

Muon transfer from hot muonic hydrogen atoms to neon

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A negative muon beam has been directed on adjacent solid layers of hydrogen and neon. Three targets differing by their deuterium concentration were investigated. Muonic hydrogen atoms can drift to the neon layer where the muon is immediately transferred. The time structure of the muonic neon X-rays follows the exponential law with a disappearance rate corresponding to the one of μ^-p atoms in each target. The rates $\lambda_{pp\mu}$ and λ_{pd} can be extracted.

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1. Introduction

This study has historically grown as an extension of a TRIUMF experiment which is already described in this volume [1]. Thus we won't repeat here most of the features which are described there, but concentrate on those details which are specific to this contribution.

By slowing down via Coulomb interaction with electrons, muons are stopped in a thin solid layer of hydrogen containing a small amount of deuterium. In this dense material, the μ^-p atoms are thermalized in a very short time ($\sim 10^{-12}$ s) compared to the μ^- lifetime. Different μ^-p disappearance channels exist : muon decay, $pp\mu$ molecular formation or muon transfer to deuterium and to impurities. After muon transfer from μ^-p , the muonic deuterium atom has an initial kinetic energy of about 45 eV. Most of them are slowed down by interacting on protons until they reach a few eV where the cross-section drops due to the Ramsauer-Townsend effect.

By using a decay e^- imaging system, the emission of these "hot" μ^-d in vacuum has been extensively studied in terms of deuterium concentration dependence and layer thickness [1]. Because of their high velocity (~ 1 cm/ μ s), these μ^-d have time to drift a few centimeters away from the hydrogen layer before muon decay occurs. Therefore the decay e^- imaging system was a very useful device for observing them. Although their velocity should be very slow compared to these μ^-d (~ 0.1 mm/ μ s), the emission of residual μ^-p atoms can not be completely excluded. The decay electron imaging system is, however, unable to extract their contribution from the target background.

The aim of this work is to study individually the μ^-p and μ^-d emission from the hydrogen layer. The basic idea is to freeze a thin neon layer on the surface of hydrogen. If muonic hydrogen atoms are emitted in this direction, these will instantaneously release their muon to form excited muonic neon atoms by transfer. Characteristic X-rays will be emitted during the cascade and observed with a Ge intrinsic semiconductor detector.

If one neglects the μ^-p and μ^-d drift time inside the hydrogen layer (this process remains to be studied in detail), the time distribution of the muonic neon X-rays emitted after muon transfer should have an exponential shape with an associated rate λ corresponding to the disappearance rate of μ^-p atoms inside the hydrogen layer:

$$\lambda = \lambda_0 + (1 - c_d)\lambda_{pp\mu} + c_d\lambda_{pd} \quad (1)$$

where the transfer to impurities has been neglected. The X-rays emitted after direct muon stop in the neon layer will be observed promptly and can be easily separated from the other events.

2. Measurement

The apparatus set-up is shown on Fig. 1. To be as careful as possible with gas mixing, and to get a maximum of information from our detectors set-up, we proceeded in two steps. The hydrogen layer was first frozen on the gold foil and the μ^-d emission observed using the imaging system. Simultaneously, the X-rays were measured to

identify the presence of impurities if any. Then, some neon was frozen on the hydrogen layer surface without warming excessively the target, and the muonic neon X-rays were collected.

This technique was applied for three targets which differed by their thickness and deuterium concentration as shown in Table 1.

3. Data analysis

The full data evaluation is not yet complete, the values reported here are also preliminary. Fig. 2 shows typical total energy spectra obtained with natural isotopic concentration of D₂ in hydrogen (target II), before (2a) and after (2b) the neon layer deposition. The absence of low Z element X-rays in Fig. 2a, and especially those from neon, indicates clearly that the impurity contamination in the hydrogen is negligible. Most of the neon line intensity of Fig. 2b comes from direct μ^- stops in the neon layer.

The specific contribution of the transfer process from muonic hydrogen can be observed by requiring a delayed time cut, as shown in Fig. 3 for the same data. It is essential to observe that the huge $\mu\text{Au}(6-5)$ peak of Fig. 2 is completely absent in these delayed X-ray spectra. Therefore one can conclude that the energy-time correlation and the time resolution were satisfactory.

