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Spin-Dependent Drell-Yan and Double Prompt Photon Production to NLO QCD †

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Abstract

We present the complete $O(\alpha_s^2)$ radiative corrections to the (non singlet) polarized Drell Yan cross section for the production of a lepton pair with a nonzero q_T . The helicity of the incoming states is arbitrary. In the case of double photon, results for the longitudinal asymmetries and on the p_T behavior of the cross section are also given (to $O(\alpha_{em}^2 \alpha_s)$) in the central rapidity region of the tagged photon.

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1 Drell Yan Lepton Pair Production at Nonzero q_T

Spin physics is a new avenue for QCD. While most of studies at hadron colliders have been focused in the past on the analysis of unpolarized collisions, more recently a large research activity in QCD has shifted toward a study of the spin structure of the nucleon in DIS at HERA, while new experiments have been planned for investigating polarized $p-p$ collisions at RHIC and at HERA-N.

RHIC is supposed to run in between 50-500 GeV center of mass energy, and, as a $p-p$ collider, will have a very high luminosity (500 pb^{-1}). Therefore, we have the hope to be able to gather from this experiments enough evidence in order to better describe the flavor and the Δg contribution in the nucleon over a wide range of x ($\equiv x_{BJ}$) and Q^2 . Given the fact that the asymmetries are usually small, unless one looks at specific intervals of x and Q^2 , it is important to focus the investigations over those processes which promise to generate large asymmetries and have a particularly clean final state. Among these processes, single and double prompt photon production together with the Drell-Yan process are surely among the most interesting to be studied.

All the models which have been proposed for the polarized structure functions [1, 2], the Δg content appears to be not well determined (for instance at $x < 0.1$ and $Q^2 \sim 100 \text{ GeV}^2$, see ref. [8] for details).

The Drell-Yan process will be crucial in order to study the distribution of transverse spin in the nucleon and unveil the transversity distribution h_1 , since the gluons decouple at leading order.

We have studied the non singlet contributions to the longitudinally polarized Drell-Yan cross section to $O(\alpha_s^2)$ [3]. The calculation is done for a nonzero q_T of the lepton pair and starts (at Born level) at order α_s . In the unpolarized case, this cross section was studied by Ellis Martinelli and Petronzio [4].

The integration of our results in $t-u$ in dimensional regularization gives the bulk of the NNLO corrections to the invariant mass distributions $d\sigma/dQ^2$ of the lepton pair.

At small p_T the cross section needs resummation. To leading order this has been done in ref. [5] following the work of refs. [6, 7].

It is important to extend completely these results to the polarized case both in the case of longitudinal and transverse asymmetries.

The next to leading order (α_s) corrections to Drell-Yan total cross section are quite simple, since the process starts at $O(1)$ through the annihilation channel, and well known [9]. They have been reconsidered by various authors (see [3] and refs. therein). They involve the usual vertex and self-energy corrections to the lowest order process ($q\bar{q}\gamma$) and can be calculated quite straightforwardly.

The parton level analysis on the nonzero q_T distributions is very involved and, unfortunately, requires a recalculation of all the radiative corrections -including the virtual corrections- considered before in ref. [4]. Our calculation is the first complete check of [4], since the helicity is an additional parameter that can be sent to zero at any stage of the calculation.

We decided to use the \overline{MS} scheme since the evolution equations are also treated in the same scheme. We show that in this scheme helicity is not conserved and one needs a finite renormalization to reenter into a physical scheme.

