

CONF

CONF-811113--40

RADIATION EFFECTS IN MAT

DE82 004099

J. L. Scott, M. L. Grossbeck, and P. J. Maziasz  
Oak Ridge National Laboratory  
P.O. Box X  
Oak Ridge, Tennessee 37830

INTER

The 14-MeV neutrons produced in a fusion reactor result in different irradiation damage than the equivalent fluence in a fast breeder reactor, not only because of the higher defect generation rate, but because of the production of significant concentrations of helium and hydrogen. Although no fusion test reactor exists, the effects of combined displacement damage plus helium can be studied in mixed-spectrum fission reactors for alloys containing nickel (e.g., austenitic stainless steels). The presence of helium appears to modify vacancy and interstitial recombination such that microstructural development in alloys differs between the fusion and fission reactor environments.

Since mechanical properties of alloys are related to the microstructure, the simultaneous production of helium and displacement damage impacts upon key design properties such as tensile, fatigue, creep, and crack growth. Through an understanding of the basic phenomena occurring during irradiation and the relationships between microstructure and properties, alloys can be tailored to minimize radiation-induced swelling and improve mechanical properties in fusion reactor service.

---

\*Research sponsored by Office of Fusion Energy, U.S. Department of Energy under contract W-7405-eng-26 with Union Carbide Corporation.

## Introduction

Fusion power together with breeder reactors and solar energy is a candidate to solve mankind's long-term energy requirements. The raw materials for fusion power are deuterium from water and lithium, which is used to produce tritium, from slat deposits. Both are sufficiently abundant to provide for projected worldwide energy demands for centuries. The product of the fusion reaction is helium, a valuable resource. There will be neutron activation of the reactor structure during operation; but the exposed structure will decay to moderate levels of activity in a few decades, so that the waste problems of a fusion economy are small in comparison to those of fission reactors. Thus, if it can be made to work economically, fusion offers enormous benefit to mankind.

At present the primary obstacle to fusion power is plasma confinement. Plasmas have been produced and confined, but not yet to the point of a self-sustained fusion process. A variety of confinement techniques are being explored, including tokamaks, tandem mirrors, and inertial confinement systems.<sup>1,2</sup> It is likely that the Tokamak Fusion Test Reactor being built at Princeton will demonstrate energy breakeven before 1985. The feasibility of mirror reactors should be established soon thereafter. It will then remain to demonstrate the system that will lead to the most economical commercial power.

The materials demands for fusion reactors are severe. An idealized blanket and shield segment of a tokamak, shown in Figure 1, illustrates the functional requirements that materials must satisfy in a fusion power system. In this design, a niobium alloy is the structural material; and

lithium is both a breeding material and coolant. Graphite is a neutron reflector and moderator. Water and lead are both shielding materials. Directly outside the shield is the superconducting coil that provides the toroidal field. This coil consists of the superconductor imbedded in a normal conductor. The temperatures at which the components operate in this design are indicated.

The blanket structure, coolant, breeding material, neutron moderator, and shielding materials are common to all fusion devices. In a few designs, neutron-multiplication systems are required for breeding. In magnetically confined systems, limiters inside the first wall both intercept the plasma under off-normal conditions and protect the blanket. Inertially confined systems do not require limiters or superconducting magnets, but they have unique components such as powerful lasers or other intense energy delivery systems, complex optical systems, and target pellets.

The first-wall structural material provides vacuum containment at  $10^{-9}$  torr (130 nPa) before fueling, and also contains the breeding material, coolant, and reflector materials. It is desirable to operate at a reasonably high temperature ( $>400^{\circ}\text{C}$ ) for good heat transfer and efficiency. Unless a protective coating is used, the wall material must be compatible with the plasma environment of photons, energetic ions, and neutral atoms. It must also be compatible with the coolant and breeding materials. That part of the structure that faces the plasma operates either in a high vacuum or in a hydrogenic environment at different stages of the burn. The external reactor structure may be exposed either to air, to an inert atmosphere, or to a vacuum.

The harshest environmental condition faced by the first-wall structure is the intense radiation field. Fusion reactions in a tokamak produce 3.5-MeV (0.55-pJ) alpha particles, which heat the plasma, and 14.1-MeV (2.25-pJ) neutrons, which heat the blanket and breed tritium. Since the positively charged alpha particles are emitted in an intense magnetic field, their orbits are confined within the plasma. After they heat ions and electrons within the plasma, their energy is transferred to the first wall primarily by bremsstrahlung and line radiation. This heat flux, 0.2 to 0.8 MW/m<sup>2</sup>, is deposited on the first wall. In contrast, the neutrons have a mean free path of 0.3 m or more in the blanket, so that their energy is deposited throughout the blanket. The energy density of the neutron current through the first wall, called the "wall loading," is in the range 1 to 4 MW/m<sup>2</sup>, or 4.43 to  $17.7 \times 10^{17}$  neutrons/(m<sup>2</sup>·s). Scattering of these neutrons adds to the flux at the first wall, creating a total flux in the range 3 to  $12 \times 10^{18}$  neutrons/(m<sup>2</sup>·s). These fluxes are lower than the peak flux in a fast breeder reactor, about  $5 \times 10^{18}$  neutrons/(m<sup>2</sup>·s), but the neutrons are more energetic.

Two consequences stem from the high-energy neutrons. First, they displace atoms in the lattice. Secondly, inelastic collisions occur, to produce high helium and hydrogen concentrations within the lattice. Damage measures for type 316 stainless steel in a typical fusion reactor and in various available irradiation test facilities are compared in Table I. The helium production rate in a fusion reactor first wall made of stainless steel is at least 40 times that in a fast breeder reactor. As Table II shows, the extent of damage assessed by each index varies

from metal to metal; but for all candidate materials for fusion reactors, the ratio of helium production to displacement per atom (dpa) production is much higher than it is in a fast reactor.

As the neutrons pass through the blanket and shield materials, atoms are knocked out of their equilibrium lattice sites, producing many vacancy-interstitial (Frenkel) pairs. Many defects are quickly annihilated, but some survive to interact with each other and with solute atoms to form dislocations and/or cavities, which affect the properties of the material. The effects are particularly important in the structural materials, which must perform well at the elevated temperatures and stresses required. The high concentrations of helium, together with atomic displacements, often lead to swelling and degradation of mechanical properties, with the result that the life of the structure is considerably shortened. It is imperative that radiation effects be understood and materials optimized for fusion service if fusion power is to be economically feasible.

In this paper the present approach to the problems and the understanding of radiation damage in fusion systems will be discussed. Some of the basic processes taking place in the metal lattice will be described. Then, ways to optimize the alloys for improved ductility and reduced swelling are proposed.

#### Approach to the Problem

The alloys for fusion service are generally selected on the basis of reactor design studies. Some of the earliest reactor studies utilized very high temperatures; consequently, niobium and molybdenum alloys were the materials of choice. Other studies focused upon the use of existing

commercial alloys, leading to the choice of austenitic stainless steels and nickel-base alloys. When emphasis was placed on the desire to minimize radioactive inventory, aluminum and vanadium alloys and SiC were selected. Recently the potential advantages of ferritic stainless steels were recognized and they were added to the set of candidate alloys being studied for fusion. Titanium alloys have also been considered for fusion because of their many desirable properties. A comparison of various alloys for an important set of reactor performance criteria is given in Table III. The specific alloys now being studied are described in a recent review of alloy development for fusion.<sup>3</sup>

The fact that no high flux fusion reactor test facility now exists hinders alloy development for fusion. At present, simulations must be used, requiring extrapolation to fusion reactor conditions. Various low-flux, 14-MeV facilities exist such as the Rotating Target Neutron Source (RTNS) at the Lawrence Livermore National Laboratory. These facilities permit the study of the evolution of the damage microstructure, and of properties sensitive to low fluences, but cannot produce fluences that match those of the fusion reactor environment. Fast-spectrum fission reactors can be used to evaluate the effects of displacement damage alone. Irradiation with dual beams of heavy ions and helium in accelerators may permit study of the swelling behavior and the microstructural evolution of various alloys, but do not provide specimens for the effective evaluation of mechanical properties.

