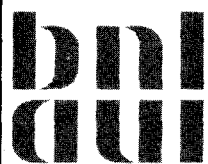


**TWO-DIMENSIONAL MODELING OF INTRA-SUBASSEMBLY  
HEAT TRANSFER AND BUOYANCY-INDUCED  
FLOW REDISTRIBUTION IN LMFBRs**

**Mohsen Khatib-Rahbar and Erik G. Cazzoli**

**Date Published — June 1984**

**CODE DEVELOPMENT, VALIDATION AND APPLICATION GROUP  
DEPARTMENT OF NUCLEAR ENERGY, BROOKHAVEN NATIONAL LABORATORY  
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## ABSTRACT

Phenomenological models and numerical techniques for prediction of coolant flow and temperature fields during forced, mixed, and free convection regimes of operation in LMFBR subassemblies are addressed. It is shown that, simplified integral solutions provide an excellent approach to assessing the importance of the intra-subassembly buoyancy induced flow redistribution, and the transverse thermal conduction and mixing effects on the assembly wide peak coolant temperatures. Furthermore, a more detailed steady-state and transient parabolic two-dimensional porous-body model, resulting in the TWIST computer code is developed. Comparison of calculated results and out-of-pile sodium and water test data indicate generally good agreement in cross-assembly temperature profiles. However, the impact of fuel pin distortion and bowing, caused by large transverse power gradients on transverse temperature distributions is found to be significant.



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## NOMENCLATURE

<u>Symbol</u>	<u>Definition</u>	<u>Unit</u>
$A_d$	Duct Wall Heat Transfer Area	$m^2$
$A_p$	Porous-Body Area	$m^2$
$A_r$	Rod Area	$m^2$
$A_w$	Wire Area	$m^2$
$B_0, B_1, B_2$	Flow Dependent Variables	-
$c_p$	Coolant Heat Capacity at Constant Pressure	$J/kg \cdot K$
$C_f$	Flow Regime Dependent Constant	-
$D_e$	Hydraulic Equivalent Diameter	$m$
$D_r$	Rod Diameter	$m$
$D_w$	Wire Diameter	$m$
$f$	Friction Factor	-
$F$	Nuclear Peak to Average Factor	-
$g$	Gravitational Acceleration	$m/s^2$
$G$	Mass Flux	$kg/m^2 \cdot s$
$Gr$	Temperature Rise Grashof Number	-
$k$	Thermal Conductivity	$W/m \cdot K$
$\ell$	Assembly Flat-to-Flat Distance	$m$
$L_h$	Rod Heated Length	$m$
$\bar{L}$	Dimensionless Assembly Length	-
$M$	Total Number of Axial Nodes	-
$M_d$	Mass of Duct Wall	$kg$
$M_s$	Mass of Interstitial Coolant	$kg$

<u>Symbol</u>	<u>Definition</u>	<u>Unit</u>
n	Exponent in the Friction Factor Equation	-
N	Total Number of Transverse Nodes	-
p	Rod Pitch	m
P	Pressure	N/m <sup>2</sup>
$\bar{P}$	Dimensionless Pressure	-
$\tilde{P}$	Modified Pressure	N/m <sup>2</sup> /kg/m <sup>2</sup>
Pe	Peclet Number (PrRe)	-
Pr	Fluid Prandtl Number	-
Q <sup>III</sup>	Volumetric Heat Generation Rate	W/m <sup>3</sup>
$\hat{Q}^{III}$	Average Volumetric Heat Generation Rate	W/m <sup>3</sup>
Re	Flow Reynolds Number	-
Ri	Richardson Number (Gr/Re <sup>2</sup> )	-
s	Even Number of Subdivisions	-
T	Temperature	k
$\hat{T}$	Average Temperature	k
T*	Reference Inlet Temperature	k
u	Superficial Velocity in the x-Direction	m/s
$\hat{u}$	Superficial Average Velocity in the x-Direction	m/s
U	Dimensionless Velocity in the x-Direction	-
v	Superficial Velocity in the y-Direction	m/s
V	Dimensional Velocity in the y-Direction	-
w	Superficial Velocity in the z-Direction	m/s
W	Weight Factor	-
x	Axial Position	m
X	Dimensionless Axial Direction	-

<u>Symbol</u>	<u>Definition</u>	<u>Unit</u>
$y$	Transverse Position	m
$Y$	Dimensionless Transverse Position	-
$Y_u$	Velocity Shape Factor	-
$Y_L$	Velocity Shape Factor for Laminar Flow	-



GREEK SYMBOLS

<u>Symbol</u>	<u>Definition</u>	<u>Unit</u>
$\alpha_i$	Fraction of the Area $A_p$ that is Occupied by Region $i$ of the Rod Bundle	-
$\alpha^*$	Temperature Peaking Factor	-
$\alpha_e$	Equivalent Thermal Diffusivity	$m^2/s$
$\beta$	Volumetric Thermal Expansion Coefficient	$K^{-1}$
$\delta$	Thermal Penetration Distance	$m$
$\delta_d$	Duct Wall Thickness	$m$
$\epsilon^*$	Mixing Parameter	-
$\epsilon_\infty^*$	Maximum Value of $\epsilon^*$	-
$\epsilon_h$	Eddy Diffusivity of Heat	$m^2/s$
$\eta^*$	Temperature Rise Hot Channel Factor	-
$\theta$	Dimensionless Temperature	-
$\kappa$	Conductivity Shape Factor	-
$\lambda_i$	Porosity for Region $i$	-
$\lambda_a$	Bundle Average Porosity	-
$\mu$	Fluid Dynamic Viscosity	$Ns/m^2$
$\nu$	Fluid Kinematic Viscosity ( $\mu/\rho$ )	$m^2/s$
$\xi$	Dimensionless Axial Distance	-
$\rho$	Fluid Density	$kg/m^3$
$\psi$	Temperature Shape Factor	-
$\psi$	Flow Intermittency Factor (Table I)	-

SUPERSCRIPT

*	Reference Inlet Condition
k	Previous Time Value
k+1	Current Time Value
^	Average Value

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## 1. INTRODUCTION

The liquid metal fast breeder reactor (LMFBR) core consists of fuel, blanket, control and shielding assemblies packed in a hexagonal configuration. The purpose of the blanket assemblies is to increase the conversion of fertile to fissile material by utilizing the neutron leakage from the fuel assemblies. For efficient breeding, a high volume fraction of the fertile material is required.

The fuel and blanket assemblies consist of cylindrical fuel elements enclosed by stainless steel cladding, arranged in a triangular pitch contained in a hexagonal duct. To minimize hot spots and provide adequate cooling, the rods are separated by spacer wires wound helically around each rod. Radial blanket assemblies occupy peripheral locations in the reactor core, and are characterized by steep power gradients; while inner blanket assemblies are located within the fuel assembly region and have higher power levels, but flatter distributions.

It is due to the presence of this radial power gradient that large sodium temperature distributions exist at full power, full flow conditions. However, at low power, low flow natural convection conditions significant heat and flow redistribution take place leading to considerable radial temperature flattening.

The prediction of assembly flow and temperature distributions must consider a wide range of coolant flow and thermal convection regimes which include laminar, transition and turbulent flow [1] during natural, mixed and forced convection conditions of reactor operation.

Analytical studies of heat transfer in rod bundles have been pursued along two separate paths - distributed parameters and lumped parameter methods [2]. In lumped parameter methods, often referred to as subchannel analyses, the rod bundle is divided into a number of subchannels whose boundaries are defined arbitrarily by surfaces of fuel elements and imaginary lines between elements and/or duct walls. Average subchannel parameters are evaluated by solving equations of continuity, momentum, and energy for each subchannel increment. Equations for each subchannel are coupled with those of its neighbors by interchannel transport of mass, momentum and energy which is treated in terms of integral transport coefficients.

Distributed parameter methods solve time-averaged Reynolds equations of momentum and energy transport for velocity and temperature distributions. These methods provide detailed information about thermal characteristics of an idealized array of fuel elements cooled by parallel flow, but have fallen short of providing overall power distribution and thermal behavior of entire fuel assemblies with spacers. Because of numerical and computer complications, as well as the lack of knowledge of local Reynolds shear stresses and turbulent heat fluxes in various directions, the application of distributed parameter methods has been rather restricted. Results of the distributed parameter method can be used to evaluate integral parameters to feed into lumped parameter analysis.

A number of models and computer codes have been developed [2-8]. These models are either limited to steady-state conditions, or are highly complex, thus requiring extensive computer storage and simulation time.

The purpose of the present study is to formulate simple, physically accurate and numerically efficient models supported by experimental data for the study of forced, mixed and natural convection flows through rod bundles using the two-dimensional porous-body approach.

Chapter 2 describes the steady-state two-dimensional porous body equations, related parameters and the governing dimensionless forms. In Chapter 3, approximate integral solutions applicable to the two limiting cases of (1) flow redistribution only; and (2) heat redistribution only are developed and discussed. A finite difference marching technique is developed for the steady-state flows in Chapter 4. The calculated results are then compared with the recent experimental data for sodium and water tests. In Chapter 5 the finite difference formulation is further extended to include transient flows; and the calculated results are also compared with transient sodium data. Finally, important results of this investigation, together with areas of future development are discussed in Chapter 6.

The computer codes and the user's guides are presented in the appendices.

## 2. STEADY-STATE POROUS-BODY MODEL

### 2.1 Governing Equations

In the porous-body treatment, the presence of fuel rods is taken into account by inclusion of a volume porosity in the governing equations. The energy generation by the rods is modeled by a continuous volumetric heat source distribution. The energy transfer in the transverse direction is modeled by molecular conduction and an empirically determined effective eddy diffusivity.

A strip section of a rod bundle is shown in Fig. 1. For simplicity, the rod bundle will be approximated by a porous slab of width  $\ell$  and height  $L_H$  corresponding to the assembly flat-to-flat and the rod heated length, respectively.

For the simulated channel described above, the following assumptions apply:

- (1) steady, two-dimensional flow,
- (2) negligible axial conduction,
- (3) negligible viscous dissipation, and
- (4) uniform pressure at any axial level.

Therefore, the conservation of mass, energy and momentum equations can be written as follows:

#### a. Conservation of Mass:

The mass continuity requires,

$$\frac{\partial(\rho u)}{\partial x} + \frac{\partial(\rho v)}{\partial y} = 0, \quad (1)$$

where  $\rho$  is the fluid density,  $u$  is the superficial velocity in the longitudinal direction and  $v$  is the superficial velocity in the transverse direction.

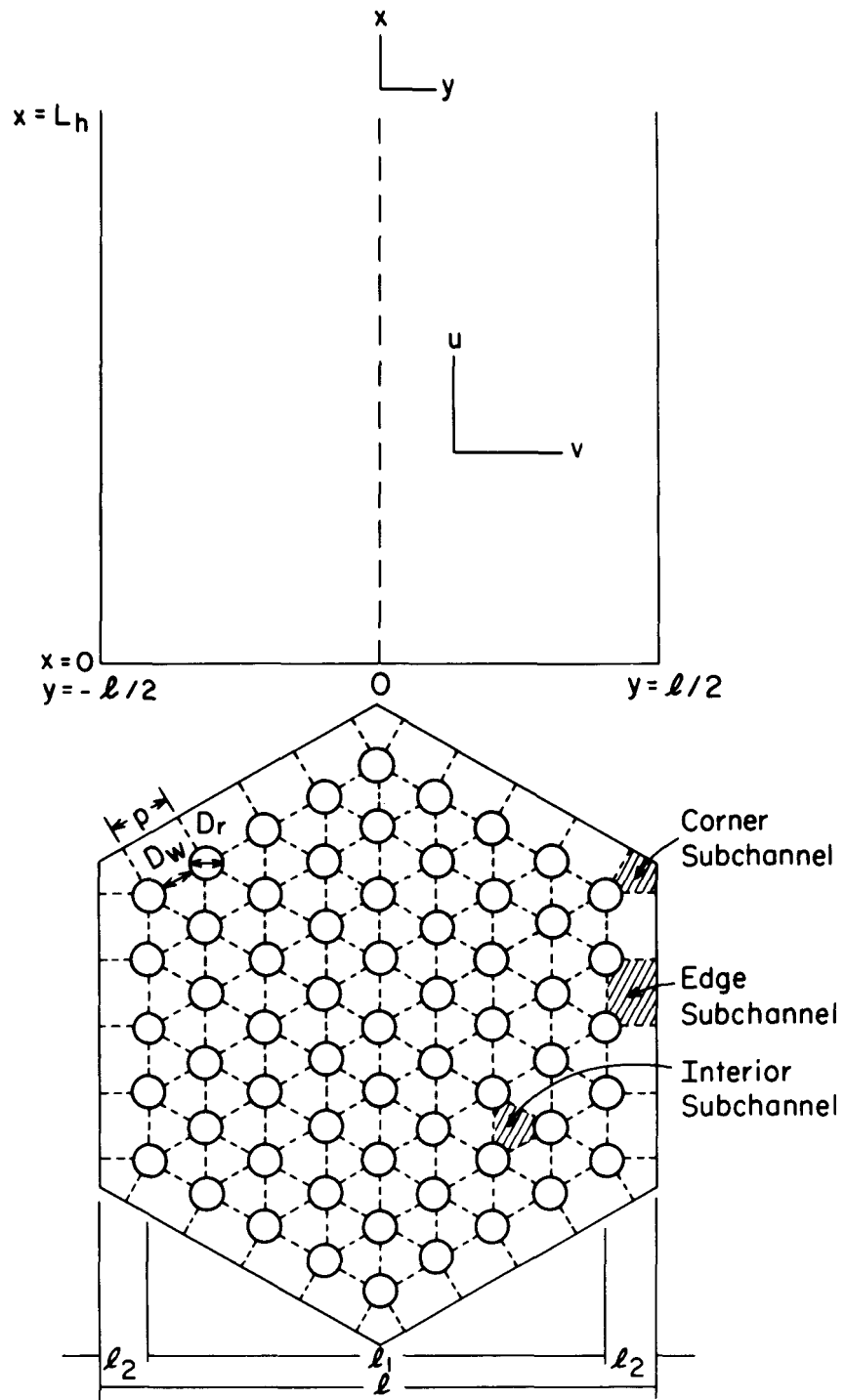


Fig. 1 Schematic of a Typical Rod Bundle and the 2-D Computational Domain

b. Conservation of Energy:

The conservation of energy requires,

$$\rho c_p \left( \frac{\partial(uT)}{\partial x} + \frac{\partial(vT)}{\partial y} \right) = \frac{\partial}{\partial y} \left( k_e \frac{\partial T}{\partial y} \right) + Q'''(x,y) \quad (2)$$

where  $c_p$  is the specific heat of sodium at constant pressure,  $T$  is the fluid temperature,  $Q'''$  is the volumetric heat generation rate and  $k_e$  is the effective thermal conductivity defined as:

$$k_e = \rho c_p \epsilon_h + \kappa k \quad (3)$$

where the eddy diffusivity of heat,  $\epsilon_h$  ( $m^2/s$ ) is defined by;

$$\epsilon_h = \hat{u} D_e \epsilon^* \quad (4)$$

Here,  $\epsilon^*$  is the mixing parameter empirically determined.

The shape factor:  $\kappa$  (also known as tortuosity) for molecular conductivity  $k$  is to account for the reduction in thermal conduction due to the presence of the fuel rods inside the sodium and is assumed to be given by\* [9];

$$\kappa = \frac{\lambda_a}{1 + 0.5 (1 - \lambda_a)} \quad (5)$$

c. Conservation of Axial Momentum:

The axial momentum is conserved according to

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = - \frac{1}{\rho} \frac{dP}{dx} - g \frac{\rho^*}{\rho} [1 - \beta(T-T^*)] + \nu \frac{\partial^2 u}{\partial y^2} - \frac{f}{2D_e \lambda^2} u^2 \quad (6)$$

---

\*Eq. (5) is only applicable to liquid metal systems (see discussions in 4.3.2)

where  $P$  is the pressure,  $T^*$  is the reference temperature,  $\rho^*$  is the density evaluated at  $T^*$ ,  $\nu$  is the kinematic viscosity,  $g$  is the gravitational acceleration,  $\beta$  is the volumetric thermal expansion coefficient,  $D_e$  is the hydraulic equivalent diameter,  $\lambda$  is the axial area fraction occupied by the coolant, and  $f$  is the skin friction factor.

The transverse momentum equation need not be formulated, since the pressure is assumed to be uniform in the transverse direction at any axial level.

## 2.2 Porous-Body Parameters

Figure 1 also shows the schematics of a rod bundle, with the porous-body parameters defined [9] as:

$$A_p = \frac{\sqrt{3}}{2} \ell^2 \quad (7)$$

$$\alpha_1 = (\ell_1/\ell)^2 \quad (8a)$$

$$\alpha_2 = (4\ell_1 \ell_2/\ell^2) \quad (8b)$$

$$\alpha_3 = 4(\ell_2/\ell)^2 \quad (8c)$$

and thus, the axial area fraction occupied by the coolant in any region of the bundle can be defined by:

$$\lambda_1 = 1 - \left[ \frac{2}{\sqrt{3}} \frac{A_r + A_w}{p^2} \right] \quad (9a)$$

$$\lambda_2 = 1 - \left[ \frac{A_r + A_w}{2 p \ell_2} \right] \quad (9b)$$

$$\lambda_3 = 1 - \left[ \frac{1}{2\sqrt{3}} \frac{A_r + A_w}{\ell_2^2} \right] \quad (9c)$$

Thus,

$$\lambda_a = \sum_{i=1}^3 (\alpha_i \lambda_i) \quad (10)$$

The hydraulic equivalent diameter in any region of the bundle is also defined by:

$$D_{e1} = \left[ \frac{\frac{2\sqrt{3}}{\pi} p^2 - (D_r^2 + D_w^2)}{(D_r + D_w)} \right] \quad (11a)$$

$$D_{e2} = \left[ \frac{\frac{8p}{\pi} \ell_2 - (D_r^2 + D_w^2)}{(D_r + D_w)} \right] \quad (11b)$$

$$D_{e3} = \left[ \frac{\frac{8\sqrt{3}}{\pi} \ell_2 - (D_r^2 + D_w^2)}{(D_r + 3D_w)} \right] \quad (11c)$$

and

$$p = D_r + D_w \quad (12)$$

where  $D_r$  is the rod outer diameter, and  $D_w$  is the diameter of the wire-wrap. Therefore, the cross-sectional area of the rod and wire can be calculated using,

$$A_r = \pi \left( \frac{D_r}{2} \right)^2 \quad (13)$$

and

$$A_w = \pi \left( \frac{D_w}{2} \right)^2 \quad (14)$$

The friction factor  $f$  is some function of the flow Reynolds number;

$$f = f(Re) \quad (15)$$

and the Reynolds number is defined as:

$$Re = \frac{\rho u D_e}{\lambda \mu} \quad (16)$$

### 2.3 Dimensionless Forms

The conservation of mass, energy and axial momentum equations (Eqs. (1), (2) and (6)) can be transformed into a dimensionless form for constant property fluids, using the following substitutions:

$$\begin{aligned} X &= \frac{x}{D_e} \\ Y &= \frac{y}{D_e} \\ U &= \frac{u}{\hat{u}} \\ V &= \frac{v}{\hat{u}} \\ \bar{p} &= \frac{p + g\rho^*x}{\frac{\rho^*\hat{u}^2}{\lambda^2}} \\ \theta &= \frac{T(x,y) - T^*}{\hat{T}(L_h) - T^*} \\ Re^* &= \frac{\rho^*\hat{u}D_e}{\lambda\mu^*} \\ Pr^* &= \frac{\mu^*c_p^*}{k_e^*} \\ Gr^* &= \frac{g \beta \rho^{*2} (\hat{T} - T^*) D_e^3}{\mu^{*2}} \end{aligned} \quad (17)$$

Therefore,

$$\frac{\partial U}{\partial X} + \frac{\partial V}{\partial Y} = 0 \quad (18)$$

$$\lambda \frac{\text{Pr}^*}{\text{Re}^*} \left( U \frac{\partial \theta}{\partial X} + V \frac{\partial \theta}{\partial Y} \right) = \frac{\partial^2 \theta}{\partial Y^2} + \frac{D_e^2 Q'''}{k_e (\hat{T} - T^*)} \quad (19)$$

and,

$$U \frac{\partial U}{\partial X} + V \frac{\partial U}{\partial Y} = - \frac{\partial \bar{P}}{\partial X} + \frac{Gr^*}{\text{Re}^{*2}} \theta + \frac{\lambda}{\text{Re}^*} \frac{\partial^2 U}{\partial Y^2} - \frac{f}{2} U^2 \quad (20)$$

alternatively,

$$Ri^* = Gr^*/\text{Re}^{*2} \quad (21)$$

where, in arriving at Equation (19), the mass continuity equation was added to the energy equation.

Equations (19) and (20) illustrate that the temperature and velocity fields in the subassembly are dependent on the magnitude of flow Reynolds number, the assembly modified Grashof number and the fluid Prandtl number. They also demonstrate, at high Reynolds number forced convection conditions the frictional forces dominate; however, at low Reynolds number conditions the buoyancy effects become increasingly significant.

### 3. APPROXIMATE INTEGRAL SOLUTIONS

In this chapter, approximate integral solutions to the conservation equations are developed for two limiting cases; namely, (a) buoyancy-induced flow redistribution in the absence of transverse thermal conduction; and (b) transverse thermal conduction in the absence of flow redistribution.

#### 3.1 Buoyancy-Induced Flow Redistribution

##### 3.1.1 Model Equations

The conservation of mass, energy and momentum equations (Eqs. (1), (2) and (3)) can be simplified assuming:

- (1) negligible transverse thermal conduction,
- (2) negligible momentum flux terms, and
- (3) constant properties.

Thus:

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0, \quad (22)$$

$$\rho c_p^* \left( u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} \right) = Q''' \quad (23)$$

and

$$-\frac{dP}{dx} = \frac{\rho^* f}{2D_e \lambda_a} u^2 + g\rho^* [1 - \beta (T - T^*)] \quad (24)$$

The solution procedure is an adaptation of the approximate integral solution proposed by Meyer [10].

Equations (22) through (24) must be solved over the slab assembly subject to the conditions:

(a) Inlet Temperature, given as,

$$T(0,y) = T^* \quad (25)$$

(b) Total Flow, the average velocity is a specified value and defined by:

$$\hat{u} = \frac{1}{\ell} \int_{-\ell/2}^{\ell/2} u(x,y) dx \quad (26)$$

(c) Assembly Edge Condition, given by,

$$v(x, \pm \ell/2) = 0 \quad (27)$$

and,

(d) Transverse Power Generation, is given by a linear distribution of the form:

$$Q'''(x,y) = \hat{Q}'''(x) \left[ 1 + (F-1) \left( \frac{2y}{\ell} \right) \right] \quad (28)$$

where F is the nuclear heat input factor, and  $\hat{Q}'''(x)$  is the average radial power distribution at a given axial location.

The friction factor can be represented by the Darcy-Weisbach relation of the form:

$$f = C_f R_e^{-n} \quad (29)$$

where  $C_f$  and n are flow dependent parameters.

Now, let us consider an axial velocity distribution which is linear in

the transverse direction at each axial position. That is,

$$u(x,y) = \hat{u} \left[ 1 + \left(\frac{2y}{\ell}\right) Y_u(x) \right] \quad (30)$$

where  $Y_u$  is the velocity shape factor, and as yet an unknown function of  $x$ .

Differentiate Eq. (24) with respect to  $y$  to obtain:

$$\frac{\partial T}{\partial y} = \frac{C_f \lambda_a^{n-2} \hat{u}^n}{g\beta D_e^{1+n} \rho^n} \left(1 - \frac{n}{2}\right) u^{1-n} \frac{\partial u}{\partial y} \quad (31)$$

and using Eq. (30) to find:

$$\frac{\partial T}{\partial y} = \frac{2B_0}{\ell} Y_u(x) u^{1-n} \quad (32)$$

where,

$$B_0 = \frac{C_f \lambda_a^{n-2} \hat{u}^n}{g\beta D_e^{1+n} \rho^n} \left(1 - \frac{n}{2}\right) \hat{u} \quad (33)$$

is the flow regime-dependent constant.

Integrate Eq. (32) over the  $y$ -direction to obtain:

$$T(x,y) = T_c(x) + \frac{B_0 \hat{u}^{1-n}}{(2-n)} \left\{ \left[ 1 + \frac{2y}{\ell} Y_u(x) \right]^{2-n} - 1 \right\} \quad (34)$$

where  $T_c(x)$  is the coolant temperature at the center of the assembly ( $y=0$ ), and is a function of axial position  $x$  only.

Now, substitute Eq. (30) back into the continuity equation (Eq. (22)) and integrate the results with respect to  $y$ , and evaluate the constant of integration by the edge condition of Eq. (27), to obtain:

$$v = \left(\frac{\ell}{4}\right) \hat{u} \frac{dY_u}{dx} \left[ 1 - \left(\frac{2y}{\ell}\right)^2 \right] \quad (35)$$

Therefore, equations (32), (34) and (35) can be used to evaluate.

$$u \frac{\partial T}{\partial y} = \frac{1}{2} B_0 \hat{u}^{2-n} Y_u \frac{dY_u}{dx} \left[ 1 + \left( \frac{2y}{\ell} \right) Y_u \right]^{1-n} \left[ 1 - \left( \frac{2y}{\ell} \right)^2 \right] \quad (36)$$

$$u \frac{\partial T}{\partial x} = \hat{u} \left[ 1 + \left( \frac{2y}{\ell} \right) Y_u \right] \frac{dT_c}{dx} + B_0 \hat{u}^{2-n} \frac{dY_u}{dx} \left[ 1 + \left( \frac{2y}{\ell} \right) Y_u \right]^{2-n} \left( \frac{2y}{\ell} \right) \quad (36b)$$

Equations (32), (34), (35) and (36) provide a complete description of the velocity and temperature fields within the assembly if  $T_c(x)$  and  $Y_u(x)$  can be determined. This solution is achieved by defining two weight functions,  $W_1(y)$  and  $W_2(y)$  [10,11], multiplying the energy equation (Eq. (23)) by each and integrating each resulting equation with respect to  $y$  across the assembly.

Choose the weight functions such that,

$$W_1(y) = 1 \quad - \ell/2 \leq y \leq \ell/2 \quad (37a)$$

and

$$W_2(y) = \begin{cases} -1 & - \ell/2 \leq y < 0 \\ +1 & 0 \leq y \leq \ell/2 \end{cases} \quad (37b)$$

Therefore,

$$\rho^* c_p^* \int_{-\ell/2}^{+\ell/2} \left( u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} \right) dy = \int_{-\ell/2}^{+\ell/2} Q'''(x,y) dy \quad (38)$$

and,

$$\begin{aligned}
 \rho^* c_p^* \left[ - \int_{-\ell/2}^0 (u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y}) dy + \int_0^{+\ell/2} (u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y}) dy \right] \\
 = - \int_{-\ell/2}^0 Q'''(x,y) dy + \int_0^{+\ell/2} Q'''(x,y) dy
 \end{aligned} \tag{39}$$

Thus, integrating Equations (38) and (39); and after simplifications one obtains:

$$\frac{dY_u}{dx} = \left( \frac{\hat{Q}'''(x)}{\rho^* c_p^* \hat{u}} \right) \left( \frac{1}{B_0 B_1 \hat{u}^{1-n}} \right) \left( 1 - \frac{F-1}{Y_u} \right) \tag{40}$$

and

$$\frac{dT_c}{dx} = \frac{\hat{Q}'''(x)}{\rho^* c_p^* \hat{u}} - B_0 B_2 \hat{u}^{1-n} \frac{dY_u}{dx} \tag{41}$$

where,

$$\begin{aligned}
 B_1 = & \left[ \frac{(1+Y_u)^{2-n} - (1-Y_u)^{2-n}}{4(2-n)} \right] + \frac{1}{4Y_u^2} \left[ \frac{(1+Y_u)^{4-n} - (1-Y_u)^{4-n}}{(4-n)} - \right. \\
 & \left. \frac{(1+Y_u)^{2-n} - (1-Y_u)^{2-n}}{(2-n)} \right] + \left( \frac{1}{Y_u^3} - \frac{1}{Y_u} \right) \left[ \frac{(1+Y_u)^{2-n} + (1-Y_u)^{2-n} - 2}{2(2-n)} \right] \\
 & - \left[ \frac{(1+Y_u)^{4-n} + (1-Y_u)^{4-n} - 2}{2Y_u^3(4-n)} \right] \tag{42}
 \end{aligned}$$

and

$$B_2 = \left[ \frac{(1+Y_u)^{2-n} - (1-Y_u)^{2-n}}{4(2-n)} \right] + \frac{1}{4Y_u^2} \left[ \frac{(1+Y_u)^{4-n} - (1-Y_u)^{4-n}}{(4-n)} - \frac{(1+Y_u)^{2-n} - (1-Y_u)^{2-n}}{(2-n)} \right] \tag{43}$$

For laminar flow ( $n=1$ ), Equations (40) through (43) reduce to the form derived by Meyer [10] that is:

$$\frac{dY_u}{d\xi} = \frac{(F-1) - Y_u}{1 - \frac{2}{3} Y_u^2} \quad (44)$$

and

$$\frac{dT_c}{d\xi} = B_o \left[ 1 - \frac{Y_u^2 - (F-1) Y_u}{Y_u^2 - \frac{3}{2}} \right] \quad (45)$$

where the dimensionless axial distance  $\xi$  is defined as,

$$\xi = \frac{\hat{Q}''' x}{B_o \rho_c^* \hat{u}_p} \quad (46)$$

Note that

$$Y_u = 0 \quad \text{at} \quad \xi = 0 \quad (47)$$

and reaches  $Y_L$  at the assembly exit, which corresponds to  $\xi = \xi_L$ , where

$$\xi_L = \frac{\hat{Q}''' L_h}{B_o \rho_c^* \hat{u}_p} \quad (48)$$

Let the average temperature rise in the assembly be  $\Delta\hat{T}$ , where

$$\Delta\hat{T} = \hat{T} - T^* = \frac{\hat{Q}''' L_h}{\rho_c^* \hat{u}_p} \quad (49)$$

Defining a Grashof number based on this temperature rise by:

$$Gr^* = g B D_e^3 \Delta\hat{T} / \nu^{*2} \quad (50)$$

Combining Equations (16), (48), (49) and (50) one obtains:

$$\xi_L = \frac{2}{C_f} \frac{Gr^*}{Re^*} \quad (51)$$

The temperature peaking factor is given by:

$$\alpha^* = (T(L_h, -l/2) - T^*)/\Delta\hat{T} \quad (52)$$

The temperature rise hot channel factor is then defined as the fractional reduction in F caused by buoyancy-induced flow redistribution. That is,

$$\eta^* = \frac{\alpha^* - 1}{F - 1} \quad (53)$$

Thus, for laminar flow, a convenient form for  $\eta^*$  can be obtained [10] as:

$$\eta^*_{\text{Laminar}} = \frac{Y_L - \frac{1}{3} Y_L^2}{\xi_L (F - 1)} \quad (54)$$

which, in a limit as  $Y_L$  approaches  $(F - 1)$ , becomes:

$$\eta^*_{\text{Laminar}} \sim \frac{1 - \frac{1}{3} (F - 1)}{\xi_L} \quad (55)$$

Equations (48) through (55) show that the hot channel factor depends not only on Reynolds number, but also on the average temperature rise (namely,  $Gr^*$ ) and the flow regime characteristics.

Equations (40) and (41), together with appropriate flow-dependent parameters, can be integrated numerically to obtain a complete characteristic of the temperature, and flow field as a function of the radial power factor F, at various average power and flow conditions in the bundle, resulting in the FREDIS code described in Appendix A.

### 3.1.2 Results

In this section, the influence of the bundle friction factor on the flow redistribution is examined. The model described in the previous sections is then used to investigate the transverse temperature flattening effects due to thermal buoyancy-induced intra-assembly flow redistribution caused by radial power gradients in the Westinghouse Blanket Test Model and the results are compared to the experimental measurements.

#### 3.1.2.1 Influence of the Friction Factor

Because of the importance of the friction factor defined in Eq. (29) previously, we will examine its influence on flow redistribution in this subsection. Table I lists some typical friction factor correlations for both fuel and blanket wire-wrapped assemblies. It is seen that the friction factor is inversely proportional to the Reynolds number in the laminar flow regime, and behaves like the Blasius approximation in the turbulent flow regime, while it is usually interpolated in the transition regime.

Figures 2 and 3 illustrate the influence of the friction factor on flow redistribution for a typical fuel and blanket assembly, respectively. It is clearly evident that a reduction in the frictional drag leads to a significantly enhanced redistribution towards the warmer side of the bundle, and thus, lower peak temperatures. Furthermore, in the transition from laminar to turbulent regime, discontinuities may occur.

In the present study, for conservatism, in the transition flow regime the friction factor is evaluated based on the turbulent flow conditions.

#### 3.1.2.2 Comparison with Experimental Measurements

The Westinghouse Blanket Test section shown schematically in Fig. 4 is a one-to-one scale model of a typical LMFBR 61-pin blanket assembly. The power level, axial and radial power distributions of the assembly are simulated using electrically heated fuel rod simulators.

Table I Friction Factor Parameters for Fuel and Blanket Assemblies

Assembly	Flow Regime	n	$C_f$	Reference
Fuel	Laminar	1	84	Spencer, [17]
	Transition & Turbulent	-	$f = (1.078 + 0.3 \times 10^6 / Re^2) f_c$ $1/f_c = -2 \log [2.51/Re f_c]$	Spencer, [17]
	Turbulent	0.25	$0.316 \left[ \frac{1.034}{(P/D)^{0.124}} + \frac{29.7(P/D)^{6.94} Re^{0.086}}{(H/D)^{2.239}} \right]^{0.885}$	Novendstern, [18]
Blanket	Laminar	1	110	Engel, [19]
	Transition	0	$\frac{110}{Re} \sqrt{1-\psi} + \frac{0.48}{Re^{25}} \sqrt{\psi} ; \psi = \frac{Re - 400}{4600}$	Engle, [19]
	Turbulent	0.25	0.48	Engel, [19]

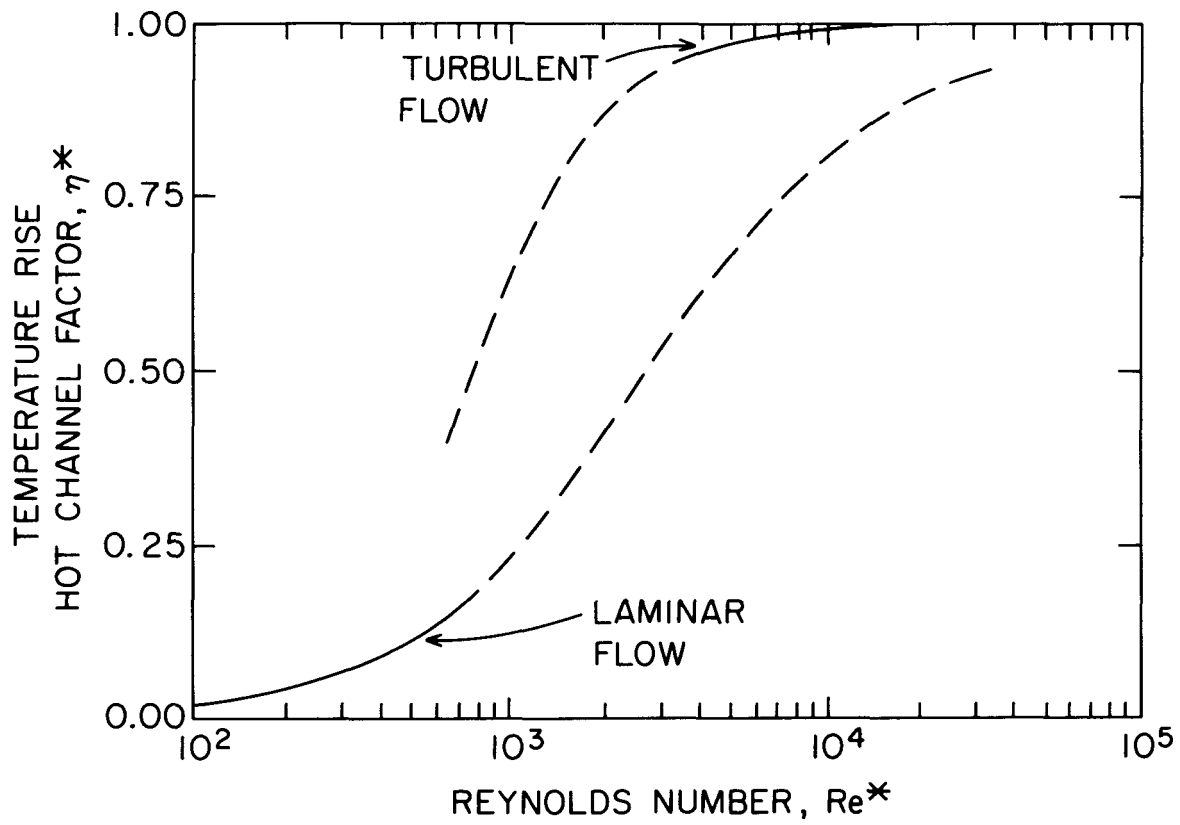


Fig. 2 Impact of Flow Regime on Calculated Fuel Assembly Temperature Rise Hot Channel Factor

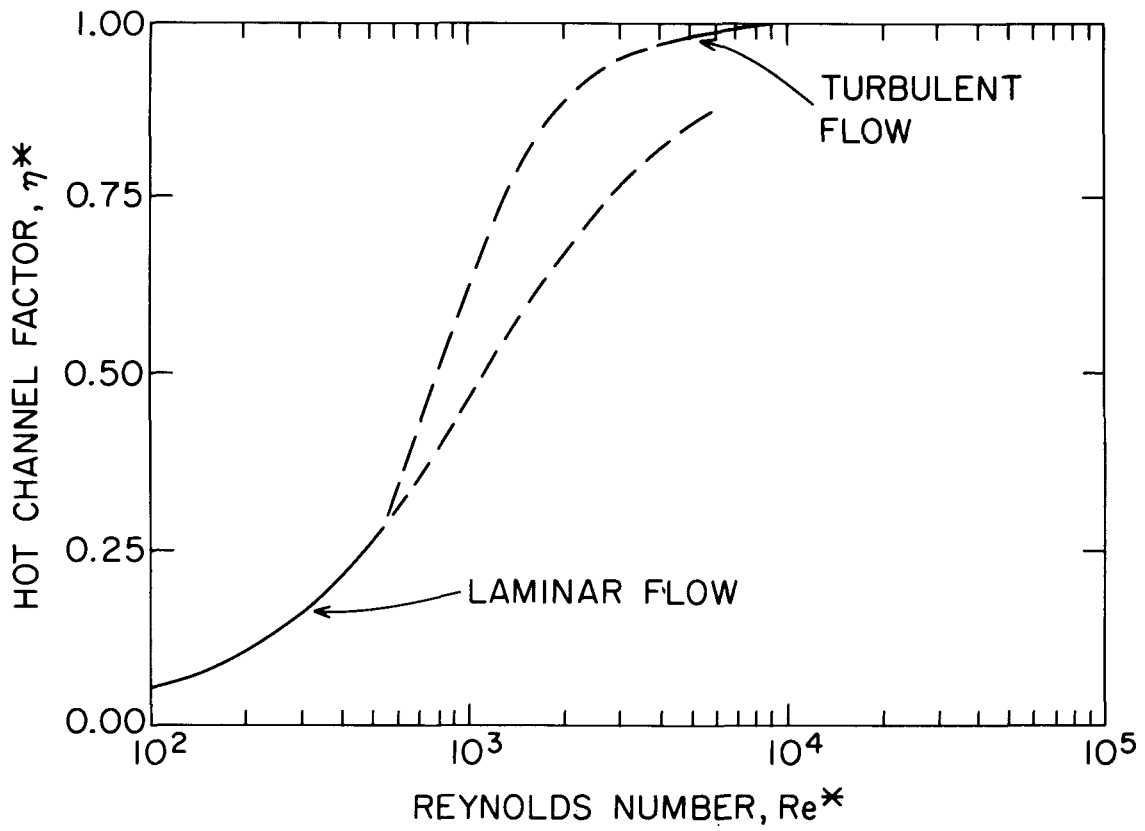


Fig. 3 Impact of Flow Regime on Calculated Blanket Assembly Temperature Rise Hot Channel Factor

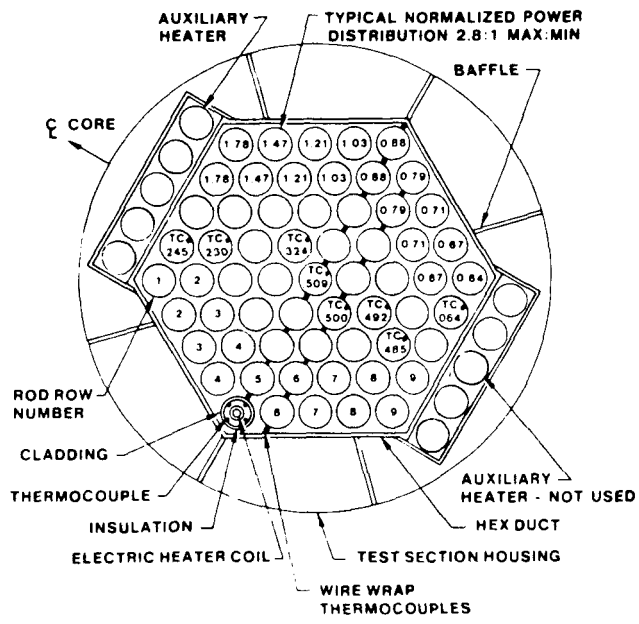
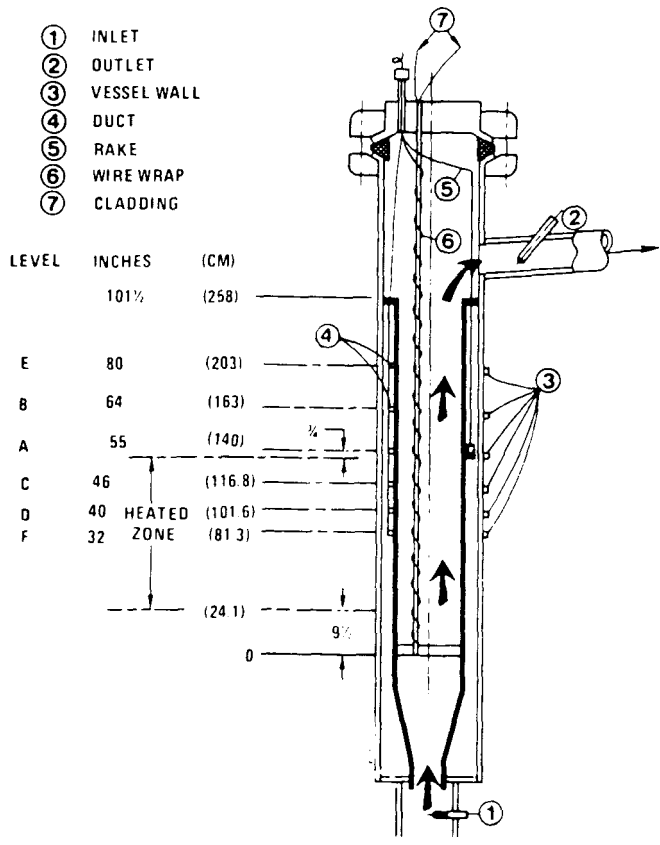


Fig. 4 Schematic of the Westinghouse 61-Pin Blanket Test Model

The test section is vertically mounted, with sodium entering at the bottom and flowing upward through the heated rod bundle, discharging from the duct near the vessel exit nozzle. The space outside the duct contains stagnant sodium and is compartmentalized by numerous horizontal and radial baffles to minimize heat loss by convective transfer in the annular space between the duct and the vessel walls.

