

Higgs-photon associated production at hadron colliders

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Abstract

MASTER

We present cross sections for the reactions $p\bar{p} \rightarrow H\gamma$ and $pp \rightarrow H\gamma$ arising from the subprocess $q\bar{q} \rightarrow H\gamma$. The calculation includes the complete one-loop contribution from all light quarks and is the main source of Higgs-photon associated production in hadron colliders. At Tevatron energies, the cross section is typically 0.1fb or less, while at LHC energies it can exceed 1.0fb.

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I. INTRODUCTION

The production of intermediate mass Higgs bosons at hadron colliders arises primarily from the gluon fusion process $gg \rightarrow H$, which is dominated by the top quark loop [1]. Depending on the collider energy, there can also be a substantial contribution from the gauge boson fusion processes $W^+W^-, ZZ \rightarrow H$ [2]. For Higgs boson masses m_H less than $2m_W$, detection of the Higgs via its dominant decay mode $H \rightarrow b\bar{b}$ must contend with a large background from $gg \rightarrow b\bar{b}$. The production of Higgs bosons in association with a gauge boson, specifically Z, W^\pm [3,4] or a photon [5,6], is potentially helpful in dealing with these backgrounds, although at a significant cost in the rate.

In this report, we examine the possibility of using Higgs-photon associated production in hadron collisions as a means of studying properties of the Higgs boson. The main source of the $H\gamma$ final state is quark-antiquark annihilation. Two gluon annihilation, which accounts for single H production, is forbidden by Furry's theorem because the gluons are in a color singlet state. For light quarks, the direct annihilation into $H\gamma$ is suppressed by the ratio m_q/m_W and is negligible. This means that one-loop electroweak corrections involving W 's, Z 's and top quarks dominate $q\bar{q} \rightarrow H\gamma$. The diagrams involved in the quark-antiquark annihilation calculation are similar to those encountered in the calculation of $e\bar{e} \rightarrow H\gamma$, which has been performed for the Standard Model [5,6] and its supersymmetric generalization [7]. In the next section, we outline the extension of our previous results [6] to $q\bar{q} \rightarrow H\gamma$. This is followed by a discussion.

II. OUTLINE OF THE CALCULATION

The amplitude for $q\bar{q} \rightarrow H\gamma$ receives contributions from pole diagrams involving virtual photon and Z exchange and from various box diagrams containing quarks, gauge bosons and/or Goldstone bosons. There are also double pole diagrams whose contribution vanishes. This is illustrated in Fig. 1. The main difference between $e\bar{e}$ annihilation and $q\bar{q}$ annihilation occurs in the crossed box diagram of Fig. 1c. Since both members of a quark doublet are charged, there are additional crossed box diagrams containing W 's. This contribution can be obtained from our Z box results by merely changing the coupling and replacing m_Z by m_W .

In the non-linear gauges we chose [6], the full amplitude consists of four separately gauge invariant terms: a photon pole, a Z pole, Z boxes and W boxes. These amplitudes can be written as

$$\mathcal{M}_{\text{pole}}^\gamma = \frac{\alpha^2 m_W}{\sin\theta_W} \bar{v}(p_2) \gamma_\mu u(p_1) \left(\frac{\delta_{\mu\nu} k \cdot (p_1 + p_2) - k_\mu (p_1 + p_2)_\nu}{s} \right) \hat{e}_\nu(k) \mathcal{A}_\gamma(s), \quad (1)$$

$$\mathcal{M}_{\text{pole}}^Z = \frac{\alpha^2 m_W}{\sin^3\theta_W} \bar{v}(p_2) \gamma_\mu (v_q + \gamma_5) u(p_1) \left(\frac{\delta_{\mu\nu} k \cdot (p_1 + p_2) - k_\mu (p_1 + p_2)_\nu}{(s - m_Z^2) + im_Z \Gamma_Z} \right) \hat{e}_\nu(k) \mathcal{A}_Z(s), \quad (2)$$

