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M.B. Brodsky, R.J. Trainor, A.T. Aldred and C.H. Sowers

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## MAGNETIC SUSCEPTIBILITY OF $\text{AnRe}_2$ COMPOUNDS\*

M.B. Brodsky, R.J. Trainor,† A.T. Aldred and C.H. Sowers  
Materials Science Division  
Argonne National Laboratory  
Argonne, IL 60439

### ABSTRACT

The compounds  $\text{AnRe}_2$ , where  $\text{An} = \text{Th, U, Np, and Pu}$  all have the  $\text{MgZn}_2$ -type hexagonal Laves phase structure. Measurements of their magnetic susceptibilities have been made from 2-300 K. Whereas  $\text{ThRe}_2$ ,  $\text{URe}_2$ , and  $\text{PuRe}_2$  have essentially temperature-independent susceptibilities,  $\text{NpRe}_2$  is ferromagnetic with  $T_c = 47$  K. The paramagnetic data for  $\text{NpRe}_2$  follow a modified Curie-Weiss law,  $\chi - \chi_0 = C/(T - \Theta)$  from 60-230 K with  $\Theta = 46.9$  K. and  $C$  equivalent to an effective moment,  $P_{\text{eff}} = 2.11 \mu_B$ . The ordered moment  $\mu_0 = 0.90 \mu_B$  and the results are consistent with other actinide compounds having good local moment behavior. The electronic specific heat of  $\text{URe}_2 = 34 \text{ mJ}/(\text{mole-K}^2)$  is large enough to indicate the possibility of spin-fluctuation behavior in this compound.

### INTRODUCTION

Many actinide intermetallic compounds with composition  $\text{AnX}_2$  form Laves phase-type structures. The largest number of these occur in the  $\text{MgCu}_2$ -type cubic phase structure. Systematic studies of the magnetic properties of  $\text{MgCu}_2$ -type compounds have been made, and reasonable correlations between magnetic ordering behavior and  $\text{An-An}$  separation have been presented for cases where  $\text{X}$  is not a  $3d$  transition element [1]. Although a number of  $\text{An}$  Laves-phase compounds also form in the hexagonal  $\text{MgZn}_2$  structure, magnetic data have been reported only for  $\text{UNi}_2$ , which orders ferromagnetically at  $T_c = 30$  K and  $\mu_0 = 0.13 \mu_B$  [2]. This paper presents the results of a magnetization study of  $\text{AnRe}_2$  compounds where  $\text{An} = \text{Th, U, Np and Pu}$ . All of these compounds form the  $\text{MgZn}_2$  structure. Although  $\text{URe}_2$  does form in the  $\text{MgZn}_2$  structure above 453 K, below that temperature it transforms to an orthorhombic structure which is a distortion of the  $\text{MgZn}_2$  structure at lower temperatures. In view of the similarity of the low temperature structure to the other  $\text{AnRe}_2$  structures, the results for  $\text{URe}_2$  are also included.

### RESULTS

Figure 1 shows the temperature dependence of the molar susceptibility for the compounds between 2-300 K.

The susceptibilities are essentially temperature-independent for  $\text{ThRe}_2$ ,  $\text{URE}_2$  and  $\text{PuRe}_2$ . However, the data for  $\text{NpRe}_2$  are indicative of ferromagnetic ordering, and these results are replotted as  $1/\chi$  versus  $T$  in Fig. 2. The results nearly fit a Curie-Weiss law. The curvature may be due to a temperature-independent term,  $\chi_0$ , or to crystal-field effects. Both approaches have been used previously to explain curvature in  $1/\chi$  versus  $T$  plots for actinides. In the absence of a crystal-field calculation for this structure, no attempt is made here to evaluate the crystal-field effects.

Figure 2 also shows the effect of assuming  $\chi_0 = 1.21 \times 10^{-3}$  emu/mole and a straight line is obtained.

The modified Curie-Weiss law,  $\chi - \chi_0 = C/(T - \theta)$ , for  $\text{NpRe}_2$  yields  $\theta = +47.4$  K and  $p_{\text{eff}} = 2.11 \mu_B$ . The existence of ferromagnetic ordering at 47 K is supported further by the magnetization data in the vicinity of  $T_c$ . These are plotted in Fig. 3 as  $\sigma^2$  versus  $H/\sigma$  and yield  $T_c = 47$  K.

