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SUMMARY OF PHYSICS FROM MEASUREMENTS WITH LONGITUDINALLY  
POLARIZED BEAMS AND TARGETS AT ZGS ENERGIES

by

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## SUMMARY OF PHYSICS FROM MEASUREMENTS WITH LONGITUDINALLY POLARIZED BEAMS AND TARGETS AT ZGS ENERGIES\*

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### ABSTRACT

An extensive amount of data were obtained from measurements of proton-proton elastic scattering from 1 to 12 GeV/c using longitudinally polarized beams and targets. We summarize physics learned from these data as well as other related experimental results. The topics will include structures observed in nucleon-nucleon scattering at lower energies and dinucleon resonances, pp scattering-amplitude measurements at 6 GeV/c, and large  $p_{\perp}$  results in pp elastic scattering.

### I. STRUCTURE IN NUCLEON-NUCLEON SYSTEM AND DINUCLEON RESONANCES

#### i) $I = 1$ System

A striking energy dependence has been observed in the difference between the nucleon-nucleon total cross-sections for pure spin states:

$$\Delta\sigma_L = (4\pi/k) \operatorname{Im}\{\phi_1(0) - \phi_3(0)\} = \sigma_{\text{Tot}}(\pm) - \sigma_{\text{Tot}}(\mp) ,$$

and

$$\Delta\sigma_T = -(4\pi/k) \operatorname{Im}\{\phi_2(0)\} = \sigma_{\text{Tot}}(\mp\mp) - \sigma_{\text{Tot}}(\mp\mp) .$$

Since then many authors have attempted to interpret the results using not only these results but other channels such as pp elastic scattering,  $\pi D$  elastic scattering,  $pp \rightarrow \pi D$ ,  $\gamma D$  reactions etc. Summary of the investigation is for instance given in Ref. 1. Figure 1 shows  $\Delta\sigma_L$  data. At present, more energy points in  $\Delta\sigma_L$  including the existing points have been measured at LAMPF in Los Alamos and simultaneously  $\Delta\sigma_T$  by a Rice group.

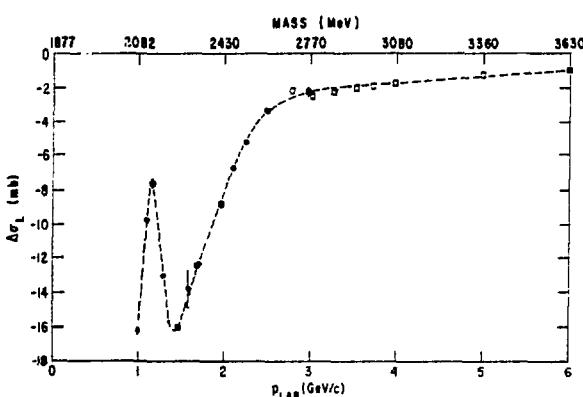
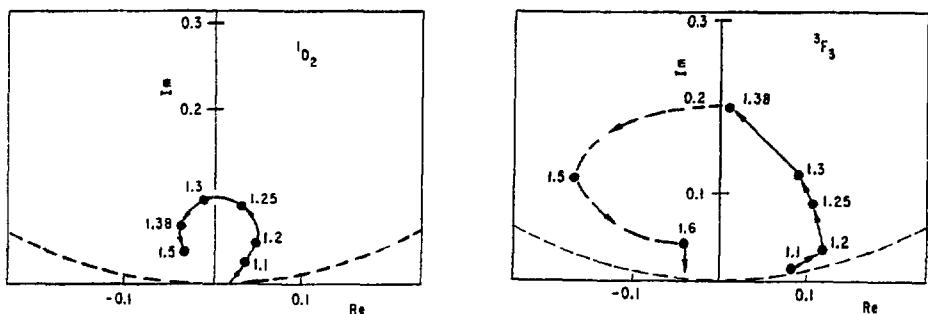


Fig. 1 Total cross-section difference  
 $\Delta\sigma_L = \sigma_{\text{Tot}}(\pm) - \sigma_{\text{Tot}}(\mp)$ .  
 The white squares are preliminary data.