A high flux deuterium-lithium source, the Fusion Materials Irradiation Test Facility (FMIT), to be built in Richland, Washington, will be available in 1985. Its primary mission will be to verify methods

developed to correlate the fusion environment with that in fission reactors. It will also be used to generate design data on a few selected high-priority materials. In the 1990s a fusion materials test reactor may be built.

Fortunately, a way exists to simulate the fusion reactor environment: nickel-containing alloys, such as austenitic stainless steels and nickel-base alloys, can yield helium by a two-step process to be discussed later.

In a mixed-spectrum fission reactor with both high-fast and high-thermal neutron fluxes, fast neutrons produce the displacements, while thermal neutrons produce helium. Typical facilities available for alloy development are the High Flux Isotope Reactor (HFIR) and the Oak Ridge Research Reactor (ORR). Since most of the results to date have been obtained on the austenitic stainless steels, especially type 316 and modifications thereof, the remainder of the paper will focus on these alloys; however, many of the principles should apply to other alloy systems as well.

### Basic Radiation Effects

Before getting into the detailed behavior of austenitic stainless steels in the mixed-spectrum fission reactor tests, it is helpful to review the basic radiation effects. A schematic diagram of a neutron interacting with a metal lattice is shown in Figure 2. An energetic neutron enters the lattice from the left and interacts elastically with a lattice atom which then becomes a primary knock-on atom. The neutron continues along its deflected path and creates other primary knock-on atoms until it no longer has sufficient energy to displace an atom, or until it is absorbed by a nucleus, creating a transmutation product. The primary

knock-on atom at first loses much of its energy through interactions with electrons. Then, as its energy decreases, nuclear cross sections increase and the atom causes many atom displacements as it traverses the lattice, creating isolated vacancies and interstitial atoms, replacement collisions, and finally a depleted zone rich in vacancies surrounded by an excess of interstitial atoms at the end of its range. In austenitic stainless steels at room temperature or above, the interstitial atoms are very mobile; and most recombine with vacancies with no lasting effect. A few interstitials are attracted to dislocations, grain boundaries, and other sinks, leaving an excess of vacancies in the lattice. Interstitials may also cluster to form two-dimensional platelets, called interstitial loops.

Vacancies are less mobile than interstitials, but are observed to migrate at significant rates at temperatures above 0.3 of the absolute melting temperature. Like interstitials, vacancies may be annihilated at a variety of sinks such as dislocations, grain boundaries, and precipitates. Vacancies may also cluster and grow into cavities, causing swelling; or they may also form dislocation loops.

#### Helium Effects

The presence of helium in the lattice can modify the response of the material. Interstitial helium, like a self-interstitial atom, is highly mobile; but it is also insoluble. Therefore, helium is quickly trapped by vacancies, dislocations, and other sinks. Vacancies containing helium atoms may be mobile. The exact mechanisms of motion are not known, but there is ample evidence that helium and vacancies flow in a coupled fashion at reactor damage rates ( $10^{-6}$  dpa/s). Since vacancy clusters are stabilized by helium, the formation of cavities normally occurs much

earlier in a fusion reactor than in a breeder reactor where little helium is produced. Also the concentration of cavities is sometimes much higher in the case of the fusion reactor. The result is that swelling should begin earlier in a fusion reactor. The rate of swelling, however, is lower at high fluences in austenitic stainless steels, particularly when a high concentration of helium-filled cavities acts as the dominant sinks for point defects in the systems.

High concentrations of helium-filled cavities in the matrix have little effect on the properties of austenitic stainless steels at room temperature; but, at temperatures above 400–500°C, helium can migrate and collect at the grain boundaries, seriously reducing ductility and causing intergranular failure. Figure 3 shows a scanning electron micrograph of the fracture surface of annealed Inconel 600 irradiated in HFIR at 650°C to a fluence that produced 1780 at. ppm He in the structure.<sup>4</sup> The sample fractured at room temperature in a brittle manner. So much helium accumulated at the grain boundaries that metallic bonding was provided only by thin webs (white in the photograph). One strategy for application of materials to fusion service is to design an alloy that reduces the helium concentration at the grain boundaries.

#### Effects of Helium on Mechanical Properties

The effects of helium on mechanical properties can be studied by introduction of helium either before or simultaneously with irradiation in much the same way that microstructural evolution is studied. The major difference lies in the fact that mechanical properties are largely bulk properties and must be studied with larger specimens than are used for electron microscopy.

As was discussed earlier, helium formed from  $(n,\alpha)$  reactions from fusion neutrons is an inescapable part of the fusion environment. The effects of helium can be studied by irradiation in 14-MeV-neutron accelerators, which better simulate the fusion environment than any other device currently available. However, this technique suffers from two serious disadvantages: the high-flux volume is so small that only a few specimens can be irradiated at a time, and the fluxes available are so low that the regime of helium embrittlement has not been attained. Even when a large high-flux neutron source becomes available (FMIT in 1985), the irradiation space will be so limited that only experiments that correlate 14-MeV-neutron effects with simulations will be performed. Simulation techniques that allow larger numbers of specimens to be studied will be necessary to develop the mechanical properties data base necessary for the design and construction of a fusion reactor.

Simulation Techniques. The techniques being used at the present time for fusion simulation may be grouped into two classes: those that introduce helium prior to irradiation and those that introduce helium simultaneously with irradiation. In the first category are the following: (1) accelerator implantation and (2) tritium trick doping. Methods introducing helium simultaneously with irradiation are: (1) dual-beam ion irradiation and (2) transmutation-produced helium from thermal neutrons in mixed-spectrum reactors.

Accelerator implantation consists of bombarding very thin ( $\sim 0.25$  mm) specimens with high-energy  $\alpha$ -particles in a manner that distributes the  $\alpha$ -particles uniformly throughout the specimen thickness. Implantation can be followed by neutron irradiation to complete the simulation of the

fusion environment. This method has the advantage that it can be used with any alloy system. It suffers from the disadvantage that only thin specimens that are subject to extraneous surface effects and corrosion can be employed.

Tritium trick doping is a clever technique that takes advantage of the decay of tritium to  $^3\text{He}$  with a 12-year half-life. Tritium is introduced into the specimens by equilibration with a tritium atmosphere. The tritium is then allowed to decay long enough to produce the desired concentration of helium, and then the excess tritium is allowed to diffuse out. The method results in a uniform distribution of helium in large specimens but suffers from the disadvantage that it applies only to metals with high solubilities for hydrogen. This limits the method to refractory metals such as niobium and vanadium and to titanium. Concentrations are limited by the decay time to tritium and the solubility.

Dual-beam ion irradiation employs a beam of heavy ions to produce atomic displacement damage and a beam of  $\alpha$ -particles to deposit helium. The helium implantation is similar to accelerator implantation previously described. The simultaneous ion bombardment provides rapid results with simultaneous helium and displacement damage in specimens of minimal radioactivity. This method is useful for fundamental studies of helium effects, and can be used with any alloy system, but is not as close a simulation to the fusion environment as neutron irradiation.

