

Columbia University
in the City of New York

NUMERICAL STUDIES OF SHOCK FOCUSING
IN A COAXIAL ELECTROMAGNETIC SHOCK TUBE

by

Hung Chang Lui

1973



Plasma Laboratory

School of Engineering and Applied Science

Columbia University

New York, N.Y. 10027

MASTER

DISCLAIMER

This report was prepared as an account of work sponsored by an agency of the United States Government. Neither the United States Government nor any agency Thereof, nor any of their employees, makes any warranty, express or implied, or assumes any legal liability or responsibility for the accuracy, completeness, or usefulness of any information, apparatus, product, or process disclosed, or represents that its use would not infringe privately owned rights. Reference herein to any specific commercial product, process, or service by trade name, trademark, manufacturer, or otherwise does not necessarily constitute or imply its endorsement, recommendation, or favoring by the United States Government or any agency thereof. The views and opinions of authors expressed herein do not necessarily state or reflect those of the United States Government or any agency thereof.

DISCLAIMER

Portions of this document may be illegible in electronic image products. Images are produced from the best available original document.

NUMERICAL STUDIES OF SHOCK FOCUSING
IN A COAXIAL ELECTROMAGNETIC SHOCK TUBE

Hung Chang Lui

NOTICE

This report was prepared as an account of work sponsored by the United States Government. Neither the United States nor the United States Atomic Energy Commission, nor any of their employees, nor any of their contractors, subcontractors, or their employees, makes any warranty, express or implied, or assumes any legal liability or responsibility for the accuracy, completeness or usefulness of any information, apparatus, product or process disclosed, or represents that its use would not infringe privately owned rights.

Submitted in Partial fulfillment
of the requirements for
The Degree of
Doctor of Engineering Science
in the Faculty of Engineering & Applied Science
Columbia University
1973

MASTER

DISTRIBUTION OF THIS DOCUMENT IS UNLIMITED

Fig

ABSTRACT

The generation of a strong transverse MHD shock wave in a coaxial electromagnetic shock tube is simulated numerically. Further plasma heating is achieved by shock reflection and focusing on the end wall. The details of the two-dimensional axisymmetric flow in the shock tube and the process of shock focusing and reflection from a curved end wall are studied. In general, the drive currents can be arbitrary functions of time, and the end wall can be of arbitrary shape.

The scheme used is the fluid-in-cell method (FLIC), modified for compressible, time-dependent magnetohydrodynamic flows to include a magnetic field perpendicular to the flow directions. The governing equations are the two-dimensional axisymmetric, ideal, single fluid MHD equations coupled with Maxwell's equations. Finite differencing is carried out for cells of variable shapes to accommodate solid walls and free surfaces. On a moving vacuum-plasma interface, a mixed Lagrangian-Eulerian procedure is used.

To optimize the wall shape for shock focusing, the location of the focused region and the wall shape are selected according to Whitham's shock-ray theory. The comparisons of the computed results between different shapes of the end wall are made to maximize the plasma heating. The most efficient shock focusing and plasma heating occurs with a quarter-elliptic end wall, giving a reflected and focused shock temperature some 8 times the incident post-shock temperature, or more than twice that obtained from a flat wall reflection. The corresponding density

in the focused region is 16 times the initial density in the shock tube.

To check the accuracy of the method, computations for a plane shock wave reflecting from a semi-circular wall in two-dimensional plane flows are also made using both the FLIC method and the two-step Lax-Wendroff scheme. The two methods agree well.

ACKNOWLEDGEMENTS

I am deeply indebted to my thesis advisors, Prof. C.K. Chu and Prof. R.T. Taussig, for their constant instruction and enlightening discussions on this research.

I am also grateful to Prof. R.A. Gross and Dr. Y.G. Chen for their help in supplying experimental data and measurements for the shock tube. To Dr. W.P. Gula and Dr. M.S. Chu, I express my thanks for their generosity in making time available for discussing computer programming at the beginning of this research. It is my pleasure to thank all members of the Plasma Physics Department for their cooperation.

I wish to thank my wife, Wenfang Wang, for her encouragement and support.

I also acknowledge the help of Mrs. Ruth Williams in correcting any grammatical errors in the original manuscript, and the expert typing of this report by Mary Black.

The research was supported by the United States Atomic Energy Commission under Contract AT(11-1)3086.

TABLE OF CONTENTS

Abstract	i
Acknowledgements	iii
I. Introduction	1
II. The FLIC MHD Code	6
2.1 The Governing Equations	6
2.2 The Computing Cells	9
2.3 Summary of the FLIC Method	14
2.4 Finite Difference Equations	16
III. Boundary Conditions	26
3.1 Introduction	26
3.2 On A Vacuum-plasma Interface	27
3.3 On A Curved Rigid Wall	36
IV. Incident Shock Waves	40
4.1 General Description	40
4.2 Constant Drive Current	43
4.3 Sinusoidal Drive Current	53
V. Shock Reflection and Convergence	60
5.1 Flat End Wall	60
5.1.1 Constant Driver	60
5.1.2 Sinusoidal Driver	70
5.2 Semi-circular End Wall	75
VI. Shock-Ray Approximation in Transverse MHD Flows	83
VII. Shock Focusing and Reflection	90

7.1 Quarter-Elliptic End Wall	90
7.2 Summary	102
VIII. Plane Flows	105
References	114
Appendice	
A. The Two-Step Lax-Wendroff Scheme	116
B. The Computer Program for Shock Focusing and Reflection from the Quarter-Elliptic End Wall	119

1. INTRODUCTION

Experimental studies of strong magnetically driven shock waves in a coaxial electromagnetic shock tube with a transverse field have been conducted for the past several years at Columbia University¹⁻⁴. At the same time, numerical simulations of the MHD shock tube with artificial dissipation⁵ and physical dissipation⁶ in one-dimensional flows were also carried out. These studies are analogous to the Hain-Roberts code for pinches⁷. However, since the nature of the flow in the coaxial shock tube is obviously two-dimensional, it is important to study the flow two-dimensionally. A numerical computation of the problem will give us interesting and useful insight into the details of the flow in the tube. It is our main purpose to investigate plasma heating numerically by shock focusing and reflection from curved end walls for temperatures of fusion interest.

The general experimental setup of the shock tube is shown schematically in Fig. 4-1. The working gas fills the space of the cylindrical annulus. A bias current flows in the central rod of the inner conductor to give an azimuthal field ahead of the shock. When the circuit is closed, the capacitor banks supply a radial current between two cylindrical electrodes passing through the working gas. This current sheet heats and pushes the gas ahead, creating a shock wave that runs down the tube. To reflect the shock wave, the end wall of the tube can be designed to form a proper profile so as to produce dense plasma heating in kilovolts by shock focusing and reflection.

A number of numerical codes (see e.g. 8,9) of multi-dimensional, time-dependent, compressible flows in hydrodynamics have been developed and applied successfully in the past. Yet it is not easy to apply them directly to MHD flows with arbitrary wall shapes or with a vacuum interface. The two-step Lax-Wendroff scheme was applied by Potter¹⁰ in plasma focus studies, in which he handled the moving vacuum-plasma interface by a minimum-density control. But this technique cannot be readily applied to MHD flows, because magnetic fields exist in both the vacuum and plasma regions, and the computed magnetic fields near the vacuum boundary will become completely inaccurate. If not properly tailored, the Lax-Wendroff scheme¹¹ may also induce numerical instabilities on a curved wall boundary. It has been suggested by many authors, e.g. Lapidus¹², to transform the curved wall into a straight boundary. Such variable transformations are usually very laborious to carry out even in simple geometries, and it may become prohibitive in more complicated geometries. The particle-in-cell method (PIC)¹³ and the grid-and-particle method (GAP)¹⁴ use particle-control in a surface cell to move the free boundary under an applied pressure, but they invariably give large fluctuations to the calculated results.

We have chosen to use the fluid-in-cell method (FLIC)¹⁵, which is originally an Eulerian finite difference scheme with mass transport differencing for solving unsteady compressible flow problems¹⁶. But, since the system of hydrodynamic equations of a perfectly conducting fluid in the presence of a

transverse magnetic field is reducible to that of ordinary fluid dynamical equations, we adapt the method rather easily. Moreover, as the magnetic field is frozen into the plasma, it can be treated in much the same way as the density.

But in the actual calculations of a plasma flow in MHD approximations, it may require some special procedures to handle boundary conditions on a curved rigid wall boundary and on a moving vacuum-plasma interface. Also for curved walls most methods to improve accuracy will result in tiny cells on the boundary. It is troublesome to handle these tiny irregular cells in the actual computations, and they cause stability problems. Since the moving vacuum-plasma interface is Lagrangian in nature, the Eulerian treatments mentioned above are less accurate. A mixed Lagrangian-Eulerian procedure is taken to handle the free boundary. Therefore, we extend the method to cells of variable shapes on solid boundaries and vacuum interfaces.

To choose an end wall shape for shock focusing, we use Whitham's shock-ray theory^{17,18}. This is an approximate procedure in which the particle paths are approximated by rays perpendicular to the shock surfaces. It has been applied to various gas dynamical problems by numerous authors such as Milton¹⁹, Lau²⁰, and Skews²¹. We extend this method to transverse MHD flows and use it to select end wall shapes. One of the best shapes of the end wall is a quarter-ellipse.

Chapter II summarizes the basic MHD equations and Maxwell's equations for the assumed physical model, and describes the modified fluid-in-cell method for MHD flows with a transverse

magnetic field. The method is designed particularly for surface cells of different shapes having irregular neighbor cells; but usually, it is also applicable to the regular inner cells located inside the computational region.

Chapter III describes the numerical boundary conditions associated with the code. The Lagrangian contour is introduced to handle the free surface and appropriate reflective boundary conditions are applied to curved walls. The cell-mixing procedures to increase stability are given as well.

In Chapter IV the results of shock formation in a coaxial shock tube are given. The shapes of the incident shocks are obtained. While the shock is almost planar, the current sheet is strongly inclined. Hence, one dimensional results are inadequate for their description.

In Chapter V the shock waves are utilized to study their reflection from flat end walls. The solution of shock convergence and reflection from a semi-circular end wall is also obtained. Usually, a shock wave will be diffracted and converged by a curved end wall when it moves along the wall surface.

In Chapter VI, Whitham's shock-ray approximation is reviewed, and extended to transverse MHD flows. End wall shapes are selected for shock focusing and reflection.

In Chapter VII, shock focusing and reflection from a quarter-elliptic end wall is presented. The shock wave will be strengthened by the elliptic end wall to maximum limit and will reflect. Of the three shapes studied, the quarter-elliptic end wall is the most efficient shape for plasma heating, giving a reflected

and focused shock temperature some 8 times the incident post-shock temperature, or more than twice that obtained from a flat wall reflection.

In Chapter VIII, we test the FLIC method by studying shock convergence and reflection from a semi-circular wall in plane transverse MHD flows, and comparing it with results from the two-step Lax-Wendroff scheme²². The agreement between the two methods is very good, which gives us some confidence in the accuracy of our basic code.

II. THE FLIC MHD CODE

2.1 The Governing Equations

The plasma in the shock tube is assumed to be a perfect gas with infinite electrical conductivity. The fluid equations under the MHD approximation are chosen for the plasma flows. Although a two-fluid model has been used in one-dimensional flows⁶, we do not feel it is easily extended to two-dimensional flows. In this study, therefore, we shall content ourselves to study only the geometrical aspects of shock focusing and plasma heating and a one fluid model is much simpler, at least as a first step.

In the simplest model, the plasma is considered as an ideal, single fluid of density ρ , pressure P , specific internal energy l , specific heat ratio γ , and velocity $\vec{U}$, carrying a current $\vec{J}$, in a magnetic field $\vec{B}$ and electrical field $\vec{E}$. The electrostatic forces and the displacement currents are negligible. The viscosity and heat conductivity are also neglected in the plasma. Maxwell's equations in Gaussian units become

$$\nabla \cdot \vec{B} = 0, \quad (2.1)$$

$$\nabla \cdot \vec{J} = 0, \quad (2.2)$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}, \quad (2.3)$$

$$\nabla \times \vec{B} = 4\pi \vec{J} / c. \quad (2.4)$$

Ohm's law is

$$\vec{E} + \frac{1}{c} \vec{U} \times \vec{B} = 0. \quad (2.5)$$

The fluid equations are

$$\left(\frac{\partial}{\partial t} + \vec{U} \cdot \nabla \right) \rho = -\rho \nabla \cdot \vec{U}, \quad (2.6)$$

$$\rho \left(\frac{\partial}{\partial t} + \vec{U} \cdot \nabla \right) \vec{U} = -\nabla P + \vec{J} \times \vec{B} / c, \quad (2.7)$$

$$\rho \left(\frac{\partial}{\partial t} + \vec{U} \cdot \nabla \right) I = -P \nabla \cdot \vec{U}, \quad (2.8)$$

$$P = (\gamma - 1) \rho I. \quad (2.9)$$

Suppose that the flow is plane ($\frac{\partial}{\partial z} = 0$) or axisymmetric ($\frac{\partial}{\partial \theta} = 0$) and the velocity is given initially to be in the (x-y) plane or (z-r) plane, and the magnetic field is initially perpendicular to it:

$$\begin{aligned} \vec{U} &= (U_x, U_y, U_z) \text{ or } (U_z, U_r, U_\theta) \\ &= (u, v, 0), \\ \vec{B} &= (B_x, B_y, B_z) \text{ or } (B_z, B_r, B_\theta) \\ &= (0, 0, B), \quad (\text{at } t=0) \end{aligned} \quad (2.10)$$

then such a flow will remain two-dimensional (plane (x-y) or axisymmetric (z-r)) at all times,

$$\begin{aligned} \vec{U} &= (u, v, 0), \\ \vec{B} &= (0, 0, B). \quad (\text{for all } t) \end{aligned} \quad (2.11)$$

Since the magnetic field is in the z-direction or θ -direction, the acceleration of the plasma and the electrical field ($\vec{E} = -\frac{1}{c} \vec{U} \times \vec{B}$) in that direction will be zero. The statement (2.11)

easily follows from Eqs. (2.3) and (2.7).

For two-dimensional plane or axisymmetric geometry, Eqs. (2.3) to (2.8) can be expressed in the form, with $\vec{U} = (u, v, 0)$ and $\vec{B} = (0, 0, B)$.

$$\frac{\partial \rho}{\partial t} = - \operatorname{div} (\rho \vec{U}), \quad (2.12)$$

$$\rho \left(\frac{\partial}{\partial t} + \vec{U} \cdot \nabla \right) u = - \frac{\partial}{\partial z} (P + B^2/8\pi), \quad (2.13)$$

$$\rho \left(\frac{\partial}{\partial t} + \vec{U} \cdot \nabla \right) v = - \frac{\partial}{\partial r} (P + B^2/8\pi) - Cb \frac{B^2}{4\pi r}, \quad (2.14)$$

$$\rho \left(\frac{\partial}{\partial t} + \vec{U} \cdot \nabla \right) I = - P \operatorname{div} (\vec{U}), \quad (2.15)$$

$$\frac{\partial}{\partial t} (B / r^m) = - \operatorname{div} (\vec{U} B / r^m), \quad (2.16)$$

and

$$\operatorname{div} (\vec{U}) = \frac{1}{r^m} \frac{\partial}{\partial r} (r^m v) + \frac{\partial u}{\partial z}, \quad (2.17)$$

where Cb and m are equal to 1 in an axisymmetric case; equal to 0 in a plane case.

We assume that the gas just in front of a strong shock wave is already preionized. The MHD equations and Maxwell's equations apply everywhere. Actually, after a shock reflection from the end wall, the temperature of plasma is about 10 Kev, and the ion cyclotron frequency becomes larger than the ion collision frequency. The fluid approximation there is poor. We shall continue to use the fluid approximation in this study, but will be aware of these possible defects.

2.2 The Computing Cells

In order to carry out the computation, the volume containing the fluid is subdivided into a number of cells. The cells in the interior region are equal in shape, and we will call them regular cells. The cells on the boundaries are irregular in shape, and we will call them irregular cells. The computational cells are shown in Fig. 2-1. The location of a cell is denoted by the indices (i, j) .

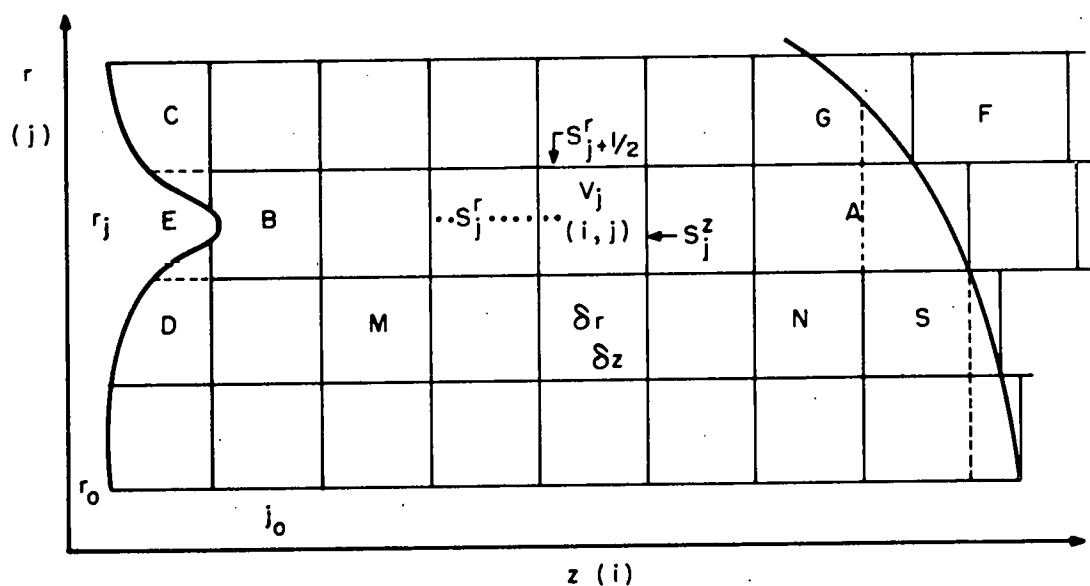


Fig. 2-1 The computing cells in $(z-r)$ coordinates

A regular cell (i, j) in plane Cartesian coordinates is a rectangle with dimension $\delta z, \delta r$. In cylindrical coordinates, it is a rectangular torus with inner and outer radii $(r_0 + (j-j_0 - 1) \delta r)$ and $(r_0 + (j-j_0) \delta r)$ respectively, and width

δz . The constants r_0 and j_0 are shown in Fig. 2-1. The regular cells have the properties shown in Table 2-1.

Table 2-1

Geometrical properties of a regular cell (i,j) in Fig. 2-1

Properties	Plane Coordinates	Cylindrical Coordinates
Volume (V_j)	$\delta r \delta z$	$2\pi (r_0 + (j-j_0 - 1/2)\delta r)\delta r\delta z$
Area (S_j^z)	δr	$2\pi (r_0 + (j-j_0 - 1/2)\delta r)\delta r$
Area (S_j^r)	δz	$2\pi (r_0 + (j-j_0 - 1/2)\delta r)\delta z$
Area ($S_{j+1/2}^r$)	δz	$2\pi (r_0 + (j-j_0)\delta r)\delta z$

Here S_j^z is the area on the right or left side of the cell (i,j) in z -direction, S_j^r is the area on the middle section of the cell (i,j) in r -direction, $S_{j+1/2}^r$ is the area on the top side of the cell (i,j) in r -direction, and V_j is the volume of the cell (i,j) .

The cells on the curved wall surface, although irregular, will not change their shapes because the wall boundary is fixed. On the other hand, the free boundary moves in space, and the cells on the free surface will change their sizes and configurations during the computational cycles. Since the motion of the free boundary must be included in the computation, these boundary cells require special treatment. The details of the resolutions on a free surface and on a curved rigid wall will be given in Chapter III.

There are a number of irregular small cells present on a

vacuum-plasma interface and on a curved wall boundary. These tiny cells will limit the time increment (δt) in the computation. To avoid this problem, we join them to their adjacent cells. The combination of a tiny cell and its adjacent cell makes the resulting cell bigger than a normal cell. The combined cell on a boundary may be in contact with two or three different neighbor cells on one of its four faces. These boundary cells are irregular in shape, they have different volumes and widths. They are still regarded as rectangular in shape, but their volumes are the geometric values of the actual configurations.

In Fig. 2-1, the left curve represents the vacuum-plasma interface, and the right curve represents the curved wall boundary. The cell (B) is in contact with three different cells (C), (D), and (E) on its left side. The cell (A) is in contact with two different cells (F) and (G) on its top side; two cells (S) and (N) on its bottom side. A cell on a boundary may have two or three neighbor cells on one of its four faces. In order to consider all geometric situations, a cell (i, j) is assumed to be in contact with eight neighbor cells $(i+1, j+1)$, $(i+1, j)$, $(i+1, j-1)$, $(i, j+1)$, $(i, j-1)$, $(i-1, j+1)$, $(i-1, j)$ and $(i-1, j-1)$ as shown in Fig. 2-2.

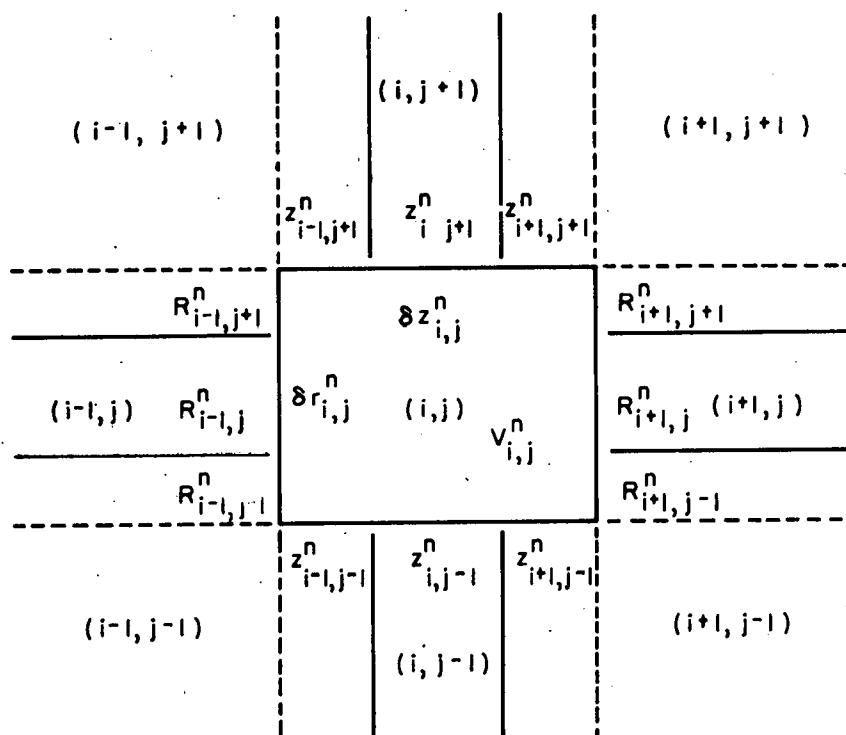


Fig. 2-2 The center cell (i, j) and its surrounding neighbors.

In Fig. 2-2, the center cell (i,j) has three neighbor cells on each of its four faces. The volumes and widths and configurations of all related cells are given on a wall boundary, or computed on a free boundary so that the geometric relations between them are known in advance. Therefore, the contacting surfaces of the surrounding neighbor cells on each face of the cell (i,j) are known. If a surrounding neighbor cell, say $(i+1, j+1)$, does not appear on a surface of the center cell, its contacting length, say $(R^n_{i+1, j+1})$, on that surface is set to zero. In case of a normal situation, the center cell (i,j) is surrounded by four regular neighbors; the following relations hold for the center cell and its surrounding neighbors:

$$V^n_{i,j} = V_j$$

$$\delta r^n_{i,j} = \delta r^n_{i+1,j} = R^n_{i+1,j} = \delta r,$$

$$= \delta r^n_{i-1,j} = R^n_{i-1,j} = \delta r,$$

$$\delta z^n_{i,j} = \delta z^n_{i,j+1} = Z^n_{i,j+1} = \delta z,$$

$$= \delta z^n_{i,j-1} = Z^n_{i,j-1} = \delta z,$$

$$Z^n_{i+s, j\pm 1} = R^n_{i\pm 1, j+t'} = 0, \quad \text{with } s = +1, -1; \\ \text{and } t' = +1, -1.$$

The lengths $R^n_{i+s, j+t'}$ and $Z^n_{i+s, i+t'}$ are shown in Fig. 2-2.

2.3 Summary of the Fluid-in-cell (FLIC) Method

A version of the FLIC method is summarized here which is originally designed for time-dependent, one-fluid problems in plane or axisymmetric compressible flows. Since the system of equations for transverse MHD flows is reducible to that for conventional compressible flows,²³ the FLIC method is modified to include a magnetic field perpendicular to the flow directions in magnetohydrodynamics.

The particle-in-cell (PIC) method¹³ is a combined Eulerian-Lagrangian scheme to solve multifluid problems in time-dependent, compressible flows. It utilizes Lagrangian particles in Eulerian cells to transport mass, momentum, and energy. The FLIC method¹⁵ is an Eulerian scheme, in which the difference equations are similar to those used in the PIC method. However, it employs a mass transport calculation which eliminates the use of particles.

We describe the axisymmetric case, the plane case being obvious and simpler. The basic flow variables to characterize the state of the plasma in transverse MHD flows are the density ρ , the z and r components of the fluid velocity, u and v , the magnetic field B , and the specific internal energy I . Other flow quantities are determined through their related equations from the given variables. The values of density, magnetic field, velocities, specific internal energy, and other flow quantities are given in each computing cell initially. Then, their values in each cell are advanced in time using the finite difference equations. The FLIC scheme consists of two steps for each cycle: First, intermediate values are calculated for the velocities

$(\tilde{u}, \tilde{v})$ and specific internal energy $(\tilde{T})$ due to the effects of fluid accelerations driven by the pressure gradients and the electromagnetic forces. Second, mass transport effects are calculated due to the movement of the fluid in each cell. The mass and the magnetic flux which flow from a cell (i,j) to a cell $(i+1,j)$ are directly proportional to the density and magnetic field of the donor cell (i,j) respectively.

The mass and magnetic flux flowing across the cell boundaries during the time increment are determined by using the intermediate values of the fluid velocities $(\tilde{u}$ and $\tilde{v})$. The mass which has crossed the cell boundaries then carries the momentum and energy corresponding to the intermediate values of the fluid velocities $(\tilde{u}$ and $\tilde{v})$ and the specific internal energy $(\tilde{T})$. The final values for ρ , B , u , v , and I are computed by redistributing the mass, magnetic flux, momentum, and energy in each cell. Throughout the steps, there is conservation of mass, momentum and energy.

2.4 Finite Difference Equations

Assume that at time $t = n \delta t$ the following values are specified for a cell (i,j) shown in Fig. 2-2.

density:	$\rho_{i,j}^n$
z-velocity component:	$u_{i,j}^n$
r-velocity component:	$v_{i,j}^n$
specific internal energy:	$l_{i,j}^n$
specific total energy:	$E_{i,j}^n = l_{i,j}^n + \frac{1}{2} (u^2 + v^2)_{i,j}^n$
magnetic field:	$B_{i,j}^n$
volume:	$V_{i,j}^n$
widths:	$\delta r_{i,j}^n$ and $\delta z_{i,j}^n$
fluid properties of its neighbors:	$\rho_{i+s,j+t'}^n$, etc. $s=1, 0, +1; t'=-1, 0, +1,$ (except $s=0$ and $t'=0$)
contacting lengths of its neighbors:	$R_{i\pm 1,j+t'}^n, \quad Z_{i+s,j\pm 1}^n$ $s=-1, 0, +1$ and $t'=-1, 0, +1$

The differencing procedures of the FLIC method for the cell (i,j) in Fig. 2-2 consist of two steps to calculate values for time $t = (n+1) \delta t$.

Step 1. The intermediate values of $\tilde{u}$, $\tilde{v}$, $\tilde{l}$ and $\tilde{E}$ due to the

effects of fluid accelerations are calculated, taking into account the pressure gradients and the magnetic forces.

$$\tilde{u}_{i,j}^n = u_{i,j}^n - \delta t A_{i,j}^z ((P^* + q)_{i+1/2,j}^n - (P^* + q)_{i-1/2,j}^n) / (\rho_{i,j}^n f_{i,j} \delta z),$$

$$\begin{aligned} \tilde{v}_{i,j}^n = & v_{i,j}^n - \delta t A_{i,j}^r ((S_{j+1/2}^r (P_T^{*n} - P_{i,j}^{*n}) - S_{j-1/2}^r (P_B^{*n} - P_{i,j}^{*n})) / (2 v_j) \\ & + (q_{i,j+1/2}^n - q_{i,j-1/2}^n) / \delta r) / (\rho_{i,j}^n f_{i,j}) - c_b (B^2)_{i,j}^n \delta t / (4 \pi r_j \rho_{i,j}^n). \end{aligned}$$

Here, $P^* = P + B^2/8\pi$,

$$A_{i,j}^z = \delta r_{i,j}^n / \delta r,$$

$$A_{i,j}^r = \delta z_{i,j}^n / \delta z,$$

$$f_{i,j} = v_{i,j}^n / v_j,$$

$$P_{i+1/2,j}^{*n} = (P_{i,j}^{*n} + P_{R,j}^{*n}) / 2,$$

$$P_{i-1/2,j}^{*n} = (P_{i,j}^{*n} + P_{L,j}^{*n}) / 2,$$

$$P_{L,j}^{*n} = \left(\sum_{s=-1}^{+1} P_{i\pm 1,j+s}^{*n} R_{i\pm 1,j+s}^n \right) / \delta r_{i,j}^n,$$

$$P_{B,j}^{*n} = \left(\sum_{s=-1}^{+1} P_{i+s,j\pm 1}^{*n} z_{i+s,j\pm 1}^n \right) / \delta z_{i,j}^n,$$

where P_{*Y}^n is the total pressure of the surrounding neighbors acting on the cell (i,j) on the Y-side, the dummy subscript $(Y = R, L, T, \text{ and } B)$ denotes the right (R), left (L), top (T), and bottom (B) side of the cell (i,j) respectively. The same notation applies to all subsequent formulae.

The time-step (δt) is chosen to meet the condition

$$| \vec{U} | \delta t / \text{Min}(\delta z_{i,j}^n, \delta r_{i,j}^n) \leq .4$$

It was introduced for accuracy in the original FLIC method. The fluid particles in a cell (i,j) will not be allowed to pass by its adjacent neighbors to farther cells during a computing cycle. It usually covers the Courant-Friedrichs-Lewy condition

$$\delta t (| \vec{U} | + C') / \text{Min}(\delta z_{i,j}^n, \delta r_{i,j}^n) \leq 1/\sqrt{2}$$

which is required for numerical stability. The notation C' represents the local magnetosonic speed.

Also,

$$\begin{aligned} \tilde{I}_{i,j}^n = & I_{i,j}^n - \delta t \{ A_{i,j}^r [P_{i,j}^n (S_{j+1/2}^r \bar{v}_{i,j+1/2}^n \\ & - S_{j-1/2}^r \bar{v}_{i,j-1/2}^n) + .5 q_{i,j+1/2}^n (S_{j+1}^r \bar{v}_{i,j+1}^n \\ & + S_j^r \bar{v}_{i,j}^n) - .5 q_{i,j-1/2}^n (S_{j-1}^r \bar{v}_{i,j-1}^n \\ & + S_j^r \bar{v}_{i,j}^n) - \bar{v}_{i,j}^n S_j^r (q_{i,j+1/2}^n - q_{i,j-1/2}^n)] \\ & + S_j^z A_{i,j}^z [\bar{u}_{i+1/2,j}^n (P_{i,j}^n + q_{i+1/2,j}^n) \\ & - \bar{u}_{i-1/2,j}^n (P_{i,j}^n + q_{i-1/2,j}^n) \end{aligned}$$

$$- \bar{u}_{i,j}^n (q_{i+1/2,j}^n - q_{i-1/2,j}^n)]] / (\rho_{i,j}^n f_{i,j} v_j) ,$$

$$\bar{E}_{i,j}^n = \bar{T}_{i,j}^n + (\bar{u}^2 + \bar{v}^2)_{i,j}^n / 2.$$

Here,

$$\bar{u}_{i,j}^n = (u_{i,j}^n + \tilde{u}_{i,j}^n) / 2,$$

$$\bar{v}_{i,j}^n = (v_{i,j}^n + \tilde{v}_{i,j}^n) / 2, \text{ etc.,}$$

$$\bar{u}_{i+1/2,j}^n = (u_{i,j}^n + \tilde{u}_{i,j}^n + u_R^n + \tilde{u}_R^n) / 4,$$

$$\bar{u}_{i-1/2,j}^n = (u_{i,j}^n + \tilde{u}_{i,j}^n + u_L^n + \tilde{u}_L^n) / 4,$$

$$\bar{v}_{i,j+1/2}^n = (v_{i,j}^n + \tilde{v}_{i,j}^n + v_T^n + \tilde{v}_T^n) / 4,$$

$$\bar{v}_{i,j-1/2}^n = (v_{i,j}^n + \tilde{v}_{i,j}^n + v_B^n + \tilde{v}_B^n) / 4,$$

The q terms that are artificial viscous pressures are of the forms.

$$q_{i+1/2,j}^n = b \ C_{i+1/2,j}^n \ \rho_{i+1/2,j}^n \ (u_{i,j}^n - u_R^n)$$

$$\text{if } K (u^2 + v^2)_{i+1/2,j}^n < (C^2)_{i+1/2,j}^n \text{ and if } u_{i,j}^n > u_R^n ,$$

$$= 0 \text{ otherwise,}$$

$$q_{i,j+1/2}^n = b \ C_{i,j+1/2}^n \ \rho_{i,j+1/2}^n \ (v_{i,j}^n - v_T^n)$$

$$\text{if } K (u^2 + v^2)_{i,j+1/2}^n < (C^2)_{i,j+1/2}^n \text{ and if } v_{i,j}^n > v_T^n ,$$

$$= 0 \text{ otherwise,}$$

$$q_{i-1/2,j}^n = b c_{i-1/2,j}^n \rho_{i-1/2,j}^n (u_L^n - u_{i,j}^n)$$

$$\text{if } K (u^2 + v^2)_{i-1/2,j}^n < (C^2)_{i-1/2,j}^n \text{ and if } u_L^n > u_{i,j}^n,$$

$$= 0 \text{ otherwise,}$$

$$q_{i,j-1/2}^n = b c_{i,j-1/2}^n \rho_{i,j-1/2}^n (v_B^n - v_{i,j}^n)$$

$$\text{if } K (u^2 + v^2)_{i,j-1/2}^n < (C^2)_{i,j-1/2}^n \text{ and if } v_B^n > v_{i,j}^n,$$

$$= 0 \text{ otherwise.}$$

The constants K and b are used to determine the magnitude of the artificial viscous pressure, e.g. $K=1$, $b \leq 0.5$. The constant (b) for cells on a wall boundary is set to be smaller than that for cells in an interior fluid region to avoid numerical viscous heating on a wall surface, say $b_{\text{wall}} \approx (1/3)b_{\text{fluid}}$. The quantity (C) is chosen to be the local sound speed in an axisymmetric case or the local magnetosonic speed in a plane case.

In a coaxial shock tube, the initial bias magnetic field is inversely proportional to the radius (r). Notice that the fluid in the tube will not flow radially under the influence of this field alone, because it is curl-free. The value (C) in the viscous pressure term must be the local sound speed in this case. Otherwise, if one chooses the value (C) to be the local magnetosonic speed in the shock tube, the viscous pressure terms will vary with respect to the radius (r). Thus, they will result in incorrect solutions to the problems. Meanwhile, the

discrete bias magnetic pressure terms ($\propto 1/r^2$) acting on the inner and outer surfaces of a toroidal cell can induce errors on the radial component of the fluid velocity. To minimize such errors, one can either use fine cells in the radial direction or set C_b to a value approximately equal to 1, say

$$[1 - (3/4)x(\delta r/r)^2 + (1/4)x(\delta r/r)^4] [1 - (\delta r/r)^2]^{-2},$$

which is obtained by setting the radial acceleration of the fluid to zero in the finite difference equation with a constant initial thermal pressure and a radially varying bias magnetic pressure ($\propto 1/r_j^2$) in the shock tube.

Let

$$X = (\rho, C, u, \tilde{u}, v, \tilde{v}),$$

then

$$X^n_{i\pm 1/2,j} = (X^n_{i,j} + X^n_{R_L}) / 2,$$

$$X^n_{i,j\pm 1/2} = (X^n_{i,j} + X^n_{T_B}) / 2,$$

$$X^n_{R_L} = \left(\sum_{s=-1}^{+1} X^n_{i\pm 1,j+s} R^n_{i\pm 1,j+s} \right) / \delta r^n_{i,j},$$

$$X^n_{T_B} = \left(\sum_{s=-1}^{+1} X^n_{i+s,j\pm 1} Z^n_{i+s,j\pm 1} \right) / \delta z^n_{i,j}.$$

Where X^n_Y is the averaged X-quantity over the surrounding neighbors on the Y-side of the cell (i,j) by the area-weighting method, the subscript Y denotes the direction of the cell surface.

Step 2. The final values due to mass transport effects are calculated.

Let $\Delta MR^n_{i\pm 1, j+s}$ be the mass flowing across the area $(R^n_{i\pm 1, j+s})$ between the cell (i, j) and the cell $(i\pm 1, j+s)$ during the time increment δt , and $\Delta MZ^n_{i+s, j\pm 1}$ be the mass flowing across the area $(Z^n_{i+s, j\pm 1})$ between the cell (i, j) and the cell $(i+s, j\pm 1)$ during the time increment δt , (with $s = -1, 0, +1$).

$$\begin{aligned} \Delta MR^n_{i\pm 1, j+s} &= R^n_{i\pm 1, j+s} S^z_j \rho^n_{i, j} (\tilde{u}^n_{i, j} + \tilde{u}^n_{i\pm 1, j+s}) \delta t / (2\delta r), \\ &\quad \text{if } (\tilde{u}^n_{i, j} + \tilde{u}^n_{i\pm 1, j+s}) \begin{matrix} > 0 \\ < 0 \end{matrix}, \\ &= R^n_{i\pm 1, j+s} S^z_{j+s} \rho^n_{i\pm 1, j+s} (\tilde{u}^n_{i, j} + \tilde{u}^n_{i\pm 1, j+s}) \delta t / (2\delta r), \\ &\quad \text{if } (\tilde{u}^n_{i, j} + \tilde{u}^n_{i\pm 1, j+s}) \begin{matrix} < 0 \\ > 0 \end{matrix}, \end{aligned}$$

$$s = -1, 0, +1$$

$$\begin{aligned} \Delta MZ^n_{i+s, j\pm 1} &= Z^n_{i+s, j\pm 1} S^r_{j\pm 1/2} \rho^n_{i, j} (\tilde{v}^n_{i, j} + \tilde{v}^n_{i+s, j\pm 1}) \delta t / (2\delta z), \\ &\quad \text{if } (\tilde{v}^n_{i, j} + \tilde{v}^n_{i+s, j\pm 1}) \begin{matrix} > 0 \\ < 0 \end{matrix}, \\ &= Z^n_{i+s, j\pm 1} S^r_{j\pm 1/2} \rho^n_{i+s, j\pm 1} (\tilde{v}^n_{i, j} + \tilde{v}^n_{i+s, j\pm 1}) \delta t / (2\delta z), \\ &\quad \text{if } (\tilde{v}^n_{i, j} + \tilde{v}^n_{i+s, j\pm 1}) \begin{matrix} < 0 \\ > 0 \end{matrix}, \end{aligned}$$

$$s = -1, 0, +1$$

$$\rho_{i,j}^{n+1} = \rho_{i,j}^n + \left(\sum_{s=-1}^{+1} \Delta MR_{i-1,j+s}^n - \sum_{s=-1}^{+1} \Delta MR_{i+1,j+s}^n + \sum_{s=-1}^{+1} \Delta MZ_{i+s,j-1}^n - \sum_{s=-1}^{+1} \Delta MZ_{i+s,j+1}^n \right) / (v_j f_{i,j}) .$$

Let $H = B/r_j^m$, the differencing equations for H are the same as those for the density ρ .

Final values for the velocity components (u,v) and for the total energy per unit mass (E) are calculated by conservation of total momentum and total energy in each cell.

$$\begin{aligned} F_{i,j}^{n+1} = & \left(\sum_{s=-1}^{+1} TR_{i-1,j+s} \tilde{F}_{i-1,j+s}^n \Delta MR_{i-1,j+s}^n \right. \\ & + \sum_{s=-1}^{+1} TZ_{i+s,j-1} \tilde{F}_{i+s,j-1}^n \Delta MZ_{i+s,j-1}^n \\ & - \sum_{s=-1}^{+1} TR_{i+1,j+s} \tilde{F}_{i+1,j+s}^n \Delta MR_{i+1,j+s}^n \\ & - \sum_{s=-1}^{+1} TZ_{i+s,j+1} \tilde{F}_{i+s,j+1}^n \Delta MZ_{i+s,j+1}^n \\ & + \tilde{F}_{i,j}^n (\rho_{i,j}^n f_{i,j} v_j \\ & + \sum_{s=-1}^{+1} (1 - TR_{i-1,j+s}) \Delta MR_{i-1,j+s}^n \\ & + \sum_{s=-1}^{+1} (1 - TZ_{i+s,j-1}) \Delta MZ_{i+s,j-1}^n \\ & - \sum_{s=-1}^{+1} (1 - TR_{i+1,j+s}) \Delta MR_{i+1,j+s}^n \end{aligned}$$

$$- \sum_{s=-1}^{+1} \left(1 - TZ_{i+s,j+1} \right) \Delta MZ_{i+s,j+1}^n \Big) /$$

$$\left(\rho_{i,j}^{n+1} f_{i,j} v_j \right).$$

Here,

$$F = (u, v, E),$$

$$TR_{i\pm 1, j+s} = 1 \text{ if fluid flows into a cell } (i, j) \text{ from}$$

$$\text{a cell } (i\pm 1, j+s), \text{ with } s=-1, 0, +1;$$

$$= 0 \text{ otherwise.}$$

$$TZ_{i+s, j\pm 1} = 1 \text{ if fluid flows into a cell } (i, j) \text{ from}$$

$$\text{a cell } (i+s, j\pm 1), \text{ with } s=-1, 0, +1;$$

$$= 0 \text{ otherwise.}$$

The specific internal energy I , gas pressure P , sound speed C , magnetic field B , and total pressure P^* are obtained by the following relations:

$$I_{i,j}^{n+1} = E_{i,j}^{n+1} - (u^2 + v^2)_{i,j}^{n+1} / 2 ,$$

$$P_{i,j}^{n+1} = (\gamma - 1) \rho_{i,j}^{n+1} I_{i,j}^{n+1} ,$$

$$C_{i,j}^{n+1} = (\gamma R I_{i,j}^{n+1} / C_v)^{\frac{1}{2}} ,$$

$$B_{i,j}^{n+1} = H_{i,j}^{n+1} (r_j^m) ,$$

$$P_{i,j}^{*n+1} = P_{i,j}^{n+1} + (B^2)_{i,j}^{n+1} / 8\pi ,$$

where the constants R and C_v are the gas constant and specific heat at constant volume respectively.

The most important feature of the present code is that it allows cells on boundaries to be different shapes and to be in contact with an arbitrary number of neighbor cells. The variable cells in the present modification can also be applied to other methods, say, the PIC method. It is also possible to include physical transport coefficients in the method by adding them directly to the finite difference equations.

III. BOUNDARY CONDITIONS

3.1 Introduction

The finite difference equations developed in Chapter II can be applied to cells on a boundary where the boundary values have been specified. There are four types of cell boundaries: inflow, outflow, free boundary, and solid wall. In all cases, boundary conditions are specified by giving the values of the flow variables in fictitious cells outside the boundaries.

On an inflow boundary, these values are those of the fluid which flows into the boundary. On an outflow boundary, these values are equal to those of the adjacent inner cells. On a free surface, the applied pressure is specified in a vacuum. On a solid wall, the normal component of the fluid velocity is equal to that of the wall surface. There are no flows of mass or energy across a wall boundary or across a non-leaky free surface. Mathematically, the number of boundary conditions is equal to the number of the characteristic lines or cones on a boundary pointing into the fluid region. Numerically, the number of boundary conditions required is equal to the number of the difference equations, and is often more than the correct number of mathematical conditions. Care must be used so that these additional conditions have as little influence as possible. In this chapter, we describe the boundary conditions for the finite difference equations when they are applied to cells on a free boundary

or on a curved wall.

In Fig. 2-1, the left curve represents a free boundary which separates the vacuum and fluid regions; the right curve represents a wall boundary which is fixed in space. The cells (A) and (C) are the surface cells on the wall and on the free boundary respectively. The cell (E) is empty in the vacuum, and the cell (F) is fictitious outside the wall boundary. The cells (M) and (N) are regular cells in the interior region. When the finite difference equations are applied to the surface cells on the boundaries, special treatment is required, which we detail below.

3.2 On a Vacuum-plasma Interface

As the free surface moves into a normal cell, that cell becomes a surface cell with a changing volume. Its cell edge in contact with the vacuum will coincide with the free surface. The free surface moves with the same velocity as the mass in the surface cell, and is a Lagrangian contour moving in the Eulerian mesh. This Lagrangian contour gives a good resolution of the movable vacuum-plasma interface if the contour does not become seriously distorted. Apparently, the surface cells on the free boundary will change their shapes, volumes, and widths when the Lagrangian contour moves into or out of them.

The finite difference equations for the center cell (i,j) in Fig. 2-2 need the following modifications when the cell (i,j) is a surface cell on the free boundary. The total pressure P^*

of an empty neighbor cell of the cell (i,j) is the applied magnetic pressure in the vacuum. Other flow variables in empty neighbor cells are set equal to those of the adjacent surface cell (i,j) . We give a simple example as follows:

- i) One of the surrounding neighbors of the cell (i,j) is empty, say the cell $(i-1,j)$ in Fig. 2-2, i.e. $\rho_{i-1,j}^n = 0$.

Then,

$$P_L^{*n} = (\underline{P_v^{*n}}_{i-1,j} R_{i-1,j}^n + P_{i-1,j+1}^{*n} R_{i-1,j+1}^n + P_{i-1,j-1}^{*n} R_{i-1,j-1}^n) / \delta r_{i,j}^n ,$$

$$X_L^n = (\underline{X}_{i,j}^n R_{i-1,j}^n + X_{i-1,j+1}^n R_{i-1,j+1}^n + X_{i-1,j-1}^n R_{i-1,j-1}^n) / \delta r_{i,j}^n ,$$

with $X = (\rho , l , u , \tilde{u} , v , \tilde{v}) ,$

$$\Delta MR_{i-1,j}^n = 0 , \text{ etc. ,}$$

where $P_v^{*n}_{i-1,j}$ is the applied magnetic pressure in the empty cell $(i-1,j)$ at the time $(n\delta t)$, and the underlined terms indicate the modification made in the usual scheme.

The values of the flow variables assigned to the empty cells do not correspond to real physical properties. It means that the flow variables on the free boundary are those of the local surface cells. For a non-leaky free surface, there are

surface cell (i,j) on the free surface are reevaluated after the movements of the cell boundaries on the Lagrangian contour. The procedures are as follows:

- 1) Calculate the new width $(\delta z_{i,j}^{n+1})$ of a surface cell (i,j) due to the movements of the vertical segments along the axial direction.

- i) Set $\delta z_{i,j}^{n+1} = \delta z_{i,j}^n$ unconditionally;

- ii) set $\delta z_{i,j}^{n+1} = \delta z_{i,j}^{n+1} \mp u_{i,j}^{n+1} \delta t$,

if the left (-sign) or right (+sign) boundary of the surface cell (i,j) is entirely in contact with the vacuum, or if the left or right boundary is in contact only partially with the vacuum, but the width $(\delta z_{i,j}^{n+1})$ has been changed after the movements of the vertical segments.

- 2) Calculate the new volume $(V_{i,j}^{n+1})$ of the surface cell (i,j) due to the axial movements of the vertical segments.

- i) Set $V_{i,j}^{n+1} = V_{i,j}^n$ unconditionally;

- ii) set $V_{i,j}^{n+1} = V_{i,j}^{n+1} - \frac{\Delta r_L}{\Delta r_R} \} u_{i,j}^{n+1} (2\pi r_j) \delta t$,

where Δr_L and Δr_R are the total lengths of the vertical segments on the left (- sign) and right (+ sign) boundaries of the surface cell (i,j) in contact with the vacuum respectively. They are computed in advance from the previous geometrical configuration of the free boundary.

- 3) Calculate the final volume ($V_{i,j}^{n+1}$) of the surface cell (i,j) due to the radial movements of the horizontal segments.

$$i) \text{ Set } V_{i,j}^{n+1} = V_{i,j}^{n+1} - \Delta z_B \} V_{i,j}^{n+1} (2\pi r_j) \delta t ,$$

where Δz_B and Δz_T are the total lengths of the horizontal segments on the bottom (- sign) and top (+ sign) boundaries of the cell (i,j) in contact with the vacuum respectively. They are calculated from the geometrical configuration of the free boundary after the movements of the vertical segments.

- 4) Calculate the new width ($\delta r_{i,j}^{n+1}$) of the surface cell (i,j) due to the movements of the horizontal segments along the radial direction.

$$i) \text{ Set } \delta r_{i,j}^{n+1} = \delta r_{i,j}^n \text{ unconditionally;}$$

$$ii) \text{ set } \delta r_{i,j}^{n+1} = \delta r_{i,j}^{n+1} - v_{i,j}^{n+1} \delta t ,$$

if the bottom (- sign) or top (+ sign) boundary of the surface cell (i,j) is entirely in contact with the vacuum after the movements of the vertical segments in all surface cells, or if the bottom or top boundary is in contact only partially with the vacuum after the movements of the vertical segments in all surface cells, but the width ($\delta r_{i,j}^{n+1}$) has been changed after the movements of the horizontal segments.

When the new value of the volume of the surface cell (i,j)

is computed, the density ($\rho_{i,j}^{n+1}$) and magnetic field ($B_{i,j}^{n+1}$) in that cell have to be reevaluated according to the continuity equations in Lagrangian form.

$$\rho_{i,j}^{n+1} = \rho_{i,j}^{n+1} V_{i,j}^n / V_{i,j}^{n+1} ,$$

$$B_{i,j}^{n+1} = B_{i,j}^{n+1} V_{i,j}^n / V_{i,j}^{n+1} .$$

When the free surface moves into a surface cell, the volume of that cell is decreasing. To get rid of a tiny surface cell (i,j) on the Lagrangian contour, we combine it with an adjacent inner cell when its volume is smaller than a half of the regular volume, i.e. $V_{i,j}^{n+1} < V_j/2$. When the free surface moves out of a surface cell, the volume of that cell is increasing. In order to keep the accuracy, the big surface cell will be divided into two cells when its volume ($V_{i,j}^{n+1}$) is larger than $1.4 V_j$.

A combination of two cells and a division of the large cell will be called "cell-mixing". The combination of two cells, and the division of a cell into two cells follow the conservation laws of total mass, momentum, and energy. The cell-mixing process is also an effective way to produce an artificial viscosity on the free surface, which helps numerical stability; it naturally will smooth the fluid properties in the mixed region. The internal energies of the mixed cells increase at the expense of the kinetic energies. On the vacuum-plasma interface, there are no viscous pressure terms between a surface cell and an empty

cell because the differences of the velocity components between them are zero. The cell-mixing process is taken to produce the same effect of the artificial viscous pressure. Since it may produce extra viscous heating, it is used on the free surface to stabilize the flow only. One can also control the use of the "cell-mixing" in the computations so that there will be no extreme viscous heating on the free boundary.

