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What is LAMPF II?

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WHAT IS LAMPF II?

by

H. A. Thiessen

ABSTRACT

The present conception of LAMPF II is a high-intensity 16-GeV synchrotron injected by the LAMPF 800-MeV H^- beam. The proton beam will be used to make secondary beams of neutrinos, muons, pions, kaons, antiprotons, and hyperons more intense than those of any existing or proposed accelerator. For example, by taking maximum advantage of a thick target, modern beam optics, and the LAMPF II proton beam, it will be possible to make a negative muon beam with nearly 100% duty factor and nearly 100 times the flux of the existing Stopped Muon Channel (SMC). Because the unique features of the proposed machine are most applicable to beams of the same momentum as LAMPF (that is, <2 GeV/c), it may be possible to use most of the experimental areas and some of the auxiliary equipment, including spectrometers, with the new accelerator. The complete facility will provide improved technology for many areas of physics already available at LAMPF and will allow expansion of medium-energy physics to include kaons, antiprotons, and hyperons. When LAMPF II comes on line in 1990 LAMPF will have been operational for 18 years and a major upgrade such as this proposal will be reasonable and prudent.

I. NUCLEAR PHYSICS AT LAMPF II

The areas of nuclear physics that can be studied with LAMPF II are listed below.

- Hypernuclei - In hypernuclei we will be able to study nuclear structure with one or more strange quarks in the nuclear bag.

- Hyperons - By studying the decays of the excited states we can determine the wave functions of the three quarks that make up the hyperons.
- Kaon-Nucleus Scattering - The kaon is the hadron with the longest mean-free path in nuclear matter and as such will be an excellent probe of nuclear structure.
- Pion-Nucleus Scattering - we can use 0.5- to 1.0-GeV pions for nuclear structure studies with the advantage that the mean-free path will be much longer and the selectivity for magnetic transitions will be significantly better than at the Energetic Pion Channel Spectrometer (EPICS).
- Hadron Resonances in Nuclear Matter - Recent work has shown that Δ -hole modes are an important part of nuclear structure. Using kaons will make it possible to excite the $Y^*(1520)$, which has a much narrower width and longer lifetime. This should be a much cleaner case to study than the Δ .
- Muon Capture and Muon Spin Rotation (μ SR) - The high flux of muons from LAMPF II will make possible significantly more sensitive experiments in muon capture, which studies the weak interaction in nuclear matter at a momentum transfer comparable to the muon mass, and in muon spin rotation, which is useful for solid-state applications.
- Neutrino-Nucleus Scattering - By using the high flux of neutrinos from LAMPF II we can study the spin and isospin structure of the charged current by scattering from selected light nuclei.
- Exotic Atoms - In addition to K^- , $\bar{\Sigma}^-$, and p^- atoms that have already been studied at existing machines, it will be possible to use intense beams of kaons and antiprotons to make Ξ^- , Ω^- , and $\bar{\Sigma}^-$ atoms that have not yet been seen. By studying the atomic spectra we can determine masses, magnetic moments, and the low-energy baryon-nucleus potential for these rare particles.

II. PARTICLE PHYSICS AT LAMPF II

The classes of particle physics that will be studied at LAMPF II are listed below.

- Rare Kaon Decays - There are many rare-decay modes that provide extremely sensitive tests of the standard Weinberg-Salam-Glashow gauge model of the electromagnetic and weak interactions. Other decay modes are sensitive to proposed extensions of this model. In many cases the measurements of the

rare-decay modes possible with LAMPF II are more sensitive than any other experiment planned at any existing or proposed accelerator.

