

MECHANICAL PROPERTIES OF NUGGET SANDSTONE

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MECHANICAL PROPERTIES OF NUGGET SANDSTONE

Abstract

The mechanical properties of Nugget sandstone have been determined from a number of independent measurements at pressures up to 30 kbar. Pressure-volume, strength, uniaxial stress, and uniaxial strain tests yield the failure surface and effective moduli as a function of stress state, i.e., mean pressure and shear stress. Acoustic velocity determinations provide the effective moduli for low amplitude dynamic waves. At atmospheric pressure the initial effective bulk modulus in tests with applied differential stresses (~40 kbar) differs from that determined hydrostatically (23 kbar). This

and a large pressure derivative of the shear modulus at low pressure are believed to be due to an abundance of cracks with low aspect ratios.

The initial effective shear modulus of about 55 kbar increases rapidly with the closing of cracks. At 1 kbar confining pressure, a shear modulus of about 120 kbar is determined in both uniaxial stress and uniaxial strain experiments. The rock has an ultimate strength comparable to granite (1.2 kbar unconfined) and exhibits brittle behavior at failure to the highest mean pressure studied (14 kbar). Failure is preceded by dilatant behavior.

Introduction

As part of a Laboratory program (Seismic Evasion) dealing with the detection of underground nuclear explosions, we have measured the physical properties of various materials selected to represent a wide range of possible responses to stress pulses. The physical properties so determined are then used in conjunction with small-scale explosion and shock-wave measurements to define the behavior of the material over a wide range of stresses, strains and strain rates. This information is then used to develop numerical techniques which seek to model stress-wave propagation. Factors such as water content and porosity, and related physical

properties such as moduli and strength, are well known to have a large effect on stress-wave propagation.

In this report we present the result of high-pressure, mechanical properties measurements on one material, the Nugget sandstone, from Utah. This sandstone is composed primarily of detrital quartz with siliceous cementing materials. Based on a reported quartz content of 99%,¹ the rock may be classified as orthoquartzite² or a quartz arenite by the more recent scheme of Pettijohn *et al.*³ The average density of the material used in this study is 2.56 g-cm^{-3} with a variation between

2.54 g-cm⁻³ and 2.57 g-cm⁻³. Most samples, however, varied by no more than 0.006 g-cm⁻³ from this average. All samples were fabricated with the cylindrical axis perpendicular to the poorly-defined bedding. No attempt was made to determine the degree of anisotropy.

We have determined the pressure-volume loading behavior to 30 kbar, the failure envelope to 14-kbar mean pressure, and acoustic velocities to 10 kbar. In addition, the three-dimensional stress-strain relationships in quasistatic loading have been measured to 13 kbar to determine such parameters as the stress and path dependence of the shear modulus and

to compare static loading moduli with those determined dynamically. The experimental techniques used have been described in detail elsewhere.⁴⁻⁸ All stresses are referred to a Cartesian coordinate system with σ_1 , σ_2 and σ_3 the maximum, intermediate, and minimum principal stresses, respectively. In all cases reported here, with the exception of the indirect tensile (Brazil) tests, σ_2 and σ_3 are equal. Stress is taken as positive and strain as negative in compression. All of our measurements are at strain rates in the range of 10^{-4} to 10^{-5} sec⁻¹, at room temperature, and on dry sandstone.

Experimental Results

PRESSURE-VOLUME

We measured the pressure-volume relationship for Nugget sandstone on six test samples. Two samples were tested to 30 kbar in a quasihydrostatic piston-cylinder device. These samples were 2.2 cm in diameter by 2.5 cm long. Tin was used as the pressure medium. The remaining four samples were tested to 14 kbar in a hydrostatic piston-cylinder apparatus using oil as the pressure fluid. Two of the samples tested to 14 kbar were 1.9 cm in diameter by 2.5 cm long, and another two were 3.1 cm in diameter by 5 cm long. The composite pressure-volume relationship from the six samples tested is shown in Fig. 1. In drawing this composite curve, the hydrostatic data were used at the low pressures (< 3 kbar). The data from all six runs agree to within 0.5% for V/V_0 in the 3 to 14-kbar region.

Over the common pressure range of 3 to 14 kbar, the bulk modulus determined by both methods agreed very well ($\pm 5\%$ of the stated value).

The composite curve shows V/V_0 to be a monotonic decreasing function over the pressure range to 30 kbar, except for the region around 8 kbar where a slight "S" shape is seen in the pressure-volume curve. This is so slight that for most purposes the bulk modulus can be considered as a monotonic increasing function of the pressure, as shown in Fig. 2. Table 1 lists the values of the bulk modulus at various pressures.