Transmutation-produced helium from thermal neutrons in mixed-spectrum reactors permits simultaneous helium formation and displacement damage production. The use of fission reactors permits large numbers of a large-size specimens to be irradiated and, therefore, lends itself to support

the mechanical properties base of a large alloy development program. Only  $^{59}\text{Ni}$  has a sufficiently large cross section for thermal-neutron absorption to be of use in helium production. The transmutation begins with  $^{58}\text{Ni}$  ( $n, \alpha$ )  $^{59}\text{Ni}$  which is followed by  $^{59}\text{Ni}$  ( $n, \alpha$ )  $^{56}\text{Fe}$ , with both reactions occurring with thermal neutrons. This method results in a homogeneous distribution of helium in any practical-size specimen (size limited only by reactor space and dissipation of nuclear heating). However, it can be used only with nickel-bearing alloys. Concentrations of several thousand parts-per-million helium can be attained, thus permitting the effects of helium at fusion reactor first-wall exposures as high as 40-50 MWyr/m<sup>2</sup>. In fact, such high concentrations can be attained that the quality of the fusion environment simulation is in question for nickel-base alloys and even in stainless steels, because of too high a helium-to-displacement per atom ratio. To more properly simulate the fusion environment, special techniques such as adjusting the ratio of fast-to-thermal neutrons (called spectral tailoring) are used to achieve the proper ratio of atomic displacements to helium concentration. Transmutation doping can, therefore, be used to study the effects of end-of-life helium concentration and study the effects of a closer simulation to the fusion environment. It suffers from the disadvantages that only nickel-bearing alloys may be studied and, as with any fission reactor irradiation, the fact that specimens are highly radioactive.

As can be concluded from the previous discussion, in order to engage in a compressive alloy development program encompassing several alloy classes, several methods of fusion irradiation simulation must be used.

Moreover, all of these methods (even accelerator-based 14-MeV-neutron generators) are only simulations of fusion reactor service. Therefore, their results must ultimately be correlated with service in an actual fusion device.

Results of Fusion Simulations. Nickel ions accompanied by  $\alpha$ -particles in a dual-beam accelerator have been used to study the effects on microstructure of preinjected helium as compared with continuously produced helium.<sup>5</sup> Significant differences between preinjection and continuous injection were observed in stainless steels in both phase stability and cavity nucleation.

Preinjection resulted in a finer microstructure than simultaneous injection; both dislocation loops and voids were smaller and more numerous. Although mechanical properties were not measured, a more refined microstructure such as results from preinjection is expected to cause higher strength accompanied by lower ductility.

Preinjection of helium followed by neutron irradiation in EBR-II was studied in V-20% Ti by Tanka, Bloom, and Horak.<sup>6</sup> They observed essentially no change in strength in the range of 400-700°C, but a substantial drop in ductility at temperatures above 600°C for 90 and 200 at. ppm He and intergranular fracture at 700°C. Since comparison samples without helium injection do not indicate any significant loss of ductility, the observed loss of ductility is attributed to helium. However, the refined microstructure observed in stainless steel in ion irradiation studies does not appear to apply to V-20% Ti, since no change in strength was observed. The loss of ductility is then not attributed to limited plastic flow resulting from dislocation pinning, but rather from the helium. Had

significant strengthening been observed, no conclusions about the effect of helium could have been drawn. It can, therefore, be seen that the results of preinjection experiments must be interpreted cautiously.

The method of tritium trick doping was used to investigate helium embrittlement in another vanadium alloy, V-15% Cr-5% Ti, by Mattas et al.<sup>7</sup> Yield strength was unaffected by helium concentrations up to 35 at. ppm, but ultimate tensile strength and elongation were decreased by the presence of helium especially at 700°C and above. Also consistent with the work of Tanaka et al., helium led to a tendency for intergranular fracture.

Mixed-spectrum fission reactor irradiations have yielded the largest amount of data on helium effects. The trends of large reductions in ductility at high temperatures continues to hold for mixed-spectrum reactor irradiations. However, strength may be increased or decreased, depending upon temperature and initial microstructure. The effect is dramatic in Inconel 600, where ductility drops from 16% to less than 1% upon irradiation at 600°C and the fracture mode transforms from ductile rupture to intergranular.<sup>8</sup> Type AISI 316 stainless steel<sup>9</sup> is more resistant to helium embrittlement but nonetheless exhibits very low ductility above 600°C after 40 dpa and approximately 4000 at. ppm He. The trends in both annealed and 20% cold-worked type 316 stainless steel have been summarized by Bloom,<sup>10</sup> who characterized the irradiation by the helium-to-displacement ratio. A high He/dpa ratio is considered to be approximately 60-80 obtained in the HFIR, while the term "low He/dpa" implies a ratio less than 1, obtained in fast reactors. Figures 4 and 5 show both yield stress and total tensile elongation for annealed and 20% cold-worked type 316

stainless steel. The increase in yield strength saturates at about 10 dpa and decreases with increasing temperature for both high and low He/dpa values. Ductility is significantly lower for the case of the high He/dpa value. In the annealed condition, ductility remains low throughout the temperature range studied, whereas a continuous decrease in ductility with increasing temperature is exhibited by material in the cold-worked condition. For type 316 stainless steel in both annealed and cold-worked condition, yield strength is lower for the high He/dpa value. Lower strength accompanied by lower ductility is characteristic of radiation and helium embrittlement.

Fatigue will be a major consideration in fusion reactors that operate in a cyclic mode. Fatigue in type 316 stainless steel in the 20%-cold-worked condition has been investigated following irradiation in the HFIR to produce helium levels as high as about 1000 at. ppm and displacement levels up to 15 dpa (ref. 11). For irradiation and test temperatures of 430°C, a reduction in fatigue life by a factor of 3 to 10 was observed (Figure 6). However, at 550°C, no significant effect of irradiation on low-cycle fatigue life was observed. At both temperatures, the irradiation reduced the strain-range level of the  $10^7$ -cycle endurance limit from 0.35 to 0.30%. However, 0.30% is sufficiently high for the normal operation of all present conceptual designs for fusion reactors. Accident conditions must still be evaluated.

Various simulation techniques have been discussed, and some results obtained with each method have been described. It remains to validate each of these techniques, even the mixed-spectrum reactor experiments and the 14-MeV-neutron accelerator experiments, the two closest simulations of

the fusion environment. A step in this direction has been taken by Vandervoort et al.,<sup>12</sup> who compared tensile properties of specimens irradiated in two types of neutron accelerators [ $\text{Be}(d,n)$  and  $\text{T}(d,n)$ ] and a fission reactor. As predicted, the high-energy neutrons from the accelerators had a larger effect on strength and ductility per unit fluence. However, when expressed in terms of damage energy, the effects on mechanical properties were found to be equivalent. It should still be kept in mind that the fluences in these experiments were too low to yield significant amounts of helium. For fluences producing appreciable helium, the helium is expected to dominate. By the time FMIT or a fusion test facility becomes available, perhaps the effects of helium will be sufficiently understood to make the outcome of verification experiments simply a necessary check.

#### Techniques for Minimizing the Effects of Helium

The metallurgist has available a number of tricks to minimize the deleterious effects of helium. One technique that works well in type 316 stainless steel is cold work. Cold work produces a high concentration of dislocations in the lattice that act as trapping sites for helium. Figure 7 shows a comparison of cold-worked and annealed steel irradiated at temperatures in the range 450–500°C so as to produce 42–53 dpa and 3000–3800 at. ppm He (ref. 13). In the cold-worked material, numerous small cavities are observed with a total of 2% swelling. In annealed steel, there are fewer, larger cavities and a total of 9% swelling. Calculations show that the amount of helium present in the samples can be accounted for in both cases by assuming that the cavities are equilibrium helium bubbles with internal pressure,  $P$ , given by

$$P = \frac{2\gamma}{r}, \quad (1)$$

where  $\gamma$  = surface tension of the metal and  $r$  = bubble radius. Smaller cavities contain helium at a higher pressure and result in less swelling for a given helium content.