The large energy dependence are seen not only in  $\Delta\sigma_L$  and  $\Delta\sigma_T$ , but in polarization,  $C_{NN} = (N,N;0,0)$ ,  $C_{LL} = (L,L;0,0)$ . The typical analysis carried out so far are the Argand plot by dispersion and relation,<sup>2</sup> elastic total cross section,<sup>3</sup> Legendre coefficient analysis,<sup>4</sup> and  $C_{NN}$  data analysis,<sup>1</sup> and phase-shift analyses. All of these analyses are consistent with the existence of diproton resonances. Particularly, strong indication of resonance in the  $1^3D_2$  and  $3^3F_3$  states are established by phase shift analyses<sup>5,6</sup> as shown in Figs. 2 and 3.



Figs. 2 Argand diagrams of the  $^1D_2$  (a) and  $^3F_3$  (b) partial waves (points are in  $\text{GeV}/c$ ); the background contributions have been subtracted.

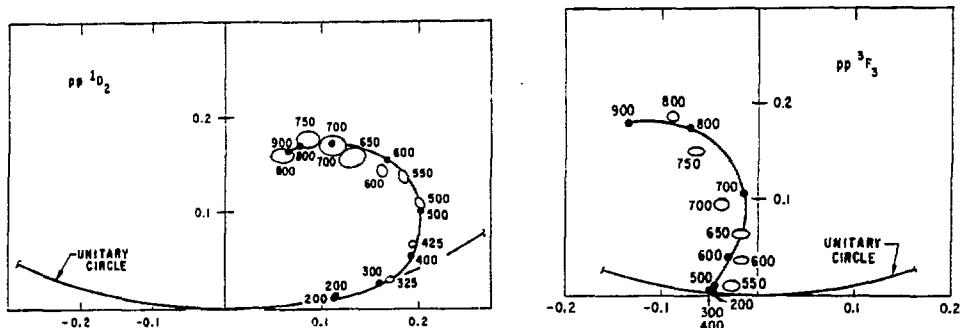
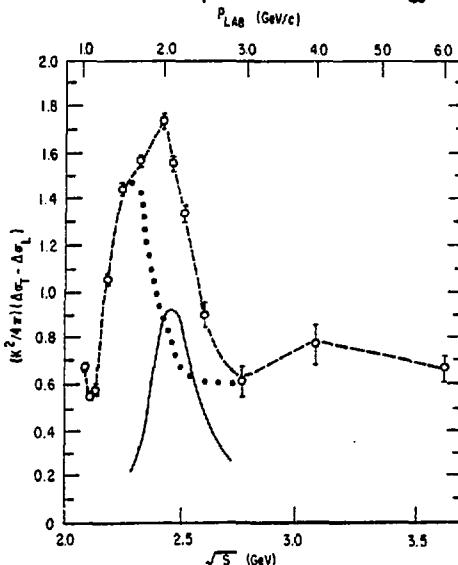


Fig. 3 Argand diagram of the  $^1D_2$  (a) and  $^3F_3$  (b) partial waves based on Arndt's phase shifts (points are in kinetic energy, MeV). The ellipses represent the errors in the real and imaginary parts of the amplitudes for energy-independent solutions. The continuous curves represent the energy-dependent solutions.



Other possible resonances in  $pp$  system include a singlet resonance in  $\Delta\sigma_T$  at 2  $\text{GeV}/c$ , a triplet resonance appearing in  $(k^2/4\pi)(\Delta\sigma_T - \Delta\sigma_1)$  as shown in Fig. 4. We also note that there is no  $^1D_2$  and  $^3F_3$  contribution to the polarization data at

Fig. 4 New triplet structure at 2.0  $\text{GeV}/c$ ; the dotted curve is deduced from  $\Delta\sigma_1$  data.

$\theta_{c.m.} = 63^\circ$ . We see an interesting structure in a plot of  $k^2 P(d\sigma/d\Omega)/\sin 2\theta_{c.m.}$  vs.  $p_{lab}$  as shown in Fig. 5. The quantity is proportional to

$$(2 \text{ Im}^3 P_0 + 3 \text{ Im}^3 P_1)(\text{Re}^3 P_2) - (2 \text{ Re}^3 P_0 + 3 \text{ Re}^3 P_1)(\text{Im}^3 P_2)$$

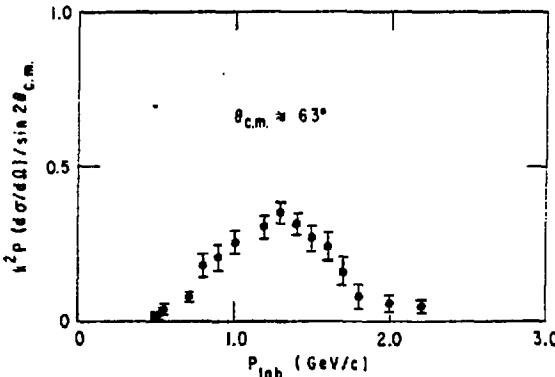


Fig. 5 Energy dependence of  $P(d\sigma/d\Omega)$  at  $\theta_{c.m.} \approx 63^\circ$ .