SHEAR STRENGTH AT HIGH PRESSURE

Jacketed cylinders (2 cm diam by 4 cm) were used to determine the strength of the dry Nugget sandstone at confining pressures ranging to 7.0 kbar. Initial

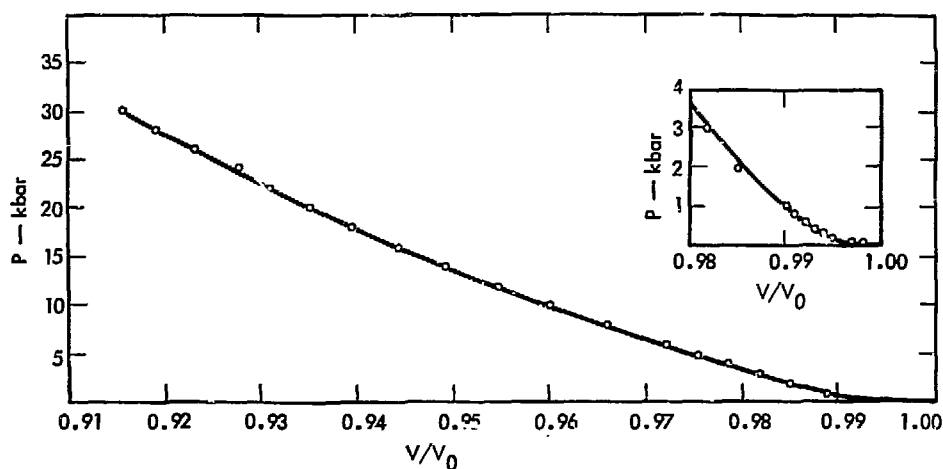


Fig. 1. Pressure-volume relationship for Nugget sandstone.

measurements were taken in the form of axial force-axial displacement traces and were reduced to yield differential stress-axial strain curves. Inspection of these

curves, together with the condition of the test sample after deformation, indicated that most, if not all of the inelastic strain had occurred by brittle fracture. Usually conjugate but occasionally single shear fracture planes, inclined $\sim 30^\circ$ to σ_1 , were obvious; cohesion across the narrow dislocation zones was usually lost. In addition, tensile fracture planes oriented parallel to σ_1 were commonly noted.

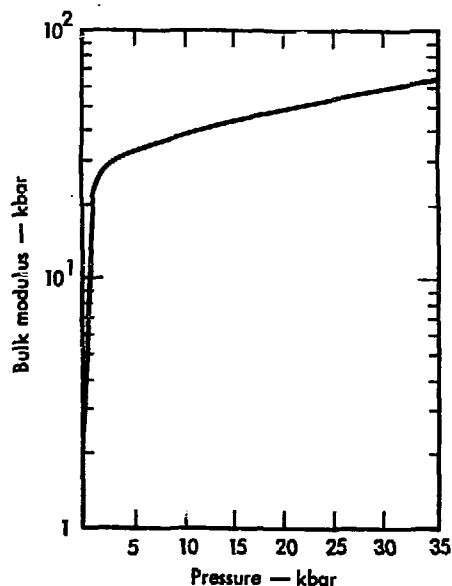


Fig. 2. Bulk modulus as a function of pressure for Nugget sandstone.

Table 1. Derived bulk moduli (K), Nugget sandstone.

Nugget Sandstone	
P (kbar)	K (kbar)
0	23
0.5	80
1	219
3	290
5	321
10	375
15	430
20	475
25	517
30	568

The ultimate (maximum) strength from these tests can be used to define a surface in shear-stress mean-pressure space. Figure 3 shows this surface or failure

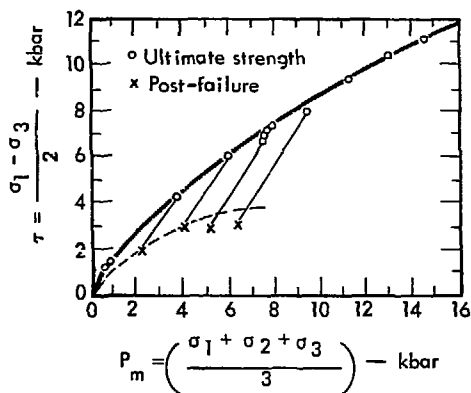


Fig. 3. Failure envelope, Nugget sandstone.

envelope along with the individual data points from each test. Table 2 summarizes the principal stresses at failure for all tests. As is usually observed in most brittle rock materials, after the initial rupture at the ultimate stress difference, the axial stress does not drop immediately to the value of the confining pressure. Instead, the fractured rock supports approximately one-half the maximum stress difference and further deformation occurs by stick-slip sliding between the rupture surfaces. As is observed in most rocks,⁹ post-failure shear strengths in Nugget sandstone remain at one-third to two-thirds of the maximum value until the jacket ruptures. Failure in the weak tygon jacket at the initial fracture prevented measurement of this phenomenon above

Table 2. Summary of uniaxial-compression and Brazil-test data, Nugget sandstone.