The procedures of the cell-mixing are described as follows:

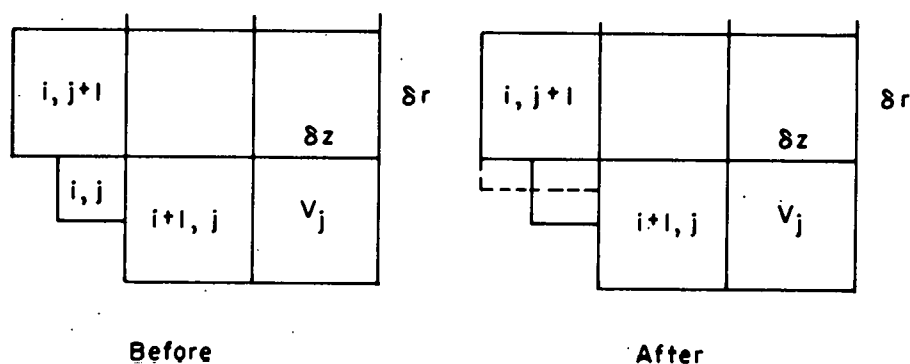


Fig. 3-2. The sample cells for cell-mixing

- 1) Combine the cell (i, j) with the cell $(i, j+1)$ if $V_{i,j}^{n+1} < V_j / 2$ and $\delta r_{i,j}^{n+1} < .54 \delta r$:

$$\underline{M}_{i,j+1}^{n+1} = \rho_{i,j+1}^{n+1} V_{i,j+1}^{n+1} + \rho_{i,j}^{n+1} V_{i,j}^{n+1}, \quad (\text{new mass of the combined cell located at } (i, j+1))$$

Similarly,

$$\underline{M}_{i,j+1}^{n+1} \underline{u}_{i,j+1}^{n+1} = \rho_{i,j+1}^{n+1} V_{i,j+1}^{n+1} u_{i,j+1}^{n+1} + \rho_{i,j}^{n+1} V_{i,j}^{n+1} u_{i,j}^{n+1},$$

$$\underline{M}_{i,j+1}^{n+1} \underline{v}_{i,j+1}^{n+1} = \rho_{i,j+1}^{n+1} V_{i,j+1}^{n+1} v_{i,j+1}^{n+1} + \rho_{i,j}^{n+1} V_{i,j}^{n+1} v_{i,j}^{n+1},$$

$$\underline{M}_{i,j+1}^{n+1} \underline{E}_{i,j+1}^{n+1} = \rho_{i,j+1}^{n+1} v_{i,j+1}^{n+1} E_{i,j+1}^{n+1} + \rho_{i,j}^{n+1} v_{i,j}^{n+1} E_{i,j}^{n+1},$$

$$\underline{H}_{i,j+1}^{n+1} = (H_{i,j+1}^{n+1} v_{i,j+1}^{n+1} + H_{i,j}^{n+1} v_{i,j}^{n+1}) / (v_{i,j+1}^{n+1} + v_{i,j}^{n+1}),$$

$$\underline{\rho}_{i,j+1}^{n+1} = \underline{M}_{i,j+1}^{n+1} / (v_{i,j+1}^{n+1} + v_{i,j}^{n+1}),$$

$$\underline{v}_{i,j+1}^{n+1} = v_{i,j+1}^{n+1} + v_{i,j}^{n+1},$$

$$\underline{\delta z}_{i,j+1}^{n+1} = \delta z_{i,j+1}^{n+1},$$

$$\underline{\delta r}_{i,j+1}^{n+1} = v_{i,j+1}^{n+1} / (2\pi r_{j+1} \underline{\delta z}_{i,j+1}^{n+1}),$$

where the underline denotes a quantity in the combined cell (i,j+1).

- 2) Divide the combined cell (i,j+1) into two cells if $\underline{v}_{i,j+1}^{n+1} > 1.4v_{j+1}$ and $\underline{\delta r}_{i,j+1}^{n+1} > 1.5 \delta r$.

Let $X = (\rho, u, v, E, H)$, then

$$x_{i,j}^{n+1} = \underline{x}_{i,j+1}^{n+1},$$

$$x_{i,j+1}^{n+1} = \underline{x}_{i,j+1}^{n+1},$$

$$v_{i,j}^{n+1} = \underline{v}_{i,j+1}^{n+1} - 2\pi r_{j+1} \underline{\delta z}_{i,j+1}^{n+1} \delta r,$$

$$\delta r_{i,j}^{n+1} = \delta r_{i,j}^{n+1},$$

$$\delta z_{i,j}^{n+1} = v_{i,j}^{n+1} / (2\pi r_j \delta r_{i,j}^{n+1}),$$

$$v_{i,j+1}^{n+1} = 2\pi r_{j+1} \frac{\delta z_{i,j+1}^{n+1}}{\delta r},$$

$$\delta z_{i,j+1}^{n+1} = \frac{\delta z_{i,j+1}^{n+1}}{\delta r},$$

$$\delta r_{i,j+1}^{n+1} = \delta r.$$

The cell (i,j) will be combined with the cell $(i+1,j)$ if $v_{i,j}^{n+1} < v_j / 2$ and $\delta z_{i,j}^{n+1} < .54 \delta z$, and the combined cell $(i+1,j)$ will be divided into two cells if $v_{i+1,j}^{n+1} > 1.4 v_j$ and $\delta z_{i+1,j}^{n+1} > 1.5 \delta z$. The formulae used in this case are similar to those in the previous case.

An additional exact condition is still needed to control the cell-mixing process so as not to produce numerical overheating on the free surface, but to insure the stability of the flow. Because the moving boundary changes its configuration during the computational cycles, the computer program for considering all the geometrical situations becomes very complicated. It simplifies the computer programming if the cells are made rectangular after cell-mixing. Any serious distortion of the free boundary will yield inaccurate solutions on the Lagrangian contour. Further improvements are still required for more accurate solutions on the moving vacuum-plasma interface.

3.3 On a Curved Rigid Wall

When the geometry of the wall shape is given, the wall boundary can be located in the computational zone. The cells on the wall surface have different shapes with the boundaries of the surface cells following the wall curve. In Fig. 3-3, the cell (S) is a surface cell on the wall boundary, which may include a regular cell and a smaller irregular cell. The cell (F) is a fictitious boundary cell. The procedures for treating the surface cells and creating the fictitious cells on a curved wall boundary are given below:

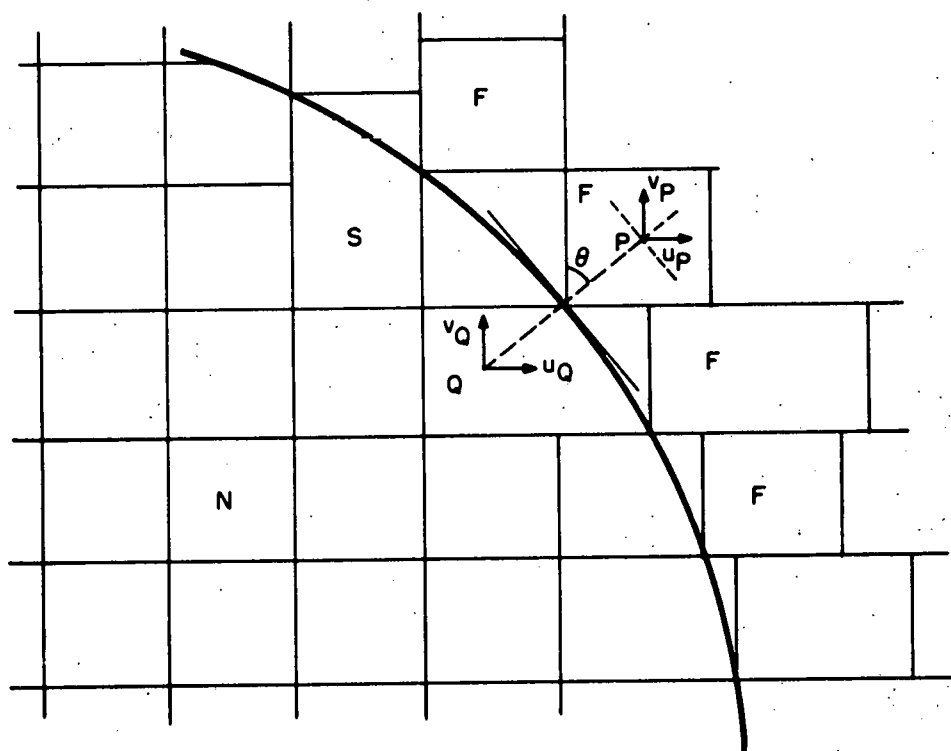


Fig. 3-3 Surface cells and fictitious cells on a curved wall.

- 1) A surface cell smaller than one half of the regular cell at the same radius is adjoined to an adjacent inner cell, making the resulting cell to be a surface cell.
- 2) Calculate its volume and widths with the wall boundary to be a part of its boundaries. But the shape of that surface cell is still regarded as a rectangle with a part outside the boundary as shown in Fig. 3-3.
- 3) Locate a fictitious cell adjacent to that surface cell, and make its shape equal to that of the adjacent surface cell in a rectangle.

The values of the flow variables for the fictitious cell are determined as follows:

- 1) Find the position of the center (P) of the fictitious cell as shown in Fig. 3-3.
- 2) Determine the position of its mirror image (Q) by reflection with respect to the tangent of the wall curve.
- 3) The values of the flow variables at the image (Q) are obtained by the area-weighting method from its surrounding cells.
- 4) The values of the flow variables at the center (P) of the fictitious cell are obtained by reflection from its image values at the point (Q).
- 5) The magnetic field at the point (P) is obtained by a special treatment, i.e., $B_p = B_Q r_Q / r_p$ in an axisymmetric case; $B_p = B_Q$ in a plane case.
- 6) There is no mass transport across the wall boundary. The

mass and magnetic flux flowing across a cell boundary between a surface cell and a fictitious cell must be set to zero.

The reflection procedure for the flow variables of the fictitious cell is as follows:

$$\rho_P = \rho_Q ,$$

$$I_P = I_Q ,$$

$$u_P = u_Q \cos(2\theta) - v_Q \sin(2\theta) ,$$

$$v_P = -u_Q \sin(2\theta) - v_Q \cos(2\theta) ,$$

$$\tilde{u}_P = \tilde{u}_Q \cos(2\theta) - \tilde{v}_Q \sin(2\theta) ,$$

$$\tilde{v}_P = -\tilde{u}_Q \sin(2\theta) - \tilde{v}_Q \cos(2\theta) ,$$

where the angle θ is shown in Fig. 3-3.

In this way, the flow variables (ρ , u , v , I) in the fictitious cell are obtained by reflection with respect to the tangent of the wall profile. The fluid velocity on the wall will be parallel to the wall surface and extended without change into the fictitious cell. The special treatment on the magnetic field is designed so that the total pressure (P^*) of the fluid near the wall will be balanced by that of the fictitious cell. In addition, we insure that there will be no mass transport between a surface cell and a fictitious cell.

The above procedures can be applied to both axisymmetric

flows and plane flows. Since the cells on the wall surface are irregular in shape, we call the method an "irregular shape" approximation. The reflective boundary conditions suggested in this section make the code powerful enough to solve problems with arbitrary smooth wall boundaries. The accuracy of the code remains high on a curved wall boundary.

IV. INCIDENT SHOCK WAVES

4.1 General Description

In the past, shock waves in coaxial shock tubes have been simplified as a one-dimensional propagation^{5,6}. It is obvious that the plasma flow in the coaxial shock tube is two-dimensional because the drive field is a function of the radial position. In this section, we apply the FLIC MHD code to study the formation and shape of the shock wave in the two-dimensional (r - z) plane. In addition, this incident shock wave and the knowledge of its properties are needed before we study shock focusing and reflection from a curved end wall.

The schematic view of the coaxial electromagnetic shock tube is shown in Fig. 4-1. First of all, a bias current is built up in the central rod of the inner conductor to set up an azimuthal bias field. When the capacitor banks are discharged, a current sheet will form in the space between the concentric electrodes to ionize the working gas in the tube. This drive current sheet and magnetic field interact and produce electromagnetic forces acting on the plasma. The magnetic pressure serves as a magnetic piston to push the plasma down the tube. The positive pressure disturbance in the plasma will eventually steepen into a shock front. Since the bias field is azimuthal, a transverse MHD shock wave will be formed when the plasma is compressed by a sufficiently strong applied pressure.

The numerical code with the Lagrangian free boundary is now applied to this problem. When the drive current sheet is formed in the plasma, the drive magnetic field is built up in the vacuum. The corresponding magnetic pressure in the vacuum is then used in the finite difference equations as the boundary value on the free surface. We prescribe the drive current as an arbitrary function of time. We first show how shock waves will be formed by constant and time increasing sinusoidal drive currents. Since there is no physical dissipation in the basic equations, numerical viscous pressure terms are included in the code to permit the computation of shock transition.

For each case of shock formation or reflection studied, the computer results will be shown on SC4060 plottings, in which there are three planes on each page (see Fig. 4-2). In each plane, the upper and lower grid lines are the outer wall and inner conductor of the shock tube respectively. The tick marks on the upper grid line denote the sizes and locations of regular computing cells. The left blank region is the vacuum; the shadow region contains the plasma. The curves in the first plane are the constant temperature contours, their values are indicated by integers listed above the plane. The density is represented by the shading points in the second plane. The darker the region is, the higher the density. The velocity and its direction in each cell is represented by a segment vector in the third plane. The details of the flow and the shock wave in the tube can be clearly seen from these planes on the

plottings at different times. In the shock transition, a number of temperature contours are present because there is a large gradient of the temperature.

In addition, we also show the profiles of magnetic field $B(z)$, temperature $T(z)$ and density $\rho(z)$ at several different radii. This is to facilitate comparison with experiments.

4.2 Constant Drive Current

In this section, the drive current in the shock tube is assumed to be a step function. The initial conditions and other required information of the tube are given in Table 4-1. The results on shock formation by a constant driver are presented in Fig. 4-2. Because of the radial effect of the drive magnetic field, the applied magnetic pressure in the tube is inversely proportional to the square of the radius (r). The plasma near the inner conductor will be driven by a larger current to move faster than that near the outer wall. Therefore, the shock wave has a slope (in this case, about 120°) with respect to the axial axis during the early stages of formation. Also, the drive current sheet on the vacuum boundary becomes inclined automatically, a radially outward force thus acts on the plasma to move it toward the outer wall. As long as the plasma behind the shock wave moves to the outer wall region continuously, compression heating occurs in the plasma in that area.

The total pressure ($P^* = P + B^2 / 8\pi$) of the plasma there becomes the highest. The increase of magnetic pressure is due to the frozen-in field moving together with the plasma; the increase of thermal pressure is due to compression heating on the outer wall. This high total pressure has two essential effects on the plasma motion and on the shock wave. First, it strengthens the shock wave near the outer wall so that the shock slope diminishes. Thus, a nearly planar shock wave will always be formed by the radially varying drive magnetic field,

provided the tube is long enough. Secondly, it slows down the axial motion of the plasma in the outer region behind the shock wave. The tail part of the plasma near the outer wall increases in length as the motion goes on. When the shock wave approaches a quasi-steady state, it makes an angle of about 110° with the axial axis, but the vacuum-plasma interface continuously increases its slope with respect to the axial axis. The inclined shock wave on the outer wall will thus be reflected.

The sequence of the shock motion, flow direction of the plasma, temperature contours, and the density distributions in the shock tube are shown in Fig. 4-2. The shock speed is about 125 cm/ μ sec, the Mach number is 867, and the corresponding Alfven Mach number at the mean radius is 4.63. At the mean radius, just behind the shock wave, the density of the plasma is about 1.8×10^{-8} gram/cm³, the magnetic field is about 28 kilo-Gauss (KG), the temperature is about 3.8 Kev, and β is 0.9.

The magnetic field, density, and temperature are also plotted with respect to axial positions at several different radii in Fig. 4-3. The humps on the curves (d) again show the effect of the compression on the outer wall. The flow conditions behind the shock wave are not uniform radially or axially.

Fig. 4-4 shows the density distribution in the coaxial shock tube which was obtained by B. Marder²⁸ using the GAP code. There is good qualitative agreement with our results.

Table 4-1

Initial conditions and other required information of the shock tube

Working gas	hydrogen
Density	$.5344 \times 10^{-8}$ gram/cm ³
Pressure	66.65 dynes/cm ²
Temperature	300 °K
Axial component of velocity	0 cm/sec
Radial component of velocity	0 cm/sec
Specific internal energy	1.87×10^{10} ergs/gram
Specific heat ratio	5/3
Gas constant	4.157×10^7 ergs/°K
Inner radius (r_i)	6.25 cm
Outer radius (r_o)	11.25 cm
Mean radius (r_m)	$(r_i r_o)^{\frac{1}{2}}$
Cell width ($\delta r = \delta z$)	0.25 cm
Tube length	32 cm
Bias magnetic field	$7000 r_m / r$ Gauss
Drive magnetic field	$4.8 \times 7000 r_m / r$ Gauss

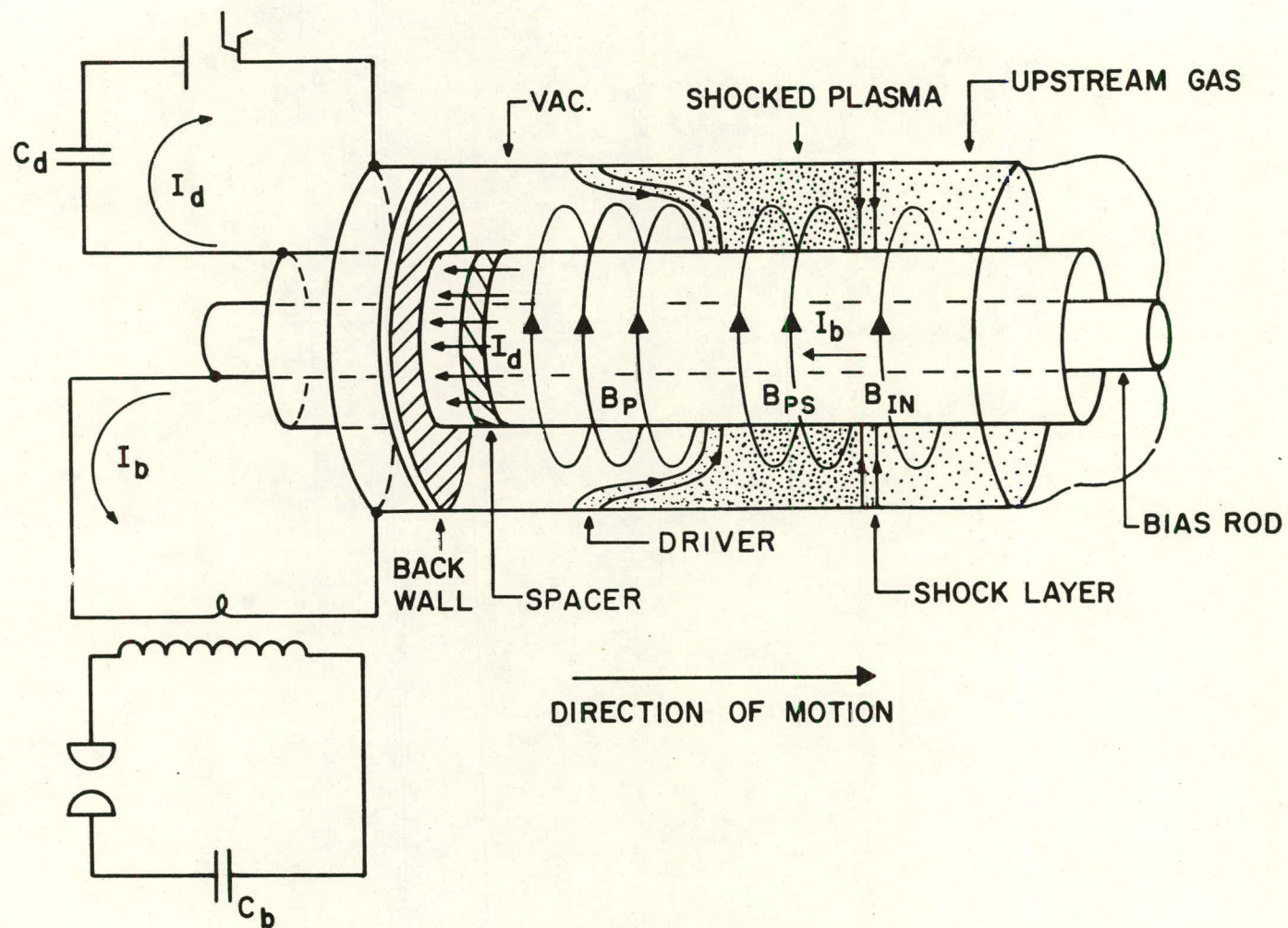


Fig. 4-1 Schematic of Shock Tube System

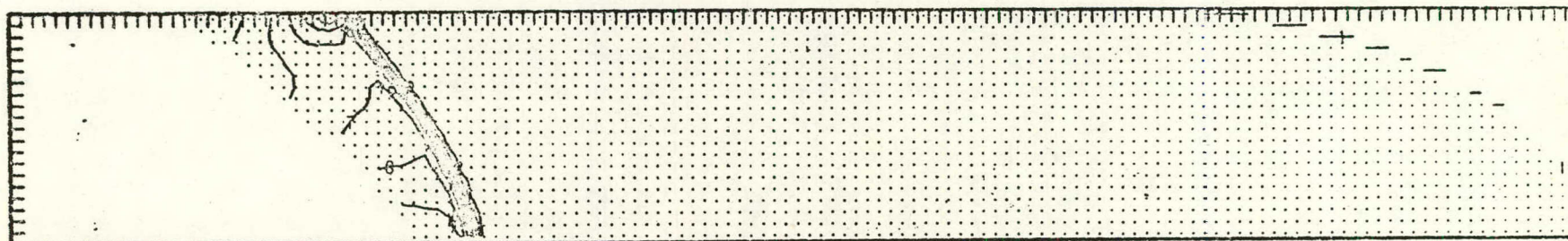
Fig. 4-2

CYCLE= 201 TIME= 0.7999691 07

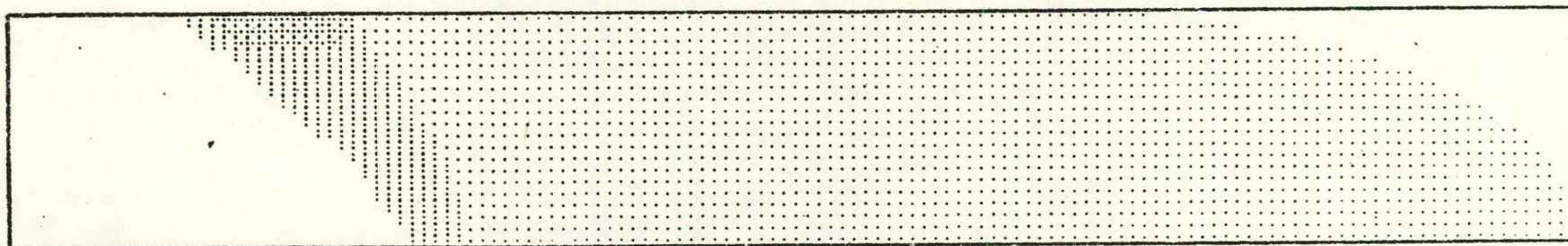
CONTOURS

1 =	0.000	2 =	0.498E+07	3 =	0.996E+07	4 =	0.149E+08	5 =	0.199E+08
6 =	0.249E+08	7 =	0.299E+08	8 =	0.349E+08	9 =	0.398E+08	10 =	0.448E+08

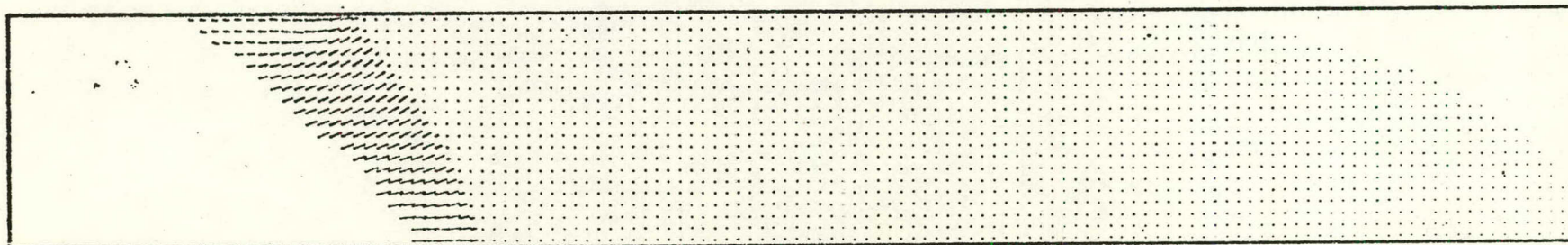
CYCLE= , TIME= DENSITY SHADING VELO



Z
TEMPERATURE CONTOURS



DENSITY SHADING



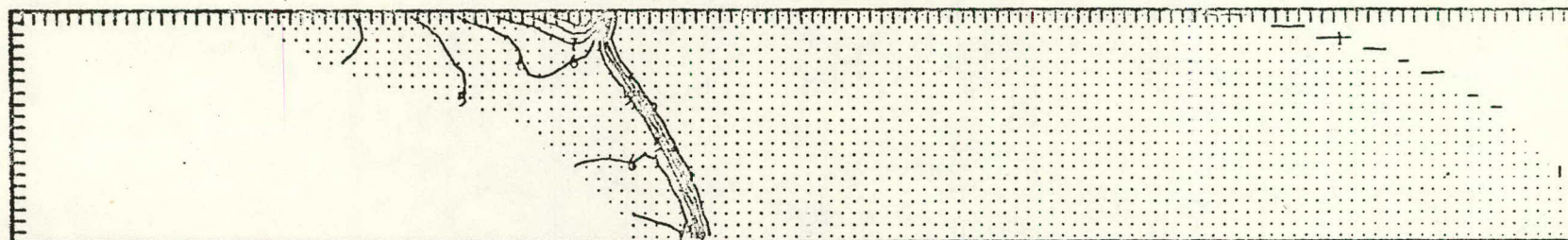
VELOCITY
NUMERICAL SHOCK WAVES SIMULATIONS AT COLUMBIA

Fig. 4-2

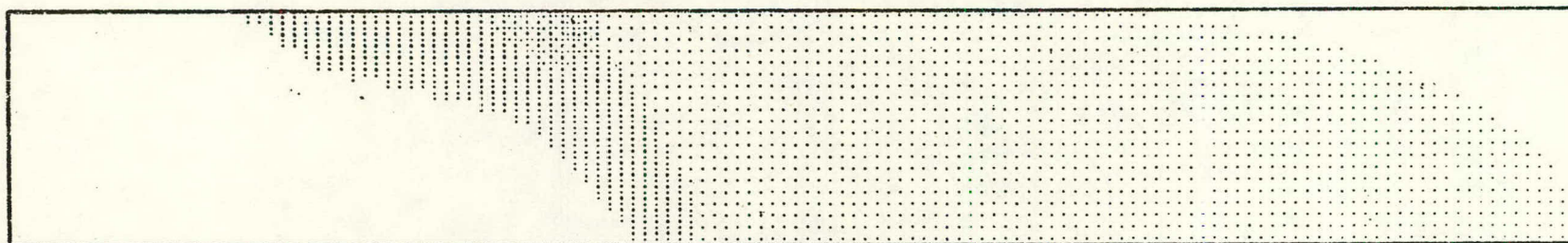
CYCLE= 301 , TIME= 0.119992E+06

CONTOURS							
1 =	0.000	2 =	0.715E+07	3 =	0.143E+08	4 =	0.214E+08
5 =	0.286E+08	6 =	0.357E+08	7 =	0.429E+08	8 =	0.500E+08
9 =	0.572E+08	10 =	0.643E+08				

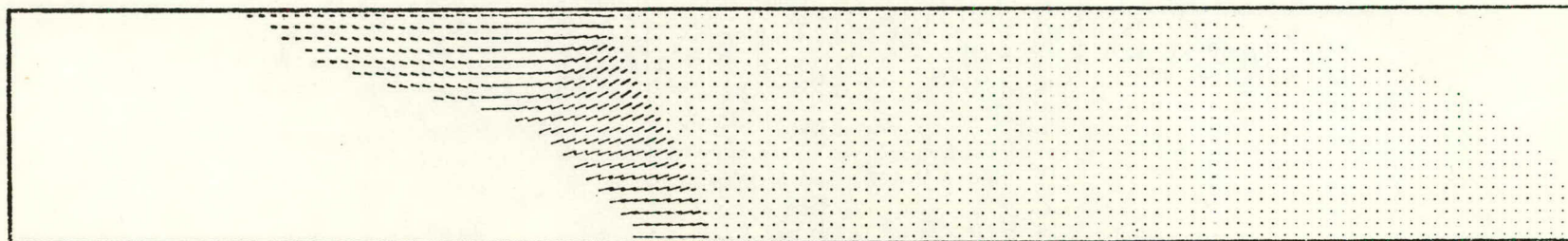
CYCLE= , TIME= DENSITY SHADING VELO



Z
TEMPERATURE CONTOURS



DENSITY SHADING



VELOCITY

NUMERICAL SHOCK WAVES SIMULATIONS AT COLUMBIA

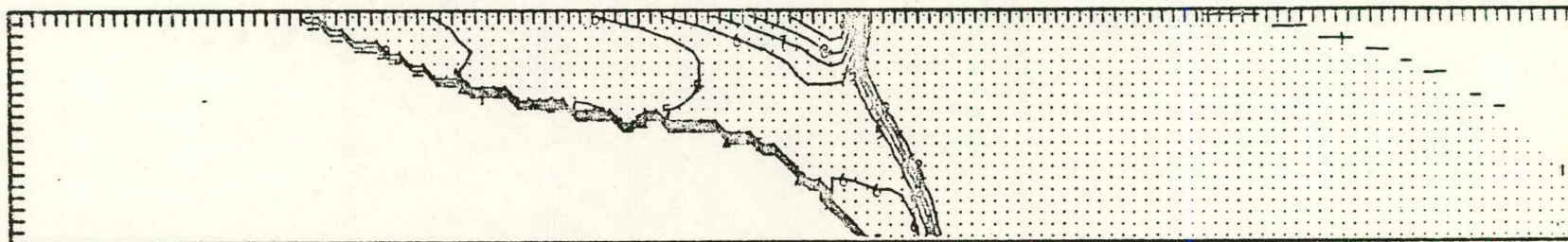
Fig. 4-2

CYCLE= 401 TIME= 0.159987E-06

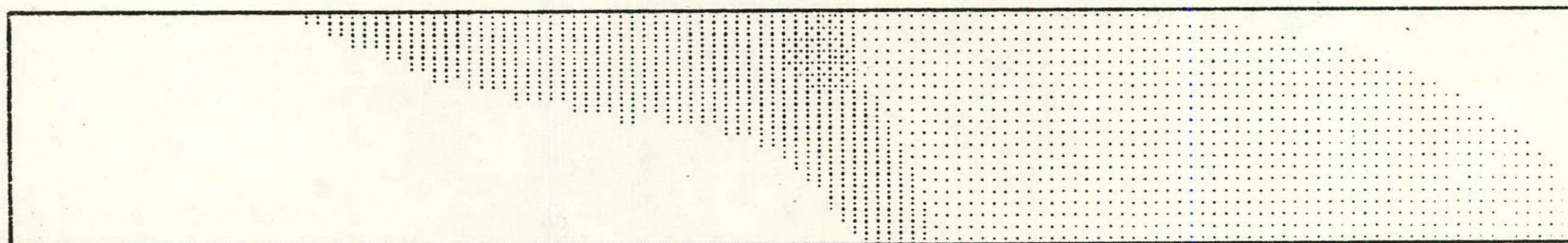
CONTOURS

1 =	0.000	2 =	0.781E+07	3 =	0.156E+08	4 =	0.234E+08	5 =	0.312E+08
6 =	0.390E+08	7 =	0.468E+08	8 =	0.546E+08	9 =	0.625E+08	10 =	0.703E+08

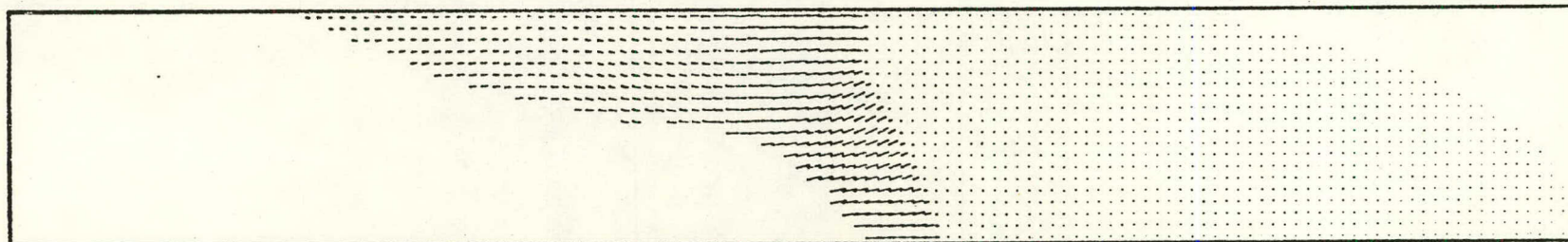
CYCLE=Y0, TIME=X% DENSITY SHADING VELOC



Z
TEMPERATURE CONTOURS



DENSITY SHADING



VELOCITY

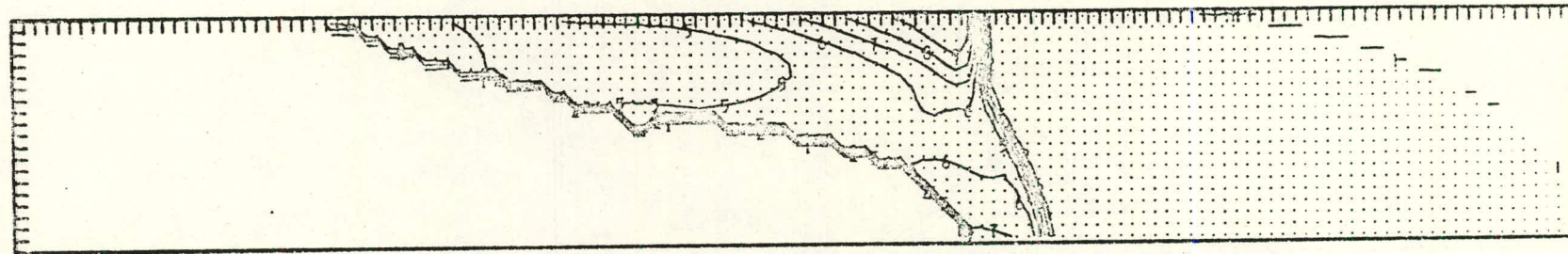
NUMERICAL SHOCK WAVES SIMULATIONS AT COLUMBIA

CYCLE= 451 , TIME= 0.179984E-06

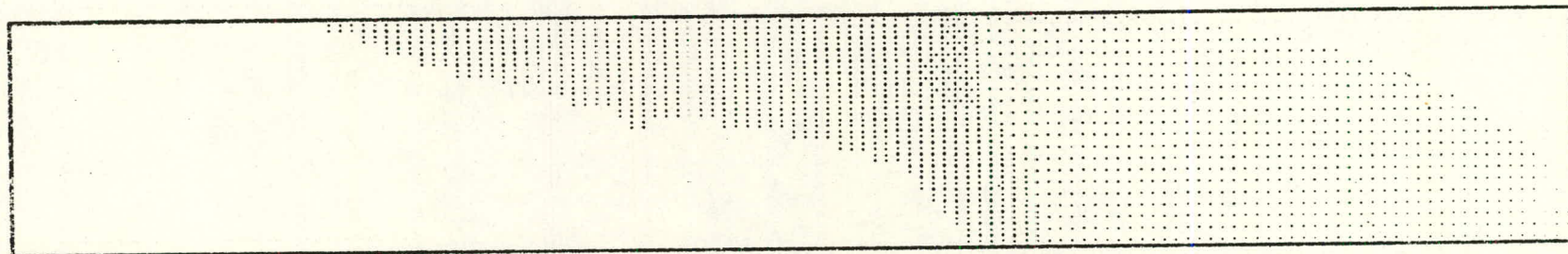
CONTOURS

1 =	0.000	2 =	0.762E+07	3 =	0.152E+08	4 =	0.229E+08	5 =	0.305E+08
6 =	0.381E+08	7 =	0.457E+08	8 =	0.533E+08	9 =	0.609E+08	10 =	0.686E+08

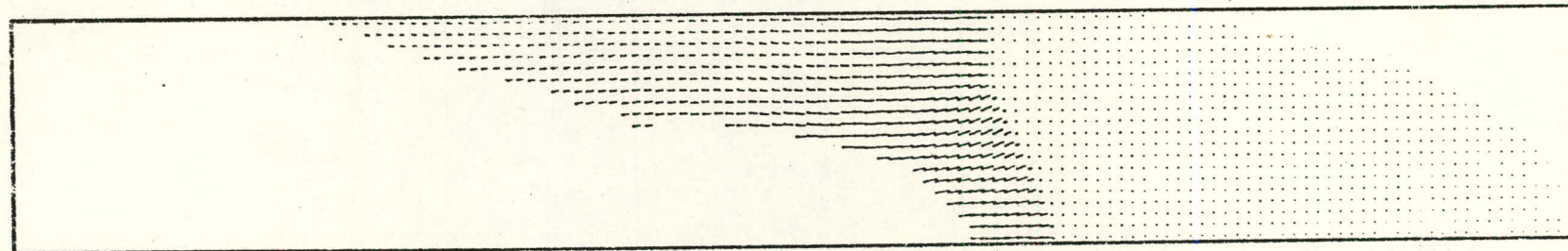
CYCLE=Y0, TIME=X0 DENSITY SHADING LEVEL



TEMPERATURE CONTOURS



DENSITY SHADING



VELOCITY

NUMERICAL SHOCK WAVES SIMULATIONS AT COLUMBIA

Fig. 4-2 The shock wave generated by a constant drive current in a coaxial shock tube

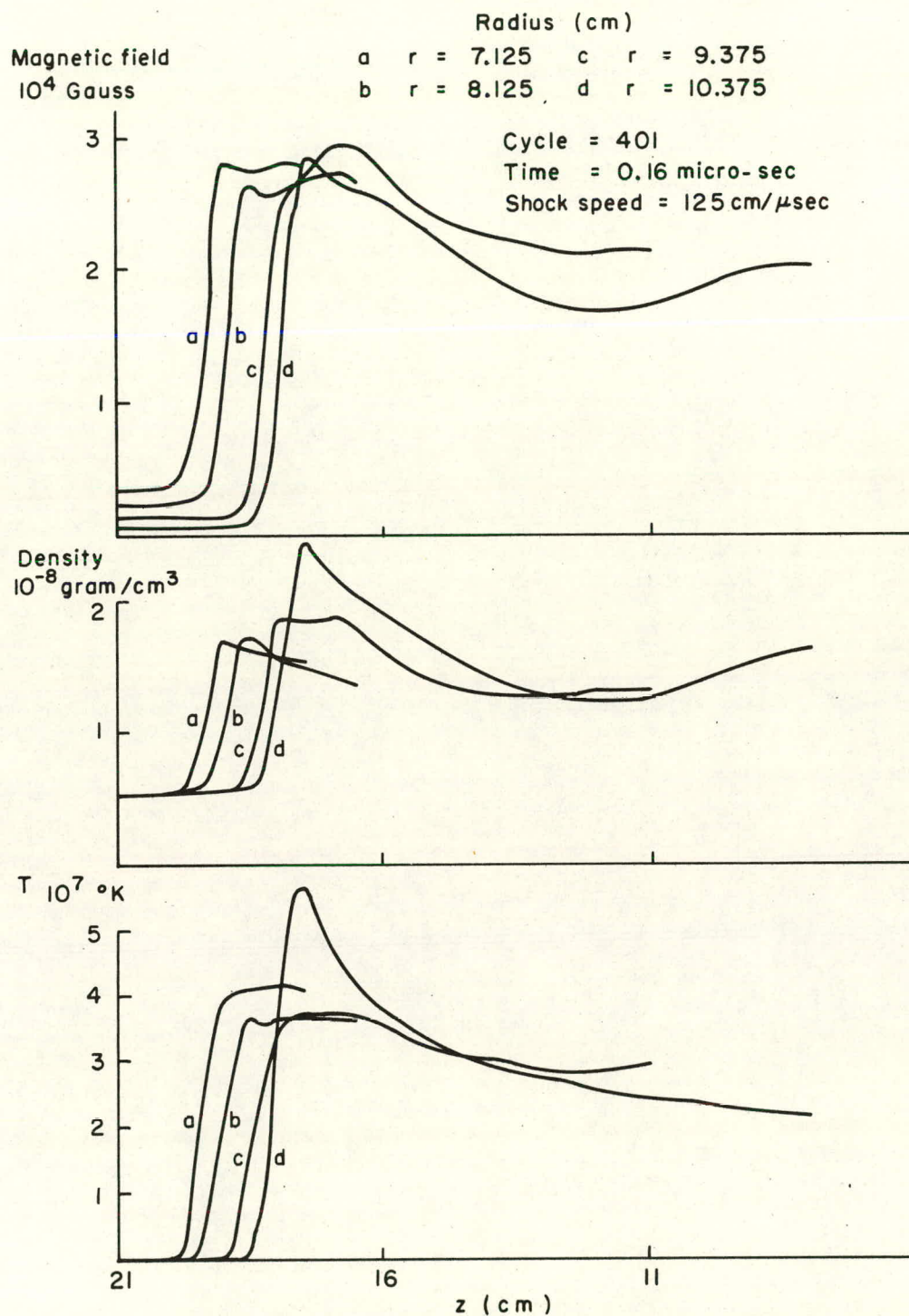


Fig. 4-3 The plots of magnetic field, density and temperature near the shock wave under a constant driver.

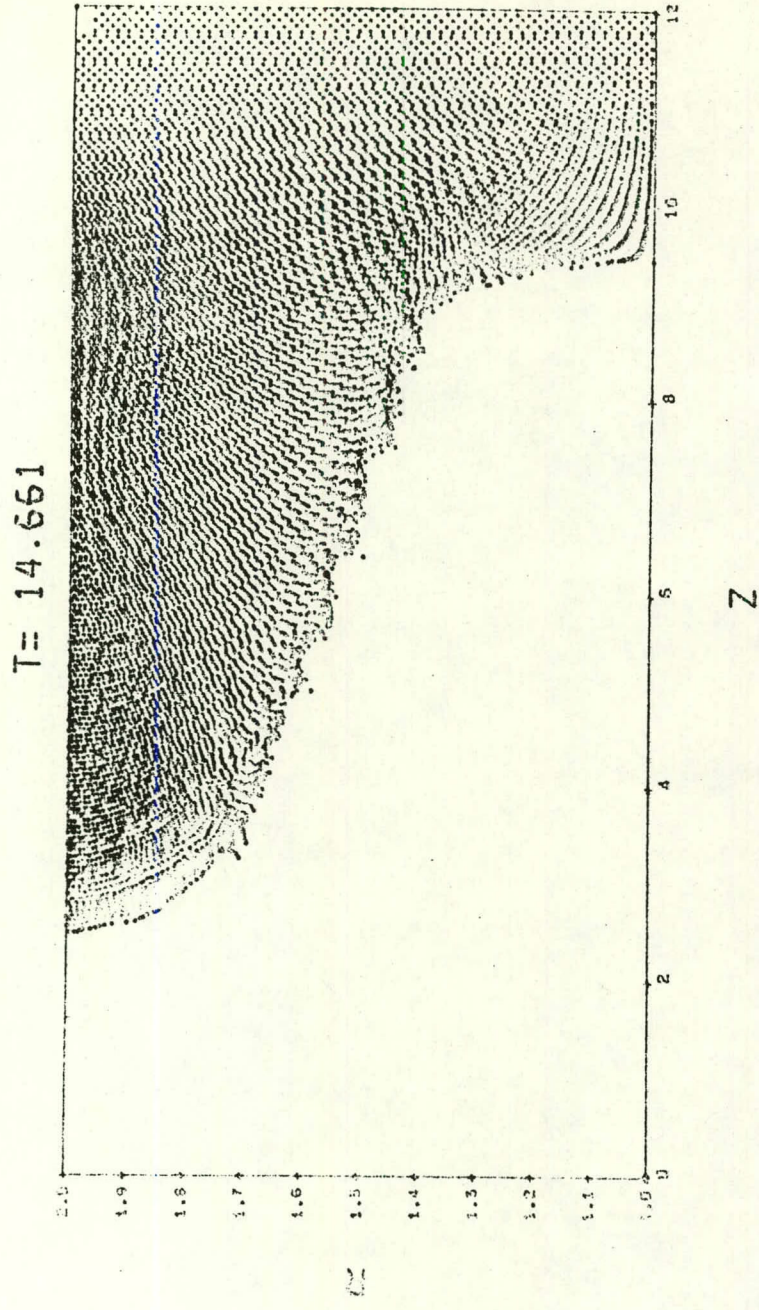


Fig. 4-4 The density distribution in the coaxial shock tube which was obtained by Marder using the GAP code. The drive magnetic field is 4 times the bias field.

4.3 Sinusoidal Drive Current

The drive current in the shock tube experiment is a sinusoidal function of time, this is a result of the LC circuit. For more realistic simulation and comparison with the experimental measurements, we therefore use the experimental data in Table 4-2 for the computation. Because the shock tube is 180 cm long, we use coarser cells (0.5 cm) in order to save computing time and core storages. The period of the sine function is 16 micro-seconds, the total time of the calculation is about 4 micro-seconds, at which time the drive current reaches maximum. The shock wave arrives at the end of the shock tube at 3.55 micro-seconds; but if the reflection of the shock wave is considered, the total time duration is 5 micro-seconds in our calculation.

The maximum amplitude (958.5 KA) of the sinusoidal current is smaller than the magnitude of the constant driver (1140 KA) in the previous section. Of course, the drive magnetic pressure induced by this current is smaller than that induced by the constant driver, and it will thus take a longer time to generate a shock wave. The fluid particles, and the shock wave will move more slowly under the influence of the sinusoidal driver than under the constant driver.

When the applied magnetic pressure in the vacuum is not high enough to produce a shock wave in the plasma, a compression wave will propagate down the tube and out toward the outer wall. The shock wave will be formed when the drive magnetic

pressure reaches a certain value at a critical moment of time. The mechanism for shock formation under the sinusoidal driver is the same as that under the constant driver. But, the shock will increase in strength with the increase of the drive current in time. The fluid particles will also increase their velocities. These are shown in Table 4-3, in which the post-shock temperatures are measured at a radius of 8.5 cm. The temperature (2.4 Kev) of the deuterium plasma measured at 3.32 micro-seconds with the sinusoidal drive current (924.6 KA) is lower than the temperature (3.8 Kev) of the hydrogen plasma with the constant driver (1140 KA) in the previous section. However, the change of the drive current with respect to time will be small enough when it is near the maximum amplitude. The flow conditions of the shock wave will not be much different from those in the case of a constant driver with the same magnitude.

Fig. 4-5 shows the magnetic field, density and temperature plotted with respect to axial position at three different radii in the tube after 3.32 micro-seconds. At a radius of 8.5 cm, the flow conditions behind the shock wave are almost uniform. The hump in the curve (c) of the magnetic field indicates the compression of the plasma on the outer wall.

Fig. 4-6 shows the time trace of the magnetic field measured at a fixed station of 100 cm with a radius of 10 cm in the tube. It agrees with the upper trace (a) in Fig. 4-7, which was obtained in our shock tube experiment.

Table 4-2

Initial conditions in a coaxial shock tube of 180 cm

Working gas	deuterium
Density	1.0688×10^{-8} gram/cm ³
Temperature	300 °K
Pressure	66.65 dynes/cm ²
Bias magnetic field	60,000/r Gauss
Drive magnetic field	$(300 + 958.5 \sin (2\pi t \times 10^6/16))$ x200/r Gauss
Inner radius	6.25 cm
Outer radius	11.25 cm
Cell width ($\delta z = \delta r$)	0.5 cm

Table 4-3

The relation between shock speed and drive current (deuterium plasma)

Time (t)	Drive Current ($958.5 \sin(2\pi t 10^6/16)$)	Shock Speed	Post-shock temp.
1.20 μsec	435.1 KA	43 cm/ μsec	$0.08 \times 10^7 \text{ } ^\circ\text{K}$
1.56	550.7	53	0.37
1.86	639.7	59.8	0.68
2.14	713.2	64.0	1.00
2.39	774.2	67.3	1.50
2.64	824.9	72.5	1.76
2.87	866.3	73.5	2.00
3.10	899.4	77.3	2.20
3.32	924.6	78.2	2.40

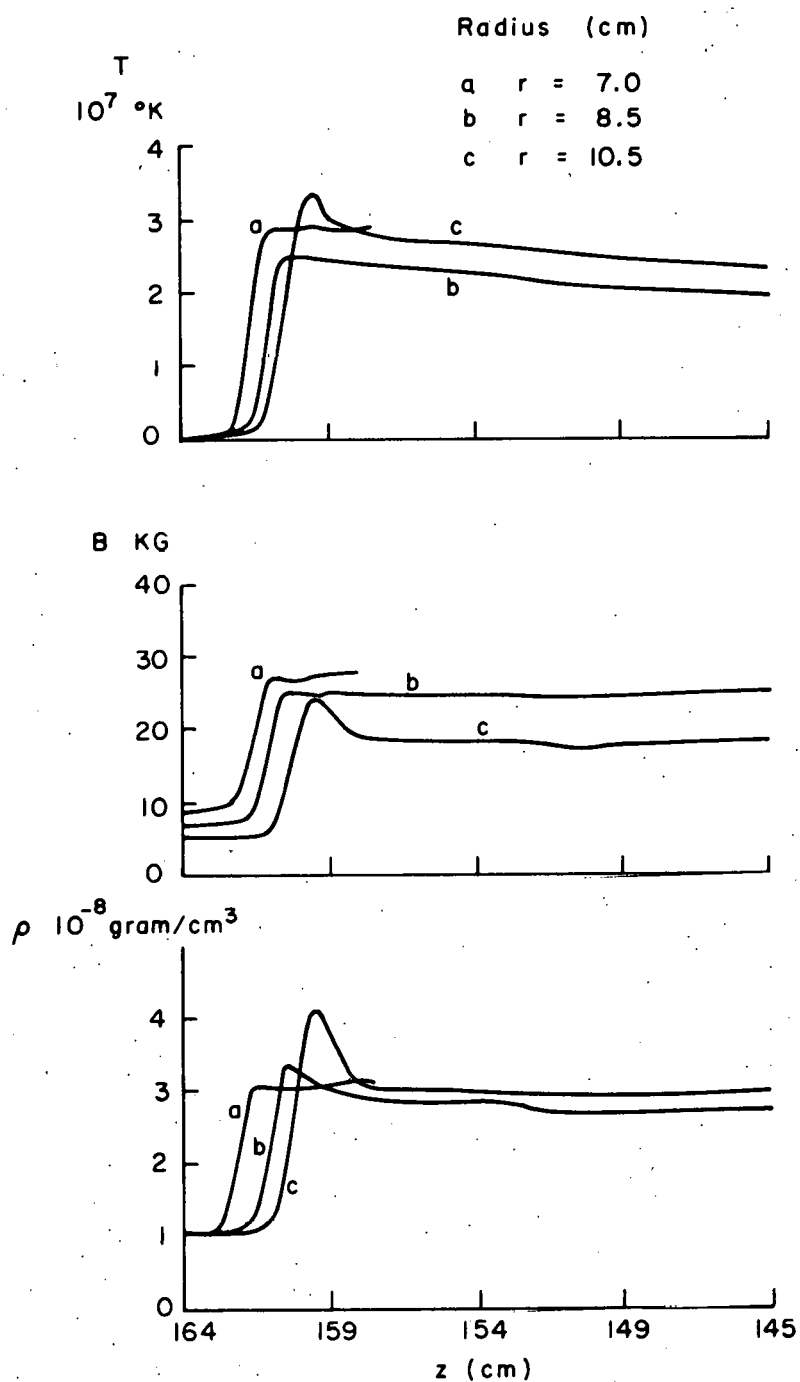


Fig. 4-5 The temperature, magnetic field and density vs. axial positions at different radii.