Table II summarizes the geometric characteristics of the test section [12-16]. The test section has been used extensively to study thermal-hydraulic characteristics of LMFBR blanket assemblies for steady state, transient, forced, mixed and free convection conditions. Due to the controlled experimental conditions, these tests can serve a very useful purpose for identification of phenomenological heat transfer and fluid dynamics issues applicable to fast reactor cores.

Figures 5 and 6 show typical steady-state transverse temperature profiles at various axial locations corresponding to different levels of power and bundle flow conditions for radial power gradients of (2.79:1) and (2.04:1), respectively. It is seen that the maximum temperature is reached away from the bundle edge in the interior of the assembly, as a result of enhanced cooling at the edge subchannels due to higher sodium flow rate. The temperature flattening at low Reynolds numbers discussed earlier is very clearly evident.

Figure 7 shows the calculated temperature rise hot channel factor as a function of inlet Reynolds number using the present flow redistribution model. Also shown are the measured data of Engel, et al., [14] corresponding to the same conditions.

At high Reynolds number forced convection conditions, the subassembly radial temperature distribution is primarily influenced by transverse thermal conduction and it is unaffected by thermal buoyancy. However, comparison with experimental data seems to indicate that this effect is counterbalanced by the larger magnitude of the velocity as assumed in the present model.

Table II Geometric Characteristics of the Westinghouse 61 Pin Test Model

$D_r$	(m)	$1.318 \times 10^{-2}$
$D_w$	(m)	$9.400 \times 10^{-4}$
$p/D_r$	(-)	1.082
$l$	(m)	0.1140
$l_1$	(m)	0.0988
$l_2$	(m)	0.0076
$L_h$	(m)	1.1430
$A_r$	(m <sup>2</sup> )	$1.3643 \times 10^{-4}$
$A_w$	(m <sup>2</sup> )	$6.9398 \times 10^{-7}$
$A_p$	(m <sup>2</sup> )	$1.1255 \times 10^{-2}$
$D_{e1}$	(m)	$3.5163 \times 10^{-3}$
$D_{e2}$	(m)	$7.1810 \times 10^{-3}$
$D_{e3}$	(m)	$4.8997 \times 10^{-3}$
$D_{ea}$	(m)	$3.7282 \times 10^{-3}$
$\alpha_1$		0.7511
$\alpha_2$		0.2311
$\alpha_3$		0.0178
$\lambda_1$		0.2214
$\lambda_2$		0.3674
$\lambda_3$		0.3077
$\lambda_a$		0.2567

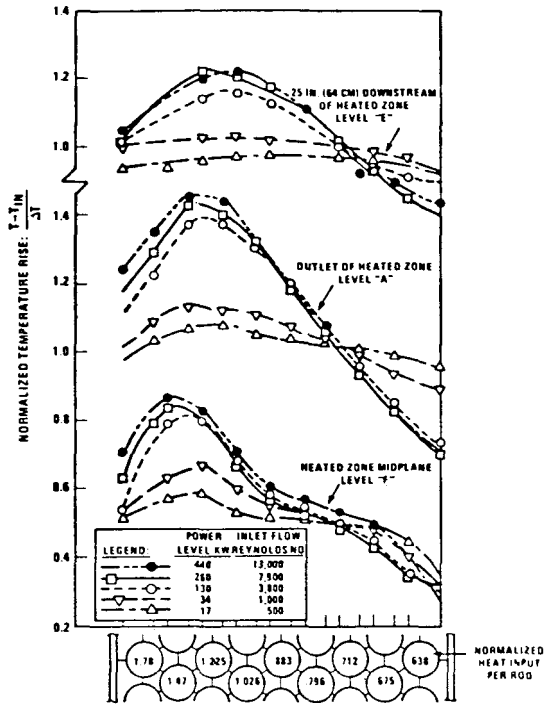


Fig. 5 Measured Transverse Temperature Distribution (2.79:1)

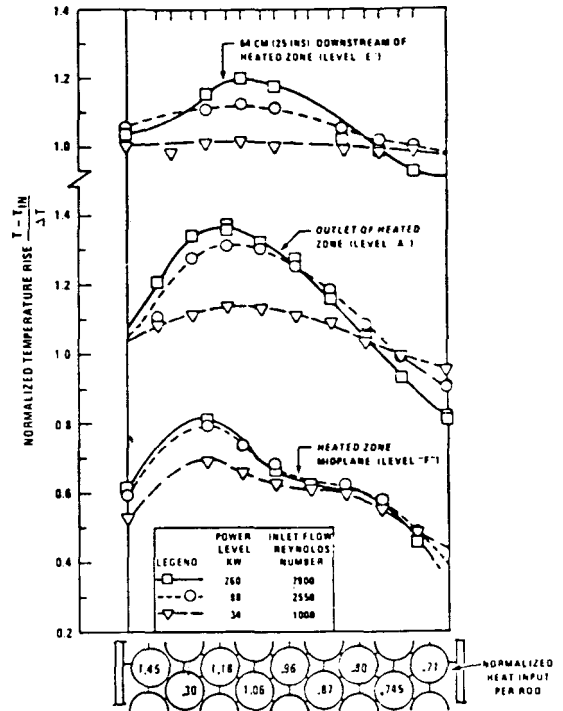


Fig. 6 Measured Transverse Temperature Distribution (2.04:1)

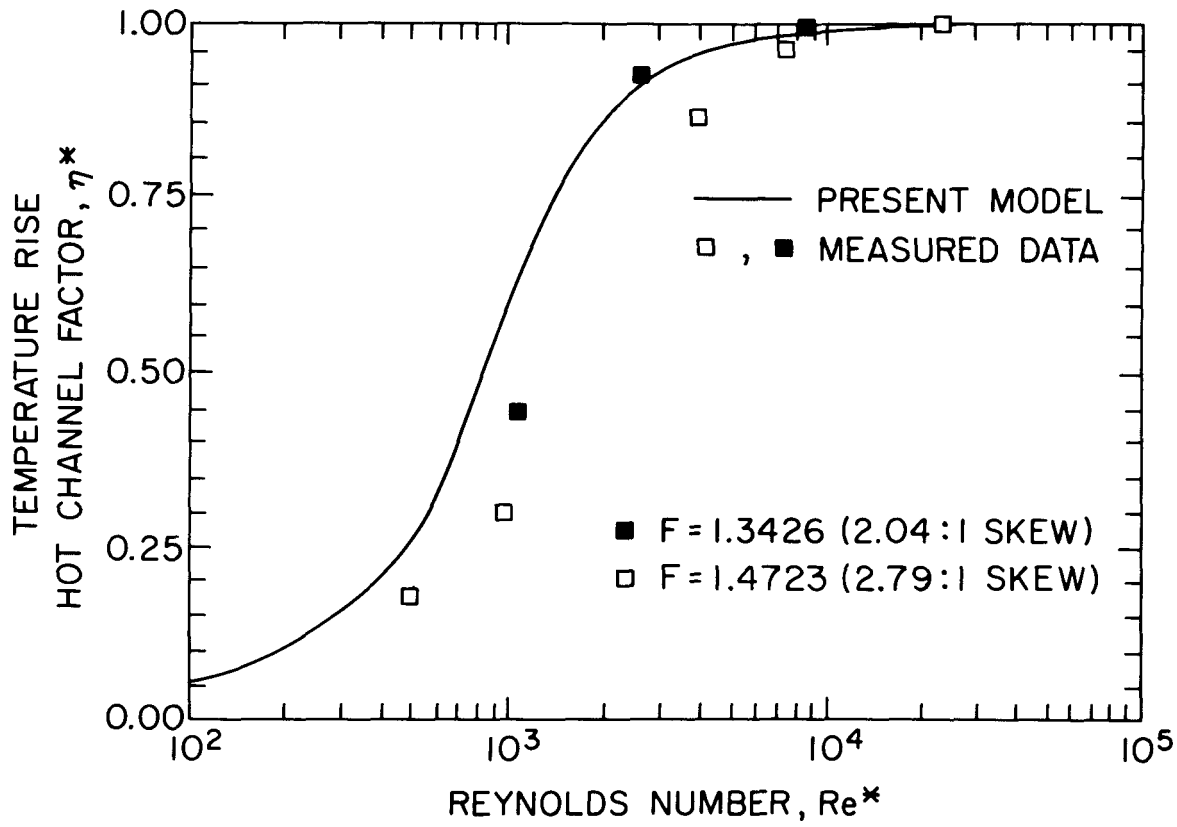


Fig. 7 Comparison of Calculated and Measured Temperature Rise Hot Channel Factors

At lower Reynolds numbers, for mixed and free convection conditions, the influence of thermal buoyancy and transverse thermal conduction become increasingly important, and thus the discrepancy between the measured and calculated values become increasingly significant.

Figure 8 shows the calculated temperature peaking factor as a function of the temperature rise Grashof number to inlet flow Reynolds number ratio ( $Gr^*/Re^*$ ) at various radial power gradients. It is seen that the temperature peaking factor remains constant in the forced convection regime, where it is only dependent on the radial power gradient. The initiation of the mixed convection regime is seen to be a unique function of  $Gr^*/Re^*$  and is found to be independent of the radial power gradient, in excellent agreement with recent experimental observation of Engel, et al., [16]. Furthermore, as the buoyancy effects become more significant, a change in shape becomes apparent beyond which free convection regime is dominant.

These calculations seem to indicate that, even though, buoyancy-induced flow redistribution has a dominant influence on transverse temperature flattening in LMFBR rod bundles, thermal conduction effects must also be included, especially during mixed and free convection regimes. The significance of molecular and turbulent diffusion is further discussed in the following section.

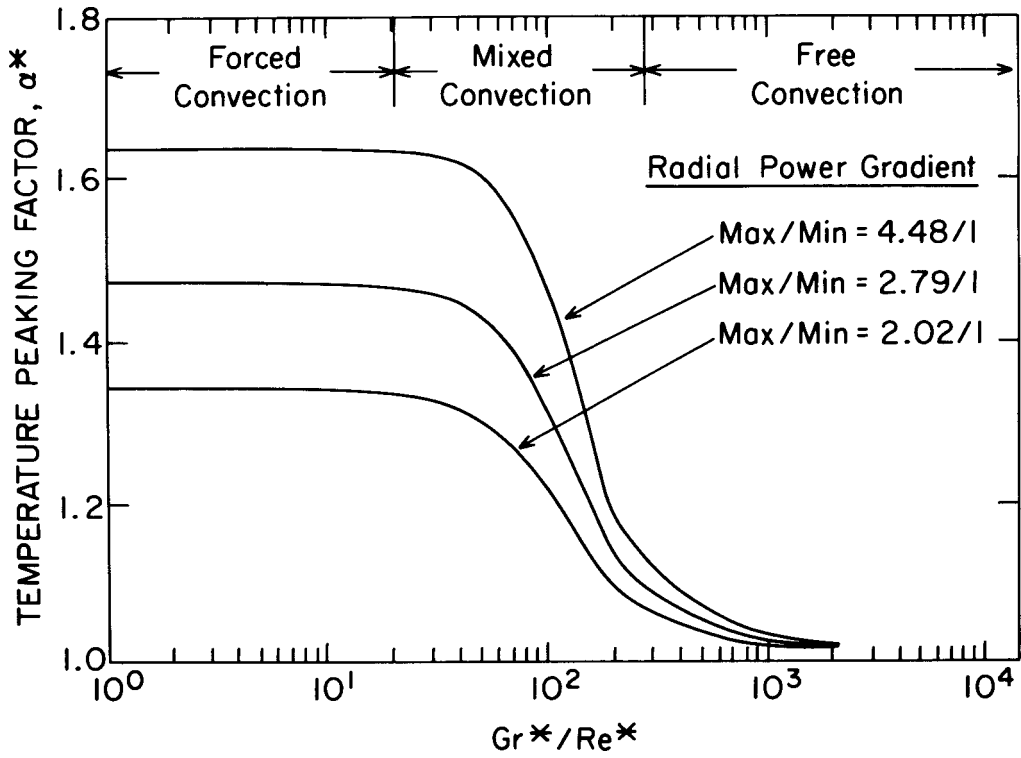


Fig. 8 Calculated Temperature Peaking Factor as a Function of Thermal Buoyancy Parameter

## 3.2 Transverse Heat Transfer

### 3.2.1 Meyer's Integral Method

In the absence of buoyancy-induced flow redistribution, the coolant flow is predominantly in the axial direction and is assumed to be independent of axial and transverse positions within the rod bundle. Therefore, the conservation of energy equation (Eq. 2) can be simplified to;

$$\rho c_p \hat{u} \frac{\partial T}{\partial x} = k_e \frac{\partial^2 T}{\partial y^2} + Q''' \quad (56)$$

where  $k_e$  is defined in equation (3) as:

$$k_e = \rho c_p \epsilon_h + \kappa k \quad (57)$$

Subject to the following boundary conditions: '

$$\left\{ \begin{array}{l} T(0, y) = T^* \\ \frac{\partial T}{\partial y}(x, \pm l/2) = 0.0 \end{array} \right. \quad (58)$$

Meyer [11] has solved this equation using an approximate integral technique. He assumed,

$$T = \hat{T}(x) \left\{ 1 + (F-1) \left(\frac{2y}{l}\right) - (F-1) \left(\frac{\delta}{l}\right) \psi_1 \right\} \quad \text{for } \delta < \frac{l}{2} \quad (59)$$

where,  $\delta$  is the distance into the assembly that the effect of insulation at the duct wall has penetrated,  $\psi_1$  is a skew symmetric function about  $y = 0$ ; that is:

$$\psi_1(-y, \delta) = -\psi_1(y, \delta) \quad (60)$$

and where, for non-negative values of  $y$ ,

$$\psi_1 = \begin{cases} 0 & 0 \leq y < (\frac{\ell}{2} - \delta) \\ \left[ \frac{2y + 2\delta - \ell}{2\delta} \right]^2 & (\frac{\ell}{2} - \delta) \leq y \leq \ell/2 \end{cases} \quad (61)$$

The temperature distribution of Equation (59) was chosen [11] to have the following nature:

- the temperature  $\hat{T}$  and the distance  $\delta$  are functions of  $x$  only;
- if at any axial location, the temperature is averaged between  $y = -\ell/2$  and  $y = \ell/2$ , the resulting average is  $\hat{T}$ ;
- The first two terms within the braces of Eq. (59) give the exact solution for the case of zero conductivity;
- The temperature shape of the penetration term (the final term within the braces in Eq. (59)) is parabolic;
- both temperature and heat flux are continuous; and
- the boundary conditions (Eq. (58)) are satisfied both at the inlet and at the duct walls if:

$$\begin{cases} \hat{T}(0) = T^* \\ \delta(0) = 0 \end{cases} \quad (62)$$

Meyer [11] obtained an integral solution using the following weighted residual technique:

Step (a) Substitute Eq. (59) into Eq. (57); multiply by the weight function  $w_1(y)$ , where;

$$w_1(y) = 1 \quad -\ell/2 \leq y \leq +\ell/2 \quad (63)$$

and integrate the resulting equation from  $y = -\ell/2$  to  $y = +\ell/2$ .

Step (b) Again substitute Eq. (59) into Eq. (57); multiply by the weight function  $w_2(y)$ , where;

$$w_2(y) = \begin{cases} 0 & -\ell/2 \leq y < 0 \\ 1 & 0 \leq y \leq \ell/2 \end{cases} \quad (64)$$

and integrate the resulting equation from  $y = \ell/2$  to  $y = +\ell/2$ .

Equations obtained by completing steps (a) and (b), can, after substituting the boundary conditions from Equation (62), be solved for  $\hat{T}$  and  $\delta$ :

$$\hat{T}(x) = T^* + \frac{\hat{Q}'''}{\rho c_p \hat{u}} x \quad (65)$$

and

$$\delta(x) = \sqrt{\frac{3 k_e x}{\rho c_p \hat{u}}} \quad (66)$$

Equation (66) ceases to apply if the value of  $\delta$  calculated is greater than  $(\ell/2)$ . That is, Eq. (66) applies only if  $x \leq x_{\max}$ ; where  $x_{\max}$  is given by:

$$x_{\max} = \left[ \frac{\rho c_p \hat{u} \ell^2}{12 k_e} \right] \quad (67)$$

An alternative equation that is equivalent to Eq. (66) is then:

$$\frac{\delta}{\ell} = \frac{1}{2} \sqrt{\frac{x}{x_{\max}}} \quad \text{for } x \leq x_{\max} \quad (68)$$

For values of  $x$  larger than  $x_{\max}$ , another integral equation for temperature was employed. Meyer [11] chose the following:

$$T(x,y) = \hat{T}(x) + T_d(x) \psi_2(y) \quad (69)$$

where  $\hat{T}$  and  $T_d$  are functions of  $x$  and  $\psi_2$  is a skew symmetric function about  $y = 0$ ; that is:

$$\psi_2(-y) = -\psi_2(y) \quad (70)$$

and where,

$$\psi_2(y) = 1 - \left[ \frac{2y - \ell}{\ell} \right]^2 \quad 0 \leq y \leq \ell/2 \quad (71)$$

The temperature distribution of Eq. (69) has the following nature:

- if at any level, the temperature is averaged between  $y = -\ell/2$  and  $y = \ell/2$ , the resulting average is  $\hat{T}$ ;
- both temperature and heat flux are continuous;
- the boundary conditions (Equation (58)) are satisfied at the duct wall; and
- the temperature can be made to match with Equation (59) at  $x = x_{\max}$

The integral equation was solved using the same steps performed earlier to obtain:

$$T_d(x) = \frac{1}{4} \hat{T}(x) \left( \frac{x_{\max}}{x} \right) \left\{ 3 - \text{Exp} [-(x - x_{\max})/x_{\max}] \right\} \quad (72)$$

Now, define a dimensionless assembly length  $\bar{L}$  as:

$$\bar{L} = L_h/x_{\max} \quad (73)$$

or using Eq. (67),

$$\bar{L} = \left[ \frac{12k_e L_h}{\rho c_p \hat{u} \ell^2} \right] \quad (74)$$

and then by using Eqs. (59), (65), (68), (72) and the definition of the temperature rise hot channel factor (Eq. (53)), one obtains:

$$\eta^* = \begin{cases} 1 - \frac{1}{2} \sqrt{\bar{L}} & \text{for } \bar{L} \leq 1 \\ \frac{1}{4\bar{L}} [3 - \text{Exp}(1-\bar{L})] & \text{for } \bar{L} > 1 \end{cases} \quad (75)$$

An alternative form for  $\bar{L}$  may be obtained using the definition of Reynolds number and Prandtl number (Eq. (17)) to obtain:

$$\bar{L} = \left( \frac{1}{\text{Re}} \right) \left( \frac{1}{\text{Pr}} \right) \left( \frac{k_e}{k} \right) \left( \frac{12 D_e L_h}{\lambda_a \ell^2} \right) \quad (76)$$

where according to Eqs. (3) and (4),

$$\left( \frac{k_e}{k} \right) = \kappa + \text{Re Pr } \lambda_a \epsilon^* \quad (77)$$

and Meyer [11] suggests the following approximation to Chiu's results [20]:

$$\left( \epsilon^*/\epsilon_{\infty}^* \right) = \begin{cases} 1 & \text{Re} \geq 7000 \\ 1 - \left[ 1 - \frac{\text{Re}}{7000} \right]^2 & \text{Re} < 7000 \end{cases} \quad (78)$$

where  $\epsilon_{\infty}^*$  is usually chosen to match the experimental data.

Therefore, Eqs. (75), (76) and (77) relate the temperature rise hot channel factor to the coolant Reynolds number, Prandtl number and the assembly geometry.

Meyer [11] also found a high flow and a low flow asymptote for  $\bar{\Gamma}$  as:

$$(\bar{\Gamma})_{\text{high flow}} \approx \left[ \frac{12 D_e L_h \epsilon_{\infty}^*}{\lambda^2} \right] \quad (79)$$

and

$$(\bar{\Gamma})_{\text{low flow}} \approx \left( \frac{1}{\text{Re}} \right) \left( \frac{1}{\text{Pr}} \right) \left[ \frac{12 D_e L_h \kappa}{\lambda_a \lambda^2} \right] \quad (80)$$

Note, that at high flow  $\bar{\Gamma}$  and thus  $\eta^*$  is independent of the flow rate.

A useful approximation to  $\eta^*$  can be obtained by taking the larger of  $\bar{\Gamma}_{\text{high flow}}$  and  $\bar{\Gamma}_{\text{low flow}}$  for a given value of Reynolds number.

### 3.2.2 Results

Results of Meyer's [11] calculations are presented in Figure 9 for CRBRP fuel and blanket assemblies. It is seen that major reduction in temperature rise hot channel factors occur below full flow values ( $\text{Re} < 3000$  for fuel assemblies and  $\text{Re} < 1500$  for blanket assemblies). However, full flow transverse thermal conduction is more significant in the blanket assemblies ( $\eta^* = 0.68$ ) as compared to the fuel assemblies ( $\eta^* = 0.90$ ). Figure 9 also shows that the asymptotic results supply an excellent representation of the complete solution.

It is therefore evident that the thermal conduction effect is quite significant at high flow forced convection; however, its influence becomes more pronounced at low Reynolds numbers, where it is of the same order as the flow redistribution observed in the previous section.

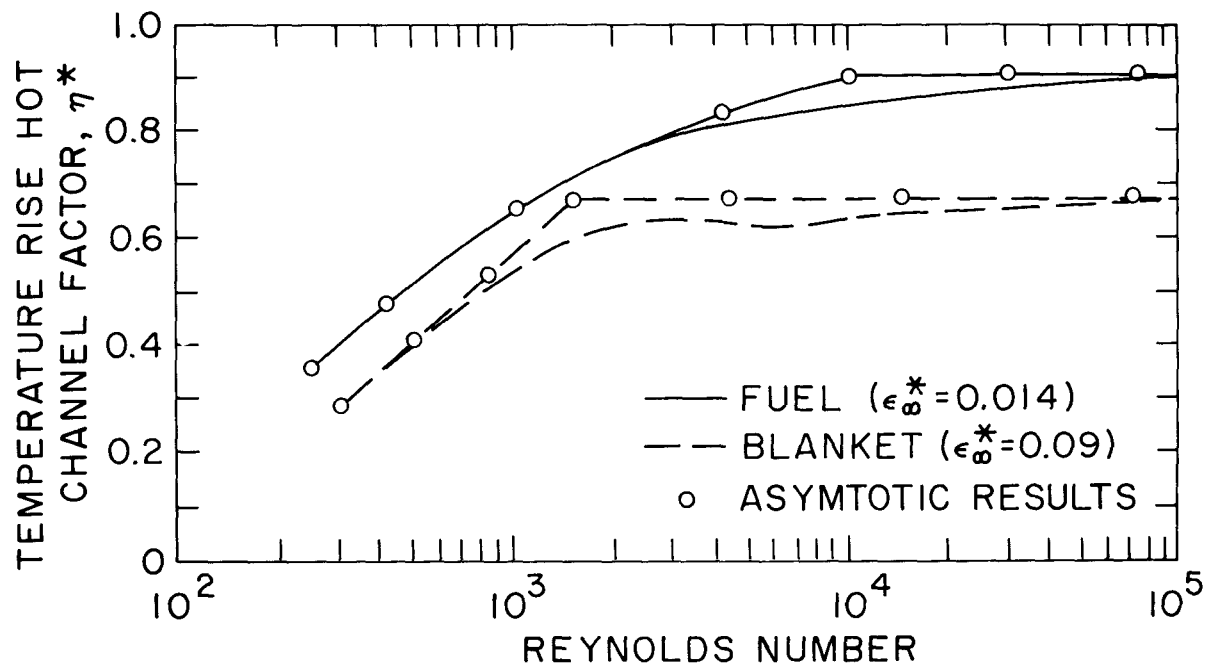


Fig. 9 Impact of Transverse Conduction and Mixing on Temperature Rise Hot Channel Factor

## 4. FINITE DIFFERENCE SOLUTION

In the previous chapter, approximate integral solutions to the conservation equations were developed. It was also demonstrated that both thermal buoyancy and transverse conduction can significantly reduce the transverse temperature gradient within an LMFBR subassembly subject to a radial power gradient.

In this chapter, finite difference solutions to a more detailed form of the conservation equations are developed, which retain both the buoyancy-induced redistribution and the transverse thermal conduction effects in the same two-dimensional slab geometry. Finally, calculated results are compared to experimental measurements for model validation.

### 4.1 Steady-State Two-Dimensional Parabolic Flow Model

In LMFBR subassemblies, the flow is predominantly in the axial direction and can therefore be treated as a two-dimensional parabolic flow. Under these conditions, convection in the axial direction dominates the axial diffusion. It is this feature that imparts the one-way character to the streamwise direction. Obviously, no reverse flow in the axial direction would be acceptable.

The solution for two-dimensional parabolic flows is obtained by starting with a known distribution of temperature and velocity at an upstream station and marching in the streamwise direction. For every forward step, the distribution of temperature and velocity in the cross-stream coordinate is calculated at one streamwise station. Thus, computationally only a one-dimensional problem needs to be solved.

Recall the conservation of mass, energy and momentum equations (Eqs. (1), (2) and (6)) of Chapter 2 where,

$$\frac{\partial(\rho u)}{\partial x} + \frac{\partial(\rho v)}{\partial y} = 0 \quad (81)$$

$$\left( u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} \right) = \alpha_e \frac{\partial^2 T}{\partial y^2} + \frac{Q'''_e}{\rho c_p} \quad (82)$$

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = - \frac{1}{\rho} \frac{dP}{dx} - g \frac{\rho^*}{\rho} [1 - \beta(T-T^*)] + v \frac{\partial^2 u}{\partial y^2} - \frac{f}{2D_e \lambda^2} u^2 \quad (83)$$

where  $\alpha_e$  is the effective thermal diffusivity and is defined as;

$$\alpha_e = k_e / \rho c_p \quad (84)$$

Equations (81) through (83) must be solved subject to the following boundary conditions (see Fig. 10):

(1) Inlet Velocity and Temperature are specified as;

$$\begin{cases} u(0,y) = F(u_{in}) \\ T(0,y) = T^* \end{cases} \quad (85)$$

(2) No Slip at the duct walls requires,

$$\begin{cases} u(x,0) = u(x,\ell) = 0 \\ v(x,0) = v(x,\ell) = 0 \end{cases} \quad (86)$$

(3) Adiabatic Duct requires,

$$\frac{\partial T}{\partial x}(x,0) = \frac{\partial T}{\partial x}(x,\ell) = 0 \quad (87)$$

## 4.2 Method of Solution

### 4.2.1 Finite Difference Equations

In finite difference form, Eqs. (81) through (83) can be expressed as follows:

#### Mass Continuity

$$\left[ \frac{(\rho u)_{i+1,j} + (\rho u)_{i+1,j+1} - (\rho u)_{i,j} - (\rho u)_{i,j+1}}{2\Delta x} \right] \quad (88)$$
$$+ \left[ \frac{(\rho v)_{i+1,j+1} - (\rho v)_{i+1,j}}{\Delta y} \right] = 0$$
$$i = 1, 2, 3, \dots, M$$
$$j = 2, 3, 4, \dots, N-2$$

Note, the somewhat unusual representation of Eq. (88) is chosen for a reason which will become clear shortly.

Now, add all of the velocities in the transverse direction, and substitute the boundary conditions of Eq. (86) to obtain:

$$\sum_{j=3}^{N-2} (\rho u)_{i+1,j} = \sum_{j=3}^{N-2} (\rho u)_{i,j} \quad (89)$$

If both sides of Eq. (89) are multiplied by  $\Delta y$ , it can be seen to be the trapezoidal rule integration form of the equation, i.e.,

$$\left[ \int_0^{\ell} \rho u \, dy \right]_{i+1} = \left[ \int_0^{\ell} \rho u \, dy \right]_i \quad (90)$$

Subsequently, dividing both sides of Eq. (90) by  $\ell$ , one obtains;

$$G = \langle \rho u \rangle_{i+1} = \langle \rho u \rangle_i \quad (91)$$

which indicates that the mass flux is constant throughout the axial location. Equation (91) is often [21] referred to as the "equation of constraint".

Energy

$$\begin{aligned}
 & u_{i,j} \left[ \frac{T_{i+1,j} - T_{i,j}}{\Delta x} \right] + v_{i,j} \left[ \frac{T_{i+1,j+1} - T_{i+1,j-1}}{2\Delta y} \right] \\
 & = \alpha_e \left[ \frac{T_{i+1,j+1} - 2T_{i+1,j} + T_{i+1,j-1}}{(\Delta y)^2} \right] + \left[ \frac{Q'''_{i,j}}{\rho c_p} \right]
 \end{aligned} \tag{92}$$

and Axial Momentum

$$\begin{aligned}
 & u_{i,j} \left[ \frac{u_{i+1,j} - u_{i,j}}{\Delta x} \right] + v_{i,j} \left[ \frac{u_{i+1,j+1} - u_{i+1,j-1}}{2\Delta y} \right] \\
 & = - \frac{\tilde{\Delta P}}{\Delta x} - \frac{f}{2D\lambda_{i,j}^2} u_{i,j} u_{i+1,j} + g_B (T_{i+1,j} - T^*) \\
 & + v_{i,j} \left[ \frac{u_{i+1,j+1} - 2u_{i+1,j} + u_{i+1,j-1}}{(\Delta y)^2} \right]
 \end{aligned} \tag{93}$$

for  $i=1, 2, 3, \dots, M$

$j=3, 4, 5, \dots, N-2$

where  $i$  and  $j$  denote  $x$  and  $y$ , respectively. The mesh points of the rectangular grid structure superimposed on the channel flow field is shown in Fig. 10. The points  $j=1$  and  $j=N$  designate the duct wall, while  $i=1$  and  $i=M$  represent the subassembly inlet and exit, respectively.

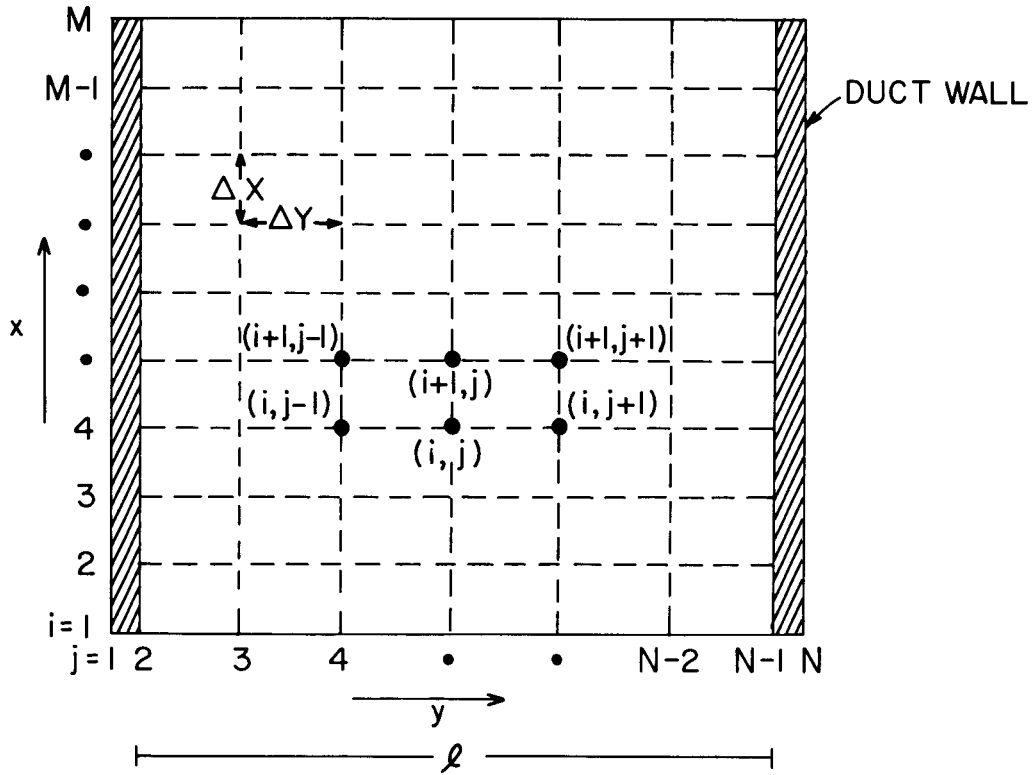


Fig. 10 Coordinate System for the Finite Difference Equations

The energy equation (Eq. (92)) together with the boundary conditions of Eq. (87) yield;

$$\begin{bmatrix}
 \phi_{i,3} & \sigma_{i,4} & 0 & 0 & \dots & -0 \\
 \xi_{i,3} & \phi_{i,4} & \sigma_{i,5} & 0 & \dots & -0 \\
 0 & \xi_{i,4} & \phi_{i,5} & \sigma_{i,6} & \dots & \\
 \vdots & \vdots & \vdots & \vdots & \ddots & \vdots \\
 0 & \dots & \dots & \dots & \dots & \sigma_{i,N-2} \\
 0 & \dots & \dots & \dots & \xi_{i,N-3} & \phi_{i,N-2}
 \end{bmatrix}
 \begin{bmatrix}
 T_{i+1,3} \\
 T_{i+1,4} \\
 \vdots \\
 T_{i+1,N-2}
 \end{bmatrix}
 =
 \begin{bmatrix}
 C_{i,3} \\
 C_{i,4} \\
 \vdots \\
 C_{i,N-2}
 \end{bmatrix}
 \quad (94a)$$

and,

$$T_{i+1,1} = T_{i+1,2} = T_{i+1,3} \quad (94b)$$

where

$$\phi_{i,j} = \left[ \frac{u_{i,j}}{\Delta x} + \frac{2\alpha_e}{(\Delta y)^2} \right] \quad (95a)$$

$$\sigma_{i,j} = \left[ \frac{v_{i,j-1}}{2\Delta y} - \frac{\alpha_e}{(\Delta y)^2} \right] \quad (95b)$$

$$\xi_{i,j} = \left[ \frac{v_{i,j+1}}{2\Delta y} + \frac{\alpha_e}{(\Delta y)^2} \right] \quad (95c)$$

$$C_{i,j} = \left[ \frac{Q_{i,j}^H}{\rho c_p} + \frac{T_{i,j} u_{i,j}}{\Delta x} \right] \quad (95d)$$

Similarly, Eqs. (89) and (93) can be written as:

$$\begin{bmatrix}
 \rho_{i+1,3} & \rho_{i+1,4} & \rho_{i+1,5} & \rho_{i+1,6} & \rho_{i+1,7} & \dots & \rho_{i+1,N-2} & 0 \\
 \gamma_{i+1,3} & \Omega_{i+1,4} & 0 & 0 & 0 & \dots & -0 & \frac{1}{\Delta x} \\
 \Gamma_{i+1,3} & \gamma_{i+1,4} & \Omega_{i+1,5} & 0 & 0 & \dots & - & \vdots \\
 0 & \Gamma_{i+1,4} & \gamma_{i+1,5} & \Omega_{i+1,6} & 0 & \dots & - & \vdots \\
 0 & 0 & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\
 0 & 0 & \vdots & \vdots & \vdots & \vdots & \vdots & \frac{1}{\Delta x} \\
 0 & 0 & \dots & -0 & \Gamma_{i+1,N-4} & \gamma_{i+1,N-3} & \Omega_{i+1,N-2} & \frac{1}{\Delta x}
 \end{bmatrix}
 \begin{bmatrix}
 u_{i+1,3} \\
 u_{i+1,4} \\
 \vdots \\
 u_{i+1,N-2} \\
 \tilde{\Delta P}_{i+1}
 \end{bmatrix}$$

$$= \begin{bmatrix}
 S_{i+1} \\
 C_{i+1,3} \\
 C_{i+1,4} \\
 \vdots \\
 \vdots \\
 C_{i+1,N-2}
 \end{bmatrix} \tag{96a}$$

and

$$u_{i+1,1} = u_{i+1,2} = u_{i+1,N-1} = u_{i+1,N} = 0 \tag{96b}$$

where:

$$v_{i+1,j} = \left[ \frac{u_{i,j}}{\Delta x} + \frac{f_{i,j} u_{i,j}}{2D_{i,j} \lambda_{i,j}^2} + \frac{2v_{i,j}}{(\Delta y)^2} \right] \quad (97a)$$

$$\Omega_{i+1,j} = \left[ \frac{v_{i,j}}{2\Delta y} - \frac{v_{i,j}}{(\Delta y)^2} \right] \quad (97b)$$

$$\Gamma_{i+1,j} = - \left[ \frac{v_{i,j}}{2\Delta y} + \frac{v_{i,j}}{(\Delta y)^2} \right] \quad (97c)$$

$$S_{i+1} = \sum_{j=3}^{N-2} (\rho_{i,j} u_{i,j}) \quad (97d)$$

$$C_{i+1,j} = \frac{u_{i,j}}{\Delta x} + g\rho [T_{i+1,j} - T^*] \quad (97e)$$

and

$$\Delta \tilde{P}_{i+1} = \left[ \frac{\Delta P + g\rho^* \Delta x}{\rho} \right]_{i+1} \quad (97f)$$

#### 4.2.2 Numerical Solution Procedure

The numerical solution procedure is as follows:

- (1) Start at row  $i=1$ , and solve Eq. (94) for  $T_{2,1}, T_{2,2}, \dots, T_{2,N}$  using the Gaussian elimination technique [22].
- (2) Calculate the physical properties as a function of  $T_{i+1,j}$ 's.

- (3) Solve Eq. (96) for  $u_{2,1}, u_{2,2}, \dots, u_{2,N}$  and  $\Delta\tilde{P}_2$  using the Gauss-Jordan reduction method [22].
- (4) Now, back substitute into Eq. (88), together with Eq. (86), to obtain  $v_{2,1}, v_{2,2}, \dots, v_{2,N}$ .
- (5) March in the stream-wise direction up to row  $i=M$ ; repeating steps (1) through (4).

### 4.3 Results and Comparison with Experimental Data

In this section, the two-dimensional parabolic flow model developed in the previous sections is applied to the Westinghouse Blanket Sodium [14] Test data and the Toshiba Water Simulation measurements [25,26].

#### 4.3.1 Comparison with Sodium Test Data

The Westinghouse Blanket Test section described previously in Section 3.1.2.2 is a highly prototypic test section simulating forced, mixed and free convection conditions typical of LMFBR blanket assemblies subject to various power and flow conditions. Table II summarizes the geometric characteristics while Table III lists the experimental conditions as given by Engel, et al., [14].

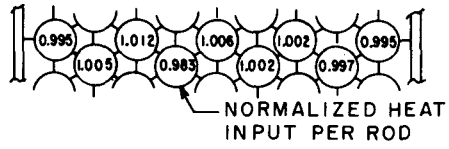
Figure 11 illustrates the transverse power profiles for the aforementioned experiments as used for the present simulation (see also Appendix B for more details).

The power distribution in the bundle is represented by the following chopped cosine axial distribution; with peak to average of 1.4; that is,

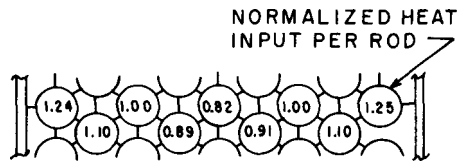
$$Q'''(x,y) = 1.40 \hat{Q}''' \bar{\psi}(y) \cos \left\{ 0.8738 \pi \left( \frac{x}{L_h} - \frac{1}{2} \right) \right\} \quad (98)$$

Table III Westinghouse Blanket Test Conditions

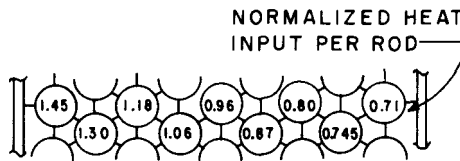
Power Skew (Max/Min)	Power Level (Kw)	Inlet Reynolds Number
1:1	34	1000
	130	3800
	260	7900
	440	13000
Inner Blanket	22	550
	44	1100
	165	4400
	440	11500
2:1	34	1000
	88	2550
	260	7900
2.8:1	17	500
	34	1000
	130	3800
	260	7900
	440	13000



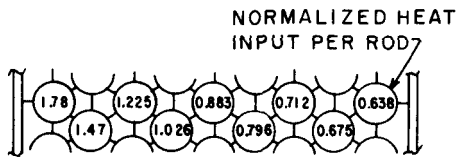
(a) UNIFORM



(b) DOUBLE HUMP



(c) 2/1 SKEW



(d) 2.8/1 SKEW

Fig. 11 Normalized Transverse Power Distribution

where  $\hat{Q}'''$  is the total average volumetric power generation in the assembly; and  $\bar{\psi}(y)$  is the radial shape factor (normalized).

Engel, et al., [14] present their temperature data by normalizing the measured local temperature rise at all elevations with respect to a calculated mixed mean sodium temperature rise at the exit of the heated section. The mixed mean sodium temperatures were calculated by mapping all measured subchannel temperatures at the exit of the heated section plane and filling-in missing subchannels (by interpolation) assuming smooth monotonic temperature profiles. The various subchannel (interior, corner or edge) temperature rises were then weighted to calculate the mixed mean outlet temperatures, using flow area and the blanket assembly flow split obtained from MIT studies as the weighting factors [15].