Test type	σ_1 (kbar)	σ_1' (kbar)	σ_3 (kbar)	σ_3' (kbar)	τ (kbar)	τ' (kbar)	P_m (kbar)	P_m' (kbar)	Behavior
Brazil	0.26		-0.09		0.17		0.06		Brittle
	0.33		-0.11		0.22		0.07		Brittle
	0.35		-0.12		0.23		0.08		Brittle
	0.45		-0.15		0.30		0.10		Brittle
Uniaxial comp.	2.30		0.001		1.15		0.77		Brittle
	2.60		0.001		1.30		0.87		Brittle
	9.60	4.76	1.00	1.06	4.30	1.85	3.87	2.25	Brittle
	14.13	8.52	2.03	2.15	6.05	3.19	6.06	4.27	Brittle
	17.52	9.44	3.01	3.19	7.26	3.13	7.85	5.27	Brittle
	16.95		3.00		6.98		7.65		Brittle ^a
	16.78		3.00		6.89		7.59		Brittle ^a
	17.50		3.00		7.25		7.83		Brittle
	17.24		3.00		7.12		7.75		Brittle
	20.04	11.12	4.02	4.29	8.01	3.42	9.36	6.57	Brittle
	23.77	5.52	5.02	5.52	9.38	0	11.27	5.52	Brittle
	26.93	6.68	6.07	6.68	10.43	0	13.02	6.68	Brittle
	29.51	7.33	7.01	7.33	11.25	0	14.51	7.33	Brittle

^a0.025-cm lead jacket.

11.2 kbar mean pressure. Both the ultimate strength and the post-failure data are quite reproducible, and each has a rather narrow scatter band. Jacketing by either tygon, lead, or copper, gave identical results.

Overall, this sandstone appears to be quite brittle and strong, at least when compared to graywacke sandstones measured in this Laboratory.¹⁰ The brittle-ductile transition for the Nugget sandstone must be in excess of 15-kbar mean pressure—at least a factor of two higher than for sandstones from the Wagon Wheel site.¹¹ The slope of the shear-stress mean-pressure envelopes (unfractured) is nearly identical (0.65) for these two sandstones. Shear strengths in the Nugget sandstone are about 10% higher, due probably to a higher initial cohesive strength and better cementation between grains and lower initial porosity.

Tensile strengths were also determined by a series of Brazil tests. The relatively high values of 90–150 bars (Table 2), are characteristic of a well-cemented aggregate. Tensile strengths in quartz-rich aggregates may vary from zero (sand pile) to at least 250 bars for a well cemented orthoquartzite.¹² The average tensile strengths for Nugget are slightly higher than for the Wagon Wheel sandstone¹¹; 120 versus 110 bars.

ULTRASONIC VELOCITY

Travel-times of both shear and compressional pulses through small cylindrical specimens (1.9 cm diam by ~2.5 cm long) were determined to 10 kbar at 1 MHz, using polycrystalline ceramic Pb (Zr,Ti)O₃ transducers.¹³ The data are plotted in

Fig. 4 as a function of pressure. Table 3 contains velocities and calculated moduli. We used Cook's method¹⁴ to estimate the change in specimen length. This method employs a step-wise integration of the calculated bulk modulus, described in detail by Schock *et al.*¹³

Both velocities are observed to increase rapidly with pressure as the effective deformation moduli increase for these low-amplitude ultrasonic waves. This behavior is characteristic of many common rock types,¹⁵ and is usually associated with microcracks having low-aspect ratios that close easily with pressure. There is little indication of the inelastic compaction which is very evident in more porous sandstones.¹⁶ Although most of the velocity change has occurred at 2 kbar, both velocities are still increasing slowly at the highest pressures obtained. This is in contrast to crystalline igneous rocks (Westerly granite) where the velocities do not increase significantly above 2 kbar because the voids are primarily long, narrow microcracks which are easily closed with pressure.¹³ Apparently a small amount of porosity remains at 10 kbar in the Nugget sandstone, which is probably more spherical, and thus is resisting collapse. Such behavior has been observed in another quartz-rich rock, the Stirling quartzite.¹⁷ The shear modulus exceeds the bulk modulus at ambient pressure, in agreement with the static results which follow. This behavior reverses at ~200 bars, and has been correlated with interlocking textures in sedimentary rock.¹⁰

The moduli calculated from the travel-times of small-amplitude waves are considerably greater than those measured