$$\begin{aligned} \mathcal{M}_{\text{box}}^Z = & -\frac{\alpha^2 m_Z}{4 \sin^3\theta_W \cos^3\theta_W} \bar{v}(p_2) \gamma_\mu (v_q + \gamma_5)^2 u(p_1) \{ [\delta_{\mu\nu} k \cdot p_1 - k_\mu (p_1)_\nu] \mathcal{B}_Z(s, t, u) \\ & + [\delta_{\mu\nu} k \cdot p_2 - k_\mu (p_2)_\nu] \mathcal{B}_Z(s, u, t) \} \hat{e}_\nu(k), \end{aligned} \quad (3)$$

$$\mathcal{M}_{\text{box}}^W = \frac{\alpha^2 m_W}{2 \sin^3 \theta_W} \bar{v}(p_2) \gamma_\mu (1 + \gamma_5)^2 u(p_1) \{ [\delta_{\mu\nu} k \cdot p_1 - k_\mu (p_1)_\nu] B_W(s, t, u) + [\delta_{\mu\nu} k \cdot p_2 - k_\mu (p_2)_\nu] B_W(s, u, t) \} \hat{e}_\nu(k), \quad (4)$$

where $s = -(p_1 + p_2)^2$, $t = -(k + p_1)^2$ and $u = -(k + p_2)^2$. Here, v_q denotes the $q\bar{q}Z$ vector coupling constant, $v_q = 1 - 4|e_q| \sin^2 \theta_W$, and e_q is the quark charge in units of the proton charge. In terms of the scalar functions defined in the appendices of our previous paper [6], we have [8]

$$\mathcal{A}_\gamma(s) = -e_q \left\{ 4 \left(6 + \frac{m_H^2}{m_W^2} \right) C_{23}(s, m_H^2, m_W^2) - 16 C_0(s, m_H^2, m_W^2) - \frac{16}{3} \frac{m_t^2}{m_W^2} \left(4 C_{23}(s, m_H^2, m_t^2) - C_0(s, m_H^2, m_t^2) \right) \right\}, \quad (5)$$

$$\mathcal{A}_Z(s) = -2I_3 \left\{ \left(5 - \tan^2 \theta_W + \frac{m_H^2}{2m_W^2} (1 - \tan^2 \theta_W) \right) C_{23}(s, m_H^2, m_W^2) + (\tan^2 \theta_W - 3) C_0(s, m_H^2, m_W^2) - \frac{1}{2} \frac{m_t^2}{m_W^2} \frac{1 - (8/3) \sin^2 \theta_W}{\cos^2 \theta_W} \left(4 C_{23}(s, m_H^2, m_t^2) - C_0(s, m_H^2, m_t^2) \right) \right\}, \quad (6)$$

$$B_Z(s, t, u) = -\tilde{e}_q A(s, t, u), \quad (7)$$

$$B_W(s, t, u) = -2I_3 [A_1(s, t, u) + A_2(s, u, t)] + \tilde{e}_q A'(s, t, u), \quad (8)$$

with I_3 denoting the third component of the *external* quark weak isospin, m_t being the top quark mass and the prime denoting the replacement $m_Z \rightarrow m_W$. \tilde{e}_q is the charge of the *internal* quark in units of the proton charge.

The differential cross section $d\sigma(q\bar{q} \rightarrow H\gamma)/d\Omega_\gamma$ is

$$\frac{d\sigma(q\bar{q} \rightarrow H\gamma)}{d\Omega_\gamma} = \frac{1}{256\pi^2} \frac{s - m_H^2}{s^2} \sum_{\text{spin}} |\mathcal{M}|^2, \quad (9)$$

where the invariant amplitude \mathcal{M} is the sum of Eqs. (1-4). To obtain the collider cross sections, we convolute Eq. (9) with the appropriate quark and antiquark structure functions using

$$\frac{d\sigma}{dm_{H\gamma}^2} = \frac{1}{s} \int_\tau^1 \frac{dx}{x} f_q(x) f_{\bar{q}}(\tau/x) \int_{-z_0}^{z_0} dz \frac{d\sigma(q\bar{q} \rightarrow H\gamma)}{dz}. \quad (10)$$

Here, $\tau = m_{H\gamma}^2/s$, $z = \cos \theta_\gamma$ and $z_0 = \cos \theta_{\min}$ is determined by the choice of the rapidity cut. We used CTEQ-1M structure functions [9] and a rapidity cut of 2.5 on both the Higgs boson and the photon.