The effect of temperature on the swelling rates of cold-worked and annealed type 316 stainless steel irradiated in HFIR is shown in Figure 8. From 375 to about 550°C the swelling of both cold-worked and annealed samples is fairly independent of temperature, but both showed significantly more swelling at 680°C. At this temperature recovery of the cold-worked microstructure occurs, so that annealed and cold-work samples behave about the same. Increased helium in HFIR often causes earlier swelling than observed in the breeder reactor (EBB-II), but void swelling in EBR-II can be greater than bubble swelling in HFIR at higher fluences.<sup>14</sup> Even though the swelling was about the same in cold-worked samples below 550°C, the bubble concentrations and sizes were different, as Table IV shows.

The use of controlled precipitation in an alloy can be more effective than cold work in mitigating the effects of helium. In austenitic stainless steels, titanium monocarbide (TiC) has been found to be an effective trap for helium.<sup>15</sup> Since the TiC is also rich in molybdenum, vanadium, and niobium, it is usually referred to as metal monocarbide (MC). Figure 9 shows a micrograph of solution-annealed type 316 stainless steel with 0.23 wt % Ti after irradiation in HFIR at 600°C to 30 dpa and 1850 at. ppm He. Although the MC phase is heavily decorated with helium, the Laves phase, a normal precipitate in type 316 stainless steel, is free of

helium. In ordinary type 316 stainless steel, the Laves phase is an effective trap for helium as Figure 1C shows; but the bubble size is larger.

The primary reason the MC phase attracts helium migrating in vacancies is that it is an oversize misfit phase, having about a 70% increase in atomic volume compared with untransformed austenite. Helium-laden vacancies are preferentially attracted by the large compressive stress around the MC particles. In addition, MC is much finer than the other precipitate phases, like Laves, reducing the migration distance. The overall result is that swelling is reduced and helium is kept out of grain boundaries.

#### Discussion

The presence of helium formed by nuclear reactions in alloys during irradiations in the fusion environment can have serious consequences and must be taken into consideration. Techniques that can be used (and illustrated with austenitic stainless steel in this paper) including cold-working and controlled precipitation of phases that act as helium traps. The same techniques can be used in other alloys, but the specific methods will vary for each alloy system. For all systems grain growth after helium accumulation must be avoided. There will also be a fluence limit beyond which the helium effects are expected to become overriding in any alloy system. Nevertheless, optimized alloys will extend the lifetime of fusion reactors, considerably, with large economic benefits.

There is, at present, a large alloy development program for fusion that was recently described elsewhere.<sup>3</sup> Many years of experimental and theoretical work will be required to fully understand the effects of

displacement damage alone and the combined effects of helium and displacement damage on the behavior of materials. Once the materials are understood, optimized alloys can then be developed. There will remain a large testing program to generate the full data base required for engineering design and code qualification. Although our present knowledge is meager, there is every reason to be optimistic that good alloys for fusion service can be developed.

#### References

1. C. C. Baker, G. A. Carlson, and R. A. Krakowski, "Trends and Developments in Magnetic Confinement Fusion Reactor Concepts," *Nucl. Technol./Fusion* 1(1): 5-78 (January 1981).
2. M. J. Monsler, J. Hovingh, D. L. Cook, T. G. Frank, and G. A. Moses, "An Overview of Inertial Fusion Reactor Design," *Nucl. Technol./Fusion* 1(3): 302-58 (July 1981).
3. R. E. Gold, E. E. Bloom, F. W. Clinard, Jr., D. L. Smith, R. D. Stevenson, and W. G. Wolfer, "Materials Technology for Fusion: Current Status and Future Requirements," *Nucl. Technol./Fusion* 1(2): 169-251 (April 1981).
4. J. L. Scott, "The Development of Advanced Structural Materials for Fusion Power," pp. 330-39 in *The 1980's - Payoff Decade For Advanced Materials*, Vol. 25, Society for the Advancement of Materials and Process Engineering, Azusa, Calif, 1980.
5. N. H. Packan and K. Farrell, "Simulation of First Wall Damage," *J. Nucl. Mater.* 85 & 86: 677-81 (December 1979).

6. M. P. Tanaka, E. E. Bloom, and J. A. Horak, "Tensile Properties and Microstructure of Helium Injected and Reactor Irradiated V-20 Ti," To be published in the proceedings of the Second Topical Meeting on Fusion Reactor Materials, ANS, Seattle, Washington, August 7-10, 1981.
7. R. F. Mattas, H. Wiedersich, D. G. Atteridge, A. B. Johnson, and J. R. Remark, "Elevated-Temperature Tensile Properties of V-15 Cr-5 Ti Containing Helium Introduced by Ion Bombardment and Tritium Decay," pp. 199-208 in *Proceedings of The Second Topical Meeting on The Technology of Controlled Nuclear Fusion CONF-760935-P1*, 1976.
8. F. W. Wiffen, "Response of Inconel 600 to Simulated Fusion Reactor Irradiation," pp. 88-106 *Effects of Radiation on Structural Materials ASTM STP 683*: American Society for Testing and Materials, Philadelphia, 1979.
9. E. E. Bloom and F. W. Wiffen, "The Effects of large Concentrations of Helium on The Mechanical Properties of Neutron-Irradiated Stainless Steel," *J. Nucl. Mater.* 58(2): 171-84 (1975).
10. E. E. Bloom, "Mechanical Properties of Materials in Fusion Reactor First Wall and Blanket Systems," *J. Nucl. Mater.* 85 & 86: 794-804 (December 1979).
11. M. L. Grossbeck and K. C. Liu, "Fatigue Behavior of Type 316 Stainless Steel Following Neutron-Irradiation-Inducing Helium," (Summary) *Trans. Am. Nucl. Soc.* 34 (TANSAO 34): 186-88 (1980).
12. R. R. Vandervoort, E. L. Raymond, and C. J. Echer, "High-Energy Neutron Irradiation Effects on The Tensile Properties and Microstructure of 316 Stainless Steel," *Radiat. Eff.* 45: 191-98 (1980).

13. F. W. Wiffen, P. J. Maziasz, E. E. Bloom, J. O. Stiegler, and M. L. Grossbeck, "The Behavior of Type 316 Stainless Steel Under Simulated Fusion Reactor Irradiation," pp. 146-59 in *The Metal Science of Stainless Steels*, The Metallurgical Society of AIME, Warrendale, Pa., 1978.
14. P. J. Maziasz, "Microstructural Development and Effects of Helium in Type 316 Stainless Steel Irradiated in HFIR and in EBR-II," *ADIP Quart. Prog. Rep. March 31, 1981*, DOE Report DOE/ER-0045/6, pp. 70-92.
15. P. J. Maziasz, "Helium Trapping and Ti-Rich MC Particles in Neutron-Irradiated Type 316 + Ti Stainless Steels," *Scr. Metall.* 14(11): 1251-56 (1980).

## List of Figures

Fig. 1. Schematic diagram of \_\_\_\_\_?

Fig. 2. Schematic diagram of neutron interactions with a metal lattice.

Fig. 3. Inconel 600, exposed in HFIR Reactor at 650°C to 8.5 displacements-per-atom level and 1780 at. ppm, fractured at 35°C.

Fig. 4. Tensile properties of 20% cold worked type 316 stainless steel. Open symbols for unirradiated samples, filled symbols for HFIR irradiated samples ( $7.1$  to  $8.7 \times 10^{26}$  neutron/m<sup>2</sup>) and EBR-II trend curves for fluences  $1.2$  to  $2.9 \times 10^{26}$  neutron/m<sup>2</sup>.