- Charge-Parity (CP) Violation in Kaon Decays - The decay of the kaon is the only verified example of violation of CP symmetry known. The origin of this CP violation has important consequences for unification. Our knowledge of all the observables of CP violation is presently limited by statistical uncertainties; hence high-intensity beams from LAMPF II will have a major impact on this field.
- Neutrino Oscillations and Neutrino-Electron Scattering - The high flux of neutrinos of variable energy from LAMPF II will make possible neutrino mass searches with 100 times the sensitivity of those performed at existing accelerators. It will also be possible to study neutrino- and antineutrino-electron scattering with sufficient precision and statistics to determine the angular distribution. Such precise experiments are important tests of the standard theory of electroweak interactions and its possible extensions.
- Muonium - The intense beam of muons from LAMPF II will make possible the study of muonium with precise tests of quantum electrodynamics at a level where weak-interaction effects become visible.
- Pion, Kaon, and Hyperon-Nucleon Scattering - Although this field has been studied in the past, many important puzzles remain. In particular, there are no spin-rotation experiments; the kaon-nucleon polarization is poorly known, especially at low energy; and the hyperon-nucleon problem is practically untouched. Much work remains to be done to clean up hadron spectroscopy.
- Antiproton Physics - It will be difficult to compete with the high-quality cooled beam of antiprotons from the low-energy antiproton ring (LEAR) at CERN. However, possibilities exist for much higher intensity beams of antiprotons at LAMPF II. An important experiment that will make use of a high-intensity, high-momentum beam is $\bar{p} + p \rightarrow e^+ + e^-$, which will measure the proton form factor in the time-like region inaccessible in electron scattering.

III. LAMPF II ACCELERATOR

The basic machine for LAMPF II is a rapid-cycling 16-GeV synchrotron. The high flux will be obtained by operating at 60 Hz, 100 times faster than existing machines. With 10^{13} protons per pulse, which has already been achieved at many laboratories, we will have an average current of 100 μ A. The major technical problems to be solved are the large amount of rf power required and the minimizing of beam losses during the acceleration cycle.

To achieve a good duty factor the rapid cycler will inject its beam into a dc-stretcher ring. Slow extraction, with nearly 100% duty factor, will be done from the stretcher ring. We also require good rf timing capability of better than 1-ns pulse width at 50 MHz or lower repetition rate. This good timing will make possible particle identification without detectors in the beam and may also make possible rf separated beams. Providing both high intensity and good timing capability simultaneously is a difficult but probably manageable task for our accelerator designers.

No decision has been made on the sharing of the primary beam from LAMPF II. For planning purposes we assume there will be four primary beams, each receiving an average of 25 μ A. One of the four beams will be a fast-extracted variable-energy beam dedicated to neutrino physics; the remaining three beams will be slow-extracted beams, which will be used to produce secondary beams. We assume that there will be three thick targets, each of more than one interaction length.

IV. LAMPF II EXPERIMENTAL AREAS

It may be possible to locate the accelerator on the site in such a way that the existing experimental areas can be provided with 16-GeV proton beams. A site layout that meets this requirement is shown in Fig. 1. The accelerator would be located in a tunnel \sim 9.14 m below the present beam elevation. We are also considering constructing completely new experimental areas, but this plan is less advanced.

A possible plan for the experimental areas is shown in Fig. 2. It is clear from such layouts that the LAMPF II experimental areas will be comparable in scope to those already in use at LAMPF. The layout shows all the presently envisaged beam lines except the neutrino line, which will be a short line from one of the straight sections of the accelerator. Of course this is only a preliminary sketch and is subject to considerable change.