#### DISCUSSION

The  $\text{NpRe}_2$  data in Fig. 2 are strongly indicative of local moment behavior. Below  $T_c$ , plots of  $\sigma$  versus  $T^{3/2}$ ,  $\sigma$  versus  $T^2$  and a Brillouin function for  $J = 1/2$  all yield  $\mu_0 = 0.899 \pm .002 \mu_B$  upon extrapolation to 0 K. Simple Russell-Saunders coupling would indicate  $0.84 \mu_B$  and  $2.70 \mu_B$  for  $p_{\text{eff}}$  of  $5f^5$  and  $5f^4$  configurations, respectively. In the well documented case of NaCl-type actinide compounds [3] intermediate coupling which involves full  $j-j$  manifold interactions, the effective moments are 1.41 and 2.59 for  $5f^5$  and  $5f^4$ , respectively. The ratio  $\mu_0/p_{\text{eff}} = 0.44$  for  $\text{NpRe}_2$  is in the range found for NaCl-type compounds. These considerations are very qualitative in the absence of crystal-field calculations for actinides with the  $\text{MgZn}_2$  structure. Whereas no attempt will be made to use the results to assign a ground state configuration to the Np atom, the results are all suggestive of good 5f local moment character in  $\text{NpRe}_2$ . This is in contrast to  $\text{NpOs}_2$  (cubic Laves phase) which is an itinerant ferromagnet [4,5].

The small break in the data for  $\text{NpRe}_2$  at 250 K (best seen in the  $1/\Delta\chi$  versus  $T$  plot in Fig. 2) may be due to a lattice distortion of the type found in  $\text{URE}_2$  at 453 K. Similar effects are found in the susceptibility data for the cubic Laves phase compounds  $\text{NpOs}_2$ ,  $\text{NpRu}_2$  and  $\text{NpIr}_2$  [4]. The room temperature x-ray diffraction pattern for the  $\text{URE}_2$  sample used here showed the presence of both the hexagonal Laves phase (major) and the low-temperature orthorhombic phase (minor). However, the absence of a temperature-dependent susceptibility or of a strongly temperature-dependent magnetic

impurity contribution indicate that both  $\text{URe}_2$  structures are nonmagnetic.

The strong magnetism of the Np compound, in comparison to the U and Pu compounds, is not uncommon. Consider the sequence  $\text{ThAl}_2$ ,  $\text{UAl}_2$ ,  $\text{NpAl}_2$  and  $\text{PuAl}_2$ .  $\text{ThAl}_2$  is nonmagnetic;  $\text{UAl}_2$  is a spin fluctuation compound [6];  $\text{NpAl}_2$  is a good ferromagnet with  $T_c = 56$  K [4]; and  $\text{PuAl}_2$  may be magnetic but only below 4 K [7]. Similarly,  $\text{USn}_3$  is likely to be a spin fluctuation compound [8];  $\text{NpSn}_3$  is an itinerant antiferromagnet [9]; and  $\text{PuSn}_3$  is nonmagnetic [8]. These results may be understood in terms of a strong spin-orbit splitting. When that splitting is larger than the 5f band-width, the 5f band splits into  $j = 5/2$  and  $j = 7/2$  sub-bands [10]. This approach can explain the lack of magnetism in Am, as well as the stronger magnetism for Np versus U and Pu in some cases.

Most certainly  $\text{ThRe}_2$  is nonmagnetic and a superconductor due to the lack of 5f electrons below the Fermi level [11]. Despite the temperature-independent paramagnetism of  $\text{URe}_2$  and  $\text{PuRe}_2$ , they may still exhibit spin-fluctuation behavior. The specific heat of the  $\text{URe}_2$  sample measured at low temperatures yields an electronic term  $\gamma = 34$  mJ/(mole-K<sup>2</sup>). It is observed that all actinide compounds having  $\gamma \geq 20$  mJ/(mole-K<sup>2</sup>) are magnetic or show spin fluctuation behavior. The simplest way to show this would be via a resistivity  $\propto T^2$  at low temperatures. However, the two-phase nature of  $\text{URe}_2$  samples currently available makes this approach suspect. The possibility of spin fluctuation effects may be tested in  $\text{PuRe}_2$ , which is single phased. However, the lower susceptibility of  $\text{PuRe}_2$ ,  $[\chi(\text{PuRe}_2)/\chi(\text{URe}_2)] = 0.5$ , makes it a less likely candidate for spin fluctuations than is  $\text{URe}_2$ .