Fig. 5. Tensile properties of solution annealed type 316 stainless steel. Open symbols for unirradiated samples, filled symbols for HFIR irradiated samples ( $5.6$  to  $8.7 \times 10^{26}$  neutron/m<sup>2</sup>) and EBR-II trend curves for fluences  $1.2$  to  $2.9 \times 10^{26}$  neutron/m<sup>2</sup>.

Fig. 6. Fatigue life of 20% cold worked type 316 stainless steel irradiated in HFIR at 430°C and tested at irradiation temperature (200–1,000 at ppm He, 5–15 dpa).

Fig. 7. Cold work reduces cavity size and swelling for irradiation in HFIR to high helium levels.

Fig. 8. A high helium content in type 316 stainless steel causes a high swelling rate.

Fig. 9. Solution-annealed type 316 stainless steel with 0.23 wt % Ti after irradiation in HFIR at 600°C to 30 dpa and 1850 at. ppm He.

Fig. 10. Helium trapping at a TiC particle in type 316 stainless steel.

Fig. 11. Annealed type 316 stainless steel after irradiation in HFIR at 600°C to 30 dpa and 1850 at. ppm Hc.

TEMPERATURE

300°K

4.2°K

COIL RESTRAINT  
SUPER CONDUCTING  
COIL ( $\text{Nb}_3\text{Sn}$ )  
GAMMA SHIELD  
(LEAD)  
NEUTRON SHIELD  
(WATER)

NIOBIUM STRUCTURE  
LITHIUM COOLANT

GRAPHITE

1300°K

NIOBIUM STRUCTURE  
LITHIUM COOLANT

$10^8$ °K

PLASMA

237 cm

Fig. 1. Schematic Diagram of Neutron Interacting with Metal Lattice.

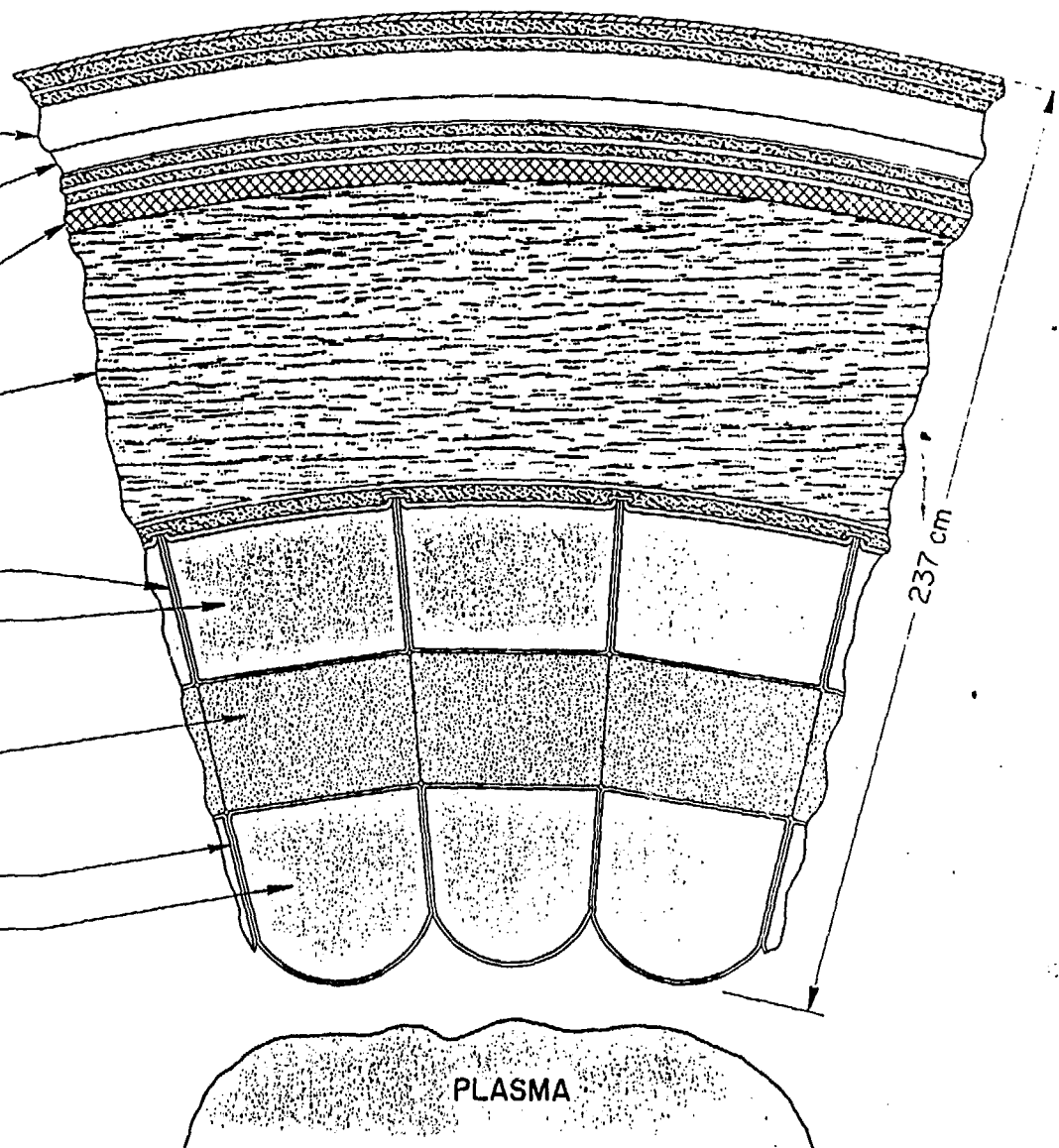


Table I. Damage production in available irradiation facilities

FACILITY	FLUX >0.1 MeV (n/m <sup>2</sup> .s)	DAMAGE PRODUCED <sup>b</sup> PER YEAR		
		dpa <sup>c</sup> <del>(No./yr)</del>	He <del>at.(ppm/yr)</del>	H <del>at.(ppm/yr)</del>
Fusion Reactor <sup>a</sup>	$2.7 \times 10^{18}$	12	144	532
Mixed-Spectrum Fission Reactor (HFIR)	15	35	1900 <sup>d</sup>	440
Fast-Breeder Reactor (EBR-II)	25	35	5	270
Rotating-Target Neutron Source	0.02	0.2	5	10

<sup>a</sup>Typical design, neutronic wall loading 1.0 MW/m<sup>2</sup>.

<sup>b</sup>Assumes 100% duty factor on all facilities.

<sup>c</sup>Displacements per atom.

<sup>d</sup>Helium production in HFIR is not linear with fluence.

Table II. **RADIATION DAMAGE IN ALLOYS TYPICAL  
OF THE ALLOY SYSTEMS PROPOSED  
FOR FUSION FIRST-WALL APPLICATIONS**

METAL OR ALLOY	PER MWy/m <sup>2</sup>		
	dpa	at. ppm H	at. ppm He
ALUMINUM	14	296	316
TI-6 AL-4 V	16	175	142
F E R R I T I C STEEL (SANDVIK HT-9)	11	450	110
AUSTENITIC SS (316)	11	532	147
NICKEL-BASE ALLOY (NIMONIC PE-16)	12	780	240
V-15 Cr-5 Ti	11	245	47
NIOBIUM	7	105	29

Table III. Comparison of candidate first-wall structural materials

Requirement	Austenitic stainless steels	Ferritic steels	High-nickel alloys	Cb (Nb)	V	Mo	Ti	Al
Availability and cost	E	E	G	F	P	F	E	E
Fabricability and weldability	E	G	E	F	F	P	G	E
High-temperature mechanical properties	E	E	E	E	E	E	G	P
Thermal-stress resistance	F	E	F	E	E	E	G	E
Compatibility with:								
Lithium	G	G	P	E	E	E	E	P
Hydrogen	E	F	E	F	F	E	P	G
Fused salts	G	G	E	P	P	G	P	P
Helium (realistic purity)	E	E	E	P	P	G	G	E
Steam	E	G	E	P	P	P	E	G
Sodium or potassium	E	E	F	E	F	E	E	P
Air	E	G	E	P	P	P	G	E
Low tritium permeability	G	F	G	P	P	E	P	G
Low neutron activation	P	F	P	P	E	P	G	E

Note: E, excellent; G, good; F, fair; and P, poor.

