA1*AK(IL)*AK(IL)*PJ1
      C
0107      A(I1,I1+3) = -EA1 + GA1*AK(IL)*AL(IL)*PJ2
0108      A(I1,I1+4) = GA1*AK(IL)*PJ1
0109      A(I1,I1+5) = -GA1*AK(IL)*AL(IL)*PJ1
0110      ANS(I1) = 4.*Q(I)/(A1*PIE)
      C
      C
      C
0111      A(I2,I1-3) = GA2*AK(IL)*PJ1
0112      A(I2,I1-2) = -GA2
0113      A(I2,I1-1) = -GA2*AK(IL)
      C
0114      A(I2,I1) = GA2*AL(IL)*PJ1 - GA1*AK(IL)*PJ1
0115      A(I2,I1+1) = CA2*PJ2 + GA2 + GA1
0116      A(I2,I1+2) = -GA2*AL(IL) + GA1*AK(IL)
      C
0117      A(I2,I1+3) = -GA2*AL(IL) * PJ1
0118      A(I2,I1+4) = -GA2
0119      A(I2,I1+5) = GA2*AL(IL)
0120      ANS(I2) = 0.
      C
      C
      C

```

```

0121      A(13,11-3) = -GA2*AK(IL)*AL(IL)*PJ1
0122      A(13,11-2) = GA2*AL(IL)
0123      A(13,11-1) = GA2*AK(IL)*AL(IL)
      C
0124      A(13,11) = -CA1*PJ3 -GA2*AL(IL)*AL(IL)*PJ1
      1 -GA1*AK(I)*AK(I)*PJ1
0125      A(13,11+1) = -GA2*AL(IL) + GA1*AK(I)
0126      A(13,11+2) = CA1*PJ2 + GM*B*H(I) + GA2*AL(IL)*AL(IL)
      1 + GA1*AK(I)*AK(I)
      C
0127      A(13,11+3) = -GA1*AK(I)*AL(I)*PJ1
0128      A(13,11+4) = -GA1*AK(I)
0129      A(13,11+5) = GA1*AK(I)*AL(I)
0130      ANS(13) = 0.
0131      GO TO 15

```

C

C

C

EQUATIONS FOR THE LAST LAYER

```

0132      14 A(11,11-3) = -EA2 + GA2*AK(IL)*AL(IL)*PJ2
0133      A(11,11-2) = -GA2*AL(IL)*PJ1
0134      A(11,11-1) = -GA2*AK(IL)*AL(IL)*PJ1
      C
0135      A(11,11) = CA1*PJ4 + EA2 + GA2*AL(IL)*AL(IL)*PJ2
0136      A(11,11+1) = GA2*AL(IL)*FJ1
0137      A(11,11+2) = -CA1*PJ3 - GA2*AL(IL)*AL(IL)*PJ1
0138      ANS(11) = 4.*Q(I)/(AJ*PIE)

```

C

C

C

```

0139      A(12,11-3) = GA2*AK(IL)*FJ1
0140      A(12,11-2) = -GA2
0141      A(12,11-1) = -GA2*AK(IL)

```

C

```

0142      A(12,11) = GA2*AL(IL)*PJ1
0143      A(12,11+1) = CA2*PJ2 + GA1
0144      A(12,11+2) = -GA2*AL(IL)
0145      ANS(12) = 0.

```

C

C

C

```

0146      A(13,11-3) = -GA2*AK(IL)*AL(IL)*PJ1
0147      A(13,11-2) = GA2*AL(IL)
0148      A(13,11-1) = GA2*AK(IL)*AL(IL)

```

C

```

0149      A(13,11) = -CA1*PJ3 - GA2*AL(IL)*AL(IL)*PJ1
0150      A(13,11+1) = -GA2*AL(IL)
0151      A(13,11+2) = -CA1*PJ2 + GM*B*H(I) + GA2*AL(IL)*AL(IL)
0152      ANS(13) = 0.

```

```

0153      15 CONTINUE

```

C

C

C

MATRIX COMPLETE - NOW PRINT MATRIX IF DEBUG REQUEST HAS BEEN MADE

```

0154      IF ( IBUG .EQ. 0 ) GO TO 160
0155      DO 16 I = 1,N3
0156      PRINT 119, I, I, ANS(I)
0157      PRINT 120, (A(I,J),J=1,N3)

```

```

0159      16  CONTINUE
0159      PRINT 121
      C
0160      160  CALL SOLVE ( A,ANS,N3,KEY )
      C
0161      IF ( KEY .NE. 0 ) GO TO 1000
0162      DO 17 I = 1,N
0163      I1 = (I-1)*3 + 1
0164      I2 = I1 + 1
0165      I3 = I1 + 2
0166      AIJ(I,JJ) = ANS(I1)
0167      BIJ(I,JJ) = ANS(I2)
0168      CIJ(I,JJ) = ANS(I3)
0169      17 20  CONTINUE
      C
      C
      C
0170      DO 21 I=1,N
0171      PRINT 130, I
0172      PRINT 131,(AIJ(I,J),J=1,M)
0173      21  CONTINUE
0174      DO 22 I = 1,N
0175      PRINT 132, I
0176      PRINT 131, ( BIJ(I,J), J=1,M)
0177      22  CONTINUE
0178      DO 24 I = 1, N
0179      PRINT 133,I
0180      PRINT 131, (CIJ(I,J),J=1,M)
0181      24  CONTINUE
      C
0182      Z(1) = -CENTZ + H(1)/2.
0183      DO 25 J = 2,N
0184      Z(J) = Z(J-1) + T(I) + ( H(J)+H(J-1))/2.
      C
0185      IF ( DABS(Z(J)) .LT. 1.0E-6 ) Z(J) = 0.
0186      25  CONTINUE
0187      DUM = IJK
0188      XDEL = ALEN/DUM
0189      XDD = -XDEL
0190      IJK = IJK + 1
      C
0191      DO 40 II = 1, IJK
0192      DO 26 I = 1,40
0193      W(I) = 0.
0194      U(I) = 0.
0195      GAMMA(I) = 0.
0196      GAMMAB(I) = 0.
0197      DGAMDX(I) = 0.
0198      DGDXX(I) = 0.
0199      DWDXX(I) = 0.
0200      D2UDX2(I) = 0.
0201      D2WDX2(I) = 0.
0202      MOM(I) = 0.
0203      VB(I) = 0.
0204      26  V(I) = 0.
0205      TOTMOM = 0.

```

```

0206      TOTVM = 0.
0207      TOTVB = 0.
0208      XDD = XDD + XDEL
0209      DO 28 I = 1,N
0210      DO 27 J = 1,M
0211      AAJ = (J-1)*2 + 1
0212      ARG = AAJ*PIE*XDD/ALEN
0213      SSS =DSIN(ARG)
0214      CCC =DCOS(ARG)
0215      ARG = AAJ*PIE/ALEN
0216      W(I) = W(I) + AIJ(I,J) * SSS
0217      U(I) = U(I) + BIJ(I,J)*CCC
0218      GAMMA(I) = GAMMA(I) + CIJ(I,J)*CCC
0219      DWDX(I) = DWDX(I) + ARG*AIJ(I,J)*CCC
0220      DUDX(I) = DUDX(I) - ARG*BIJ(I,J) * SSS
0221      DGAMDX(I) = DGAMDX(I) - ARG*CIJ(I,J) * SSS
0222      D2WDX2(I) = D2WDX2(I) - ARG*ARG*AIJ(I,J) * SSS
0223      D2UDX2(I) = D2UDX2(I) - ARG*ARG*BIJ(I,J) * CCC
0224      ST1(I) = DUDX(I) + H(I)*(D2WDX2(I)-DGAMDX(I))/2.
0225      ST2(I) = DUDX(I) - H(I)*(D2WDX2(I)-DGAMDX(I))/2.
0226      A1 = I
0227      Z1 = Z(I) - H(I)/2.
0228      Z2 = Z(I) + H(I)/2.
0229      ZAVG = (Z1+Z2)/2.
0230      MOM(I) =EM*B*(DUDX(I)*ZAVG*H(I) - (D2WDX2(I)-DGAMDX(I))*H(I)*
0231      1 H(I)*H(I) /12.)
0232      TOTMOM = TOTMOM + MOM(I)
0233      V(I) = GAMMA(I) * GM * B * H(I)
0234      TOTVM = TOTVM + V(I)
0235      IF ( I .EQ. 1 ) GO TO 28
0236      IL = I-1
0237      GAMMAB(I) = (U(IL)-U(I) -AK(IL)*(DWDX(IL)-GAMMA(IL)) -AL(IL)*
0238      1 (DWDX(I)-GAMMA(I)))/T(IL)
0239      VB(I) = -GAMMAB(I)*GX*B *T(IL)
0240      TOTVB = TOTVB + VB(I)
0241      28 CONTINUE
0242      RV = QT *ALEN*(1.-2.*XDD/ALEN)/2.
0243      RMOD = QT * ALEN*(XDD-XDD*XDD/ALEN)/2.
0244      VTOT = TOTVM + TOTVB
0245      PRINT 140, XDD
0246      DO 30 I = 1, N
0247      PRINT 142,I,W(I),U(I),GAMMA(I),DWDX(I),DUDX(I),DGAMDX(I),D2WDX2(I)
0248      1 , D2UDX2(I),ST1(I),ST2(I),GAMMAB(I)
0249      30 CONTINUE
0250      PRINT 141
0251      DO 31 I = 1 , N
0252      PRINT 143, MOM(I),V(I),VB(I)
0253      31 CONTINUE
0254      PRINT 144, RMOD,TOTMOM
0255      PRINT 145, RV, VTOT, TOTVM, TOTVB
0256      40 CONTINUE
0257      GO TO 1
0258      1000 CALL EXIT
0259      100 FORMAT(20A4)
0260      101 FORMAT(10I5)
0261      102 FORMAT(4F20.0)

```