The complete absence of the delayed muonic gold X-rays is *a priori* very intriguing because one would expect, due to the isotropic μ^- d emission, that delayed muonic gold X-rays would be apparent. In addition, no muonic carbon X-rays are observed.

This allows us to rule out the presence of a thin carbon layer at the surface of the gold foil, as discussed for an analogous experiment in Ref. 2. The intensity of the muonic transition $\mu\text{Au}(6-5)$ by direct μ^- stop is about twenty times more than via μ^- transfer, because the circular transitions are much more favoured for the former.

One can then assert that the delayed muonic neon X-rays of Fig. 3 follow the emission of muonic hydrogen into the adjacent neon layer where muon transfer eventually occurs.

Unfortunately, the statistics of the delayed events in the neon X-rays (and especially the 2-1), and the low signal to background ratio are insufficient to get a time spectrum of these events. We have proceeded by grouping the counts in four time bins: from 40 to 120 ns after μ^- signal (this spectrum will be called Del I), from 121 to 339 ns (Del II), from 340 to 1539 ns (Del III) and from 1540 to 7931 ns (Del IV). The fitted intensities of $\mu\text{Ne}(2-1)$ in each time bin are reported in Table 2.

The width of each time bin was chosen to have almost the same statistics in each. As previously stated, the delayed time structure of the neon X-rays should be exponential. This statement was confirmed by fitting the intensity structure as a function of the time window. The lifetimes corresponding to each target are also reported in Table 2.

Using Eq. 1, one obtains a set of three linear equations, each one corresponding to one value of deuterium concentration c_d , with two free parameters $\lambda_{pp\mu}$ and λ_{pd} . One finally gets :

$$\lambda_{pp\mu} = (2.18 \pm 0.32) \cdot 10^6 s^{-1}$$

$$\lambda_{pd} = (0.98 \pm 0.19) \cdot 10^{10} s^{-1}$$

These values are in very good agreement with those obtained at comparable experimental conditions i.e. with liquid hydrogen in the sixties [3-5]. The rate λ_{pd} however differs from the value which is generally accepted and measured at room temperature [6] of $(1.68 \pm 0.26) \cdot 10^{10} s^{-1}$, although recent calculations [7] do not predict such energy dependence.

4. Acknowledgements

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References

- [1] See G.M. Marshall *et al.* contribution in this volume.
- [2] J.B. Kraiman, W.H. Breunlich, M. Cargnelli, G. Chen, P.P. Guss, F.J. Hartmann, P. Kammel, J. Marton, J.J. Reidy, R.T. Siegel, W.F. Vulcan, R.E. Welsh, H. Woolverton, A. Zehnder, and J. Zmeskal, Phys. Rev. Lett. **63**, 1942 (1989).
- [3] V.P. Dzhelepov, P.F. Ermolov, E.A. Kushnirenko, V.I. Moskalev, and S.S. Gershtein, Sov. Phys. JETP **15**, 306 (1962).
- [4] E.J. Bleser, E.W. Anderson, L.M. Lederman, S.L. Meyer, J.L. Rosen, J.E. Rothberg, and I-T. Wang, Phys. Rev. **132**, 2679 (1963).
- [5] G. Conforto, C. Rubbia, E. Zavattini, and S. Focardi, Il Nuovo Cim. **33**, 4281 (1964).
- [6] A. Bertin, M. Bruno, A. Vitale, A. Placci, and E. Zavattini, Phys. Rev. **A7**, 462 (1973).
- [7] C. Chiccoli, V.I. Korobov, V.S. Melezhik, P. Pasini, L.I. Ponomarev, and J. Wosniak, Muon Cat. Fusion **7**, 87 (1992).

Table 1: Composition of the three targets used. The \oplus sign indicates that the material was added to the surface and not mixed. 1000 torr \cdot l H_2 is $4.2 \text{ mg} \cdot \text{cm}^{-2}$.

