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HYPERFINE MAGNETIC FIELD MEASUREMENTS IN ErRh_4B_4 *

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We report measurements of the hyperfine magnetic field in ErRh_4B_4 in zero applied field in a range of temperatures around the reentrant temperature T_{C2} . Using the Mössbauer effect of ^{57}Fe impurities as a hyperfine field microprobe, we find that the low temperature ferromagnetic order persists well above T_{C2} into the superconducting state in this compound, thereby providing direct evidence of coexistence between superconductivity and magnetic order.

In the late 1950's, Matthias et al.[1] discovered alloy systems which gave strong indications of coexistence of superconducting and magnetically ordered states. Subsequent studies employing Mössbauer effect microprobes (^{57}Fe nuclei)[2,4] confirmed the earlier findings and proved coexistence of superconductivity and some kind of spontaneous magnetic order.

In 1977, the superconducting and magnetic properties of the ternary compound ErRh_4B_4 were documented[3]. Among the most interesting of these properties is the reentrant behavior at the so-called reentrant temperature T_{C2} . At temperatures below T_{C2} , the superconductivity of ErRh_4B_4 is destroyed by a competing ferromagnetically ordered phase. Neutron scattering experiments indicate that the magnetic order at temperatures around T_{C2} is of the ferromagnetic spin density wave type, which is consistent with the magnetic structure of ErRh_4B_4 determined by x-ray diffraction.

We report here the results of our recent measurements of the hyperfine field in the Mössbauer effect ^{57}Fe nuclei as a hyperfine magnetic field microprobe in ErRh_4B_4 . $A^{57}\text{Fe}$ nuclei substitute at an interstitial site in the ErRh_4B_4 lattice at $1/4, 1/2, 3/4, 1/2, 1/4$ in the ab plane. A small hyperfine field is found, with a temperature dependence consistent with the reentrant behavior. The hyperfine field splits into two parts at the superconducting transition temperature, T_{C2} , and the splitting is only slightly greater at T_{C2} than at the lowest temperature investigated. The $A^{57}\text{Fe}$ is expected to substitute for the Er in ErRh_4B_4 rather than the Rh , the average concentration of $A^{57}\text{Fe}$ is estimated to be less than 10 ppm.

The reentrant temperature of the doped sample was determined from measurements of the ac susceptibility. Hysteresis was observed in the data depending on whether the sample was

warmed or cooled through T_{C2} . We also observed a trend in the data toward successively higher values of T_{C2} on cooling and successively lower values on warming, as the sample was repetitively cycled through the transition region. For this reason we quote ranges in which T_{C2} was measured on warming and cooling, respectively, 0.74 K to 0.75 K and 0.66 K to 0.71 K.

MF measurements were performed with the source cooled in a ^4He cryostat. A constant-acceleration-mode MF spectrometer was employed with a room temperature single-line potassium ferrioxalate absorber.

At all temperatures the MF spectrum is comprised chiefly of a slightly asymmetrical quadrupole doublet. A simple model, based on a single site, symmetric hyperfine split quadrupole sites, is used to characterize the data at all sites. Indeed, the available Mössbauer data at various temperatures do show that there are two non-equivalent sites present in the ratio $A^{57}\text{Fe} = 1/2$. One assumption made in the MF data is therefore based on a model characterized by two nonequivalent sites A^1 and A^2 in the ratio $A^1/A^2 = 1/2$. For the quadrupole probe, the model also reflects our assumption that only one of the sites experiences a measurable hyperfine field as a result of the ferromagnetic ordering of the Er atoms in the host compound.

The details of the model are determined by suitable constraints to fit eq. (1) to the MF spectra accumulated at the extremes of the temperature range investigated. The $A^{57}\text{Fe}$ hyperfine data determines the center A^1 and quadrupole splitting of each of the two sites. Fits to spectra accumulated at the lowest temperature show that A^1 is the magnetically influenced site. The model can only be fully determined, however, with the introduction of the relative directions of the principal axes of the electric field gradient and the internal magnetic field at the A^1

site. As these directions are not known, we take them to be colinear.