In the present study, the mean sodium temperature at the exit is calculated as,

$$\hat{T}(L_h) = \frac{1}{\hat{u}_\ell} \int_0^\ell T(L_h, y) u(L_h, y) dy \quad (99)$$

where the integration is carried out numerically by the Simpson's rule:

$$\begin{aligned} \hat{T}(L_h) \approx \frac{1}{3\hat{u}_s} [ & (uT)_0 + 4(uT)_1 + 2(uT)_2 + \dots \\ & + 2(uT)_{s-2} + 4(uT)_{s-1} + (uT)_s ] \end{aligned} \quad (100)$$

where  $s$  is an even number of equal subdivisions in the transverse direction.

It must be noted that in the present analysis the effect of heat losses from the test assembly has been neglected.

Parameters which can have significant effects on calculated temperature distributions include, inlet velocity distribution, bundle tolerance and especially the bundle bowing, pin distortion and the bundle housing thermo-mechanical behavior caused by heating effects.

In the present application the effect of bundle bowing and distortion are not taken into account. However, these effects can readily be included through adjustments in the hydraulic diameters and the porosity to match the experimental data.

In a typical fuel and blanket assembly, the flow area of the channels in the wall region is nearly twice as large as the flow area of the interior channels. As a result, the coolant temperature rise in the wall region is lower than in the interior channels. The wall channels act as a heat sink to the interior regions of the assembly. At low Reynold numbers (or at low Peclet number;  $Pe = Re Pr$ ), the flattening of the transverse temperature profile occurs as a result of the two distinct processes described earlier, namely; first, by energy transport due to increased conduction, combined with wire wrap induced mixing, and secondly, by buoyancy-induced flow redistribution. As the wall channels are cooler than the interior channels, there will be density differences between the wall and the interior regions of the bundle, even in the absence of radial power skew. The effect of higher density, and to a lesser extent the higher viscosity in the wall region compared to the interior region, results in higher frictional forces in the wall channels. The combined effect of frictional and static head forces results in a net diversion of flow from the wall to the interior regions of the bundle. This buoyancy-induced cross flow tends to further flatten the transverse temperature gradient. This cross flow effect can be further enhanced in the presence of radial power skew across the bundle [23].

The following analysis of simulation experiments demonstrates the importance of intra-subassembly processes as they influence the sodium temperatures during forced, mixed and free convection conditions.

#### A. Uniform Power Distribution

Figure (11a) depicts the normalized transverse power distribution in the assembly. It is seen that the distribution is nearly uniform (flat), thus the expected temperature distribution is inversely proportional to the velocity distribution shown in Figure 12 (high Reynolds number).

Figures 13 through 17 illustrate the calculated transverse coolant temperature distribution corresponding to the power and flow conditions of Table III for two different mixing parameters ( $\epsilon_{\infty}^* = 0.09$  and  $\epsilon_{\infty}^* = 0.03$ ).

Also shown are the experimental data of Engel, et al. [14]. It is seen that the agreement between the calculated and measured data improves in going from the forced to the free convection regime. Furthermore, the quantitative differences in the measured and calculated temperatures are believed to be in part due to the bundle distortions and inlet velocity distribution. Similar observations were made by Juneau and Khan [23] using the ENERGY and COBRA computer codes.

Figures 13 through 17 also demonstrate the influence of the mixing parameter  $\epsilon_{\infty}^*$  on the bundle temperature profile calculations. It is seen that a reduction in turbulent mixing reduces the transverse temperature flattening at high Reynolds numbers, leading to higher temperatures in the interior region, and subsequently lower sodium temperatures in the side (wall) region.

Figure 15 illustrates the temperature distribution downstream of the heated section corresponding to the exit of the subassembly. It is seen that significant flattening of the radial temperature distribution occurs as a result of turbulent and molecular diffusion, in close agreement with experimental measurements.

#### B. Double-Hump Power Distribution

The double-humped power distribution depicted in Fig. (11b) exemplifies the power distribution for LMFBR inner blanket subassemblies.

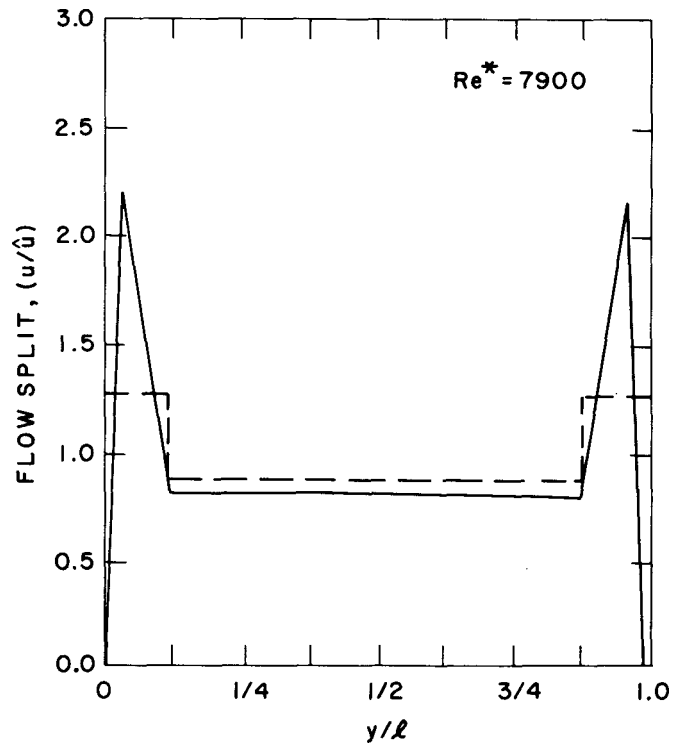


Fig. 12 Calculated Flow Split at the Exit of the Heated Section

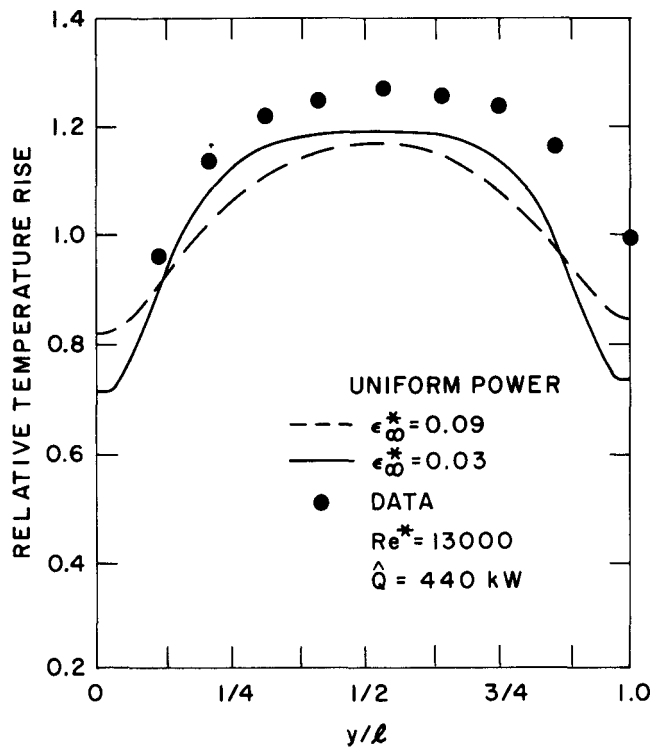


Fig. 13 Comparison of Calculated and Measured Transverse Temperature Distribution (Uniform)

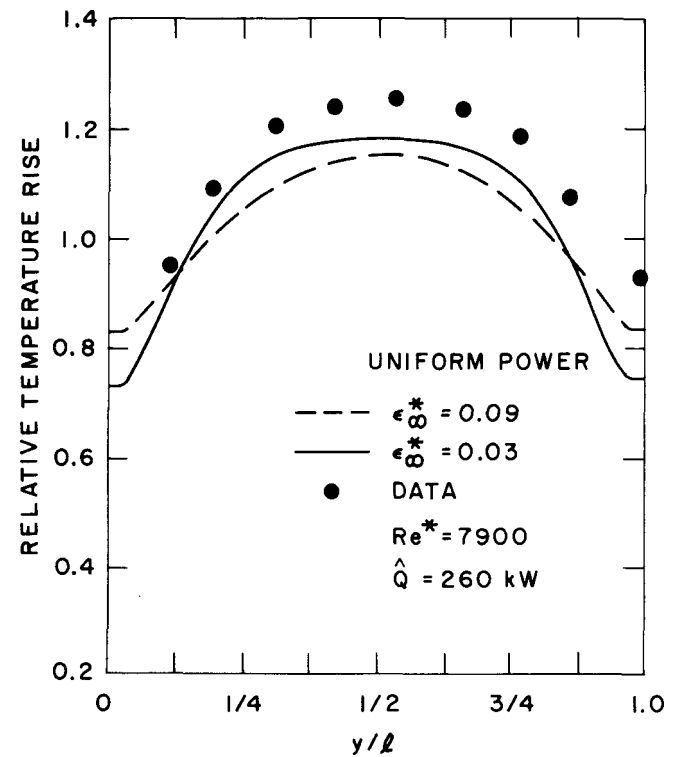


Fig. 14 Comparison of Calculated and Measured Transverse Temperature Distribution (Uniform)

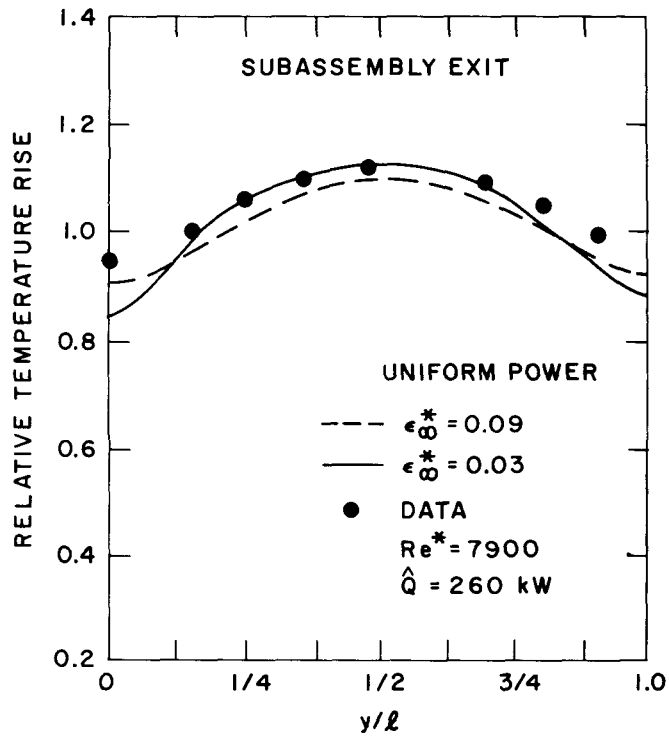


Fig. 15 Comparison of Calculated and Measured Transverse Temperature Distribution at the Subassembly Exit (Uniform)

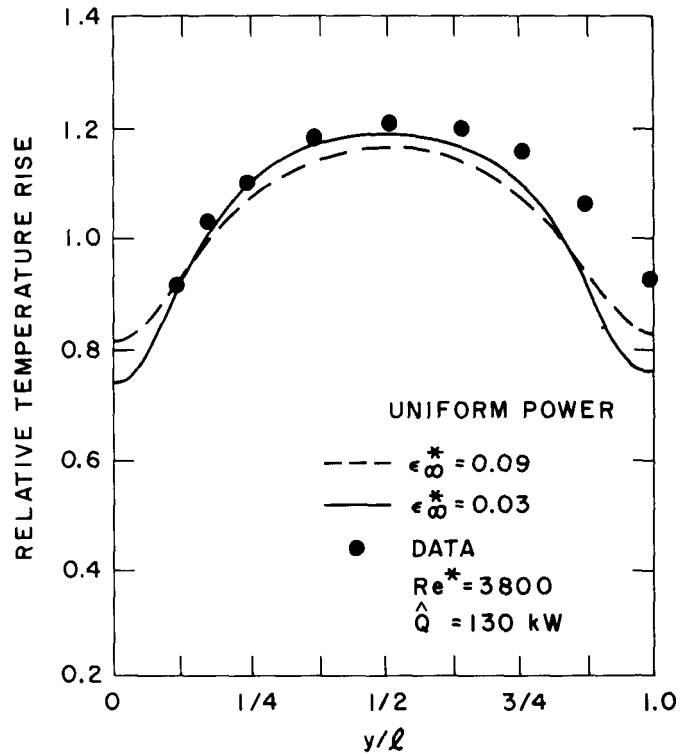


Fig. 16 Comparison of Calculated and Measured Transverse Temperature Distribution (Uniform)

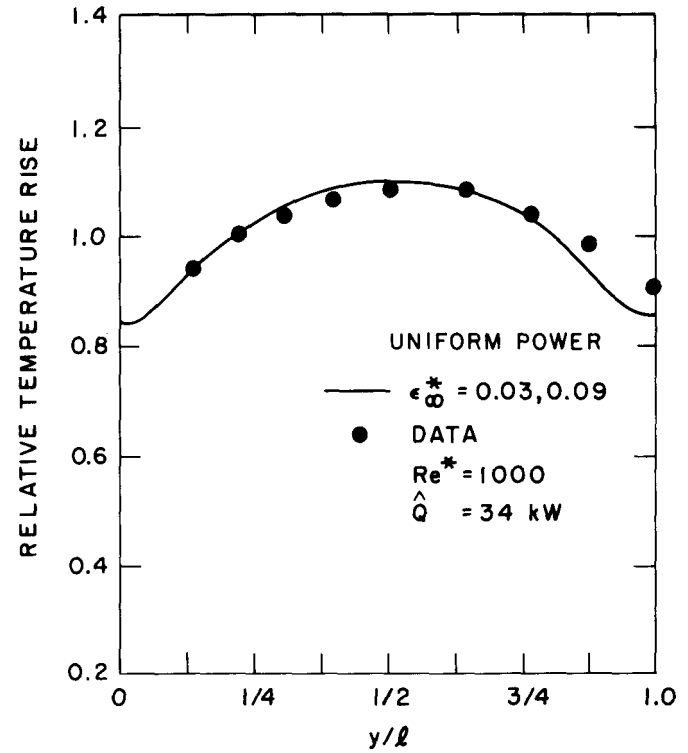


Fig. 17 Comparison of Calculated and Measured Transverse Temperature Distribution (Uniform)

Figures 18 through 21 show the comparison of the measured and calculated temperature distribution at high, intermediate and low Reynolds numbers corresponding to forced, mixed and free convection conditions. The agreement between the calculated and measured values improves at low flow natural convection conditions; a trend also observed for the uniform power distribution case discussed earlier.

### C. Steep Power Gradients

Figures (11c) and (11d) show the normalized transverse power gradients corresponding to 2:1 and 2.8:1 power skews, respectively, while Table III summarizes the corresponding test conditions.

Figures 22 through 24 show the calculated and measured transverse temperature distribution at 2:1 power skew across the assembly. It is seen that the calculated results are in agreement with the experimental measurements, especially in the higher power side of the assembly. Again, the agreement between the calculated and measured temperatures improve at lower Reynolds numbers.

The larger discrepancy in the low-power side of the bundle is believed to be influenced by the reduction in the flow area caused by pin distortion and bowing as a result of large transverse power (temperature) gradients. Similar observations have been made using COTEC, ENERGY-II and COBRA-IV computer codes [14,23,24].

Figures 25 through 29 show the calculated and measured transverse temperature distribution at 2.8:1 power skew across the assembly. The behaviors are quite similar to the 2:1 power skew case, however, the maximum temperatures are higher as a consequence of the larger power gradients.

Figures 25 and 29 also show the transverse temperature distribution corresponding to the midplane of the heated section. It is seen that the agreement between the calculated and measured temperature distributions is good. Furthermore, the better agreement with experimental measurements, as compared to the exit of the heated section, suggests that differential bowing seems to be more pronounced towards the top of the assembly.

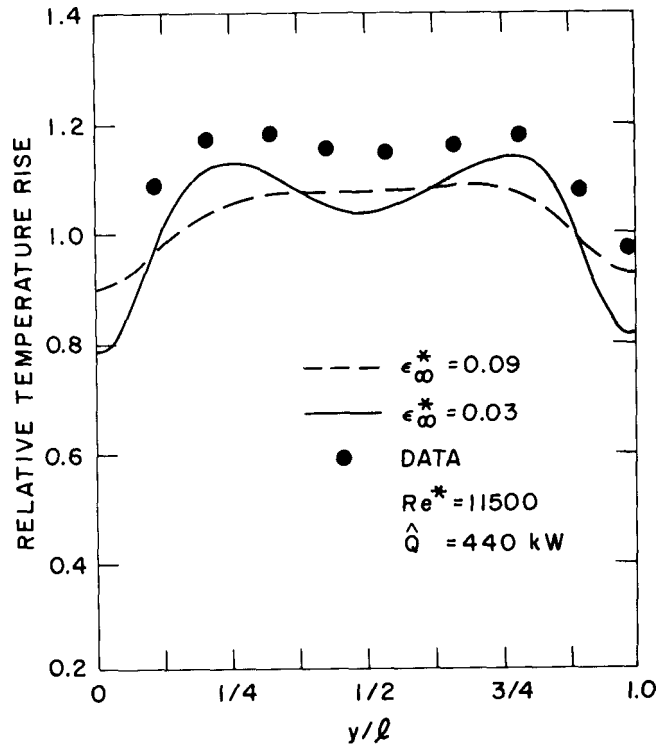


Fig. 18 Comparison of Calculated and Measured Transverse Temperature Distribution (Internal Blanket)

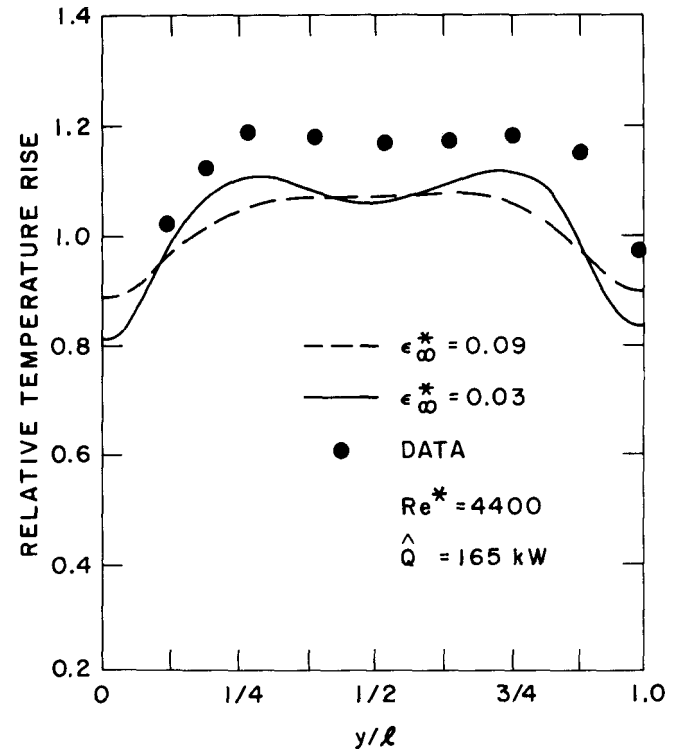


Fig. 19 Comparison of Calculated and Measured Transverse Temperature Distribution (Internal Blanket)

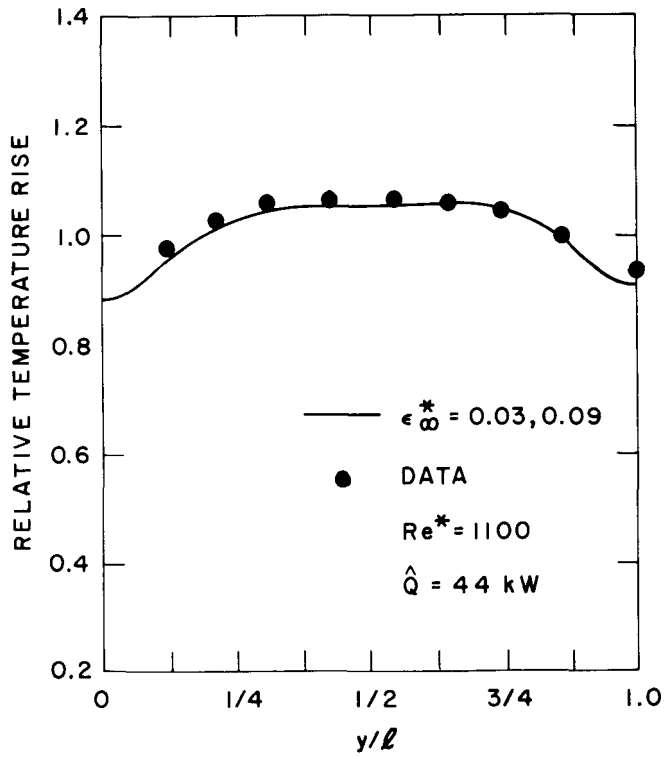


Fig. 20 Comparison of Calculated and Measured Transverse Temperature Distribution (Internal Blanket)

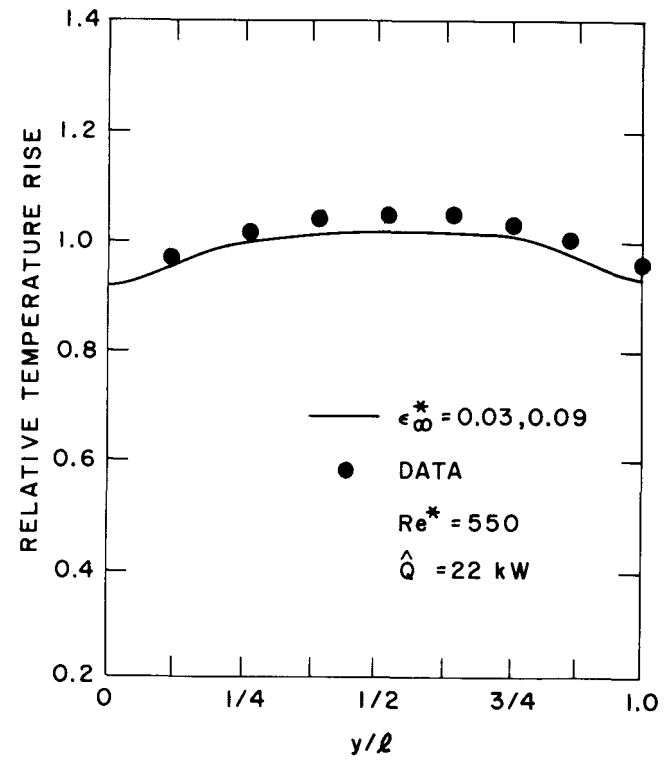


Fig. 21 Comparison of Calculated and Measured Transverse Temperature Distribution (Internal Blanket)

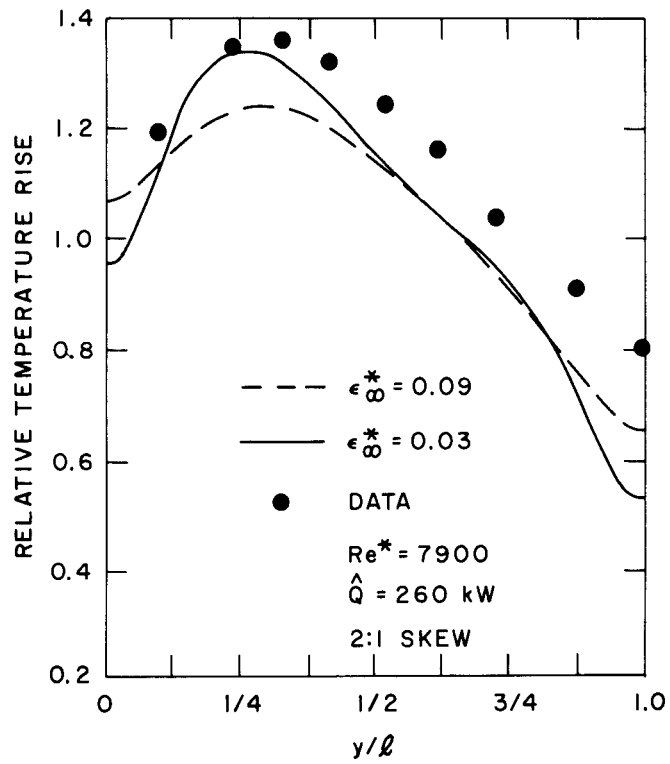


Fig. 22 Comparison of Calculated and Measured Transverse Temperature Distribution (2:1 Skew)

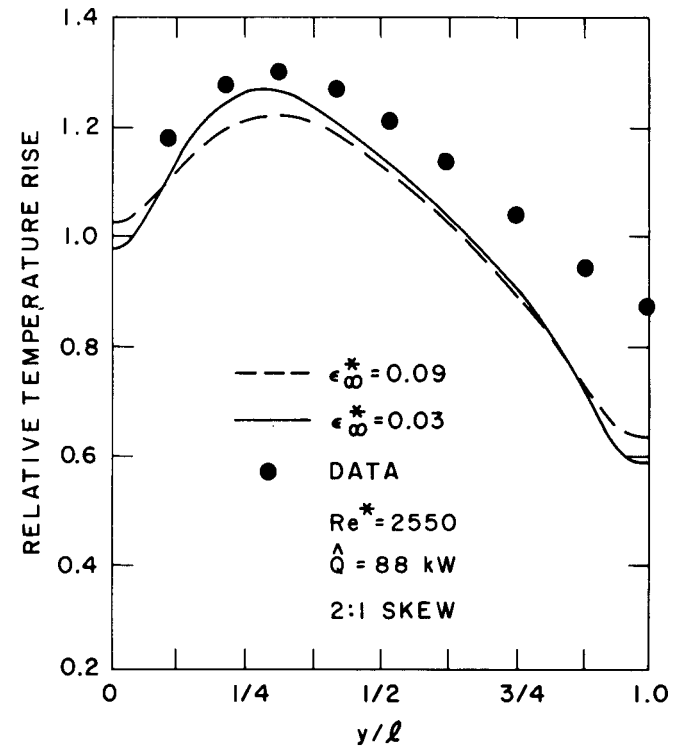


Fig. 23 Comparison of Calculated and Measured Transverse Temperature Distribution (2:1 Skew)

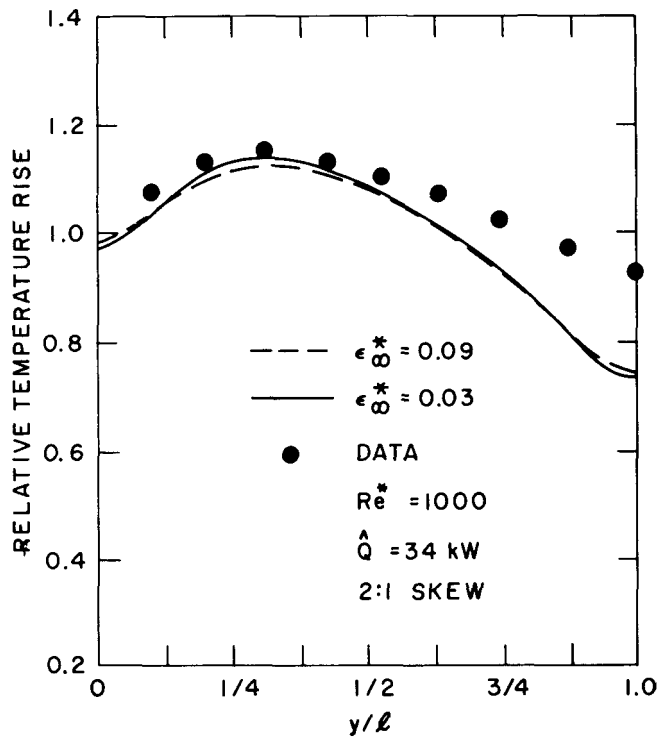


Fig. 24 Comparison of Calculated and Measured Transverse Temperature Distribution (2:1 Skew)

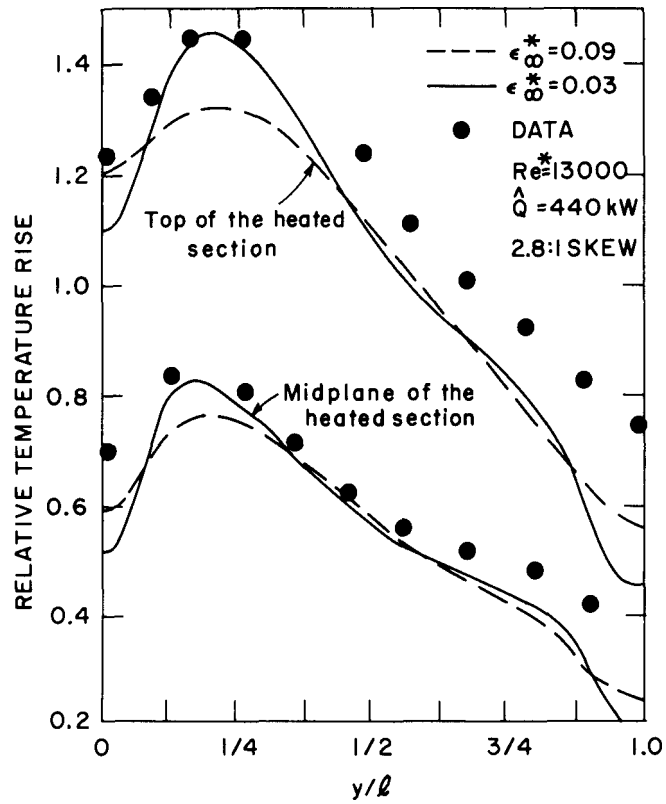


Fig. 25 Comparison of Calculated and Measured Transverse Temperature Distribution (2.8:1 Skew)

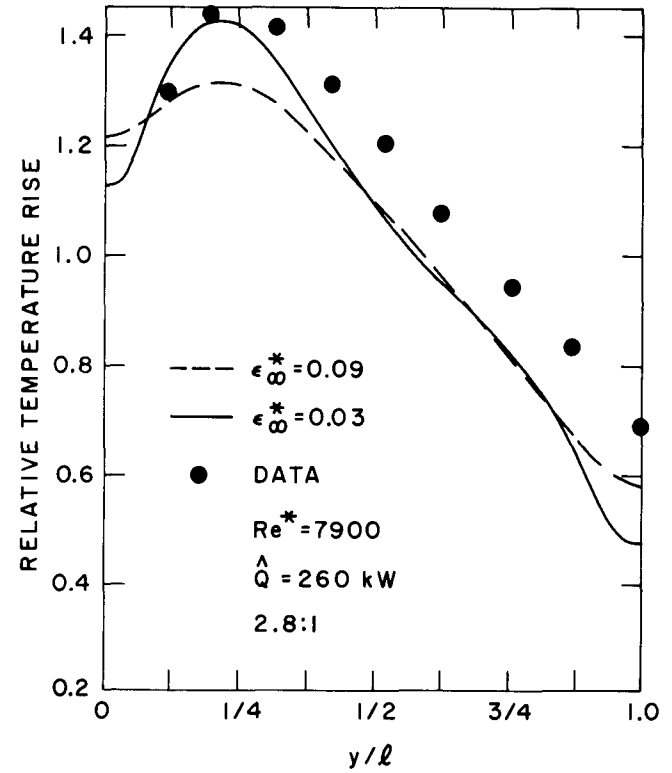


Fig. 26 Comparison of Calculated and Measured Transverse Temperature Distribution (2.8:1 Skew)

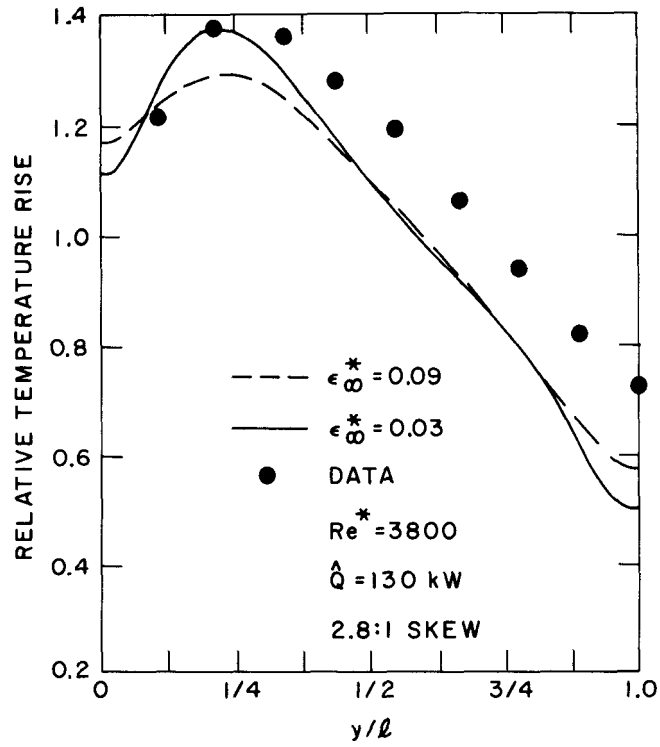


Fig. 27 Comparison of Calculated and Measured Transverse Temperature Distribution (2.8:1 Skew)

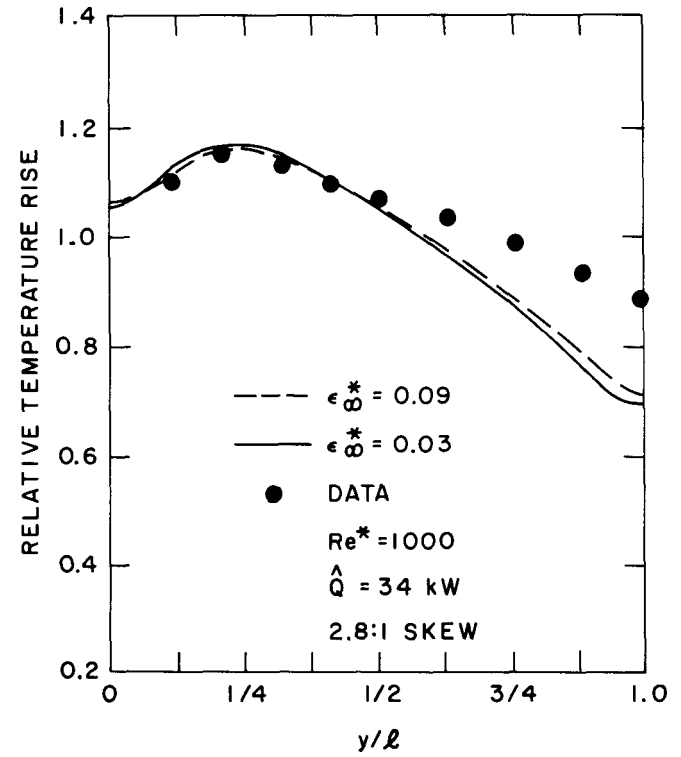


Fig. 28 Comparison of Calculated and Measured Transverse Temperature Distribution (2.8:1 Skew)

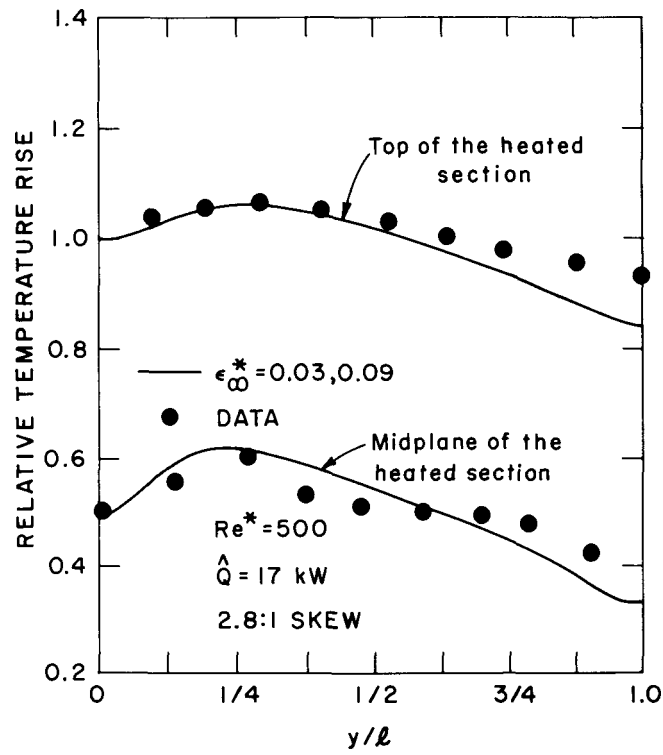


Fig. 29 Comparison of Calculated and Measured Transverse Temperature Distribution (2.8:1 Skew)

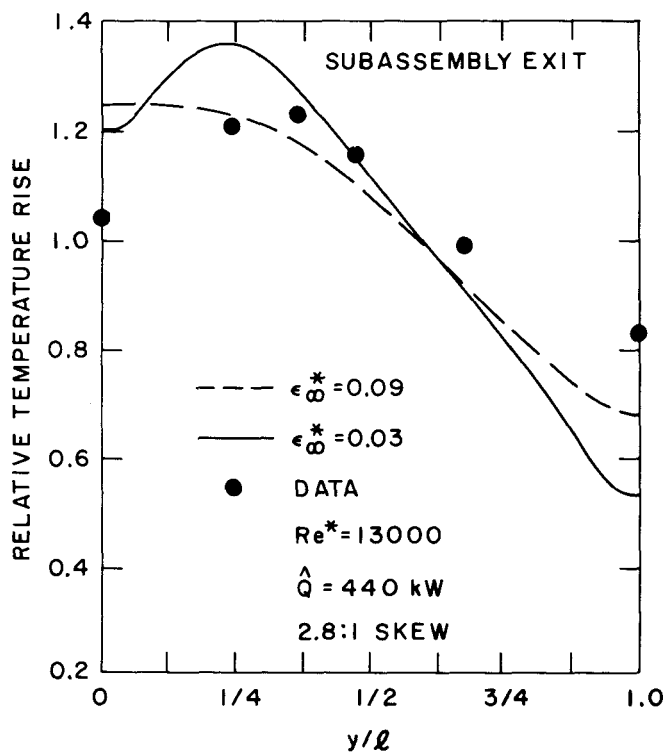


Fig. 30 Comparison of Calculated and Measured Transverse Temperature Distribution at the Subassembly Exit (2.8:1 Skew)

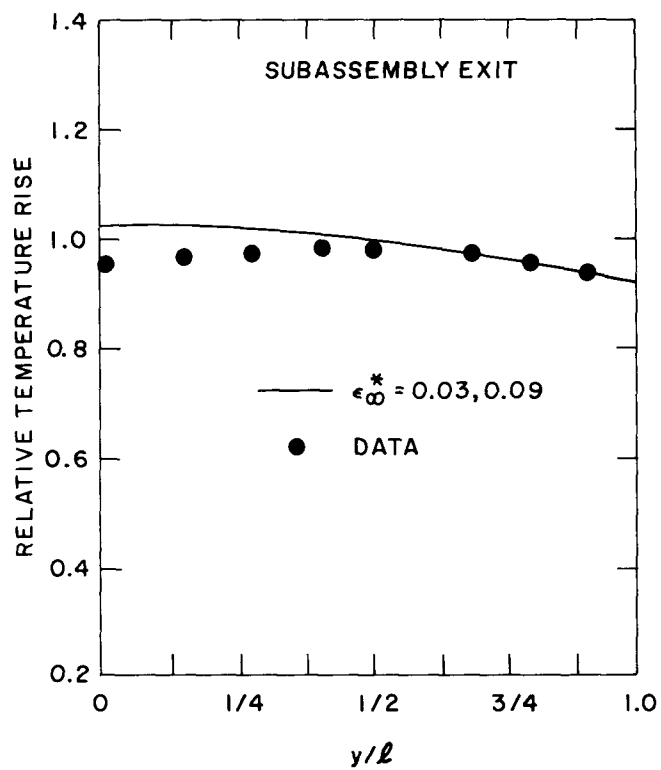


Fig. 31 Comparison of Calculated and Measured Transverse Temperature Distribution at the Subassembly Exit (2.8:1 Skew)

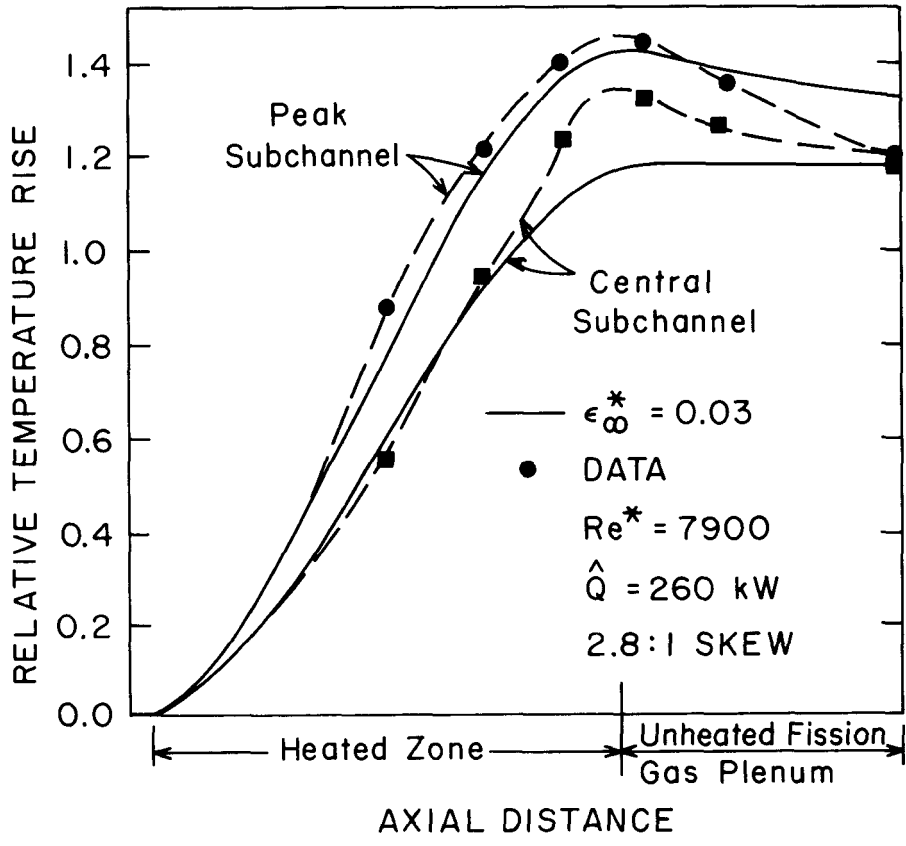


Fig. 32 Comparison of Calculated and Measured Axial Temperature Distribution for Peak and Central Subchannels (2.8:1 Skew)

Figures 30 and 31 illustrate the transverse temperature profiles at the exit of the subassembly (top of the unheated fission gas plenum region). At high Reynolds number (Fig. 30), there seems to be considerably more mixing between the side and interior channels as predicted by the present model, also in agreement with COTEC calculations [14]. At low Reynolds number (Fig. 31), the predicted profile is in excellent agreement to the measured data.