Target #	Composition		
I.	1506 torr \cdot l H_2	\oplus	194 torr \cdot l Ne
II.	1454 torr \cdot l (H_2 + 145 ppm D_2)	\oplus	193 torr \cdot l Ne
III.	725 torr \cdot l (H_2 + 1150 ppm D_2)	\oplus	175 torr \cdot l Ne

Table 2: Intensity of the $\mu\text{Ne}(2-1)$ line measured for different time windows. By assuming an exponential dependence of the delayed intensities, the lifetime $\tau = 1/\lambda$ of the μ^-p is deduced.

Target	c_d (ppm)	$\mu\text{Ne}(2-1)$ intensity for each time cut					Lifetime τ (ns)
		Prompt	Del I	Del II	Del III	Del IV	
I	0	9116	170	260	246	33	303
		(125)	(42)	(37)	(33)	(33)	(40)
II	145	5032	228	311	171	27	210
		(91)	(34)	(32)	(26)	(24)	(21)
III	1150	7057	365	110	20		56
		(100)	(36)	(28)	(23)		(9)

Figure 1: Experimental apparatus

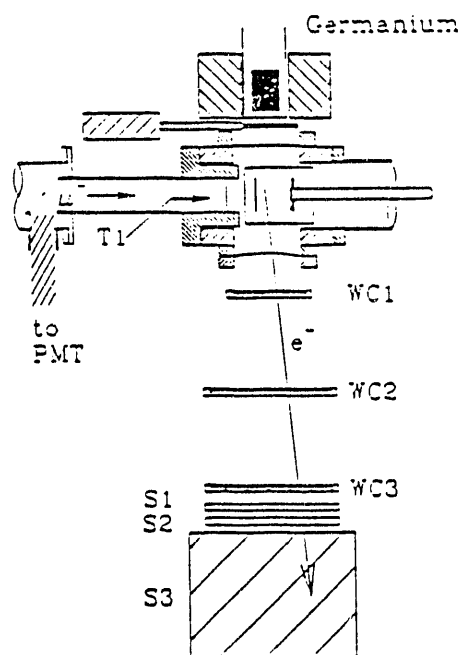


Figure 2: Part of the X-ray energy spectra without any time cut measured with target II (cf. Table 1). Upper (a) and lower (b) spectra have been obtained before and after neon deposition, respectively.

