

DISCLAIMER

This document was prepared pursuant to a contract with the United States Government. Neither the United States Government nor any of its employees make any representations about the accuracy of the information contained in this document. Any opinions, findings, and conclusions or recommendations expressed herein are those of the author(s) and do not necessarily reflect the views of the United States Government. The United States Government does not make any express or implied representations about the accuracy of the information contained in this document. The document contains neither recommendations nor conclusions of the United States Government. The United States Government is not responsible for typographical or editorial errors that may appear in this document.

Los Alamos National Laboratory is operated by the University of California for the United States Department of Energy under contract W-7405-ENG-36

LA-UR-82-3397

DE83 003559

TITLE: Special Features of the Alpha Induced d and \bar{d} Breakup

AUTHOR(S): J. M. Lambert and P. A. Treado, Georgetown U
 Y. Koike, Kyoto University, Japan
 I. Slavus, Rudjer Boskovic Institute, Yugoslavia
 F. D. Correll, U.S. Naval Academy
 Ronald E. Brown, R. A. Hardkopf, Nelson Jarmie of the
 Los Alamos National Laboratory

SUBMITTED TO: Contributed paper 7th Conf. on Applications of Accelerators
 in Research and Industry, November 8-10, Denton, TX.

NOTICE

**PORTIONS OF THIS REPORT ARE ILLEGIBLE. It
 has been reproduced from the best available
 copy to permit the broadest possible avail-
 ability.**

DISTRIBUTION BY THE U. S. GOVERNMENT

u.s. g. o.

By acceptance of this article for publication [yellow box] recognizes that the U.S. Government retains a non-exclusive, royalty-free license to publish or reproduce the published form of this contribution, or to allow others to do so, for U.S. Government purposes.

The Los Alamos National Laboratory requests that the publisher identify this article as work performed under the auspices of the U.S. Department of Energy.

MASTER

[Los Alamos] Los Alamos National Laboratory
 Los Alamos, New Mexico 87545

Paper to be submitted to the "Conference on the Application of Accelerators in Research and Industry" for publication in the IEEE Transactions in Nuclear Science

SPECIAL FEATURES OF THE ALPHA INDUCED d AND \bar{d} BREAKUP

J.M. Lambert and P.A. Treado
Department of Physics, Georgetown University.
Washington, DC 20057 USA

I. Slaus
Rudjer Boskovic Institute and Georgetown University
41001 Zagreb, Yugoslavia

Y. Koike
Department of Nuclear Engineering, Kyoto University
Kyoto, Japan

F.D. Correll
Department of Physics, U.S. Naval Academy
Annapolis, Maryland 21402 USA

Ronald E. Brown, R.A. Hardekopf and Nelson Jarmie
Los Alamos National Laboratory
Los Alamos, New Mexico 87545 USA

Summary

The alpha induced deuteron breakup reactions have been of interest since the three-nucleon, nucleon induced deuteron breakup, problem was addressed with tractable and predictive codes based on the Faddeev formalism and with appropriate nucleon-nucleon forces. In this paper we discuss a few special features of the alpha induced deuteron breakup reactions. Specifically, we point out the importance of the n-p tensor force in the predictions of the three-body model that fit the tensor analyzing powers better and the concomitant deterioration of the fit to the vector analyzing power caused by the inclusion of the tensor force. We suggest that there exists evidence for the 1S_0 n-p interaction, which is isospin forbidden, in both the cross section and spin observables data. Finally, we discuss certain characteristics of the p- α quasi-free-scattering and n- α final-state-interaction processes.

Introduction

Alpha induced deuteron breakup has been studied over wide range of energies, 4 - 55 MeV, for p- α quasifree scattering (QFS) and p- α , n- α and n-p final-state interaction (FSI) kinematic conditions¹⁻³. Data recently have been compared with the predictions of the three-body model⁴ which assumes that the alpha particle is a structureless boson and which uses separable potentials for the input N- α and n-p forces. Specifically, the model employs N- α $s_{1/2}$, $p_{3/2}$ and $p_{1/2}$ forces with Yamaguchi form factors. The Coulomb force is not explicitly included in the three-body equations, but only in the two-body p- α t-matrix so as to reproduce experimental p- α phase shifts. For the n-p system only a 3S_1 rank-one Yamaguchi potential is used. While the three-body model provided satisfactory fits to most of the data, some significant discrepancies exist. Therefore, it is necessary to investigate whether, and in which way, more realistic forces affect the predictions from the model. This paper discusses some special features of the alpha induced deuteron breakup and information derived from these features.