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†Permanent Address: Lawrence Livermore Laboratory, Livermore, CA 94550.

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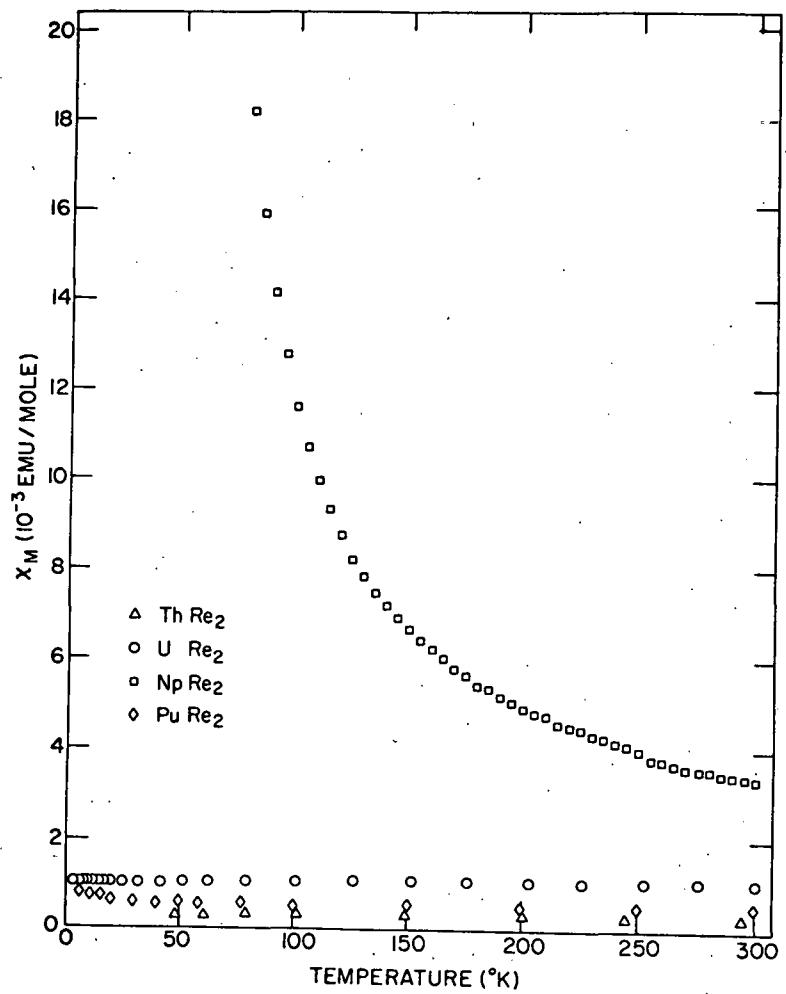


Fig. 1. Molar susceptibilities vs. temperature for AnRe<sub>2</sub> compounds.