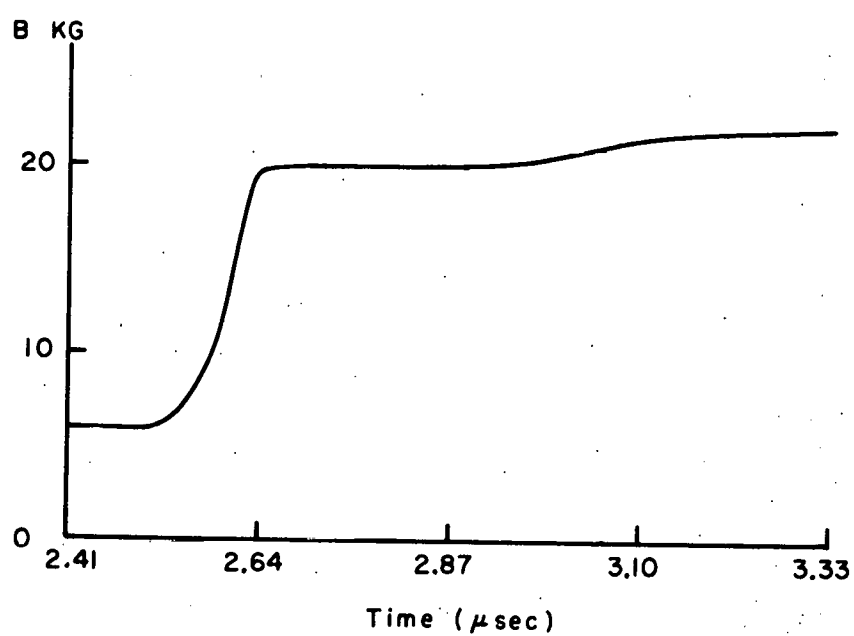


Fig. 4-6 The trace of the magnetic field measured at a fixed station (100 cm) with a radius of 10 cm .

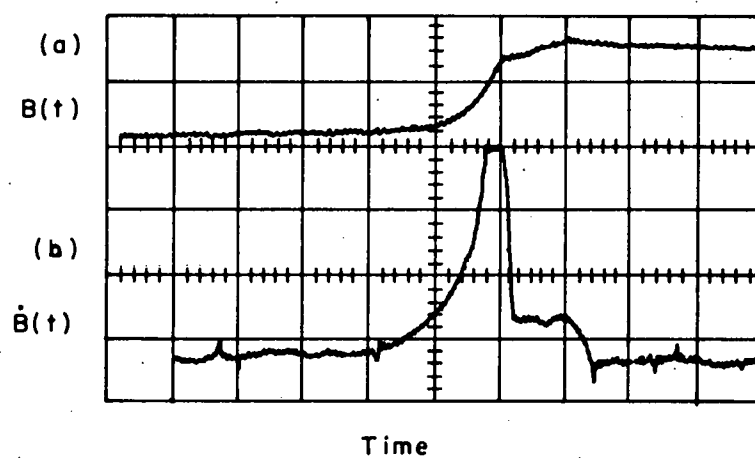


Fig. 4-7 A typical magnetic probe trace from the Columbia University shock tube, taken from a station 100cm down the shock tube.

4-501- MC-2151

V. SHOCK REFLECTION AND CONVERGENCE

5.1 Flat End Wall

5.1.1 Constant Driver

To reflect a shock wave in the simplest way, we locate a flat end wall in the shock tube. The shock wave in Section 4.2, with a constant drive current of 1140 KA, is employed to study its reflection from the end wall. The fluid particles behind the shock wave travel in the range of 110 cm/ μ sec, the incident shock speed is about 125 cm/ μ sec. Other flow conditions of the incident shock wave are given in Section 4.2.

Fig. 5-1 shows the sequence of the shock wave moving down the tube and reflecting from the flat end wall, which is at the right end. At cycle=651, the shock wave arrives at the end wall. At cycle=701, the shock wave just reflects off the wall surface. After shock reflection, the plasma near the outer wall again remains the hottest with the temperature 14.5×10^7 °K, density 5.1×10^{-8} gram/cm³, and magnetic field 5.3×10^4 Gauss. At cycle=801, the reflected shock wave travels downstream, and the hot plasma expands into the vacuum near the inner conductor.

Fig. 5-2 shows the traces of the magnetic field, density and temperature with respect to time at a fixed station, which is 5 cm in front of the flat end wall with a radius of 9.375 cm. The first jump in the trace of the magnetic field indicates the arrival of the incident shock wave at the fixed station; the

second jump shows the return of the reflected shock wave.

The total pressure ($P^* = 2.4 \times 10^8$ dynes/cm²) in the hot plasma after shock reflection is much higher than the applied magnetic pressure (0.4×10^8 dynes/cm² at the mean radius) in the vacuum so that the reflected shock and particles are accelerated to expand into the vacuum quickly. The magnetic field, density and temperature decrease after the reflected shock wave moves away from the fixed station.

The relation between the incoming shock and the reflected shock along a radius of 9.375 cm in the tube is also demonstrated by the plots of the magnetic field, density and temperature with respect to axial position at different times in Fig. 5-3. The flat end wall is located at 32.5 cm. There are small numerical oscillations in the density and magnetic field just behind the reflected shock wave. The reflected shock speed is about 100 cm/μsec, slower than the incident shock speed of 125 cm/μsec.

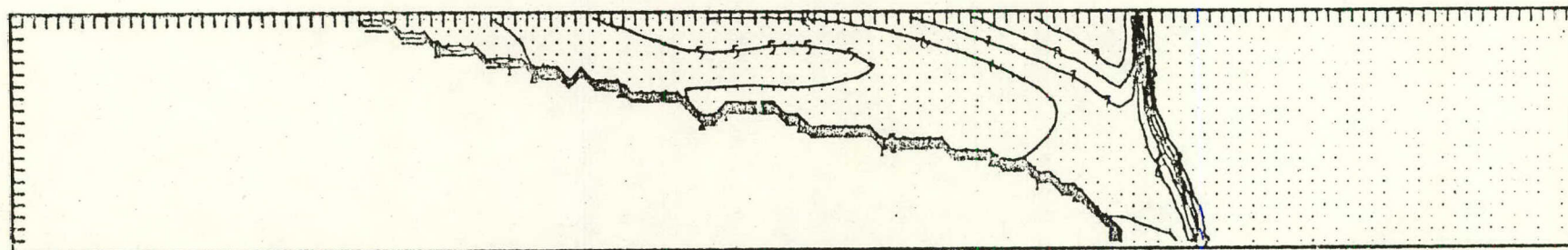
Fig. 5-1

CYCLE= 501 , TIME= 0.199982E-06

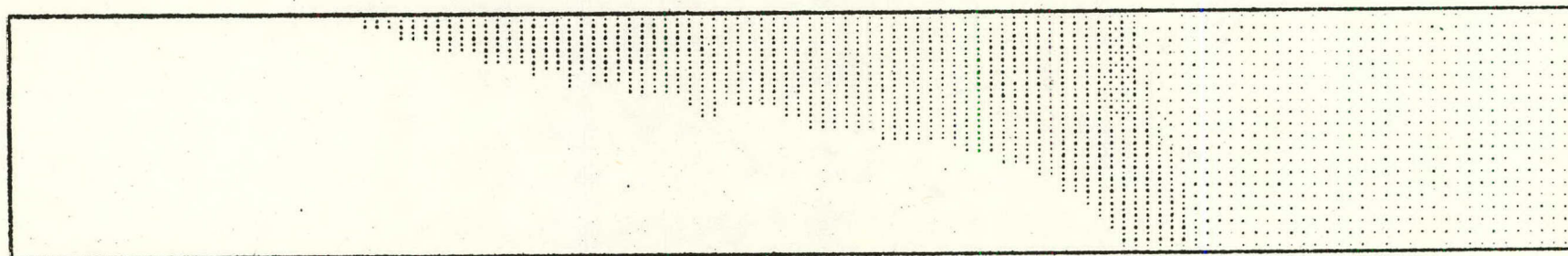
CONTOURS

1 =	0.000	2 =	0.734E+07	3 =	0.147E+08	4 =	0.220E+08	5 =	0.294E+08
6 =	0.367E+08	7 =	0.441E+08	8 =	0.514E+08	9 =	0.587E+08	10 =	0.661E+08

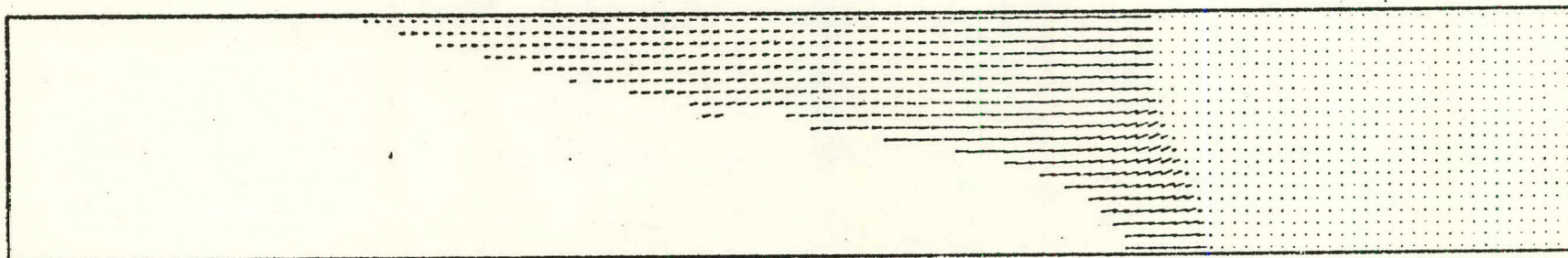
CYCLE=Y , TIME=Z DENSITY SHADING



Z
TEMPERATURE CONTOURS



DENSITY SHADING



VELOCITY
NUMERICAL SHOCK WAVES SIMULATIONS AT COLUMBIA

Fig. 5-1

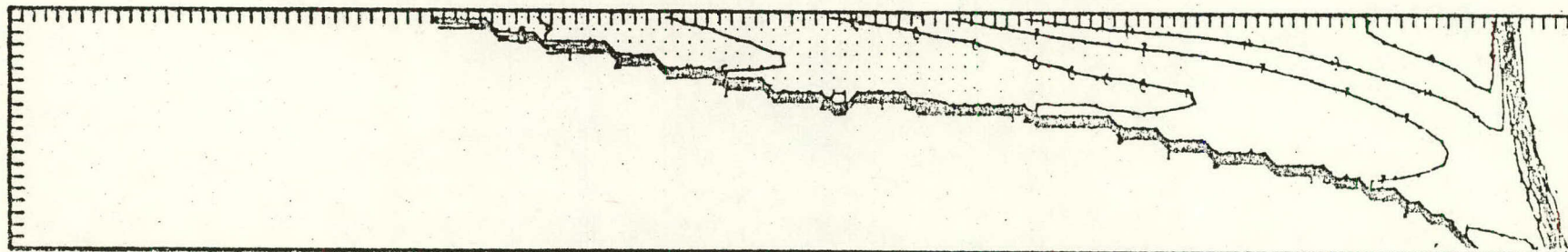
CYCLE= 651
CONTOURS

, TIME= 0.259974E-06

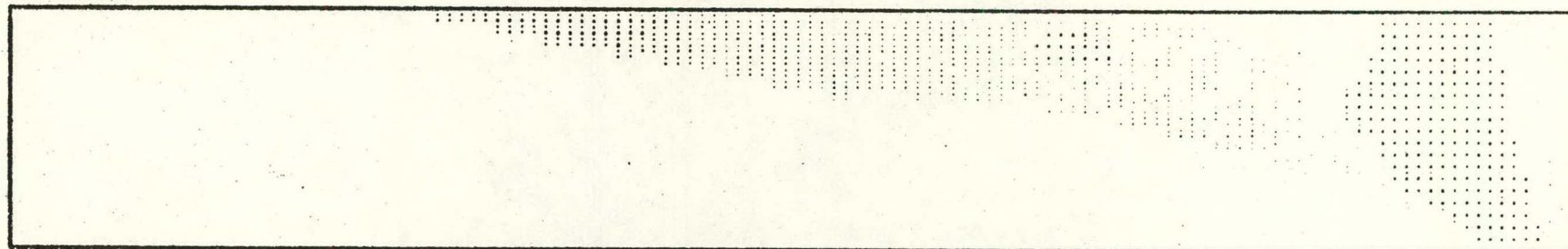
Page

1 =	0.000	2 =	0.652E+07	3 =	0.130E+08	4 =	0.196E+08	5 =	0.261E+08
6 =	0.326E+08	7 =	0.391E+08	8 =	0.457E+08	9 =	0.522E+08	10 =	0.587E+08

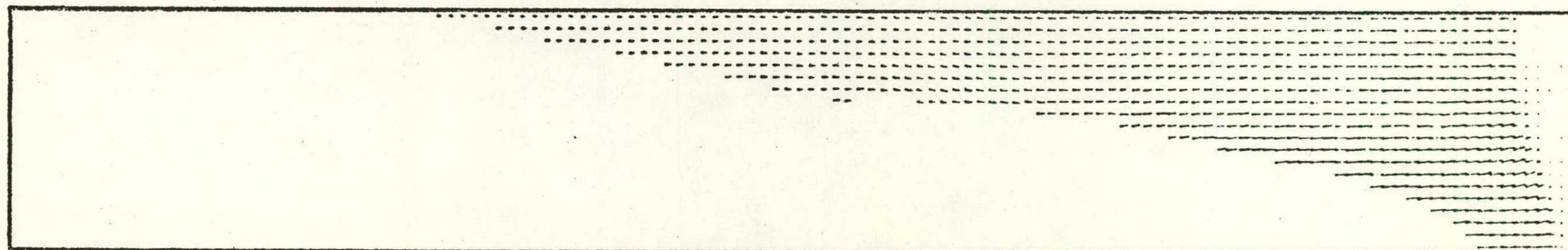
CYCLE=, TIME= DENSITY SHADING



Z
TEMPERATURE CONTOURS



DENSITY SHADING



VELOCITY

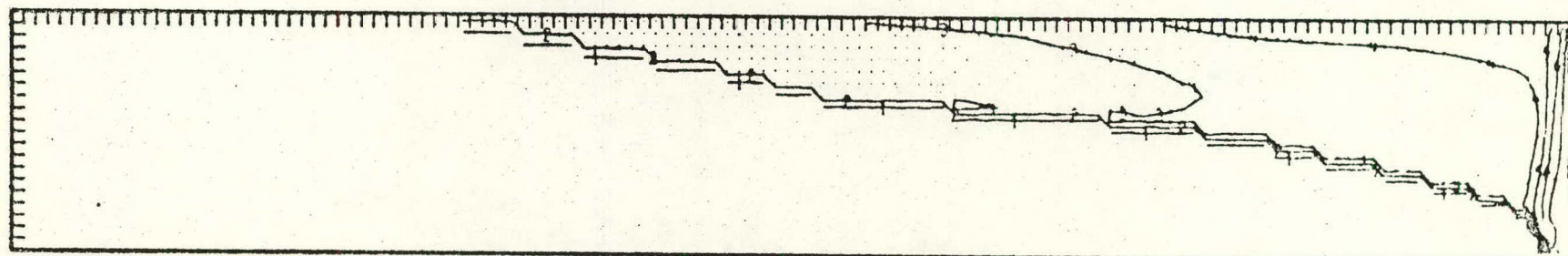
NUMERICAL SHOCK WAVES SIMULATIONS AT COLUMBIA

Fig. 5-1

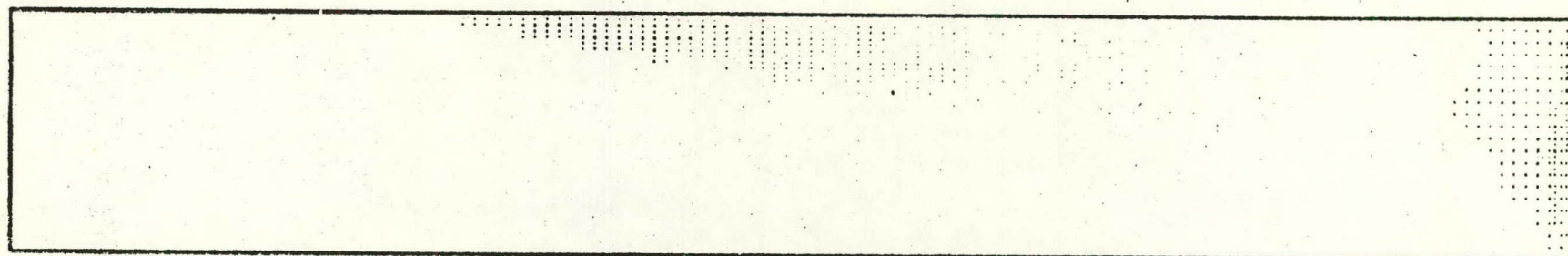
CYCLE= 701 , TIME= 0.279972E-06
CONTOURS

1 =	0.000	2 =	0.164E+08	3 =	0.328E+08	4 =	0.492E+08	5 =	0.656E+08
6 =	0.820E+08	7 =	0.984E+08	8 =	0.115E+09	9 =	0.131E+09	10 =	0.148E+09

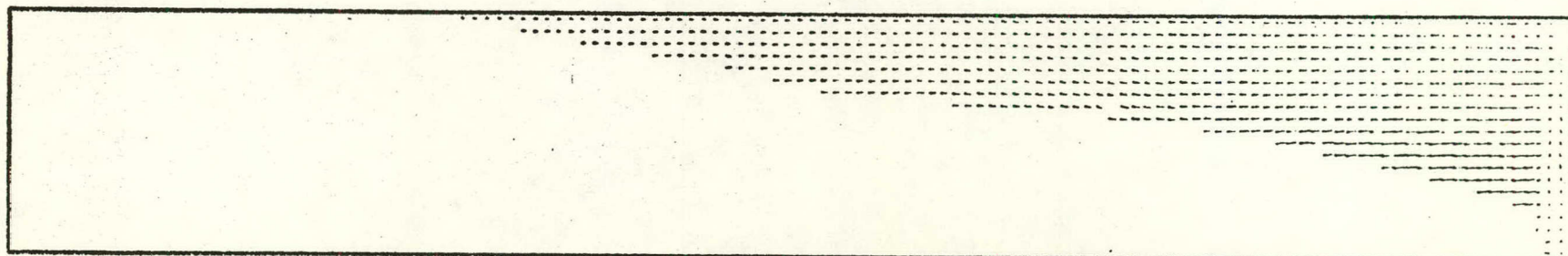
CYCLE-Y TIME DENSITY SHADING



TEMPERATURE CONTOURS



DENSITY SHADING



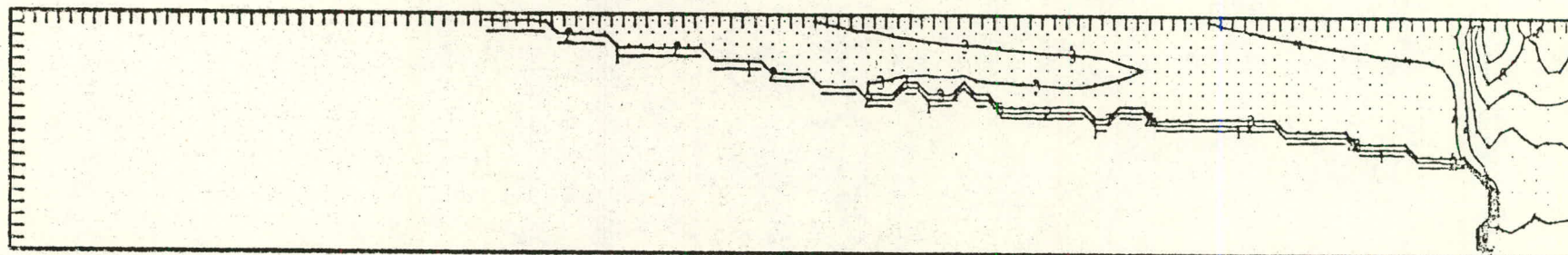
VELOCITY

NUMERICAL SHOCK WAVES SIMULATIONS AT COLUMBIA

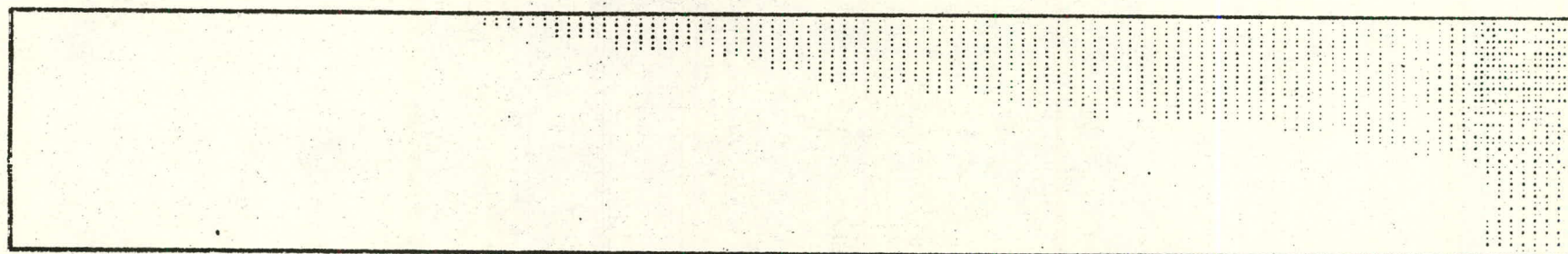
CYCLE= 751 , TIME= 0.299969E-06
CONTOURS

1 =	0.000	2 =	0.149E+08	3 =	0.297E+08	4 =	0.446E+08	5 =	0.594E+08
6 =	0.743E+08	7 =	0.892E+08	8 =	0.104E+09	9 =	0.119E+09	10 =	0.134E+09

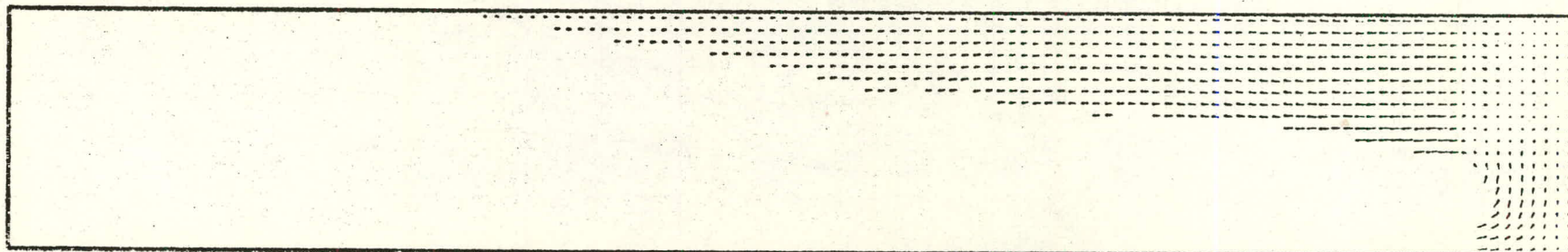
CYCLE=Y, TIME=Z DENSITY SHADING



Z
TEMPERATURE CONTOURS



DENSITY SHADING



VELOCITY
NUMERICAL SHOCK WAVES SIMULATIONS AT COLUMBIA

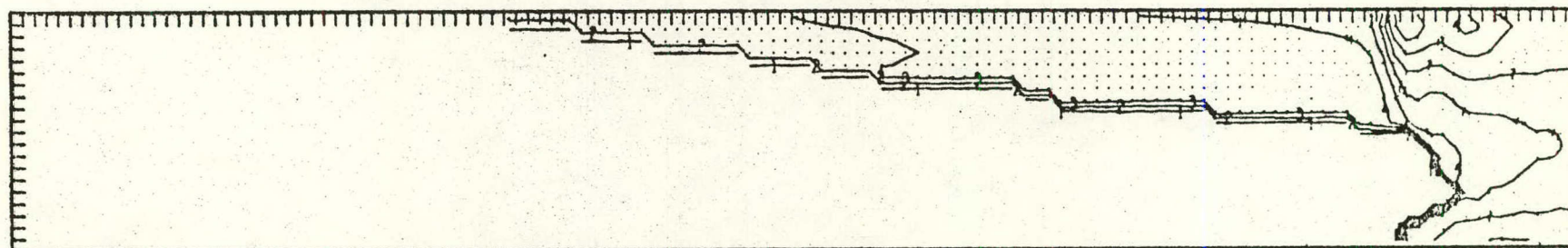
Fig. 5-1

Page

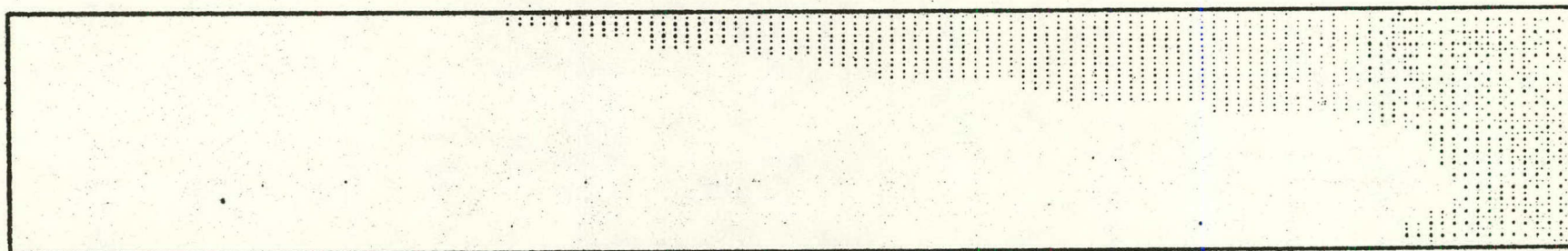
CYCLE= 801 , TIME= 0.319967E-06
CONTOURS

1 =	0.000	2 =	0.136E+08	3 =	0.271E+08	4 =	0.407E+08	5 =	0.543E+08
6 =	0.678E+08	7 =	0.814E+08	8 =	0.950E+08	9 =	0.109E+09	10 =	0.122E+09

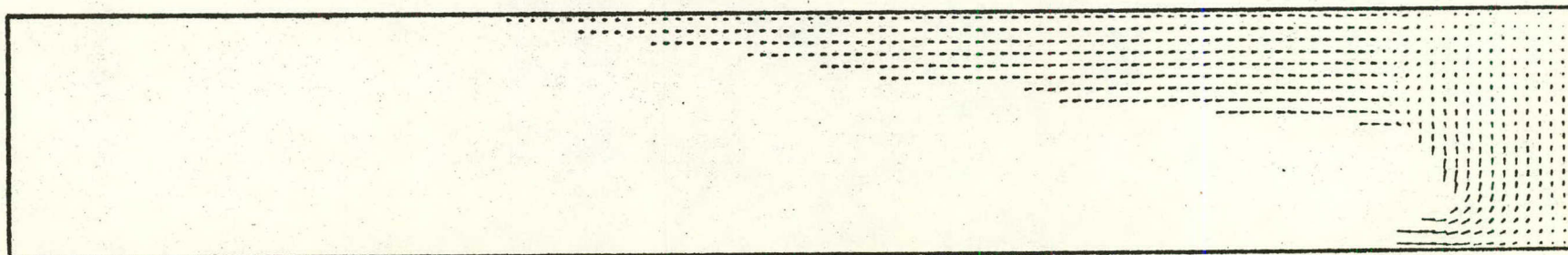
CYCLE=Y , TIME= DENSITY SHADING



Z
TEMPERATURE CONTOURS



DENSITY SHADING



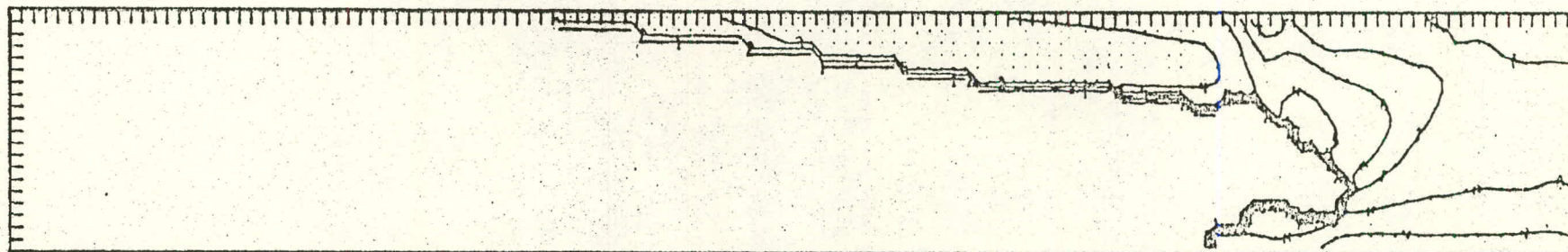
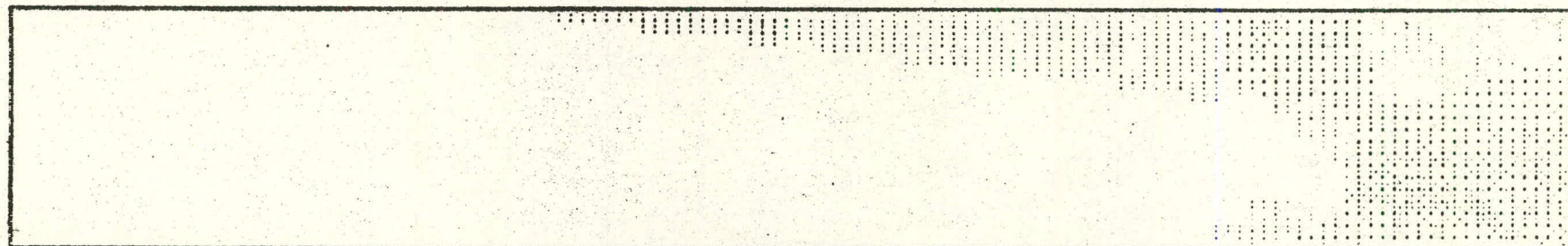
VELOCITY
NUMERICAL SHOCK WAVES SIMULATIONS AT COLUMBIA

CYCLE= 901 , TIME= 0.359962E-06

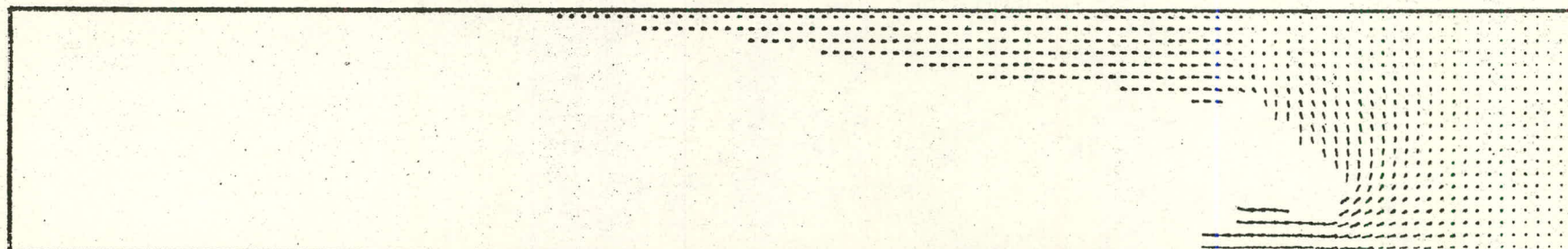
CONTOURS

1 =	0.000	2 =	0.107E+08	3 =	0.215E+08	4 =	0.322E+08	5 =	0.429E+08
6 =	0.537E+08	7 =	0.644E+08	8 =	0.751E+08	9 =	0.859E+08	10 =	0.966E+08

CYCLE=Y , TIME= DENSITY TRAINING


2
TEMPERATURE CONTOURS


DENSITY SHADING



VELOCITY

NUMERICAL SHOCK WAVES SIMULATIONS AT COLUMBIA

Fig. 5-1 The incident shock wave and its reflection from a flat end wall (constant driver)

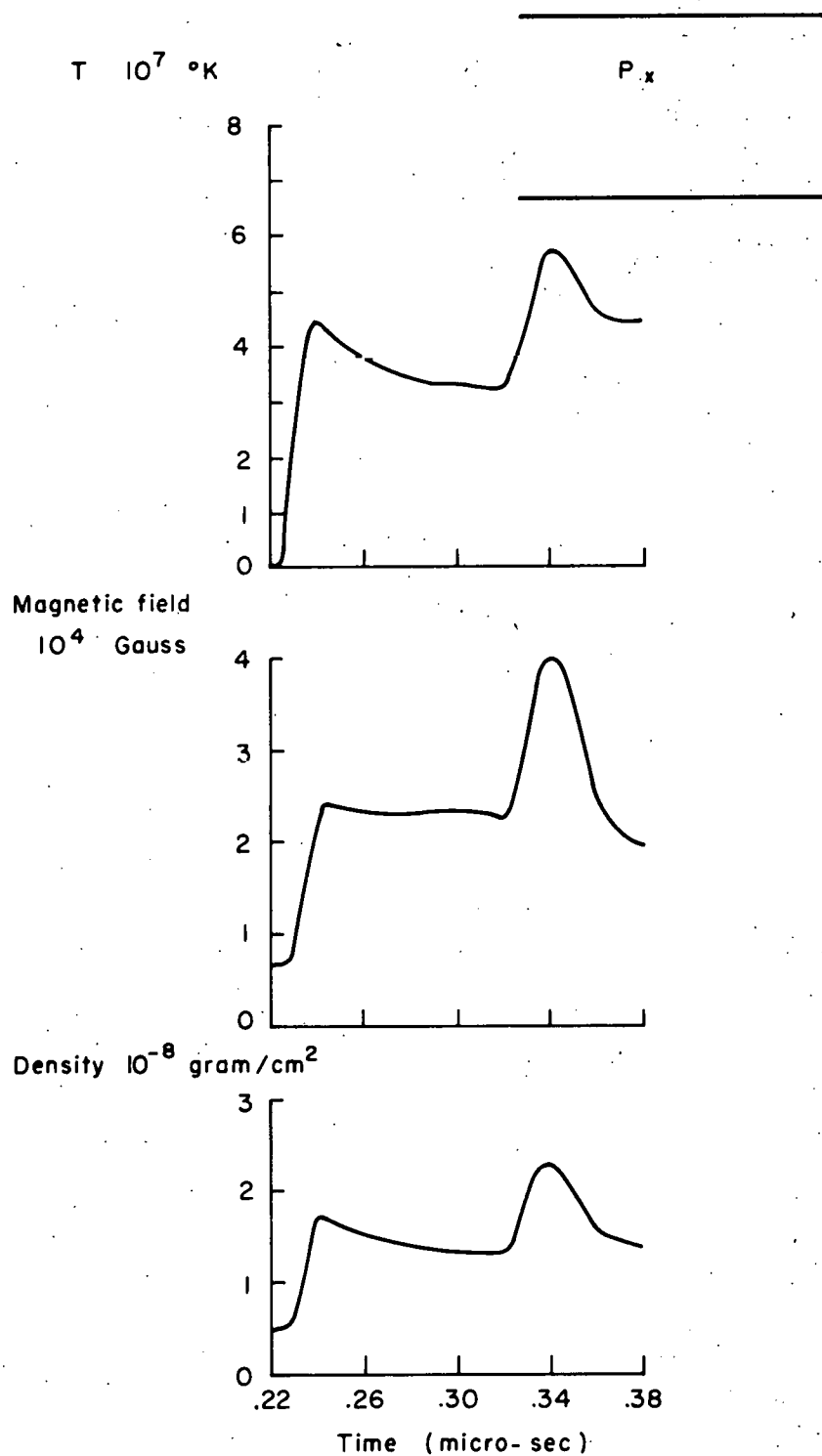


Fig. 5-2 The traces of the temperature, magnetic field and density with respect time at the fixed station (P) located at 5 cm in front of the flat end wall with the radius equal to 9.375 cm.

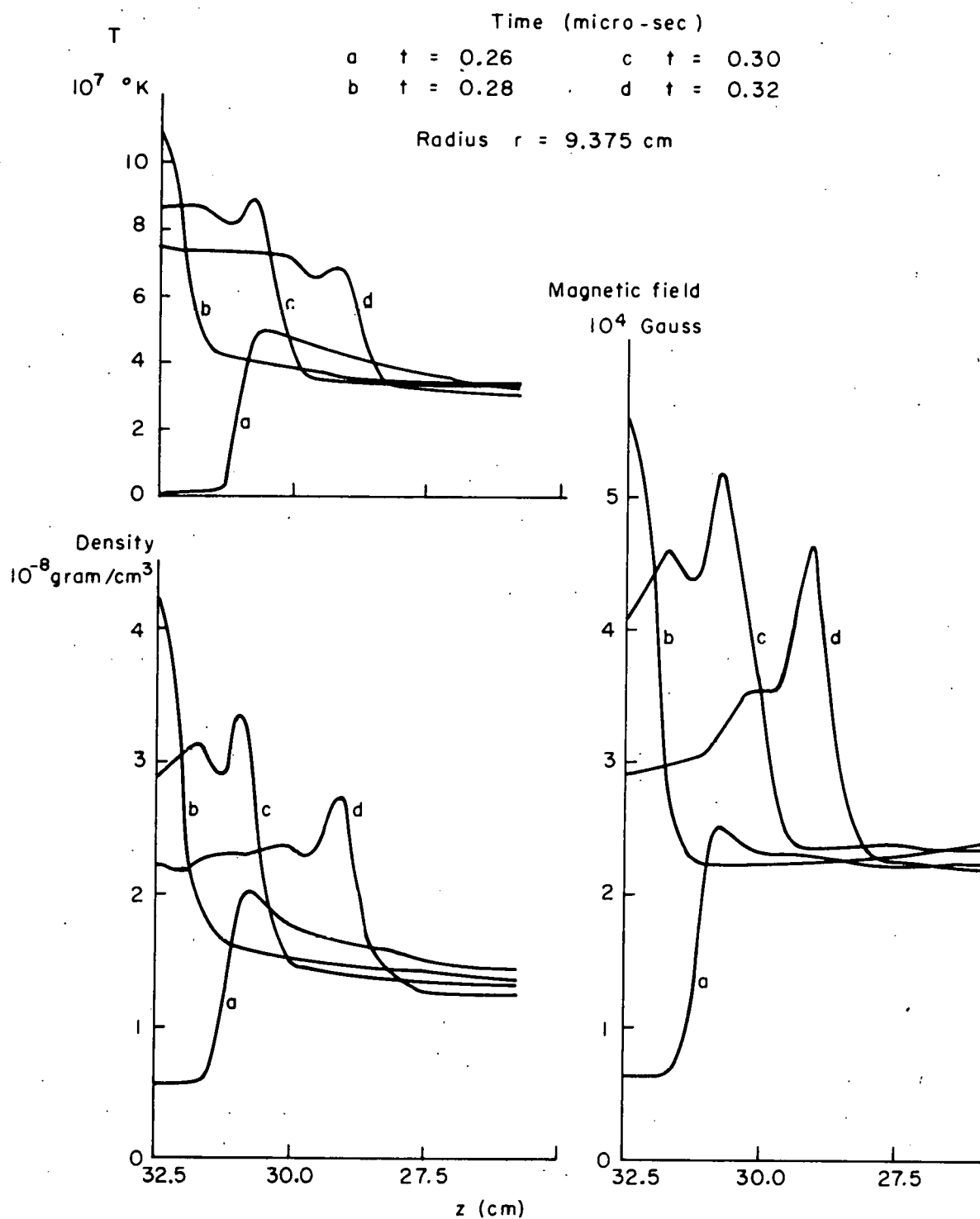


Fig. 5-3 The plots of the magnetic field, density and temperature with respect to the axial positions near the flat end wall with the constant driver.

5.1.2 Sinusoidal Driver

The shock wave in Section 4.3 at 3.32 micro-seconds is utilized so that we can compare its reflection from a flat end wall with experimental measurements. The flow conditions of the incident shock wave are given in Section 4.3. The shock wave runs down the tube with a speed of about 78.2 cm/ μ sec.

Owing to the slope of the incident shock wave, the impact of the shock wave on the outer corner of the end wall produces a hot plasma with a temperature of about 7.8 Kev. Although the plasma in the tube may flow two-dimensionally, the reflected shock wave travels downstream almost like a simple one-dimensional reflection.

Fig. 5-4 shows the plots of the temperature, magnetic field, and density with respect to axial position in front of the flat end wall, which is at 180 cm, at a radius of 9 cm. The jumps of the temperature, magnetic field, and density between the reflected shock wave are from 2 Kev, 22 KG, and 2.6×10^8 gram/cm³ to 5 Kev, 34 KG, and 4.6×10^8 gram/cm³ respectively. The reflected shock speed is about 79 cm/ μ sec at a radius of 9 cm, slightly greater than the incident shock speed 78.2 cm/ μ sec. The incoming shock wave is in front of the flat end wall at 3.538 μ sec. At 3.644 μ sec, the shock wave reflects from the end wall and travels downstream. The motion of the shock wave and the plasma expansion are similar to those in the case of constant driver respectively.

Fig. 5-5 shows the time trace of the computed magnetic

field at a fixed position of 175 cm with a radius of 9 cm in the shock tube. Qualitatively, it agrees with the experimental measurement in Fig. 5-6. The quantitative comparison between them is given below.

	Magnetic Field (KG)		
	bias	post-shock	maximum after reflection
computed (Fig. 5-5)	6.6	24	36.8
experimental (Fig. 5-6)	6.6	19.8	42.9

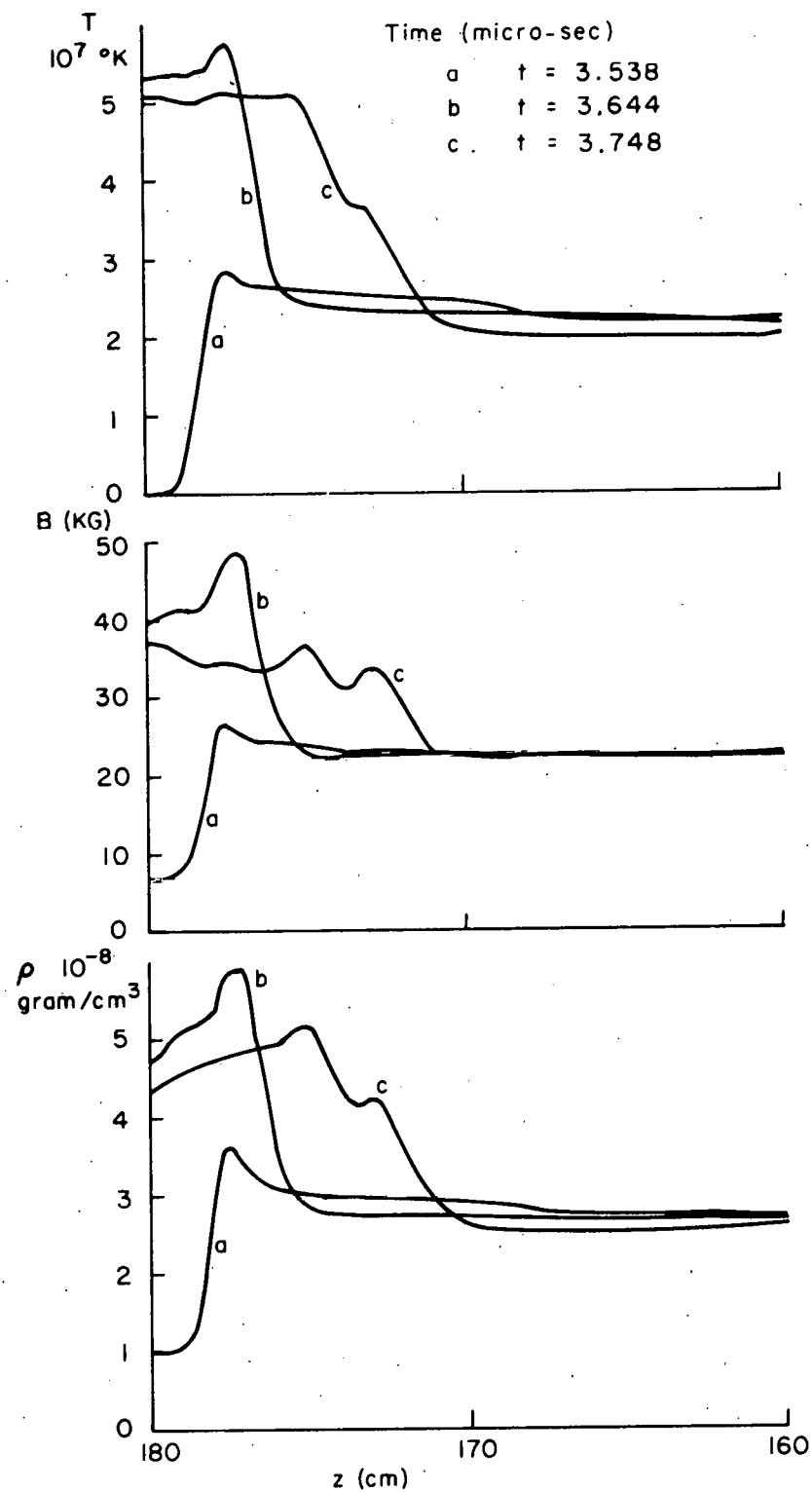


Fig. 5-4 The temperature, magnetic field and density at a radius of 9cm plotted versus axial positions in front of the flat end wall at different times.

Tube length : 180 cm
 Incoming shock speed : 78 cm/micro-sec
 End wall : flat
 Magnetic probe : located at the position P
 Other flow variables :

Time	T (°K)	ρ (gram/cm ³)
A	600	$.10688 \times 10^{-7}$
B	2.59×10^7	$.3036 \times 10^{-7}$
C	5.04×10^7	$.5194 \times 10^{-7}$
D	3.89×10^7	$.323 \times 10^{-7}$

Magnetic field

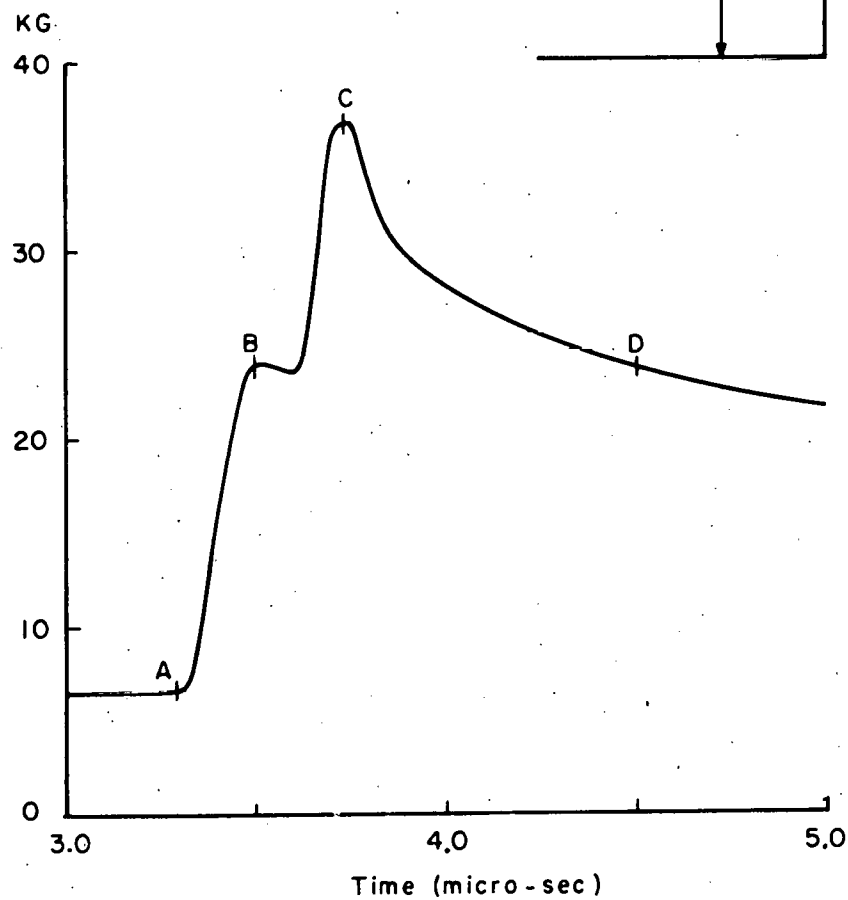


Fig. 5-5 The magnetic field vs. time measured at the position P.

W. L. A-907-ME-27

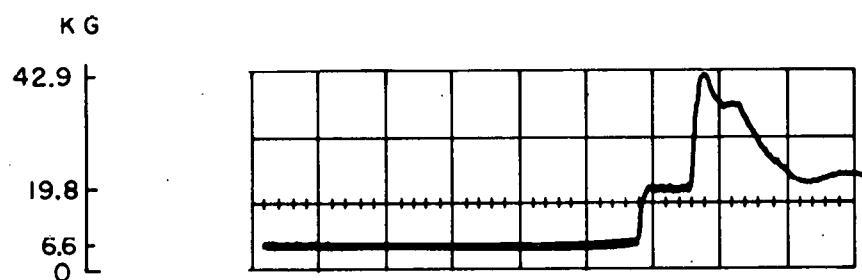


Fig. 5-6 The magnetic trace measured experimentally at the fixed station (175cm) in front of the flat end wall in the shock tube which is 180cm long.

5.2 Semi-circular End Wall

In order to achieve better plasma heating by shock convergence and reflection, we consider a curved end wall instead of a flat end wall. From a practical point of view, it would be advantageous to use a semi-circular end wall instead of other curves, since it is easiest to manufacture. The same shock wave in Section 5.1.1 is again employed, so that we may study its convergence and reflection from such a semi-circular end wall, and compare the results with those of the flat end wall.

Fig. 5-7 shows the sequence of the shock flow in the tube with the wall profile at the right end. When the shock wave moves along the circular wall surface, it is strengthened near the side wall surfaces. The shock wave does not fully converge when it arrives at the end of the tube. Because the incident shock wave has a slope greater than 90° with respect to the axial axis, the converged and reflected shock position is at a radius of 9.375 cm near the end wall surface.

At cycle = 701, the shock wave reflects off the wall surface. The plasma in front of the reflected shock still moves toward the focused position, but the hot plasma behind the reflected shock wave moves in the opposite direction. The reflected shock wave is approximately circular in shape in the focused region. Later on, its shape becomes flat when it reflects to the straight part of the tube. The maximum value of the temperature at the focused position is about 16.7 Kev, and the corresponding density and magnetic field are 5.36×10^{-8} gram/cm³

and 6.8×10^4 Gauss respectively. This temperature is only slightly higher than the 14.5 Kev obtained by flat wall reflection. Thus, the semi-circular end wall is not a very effective shape for plasma heating. (We shall see later that in plane geometry, it is actually more effective).

From the dynamical point of view, the plasma in a high pressure region will be driven to a lower pressure region. At cycle = 751, the reflected shock wave travels downstream, and the hot plasma expands into the vacuum near the inner conductor. The temperature contours indicate that there are recompressions on the side walls when the reflected shock wave reaches the straight part of the shock tube.

Fig. 5-8 shows the reflection of the shock wave from the semi-circular wall with flow variables plotted versus the axial position in front of the end wall, at a radius of 9.375 cm in the shock tube. The curve (a) of the temperature shows the incident shock wave; the curve (b) shows the reflection of the converged shock wave; and the curve (c) shows the reflected shock wave travelling downstream. When the hot plasma expands into the vacuum and the reflected shock wave moves away from the focused region, the magnetic field, density and temperature there decrease quickly.

Since the semi-circular end wall is not a good shape for plasma heating, in the next chapter, we shall find a wall shape which is more successful in converging shock waves using Whitham's shock-ray approximation.