Figure 32 shows a longitudinal temperature distribution at Reynolds number of 7900 corresponding to 2.8:1 power skew. The agreement between predictions and measured data is good in the heated section with prediction of peak temperature becoming more conservative near the subassembly outlet. In the upper unheated region as mentioned earlier, in high Reynolds number range, mixing appears to be much stronger than predicted, leading to flatter measured transverse profile.

#### 4.3.2 Comparison With Water Test Data

The Toshiba Water simulation test assembly shown schematically in Fig. 33, consists of 91 helically-wound wire spacer rods arranged in a triangular pitch configuration. Each simulator rod contains an electrically heated element to simulate nuclear heating with a chopped cosine axial power distribution [25, 26]. Table IV summarizes the geometric characteristics of the test assembly. Figure 33 also shows a cross sectional view of the rod bundle and the thermocouple locations placed at vertical planes L2 and L3.

Experimentally, bundle operating parameters were adjusted to achieve the desired flow rate and inlet temperature, then the power was gradually increased to the selected level. Measurements were taken following attainment of the steady-state operating conditions.

An interesting feature of these water simulation experiments is that due to the small thermal conductivity of the coolant, the impact of intra-assembly heat transfer is quite insignificant as compared to the buoyancy-induced flow redistribution. Thus, the data provide an excellent data base for confirmation of buoyancy-induced flow redistribution in rod bundle geometries.

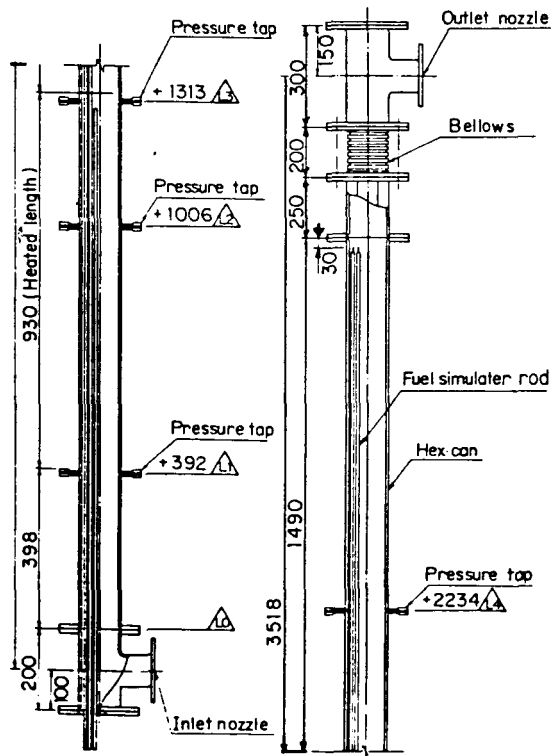
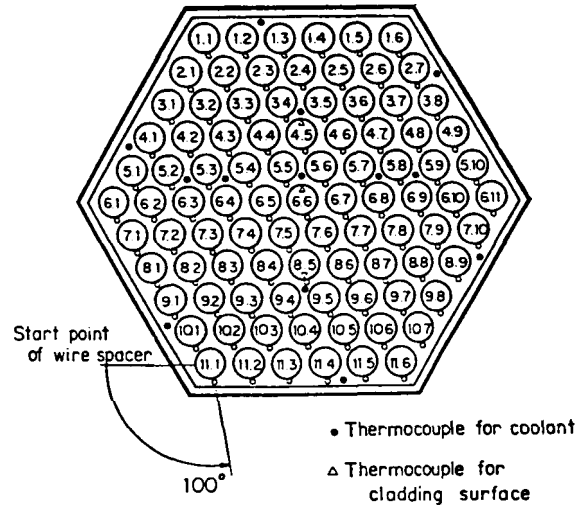


Fig. 33 Schematic of the Toshiba 91-Pin Test Model

Table IV Geometric Characteristics of the  
TOSHIBA 91 Pin Test Model

$D_r$	(m)	$6.5 \times 10^{-3}$
$D_w$	(m)	$1.3 \times 10^{-3}$
$p/D_r$	(-)	1.21
$\ell$	(m)	0.07730
$\ell_1$	(m)	0.06815
$\ell_2$	(m)	0.00457
$L_h$	(m)	0.930
$A_r$	(m <sup>2</sup> )	$3.32 \times 10^{-5}$
$A_w$	(m <sup>2</sup> )	$1.33 \times 10^{-6}$
$A_p$	(m <sup>2</sup> )	$5.17 \times 10^{-3}$
$D_{e1}$	(m)	$3.12 \times 10^{-3}$
$D_{e2}$	(m)	$6.12 \times 10^{-3}$
$D_{e3}$	(m)	$4.33 \times 10^{-3}$
$D_{ea}$	(m)	$4.01 \times 10^{-3}$
$\alpha_1$		0.7774
$\alpha_2$		0.2086
$\alpha_3$		0.0140
$\lambda_1$		0.3566
$\lambda_2$		0.5204
$\lambda_3$		0.5234
$\lambda_a$		0.3931

Recall in Eq. (3),  $\kappa$ , the molecular conduction correction factor was introduced to account for the reduction in effective conductivity of sodium due to the presence of the fuel rods. However, in the present experiment, the thermal conductivity of fuel rod simulators is much greater than the coolant (water). Thus, at low flow rates, heat tends to conduct much easier through the steel cladding and the wire-wrap than water (coolant), thus  $\kappa$  may be greater than unity.

The mixing parameter  $\epsilon_{\infty}^*$  is not very important for the Reynolds number regimes of these experiments and thus, the existing model is assumed to be valid.

In order to simulate the hydro-dynamic characteristics of the Toshiba assembly, the following friction factor correlations are used:

$$f = 37.4 \text{ Re}^{-0.872} \quad \text{Re} \leq 1200 \quad (101a)$$

$$f = 41.7/\text{Re} + 0.12 \text{ Re}^{-0.159} \quad \text{Re} > 1200 \quad (101b)$$

These relations compare favorably with the unheated test data as recommended by Mawatari, et al. [25].

#### A. Uniform Power Distribution

Figure 34 shows the comparison between measured [25,26] and calculated transverse temperature distribution corresponding to uniform transverse power distribution.

It is evident that good agreement exists between the measured and predicted results. At high Reynolds number the calculated relative temperature rise is within 9% of the experimentally measured values (less than 4°K in actual temperature).

At low Reynolds number the maximum discrepancy is seen to occur near the edges where the calculated relative temperature rise is within about 15% of the measured data (< 6°K in actual temperature).

Experiment	●	○
Calculation	—	---
Test No.	HT303-1	HT303-2
Re*	336	3330
$\hat{\Delta T}$ (K)	38.6	36.1
T* (K)	314	318

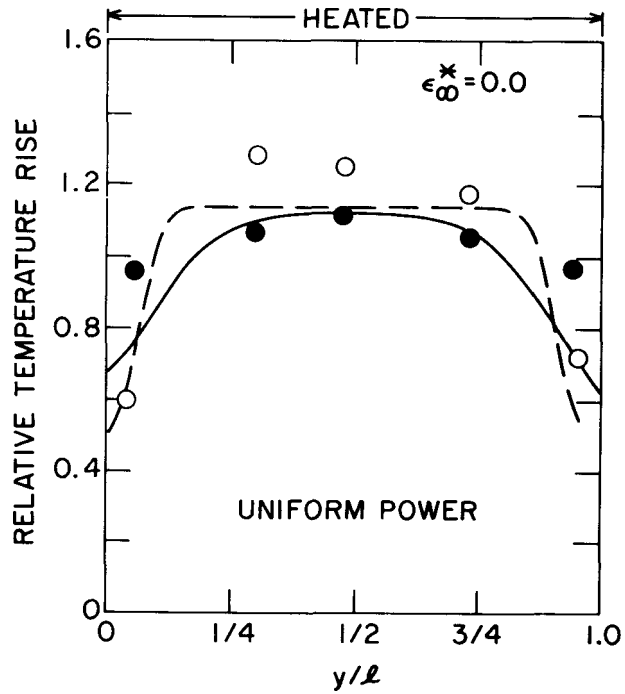


Fig. 34 Comparison of Calculated and Measured Transverse Temperature Distribution for Toshiba Experiments (Uniform)

## B. Non-Uniform Power Distribution

Figure 35 shows the comparison between measured and calculated transverse temperature distribution corresponding to a non-uniform transverse power distribution.

In this case, one half of the bundle is heated uniformly (51 rods) while the other half (remaining 40 rods) are unheated. This gives rise to a large temperature gradient between the heated and unheated halves of the test section.

It is clearly evident that at high Reynolds number the agreement between the calculated and measured data is excellent. However, at low Reynolds number natural convection conditions, experimental measurements indicate significant energy redistribution to the unheated side of the bundle, while the present calculations show a much smaller flattening effect.

This may be attributed to three effects, namely; (1) continued swirl flow mixing at low flows, (2) enhanced molecular conduction through the cladding and the wire-wraps, and (3) bowing and bundle distortion caused by large transverse power gradient across the assembly.

Experiment	●	○
Calculation	—	- - -
Test No.	HT40I-I	HT40I-O
Re*	334	2920
$\Delta T$ (K)	40.2	38.3
T* (K)	314	319

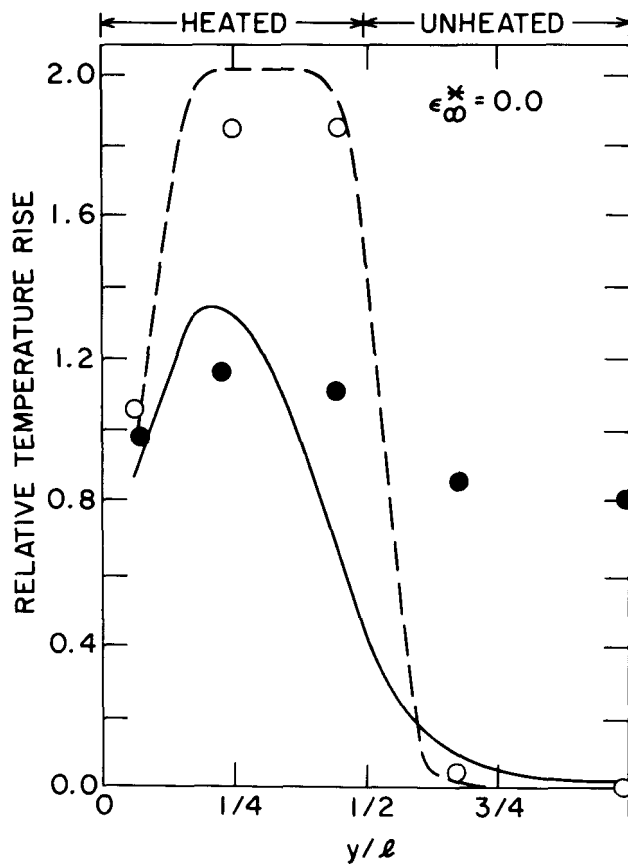


Fig. 35 Comparison of Calculated and Measured Transverse Temperature Distribution for Toshiba Experiments (Partially Uniform)

## 5. TRANSIENT POROUS-BODY MODEL

The model described in the previous chapter is applicable to steady-steady or quasi-steady processes. However, for the study of protected loss-of-flow events in LMFBR subassemblies, where the flow regime changes quickly from forced to mixed and eventually free convection, a time dependent model is desirable.

In this chapter, the previously developed (see 4.1) parabolic flow model is extended to include transient processes. The model and the resulting computer code is then used to post-predict the 61-pin rod bundle transient experiments conducted at Westinghouse [27].

### 5.1 Governing Equations

The basic conservation equations of section 2.1 can be easily extended to include the time-dependent terms assuming that the Boussinesq approximation for incompressible fluids holds (density is a constant of motion, i.e.,  $\frac{\partial \rho}{\partial t} = 0$ ), thus:

a) Conservation of Mass:

$$\frac{\partial(\rho u)}{\partial x} + \frac{\partial(\rho v)}{\partial y} = 0 \quad (102)$$

b) Conservation of Energy:

$$\frac{\partial T}{\partial t} + u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \alpha_e \frac{\partial^2 T}{\partial y^2} + \frac{Q'''(x,y)}{\rho c_p} \quad (103)$$

c) Conservation of Axial Momentum:

$$\begin{aligned} \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} &= -\frac{1}{\rho} \frac{\partial P}{\partial x} - g \frac{\rho^*}{\rho} [1 - \beta(T-T^*)] \\ &+ v \frac{\partial^2 u}{\partial y^2} + \frac{f}{2D_e \lambda^2} u^2 \end{aligned} \quad (104)$$

Subject to the following initial and boundary conditions:

Initial Conditions

$$\left. \begin{aligned} u(x,y,0) &= \text{known} \\ v(x,y,0) &= \text{known} \\ T(x,y,0) &= \text{known} \end{aligned} \right\} \text{Through Steady-State Solution,} \quad (105)$$

Boundary Conditions

$$\begin{aligned} u(0,y,t) &= \text{specified} \\ T(0,y,t) &= \text{specified} \\ Q''(x,y,t) &= \text{specified} \end{aligned} \quad (106)$$

Furthermore, the heat flux at the duct walls is assumed to be continuous.

The energy equations for the stagnant sodium region adjacent to the duct wall and the duct wall structure can be written as:

Stagnant Sodium

$$M_s C_s \frac{dT_s}{dt} = \frac{k_d}{\delta_d} A_d (T_d - T_s) \quad \begin{matrix} s=1,N \\ d=2,N-1 \end{matrix} \quad (107)$$

Duct Wall

$$M_d C_d \frac{dT_d}{dt} = \frac{k_d}{\delta_d} A_d (T_s - T_d) + \frac{k_s}{\Delta y} A_d (T_{d'} - T_d) \quad (108)$$

$$d' = 3 \quad \text{for } d = 2$$

$$d' = N-2 \quad \text{for } d = N-1$$

where M is the mass, C is the specific heat, k is the thermal conductivity,  $\delta$  is the wall thickness, A is the heat transfer area, and the subscripts d and s refer to the duct wall and the sodium, respectively.

The forms of equations (107) and (108) indicate that the duct wall outer surface is at equilibrium with the stagnant sodium layer (usually present between subassemblies).

Initially ( $t = 0$ ) the adiabatic boundary condition is identically satisfied, that is:

$$\begin{cases} T_1 = T_2 = T_3 \\ T_N = T_{N-1} = T_{N-2} \end{cases} \quad (109)$$

## 5.2 Method of Solution

### 5.2.1 Finite Difference Equations

In finite difference form, Equations (102) through (108) can be expressed as follows:

#### Mass Continuity

$$\left[ \frac{(\rho u)_{i+1,j} + (\rho u)_{i+1,j+1} - (\rho u)_{i,j} - (\rho u)_{i,j+1}}{2\Delta x} \right]^{k+1} + \left[ \frac{(\rho v)_{i+1,j+1} - (\rho v)_{i+1,j}}{\Delta y} \right]^{k+1} = 0 \quad (110)$$

#### Energy

$$\left[ \frac{T_{i+1,j}^{k+1} - T_{i+1,j}^k}{\Delta t} \right] + u_{i,j}^{k+1} \left[ \frac{T_{i+1,j}^{k+1} - T_{i,j}^{k+1}}{\Delta x} \right] + v_{i,j}^{k+1} \left[ \frac{T_{i+1,j+1}^{k+1} - T_{i+1,j-1}^{k+1}}{2\Delta y} \right] = \alpha_e \left[ \frac{T_{i+1,j+1}^{k+1} - 2T_{i+1,j}^{k+1} + T_{i+1,j-1}^{k+1}}{(\Delta y)^2} \right] + \frac{Q_{i,j}^{k+1}}{\rho c_p} \quad (111)$$

Axial Momentum

$$\begin{aligned}
 & \left[ \frac{u_{i+1,j}^{k+1} - u_{i+1,j}^k}{\Delta t} \right] + u_{i,j}^{k+1} \left[ \frac{u_{i+1,j}^{k+1} - u_{i,j}^{k+1}}{\Delta x} \right] + v_{i,j}^{k+1} \left[ \frac{u_{i+1,j+1}^{k+1} - u_{i+1,j-1}^{k+1}}{2\Delta y} \right] \\
 & = - \left[ \frac{\Delta \tilde{p}}{\Delta x} \right]_{i+1}^{k+1} - \frac{f_{i,j}^{k+1}}{2D\lambda_{i,j}^2} u_{i,j}^{k+1} u_{i+1,j}^{k+1} + g \beta (T_{i+1}^{k+1} - T^*{}^{k+1}) + \\
 & v_{i,j} \left[ \frac{u_{i+1,j+1}^{k+1} - 2u_{i+1,j}^{k+1} + u_{i+1,j-1}^{k+1}}{(\Delta y)^2} \right] \tag{112}
 \end{aligned}$$

Stagnant Sodium

$$\left[ \frac{M_s C_s}{\Delta t} + \frac{k_d}{\delta_d} A_d \right] T_{i+1,s}^{k+1} - \frac{k_d}{\delta_d} A_d T_{i+1,d}^{k+1} = \frac{M_s C_s}{\Delta t} T_{i+1,s}^k \tag{113}$$

Duct Wall

$$\begin{aligned}
 & \left[ \frac{M_d C_d}{\Delta t} + \frac{k_d}{\delta_d} A_d + \frac{k_s}{\Delta y} A_d \right] T_{i+1,d}^{k+1} - \frac{k_d}{\delta_d} A_d T_{i+1,s}^{k+1} - \frac{k_s}{\Delta y} A_d T_{i+1,d}^{k+1} \\
 & = \frac{M_d C_d}{\Delta t} T_{i+1,d}^k \tag{114}
 \end{aligned}$$

where  $\Delta t$  is the time-step, and the superscripts  $k$  and  $k+1$  denote the previous and current time values, respectively.

The time-differencing forms of Eqs. (110) through (114) are based on the fully implicit formulation, which are unconditionally stable [28] for any value of  $\Delta t$ . However, the time-step size is governed by the desired degree of accuracy which is discussed in the following subsection.

Equations (110) through (114) can be recasted into the familiar matrix forms described in section 4.2.1 to obtain:

$$\begin{bmatrix}
 \phi_{i,1} & \sigma_{i,2} & 0 & 0 & 0 & \dots & 0 \\
 \zeta_{i,1} & \phi_{i,2} & \sigma_{i,3} & 0 & 0 & \dots & 0 \\
 0 & \zeta_{i,2} & \phi_{i,3} & \sigma_{i,4} & 0 & \dots & 0 \\
 \vdots & \vdots & \vdots & \vdots & \vdots & \ddots & \vdots \\
 \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & 0 \\
 \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \sigma_{i,N} \\
 \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\
 0 & \dots & \dots & \dots & \dots & \dots & \zeta_{i,N-1} & \phi_{i,N}
 \end{bmatrix}
 \begin{bmatrix}
 T_{i+1,1}^{k+1} \\
 T_{i+1,2}^{k+1} \\
 T_{i+1,3}^{k+1} \\
 \vdots \\
 \vdots \\
 T_{i+1,N-1}^{k+1} \\
 T_{i+1,N}^{k+1}
 \end{bmatrix}
 =
 \begin{bmatrix}
 C_{i,1} \\
 C_{i,2} \\
 C_{i,3} \\
 \vdots \\
 \vdots \\
 C_{i,N-1} \\
 C_{i,N}
 \end{bmatrix}
 \quad (115)$$

where:

$$\phi_{i,1} = \phi_{i,N} = \left[ \frac{M_s C_s}{\Delta t} + \frac{k_d}{\delta_d} A_d \right] \quad (116a)$$

$$\phi_{i,2} = \phi_{i,N-1} = \left[ \frac{M_d C_d}{\Delta t} + \frac{k_d}{\delta_d} A_d + \frac{k_s}{\Delta y} A_d \right] \quad (116b)$$

$$\phi_{i,j} = \left[ \frac{1}{\Delta t} + \frac{u_{i,j}^{k+1}}{\Delta x} + \frac{2\alpha_e}{(\Delta y)^2} \right] \quad (116c)$$

$$\sigma_{i,2} = \sigma_{i,N} = - \left[ \frac{k_d}{\delta_d} A_d \right] \quad (116d)$$

$$\sigma_{i,3} = - \left[ \frac{k_s}{\Delta y} A_d \right] \quad (116e)$$

$$\sigma_{i,j} = \left[ \frac{v_{i,j-1}^{k+1}}{2\Delta y} - \frac{\alpha_e}{(\Delta y)^2} \right] \quad (116f)$$

$$\zeta_{i,1} = \zeta_{i,N-1} = - \left[ \frac{k_d}{\delta_d} A_d \right] \quad (116g)$$

$$\zeta_{i,N-2} = - \left[ \frac{k_s}{\Delta y} A_d \right] \quad (116h)$$

$$\zeta_{i,j} = \left[ \frac{v_{i,j+1}^{k+1}}{2\Delta y} + \frac{\alpha_e}{(\Delta y)^2} \right] \quad (116i)$$

$$C_{i,1} = \left[ \frac{M_s C_s}{\Delta t} \right] T_{i,1}^k \quad (116j)$$

$$C_{i,2} = \left[ \frac{M_d C_d}{\Delta t} \right] T_{i,2}^k \quad (116k)$$

$$C_{i,j} = \left[ \frac{Q_{ij}^{k+1}}{\rho c_p} + \frac{T_{i,j}^k u_{i,j}^{k+1}}{\Delta x} + \frac{T_{i,j}^k}{\Delta t} \right] \quad j = 3, 4, \dots, N-2 \quad (116l)$$

$$C_{i,N-1} = \left[ \frac{M_d C_d}{\Delta t} \right] T_{i,N-1}^k \quad (116m)$$

$$C_{i,N} = \left[ \frac{M_s C_s}{\Delta t} \right] T_{i,N}^k \quad (116n)$$

Similarly, Eqs. (89) and (112) can be written as:

$$\begin{bmatrix} \rho_{i+1,3} & \rho_{i+1,4} & \rho_{i+1,5} & \rho_{i+1,6} & \rho_{i+1,7} & \dots & \rho_{i+1,N-2} & 0 \\ \gamma_{i+1,3} & \Omega_{i+1,4} & 0 & 0 & 0 & \dots & 0 & \frac{1}{\Delta x} \\ \Gamma_{i+1,3} & \gamma_{i+1,4} & \Omega_{i+1,5} & 0 & 0 & \dots & \vdots & \vdots \\ 0 & \Gamma_{i+1,4} & \gamma_{i+1,5} & \Omega_{i+1,6} & 0 & \dots & \vdots & \vdots \\ 0 & 0 & \gamma_{i+1,6} & \Omega_{i+1,7} & 0 & \dots & \vdots & \vdots \\ \vdots & \vdots & \vdots & \vdots & \vdots & \dots & \vdots & \vdots \\ 0 & 0 & \dots & 0 & \Gamma_{i+1,N-4} & \gamma_{i+1,N-3} & \Omega_{i+1,N-2} & \frac{1}{\Delta x} \end{bmatrix} \begin{bmatrix} u_{i+1,3}^{k+1} \\ u_{i+1,4}^{k+1} \\ \vdots \\ \vdots \\ \vdots \\ \vdots \\ u_{i+1,N-2}^{k+1} \\ \tilde{\Delta P}_{i+1}^{k+1} \end{bmatrix}$$

$$= \begin{bmatrix} S_{i+1} \\ C_{i+1,3} \\ C_{i+1,4} \\ \vdots \\ \vdots \\ \vdots \\ C_{i+1,N-2} \end{bmatrix} \quad (117a)$$

and

$$u_{i+1,1}^{k+1} = u_{i+1,2}^{k+1} = u_{i+1,N-1}^{k+1} = u_{i+1,N}^{k+1} = 0 \quad (117b)$$

where,

$$\gamma_{i+1,j} = \left[ \frac{1}{\Delta t} + \frac{u_{i,j}^{k+1}}{\Delta x} + \frac{f_{i,j}^{k+1}}{2D_{i,j}} \frac{u_{i,j}^{k+1}}{\lambda_{i,j}^2} + \frac{2v_{i,j}}{(\Delta y)^2} \right] \quad (118a)$$

$$\Omega_{i+1,j} = \left[ \frac{v_{i,j}^{k+1}}{2\Delta y} - \frac{v_{i,j}}{(\Delta y)^2} \right] \quad (118b)$$

$$\Gamma_{i+1,j} = - \left[ \frac{v_{i,j}^{k+1}}{2\Delta y} + \frac{v_{i,j}}{(\Delta y)^2} \right] \quad (118c)$$

$$S_{i+1} = \sum_{j=3}^{N-2} \left[ \rho_{i,j} u_{i,j} \right]^{k+1} \quad (118d)$$

$$C_{i+1,j} = \frac{1}{\Delta t} u_{i+1,j}^k + \frac{u_{i,j}^{k+1}}{\Delta x} + g\beta \left[ T_{i+1,j}^{k+1} - T^{*k+1} \right] \quad (118e)$$

### 5.2.2 Numerical Solution Procedure

The numerical solution procedure is somewhat similar to the steady-state calculations described in Section 4.2.2. The steps are to:

- 1) Carry out the steady-state initialization ( $t = 0$ ) and set the time step.
- 2) Calculate the forcing functions (i.e., time-dependent  $Q'''$ ,  $u_{in}$ ,  $T^*$ ).
- 3) Start at  $i=1$  row and solve Eq. (115) for  $T_{2,1}^{k+1}$ ,  $T_{2,2}^{k+1}$ , .....,  $T_{2,N}^{k+1}$  using the Gaussian elimination technique [22].
- 4) Calculate the physical properties as a function of  $T_{i+1,j}$ 's.
- 5) Solve Eq. (117) for  $u_{2,1}^{k+1}$ ,  $u_{2,2}^{k+1}$ , .....,  $u_{2,N}^{k+1}$  and  $\Delta \tilde{P}_2^{k+1}$  using the Gauss-Jordan reduction method [22].
- 6) Back substitute into Eq. (110) to obtain  $v_{2,1}^{k+1}$ ,  $v_{2,2}^{k+1}$ , .....,  $v_{2,N}^{k+1}$ .
- 7) March in the stream-wise direction up to  $i=M$  row; repeating steps (3) through (6).

8) Adjust the time step using:

$$\text{if } \left\{ \begin{array}{l} \left| \frac{T_{i,j}^{k+1} - T_{i,j}^k}{T_{i,j}^k} \right| \leq \epsilon ; \Delta t^{k+1} = 2\Delta t^k \\ \text{otherwise} ; \Delta t^{k+1} = \Delta t^k / 2 \end{array} \right. \quad (119)$$

However the time step must be limited to

$$\Delta t_{\min} \leq \Delta t \leq \Delta t_{\max} \quad (120)$$

9) Advance the time and repeat steps (2) through (8) above.

The present steady-state and transient model has been assembled into a computer code, named TWIST.

The user's manual for the TWIST code along with the source listing is given in Appendix B. It must be noted, that code modifications to include other boundary conditions can be readily made; therefore the program structure is sufficiently generalized to perform a variety of steady-state and transient calculations.

### 5.3 Results and Comparison with Experimental Data

The transient porous-body model described in the previous sections is used to simulate the heat transfer tests performed at the Westinghouse blanket test model described earlier in sections 3.1.2.2 and 4.3.1.

The experiments were performed by first establishing steady state conditions (power input, flow, temperature) and then decreasing the power input and sodium flow at pre-programmed rates [27]. The test series included tests with heat input to all rods, and also to four rows (26 rods) only. The transients simulated both undercooling and overcooling events.

## A. Undercooling Event

Engel, et al., [27] report results for four transient undercooling events corresponding to various heat-input gradients.

For the purpose of the current simulation, Test Run 613 is chosen. This test corresponds to 1:0 maximum-to-minimum heat input, where only 26 of the 61 rods were uniformly heated giving rise to large temperature gradients between the heated and unheated halves of the test section.

Initially, the bundle was operating at 188 kw power, 56 gpm flow and 588K sodium inlet temperature condition. The bundle power and sodium flow rates were changed in a pre-programmed manner to achieve the rates displayed in Figs. 36 and 37.

In steady-state operations, the power-to-flow ratio is a measure of the mean sodium temperature rise through the test section (see Fig. 38). During the transient, the temperature rise is also affected by the stored heat released from the heater rods and/or the duct walls. It must be noted that the present model assumes that the heat input is directly deposited into the coolant, and thus neglects the heat storage effects which can influence the early stages of the transient. However, modification of the model to include fuel pin to coolant heat transfer is straightforward.

Figures 39 through 41 illustrate the effect of transient operation on cross-assembly temperature profiles. Also shown are the corresponding measured values.

It is seen that a significant profile flattening takes place, similar to the steady state cases studied earlier in Chapter 4. Furthermore, maximum discrepancies between the calculated and measured temperatures occur in the unheated side of the bundle, similar to the earlier steady-state observations. However, the peak sodium temperatures are in excellent agreement with the test data.

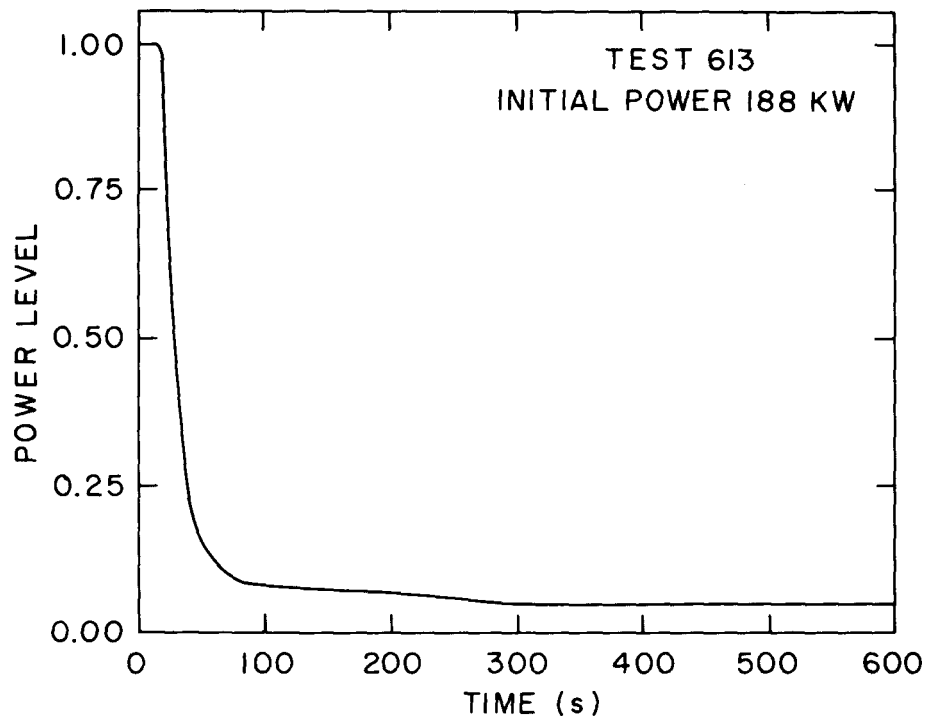


Fig. 36 Input Power Level (Test 613)

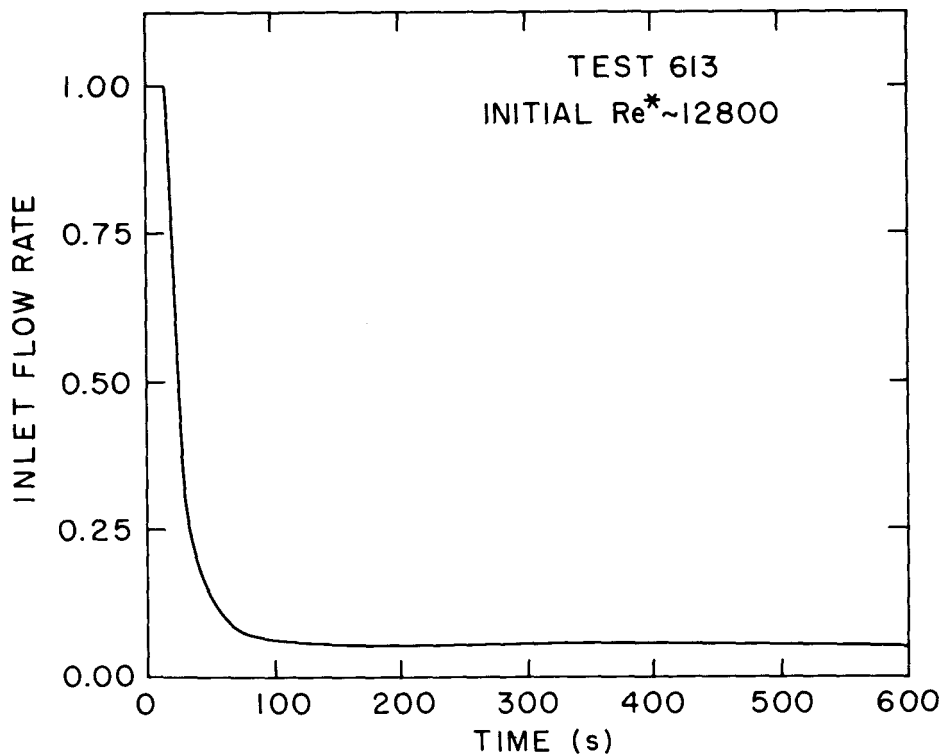


Fig. 37 Inlet Flow Rate (Test 613)

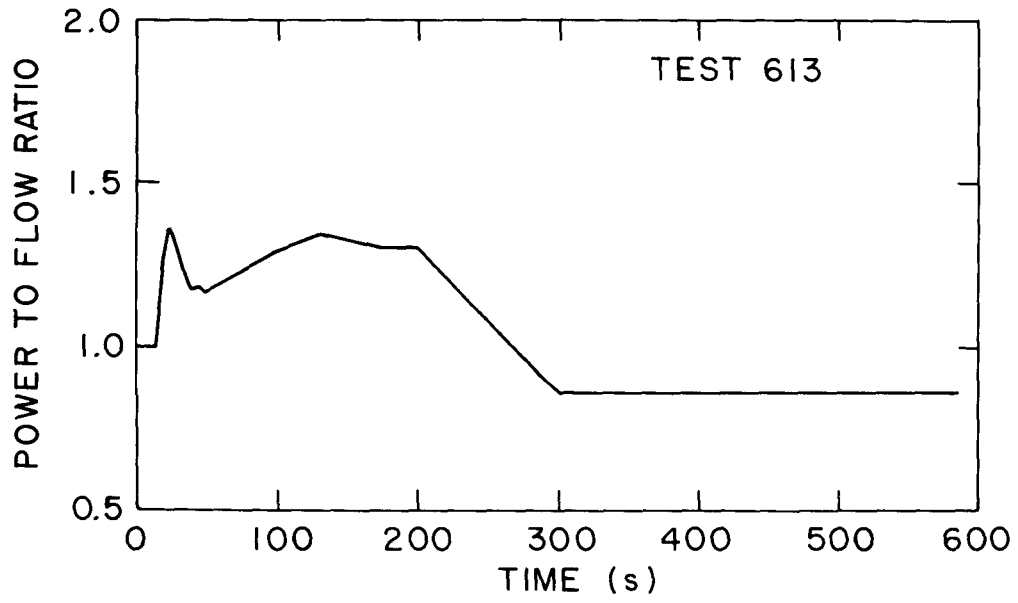


Fig. 38 Power to Flow Ratio (Test 613)

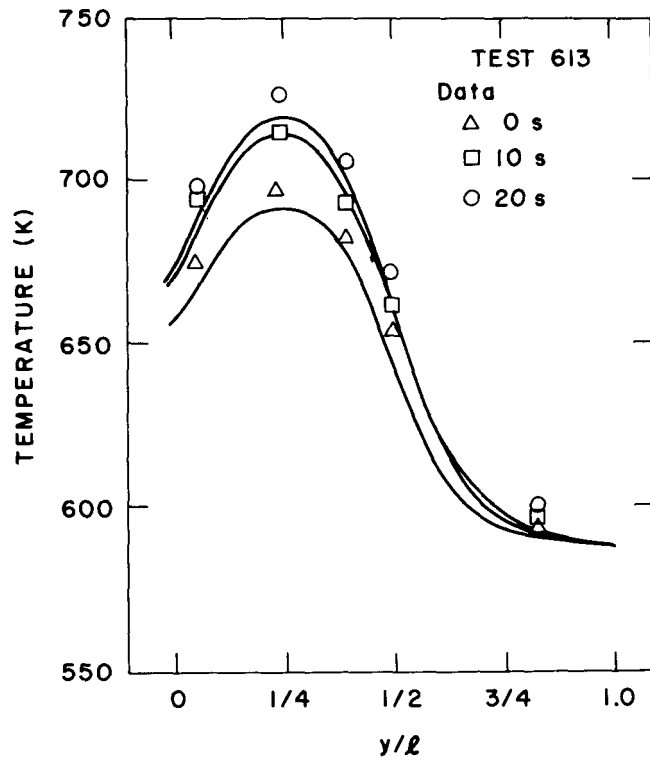


Fig. 39 Comparison of Calculated and Measured Transient Temperature Distributions for Test 613 (Early)

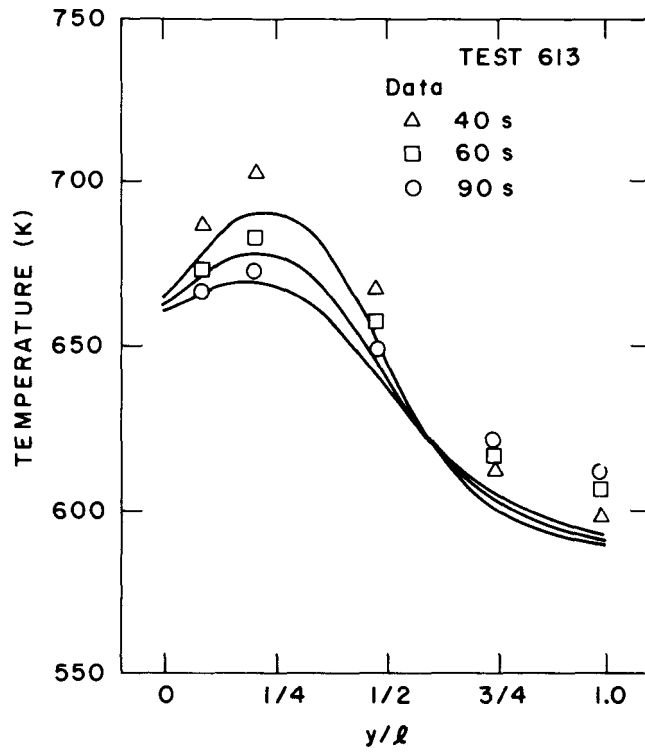


Fig. 40 Comparison of Calculated and Measured Transient Temperature Distribution for Test 613 (Intermediate)

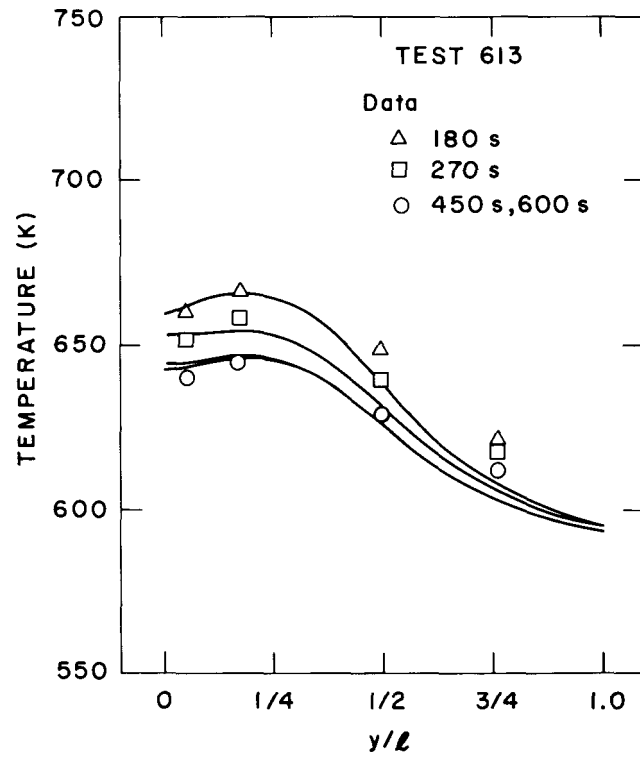


Fig. 41 Comparison of Calculated and Measured Transient Temperature Distributions for Test 613 (Late)

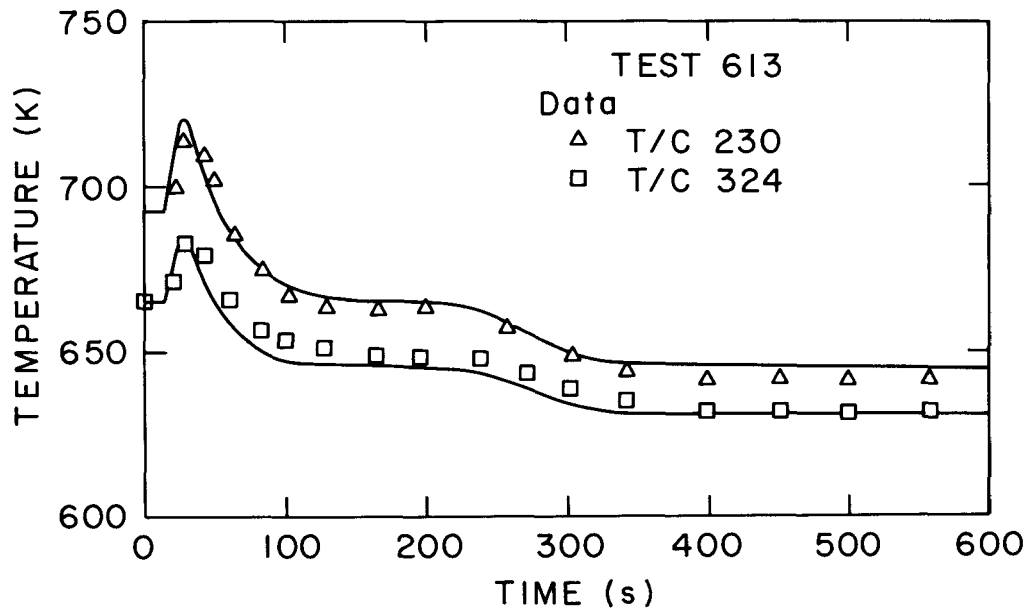


Fig. 42 Comparison of Calculated and Measured Temporal Variations of Sodium Temperature at T/C #230 and #324 for Test 613

Figure 42 shows that peak sodium temperatures are in excellent agreement with the measured data where the temperatures follow the behavior of the power to flow ratio (Fig. 38). However, the peak temperatures are higher during the initial high flow part of the transient even when the power-to-flow ratio peaked back to near initial high values later in the transient. This reflects contribution of buoyancy-induced flow redistribution and transverse thermal conduction effects which had been observed in low flow steady state tests. Heat transfer to the surrounding sodium and structures is also a significant contributing effect.