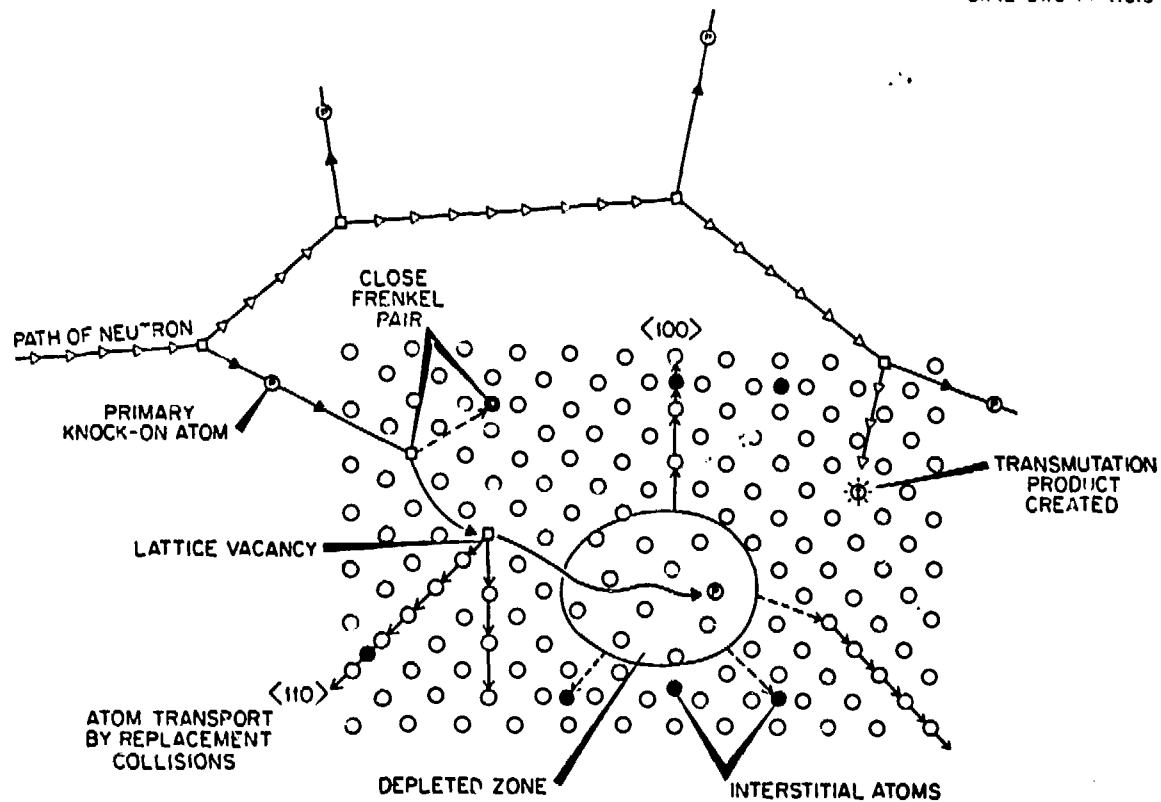


Fig. 2. Schematic diagram of neutron interactions with a metal lattice.

Fig. 3. INCONEL 600, EXPOSED IN HFIR REACTOR AT 650°C TO 8.5 DISPLACEMENTS-  
PER-ATOM LEVEL AND 1780 ATOM PARTS PER MILLION, FRACTURED AT ~~350°C~~

35°C



INCONEL 600

300 X

1780 appm He  
PRODUCED AT 650°C  
FRACTURE AT 35°C

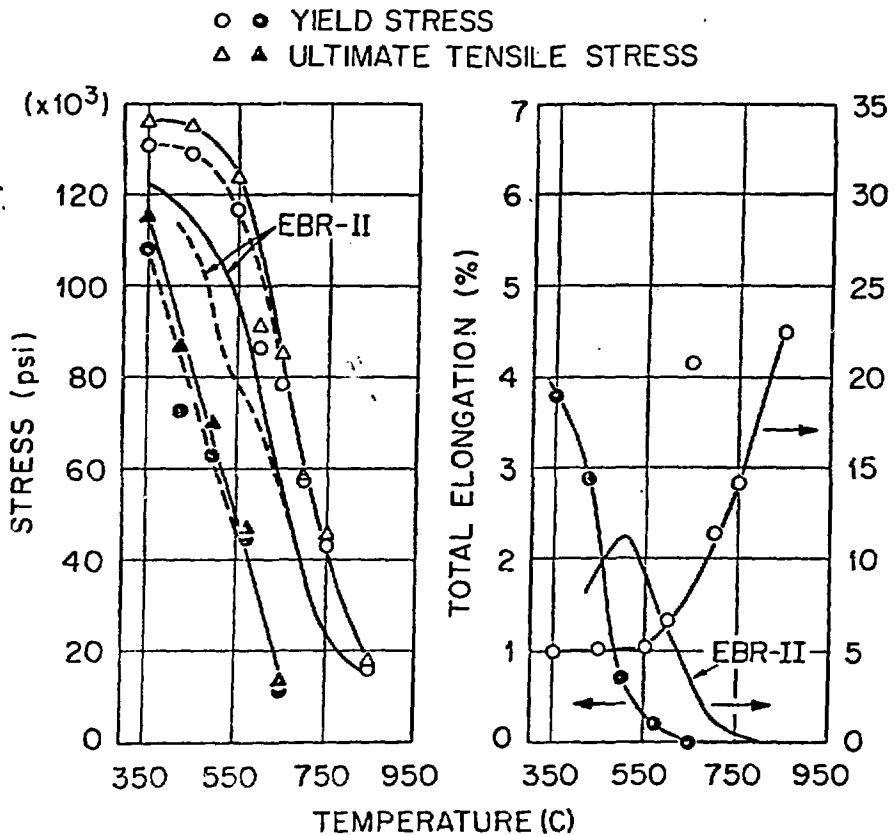


Fig. 4. Tensile Properties of 20 % Cold Worked Type 316 Stainless Steel. Open Symbols for Unirradiated Samples, Filled Symbols for HFIR Irradiated Samples ( $7.1$  to  $8.7 \times 10^{26} \text{ n/m}^2$ ) and EBR-II Trend Curves for Fluences  $1.2$  to  $2.9 \times 10^{26} \text{ n/m}^2$ .