III. DISCUSSION

The total cross section for $p\bar{p}(pp) \rightarrow H\gamma$ is illustrated in Fig.2 for both Tevatron and LHC energies. At the Tevatron, it is clear that an upgrade in luminosity to 20 - 30 fb⁻¹/year

is necessary for $H\gamma$ to be seen. Luminosities of this order are contemplated in the TeV33 working group study [10]. For the LHC with a proposed luminosity of $100\text{fb}^{-1}/\text{year}$, associated production for $m_H < 200$ GeV will be observable.

In addition to $H\gamma$ production from light quark annihilation by means of one-loop amplitudes, there is the possibility of heavy quark annihilation at the tree level. As mentioned in the introduction, the tree level amplitudes are suppressed by a factor of m_q/m_W , which, while negligible for u and d quarks, could be significant for s and c quarks [11,12]. We specifically checked the contribution of the $s\bar{s}$ tree level amplitude and found that for $m_H < m_Z$, the one-loop $s\bar{s}$ amplitude dominates because of the Z pole. When $m_H > m_Z$, the tree level contribution exceeds the one-loop contribution. However, in this case, the s quark luminosity is so small that both contributions are negligible compared to that of the light quarks.

Finally, to determine the sensitivity of associated production to changes in the $t\bar{t}H$ coupling, we computed the cross section including a factor λ multiplying the Standard Model coupling [13]. For a 2 TeV Tevatron upgrade, this effect is illustrated in Fig. 3 for $0.0 \leq \lambda \leq 4.0$ as a function of m_H , and in Fig. 4 as a function of λ . The most obvious effect is the uniform decrease in the cross section as λ varies from $\lambda = 0$ to the Standard Model value at $\lambda = 1.0$. With increasing λ , the cross sections eventually exceed the Standard Model result. For larger values of m_H , the observation of $H\gamma$ events at a 2 TeV machine with a luminosity of $20\text{-}30\text{fb}^{-1}/\text{year}$ would necessitate a non Standard Model coupling.

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FIGURES

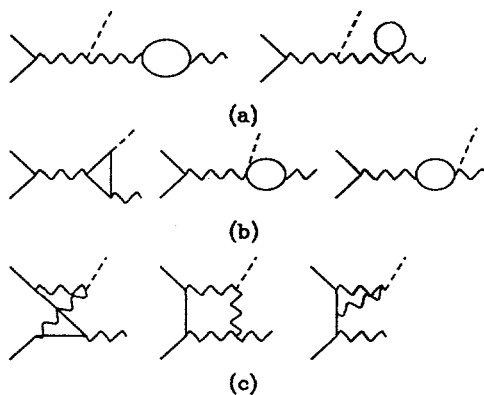


FIG. 1. Typical diagrams for the double pole (a), single pole (b) and box (c) corrections are shown. An external solid line represents a quark, a wavy line a gauge boson, a dashed line a Higgs boson and an internal solid line a quark, gauge boson, Goldstone boson or ghost.