Fig. 5-7

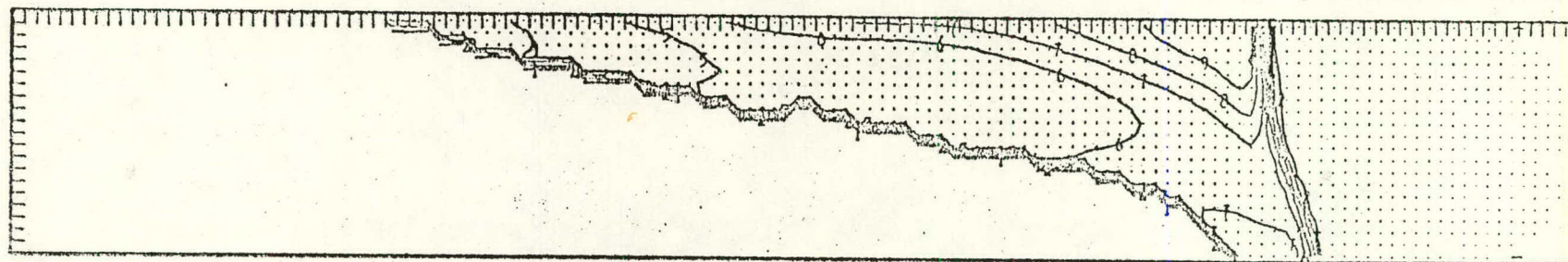
CYCLE= 551 , TIME= 0.219979E-06

CONTOURS

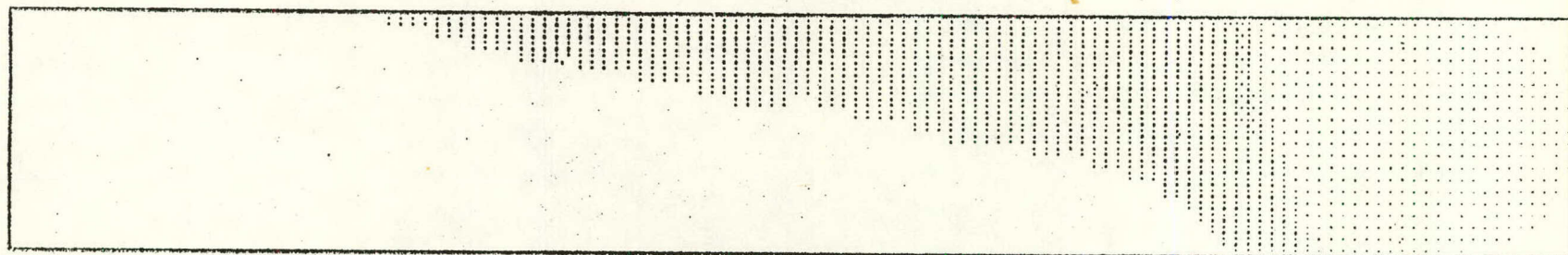
1 =	0.000	2 =	0.702E+07	3 =	0.140E+08	4 =	0.210E+08	5 =	0.281E+08
6 =	0.351E+08	7 =	0.421E+08	8 =	0.491E+08	9 =	0.561E+08	10 =	0.631E+08

PAGE 1

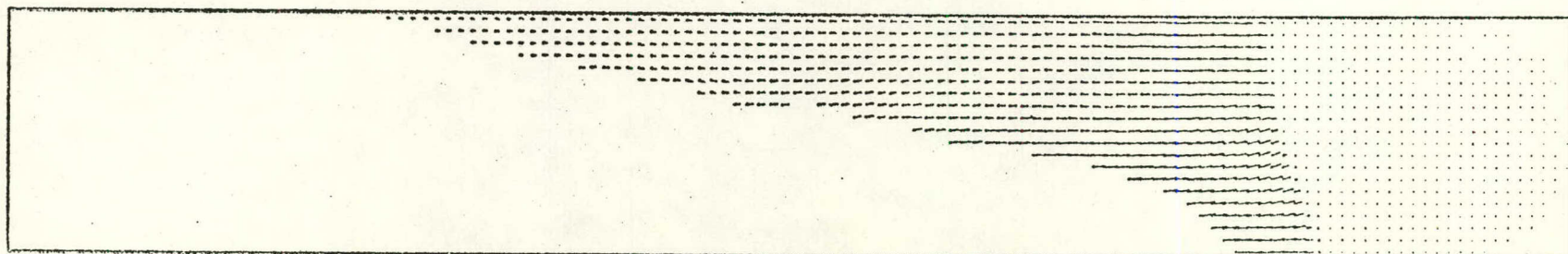
CYCLE= , TIME= DENSITY SHADING



TEMPERATURE CONTOURS



DENSITY SHADING



VELOCITY

Fig. 5-7

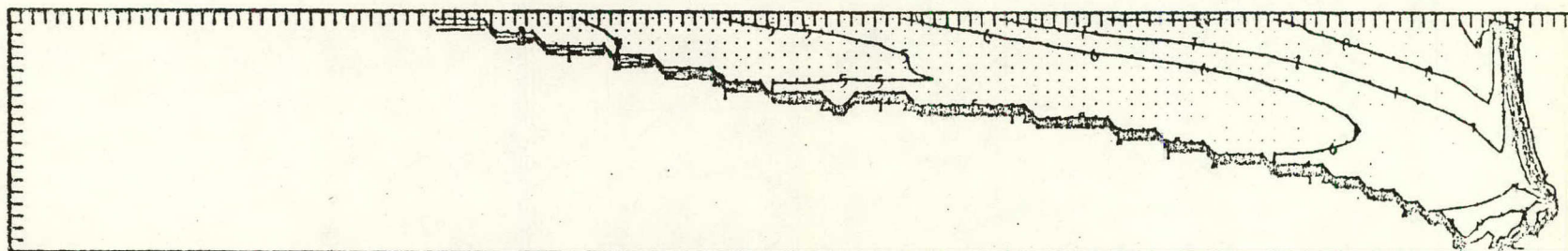
Page

CYCLE= 651 , TIME= 0.259974E-06

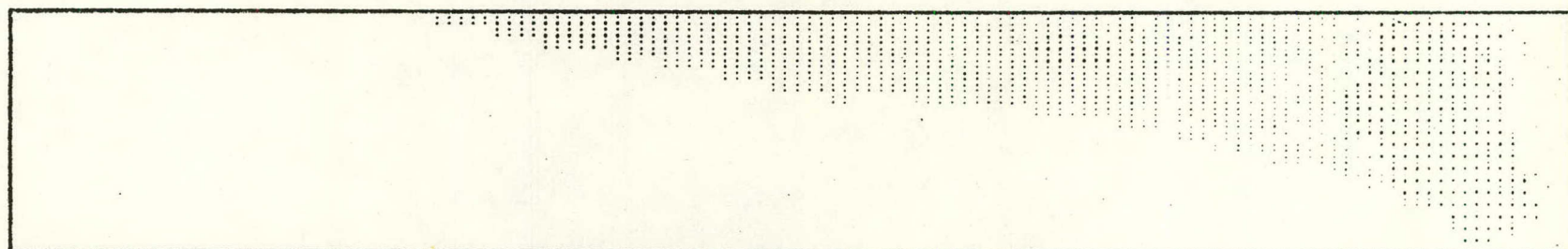
CONTOURS

1 =	0.000	2 =	0.719E+07	3 =	0.144E+08	4 =	0.216E+08	5 =	0.268E+08
6 =	0.360E+08	7 =	0.432E+08	8 =	0.504E+08	9 =	0.576E+08	10 =	0.648E+08

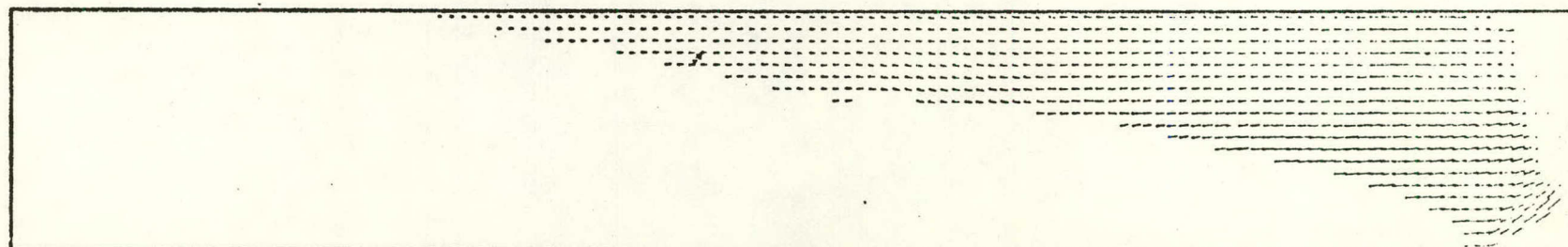
CYCLE= , TIME= DENSITY SHADING



TEMPERATURE CONTOURS



DENSITY SHADING



VELOCITY
NUMERICAL SHOCK WAVES SIMULATIONS AT COLUMBIA

Fig. 5-7

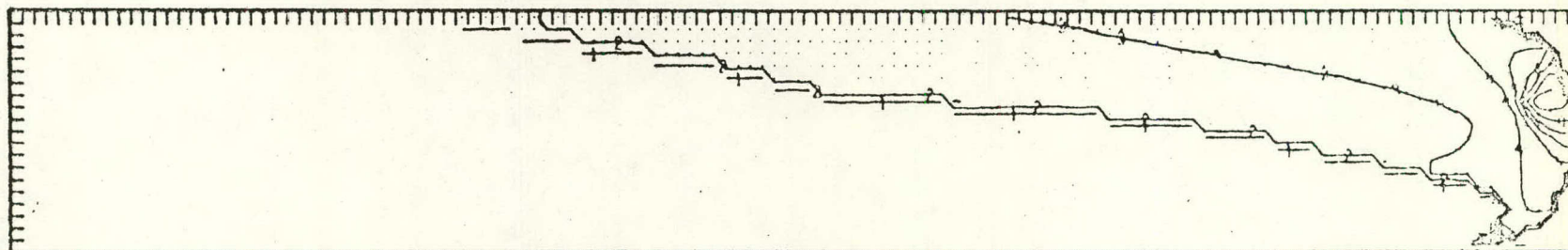
Page

CYCLE= 701 , TIME= 0.279972E-06

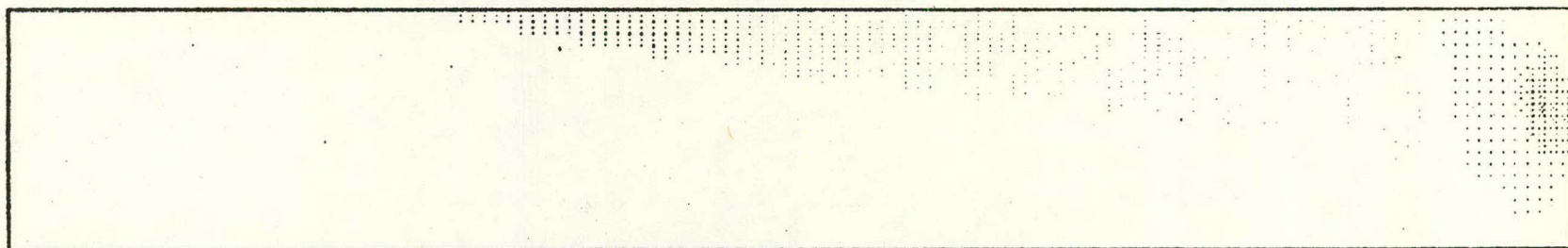
CONTOURS

1 =	0.000	2 =	0.185E+08	3 =	0.371E+08	4 =	0.556E+08	5 =	0.742E+08
6 =	0.927E+08	7 =	0.111E+09	8 =	0.130E+09	9 =	0.148E+09	10 =	0.167E+09

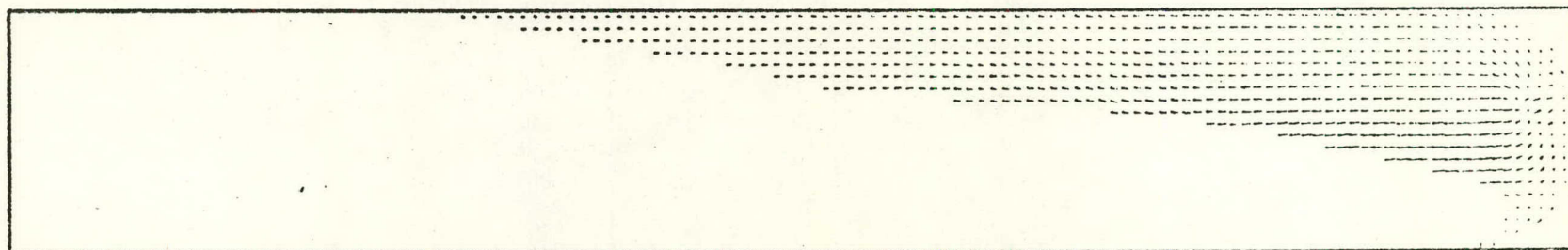
CYCLE= , TIME= DENSITY SHADING



Z
TEMPERATURE CONTOURS



DENSITY SHADING



VELOCITY
NUMERICAL SHOCK WAVES SIMULATIONS AT COLUMBIA

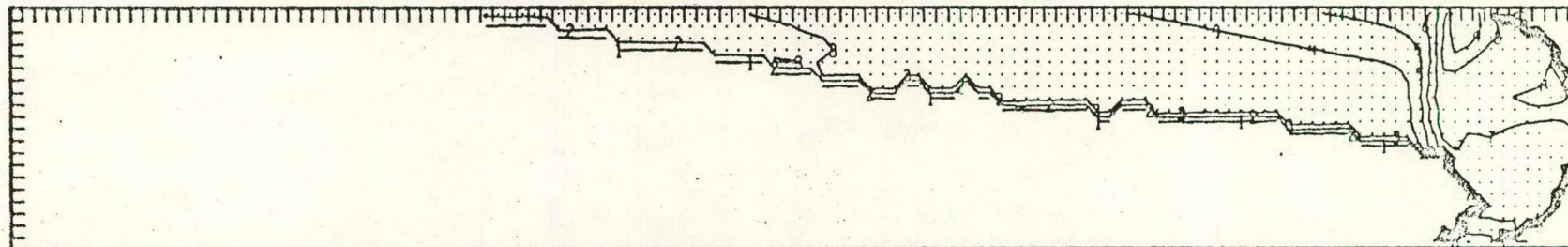
Fig. 5-7

CYCLE= 751 , TIME= 0.299969E-06

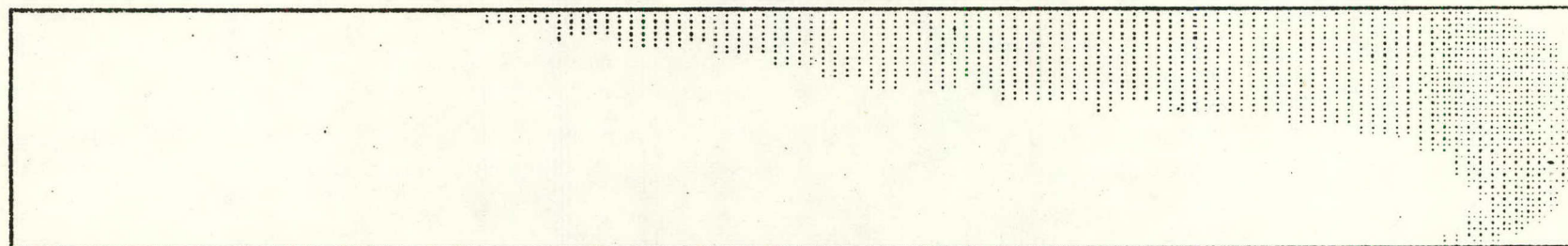
CONTOURS

1 =	0.000	2 =	0.131E+08	3 =	0.262E+08	4 =	0.392E+08	5 =	0.523E+08
6 =	0.654E+08	7 =	0.785E+08	8 =	0.916E+08	9 =	0.105E+09	10 =	0.118E+09

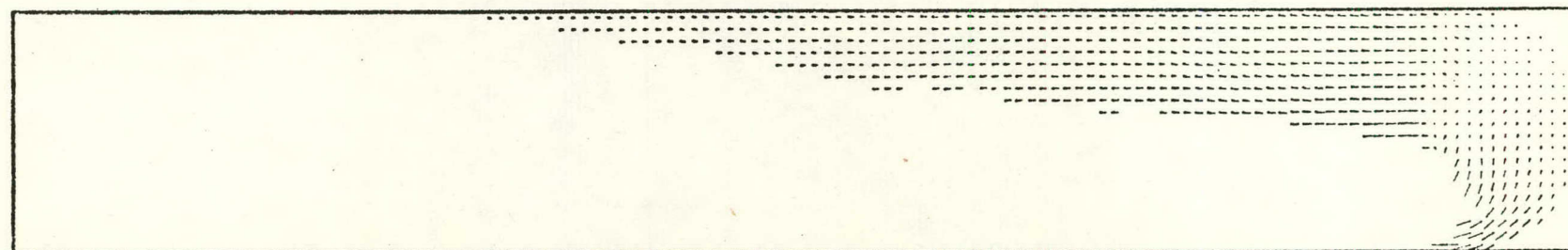
CYCLE=- , TIME=n DENSITY SHADING



Z
TEMPERATURE CONTOURS



DENSITY SHADING



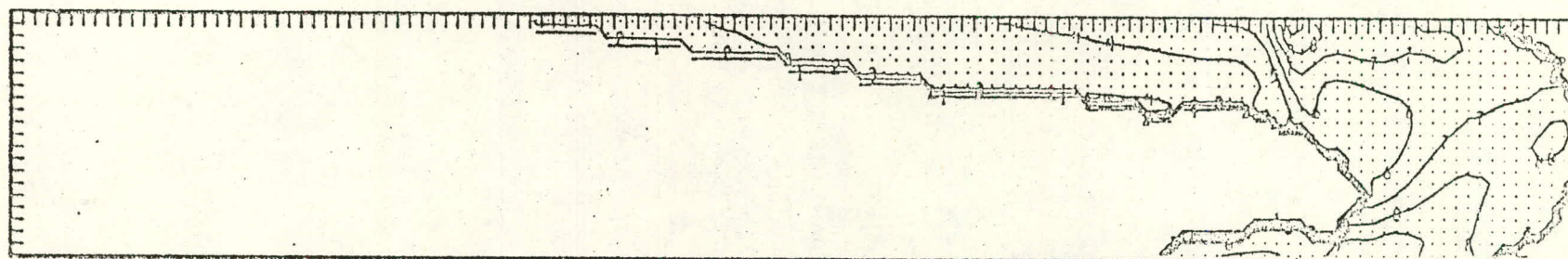
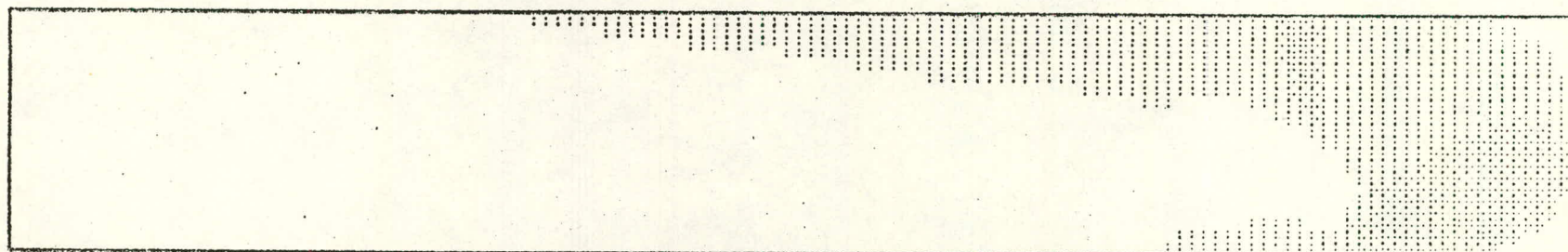
VELOCITY

CYCLE= 851 , TIME= 0.339964E-06

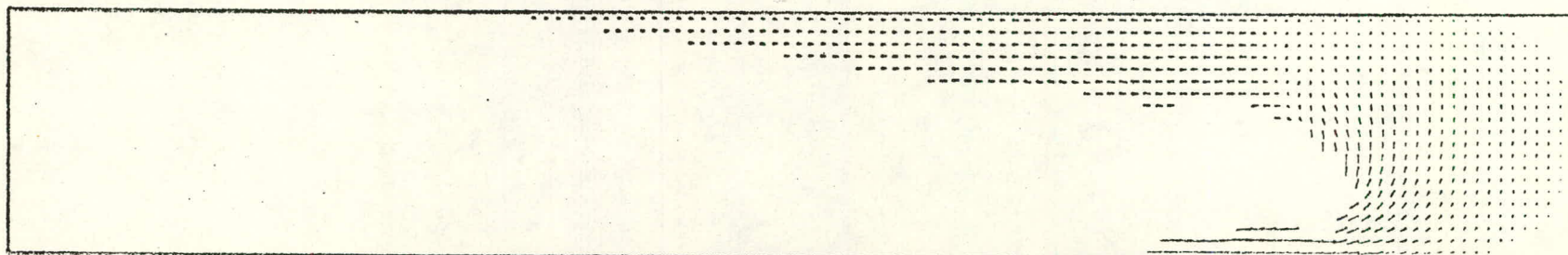
CONTOURS

1 =	0.000	2 =	0.109E+08	3 =	0.219E+08	4 =	0.328E+08	5 =	0.438E+08
6 =	0.547E+08	7 =	0.657E+08	8 =	0.766E+08	9 =	0.876E+08	10 =	0.985E+08

CYCLE=- , TIME=n DENSITY SHADING

Z
TEMPERATURE CONTOURS

DENSITY SHADING



VELOCITY

NUMERICAL SHOCK WAVES SIMULATIONS AT COLUMBIA

Fig. 5-7 The incident shock wave and its reflection from a semi-circular end wall

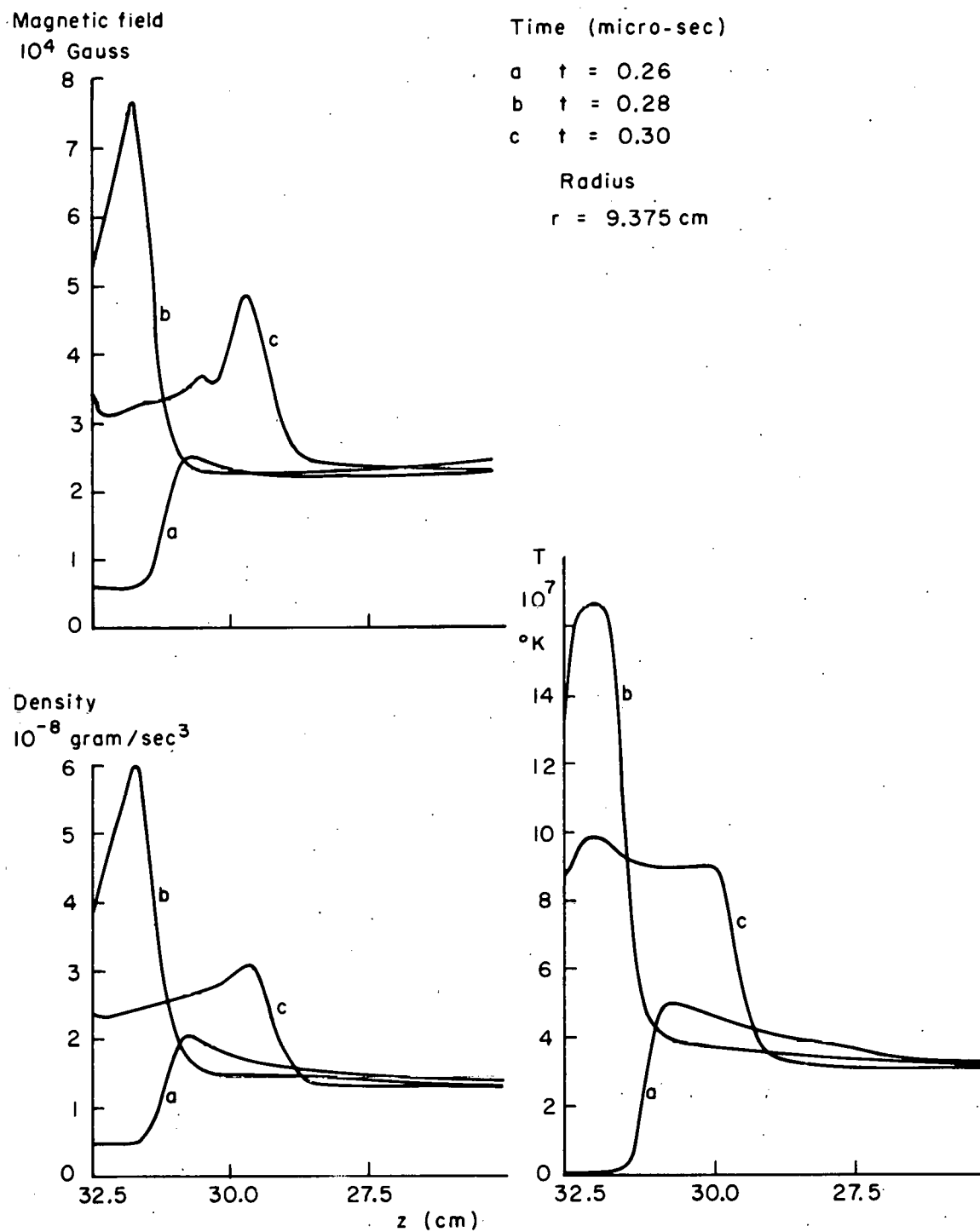


Fig. 5-8 The plots of the magnetic field, density and temperature with respect to the axial positions near the circular end wall.

VI. SHOCK-RAY APPROXIMATION IN TRANSVERSE MHD FLOWS

In order to select a wall profile and to determine the location for shock focusing and reflection from a curved end wall in the coaxial shock tube, Whitham's shock-ray approximation is reviewed and extended to plane transverse MHD flows. The system of equations in a two-dimensional transverse MHD flow is reducible to the case of a conventional two-dimensional unsteady fluid flow²³. Therefore, we can apply Whitham's formulations and Lau's shaping method²⁰ to design the end wall of the shock tube for shock convergence in transverse MHD flows.

According to the theory of sound, a wave front carrying a disturbance from a surface of arbitrary shape moves with the local speed of sound along the rays which are normal to the surface. Rays are orthogonal trajectories of successive positions of the wave front. They may be considered as carrying the discontinuities. Similarly, the particles immediately behind a shock wave move perpendicularly to the shock front, i.e., along the ray directions. In particular, a wall profile is a ray.

Whitham extended the theory of geometrical acoustics²⁴ to the shock-ray approximation^{17,18} by constructing shock-ray coordinates in the Cartesian space and by formulating the kinematic relations. He made the assumption that the shock Mach number and the area between ray tubes were functionally related. A smoothly converging wall will increase the shock Mach number

and generate a geometrical wave along the shock surface. For a properly shaped wall, focusing of the shock into a small region will occur.

We now extend this approximation to the transverse MHD flows. Let the shock position be defined by the curves $\alpha = \text{const}$ and the rays by $\beta = \text{const}$, with $\alpha(x, y, z) = C_0 t$. C_0 is the local magnetosonic speed in front of the shock wave. The $\alpha - \beta$ curves in two-dimensional axisymmetric ($r - z$) coordinates are shown in Fig. 6-1.

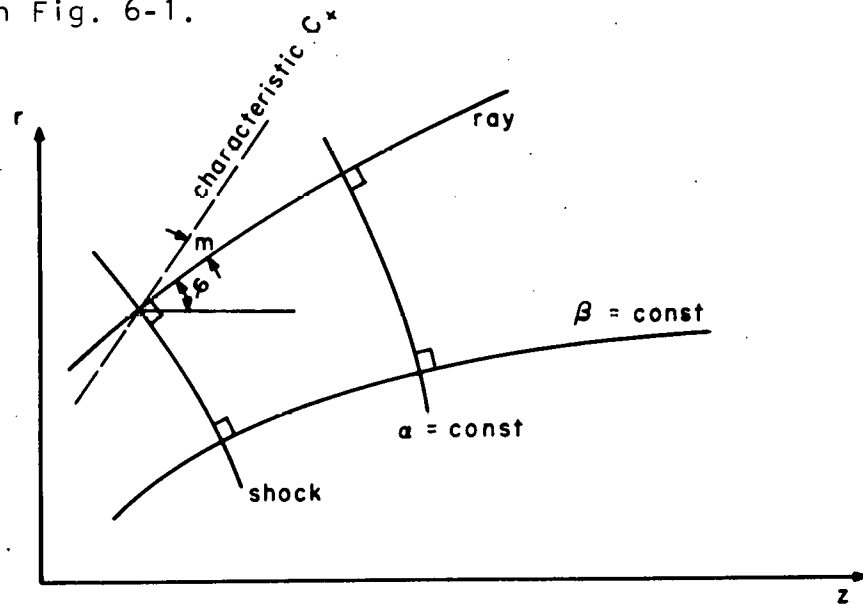


Fig. 6-1 Orthogonal shock-ray curves.

Using Whitham's formulations¹⁸, the following kinematic relations can be obtained for transverse MHD shock flows in axisymmetric geometry.

$$\alpha_z = \cos \varphi / M^*, \quad \alpha_r = \sin \varphi / M^*, \quad (6.1)$$

$$\frac{\partial}{\partial z} \left(\frac{\sin \varphi}{M^*} \right) - \frac{\partial}{\partial r} \left(\frac{\cos \varphi}{M^*} \right) = 0, \quad (6.2)$$

$$\frac{\partial}{\partial z} \left(\frac{r \cos \varphi}{A} \right) + \frac{\partial}{\partial r} \left(\frac{r \sin \varphi}{A} \right) = 0, \quad (6.3)$$

where the subscripts denote derivatives, M^* is the magnetosonic Mach number and A is the area between ray tubes. φ is the angle between a ray and the z -axis as shown in Fig. 6-1. The following derivation of the dynamical equation, $A = A(M^*)$, is similar to Whitham's work²⁵.

In a transverse MHD flow, the equations of motion in a non-uniform tube are

$$\rho_t + u \rho_z + \rho \left(u_z + \frac{u}{A} \frac{dA(z)}{dz} \right) = 0, \quad (6.4)$$

$$u_t + u u_z + \frac{1}{\rho} P_z^* = 0, \quad (6.5)$$

$$P_t^* + u P_z^* - C^2 (\rho_t + u \rho_z) = 0, \quad (6.6)$$

$$\frac{B}{\rho} = \text{const}, \quad (6.7)$$

$$P^* = P + \frac{B^2}{8\pi}, \quad (6.8)$$

$$C^2 = \frac{\partial P^*}{\partial \rho}, \quad (6.9)$$

where C is the magnetosonic speed, $C^2 = a^2 + b'^2$, with a being the sound speed and $b' \left(\frac{B}{\sqrt{4\pi\rho}} \right)$ being the Alfvén speed. The pressure term due to the area change does not appear in Eq. (6.5) because it can be neglected in a smoothly converging tube.

The characteristic form in $(z - t)$ plane is

$$\left(\frac{P_t^*}{u+c} + P_z^* \right) \pm \rho c \left(\frac{u_t}{u+c} + u_z \right) + \frac{\rho c^2 u}{(u+c)A} \frac{dA(z)}{dz} = 0, \quad (6.10)$$

$$\left(\frac{P_t^*}{u} + P_z^* \right) - c^2 \left(\frac{\rho_t}{u} + \rho_z \right) = 0. \quad (6.11)$$

The equations (6.10) and (6.11) can be rewritten as follows:

$$dP^* + \rho c du + \frac{\rho c^2 u}{u+c} \frac{dA}{A} = 0 \quad \text{on } C_+ : \frac{dz}{dt} = u + c, \quad (6.12)$$

$$dP^* - \rho c du + \frac{\rho c^2 u}{u-c} \frac{dA}{A} = 0 \quad \text{on } C_- : \frac{dz}{dt} = u - c, \quad (6.13)$$

$$dP^* - c^2 d\rho = 0 \quad \text{on } P : \frac{dz}{dt} = u. \quad (6.14)$$

The transverse MHD shock jump conditions with an infinite shock strength ($M^* \approx \infty$) are

$$P^* = P_o^* + \rho_o c_o^2 M^{*2} \frac{2}{\gamma+1}, \quad (6.15)$$

$$\rho = \rho_o \frac{\gamma+1}{\gamma-1}, \quad (6.16)$$

$$u = M^* \frac{2c_o}{\gamma+1}, \quad (6.17)$$

$$c = c_o M^* \sqrt{2\gamma(\gamma-1)} / (\gamma+1), \quad (6.18)$$

where the subscript (o) denotes a quantity in front of the shock wave.

Since the shock wave is along C_+ , we get the following dynamical relation by substituting Eqs. (6.15), (6.16), (6.17),

and (6.18) into Eq. (6.12).

$$\frac{dA}{A} \frac{M^*}{dM^*} = f(\gamma) = -\left\{ \frac{2}{\gamma} + 1 + \sqrt{\frac{1}{\gamma(\gamma-1)}} + \sqrt{\left(1 - \frac{1}{\gamma}\right)} \right\} \quad (6.19)$$

The result is similar to Whitham's formula for very strong shocks²⁵. Although the dynamical relation (6.19) is derived from the equations of plane transverse MHD flows, we still use it as an approximation to the coaxial shock tube by treating the propagation of the shock wave down each ray tube separately.

With the dynamical equation (6.19), we can rewrite hyperbolic Eqs. (6.2) and (6.3) in the characteristic forms.

$$d\varphi + \frac{1}{\tan(\pm m)} \frac{dM^*}{M^*} + \frac{\tan \varphi \tan(\pm m)}{\tan \varphi + \tan(\pm m)} \frac{dr}{r} = 0, \\ \text{along } \frac{dr}{dz} = \tan(\varphi \pm m) \quad (6.20)$$

$$\tan m = \sqrt{\frac{-A}{M^*} \frac{dM^*}{dA}} = \frac{1}{2.15} \quad (\text{with } \gamma = 5/3), \quad (6.21)$$

where m is the angle between a ray and a characteristic line C_+ . These equations are of the same form as the gas dynamical equations, with their characteristics, Riemann invariants, shocks, etc. The propagation speed of the geometrical wave,

$$\frac{d\beta}{d\alpha} = \pm \left(- \frac{dM^{*2}}{dA^2} \right)^{\frac{1}{2}} r$$

in an axisymmetric case, is an increasing function of magnetosonic Mach number M^* . The geometrical waves with increased magnetosonic Mach number will eventually break and form a shock, in this case, a "shock-shock" called by Whitham. Such a "shock-shock" which propagates on the shock surface is a discontinuity

in magnetosonic Mach number (M^*) and shock slope (φ).

With Eqs. (6.20) and (6.21), we follow Lau's wall shaping method using the theory of characteristics to find the end wall profile for shock focusing in the coaxial shock tube. The incident shock wave is assumed to be a plane shape. Because the Eqs. (6.20) and (6.21) are completely similar to those used by Lau, the wall shape for shock focusing in our case will be exactly the same as that found by him. It is approximately a quarter of the ellipse with the ratio of the major axis to the minor axis equal to 2.25 [= (2.15 + 0.1)]. This is the optimal ratio, as the shock focuses to a point in the $r-z$ plane, or actually to a ring. The difference between two slightly different wall shapes will eventually become negligible when a shock wave is converged into a same location of the bottom of the end wall. The elliptic shape makes the numerical computation particularly easy. In Fig. 6-2, the line BC is the first positive characteristic due to the wall curvature. The dimension of the triangle ABC is given by Eq. (6.21), i.e., $\tan(m) = AC/AB = 1/2.15$. The numerical value (0.1) is chosen to make the ellipse following Lau's wall shape fairly. It can be changed to other small values. The focused location is arranged to be near the inner conductor according to Eq. (6.20). The end wall (CE) in the shock tube is shown in Fig. 6-2.

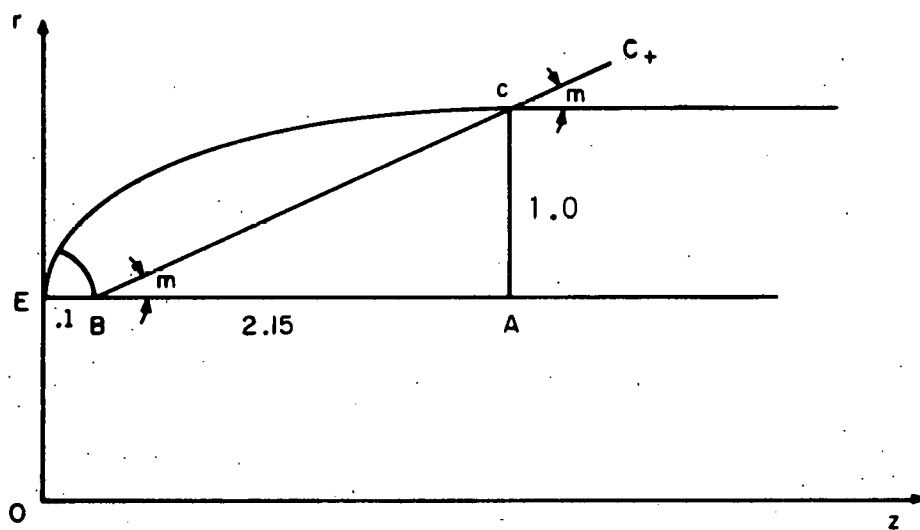


Fig. 6-2 The Selected Elliptic End Wall.

VII. SHOCK FOCUSING AND REFLECTION

7.1 Quarter-elliptic End Wall

The elliptic wall designed in the previous chapter will now be checked out by computation. The shock wave of Section 4.2 is used again so that we can study its focusing and reflection. The solution is presented in the plottings such as Figs. 7-1, 7-2, and 7-3.

Fig. 7-1 shows the sequence of the shock flow in the tube with the elliptic end wall located at the right end. These pictures are helpful to us in understanding the details of the flow, the process of shock focusing and reflection, and plasma heating. At cycle = 501, the incident shock wave runs down the tube with a speed of 125 cm/ μ sec. At cycle = 601, the shock wave moving through the converging channel is diffracted and strengthened on the side wall, and the plasma behind the wall shock is further compressed. The fluid particles behind the wall shock travel along the wall surface.

At cycle = 651, the wall shock and the incoming main shock wave make an angle of about 110° . The geometrical wave on the shock surface now breaks into a geometric shock (Mach reflection). Whitham called such a geometric discontinuity on the shock surface a "shock-shock". The magnetosonic Mach number and the angle (ϕ) of the ray on the wall shock differ with those on the main shock wave respectively. No matter what the shock shape is at the bottom of the end wall, the shock will inevitably

reach a maximum in strength, i.e., shock focusing. Then, the high kinetic energy of a fast moving particle suddenly will change into plasma thermal energy in the focused region.

At cycle = 701, the reflected shock wave has indeed focused and produces a small region of hot plasma with a temperature of 32 Kev, or some 8 times the incident post-shock temperature. The maximum density in the focused region is 8.12×10^{-8} gram/cm³, or 16 times the initial density in the shock tube. The duration of the hot plasma is quite short, being equal to the transit time of the reflected shock wave across the major axis of the ellipse.

At cycle = 751, the shock again reflects from the focused region, and behind the reflected shock is a rarefaction wave. The velocities of the incoming particles and the reflected particles between the reflected shock wave are in opposite directions. Then, the reflected shock wave propagates downstream near the outer wall, and moves into the vacuum near the inner conductor. The plasma expands very rapidly along the surface of the inner conductor.

Fig. 7-2 shows the plots of the magnetic field, density and temperature with respect to axial position in the focused region with a radius of 6.625 cm at different times. The bottom of the elliptic end wall is at 33.75 cm. At 0.28 microseconds, shock focusing and reflection make all the flow variables peak at the focused position.

Fig. 7-3 shows the traces of the magnetic field, density and temperature with respect to time at a fixed station near

the focused region. The station is shown on the first page of Fig. 7-1 as the point (G). A shock compression raises the total pressure (P^*) due to the arrival of the incoming shock wave. When the shock wave moves away from the station, the total pressure drops since the magnetic field and density of the plasma just behind the incident shock wave are higher than those of the downstream plasma; this is due to the effects of the side wall and drive current. A further shock compression occurs when the reflected shock wave returns to the station. Finally, the total pressure goes down quickly as the reflected shock wave moves away from the station. Milton¹⁹ pointed out that ordinary gas shock convergence and reflection by a curved wall in plane flows follows purely the cylindrical implosion cycle^{26,27}, and that the flow conditions in the area contraction of the shock wave are not uniform. In our case, the magnetic field in front of the incident shock wave varies radially. However, some features of cylindrical implosion appear in this case.

To keep the temperature of the plasma high, it is necessary to keep the plasma from expanding. An externally applied magnetic field after shock reflection at a suitable downstream position will slow the plasma expansion resulting in extended heating. The plasma under the applied magnetic fields will not expand when the applied magnetic energy density is much greater than the total gas energy density.

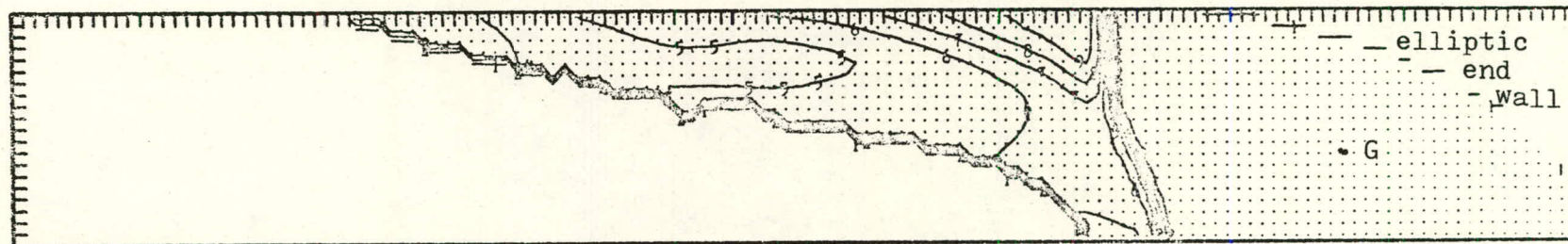
Fig. 7-1

CYCLE= 501 , TIME= 0.199982E-06

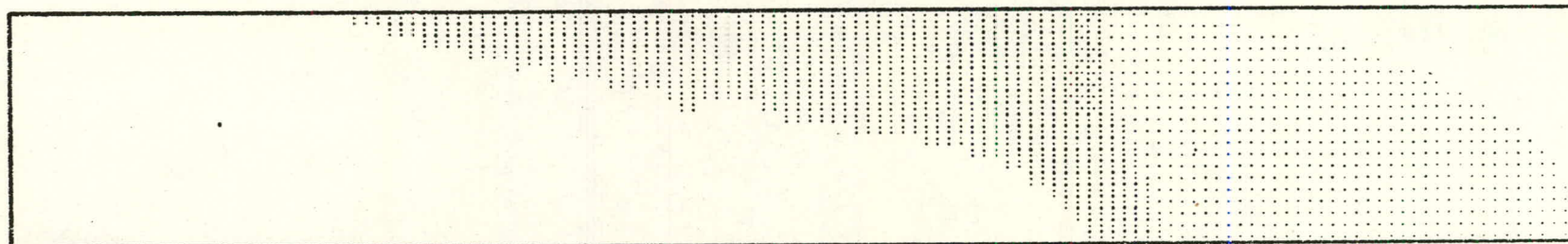
CONTOURS

1 =	0.000	2 =	0.739E+07	3 =	0.148E+08	4 =	0.222E+08	5 =	0.296E+08
6 =	0.369E+08	7 =	0.443E+08	8 =	0.517E+08	9 =	0.591E+08	10 =	0.665E+08

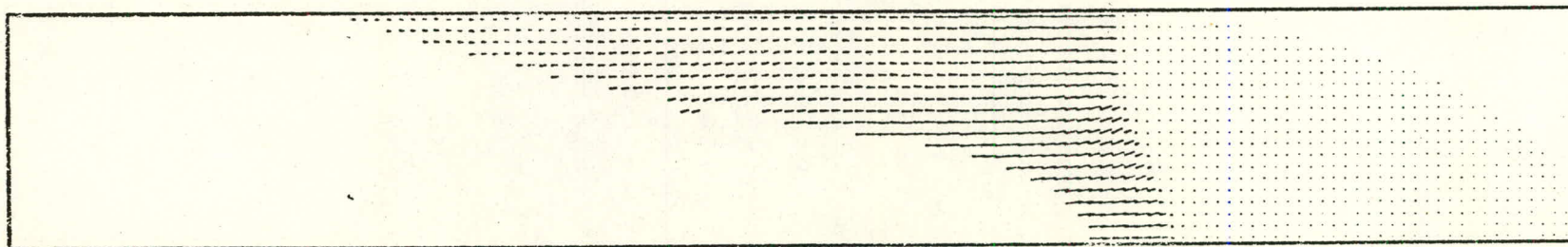
CYCLE=Y0, TIME=X0 DENSITY SHADING AVELO



Z
TEMPERATURE CONTOURS



DENSITY SHADING



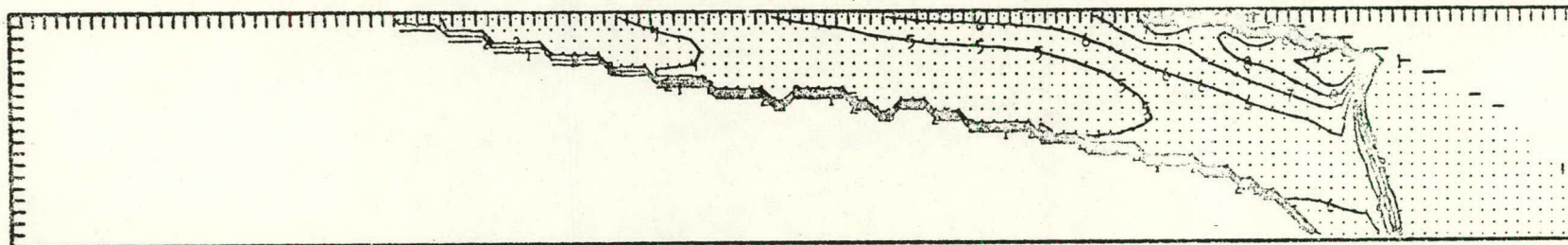
VELOCITY

NUMERICAL SHOCK WAVES SIMULATIONS AT COLUMBIA

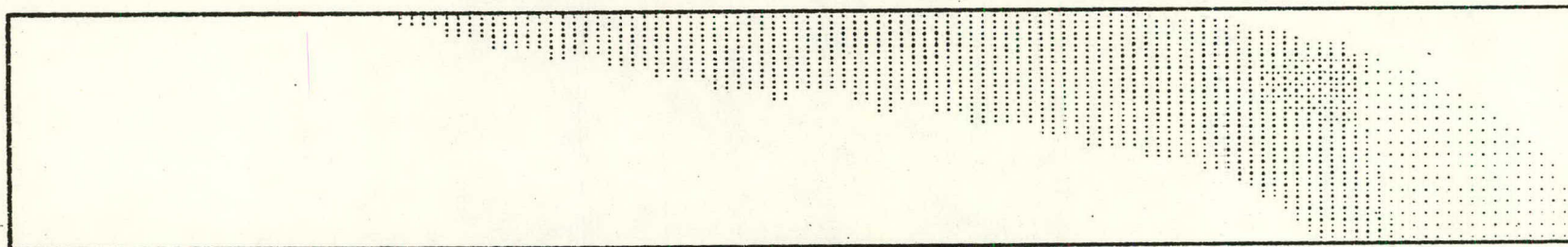
Fig. 7-1

CYCLE= 601 , TIME= 0.239771 06
 CONTOURS
 1 = 0.000 2 = 0.867E+07 3 = 0.173E+08 4 = 0.260E+08 5 = 0.347E+08
 6 = 0.433E+08 7 = 0.520E+08 8 = 0.607E+08 9 = 0.694E+08 10 = 0.780E+08

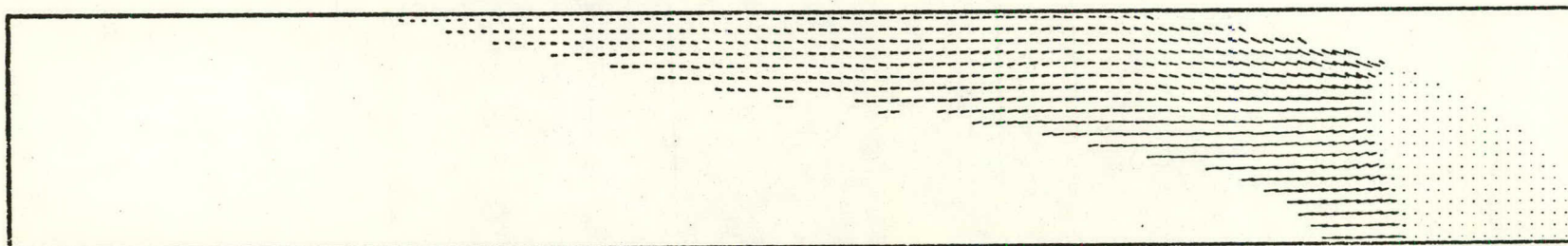
CYCLE=Y0, TIME=X0 DENSITY SHADING AVELD



Z
TEMPERATURE CONTOURS



DENSITY SHADING



VELOCITY
NUMERICAL SHOCK WAVES SIMULATIONS AT COLUMBIA

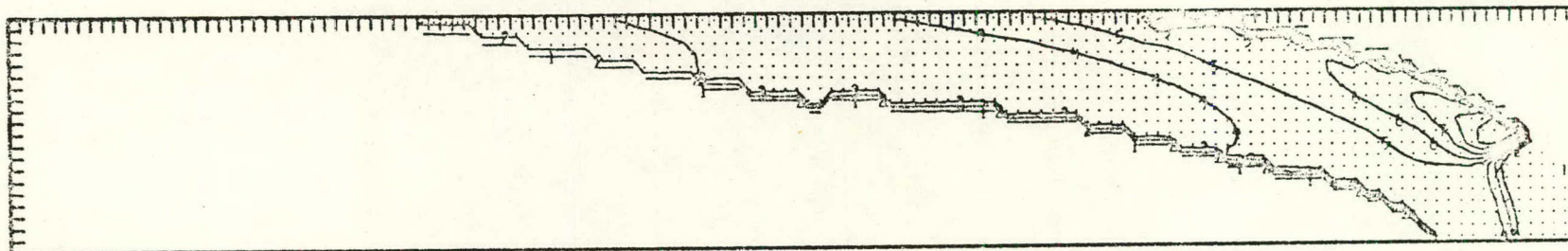
Fig. 7-1

CYCLE= 651 , TIME= 0.259974E+06

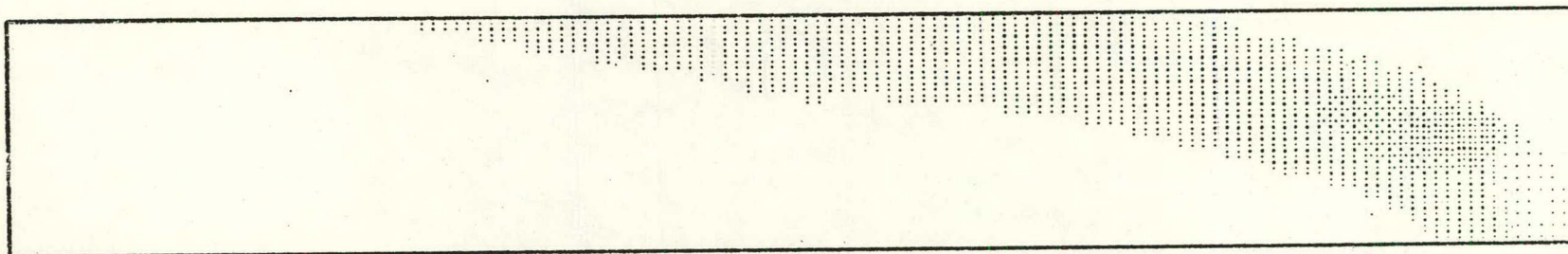
CONTOURS

1 =	0.000	2 =	0.124E+08	3 =	0.248E+08	4 =	0.371E+08	5 =	0.495E+08
6 =	0.619E+08	7 =	0.743E+08	8 =	0.867E+08	9 =	0.990E+08	10 =	0.111E+09