#### B. Overcooling Event

Engel, et al. [27] also performed two transient overcooling events; namely, run numbers 609 and 611, corresponding to heat input gradients of 2.8 to 1, and uniform distribution across the bundle.

Figure 43 shows the transient power-to-flow ratio forcing function approximating the measured values of the run number 611. Initially, the bundle was operating at 440 kw power, 56 gpm flow and 588 K sodium inlet temperature condition.

Following the transient initiation the power and flow rates are changed in order to achieve a rapid undercooling of the subassembly as shown by the power-to-flow ratio in Fig. 43.

Figures 44 and 46 illustrate the cross-assembly temperature profiles as calculated by the present model. Clearly, the overcooling effect is quite evident early in the transient, where the impact of heat storage in the duct and surrounding sodium region is also illustrated (it must be noted that cross-assembly temperature data were not reported for Test 611 in Ref. [27]).

Following the increase in power-to-flow ratio, caused by reduction in sodium flow rate, the sodium temperatures across the assembly increase as shown in Fig. 46. Here again, the effect of temperature flattening at low Reynolds number is clearly evident.

Figures 47 and 48 show the comparisons at selected thermocouple locations with the experimental measurements of Engel, et al. [27].

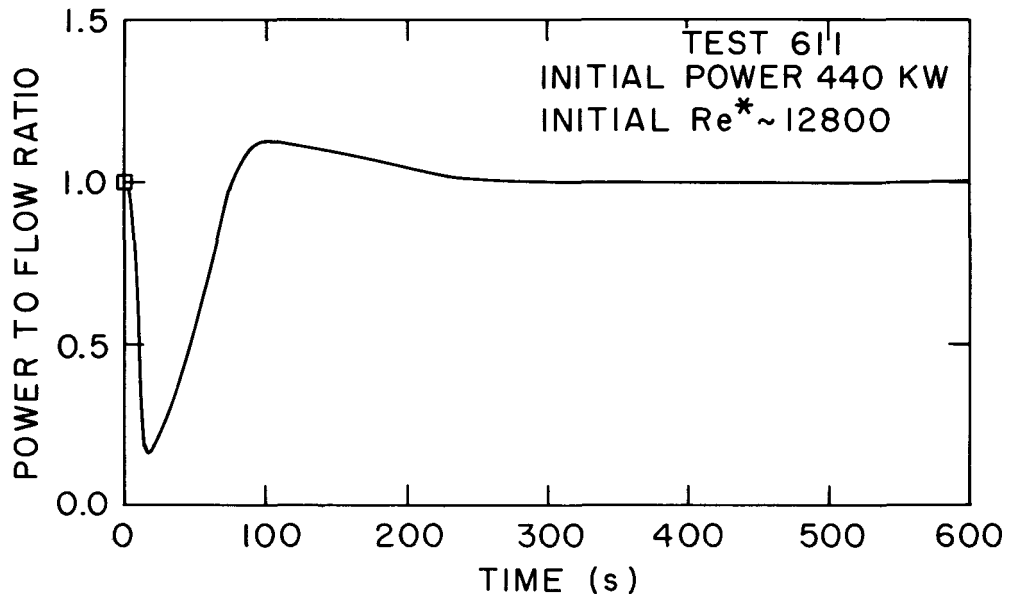


Fig. 43 Power to Flow Ratio (Test 611)

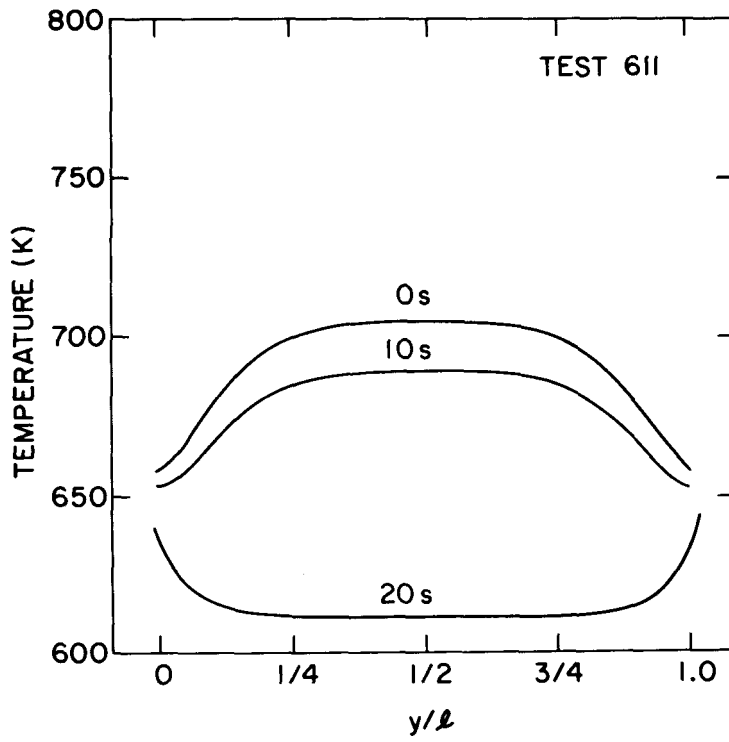


Fig. 44 Calculated Transient Temperature Distributions for Test 611 (Early)

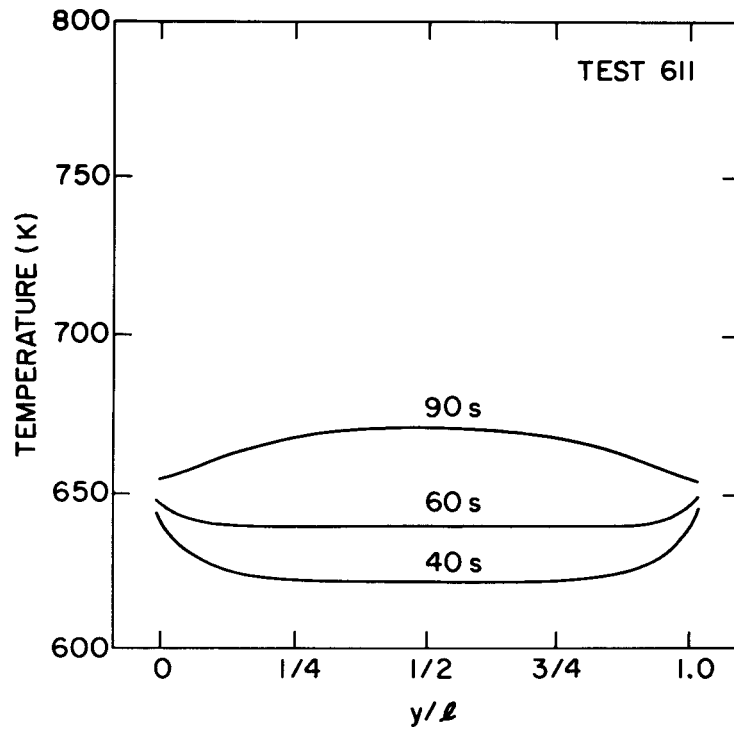


Fig. 45 Calculated Transient Temperature Distributions for Test 611 (Intermediate)

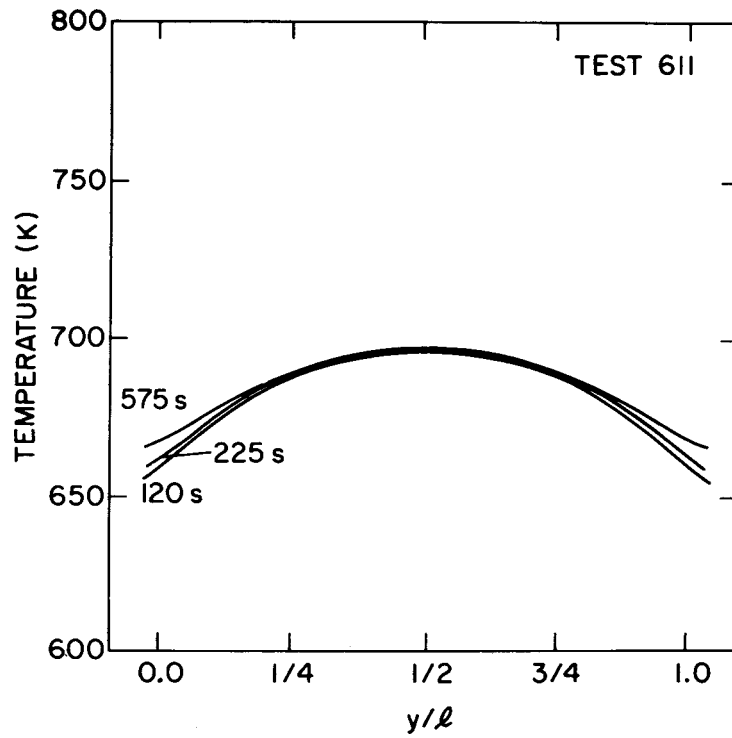


Fig. 46 Calculated Transient Temperature Distributions for Test 611 (Late)

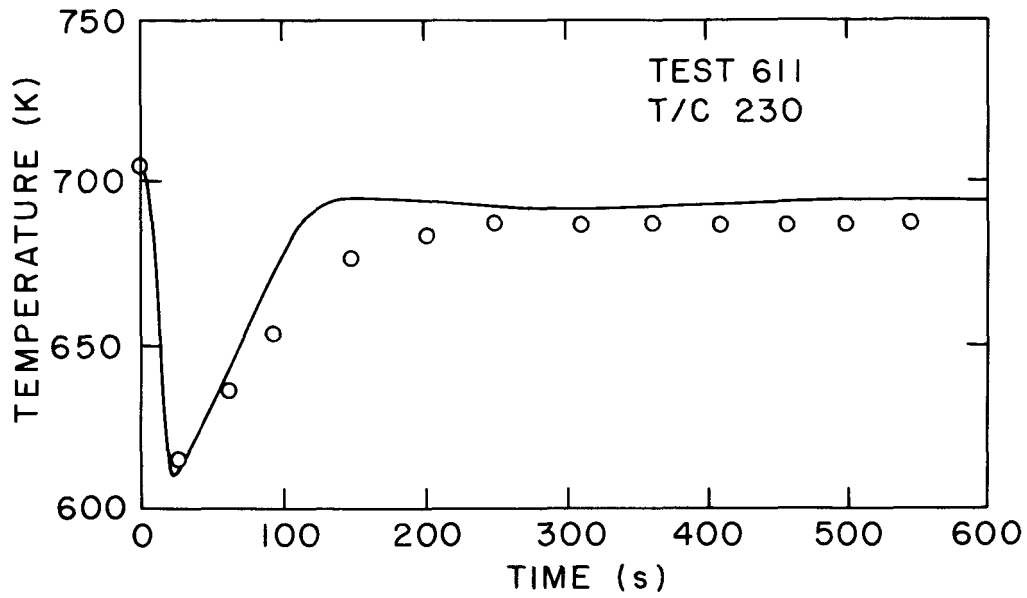


Fig. 47 Comparison of Calculated and Measured Temporal Variations of Sodium Temperature at T/C #230 for Test 611

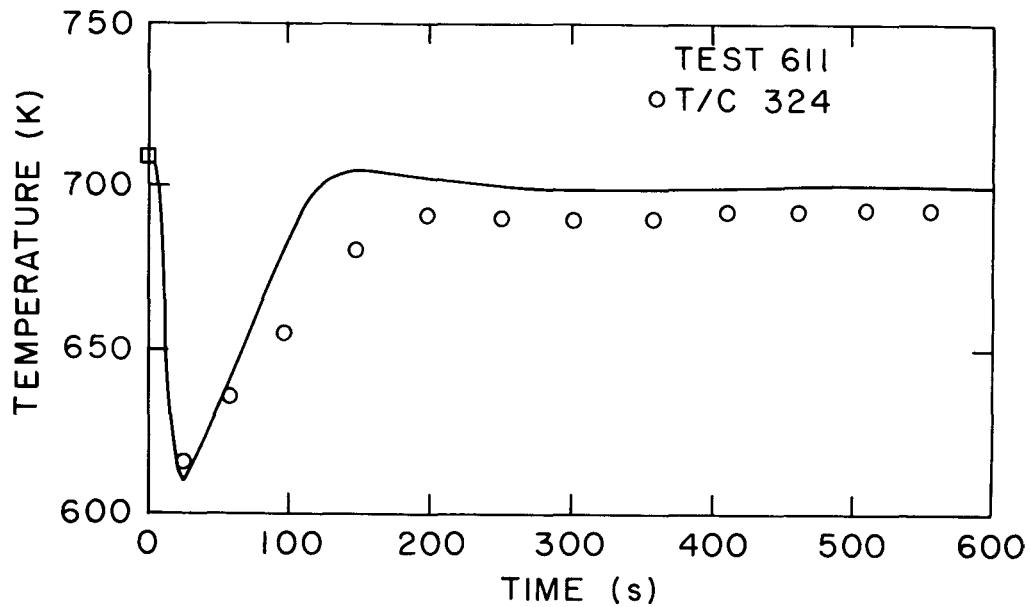


Fig. 48 Comparison of Calculated and Measured Temporal Variations of Sodium Temperature at T/C #324 for Test 611

It is seen that good agreement is obtained except between 100 to 200 seconds where the present model tends to overpredict the temperature as a result of the absence of fuel pin modeling, and perhaps inaccurate power and flow forcing functions as they were taken from a rather illegible figure of the original publication [27].

It is important to note, that the machine time requirements, for the aforementioned transients ranged from 5 to 10 times faster than real time depending on the accuracy criterion and the maximum allowable time step.

## 6. SUMMARY, CONCLUSIONS AND FUTURE WORKS

### 6.1 Summary

Two-dimensional porous-body models for study of the temperature flattening effect caused by buoyancy-induced intra-subassembly flow redistribution and transverse thermal conduction were proposed.

In Chapter 2, governing equations and porous-body parameters were set forth. In Chapter 3, approximate integral solutions for two limiting cases of flow redistribution and transverse conduction were developed. The approximate solutions provide excellent insight into the problem, while at the same time, the flow redistribution model can be used to determine maximum sodium temperatures within representative LMFBR assemblies under forced, mixed and natural convection conditions.

Detailed finite difference solutions based upon a parabolic flow model were developed in Chapter 4. The steady state model was then used to post-predict a number of sodium and water experiments.

Finally, a transient two-dimensional porous-body model applicable to forced, mixed and free convection regimes was developed in Chapter 5. This model was also applied to recent experimental measurements.

### 6.2 Conclusions

The importance of heat transfer and flow regime changes on LMFBR subassembly sodium flow behavior was demonstrated through approximate and detailed steady state and transient solutions to two-dimensional porous-body equations, and comparisons with experimental data.

Numerical results supported by experimental data indicate that,

- 1) Thermal buoyancy-induced flow redistribution is the dominant mode of the intra-assembly temperature flattening.
- 2) Flow regime change is significant.
- 3) Transverse thermal conduction can become quite significant at low flow natural convection conditions.

- 4) The proposed integral solution for calculation of the maximum sodium temperature within an LMFBR subassembly may be adequate under most of the natural circulation conditions.
- 5) The detailed two-dimensional model is quite accurate and highly efficient, and can therefore be used either directly or in conjunction with a system code for detailed calculation of intra-subassembly dynamics.

### 6.3 Future Works

The two-dimensional model described in this report is sufficient for representation of intra-subassembly behavior. However, in order to include inter-subassembly heat transfer processes as well, one needs to obtain a more complete representation of temperature and velocity fields within the entire assembly.

Therefore, for a complete three-dimensional analysis of the assembly; the Z-component of the velocity can be incorporated into the conservation equations as:

#### a) Conservation of Mass

$$\frac{\partial(\rho u)}{\partial x} + \frac{\partial(\rho v)}{\partial y} + \frac{\partial(\rho w)}{\partial z} = 0 \quad (121)$$

#### b) Conservation of Energy

$$\frac{\partial T}{\partial t} + u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} + w \frac{\partial T}{\partial z} = \alpha_e \left( \frac{\partial^2 T}{\partial y^2} + \frac{\partial^2 T}{\partial z^2} \right) + \frac{Q'''(x,y,z)}{\rho c_p} \quad (122)$$

c) Conservation of Axial Momentum

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} = - \frac{1}{\rho} \frac{\partial P}{\partial x} - g \frac{\rho^*}{\rho} \left[ 1 - \beta (T - T^*) \right] + v \left[ \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2} \right] + \frac{f}{2D\lambda^2} u^2 \quad (123)$$

This model corresponds to the usual assumptions that transverse momentum flux is small, that gradients in the x-direction are smaller than those in the other directions, and  $P = P(x)$  only. This model is as yet incomplete, because, Equations (121) through (123) represent three equations in four unknowns,  $u$ ,  $v$ ,  $w$ , and  $T$ . (As in the previous 2-D case considered earlier,  $P$  is not a true unknown in the same sense as  $u$ ,  $v$ ,  $w$ , and  $T$  since  $P = P(x)$  only.) An additional relation between  $v$  and  $w$  must be specified in order to have a complete set.

One such additional relation between  $v$  and  $w$  is suggested by Carlson [21, 29] as:

$$\frac{w}{z} = \frac{v}{y} \quad (124)$$

This relation, originally proposed for square ducts, assumes that at each point in the duct cross section the transverse velocity vector is directed towards the direct center line. Such an approximation should of course be verified and/or modified for application to large aspect ratio channels (sub-assemblies). Modifications include fomulation of transverse momentum equations, while retaining the parabolic nature of the equations ( $\partial^2/\partial x^2 = 0$ ).

The present two-dimensional model should also be incorporated into the Super System Code (SSC) [30] for best-estimate prediction of protected loss-of-flow transients in LMFBRs.

## REFERENCES

1. M. Khatib-Rahbar and K.B. Cady, "Dynamical Models and Numerical Simulation of System-Wide Transients in Loop-Type LMFBRs," Nucl. Eng. & Design, 64, 259-281, (1981)
2. E.U. Khan, "LMFBR In-Core Thermal-Hydraulics: The State of the Art and U.S. Research and Development Needs," Pacific Northwest Laboratory, PNL-3337/UC-32 (April 1980).
3. E.U. Khan, W.M. Rohsenow, A.A. Sonin and N.E. Todreas, "Manual for ENERGY I, II, III Computer Programs," Massachusetts Institute of Technology, Report COO-2245-18TR, (May 1975).
4. C.L. Wheeler, et al., "COBRA-IV-I: An Interim Version of COBRA for Thermal-Hydraulic Analysis of Rod Bundle Nuclear Fuel Element and Cores", Pacific Northwest Laboratory, BNWL-1962, (March 1976).
5. T.L. George, et al., "COBRA-WC: A Version of COBRA for Single-Phase Multiassembly Thermal-Hydraulic Transient Analysis," Pacific Northwest Laboratory, PNL-3259 (1980). Also see AIChE Symposium Series, No. 199, Vol. 76, 205-214.
6. W.L. Chen, M.A. Grolmes and M. Ishii, "A Simple Forced Diversion Model for Study of Thermal-Hydraulic Transients in LMFBR Subassembly," Nucl. Eng. & Design, 45, 53-66 (1978).
7. H.G. Johnson, "CORA - A Computer Code for Thermal and Hydraulic Coupling of Reactor Core Assemblies," Hanford Engineering Development Laboratory, TC-1505, (September 1979).
8. W.T. Sha, et al., "COMMIX-1: A Three-Dimensional Transient Single-Phase Component Computer Program for Thermal-Hydraulic Analysis," NUREG/CR-0785, ANL-77-96, (September 1978).
9. Informal information obtained from J.E. Meyer on "Transverse Heat Transfer in an LMFBR Core," (January 1979).
10. Informal information obtained from J.E. Meyer on "Intra-Assembly Buoyancy Effects in an LMFBR Core," (June 1979).
11. Informal information obtained from J.E. Meyer on "Heat Transfer Within the Coolant of an LMFBR Core," (April 1979).
12. F.C. Engel, R.A. Markley and B. Minushkin, "Buoyancy Effects on Sodium Coolant Temperature Profiles Measured in an Electrically Heated Mock-up of a 61-Rod Breeder Reactor Blanket Assembly," ASME paper 78-WA/HT-25, (1978).

13. F.E. Engel, and R.A. Markley, "ARD 61-Rod Electrically Heated Blanket Mock-up Heat Transfer Test Results for Turbulent Flow Operation and Comparison with COTEC Predictions," Paper presented at the Core T & H Technical Group Meeting, Argonne National Laboratory, April 24-25, (1979).
14. F.C. Engel, B. Minushkin, R.J. Atkins and R.A. Markley, "Characterization of Heat Transfer and Temperature Distributions in an Electrically Heated Model of an LMFBR Blanket Assembly," Nucl. Eng. & Design, 62, 335-347, (1980).
15. F.C. Engel, R.A. Markley and B. Minushkin, "Temperature Profiles in Natural and Forced Circulation of Sodium Through a Vertical LMFBR Blanket Test Assembly," American Institute of Chemical Engineers Symposium Series, 208, Milwaukee, (1981).
16. F.C. Engel, R.A. Markley and A.A. Bishop, "The Effects of Radial Heat Flux Gradients and Flow Regimes on the Peak Sodium Temperature Rise in Wire Wrapped Rod Bundles," Proceedings of the Topical Meeting on Advances in Reactor Physics and Core Thermal Hydraulics, NUREG/CP-0034, Vol. 2, (1982).
17. D.R. Spencer and R.A. Markley, "Friction Factor Correlation for 217-Pin Wire-Wrap-Spaced LMFBR Fuel Assemblies," Trans. Am. Nucl. Soc. 39, 1014, (1981).
18. E.H. Novendstern, "Pressure Drop Model for Wire-Wrapped Fuel Assemblies," Trans. Am. Nucl. Soc. 14, 660 (1971).
19. F.C. Engel, R.A. Markley and A.A. Bishop, "Laminar, Transition and Turbulent Parallel Flow Pressure Drop Across Wire Wrap Spaced Rod Bundles, Nucl. Sci. & Eng. 69, 290-296 and 74, 226, (1977).
20. C. Chiu, "Investigation of Coolant Mixing and Flow Split for LMFBR Wire-Wrapped Assemblies," PhD Thesis, Massachusetts Institute of Technology, (1977).
21. R.W. Hornbeck, "Numerical Marching Techniques for Fluid Flows With Heat Transfer," National Aeronautics and Space Administration, NASA-SP-297, Washington, D.C. (1973).
22. B. Carnahan, H.A. Luther, and J.O. Wilkes, Applied Numerical Methods, John Wiley & Sons, Inc. (1969).
23. J. Juneau and E.U. Khan, "Analysis of Steady-State Combined Forced and Free Convection Data in Rod Bundles," Argonne National Laboratory, FRA-TM-116, (January 1979).
24. A.K. Agrawal and M. Khatib-Rahbar, "Dynamic Simulation of LMFBR Systems," Atomic Energy Review 18, No. 2, 329, IAEA, Vienna (1980).

25. K. Mawatari, F. Namekawa, N. Handa, F. Kasahara, and Y. Ishida, "Natural Circulation Decay Heat Removal Experiments and Analysis in An LMFBR Fuel Assembly," Proceedings of International Topical Meeting on Fast Reactor Safety, Lyon, France, July 18-21 (1982).
26. F. Namekawa to M. Khatib-Rahbar, Personal Communication, Toshiba Nuclear Engineering Laboratory (April 1983)
27. F.C. Engel, R.A. Markley and B. Minushkin, "Loss of Flow Transient Heat Transfer Tests of a Full Size LMFBR Blanket Model," ASME Paper No. 82-WA/HT-35 (1982).
28. D. Potter, Computational Physics, John Wiley & Sons, London (1973).
29. G.A. Carlson, "Laminar Entrance Flow in a Square Duct," Ph.D. Thesis, Carnegie Institute of Technology (1966).
30. J.G. Guppy, et al., "An Advanced Thermohydraulic Simulation Code for Transients in LMFBRs, SSC, Rev. 0," Brookhaven National Laboratory Report, BNL-NUREG-51650 (April 1983).

## APPENDIX A

### USER'S MANUAL FOR FREDIS

The approximate integral method described in Section 3.1 is implemented in the FREDIS (Flow REDIstribution Model) code.

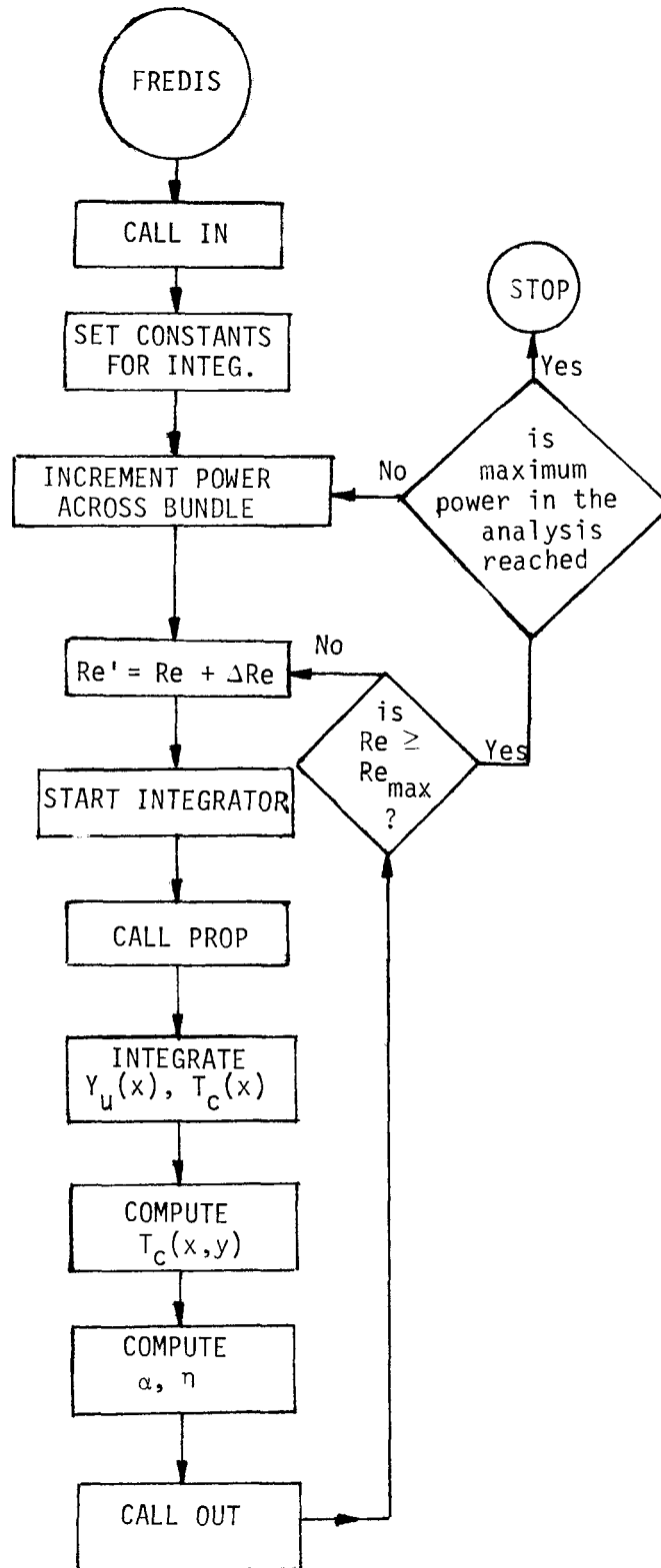
FREDIS integrates Equations (40) and (41) using the user supplied geometry and bundle inlet conditions.

The integration procedure is based upon the Gear method. However, other numerical integration techniques may be readily substituted.

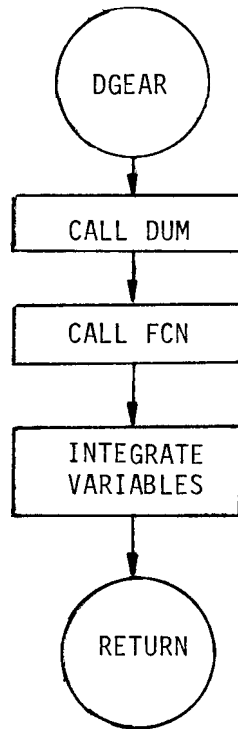
In this Appendix, the code flow chart, together with the description of various subroutines, code input, code output and limitations are discussed. The FORTRAN source listing is also included.

A.1 PROGRAM FLOWCHART AND CODE STRUCTURE

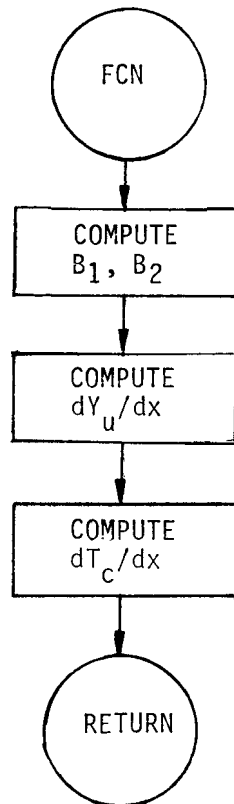
a) MAIN DRIVER



b) INTEGRATOR DRIVER



c) DIFFERENTIAL EQUATIONS



```

PROGRAM FREDIS(INPUT,OUTPUT,TAPE1,TAPE2=OUTPUT,TAPE98)
C
C PROGRAM TO COMPUTE FLOW REDISTRIBUTION PARAMETERS INSIDE
C A ROD BUNDLE DUE TO THERMAL BUOYANCY
C
C
COMMON /NEW/ .DH
COMMON /DATA/ TIN,YDE, ALAM,ALH,AL2,F,QPPP,WBAR,RE,AMU
COMMON /UNIT/ INP,IOUP,ISAVE
COMMON /WHICH/ I1,YYY(3),TTT(3)
COMMON /SAVE/ NXRE,XRE(50),XQ(50),TEMP(11,3,50),DIST(11)
C
COMMON /INTEG/ NEQ, Z, H, ZEND, Y(2),INDEX
COMMON /NEW2/ Q
COMMON /NEW4/ DELTAT
C
DIMENSION IWK(2), WK(30)
C
EXTERNAL FCN
EXTERNAL DUM
C
C
C READ INPUT DATA
C
CALL IN
C
SET CONSTANT DATA FOR INTEGRATOR
C METH = 1 ADAMS METHOD
C = 2 GEAR METHOD
C MITER = 2 CHORD METH,JACOBIAN COMP. INTERNALLY
C = 3 " " " , DIAGONAL JACOBIAN COMP.
C
TOLERANCE SHOULD BE LESS THAN OR EQUAL TO 1.E-3
C
DO 1 L=1,3
TTT(L) = 0.0
YYY(L) = 0.0
1 CONTINUE
NEQ = 2
METH = 2
MITER = 2
IEND= 0
TOL = 5.E-04
C
REYNOLDS NUMBER STEPS
C
DELTR1 = 50.
DELTR2 = 100.
DELTR3 = 500.
DELTR4 = 2000.
C
STEPS FOR DELTAT
C
DT1 = 25.
DT2 = 75.
DO 7777 M1=1,8
IF(M1.LE.3) DELTAT = DELTAT + DT1
IF(M1.GT.3) DELTAT = DELTAT + DT2
C
PERFORM CALCULATIONS FOR INCREASING REYNOLDS NUMBERS
C
DO 1000 I=1,NXRE
Q = 0.0
IF(XRE(1).EQ.0.0) GOTO 1003
C
REYNOLDS NUMBERS WERE PRESCRIBED IN INPUT
C
RE = XRE(I)
GOTO 1004
1003 IF(RE.LT.600.) RE = RE + DELTR1
IF(RE.GE.1000.0.AND.RE.LE.7000.) RE=RE+DELTR3
IF(RE.GE.600.0.AND.RE.LT.1000.)
+ RE = RE + DELTR2
IF(RE.GT.7000.) RE = RE + DELTR4
1004 I1 = I
C
START UP INTEGRATOR
C
CALL START
C
COMPUTE PROPERTIES, FRICTION FACTOR, AVERAGE INLET
VELOCITY AND POWER ACROSS THE BUNDLE
C
CALL PROP
C
INTEGRATE Y(Z) AND TCOOL(Z)
C
CALL DGEAR(NEQ,FCN,DUM,Z,H,Y,ZEND,TOL,METH,MITER,
+ INDEX,IWK,WK,IER)
C
COMPUTE T(X,Z)
C
CALL TXZ
C
COMPUTE THE TEMPERATURE RISE HOT CHANNEL FACTOR
C
CALL FF
C
OUTPUT RESULTS
C
CALL OUT

```

```

C
C
C
1000    CONTINUE
7777    CONTINUE
C
      STOP
      END
      SUBROUTINE DUM(N,X,Y,PD)
C
C
C
      DUMMY ROUTINE NEEDED BY INTEGRATOR
C
      INTEGER N
      REAL Y(N),PD(N,N),X
      RETURN
C
      END
      SUBROUTINE FCN(NEQ,Z,Y,YPRIME)
C
C
C
      THIS SUBROUTINE COMPUTES DY/DZ AND DTCOOL/DZ
C
C
C
      COMMON /NEW1/  AKE
      COMMON /DATA/  TIN,YDE,ALAM,ALH,AL2,F,QPPP,WBAR,RE,AMU
      COMMON /WHICH/  II,YYY(3),TTT(3)
      COMMON /OTHER/  BZERO,EXPON,T,DELTAH,DHBAR,FDELTA,FSUBF,
C
C
C
      CP, RHO,BTWO,BONE
      REAL Y(NEQ),YPRIME(NEQ)
C
C
C
      COMPUTE COEFFICIENTS B1 AND B2 AS FUNCTIONS OF
      Y(Z) AND EXPON
C
      GEX = 1. - EXPON
      AMULT = 2./AL2**2*AKE/(RHO*CP)/WBAR**EXPON
      Y1 = Y(1)*(1.+Y(1))**GEX-(1.-Y(1))**GEX
      Y2 = 2.*(2.-(1.+Y(1))**GEX-(1.-Y(1))**GEX)
      R1 = AMULT * (Y1+Y2)
      R2 = AMULT * Y1
      IF (RE.LE.400.) R1 = R2 = 0.0
      CALL COMPB(Y,NEQ,EXPON,BONE,BTWO)
C
C
C
      COMPUTE DERIVATIVES
C
      FACT = QPPP / ( RHO*CP*WBAR )
      DYDZ = FACT * (1.-(F-1.)/Y(1)) / ( BZERO * BONE *
+
      WBAR**(1.-EXPON))
      DYDZ = DYDZ + R1/(BZERO*BONE*WBAR**GEX)
C
      DTCOZ = FACT - BZERO*BTWO*(WBAR**(1.-EXPON)) *
+
      DYDZ
      DTCOZ = DTCOZ + R2
C
C
C
      TRANSFER DERIVATIVES TO INTEGRATOR ARRAYS
C
      YPRIME(1) = DYDZ
      YPRIME(2) = DTCOZ
C
C
C
      SAVE DATA FOR Z=1/4 AND Z=1/2 OF HEATED LENGTH
C
      ZTEST1 = ALH/4.
      ZTEST2 = ALH/2.
      IF (ABS(Z-ZTEST1).GT.10.0E-02) GOTO 1
      YYY(1) = Y(1)
      TTT(1) = Y(2)
      GOTO 2
1
      IF (ABS(Z-ZTEST2).GT.10.0E-02) GOTO 2
      YYY(2) = Y(1)
      TTT(2) = Y(2)
2
      CONTINUE
C
      RETURN
      END
      SUBROUTINE COMPB(Y,N,EX,BONE,BTWO)
C
C
C
      THIS SUBROUTINE COMPUTES COEFFICIENTS B1 AND B2
      FUNCTIONS OF Y(Z) AND EXPON NEEDED BY DIFF. EQ.S
C
      DIMENSION Y(N)
C
C
C
      YY = Y(1)
      YYP2 = ( YY + 1. ) ** (2. - EX)
      YYP4 = ( YY + 1. ) ** (4. - EX)
      YYM2 = ( 1. - YY ) ** (2. - EX)
      YYM4 = ( 1. - YY ) ** (4. - EX)
C
      B11 = (YYP2-YYM2)/(8.-4.*EX)
      B12 = ((YYP4-YYM4)/(4.-EX)-(YYP2-YYM2)/(2.-EX))/
+
      (4.*YY*YY)
      B13 = -((YYP2+YYM2-2.)/(4.-2.*EX))*(1./YY**3 -
+
      1./YY)
      B14 = (YYP4+YYM4-2.)/(8.*YY**3-2.*EX*YY**3)
C
      BONE = B11 + B12 - B13 - B14
C
      B21 = (YYP2-YYM2)/(2.-EX)
      B22 = (YYP4-YYM4)/(4.-EX)
C

```

```

BTWO = (B21 + (B22-B21)/(YY**2))/4.0
C
C
RETURN
END
SUBROUTINE IN
C
C INPUT SUBROUTINE
C
COMMON /NEW/ DH
COMMON /NEW4/ DELTAT
COMMON /DATA/ TIN,YDE,ALAM,ALH,AL2,F,QPPP,WBAR,RE,AMU
COMMON /SAVE/ NXRE,XRE(50),XQ(50)
COMMON /UNIT/ INP,IOUP,ISAVE
C
C
NAMELIST /CONTR/ NXRE
NAMELIST /GEOM1/ YDE,ALAM
NAMELIST /GEOM2/ ALH,AL2
NAMELIST /DATIN/ TIN,F,RE
NAMELIST /REYNOL/ XRE
C
C
INP = 1
IOUP = 2
ISAVE = 98
READ(INP,CONTR)
READ(INP,GEOM1)
READ(INP,GEOM2)
READ(INP,DATIN)
READ(INP,REYNOL)
C
C
FORMAT(I3)
IF(XRE(1).EQ.0.0) DELTAT = 0.
C
WRITE(ISAVE,1) NXRE
C
C
RETURN
END
SUBROUTINE START
C
THIS SUBROUTINE FIRES UP DGEAR
C
COMMON /INTEG/ NEQ, Z, H, ZEND, Y(2),INDEX
COMMON /DATA/TIN,YDE,ALAM,ALH,AL2,F,QPPP,WBAR,RE,AMU
C
C
NEQ = 2
Z = 0.0
H = 1.E-04
C
SET Y(1) TO A VERY SMALL NUMBER (NEARLY ZERO)
C
Y(1) = 1.E-06
Y(2) = TIN
ZEND = ALH
INDEX = 1
C
RETURN
END
SUBROUTINE PROP
C
THIS SUBROUTINE COMPUTES PROPERTIES AND FRICTION FACTOR:
CF(RE), EXPON(RE)
C
COMMON /NEW/ DH
COMMON /NEW1/ AKE
COMMON /NEW2/ Q
COMMON /NEW10/ C9GRAS,R1,BUOY,FRIC
COMMON /NEW4/ DELTAT
COMMON /DATA/ TIN,YDE,ALAM,ALH,AL2,F,QPPP,WBAR,RE,AMU
COMMON /OTHER/ BZERO,EXPON,T,DELTAH,DHBAR,FDELTA,FSUBF,
C CP,RHO,BTWO,BONE
C
DATA C1,C2,C3/-2.4892,220.65,-.21389/
DATA R0,R1,R2,R3/1011.597,-.22051,-1.922243E-05,
D 5.63769E-09/
C
C FRICTION FACTOR IS FOR BLANKET ASSEMBLIES
C
TT = TIN
EXPON = 1.
CF = 110.
IF(RE.GT.550.) EXPON = .25
IF(RE.GT.550.) CF = .48
FRIC = CF/RE**EXPON
C
RHO = R0+R1*TT+R2*TT*TT+R3*TT*TT*TT
C
C POROUS BODY AREA
C
AP = SQRT(3.) * AL2 * AL2 / 2.0
C
C SET THE THERMAL EXPANSION COEFFICIENT
C
BETA = 2.91E-04

```

```

C
C      OTHER SODIUM PROPERTIES
C
CP = 1630.22 - .83354*TT + 4.62838E-04*TT*TT
RHS = C1 + C2/TT + C3*ALOG(TT)
AMU = 10.**(RHS)
AKE = 109.7 - 6.4499E-02*TT + 1.1728E-05*TT*TT
C
C      COMPUTE INLET VELOCITY AND POWER
C
IF(Q.EQ.0.0) GOTO 10
DH = Q/(RE*AMU*AP*ALAM/YDE)
GOTO 20
10  TFIN = TIN + DELTAT
DH = SUM(TIN,TFIN)
20  CONTINUE
WBAR = RE*ALAM*AMU/(RHO*YDE)
QPPP = RHO*WBAR*DH/ALH
C
C      COMPUTE COEFFICIENT BZERO
C
BZERO = CF*ALAM**(EXPON-2.)*AMU**EXPON*RHO**(1.-EXPON)/
+ (9.81*BETA*YDE**(1.+EXPON)*RHO) *
+ (1. - EXPON/2.)*WBAR
IF(EXPON.GE.1.0) BZERO = CF*AMU*WBAR/(19.62*ALAM*
+ BETA*YDE*YDE*RHO)
C
C      GRASHOFF NUMBER
C
C9GRAS = 9.81*BETA*RHO**2*YDE**3*DELTAT/
+ (AMU**2)
C
C      RICHARDSON NUMBER
C
RI = C9GRAS/RE/RE
C
C      BUOYANCY PARAMETER
C
BUOY = 2.*RI / FRIC
C
C      RETURN
END
SUBROUTINE TXZ
C
C      THIS SUBROUTINE COMPUTES T(X,Z), NAMELY FOR:
C      X=AL2 AND
C      Z=ALH
C
C      EQUATION IN GENERAL IS :
C      T = TCOOL(Z) + G(WBAR,EXPON)*((1.+X/AL2*Y(Z))**(2.-EXPON)
C      -1.)
C
C      COMMON /DATA/ TIN,YDE,ALAM,ALH,AL2,F,QPPP,WBAR,RE,AMU
C      COMMON /INTEG/ NEQ,Z,H,ZEND,Y(2),INDEX
C      COMMON /OTHER/ BZERO,EXPON,T,DELTAH,DHBAR,FDELTA,FSUBF,
C      CP,RHO,BTWO,BONE
C      COMMON /SAVE/ NXRE,XRE(50),XQ(50),TEMP(11,3,50),
C      + DIST(11)
C      COMMON /WHICH/ II,YYY(3),TTT(3)
C
AL = AL2
AL2P = AL / 2.
DX = AL/10.
YYY(3) = Y(1)
TTT(3) = Y(2)
C
T = Y(2) + BZERO*(WBAR)**(1.-EXPON)/(2.-EXPON)*
+ ((1. + Y(1))**(2.-EXPON) -1.)
C
C      SAVE RESULTS FOR SLICE
C
DO 2 JJ=1,3
X = -AL2P - DX
DO 1 KK=1,11
X = X + DX
TX=TTT(JJ)+BZERO*WBAR**(1.-EXPON)/(2.-EXPON)*
+ ((1.+X/AL2P*YYY(JJ))**(2.-EXPON)-1.)
TEST = (TX-TIN)/(Y(2)-TIN)
DIST(KK) = X
TEMP(KK,JJ,II) = TEST
1  CONTINUE
2  CONTINUE
C
RETURN
END
SUBROUTINE FF
C
C      THIS SUBROUTINE COMPUTES THE TEMPERATURE RISE HOT
C      CHANNEL FACTOR AND TEMPERATURE PEAKING FACTOR
C
COMMON /NEW/ DH
COMMON /DATA/ TIN,YDE,ALAM,ALH,AL2,F,QPPP,WBAR,RE,AMU
COMMON /INTEG/ NEQ,Z,H,ZEND,Y(2),INDEX
COMMON /OTHER/ BZERO,EXPON,T,DELTAH,DHBAR,FDELTA,FSUBF,
C      CP,RHO,BTWO,BONE
C
C      TEMPERATURE PEAKING FACTOR

```

```

C      FDELTA = (T-TIN)/(Y(2)-TIN)
C      HOT CHANNEL FACTOR
C      FSUBF = (FDELTA - 1.) / (F - 1.)
C
C      RETURN
C      END
C      FUNCTION SUM(T1,T2)
C      THIS FUNCTION COMPUTES DELTAH(T1,T2)
C      C0 = 1630.22
C      C1 = -0.83354
C      C2 = 4.62838E-04
C      SUM = C0*(T2-T1) + C1/2.*(T2**2-T1**2) +
C      + C2/3.*(T2**3-T1**3)
C
C      RETURN
C      END
C      SUBROUTINE OUT
C      THIS SUBROUTINE IS THE OUTPUT ROUTINE
C
C      COMMON /NEW2/  Q
C      COMMON /NEW10/ C9GRAS,R1,BUOY,FRIC
C      COMMON /DATA/  TIN,YDE,ALAM,ALH,AL2,F,QPPP,WBAR,RE,AMU
C      COMMON /UNIT/  INP,IOUP,ISAVE
C      COMMON /WHICH/  I1
C      COMMON /SAVE/  NXRE,XRE(50),XQ(50),TEMP(11,3,50),DIST(11)
C      COMMON /OTHER/ BZERO,EXPON,T,DELTAH,DHBAR,FDELTA,FSUBF,
C      CP,RHO,BTWO,BONE
C
C      WRITE(2,2)
C      WRITE(2,100)
100  FORMAT(6H  TIN ,15X,2H Q,15X,2HRE,15X,5HFSUBF,12X,4HWBAR)
C      WRITE(2,1)  TIN,Q,RE,FSUBF,WBAR
1    FORMAT(5E15.7)
2    FORMAT(2X,50(1H*),//)
C      WRITE(2,102) C9GRAS,R1,BUOY,FRIC
102  FORMAT(3X,6HGRASHO,2X,E15.7,4X,4HFRIC,2X,E15.7,
+    4X,4HBUOY,2X,E15.7,4X,4HFRIC,2X,E15.7)
C      WRITE(1SAVE,1000) I1,RE
1000 FORMAT(I3,E16.6)
C
C      SAVE DATA FOR PLOTTING
C
C      DO 10 J=1,3
C      WRITE(1SAVE,1001) ((DIST(K),TEMP(K,J,I1));K=1,11)
10   CONTINUE
1001  FORMAT(11E10.3)
C      WRITE(1SAVE,1002)  FSUBF,BUOY
1002  FORMAT(2E16.6)
C
C      RETURN
C      END

```

## A.2 ALPHABETICAL DESCRIPTION OF ASSOCIATED SUBROUTINES

COMPB: This subroutine computes coefficients  $B_1$  and  $B_2$  as functions of  $Y_u(x)$ .