```

0259      103  FORMAT(12,F25.0)
0260      110  FORMAT(1H1,/,20X,20A4)
0261      111  FORMAT(/,10X,'N = ',15,10X,'M = ',15,10X,'IJK = ',15)
0262      112  FORMAT(/,10X,'LENGTH OF BEAM(INCHES) = ',F15.2,
           1 /10X,'WIDTH OF BEAM (INCHES) = ',F15.2,
           2 /10X,'POSITION OF NEUTRAL AXIS = ',F15.2)
0263      113  FORMAT(/,10X,'EM = ',D16.8,/,10X,'GM = ',D16.8,/,10X,'EX = ',D16.8,
           1 /10X,'GX = ',D16.8)
0264      114  FORMAT(/,10X,'CROSS SECTION DESCRIPTION',/,10X,'LAYER NO.', 2X,'LOAD
           1DING INTENSITY (LBS/IN)',5X,'H(I)',20X,'T(I)' )
0265      115  FORMAT(10X,16,3D20.8 )
0266      116  FORMAT(5X,'NET LOAD = ',D20.8)
0267      119  FORMAT(/,5X,'ROW = ',15, '   'ANS('I3.') = ',D20.8 )
0268      120  FORMAT( 1X,6D20.8 )
0269      121  FORMAT(1X,' - - - - - ')
           1 - - - - - ')
0270      130  FORMAT(/,10X,'AIJ      I = ',15 )
0271      131  FORMAT( 10X,5D20.8 )
0272      132  FORMAT(/,10X,'BIJ      I = ',15 )
0273      133  FORMAT(/,10X,'CIJ      I = ',15 )
0274      140  FORMAT(/,10X,'X = ',F16.3,/,5X,'      W      U      GAMMA      DWDX
           1      DUDX      DGAMDY D2WDX2      D2UDX2      ST1      ST2
           2 GAMMAB')
0275      141  FORMAT(12X,'MOM-M',10X,'V - M',13X,'V - B')
0276      142  FORMAT(1X,12,11F10.6 )
0277      143  FORMAT(10X,F10.2,5X,F10.2,5X,F10.2 )
0278      144  FORMAT(1X,'MOMENT SHOULD = ',F10.2,5X,'MOMENT = ',F10.2 )
0279      145  FORMAT(1X,'SHEAR SHOULD = ',F10.2,5X,'SHEAR = ',F10.2,5X,
           1' = (' , F10.2,' + ', F10.2,' ' )' /)
0280      END

```