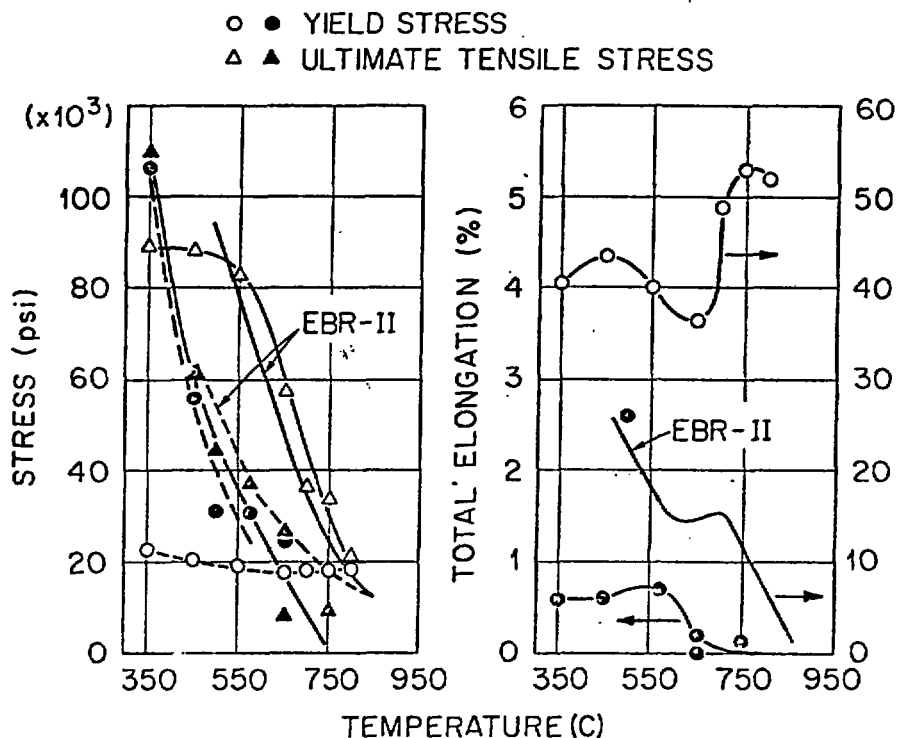


Fig. 5. Tensile Properties of Solution Annealed Type 316 Stainless Steel. Open Symbols for Unirradiated Samples, Filled Symbols for HFIR Irradiated Samples ( $5.6$  to  $8.7 \times 10^{26} \text{ n/m}^2$ ) and EBR-II Trend Curves for Fluences  $1.2$  to  $2.9 \times 10^{26} \text{ n/m}^2$ .

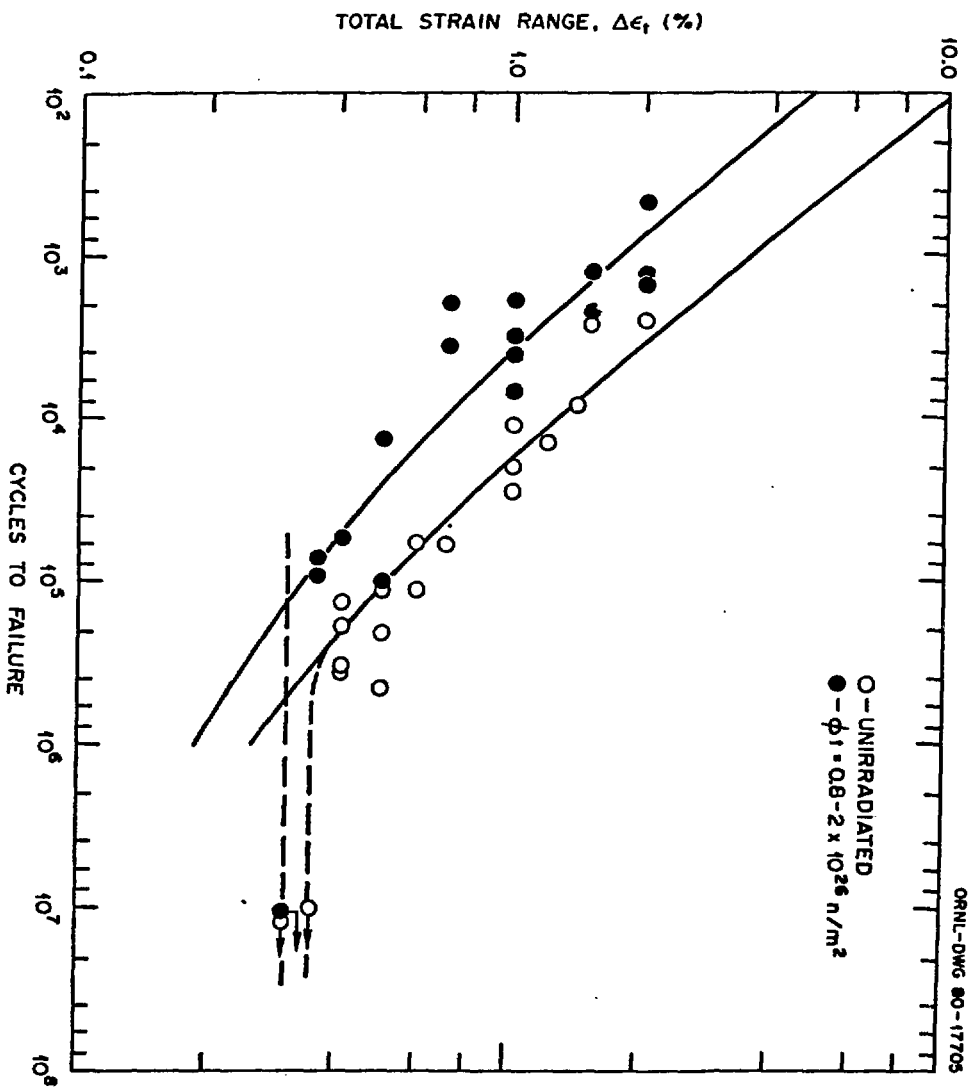


Fig. 6. Fatigue Life Of 20% C.W. Type 316 Stainless Steel Irradiated in HFIR At 430 °C And Tested At Irradiation Temperature (200-4,000 appm He, 5-15 dpa).

THE EFFECT OF COLD WORK  
316 Stainless Steel  
HFIR Irradiation

Y-143434



- 20% C.W. 316 SS
- Irradiated at 450°C, 3800 appm He, 53 dpa
- Swelling - 2.0%



- Annealed 316 SS
- Irradiated at 480°C, 3000 appm He, 42 dpa
- Swelling - 9.0%

0.1  $\mu\text{m}$

Fig. 7. Cold Work Reduces Cavity Size and Swelling for Irradiation in HFIR to High Helium Levels.

Table IV. Microstructural swelling data on irradiated type 316 stainless steels

Condition	Irradiation temperature (°C)	Displacement damage (dpa) <sup>a</sup>	Helium content (appm) <sup>b</sup>	Swelling (%)	Cavity concentration (No./cm <sup>3</sup> )	Diameter (Å)
20% Cold-worked	380	49	3320	2.2	$1.8 \times 10^{16}$	95
	450	54	3660	2.0	$6.6 \times 10^{15}$	170
	550	42	2990	1.4	$2.4 \times 10^{15}$	210
	600	60	4070	5.0	$3.3 \times 10^{14}$	647
	680	61	4140	16.8	$6.3 \times 10^{13}$	1100
Annealed	480	42	2950	8.8	$1.4 \times 10^{15}$	386
	550	42	2990	8.5	$4.4 \times 10^{15}$	500
	680	61	4140	15.2	$4.6 \times 10^{13}$	1083

<sup>a</sup>Displacements per atom.<sup>b</sup>Atom parts per million.