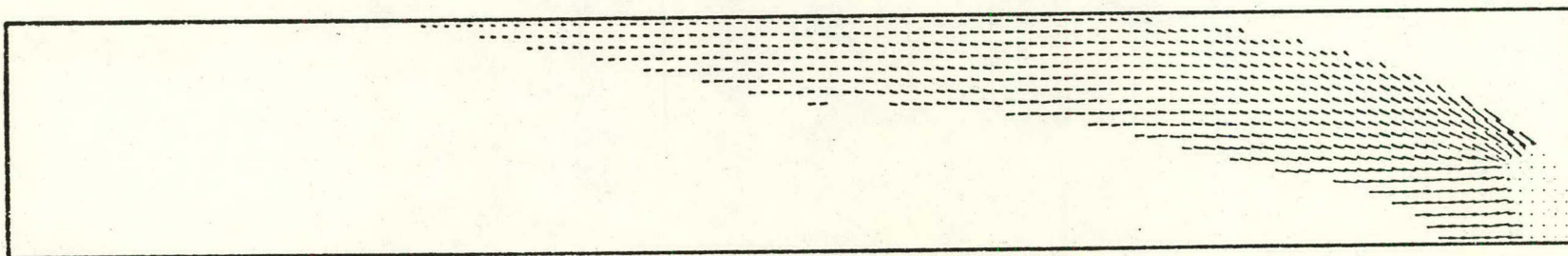
CYCLE=Y0, TIME=X0 DENSITY SHADING VELO



TEMPERATURE CONTOURS



DENSITY SHADING



VELOCITY

NUMERICAL SHOCK WAVES SIMULATIONS AT COLUMBIA

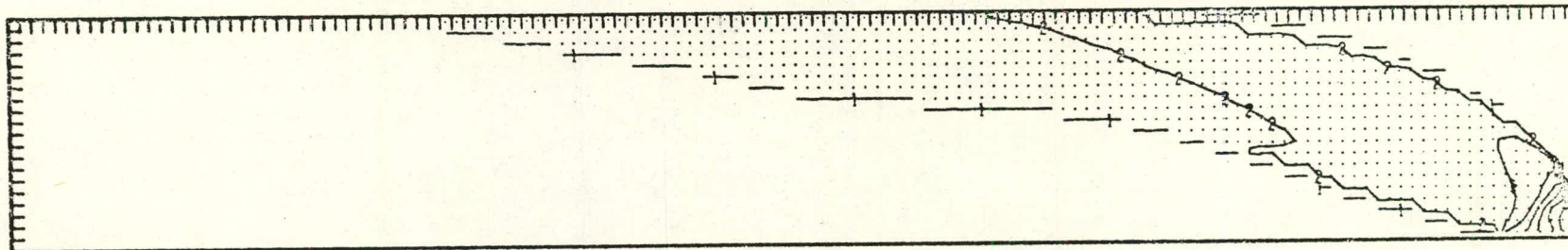
Fig. 7-1

CYCLE= 701 , TIME= 0.279972E-06

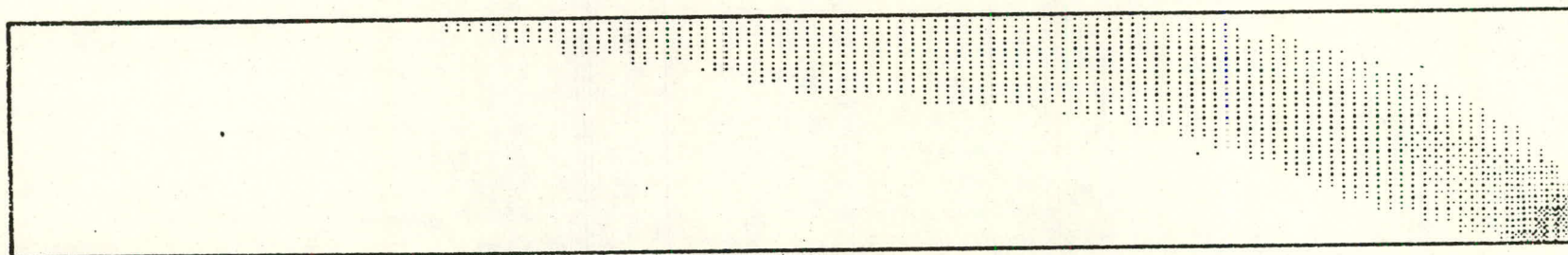
CONTOURS

1 =	0.000	2 =	0.381E+08	3 =	0.761E+08	4 =	0.114E+09	5 =	0.152E+09
6 =	0.190E+09	7 =	0.228E+09	8 =	0.266E+09	9 =	0.304E+09	10 =	0.342E+09

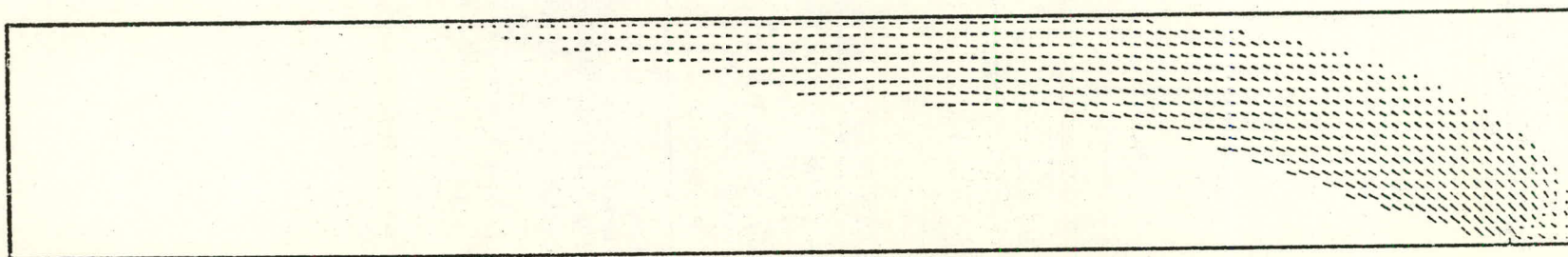
CYCLE=Y0, TIME=X0 DENSITY SHADING VELO



Z
TEMPERATURE CONTOURS



DENSITY SHADING



VELOCITY

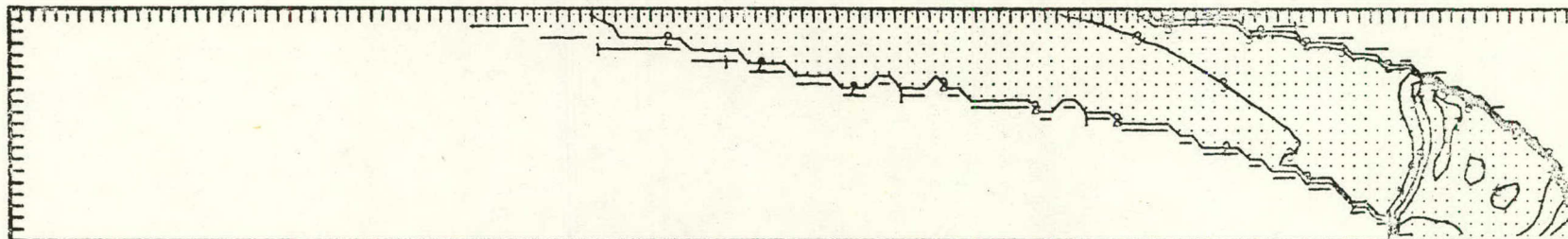
NUMERICAL SHOCK WAVES SIMULATIONS AT COLUMBIA

Fig. 7-1

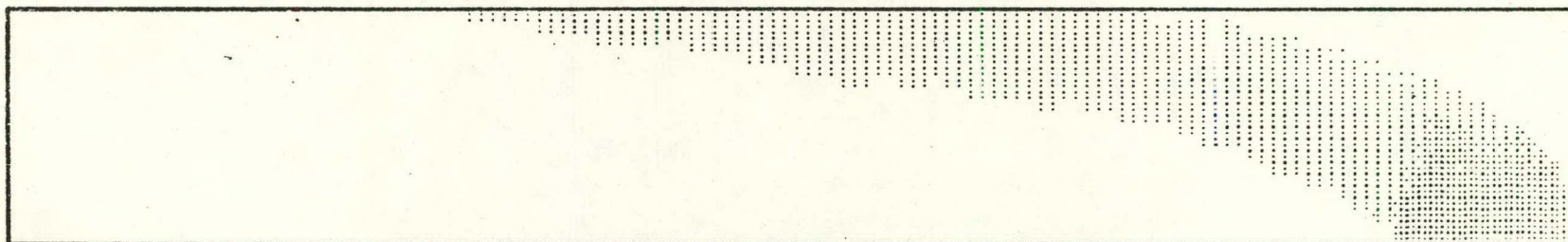
CYCLE= 751 , TIME= 0.2995651 06

CONTOURS
 1 = 0.000 2 = 0.203E+08 3 = 0.407E+08 4 = 0.610E+08 5 = 0.813E+08
 6 = 0.102E+09 7 = 0.122E+09 8 = 0.142E+09 9 = 0.163E+09 10 = 0.183E+09

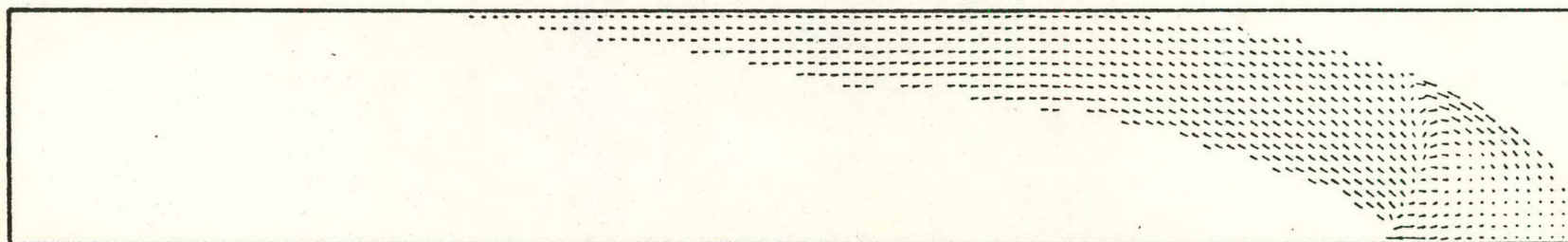
CYCLE=Y0, TIME=X0 DENSITY SHADING AVELO



Z
TEMPERATURE CONTOURS



DENSITY SHADING



VELOCITY

NUMERICAL SHOCK WAVES SIMULATIONS AT COLUMBIA

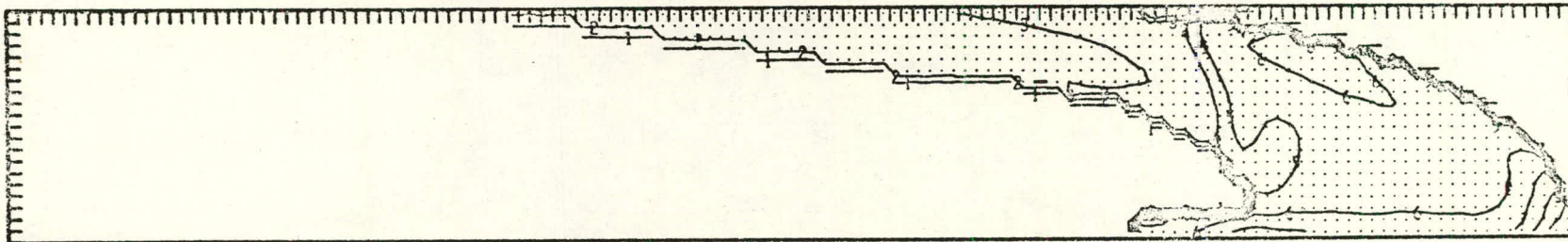
Fig. 7-1

CYCLE= 851 , TIME= 0.339964E 06

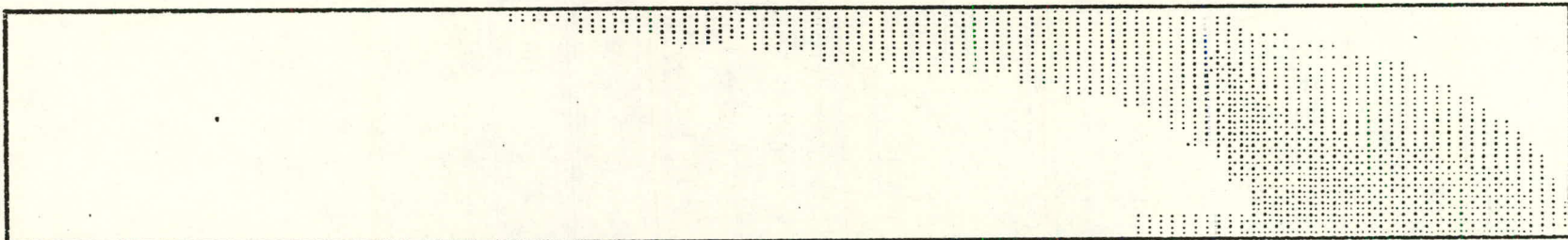
CONTOURS

1 =	0.000	2 =	0.166E+08	3 =	0.332E+08	4 =	0.498E+08	5 =	0.665E+08
6 =	0.831E+08	7 =	0.997E+08	8 =	0.116E+09	9 =	0.133E+09	10 =	0.150E+09

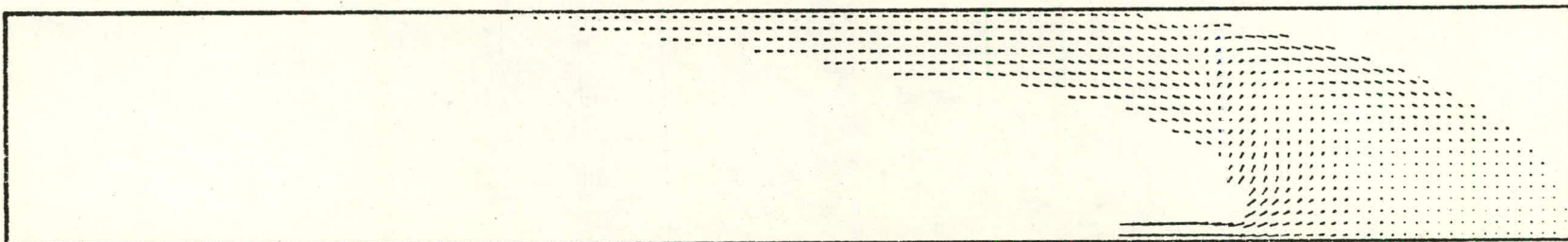
CYCLE=90, TIME=X, DENSITY SHADING



Z
TEMPERATURE CONTOURS



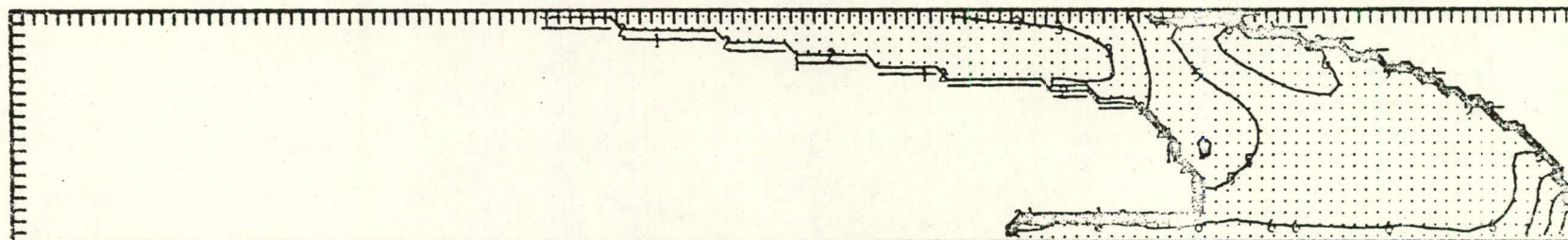
DENSITY SHADING



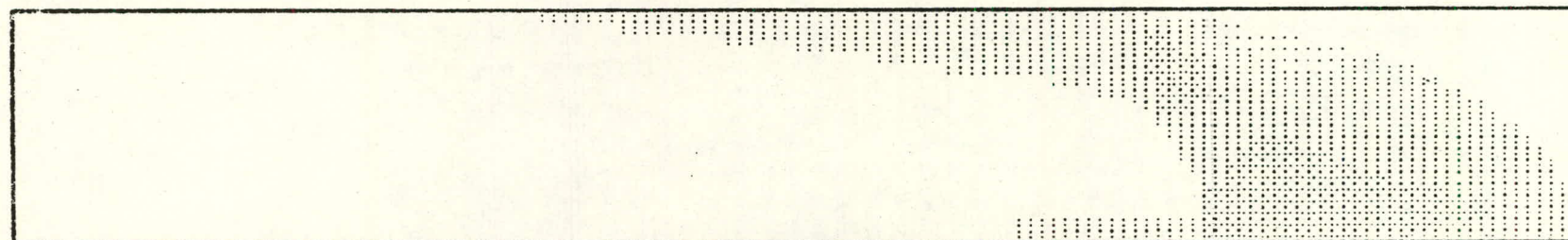
VELOCITY

CYCLE= 901 TIME= 0.359962E-06
 CONTOURS
 1 = 0.000 2 = 0.158E+08 3 = 0.315E+08 4 = 0.473E+08 5 = 0.631E+08
 6 = 0.788E+08 7 = 0.946E+08 8 = 0.110E+09 9 = 0.126E+09 10 = 0.142E+09

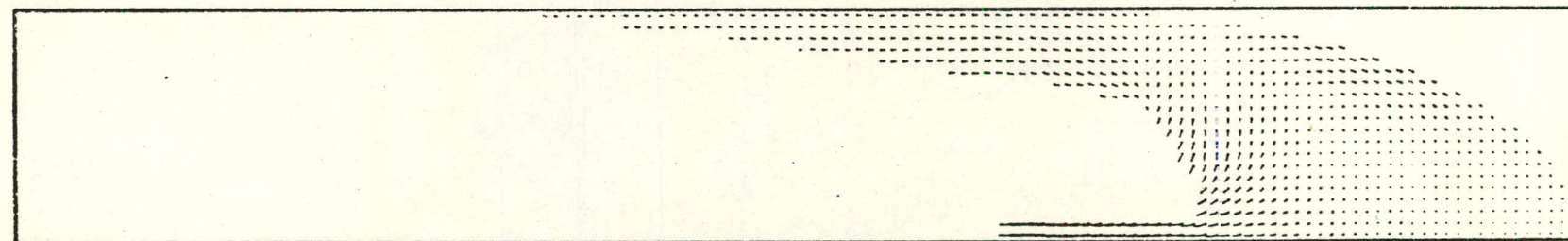
CYCLE=90, TIME=8, DENSITY SHADING/VELO



Z
 TEMPERATURE CONTOURS



DENSITY SHADING

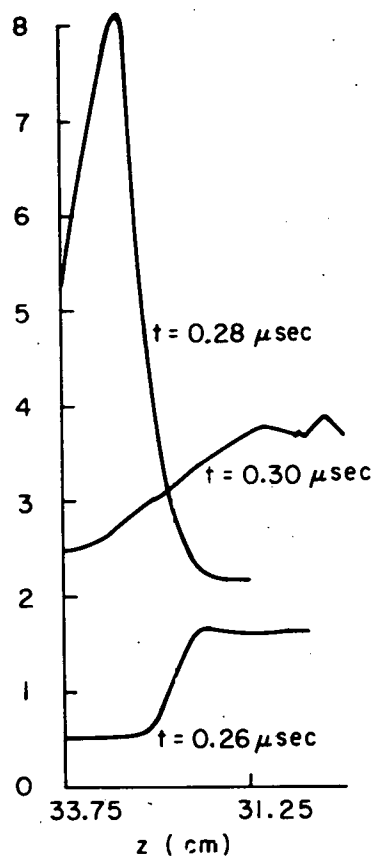


VELOCITY

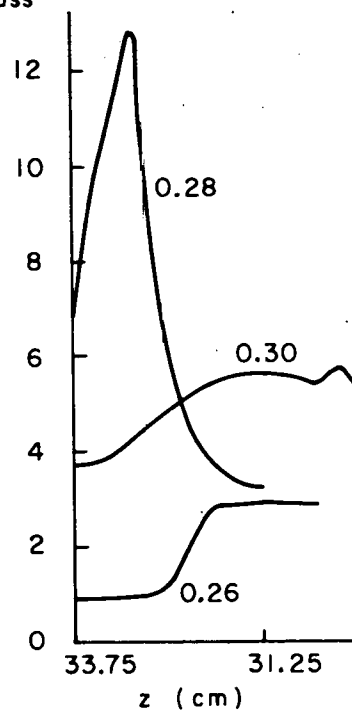
NUMERICAL SHOCK WAVES SIMULATIONS AT COLUMBIA

Fig. 7-1 Shock focusing and reflection from the selected quarter-elliptic end wall

Density
 10^{-8}
 Gram/cm³



Magnetic field
 10^4
 Gauss



T
 10^7
 $^{\circ}\text{K}$

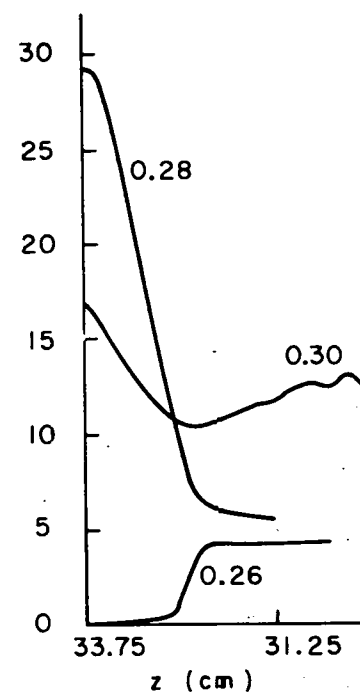


Fig. 7-2 Plots of the magnetic field, the density and the temperature with respect to the axial positions near the focused region with the radius $r = 6.625$ cm at different sequences. The incident shock speed is 125 cm/ μsec .

The variables are measured at the station (G) shown on the first page in Fig. 7-1

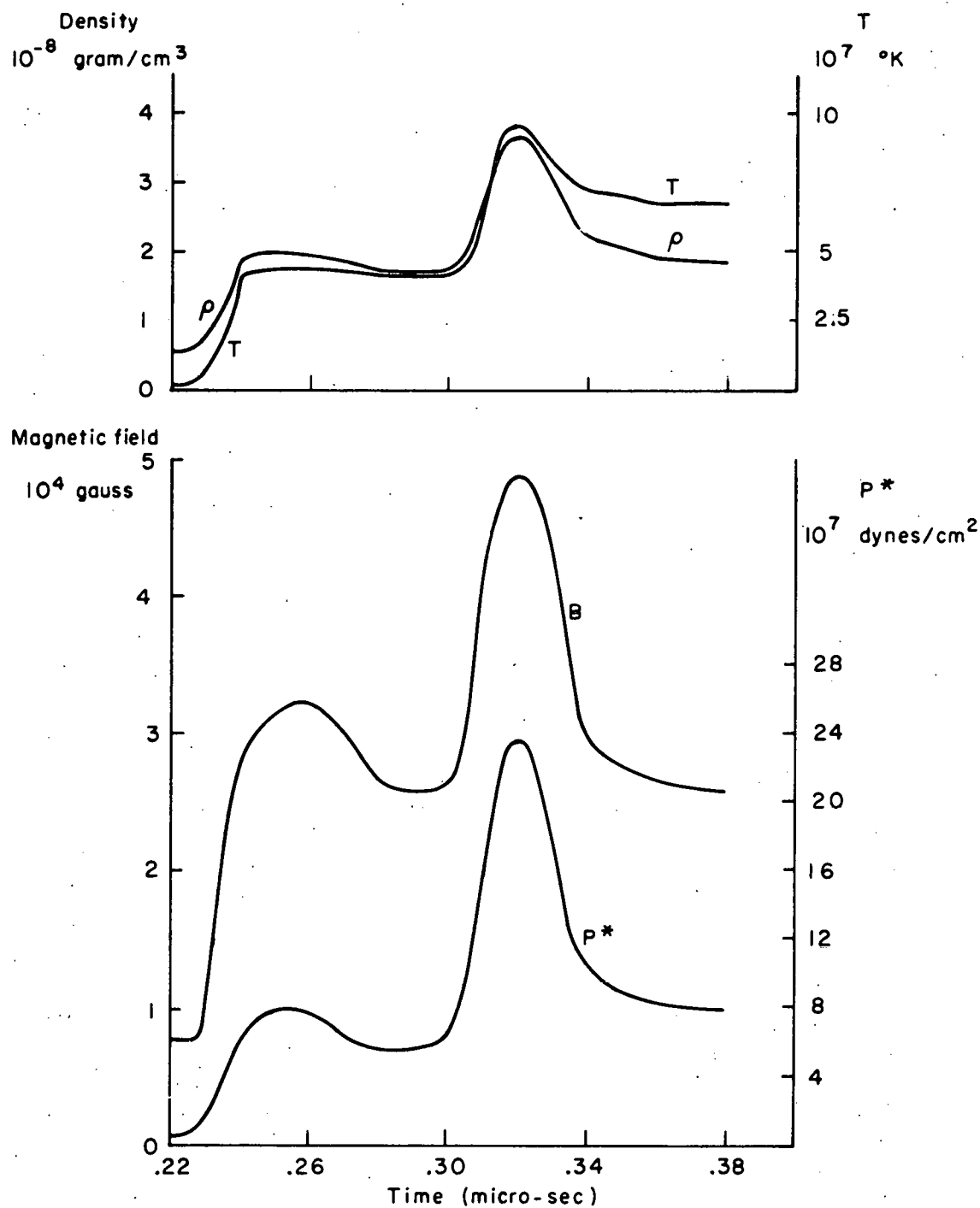


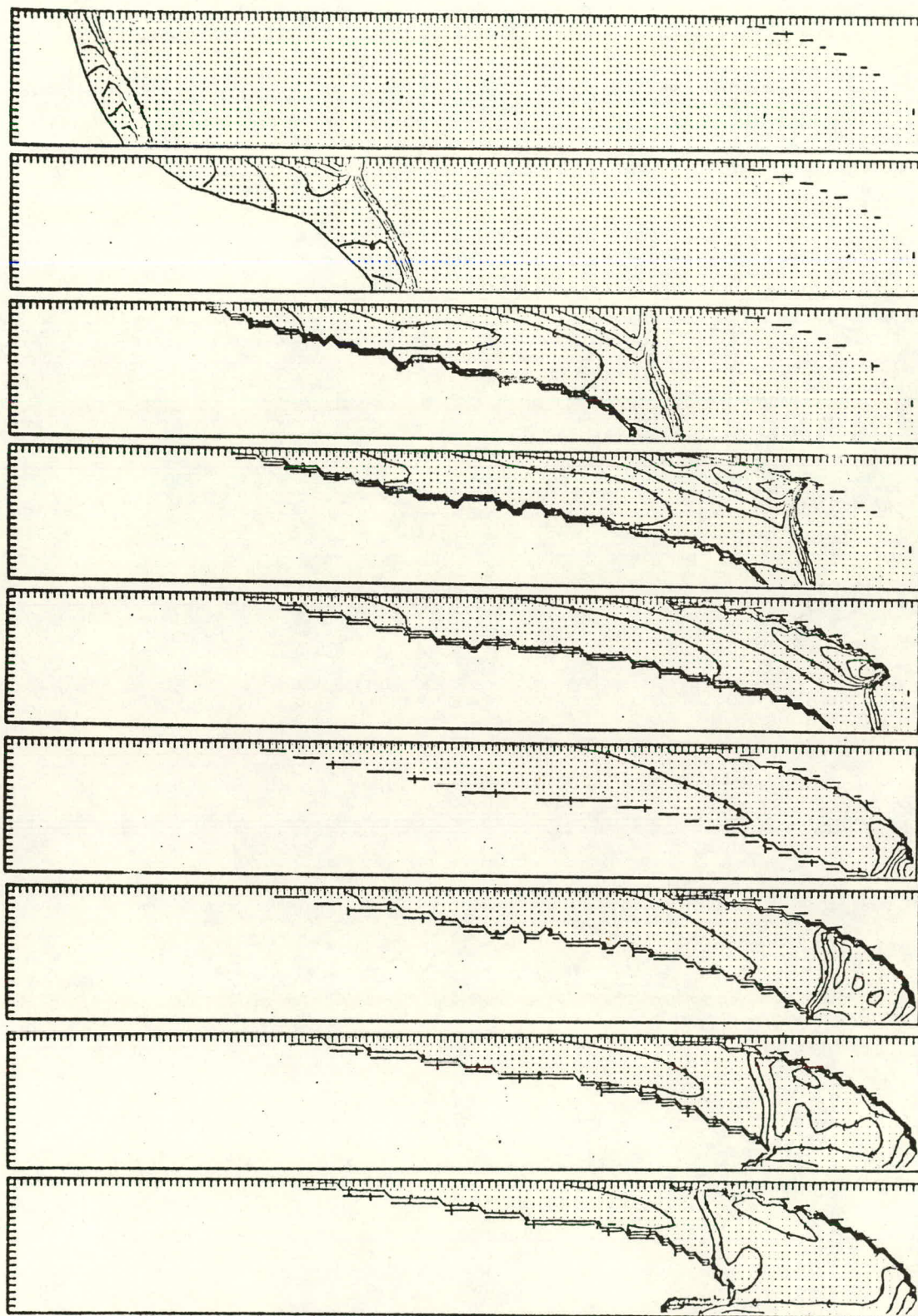
Fig. 7-3 The traces of the temperature, density, magnetic field and total pressure (P^*) at a fixed position.

7.2 Summary

As we have seen, the end wall profile in a coaxial shock tube obviously plays an important rule in plasma heating. The convergence of the incident shock wave, the motions of fluid particles, and shock flows depend strongly on the wall profile. Generally speaking, the flat wall does not converge a plane shock; the semi-circular wall converges an incident shock to some degree; and the selected quarter-elliptic end wall makes an incident shock wave fully convergent. The maximum temperature and density of plasma achieved during shock focusing and reflection from the quarter-elliptic end wall, are much higher than those due to shock reflection from the semi-circular and flat end walls respectively. The maximum temperature and the corresponding density and magnetic field are listed below for the shock wave which has been reflected from three different wall shapes. They are measured at the focused locations of the quarter-elliptic ($r = 6.625$ cm), semi-circular ($r = 9.375$ cm), and flat ($r = 11$ cm) end walls at the moment of maximum plasma heating.

wall shape	flow variables measured in the focused region at the moment of maximum plasma heating		
	maximum temp. ($^{\circ}\text{K}$)	corresponding density (gram/cm ³)	corresponding magnetic field (KG)
elliptic	29.0×10^7	5.30×10^{-8}	78.7
circular	16.7×10^7	5.36×10^{-8}	68
flat	14.5×10^7	5.1×10^{-8}	53

For the convenience to see the motions of the incident and reflected shock waves, Fig. 7-4 shows the whole sequence of the temperature contours for shock focusing and reflection from the selected quarter-elliptic end wall in the coaxial shock tube.



2
TEMPERATURE CONTOURS

Fig. 7-4 Shock focusing and reflection from the quarter-elliptic end wall in the coaxial shock tube

VIII. PLANE FLOWS

We did some preliminary work on plane flows to check the accuracy of the FLIC method. The code was applied to shock reflection in plane flows, and the results were compared with those obtained by the two-step Lax-Wendroff scheme²², which was given in Appendix A.

The treatment of boundary conditions using irregular cells on the curved wall described in Chapter III is designed for the FLIC method only. For the Lax-Wendroff scheme, we use a simpler approximation, using only regular cells. This same approximation is also used for the FLIC method; and in plane flows, it does not give much difference in results from the more accurate method of Chapter III. More specifically, on a curved wall boundary, all surface cells and fictitious boundary cells are treated in the same way as the interior cells, in spite of the fact that they are split by the wall boundary. They are kept in the same shape as the interior regular cells. The values of the flow variables in a fictitious cell, which has at least one half of its area outside the wall boundary, are calculated by reflection with respect to the tangent of the wall profile. In this way, there are no flows of mass, momentum and energy across the wall boundary as the normal component of the fluid velocity on a fixed wall vanishes. This method makes it easy to handle boundary cells on a curved wall in plane flows. We call it a "regular shape" approximation.

The problem, with which we compare the two methods, is the reflection of a plane shock wave from a semi-circular wall in a transverse MHD flow. In Table 8-1, we give the required parameters for both schemes. The downstream conditions are obtained from the upstream conditions in front of the shock wave. The initial shock position of the incident shock wave is given to be at 3.2 cm in front of the center of the semi-circular wall as shown in Fig. 8-1. The radius of the wall curvature is 2.5 cm. Then, the computations of the shock flow can be made by using a continuous inflow boundary at the left end. Actually, reflection is applied to the center line because of the symmetry, so that only one half of the flow region is in the computations.

To select the right amount of numerical dissipation, for this set of tests as well as for Chapters IV, V and VII, we test the FLIC method on shock reflection from a flat wall using various values for the coefficient of the numerical dissipation, until it is large enough to dampen significantly the numerical oscillations. The coefficient (b) in the viscous pressure term, say $q_{i+\frac{1}{2},j}^n$, is chosen to be 0.5. The detailed relation between b and $q_{i+\frac{1}{2},j}^n$ is given in Chapter II. Similar tests are carried out for the two-step Lax-Wendroff scheme, for which we need a coefficient (k=1.2) in a viscosity term which is shown in Appendix A. These values of the coefficients (b and k) are used in the computations of shock reflection from the semi-circular wall.

Fig. 8-2 shows the results obtained by the FLIC method for a shock wave reflecting from a flat wall using the initial conditions given in Table 8-1. The reflected shock wave is shown on the lower curves. The density profile is on the left, the temperature profile is on the right. The solid curves are a large cell result ($\delta x = \delta y = 0.2$ cm), the dotted curves are a fine cell result ($\delta x = \delta y = 0.1$ cm). The shock thickness in the fine cell result is narrower than that of the large cell result since the numerical diffusion is proportional to the cell size. An ordinary numerical scheme will spread out a step function into several cells. But the amount of inherent numerical dissipation in the large cell result is greater than that of the fine cell result. These are the common nature of the numerical schemes. However, these computed results agree well with the analytical solution.

The sequence of temperature contours obtained by the FLIC method for the plane shock wave reflecting from the semi-circular wall is shown in Fig. 8-3. The comparison of the computed results between the two methods in this problem is given in Fig. 8-4 and Table 8-2. Fig. 8-4 shows the temperature and density contours obtained from three different calculations at the same instant of time: a large grid FLIC method ($\delta x = \delta y = 0.2$ cm), a fine grid FLIC method ($\delta x = \delta y = 0.1$ cm), and a large grid two-step Lax-Wendroff scheme ($\delta x = \delta y = 0.2$ cm). The agreement between the three results is good. In addition, we also tabulate the density and temperature measured at the position B (shown in Fig. 8-1) with

respect to time under the two methods in Table 8-2. The average errors between them based on the FLIC result is about 2.34% in density and about 1.36% in temperature. This gives us some confidence on the basic soundness of the FLIC method for this class of problems.

Table 8-1

The required parameters and initial conditions for a plane shock wave reflecting from a semi-circular wall in plane flows

Gas	hydrogen
Mach number	1000
Shock speed	144 cm/ μ sec
Cell width ($\delta x = \delta y$)	0.2 cm

	<u>Upstream</u>	<u>Downstream</u>
Density	$.5344 \times 10^{-8}$ gram/cm ³	1.86×10^{-8} gram/cm ³
Pressure	66.65 dynes/cm ²	5.74×10^7 dynes/cm ²
Temperature	300 °K	7.43×10^7 °K
Transverse field	7000 Gauss	24364.7 Gauss
Axial component of fluid velocity	0 cm/ μ sec	102.63 cm/ μ sec

Table 8-2

The density and temperature measured at the position B (shown in Fig. 8-1) with respect to time under both schemes

Time μsec	FLIC ($\delta x = \delta y = 0.2 \text{ cm}$)		Lax-Wendroff ($\delta x = \delta y = 0.2 \text{ cm}$)	
	Density $10^{-8} \text{ gram/cm}^3$	Temp. $10^8 \text{ }^\circ\text{K}$	Density $10^{-8} \text{ gram/cm}^3$	Temp. $10^8 \text{ }^\circ\text{K}$
0	.53	0	.53	0
.0155	1.93	.78	2.03	.81
.0255	1.83	.70	1.87	.72
.0355	5.04	2.33	4.67	2.44
.0455	4.18	2.00	3.91	2.07
.0555	3.97	1.89	3.92	1.88
.0655	3.95	1.87	4.01	1.78

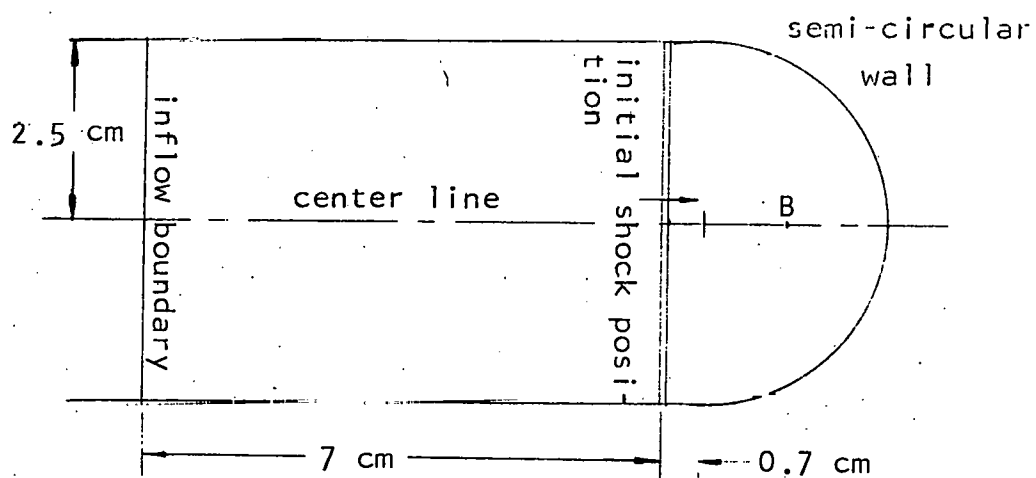


Fig. 8-1 The initial shock position in the flow region

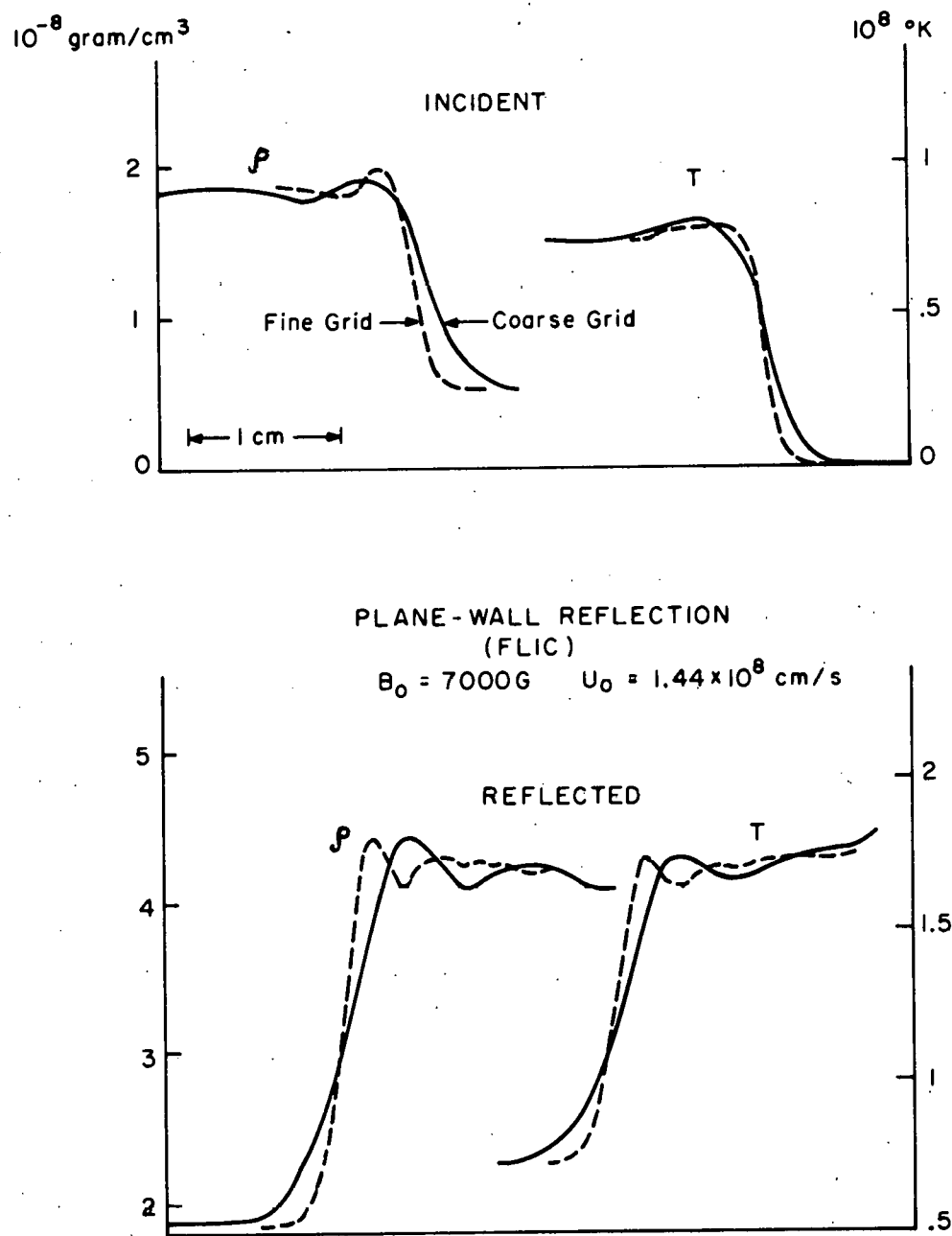


Fig. 8-2 The computer results of an incident shock wave reflecting from a flat wall in a plane flow.

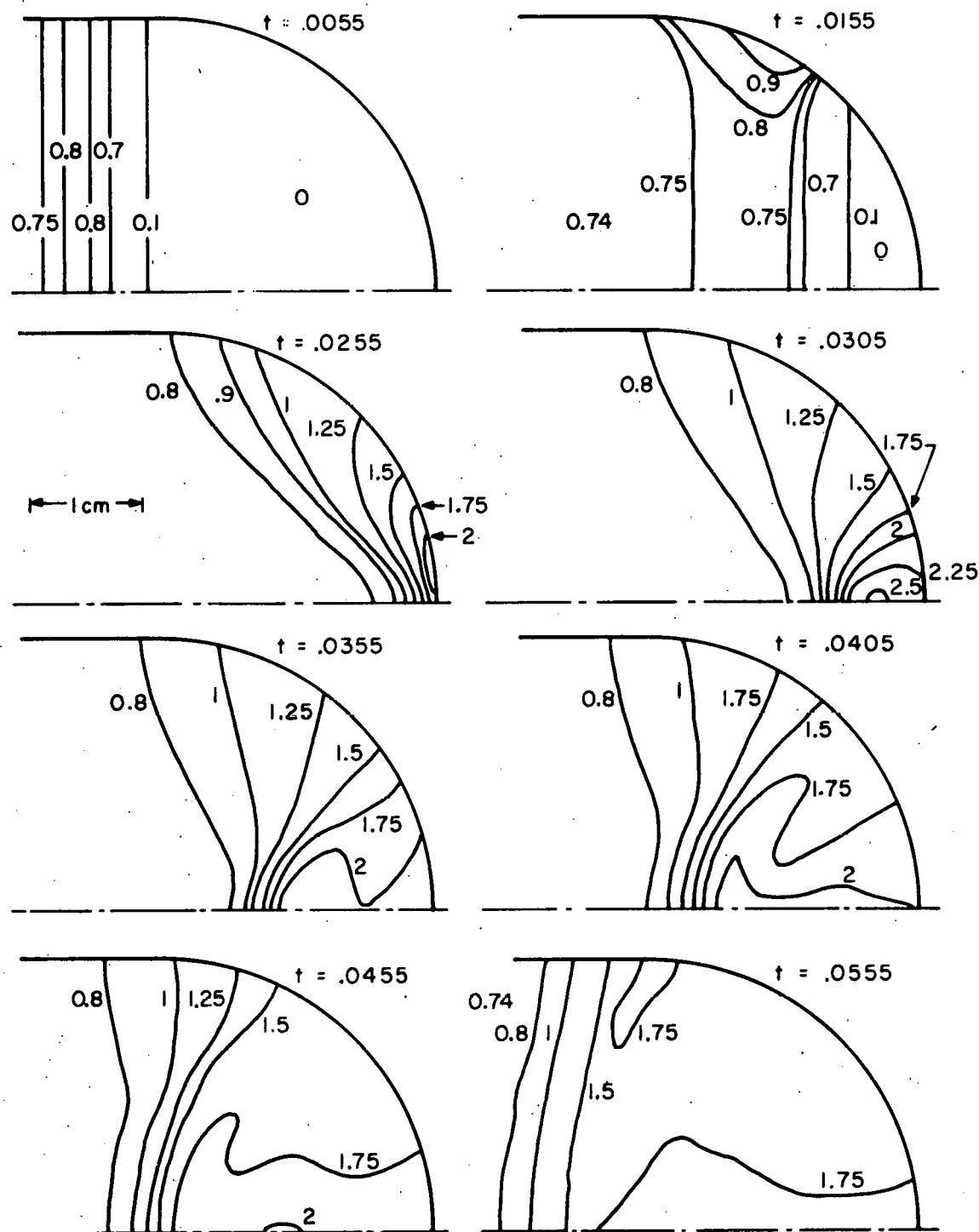


Fig. 8-3 The sequence of temperature contours for a plane shock wave reflecting from a semi-circular wall in a plane flow, with $B_0 = 7000G$, $U_s = 1.44 \times 10^8$ cm/sec, t in μ sec, T in 10^8 °K

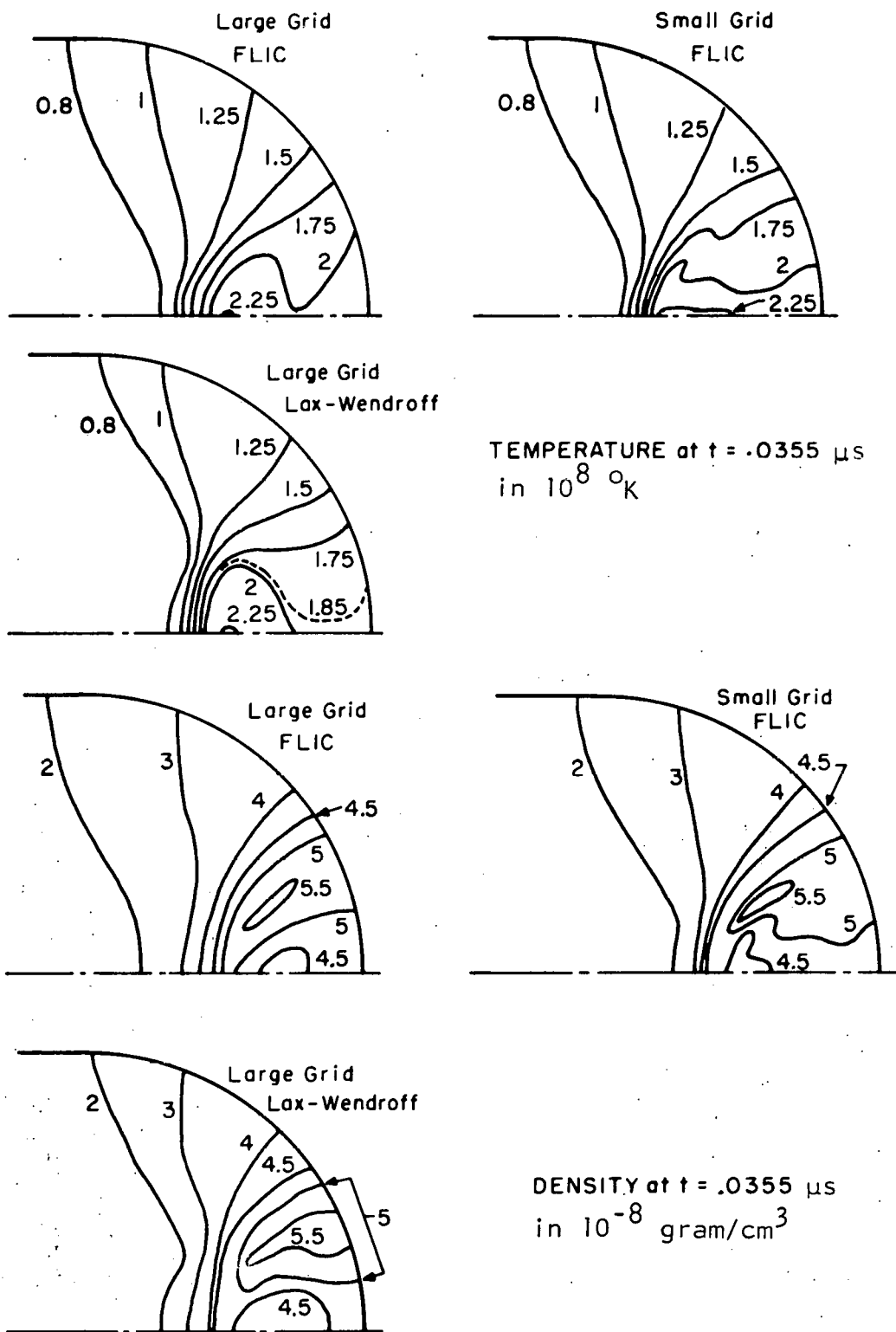


Fig. 8-4. The temperature and density contours obtained from three different calculations for the same instant of time.

REFERENCES

1. Meehan, R. J., Ph. D. Thesis, Report No. 44, Plasma Lab., Columbia Univ., (1969).
2. Halmoy, E., Phys. Fluids, 14, 2134, (1971).
3. Moriette, P. R., Phys. Fluids, 15, 51, (1972).
4. Gross, R. A., Chen, Y. G., Halmoy, E., and Moriette, P. R., Phys. Rev. Letters, 25, 575, (1970).
5. Schneider, S. H., Chu, C. K., and Leonard, B. P., Phys. Fluids, 14, 1103, (1971).
6. Schneider, S. H., Phys. Fluids, 15, 805, (1972).
7. Hain, K., Hain, G., Roberts, K. V., Roberts, S. J., and Koppendorfer, W., Z. Naturforsch, 15a, 1039, (1960).
8. Alder, B., Methods in Computational Physics, Vol. 3, Academic Press, (1964).
9. Chu, C. K., Computational Fluid Dynamics, AIAA Selected Reprint Series.
10. Potter, D. E., Phys. Fluids, 14, 1971, (1971).
11. Burstein, S. Z., AIAA Journal, 2, 2111, (1964).
12. Lapidus, A., J. Computational Phys., 2, 154, (1967).
13. Amsden, A. A., Los Alamos Rep. La-3466, (1966).
14. Marder, B. M., Los Alamos Rep. LA-DC 13103, (1971).
15. Gentry, R. A., Martin, R. E., and Daly, B. J., J. Computational Phys., 1, 87, (1966).
16. Butler, T. D., Phys. Fluids, 10, 1205, (1967).
17. Whitham, G. B., J. Fluid Mech., 2, 145, (1957).
18. Whitham, G. B., J. Fluid Mech., 5, 369, (1959).