```

0001      SUBROUTINE SOLVE(A,ANS,N,KEY)
0002      REAL*8 A(120,120),ANS(1),C
0003      NM1 = N - 1
0004      KEY = 0
0005      DO 10 IR = 1, NM1
0006      C = A(IR,IR)
0007      IF ( C .EQ. 0.0 ) GO TO 200
0008      DUM = DABS(C)
0009      IF ( DUM .LT. 1.0E-10 ) PRINT 400, IR
0010      C PRINT 300, IR,C
0011      C PRINT 301,(A(IR,IC),IC=1,N)
0012      IRP1 = IR + 1
0013      C PRINT 301, ( A(IRP1,IC),IC = 1,N)
0014      IRP2 = IR + 2
0015      C PRINT 301, ( A(IRP2,IC),IC = 1,N)
0016      C PRINT 301, ANS(IR),ANS(IRP1),ANS(IRP2)
0017      DO 9 I = IRP1, N
0018      IF ( A(I,IR) .EQ. 0. ) GO TO 9
0019      C = A(I,IR)/A(IR,IR)
0020      A(I,IR) = 0.
0021      DO 8 J = IRP1, N
0022      A(I,J) = A(I,J) - C*A(IR,J)
0023      ANS(I) = ANS(I) - C*ANS(IR)
0024      9 CONTINUE
0025      C PRINT 301, (A(IR,IC),IC = 1, N )
0026      C PRINT 301, ( A(IRP1,IC),IC=1,N)
0027      C PRINT 301,(A(IRP2,IC),IC=1,N)
0028      C PRINT 301, ANS(IR),ANS(IRP1),ANS(IRP2)
0029      10 CONTINUE
0030      DO 15 I = 1,N
0031      C PRINT 500, I,ANS(I)
0032      C PRINT 301,(A(I,J),J=1,N)
0033      C15 CONTINUE
0034      C PRINT 301, (ANS(IC),IC = 1,N)
0035      ANS(N) = ANS(N) / A(N,N)
0036      DO 20 I = 2,N
0037      IR = N+1 - I
0038      C = ANS(IR)
0039      IRP1 = IR + 1
0040      DO 18 J = IRP1,N
0041      C = C - A(IR,J)*ANS(J)
0042      20 ANS(IR) = C / A(IR,IR)
0043      C PRINT 301, ( ANS(IC),IC=1,N)
0044      GO TO 9000
0045      200 KEY = IR
0046      PRINT 401, KEY
0047      PRINT 301,(ANS(IC),IC=1,N)
0048      9000 RETURN
0049      300 FORMAT(//10X,I10,10X,D20.8 )
0050      301 FORMAT (1X,5D20.8 )
0051      400 FORMAT(//10X,'WARNING - DIAGONAL TERM ON ROW',I5,' IS SMALLER THAN
0052      1AN 1.0E-10 ')
0053      401 FORMAT(//10X,'***** SINGULAR MATRIX - DIAGONAL TERM IS ZERO',10X,