Parameters: Y      dependent variables  
              N      number of dependent variables  
              EX      friction factor exponent  
              BONE    coefficient  $B_1$ , Eq. 42  
              BTWO    coefficient  $B_2$ , Eq. 43

DUM: Dummy Routine needed by integrator package.

Parameters: N      number of dependent variables  
              X      unused  
              Y      dependent variables  
              PD      unused

FCN: This subroutine computes  $\frac{dY_u}{dx}$  and  $\frac{dT_c}{dx}$

Parameters: NEQ    number of differential equations  
              Z      independent variable  $x$   
              Y      dependent variables;  $Y_u(x)$  and  $T_c(x)$   
              YPRIME value of  $dY_u/dx$  and  $dT_c/dx$

FF: This subroutine computes  $\alpha^* = (T - T^*)/(T - T^*)$  and  $\eta^* = (\alpha^*-1)/(F-1)$

IN: Input subroutine

OUT: Output subroutine

PROP: This subroutine computes physical properties and friction factor based on average conditions.

START: This subroutine sets initial values and constants needed by integrator.

TXZ: This subroutine computes  $T(x,y)$

WREDIS: Main program driver

Auxiliary Routines:

DGEAR: Integrator package from IMSLX Program Library (version 1981); can use Adam's method or Gear's method to integrate system of differential equations;

Parameters: NEQ number of differential equations  
FCN external routine which computes R.H.S. of differential equations  
DUM dummy external routine  
Z independent variable  
H initial step on independent variable  
Y array of dependent variables  
ZEND final value of independent variable  
TOL tolerance  
METH method used to integrate  
MITER defines which chord method used in integration  
INDEX unused  
IWK internal dimensions of working area for dependent variables  
WK working area for integration

### A.3 DESCRIPTION OF INPUT

TAPE 1 is defined as the input file; input consists of 5 NAMELIST groups as follows:

NAMELIST/CONTR/	NXRE	number of Reynolds numbers for which steady-state calculations are desired.
NAMELIST/GEOM1/	YDE	subassembly hydraulic diameter, m
	ALAM	subassembly channel porosity
NAMELIST/GEOM2/	ALH	subassembly heated length, m
	AL2	subassembly flat-to-flat, m
NAMELIST/DATIN/	TIN	inlet temperature, K
	F	peak to average power
	RE	initial Reynolds number
NAMELIST/REYNOL/	XRE	specific Reynolds numbers for which steady-state calculations is desired; if the first is = 0.0, the program will automatically step in Re and power (up to 30 Re values can be specified).

The following is a sample input file

```
$CONTR  NXRE=40  $END
$GEOM1  YDE=3.7713E-03, ALAM=.2546  $END
$GEOM2  ALH=1.143, AL2=.11312  $END
$DATIN  TIN=590.23, F=1.4723, RE=50.  $END
$REYNOL XRE=0.,0.,0.,0.  $END
```

This run will perform calculations for a power skew case of 2:1, inlet T of 590.23 K, starting at Re = 50, stepping automatically in Re and power for 40 progressively increasing Reynolds numbers.

#### A.4 DESCRIPTION OF OUTPUT

TAPE 2 is defined as the output file, while TAPE 99 is defined as the result storage file for eventual plotting and/or display.

On output the following data is displayed:

TIN	-	inlet temperature, K
Q	-	average power, watts/m <sup>2</sup>
RE	-	Reynolds number
FSUBF	-	hot channel factor, $\eta^*$
WBAR	-	average velocity, m/s
C9GRAS	-	Grashof number
RI	-	Richardson number
BUOY	-	buoyancy parameter (2RI/FRIC)
FRIC	-	friction factor

The following data is stored:

NXRE	-	number of Reynolds numbers for which steady-state calculations are to be performed
II	-	Reynolds number index, $1 \leq II \leq NXRE$
RE	-	current Reynolds number
DIST(k)	-	node position, (k = 1,11), m
TEMP(k,J,II)	-	Temperature at position k for Reynolds number index II at J = 1 (one fourth of heated length); J = 2 (one half of heated length); J = 3 (exit of heated section), K
FSUBF	-	temperature peaking factor $\alpha^*$
BUOY	-	buoyancy parameter

Following is a sample output as generated by FREDIS:

\*\*\* WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR  
\*\*\*\*\*

TIN	Q	RE	FSUBF	WBAR		
.5902300E+03	0.	.1000000E+03	.2081656E+00	.2551634E-02		
GRASHO	.2679620E+05	RICH .2679620E+01	BUOY .4872036E+01		FRIC	.1100000E+01

\*\*\* WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR  
\*\*\*\*\*

TIN	Q	RE	FSUBF	WBAR		
.5902300E+03	0.	.1500000E+03	.3066541E+00	.3827451E-02		
GRASHO	.2679620E+05	RICH .1190942E+01	BUOY .3248024E+01		FRIC	.7333333E+00

\*\*\* WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR  
\*\*\*\*\*

TIN	Q	RE	FSUBF	WBAR		
.5902300E+03	0.	.2000000E+03	.3917585E+00	.5103268E-02		
GRASHO	.2679620E+05	RICH .6699050E+00	BUOY .2436018E+01		FRIC	.5500000E+00

\*\*\* WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR  
\*\*\*\*\*

TIN	Q	RE	FSUBF	WBAR		
.5902300E+03	0.	.2500000E+03	.4625741E+00	.6379084E-02		
GRASHO	.2679620E+05	RICH .4287392E+00	BUOY .1948815E+01		FRIC	.4400000E+00

\*\*\* WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR  
\*\*\*\*\*

TIN	Q	RE	FSUBF	WBAR		
.5902300E+03	0.	.3000000E+03	.5206782E+00	.7654901E-02		
GRASHO	.2679620E+05	RICH .2977355E+00	BUOY .1624012E+01		FRIC	.3666667E+00

\*\*\* WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR  
\*\*\*\*\*

TIN	Q	RE	FSUBF	WBAR		
.5902300E+03	0.	.3500000E+03	.5683969E+00	.8930718E-02		
GRASHO	.2679620E+05	RICH .2187445E+00	BUOY .1392010E+01		FRIC	.3142857E+00

\*\*\* WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR  
\*\*\*\*\*

TIN	Q	RE	FSUBF	WBAR		
.5902300E+03	0.	.4000000E+03	.6080088E+00	.1020654E-01		
GRASHO	.2679620E+05	RICH .1674762E+00	BUOY .1218009E+01		FRIC	.2750000E+00

\*\*\* WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR  
\*\*\*\*\*

TIN	Q	RE	FSUBF	WBAR		
.5902300E+03	0.	.4500000E+03	.6412715E+00	.1148235E-01		
GRASHO	.2679620E+05	RICH .1323269E+00	BUOY .1082675E+01		FRIC	.2444444E+00

\*\*\* WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR  
\*\*\*\*\*

TIN	Q	RE	FSUBF	WBAR		
.5902300E+03	0.	.5000000E+03	.6694783E+00	.1275817E-01		
GRASHO	.2679620E+05	RICH .1071848E+00	BUOY .9744073E+00		FRIC	.2200000E+00

\*\*\* WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR  
\*\*\*\*\*

TIN	Q	RE	FSUBF	WBAR		
.5902300E+03	0.	.5500000E+03	.6936627E+00	.1403399E-01		
GRASHO	.2679620E+05	RICH .8858248E-01	BUOY .8858248E+00		FRIC	.2000000E+00

\*\*\* WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR  
\*\*\*\*\*

TIN	Q	RE	FSUBF	WBAR		
.5902300E+03	0.	.7000000E+03	.8265928E+00	.1786144E-01		
GRASHO	.2679620E+05	RICH .5468612E-01	BUOY .1172034E+01		FRIC	.9331828E-01

\*\*\* WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR  
\*\*\*\*\*

TIN	Q	RE	FSUBF	WBAR		
.5902300E+03	0.	.8000000E+03	.8645263E+00	.2041307E-01		
GRASHO	.2679620E+05	RICH .4186906E-01	BUOY .9278002E+00		FRIC	.9025447E-01

\*\*\* WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR  
\*\*\*\*\*

TIN	Q	RE	FSUBF	WBAR		
.5902300E+03	0.	.9000000E+03	.8913175E+00	.2296470E-01		
GRASHO	.2679620E+05	RICH .3308173E-01	BUOY .7549837E+00		FRIC	.8763561E-01

\*\*\* WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR  
\*\*\*\*\*

TIN	Q	RE	FSUBF	WBAR		
.5902300E+03	0.	.1000000E+04	.9108358E+00	.2551634E-01		
GRASHO	.2679620E+05	RICH .2679620E-01	BUOY .6278588E+00		FRIC	.8535741E-01

\*\*\* WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR  
\*\*\*\*\*

TIN	Q	RE	FSUBF	WBAR		
.5902300E+03	0.	.1500000E+04	.9582280E+00	.3827451E-01	FRIC	.7712913E-01
GRASHO	.2679620E+05	RICH .1190942E-01	BUOY .3098178E+00			
*** WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR						
*****						
.5902300E+03	0.	.2000000E+04	.9753961E+00	.5103268E-01	FRIC	.7177674E-01
GRASHO	.2679620E+05	RICH .6699050E-02	BUOY .1866635E+00			
*** WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR						
*****						
.5902300E+03	0.	.2500000E+04	.9836270E+00	.6379084E-01	FRIC	.6788225E-01
GRASHO	.2679620E+05	RICH .4287392E-02	BUOY .1263185E+00			
*** WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR						
*****						
.5902300E+03	0.	.3000000E+04	.9882004E+00	.7654901E-01	FRIC	.6485761E-01
GRASHO	.2679620E+05	RICH .2977355E-02	BUOY .9181207E-01			
*** WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR						
*****						
.5902300E+03	0.	.3500000E+04	.9910750E+00	.8930718E-01	FRIC	.6240570E-01
GRASHO	.2679620E+05	RICH .2187445E-02	BUOY .7010401E-01			
*** WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR						
*****						
.5902300E+03	0.	.4000000E+04	.9929867E+00	.1020654E+00	FRIC	.6035680E-01
GRASHO	.2679620E+05	RICH .1674762E-02	BUOY .5549540E-01			
*** WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR						
*****						
.5902300E+03	0.	.4500000E+04	.9943355E+00	.1148235E+00	FRIC	.5860546E-01
GRASHO	.2679620E+05	RICH .1323269E-02	BUOY .4515856E-01			
*** WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR						
*****						
.5902300E+03	0.	.5000000E+04	.9953490E+00	.1275817E+00	FRIC	.5708194E-01
GRASHO	.2679620E+05	RICH .1071848E-02	BUOY .3755471E-01			
*** WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR						
*****						
.5902300E+03	0.	.5500000E+04	.9960936E+00	.1403399E+00	FRIC	.5573790E-01
GRASHO	.2679620E+05	RICH .8858248E-03	BUOY .3178537E-01			
*** WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR						
*****						
.5902300E+03	0.	.6000000E+04	.9966682E+00	.1530980E+00	FRIC	.5453853E-01
GRASHO	.2679620E+05	RICH .7443389E-03	BUOY .2729589E-01			
*** WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR						
*****						
.5902300E+03	0.	.6500000E+04	.9970937E+00	.1658562E+00	FRIC	.5345802E-01
GRASHO	.2679620E+05	RICH .6342296E-03	BUOY .2372813E-01			
*** WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR						
*****						
.5902300E+03	0.	.7000000E+04	.9975538E+00	.1786144E+00	FRIC	.5247673E-01
GRASHO	.2679620E+05	RICH .5468612E-03	BUOY .2084205E-01			
*** WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR						
*****						
.5902300E+03	0.	.9500000E+04	.9987790E+00	.2424052E+00	FRIC	.4861948E-01
GRASHO	.2679620E+05	RICH .2969108E-03	BUOY .1221366E-01			
*** WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR						
*****						
.5902300E+03	0.	.1150000E+05	.9993048E+00	.2934379E+00	FRIC	.4635182E-01
GRASHO	.2679620E+05	RICH .2026178E-03	BUOY .8742603E-02			
*** WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR						
*****						
.5902300E+03	0.	.1350000E+05	.9996547E+00	.3444706E+00	FRIC	.4453052E-01
GRASHO	.2679620E+05	RICH .1470299E-03	BUOY .6603556E-02			
*** WARNING WITH FIX ERROR (IER = 67) FROM IMSL ROUTINE DGEAR						
*****						

## A.6 LIMITATIONS

The following limitations must be kept in mind:

1. For  $F=1$  (Uniform Distribution)  $\eta^*$  is set to unity.
2. Appropriate flow-dependent friction factor parameters must be inputted. The present version uses the blanket assembly correlations given in Table I as default values.
3. The integrator subroutine DGEAR should be substituted with an equivalent integration routine, if not available.



## APPENDIX B

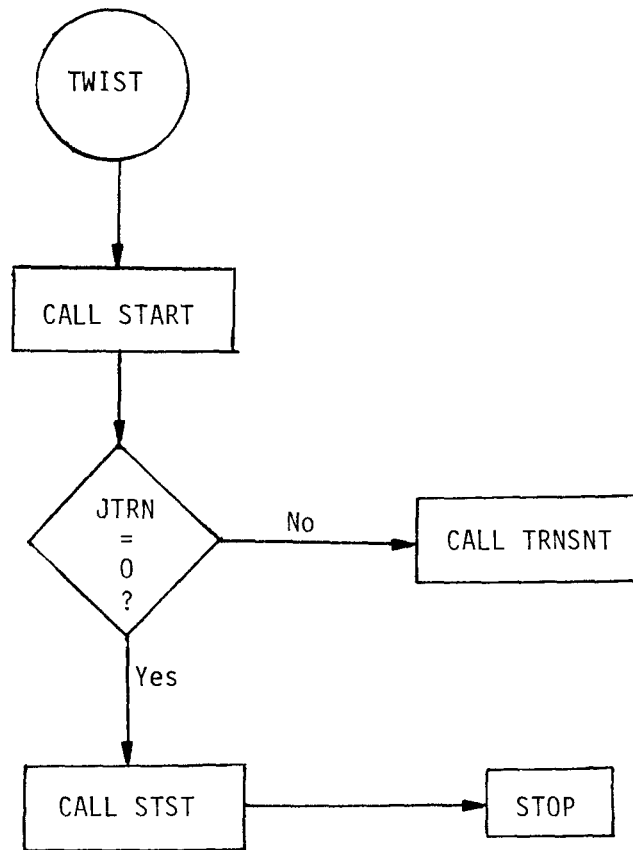
### USER'S MANUAL FOR TWIST

The steady-state and transient parabolic flow model developed in Chapters 4 and 5 have been implemented into the Two-dimensional Intra-Subassembly Thermal-hydraulics (TWIST) code.

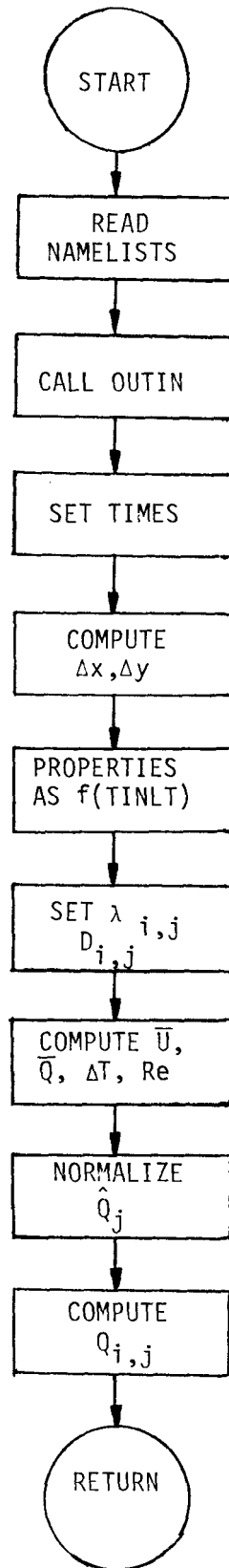
TWIST is a highly modular computer program capable of calculating a number of steady-states sequentially, or as initial conditions to the transient calculations, using user supplied boundary conditions and transient forcing functions.

B.1 PROGRAM FLOW CHART AND CODE STRUCTURE

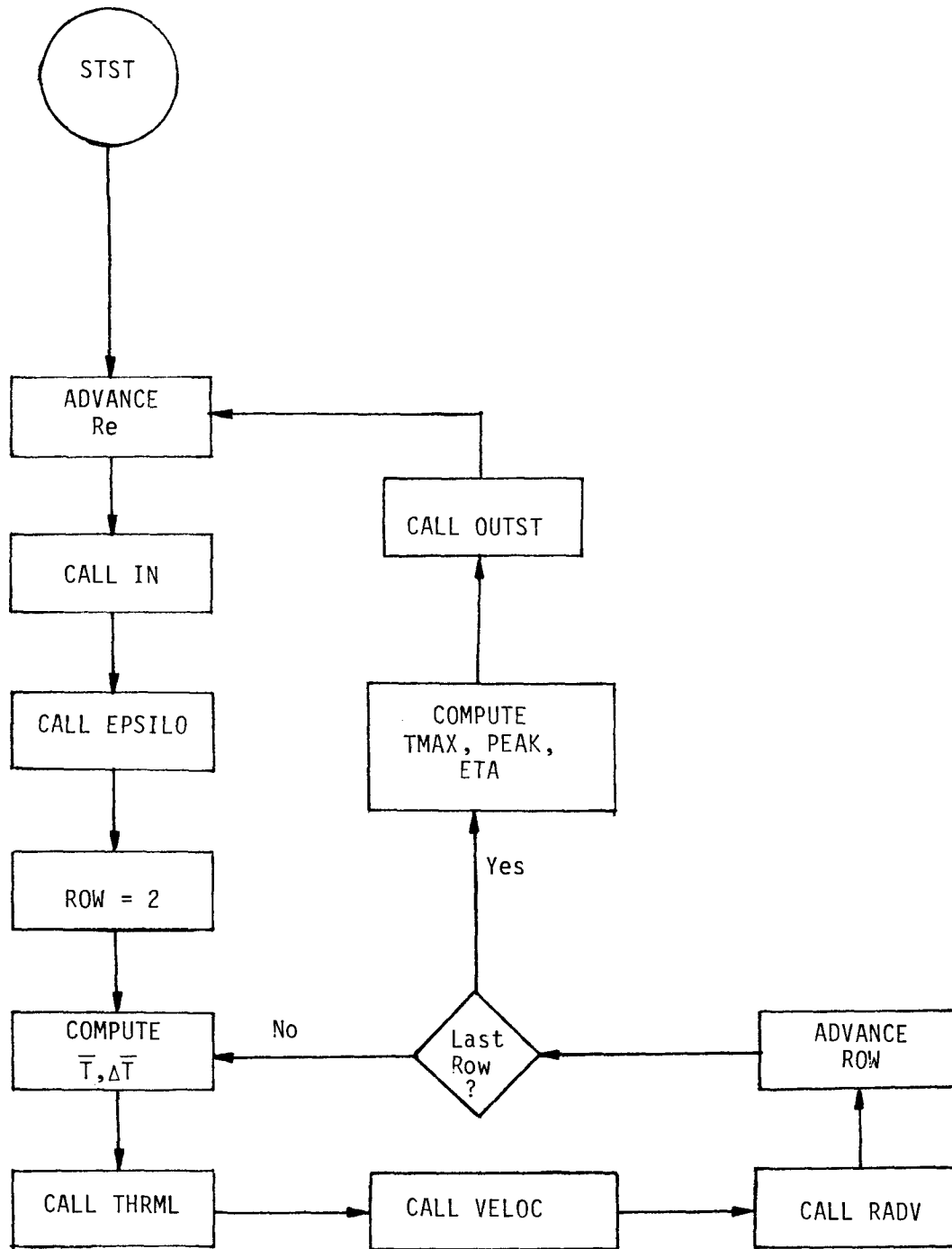
1) MASTER DRIVER



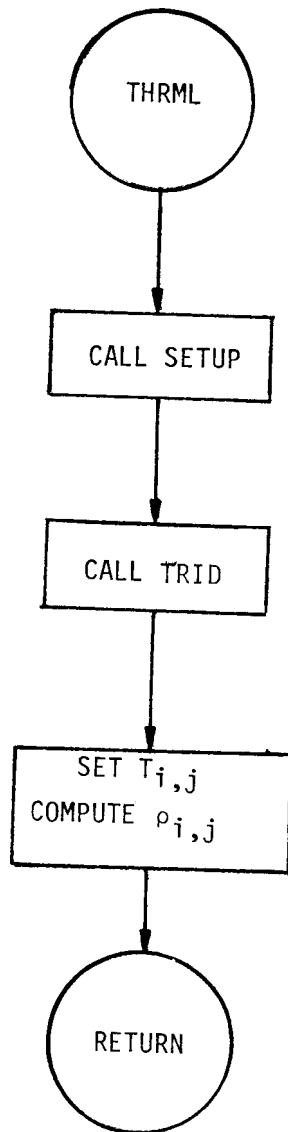
2) START MODULE



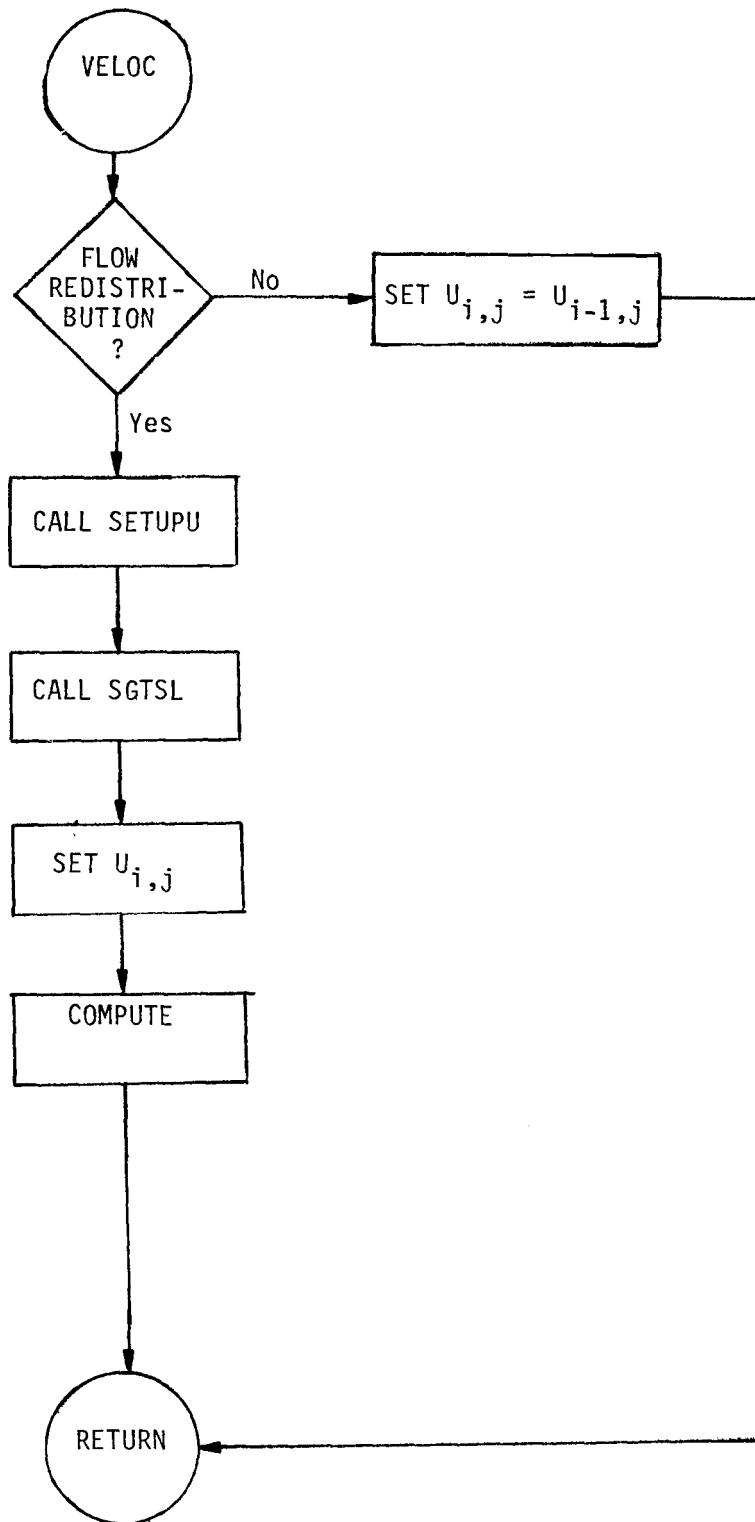
3) STEADY-STATE DRIVER



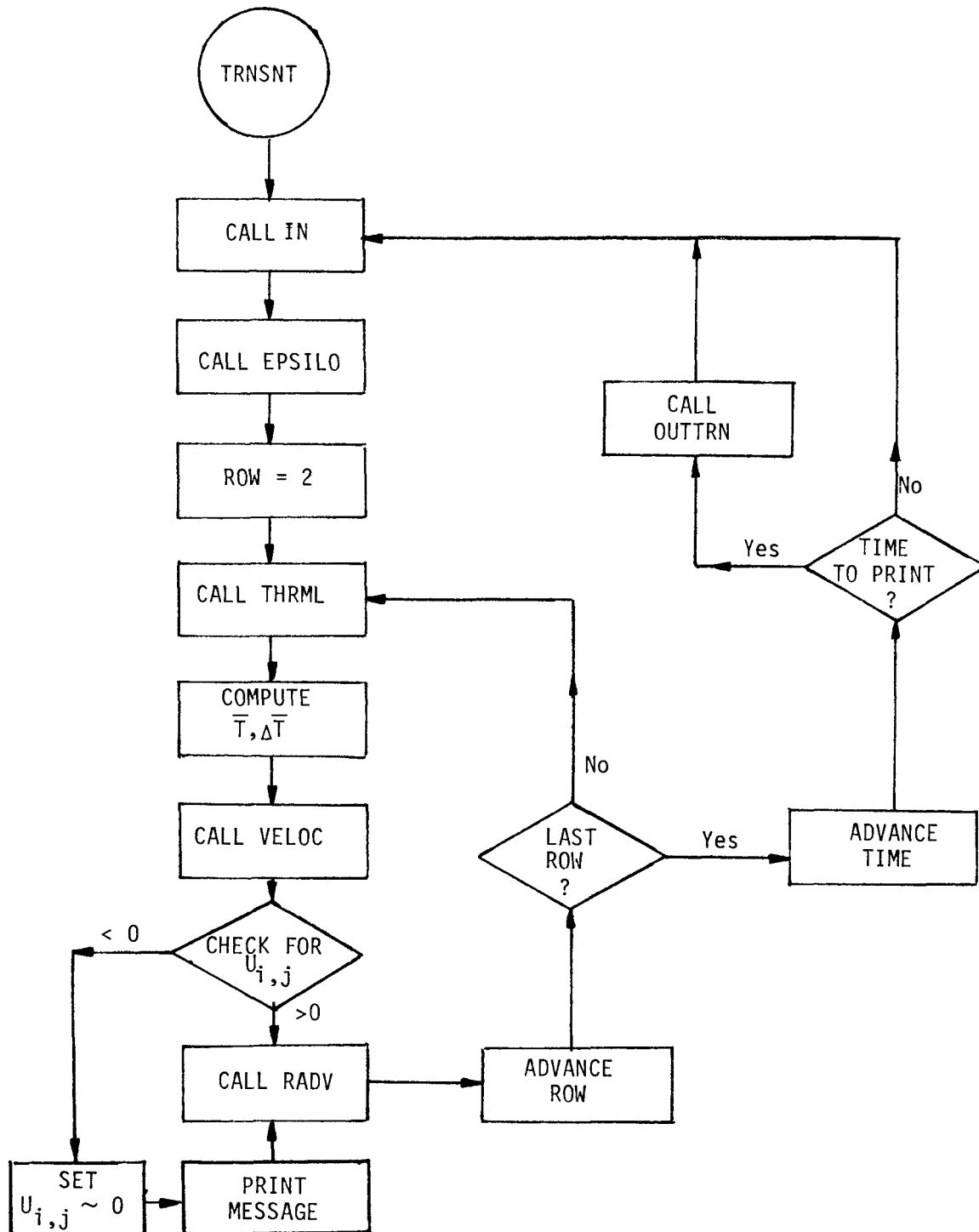
5) DRIVER FOR ENERGY EQUATION



6) DRIVER FOR AXIAL MOMENTUM EQUATION



4) TRANSIENT DRIVER



```

PROGRAM TWIST(INPUT,OUTPUT,TAPE1,TAPE2=OUTPUT,TAPE98)
C
C   PROGRAM TO COMPUTE TWO DIMENSIONAL TEMP. AND VELOCITY
C   DISTRIBUTION IN A ROD BUNDLE.
C   X IS AXIAL DIRECTION ( U )
C   Y IS TRANSVERSE DIRECTION ( V )
C   U IS AXIAL VELOCITY
C   V IS TRANSVERSE VELOCITY
C
C   NEEDS AS BOUNDARY CONDITIONS
C
C   T(I,J)   T DISTRIB.  $\leq$  INLET
C   U(I,J)   U DISTRIB.  $\leq$  INLET
C   V(I,J)   V DISTRIB.  $\leq$  INLET
C
C   OTHER DEFINITIONS :
C   XL  AXIAL LENGTH OF ASSEMBLY
C   YL  ASSEMBLY FLAT-TO-FLAT
C   N   NUMBER OF RADIAL NODES
C   M   NUMBER OF AXIAL NODES
C
COMMON /TRNOP/  JTRN
C
C   READ INPUT DATA
C
CALL START
C
C   IF JTRN = 0 COMPUTE STEADY STATE ONLY
C
IF(JTRN.EQ.0)  CALL STST
C
C   IF JTRN = 1 COMPUTE TRANSIENT
C
IF(JTRN.EQ.1)  CALL TRNSNT
C
C
C   STOP
C   END
C   SUBROUTINE START
C
C   THIS SUBROUTINE READS INPUT DATA
C   THROUGH 12 NAMELIST GROUPS
C
COMMON /TRNOP/  JTRN
C
COMMON /DATIN/  IDETAI,NRE,MODE,ALINT,ALEXT,YDINT,YDEXT,REEND,
+  QNORM(40)
C
COMMON /MORE/  AMU,F,FRIC,DTAV,YDE,BETA,TINLT,ALAM,RE,
+  DELTAT,PEAK,ETA,QBAR,TBAR,XKF,ICOND,IRED5
C
COMMON /KNOW/  XL,YL,N,M,RHO,CP,ALPHA,ALK,DX,DY,UBAR,
+  CF,EPSII,Q(40,40)
C
COMMON /AXIA/  XLLFGP,XLH,XLUFGP,MLFGP,MH,MUFGP,
+  DXLFGP,DXH,DXUFGP
C
COMMON /LAMBDA/  ALAJ(40),YD(40),USTAR,UBOLD,TBOLD
C
COMMON /TIMER/  STIME,SDELTA,CTIME,CTIMS,DTTRN,SEND,SDMAX,SDMIN
C
COMMON /PRINTR/  SPRNT,SNEXT,DTPRNT
C
COMMON /STRC/  XMS,CS,DELTAW,XMC1,XMCN
C
COMMON /SS/  UZERO,QZERO
C
COMMON /U/  TIME(40),URAT(40),NIME
C
COMMON /P/  TIM(40),PW(40),NTIM
C
C   NAMELIST GROUPS
C
C   NAMELIST /OPTION/  JTRN
C
C   NAMELIST /CONTR/  IDETAI,NRE,IRED5,ICOND,N,
+  MLFGP,MH,MUFGP
C
C   NAMELIST /GEOM1/  XLLFGP,XLH,XLUFGP,YL
C
C   NAMELIST /GEOM2/  ALAM,ALINT,ALEXT
C
C   NAMELIST /GEOM3/  YDE,YDINT,YDEXT
C
C   NAMELIST /INITC/  RE,REEND,TINLT,EPSII,QBAR,
+  DELTAT,UBAR
C
C   NAMELIST /PROP/  BETA
C
C   NAMELIST /RADQ/  QNORM
C
C   NAMELIST /TIMES/  SDMIN,SDMAX,DTPRNT,SEND
C
C   NAMELIST /WALL/  XMS,CS,DELTAW,XMC1,XMCN
C
C   NAMELIST /UTIME/  TIME,URAT
C
C   NAMELIST /PTIME/  TIM , PW
C
C
C   READ INPUT DATA FROM NAMELISTS