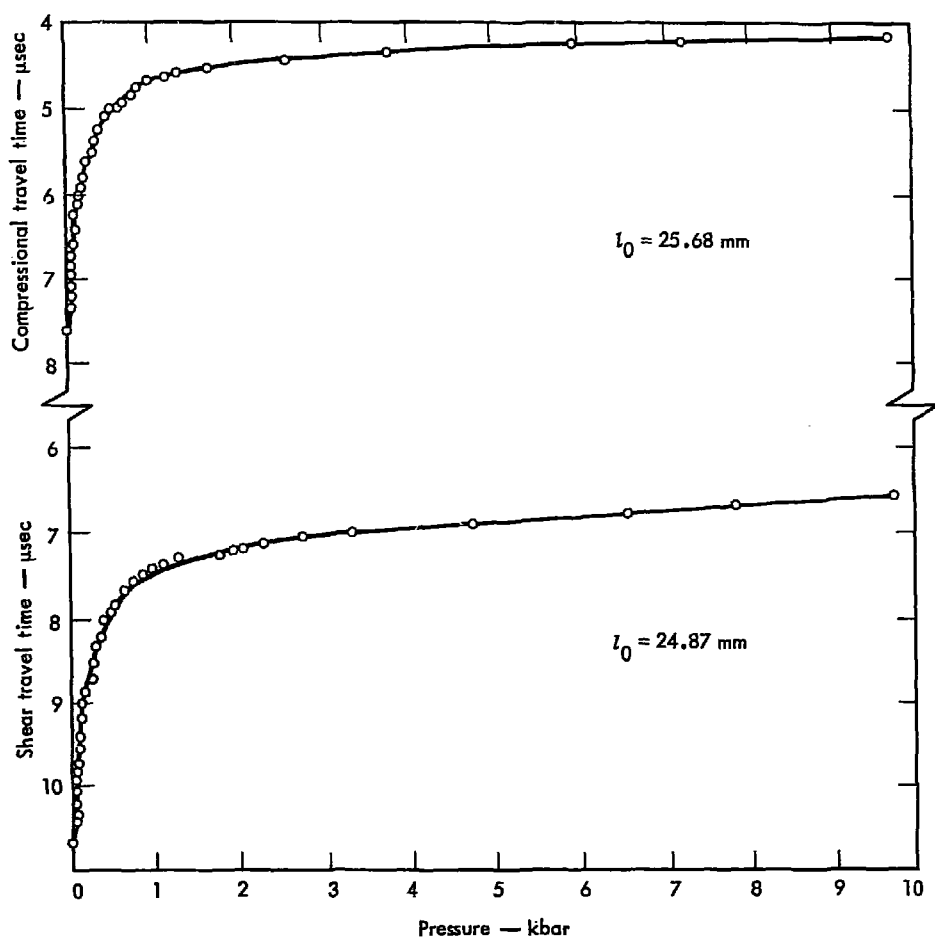


Fig. 4. Measured travel-times for small-amplitude strain waves as a function of hydrostatic pressure for Nugget sandstone.

during the static experiments (compare Tables 1 and 3, and the uniaxial stress and uniaxial strain results). This discrepancy is typical of porous rock, particularly those containing long narrow cracks.¹⁸ Rock moduli are usually strong functions of stress amplitude and strain rate. This, together with the effect of wave-scattering around arrays of micro-

cracks, offers alternative explanations of these differences in moduli.

THREE-DIMENSIONAL STRESS-STRAIN

The three-dimensional stress-strain relationship was measured over a variety of loading paths in shear stress and mean stress space, using jacketed cylindrical

Table 3. Calculated velocities and moduli for Nugget sandstone, perpendicular to the bedding.

Pressure (kbar)	ρ (g/cm ³)	V_p (km/sec)	V_s (km/sec)	K (Mbar)	μ (Mbar)	Poisson's ratio	Young's modulus (Mbar)
0	2.560	3.44	2.34	0.117	0.140	0.07	0.300
0.1	2.562	4.14	2.73	0.184	0.191	0.11	0.426
0.2	2.563	4.46	2.89	0.224	0.215	0.14	0.488
0.4	2.565	4.92	3.08	0.297	0.244	0.18	0.574
0.6	2.567	5.16	3.18	0.337	0.260	0.19	0.620
0.8	2.568	5.30	3.25	0.358	0.272	0.20	0.651
1.0	2.570	5.45	3.41	0.366	0.299	0.18	0.705
2.0	2.576	5.71	3.45	0.429	0.307	0.21	0.744
3.0	2.582	5.82	3.52	0.448	0.319	0.21	0.774
4.0	2.588	5.91	3.54	0.470	0.325	0.22	0.792
5.0	2.593	5.99	3.59	0.485	0.333	0.22	0.813
6.0	2.599	6.05	3.62	0.499	0.340	0.22	0.831
7.0	2.604	6.08	3.64	0.503	0.344	0.22	0.841
8.0	2.609	6.12	3.67	0.508	0.352	0.22	0.857
9.0	2.614	6.14	3.69	0.512	0.356	0.22	0.867
10.0	2.619	6.15	3.72	0.509	0.362	0.21	0.878

samples. For work at axial stresses above 5 kbar, the samples (2-cm-diam \times 3 cm) were loaded with the piston of a piston-cylinder pressure vessel in which the confining pressure could be controlled independent of the advance of the piston. In this case, the samples were jacketed with either 0.012-cm-thick or 0.025-cm-thick lead, and strain gages were bonded to the jacket. For work at low axial stresses at confining pressures to 4 kbar, a fluid loading system was used and the samples (2.5-cm-diam \times 6 cm) were jacketed with a thin film of epoxy.⁸ Strain-gage output was corrected for the effect of pressure.¹⁹

Uniaxial Stress

Samples were loaded to failure in uniaxial stress at confining pressures to 5 kbar. Figure 5 shows the measured

axial and circumferential strains on loading to failure at atmospheric confining pressure. In contrast to the circumferential strain, the initial increase in axial

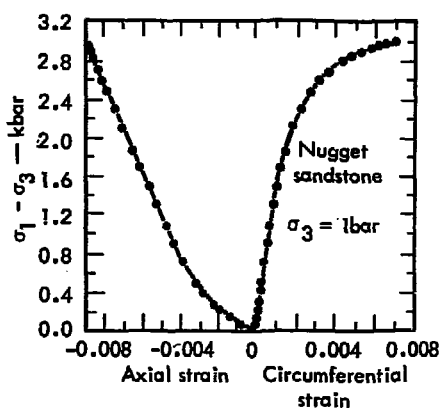


Fig. 5. Axial and circumferential strain as a function of axial stress difference at atmospheric confining pressure.