0054      1 I5 )
0055      500 FORMAT(/10X,I5,10X,D20.8 )
0056      END

```

REFERENCES

¹R. Marino, J. Citrolo, J. Frankenberg, PLT Toroidal Field Coil Power Tests, Proceedings of the Sixth Symposium on Engineering Problems of Fusion Research, San Diego, Cal. November 18-21, 1975 (p. 489-492).

²J. Frankenberg, PLT Power Test II, Private Communication, Plasma Physics Laboratory, Princeton, N. J., February, 1975.

³R. Marino, Shear Properties of a PLT Toroidal Field Coil Laminate, Private Communication, Plasma Physics Laboratory, Princeton, N. J., December, 1973.

⁴J. Bialek, Analysis of Laminated Beams for Use in Coil Design, Proceedings of the Sixth Symposium on Engineering Problems of Fusion Research, San Diego, Cal., November 18-21, 1975 (p. 950-954).

TABLE 1

$$nh + (n-1)t = H \quad R = \frac{(n-1)t}{H}$$

H = 5 inch

R = 0.05

R = 0.10

R = 0.15

n	h (in)	t (in)	h (in)	t (in)	H (in)	t (in)
2	2.375	0.25	2.25	0.5	2.125	0.75
3	1.583333	0.125	1.5	0.25	1.416666	0.375
4	1.1875	0.083333	1.125	0.166666	1.0625	0.25
5	0.95	0.0625	0.9	0.125	0.85	0.1875
6	0.791666	0.05	0.75	0.10	0.708333	0.15
8	0.59375	0.035714	0.5625	0.07142857	0.53125	0.107142857
12	0.3958333	0.0227272	0.375	0.04545454	0.35416666	0.068181818
20	0.2375	0.01315789	0.225	0.0263157	0.2125	0.03947368
30	0.158333	0.008620689	0.15	0.017241379	0.1465517	0.025862069

TABLE 2.

R = 0.05

Max Deflection vs. G_e and n (deflection in inches)

$n \rightarrow$ $G \downarrow$	2	3	4	5	6	8	12	20	30
375000.	0.125488	0.127656	0.128725	0.129366	0.129794	0.130330	0.130867	0.131298	0.131515
100000.	0.126306	0.128295	0.129299	0.129906	0.130314	0.130827	0.131343	0.131759	0.131968
10000.	0.136154	0.136096	0.136319	0.136523	0.136686	0.136918	0.137179	0.137408	0.137528