```

```

C
READ(1,OPTION)
READ(1,CONTR)
READ(1,GEOM1)
READ(1,GEOM2)
READ(1,GEOM3)
READ(1,INITC)
READ(1,PROP)
READ(1,RADQ)
READ(1,TIMES)
READ(1,WALL)
READ(1,UTIME)
READ(1,PTIME)

C
C      SET TRANSIENT TIME AND INTERVALS
C
C      INITIAL TIME
C
C      STIME = 0.0
C      PRINTOUT TIME
SPRNT = STIME
C      NEXT PRINTOUT TIME
SNEXT = SPRNT
IF(SDMAX.EQ.0.0)   SDMAX = 1.E-05
C      THE LATTER IS NECESSARY TO AVOID DIVISION BY 0
C
C      INITIAL TIME STEP
SDELT = SDMAX / 32.0

C      MODE=0 MEANS THAT UBAR IS NOT SPECIFIED,HENCE WILL BE COMPUTED
C      FROM REYNOLDS NUMBER
C
MODE = 0
IF(UBAR.GT.0.0)   MODE = 1

C
C      DEFINE ASSEMBLY GEOMETRY
C
XL = XLH + XLUFGP + XLLFGP
DXH = XLH / FLOAT(MH-1)
DXLFGP = 0.0
DXUFGP = 0.0
IF(MLFGP.EQ.0)   GOTO 1
DXLFGP = XLLFGP / FLOAT(MLFGP-1)
DXH = XLH / FLOAT(MH)
1   IF(MUFGP.EQ.0)   GOTO 2
DXUFGP = XLUFGP / FLOAT(MUFGP)
2   M = MLFGP + MH + MUFGP
DY = YL / FLOAT(N-3)

C
C      COMPUTE PROPERTIES AS FUNCTION OF INLET CONDITIONS
C
RHO = DENS(TINLT)
CP = CAP(TINLT)
AMU = VISC(TINLT)

C
C      SET CONDUCTION AND DISTRIBUTION PARAMETERS
C
C      NO CONDUCTION
IF(ICOND.EQ.0)   EPSII = 0.0
C      CALCULATE TORTUOSITY
XKF = ALAM / (1. + .5*(1.-ALAM))
C      NO FLOW REDISTRIBUTION
IF(IREDS.EQ.0)   BETA = 0.0

C
C      SET ASSEMBLY POROSITIES AND HYDR. DIAMETERS
C
YD(1) = YD(N) = YDEXT
YD(2) = YD(N-1) = YDEXT
YD(3) = YD(N-2) = YDEXT
ALAJ(1) = ALAJ(N) = ALEXT
ALAJ(2) = ALAJ(N-1) = ALEXT
ALAJ(3) = ALAJ(N-2) = ALEXT
DELY = (YDEXT - YDINT) / 3.
YD(4) = YD(N-3) = YDEXT - DELY
YD(5) = YD(N-4) = YDEXT - 2.*DELY
DELA = (ALEXT - ALINT) / 3.
ALAJ(4) = ALAJ(N-3) = ALEXT - DELA
ALAJ(5) = ALAJ(N-4) = ALEXT - 2.*DELA
NEND = N - 5
DO 90 JJ=6,NEND
ALAJ(JJ) = ALINT
YD(JJ) = YDINT
90   CONTINUE

C
C      DEFINE INLET AVER. VELOCITY,TOTAL POWER AND AVERAGE TEMP. RISE
C
IF(MODE.EQ.1)   GOTO 91
UBAR = RE *ALAM * AMU / RHO / YDE
GOTO 92
91   RE = UBAR * RHO * YDE / ALAM / AMU
92   CONTINUE
IF(QBAR.EQ.0.0)   GOTO 100
AP = SQRT(3.) / 2. * YL*YL
QBAR = QBAR / XLH / AP
DELTAT = QBAR * XLH / RHO / CP / UBAR
GOTO 101
100   QBAR = DELTAT * RHO * CP * UBAR / XLH
101   CONTINUE
UZERO = UBAR
QZERO = QBAR
C

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```

C          MASS OF STRUCTURE PER AXIAL NODE
C
C          XMS = XMS / FLOAT(M-1)
C
C          NUMBER OF DATA POINTS IN FORCING FUNCTIONS TABLES
C
C          DO 200 KT = 2,40
C          IF (TIME(KT).LE.0.0) GOTO 200
C          NTIME = KT
200      CONTINUE
C
C          DO 201 KT=2,40
C          IF (TIM(KT).LE.0.0) GOTO 201
C          NTIM = KT
201      CONTINUE
C
C          NORMALIZE RADIAL POWER DISTRIBUTION, USING SIMPSON'S RULE
C
C          NMM = N - 3
C          QTOT = QNORM(3) + QNORM(N-2)
C          DO 400 J=4,NMM
C          COEF = 4.
C          JL = J - (J/2) * 2
C          IF (JL.NE.0) COEF = 2.
C          QTOT = QTOT + QNORM(J) * COEF
400      CONTINUE
C          QTOT = QTOT / 3. / FLOAT(N-5)
C          DO 410 J=1,N
410      QNORM(J) = QNORM(J) / QTOT
C
C          COMPUTE POWER SKEW (PEAK-TO-AVERAGE)
C
C          F = QNORM(3) / ((QNORM(3)+QNORM(N-2))/2.)
C
C          AXIAL POWER DISTRIBUTION IS IN THE SHAPE OF A CHOPPED
C          COSINE CURVE, WITH PEAK-TO-AVERAGE OF 1.4
C
C          X = 0.0
C          DO 500 I=1,M
C          QX = 1.4 * COS(.8738 * 3.14156 * ((X-XLLFGP)/XLH - .5))
C          IF (I.LE.MLFGP.OR.I.GT.(MLFGP+MH)) QX = 0.0
C          DX = DXLFGP
C          IF (I.GE.MLFGP) DX = DXH
C          X = X + DX
C          DO 600 J=1,N
C          Q(I,J) = QBAR * QNORM(J) * QX
C          IF (J.EQ.1.OR.J.EQ.N) Q(I,J) = 0.0
C          IF (J.EQ.2.OR.J.EQ.(N-1)) Q(I,J) = 0.0
600      CONTINUE
500      CONTINUE
C
C          CALL OUTIN
C
C          RETURN
C          END
C          SUBROUTINE STST
C
C          THIS SUBROUTINE IS THE STEADY-STATE DRIVER
C
C          COMMON /KNOW/ XL, YL, N, M, RHO, CP, ALPHA, ALK, DX, DY, UBAR,
+          CF, EPSII,
+          Q(40,40)
C
C          COMMON/AXIA/ XLLFGP,XLH,XLFGP,MLFGP,MH,MUFGP,
+          DXLFGP,DXH,DXUFGP
C
C          COMMON /DATIN/ IDETAI,NRE,MODE,ALINT,ALEXT,YDINT,YDEXT,REEND,
+          QNORM(40)
C
C          COMMON /UNKN/ T(40,40), U(40,40), V(40,40),DEN(40,40)
C
C          COMMON /MORE/ AMU,F,FRIC,DTAV,YDE,BETA,TINLT,ALAM,RE
+          ,DELTAT,PEAK,ETA,QBAR,TBAR,XKF,ICOND,IRED5
C
C          COMMON /LAMBDA/ ALAJ(40),YD(40),USTAR,UBOLD,TBOLD
C
C          SET REYNOLDS NUMBER STEPS
C
C          DELRE1 = 100.
C          DELRE2 = 500.
C          DELRE3 = 5000.
C
C          BEGIN COMPUTATIONS
C
C          DO 700 IRE = 1,NRE
C
C          GET INLET CONDITIONS
C          UBOLD = UBAR
C
C          CALL IN
C
C          COMPUTE EDDY DIFFUSIVITY OF HEAT
C          CALL EPSILO

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TBAR = TINLT
C
C      NOW MARCH FROM ROW TWO TO TOP OF ASSEMBLY
C
X = 0.0
DO 100 I=2,M
C
C      FIND WHETHER IT IS A PLENA SLICE OR HEATED
C      COMPUTE TBAR FOR SLICE FROM DELTAT
C
DX = DXLFGP
IF(I.GT.MLFGP) DX = DXH
IF(I.GT.(MLFGP+MH)) DX = DXUFGP
X = X + DX
TBOLD = TBAR
TBAR = TINLT + 1.4*QBAR/RHO/CP/UBAR*XLH/
+ (3.14156*.8738)*(SIN(.8738*3.14156*
+ ((X-XLLFGP)/XLH-.5)) + .98)
IF(I.LE.MLFGP.OR.I.GT.(MLFGP+MH)) TBAR = TBOLD
DTAV = TBAR - TBOLD
C
C      THERMAL CALCULATIONS FOR ROW I
C
CALL THRML(I)
C
C      NOW SOLVE FOR U(I,J)
C
CALL VELOC(I)
C
C      FINALLY COMPUTE V(I,J)
C
CALL RADV(I)
C
C      DETAILED PRINTOUT ACTIVATED
C
IF(IDETAI.EQ.0) GOTO 600
C
II = I - 1
CALL OUTROW(II)
600 CONTINUE
C
C      CALCULATIONS DONE FOR ROW I
C
100 CONTINUE
C
C      DISPLAY RESULTS FOR EXIT ROW
C
CALL OUTROW(M)
C
C      FIND TMAX IN LAST ROW
C      COMPUTE HOT CHANNEL FACTOR AND TEMP. PEAKING FACTOR
C
TMAX = T(M,1)
DO 400 J=2,N
IF(T(M,J).GT.TMAX) TMAX = T(M,J)
400 CONTINUE
PEAK = (TMAX - TINLT) / (TBAR - TINLT)
ETA = 1.
IF(F.EQ.1.) GOTO 477
ETA = (PEAK - 1.) / (F - 1.)
477 CONTINUE
CALL OUTST
C
C      ADVANCE REYNOLDS NUMBER
C
DELRE = DELRE1
IF(RE.GE.1500.) DELRE = DELRE2
IF(RE.GE.6000.) DELRE = DELRE3
RE = RE + DELRE
IF(RE.GT.REEND) STOP
700 CONTINUE
C
C      RETURN
C      END
C      SUBROUTINE TRNSNT
C
C      THIS SUBROUTINE IS THE TRANSIENT DRIVER
C
COMMON /KNOW/ XL, YL, N, M, RHO, CP, ALPHA, ALK, DX, DY, UBAR,
+ CF, EPSII, Q(40,40)
C
COMMON /DATIN/ IDETAI, NRE, MODE, ALINT, ALEXT, YDINT, YDEXT, REEND,
+ QNORM(40)
C
COMMON /AXIA/ XLLFGP, XLH, XLUFGP, MLFGP, MH, MUFGP,
+ DXLFGP, DXH, DXUFGP
C
COMMON /UNKN/ T(40,40), U(40,40), V(40,40), DEN(40,40)
C
COMMON /MORE/ AMU, F, FRIC, DTAV, YDE, BETA, TINLT, ALAM, RE
+ , DELTAT, PEAK, ETA, QBAR, TBAR, XKF, ICOND, IREDS
C
COMMON /LAMBDA/ ALAJ(40), YD(40), USTAR, UBOLD, TBOLD
C
COMMON /TIMER/ STIME, SDEL, CTIME, CTIMS, DTTRN, SEND, SDMAX, SDMIN
C
COMMON /PRINTR/ SPRNT, SNEXT, DTPRNT
C
DIMENSION TOLD(40,40)

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C
C
C      SAVE DATA FROM PREVIOUS TIME STEP;
C      GET INLET CONDITIONS AND COMPUTE EDDY DIFFUSIVITY
C
C      JNEG = 0
C      JPRNT = 0
700     CONTINUE
C
C      UBOLD = UBAR
C      CALL IN
C      CALL EPSILO
C      DO 500 J=1,N
C      DO 500 I=1,M
C      TOLD(I,J) = T(I,J)
500     CONTINUE
C
C      TBAR = TINLT
C
C      MARCH UP THE ASSEMBLY
C
C      DO 100 I=2,M
C
C      TBOLD = TBAR
C      DX = DXLFGP
C      IF(I.GT.MLFGP) DX = DXH
C      IF(I.GT.(MLFGP+MH)) DX = DXUFGP
C
C      THERMAL CALCULATIONS FOR ROW I
C
C      CALL THRML(I)
C
C      COMPUTE TBAR FOR ROW I, USING SIMPSON'S RULE
C
C      NLAST = N - 2
C      TBAR = 0.0
C      DO 121 J=3,NLAST
C      COEF = 2.0
C      JJ = J - 2*(J/2)
C      IF(JJ.NE.0) COEF = 4.0
C      TBAR = TBAR + COEF*T(I,J)*U(I,J)
121     CONTINUE
C      TBAR = TBAR / 3. / FLOAT(N-3)
C      TBAR = TBAR / UBAR
C      DTRN = TBAR - TBOLD
C
C      NOW SOLVE FOR U(I,J)
C
C      CALL VELOC(I)
C
C      CHECK FOR NEGATIVE VELOCITY
C      IF ANY U IS < 0 , SET IT TO 1.E-5 UNTIL IT RECOVERS
C      ALSO, PRINT A MESSAGE^
C
C      JNEG = 0
C      DO 261 JJ=3,NLAST
C      IF(U(I,JJ).LT.0.0) JNEG = JJ
C      IF(U(I,JJ).LT.0.0) U(I,JJ) = UBAR * 1.0E-05
261     CONTINUE
C      IF(JNEG.EQ.0) GOTO 262
C      IF(JPRNT.EQ.1) GOTO 262
C      JPRNT = 1
C      PRINT*," AT TIME = ",STIME," IN ROW ",I," AT NODE ",
+      JNEG," U WAS NEGATIVE "
262     CONTINUE
C      IF(JNEG.GT.0) GOTO 263
C      IF(JPRNT.EQ.0) GOTO 263
C      PRINT *," U RECOVERED AT TIME = ",STIME
C      JPRNT = 0
263     CONTINUE
C
C      FINALLY COMPUTE V(I,J)
C
C      CALL RADV(I)
C
C      IF(IDETA1.EQ.0) GOTO 600
C
C      DETAILED PRINTOUT ACTIVATED
C
C      II = I - 1
C      CALL OUTROW(II)
600     CONTINUE
C
C      END CALCULATIONS FOR SLICE I
C
100     CONTINUE
C
C      DTRN = TBAR - TINLT
C      IF(STIME.NE.SPRNT) GOTO 333
C
C      IT IS TIME TO PRODUCE GENERAL PRINTOUT
C
C      SPRNT = SPRNT + DTPRNT
C      SNEXT = SNEXT + DTPRNT
C
C      CALL OUTTRN
C      CALL OUTROW(M)
333     CONTINUE
C

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C      ADVANCE TIME AND RESET INTERVALS
C
      DTMAX = 0.0
      DO 555 I=1,M
      DO 555 J=1,N
      DT = ABS((TOLD(I,J)-T(I,J))/T(I,J))
      IF(DT.GT.DTMAX) DTMAX = DT
555    CONTINUE
      IF(DTMAX.LE.0.005) SDELTA = 2.0*SDELTA
      IF(DTMAX.GT.0.01) SDELTA = SDELTA/2.0
      IF(SDELTA.GT.SDMAX) SDELTA = SDMAX
      IF(SDELTA.LT.SDMIN) SDELTA = SDMIN
      IF((STIME+SDELTA).GT.SNEXA)
+     SDELTA = SNEXA - STIME
      STIME = STIME + SDELTA
      IF(STIME.GT.SEND) STOP
      GOTO 700
C
C      END
C      SUBROUTINE THRML(I)
C
C      THIS SUBROUTINE SOLVES THE ENERGY EQUATION
C
      COMMON /UNKN/ T(40,40),U(40,40),V(40,40),DEN(40,40)
C
      COMMON /WORK/ C(40), D(40), E(40), B(40)
C
      COMMON /KNOW/ XL,YL,N,M,RHO,CP,ALPHA,ALK,DX,DY,UBAR,CF,EPSII,
+     Q(40,40)
C
C      COMPUTE MATRIX TERMS FOR ROW I AND R.H.S.
C
      CALL SETUP(I)
C
C      SOLVE A(C,D,E) * T = B
C
      CALL TRID(C,D,E,B,N)
C
C      B NOW CONTAINS T(I,J)
C
      DO 100 J=1,N
      T(I,J) = B(J)
      DEN(I,J) = DENS(T(I,J))
100    CONTINUE
C
C      RETURN
C      END
C      SUBROUTINE RADV(I)
C
C      THIS SUBROUTINE COMPUTES THE TRANSVERSE VELOCITY
C      DISTRIBUTION USING THE CONTINUITY EQUATION
C
      COMMON /UNKN/ T(40,40),U(40,40),V(40,40),DEN(40,40)
C
      COMMON /KNOW/ XL,YL,N,M,RHO,CP,ALPHA,ALK,DX,DY,UBAR,CF,EPSII,
+     Q(40,40)
C
C      NLP = N - 2
C
      DO 100 J=2,NLP
      V(I,J) = (DEN(I,J-1)*V(I,J-1) - DY/2./DX*
+     (DEN(I,J-1)*U(I,J-1) + DEN(I,J)*U(I,J) -
+     DEN(I-1,J-1)*U(I-1,J-1) -
+     DEN(I-1,J)*U(I-1,J) ) ) /DEN(I,J)
100    CONTINUE
C
C      RETURN
C      END
C      SUBROUTINE TRID(A,B,C,R,N)
C
C      SUBROUTINE TO SOLVE A SET OF LINEAR EQUATIONS WITH
C      TRIDIAGONAL MATRIX OF COEFFICIENTS BY GAUSSIAN
C      ELIMINATION OF THE FORM :
C      MAT(A,B,C) * X = R
C      ORDER OF MAT IS N*N
C      RESULT IS RETURNED IN R
C
      DIMENSION A(40),B(40),C(40),R(40)
C
      DIMENSION X(40)
C
      IF(B(N).EQ.0.0) GOTO 400
      A(N) = A(N)/B(N)
      R(N) = R(N)/B(N)
C
      DO 100 I=2,N
      I1 = -I + N + 2
      BN = 1./(B(I1-1) - A(I1)*C(I1-1))
      IF(I1.GT.2) A(I1-1) = A(I1-1) * BN
      R(I1-1) = (R(I1-1)-C(I1-1)*R(I1)) * BN
100    CONTINUE
C
      X(1) = R(1)
      DO 200 I=2,N
      X(I) = R(I) - A(I)*X(I-1)

```

```

200          CONTINUE
C
DO 300 I=1,N
R(I) = X(I)
300          CONTINUE
C
RETURN
C
400          CONTINUE
PRINT *, " STOP CALCULATIONS BECAUSE ELEMENT IN DIAGONAL = 0."
C
STOP
C
END
SUBROUTINE EPSILO
C
      THIS SUBROUTINE COMPUTES THE THERMAL CONDUCTIVITY
      OF THE COOLANT AND SUBSEQUENTLY EDDY
      DIFFUSIVITY OF HEAT AND EFFECTIVE THERMAL DIFFUSIVITY
C
COMMON /KNOW/ XL,YL,N,M,RHO,CP,ALPHA,ALK,DX,DY,UBAR,
+ CF,EPSII, Q(40,40)
C
COMMON /MORE/ AMU,F,FRIC,DTAV,YDE,BETA,TINLT,ALAM,RE,
+ DELTAT,PEAK,ETA,QBAR,TBAR,XKF,ICOND,IREDIS
C
C
ALK = CON( TINLT )
C
IF(ICOND.EQ.0) ALK = 1.0E-05
C
EPSI = EPSII*( 1. - (1.-RE/7000.))**2 )
IF(RE.GE.7000.) EPSI = EPSII
EPSIH = UBAR * YDE * EPSI
C
ALPHA = EPSIH + XKF*ALK/RHO/CP
C
C
RETURN
END
SUBROUTINE SETUPU(I)
C
      THIS SUBROUTINE COMPUTES TERMS OF EQUATIONS
      A(C,D,E) * U = B
      FOR AXIAL VELOCITY COMPONENT U
C
COMMON /KNOW/ XL,YL,N,M,RHO,CP,ALPHA,ALK,DX,DY,UBAR,CF,
+ EPSII, Q(40,40)
C
COMMON /UNKN/ T(40,40),U(40,40),V(40,40),DEN(40,40)
C
COMMON /WORK/ C(40),D(40),E(40),B(40)
C
COMMON /NEWK/ A(40,40),WK(40),IWK(40)
C
COMMON /MORE/ AMU,F,FRIC,DTAV,YDE,BETA,TINLT,ALAM,RE,
+ DELTAT,PEAK,ETA,QBAR,TBAR,XKF,ICOND,IREDIS
C
COMMON /LAMBDA/ ALAJ(40),YD(40),USTAR,UBOLD,TBOLD
C
COMMON /TIMER/ STIME,SDELCT,CTIME,CTIMS,DTTRN,SEND,SDMAX,SDMIN
C
      COMPUTATIONS ARE PERFORMED FOR COOLANT ONLY;
      AT WALLS AND FOR STAGNANT COOLANT U IS = 0
C
      II = I
      DELTIM = CTIME / SDELCT
      NLAST = N - 2
      NFIRST = 3
      NU = N - 3
      RHOBAR = DENS( TBAR )
      ANU = AMU/RHOBAR/DY/DY
      DO 77 III=1,40
      DO 77 JJJ=1,40
      A(III,JJJ) = 0.0
77          CONTINUE
C
      FIRST ROW OF A(I,J)
C
      DO 1 J=NFIRST,NLAST
      A(1,J-2) = DEN(II,J)
      CONTINUE
C
      LAST COLUMN OF A(I,J)
C
      DO 2 JI=NFIRST,NLAST
      A(JI-1,NU) = 1. / DX
      CONTINUE
C
      COEFFICIENTS OF U(1,J)
C
      DO 100 JJ=NFIRST,NLAST
      JD = JJ - 1
      AL = ALAJ(JJ)
      IF(II.LE.2) AL = ALAM
      YDIAM = YD(JJ)
      IF(II.LE.2) YDIAM = YDE
      A(JD,JD-1) = U(II-1,JJ)/DX + 2.*ANU
      REIJ=DEN(II-1,JJ)*YDIAM*U(II-1,JJ)/AL/AMU

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F1J = FRC(RE1J)
A(JD,JD-1) = A(JD,JD-1) + U(II-1,JJ)*F1J/2./YDIAM/AL**2
+ DELTIM
100 CONTINUE
C
C COEFFICIENTS OF U(I,J-1)
C
NSUB = NFIRST + 1
DO 200 JJ=NSUB,NLAST
JC = JJ - 1
A(JC,JC-2) = -V(II-1,JJ-1)/2./DY - ANU
200 CONTINUE
C
C COEFFICIENTS OF U(I,J+1)
C
NSUP = NLAST - 1
DO 300 JJ=NFIRST,NSUP
JE = JJ + 1
A(JE,JE) = V(II-1,JJ+1)/2./DY - ANU
300 CONTINUE
C
C
C R. H. S.
C
B(1) = 0.0
DO 400 J=NFIRST,NLAST
B(1) = B(1) + DEN(II-1,J)*U(II-1,J)
400 CONTINUE
C
DO 500 KB = 2,NU
B(KB) = U(II-1,KB+1)*U(II-1,KB+1)/DX +
+ 9.81*BETA*(T(II,KB+1) - TINLT) +
+ DELTIM * U(II,KB+1)
500 CONTINUE
C
C AUGMENT MATRIX A WITH R.H.S. COLUMN
C
DO 600 KK=1,NU
A(KK,NU+1) = B(KK)
600 CONTINUE
C
C
C RETURN
C END
C SUBROUTINE VELOC(I)
C
C THIS SUBROUTINE COMPUTES AXIAL VELOCITY IN ROW I
C
COMMON /LAMBDA/ ALAJ(40),YD(40),USTAR,UBOLD,TBOLD
C
COMMON /MORE/ AMU,F,FRIC,DTAV,YDE,BETA,TINLT,ALAM,RE,
+ DELTAT,PEAK,ETA,QBAR,TBAR,XKF,ICOND,IREDS
C
COMMON /WORK/ C(40),D(40),E(40),B(40)
C
COMMON /NEWK/ A(40,40),WK(40),IWK(40)
C
COMMON /KNOW/ XL,YL,N,M,RHO,CP,ALPHA,ALK,DX,DY,UBAR,
+ CF,EPSIL,Q(40,40)
C
COMMON /UNKN/ T(40,40),U(40,40),V(40,40),DEN(40,40)
C
C
C USTAR = 0.0
C NLAST = N - 2
C
C IN THE ABSENCE OF BUOYANCY INDUCED FLOW REDISTRIBUTION,
C SET THE AXIAL VELOCITY TO THE INLET VELOCITY
C
IF(IREDS.EQ.0) GOTO 130
C
C COMPUTE TERMS OF THE EQUATION
C A(C,D,E) * U = B
C
CALL SETUPU(I)
C
C INVERT EQUATION
C
NU = N - 3
NUP1=NU + 1
CALL SOLVEU(A,B,NU,NUP1)
C
C NOW B CONTAINS U
C
DO 125 J=3,NLAST
U(1,J) = B(J-2)
125 CONTINUE
GOTO 132
C
130 CONTINUE
DO 131 J=1,N
U(1,J) = U(II-1,J)
131 CONTINUE
132 CONTINUE
C
C COMPUTE AVERAGE U FOR ROW I,USING SIMPSON'S RULE
C
USTAR = 0.0
DO 727 JJ=3,NLAST
COEF = 4.
JJL = JJ - 2 * (JJ/2)
IF(JJL.NE.0) COEF = 2.

```

```

727  USTAR = USTAR + COEF * U(I,JJ)
      CONTINUE
      USTAR = USTAR * DY / 3. / YL
C
C
      RETURN
      END
      SUBROUTINE SOLVEU(A,X,M,N)
C
C      SUBROUTINE TO INVERT EQUATION A(M,M)*X = A(N)
C      USING THE GAUSS-JORDAN REDUCTION METHOD
C
      DIMENSION A(40,40),X(40)
C
C      A IS NON-SINGULAR
C
      DO 100 K=1,M
      KP1 = K + 1
C
C      NORMALIZE BY PIVOT A(K,K)
C
      PIVOT = A(K,K)
      DO 200 J=KP1,N
      A(K,J) = A(K,J) / PIVOT
200  CONTINUE
      A(K,K) = 1.0
C
C      REDUCE OFF-DIAGONAL TERMS
C
      DO 100 I1=1,M
      IF(I1.EQ.K) GOTO 100
      IF(A(I1,K).EQ.0.0) GOTO 100
      DO 310 JJ=KP1,N
      A(I1,JJ) = A(I1,JJ) - A(I1,K)*A(K,JJ)
310  CONTINUE
      A(I1,K) = 0.0
100  CONTINUE
C
C      SOLUTION IS IN A(I,N)
C
      DO 400 K=1,M
      X(K) = A(K,N)
400  CONTINUE
C
C
      RETURN
      END
      SUBROUTINE IN
C
C      THIS SUBROUTINE INITIALIZES ARRAYS AND SETS INLET
C      BOUNDARY CONDITIONS
C
      COMMON /KNOW/ XL,YL,N,M,RHO,CP,ALPHA,ALK,DX,DY,UBAR,CF
+      , EPSII, Q(40,40)
C
      COMMON /UNKN/ T(40,40),U(40,40),V(40,40),DEN(40,40)
C
      COMMON /MORE/ AMU,F,FRIC,DTAV,YDE,BETA,TINLT,ALAM
+      ,RE,DELTAT,PEAK,ETA,QBAR,TBAR,XKF,ICOND,IREDS
C
      COMMON /LAMBDA/ ALAJ(40),YD(40),USTAR,UBOLD,TBOLD
C
      COMMON /TIMER/ STIME,SDELT,CTIME,CTIMS,DTTRN,SEND,SDMAX,SDMIN
C
      COMMON /SS/ UZERO,QZERO
C
      COMMON /STRC/ XMS,CS,DELTAW,XMC1,XMCN
C
      COMMON/AXIA/ XLLFGP,XLH,XLUFGP,MLFGP,MH,MUFGP,
+      DXLFGP,DXH,DXUFGP
C
      COMMON /DATIN/ IDETA1,NRE,MODE,ALINT,ALEXT,YDINT,YDEXT,REEND,
+      QNORM(40)
C
      COMMON /TRNOP/ JTRN
C
C
C
      DO 80 J=1,N
      DEN(1,J) = RHO
80  CONTINUE
C
      IF(MODE.EQ.1) GOTO 91
      UBAR = RE * ALAM * AMU / RHO / YDE
      GOTO 92
91  RE = UBAR * RHO * YDE / ALAM / AMU
92  CONTINUE
      FRIC = FRC( RE )
C
C      FOR TRANSIENT CALCULATIONS COMPUTE ADVANCED UBAR AND POWER
C
      URT = UR(STIME)
      UBAR = UZERO * URT
      RE = UBAR*RHO*YDE/ALAM/AMU
C
      CTIME = 0.0
      IF(STIME.GT.0.0) CTIME = 1.0
      CTIMS = 0.0
      IF(STIME.GT.0.0) CTIMS = 1.0
      P = POVERW(STIME)
C

```

```

C      QBAR = P * QZERO* UBAR / UZERO
C
C      U(I,J) AT INLET IS = 0 AT WALLS AND FOR STAGNANT COOLANT,
C      CALCULATE U(I,J) SO THAT THE AVERAGE VALUE BECOMES UBAR
C      USING THE SIMPSON'S RULE
C      IN OTHER NODES
C
C      DO 100 I=1,M
C      DO 100 J = 1,N
C      V(I,J) = 0.0
C      U(I,J) = UBAR*3.*FLOAT(N-3)/(6.*FLOAT((N-3)/2)-2.)
C      IF(J.EQ.1.OR.J.EQ.N) U(I,J) = 0.0
C      IF(J.EQ.2.OR.J.EQ.(N-1)) U(I,J) = 0.0
100    CONTINUE
C      DO 200 I=1,M
C      DO 200 J=1,N
C      IF(STIME.GT.0.0) GOTO 201
C      T(I,J) = TINLT
201    CONTINUE
200    CONTINUE
C      X = 0.0
C      DO 500 I=1,M
C      QX = 1.4*COS(.8738*3.14156*((X-XLLFGP)/XLH-.5))
C      IF(I.LE.MLFGP.OR.I.GT.(MLFGP+MH)) QX = 0.0
C      DX = DXLFGP
C      IF(I.GE.MLFGP) DX = DXH
C      X = X + DX
C      DO 600 J=1,N
C      Q(I,J) = QBAR * QNORM(J) * QX
C      IF(J.EQ.1.OR.J.EQ.N) Q(I,J) = 0.0
C      IF(J.EQ.2.OR.J.EQ.(N-1)) Q(I,J) = 0.0
600    CONTINUE
500    CONTINUE
C
C      RETURN
C      END
C      SUBROUTINE SETUP(I)
C
C      THIS SUBROUTINE DEFINES, FOR I=2 TO M, THE TRIDIAGONAL
C      TERMS OF A*T = B
C
C      B = R.H.S.
C      C = SUBDIAGONAL TERMS OF A
C      D = DIAGONAL " "
C      E = SUPERDIAGONAL " "
C
C      COMMON /KNOW/ XL,YL,N,M,RHO,CP,ALPHA,ALK,DX,DY,UBAR,CF,
+      EPSII, Q(40,40)
C
C      COMMON /MORE/ AMU,F,FRIC,DTAV,YDE,BETA,TINLT,ALAM,RE,
+      DELTAT,PEAK,ETA,QBAR,TBAR,XKF,ICOND,IRED5
C
C      COMMON /UNKN/ T(40,40),U(40,40),V(40,40),DEN(40,40)
C
C      COMMON /WORK/ C(40),D(40),E(40),B(40)
C
C      COMMON /TIMER/ STIME,SDELT,CTIME,CTIMS,DTRN,SEND,SDMAX,SDMIN
C
C      COMMON /STRC/ XMS,CS,DELTAW,XMC1,XMCN
C
C      COMMON /LAMBDA/ ALAJ(40),YD(40),USTAR,UBOLD,TBOLD
C
C      I1 = 1
C      CALPHA = ALPHA / DY / DY
C      DTIME = CTIME / SDELT
C      AREA = SQRT(3.)*YL*DX
C      AC2 = ACN1 = AREA
C      LSUB = N - 2
C
C      DEFINE DIAGONAL TERMS
C
C      DO 100 L=3,LSUB
100    D(L) = U(I1-1,L)/DX + 2.*CALPHA + DTIME
C
C      THESE CALCULATIONS ARE FOR THE WALL AND THE
C      EXTERNAL COOLANT
C
C      AKW1 = AKW( T(I1,2) )
C      AKWN = AKW( T(I1,N-1) )
C      C1 = CAP(T(I1,1))
C      CN = CAP(T(I1,N))
C      D(1) = XMC1*DTIME*C1 + AKW1*AREA/DELTAW
C      D(N) = XMCN*DTIME*CN + AKWN*AREA/DELTAW
C
C      AK2 = CON(T(I1,2))
C      AKN1 = CON(T(I1,N-1))
C
C      D(2) = XMS*CS*DTIME + AKW1*AREA/DELTAW +
+      AK2 * AC2 / DY
C      D(N-1) = XMS*CS*DTIME + AKWN*AREA/DELTAW +
+      AKN1 * ACN1 / DY
C
C      SUBDIAGONAL TERMS
C
C      DO 200 L=3,LSUB
200    C(L) = -(V(I1-1,L)/2./DY + CALPHA )
C      C(N) = -AKWN * AREA / DELTAW
C      E(2) = - AK2 * AC2 / DY
C      C(N-1) = - AKN1 * ACN1 / DY

```

```

C
C SUPERDIAGONAL TERMS
C
E(1) = -AKW1 * AREA / DELTAW
C(2) = -AKW1 * AREA / DELTAW
DO 300 L=3,LSUB
300 E(L) = V(II-1,L)/2./DY - CALPHA
E(N-1) = -AKWN * AREA / DELTAW
C
C R.H.S.
C
C DO 400 L=3,LSUB
B(L) = T(II-1,L) * U(II-1,L) / DX + Q(II,L) / RHO / CP
+ DTIME * T(II,L)
400 CONTINUE
B(1) = XMC1*C1*DTIME*T(II,1)
B(2) = XMS*CS*DTIME*T(II,2)
B(N-1) = XMS*CS*DTIME*T(II,N-1)
B(N) = XMCN*CN*DTIME*T(II,N)
C
C RETURN
END
SUBROUTINE OUTTRN
C
C TRANSIENT OUTPUT ROUTINE
C
COMMON /KNOW/ XL,YL,N,M,RHO,CP,ALPHA,ALK,DX,DY,UBAR,CF,
+ EPSII,Q(40,40)
C
COMMON /UNKN/ T(40,40),U(40,40),V(40,40),DEN(40,40)
C
COMMON /MORE/ AMU,F,FRIC,DTAV,YDE,BETA,TINLT,ALAM,RE,
+ DELTAT,PEAK,ETA,QBAR,TBAR,XKF,ICOND,IREDS
C
COMMON /TIMER/ STIME,SDELT,CTIME,CTIMS,DTTRN,SEND,SDMAX,SDMIN
C
COMMON /SS/ UZERO,QZERO
C
C
C WRITE(2,99)
99 FORMAT(////)
C
WRITE(2,666) STIME,SDELT
666 FORMAT(1H,10(1H*)," TIME = ",2F10.4)
WRITE(2,4) RE
4 FORMAT(" RE = ",E12.5)
WRITE(2,7) TINLT
7 FORMAT(//," INLET TEMPERATURE = ",F10.3)
WRITE(2,8) DTTRN
8 FORMAT(" AVERAGE TEMPERATURE RISE = ",F10.3)
WRITE(2,10) TBAR
10 FORMAT(" AVERAGE EXIT TEMPERATURE = ",F10.3)
WRITE(2,9) QBAR
9 FORMAT(" AVERAGE POWER DENSITY = ",E12.5)
C
C TAPE9B STORES DATA FOR PLOTTING
C
UU = UBAR / UZERO
QQ = QBAR / QZERO
QTOU = QQ / UU
1 WRITE(98,1) STIME,UU,DTTRN,QQ,QTOU
FORMAT(5E13.6)
DO 300 L=1,M
300 WRITE(98,1) (T(L,J),J=1,N)
CONTINUE
C
RETURN
END
SUBROUTINE OUTROW(1)
C
C THIS SUBROUTINE DISPLAYS RESULTS FOR ROW 1
C
COMMON /TRNOP/ JTRN
C
COMMON /KNOW/ XL,YL,N,M,RHO,CP,ALPHA,ALK,DX,DY,UBAR,CF,
+ EPSII,Q(40,40)
C
COMMON /UNKN/ T(40,40),U(40,40),V(40,40),DEN(40,40)
C
COMMON /MORE/ AMU,F,FRIC,DTAV,YDE,BETA,TINLT,ALAM,RE,DELTAT,
+ PEAK,ETA,QBAR,TBAR,XKF,ICOND,IREDS
C
COMMON /TIMER/ STIME,SDELT,CTIME,CTIMS,DTTRN,SEND,SDMAX,SDMIN
C
C DIMENSION J(40)
C
IF(JTRN.EQ.0) DTTRN = DTAV
WRITE(2,90)
90 FORMAT(//)
DO 100 JJ=1,N
100 J(JJ) = JJ
TBARP = TBAR - DTTRN
IF(I.EQ.M) TBARP = TBAR
TCOMP = TBAR
GIN = RHO*UBAR
GOUT = DENS(TBARP) * UBAR
WRITE(2,76) J,TCOMP,GIN,GOUT
76 FORMAT(" IN ROW ",15," AVER TEMP = ",F12.4,

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+ " MASS FLOW IN = ",E13.6," MASS FLOW OUT = ",
+ E13.6)
DO 220 JJ=1,N
Q(I,JJ) = Q(I,JJ)/QBAR
220 CONTINUE
JMANY = N / 6
JP5 = 6
J1 = -5
DO 200 JJ=1,JMANY
WRITE(2,91)
91 FORMAT( / )
J1 = J1 + JP5
J5 = J1 + 5
WRITE(2,77) (J(JJJ),JJJ=J1,J5)
77 FORMAT(" J ",6I20)
WRITE(2,78) (U(I,JU),JU=J1,J5)
78 FORMAT(" U ",6E20.4)
WRITE(2,79) (V(I,JV),JV=J1,J5)
79 FORMAT(" V ",6E20.4)
WRITE(2,80) (T(I,JT),JT=J1,J5)
80 FORMAT(" T ",6F19.3)
WRITE(2,81) (Q(I,JQ),JQ=J1,J5)
81 FORMAT(" Q ",6E20.4)
200 CONTINUE
C
ILEFT = JMANY*6 - N
IF(ILEFT.EQ.0) GOTO 300
JMP1 = JMANY*6 + 1
WRITE(2,77) (J(JJJ),JJJ=JMP1,N)
WRITE(2,78) (U(I,JU),JU=JMP1,N)
WRITE(2,79) (V(I,JV),JV=JMP1,N)
WRITE(2,80) (T(I,JT),JT=JMP1,N)
WRITE(2,81) (Q(I,JQ),JQ=JMP1,N)
300 CONTINUE
DO 400 JJ=1,N
Q(I,JJ) = Q(I,JJ)*QBAR
400 CONTINUE
C
C
RETURN
END
SUBROUTINE OUTIN
C
C THIS SUBROUTINE PRINTS INPUT DATA
C
C
COMMON /TRNOP/ JTRN
C
COMMON /STRC/ XMS,CS,DELTAW,XMC1,XMCN
C
COMMON /TIMER/ STIME,SDELTA,CTIME,CTIMS,DTRN,SEND,SDMAX,SDMIN
C
COMMON /PRINTR/ SPRNT,SNEXT,DTPRNT
C
COMMON /KNOW/ XL,YL,N,M,RHO,CP,ALPHA,ALK,DX,DY,UBAR,
+ CF,EPS11,Q(40,40)
C
COMMON /AXIA/ XLLFGP,XLH,XLUFGP,MLFGP,MH,MUFGP,
+ DXLFGP,DXH,DXUFGP
C
COMMON /DATIN/ IDETA1,NRE,MODE,ALINT,ALEXT,YDINT,YDEXT,REEND,
+ QNORM(40)
C
COMMON /MORE/ AMU,F,FRIC,DTAV,YDE,BETA,TINLT,ALAM,RE,
+ DELTAT,PEAK,ETA,QBAR,TBAR,XKF,ICOND,IREDS
COMMON /U/ TIME(40),URAT(40),NTIME
C
COMMON /P/ TIM(40),PW(40),NTIM
C
C
WRITE(2,100)
100 FORMAT(1H1,10(2H*),18H OPTIONS SELECTED ,10(2H*))
WRITE(2,200)
200 FORMAT(10(1H),/)
WRITE(2,101) JTRN,IREDS,ICOND,IDETA1
101 FORMAT(20H TRANSIENT OPTION = ,15,/,
+ 23H FLOW REDISTRIBUTION = ,15,/,
+ 13H CONDUCTION = ,15,/,
+ 21H DETAILED PRINTOUT = ,15,/)
WRITE(2,102)
102 FORMAT(10(2H*),16H GEOMETRIC DATA ,10(2H*),/)
WRITE(2,103) N,MLFGP,MH,MUFGP
103 FORMAT(26H NUMBER OF RADIAL NODES = ,15,/,
+ 24H NUMBER OF LF GP NODES = ,15,/,
+ 26H NUMBER OF HEATED NODES = ,15,/,
+ 24H NUMBER OF UF GP NODES = ,15,/)
WRITE(2,1) XLLFGP,XLH,XLUFGP,YL
1 FORMAT(28H LOWER FISSION GAS PLENUM = ,F10.4,4H M /,
+ 26H ASSEMBLY HEATED LENGTH = ,F10.4,4H M /,
+ 28H UPPER FISSION GAS PLENUM = ,F10.4,4H M /,
+ 16H FLAT-TO-FLAT = ,F10.4,4H M /)
WRITE(2,2) ALAM,ALINT,ALEXT
2 FORMAT(36H LAMBDA AVERAGE, INTERIOR, EXTERIOR = ,
+ 3F10.4)
WRITE(2,3) YDE,YDINT,YDEXT
3 FORMAT(48H HYDRAULIC DIAMETER AVERAGE, INTERIOR, EXTERIOR =
+ 3E15.4,4H M )
WRITE(2,4)
4 FORMAT(///,10(2H*),20H INITIAL CONDITIONS ,10(2H*))

```



```

WRITE(98,1) (T(L,J),J=1,N)
WRITE(98,1) (U(L,JU),JU=1,N)
300 CONTINUE
C
RETURN
END
FUNCTION UR(T)
C
THIS FUNCTION INTERPOLATES BETWEEN TIMES TO
RETURN NORMALIZED VELOCITY AT TIME T
C
COMMON /U/ TIME(40),URAT(40),NTIME
C
ITIME = 0
DO 100 IT = 1,NTIME
ITIME = ITIME + 1
IF(T.LT.TIME(IT)) GOTO 101
100 CONTINUE
101 IF(ITIME.EQ.NTIME) GOTO 200
UR = URAT(ITIME-1) - (T-TIME(ITIME-1))*(URAT(ITIME)-
+ URAT(ITIME-1))/(TIME(ITIME-1)-TIME(ITIME))
GOTO 210
200 UR = URAT(ITIME)
210 CONTINUE
RETURN
C
END
FUNCTION POWERW(T)
C
THIS FUNCTION INTERPOLATES BETWEEN TIMES TO
RETURN POWER/FLOW RATIO AT TIME T
C
COMMON /P/ TIM(40),PW(40),NTIM
C
ITIME = 0
DO 100 IT = 1,NTIM
ITIME = ITIME + 1
IF(T.LT.TIM(IT)) GOTO 101
100 CONTINUE
101 IF(ITIME.EQ.NTIM) GOTO 200
POVERW = PW(ITIME-1) - (T-TIM(ITIME-1))*(PW(ITIME) -
+ PW(ITIME-1))/(TIM(ITIME-1)-TIM(ITIME))
GOTO 210
200 POWERW = PW(ITIME)
210 CONTINUE
RETURN
END
FUNCTION AKW(TT)
C
THERMAL CONDUCTIVITY FOR WALL
C
T = TT
AKW = 9.01748 + 1.62997E-02 * T -
+ 4.80329E-06 * T * T + 2.18422E-09 * T ** 3
C
RETURN
END
FUNCTION CON(T)
C
SODIUM THERMAL CONDUCTIVITY
C
CON = 109.7 + (-6.4499E-02 + 1.1728E-05*T)
+ * T
C
RETURN
END
FUNCTION VISC(T)
C
SODIUM VISCOSITY
C
RHS = -2.4892 + 220.65/T - .21389*ALOG(T)
VISC = 10. ** RHS
C
RETURN
END
FUNCTION DENS(T)
C
SODIUM DENSITY
C
DENS = 1011.597 - .22051*T - 1.922243E-5*T*T +
+ 5.63769E-9*T*T*T
C
RETURN
END
FUNCTION CAP(T)
C
SODIUM HEAT CAPACITY
C
CAP = 1630.22 - .83354*T + 4.62838E-4*T*T
C
RETURN
END
FUNCTION FRC(RE)
C
FRICITION FACTOR COEFFICIENT FOR 61-PIN

```

```
C          BLANKET ASSEMBLY
C
FRC = 110. / RE
IF (RE.LT.400.) RETURN
IF (RE.GT.5000.0) GOTO 1
PSI = (RE - 400.) / 4600.
FRC = 110./RE*SQRT(1.-PSI) + .48/(RE**.25)*SQRT(PSI)
RETURN
1  CONTINUE
C  FRC = .48 * RE**(-.25)
C
RETURN
END
```

## B.2 ALPHABETICAL DESCRIPTION OF THE ASSOCIATED SUBROUTINES

- AKW: This function computes the thermal conductivity of structure.  
Parameter: T Temperature, K
- CAP: This function computes the Heat Capacity of Sodium.  
Parameter: T Temperature, K
- CON: This function computes the Thermal Conductivity of Sodium.  
Parameter: T Temperature, K
- DENS: This function computes Sodium Density.  
Parameter: T Temperature, K
- EPSILO: This subroutine computes Eddy Diffusivity of Heat.
- FRC: This function computes Sodium Frictional Coefficient.  
Parameter: Re Reynolds number
- IN: This subroutine initializes arrays at channel inlet at each iteration.
- OUTIN: This subroutine displays input data.
- OUTROW: This subroutine displays results for axial slice I.  
Parameter: I axial slice number
- OUTST: Output subroutine for steady-state calculations.
- OUTTRN: Output routine for transient calculations.

POVERW: This function interpolates between times to compute Power/Flow ratio at time t.  
Parameter: T time, s

RADV: This subroutine computes the transverse velocity components from the continuity equation.  
Parameter: I axial slice number

SETUP: This subroutine computes tridiagonal terms of the coefficient matrix of the energy equation for temperature calculations.  
Parameter: I axial slice number

SETUPU: This subroutine computes terms of the coefficient matrix of the axial momentum equation.  
Parameter: I axial slice number

SOLVEU: Matrix inverter subroutine, uses the Gauss-Jordan reduction method  
Parameters: A augmented matrix of coefficients and R.H.S.  
X vector of results  
M leading dimension of A  
N M+1

START: This subroutine reads input data and computes constants and properties.

STST: Driver for steady-state calculations.

THRML: This subroutine drives the energy calculations for axial slice I.  
Parameter: I axial slice number

TRID: Tridiagonal matrix inverter; uses the Gaussian elimination method.  
Parameters: A vector of subdiagonal terms  
B vector of diagonal terms  
C vector of superdiagonal terms  
R R.H.S., on output, result  
N leading dimension of vectors

TRNSNT: This subroutine drives the transient calculations.

TWIST: Main driver.

Parameters defining files:

TAPE1 (input)

TAPE2 (output)

TAPE98 (storage file)

UR: This function interpolates between times to compute normalized inlet average velocity at time t.

Parameters: t time, s

VELOC: This subroutine computes axial velocities from the momentum equation.

Parameter: I axial slice number

VISC: This function computes viscosity of sodium.