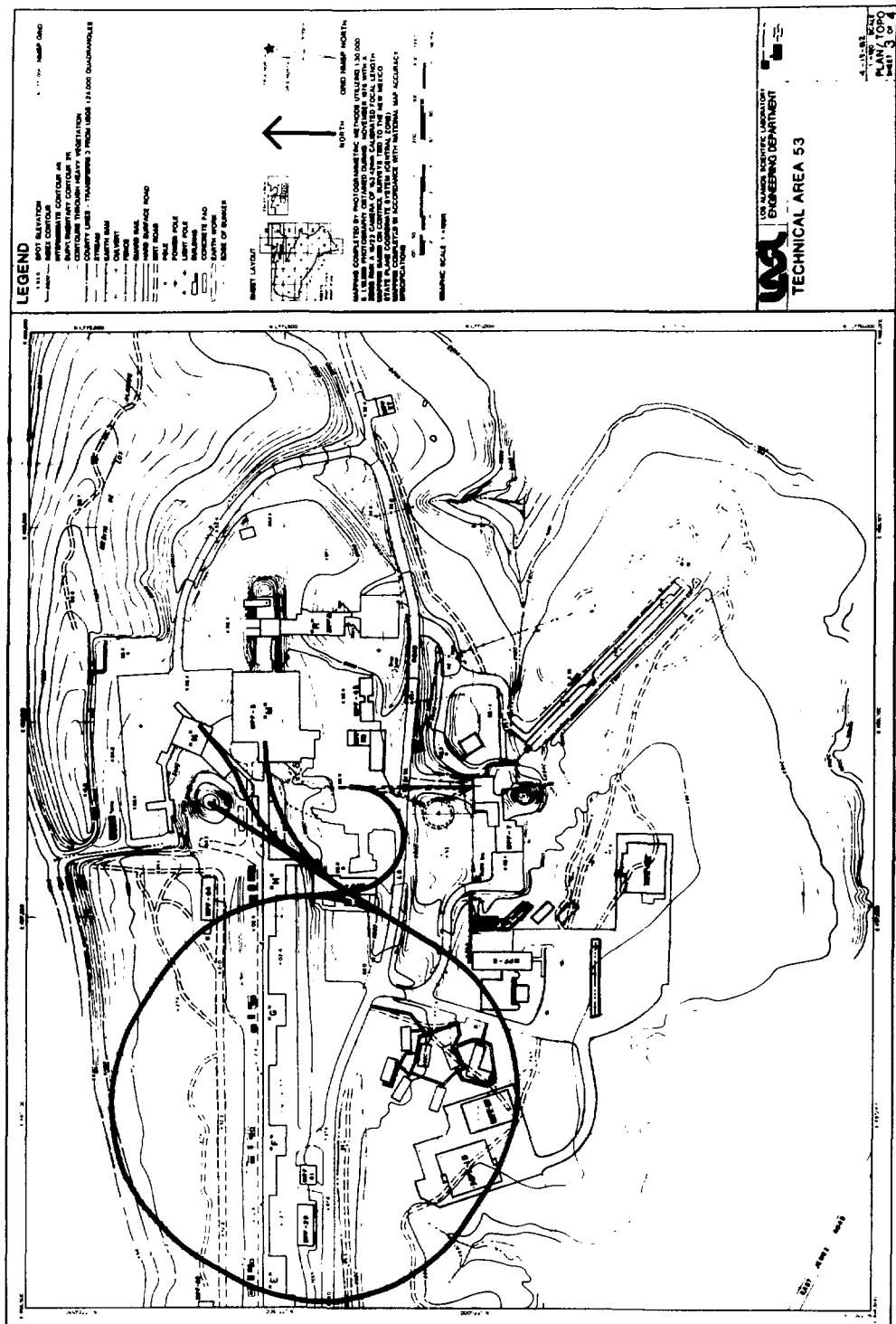


Fig. 1.
Proposed site layout for LAMPF II.