19. Milton, B. E., and Archer, R. D., AIAA Journal, 7, 779, (1969).
Milton, B. E., and Archer, R. D., 8th International Symposium,
London, (1971).
20. Lau, J., CASI Transaction, 4, 13, (1971).
21. Skews, B. W., AIAA Journal, 10, 839, (1972).
22. Burstein, S. Z., J. Computational Phys., 2, 198, (1967).
23. Chu, C. K., and Harold Grad, in Research Frontiers in Fluid Dynamics, (R. Seeger and G. Temple, eds.), Wiley-Interscience, pp. 308-349, (1964).
24. Jeffrey, A., and Taniuti, T., Non-Linear Wave Propagation, Academic Press, p. 159, (1964).
25. Whitham, G. B., J. Fluid Mech., 2, 185, (1958).
26. Payne, R. B., J. Fluid Mech., 2, 185, (1957).
27. Lapidus, A., J. Computational Phys., 8, 106, (1971).
28. Private communication.
29. Richtmyer, R. D., and Morton, K. W., Difference Methods for Initial-Value Problems, second edition, Interscience Publishers, pp. 365-368, (1967).

APPENDIX A

The Two-Step Lax-Wendroff Scheme

The basic equations of a two-dimensional transverse MHD flow in Cartesian coordinates (x, y) can be expressed as follows:

$$\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x} (m) + \frac{\partial}{\partial y} (n) = 0, \quad (\text{A-1})$$

$$\frac{\partial m}{\partial t} + \frac{\partial}{\partial x} \left(\frac{m^2}{\rho} + P^* \right) + \frac{\partial}{\partial y} \left(\frac{mn}{\rho} \right) = 0, \quad (\text{A-2})$$

$$\frac{\partial n}{\partial t} + \frac{\partial}{\partial x} \left(\frac{mn}{\rho} \right) + \frac{\partial}{\partial y} \left(\frac{n^2}{\rho} + P^* \right) = 0, \quad (\text{A-3})$$

$$\frac{\partial E}{\partial t} + \frac{\partial}{\partial x} \left(\frac{m}{\rho} (E + P^*) \right) + \frac{\partial}{\partial y} \left(\frac{n}{\rho} (E + P^*) \right) = 0, \quad (\text{A-4})$$

$$\frac{\partial B}{\partial t} + \frac{\partial}{\partial x} \left(\frac{mB}{\rho} \right) + \frac{\partial}{\partial y} \left(\frac{nB}{\rho} \right) = 0. \quad (\text{A-5})$$

Here,

$$m = \rho u,$$

$$n = \rho v,$$

$$P^* = P + B^2/8\pi,$$

$$= (\gamma - 1) \left(E - \frac{1}{2}(m^2 + n^2)/\rho \right)$$

$$+ (2 - \gamma) B^2/8\pi,$$

$$P = (\gamma - 1) \rho I,$$

$$E = \rho \left(1 + \frac{1}{2}(u^2 + v^2) \right) + \frac{B^2}{8\pi},$$

where ρ , u , v , P , B and l are the density, x-component of velocity, y-component of velocity, thermal pressure, magnetic field and specific internal energy respectively.

These equations fit in the following form,

$$\frac{\partial w}{\partial t} = \frac{\partial}{\partial x} [f] + \frac{\partial}{\partial y} [g] . \quad (A-6)$$

The difference schemes²² for Eq. (A-6) consists of two steps as follows:

$$\begin{aligned} \text{Step 1: } w_{h,l}(t+\delta t) = & \frac{1}{4}(w_{h+\frac{1}{2},l+\frac{1}{2}} + w_{h+\frac{1}{2},l-\frac{1}{2}} \\ & + w_{h-\frac{1}{2},l+\frac{1}{2}} + w_{h-\frac{1}{2},l-\frac{1}{2}}) + \frac{\lambda}{2} (f_{h+\frac{1}{2},l+\frac{1}{2}} - f_{h-\frac{1}{2},l+\frac{1}{2}} \\ & + f_{h+\frac{1}{2},l-\frac{1}{2}} - f_{h-\frac{1}{2},l-\frac{1}{2}}) + \frac{\lambda}{2} (g_{h+\frac{1}{2},l+\frac{1}{2}} - g_{h+\frac{1}{2},l-\frac{1}{2}} \\ & + g_{h-\frac{1}{2},l+\frac{1}{2}} - g_{h-\frac{1}{2},l-\frac{1}{2}}) , \end{aligned}$$

$$\text{with } \delta x = \delta y , \quad \lambda = \delta t / \delta x ,$$

$$h = i \pm \frac{1}{2} , \quad l = j \pm \frac{1}{2} .$$

$$\begin{aligned} \text{Step 2: } w_{i,j}(t+\delta t) = & w_{i,j}(t) + \frac{\lambda}{4} (f_{i+1,j} - f_{i-1,j} + \underline{f}_x^{t+\delta t}) \\ & + \frac{\lambda}{4} (g_{i,j+1} - g_{i,j-1} + \underline{g}_y^{t+\delta t}) + q_{i,j}^t , \end{aligned}$$

$$\underline{f}_x^{t+\delta t} = f_{i+\frac{1}{2},j+\frac{1}{2}}^{t+\delta t} - f_{i-\frac{1}{2},j+\frac{1}{2}}^{t+\delta t} + f_{i+\frac{1}{2},j-\frac{1}{2}}^{t+\delta t} - f_{i-\frac{1}{2},j-\frac{1}{2}}^{t+\delta t} ,$$

$$\underline{g}_y^{t+\delta t} = g_{i+\frac{1}{2},j+\frac{1}{2}}^{t+\delta t} - g_{i+\frac{1}{2},j-\frac{1}{2}}^{t+\delta t} + g_{i-\frac{1}{2},j+\frac{1}{2}}^{t+\delta t} - g_{i-\frac{1}{2},j-\frac{1}{2}}^{t+\delta t} ,$$

$$q_{i,j}^t = \frac{0.5}{4} \lambda \cdot k \cdot (|u_{i+1,j}^t - u_{i,j}^t| \times (w_{i+1,j}^t - w_{i,j}^t))$$

$$\begin{aligned}
& -|u_{i,j}^t - u_{i-1,j}^t| \times (w_{i,j}^t - w_{i-1,j}^t) \\
& + |v_{i,j+1}^t - v_{i,j}^t| \times (w_{i,j+1}^t - w_{i,j}^t) \\
& - |v_{i,j}^t - v_{i,j-1}^t| \times (w_{i,j}^t - w_{i,j-1}^t) \quad ,
\end{aligned}$$

with

$$k = \frac{1.2}{2} \left[1 + \exp \left(1 - \frac{\sqrt{u^2 + v^2}}{U_2} \right) \right] \text{ if } \sqrt{u^2 + v^2} / U_2 < 1 ;$$

= 1.2 otherwise.

where $q_{i,j}^t$ is the viscosity term²⁹, U_2 is the speed of fluid behind an incoming shock wave.

APPENDIX B

The computer program for shock focusing and reflection from the selected quarter-elliptic end wall in a coaxial shock tube with a constant drive current


```

0244 200111,J-11*(DSX(1,1)-CSX(1,1)-CSX(1,1)-J-11)/CSX(1,1)
0247 62 IF (DSX(1,1).GT.CSX(1,1).AND. CSX(1,1).GT.O..AND.CP(1,1).J-11)
1.GT.O.) TPR=(PP(1,1,1)-CSX(1,1)-PP(1,1,1)-J-11)*DSX(1,1)-
2.CSX(1,1)-J-11)/DSX(1,1)
0248 1F (DSX(1,1).GT.C..AND.DSX(1,1).GT.DSX(1,1).AND.OP(1,1).J-11)
1.GT.O.) TPR=(PP(1,1,1)-CSX(1,1)-PP(1,1,1)-J-11)*DSX(1,1)-
2.CSX(1,1)-J-11)/DSX(1,1)
0249 1F (CSX(1,1).LT.CSX(1,1).AND.CP(1,1).J-11).GT.C..AND.
1.CSX(1,1).GT.O.) TPR=(PP(1,1,1)-J-11)*DSX(1,1)-CSX(1,1)-J-11)
2.PP(1,1,1)*DSX(1,1)+CSX(1,1)-J-11)/DSX(1,1)
0250 6 CONTINUE
0251 GUL(1,1)=FUI(1,1)-DT*(XJ*(L.F*(TPR-TPL)+C1-Q7)
3 /CP(1,1,1)+X*SF(1,1)
0252 GVL(1,1)=FV(1,1,1)-DT*(YIJ*(L.F*YALJ)+TPH-PP(1,1,1)-SYA(J-11)
1.TPR-PP(1,1,1))/2.+VCL(1,1)+C3-C4)/CP(1,1,1)+SF(1,1)
2-PP(1,1,1)*2/2.+XSR(J)/DX(1,1)+CP(1,1,1)+CP(1,1,1)
0253 1F (ARSIGUL(1,1)) .LE. EPS1) GUL(1,1)=0.
0254 1F (ARSIGVL(1,1)) .LE. EPS1) GVL(1,1)=0.
0255 21 CONTINUE
0256 DO 22 J=15,1F
0257 GUL(1,1,1)=GUL(1,1,1)
0258 GVL(1,1,1)=GVL(1,1,1)
0259 GUL(1,1,1)=GUL(1,1,1)
0260 GVL(1,1,1)=GVL(1,1,1)
0261 22 CONTINUE
0262 1F (1F.GT.1M) CALL CIRCLE(DX,1M,1C,JC,JP,RA,PEAN)
0263 NPA2
0264 DO 23 J=JB,JE
0265 DO 23 I=15,1F
0266 1F (ZSX(1,1).EQ.O..AND.RSY(1,1).EQ.O.) GO TO 23
0267 C=0.5
0268 1F (J.EQ.JE .OR. J.EQ.JR) C=C-18
0269 1F (1.LE.1M) GO TO 2223
0270 CALL DISTANCE(J,DA,CB,CFL,FX)
0271 1F (ABS(DEL).LE.1.20*PX) CR=0.1P
0272 223 CONTINUE
0273 CALL FRAC(SF(1,1),AXIJ,FIJ, FX,DY)
0274 1F (1F)
VISCOS(CB,CK,VP,FUR,FUL,GUR,GUL,FVT,FVB,GVT,GVB,DT,DX,
1F(J))
0275 G1(1,1)=F1(1,1,1)-DT*(AYIJ*(P1(1,1)+SYALJ)+FV1+
1.GVT
+VVL(1,1,1)+GVL(1,1,1)/4.-SYA(J-11)*(FVB+
2.GVB
+VVL(1,1,1)+GVL(1,1,1)/4.)+C2+.5*(SYE(J+1)+
3.FVT
+GVT
+SYR(J)+FV(1,1,1)+GVL(1,1,1)+GVL(1,1,1)-.5*C4+.5*
4.(SYR(J)+FV(1,1,1)+GVL(1,1,1)+SYE(J-11)*(FVB
+GVR))
5-1/4*(FV(1,1,1)+GVL(1,1,1)+SYR(J)+C3-C4)+XSR(J)*XIJ*.25*
6.FUR
+GUR
+FUL(1,1,1)+GUL(1,1,1)+P1(1,1,1)+Q1)
7-.25*(FUL
+GUL
+FUL(1,1,1)+GUL(1,1,1)+P1(1,1,1)+Q1)
8-1/4*(FUL(1,1,1)+GUL(1,1,1)+C1-Q7))/CP(1,1,1)+SF(1,1)+VOL(1,1,1)
9.F1(1,1,1)+G1(1,1,1)+F*(GUL(1,1,1)+2*GVL(1,1,1)+Q2)
0276 23 CONTINUE
0277 CALL TRANSP(DX,DY,DT,JB,JE,1F,EPST,15,PEAN,1M,1M)
0278 DO 477 J=JB,JE
0279 DO 477 I=15,1F
0280 1F (ZSX(1,1).EQ.O..AND.RSY(1,1).EQ.O.) GO TO 477
0281 1F (PSY(1,1).EQ.PSY(1+1,1).AND.(PSY(1,1).EQ.PSY(1-1,1).AND.

```

```

0282 1F (ZSX(1,1).EQ.ZSX(1,1+1).AND.(ZSX(1,1).EQ.ZSX(1,1-1)) GO TO 477
1F (DR(1,1-1,1).EQ.O..AND.RSY(1,1).EQ.O.) LE.PSY(1,1).AND.
1PSY(1,1).LE.PSY(1,1-1) GO TO 9102
0284 1F (PR(1,1-1,1).EQ.O..AND.FUI(1,1).LT.O..AND.VCL(1,1).GF.
1PSY(1,1)+SYR(J)) GO TO 9102
0285 GO TO 9105
0286 9102 ZSX(1,1)=ZSX(1,1)-FUI(2,1,1)*DT
0287 CSX(1,1)=ZSX(1,1)
0288 DSX(1,1)=ZSX(1,1)
0289 9105 CONTINUE
0290 1F (DR(1,1+1,1).EQ.O..AND.RSY(1,1).LE.RSY(1,1+1).AND.
1PSY(1,1).LE.RSY(1,1+1) GO TO 9104
0291 1F (PR(1,1+1,1).EQ.O..AND.FUI(1,1).GT.O..AND.VCL(1,1).GF.
1PSY(1,1)+SYR(J)) GO TO 9104
0292 GO TO 9107
0293 9104 ZSX(1,1)=ZSX(1,1)+FUI(2,1,1)*DT
0294 CSX(1,1)=ZSX(1,1)
0295 DSX(1,1)=ZSX(1,1)
0296 9107 CONTINUE
0297 477 CONTINUE
0298 DO 25 J=JB,JE
0299 DO 25 I=15,1F
0300 1F (ZSX(1,1).EQ.O..AND.RSY(1,1).EQ.O.) GO TO 25
0301 SFV=VCL(1,1)
0302 1F (PSY(1,1).EQ.PSY(1+1,1).AND.(PSY(1,1).EQ.PSY(1-1,1).AND.
1ZSX(1,1).EQ.ZSX(1,1+1).AND.(ZSX(1,1).EQ.ZSX(1,1-1)) GO TO 8551
0303 1F (PR(1,1-1,1).EQ.O..AND.CSY(1,1).LE.DSY(1,1).AND.CSY(1,1-1)
1J-1).LE.DSY(1,1-1)) VCL(1,1)=VCL(1,1)-PSY(1,1)*FUI(1,1)*DT+XSR(J)
2/DX
0304 1F (DR(1,1+1,1).EQ.O..AND.CSY(1,1).LE.CSY(1,1+1).AND.DSY(1,1)
1J-1).LE.CSY(1,1+1)) VCL(1,1)=VCL(1,1)+PSY(1,1)*FUI(1,1)*DT+XSR(J)
2/DX
0305 1F (PR(1,1-1,1).EQ.O..AND.CSY(1,1).GT.DSY(1,1).AND.CSY(1,1-1)
1J-1).GT.DSY(1,1-1)) VCL(1,1)=VCL(1,1)-(PSY(1,1)+DSY(1,1)+DSY(1,1-1)
2J-1)-CSY(1,1)+PSY(1,1)+FUI(1,1)*DT+XSR(J)/DX
0306 1F (DR(1,1+1,1).EQ.O..AND.CSY(1,1).GT.CSY(1,1+1).AND.DSY(1,1+1)
1J-1).GT.CSY(1,1+1)) VCL(1,1)=VCL(1,1)+(PSY(1,1)+DSY(1,1)+DSY(1,1+1)
2FV(1,1)-DSY(1,1)+PSY(1,1)+FUI(1,1)*DT+XSR(J)/DX
0307 1F (DR(1,1-1,1).EQ.O..AND.CSY(1,1).GT.CSY(1,1-1).AND.
1FV(1,1)-DSY(1,1)+PSY(1,1)+FUI(1,1)*DT+XSR(J)/DX
2FV(1,1)-DSY(1,1)+PSY(1,1)+FUI(1,1)*DT+XSR(J)/DX
0308 1F (DR(1,1+1,1).EQ.O..AND.CSY(1,1).GT.CSY(1,1+1).AND.
1FV(1,1)-DSY(1,1)+PSY(1,1)+FUI(1,1)*DT+XSR(J)/DX
2FV(1,1)-DSY(1,1)+PSY(1,1)+FUI(1,1)*DT+XSR(J)/DX
0309 1F (DR(1,1-1,1).EQ.O..AND.CSY(1,1).GT.CSY(1,1-1).AND.
1FV(1,1)-DSY(1,1)+PSY(1,1)+FUI(1,1)*DT+XSR(J)/DX
2FV(1,1)-DSY(1,1)+PSY(1,1)+FUI(1,1)*DT+XSR(J)/DX
0310 1F (DR(1,1+1,1).EQ.O..AND.CSY(1,1).GT.CSY(1,1+1).AND.
1FV(1,1)-DSY(1,1)+PSY(1,1)+FUI(1,1)*DT+XSR(J)/DX
2FV(1,1)-DSY(1,1)+PSY(1,1)+FUI(1,1)*DT+XSR(J)/DX
0311 1F (DR(1,1-1,1).EQ.O..AND.CSY(1,1).GT.CSY(1,1-1).AND.
1FV(1,1)-DSY(1,1)+PSY(1,1)+FUI(1,1)*DT+XSR(J)/DX
2FV(1,1)-DSY(1,1)+PSY(1,1)+FUI(1,1)*DT+XSR(J)/DX
0312 1F (DSY(1,1).GT.O..AND.CSY(1,1).LE.DSY(1,1).AND.FUI(1,1).LT.
1O..AND.VCL(1,1).GF.PSY(1,1)+YALJ) GO TO 9061
0313 GO TO 35A2
0314 CSY(1,1)=DSY(1,1)
0315 DSY(1,1)=PSY(1,1)

```



```

C390      IF (ABS(FV(1,1,J)) .LE. EPS) FV(1,1,J)=0.
C400      IF (ABS(FV(1,1,J)) .LE. EPS) FV(1,1,J)=0.
C401      IF (I.E. 1) GO TO 5575
C402      CALL DISTAN(1,1,CA,CP,CFL,CX)
C403      IF (ABS(DEL).LE.1.74*CX) GO TO 5558
C404      9549 CONTINUE
C405      IF (DEL(1,1,J).GT.0..AND.CR(1,1,J-1).GT.0..AND.RSV(1,1,NE.DY)
      100 TO 9109
      GO TO 9109
C406      RSV(1,1,J)=DY
C407      CSV(1,1,J)=DY
C408      RSV(1,1,J)=DY
C409      CSV(1,1,J)=DY
C410      9109 CONTINUE
C411      IF (DEL(1,1,J).GT.0..AND.CR(1,1,J-1).GT.0..AND.ZSX(1,1,NE.DX) GO
      1 TO 9208
      GO TO 9208
C412      9208 ZSX(1,1,J)=DX
C413      CSX(1,1,J)=DX
C414      DSX(1,1,J)=DX
C415      9209 CONTINUE
C416      IF (RSV(1,1,J).GT.DY..AND.VOL(1,1,LY. VOL(2,1)/2.) GO TO 9551
C417      GO TO 9552
C418      RSV(1,1,J)=DY
C419      CSV(1,1,J)=DY
C420      RSV(1,1,J)=DY
C421      CSV(1,1,J)=DY
C422      9552 IF (ZSX(1,1,J).GT.DX .AND.VOL(1,1,LY.VOL(2,1)/2.) GO TO 9553
C423      GO TO 9554
C424      ZSX(1,1,J)=DX
C425      CSX(1,1,J)=DX
C426      DSX(1,1,J)=DX
C427      9554 IF (RSV(1,1,J).GT. 1.24*DY) GO TO 9555
C428      GO TO 9556
C429      RSV(1,1,J)=VOL(1,1)+DX/DSX(1,1)/ZSX(1,1)
C430      CSV(1,1,J)=RSV(1,1,J)
C431      DSX(1,1,J)=RSV(1,1,J)
C432      9556 IF (ZSX(1,1,J).GT. 1.24*DX) GO TO 9557
C433      GO TO 9558
C434      ZSX(1,1,J)=VOL(1,1)+DX/DSX(1,1)/RSV(1,1,J)
C435      CSX(1,1,J)=ZSX(1,1,J)
C436      DSX(1,1,J)=ZSX(1,1,J)
C437      9558 CONTINUE
C438      95 CONTINUE
C439      TIME=TIME+DT
C440      IF (M .EQ. 400) GO TO 1227
C441      GO TO 1238
C442      1237 CONTINUE
C443      PRINTD 8
C444      WRITE (R,1229) IF .A. TIME,DT
C445      1229 FORMAT (14,14,E15.8,E15.8/)
C446      DO 1106 J=1,JD
C447      DO 1106 I=1,IE
C448      WRITE (R,1109) DEL(1,1,J),FV(1,1,J),FV(1,1,J),DEL(1,1,J),FV(1,1,J),
      1109 1TC(1,1,J),PP(1,1,J),P(1,1,J),PB(1,1,J),IMP(1,1,J),VOL(1,1,J),
      1110 ZSX(1,1,J),CSX(1,1,J),ZSX(1,1,J),CSV(1,1,J),RSV(1,1,J),EV(1,1,J),
      1111 RSV(1,1,J),FV(1,1,J)
C449      1106 CONTINUE
C450      1106 CONTINUE

```

```

C450      1105 FORMAT (5F15.8/)
C451      END FILE 8
C452      1238 CONTINUE
C453      GO TO 8
C454      21 CONTINUE
C455      CALL EXITG(2)
C456      STOP
C457      END

```



```

0020      IE=16
0030      C LABELS AXES AND TITLES GRID
0031      IF(IITICK.NE.1)GO TO 1
0032      CALL SETSMG(2,102,.01)
0033      CALL SETSMG(2,103,.01)
0034      C HEIGHT OF TICK MARKS ABOUT 30 RASTERS
0035      DX=DX/(INX-2)
0036      DY=DY/(INY-2)
0037      CALL LABELG(2,0,CX,1,IEK)
0038      IE=17
0039      WRITE (6,999) IE
0040      CALL LABELG(2,1,CY,1,IEK)
0041      IE=18
0042      WRITE (6,999) IE
0043      IF(ICRS.NE.1)GO TO 2
0044      CALL SETSMG(2,14,C.)
0045      C SCALING IN SUBJECT SPACE UNITS
0046      CALL SEGHTG(2,1,XPIR,0,XMAX,0.)
0047      CALL SEGHTG(2,1,0,YMIN,C,YMAX)
0048      C DRAWS CIRCLES LINES
0049      CONTINUE
0050      DCNT=(CNTPX-CNTMNI)/(NCNT-1)
0051      CNT(1)=CNTMNI
0052      DO 3 I=2,NCNT
0053      CNT(I)=CNT(I-1)+DCNT
0054      WRITE(6,103)(I,CNT(I),I=1,NCNT)
0055      CALL OBJCTG(2,0.1,C.849,1.33,1.)
0056      IE=19
0057      WRITE (6,999) IE
0058      C TWO OBJECT SPACE IS TO RIGHT OF GRID WHERE WE WILL LIST CONTOUR
0059      C VALUES. MODE NO. 14 ALREADY 0 FROM A PREVIOUS CALL (SUBJECT SPACE UNITS)
0060      C CHANGE IT BACK TO 39 (OBJECT SPACE UNITS)
0061      CALL SETSMG(2,14,3.)
0062      IE=20
0063      WRITE (6,999) IE
0064      CALL LEGNDG(2,0.1,C.98,6,'CYCLE=')
0065      IE=21
0066      WRITE (6,999) IE
0067      CALL NUMPRG(2,0.2,C.98,4,M)
0068      IE=22
0069      WRITE (6,999) IE
0070      CALL LEGNDG(2,0.3,C.98,6,'TIME=')
0071      IE=23
0072      WRITE (6,999) IE
0073      CALL NUMPRG(2,0.4,C.98,-14.6,TIME)
0074      IE=24
0075      WRITE (6,999) IE
0076      CALL LEGNDG(2,0.1,C.96,8,CNTP)
0077      IE=25
0078      WRITE (6,999) IE
0079      YY=0.94
0080      XX=0.1
0081      DO 4 I=1,NCNT
0082      IF (I.EQ. 6) YY=0.1
0083      IF (I.EQ. 6) YY=0.92

```

```

0078      CALL FMTSG(2,1,4,0.1,CNTR(1))
0079      CALL FMTSG(2,3,12,3,CNTR(1),CNTR(2))
0080      CALL LEGNDG(2,XX,YY,20,CNTR(1))
0081      XX=XX*0.2
0082      CONTINUE
0083      CALL OBJCTG(2,0.1,C.84,1.3,YEN)
0084      IE=26
0085      WRITE (6,999) IE
0086      C RESETS OBJECT SPACE TO SQUARE
0087      C NOW START WORK DRAWING CONTOURS
0088      CALL SETSMG(2,14,C.)
0089      IE=27
0090      WRITE (6,999) IE
0091      NX=NX-2
0092      NY=NY-2
0093      DO 5 I=1,NCNT
0094      CALL FMTSG(2,1,1,0.1,ICNT)
0095      ICCUNT=0
0096      X=XMIN+DX/2.
0097      C WHEN ICOUNT IS A MULTIPLE OF 10, PRINT THE COUNTER NUMBER ON THE CONTOUR
0098      DO 6 I=2,NXI
0099      Y=YMIN+DY/2.
0100      DO 7 JY=2,NYI
0101      DO 8 JJ=1,4
0102      MCC(JJ)=0.
0103      MCC(JJ)=0.
0104      IF(IIDBUG.EQ.1)WRITE(6,104)IX,JY,X,Y
0105      IF(IIX,JY).EQ.CNTR(1) MCC(1)=1
0106      IF(IIX,JY+1).EQ.CNTR(1) MCC(2)=1
0107      IF(IIX+1,JY).EQ.CNTR(1) MCC(3)=1
0108      IF(IIX+1,JY+1).EQ.CNTR(1) MCC(4)=1
0109      A=CNT(1)-F(IIX,JY)
0110      B=CNT(1)-F(IIX,JY+1)
0111      C=CNT(1)-F(IIX+1,JY)
0112      D=CNT(1)-F(IIX+1,JY+1)
0113      IF(IIDBUG.EQ.1)WRITE(6,105)A,B,C,D
0114      IF(A<0.LT.0)MCS(1)=1
0115      IF(A<0.LT.0)MCS(2)=1
0116      IF(B<0.LT.0)MCS(3)=1
0117      IF(C<0.LT.0)MCS(4)=1
0118      C THAT IS, IF THE PRODUCT OF THE DIFFERENCES IS NEGATIVE, THE
0119      C TWO ENDS LIE ON OPPOSITE SIDES OF CNTR(1), OR THE LINE INTERSECTS
0120      C THE GRID BOX SIDE. A,B,C,AND D WILL BE USED LATER TO CALCULATE
0121      C INTERPOLATIONS
0122      IPCC=MCC(1)+MCC(2)+MCC(3)+MCC(4)
0123      IMCS=MCS(1)+MCS(2)+MCS(3)+MCS(4)
0124      IF(IMCC.GF.3)GO TO 7
0125      C IF MORE THAN 2 CORNERS EQUAL TO THE CONTOUR, CONT DRAW ANY LINES IS THIS
0126      C BOX.
0127      IF(IMCS.NE.1)AND.1MCS.NE.3)GO TO 9
0128      IF(IMCC.NE.0)GO TO 9
0129      WRITE(6,100)IMCS
0130      C THE LOGIC OF THE SITUATION MAKES IT OBVIOUS THAT THE CONTOUR LINE
0131      C MAY ONLY INTERSECT AN EVEN NUMBER OF SIDES, UNLESS IT ALSO INTERSECTS
0132      C A CORNER
0133      RETURN

```

```

0127 9 IF(1MCS .EQ. 0 .AND. 1MCC .EQ. 0) GO TO 7
0128 C LINE DOES NOT CROSS THIS BOX, GO ON TO NEXT
0129 IF(1MCS .EQ. 2) GO TO 11
0130 C GO TO 11 IF 2 SIDES INTERSECTED. IF THE OPPOSITE CORNER IS ALSO
0131 C INTERSECTED, IGNORE IT
0132 IF(1MCS .EQ. 4) GO TO 12
0133 C GO TO STATEMENT 12 IF ALL FOUR SIDES INTERSECTED
0134 IF(1MCS .EQ. 1 .AND. 1MCC .EQ. 1) GO TO 13
0135 C TO 13 IF ONE SIDE AND ONE CORNER INTERSECTED
0136 IF(1MCC .EQ. 1) GO TO 7
0137 C IF ONE CORNER INTERSECTED ONLY, IGNORE IT. IT WILL BE PICKED UP
0138 C IN ANOTHER BOX.
0139 C AT THIS POINT IN THE CODE, 1MCC=2, 1MCS=C, OR TWO CORNERS INTERSECTED,
0140 C NO SIDES
0141 C IF THE TWO CORNERS ARE OPPOSITE, IGNORE THEM IF THE OTHER TWO
0142 C CORNERS ARE ON THE SAME SIDE OF THE CONTOUR
0143 IF(1MCC(1) .EQ. 1MCC(3) .AND. 1MCC(2) .EQ. 1MCC(4)) GO TO 71
0144 X1=X
0145 Y1=Y
0146 X2=X+DX
0147 Y2=Y
0148 IF(1MCC(1) .EQ. 1) GO TO 14
0149 IF(1MCC(2) .NE. 1) GO TO 15
0150 V1=Y1+DY
0151 Y2=Y2+DY
0152 X2=X2+DX
0153 GO TO 17
0154 15 IF(1MCC(3) .NE. 1) GO TO 14
0155 X1=X1+DX
0156 Y1=Y1+DY
0157 X2=X2+DX
0158 GO TO 17
0159 14 IF(1MCC(4) .NE. 1) GO TO 16
0160 X2=X2+DX
0161 GO TO 17
0162 16 Y2=Y2+DY
0163 CONTINUE
0164 C
0165 IF(1MCC(1) .EQ. 1) WRITE(6,100) X1,Y1,X2,Y2
0166 C
0167 CALL SEGMENT(2,X1,Y1,X2,Y2)
0168 ICOUNT=ICOUNT+1
0169 IF(1MCC(1) .NE. 1) GO TO 7
0170 ICOUNT=C
0171 CALL LEGNDG(2,X1,Y1,1,ICOUNT)
0172 GO TO 7
0173 C HERE TWO OPPOSITE CORNERS HAVE A CONTOUR GOING THROUGH THEM.
0174 C TEST IF OTHER TWO CORNERS ARE ON SAME SIDE OF THE CONTOUR. IF SO
0175 C SKIP OUT. THE CORNERS WILL BE PICKED UP IN OTHER BOXES.
0176 71 IF(1MCC(2) .EQ. 1) GO TO 72
0177 IF(1MCC(4) .GT. 0) GO TO 7
0178 C IF CORNERS 2 AND 4 ARE ON SAME SIDE OF CONTOUR GO OUT, OTHERWISE
0179 C CONNECT 1 AND 3
0180 X1=X
0181 Y1=Y
0182 X2=X+DX

```

```

0160 Y2=Y+DY
0161 GO TO 17
0162 72 IF(A=C .GT. 0) GO TO 7
0163 C IF 1 AND 3 ON SAME SIDE GO OUT, OTHERWISE CONNECT 2 AND 4
0164 X1=X
0165 Y1=Y+DY
0166 X2=X+DX
0167 Y2=Y
0168 GO TO 17
0169 C NOW TAKE 1MCS=2, 1MCC=1, C, BUT DOESN'T PLAY A PART
0170 11 IC=1
0171 IF(1MCS(1) .NE. 1) GO TO 151
0172 XPT(10)=X+DX*A/(F(IX+1,JY)-F(IX,JY))
0173 YPT(10)=Y
0174 IO=IO+1
0175 IF(1MCS(2) .NE. 1) GO TO 161
0176 XPT(10)=X
0177 YPT(10)=Y+DY*A/(F(IX,JY+1)-F(IX,JY))
0178 IO=IO+1
0179 IF(10 .EQ. 3) GO TO 18
0180 IF(1MCS(3) .NE. 1) GO TO 171
0181 XPT(10)=X+DX*2/(F(IX+1,JY+1)-F(IX,JY+1))
0182 YPT(10)=Y+DY
0183 IO=IO+1
0184 IF(10 .EQ. 3) GO TO 18
0185 IF(1MCS(4) .NE. 1) GO TO 19
0186 XPT(10)=X+DX
0187 YPT(10)=Y+DY*D/(F(IX+1,JY+1)-F(IX+1,JY))
0188 CONTINUE
0189 18
0190 C IF(1MCC(1) .EQ. 1) WRITE(6,107) (XPT(10),YPT(10),10-1,2)
0191 C
0192 CALL LINES(2,XPT,YPT)
0193 ICOUNT=ICOUNT+1
0194 IF(1MCC(1) .NE. 1) GO TO 7
0195 ICOUNT=C
0196 CALL LEGNDG(2,XPT(1),YPT(1),1,ICOUNT)
0197 GO TO 7
0198 19 WRITE(6,102)
0199 GO TO 7
0200 C NOW CASE OF ALL SIDES INTERSECTED, 1MCC=0
0201 12 DX1=DX*A/(F(IX+1,JY)-F(IX,JY))
0202 DX2=DX*B/(F(IX+1,JY+1)-F(IX,JY+1))
0203 CY1=C*A/(F(IX,JY+1)-F(IX,JY))
0204 DY2=DY*D/(F(IX+1,JY+1)-F(IX+1,JY))
0205 C NOW TEST NEXT BOX TO SEE WHICH PAIRS TO CONNECT.
0206 C IF F(IX,JY+2) AND F(IX,JY+1) ARE ON THE SAME SIDE OF CNT(1),
0207 C THEN CONNECT SIDES 2 AND 3 AND SIDES 1 AND 4, OTHERWISE 1-2 AND 3-4.
0208 C IF JY+2 .GT. NY, USE F(IX-1,JY) AND F(IX,JY). IF THEY ARE ON THE SAME
0209 C SIDE OF CNT(1), CONNECT 1-2,3-4, OTHERWISE 1-4,2-3
0210 C
0211 IF (JY+2 .GT. NY-1) GO TO 20
0212 XX=B*(CAT(1)-F(IX,JY+2))
0213 IF(XX .LT. 0) GO TO 21
0214 XPT(1)=X+DX1
0215 YPT(1)=Y

```

```

C205      XPT(2)=X
C206      YPT(2)=Y+DY1
C207      XPT(1)=X+DX
C208      YPT(1)=Y+DY2
C209      XPT(2)=X+DX2
C210      YPT(2)=Y+DY
C211      GO TO 22
      C THAT CONNECTS 1-4,2-3
C212      21 XPT(1)=X+CX1
C213      YPT(1)=Y+CY1
C214      XPT(2)=X+CX2
C215      YPT(2)=Y+CY
C216      XPT(1)=X
C217      YPT(1)=Y+DY1
C218      XPT(2)=X+DX
C219      YPT(2)=Y+DY2
      C THAT CONNECTED 1-2,3-4
C220      22 CONTINUE
C221      C
      IF(IDEBUG.EQ.1)WRITE(6,1C8)(XPT(1),YPT(1),XPT(2),YPT(2),
      IIT=1,2)
      C
C222      CALL SEGMTG(2,2,XPT,YPT,XPT1,YPT1)
C223      ICOUNT=ICOUNT+1
C224      IF(ICOUNT.NE.10)GO TO 7
C225      ICOUNT=0
C226      CALL LEGNDG(2,XPT(1),YPT(1),1,ICHT)
C227      GO TO 7
C228      20 XX=A*(CAT(1)-F(IX,JY-1))
C229      IF(XX.LT.C)GO TO 23
C230      GO TO 21
      C NOW WE TAKE CASE OF ONE SIDE AND ONE CORNER
C231      13 X1=X
C232      Y1=Y
C233      IF(MCC(2).EQ.1)Y1=Y1+CY
C234      IF(MCC(3).EQ.1)X1=X1+DX
C235      IF(MCC(1).EQ.1)GC TO 24
C236      IF(MCC(4).NE.1)GC TO 24
C237      X1=X+DX
C238      Y1=Y+DY
C239      24 IF(MCS(1).NE.1)GC TO 25
C240      X2=X+DX*A/(F(IX+1,JY)-F(IX,JY))
C241      Y2=Y
C242      GO TO 28
C243      25 IF(MCS(2).NE.1)GC TO 26
C244      X2=X
C245      Y2=Y+DY*A/(F(IX,JY+1)-F(IX,JY))
C246      GO TO 28
C247      26 IF(MCS(3).NE.1)GC TO 27
C248      X2=X+DX*B/(F(IX+1,JY+1)-F(IX,JY+1))
C249      Y2=Y+CY
C250      GO TO 28
C251      27 X2=X+DX
C252      Y2=Y+DY*D/(F(IX+1,JY+1)-F(IX+1,JY))
C253      28 CONTINUE
      C

```

```

C254      C
      IF(IDEBUG.EQ.1)WRITE(6,1C9)(X1,Y1,X2,Y2)
      C
C255      CALL SEGMTG(2,1,X1,Y1,X2,Y2)
C256      ICOUNT=ICOUNT+1
C257      IF(ICOUNT.NE.1)GC TO 7
C258      ICOUNT=0
C259      CALL LEGNDG(2,X1,Y1,1,ICHT)
C260      Y=Y+DY
C261      6 X=X+DX
C262      5 CONTINUE
C263      DO 201 I=2,NXX
C264      XF=I
C265      DO 201 J=2,NYY
C266      YG=J
C267      XAX(I,J)=XMIN+(XF-2.-3.)*CX/2.
C268      YAY(I,J)=YMIN+(YG-2.-3.)*DY/2.
C269      XPX(I,J)=XAX(I,J)+FV(1,1,J)/0.8E+08*CX
C270      YPY(I,J)=YAY(I,J)+FV(1,1,J)/0.8E+08*DY
C271      201 CONTINUE
C272      DO 209 I=2,NXX
C273      DO 209 J=2,NYY
C274      IF (RSV(I,J).EQ.0.) GO TO 209
C275      CALL POINTG(2,1,XAX(I,J),YAY(I,J))
C276      209 CONTINUE
C277      CALL RSETMG(2)
C278      IF=6
C279      WRITE (6,999) IE
C280      YEN=0.36*YDIS
C281      CALL OBJCTG(2,0.1,C.36,1.3,YEN)
C282      IF=9
C283      WRITE (6,999) IE
C284      CALL SUBJEG(2,XMIN,YMIN,XPAX,YMAX)
C285      IF=2
C286      WRITE (6,999) IE
C287      DX=XMAX-XMIN
C288      DY=YMAX-YMIN
C289      CALL GRIDG(2,DX,DY,0,0)
C290      IF=10
C291      WRITE (6,999) IE
C292      CALL TITLEG(2,15,15+DEASITY SPACING,8,LY,0,0)
C293      IF=11
C294      WRITE (6,999) IE
C295      DX=DX/(NX-2)
C296      DY=DY/(NY-2)
C297      CALL SETSPG(2,14,0.)
C298      IF=5
C299      WRITE (6,999) IE
C300      DO 221 I=2,NXX
C301      DO 221 J=2,NYY
C302      IF (RSV(I,J).EQ.0.) GO TO 221
C303      NN=DR(1,1,J)/0.05CCOE-C7
C304      IF (NN.GT.17) NN=17
C305      PX(I)=XAX(I,J)
C306      PY(I)=YAY(I,J)
C307      PX(2)=XAX(I,J)
C308      PY(2)=YAY(I,J)+DY/2.22

```

```
0309      PX(3)=XAX(I,J)
0310      PY(3)=YAY(I,J)-DY/2.22
0311      PX(4)=XAX(I,J)+DX/2.22
0312      PY(4)=PY(1)
0313      PX(5)=PX(4)
0314      PY(5)=PY(2)
0315      PX(6)=PX(4)
0316      PY(6)=PY(3)
0317      PX(7)=XAX(I,J)-DX/2.22
0318      PY(7)=PY(1)
0319      PX(8)=PX(7)
0320      PY(8)=PY(2)
0321      PX(9)=PX(7)
0322      PY(9)=PY(3)
0323      PX(10)=PX(1)+DX/4.6
0324      PY(10)=PY(1)+DY/4.6
0325      PX(11)=PX(10)
0326      PY(11)=PY(10)-DY/4.6
0327      PX(12)=PX(11)-DX/4.6
0328      PY(12)=PY(10)
0329      PX(13)=PX(12)
0330      PY(13)=PY(12)
0331      PX(14)=PX(14)-DX/4.6
0332      PY(14)=PY(12)
0333      PX(15)=PX(14)
0334      PY(15)=PY(11)
0335      PX(16)=PX(17)+DX/4.6
0336      PY(16)=PY(12)
0337      PX(17)=PX(16)
0338      PY(17)=PY(13)
0339      DD 222 NPT=1,NN
0340      CALL POINTG(2,1,PX(INPT),PY(RPT))
0341      222 CONTINUE
0342      221 CONTINUE
0343      YEN=0.1+YDIS
0344      CALL OBJCTG(2,C11,C:211.3,YEN)
0345      IE=1
0346      WRITE (6,999) IE
0347      DX=XMAX-XMIN
0348      DY=YMAX-YMIN
0349      CALL GRIDG(2,DX,DY,0,0)
0350      IE=3
0351      WRITE (6,999) IE
0352      CALL TITLEG(2,8,8+VELCC(1),8,LY,47,47) NUMERICAL SHOCK WAVES SIMUL
    1ATIONS AT COLUMBIA
0353      IE=4
0354      WRITE (6,999) IE
0355      DX=DX/INX-2)
0356      DY=DY/INY-2)
0357      CALL SETSMG(2,30,CJ9)
0358      IE=7
0359      WRITE (6,999) IE
0360      DD 312 I=2,NX
0361      DD 312 J=2,NY
0362      IF (RSY(I,J).EQ.0.) GC TC 312
0363      CALL SEGMENT(2,1,XAX(I,J),YAY(I,J),XPX(I,J),YPY(I,J))
```

```
0364      CALL POINTG(2,1,XAX(I,J),YAY(I,J))
0365      312 CONTINUE
0366      CALL PAGEG(2,0,1,1)
0367      IE=230
0368      WRITE (6,999) IE
0369      CALL RSETHG(2)
0370      IE=31
0371      WRITE (6,999) IE
0372      RETURN
0373      100 FORMAT ('OERROR, IPCS= ',I2)
0374      101 FORMAT ('OERROR,NCNT= ',I5)
0375      102 FORMAT ('OERROR, IPCS=2, BUT LESS THAN TWO PCS =1')
0376      103 FORMAT ('O1',1CX,'CNT(1)'/(15,5X,1PE12.3))
0377      104 FORMAT ('O1X= ',I5,' JY= ',I5,' X= ',1PE12.3,' Y= ',1PE12.3)
0378      105 FORMAT ('O1X= ',1PE12.3,' P= ',1PE12.3,' C= ',1PE12.3,' D= ',1PE12.3)
0379      106 FORMAT ('O2 CORNERS/' X1= ',1PE12.3,' Y1= ',1PE12.3,' X2= ',1PE12.3,
    13,' Y2= ',1PE12.3)
0380      107 FORMAT ('O2 SIDES/' XPT=1CX,'YPT'/(12PE12.3))
0381      108 FORMAT ('OFCUR SIDES/' XPT=1CX,'YPT=1CX,'XPT1=1CX,'YPT1'/(
    11PE12.3))
0382      109 FORMAT ('ODONE SIDE,CNE CORNER/' X1= ',1PE12.3,' Y1= ',1PE12.3,
    1' X2= ',1PE12.3,' Y2= ',1PE12.3)
0383      999 FORMAT (1H,'IE=',I4)
0384      END
```

```

0001 SUPROUTINE DIVINCEK,CY,FK,FR,CV,JP,JE,IS,IE,IP,JP,ICJ
0002 DIMENSION CM(2,137,22),FU(2,137,22),FV(2,137,22),FI(2,137,22),
1FE(2,137,22),TC(2,137,22),PM(2,137,22),PI(2,137,22),PM(2,137,22),
2GUL(1,137,22),GV(1,137,22),GIL(1,137,22),GFI(1,137,22),TMF(2,137,22),
3SYA(22),SAP(22),SYH(22),VCL(137,22),S1(22),CPH(22),OVI(137,22),
4CSX(137,22),DSX(137,22),CSY(137,22),CSY(137,22),RSY(137,22),
5ZSX(137,22)
0003 COMMON DR,FU,FV,FI,FE,TC,PP,P,RR,CU,GV,GT,GF,TMP,SYA,SRB,SYR,VOL,
1SI,CPB,NV,CSA,USX,CSY,DSY,PSY, ZSX,C1,C2,Q3,C4,P,I,J
0004 DC 744 J=JP,JE
0005 DC 744 I=IS,IE
0006 IF (ZSX(I,J).EQ.O..AND.RSY(I,J).EQ.O.) GC TC 744
0007 IF (RSY(I,J).EQ.OX..AND.ZSX(I,J).EQ.OX..AND.VOL(I,J).EQ.VOL(2,J))
1 GO TO 744
0008 L=O
0009 VY=O.
0010 VZ=O.
0011 IF (DR(I,I+1,J).EQ.C1) L=-1
0012 IF (DR(I,I-1,J).EQ.O.) L=1
0013 IF (L.EQ.C1) GO TC 744
0014 IF (ZSX(I,J).GT.1.5*CX..AND.VOL(I,J).GT.VOL(2,J)..AND.
1VOL(I,J).GT.1.4*RSY(I,J)+SYH(J)) GC TC 761
0015 IF (RSY(I-L,J+1).GT.1.5*CX..AND.VCL(I-L,J+1).GE.VCL(2,J+1))
1..AND.VOL(I-L,J+1).GT.1.4*ZSX(I-L,J+1)+SYH(J+1)) GO TO 762
0016 IF (RSY(I-L,J-1).GT.1.5*CX..AND.VOL(I-L,J-1).GE.VOL(2,J-1))
1..AND.VOL(I-L,J-1).GT.1.4*ZSX(I-L,J-1)+SYH(J-1)) GO TC 762
0017 GC TO 744
0018 762 CONTINUE
0019 IF (ZSX(I,J).GT.CX..AND.VOL(I,J).GT.RSY(I,J)+SYH(J)) GC TC861
0020 GC TC 744
0021 761 IF (ZSX(I,J+1).GT.1.5*CX..CP.ZSX(I,J-1).GT.1.5*CX) GC TC 861
0022 IF (DR(I,I-L,J+1).EQ.O..AND.DR(I,I-L,J-1).EQ.O.) GO TO 744
0023 861 CCATINUE
0024 VY=VOL(I,J)-RSY(I,J)+SYE(J)
0025 IF (VOL(I-L,J+1).GT.ZSX(I-L,J+1)+SYH(J+1)..AND.RSY(I-L,J+1)
1..GT.DY) VY=VCL(I-L,J+1)-ZSX(I-L,J+1)+SYE(J+1)
0026 IF (VOL(I-L,J-1).GT.ZSX(I-L,J-1)+SYE(J-1)..AND.RSY(I-L,J-1)
1..GT.DY) VZ=VOL(I-L,J-1)-ZSX(I-L,J-1)+SYE(J-1)
0027 VOL(I-L,J)=VY+VZ
0028 IF (VY.EQ.O..AND.VZ.EQ.O.) GO TC 861
0029 GC TO 882
0030 881 DSX(I-L,J)=DSX(I,J)-DX
0031 DSX(I-L,J)=DSX(I,J)-DX
0032 ZSX(I-L,J)=APAXI(CSX(I-L,J),CSX(I-L,J))
0033 GO TO 981
0034 882 IF (VY.GT.O..AND.VZ.EQ.O.) GC TO 882
0035 GC TO 884
0036 883 CSX(I-L,J)=ZSX(I-L,J+1)
0037 DSX(I-L,J)=ZSX(I-L,J+1)
0038 ZSX(I-L,J)=ZSX(I-L,J+1)
0039 VOL(I-L,J+1)=VOL(I-L,J+1)-VY
0040 DSY(I-L,J+1)=DY
0041 CSY(I-L,J+1)=DY
0042 RSY(I-L,J+1)=DY
0043 CSX(I-L,J+1)=ZSX(I-L,J+1)
0044 DSX(I-L,J+1)=ZSX(I-L,J+1)

```

```

0045 GO TO 981
0046 884 IF (VY.EQ.O..AND.VZ.GT.O.) GO TO 885
0047 GO TO 886
0048 885 CSX(I-L,J)=ZSX(I-L,J-1)
0049 DSX(I-L,J)=ZSX(I-L,J-1)
0050 ZSX(I-L,J)=ZSX(I-L,J-1)
0051 VOL(I-L,J-1)=VOL(I-L,J-1)-VZ
0052 CSY(I-L,J-1)=DY
0053 DSY(I-L,J-1)=DY
0054 RSY(I-L,J-1)=DY
0055 CSX(I-L,J-1)=ZSX(I-L,J-1)
0056 DSX(I-L,J-1)=ZSX(I-L,J-1)
0057 GO TO 981
0058 886 IF (VY.GT.O..AND.VZ.GT.O.) GC TC 887
0059 GO TO 981
0060 887 ZSX(I-L,J)=VCL(I-L,J)/SYE(J)
0061 CSX(I-L,J)=ZSX(I-L,J)
0062 DSX(I-L,J)=ZSX(I-L,J)
0063 RSY(I-L,J)=DY
0064 CSY(I-L,J+1)=DY
0065 DSY(I-L,J+1)=DY
0066 VOL(I-L,J+1)=VOL(I-L,J+1)-VY
0067 DSY(I-L,J+1)=DY
0068 RSY(I-L,J+1)=DY
0069 CSY(I-L,J+1)=DY
0070 VOL(I-L,J-1)=VOL(I-L,J-1)-VZ
0071 CSX(I-L,J+1)=ZSX(I-L,J+1)
0072 DSX(I-L,J+1)=ZSX(I-L,J+1)
0073 CSX(I-L,J-1)=ZSX(I-L,J-1)
0074 DSX(I-L,J-1)=ZSX(I-L,J-1)
0075 981 CONTINUE
0076 CSY(I,J)=PSY(I,J)
0077 DSY(I,J)=RSY(I,J)
0078 VOL(I,J)=VOL(I,J)-VX
0079 DSX(I,J)=DX
0080 CSX(I,J)=DX
0081 ZSX(I,J)=OX
0082 RSY(I-L,J)=VCL(I-L,J)*CX/ZSX(I-L,J)/SX(I,J)
0083 CSY(I-L,J)=RSY(I-L,J)
0084 DSY(I-L,J)=RSY(I-L,J)
0085 DR(I,I-L,J)=DR(I,I,J)+VX*CR(I,I-L,J)+VY*CR(I,I-L,J)+VZ
1..AB(I,I-L,J)=ER(I,I,J)+VX*RB(I,I-L,J)+VY*SYB(I,J)/SYB(J+1)+
1..RB(I,I-L,J)+VZ*SYE(J)/SYE(J-1)
0087 FE(I,I-L,J)=CR(I,I,J)+VX*FE(I,I,J)+CR(I,I-L,J)+VY*FE(I,I-L,J)+
1..CP(I,I-L,J)+VZ*FE(I,I-L,J-1)
0088 FU(I,I-L,J)=CR(I,I,J)+VX*FU(I,I,J)+CR(I,I-L,J)+VY*FU(I,I-L,J)+
1..DR(I,I-L,J-1)+VZ*FU(I,I-L,J-1)
0089 FV(I,I-L,J)=CR(I,I,J)+VX*FV(I,I,J)+CR(I,I-L,J)+VY*FV(I,I-L,J)+
1..PI(I,I-L,J)+VZ*FV(I,I-L,J-1)
0090 FE(I,I-L,J)=FE(I,I-L,J)/CR(I,I-L,J)
0091 FU(I,I-L,J)=FU(I,I-L,J)/CR(I,I-L,J)
0092 FV(I,I-L,J)=FV(I,I-L,J)/CR(I,I-L,J)
0093 CR(I,I-L,J)=CR(I,I-L,J)/VCL(I-L,J)
0094 DR(I,I-L,J)=LE(I,I-L,J)/VCL(I-L,J)
0095 FI(I,I-L,J)=FE(I,I-L,J)-S*FU(I,I-L,J)+2*FV(I,I-L,J)+2)
0096 TC(I,I-L,J)=FK*FR*FI(I,I-L,J)/CVI*O.5