1000.	0.216330	0.209402	0.204369	0.201387	0.199484	0.197234	0.195144	0.193593	0.192858
100.	0.452480	0.636164	0.709579	0.736914	0.746631	0.749244	0.742575	0.732034	0.725423
50.	0.507046	0.839725	1.047946	1.153260	1.226186	1.280850	1.303911	1.300135	1.291920
10.	0.565747	1.173679	1.850724	2.500332	3.070120	3.927623	4.822034	5.359399	5.504592

TABLE 3.
Max Deflection vs. G_e and n (deflection in inches)

$R = 0.10$

$G_e \rightarrow$ $n \downarrow$	2	3	4	5	6	8	12	20	30
375000.	0.125897	0.130462	0.132686	0.134019	0.134911	0.136029	0.137154	0.138060	0.138515
100000.	0.127509	0.131714	0.133812	0.135083	0.135936	0.137012	0.138099	0.138976	0.139418
10000.	0.146675	0.146943	0.147571	0.148090	0.148492	0.149053	0.149675	0.150214	0.150497
1000.	0.285919	0.284303	0.278109	0.273811	0.270914	0.267378	0.264017	0.261489	0.260284
100.	0.574668	0.911902	1.104365	1.203787	1.255214	1.297155	1.312042	1.306132	1.298246
50.	0.624237	1.137125	1.548667	1.832610	2.018578	2.219217	2.352114	2.394436	2.394517
10.	0.672494	1.437954	2.366305	3.353189	4.316438	5.992962	8.169095	9.856207	10.437098

TABLE 4.
Max. deflection vs. G_e and n (deflection in inches)

$R = 0.15$

$n \rightarrow$ $G_{\downarrow}$	2	3	4	5	6	8	12	20	30
375000.	0.126432	0.133669	0.137150	0.139238	0.140637	0.142395	0.144168	0.145599	0.133568
100000.	0.128849	0.135507	0.138808	0.140808	0.142153	0.143852	0.145573	0.146967	0.134843
10000.	0.156805	0.157803	0.159033	0.159987	0.160709	0.161704	0.162796	0.163737	0.150475
1000.	0.349154	0.354363	0.348477	0.343720	0.340369	0.336195	0.332189	0.329173	0.304777
100.	0.696153	1.155737	1.453212	1.624639	1.722160	1.812594	1.859371	1.865118	1.748101
50.	0.749723	1.410151	1.988111	2.421609	2.726462	3.082598	3.348554	3.456865	3.269109