strain is nonlinear. Similar results are well known in granitic rocks,^{20,21} and have been attributed to the influence of cracks with low-aspect ratios. In response to an applied axial stress, those cracks whose major axes lie in or nearly in the plane normal to the applied stress offer little resistance and are closed at low stresses. In contrast, those cracks whose major axes are parallel to σ_1 cannot be closed by this stress. Furthermore they cannot be opened at low stresses. The results of these interactions is an increase in axial strain larger than that for the rock without cracks during the first several hundred bars of axial compression. After the cracks are closed, the continued strain response is essentially that of the rock without cracks until higher stress levels where microfracturing is initiated. Thus, one may conclude from the data in Fig. 5 below 1 kbar that the Nugget sandstone contains abundant cracks with low aspect ratios.

As the axial stress is increased to failure, the circumferential strain now begins to increase in a nonlinear fashion while the axial strain increase is nearly linear. This behavior corresponds to an increase in the macroscopic volume of the rock (Fig. 6). The volume increase before failure (dilatancy) is well known in brittle rocks²² and is thought to be the effect of microfracturing²³ which precedes ultimate failure. Since the predominant strain is circumferential, the microfracturing must be the result of the opening of cracks whose major axes are nearly parallel to σ_1 . In Fig. 7 we show the shear strain with increasing axial-stress difference. The initial shear modulus ($-1/2$ slope in Fig. 7) is found to

be 54 kbar. After the cracks are closed, the shear modulus has increased to 155 kbar. This represents the shear modulus of the rock without the influence of cracks. As the rock dilates the shear modulus decreases, gradually approaching zero as the rock breaks. The initial bulk

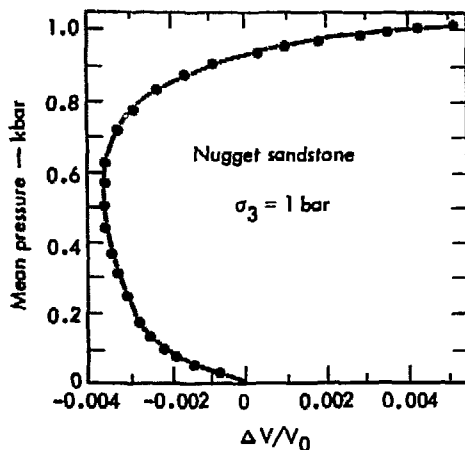


Fig. 6. Volume strain as a function of mean pressure on uniaxial stress loading at atmospheric confining pressure (data from Fig. 5).

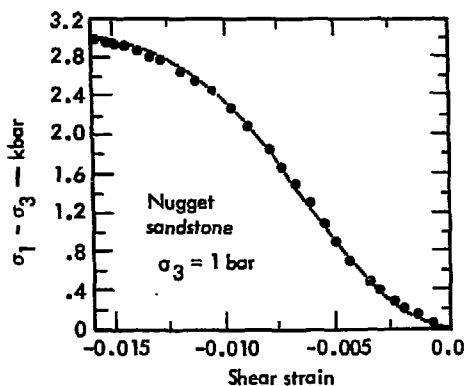


Fig. 7. Shear strain as a function of axial stress difference at atmospheric confining pressure (data from Fig. 5).

modulus, determined from the data in Fig. 6, is 35 kbar which is lower than the initial effective shear modulus. This is characteristic of sandstones with appreciable detrital quartz contents,¹⁰ and may be due to the interlocking nature of the grains themselves. An interlocking texture would serve to inhibit sliding of grains while not affecting compaction.⁷

If we examine the data obtained in a uniaxial stress test at 1-kbar confining pressure we see that compaction of cracks oriented primarily normal to σ_1 is responsible for the nonlinear behavior. At 1 kbar the confining pressure itself is enough to close these cracks, and thus a plot of axial strain (Fig. 8) does not show the initial nonlinearity noted in the same test at 1-bar confining pressure. The shear strain (Fig. 9) also shows no evidence of the influence of cracks at low σ_1 . At high σ_1 , the rock still dilates before failure (Fig. 10).

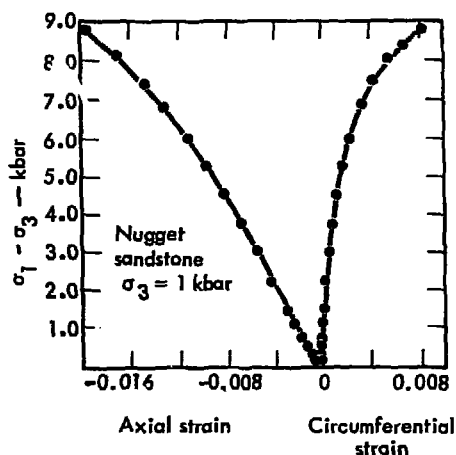


Fig. 8. Axial and circumferential strain as a function of axial stress difference at 1-kbar confining pressure.