```

```

CC97      P(1,1-L,J)=(FK-1,1)*DR(1,1-L,J)*F(1,1-L,J)
CC98      PP(1,1-L,J)=P(1,1-L,J)*PP(1,1-L,J)**2/25.1328
CC99      TMP(1,1-L,J)=F(1,1-L,J)/CV
C100      BV(1,1-L,J)=0.
C101      SI(J)-SI(J)-1.
C102      744 CONTINUE
C103      DC 748 J=JB,JE
C104      DO 748 I=1,IE
C105      IF (ZSX(I,J).EQ.0..AND.RSY(I,J).EQ.0.) GO TO 748
C106      IF (RSY(I,J).EQ.0X..AND.ZSX(I,J).EQ.0X..AND.VUL(I,J).EQ.VOL(2,J))
      1 GO TO 748
C107      L=0
C108      IF (DR(1,1,J-1).EQ.0.) L=-1
C109      IF (DR(1,1,J-1).EQ.0.) L=+1
C110      IF (L.EQ.0) GO TO 748
C111      IF (RSY(I,J).GT. 1.5*DY .AND. VCL(1,J).GE. VCL(2,J).AND.
      1VUL(1,J).GT. 1.4*ZSX(1,J)*SYB(J)) GO TO 743
C112      GC TO 748
C113      763 IF (RSY(I+1,J).GT. 1.5*DY .OR. RSY(I-1,J).GT. 1.5*DY) GO TO 764
C114      IF (DR(1,1+J-L).CC.0..AND.DR(1,1-J-L).EQ.0.) GO TO 748
C115      764 CONTINUE
C116      CSX(I,J-L)=ZSX(I,J)
C117      DSX(I,J-L)=ZSX(I,J)
C118      ZSX(I,J-L)=ZSX(I,J)
C119      VOL(I,J-L)=VCL(I,J)-ZSX(I,J)*SYB(J)
C120      RSY(I,J-L)=VCL(I,J-L)*CX/SX(I,J-L)/ZSX(I,J-L)
C121      CSY(I,J-L)=RSY(I,J-L)
C122      OSY(I,J-L)=RSY(I,J-L)
C123      VOL(I,J)=ZSX(I,J)*SYB(J)
C124      CSY(I,J)=OY
C125      OSY(I,J)=OY
C126      RSY(I,J)=OY
C127      CSX(I,J)=ZSX(I,J)
C128      DSX(I,J)=ZSX(I,J)
C129      FE(1,1,J-L)=FE(1,1,J)
C130      FU(1,1,J-L)=FU(1,1,J)
C131      FV(1,1,J-L)=FV(1,1,J)
C132      FI(1,1,J-L)=FI(1,1,J)
C133      DR(1,1-L)=CR(1,1,J)
C134      TC(1,1,J-L)=TC(1,1,J)
C135      P (1,1,J-L)= P(1,1,J)
C136      BB(1,1,J-L)=BB(1,1,J)*SYB(J-L)/SYB(J)
C137      PP(1,1,J-L)=P(1,1,J-L)*BB(1,1,J-L)**2/25.1328
C138      TMP(1,1,J-L)=TMP(1,1,J)
C139      BV(1,1-L)=0.
C140      SI(J-L)-SI(J-L)-1.
C141      748 CONTINUE
C142      RETURN
C143      END

```

```

0001      SUBROUTINE DISTAN(I,J,CA,CB,DEL,DX)
0002      CA=5.
0003      U=1.
0004      XIM1=92.
0005      XJP=2.
0006      XI=1
0007      XJ=J
0008      DA=(XI-XIM1+.5)*DX
0009      DB=(XJ-XJP+.5)*DX
0010      DAB=SQRT(CA**2+DB**2)
0011      ZN=ATAN(DB/CB)
0012      OS=RA+SQRT(U/(COS(ZN)**2+(SIN(ZN)/2.25)**2))
0013      DEL=CAB-DS
0014      RETURN
0015      END

```


FOOTRAN IV G LEVEL 19

UPPER

DATE = 72173

19/11/47

PAGE 0001

```

0001 SUBROUTINE UPPER(DX, I, K, IR, CV, JR, JF, P, H, AN)
0002 G14, KSI(19, D12, 137, 221, F12, 137, 221, F12, 137, 221, F12, 137, 221,
137, 221, 137, 221, F12, 137, 221, F12, 137, 221, F12, 137, 221, F12, 137, 221,
2001, 137, 221, CV11, 137, 221, G11, 137, 221, G11, 137, 221, G11, 137, 221,
35VA(221, SX1(221, SYH(221, VOL(137, 221, S1(221, CPH(221, BV1(137, 221,
40SX(137, 221, OSX(137, 221, CSY(137, 221, OSY(137, 221, PSY(137, 221,
52SX(137, 221)
0003 CCM, MCN, DR, DV, IV, F1, F2, TC, PP, P, BR, CU, W, G1, G2, TMP, SYA, SKD, SYB, VOL,
451, CPH, BV, CSX, OSX, CSY, OSY, HSY, Z5, Q1, Q2, Q3, Q4, M, I, J
0004 L=0
0005 IF (PSY(I, J-1), GT, .1E-02) L=L+1
0006 IF (PSY(I, J+1), GT, .1E-02) L=L+1
0007 IF (L, EQ, 0) GO TO 1
0008 IF (I, J, L), GT, J) GO TO 1
0009 IF (I, J, L), LT, J) GO TO 1
0010 VS=0.
0011 K=0
0012 IF (DR(I, I+1, J), .EQ, 0.) K=1
0013 IF (DR(I, I-1, J), .EQ, 0.) K=1
0014 IF (K, .EQ, 0, .OR, (DR(I, I+1, J), .EQ, 0, .AND, DR(I, I-1, J), .EQ, 0, .AND, TO2
0015 IF (ZSX(I, J), GT, ZSX(I, J+1), .AND, PSY(I, J+1), GT, 0, .AND, VOL(I, J)
1, GT, ZSX(I, J+1), .AND, CSY(I, J+1), .AND, SYA(I, J+1), .AND, SYB(I, J+1), .AND, VOL(I, J)
1, ZSX(I, J+1), .AND, CSY(I, J+1), .AND, SYA(I, J+1), .AND, SYB(I, J+1), .AND, VOL(I, J)
VOL(I, J)=VOL(I, J)-VS
0016 2 CONTINUE
0017 FU(I, I, J+1)=(VOL(I, J)+DR(I, I, J)+FU(I, I, J)+VOL(I, J+1)+DR(I, I, J+1)
0018 1, FU(I, I, J+1)/(VOL(I, J)+DR(I, I, J)+FU(I, I, J)+VOL(I, J+1)+DR(I, I, J+1)
0019 FU(I, I, J+1)=(VOL(I, J)+DR(I, I, J)+FU(I, I, J)+VOL(I, J+1)+DR(I, I, J+1)
1, FU(I, I, J+1)/(VOL(I, J)+DR(I, I, J)+FU(I, I, J)+VOL(I, J+1)+DR(I, I, J+1)
0020 FU(I, I, J+1)=(VOL(I, J)+DR(I, I, J)+FU(I, I, J)+VOL(I, J+1)+DR(I, I, J+1)
1, FU(I, I, J+1)/(VOL(I, J)+DR(I, I, J)+FU(I, I, J)+VOL(I, J+1)+DR(I, I, J+1)
0021 DR(I, I, J+1)=(VOL(I, J)+DR(I, I, J)+VOL(I, J+1)+DR(I, I, J+1)/
1(VOL(I, J)+VOL(I, J+1))
0022 BR(I, I, J+1)=(VOL(I, J)+SX(I, J+1)/SX(I, J)+VOL(I, J+1)+
1BR(I, I, J+1)/(VOL(I, J)+VOL(I, J+1))
0023 FU(I, I, J+1)=FU(I, I, J+1)/DR(I, I, J+1)
0024 FU(I, I, J+1)=FU(I, I, J+1)/DR(I, I, J+1)
0025 FE(I, I, J+1)=FE(I, I, J+1)/DR(I, I, J+1)
0026 FI(I, I, J+1)=FI(I, I, J+1)-.5*(FU(I, I, J+1)+2*FU(I, I, J+1)+2)
0027 TC(I, I, J+1)=(FK+FR+FE(I, I, J+1)/CV)**.5
0028 P(I, I, J+1)=(FK+FR+FE(I, I, J+1)/CV)**.5
0029 P(I, I, J+1)=P(I, I, J+1)+BR(I, I, J+1)**2/25, 1328
0030 TMP(I, I, J+1)=FI(I, I, J+1)/CV
0031 VOL(I, J+1)=VOL(I, J)+VOL(I, J+1)
0032 RSY(I, J+1)=VOL(I, J+1)*DX/SYF(J+1)/ZSX(I, J+1)
0033 CSY(I, J+1)=PSY(I, J+1)
0034 DSY(I, J+1)=PSY(I, J+1)
0035 ST(I, J)=ST(I, J)+1
0036 IF (VS, .LE, 0, .OR, K, .EQ, 0) GO TO 11
0037 FU(I, I, K, J+1)=(VS+DR(I, I, J)+FU(I, I, J)+VOL(I, I, K, J+1)+DR(I, I, K, J+1)
1*FU(I, I, K, J+1)/(VOL(I, I, K, J+1)+VS)
0038 FU(I, I, K, J+1)=(VS+DR(I, I, J)+FU(I, I, J)+VOL(I, I, K, J+1)+DR(I, I, K, J+1)
1*FU(I, I, K, J+1)/(VOL(I, I, K, J+1)+VS)
0039 FE(I, I, K, J+1)=(VS+DR(I, I, J)+FE(I, I, J)+VOL(I, I, K, J+1)+DR(I, I, K, J+1)
1*FE(I, I, K, J+1)/(VOL(I, I, K, J+1)+VS)
0040 DR(I, I, K, J+1)=(VS+DR(I, I, J)+VOL(I, I, K, J+1)+DR(I, I, K, J+1)/

```

FOOTRAN IV G LEVEL 19

UPPER

DATE = 72173

19/11/47

PAGE 0002

```

0041 1(VOL(I, I, K, J+1)+VS)
0042 BR(I, I, K, J+1)=(VOL(I, I, K, J+1)+VS)
1BR(I, I, K, J+1)/(VOL(I, I, K, J+1)+VS)
0043 FU(I, I, K, J+1)=FU(I, I, K, J+1)/DR(I, I, K, J+1)
0044 FE(I, I, K, J+1)=FE(I, I, K, J+1)/DR(I, I, K, J+1)
0045 FI(I, I, K, J+1)=FI(I, I, K, J+1)-.5*(FU(I, I, K, J+1)+2*FU(I, I, K, J+1)+2)
0046 TC(I, I, K, J+1)=(FK+FR+FE(I, I, K, J+1)/CV)**.5
0047 P(I, I, K, J+1)=(FK+FR+FE(I, I, K, J+1)/CV)**.5
0048 P(I, I, K, J+1)=P(I, I, K, J+1)+BR(I, I, K, J+1)**2/25, 1328
0049 TMP(I, I, K, J+1)=FI(I, I, K, J+1)/CV
0050 VOL(I, I, K, J+1)=VOL(I, I, K, J+1)+VS
0051 RSY(I, I, K, J+1)=VOL(I, I, K, J+1)*DX/SYF(J+1)/ZSX(I, I, K, J+1)
0052 CSY(I, I, K, J+1)=PSY(I, I, K, J+1)
0053 DSY(I, I, K, J+1)=PSY(I, I, K, J+1)
0054 11 CALL VACUUM(MEAN, DX)
0055 1 RETURN
0056 END

```

```

0001      SUBROUTINE TRANSP(IX,IY,IT,JK,JF,IF,IPSI,IS,IM,AN,IM,JM)
0002      DIMENSION IP(2,137,22),IC(2,137,22),IUI(2,137,22),IV(2,137,22),I(2,137,22),
      IF(2,137,22),IC(2,137,22),PP(2,137,22),PE(2,137,22),BU(2,137,22),
      2GUIL(137,22),GV(1,137,22),G(1,137,22),GL(1,137,22),IMPE(137,22),
      3SYA(22),SX(22),SY(22),VOL(137,22),SI(22),CPR(22),JVI(137,22),
      4CSX(137,22),DSX(137,22),CSY(137,22),OSY(137,22),RSY(137,22),
      5ZSX(137,22)
      COMMON DR,FV,FV,FE,FC,PP,P,IM,GO,GV,G,GE,IMP,SYA,SXB,SYB,VOL,
      1SI,CPR,IV,CSX,DSX,CSY,OSY,RSY,ZSA,CI,Q2,Q3,Q4,M,I,J
      0003      DO 25 J=JK,JF
      0004      DO 25 I=IS,IE
      0005      IF (ZSX(I,J).EQ.0.)AND.(RSY(I,J).EQ.0.) GO TO 25
      0006      SFI=VOL(I,J)/VUL(I,J)
      0007      IF (RSY(I,J).EQ.RSY(I+1,J))AND.(RSY(I,J).EQ.RSY(I-1,J))AND.
      0008      1(ZSX(I,J).EQ.ZSX(I,J+1))AND.(ZSX(I,J).EQ.ZSX(I,J-1)) GO TO 6
      0009      IF (C(1,I-1,J+1).GT.0.)AND.(C(1,I-1,J-1).GT.0.) GO TO 153
      0010      IF (CSY(I-1,J).GT.0.)AND.(CSY(I-1,J).LT.OSY(I,J))AND.
      1IGUI(1,I,J)+GUI(1,I-1,J)).LT.0.)AND.(VUL(I,J).GE.RSY(I,J)*SYB(J))
      2GO TO 9571
      0011      GO TO 153
      0012      9571 CSY(I-1,J)=DSY(I,J)
      0013      OSY(I-1,J)=DSY(I,J)
      0014      RSY(I-1,J)=DSY(I,J)
      0015      153 CONTINUE
      0016      IF (C(1,I+1,J+1).GT.0.)AND.(C(1,I+1,J-1).GT.0.) GO TO 155
      0017      IF (OSY(I+1,J).GT.0.)AND.(OSY(I+1,J).LT.CSY(I,J))AND.
      1IGUI(1,I,J)+GUI(1,I+1,J)).GT.0.)AND.(VUL(I,J).GE.RSY(I,J)*SYB(J))
      2GO TO 9573
      0018      GO TO 155
      0019      9573 DSY(I+1,J)=CSY(I,J)
      0020      CSY(I+1,J)=CSY(I,J)
      0021      PSY(I+1,J)=CSY(I,J)
      0022      155 CONTINUE
      0023      IF (C(1,I-1,J-1).GT.0.)AND.(C(1,I+1,J-1).GT.0.) GO TO 152
      0024      IF (CSX(I,J-1).GT.0.)AND.(CSX(I,J-1).LT.DSX(I,J))AND.
      1IGV(1,I,J)+GV(1,I,J-1)).LT.0.)AND.(VUL(I,J).GE.ZSX(I,J)*SYB(J))
      2GO TO 9575
      0025      GO TO 152
      0026      9575 CSX(I,J-1)=DSX(I,J)
      0027      DSX(I,J-1)=DSX(I,J)
      0028      ZSX(I,J-1)=DSX(I,J)
      0029      152 CONTINUE
      0030      IF (C(1,I-1,J+1).GT.0.)AND.(C(1,I+1,J+1).GT.0.) GO TO 156
      0031      IF (CSX(I,J+1).GT.0.)AND.(CSX(I,J+1).LT.CSX(I,J))AND.
      1IGV(1,I,J)+GV(1,I,J+1)).GT.0.)AND.(VUL(I,J).GE.ZSX(I,J)*SYB(J))
      2GO TO 9577
      0032      GO TO 156
      0033      9577 DSX(I,J+1)=CSX(I,J)
      0034      ZSX(I,J+1)=CSX(I,J)
      0035      CSX(I,J+1)=CSX(I,J)
      0036      156 CONTINUE
      0037      6 CONTINUE
      0038      AA=AMIN(CSY(I,J),DSY(I+1,J))/DX
      0039      AH=AMIN(DSY(I,J),CSY(I-1,J))/DX
      0040      AC=AMIN(CSX(I,J),DSX(I+1,J))/DX
      0041      AD=AMIN(DSX(I,J),CSX(I-1,J))/DX

```

```

0042      TTR=C.
0043      TTL=0.
0044      TTT=C.
0045      TTG=C.
0046      TTA=0.
0047      TTB=C.
0048      TTC=0.
0049      TTD=0.
0050      TTE=C.
0051      TT2=0.
0052      TT3=C.
0053      TT4=C.
0054      DPR=C.
0055      DML=C.
0056      DMT=0.
0057      DMG=C.
0058      DRR=C.
0059      DBL=C.
0060      DST=C.
0061      DRG=C.
0062      DMI=C.
0063      DM2=C.
0064      DM3=C.
0065      DM4=C.
0066      DMA=C.
0067      DMD=C.
0068      DMC=C.
0069      DMD=C.
0070      DBI=0.
0071      DB2=C.
0072      DB3=C.
0073      DB4=C.
0074      DBA=C.
0075      DBB=C.
0076      DDC=0.
0077      DHD=C.
0078      IF (1.5*(GUI(1,I,J)+GUI(1,I+1,J)).GT.0.) DM1=SXB(J)*DR(1,I,J)*.5*
      1 (GUI(1,I,J)+GUI(1,I+1,J))*DT*AA
0079      IF (1.5*(GUI(1,I,J)+GUI(1,I+1,J)).LT.0.) DM1=SXB(J)*DR(1,I,J)*.5*
      1 (GUI(1,I,J)+GUI(1,I+1,J))*DT*AA
0080      IF (1.5*(GV(1,I,J)+GV(1,I+1,J)).GT.0.) DM3=SYA(J)*DR(1,I,J)*.5*
      1 (GV(1,I,J)+GV(1,I+1,J))*DT*AC
0081      IF (1.5*(GV(1,I,J)+GV(1,I+1,J)).LT.0.) DM3=SYA(J)*DR(1,I,J)*.5*
      1 (GV(1,I,J)+GV(1,I+1,J))*DT*AC
0082      IF (1.5*(GUI(1,I,J)+GUI(1,I-1,J)).GT.0.) DM2=SXB(J)*DR(1,I-1,J)*
      1 (GUI(1,I,J)+GUI(1,I-1,J))*DT*.5*AB
0083      IF (1.5*(GUI(1,I,J)+GUI(1,I-1,J)).LT.0.) DM2=SXB(J)*DR(1,I,J)*
      1 (GUI(1,I,J)+GUI(1,I-1,J))*DT*.5*AB
0084      IF (1.5*(GV(1,I,J)+GV(1,I-1,J)).GT.0.) DM4=SYA(J-1)*DR(1,I,J-1)*
      1.5*(GV(1,I,J)+GV(1,I-1,J))*DT*AD
0085      IF (1.5*(GV(1,I,J)+GV(1,I-1,J)).LT.0.) DM4=SYA(J-1)*DR(1,I,J)*
      1.5*(GV(1,I,J)+GV(1,I-1,J))*DT*AD
0086      IF (1.5*(GUI(1,I,J)+GUI(1,I+1,J)).GT.0.) DBI=SXB(J)*DB(1,I,J)*.5*
      1 (GUI(1,I,J)+GUI(1,I+1,J))*DT*2*.5*1416*DX/SXB(J)*AA
0087      IF (1.5*(GUI(1,I,J)+GUI(1,I+1,J)).LT.0.) DBI=SXB(J)*DB(1,I,J)*.5*
      1 (GUI(1,I,J)+GUI(1,I+1,J))*DT*2*.5*1416*DX/SXB(J)*AA

```



```

0162 9325 DMC=SYA(J)*DR(I,I,J)*.5*(GV(I,I,J)+GV(I,I-1,J-1))*DT*(CSX(I,I,J)-
1 DSX(I,I,J))/DX
0163 DMC=SYA(J)*BR(I,I,J)*.5*(GV(I,I,J)+GV(I,I-1,J-1))*DT*(CSX(I,I,J)-
1 DSX(I,I,J))/DX
0164 GO TO 9328
0165 9326 IF (DR(I,I-1,J-1).EQ.0.)AND.(CSX(I,I,J).GT.DSX(I,I,J).AND.(GV(I,I,J)+
1 GV(I,I-1,J-1)).LT.0.) GO TO 9327
0166 GO TO 9328
0167 9327 DMC=SYA(J)*DR(I,I-1,J-1)*.5*(GV(I,I,J)+GV(I,I-1,J-1))*DT*(CSX(I,I,J)-
1 DSX(I,I,J))/DX
0168 DMC=SYA(J)*BR(I,I-1,J-1)*.5*(GV(I,I,J)+GV(I,I-1,J-1))*DT*(CSX(I,I,J)-
1 DSX(I,I,J))/DX
0169 TTC=1
0170 9328 CONTINUE
0171 IF (DR(I,I-1,J-1).EQ.0.) GO TO 9334
0172 IF (DR(I,I,J-1).EQ.0.)AND.(CSX(I,I-1,J-1).GT.DSX(I,I-1,J-1).AND.
1 (GV(I,I,J)+GV(I,I-1,J-1)).GT.0.) GO TO 9331
0173 GO TO 9332
0174 9331 DMC=SYA(J-1)*DR(I,I-1,J-1)*.5*(GV(I,I,J)+GV(I,I-1,J-1))*DT*
1 (CSX(I,I-1,J-1)-DSX(I,I-1,J-1))/DX
0175 DMC=SYA(J-1)*BR(I,I-1,J-1)*.5*(GV(I,I,J)+GV(I,I-1,J-1))*DT*
1 (CSX(I,I-1,J-1)-DSX(I,I-1,J-1))/DX
0176 TTD=1
0177 GO TO 9334
0178 9332 IF (DR(I,I,J-1).EQ.0.)AND.(CSX(I,I-1,J-1).GT.DSX(I,I-1,J-1).AND.
1 (GV(I,I,J)+GV(I,I-1,J-1)).LT.0.) GO TO 9333
0179 GO TO 9334
0180 9333 DMC=SYA(J-1)*DR(I,I,J-1)*.5*(GV(I,I,J)+GV(I,I-1,J-1))*DT*
1 (CSX(I,I-1,J-1)-DSX(I,I-1,J-1))/DX
0181 DMC=SYA(J-1)*BR(I,I,J-1)*.5*(GV(I,I,J)+GV(I,I-1,J-1))*DT*
1 (CSX(I,I-1,J-1)-DSX(I,I-1,J-1))/DX
0182 9334 CONTINUE
0183 IF (DR(I,I-1,J-1).EQ.0.) GO TO 9338
0184 IF (DR(I,I,J-1).EQ.0.)AND.(CSX(I,I-1,J-1).GT.DSX(I,I-1,J-1).AND.
1 (GV(I,I,J)+GV(I,I-1,J-1)).GT.0.) GO TO 9335
0185 GO TO 9336
0186 9335 DMC=SYA(J-1)*DR(I,I-1,J-1)*.5*(GV(I,I,J)+GV(I,I-1,J-1))*DT*
1 (CSX(I,I-1,J-1)-DSX(I,I-1,J-1))/DX
0187 DMC=SYA(J-1)*BR(I,I-1,J-1)*.5*(GV(I,I,J)+GV(I,I-1,J-1))*DT*
1 (CSX(I,I-1,J-1)-DSX(I,I-1,J-1))/DX
0188 TTC=1
0189 GO TO 9338
0190 9336 IF (DR(I,I,J-1).EQ.0.)AND.(CSX(I,I-1,J-1).GT.DSX(I,I-1,J-1).AND.
1 (GV(I,I,J)+GV(I,I-1,J-1)).LT.0.) GO TO 9337
0191 GO TO 9338
0192 9337 DMC=SYA(J-1)*DR(I,I,J-1)*.5*(GV(I,I,J)+GV(I,I-1,J-1))*DT*
1 (CSX(I,I-1,J-1)-DSX(I,I-1,J-1))/DX
0193 DMC=SYA(J-1)*BR(I,I,J-1)*.5*(GV(I,I,J)+GV(I,I-1,J-1))*DT*
1 (CSX(I,I-1,J-1)-DSX(I,I-1,J-1))/DX
0194 9338 CONTINUE
0195 IF (DR(I,I-1,J-1).EQ.0.) GO TO 8404
0196 IF (DR(I,I,J-1).EQ.0.)AND.(DSY(I,I,J).GT.CSY(I,I-1,J).AND.(GV(I,I,J)+
1 GV(I,I-1,J-1)).LT.0.) GO TO 8301
0197 GO TO 8302
0198 8301 DMB=SR(I,J)*DR(I,I,J)*.5*(GV(I,I,J)+GV(I,I-1,J-1))*DT*(DSY(I,I,J)-
1 CSY(I,I-1,J))/DX

```

```

0199 DBB=BR(I,I,J)*.5*(GV(I,I,J)+GV(I,I-1,J-1))*DT*(DSY(I,I,J)-CSY(I,I-1,J))
1 *.6.2832
0200 GO TO 8404
0201 8302 IF (DR(I,I,J-1).EQ.0.)AND.(DSY(I,I,J).GT.CSY(I,I-1,J).AND.(GV(I,I,J)+
1 GV(I,I-1,J-1)).GT.0.) GO TO 8303
0202 GO TO 8404
0203 8303 DMB=SR(I,J)*DR(I,I-1,J-1)*.5*(GV(I,I,J)+GV(I,I-1,J-1))*DT*
1 (DSY(I,I-1,J-1)-CSY(I,I-1,J-1))/DX
0204 DMB=SR(I,J)*BR(I,I-1,J-1)*.5*(GV(I,I,J)+GV(I,I-1,J-1))*DT*(DSY(I,I-1,J-1)-
1 CSY(I,I-1,J-1))/DX
0205 TTS=1
0206 8404 CONTINUE
0207 IF (DR(I,I-1,J-1).EQ.0.) GO TO 8405
0208 IF (DR(I,I,J-1).EQ.0.)AND.(DSY(I,I,J).GT.CSY(I,I-1,J).AND.(GV(I,I,J)+
1 GV(I,I-1,J-1)).LT.0.) GO TO 8407
0209 GO TO 8406
0210 8405 DMB=SR(I,J)*DR(I,I,J)*.5*(GV(I,I,J)+GV(I,I-1,J-1))*DT*(DSY(I,I,J)-
1 CSY(I,I-1,J))/DX
0211 DBL=BR(I,I,J)*.5*(GV(I,I,J)+GV(I,I-1,J-1))*DT*(DSY(I,I,J)-CSY(I,I-1,J))
1 *.6.2832
0212 GO TO 8404
0213 8406 IF (DR(I,I,J-1).EQ.0.)AND.(DSY(I,I,J).GT.CSY(I,I-1,J).AND.(GV(I,I,J)+
1 GV(I,I-1,J-1)).GT.0.) GO TO 8407
0214 GO TO 8408
0215 8407 DMB=SR(I,J)*DR(I,I-1,J-1)*.5*(GV(I,I,J)+GV(I,I-1,J-1))*DT*
1 (DSY(I,I-1,J-1)-CSY(I,I-1,J-1))/DX
0216 DBL=BR(I,I-1,J-1)*.5*(GV(I,I,J)+GV(I,I-1,J-1))*DT*(DSY(I,I-1,J-1)-
1 CSY(I,I-1,J-1))/DX
0217 TTD=1
0218 8408 CONTINUE
0219 IF (DR(I,I-1,J-1).EQ.0.) GO TO 8315
0220 IF (DR(I,I,J-1).EQ.0.)AND.(DSY(I,I,J).GT.CSY(I,I-1,J).AND.
1 (GV(I,I,J)+GV(I,I-1,J-1)).LT.0.) GO TO 8312
0221 GO TO 8315
0222 8312 DMB=SR(I,J)*DR(I,I-1,J-1)*.5*(GV(I,I,J)+GV(I,I-1,J-1))*DT*
1 (DSY(I,I-1,J-1)-CSY(I,I-1,J-1))/DX
0223 DBA=BR(I,I,J)*.5*(GV(I,I,J)+GV(I,I-1,J-1))*DT*(DSY(I,I,J-1)-
1 CSY(I,I-1,J-1))/DX
0224 TTD=1
0225 GO TO 8315
0226 8313 IF (DR(I,I-1,J-1).EQ.0.)AND.(DSY(I,I-1,J-1).GT.CSY(I,I-1,J-1).AND.
1 (GV(I,I,J)+GV(I,I-1,J-1)).GT.0.) GO TO 8314
0227 GO TO 8315
0228 8314 DMB=SR(I,J)*DR(I,I,J)*.5*(GV(I,I,J)+GV(I,I-1,J-1))*DT*(DSY(I,I,J-1)-
1 CSY(I,I-1,J-1))/DX
0229 DBA=BR(I,I,J)*.5*(GV(I,I,J)+GV(I,I-1,J-1))*DT*(DSY(I,I,J-1)-
1 CSY(I,I-1,J-1))/DX
0230 8315 CONTINUE
0231 IF (DR(I,I-1,J-1).EQ.0.) GO TO 8319
0232 IF (DR(I,I,J-1).EQ.0.)AND.(DSY(I,I,J).GT.CSY(I,I-1,J).AND.
1 (GV(I,I,J)+GV(I,I-1,J-1)).LT.0.) GO TO 8316
0233 GO TO 8317
0234 8316 DMB=SR(I,J)*DR(I,I-1,J-1)*.5*(GV(I,I,J)+GV(I,I-1,J-1))*DT*
1 (DSY(I,I-1,J-1)-CSY(I,I-1,J-1))/DX
0235 DBA=BR(I,I-1,J-1)*.5*(GV(I,I,J)+GV(I,I-1,J-1))*DT*(DSY(I,I-1,J-1)-
1 CSY(I,I-1,J-1))/DX

```

```
0236 TTR=1.  
0237 GO TO 8319  
0238 8317 IF (CR(I,I-1,J),EQ.0.,AND.DSY(I-1,J-1),GT. CSY(I,J-1),AND.  
1 (GV(I,I-1,J)+GV(I,I-1,J-1)),GT.0.) GO TO 8318  
0239 GO TO 8319  
0240 8318 DMG=SYA(I-1,J-1)*DBR(I,I-1,J)*.5*(GV(I,I-1,J)+GV(I,I-1,J-1))*DT*  
1 (DSY(I-1,J-1)-CSY(I,J-1))/DX  
0241 DBR=DM(I,I-1,J)*.5*(GV(I,I-1,J)+GV(I,I-1,J-1))*DT*(DSY(I-1,J-1)-  
1 CSY(I-1,J-1))*6.2832  
0242 8319 CONTINUE  
0243 IF (CR(I,I-1,J-1),EQ.0.) GO TO 8324  
0244 IF (CR(I,I-1,J),EQ.0.,AND.DSX(I-1,J),GT.CSX(I,J-1),AND.(GV(I,I-1,J)+  
1 GV(I,I-1,J-1)),LT.0.) GO TO 8321  
0245 GO TO 8322  
0246 8321 DMG=SYA(I-1,J)*DBR(I,I-1,J)*.5*(GV(I,I-1,J)+GV(I,I-1,J-1))*DT*(DSX(I,J)-  
1 CSX(I-1,J-1))/DX  
0247 DBR=SYA(I-1,J)*DBR(I,I-1,J)*.5*(GV(I,I-1,J)+GV(I,I-1,J-1))*DT*(DSX(I,J)-  
1 CSX(I-1,J-1))*6.2832/SYR(J)  
0248 GO TO 8324  
0249 8322 IF (CR(I,I-1,J),EQ.0.,AND.DSX(I-1,J),GT.CSX(I,J-1),AND.(GV(I,I-1,J)+  
1 GV(I,I-1,J-1)),GT.0.) GO TO 8323  
0250 GO TO 8324  
0251 8323 DMG=SYA(I-1,J)*DBR(I,I-1,J)*.5*(GV(I,I-1,J)+GV(I,I-1,J-1))*DT*  
1 (DSX(I,J)-CSX(I,J-1))/DX  
0252 DBR=SYA(I-1,J)*DBR(I,I-1,J)*.5*(GV(I,I-1,J)+GV(I,I-1,J-1))*DT*  
1 (DSX(I,J)-CSX(I,J-1))*6.2832/SYR(J-1)  
0253 TTR=1.  
0254 8324 CONTINUE  
0255 IF (CR(I,I-1,J-1),EQ.0.) GO TO 8328  
0256 IF (CR(I,I-1,J),EQ.0.,AND.DSX(I-1,J),GT.CSX(I,J-1),AND.(GV(I,I-1,J)+  
1 GV(I,I-1,J-1)),LT.0.) GO TO 8325  
0257 GO TO 8326  
0258 8325 DMG=SYA(I-1,J)*DBR(I,I-1,J)*.5*(GV(I,I-1,J)+GV(I,I-1,J-1))*DT*(DSX(I,J)-  
1 CSX(I-1,J-1))/DX  
0259 DBR=SYA(I-1,J)*DBR(I,I-1,J)*.5*(GV(I,I-1,J)+GV(I,I-1,J-1))*DT*(DSX(I,J)-  
1 CSX(I-1,J-1))*6.2832/SYR(J)  
0260 GO TO 8328  
0261 8326 IF (CR(I,I-1,J),EQ.0.,AND.DSX(I-1,J),GT.CSX(I,J-1),AND.(GV(I,I-1,J)+  
1 GV(I,I-1,J-1)),GT.0.) GO TO 8327  
0262 GO TO 8328  
0263 8327 DMG=SYA(I-1,J)*DBR(I,I-1,J)*.5*(GV(I,I-1,J)+GV(I,I-1,J-1))*DT*  
1 (DSX(I,J)-CSX(I,J-1))/DX  
0264 DBR=SYA(I-1,J)*DBR(I,I-1,J)*.5*(GV(I,I-1,J)+GV(I,I-1,J-1))*DT*  
1 (DSX(I,J)-CSX(I,J-1))*6.2832/SYR(J-1)  
0265 TTR=1.  
0266 8328 CONTINUE  
0267 IF (CR(I,I-1,J-1),EQ.0.) GO TO 8334  
0268 IF (CR(I,I-1,J),EQ.0.,AND.DSX(I-1,J),GT.CSX(I-1,J),AND.  
1 (GV(I,I-1,J)+GV(I,I-1,J-1)),LT.0.) GO TO 8331  
0269 GO TO 8332  
0270 8331 DMG=SYA(I-1,J)*DBR(I,I-1,J)*.5*(GV(I,I-1,J)+GV(I,I-1,J-1))*DT*  
1 (DSX(I-1,J)-CSX(I-1,J))/DX  
0271 DBR=SYA(I-1,J)*DBR(I,I-1,J)*.5*(GV(I,I-1,J)+GV(I,I-1,J-1))*DT*  
1 (DSX(I-1,J)-CSX(I-1,J))*6.2832/SYR(J-1)  
0272 TTR=1.  
0273 GO TO 8334
```

```
0274 8332 IF (CR(I,I-1,J-1),EQ.0.,AND.DSX(I-1,J-1),GT.CSX(I-1,J),AND.  
1 (GV(I,I-1,J)+GV(I,I-1,J-1)),GT.0.) GO TO 8333  
0275 GO TO 8334  
0276 8333 DMG=SYA(I-1,J)*DBR(I,I-1,J)*.5*(GV(I,I-1,J)+GV(I,I-1,J-1))*DT*  
1 (DSX(I-1,J-1)-CSX(I-1,J-1))/DX  
0277 DBR=SYA(I-1,J)*DBR(I,I-1,J)*.5*(GV(I,I-1,J)+GV(I,I-1,J-1))*DT*  
1 (DSX(I-1,J-1)-CSX(I-1,J-1))*6.2832/SYR(J)  
0278 8334 CONTINUE  
0279 IF (CR(I,I-1,J-1),EQ.0.) GO TO 8338  
0280 IF (CR(I,I-1,J),EQ.0.,AND.DSX(I-1,J),GT.CSX(I-1,J),AND.  
1 (GV(I,I-1,J)+GV(I,I-1,J-1)),LT.0.) GO TO 8335  
0281 GO TO 8336  
0282 8335 DMG=SYA(I-1,J)*DBR(I,I-1,J)*.5*(GV(I,I-1,J)+GV(I,I-1,J-1))*DT*  
1 (DSX(I-1,J-1)-CSX(I-1,J-1))/DX  
0283 DBR=SYA(I-1,J)*DBR(I,I-1,J)*.5*(GV(I,I-1,J)+GV(I,I-1,J-1))*DT*  
1 (DSX(I-1,J-1)-CSX(I-1,J-1))*6.2832/SYR(J-1)  
0284 TTR=1.  
0285 GO TO 8338  
0286 8336 IF (CR(I,I-1,J-1),EQ.0.,AND.DSX(I-1,J-1),GT.CSX(I-1,J),AND.  
1 (GV(I,I-1,J)+GV(I,I-1,J-1)),GT.0.) GO TO 8337  
0287 GO TO 8338  
0288 8337 DMG=SYA(I-1,J)*DBR(I,I-1,J)*.5*(GV(I,I-1,J)+GV(I,I-1,J-1))*DT*  
1 (DSX(I-1,J-1)-CSX(I-1,J-1))/DX  
0289 DBR=SYA(I-1,J)*DBR(I,I-1,J)*.5*(GV(I,I-1,J)+GV(I,I-1,J-1))*DT*  
1 (DSX(I-1,J-1)-CSX(I-1,J-1))*6.2832/SYR(J)  
0290 8338 CONTINUE  
0291 IF (I.LE.1M) GO TO 654  
0292 IF (I.GT. JM) GO TO 653  
0293 IF (DSY(I-1,J),GT.CSY(I-1,J),AND.CSY(I-1,J),GT.0.,AND.DR(I,I-1,J-1)  
1 ,GT.0.,AND.DSX(I-1,J-1),EQ.0.,AND.(GV(I,I-1,J)+GV(I,I-1,J-1)),GT.0.)  
2 GO TO 7441  
0294 GO TO 7442  
0295 7441 DMG=SYA(I-1,J)*DBR(I,I-1,J)*.5*(GV(I,I-1,J)+GV(I,I-1,J-1))*DT*  
1 (DSY(I-1,J)-CSY(I-1,J))/DX  
0296 DBR=SYA(I-1,J)*DBR(I,I-1,J)*.5*(GV(I,I-1,J)+GV(I,I-1,J-1))*DT*  
1 (DSY(I-1,J)-CSY(I-1,J))*6.2832/SYR(J-1)  
0297 TTR=1.  
0298 GO TO 7444  
0299 7442 IF (DSY(I-1,J),GT.CSY(I-1,J),AND.CSY(I-1,J),GT.0.,AND.DR(I,I-1,J-1)  
1 ,GT.0.,AND.DSX(I-1,J-1),EQ.0.,AND.(GV(I,I-1,J)+GV(I,I-1,J-1)),GT.0.)  
2 GO TO 7443  
0300 GO TO 7444  
0301 7443 DMG=SYA(I-1,J)*DBR(I,I-1,J)*.5*(GV(I,I-1,J)+GV(I,I-1,J-1))*DT*(DSY(I-1,J)  
1 -CSY(I-1,J))/DX  
0302 DBR=SYA(I-1,J)*DBR(I,I-1,J)*.5*(GV(I,I-1,J)+GV(I,I-1,J-1))*DT*(DSY(I-1,J)  
1 -CSY(I-1,J))*6.2832/SYR(J)  
0303 7444 CONTINUE  
0304 IF (DSY(I-1,J-1),GT.CSY(I-1,J-1),AND.CSY(I-1,J-1),GT.0.,AND.  
1 DR(I,I-1,J),GT.0.,AND.DSX(I-1,J),EQ.0.,AND.(GV(I,I-1,J)+  
2 GV(I,I-1,J-1)),GT.0.) GO TO 7445  
0305 GO TO 7446  
0306 7445 DMG=SYA(I-1,J)*DBR(I,I-1,J)*.5*(GV(I,I-1,J)+GV(I,I-1,J-1))*DT*(DSY(I-1,J-1)  
1 -CSY(I-1,J-1))/DX  
0307 DBR=SYA(I-1,J)*DBR(I,I-1,J)*.5*(GV(I,I-1,J)+GV(I,I-1,J-1))*DT*(DSY(I-1,J-1)  
1 -CSY(I-1,J-1))*6.2832/SYR(J-1)  
0308 GO TO 7446
```

```
0109 7446 IF (DSY(I-1,J-1).GT.CSY(I-1,J).AND.CSY(I-1,J).GT.O..AND.  
1 DR(I-1,J-1).GT.O..AND.ZSX(I-1,J).EQ.O..AND.(GV(I-1,J)+GV(I-1,J-1)).GT.O.)  
2 GUL(I-1,J-1).LT.O.) GO TO 7447  
0310 GO TO 7448  
0311 7447 DBA=SR(I-1,J-1)*DB(I-1,J-1)*.5*(GV(I-1,J)+GV(I-1,J-1))*DT*  
1 (DSY(I-1,J-1)-CSY(I-1,J-1))/DX  
DBA=SR(I-1,J)*DB(I-1,J-1)*.5*(GV(I-1,J)+GV(I-1,J-1))*DT*  
0312 1 (DSY(I-1,J-1)-CSY(I-1,J-1))*6.2832/SXB(J-1)  
TTA=1.  
0313 7448 CONTINUE  
0314 IF (CSX(I-1,J).GT.DSX(I-1,J).AND.DSX(I-1,J).GT.O..AND.DR(I-1,J-1)  
1.GT.O..AND.RSY(I-1,J).EQ.O..AND.(GV(I-1,J)+GV(I-1,J-1)).GT.O.)  
2GO TC 7451  
0316 GO TC 7452  
0317 7451 DBA=SYA(J-1)*DR(I-1,J-1)*.5*(GV(I-1,J)+GV(I-1,J-1))*DT*(CSX(I-1,J)-  
1DSX(I-1,J-1))/DX  
0318 DBA=SYA(J-1)*DB(I-1,J-1)*.5*(GV(I-1,J)+GV(I-1,J-1))*DT*(CSX(I-1,J)-  
1DSX(I-1,J-1))*6.2832/SXB(J)  
0319 7452 IF (CSX(I-1,J).GT.DSX(I-1,J).AND.DSX(I-1,J).GT.O..AND.DR(I-1,J-1)  
1.GT.O..AND.RSY(I-1,J).EQ.O..AND.(GV(I-1,J)+GV(I-1,J-1)).GT.O.)  
2GO TC 7453  
0321 GO TC 7454  
0322 7453 DBA=SYA(J-1)*DR(I-1,J-1)*.5*(GV(I-1,J)+GV(I-1,J-1))*DT*  
1(CSX(I-1,J)-DSX(I-1,J-1))/DX  
0323 DBA=SYA(J-1)*DB(I-1,J-1)*.5*(GV(I-1,J)+GV(I-1,J-1))*DT*  
1(CSX(I-1,J)-DSX(I-1,J-1))*6.2832/SXB(J-1)  
0324 TTA=1.  
0325 7454 CONTINUE  
0326 IF (CSX(I-1,J).LT.CSX(I-1,J-1).AND.DSX(I-1,J).GT.O..AND.  
1DR(I-1,J-1).GT.O..AND.RSY(I-1,J).EQ.O..AND.(GV(I-1,J)+  
2GV(I-1,J-1)).GT.O.) GO TO 7461  
0327 GO TC 7462  
0328 7461 DBA=SYA(J-1)*DR(I-1,J-1)*.5*(GV(I-1,J)+GV(I-1,J-1))*DT*  
1(CSX(I-1,J-1)-DSX(I-1,J-1))/DX  
0329 DBA=SYA(J-1)*DB(I-1,J-1)*.5*(GV(I-1,J)+GV(I-1,J-1))*DT*  
1(CSX(I-1,J-1)-DSX(I-1,J-1))*6.2832/SXB(J-1)  
0330 TTA=1.  
0331 GO TC 7464  
0332 7462 IF (DSX(I-1,J).LT.CSX(I-1,J-1).AND.DSX(I-1,J).GT.O..AND.  
1DR(I-1,J-1).GT.O..AND.RSY(I-1,J).EQ.O..AND.(GV(I-1,J)+  
2GV(I-1,J-1)).GT.O.) GO TO 7463  
0333 GO TO 7464  
0334 7463 DBA=SYA(J-1)*DR(I-1,J-1)*.5*(GV(I-1,J)+GV(I-1,J-1))*DT*  
1(CSX(I-1,J-1)-DSX(I-1,J-1))/DX  
0335 DBA=SYA(J-1)*DB(I-1,J-1)*.5*(GV(I-1,J)+GV(I-1,J-1))*DT*  
1(CSX(I-1,J-1)-DSX(I-1,J-1))*6.2832/SXB(J)  
0336 7464 CONTINUE  
0337 GO TC 654  
0338 653 CONTINUE  
0339 1E.DSY(I-1,J).GT.CSY(I-1,J).AND.CSY(I-1,J).GT.O..AND.DR(I-1,J-1)  
1.GT.O..AND.ZSX(I-1,J).EQ.O..AND.(GV(I-1,J)+GV(I-1,J-1)).GT.O.)  
2 GO TO 6441  
0340 GO TC 6442  
0341 6441 DBA=SR(I-1,J)*DB(I-1,J-1)*.5*(GV(I-1,J)+GV(I-1,J-1))*DT*  
1 (DSY(I-1,J)-CSY(I-1,J-1))/DX
```

```
0342 DBA=SR(I-1,J)*DB(I-1,J-1)*.5*(GV(I-1,J)+GV(I-1,J-1))*DT*  
1 (DSY(I-1,J)-CSY(I-1,J-1))*6.2832/SXB(J-1)  
0343 TTA=1.  
0344 GO TO 6444  
0345 6442 IF (DSY(I-1,J).GT.CSY(I-1,J).AND.CSY(I-1,J).GT.O..AND.DR(I-1,J-1)  
1.GT.O..AND.ZSX(I-1,J).EQ.O..AND.(GV(I-1,J)+GV(I-1,J-1)).GT.O.)  
2 GO TO 6443  
0346 GO TO 6444  
0347 6443 DBA=SR(I-1,J)*DB(I-1,J-1)*.5*(GV(I-1,J)+GV(I-1,J-1))*DT*(DSY(I-1,J)  
1-CSY(I-1,J-1))/DX  
0348 DBA=SR(I-1,J)*DB(I-1,J-1)*.5*(GV(I-1,J)+GV(I-1,J-1))*DT*(DSY(I-1,J)  
1-CSY(I-1,J-1))*6.2832/SXB(J)  
0349 6444 CONTINUE  
0350 IF (CSY(I-1,J-1).GT.CSY(I-1,J-1).AND.CSY(I-1,J-1).GT.O..AND.  
1DR(I-1,J-1).GT.O..AND.ZSX(I-1,J).EQ.O..AND.(GV(I-1,J)+  
2GV(I-1,J-1)).GT.O.) GO TO 6445  
0351 GO TC 6446  
0352 6445 DBA=SR(I-1,J)*DB(I-1,J-1)*.5*(GV(I-1,J)+GV(I-1,J-1))*DT*(DSY(I-1,J-1)  
1-CSY(I-1,J-1))/DX  
0353 DBA=SR(I-1,J)*DB(I-1,J-1)*.5*(GV(I-1,J)+GV(I-1,J-1))*DT*(DSY(I-1,J-1)  
1-CSY(I-1,J-1))*6.2832/SXB(J)  
0354 GO TO 6449  
0355 6446 IF (CSY(I-1,J-1).GT.CSY(I-1,J-1).AND.CSY(I-1,J-1).GT.O..AND.  
1DR(I-1,J-1).GT.O..AND.ZSX(I-1,J).EQ.O..AND.(GV(I-1,J)+  
2GV(I-1,J-1)).GT.O.) GO TO 6447  
0356 GO TO 6449  
0357 6447 DBA=SR(I-1,J)*DB(I-1,J-1)*.5*(GV(I-1,J)+GV(I-1,J-1))*DT*  
1 (DSY(I-1,J-1)-CSY(I-1,J-1))/DX  
0358 DBA=SR(I-1,J)*DB(I-1,J-1)*.5*(GV(I-1,J)+GV(I-1,J-1))*DT*  
1 (DSY(I-1,J-1)-CSY(I-1,J-1))*6.2832/SXB(J-1)  
0359 TTA=1.  
0360 6448 CONTINUE  
0361 IF (DSX(I-1,J).GT.CSX(I-1,J-1).AND.CSX(I-1,J-1).GT.O..AND.DR(I-1,J-1)  
1.GT.O..AND.RSY(I-1,J).EQ.O..AND.(GV(I-1,J)+GV(I-1,J-1)).GT.O.)  
2GO TC 6451  
0362 GO TC 6452  
0363 6451 DBA=SYA(J-1)*DR(I-1,J-1)*.5*(GV(I-1,J)+GV(I-1,J-1))*DT*(DSX(I-1,J)-  
1DSX(I-1,J-1))/DX  
0364 DBA=SYA(J-1)*DB(I-1,J-1)*.5*(GV(I-1,J)+GV(I-1,J-1))*DT*(DSX(I-1,J)-  
1DSX(I-1,J-1))*6.2832/SXB(J)  
0365 GO TO 6454  
0366 6452 IF (DSX(I-1,J).GT.CSX(I-1,J-1).AND.CSX(I-1,J-1).GT.O..AND.DR(I-1,J-1)  
1.GT.O..AND.RSY(I-1,J).EQ.O..AND.(GV(I-1,J)+GV(I-1,J-1)).GT.O.)  
2GO TC 6453  
0367 GO TC 6454  
0368 6453 DBA=SYA(J-1)*DR(I-1,J-1)*.5*(GV(I-1,J)+GV(I-1,J-1))*DT*  
1(DSX(I-1,J)-DSX(I-1,J-1))/DX  
0369 DBA=SYA(J-1)*DB(I-1,J-1)*.5*(GV(I-1,J)+GV(I-1,J-1))*DT*  
1(DSX(I-1,J)-DSX(I-1,J-1))*6.2832/SXB(J-1)  
0370 TTA=1.  
0371 6454 CONTINUE  
0372 IF (CSX(I-1,J).LT.DSX(I-1,J-1).AND.CSX(I-1,J).GT.O..AND.  
1DR(I-1,J-1).GT.O..AND.RSY(I-1,J).EQ.O..AND.(GV(I-1,J)+  
2GV(I-1,J-1)).GT.O.) GO TO 6461  
0373 GO TC 6462  
0374 6461 DBA=SYA(J-1)*DR(I-1,J-1)*.5*(GV(I-1,J)+GV(I-1,J-1))*DT*
```

```

0375      I1DSX(I-1,J)-CSX(I-1,J)/DX
0376      DUC=SYA(I,J)*PP(I-1,J)+.5*(GV(I-1,J)+GV(I-1,J+1))*DT*
0377      I1DSX(I-1,J)-CSX(I-1,J))*6.2832/SXB(I,J)
0378      TTC=1.
0379      GO TO 6464
0380      6462 IF (CSX(I-1,J).LT.DSX(I-1,J).AND.CSX(I-1,J).GT.0..AND.
0381      I1R(I-1,J+1).GT.0..AND.RSY(I-1,J).EQ.0..AND.(GV(I-1,J)+
0382      2GV(I-1,J+1)).GT.0.) GO TO 6463
0383      GO TO 6464
0384      6463 DMC=SYA(I,J)*DRI(I-1,J)+.5*(GV(I-1,J)+GV(I-1,J+1))*DT*
0385      I1DSX(I-1,J)-CSX(I-1,J)/DX
0386      DUC=SYA(I,J)*PP(I-1,J)+.5*(GV(I-1,J)+GV(I-1,J+1))*DT*
0387      I1DSX(I-1,J)-CSX(I-1,J))*6.2832/SXB(I,J)
0388      6464 CONTINUE
0389      654 CONTINUE
0390      39 CONTINUE
0391      DRI(2,I,J)=DRI(1,I,J)*(DMA+DM2+DM3+DM4+DM5+DML+DM3-DPI-DMC-DMA-
0392      1DMR-CMT)/(VOL(I,J)*SF(I,J))
0393      DRI(2,I,J)=DRI(1,I,J)*(DM4+DM2+DM3+DM4+DM5+DML+DM3-DPI-DMC-DMA-
0394      1DMR-CMT)/(VOL(I,J)*SF(I,J)*SYB(I,J)/6.2832*DX)
0395      0397 FUI(2,I,J)=(TT2*DM2*GV(I-1,I-1,J)+TT4*DM4*GV(I-1,I-1,J)-TT1*DM1*GV(I-
0396      1,I-1,J)-TT3*DM3*GV(I-1,I-1,J)+GV(I-1,I-1,J+1))*(TT4*DM4-TTC*DMC)
0397      2*GV(I-1,I-1,J)-1)*(TT2*DM2-TTA*DMA)-GV(I-1,I-1,J+1)*(TT4*DM4+TTN*DMR)
0398      3*GV(I-1,I-1,J)-1)*(TT2*DM2-TTA*DMA)-GV(I-1,I-1,J+1)*VOL(I-1,J)+
0399      4SF(I-1,I-1,J)-TT4*DM4*(1-TT2)*DM2*(1-TT3)*DM4*(1-TT1)*DM1*(1-TT3)*
0400      5DMC*(1-TT1)*DM1*(1-TT3)*DM3*(1-TT1)*DM3*(1-TT1)*DMA*(1-TTC)*
0401      6DMC*(1-TT1)*DMT*(1-TT3)*DMR)/(VOL(I,J)*DRI(2,I,J)*SF(I,J))
0402      0398 FV(2,I,J)=(TT2*DM2*GV(I-1,I-1,J)+TT4*DM4*GV(I-1,I-1,J)-TT1*DM1*GV(I-
0403      1,I-1,J)-TT3*DM3*GV(I-1,I-1,J)+GV(I-1,I-1,J+1))*(TT4*DM4-TTC*DMC)
0404      2*GV(I-1,I-1,J)-1)*(TT2*DM2-TTA*DMA)-GV(I-1,I-1,J+1)*(TT4*DM4+TTN*DMR)
0405      3*GV(I-1,I-1,J)-1)*(TT2*DM2-TTA*DMA)-GV(I-1,I-1,J+1)*VOL(I-1,J)+
0406      4SF(I-1,I-1,J)-TT4*DM4*(1-TT2)*DM2*(1-TT3)*DM4*(1-TT1)*DM1*(1-TT3)*
0407      5DMC*(1-TT1)*DM1*(1-TT3)*DM3*(1-TT1)*DM3*(1-TT1)*DMA*(1-TTC)*
0408      6DMC*(1-TT1)*DMT*(1-TT3)*DMR)/(VOL(I,J)*DRI(2,I,J)*SF(I,J))
0409      FE(2,I,J)=(TT2*DM2*GV(I-1,I-1,J)+TT4*DM4*GV(I-1,I-1,J)-TT1*DM1*GV(I-
0410      1,I-1,J)-TT3*DM3*GV(I-1,I-1,J)+GV(I-1,I-1,J+1))*(TT4*DM4-TTC*DMC)
0411      2*GV(I-1,I-1,J)-1)*(TT2*DM2-TTA*DMA)-GV(I-1,I-1,J+1)*(TT4*DM4+TTN*DMR)
0412      3*GV(I-1,I-1,J)-1)*(TT2*DM2-TTA*DMA)-GV(I-1,I-1,J+1)*VOL(I-1,J)+
0413      4SF(I-1,I-1,J)-TT4*DM4*(1-TT2)*DM2*(1-TT3)*DM4*(1-TT1)*DM1*(1-TT3)*
0414      5DMC*(1-TT1)*DM1*(1-TT3)*DM3*(1-TT1)*DM3*(1-TT1)*DMA*(1-TTC)*
0415      6DMC*(1-TT1)*DMT*(1-TT3)*DMR)/(VOL(I,J)*DRI(2,I,J)*SF(I,J))
0416      IF (ABS(FUI(2,I,J))-LE.EPS) FUI(2,I,J)=0.
0417      IF (ABS(FV(2,I,J))-LE.EPS) FV(2,I,J)=0
0418      FV2=FV(2,I,J)*2
0419      GO TO 1131
0420      1130 FV2=C.
0421      0395 1131 FUI(2,I,J)=FUI(2,I,J)+.5*(FUI(2,I,J)+FV2)
0422      0396 1E=(DP(I-1,I-1,J)+EQ.0..AND.DP(I-1,I-1,J).EQ.0..AND.DRI(1,I-1,J).EQ.
0423      10..AND.DRI(1,I-1,J).EQ.0.) CALL VACUUMIRMEAN,DX)
0424      0397 25 CONTINUE
0425      0398 RETURN
0426      0399 END