10.	0.800331	1.726827	2.880548	4.149118	5.434599	7.800255	11.167997	14.096046	14.299311

TABLE 5.

Key

R = 0.05

Maximum tensile and shear strains

Applied moment = 5000. in lbs.

Applied shear = 100. lbs.

Tensile (copper)

Shear (epoxy)

G	n	2	3	4	5	6	8	12	20	30
375000.		0.000075	0.000076	0.000077	0.000077	0.000078	0.000078	0.000078	0.000079	0.000079
		0.000076	0.000067	0.000075	0.000072	0.000075	0.000076	0.000076	0.000076	0.000076
100,000.0		0.000075	0.000076	0.000077	0.000077	0.000078	0.000078	0.000078	0.000070	0.000079
		0.000282	0.000250	0.000282	0.000271	0.000283	0.000283	0.000283	0.000284	0.000284
10000.		0.000077	0.000077	0.000078	0.000078	0.00-078	0.000078	0.000079	0.000079	0.000070
		0.002668	0.002483	0.002777	0.002683	0.002799	0.002898	0.002817	0.002823	0.002826
1000.		0.000092	0.000087	0.000085	0.000084	0.000084	0.000084	0.000084	0.000084	0.000084
		0.020385	0.021719	0.025298	0.025089	0.026339	0.026748	0.027071	0.027262	0.027334
100.		0.000139	0.000148	0.000143	0.000138	0.000133	0.000128	0.000125	0.000123	0.000122
		0.070400	0.125477	0.167338	0.185013	0.200591	0.215236	0.228324	0.237119	0.240554
50.		0.000150	0.000179	0.000183	0.000179	0.000174	0.000167	0.000159	0.000155	0.000154
		0.081890	0.174158	0.233592	0.301711	0.337269	0.376825	0.413188	0.438898	0.449638
10.		0.000162	0.000228	0.000276	0.000307	0.000325	0.000337	0.000329	0.000310	0.000302
		0.094242	0.253808	0.452708	0.651362	0.835059	1.128399	1.476938	1.761533	1.895361

TABLE 6.

key

R = 0.10

Maximum tensile and shear strains

Applied moment = 5000. in #

Applied shear = 100. #

Tensile (copper)

Shear (epoxy)

n	2	3	4	5	6	8	12	20	30
G 375000.	0.000075	0.000078	0.000079	0.000080	0.000081	0.000081	0.000082	0.000082	0.000083
	0.000075	0.000066	0.000075	0.000072	0.000075	0.000075	0.000075	0.000076	0.000076