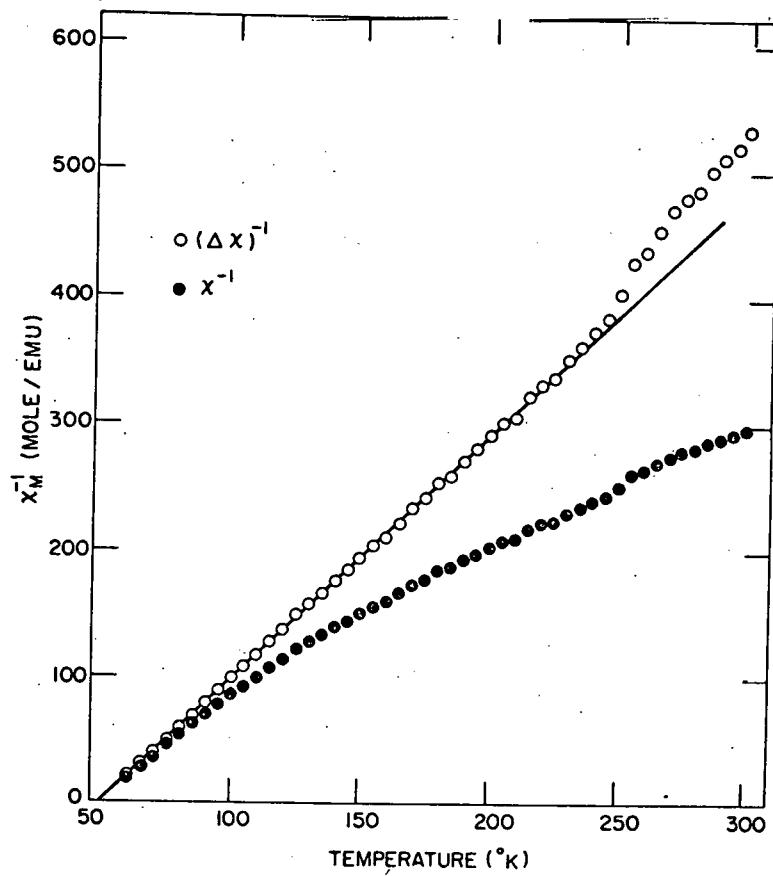


Fig. 2. Inverse molar susceptibility and corrected susceptibility vs. temperature for  $\text{NpRe}_2$ .

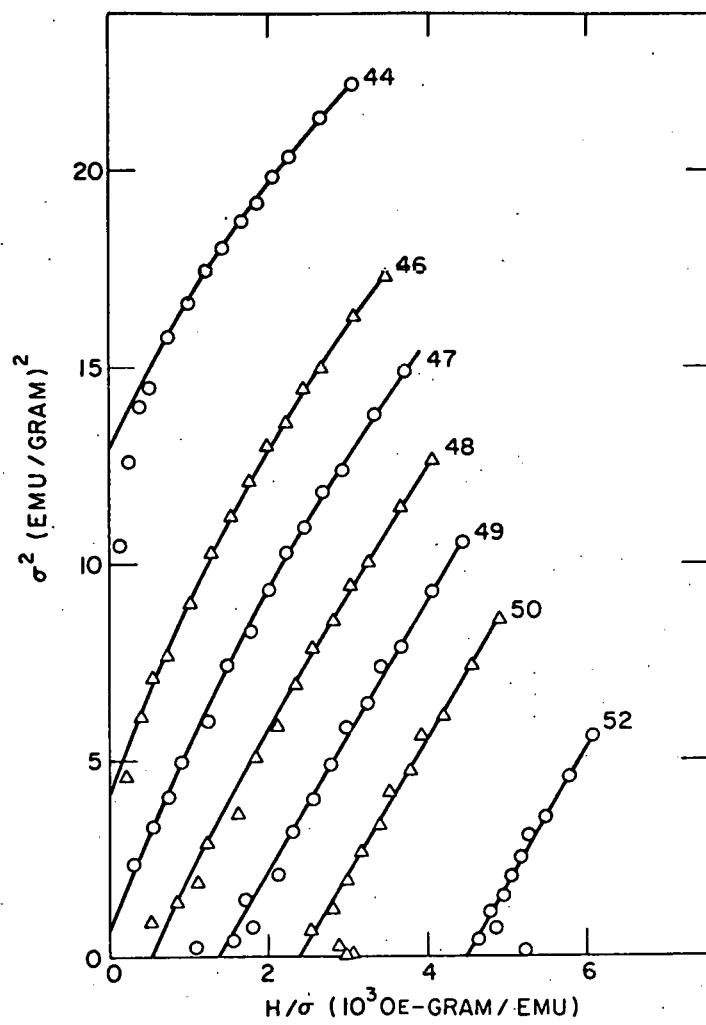


Fig. 3. Arrott plots for NpRe<sub>2</sub>.