Quasifree Scattering

The importance of the two-particle input force for the QFS is well established from the $D(p,pp)n$ studies⁵. Similarly, in order to fit the p- α QFS data it is necessary to employ a good p- α input force¹. However, even at the c.m. energy of 47 MeV, though the impulse approximation term, $I_{\alpha p}$, is the largest one, other terms in the multiple scattering series, particularly those which terminate in the n-p FSI, are very important.

The inadequacy of the plane-wave-impulse-approximation (PWIA) approach is dramatically demonstrated through the study of polarization observables. For the incident deuteron energy of 17 MeV, the measured Λ_y of the p- α QFS is 0.50 ± 0.03 , while at the corresponding conditions, i.e., $T_p^{\text{lab}} = 5.7$ MeV and $\theta_{\text{c.m.}} = 45^\circ$, the Λ_y measured for free p- α scattering is -0.6 . However, in the QFS kinematic region the measured tensor-analyzing powers are consistent with zero and intuitively

expected from the QFS picture.

This information is useful when one plans to use the $^4\text{He}(\vec{d},\text{ab})\text{cd}$ reactions to study two-spectator QFS processes.

Tensor n-p Interaction

The tensor n-p force recently has been included in the three-body model⁶⁾. The comparison of the three-body model predictions, when the tensor force is included, with our 17 MeV $^4\text{He}(\vec{d},\text{p}\alpha)\text{n}$ data²⁾ shows that the inclusion of the tensor force which predicts a deuteron D-state probability of 7% significantly improves the fit to the tensor-analyzing power data, as shown in Figure 1. On the contrary, the fit to the vector-analyzing power, A_y , becomes somewhat worse. We also have done the calculation with the tensor force which results in D-state probability of 4% and, as expected, the results are in between those for 7% and 0% D-state probabilities. Similar indications have been obtained from the kinematically incomplete measurements of the $^4\text{He}(\vec{d},\text{p})^5\text{He}$ reaction at 21 MeV⁷⁾. Thus, we conclude that the tensor-analyzing powers are sensitive to the n-p tensor force. The vector-analyzing power is influenced by the hitherto unincluded forces. Thus, it is worthwhile to investigate whether modifications of the N- α force, notably the $P_{1/2}$ force, would restore agreement between the predictions of the three-body model and the data. Likewise, it is important to investigate the effect of the tensor force at higher energies.

The $^5\text{He}_{\text{gs}}$ FSI Structure

The cross section enhancement due to the n- α FSI, $^5\text{He}_{\text{gs}}$, seems to vary rapidly as a function of the relative momentum angle. Figure 2 shows that the inclusion of the tensor force modifies the predicted $^5\text{He}_{\text{gs}}$ spectral structures rather significantly and in different ways. I.e., one peak increases in magnitude while the other decreases. With the limited data available, we suggest that alpha-induced deuteron breakup spectra which include the $^5\text{He}_{\text{gs}}$ structure should be acquired.

This should be done for both low and high energies.

The 1S_0 n-p Force

The 1S_0 n-p FSI is isospin forbidden in the d- α breakup. The present three-body model ignores the 1S_0 n-p force. The possible contribution of the 1S_0 n-p interaction has been studied experimentally⁸⁾ and theoretically in the framework of the impulse approximation⁹⁾.