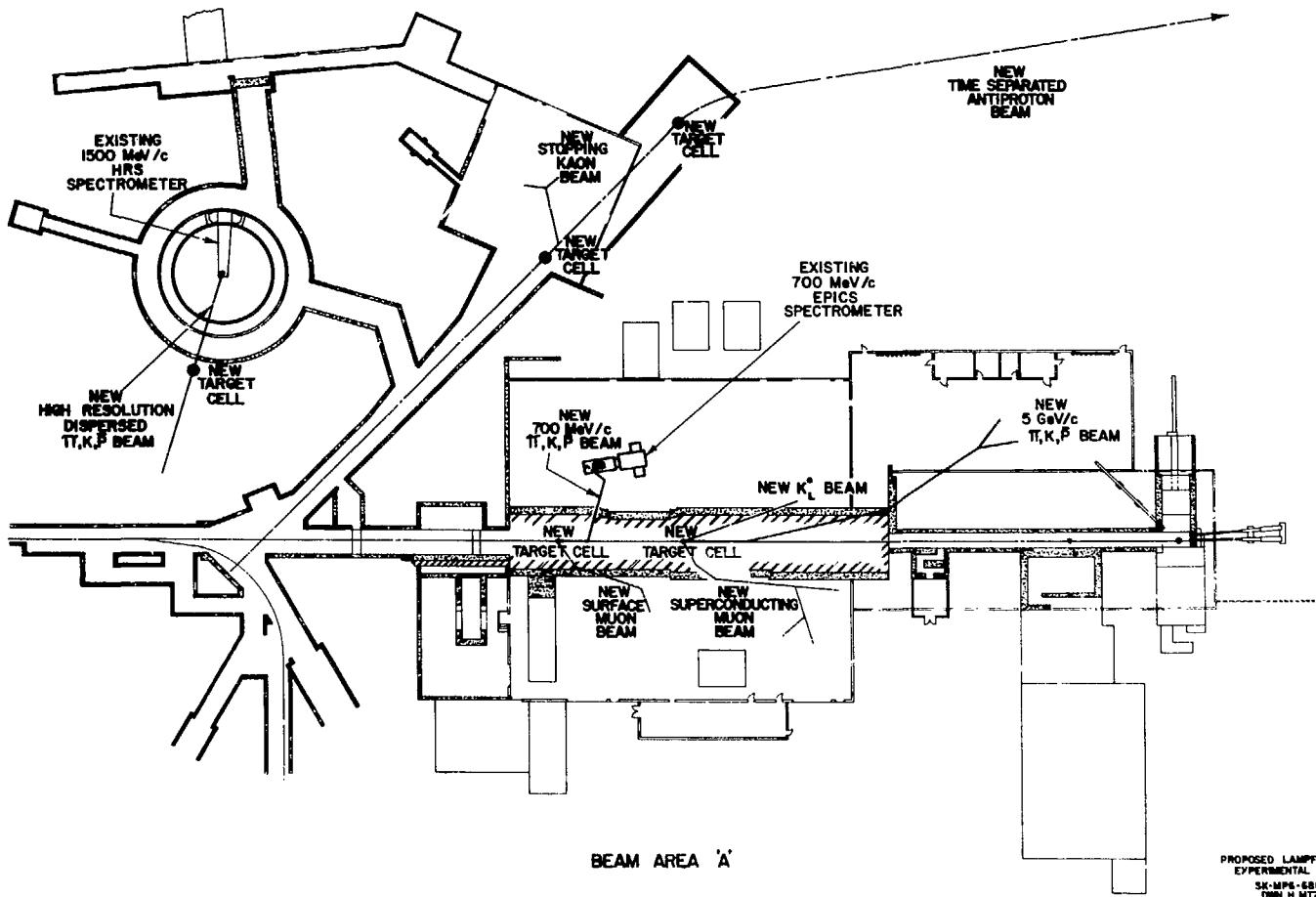


FIG. 2.
Proposed plan for new experimental areas.

V. EXTENSIONS TO THE BASIC PLAN

We are actively pursuing two directions in addition to the plans mentioned above. First, we are looking at the possibility of providing a cooled antiproton beam such as LEAR at CERN. It will be difficult to provide a competitive proposal unless we come up with a new-idea antiproton beam. The Nucleon-Antinucleon Working Group is studying this problem.

A second possibility is that we construct a colliding-beam facility for polarized protons and heavy ions. Such a facility would provide a unique capability if we can achieve adequate luminosity. If the interference with the operation of the other facilities can be kept to a few hours a day we would consider this a feasible option. Nick DiGiacomo of P Division is looking at the feasibility of a collider facility.

VI. COST AND TIME SCALE

We hope to complete a plan for the accelerator and experimental areas during 1983; a proposal could be ready for submission at the end of 1983. If approved, construction could start in FY 1986 and operation could occur in 1990. A detailed cost estimate must await completion of our plan; however, early estimates of \$75M for the accelerator and \$75M for experimental areas have been confirmed by several consultants.

VII. USER INPUT FOR LAMPF II

Several working groups of prospective LAMPF II users have been organized. Those presently active are as follows.

1. Nuclear Physics
2. Muons
3. Rare Kaon Decays
4. CP Violation
5. Hyperons
6. Nucleon/Antinucleon
7. Nuclear Chemistry
8. Neutrinos

Each working group will hold several meetings during the next year to select the most interesting few experiments in each experiment area and to study each experiment carefully enough to demonstrate its feasibility. The list of working groups is arbitrary -- changes can be made if desired. The next meetings of the

working groups will be parallel sessions at the LAMPF II Workshop, July 19-22, 1982.

VIII. NUCLEAR PHYSICS WORKING GROUP

At the February meeting Carl Dover presented a review of kaon-nucleus physics, Lee Roberts reviewed kaonic atoms, and Jim Carr discussed his first results on the pion-nucleus effective interaction. The pion-nucleus scattering calculations were the most important new results shown at this session and indicated that significant sensitivity to neutron or proton particle-hole states is possible, with magnetic transitions enhanced near 500 MeV and natural-parity transitions dominating near 1000 MeV. H. A. Thiessen reviewed the proposed 700-MeV/c dispersed kaon beam for use with the EPICS spectrometer, Harald Enge presented his proposal for a low-energy kaon beam separated by an absorber, and Donald Lobb showed his design for an achromatic stopping pion beam proposed for the TRIUMF kaon factory.