The initial shear modulus at 1 kbar is 223 kbar, well above the initial value at 1 bar and higher than the value obtained at 1 bar after the axial stress-strain curve became linear. This is due to the increased confining pressure which causes

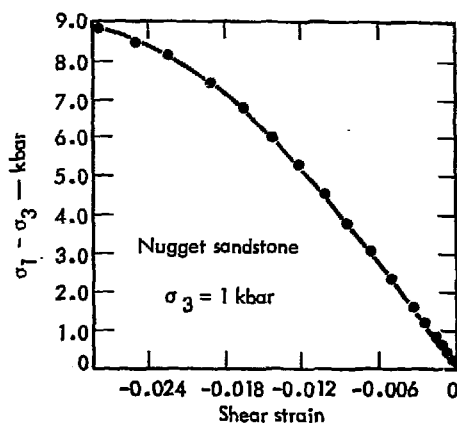


Fig. 9. Shear strain as a function of axial stress difference at 1-kbar confining pressure.

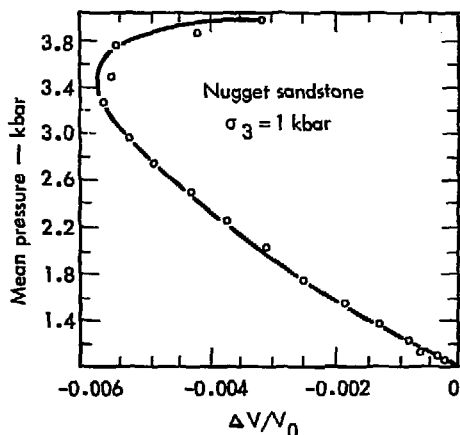


Fig. 10. Volume strain as a function of mean pressure on uniaxial stress loading at 1-kbar confining pressure.

increased friction between constituent grains and cement. The initial bulk modulus at 1 kbar is 280 kbar, higher than that determined under hydrostatic conditions. This surprising behavior persists up to 5 kbar which is the limit of the uniaxial stress tests performed here and is probably due to the influence of cracks, as discussed below.

The behavior of this rock in uniaxial stress is remarkably similar to that of granitic rocks. All of the features just noted as being due to the closing of cracks and the dilatant behavior are found in granitic-type rocks.^{7,20} However, both the shear and bulk moduli are lower than those for granites, a feature characteristic of quartzitic sandstones.⁷

A typical granite shows a smaller amount of compression in the axial direction due to the closing of cracks. The Hoggar granite²⁰ exhibits <0.05% strain due to the initial nonlinear portion of its stress-strain curve, in comparison to the 0.2% shown by the Nugget sandstone in Fig. 5. Nevertheless, this sandstone is somewhat stronger than most granites studied in this Laboratory.²⁰

No evidence of a decrease in dilatant behavior is observed to 5-kbar confining pressure. This is in agreement with the lack of any evidence of ductile behavior in the failure response.⁷

Uniaxial Strain

The uniaxial strain test is accomplished by varying the confining pressure as axial stress is increased. Uniaxial strain is assumed to be the condition under which deformation by a plane-shock wave takes place. The condition arises primarily due to restrictions of symmetry, ne-

glecting boundary effects. In principle, the high strain rates involved in the shock-wave experiments reduce the possibility of boundary effects through inertial considerations. During the passage of a plane shock wave no strain normal to the propagation direction occurs, and the pressure surrounding the sample will increase with the axial stress according to the effective moduli of the material, viz.

$$\frac{\sigma_1}{\sigma_3} = \frac{3K + 4\mu}{3K - 2\mu},$$

where K is the effective bulk modulus and μ the effective shear modulus. The uniaxial strain test therefore accomplishes the same result, but at a much lower strain rate.

Figure 11 shows the loading path in uniaxial strain to 7-kbar confining pressure compared to the failure envelope (Fig. 3). Several unloading paths in uniaxial strain are also shown. The loading path, after having started in the direction of the failure surface, begins to curve toward higher confining pressures and seems to parallel the failure envelope. In a recent report²⁴ it was shown that in granites the uniaxial strain loading path coincides with the lower limit of the region which describes dilatant behavior. This is because the onset of dilatant behavior is path-independent, and the nonlinear increase in circumferential strain (which is a necessary feature of dilatant behavior) is precluded by definition in a uniaxial strain test. In Fig. 11 we have plotted three points at which the Nugget sandstone dilates on loading in uniaxial stress.

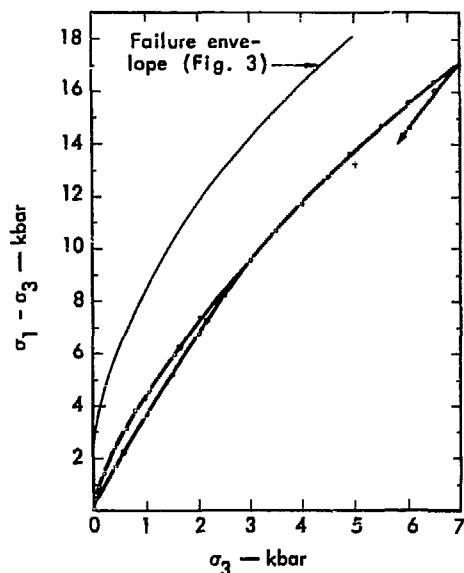


Fig. 11. Axial differential stress as a function of confining pressure for loading in uniaxial strain and on unloading from several pressures. (+) represents the onset of dilatant behavior in a uniaxial stress test.