Our 17 MeV data²⁾, analyzed in a higher resolution mode than previously considered, are shown in Figure 3 (3a: $\theta_\alpha = 20^\circ$, $\theta_p = 120^\circ$ and 3b: $\theta_\alpha = 30^\circ$, $\theta_p = 82.8^\circ$) and are compared with the three-body model predictions. First, we note that there was an error in our code related to the n-p FSI enhancement. The error has been corrected⁶⁾ and, in order to illustrate its effect we show the old incorrect calculation, the new correct, but not including the tensor force, calculation, and the new calculation which includes the tensor force. The error appreciably influenced only the region where the n-p relative energy is close to zero. In the region far from the n-p FSI the inclusion of the tensor force improves the fit to the cross section as shown in Figure 3a. However, in the n-p FSI region, as shown in Figure 3b, the inclusion of the tensor force makes the fit worse, underlining the need to incorporate the 1S_0 n-p force also. The effect is not large, but it is apparent. The discrepancy at E_p approximately 4 MeV could be a reflection of the same 1S_0 n-p importance of an inadequate $P_{1/2}$ n- α force.

More insight into the role of the 1S_0 n-p force can be obtained from the spin observables. In Table 1 we compare our measurements at 17 MeV for the reaction $^4\text{He}(\text{d},\text{d}^*)\alpha$ at $\theta_{\text{c.m.}} = 116.5^\circ$ with the previously measured elastic scattering data¹⁰⁾.

Table I

Observable	Comparison between ${}^4\text{He}(\vec{d},d)\alpha$ and ${}^4\text{He}(\vec{d},d^*)\alpha$ data	
	${}^4\text{He}(\vec{d},d)\alpha$, ref. 10	${}^4\text{He}(\vec{d},d^*)\alpha$, ours
A_y	-0.40	-0.3 ± 0.1
A_{xx}	+0.25	$+0.6 \pm 0.1$
A_{yy}	-0.65	-0.8 ± 0.1

The good agreement for the A_y and A_{yy} suggests that the 3S_1 n-p force dominates the n-p FSI interaction. However, the fact that the A_{xx} results differ markedly suggest that this observable may be sensitive to the details of the n-p force, particularly to the 1S_0 component.

References.

1. J.M. Lambert, P.A. Treado, P.G. Roos, N.S. Chant, A. Nadasen, I. Slaus and Y. Koike, Phys. Rev. C26, 357 (1982).
2. I. Slaus, J.M. Lambert, P.A. Treado, F.D. Correll, R.E. Brown, R.A. Hardekopf, N. Jarmie, Y. Koike and W. Gruebler, submitted to Nucl. Phys., and references therein.
3. H. Paetz gen. Schieck, private communication to Y. Koike.
4. Y. Koike, Nucl. Phys. A301, 411 (1978).
5. I. Slaus, J. Phys. Soc. Japan 44, 378 (1978).
6. Y. Koike, to be published.
7. M. Ishikawa, S. Seki, K. Furuno, Y. Tagishi, M. Sawada, T. Sugiyama, K. Matsuda, J. Sanada, and Y. Koike, J. Phys. Soc. Japan 51, 1327 (1982).
8. M. Bruno, F. Cannata, M. D'Agostino, G. Vannini, M. Lombardi and Y. Koike, Lett. Nuovo Cim. 29 385 (1980).
9. C. Werntz and F. Cannata, Phy. Rev. C24, 349 (1981).
10. W. Gruebler, V. Knig, P. Schmelzbach, B. Jenny, H.R. Bürge, R.A. Hardekopf, J. Nurzynski and R. Risler, Nucl. Phys. A334, 365 (1980).

Figure Captions.

1. The tensor analyzing power, A_{yy} , for the reaction $^4\text{He}(\vec{d},p\alpha)n$ at 17 MeV and $\theta_\alpha = 30^\circ$, $\theta_p = 82.8^\circ$. The solid curve is for the three-body model prediction with no n-p tensor force and the dashed curve is for the three-body model prediction including the n-p tensor force with the deuteron D-state probability of 7%. Kinematic conditions for the n-p and $^5\text{He}_{\text{gs}}$ FSI, as well as for collinearity, are indicated.
2. The cross section for the reaction $^4\text{He}(\vec{d},p\alpha)n$ with $\theta_\alpha = 20^\circ$, $\theta_p = 120^\circ$, and $E_d = 17$ MeV. The curves are: W = a previous calculation²⁾ without the n-p tensor force and with an error in the n-p FSI enhancement region; C = no tensor force and no error; and C-T = with the tensor force and $P_d = 7\%$, and no error.
3. The cross section, in the higher resolution mode, for the $^4\text{He}(\vec{d},p\alpha)n$; (a) $\theta_\alpha = 20^\circ$, $\theta_p = 120^\circ$, and (b) $\theta_\alpha = 30^\circ$, $\theta_p = 82.8^\circ$. The curves are labeled as in Figure 2.

