

MASTER

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PHOTOIONIZATION CROSS SECTION FOR He IN THE
HYPERSPHERICAL COORDINATE METHOD

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In order to more fully explore the role of electron correlations in the photoionization process, we have employed the hyperspherical coordinate method of Macek¹ in calculating photoionization cross sections of He. While this system of coordinates has long been recognized as particularly suitable for the description of two-electron dynamics² it is only recently that quantitative calculations have been made. Previous applications have included predictions of doubly-excited-state energies in He and H⁻,^{3,5} and the treatment of two-electron atomic scattering.^{3,6} Ours is the first application of the method to a dipole process.

The hyperspherical coordinate system expresses the six-dimensional two-electron space in terms of a hyperspherical radius $R = (r_1^2 + r_2^2)^{1/2}$ and five angular variables. Following Macek,¹ we take these to be $\alpha = \tan^{-1}(r_2/r_1)$, \hat{r}_1 and \hat{r}_2 . In this set of coordinates, the two-electron Schrödinger equation is nearly separable. The wavefunction is written as

$$\Psi(R, \alpha, \hat{r}_1, \hat{r}_2) = (R^{5/2} \sin \alpha \cos \alpha)^{-1} \sum_u F_u(R) \phi_u(R; \alpha, \hat{r}_1, \hat{r}_2). \quad (1)$$

The eigenvalues of the angular equation, which are parametrically dependent upon R , yield channel potentials for the radial equation, and the eigenfunctions, ϕ_u , of the angular equation are used to obtain the channel-channel coupling for the radial equation.

Our first application of the hyperspherical method has been to calculate cross sections for photoionization of He using a single-channel, adiabatic approach, from threshold to 1 a.u. above threshold. In this energy range the residual ion is in its ground state. Our initial (final) -state wavefunction was obtained by using the lowest $^1S^e(^1P^0)$ potential curve along with the diagonal coupling-matrix element in the radial equation. As a measure of the degree of correlation contained in our wavefunctions, we observe that our ground state

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energy (-2.89517 a.u.) is quite close to the essentially exact result of Pekeris⁷ (-2.90372 a.u.) in comparison with the Hartree-Fock⁸ ground state energy (-2.86168 a.u.).

In Table 1 we compare our cross section results with those of other detailed calculations⁹⁻¹¹ as well as with revised experimental data of Samson,¹² which have error bars of $\pm 3\%$. Within the kinetic energy range $0.0 \leq \epsilon \leq 0.25$ a.u. our results are comparable to or slightly better than these other calculations. Above $\epsilon = 2.5$ a.u. the RPA⁹ and polarized orbital¹¹ calculations tend higher than experiment, while our results tend lower ($\sim 12\%$ lower at $\epsilon = 0.9$ a.u.). Close-coupling results¹⁰ are not given below $\epsilon = 0.1$ a.u. because of numerical problems near threshold, but above $\epsilon = 0.25$ a.u. they are in best agreement with experiment. Preliminary calculations indicate that when the hyperspherical channel corresponding to the $n = 2+$ state is included the cross section is practically unchanged near threshold, but is increased by about 10% over the single channel result at $\epsilon = 0.9$ a.u.

Table 1. Comparison of the Present Results with Some of the More Detailed Theoretical Calculations of the Photoionization Cross Section for the Process $\text{He} + \omega \rightarrow \text{He}^+(1s) + e^-$

Photoelectron Energy(a.u.) ^a	$\sigma(\text{Mb})^b$				
	Experimental results ^c	Present results ^d	RPA Method ^e	Close-Coupling Method ^f	Polarized Orbital Method ^g
0.0	7.56 ± 0.23	7.65	7.55	—	7.56 (7.84)
0.05	6.89 ± 0.21	6.96	7.00	—	7.00 (7.24)
0.1	6.41 ± 0.19	6.36	6.45	6.32 (6.25)	6.47 (6.69)
0.15	5.92 ± 0.18	5.81	6.00	—	—
0.2	5.52 ± 0.17	5.33	5.55	5.38 (5.34)	5.55 (5.72)
0.25	5.02 ± 0.15	4.88	5.18	—	—
0.3	4.63 ± 0.14	4.48	4.85	4.61 (4.59)	4.79 (4.92)
0.4	3.94 ± 0.12	3.81	4.28	3.98 (3.97)	4.14 (4.25)
0.5	3.41 ± 0.10	3.26	3.75	3.47 (3.46)	3.61 (3.70)
0.7	2.57 ± 0.08	2.43	2.87	2.68 (2.67)	2.79 (2.85)
0.9	2.10 ± 0.06	1.85	2.23	2.13 (2.11)	2.19 (2.23)

^a1 a.u. = 27.2108 eV. ^bTheoretical results are given in dipole length approximation. When available, dipole velocity results are given in parenthesis.

^cRef. 12. Linear interpolation of the densely spaced experimental results (see Fig. 1) is used to provide values at the same energies for which theoretical

results are available. ^dThe diagonal non-adiabatic coupling terms are included in calculating both initial and final state wavefunctions. ^eRef. 9. Numerical values were supplied to us by the author. ^fRef. 10, Table IV. Final state is obtained from a $1s - 2\bar{s} - 2\bar{p}$ expansion. Initial state is the 56-term Pekeris variational wavefunction. ^gRef. 11, Table 2. Initial state is the 2-term Hart-Herzberg variational wavefunction.

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