The onset of dilatancy is defined here as the point at which the mean pressure-volume strain curve in uniaxial stress loading becomes significantly steeper than the hydrostat at corresponding mean pressures. In other words, the bulk modulus becomes greater. (Compare Figs. 6 and 9 with Fig. 1.) These points (Fig. 11) support the hypothesis that the low strain rate uniaxial strain loading curve is controlled by the onset of dilatant behavior. However, available evidence²⁴ indicates that under shock-loading conditions dilatant behavior may take place although at higher stress levels due to effects of strain rate on crack propagation. Since the failure envelope

shown in Figs. 3 and 11 is comparable to that found for three granitic rocks,²⁰ we can anticipate a Hugoniot-elastic-limit (HEL) for this rock at stress levels of at least 40-50 kbar. The exact value will be influenced by the strain-rate dependence of failure and inelastic processes which occur before failure. HEL values for polycrystalline quartz aggregates range from 40-97 kbar in the case of pure material and from 40-59 kbar for coarse-grained natural aggregates.²⁵ With the extremely high quartz content of this sandstone, one may expect the HEL to be in the 40- to 60-kbar range. Since there is no evidence of compaction in these rocks, as in other sandstones,^{7,9} we do not anticipate a relaxation at stresses below those associated with brittle failure.

In Fig. 12 we show the strain measured during uniaxial strain loading. Both the axial stress and the mean pressure are shown for comparison. Also shown are six points determined by Larson²⁶ in a series of low-pressure gas gun experiments. In general, these points agree well with the static uniaxial-strain data. There is a tendency toward steeper loading in the static results which might be attributed to strain rate. However, processes which involve compaction, crack propagation, or other inelastic processes are generally strain-rate-dependent in the opposite sense, i.e., they show less compaction at higher strain rates. We have no explanation for this discrepancy.

The initial slope (σ_1 versus $\Delta V/V_0$) in Fig. 12 yields a velocity of 2.0 km-sec^{-1} . This reflects the low initial slope shown in the Fig. 12 insert. Beyond 1-kbar axial stress the curve is much steeper, and yields a velocity of 4.6 km-sec^{-1} for the