At a second meeting held in March, Leonard Kisslinger discussed the University of Washington model for short-range effects in hadron-nucleus scattering based on a quark model, Ben Gibson reviewed the field of low-mass hypernuclei, and Peter Mulders presented a discussion of quark-model predictions for hadron spectroscopy. On the experimental side H. A. Thiessen presented first results for a high-momentum dispersed kaon beam proposed for the High-Resolution Spectrometer (HRS) and Ed Hungerford began a discussion of the experimental problems of observing cascade hypernuclei. Morgan May reviewed the earlier Thiessen proposal for a low-momentum kaon beam and, using a good deal of experience from Brookhaven, concluded that this proposal is practical.

The problem facing this working group is to narrow the range of experiments to the point where they can be accomplished with a small number (say, two) of beam lines. The problem of separators for a high-momentum beam is still under discussion, as is the question of the origin of backgrounds in a separated beam. The issue of coincidence experiments has not yet been addressed.

IX. \bar{N} AND \bar{N} ACTIVITIES

Can we make an antiproton beam that would compete with LEAR? LAMPF II could produce more antiprotons than CERN (especially if we had 32-GeV protons instead of 16 GeV), but could we cool them? CERN and Fermi National Accelerator Laboratory (FNAL) are probably close to the limit of stochastic cooling (10^{11}

\bar{p}/day), but Peter McIntyre (Texas A&M University) and Billy Bonner (Los Alamos) are investigating possible electron-cooling schemes to cool $10^{12} \bar{p}/\text{day}$.

A good beam of polarized antiprotons would give LAMPF II a major advantage, but how do we do this? Our polarized proton beam is of no help unless the spin-transfer coefficients are large. We will measure them, but they are expected to be small. LEAR plans to make antiprotons and then polarize them by scattering, but this destroys the cooled-beam quality. Perhaps we can scatter first and then cool.

If all grand schemes prove impractical, however, there is still valuable work that we could do with a conventional (uncooled) \bar{p} secondary beam line.

Colliding beams of polarized protons would be exciting ($16 + 16 \text{ GeV} \approx 500 \text{ GeV}$ on a fixed target). The technology is well understood, but if we are to have variable energy without interfering with kaon production we would need two additional rings.

Finally, we could turn our existing polarization expertise to πp , $K p$, and Λp experiments. These fields have been discussed by Kelley and by Cutkosky (the Isgur-Karl catastrophe) in the proceedings of the 31-GeV workshop [Los Alamos National Laboratory report LA-8775-C (March 1981), pp. 166 and 185].

X. RARE KAON DECAY WORKING GROUP SUMMARY

Some of the rare decays considered are discussed below.

The standard Weinberg-Salam-Glashow model offers no explanation or differentiation among the three generations of leptons and quarks observed in nature. New interactions, not contained in this model, have been proposed to link the generations. Rare lepton-flavor-violating decays, such as $K^+ \rightarrow \pi^+ \mu e$ and $K^0 \rightarrow \mu e$, test for the presence of these interactions even if the mass of the particle responsible for these interactions is many tera-electron volts. There is no direct way to search for such massive objects.

The decay $K^+ \rightarrow \pi^+ \bar{v}v$ is suppressed by the Glashow-Iliopoulos-Maiani (GIM) mechanism. The particular interest in this process is that the decay rate is proportional to the number of (light) neutrino types, thus making it possible to determine the number of lepton generations. The high flux available at LAMPF II is necessary to improve the experimental sensitivity.

The decay $K_L^0 \rightarrow K^{\pm} e^{\mp} \bar{\nu}$ is expected to have a branching ratio of 3×10^{-9} . This decay provides a sensitive test of the conserved vector current (CVC) hypothesis in the presence of a strange quark. CVC, which is a cornerstone of

all present theories, has never been tested in this domain. Other rare decays, which have been discussed and are of great interest, are $K_L^0 \rightarrow \mu^+ \mu^-$, $K^+ \rightarrow \pi^+ e^+ e^-$, $K^0 \rightarrow \ell^+ \ell^- \gamma$, and $K^+ \rightarrow e^+ \nu$.

