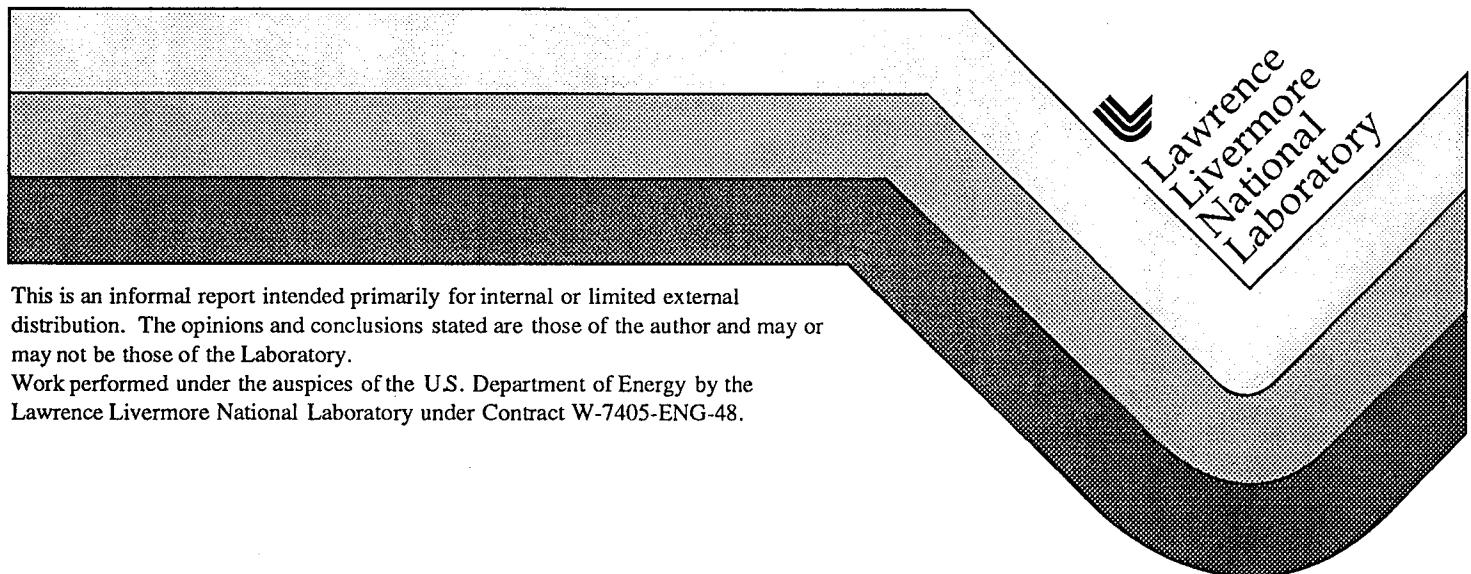


U235: A Gamma Ray Analysis Code for Uranium Isotopic Determination

DeLynn Clark

December 1996



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U235: A Gamma Ray Analysis Code for Uranium Isotopic Determination

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Abstract:

A ^{235}U analysis code, U235, has been written that can nondestructively determine the percentage of ^{235}U in a uranium sample from the analysis of the emitted gamma rays. The code is operational and work is now underway to improve the accuracy of the calculation, particularly at the high ($>90\%$) and low ($<0.7\%$) ^{235}U concentrations. A technique has been found to evaluate low ^{235}U concentrations that works well on the existing standards. Work is now under way to evaluate this technique for other detectors and other types of samples. Work is also proceeding on: (1) ways to better determine gamma backgrounds, (2) techniques to determine the equivalent thickness of the sample to correct for gamma attenuation, (3) evaluation of the existing data base of branching ratios of ^{235}U , ^{238}U and their daughters gamma rays to allow better results and (4) evaluation of the existing data base on the emission ratios for uranium, thorium, and protactinium x-rays.