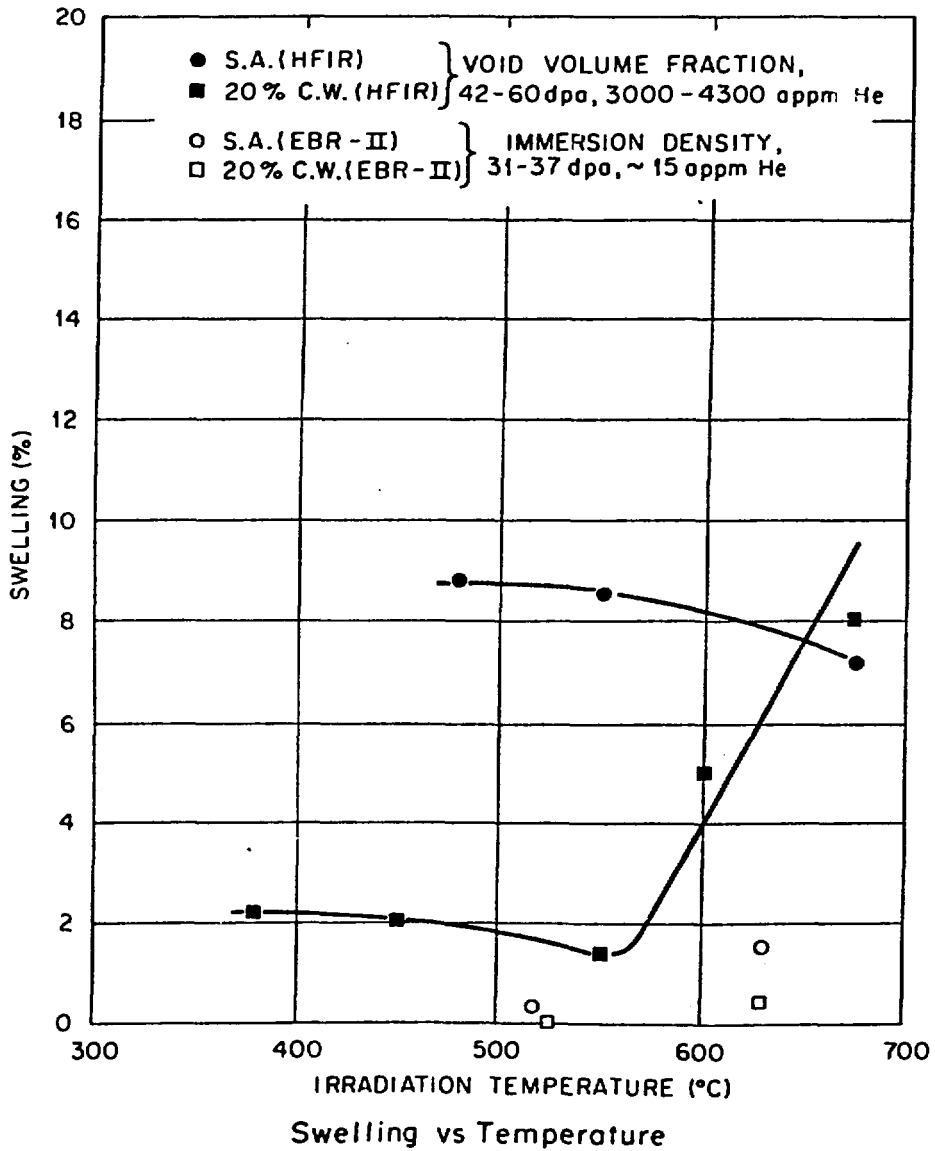


Fig. 8. A High Helium Content in 316 Stainless Steel Causes a High Swelling Rate.

E-7675  
YE-11482



Fig. 9. Solution-Annealed Type 316 Stainless Steel with 0.23 wt % Ti after Irradiation in HFIR at 600°C to 30 dpa and 1850 appm He.

Fig. 10. Helium Trapping at a TiC Particle in Type 316 Stainless Steel.

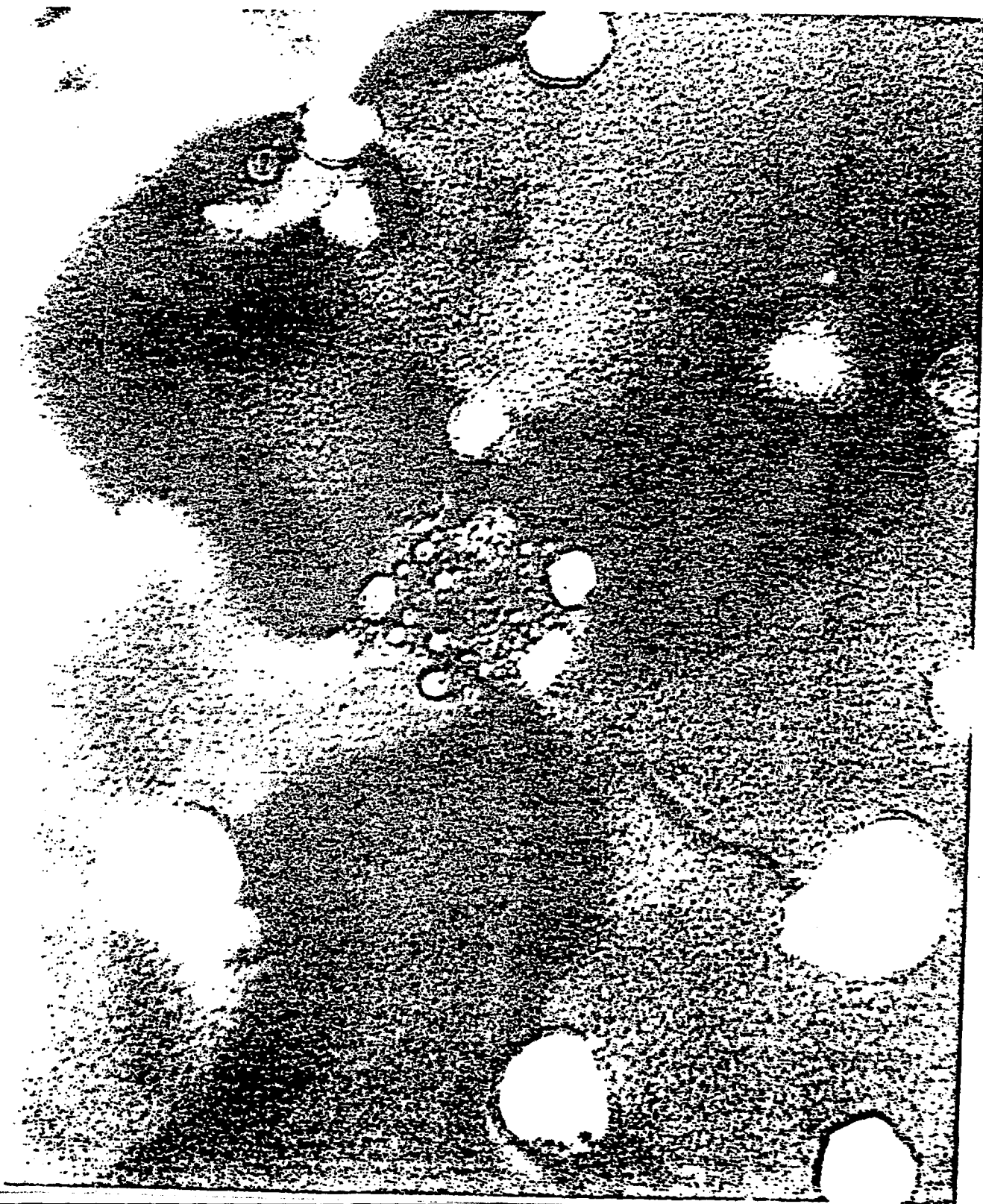


Fig. 11. Annealed Type 316 Stainless Steel after Irradiation in HEIR at 600°C to 30 dpa and 1850 appm He.

B-19172 YE-11481 3x