XI. CP VIOLATION WORKING GROUP SUMMARY

Since the discovery of CP violation in 1964, many difficult and elegant experiments have been performed to study this phenomenon. However, the most important issue, namely what interaction is the underlying cause of the observed effects, remains a mystery. CP violation has been observed only in the neutral kaon system, but more precise experiments are required to distinguish between various theoretical models of the effect.

The uncertainty in the experimental determination of every CP-violation parameter is dominated by statistical errors. The large increase in kaon intensity available at LAMPF II will be clearly invaluable in improving these measurements. We could also design K^0 beams with improved properties (smaller beam spot and lower neutron contamination, for example) that would permit experiments with smaller systematic errors.

The experiments that have been identified as being particularly important are

- a precise measurement of the rate for $K_L^0 \rightarrow \pi^0 \pi^0$,
- a measurement of the μ^+ polarization transverse to the decay plane in $K_L^0 \rightarrow \pi^- \mu^+ \nu_\mu$ and $K^+ \rightarrow \pi^0 \mu^+ \nu_\mu$,
- a search for CP-violating effects in the rare-decay modes $K^0 \rightarrow \mu^+ \mu^-$ and $K^0 \rightarrow \gamma \gamma$,
- a search for the CP-violating decay $K_S^0 \rightarrow \pi^+ \pi^- \pi^0$, and
- a comparison of the decays $K^+ \rightarrow \pi^+ \pi^+ \pi^-$ and $K^- \rightarrow \pi^- \pi^+ \pi^-$.

XII. NEUTRINO WORKING GROUP SUMMARY

It is fairly straightforward to identify the physics one wants to address with a neutrino facility at LAMPF II. These are

- neutrino oscillation searches with improved sensitivities. In particular, we could search for ν_μ disappearance and ν_μ oscillation into ν_e with sensitivities not matched by any other facility.
- neutrino-electron elastic scattering. The $\nu_\mu - e$ and $\bar{\nu}_\mu - e$ elastic scattering are the simplest weak neutral-current interactions that can be studied. Data of unprecedented accuracy could be obtained at LAMPF II that would

provide the most stringent test of the electroweak theories and the most accurate measurement possible of the Weinberg angle. The Neutrino Working Group is investigating how to minimize backgrounds and systematic errors.

We could also measure $\nu_e - e$ and $\bar{\nu}_e - e$ elastic scattering in a neutrino beam derived from K_L^0 decay, and we could separate ν_e and $\bar{\nu}_e$ contributions from their different distributions in $y = E_e/E_\nu$.

- neutrino-proton elastic scattering. This is the simplest weak neutral-current reaction involving hadrons. As such, the interpretation of any results is subject to some quantum chromodynamics (QCD) model dependence. Coupled with precise neutrino-electron data, this reaction can be used to test QCD assumptions. Low- q^2 neutrino-proton data can also be used to test the space-time properties of the weak neutral current.
- other physics. Enormous numbers of other types of events can be collected. Here we include quasi-elastic scattering and single-pion production. It is not clear that the vastly improved statistical sample will be useful, as systematic errors and model dependences in the interpretation are serious problems. On the other hand, the large fluxes imply that one can contemplate a relatively small, heavily instrumented detector to study these processes and minimize these problems.

The group is also addressing the questions of what a detector and a neutrino facility might look like, what proton energy is optimum, and what fluxes are to be expected.

XIII. MUON WORKING GROUP

A workshop entitled "Muon Science and Facilities at Los Alamos" was held at LAMPF March 15-18, 1982 as part of the continuing series of LAMPF II planning sessions. About 45 workshop participants from the United States and abroad discussed current worldwide facilities and scientific activities using stopping muon beams with a view toward delineating future directions for such work. The workshop was organized around five working groups:

1. Particle Physics - V. W. Hughes (Yale University), Chairman;
2. Nuclear Physics - J. D. Walecka (Stanford University), Chairman;
3. Solid-State Physics - A. Schenck (ETH Zürich), Chairman;

4. Chemistry - D. Fleming (University of British Columbia, TRIUMF), Chairman; and

5. Facilities and Costs - J. Bradbury (Los Alamos), Chairman.

In addition, there were the following keynote speakers:

K. Nagamini (Tokyo) - "Muon Physics Program and Facilities at the KEK-BOOM,"

V. W. Hughes (Yale University) - "Weak and Electromagnetic Interaction Studies with Muons," and

J. H. Brewer (University of British Columbia, TRIUMF) - "Use of Muons in the Study of Solid-State Physics and Chemistry."