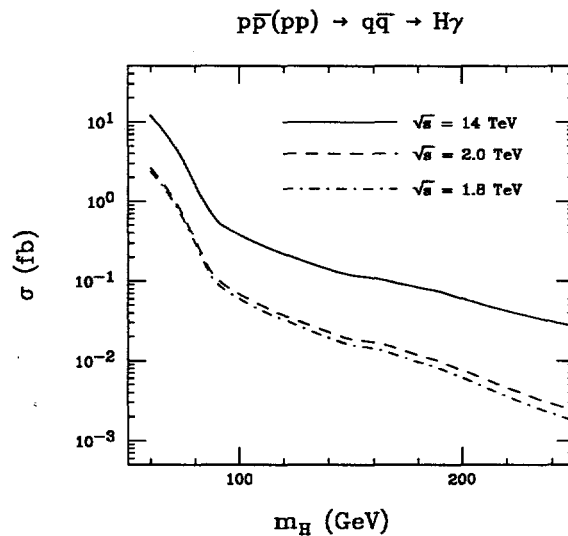


FIG. 2. The total cross section for $H\gamma$ production from $q\bar{q}$ annihilation in $p\bar{p}$ and pp scattering is shown for various collider energies. The dot-dashed line corresponds to the current Tevatron energy, the dashed line to an upgraded Tevatron energy and the solid line to the proposed LHC energy.

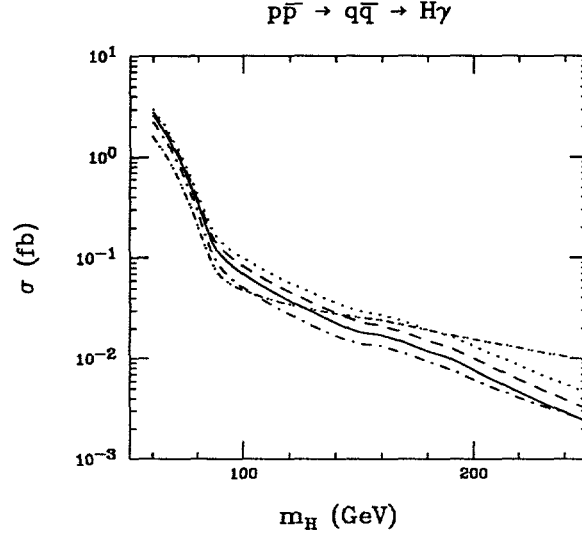


FIG. 3. The total cross section for $H\gamma$ production from $q\bar{q}$ annihilation in $p\bar{p}$ scattering is shown for varying strengths of the $t\bar{t}H$ coupling (in multiples λ of the Standard Model coupling) at an upgraded Tevatron. The solid line is $\lambda = 1$, the dot-dashed line $\lambda = 2$, the dot-dot-dashed line $\lambda = 4$, the dashed line $\lambda = .5$, and the dotted line $\lambda = 0$.

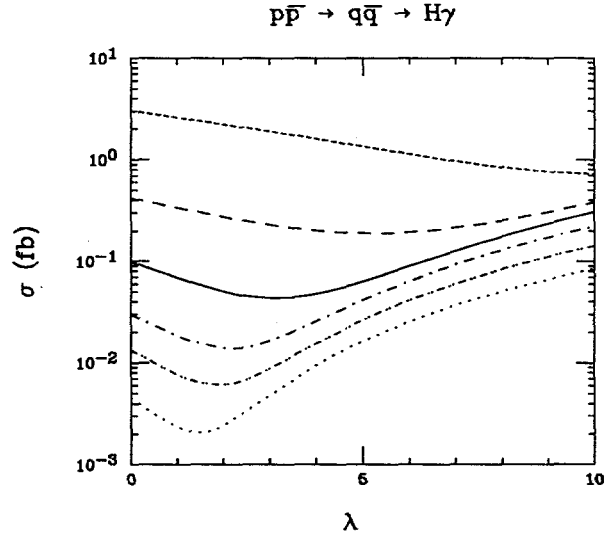


FIG. 4. The total cross section for $H\gamma$ production from $q\bar{q}$ annihilation in $p\bar{p}$ scattering is shown as a function of the $t\bar{t}H$ coupling (in multiples λ of the Standard Model coupling) at an upgraded Tevatron for several values of m_H . The short dashed line is $m_H = 60$ GeV, the dashed line $m_H = 80$ GeV, the solid line $m_H = 100$ GeV, the dot-dashed line $m_H = 150$ GeV, the dot-dot-dashed line $m_H = 200$ GeV and the dotted line $m_H = 250$ GeV.