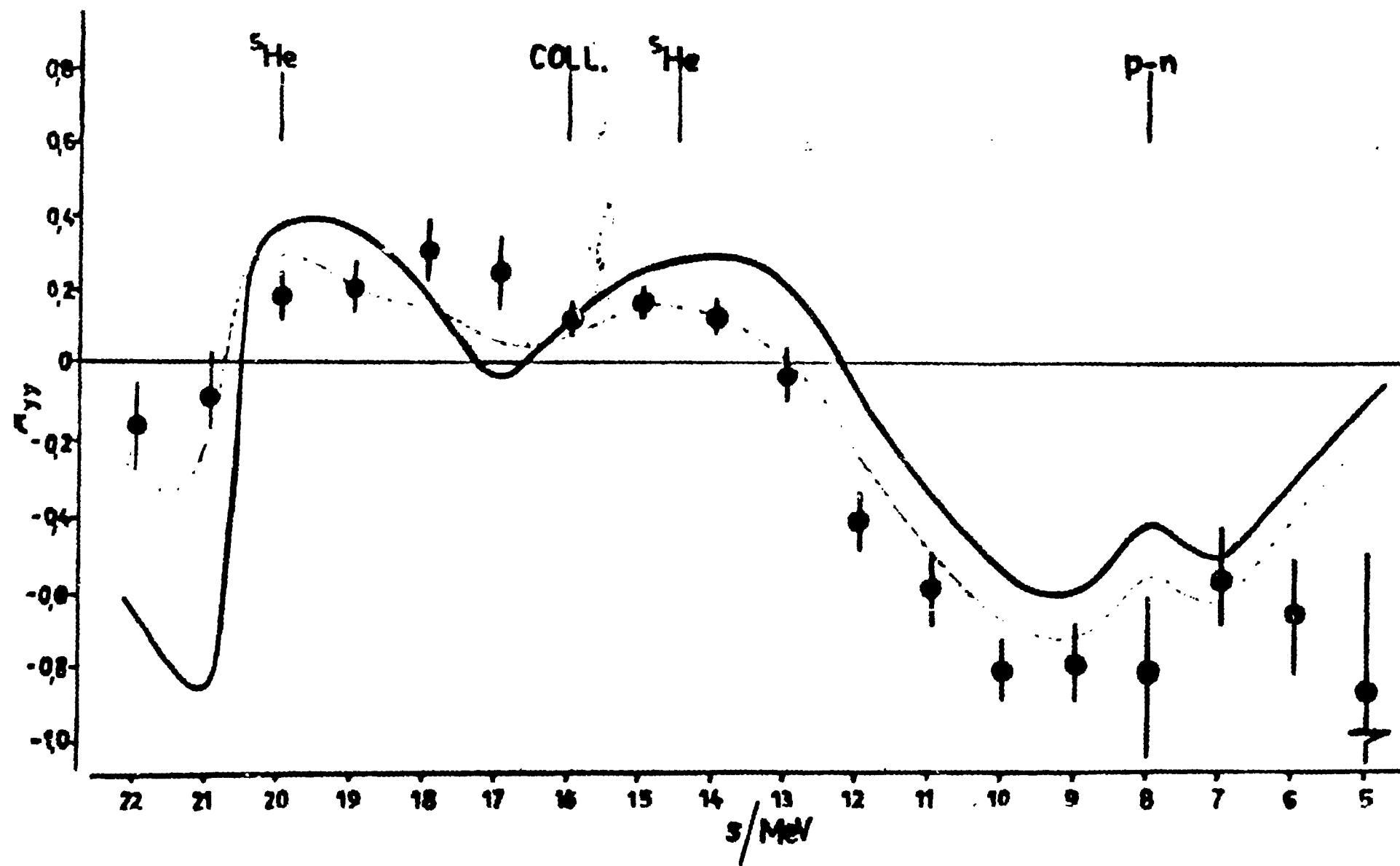


Figure 1.