The working group discussions focused on future facilities that would be possible at Los Alamos with the planned Proton Storage Ring (PSR) and LAMPF II. The PSR will produce protons with a planned maximum intensity on target of about 100 μ A in two pulse modes, a 1-ns-wide burst at 720 Hz or a 270-ns-wide burst at 12 Hz. Pulsed beams are important for the following types of experiments:

- studies of delayed processes where beam-associated background can be reduced by a prompt timing cut;
- experiments using intrinsically pulsed environments, such as laser pumping of muonic atoms or applications of rf fields;
- muonium hyperfine structure studies with line-narrowing techniques; and
- muon-spin-rotation experiments wherein all muons in a given pulse arrive at the target within a few nanoseconds, thus circumventing pile-up timing problems.

The working model for LAMPF II was about 25 μ A in a dc mode at about 16-GeV energy. The principal advantage LAMPF II would afford is a dc beam, providing an order of magnitude lower instantaneous rates than are currently achieved at LAMPF with its 6% duty factor and a much higher muon intensity than is now available at the Stopped Muon Channel (SMC).

Before the workshop a preliminary design for a superconducting-solenoidal decay beam channel was worked out by G. H. Sanders, R. Werbeck, and R. H. Heffner of Los Alamos. Table I shows the calculated values for the SMC and for the new channel at both the PSR and LAMPF II. The LAMPF II production target will be thinner than at the PSR to accommodate the increased heat production at 16 GeV.

The most striking feature of the new solenoidal decay channel is that it produces 5-10 times the muon/proton ratio compared to the SMC. Furthermore, at the PSR a relatively thick production target can be used because the proton

energy is modest and the proton beam need not be passed to a downstream target. Thus, although the PSR has a substantially lower primary proton current, the final muon stopping rates would be about the same. The PSR, of course, provides a unique duty factor. At LAMPF II the muon yield per proton is about two orders of magnitude greater than at LAMPF/SMC because of the solenoid decay channel and the increased production cross section at the higher energy. Therefore, a LAMPF II channel could produce a dc μ^- stopping beam nearly 30 times more intense than at the SMC; the μ^+ stopping rate would be about 6 times larger. In addition to these advantages this new solenoidal channel would provide significantly better rate and phase-space control.

A possible location for a muon channel coupled to the PSR was found south of the Weapons Neutron Research (WNR) facility. Layouts and designs of the proton transport line, target cell, and shielding were studied and a rough cost estimate was obtained.

One of the most important conclusions reached at the workshop regarding beam characteristics was the need for high-intensity, low-momentum (<50 -MeV/c) μ^+ and μ^- beams. Consequently, future beam-line designs for either the PSR or LAMPF II will focus on these low-momentum beams in addition to the decay beam discussed above.

In general there was great enthusiasm generated at the workshop for the future of physics with stopping muon beams. No other particle probe at any of the world's meson facilities provides the extraordinary range of scientific study, from chemistry to particle physics, as does the muon. Equally great was the enthusiasm shown for possible new muon facilities at Los Alamos.

TABLE I
COMPARISON OF CHANNEL FLUXES

		Columns					
(1)	(2)	(3) ^a	(4)	(5) ^b	(6) ^c	(7)	(8) ^d
Location	Target	I_p (μ A)	$\Delta p_{\pi}/p_{\pi}$ (FWHM)	$\Delta\Omega_{\pi}$ (str)	μ Spot Size (cm^2)	$\mu^-/\mu\text{A}\cdot\text{s}$	$\mu^-/\mu\text{A}\cdot\text{s}\cdot\text{q}$
LAMPF II	8-cm Be	25	0.158	0.196	100	1.6×10^7	5.3×10^4
PSR	30-cm Be	15-100	0.158	0.196	100	8.8×10^5	3.0×10^3
SMC	5-cm C	500	0.134	0.456	50	1.4×10^4	0.93×10^2

^aColumn 3 contains estimates of proton intensities for the PSR and LAMPF II. The SMC value is a nominal one as of February 1982.

^bColumn 5 is the geometrical solid angle subtended by the front aperture of the channel.

^cColumn 6 is for the FWHM and thus contains about 50% of the beam.