100000.	0.000075	0.000078	0.000079	0.000080	0.000081	0.000081	0.000082	0.000083	0.000083
	0.000279	0.000247	0.000210	0.000269	0.000280	0.000281	0.000282	0.000283	0.000284
10000.	0.000079	0.000080	0.000081	0.000081	0.000082	0.000082	0.000083	0.000083	0.000084
	0.002577	0.002398	0.002723	0.002642	0.002760	0.002776	0.002791	0.002803	0.002808
1000.	0.000105	0.000097	0.000094	0.000093	0.000093	0.000092	0.000092	0.000092	0.000093
	0.017837	0.020241	0.023819	0.023948	0.025219	0.025807	0.026297	0.026602	0.026721
100.	0.000163	0.000186	0.000187	0.000181	0.000176	0.000168	0.000161	0.000159	0.000158
	0.048541	0.096736	0.136526	0.158287	0.174895	0.192919	0.209522	0.221292	0.226198
50.	0.000173	0.000219	0.000236	0.000239	0.000236	0.000227	0.000215	0.000208	0.000206
	0.053793	0.123889	0.192793	0.242618	0.280679	0.327962	0.313377	0.406742	0.421496
10.	0.000182	0.000262	0.000328	0.000378	0.000414	0.000452	0.000463	0.000440	0.000424
	0.058903	0.160104	0.293950	0.441490	0.590683	0.861168	1.241119	1.595153	1.770150

TABLE 7.

key

R = 0.15

Maximum tensile and shear strains

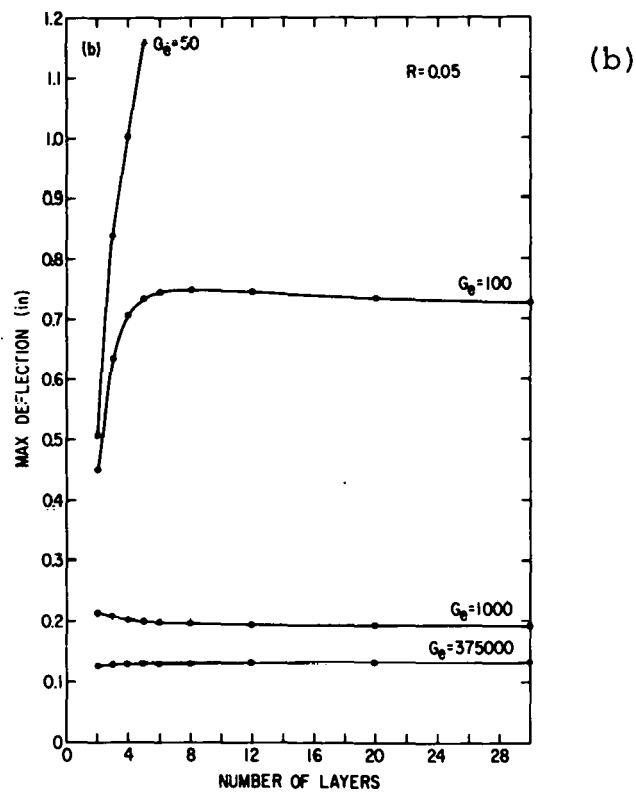
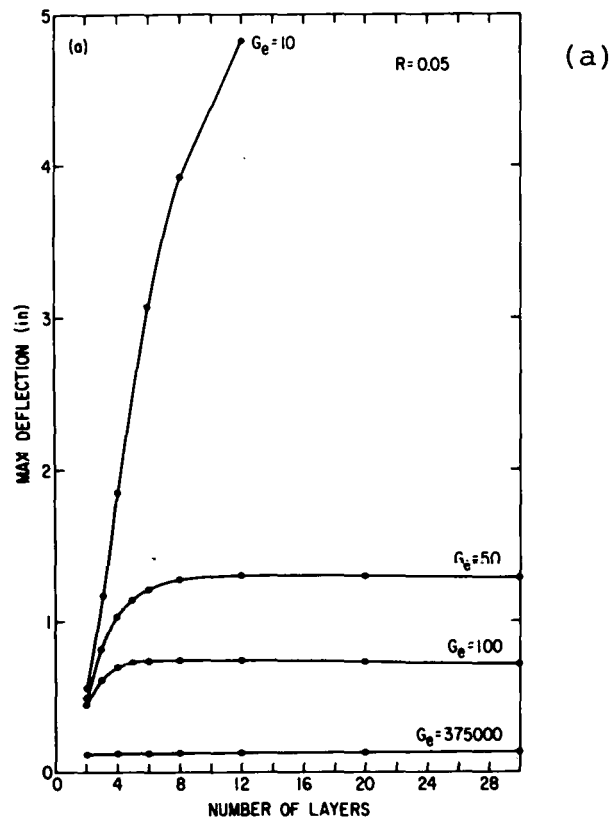
Applied moment = 5000. in lbs.

Applied shear = 100. lbs

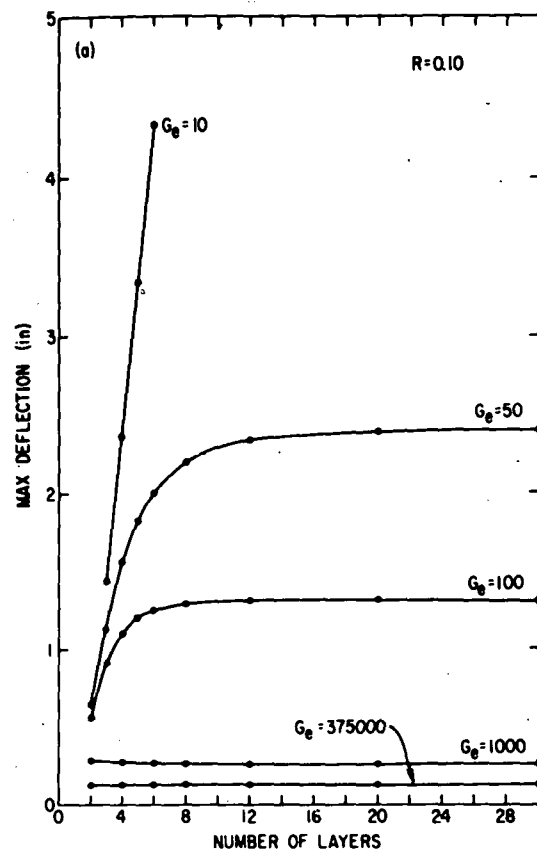
Tensile (copper)

Shear (epoxy)

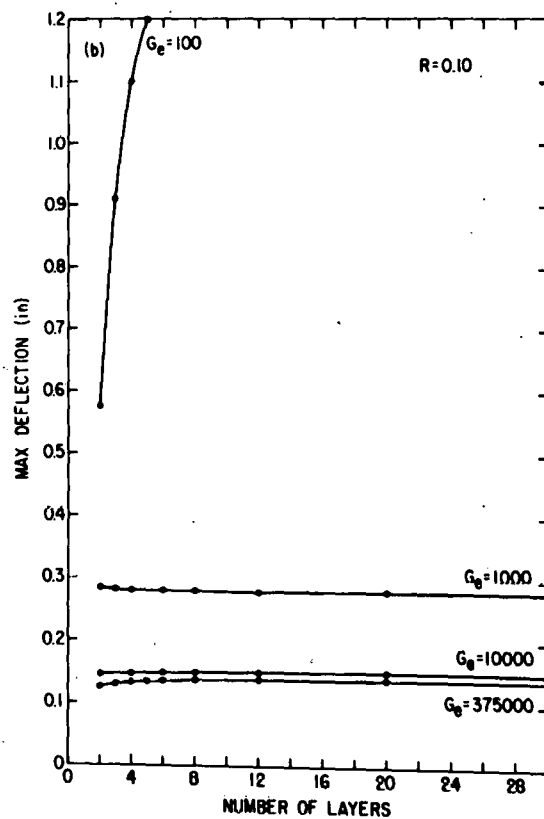
h	2	3	4	5	6	8	12	20	30
G 375000.	0.000075	0.000080	0.000082	0.000083	0.000084	0.000085	0.000086	0.000087	0.000082
	0.000074	0.000065	0.000074	0.000071	0.000074	0.000075	0.000075	0.000075	0.000073
100000.	0.000076	0.000080	0.000082	0.000083	0.000084	0.000085	0.000086	0.000087	0.000082
	0.000275	0.000245	0.000276	0.000266	0.000278	0.000280	0.000281	0.000282	0.000275
10000.	0.000080	0.000082	0.000084	0.000085	0.000086	0.000086	0.000087	0.000088	0.000083