```

```

0001      SUBROUTINE PLOTS(I8,IE,J8,JE,DX,DT,TIME,N,I,J,JD,TD)
0002      DIMENSION DRI(2,137,22),FUI(2,137,22),FV(2,137,22),F(2,137,22),
0003      1FE(2,137,22),TC(2,137,22),PP(2,137,22),P(2,137,22),R(2,137,22),
0004      2GV(1,137,22),GV(1,137,22),G(1,137,22),GE(1,137,22),TMP(2,137,22),
0005      3SYA(22),SXB(22),SYB(22),VOL(137,22),SI(22),CPB(22),RV(137,22),
0006      4CSX(137,22),DSX(137,22),CSY(137,22),DSY(137,22),RSY(137,22),
0007      5ZSX(137,22)
0008      DIMENSION DRI(2,137,22),PDR(2,137,22),PTMP(2,137,22)
0009      COMMON OR,FU,EV,FI,FE,TC,PP,P,BG,GU,GV,GI,GE,TMP,SYA,SXB,SYB,VOL,
0010      1SI,CPB,BV,CSX,DSX,CSY,DSY,RSY, ZSX,CL,02,03,04,N,I,J
0011      SCR=1.E+04
0012      TOR=1.OF+09
0013      IRR=IR-1
0014      IF (N.EQ.1) NN=2
0015      IF (N.EQ.2) NN=1
0016      DO 77 J=JB,JE
0017      DO 77 I=IRN,IE
0018      IF (FV(I,N,I,J).NE.0..AND..OR(I,I,J).EQ.0.) FV(I,N,I,J)=0.
0019      IF (FV(I,N,I,J).NE.0..AND..DRI(1,I,J).EQ.0.) FV(I,N,I,J)=0.
0020      PDR(I,N,I,J)=DRI(I,N,I,J)/TOR
0021      PTMP(I,N,I,J)=TMP(I,N,I,J)/1000000.
0022      PDR(I,N,I,J)=FV(I,N,I,J)/TOR
0023      77 CONTINUE
0024      WRITE (6,64)
0025      65 FORMAT (1H0,'DISTRIBUTION OF TEMPERATURE IN SHOCK TUBE')
0026      DO 64 I=IRN,IE
0027      WRITE (6,63) (PTMP(I,N,I,J),J=JB,JE)
0028      64 CONTINUE
0029      63 FORMAT (1H0,'33F4.2)
0030      WRITE (6,65)
0031      65 FORMAT (1H0,'DENSITY VALUES')
0032      DO 65 I=IRN,IE
0033      WRITE (6,63) (PDR(I,N,I,J),J=JB,JE)
0034      65 CONTINUE
0035      63 CONTINUE
0036      WRITE (6,66)
0037      66 FORMAT (1H0,'U-COMPONENT')
0038      DO 66 I=IRN,IE
0039      WRITE (6,53) (PDR(I,N,I,J),J=JB,JE)
0040      66 CONTINUE
0041      WRITE (6,67)
0042      67 FORMAT (1H0,'V-COMPONENT')
0043      DO 67 I=IRN,IE
0044      WRITE (6,53) (PDR(I,N,I,J),J=JB,JE)
0045      67 CONTINUE
0046      53 FORMAT (1H0,'26F5.2)
0047      RETURN
0048      END

```

```

0001 SURROUTNF VISCOS ECR,CK,MP,FUR,FUL,GUP,GUL,FVT,FVB,GVT,GVB,DT,DX,
      11P,J1
0002 DIMENSION DP(2,137,22),FUI(2,137,22),FV(2,137,22),F(12,137,22),
      1FF(2,137,22),TC(2,137,22),PP(2,137,22),P(2,137,22),RR(2,137,22),
      2GU(1,137,22),GV(1,137,22),GT(1,137,22),GE(1,137,22),TMP(2,137,22),
      3SYA(22),SX(22),SYR(22),VOL(137,22),SI(22),CPR(22),HVL(137,22),
      4CSX(137,22),DSX(137,22),CSY(137,22),DSY(137,22),RSY(137,22),
      5ZSX(137,22)
0003 COMMON DP,FU,FV,FI,FE,TC,PP,P,RR,GU,GV,GT,GE,TMP,SYA,SXR,SYR,VOL,
      1SI,CPR,BV,CSX,DSX,CSY,DSY,RSY,ZSX,Q1,Q2,Q3,Q4,M,I,J
0004 C1=0.
0005 C2=0.
0006 C3=0.
0007 C4=0.
0008 DRR=CR(1,1,J)
0009 DPL=CR(1,1,J)
0010 DRT=DR(1,1,J)
0011 DRB=DR(1,1,J-1)
0012 IF(DR(1,1,J).EQ.0.) GO TO 1
0013 GO TC 2
0014 1 CONTINUE
0015 FUI(1,1,J)=FUI(1,1,J)
0016 FV(1,1,J)=FV(1,1,J)
0017 GUL(1,1,J)=GUL(1,1,J)
0018 TC(1,1,J)=TC(1,1,J)
0019 DRR=DR(1,1,J)
0020 2 CONTINUE
0021 IF(DR(1,1,J).EQ.0.) GO TO 3
0022 GO TC 4
0023 3 CONTINUE
0024 FUI(1,1,J)=FUI(1,1,J)
0025 FV(1,1,J)=FV(1,1,J)
0026 GUL(1,1,J)=GUL(1,1,J)
0027 TC(1,1,J)=TC(1,1,J)
0028 DPL=DR(1,1,J)
0029 4 CONTINUE
0030 IF(DR(1,1,J-1).EQ.0.) GO TO 5
0031 GO TO 6
0032 5 CONTINUE
0033 FUI(1,1,J-1)=FUI(1,1,J)
0034 FV(1,1,J-1)=FV(1,1,J)
0035 GUL(1,1,J-1)=GUL(1,1,J)
0036 TC(1,1,J-1)=TC(1,1,J)
0037 DRT=DR(1,1,J)
0038 6 CONTINUE
0039 IF(DR(1,1,J-1).EQ.0.) GO TO 7
0040 GO TC 8
0041 7 CONTINUE
0042 FUI(1,1,J-1)=FUI(1,1,J)
0043 FV(1,1,J-1)=FV(1,1,J)
0044 GUL(1,1,J-1)=GUL(1,1,J)
0045 TC(1,1,J-1)=TC(1,1,J)
0046 DRB=CR(1,1,J)
0047 8 CONTINUE
0048 FUR=FUI(1,1,J)
0049 FUL=FUI(1,1,J)

```

```

0050 FUI=FUI(1,1,J-1)
0051 FVB=FUI(1,1,J-1)
0052 FVR=FV(1,1,J-1)
0053 FVL=FV(1,1,J-1)
0054 FVT=FV(1,1,J-1)
0055 FVB=FV(1,1,J-1)
0056 GUR=GU(1,1,J-1)
0057 GUL=GU(1,1,J-1)
0058 GVT=GV(1,1,J-1)
0059 GVB=GV(1,1,J-1)
0060 TCR=TC(1,1,J-1)
0061 TCL=TC(1,1,J-1)
0062 TCR=TC(1,1,J-1)
0063 TCR=TC(1,1,J-1)
0064 IF (RSY(1,1,J).EQ.RSY(1,1,J-1).AND.(RSY(1,1,J).EQ.RSY(1,1,J-1).AND.
      1ZSX(1,1,J).EQ.ZSX(1,1,J-1).AND.(ZSX(1,1,J).EQ.ZSX(1,1,J-1))) GO TO 59
0065 IF (CSY(1,1,J).GT.0. .AND. DR(1,1,J-1).EQ.0. .AND. RSY(1,1,J)
      1.GT. DSY(1,1,J)) GO TO 11
0066 IF (CSY(1,1,J).GT.0. .AND. DR(1,1,J-1).EQ.0. .AND. RSY(1,1,J)
      1.GT. DSY(1,1,J)) GO TO 11
0067 GO TO 41
0068 11 CONTINUE
0069 FUR=(FUI(1,1,J)+DSY(1,1,J)+FUI(1,1,J)+RSY(1,1,J)-DSY(1,1,J))
      1/RSY(1,1,J)
0070 FVR=(FV(1,1,J)+DSY(1,1,J)+FV(1,1,J)+RSY(1,1,J)-DSY(1,1,J))
      1/RSY(1,1,J)
0071 DRR=(DR(1,1,J)+DSY(1,1,J)+DR(1,1,J)+RSY(1,1,J)-DSY(1,1,J))
      1/RSY(1,1,J)
0072 TCR=(TC(1,1,J)+DSY(1,1,J)+TC(1,1,J)+RSY(1,1,J)-DSY(1,1,J))
      1/RSY(1,1,J)
0073 IF (PP.EQ.1) GO TO 41
0074 GUR=(GU(1,1,J)+DSY(1,1,J)+GU(1,1,J)+RSY(1,1,J)-DSY(1,1,J))
      1/RSY(1,1,J)
0075 41 CONTINUE
0076 IF (CR(1,1,J).EQ.0. .AND. DSY(1,1,J-1).GT. CSY(1,1,J) .AND.
      1DSY(1,1,J-1).GT. CSY(1,1,J)) GO TO 12
0077 GO TC 42
0078 12 CONTINUE
0079 FUR=(FUI(1,1,J-1)+DSY(1,1,J-1)+CSY(1,1,J-1)+FUI(1,1,J-1)+
      1DSY(1,1,J-1)+CSY(1,1,J-1)+FUI(1,1,J-1)+DSY(1,1,J-1)+CSY(1,1,J-1)+
      2J-1)+DSY(1,1,J-1)+DSY(1,1,J-1)+FUI(1,1,J-1)+DSY(1,1,J-1)+
      3FVR=(FV(1,1,J-1)+DSY(1,1,J-1)+GV(1,1,J-1)+FV(1,1,J-1)+DSY(1,1,J-1)+
      4DSY(1,1,J-1)+CSY(1,1,J-1)+FV(1,1,J-1)+DSY(1,1,J-1)+CSY(1,1,J-1)+
      52J-1)+DSY(1,1,J-1)+DSY(1,1,J-1)+FV(1,1,J-1)+DSY(1,1,J-1)+
      6DRR=(DR(1,1,J-1)+DSY(1,1,J-1)+CSY(1,1,J-1)+DR(1,1,J-1)+
      7DSY(1,1,J-1)+CSY(1,1,J-1)+DR(1,1,J-1)+DSY(1,1,J-1)+CSY(1,1,J-1)+
      82J-1)+DSY(1,1,J-1)+DSY(1,1,J-1)+FV(1,1,J-1)+DSY(1,1,J-1)+
      9TCR=(TC(1,1,J-1)+DSY(1,1,J-1)+CSY(1,1,J-1)+TC(1,1,J-1)+
      10DSY(1,1,J-1)+CSY(1,1,J-1)+TC(1,1,J-1)+DSY(1,1,J-1)+CSY(1,1,J-1)+
      112J-1)+DSY(1,1,J-1)+DSY(1,1,J-1)+FV(1,1,J-1)+DSY(1,1,J-1)+
      12IF (PP.EQ.1) GO TO 42
0084 GUR=(GU(1,1,J-1)+DSY(1,1,J-1)+CSY(1,1,J-1)+GU(1,1,J-1)+
      1DSY(1,1,J-1)+CSY(1,1,J-1)+GU(1,1,J-1)+DSY(1,1,J-1)+CSY(1,1,J-1)+
      2J-1)+DSY(1,1,J-1)+DSY(1,1,J-1)+FV(1,1,J-1)+DSY(1,1,J-1)+
      342 CONTINUE
0086 IF (CR(1,1,J).EQ.0. .AND. DSY(1,1,J-1).GT. CSY(1,1,J) .AND.

```



```
0151 DRT = (DRT(1,I,J+1)*DSX(I,I,J+1)+DR(1,I,J))/ZSX(I,I,J)-DSX(I,I,J+1)
0154 TCT = (TCT(1,I,J+1)*DSX(I,I,J+1)+TC(1,I,J))/ZSX(I,I,J)-DSX(I,I,J+1)
0155 IF (MP.EQ.1) GO TO 49
0156 GVT = (GVT(1,I,J+1)*DSX(I,I,J+1)+GV(1,I,J))/ZSX(I,I,J)-DSX(I,I,J+1)
0157 1/ZSX(I,I,J)
0157 49 CONTINUE
0159 IF (CR(1,I,J+1).EQ.0..AND.DSX(I,I,J+1).GT.CSX(I,I,J).AND.
0159 DSX(I-I,J+1).GT.CSX(I-I,J)) GO TO 21
0160 GO TO 50
0161 21 CONTINUE
0161 FUT = (FUT(1,I+1,J+1)*DSX(I+1,J+1)-CSX(I+1,J+1)+FUT(1,I,J+1)+
0161 1DSX(I-I,J+1)-CSX(I-I,J+1)+FUT(1,I,J+1)*ZSX(I,I,J)+CSX(I+1,J+1)+
0161 2CSX(I-I,J+1)-DSX(I-I,J+1)-DSX(I-I,J+1))/ZSX(I,I,J)
0162 FVT = (FVT(1,I+1,J+1)*DSX(I+1,J+1)-CSX(I+1,J+1)+FVT(1,I,J+1)+
0162 1DSX(I-I,J+1)-CSX(I-I,J+1)+FVT(1,I,J+1)*ZSX(I,I,J)+CSX(I+1,J+1)+
0162 2CSX(I-I,J+1)-DSX(I-I,J+1)-DSX(I-I,J+1))/ZSX(I,I,J)
0163 DRT = (DRT(1,I+1,J+1)*DSX(I+1,J+1)-CSX(I+1,J+1)+DRT(1,I,J+1)+
0163 1DSX(I-I,J+1)-CSX(I-I,J+1)+DRT(1,I,J+1)*ZSX(I,I,J)+CSX(I+1,J+1)+
0163 2CSX(I-I,J+1)-DSX(I-I,J+1)-DSX(I-I,J+1))/ZSX(I,I,J)
0164 TCT = (TCT(1,I+1,J+1)*DSX(I+1,J+1)-CSX(I+1,J+1)+TCT(1,I,J+1)+
0164 1DSX(I-I,J+1)-CSX(I-I,J+1)+TCT(1,I,J+1)*ZSX(I,I,J)+CSX(I+1,J+1)+
0164 2CSX(I-I,J+1)-DSX(I-I,J+1)-DSX(I-I,J+1))/ZSX(I,I,J)
0165 IF (MP.EQ.1) GO TO 50
0166 GVT = (GVT(1,I+1,J+1)*DSX(I+1,J+1)-CSX(I+1,J+1)+GVT(1,I,J+1)+
0166 1DSX(I-I,J+1)-CSX(I-I,J+1)+GVT(1,I,J+1)*ZSX(I,I,J)+CSX(I+1,J+1)+
0166 2CSX(I-I,J+1)-DSX(I-I,J+1)-DSX(I-I,J+1))/ZSX(I,I,J)
0167 50 CONTINUE
0168 IF (CR(1,I,J+1).EQ.0..AND.DSX(I+1,J+1).GT.CSX(I+1,J).AND.
0168 1DSX(I-I,J+1).LE.CSX(I-I,J)) GO TO 22
0169 GO TO 51
0170 22 CONTINUE
0171 FUT = (FUT(1,I+1,J+1)*DSX(I+1,J+1)-CSX(I+1,J+1)+FUT(1,I,J+1)+
0171 1ZSX(I,I,J)+CSX(I+1,J+1)-DSX(I+1,J+1))/ZSX(I,I,J)
0172 FVT = (FVT(1,I+1,J+1)*DSX(I+1,J+1)-CSX(I+1,J+1)+FVT(1,I,J+1)+
0172 1ZSX(I,I,J)+CSX(I+1,J+1)-DSX(I+1,J+1))/ZSX(I,I,J)
0173 DRT = (DRT(1,I+1,J+1)*DSX(I+1,J+1)-CSX(I+1,J+1)+DRT(1,I,J+1)+
0173 1ZSX(I,I,J)+CSX(I+1,J+1)-DSX(I+1,J+1))/ZSX(I,I,J)
0174 TCT = (TCT(1,I+1,J+1)*DSX(I+1,J+1)-CSX(I+1,J+1)+TCT(1,I,J+1)+
0174 1ZSX(I,I,J)+CSX(I+1,J+1)-DSX(I+1,J+1))/ZSX(I,I,J)
0175 IF (MP.EQ.1) GO TO 51
0176 GVT = (GVT(1,I+1,J+1)*DSX(I+1,J+1)-CSX(I+1,J+1)+GVT(1,I,J+1)+
0176 1ZSX(I,I,J)+CSX(I+1,J+1)-DSX(I+1,J+1))/ZSX(I,I,J)
0177 51 CONTINUE
0178 IF (CR(1,I,J+1).EQ.0..AND.DSX(I+1,J+1).LE.CSX(I+1,J).AND.
0178 1DSX(I-I,J+1).GT.CSX(I-I,J)) GO TO 23
0179 GO TO 52
0180 23 CONTINUE
0181 FUT = (FUT(1,I+1,J+1)*DSX(I+1,J+1)-CSX(I+1,J+1)+FUT(1,I,J+1)+
0181 1ZSX(I,I,J)+CSX(I+1,J+1)-DSX(I+1,J+1))/ZSX(I,I,J)
0182 FVT = (FVT(1,I+1,J+1)*DSX(I+1,J+1)-CSX(I+1,J+1)+FVT(1,I,J+1)+
0182 1ZSX(I,I,J)+CSX(I+1,J+1)-DSX(I+1,J+1))/ZSX(I,I,J)
0183 DRT = (DRT(1,I+1,J+1)*DSX(I+1,J+1)-CSX(I+1,J+1)+DRT(1,I,J+1)+
0183 1ZSX(I,I,J)+CSX(I+1,J+1)-DSX(I+1,J+1))/ZSX(I,I,J)
0184 TCT = (TCT(1,I+1,J+1)*DSX(I+1,J+1)-CSX(I+1,J+1)+TCT(1,I,J+1)+
0184 1ZSX(I,I,J)+CSX(I+1,J+1)-DSX(I+1,J+1))/ZSX(I,I,J)
```

```
0185 1ZSX(I,I,J)+CSX(I-I,J)-DSX(I-I,J+1))/ZSX(I,I,J)
0186 IF (MP.EQ.1) GO TO 52
0186 GVT = (GVT(1,I+1,J+1)*DSX(I+1,J+1)-CSX(I+1,J+1)+GVT(1,I,J+1)+
0186 1ZSX(I,I,J)+CSX(I+1,J+1)-DSX(I+1,J+1))/ZSX(I,I,J)
0187 52 CONTINUE
0188 IF (CSX(I,I,J).GT.0..AND.DR(1,I+1,J-1).EQ.0..AND.ZSX(I,I,J).GT.
0188 1CSX(I,I,J-1) GO TO 24
0189 IF (CSX(I,I,J).GT.0..AND.DR(1,I+1,J-1).EQ.0..AND.ZSX(I,I,J).GT.
0189 1CSX(I,I,J-1) GO TO 25
0190 GO TO 53
0191 24 CONTINUE
0192 FUB = (FUB(1,I,J-1)*CSX(I,I,J-1)+FUT(1,I,J-1)*ZSX(I,I,J)-CSX(I,I,J-1)
0192 1/ZSX(I,I,J)
0193 FVB = (FVB(1,I,J-1)*CSX(I,I,J-1)+FVT(1,I,J-1)*ZSX(I,I,J)-CSX(I,I,J-1)
0193 1/ZSX(I,I,J)
0194 DRB = (DRB(1,I,J-1)*CSX(I,I,J-1)+DRT(1,I,J-1)*ZSX(I,I,J)-CSX(I,I,J-1)
0194 1/ZSX(I,I,J)
0195 TCB = (TCB(1,I,J-1)*CSX(I,I,J-1)+TCT(1,I,J-1)*ZSX(I,I,J)-CSX(I,I,J-1)
0195 1/ZSX(I,I,J)
0196 IF (MP.EQ.1) GO TO 53
0197 GVB = (GVB(1,I,J-1)*CSX(I,I,J-1)+GVT(1,I,J-1)*ZSX(I,I,J)-CSX(I,I,J-1)
0197 1/ZSX(I,I,J)
0198 53 CONTINUE
0199 IF (CR(1,I,J-1).EQ.0..AND.CSX(I,I,J-1).GT.DSX(I+1,J).AND.CSX(I-I,J
0199 1-I).GT.DSX(I-I,J)) GO TO 26
0200 GO TO 54
0201 26 CONTINUE
0202 FUB = (FUB(1,I+1,J-1)*CSX(I+1,J-1)-DSX(I+1,J-1)+FUT(1,I+1,J-1)+
0202 1CSX(I-I,J-1)-DSX(I-I,J-1)+FUT(1,I,J-1)*ZSX(I,I,J)+DSX(I+1,J-1)+
0202 2DSX(I-I,J-1)-CSX(I-I,J-1)-CSX(I-I,J-1))/ZSX(I,I,J)
0203 FVB = (FVB(1,I+1,J-1)*CSX(I+1,J-1)-DSX(I+1,J-1)+FVT(1,I+1,J-1)+
0203 1CSX(I-I,J-1)-DSX(I-I,J-1)+FVT(1,I,J-1)*ZSX(I,I,J)+DSX(I+1,J-1)+
0203 2DSX(I-I,J-1)-CSX(I-I,J-1)-CSX(I-I,J-1))/ZSX(I,I,J)
0204 DRB = (DRB(1,I+1,J-1)*CSX(I+1,J-1)-DSX(I+1,J-1)+DRT(1,I+1,J-1)+
0204 1CSX(I-I,J-1)-DSX(I-I,J-1)+DRT(1,I,J-1)*ZSX(I,I,J)+DSX(I+1,J-1)+
0204 2DSX(I-I,J-1)-CSX(I-I,J-1)-CSX(I-I,J-1))/ZSX(I,I,J)
0205 TCB = (TCB(1,I+1,J-1)*CSX(I+1,J-1)-DSX(I+1,J-1)+TCT(1,I+1,J-1)+
0205 1CSX(I-I,J-1)-DSX(I-I,J-1)+TCT(1,I,J-1)*ZSX(I,I,J)+DSX(I+1,J-1)+
0205 2DSX(I-I,J-1)-CSX(I-I,J-1)-CSX(I-I,J-1))/ZSX(I,I,J)
0206 IF (MP.EQ.1) GO TO 54
0207 GVB = (GVB(1,I+1,J-1)*CSX(I+1,J-1)-DSX(I+1,J-1)+GVT(1,I+1,J-1)+
0207 1CSX(I-I,J-1)-DSX(I-I,J-1)+GVT(1,I,J-1)*ZSX(I,I,J)+DSX(I+1,J-1)+
0207 2DSX(I-I,J-1)-CSX(I-I,J-1)-CSX(I-I,J-1))/ZSX(I,I,J)
0208 54 CONTINUE
0209 IF (CR(1,I,J-1).EQ.0..AND.CSX(I+1,J-1).GT.DSX(I+1,J).AND.
0209 1CSX(I-I,J-1).LE.DSX(I-I,J)) GO TO 27
0210 GO TO 55
0211 27 CONTINUE
0212 FUB = (FUB(1,I+1,J-1)*CSX(I+1,J-1)-DSX(I+1,J-1)+FUT(1,I+1,J-1)+
0212 1ZSX(I,I,J)+DSX(I+1,J-1)-CSX(I+1,J-1))/ZSX(I,I,J)
0213 FVB = (FVB(1,I+1,J-1)*CSX(I+1,J-1)-DSX(I+1,J-1)+FVT(1,I+1,J-1)+
0213 1ZSX(I,I,J)+DSX(I+1,J-1)-CSX(I+1,J-1))/ZSX(I,I,J)
0214 DRB = (DRB(1,I+1,J-1)*CSX(I+1,J-1)-DSX(I+1,J-1)+DRT(1,I+1,J-1)+
0214 1ZSX(I,I,J)+DSX(I+1,J-1)-CSX(I+1,J-1))/ZSX(I,I,J)
0215 TCB = (TCB(1,I+1,J-1)*CSX(I+1,J-1)-DSX(I+1,J-1)+TCT(1,I+1,J-1)+
0215 1ZSX(I,I,J)+DSX(I+1,J-1)-CSX(I+1,J-1))/ZSX(I,I,J)
```



```
0275 IF (ZSX(I,J).GT.CSX(I,J-1).AND. CSX(I,J-1).GT.O..AND.  
0286 IMR(I,I-1,J-1).GT.O..AND.DR(I,I-1,J).EQ.O.) GO TO 65  
0287 GO TO 85  
0288 65 CONTINUE  
0289 FUB=(FV(I,I-1,J-1)*(ZSX(I,J)-CSX(I,J-1))+FV(I,I,J-1)*CSX(I,J-1))  
0291 1/ZSX(I,J)  
0291 FVB=(FV(I,I-1,J-1)*(ZSX(I,J)-CSX(I,J-1))+FV(I,I,J-1)*CSX(I,J-1))  
0291 1/ZSX(I,J)  
0291 DHR=(DR(I,I-1,J-1)*(ZSX(I,J)-CSX(I,J-1))+DR(I,I,J-1)*CSX(I,J-1))  
0291 1/ZSX(I,J)  
0291 TCR=(TC(I,I-1,J-1)*(ZSX(I,J)-CSX(I,J-1))+TC(I,I,J-1)*CSX(I,J-1))  
0291 1/ZSX(I,J)  
0292 GVB=(GV(I,I-1,J-1)*(ZSX(I,J)-CSX(I,J-1))+GV(I,I,J-1)*CSX(I,J-1))  
0292 1/ZSX(I,J)  
0293 85 CONTINUE  
0294 IF (ZSX(I,J).GT.CSX(I,J-1).AND. CSX(I,J-1).GT.O..AND.DR(I,I-1,J-1)  
0295 1.GT.O..AND.DR(I,I-1,J).EQ.O.) GO TO 32  
0295 GO TO 159  
0296 32 CONTINUE  
0297 FUB=(FV(I,I-1,J-1)*(ZSX(I,J)-CSX(I,J-1))+FV(I,I,J-1)*CSX(I,J-1))  
0297 1/ZSX(I,J)  
0297 FVB=(FV(I,I-1,J-1)*(ZSX(I,J)-CSX(I,J-1))+FV(I,I,J-1)*CSX(I,J-1))  
0297 1/ZSX(I,J)  
0297 DHR=(DR(I,I-1,J-1)*(ZSX(I,J)-CSX(I,J-1))+DR(I,I,J-1)*CSX(I,J-1))  
0297 1/ZSX(I,J)  
0297 TCR=(TC(I,I-1,J-1)*(ZSX(I,J)-CSX(I,J-1))+TC(I,I,J-1)*CSX(I,J-1))  
0297 1/ZSX(I,J)  
0301 IF (MP.FO.1) GO TO 159  
0302 GVB=(GV(I,I-1,J-1)*(ZSX(I,J)-CSX(I,J-1))+GV(I,I,J-1)*CSX(I,J-1))  
0302 1/ZSX(I,J)  
0303 159 CONTINUE  
0304 IF (I.LE.14) GO TO 777  
0305 IF (J.GT.JM).K=1  
0306 IF (J.LE.JM).K=1  
0307 IF (RSY(I,J).GT.CSY(I-1,J).AND.CSY(I-1,J).GT.O..AND.DR(I,I-1,J-K)  
0308 1.GT.O..AND.ZSX(I,J-K).EQ.O.) GO TO 101  
0308 GO TO 102  
0309 101 FUL=(FV(I,I-1,J)*CSY(I-1,J)+FV(I,I-1,J-K)*(RSY(I,J)-CSY(I-1,J)))  
0310 1/RSY(I,J)  
0310 FVL=(FV(I,I-1,J)*CSY(I-1,J)+FV(I,I-1,J-K)*(RSY(I,J)-CSY(I-1,J)))  
0310 1/RSY(I,J)  
0311 DRL=(DR(I,I-1,J)*CSY(I-1,J)+DR(I,I-1,J-K)*(RSY(I,J)-CSY(I-1,J)))  
0311 1/RSY(I,J)  
0312 TCL=(TC(I,I-1,J)*CSY(I-1,J)+TC(I,I-1,J-K)*(RSY(I,J)-CSY(I-1,J)))  
0312 1/RSY(I,J)  
0313 GUL=(GV(I,I-1,J)*CSY(I-1,J)+GV(I,I-1,J-K)*(RSY(I,J)-CSY(I-1,J)))  
0313 1/RSY(I,J)  
0314 102 IF (RSY(I,J).GT.DSY(I+1,J).AND.DSY(I+1,J).GT.O..AND.DR(I,I-1,J-K)  
0315 1.GT.O..AND.ZSX(I,J-K).EQ.O.) GO TO 103  
0315 GO TO 104  
0316 103 FUR=(FV(I,I+1,J)*DSY(I+1,J)+FV(I,I+1,J-K)*(RSY(I,J)-DSY(I+1,J)))  
0317 1/RSY(I,J)  
0317 FVR=(FV(I,I+1,J)*DSY(I+1,J)+FV(I,I+1,J-K)*(RSY(I,J)-DSY(I+1,J)))  
0317 1/RSY(I,J)  
0318 DRR=(DR(I,I+1,J)*DSY(I+1,J)+DR(I,I+1,J-K)*(RSY(I,J)-DSY(I+1,J)))  
0318 1/RSY(I,J)
```

```
0319 TCR=(TC(I,I+1,J)*DSY(I+1,J)+TC(I,I+1,J-K)*(RSY(I,J)-DSY(I+1,J)))  
0320 1/RSY(I,J)  
0320 GVB=(GV(I,I+1,J)*DSY(I+1,J)+GV(I,I+1,J-K)*(RSY(I,J)-DSY(I+1,J)))  
0320 1/RSY(I,J)  
0321 104 IF (DSY(I+1,J-K).GT.CSY(I+1,J).AND.CSY(I+1,J).GT.O..AND.  
0322 DR(I,I+1,J).GT.O..AND.ZSX(I+1,J).EQ.O.) GO TO 105  
0322 GO TO 106  
0323 105 FUR=(FV(I,I+1,J-K)*(DSY(I+1,J)-CSY(I+1,J-K))+FV(I,I+1,J)*(RSY(I,J)  
0324 1.CSY(I+1,J-K).DSY(I+1,J-K))/RSY(I,J)  
0324 FVR=(FV(I,I+1,J-K)*(DSY(I+1,J-K)-CSY(I+1,J-K))+FV(I,I+1,J)*(RSY(I,J)  
0325 1.CSY(I+1,J-K).DSY(I+1,J-K))/RSY(I,J)  
0325 DRR=(DR(I,I+1,J-K)*(DSY(I+1,J-K)-CSY(I+1,J-K))+DR(I,I+1,J)*(RSY(I,J)  
0326 1.CSY(I+1,J-K).DSY(I+1,J-K))/RSY(I,J)  
0326 TCR=(TC(I,I+1,J-K)*(DSY(I+1,J-K)-CSY(I+1,J-K))+TC(I,I+1,J)*(RSY(I,J)  
0327 1.CSY(I+1,J-K).DSY(I+1,J-K))/RSY(I,J)  
0327 GUR=(GV(I,I+1,J-K)*(DSY(I+1,J-K)-CSY(I+1,J-K))+GV(I,I+1,J)*(RSY(I,J)  
0328 1.CSY(I+1,J-K).DSY(I+1,J-K))/RSY(I,J)  
0329 106 IF (J.GT.JM) GO TO 778  
0329 IF (ZSX(I,J).GT.DSX(I+1,J).AND.DSX(I+1,J).GT.O..AND.DR(I,I+1,J)  
0330 1.GT.O..AND.RSY(I+1,J).EQ.O.) GO TO 107  
0330 GO TO 108  
0331 107 FUL=(FV(I,I+1,J)*DSX(I+1,J)+FV(I,I+1,J-K)*(ZSX(I,J)-DSX(I+1,J)))  
0332 1/ZSX(I,J)  
0332 FVL=(FV(I,I+1,J)*DSX(I+1,J)+FV(I,I+1,J-K)*(ZSX(I,J)-DSX(I+1,J)))  
0332 1/ZSX(I,J)  
0333 DRL=(DR(I,I+1,J)*DSX(I+1,J)+DR(I,I+1,J-K)*(ZSX(I,J)-DSX(I+1,J)))  
0333 1/ZSX(I,J)  
0334 TCL=(TC(I,I+1,J)*DSX(I+1,J)+TC(I,I+1,J-K)*(ZSX(I,J)-DSX(I+1,J)))  
0334 1/ZSX(I,J)  
0335 GUL=(GV(I,I+1,J)*DSX(I+1,J)+GV(I,I+1,J-K)*(ZSX(I,J)-DSX(I+1,J)))  
0335 1/ZSX(I,J)  
0336 108 IF (ZSX(I,J).GT.CSX(I,J-1).AND.CSX(I,J-1).GT.O..AND.DR(I,I-1,J)  
0337 1.GT.O..AND.RSY(I+1,J).EQ.O.) GO TO 109  
0337 GO TO 110  
0338 109 FUB=(FV(I,I-1,J-1)*CSX(I,J-1)+FV(I,I-1,J-1)*(ZSX(I,J)-CSX(I,J-1)))  
0339 1/ZSX(I,J)  
0339 FVB=(FV(I,I-1,J-1)*CSX(I,J-1)+FV(I,I-1,J-1)*(ZSX(I,J)-CSX(I,J-1)))  
0339 1/ZSX(I,J)  
0340 DHR=(DR(I,I-1,J-1)*CSX(I,J-1)+DR(I,I-1,J-1)*(ZSX(I,J)-CSX(I,J-1)))  
0341 1/ZSX(I,J)  
0341 TCR=(TC(I,I-1,J-1)*CSX(I,J-1)+TC(I,I-1,J-1)*(ZSX(I,J)-CSX(I,J-1)))  
0342 1/ZSX(I,J)  
0342 GVB=(GV(I,I-1,J-1)*CSX(I,J-1)+GV(I,I-1,J-1)*(ZSX(I,J)-CSX(I,J-1)))  
0343 1/ZSX(I,J)  
0343 110 IF (CSX(I-1,J-1).GT.DSX(I-1,J).AND.DSX(I-1,J).GT.O..AND.  
0344 DR(I,I-1,J).GT.O..AND.RSY(I+1,J).EQ.O.) GO TO 111  
0344 GO TO 112  
0345 111 FUB=(FV(I,I-1,J-1)*CSX(I-1,J-1)+FV(I,I-1,J-1)*(DSX(I-1,J)  
0346 1.J)*ZSX(I-1,J)-CSX(I-1,J-1))/ZSX(I,J)  
0346 FVB=(FV(I,I-1,J-1)*CSX(I-1,J-1)+FV(I,I-1,J-1)*(DSX(I-1,J)  
0347 1.J)*ZSX(I-1,J)-CSX(I-1,J-1))/ZSX(I,J)  
0347 DHR=(DR(I,I-1,J-1)*CSX(I-1,J-1)+DR(I,I-1,J-1)*(DSX(I-1,J)  
0348 1.J)*ZSX(I-1,J)-CSX(I-1,J-1))/ZSX(I,J)  
0348 TCR=(TC(I,I-1,J-1)*CSX(I-1,J-1)+TC(I,I-1,J-1)*(DSX(I-1,J)  
0349 1.J)*ZSX(I-1,J)-CSX(I-1,J-1))/ZSX(I,J)  
0349 GVB=(GV(I,I-1,J-1)*CSX(I-1,J-1)+GV(I,I-1,J-1)*(DSX(I-1,J)  
0349 1.J)*ZSX(I-1,J)-CSX(I-1,J-1))/ZSX(I,J)
```

```
0350      1J1=ZSX(I,J)-CSX(I-1,J-1))/ZSX(I,J)
0351      GO TO 122
0352      778 CONTINUE
0352      IF (ZSX(I,J).GT.CSX(I,J-1).AND.CSX(I,J-1).GT.O..AND.DR(I,1,J)
0353      1.GT.O..AND.RSY(I+1,J).EQ.O.) GO TO 117
0353      GO TO 119
0354      117 FUR=(FUI(I,J-1)*CSX(I,J-1)+FUI(I+1,J-1)*(ZSX(I,J)-CSX(I,J-1))
0355      1/ZSX(I,J)
0355      FVB=(FVI(I,J-1)*CSX(I,J-1)+FVI(I+1,J-1)*(ZSX(I,J)-CSX(I,J-1))
0356      1/ZSX(I,J)
0356      DRB=(DR(I,1,J-1)*CSX(I,J-1)+DR(I+1,J-1)*(ZSX(I,J)-CSX(I,J-1))
0357      1/ZSX(I,J)
0357      TCB=(TC(I,1,J-1)*CSX(I,J-1)+TC(I+1,J-1)*(ZSX(I,J)-CSX(I,J-1))
0358      1/ZSX(I,J)
0358      GVB=(GVI(I,1,J-1)*CSX(I,J-1)+GVI(I+1,J-1)*(ZSX(I,J)-CSX(I,J-1))
0359      1/ZSX(I,J)
0359      118 IF (ZSX(I,J).GT.DSX(I,J+1).AND.DSX(I,J+1).GT.O..AND.DR(I,1+1,J)
0360      1.GT.O..AND.RSY(I+1,J).EQ.O.) GO TO 119
0360      GO TO 120
0361      119 FUI=(FUI(I,1,J+1)*DSX(I,J+1)+FUI(I+1,J+1)*(ZSX(I,J)-DSX(I,J+1))
0362      1/ZSX(I,J)
0362      FVI=(FVI(I,1,J+1)*DSX(I,J+1)+FVI(I+1,J+1)*(ZSX(I,J)-DSX(I,J+1))
0363      1/ZSX(I,J)
0363      DRI=(DRI(I,1,J+1)*DSX(I,J+1)+DRI(I+1,J+1)*(ZSX(I,J)-DSX(I,J+1))
0364      1/ZSX(I,J)
0364      TCI=(TC(I,1,J+1)*DSX(I,J+1)+TC(I+1,J+1)*(ZSX(I,J)-DSX(I,J+1))
0365      1/ZSX(I,J)
0365      GVI=(GVI(I,1,J+1)*DSX(I,J+1)+GVI(I+1,J+1)*(ZSX(I,J)-DSX(I,J+1))
0366      1/ZSX(I,J)
0366      120 IF (DSX(I-1,J+1).GT.CSX(I-1,J).AND.CSX(I-1,J).GT.O..AND
0367      1.DRI(I,1,J).GT.O..AND.RSY(I+1,J).EQ.O.) GO TO 121
0367      GO TO 122
0368      121 FUI=(FUI(I-1,J+1)*(DSX(I-1,J+1)-CSX(I-1,J))+FUI(I,J+1)*CSX(I-1,
0369      1J)+ZSX(I,J)-DSX(I-1,J+1))/ZSX(I,J)
0369      FVI=(FVI(I-1,J+1)*(DSX(I-1,J+1)-CSX(I-1,J))+FVI(I,J+1)*CSX(I-1,
0370      1J)+ZSX(I,J)-DSX(I-1,J+1))/ZSX(I,J)
0370      DRI=(DRI(I-1,J+1)*(DSX(I-1,J+1)-CSX(I-1,J))+DRI(I,J+1)*CSX(I-1,
0371      1J)+ZSX(I,J)-DSX(I-1,J+1))/ZSX(I,J)
0371      TCI=(TC(I-1,J+1)*(DSX(I-1,J+1)-CSX(I-1,J))+TC(I,J+1)*CSX(I-1,
0372      1J)+ZSX(I,J)-DSX(I-1,J+1))/ZSX(I,J)
0372      GVI=(GVI(I-1,J+1)*(DSX(I-1,J+1)-CSX(I-1,J))+GVI(I,J+1)*CSX(I-1,
0373      1J)+ZSX(I,J)-DSX(I-1,J+1))/ZSX(I,J)
0373      122 CONTINUE
0374      777 CONTINUE
0375      59 CONTINUE
0376      IF (1/5*(FUR +FUI(I,J)))*2*(1/5*(FVR +FVI(I,J))
0377      1*(2)*CK.LT.(1/5*(TCR +TC(I,1,J)))*2 .AND. FUI(I,J) .GT.
0378      2FUR ) Q1=.25*CB*(TCR +TC(I,1,J))*IDRR
0377      3DR(I,1,J)=(FUI(I,1,J)-FUR )
0377      IF (1/5*(FUI +FUI(I,J)))*2*(1/5*(FVI +FVI(I,J))
0378      1*(2)*CK.LT.(1/5*(TCI +TC(I,1,J)))*2 .AND. FVI(I,J)
0379      2.GT. FVT ) Q3=.25*CB*(TCI +TC(I,1,J))*IDRT
0378      3DR(I,1,J)=(FVI(I,1,J)-FVT )
0378      IF (1/5*(FUI +FUI(I,J)+FUI(I,1,J)+FUI(I,1,J)))*2*(1/5*(FVI +FVI(I,1,J)+FVI(I,1,J)+FVI(I,1,J))
0379      1*(2)*CK.LT.(1/5*(TCI +TC(I,1,J)+TCI(I,1,J)+TCI(I,1,J)))*2 .AND. FUL
0380      2.GT. FUI(I,1,J) ) Q2=.25*CB*(TCI(I,1,J)+TCI(I,1,J))*IDRI
0379      3DR(I,1,J)=
```

```
0379      3DR(I,1,J)=(FUI +FUI(I,1,J)+FUI(I,1,J)+FUI(I,1,J))
0379      IF (1/5*(FUI(I,1,J)+FUI(I,1,J)+FUI(I,1,J)+FUI(I,1,J)))*2*(1/5*(FVI(I,1,J)+FVI(I,1,J)+FVI(I,1,J)+FVI(I,1,J))
0380      1*(2)*CK.LT.(1/5*(TCI(I,1,J)+TCI(I,1,J)+TCI(I,1,J)+TCI(I,1,J)))*2 .AND. FVB
0381      2.GT. FV(I,1,J) ) Q4=.25*CB*(TCI(I,1,J)+TCI(I,1,J)+TCI(I,1,J)+TCI(I,1,J))*IDRB
0380      3DRB(I,1,J)=(FVB +FVI(I,1,J))
0380      RETURN
0381      END
```


FORTRAN IV G LEVEL 21

AREAV

DATE = 72210

11/03/72

PAGE 0003

```

0078      W3=(A1+BV(1,1W+1,JW+1)+A2+CV(1,1W,JW+1)+A3+GV(1,1W,JW)+A4+FV(1,
0079      W4=(A1+GV(1,1W+1,JW+1)+A2+GU(1,1W,JW+1)+A3+GU(1,1W,JW)+A4+GU(1,
0080      W5=(A1+CV(1,1W+1,JW+1)+A2+CV(1,1W,JW+1)+A3+GV(1,1W,JW)+A4+GV(1,
0081      W6=(A1+P(1,1W+1,JW+1)+A2+P(1,1W,JW+1)+A3+P(1,1W,JW)+A4+P(1,
0082      W7=(A1+BV(1,1W+1,JW+1)+A2+BV(1,1W,JW+1)+A3+BV(1,1W,JW)+A4+BV(1,
0083      W9=(A1+TC(1,1W+1,JW+1)+A2+TC(1,1W,JW+1)+A3+TC(1,1W,JW)+A4+TC(1,
0084      W9=(A1+CE(1,1W+1,JW+1)+A2+CE(1,1W,JW+1)+A3+CE(1,1W,JW)+A4+CE(1,
0085      5 RETURN
0086      END

```

FORTRAN IV G LEVEL 21

FRAC

DATE = 72210

11/03/72

PAGE C001

```

0001      SUBROUTINE FRAC(SFIJ,AXIJ,AYIJ, CX,DY)
0002      DIMENSION DR(2,137,22),FU(2,137,22),FV(2,137,22),FI(2,137,22),
0003      1FE(2,137,22),TC(2,137,22),PP(2,137,22),P(2,137,22),BE(2,137,22),
0004      2GU(1,137,22),GV(1,137,22),GI(1,137,22),GE(1,137,22),TMP(2,137,22),
0005      3SYA(22),SYB(22),VCL(137,22),SI(22),CPB(22),BV(137,22),
0006      4CSX(137,22),DSX(137,22),CSY(137,22),DSY(137,22),RSY(137,22),
0007      5ZSX(137,22)
0008      COMMON DR,FU,FV,FI,FE,TC,PP,P,BE,GU,GV,GI,GE,TMP,SYA,SYB,VOL,
0009      1SI,CPB,BV,CSX,DSX,CSY,DSY,RSY, ZSX,C1,C2,C3,C4,P,I,J
0010      SFIJ=VCL(I,J)/VCL(2,J)
0011      AXIJ=RSY(I,J)/DY
0012      AYIJ=ZSX(I,J)/DX
0013      RETURN
0014      END

```

FORTRAN IV G LEVEL 19

VACUUM

DATE = 72173

19/11/77

PAGE 0001

```

0001      SUBROUTINE VACUUM(RMEAN,DX)
0002      DIMENSION DP(2,137,22),FU(2,137,22),FV(2,137,22),FI(2,137,22),
0003      1FE(2,137,22),TC(2,137,22),PP(2,137,22),P(2,137,22),BE(2,137,22),
0004      2GU(1,137,22),GV(1,137,22),GI(1,137,22),GE(1,137,22),TMP(2,137,22),
0005      3SYA(22),SYB(22),VOL(137,22),SI(22),CPB(22),BV(137,22),
0006      4CSX(137,22),DSX(137,22),CSY(137,22),DSY(137,22),RSY(137,22),
0007      5ZSX(137,22)
0008      COMMON DR,FU,FV,FI,FE,TC,PP,P,BE,GU,GV,GI,GE,TMP,SYA,SYB,VOL,
0009      1SI,CPB,BV,CSX,DSX,CSY,DSY,RSY, ZSX,Q1,Q2,Q3,Q4,M,I,J
0010      BV(I,J)=7000.*6.2832*DX*RMEAN/SX(I,J)*4.*
0011      PP(1,I,J)=PV(1,I,J)**2/25.1328
0012      BR(1,I,J)=0.
0013      BR(2,I,J)=0.
0014      GE(1,I,J)=0.
0015      FE(1,I,J)=0.
0016      P(1,I,J)=0.
0017      FI(1,I,J)=0.
0018      FI(2,I,J)=0.
0019      GU(1,I,J)=0.
0020      GV(1,I,J)=0.
0021      FV(1,I,J)=0.
0022      FU(1,I,J)=0.
0023      TC(1,I,J)=0.
0024      DR(2,I,J)=0.
0025      DR(1,I,J)=0.
0026      TMP(1,I,J)=0.
0027      TMP(2,I,J)=0.
0028      DSX(1,I,J)=0.
0029      DSY(1,I,J)=0.
0030      CSX(1,I,J)=0.
0031      CSY(1,I,J)=0.
0032      ZSX(1,I,J)=0.
0033      RSY(1,I,J)=0.
0034      RETURN
0035      END

```

//CO.FTCR001 DD CEN=H.LUT,VOL=SER=H10467,UNIT=2400-9,
// DISP=(NEW,KEEP),DCB=(RECFM=FB,LRECL=80,BLKSTZE=3200),
// LABEL=(1,SL)
//CO.SC406077 DD DISP=NEW,UNIT=2400-7,LABEL=(1,ML),
// VOL=SER=CSOJ37,ORCL=DFN=2,TRTCH=0
//CC.SYSIN DD *
/*