^dColumn 8 is for the CH_2 stopping target (range spread $\approx 3 \text{ g/cm}^2$) and uses the nominal spot size in Column 6. As an example, consider a 1-cm-thick, 50- cm^2 CH_2 target:

		$\mu^- \text{ stops/s}$	$\mu^+ \text{ stops/s}$
LAMPF II	(25 μ A)	6.6×10^7	6.6×10^7
PSR	(15 μ A) ^{aa}	2.3×10^6	1.1×10^7
SMC	(500 μ A)	2.4×10^6	1.2×10^7

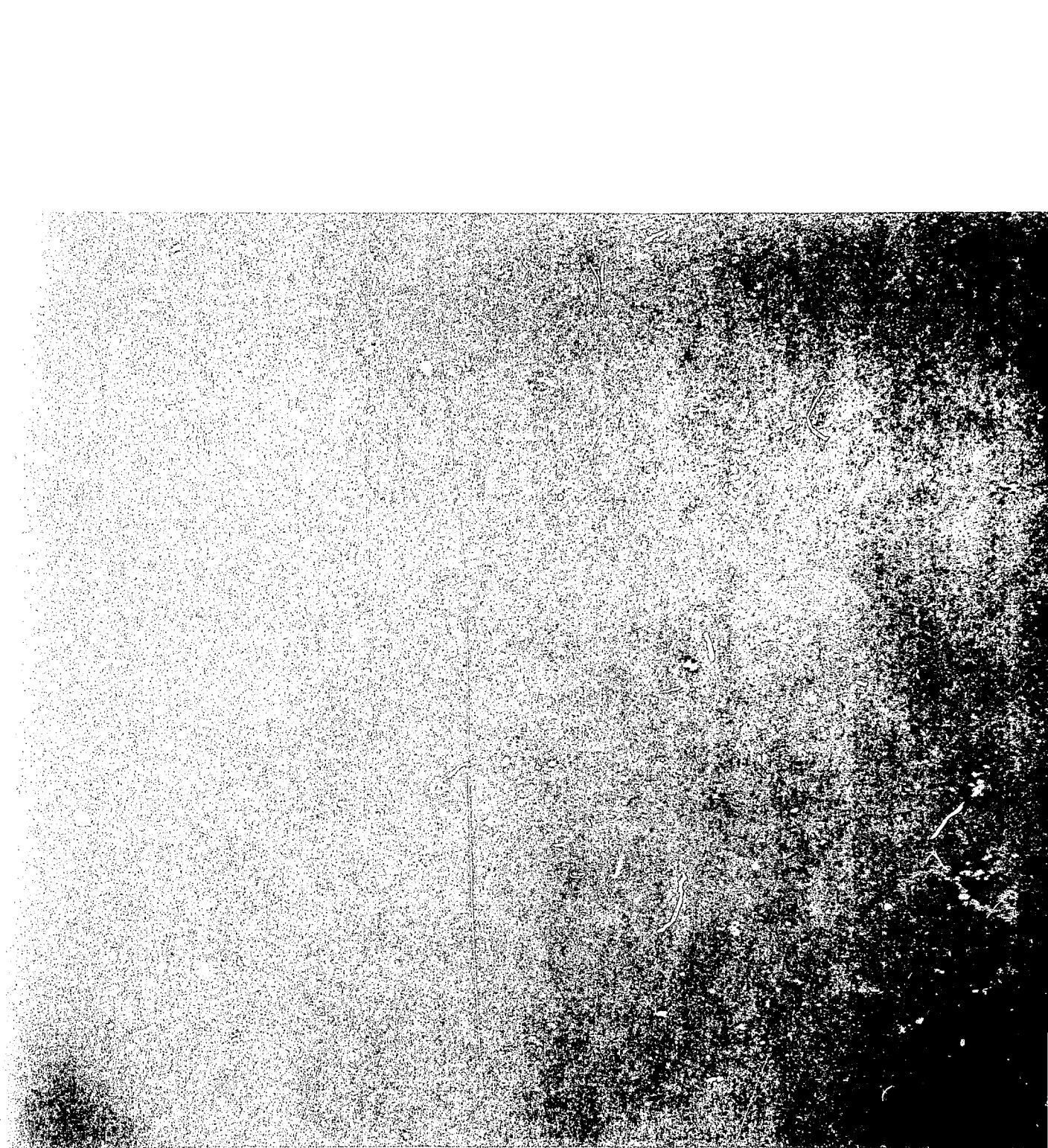
^{aa}PSR rates corrected for proton absorption in the thick target.

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