	0.002517	0.002357	0.002680	0.002608	0.002727	0.002749	0.002770	0.002786	0.002713
1000.	0.000116	0.000107	0.000103	0.000101	0.000101	0.000101	0.000101	0.000101	0.000096
	0.016652	0.019422	0.022914	0.023218	0.022547	0.025159	0.025747	0.026128	0.025494
100.	0.000184	0.000217	0.000221	0.000216	0.000210	0.000200	0.000192	0.000189	0.000180
	0.041488	0.085109	0.122584	0.145359	0.162240	0.181603	0.199737	0.212907	0.211611
50.	0.000195	0.000252	0.000279	0.000287	0.000285	0.000276	0.000260	0.000251	0.000238
	0.045310	0.105772	0.167908	0.216627	0.254807	0.304734	0.354211	0.391196	0.395030
10.	0.000205	0.000296	0.000273	0.000437	0.000484	0.000541	0.000569	0.000544	0.000499
	0.048924	0.131462	0.242337	0.368427	0.500079	0.750871	1.131278	1.513544	1.650736



764770
Fig. 4(a) and (b).

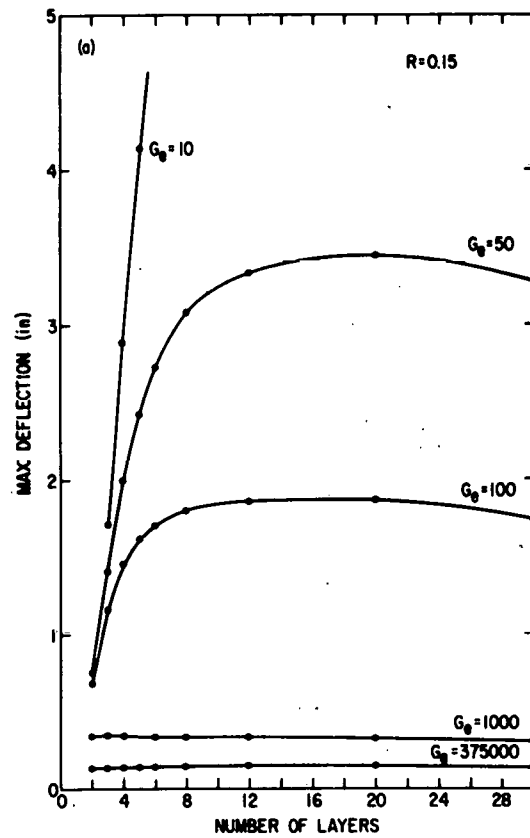


(a)

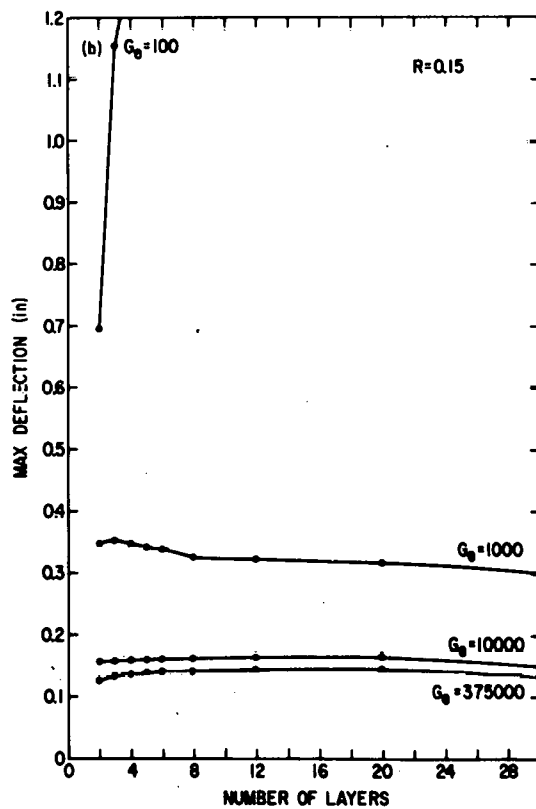


(b)

764789
Fig. 5(a) and (b).

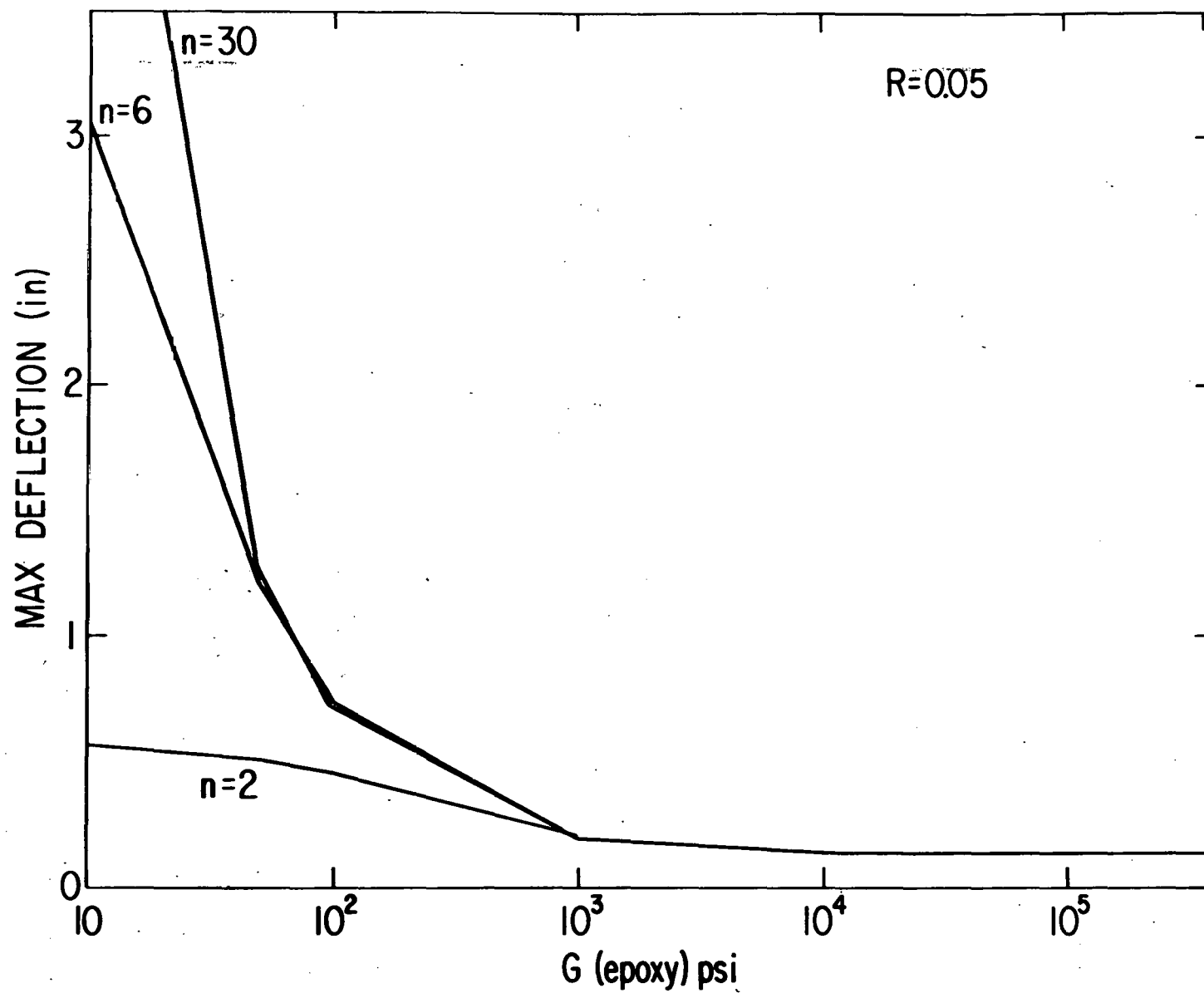


(a)

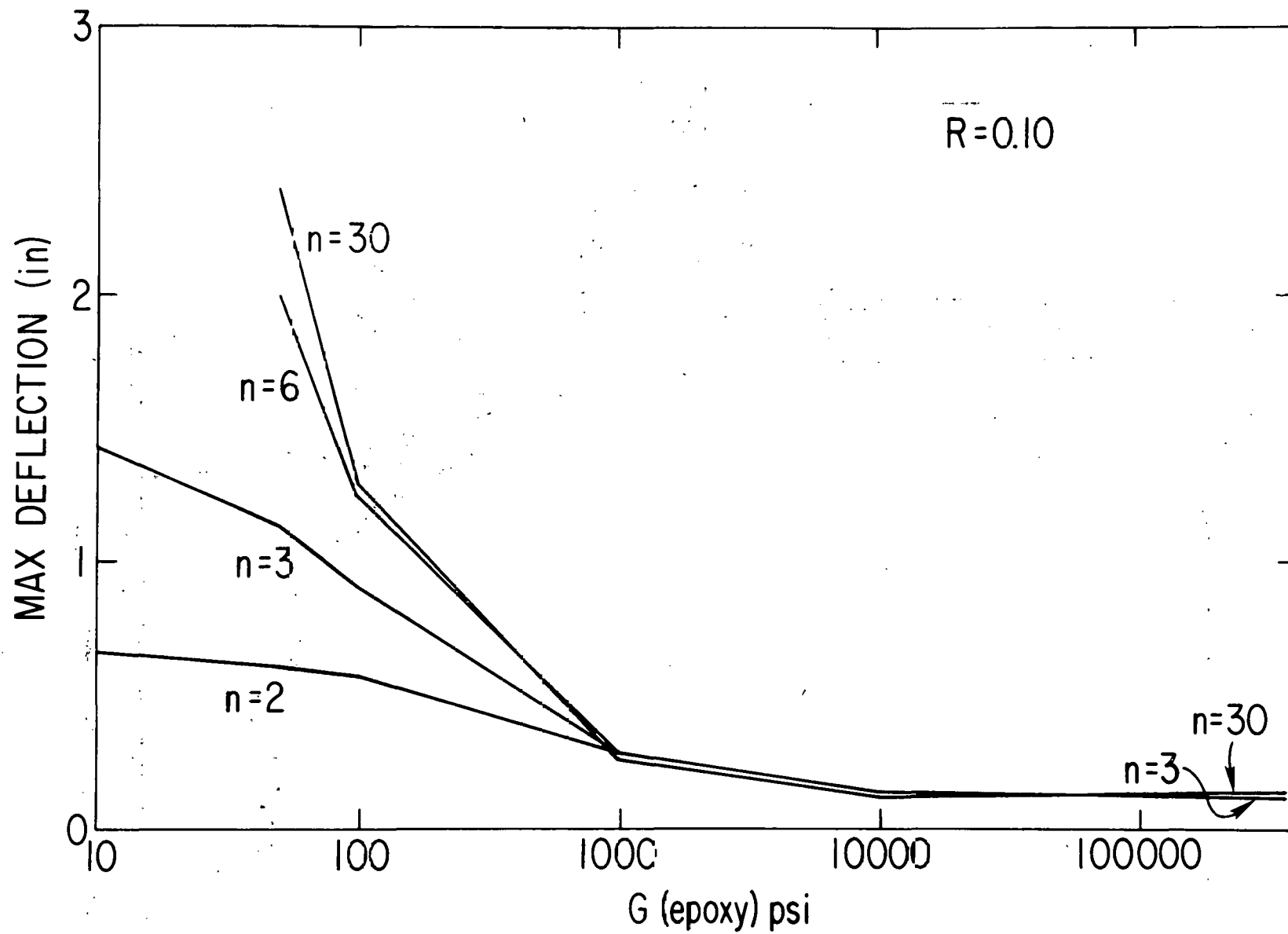


(b)

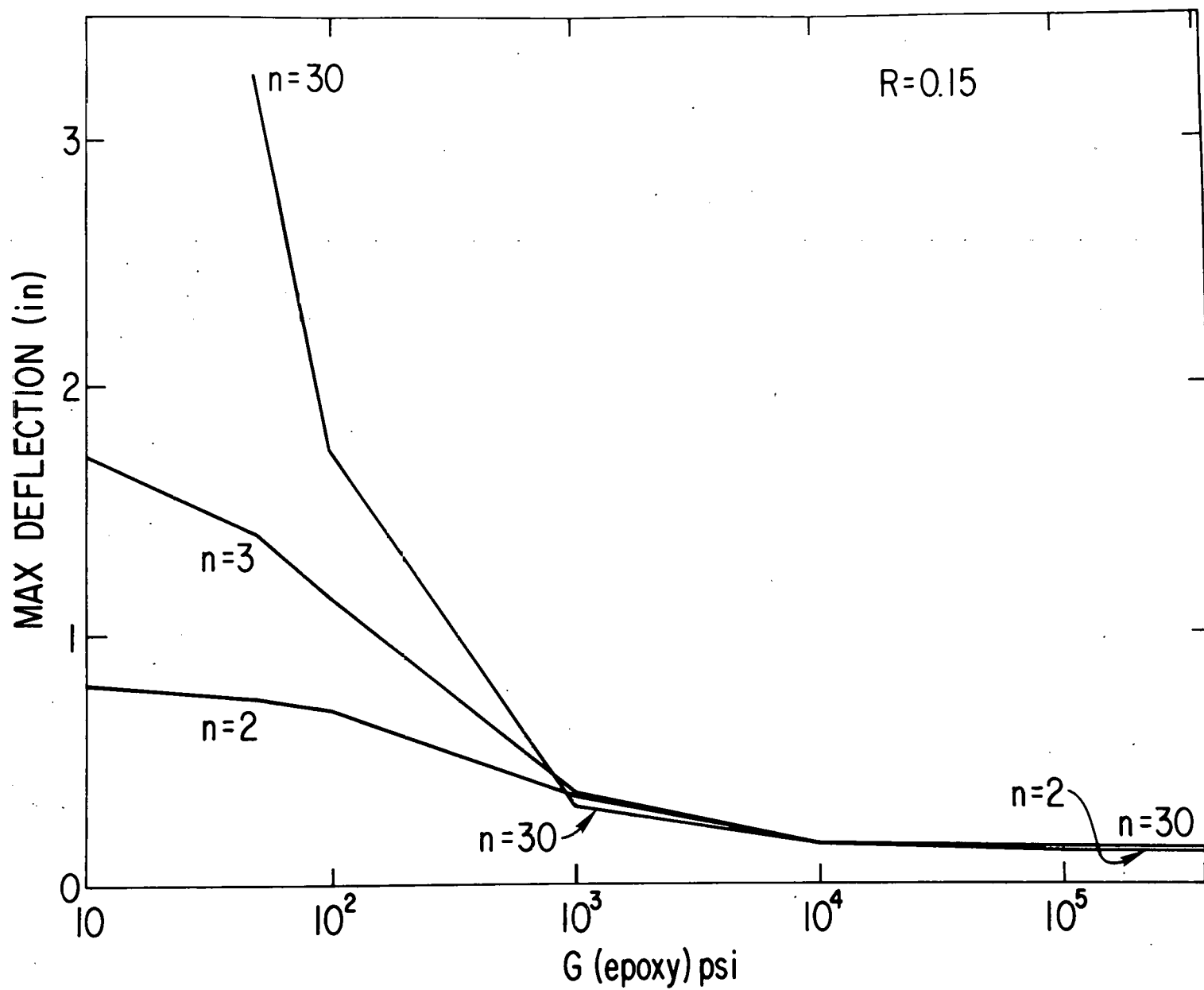
764768
Fig. 6(a) and (b).



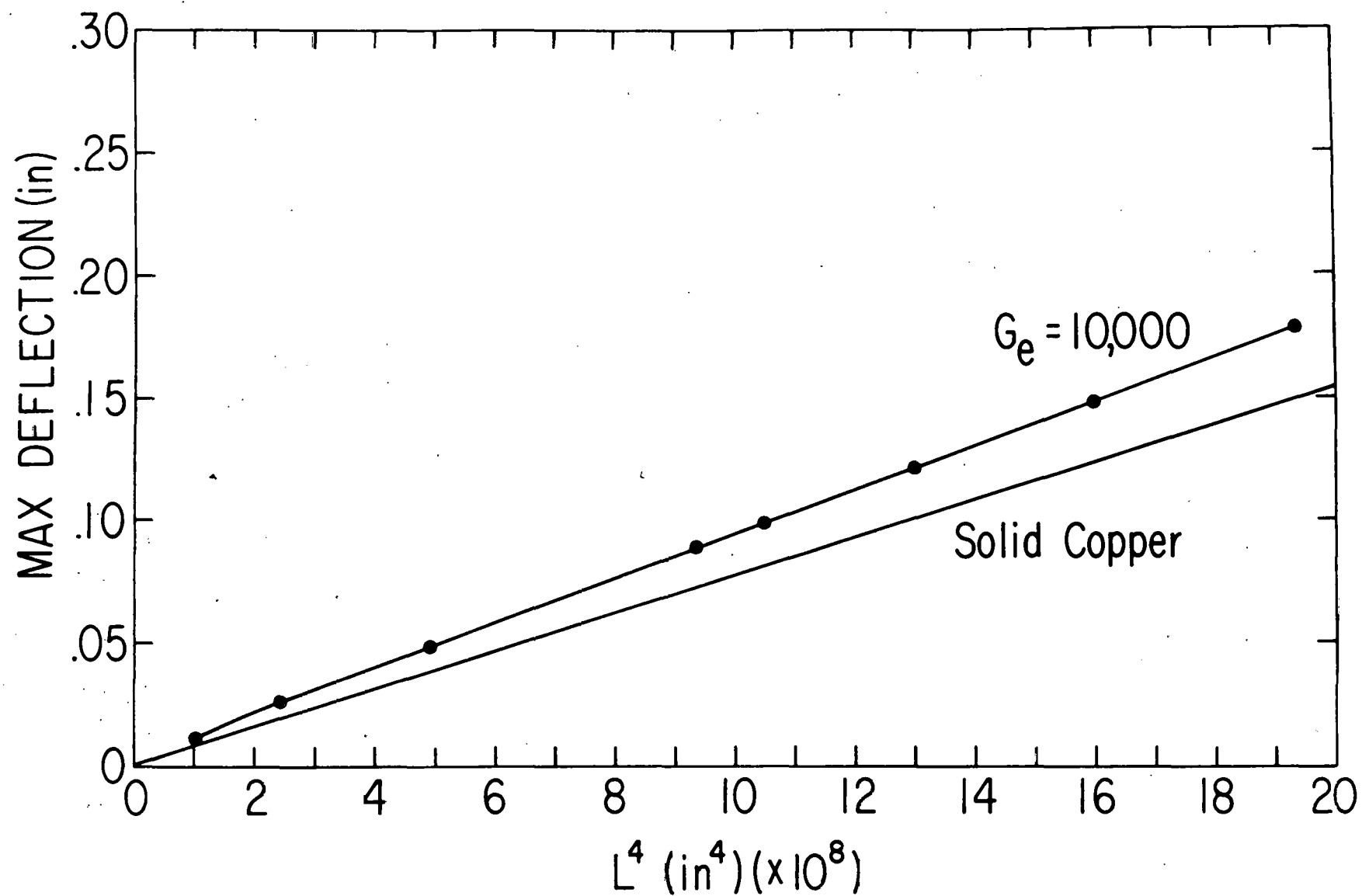
764780
Fig. 7.



764777
Fig. 8.



764771
Fig. 9.



764779
Fig. 10.

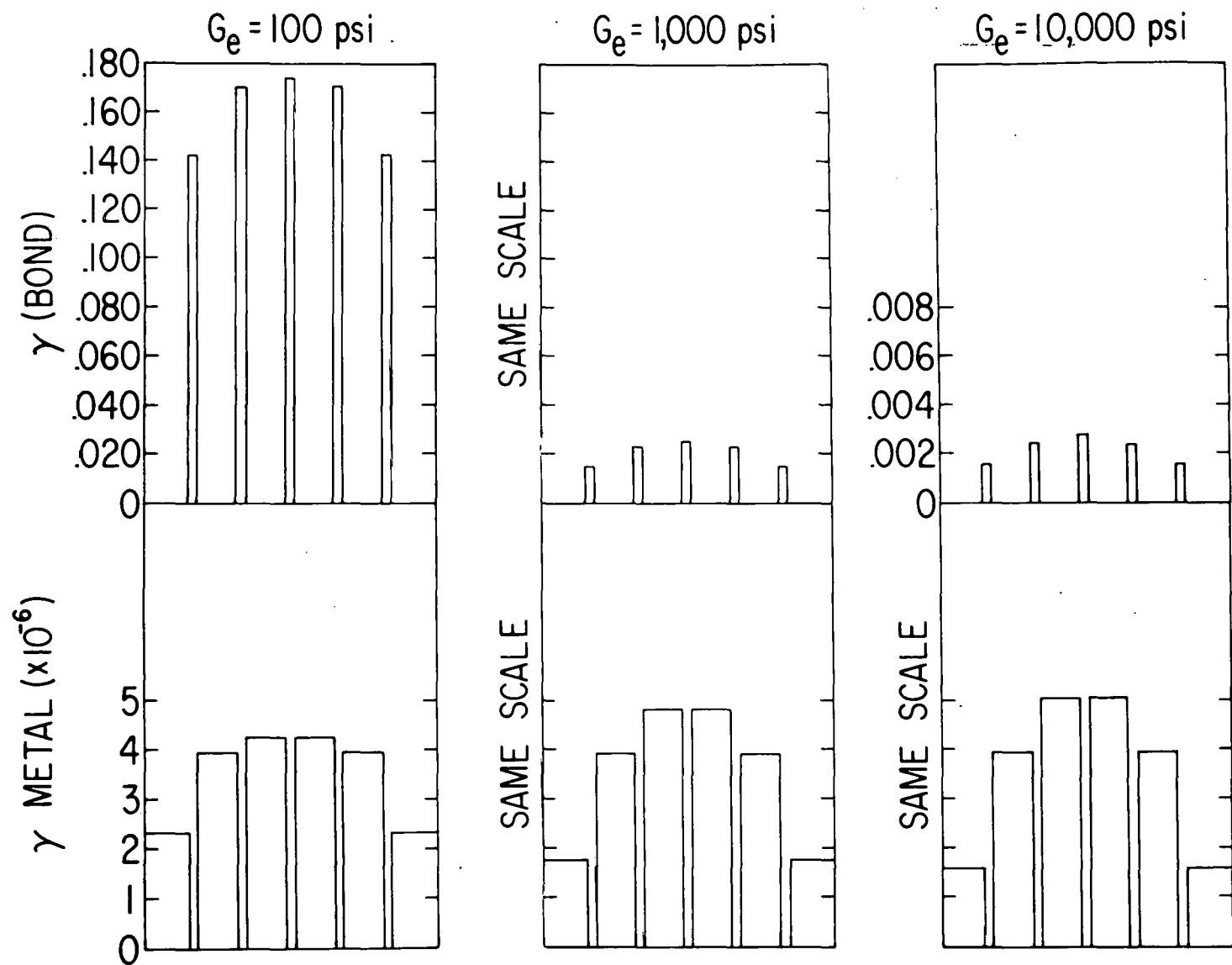
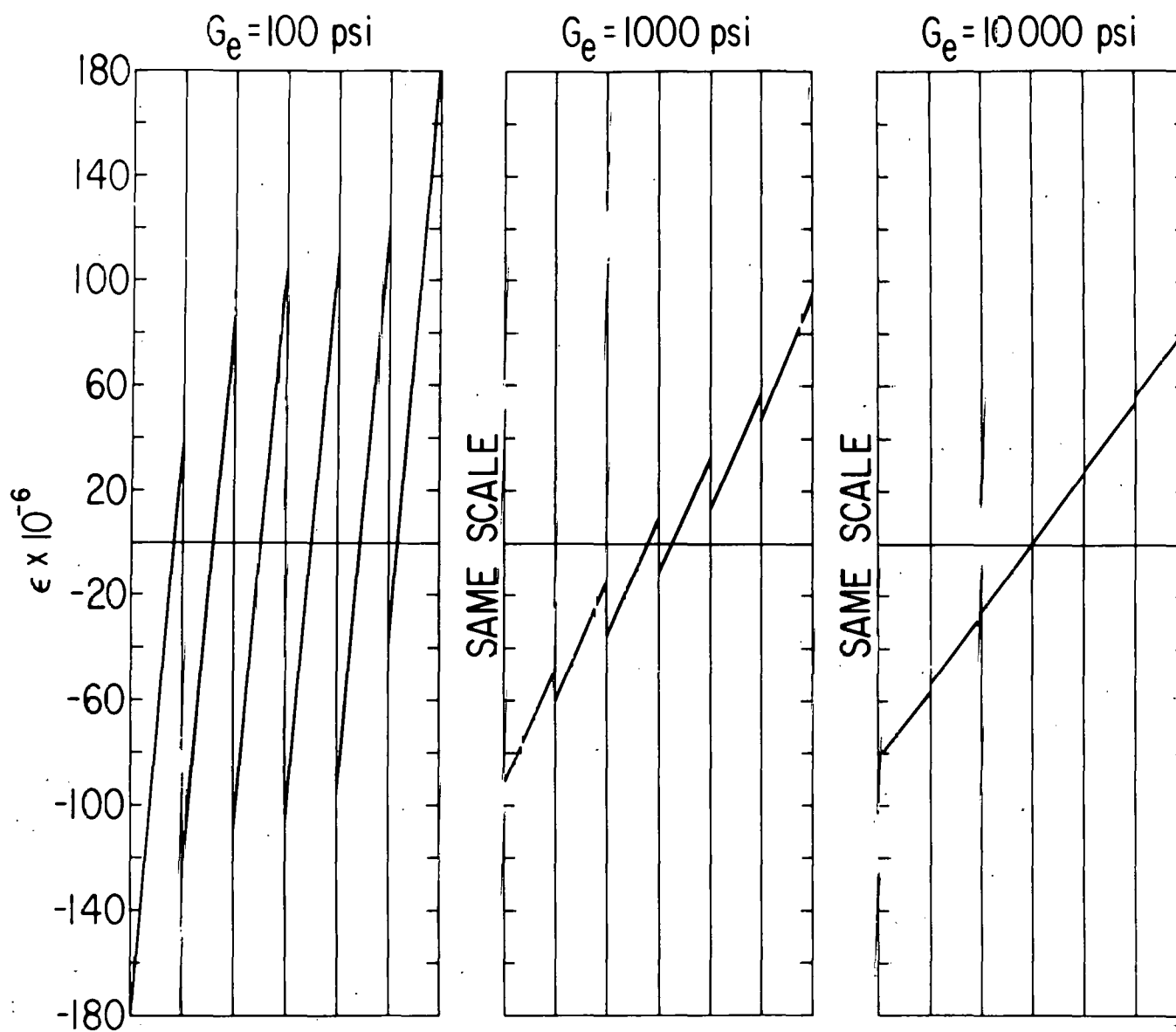


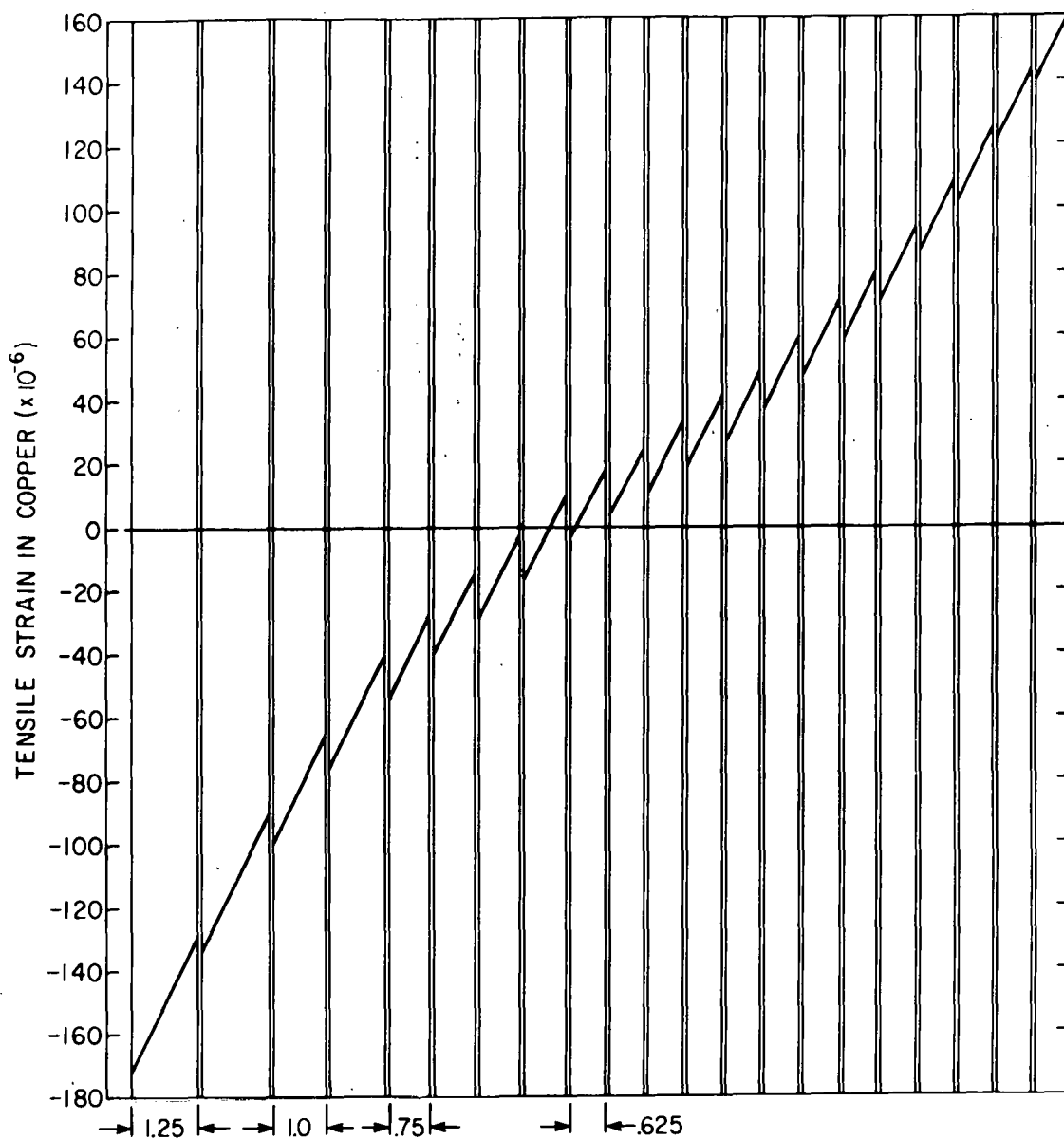
Fig. 11. Shear strain in a six layer beam ($R = 0.10$).



764775
Fig. 12. Tensile strain in a six layer beam ($R = 0.10$).

Applied Moment
is 105000 in #

Moment From
Tensile Strains
is ≈ 104902 in #

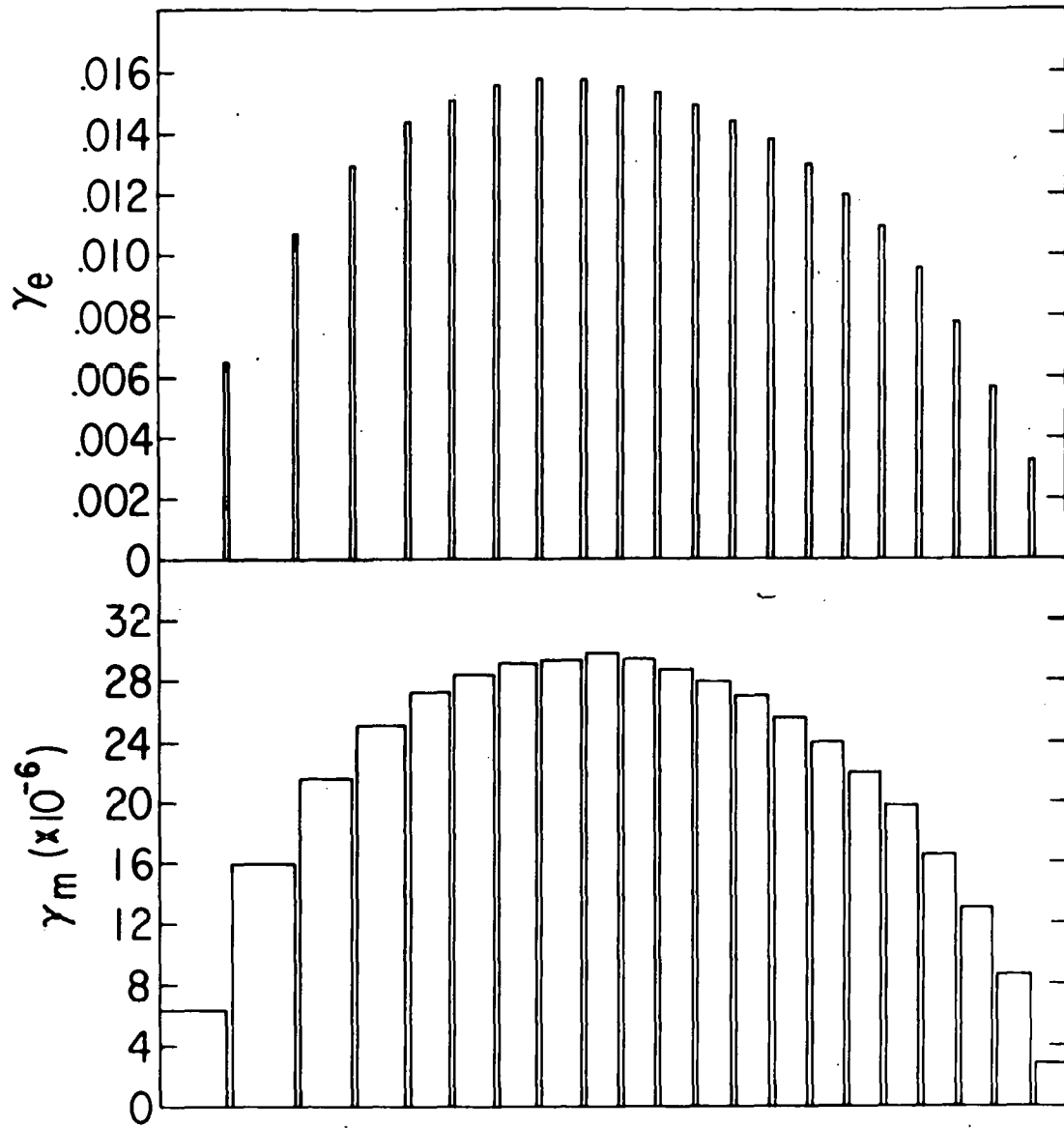


764773

Fig. 13. Tensile strain vs. layer in the PLT cross section.

Applied Shear
is 2100 #

Shear From
Total Strain
is 2201 #



764776

Fig. 14. Shear strain vs. layer in the PLT cross section.