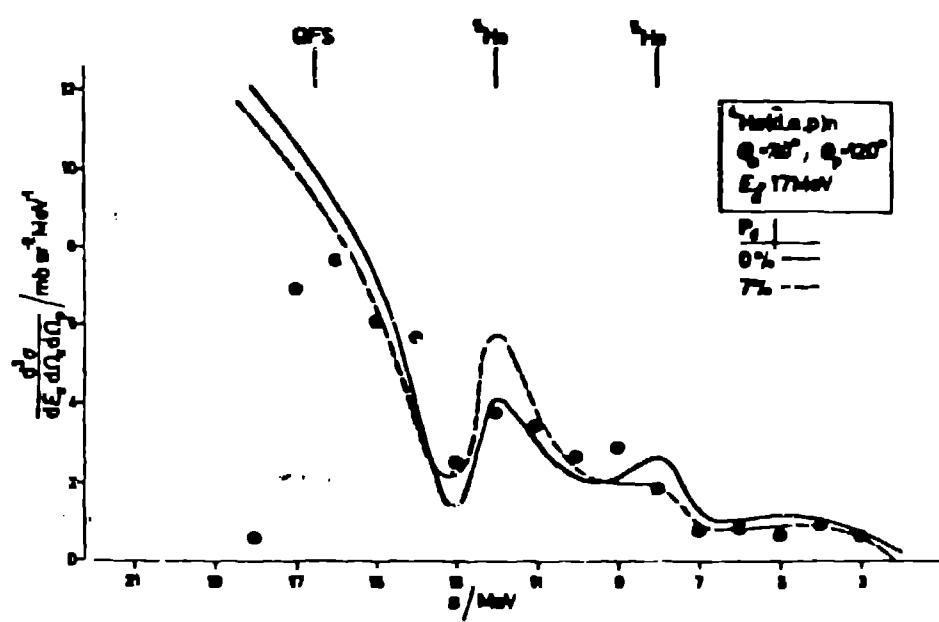


Figure 2.