At intermediate temperatures, a fitting procedure determines the optimum hyperfine-magnetic splitting of the B' site quadrupole which is consistent with our fully determined dual-site model. The introduction of an additional B' site linewidth broadening allows a second degree of freedom to the fit. This eases the constraint imposed by our simple assumption regarding the relative directions of the hyperfine fields and also reflects the possibility that all B' sites are not precisely equivalent.

The results of fits to the intermediate-temperature ME spectra exhibit several interesting features. Most notable, the B' site remains non-magnetic at all temperatures above approximately 1 K. Below this temperature, however, the magnetic splitting increases significantly, reaching a saturation value of 0.067 mm s^{-1} (corresponding to $0.21 \pm 0.02 \text{ T}$) at a temperature of approximately 0.6 K. Linewidth broadening of the B' site component lines exhibits a quite similar temperature dependence. The linewidth increases from a nominal value of 0.36 mm s^{-1} at high temperatures to approximately 0.40 mm s^{-1} at saturation; this broadening roughly corresponds to 0.2 T.

Figure 1 is a plot of the hyperfine magnetic field at the B' impurity probe sites as a function of temperature. Hyperfine magnetic fields were calculated by calibrating the spectrometer against a standard ^{57}Fe source. Although relatively small at saturation, the hyperfine magnetic fields and presumably the hyperfine magnetic ordering parameters, increase well above 1 K. In contrast to the general behavior observed with other measurement techniques (not only in superconductors, but also in normal metals), neutron scattering and heat capacity measurements were unable to resolve any appreciable hysteresis in the hyperfine field and its cooling or warming the sample through the transition region.

However these ME measurements have marginal sensitivity for observing hysteresis effects. That the hyperfine magnetic fields reported in this work are due to impurity-impurity magnetic ordering of the Fe atoms rather than to the host Er atoms seems doubtful in view of the extremely low Fe concentration. Relaxation time effects can also be ruled out for ^{57}Fe in this metallic system.

In conclusion, we have measured a temperature dependent hyperfine magnetic field at ME impurity probe sites in ErRh_4B_4 . This field persists above the reentrant temperature of the compound into the superconducting state, thereby providing direct evidence of coexistence of a magnetically-ordered state and superconductivity in this material in a small temperature interval above T_{C2} .

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REFERENCES

- [1] Matthias, B. T., Suhl, H. and Corenzweig, E., Phys. Rev. Lett. **1** (1958) 449.
- [2] Erickson, D. J., Olsen, C. E. and Taylor, R. D., in Gruverman, I. J. and Seidel, G. W. (eds.) Mossbauer Effect Methodology, Vol. 8, (Plenum Press, New York, 1973) p. 73.
- [3] Taylor, R. D., Decker, W. R., Erickson, D. J., Gierke, A. L., Matthias, B. T., Olsen, C. E. and Sklarz, L. G., in Low Temperature Physics--LT-14, Vol. 2, (Plenum Press, New York, 1974) p. 605.
- [4] Fert, M. A., Johnston, D. C., Pelong, L. J., McCallum, R. S., Maple, M. B., and Matthias, B. T., Phys. Rev. Lett. **38** (1977) 198.
- [5] Moncton, D. C., McCallum, R. S., Schmidt, P. H., Shrieve, G., Howlanson, W., Maple, M. B., MacKay, H. B., Kozl, L. D., Fisk, Z. and Johnston, D. C., Phys. Rev. Lett. **40** (1980) 2200.
- [6] Vandenberg, J. M. and Matthias, B. T., Science **184** (1973) 194.