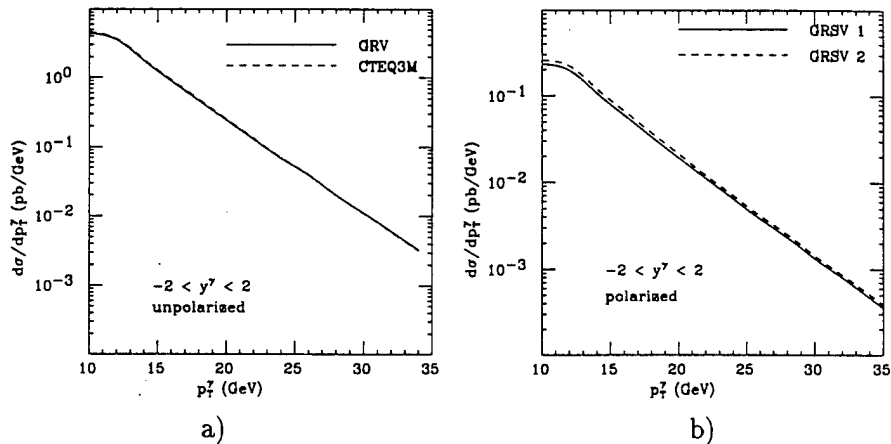


Figure 1:

The non singlet sector of Drell-Yan is defined as a difference between scattering and annihilation diagrams of quarks and antiquarks.

At the parton level we define

$$\sigma_{NS} \equiv \sum_{f,f'} (\sigma_{q_f \bar{q}_{f'}} - \sigma_{q_f q_{f'}}), \quad (1)$$

where we have included a sum over all the flavors f, f' . A careful diagrammatic analysis shows that the cross section so defined is diagonal in flavor.

We refer to our works for more details. Here we just quote the result. In the \overline{MS} , to order α_s^2 the non-singlet parton level spin asymmetry becomes

$$s \frac{d\hat{\sigma}_{NS}^{LL}}{dtdu}(s, t, u) = -s \frac{d\hat{\sigma}_{NS}^{\Sigma}}{dtdu} + s \frac{d\hat{\sigma}_0^{hat}}{dtdu} - 2s \frac{d\hat{\sigma}_4^{\Sigma}}{dtdu}, \quad (2)$$

where $d\hat{\sigma}_{NS}^{\Sigma}$ is the non-singlet unpolarized spin averaged cross section of ref. [4], while $d\hat{\sigma}_0^{hat}$ is a regularization scheme dependent helicity non-conserving piece. Similarly, $d\hat{\sigma}_4^{\Sigma}$ is the unpolarized contributions coming from the qq scattering sector.

In this scheme $P_{qq}(z) = \Delta P_{qq}(z)$ in D -dimensions and the helicity violating terms are *explicitly* needed in the convolutions of the hard scattering with the NLO evolved structure functions. The evolution DGLAP kernel has obviously to be evaluated in this same scheme.

2 Double Photon Production in Polarized Collisions to NLO

The analysis of Double Photon production can be found in ref. [8]. Here we quote briefly some of the results.

In Figs.1a and 1b we display the unpolarized and the polarized cross sections. The results indicate that it will be difficult to measure this cross section beyond about $p_T = 20$ GeV, even with the planned high luminosities at RHIC. More interestingly, Fig.1b suggests, as we expected, that GRSV distributions give very similar results and cannot be separated by this process.

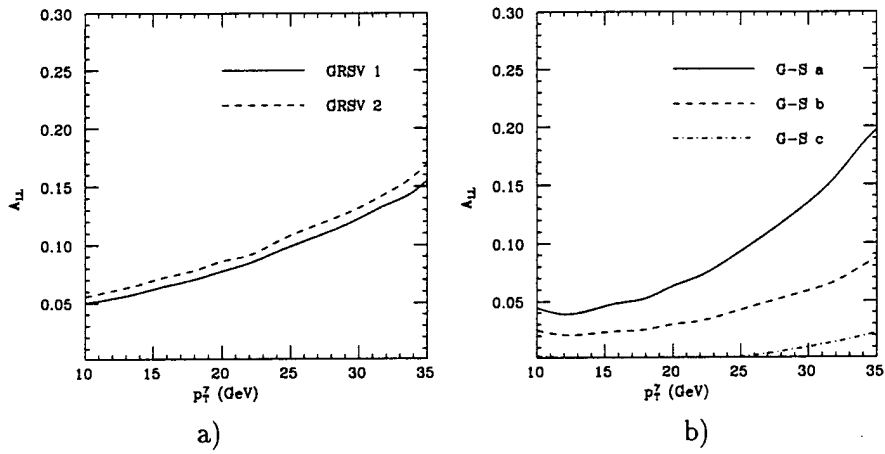


Figure 2:

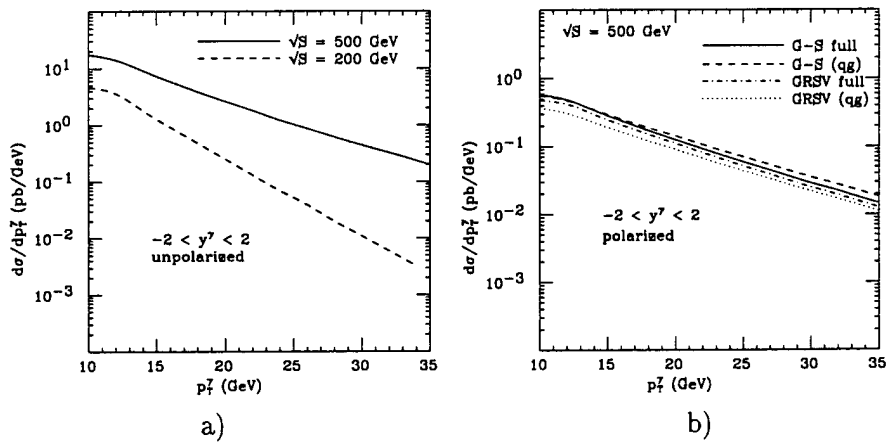


Figure 3:

In Figs.2a and 2b we show the longitudinal asymmetries as predicted by the various parameterizations. The asymmetry, A_{LL} , is defined by the relation

$$A_{LL} = \frac{\frac{d\Delta\sigma}{dp_T}}{\frac{d\sigma}{dp_T}}, \quad (3)$$

the ratio of the polarized to the unpolarized cross section, and gives a measure of the spin dependence or spin sensitivity of the process. Again the two GRSV distributions, as expected, give similar results, while the G-S ones give clearly distinguishable results. The predictions also indicate that the asymmetries are not very large in the measurable region of the cross section, varying between 5 and 8% for the GRSV and 4 and 6% for the G-S a parameterization.

Fig.2b suggests, on the other hand, that the three G-S distributions give substantially different predictions, and will probably be distinguishable. The G-S a parameterization gives similar results to the two GRSV versions.

It will require very high statistics measurements to measure this asymmetry, but this may be achievable at RHIC, as long as the cross section can be measured, given the planned detectors.

At $\sqrt{S} = 500$ GeV the situation improves somewhat with regards to the size of the cross section. Fig.3a compares the p_T distributions at $\sqrt{S} = 500$ and 200 GeV for the unpolarized cross section. As expected, at higher cms energies the cross section is substantially larger. In Fig.3b we show the polarized cross section at $\sqrt{S} = 500$ GeV. On the figure we also show the full prediction given by the G-S a distribution as well as the contribution from qg initiated subprocesses. In the case of G-S a the contribution for the $q\bar{q}$ initiated subprocess (not shown) is negative. Corresponding distributions are shown for the GRSV 1 parameterization. In this case the $q\bar{q}$ process gives a positive contribution but it is substantially smaller than the qg one. The G-S a and GRSV 1 distributions predict similar cross sections, but the relative importance of the subprocesses is clearly very different. The most interesting aspect of this result from the point of view of sensitivity to ΔG is the fact that in both cases, the qg initiated process dominates. We should mention that in this calculation we do not include contributions for the higher order process ($O(\alpha_s^2)$) $gg \rightarrow \gamma\gamma$, preferring to keep consistently to $O(\alpha_s)$.

3 Conclusions

We have presented a complete parton level study of the NLO radiative corrections to Drell Yan for a non zero q_T of the lepton pair, thereby extending the result of ref. [4] which discussed the unpolarized case. Our results completely overlap with these former calculations, except for the extraction of the helicity violating contributions. The derivation of these terms requires a complete recalculation of all the diagrams in a general way. Numerical results in the case of double photon have been given. The cross section is one power of α_{em} suppressed compared to the single photon case. Around $\sqrt{s} \sim 400 - 500$ GeV the corrections are sizeable with a K -factor $K_f \sim 2$.

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