Recently Borisov et al.<sup>7</sup> investigated the energy dependence of  $(C_{NN} - C_{LL})$  at  $\theta_{c.m.} = 90^\circ$  in where neither  $^1D_2$  nor  $^3F_3$  contribution exist. As shown in Fig. 6, there is a sharp structure around 1.3 GeV/c. One may conclude  $^3P_0$  partial wave is responsible for the structure.

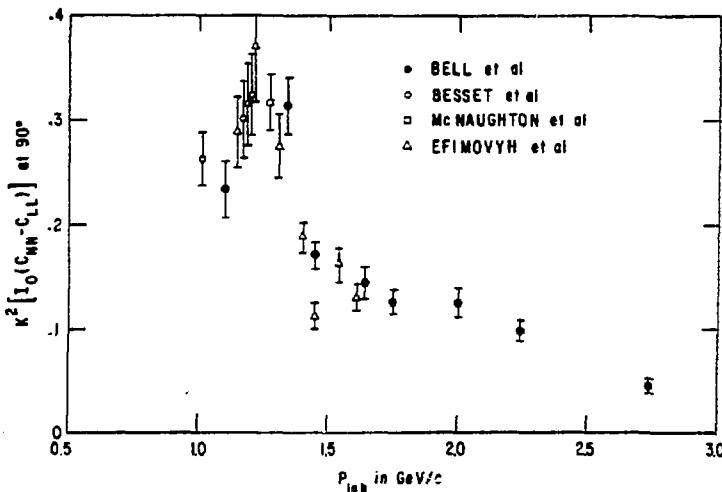


Fig. 6 A plot of  $k^2 [I_0 (C_{NN} - C_{LL})]$  at  $90^\circ$  versus  $p_{lab}$  in GeV/c.

### ii) $\Delta \sigma$ ( $I = 0$ ) Measurements

The Argonne PPT group has recently measured the difference between isoscalar nucleon-nucleon total cross sections for pure longitudinal

initial spin states,  $\Delta\sigma_L$ (pd), using a polarized proton beam and a polarized deuteron target. In the simplest approximation,  $\Delta\sigma_L$ (pd)  $\approx$   $\Delta\sigma_L$ (pp) +  $\Delta\sigma_L$ (pn). One can extract  $\Delta\sigma_L$ (I=0) data using both  $\Delta\sigma_L$ (pd) and  $\Delta\sigma_L$ (pp) as shown in Fig. 7; a significant structure is observed around 1.5 GeV/c. This seems to suggest the existence of a new isoscalar spin-singlet dinucleon resonance. We note here that there exists a clear shoulder in the np total cross-section data<sup>9</sup> in the vicinity of 1.5 GeV/c.

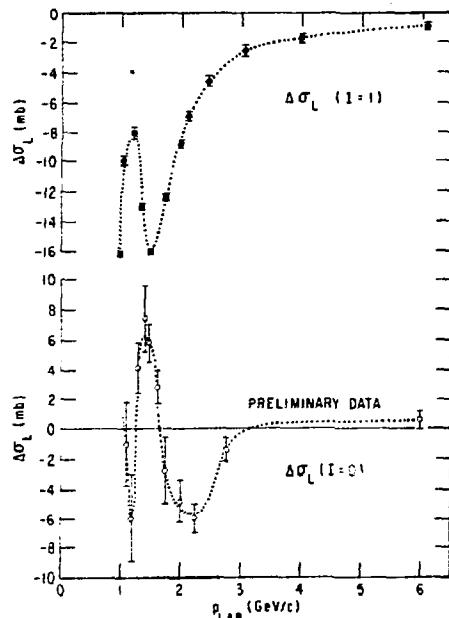


Fig. 7  $\Delta\sigma_L$ (I = 0) together with  $\Delta\sigma_L$ (I = 1).