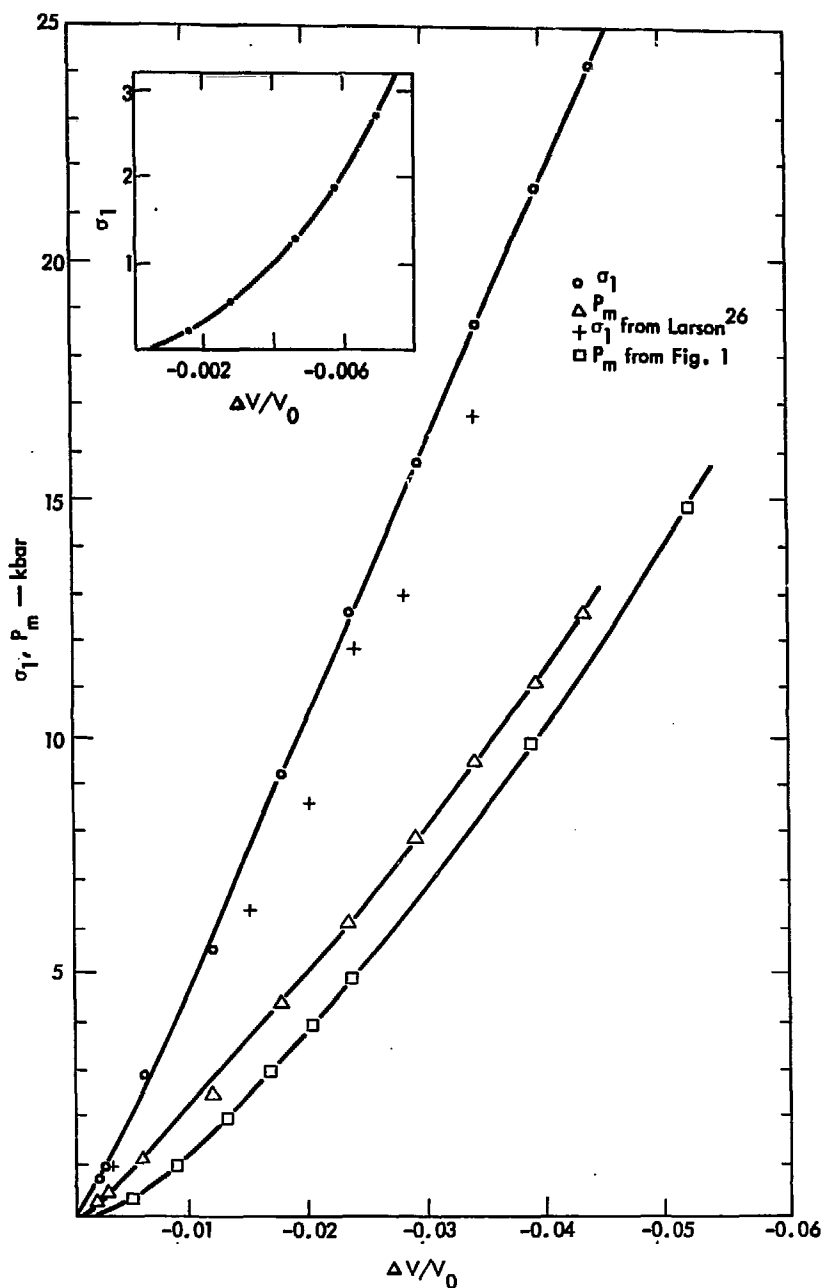


Fig. 12. Axial stress (σ_1) and mean pressure (P_m) as a function of volume strain for loading in uniaxial strain. Shock-wave data obtained by Larson²¹ are shown for comparison.

propagation of an elastic wave. This is the first material in which we have been able to observe such a steepening in a slope at pressures below 1 kbar in a uniaxial strain test. The observed behavior at low stresses in the uniaxial stress tests and the behavior of cracks suggests that the effect may have been masked in the

granites by the relatively low strains due to compression of cracks. We observe no discontinuity in the σ_1 versus $\Delta V/V_0$ loading curve attributable to pore collapse, as observed in more porous sandstones.¹⁰ The low porosity of this sandstone apparently precludes this.

Discussion

Comparison of the hydrostatic pressure-volume curve with the mean pressure-volume relation obtained in uniaxial strain (Fig. 12) reveals that the uniaxial strain data show less compression than the hydrostatic data at corresponding mean pressures. Inspection of the uniaxial stress data in Figs. 6 and 9 shows that this loading path is also above the hydrostat. Granites load along the hydrostat to the highest stress attained in uniaxial strain and to the onset of dilatancy in uniaxial stress.¹⁷ Other sandstones¹⁰ load along the hydrostat until they either dilate or undergo compaction due to collapse of their pore space. The disagreement between the hydrostat and mean pressure-volume curves in nonhydrostatic loading is probably due to cracks. In the uniaxial strain test, the maximum principal stress is always significantly greater (increases faster) than the confining pressure, causing those cracks oriented parallel to the principal applied stress to remain open. However, when the rock is compressed hydrostatically the cracks close without regard to orientation. This behavior has not been observed in the uniaxial strain loading paths in granites, and is probably the result of the much smaller strains.¹⁹

For granitic rocks, the computed volume strains have an uncertainty which approaches the theoretical difference (0.001 in strain) in the pressure-volume curves.

The data in Fig. 12 support the above argument. At pressures above 1 kbar, the offset between the uniaxial strain data and the hydrostat is about 0.004 in strain. If one considers a three-dimensional array of homogeneously distributed and oriented cracks, on the average only one-third will close when a uniaxial stress is applied. Examination of the data in Fig. 12 shows that the compression at low pressures before the knee in the uniaxial strain test is about 0.0015, or about one-third to one-fourth the total before the knee in the hydrostat (0.005). Of course, this analysis is an oversimplification. In reality cracks will exhibit a range of aspect ratios, external stresses will be transmitted internally through the rock matrix in a complex fashion, and perfect agreement to the above model should not be expected. The theoretical treatments of Walsh²¹ and Walsh and Brace²⁷ predict such an effect on cracks at low stresses.

The initial loading data in Fig. 11 and 12 yield an effective shear modulus of

56 kbar and an effective bulk modulus of 43 kbar. These are in good agreement with the values obtained on uniaxial stress loading at 1 bar. However, the bulk modulus is higher than that calculated from hydrostatic loading, presumably due to the effect just analyzed. At axial stress values above 1 kbar, the effective shear modulus becomes 219 kbar and remains relatively constant with increasing stress. The effective bulk modulus above the knee in Fig. 11 is 268 kbar and increases with increasing stress. These values agree very well with those from the initial loading in uniaxial stress at 1-kbar confining pressure.

On unloading, the mean pressure-volume strain data in Fig. 12 do not show any detectable departure from the loading values. However, the data in Fig. 11 show an initial stiffening in shear response on unloading. This is believed to result from friction on surfaces which slide on loading and become

temporarily locked when stresses are relaxed.

In summary, we have determined the mechanical properties of samples from a block of dry Nugget sandstone ($\rho_0 = 2.56 \text{ g-cm}^{-3}$) both at atmospheric and under high confining pressures. The initial loading moduli are found to be strongly influenced by an inelastic phenomenon at stress levels below 1 kbar. This phenomenon is believed to represent the closing of cracks with low-aspect ratios. At confining pressures of 1 kbar and above, these cracks appear to be closed and the rock exhibits a much stiffer response to both shear and simple compressive stresses. No evidence of an irreversible collapse of pore space is found to pressures of 30 kbar. The rock dilates before failure in the same manner as granitic rocks. Failure in uniaxial stress is by brittle fracture to the highest confining pressures attained (7.0 kbar).

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