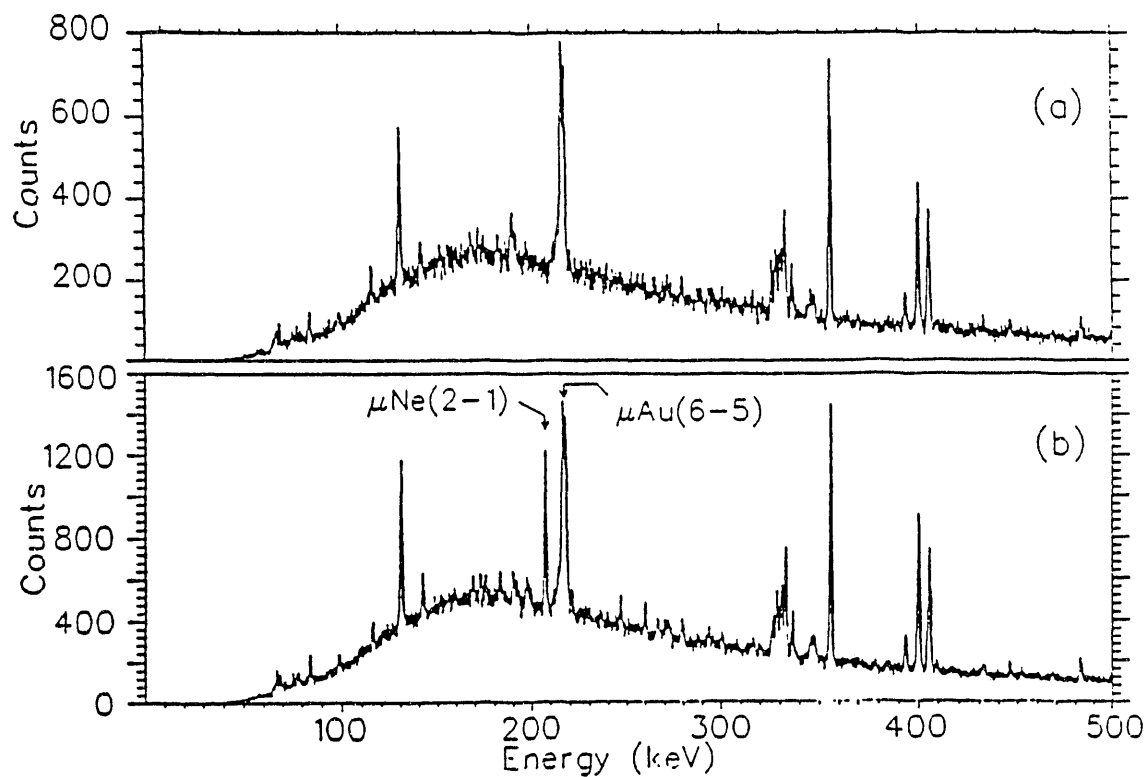
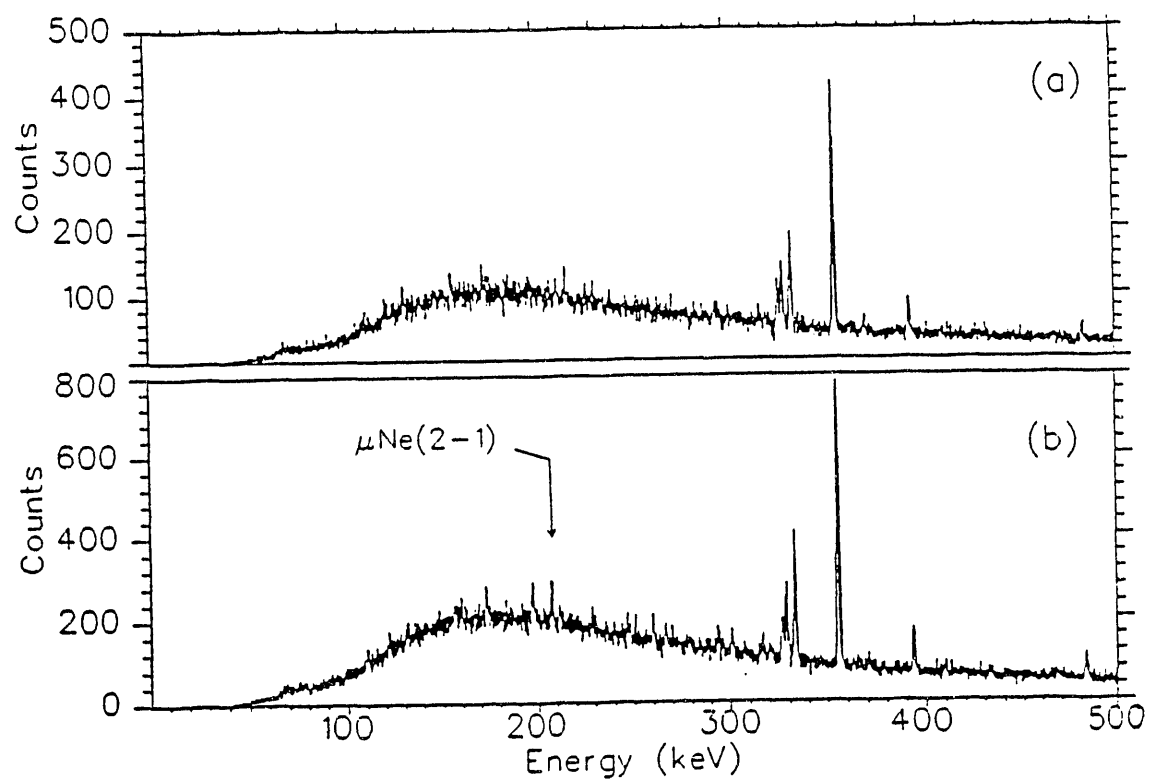


Figure 3: Part of the delayed X-ray energy spectra measured with target II (cf. Table 1). Upper (a) and lower (b) spectra have been obtained before and after neon deposition, respectively.



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