A recent phase-shift analysis using these preliminary data by Hoshizaki et al.<sup>10</sup> suggests that there exists a partial wave whose behavior is consistent with the Breit-Wigner Resonance formula, namely, the spin singlet  $^1F_3$  wave. From the dispersion analysis of a forward I = 0 scattering amplitude using the data on  $\Delta\sigma_L$ (I = 0), Grein and Kroll<sup>11</sup> showed that the Argand plot of the amplitude has a resonancelike behavior around 1.5 GeV/c, and that suggests the existence of a spin-singlet dibaryon resonance.

### iii) Conclusion on I = 0 and I = 1 Resonances

Candidates for dibaryon resonances that can couple to nucleon-nucleon systems are summarized in the table below.

#### Candidates for Dinucleon Resonances

##### i) I = 1 Isospin State

	$B_1^2$ (2.14)	$B_1^2$ (2.18)	$B_1^2$ (2.22)	$B_1^2$ (2.43)	$B_1^2$ (2.43)
Mass, GeV	2.14 - 2.17	2.18 - 2.20	2.20 - 2.25	2.43 - 2.50	2.43 - 2.50
Width, MeV	50 - 100	100 - 200	100 - 200	$\sim$ 150	$\sim$ 150
Quantum State	$^1D_2$	Triplet P ?	$^3F_3$	probably $^1G_4$	Triplet $R_{JJ}$ ?

##### ii) I = 0 Isospin State

	$B_0^2$ (2.14)	$B_0^2$ (2.22)	$B_0^2$ (2.43)
Mass, GeV	2.14 - 2.17	2.20 - 2.26	2.40 - 2.50
Width, MeV	50 - 100	100 - 200	
Quantum State	Triplet ?	$^1F_3$	Triplet ?

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## II. SCATTERING-AMPLITUDE MEASUREMENTS

The present state of measured observables toward  $p\bar{p}$  scattering-amplitude was summarized in a report.<sup>12</sup> Here we will not document the details.

## III. $p\bar{p}$ SPIN CORRELATIONS AT HIGH $p_{\perp}$

Elastic scattering at high  $p_{\perp}$  has been theoretically studied by Chen<sup>13</sup>, Farrar et al.,<sup>14</sup> Brodsky et al.,<sup>15</sup> and Wolters.<sup>16</sup>

It is our major aim to observe if the quark would carry the spin of the proton. The prediction of QIM assumes helicity conservation among quarks, that is, quark mass is nearly zero. In order to test the prediction at 12 GeV/c (highest momenta at ZGS), we need to measure only two spin-spin correlation parameters so that we test if the s-channel helicity amplitude  $\phi_2 = \langle \dots | + + \rangle = 0$  or not; experimentally we need to check if  $C_{SS} = -C_{NN}$  where a value of  $C_{SS}$  can be obtained by measuring  $C_{LL}$ , because at  $\theta_{c.m.} = 90^\circ$  the following relationship holds:

$$C_{SS} = C_{NN} - C_{LL} - 1$$

The parameter  $C_{NN}$  is already measured by the Michigan Group.<sup>17</sup> We have measured  $C_{LL}$  at  $\theta_{c.m.} = 90^\circ$ , and preliminary data (about 1/3 of the existing data) are shown in Fig. 8. Together with these data the quark helicity conservation is verified within our statistical error.

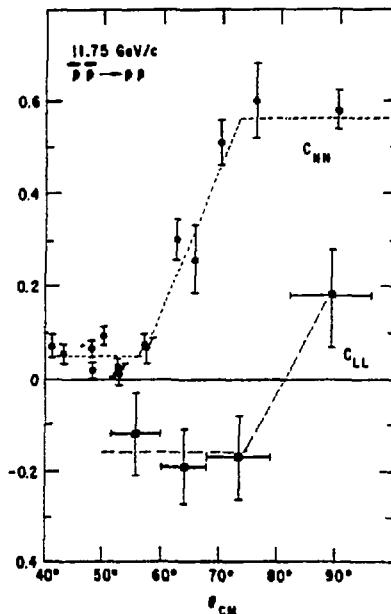


Fig. 8 Spin correlation parameters  $C_{LL}$  and  $C_{NN}$  for  $p\bar{p}$  elastic scattering at 11.75 GeV/c near  $\theta_{c.m.} = 90^\circ$ . The curves shown are to guide the eye.

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