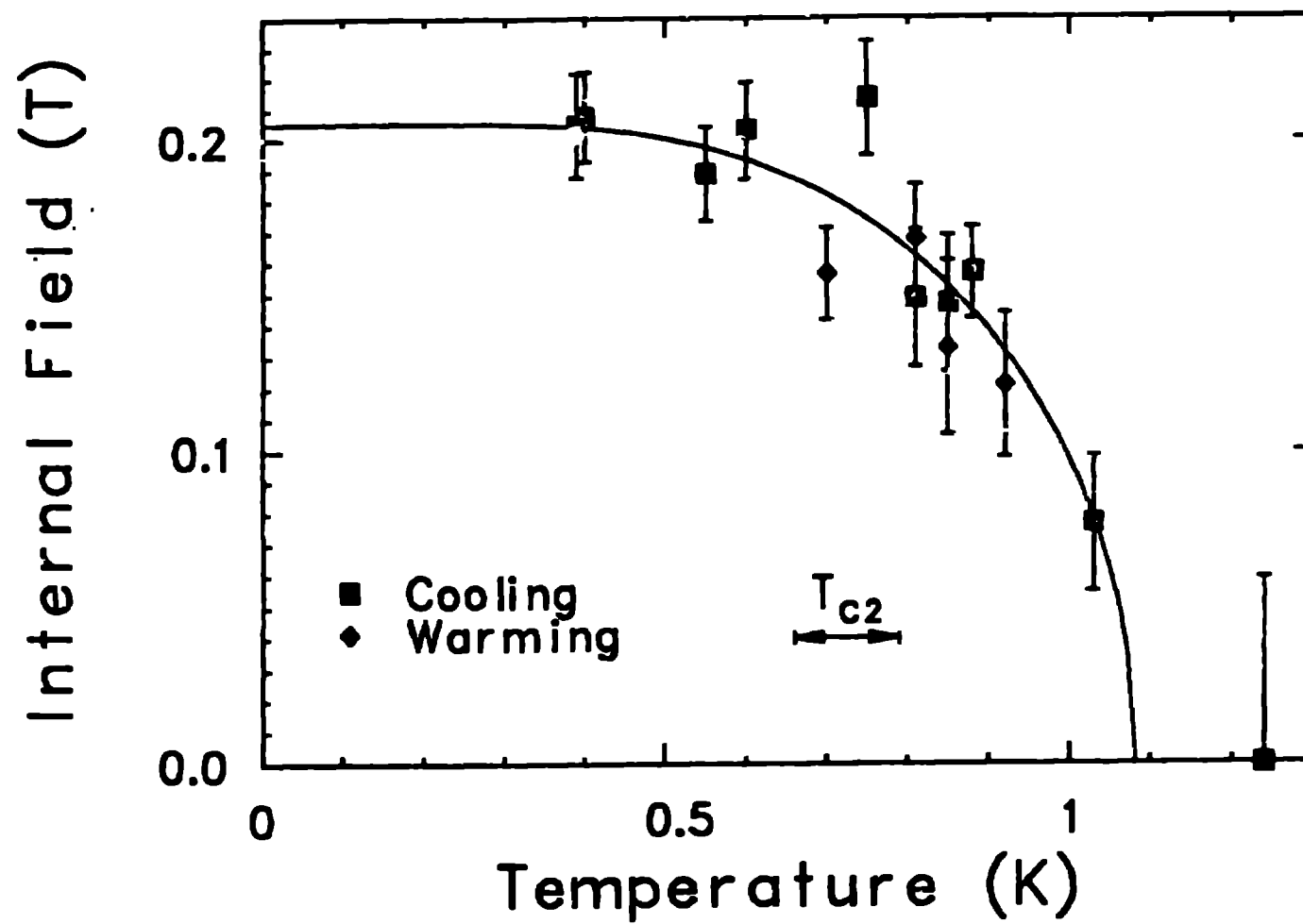


Figure 1. Internal magnetic field in ErRh_4B_4 : the solid line is a molecular field fit to the data.