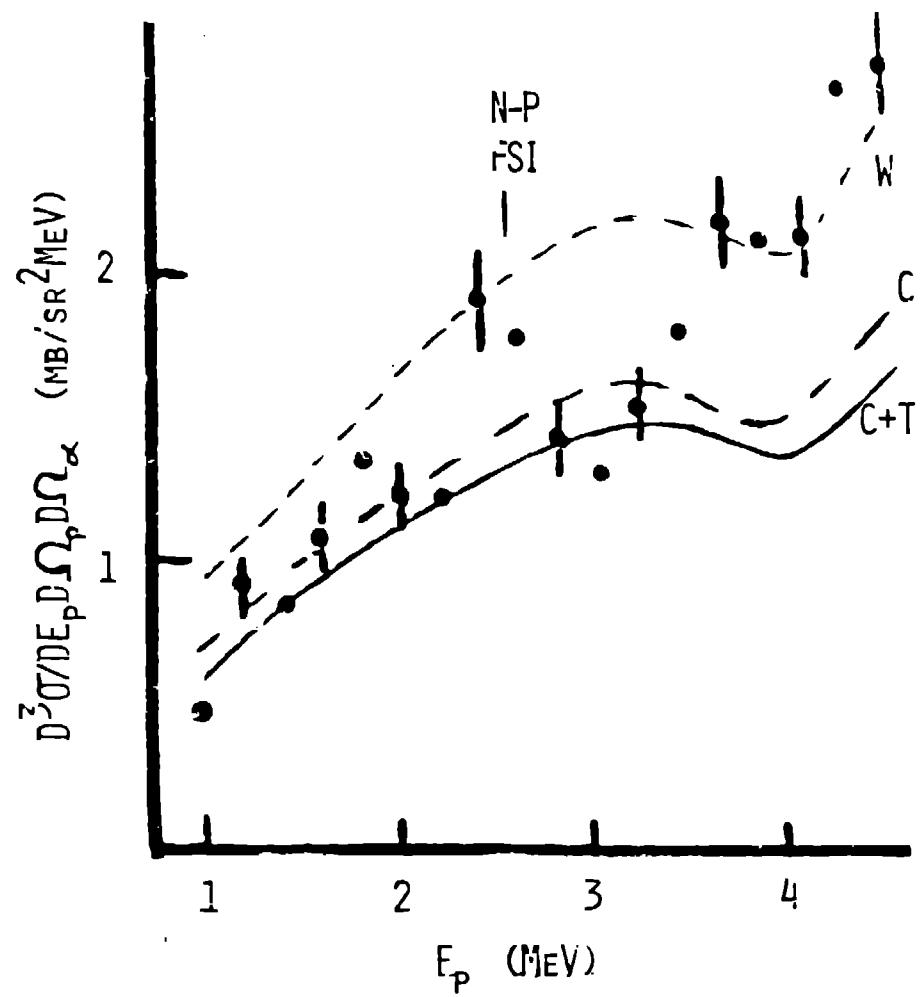


Figure 3b.

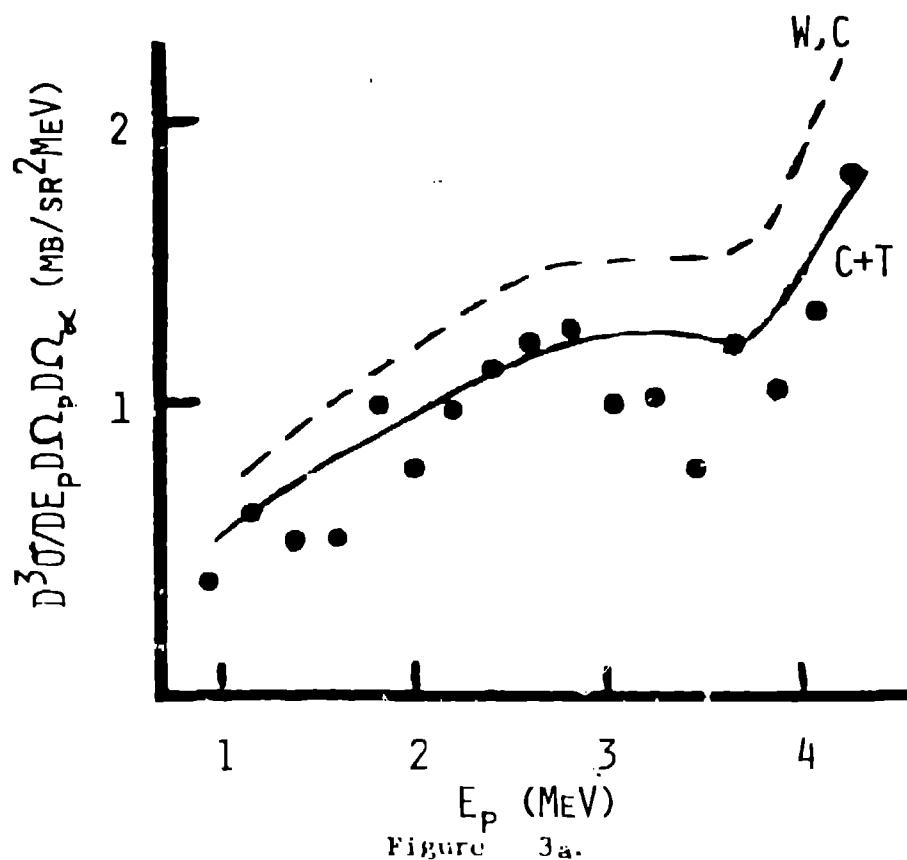


Figure 3a.