Parameter: T Temperature, K

### B.3 DESCRIPTION OF INPUT

Input data resides in file "TAPE1", and is organized in 12 NAMELIST groups as follows:

NAMELIST/OPTION/	JTRN	- Integer variable which defines whether steady-state or transient* calculations will be performed; { 0, steady-state { 1, transient
NAMELIST/CONTR/	IDETAI	- Integer control variable for printout { 0, Summary printouts will be produced at end of each calculational step. { 1, Detailed printout will be produced for each axial slice at the end of each calculational step.
	NRE	- total number of Reynolds numbers upon which to perform calculations (For steady-state calculations only)
	IREDS	- Integer control variable for flow redistribution option { 0, flow redistribution ignored { 1, flow redistribution included
	ICOND	- Integer control variable for conduction model { 0, conduction and mixing effects ignored { 1, conduction and mixing effects included
	N	- Number of radial nodes; N should be equal to $N_{\text{heated}} + 2$ (walls nodes) + 2 (stagnant coolant nodes) and $N_{\text{heated}}$ must be <u>odd</u> number; a maximum of 39 nodes is allowed in the present version of the code.

\*A steady state initialization is always performed prior to transient calculations.

NAMelist/CONTR/            MLFGP - Number of axial Lower Fission Gas Plenum nodes  
                              MH     - Number of axial heated nodes  
                              MUFGP - Number of axial Upper Fission Gas Plenum nodes

NOTE: The sum MLFGP + MH + MUFGP should not exceed 40 in the present version of the code.

NAMelist/GEOM1/            XLLFGP - Lower Fission Gas Plenum Length, m  
                              XLH     - Heated channel length, m  
                              XLUFGP - Upper Fission Gas Plenum Length, m  
                              YL       - Subassembly flat-to-flat, m

NAMelist/GEOM2/            ALAM   - Average porosity  
                              ALINT   - Porosity, interior region of the assembly  
                              ALEXT   - Porosity, wall region of the assembly

NAMelist/GEOM3/            YDE    - Average hydraulic diameter, m  
                              YDINT   - Hydraulic diameter, interior region of the assembly, m  
                              YDEXT   - Hydraulic diameter, wall region of the assembly, m

NAMelist/INITC/            Re     - Initial Reynolds number  
                              TINLT   - Inlet temperature, K  
                              EPSII   - Value of  $\epsilon_{\infty}^*$   
                              QBAR     - Total power in the assembly, watts  
                              DELTAT - Total assembly average  $\Delta\hat{T}$ , K  
                              UBAR     - Average inlet velocity, m/s

NOTE: Of the parameters RE and UBAR, one needs to be defined, the other should be set to 0.0, and the program will compute it; similarly for the parameters QBAR and DELTAT, where only one needs to be specified.

NAMelist/PROP/            BETA        -   Thermal expansion coefficient, K<sup>-1</sup>

NAMelist/RADQ/            QNOR        -   Radial power distribution in each axial slice; up to  
40 values allowed in the present version of the code.

NOTE:                      Need not be normalized; its real dimensions are QNORM(N), and the first and  
last two values should be equal to zero (wall and stagnant coolant nodes)

NAMelist/TIMES/            SDMIN       -   Minimum time step, s  
SDMAX       -   Maximum time step, s  
DTPRNT      -   Printout time, s  
SEND        -   Duration of transient (t<sub>max</sub>), s

NAMelist/WALL/            XMS        -   mass of wall on either side of the subassembly, Kg  
CS           -   heat capacity of duct wall, J/(Kg·K)  
DELTAW      -   wall thickness, m  
XMC1        -   mass of stagnant coolant adjacent to the left wall, Kg  
XMCN        -   mass of stagnant coolant adjacent to the right wall, Kg

LIST/UTIME/                Transient normalized velocity forcing function  
TIME, URAT -   values of normalized inlet average velocity at time,  
TIME

NOTE:                      Up to 40 values are allowed in the present version of the code. First all  
the values of TIME are inputted, then all values of URAT. An additional in-  
put is required, equal to -1, at the end of the TIME table.

NAMelist/PTIME/            Transient power/flow ratio forcing function  
TIM, PW     -   values of normalized power to flow ratio at time, TIM.

NOTE:                      See above note.



## B.4 DESCRIPTION OF OUTPUT

### a) With IDETAI = 1 (Detailed Printout)

For each axial slice in calculational step the following results will be displayed:

- i) Slice number, average temperature in the slice, mass flow rate in, and mass flow rate out of the slice.
- ii) For each radial node  $j$  in slice  $i$ 
  - node number  $j$
  - axial velocity m/s
  - radial velocity m/s
  - temperature K
  - Heat flux  $Q_{i,j}/\hat{Q}$

Summary printout will also be displayed at the end of calculation step.

### b) With IDETAI = 0 (Summary Printout)

- i) Same as a), only for exit of channel, and in addition:
- ii) for steady-state calculations:
  - F peak to average power
  - Re Reynolds number
  - PEAK temperature peaking factor
  - ETA hot channel factor
  - TINLT inlet temperature
  - DELTAT average channel temperature rise
  - QBAR average power density per axial slice
  - GRA Grashof number,  $Gr^*$
  - Ri Richardson number
  - GRRE =  $Gr^*/Re^*$
  - BUOY buoyancy parameter =  $2 Ri/f$

in addition, the following data is stored on file Tape 98 (for plotting purposes):

RE, ETA, BUOY, PEAK, DELTAT (5E13.6)  
 $T_{i,j}, U_{i,j}$  ( $i = 1, M; j = 1, N$ ) (5E13.6)

iib) for transient calculations:

STIME, SDEL	current time and time step	s
RE	Reynolds number	
TINLT	inlet temperature	K
DTRN	average temperature rise from inlet to exit of assembly	K
TBAR	average exit temperature	K
QBAR	average power density	watts/m <sup>2</sup>

while on file TAPE 98 (plotting file) the following results are stored:

Time, normalized velocity, average temperature rise,  
normalized power and power-to-flow ratio (5E13.6)

$T_{i,j}$  ( $i = M; j = 1,N$ ) (5E13.6)

c) Regardless of the value of 1DETAI, all input data is displayed at the beginning of execution.

Sample Output:

\*\*\*\*\* OPTIONS SELECTED \*\*\*\*\*

TRANSIENT OPTION = 1  
FLOW REDISTRIBUTION = 1  
CONDUCTION = 1  
DETAILED PRINTOUT = 0

\*\*\*\*\* GEOMETRIC DATA \*\*\*\*\*

NUMBER OF RADIAL NODES = 37  
NUMBER OF LFDP NODES = 0  
NUMBER OF HEATED NODES = 20  
NUMBER OF UFGP NODES = 0

LOWER FISSION GAS PLENUM = 0.0000 M  
ASSEMBLY HEATED LENGTH = 1.1430 M  
UPPER FISSION GAS PLENUM = 0.0000 M  
FLAT-TO-FLAT = .1140 M

LAMBDA AVERAGE, INTERIOR, EXTERIOR = .2567 .2214 .3674  
HYDRAULIC DIAMETER AVERAGE, INTERIOR, EXTERIOR = .3728E-02 .3516E-02 .7181E-02 M

\*\*\*\*\* INITIAL CONDITIONS \*\*\*\*\*

REYNOLD NUMBER = .128681E+05  
FINAL REYNOLD NUMBER = 0.  
INLET TEMPERATURE = 588.000 K  
OVERALL ASSEMBLY DELTAT = 101.989 K  
AVERAGE POWER = .342032E+08 WATTS  
AVERAGE INLET VELOCITY = .336400E+00 M/S

\*\*\*\*\* PROPERTIES \*\*\*\*\*

EPSILON STAR INFINITY = .300000E-01  
BETA = .291000E-03

\*\*\*\*\* NORMALIZED RADIAL POWER DENSITY \*\*\*\*\*

0.	0.	.1000E+01	.1000E+01	.1000E+01	.1000E+01	.1000E+01	.1000E+01	.1000E+01	.1000E+01	.1000E+01
.1000E+01	.1000E+01	.1000E+01	.1000E+01	.1000E+01	.1000E+01	.1000E+01	.1000E+01	.1000E+01	.1000E+01	.1000E+01
.1000E+01	.1000E+01	.1000E+01	.1000E+01	.1000E+01	.1000E+01	.1000E+01	.1000E+01	.1000E+01	.1000E+01	.1000E+01
.1000E+01	.1000E+01	.1000E+01	.1000E+01	.1000E+01	0.	0.	.1000E+01	.1000E+01	.1000E+01	.1000E+01

\*\*\*\*\* WALL PROPERTIES \*\*\*\*\*

MASS OF STRUCTURE = .316 KG PER SIDE  
MASS OF SURROUNDING SODIUM = 22.000 22.000 KG AT LEFT AND RIGHT  
WALL THICKNESS = .3810E-03 M  
STRUCTURE CP = 560.000 J/KG-K

\*\*\*\*\* TRANSIENT DATA \*\*\*\*\*

FINAL TIME = 600.00 SEC  
MINIMUM AND MAXIMUM TIME STEPS = .1000E-04 .1000E+01 SEC  
PRINTOUT EVERY = 5.00 SEC

FORCING FUNCTIONS

NORMALIZED VELOCITY VS TIME

1.0000	0.0000
.9700	10.0000
.7000	18.0000
.5400	21.0000
.4000	25.0000
.2100	40.0000
.1300	50.0000
.0650	70.0000
.0620	100.0000
.0700	400.0000
.0700	1000.0000

POWER TO FLOW RATIO VS TIME

1.0000	0.0000
1.0000	7.0000
.2000	15.0000
.1800	17.0000
.1500	19.0000
.2000	25.0000
.5000	50.0000
.9800	75.0000
1.1200	95.0000
1.1200	100.0000
1.0000	250.0000
1.0000	1000.0000

\*\*\*\*\* TIME = 0.0000 .0313  
 RE = .12868E+05

INLET TEMPERATURE = 588.000  
 AVERAGE TEMPERATURE RISE = 103.721  
 AVERAGE EXIT TEMPERATURE = 691.721  
 AVERAGE POWER DENSITY = .34203E+08

IN ROW 20 AVER TEMP = 691.7213 MASS FLOW IN = .294833E+03 MASS FLOW OUT = .286523E+03

J		1		2		3		4		5		6
U	0.		0.		.8054E+00		.6130E+00		.4413E+00		.2927E+00	
V	0.		0.		.7102E-06		.1741E-05		.1564E-05		.3164E-05	
T	658.000		658.000		658.000		662.857		670.143		677.784	
Q	0.		0.		.2757E+00		.2757E+00		.2757E+00		.2757E+00	
J		7		8		9		10		11		12
U	.2928E+00		.2929E+00		.2931E+00		.2932E+00		.2933E+00		.2933E+00	
V	-.1347E-05		-.3121E-05		-.4661E-05		-.5768E-05		-.6381E-05		-.6522E-05	
T	684.484		690.032		694.429		697.791		700.285		702.086	
Q	.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00	
J		13		14		15		16		17		18
U	.2934E+00		.2934E+00		.2934E+00		.2934E+00		.2934E+00		.2934E+00	
V	-.6251E-05		-.5646E-05		-.4785E-05		-.3735E-05		-.2556E-05		-.1297E-05	
T	703.356		704.230		704.817		705.199		705.434		705.560	
Q	.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00	
J		19		20		21		22		23		24
U	.2934E+00		.2934E+00		.2934E+00		.2934E+00		.2934E+00		.2934E+00	
V	.2769E-13		.1297E-05		.2556E-05		.3735E-05		.4785E-05		.5646E-05	
T	705.599		705.560		705.434		705.199		704.817		704.230	
Q	.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00	
J		25		26		27		28		29		30
U	.2934E+00		.2933E+00		.2933E+00		.2932E+00		.2931E+00		.2929E+00	
V	.6251E-05		.6522E-05		.6381E-05		.5768E-05		.4661E-05		.3121E-05	
T	703.356		702.086		700.285		697.791		694.429		690.032	
Q	.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00	
J		31		32		33		34		35		36
U	.2928E+00		.2927E+00		.4413E+00		.6130E+00		.8054E+00		0.	
V	.1347E-05		-.3164E-05		-.1564E-05		-.1741E-05		-.7102E-06		0.	
T	684.484		677.784		670.143		662.857		658.000		658.000	
Q	.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00		0.	
J		37										
U	0.											
V	0.											
T	658.000											
Q	0.											

\*\*\*\*\* TIME = 5.0000 .0156  
 RE = .12675E+05

INLET TEMPERATURE = 588.000  
 AVERAGE TEMPERATURE RISE = 103.260  
 AVERAGE EXIT TEMPERATURE = 691.260  
 AVERAGE POWER DENSITY = .33690E+08

IN ROW 20 AVER TEMP = 691.2598 MASS FLOW IN = .290411E+03 MASS FLOW OUT = .282262E+03

J		1		2		3		4		5		6
U	0.		0.		.4733E+00		.4430E+00		.3964E+00		.3274E+00	
V	0.		0.		.4842E-06		.1088E-05		.6929E-06		-.6402E-06	
T	658.011		658.090		658.417		663.140		670.221		677.678	
Q	0.		0.		.2757E+00		.2757E+00		.2757E+00		.2757E+00	
J		7		8		9		10		11		12
U	.3274E+00		.3275E+00		.3276E+00		.3276E+00		.3276E+00		.3277E+00	
V	-.2643E-05		-.4994E-05		-.7084E-05		-.8601E-05		-.9435E-05		-.9603E-05	
T	684.229		689.646		693.932		697.202		699.623		701.367	
Q	.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00	
J		13		14		15		16		17		18
U	.3277E+00		.3277E+00		.3277E+00		.3277E+00		.3277E+00		.3277E+00	
V	-.9186E-05		-.8289E-05		-.7020E-05		-.5479E-05		-.3749E-05		-.1	

T	702.593	703.435	703.999	704.365	704.589	704.709
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00
J	19	20	21	22	23	24
U	.3277E+00	.3277E+00	.3277E+00	.3277E+00	.3277E+00	.3277E+00
V	.3321E-13	.1903E-05	.3749E-05	.5479E-05	.7020E-05	.8289E-05
T	704.747	704.709	704.589	704.365	703.999	703.435
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00
J	25	26	27	28	29	30
U	.3277E+00	.3277E+00	.3276E+00	.3276E+00	.3276E+00	.3275E+00
V	.9186E-05	.9603E-05	.9435E-05	.8601E-05	.7084E-05	.4994E-05
T	702.593	701.367	699.623	697.202	693.932	689.646
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00
J	31	32	33	34	35	36
U	.3274E+00	.3274E+00	.3964E+00	.4430E+00	.4733E+00	0.
V	.2643E-05	.6402E-06	-.6929E-06	-.1088E-05	-.4842E-06	0.
T	684.229	677.678	670.221	663.140	658.417	658.090
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	0.
J	37					
U	0.					
V	0.					
T	658.011					
Q	0.					

\*\*\*\*\* TIME = 10.0000 .0313  
RE = .12482E+05

INLET TEMPERATURE = .588.000  
AVERAGE TEMPERATURE RISE = 90.472  
AVERAGE EXIT TEMPERATURE = 678.472  
AVERAGE POWER DENSITY = .23224E+08

IN ROW 20 AVER TEMP = 678.4716 MASS FLOW IN = .285988E+03 MASS FLOW OUT = .278961E+03

J	1	2	3	4	5	6
U	0.	0.	.5346E+00	.4782E+00	.4039E+00	.3116E+00
V	0.	0.	-.2223E-05	-.5863E-05	-.8323E-05	-.1001E-04
T	657.972	656.557	653.138	654.994	660.094	666.241
Q	0.	0.	.2757E+00	.2757E+00	.2757E+00	.2757E+00
J	7	8	9	10	11	12
U	.3117E+00	.3118E+00	.3118E+00	.3119E+00	.3119E+00	.3120E+00
V	-.1143E-04	-.1287E-04	-.1394E-04	-.1446E-04	-.1437E-04	-.1371E-04
T	671.791	676.469	680.228	683.136	685.313	686.897
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00
J	13	14	15	16	17	18
U	.3120E+00	.3120E+00	.3120E+00	.3120E+00	.3120E+00	.3120E+00
V	-.1255E-04	-.1099E-04	-.9117E-05	-.7016E-05	-.4758E-05	-.2402E-05
T	688.021	688.799	689.324	689.666	689.877	689.990
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00
J	19	20	21	22	23	24
U	.3120E+00	.3120E+00	.3120E+00	.3120E+00	.3120E+00	.3120E+00
V	.2754E-13	.2402E-05	.4758E-05	.7016E-05	.9117E-05	.1099E-04
T	690.026	689.990	689.877	689.666	689.324	688.799
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00
J	25	26	27	28	29	30
U	.3120E+00	.3120E+00	.3119E+00	.3119E+00	.3118E+00	.3118E+00
V	.1255E-04	.1371E-04	.1437E-04	.1446E-04	.1394E-04	.1287E-04
T	688.021	686.897	685.313	683.136	680.228	676.469
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00
J	31	32	33	34	35	36
U	.3117E+00	.3116E+00	.4039E+00	.4782E+00	.5346E+00	0.
V	.1143E-04	.1001E-04	.8323E-05	.5863E-05	.2223E-05	0.
T	671.791	666.241	660.094	654.994	653.138	656.557
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	0.
J	37					
U	0.					
V	0.					
T	657.972					
Q	0.					



J		19		20		21		22		23		24
U		.2264E-01		.2262E-01		.2256E-01		.2245E-01		.2230E-01		.2210E-01
V		.3751E-14		.2589E-05		.5302E-05		.8256E-05		.1155E-04		.1527E-04
T		695.651		695.558		695.280		694.809		694.136		693.249
Q		.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00
J		25		26		27		28		29		30
U		.2185E-01		.2154E-01		.2116E-01		.2073E-01		.2022E-01		.1964E-01
V		.1946E-04		.2410E-04		.2913E-04		.3438E-04		.3957E-04		.4425E-04
T		692.132		690.768		689.142		687.238		685.046		682.562
Q		.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00
J		31		32		33		34		35		36
U		.1899E-01		.1829E-01		.2993E-01		.4091E-01		.5029E-01		0.
V		.4777E-04		.4926E-04		.4638E-04		.3509E-04		.1340E-04		0.
T		679.791		676.752		673.484		670.293		667.557		665.611
Q		.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00		0.
J		37										
U		0.										
V		0.										
T		664.750										
Q		0.										

\*\*\*\*\* TIME = 570.0000 1.0000  
 RE = .90076E+03

INLET TEMPERATURE = 588.000  
 AVERAGE TEMPERATURE RISE = 97.630  
 AVERAGE EXIT TEMPERATURE = 685.630  
 AVERAGE POWER DENSITY = .23942E+07

IN ROW	20	AVER TEMP =	685.6296	MASS FLOW IN =	.206383E+02	MASS FLOW OUT =	.200909E+02					
J		1		2		3		4		5		6
U		0.		0.		.5032E-01		.4092E-01		.2994E-01		.1830E-01
V		0.		0.		-.1341E-04		-.3513E-04		-.4645E-04		-.4935E-04
T		664.837		665.692		667.625		670.351		673.534		676.795
Q		0.		0.		.2757E+00		.2757E+00		.2757E+00		.2757E+00
J		7		8		9		10		11		12
U		.1899E-01		.1964E-01		.2022E-01		.2072E-01		.2116E-01		.2153E-01
V		-.4786E-04		-.4435E-04		-.3967E-04		-.3449E-04		-.2923E-04		-.2419E-04
T		679.828		682.594		685.074		687.263		689.164		690.787
Q		.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00
J		13		14		15		16		17		18
U		.2184E-01		.2210E-01		.2230E-01		.2245E-01		.2255E-01		.2262E-01
V		-.1954E-04		-.1535E-04		-.1161E-04		-.8303E-05		-.5334E-05		-.2605E-05
T		692.149		693.254		694.150		694.822		695.292		695.570
Q		.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00
J		19		20		21		22		23		24
U		.2264E-01		.2262E-01		.2256E-01		.2245E-01		.2230E-01		.2210E-01
V		.2070E-14		.2605E-05		.5334E-05		.8303E-05		.1161E-04		.1535E-04
T		695.662		695.570		695.292		694.822		694.150		693.264
Q		.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00
J		25		26		27		28		29		30
U		.2184E-01		.2153E-01		.2116E-01		.2072E-01		.2022E-01		.1964E-01
V		.1954E-04		.2419E-04		.2923E-04		.3449E-04		.3967E-04		.4435E-04
T		692.149		690.787		689.164		687.263		685.074		682.594
Q		.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00
J		31		32		33		34		35		36
U		.1899E-01		.1830E-01		.2994E-01		.4092E-01		.5032E-01		0.
V		.4786E-04		.4935E-04		.4645E-04		.3513E-04		.1341E-04		0.
T		679.828		676.795		673.484		670.351		667.625		665.692
Q		.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00		0.
J		37										
U		0.										
V		0.										
T		664.837										
Q		0.										

\*\*\*\*\* TIME = 575.0000 1.0000  
 RE = .90076E+03

INLET TEMPERATURE = 588.000  
 AVERAGE TEMPERATURE RISE = 97.658  
 AVERAGE EXIT TEMPERATURE = 685.658

AVERAGE POWER DENSITY = .23942E+07

IN ROW 20 AVER TEMP = 685.6583 MASS FLOW IN = .206383E+02 MASS FLOW OUT = .200907E+02

J		1		2		3		4		5		6
U	0.		0.		.5034E-01		.4093E-01		.2995E-01		.1830E-01	
V	0.		0.		-.1342E-04		-.3517E-04		-.4651E-04		-.4943E-04	
T	664.923		665.772		667.692		670.408		673.583		676.838	
Q	0.		0.		.2757E+00		.2757E+00		.2757E+00		.2757E+00	
J		7		8		9		10		11		12
U	.1899E-01		.1964E-01		.2022E-01		.2072E-01		.2116E-01		.2153E-01	
V	-.4796E-04		-.4445E-04		-.3978E-04		-.3459E-04		-.2933E-04		-.2429E-04	
T	679.865		682.626		685.102		687.287		689.185		690.806	
Q	.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00	
J		13		14		15		16		17		18
U	.2184E-01		.2209E-01		.2229E-01		.2244E-01		.2255E-01		.2261E-01	
V	-.1963E-04		-.1542E-04		-.1168E-04		-.8351E-05		-.5366E-05		-.2622E-05	
T	692.165		693.279		694.164		694.835		695.304		695.582	
Q	.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00	
J		19		20		21		22		23		24
U	.2263E-01		.2261E-01		.2255E-01		.2244E-01		.2229E-01		.2209E-01	
V	.3958E-14		.2622E-05		.5366E-05		.8351E-05		.1168E-04		.1542E-04	
T	695.674		695.582		695.304		694.835		694.164		693.279	
Q	.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00	
J		25		26		27		28		29		30
U	.2184E-01		.2153E-01		.2116E-01		.2072E-01		.2022E-01		.1964E-01	
V	.1963E-04		.2429E-04		.2933E-04		.3459E-04		.3978E-04		.4445E-04	
T	692.165		690.806		689.185		687.287		685.102		682.626	
Q	.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00	
J		31		32		33		34		35		36
U	.1899E-01		.1830E-01		.2995E-01		.4093E-01		.5034E-01		0.	
V	.4796E-04		.4943E-04		.4651E-04		.3517E-04		.1342E-04		0.	
T	679.865		676.838		673.583		670.408		667.692		665.772	
Q	.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00		0.	
J		37										
U	0.											
V	0.											
T	664.923											
Q	0.											

\*\*\*\*\* TIME = 580.0000 1.0000  
RE = .90076E+03

INLET TEMPERATURE = 588.000  
AVERAGE TEMPERATURE RISE = 97.687  
AVERAGE EXIT TEMPERATURE = 685.687  
AVERAGE POWER DENSITY = .23942E+07.

IN ROW 20 AVER TEMP = 685.6867 MASS FLOW IN = .206383E+02 MASS FLOW OUT = .200906E+02

J		1		2		3		4		5		6
U	0.		0.		.5036E-01		.4094E-01		.2995E-01		.1830E-01	
V	0.		0.		-.1343E-04		-.3521E-04		-.4658E-04		-.4951E-04	
T	665.009		665.853		667.759		670.465		673.632		676.880	
Q	0.		0.		.2757E+00		.2757E+00		.2757E+00		.2757E+00	
J		7		8		9		10		11		12
U	.1899E-01		.1964E-01		.2021E-01		.2072E-01		.2116E-01		.2153E-01	
V	-.4805E-04		-.4455E-04		-.3988E-04		-.3469E-04		-.2943E-04		-.2438E-04	
T	679.902		682.658		685.130		687.311		689.206		690.825	
Q	.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00	
J		13		14		15		16		17		18
U	.2183E-01		.2209E-01		.2229E-01		.2244E-01		.2254E-01		.2261E-01	
V	-.1971E-04		-.1550E-04		-.1174E-04		-.8399E-05		-.5399E-05		-.2638E-05	
T	692.182		693.294		694.178		694.848		695.316		695.594	
Q	.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00	
J		19		20		21		22		23		24
U	.2263E-01		.2261E-01		.2254E-01		.2244E-01		.2229E-01		.2209E-01	
V	.4107E-14		.2638E-05		.5399E-05		.8399E-05		.1174E-04		.1542E-04	
T	695.686		695.594		695.316		694.848		694.178		693.294	
Q	.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00		.2757E+00	
J		25		26		27		28		29		30
U	.2183E-01		.2153E-01		.2116E-01		.2072E-01		.2021E-01		.1964E-01	

V	.1971E-04	.2438E-04	.2943E-04	.3469E-04	.3988E-04	.4455E-04
T	692.182	690.825	689.206	687.311	685.130	682.658
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00

J	31	32	33	34	35	36
U	.1899E-01	.1830E-01	.2995E-01	.4094E-01	.5036E-01	0.
V	.4805E-04	.4951E-04	.4658E-04	.3521E-04	.1343E-04	0.
T	679.902	676.880	673.632	670.465	667.759	665.853
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	0.

J	37
U	0.
V	0.
T	665.009
Q	0.

\*\*\*\*\* TIME = 585.0000 1.0000  
 RE = .90076E+03

INLET TEMPERATURE = 588.000  
 AVERAGE TEMPERATURE RISE = 97.715  
 AVERAGE EXIT TEMPERATURE = 685.715  
 AVERAGE POWER DENSITY = .23942E+07

IN ROW 20 AVER TEMP = 685.7149 MASS FLOW IN = .206383E+02 MASS FLOW OUT = .200904E+02

J	1	2	3	4	5	6
U	0.	0.	.5039E-01	.4096E-01	.2996E-01	.1830E-01
V	0.	0.	-.1344E-04	-.3525E-04	-.4664E-04	-.4960E-04
T	665.095	665.932	667.826	670.521	673.680	676.922
Q	0.	0.	.2757E+00	.2757E+00	.2757E+00	.2757E+00

J	7	8	9	10	11	12
U	.1899E-01	.1964E-01	.2021E-01	.2072E-01	.2115E-01	.2152E-01
V	-.4814E-04	-.4465E-04	-.3999E-04	-.3480E-04	-.2953E-04	-.2448E-04
T	679.938	682.689	685.157	687.335	689.227	690.843
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00

J	13	14	15	16	17	18
U	.2183E-01	.2208E-01	.2228E-01	.2243E-01	.2254E-01	.2260E-01
V	-.1980E-04	-.1557E-04	-.1180E-04	-.8447E-05	-.5431E-05	-.2654E-05
T	692.198	693.309	694.191	694.860	695.328	695.605
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00

J	19	20	21	22	23	24
U	.2262E-01	.2260E-01	.2254E-01	.2243E-01	.2228E-01	.2208E-01
V	.4610E-14	.2654E-05	.5431E-05	.8447E-05	.1180E-04	.1557E-04
T	695.697	695.605	695.328	694.860	694.191	693.309
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00

J	25	26	27	28	29	30
U	.2183E-01	.2152E-01	.2115E-01	.2072E-01	.2021E-01	.1964E-01
V	.1980E-04	.2448E-04	.2953E-04	.3480E-04	.3999E-04	.4465E-04
T	692.198	690.843	689.227	687.335	685.157	682.689
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00

J	31	32	33	34	35	36
U	.1899E-01	.1830E-01	.2996E-01	.4096E-01	.5039E-01	0.
V	.4814E-04	.4960E-04	.4664E-04	.3525E-04	.1344E-04	0.
T	679.938	676.922	673.680	670.521	667.826	665.932
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	0.

J	37
U	0.
V	0.
T	665.095
Q	0.

\*\*\*\*\* TIME = 590.0000 1.0000  
 RE = .90076E+03

INLET TEMPERATURE = 588.000  
 AVERAGE TEMPERATURE RISE = 97.743  
 AVERAGE EXIT TEMPERATURE = 685.743  
 AVERAGE POWER DENSITY = .23942E+07

IN ROW 20 AVER TEMP = 685.7428 MASS FLOW IN = .206383E+02 MASS FLOW OUT = .200903E+02

J	1	2	3	4	5	6
U	0.	0.	.5041E-01	.4097E-01	.2996E-01	.1831E-01
V	0.	0.	-.1346E-04	-.3529E-04	-.4671E-04	-.4968E-04

T	665.180	666.011	667.892	670.577	673.728	676.964
Q	0.	0.	.2757E+00	.2757E+00	.2757E+00	.2757E+00
J	7	8	9	10	11	12
U	.1899E-01	.1964E-01	.2021E-01	.2072E-01	.2115E-01	.2152E-01
V	-.4824E-04	-.4475E-04	-.4009E-04	-.3490E-04	-.2963E-04	-.2447E-04
T	679.974	682.721	685.184	687.359	689.247	690.801
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00
J	13	14	15	16	17	18
U	.2183E-01	.2208E-01	.2228E-01	.2243E-01	.2253E-01	.2260E-01
V	-.1989E-04	-.1565E-04	-.1186E-04	-.8494E-05	-.5464E-05	-.2671E-05
T	692.214	693.323	694.205	694.873	695.340	695.617
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00
J	19	20	21	22	23	24
U	.2262E-01	.2260E-01	.2253E-01	.2243E-01	.2228E-01	.2208E-01
V	.2629E-14	.2671E-05	.5464E-05	.8494E-05	.1186E-04	.1565E-04
T	695.708	695.617	695.340	694.873	694.205	693.323
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00
J	25	26	27	28	29	30
U	.2183E-01	.2152E-01	.2115E-01	.2072E-01	.2021E-01	.1964E-01
V	.1989E-04	.2457E-04	.2963E-04	.3490E-04	.4009E-04	.4475E-04
T	692.214	690.861	689.247	687.359	685.184	682.721
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00
J	31	32	33	34	35	36
U	.1899E-01	.1831E-01	.2996E-01	.4097E-01	.5041E-01	0.
V	.4824E-04	.4968E-04	.4671E-04	.3529E-04	.1346E-04	0.
T	679.974	676.964	673.728	670.577	667.892	666.011
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	0.
J	37					
U	0.					
V	0.					
T	665.180					
Q	0.					

\*\*\*\*\* TIME = 595.0000 1.0000  
RE = .90076E+03

INLET TEMPERATURE = 588.000  
AVERAGE TEMPERATURE RISE = 97.770  
AVERAGE EXIT TEMPERATURE = 685.770  
AVERAGE POWER DENSITY = .23942E+07

IN ROW 20 AVER TEMP = 685.7704 MASS FLOW IN = .206383E+02 MASS FLOW OUT = .200901E+02

J	1	2	3	4	5	6
U	0.	0.	.5043E-01	.4098E-01	.2997E-01	.1831E-01
V	0.	0.	-.1347E-04	-.3533E-04	-.4678E-04	-.4977E-04
T	665.264	666.089	667.957	670.632	673.775	677.005
Q	0.	0.	.2757E+00	.2757E+00	.2757E+00	.2757E+00
J	7	8	9	10	11	12
U	.1899E-01	.1964E-01	.2021E-01	.2071E-01	.2115E-01	.2152E-01
V	-.4834E-04	-.4486E-04	-.4020E-04	-.3501E-04	-.2974E-04	-.2467E-04
T	680.009	682.751	685.211	687.382	689.268	690.879
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00
J	13	14	15	16	17	18
U	.2182E-01	.2207E-01	.2227E-01	.2242E-01	.2253E-01	.2259E-01
V	-.1997E-04	-.1572E-04	-.1193E-04	-.8543E-05	-.5497E-05	-.2687E-05
T	692.230	693.338	694.218	694.885	695.352	695.628
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00
J	19	20	21	22	23	24
U	.2261E-01	.2259E-01	.2253E-01	.2242E-01	.2227E-01	.2207E-01
V	.4439E-14	.2687E-05	.5497E-05	.8543E-05	.1193E-04	.1572E-04
T	695.720	695.628	695.352	694.885	694.218	693.338
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00
J	25	26	27	28	29	30
U	.2182E-01	.2152E-01	.2115E-01	.2071E-01	.2021E-01	.1964E-01
V	.1997E-04	.2467E-04	.2974E-04	.3501E-04	.4020E-04	.4486E-04
T	692.230	690.879	689.268	687.382	685.211	682.751
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00
J	31	32	33	34	35	36
U	.1899E-01	.1831E-01	.2997E-01	.4098E-01	.5043E-01	0.
V	.4834E-04	.4977E-04	.4678E-04	.3533E-04	.1347E-04	0.
T	680.009	677.005	673.775	670.632	667.957	666.089
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	0.

J 37  
 U 0.  
 V 0.  
 T 665.264  
 Q 0.

\*\*\*\*\* TIME = 600.0000 1.0000  
 RE = .90076E+03

INLET TEMPERATURE = 588.000  
 AVERAGE TEMPERATURE RISE = 97.798  
 AVERAGE EXIT TEMPERATURE = 685.798  
 AVERAGE POWER DENSITY = .23942E+07

IN ROW 20 AVER TEMP = 685.7978 MASS FLOW IN = .206383E+02 MASS FLOW OUT = .200900E+02

J	1	2	3	4	5	6
U	0.	0.	.5045E-01	.4099E-01	.2998E-01	.1831E-01
V	0.	0.	-.1348E-04	-.3537E-04	-.4685E-04	-.4985E-04
T	665.348	666.167	668.022	670.686	673.822	677.045
Q	0.	0.	.2757E+00	.2757E+00	.2757E+00	.2757E+00
J	7	8	9	10	11	12
U	.1899E-01	.1964E-01	.2021E-01	.2071E-01	.2114E-01	.2151E-01
V	-.4843E-04	-.4496E-04	-.4030E-04	-.3511E-04	-.2984E-04	-.2476E-04
T	680.045	682.782	685.237	687.405	689.288	690.897
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00
J	13	14	15	16	17	18
U	.2182E-01	.2207E-01	.2227E-01	.2242E-01	.2252E-01	.2259E-01
V	-.2006E-04	-.1580E-04	-.1199E-04	-.8591E-05	-.5529E-05	-.2704E-05
T	692.246	693.352	694.231	694.897	695.363	695.639
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00
J	19	20	21	22	23	24
U	.2261E-01	.2259E-01	.2252E-01	.2242E-01	.2227E-01	.2207E-01
V	.3256E-14	.2704E-05	.5529E-05	.8591E-05	.1199E-04	.1580E-04
T	695.731	695.639	695.363	694.897	694.231	693.352
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00
J	25	26	27	28	29	30
U	.2182E-01	.2151E-01	.2114E-01	.2071E-01	.2021E-01	.1964E-01
V	.2006E-04	.2476E-04	.2984E-04	.3511E-04	.4030E-04	.4496E-04
T	692.246	690.897	689.288	687.405	685.237	682.782
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00
J	31	32	33	34	35	36
U	.1899E-01	.1831E-01	.2998E-01	.4099E-01	.5045E-01	0.
V	.4843E-04	.4985E-04	.4685E-04	.3537E-04	.1348E-04	0.
T	680.045	677.045	673.822	670.686	668.022	666.167
Q	.2757E+00	.2757E+00	.2757E+00	.2757E+00	.2757E+00	0.
J	37					
U	0.					
V	0.					
T	665.348					
Q	0.					

SYS DEVICES 844/15/PF FLS=200K FLL=1747K MXS=160K MXL=1321K MXB=1321B

HH.MM.SS CPU SECOND ORIGIN

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10.19.41.MFB. QL NOS/BE 1.4 MFB 518CK 02/06/84
10.33.13 00000.004 MFZ. -ERIK,STMFZ,T477.
10.33.13 00000.005 JOB. -ACCOUNT,EGC,1530,*.:
10.33.17 00000.026 JOB. -ATTACH,P,TWODNEW,ID=EGC,ST=MFB.
10.33.17 00000.027 JOB. -FILE,P,RT=Z,FL=80.
10.33.17 00000.028 JOB. -ATTACH,TAPE1,FLATDISTRIB,ID=EGC,ST=STP.
10.33.17 00000.029 JOB. -FILE,TAPE1,RT=Z,FL=80.
10.33.17 00000.030 LOD. -FTN,I=P,R=3.
10.33.18 00000.038 MFZ. JM262 - STAGE PF ST=MFB LFN=P
10.33.33 MFB. 10.31.30.Z0005RJ 02/06/84
10.33.33 MFB. 10.31.31.-SP- ATTACHING PFN=TWODNEW
10.33.33 MFB. 10.31.35.AT CY= 035 SN=SHARE
10.33.33 MFB. 10.31.38.NO. WORDS= 11424 NO. OF EOR = 0
10.33.34 MFB. JM511 - WORDS READ- 26240B
10.36.19 00002.912 USR. 2.879 CP SECONDS COMPILATION TIME
10.36.19 00002.912 LOD. -LGO,PL=77777.
10.36.20 00003.113 MFZ. LD610 - FLS REQUIRED TO LOAD - 0015624 OU.COG
10.36.20 00003.114 MFZ. LD603 - EXECUTION INITIATED OS.EXP
10.36.20 00003.114 USR. FORTRAN LIBRARY 552 05/02/83
10.36.20 00003.118 MFZ. JM262 - STAGE PF ST=STP LFN=TAPE1
10.36.29 MFA. 10.34.34.Z00055G 02/06/84
10.36.29 MFA. 10.34.34.-SP- ATTACHING PFN=FLATDISTRIB
10.36.30 MFA. 10.34.34.AT CY= 010 SN=SHARE
10.36.30 MFA. 10.34.35.NO. WORDS= 149 NO. OF EOR = 0
10.36.31 STP. JM511 - WORDS READ- 225B
10.43.32 00161.407 USR. STOP
10.43.32 00161.407 USR. 043000 FINAL EXECUTION FL.
10.43.32 00161.407 USR. 158.292 CP SECONDS EXECUTION TIME.
10.43.32 00161.409 MFZ. JM166 - MAXIMUM USER SCM 47000B WORDS
10.43.32 00161.409 MFZ. JM167 - MAXIMUM USER LCM 0B WORDS
10.43.32 00161.410 MFZ. JM170 - MAXIMUM JS+10 LCM 110B BUFFERS
10.43.32 00161.410 MFZ. RM770 - MAXIMUM ACTIVE FILES 5
10.43.32 00161.410 MFZ. RM771 - OPEN/CLOSE CALLS 26
10.43.32 00161.410 MFZ. RM772 - DATA TRANSFER CALLS 42,217
10.43.32 00161.410 MFZ. RM773 - CONTROL/POSITIONING CALLS 150
10.43.32 00161.410 MFZ. RM774 - BM DATA TRANSFER CALLS 2,865
10.43.32 00161.410 MFZ. RM775 - BM CONTROL/POSITIONING CALLS 227
10.43.32 00161.411 MFZ. RM776 - QUEUE MANAGER CALLS 528
10.43.32 00161.411 MFZ. RM777 - RECALL CALLS 437
10.43.32 00161.411 MFZ. SCM 3 019.224 KWS
10.43.32 00161.411 MFZ. I/O 0.266 MW
10.43.32 00161.412 MFZ. RMS 27.073 MWS
10.43.32 00161.412 MFZ. USER 157.465 SEC
10.43.32 00161.412 MFZ. JOB 161.414 SEC
10.43.32 00161.412 MFZ. DIO 1 605.577 KW
10.43.32 00161.412 MFZ. SIO 11.573 KW
10.43.32 00161.412 MFZ. RU 348.351
10.43.32 00161.413 MFZ. BNL CCUS 11.950
10.43.32 00161.413 MFZ. SC050 - 007420 SC/LC SWAPS
10.43.32 00161.413 MFZ. EJ END OF JOB - PN001530

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B.5 CODE LIMITATIONS

1. The present version of the code limits the total number of axial and radial nodes to 40 each; if a larger number of nodes is desired in the analysis, all arrays of the form  $X(M_{axial}, N_{radial})$  and  $Y(N_{radial})$  should be properly redimensioned, bearing in mind that the number of radial nodes, for computational reasons, should be odd in input.
2. Due to the presence of discontinuity in the slope of the friction factor vs.  $Re$  for blanket assemblies, inaccuracies in computed velocities may be present in the vicinity of  $4500 < Re \leq 5500$ .
3. Flow-reversal in channel has not been considered; hence, if axial velocity becomes negative, a message is printed to this effect, giving time and position of occurrence, and axial velocity is reset to  $10^{-5} \times \hat{U}$ ; a similar message is displayed when velocity recovers to positive values.
4. There is no restart capability at present.
5. Axial power distribution is in the shape of a chopped cosine, with peak to average value at 1.4 while the upper and lower fission gas plena have no heat fluxes associated with them.
6. At present, friction factors for radial blankets are included as default, for fuel assembly calculations appropriate friction factors must be provided.
7. Porosity and hydraulic diameters are assigned as follows:

radial node number	$j = 1, 2, 3;$	and $N, N-1, N-2$	are	External
	$6 \leq j \leq N-5$		are	Internal
	others			linearly interpolate

8. The number of radial nodes,  $N$ , must be greater or equal to 11 ( $N$  must be an odd number).



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7. AUTHOR(S) Mohsen Khatib-Rahbar and Erik G. Cazzoli		3. RECIPIENT'S ACCESSION NO.	
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16. ABSTRACT (200 words or less)  Phenomenological models and numerical techniques for prediction of coolant flow and temperature fields during forced, mixed, and free convection regimes of operation in LMFBR subassemblies are addressed. It is shown that, simplified integral solutions provide an excellent approach to assessing the importance of the intra-subassembly buoyancy induced flow redistribution, and the transverse thermal conduction and mixing effects on the assembly wide peak coolant temperatures. Furthermore, a more detailed steady-state and transient parabolic two-dimensional porous-body model, resulting in the TWIST computer code is developed. Comparison of calculated results and out-of-pile sodium and water test data indicate generally good agreement in cross-assembly temperature profiles. However, the impact of fuel pin distortion and bowing, caused by large transverse power gradients on transverse temperature distributions are found to be significant.			
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