Introduction:

Gamma ray spectrometry can be used to analyze uranium isotopic abundance ratios. The "standard" uranium enrichment meter relies on making standards of the various sample types of interest. Analyzing these standards with mass spectrometry to find the appropriate calibration factors is then done to calibrate out all the unknowns in the counting scheme. Then the strong ^{235}U gamma peak at 185.712 keV can be counted with a "simple" two channel analyzer to find the peak counts and background. The net 185.715 counts are used to calculate the enrichment. This technique works well but has the draw back that new "standards" have to be made for each different geometry and analyzed by mass spectrometry. This calibration process is often very time consuming and costly as well as being limited to "calibrated" geometries.

Accurate analysis of a radioactive sample by high resolution Germanium detector spectrometry requires correct information on the gamma ray and x-ray branching ratios for the radionuclides in the sample. ^{235}U and ^{238}U sample analysis is complicated in that the gammas observed often come from their radioactive daughters produced by successive alpha and beta decays. In addition to gamma decay these elements decay by internal conversion, IC, and subsequent

emission of daughter product x-rays. For example, when ^{235}U alpha decays the result is a radioactive ^{231}Th nucleus. This thorium isotope decays by both gamma emission and internal conversion (IC). Internal conversion results in an electron being ejected; usually from the K-shell; but L, M etc. shell conversion are also possible. This ejected electron leaves a hole in the atomic orbit that is filled by the fall of another electron from a higher electron orbital, giving rise to a thorium x-ray spectrum associated with the decay. In addition, x-rays are also produced by gammas interacting (via the photoelectric effect) in the material itself—so called fluorescent x-rays. In the case of a pure uranium sample these will be uranium x-rays. Internal conversion processes give rise to characteristic x-rays of the daughter product (not the parent) and are not proportional to the amount of material (the amount of thorium in a decaying sample of purified uranium is very small) in the sample. Internally conversion induced x-rays are proportional to the number of decays; i.e. each decay has a fractional output of x-rays of the daughter product regardless of the parent material present in the sample. This fact makes these x-rays usable for isotopic analysis if the sample has a very low concentration of daughter material (Th). To accurately use these IC x-ray peaks requires that the thorium present in very old “natural” uranium samples be removed. X-rays induced by the photoelectric effect (fluorescent x-rays) have energies characteristic of the bulk material—being proportional to the mass of material present in the source. The observed x-rays, from both fluorescent and internally converted sources, must originate within a mean free path of the surface to be easily observed.

Branching ratio and gamma, x-ray energy data have been published in various places^{1,2,3,4} for ^{235}U and ^{238}U and some of their daughter products; but this data is sometimes incomplete, or of inadequate accuracy. The current status of this data are summarized in Appendix A.

Figure 1 shows the main decay scheme for ^{235}U and ^{238}U . Pure ^{238}U emits only a 49.55 keV gamma that is so weak as to be almost useless for analysis (see Figure 2.). Fortunately, ^{238}U alpha (and beta) decays so that in within a few months it is in equilibrium with the ^{238}U decay and there are gammas from ^{234}Th , ^{234}Pa and ^{234}U available for analysis.(see Appendix E for a discussion of equilibrium). The small percentage (.0057%) of natural ^{234}U typically observed is due to the constant decay of ^{238}U . Similarly ^{235}U relatively quickly decays to equilibrium with its daughters, ^{231}Th and ^{231}Pa . Samples of uranium that have been enriched or separated can be analyzed for their ^{235}U concentrations by using these daughter product decays in all cases except very fresh (<2 month old) samples. At present the only way to accurately measure "fresh" samples before equilibrium is established is to use mass spectrometry.

Alternately, using high resolution gamma spectrometers, the spectra can be measured and the $^{235}\text{U}/^{238}\text{U}$ ratio determined by finding the peak intensities of neighboring gamma (or x-ray)

¹ Firestone, B. F.. ed., "Table of Isotopes", 8th Edition Lawrence Berkeley Laboratory, John Wiley & Sons, 1996

² Decay Data of the Transactinium Nuclides, Technical Report 261, IAEA, 1986

³Roy, J. C., et. al. *Int. J. Appl. Radiation Isotopes* , 35, pg 899, 1984

⁴Lammer, M. and Schwerer, O. *Handbook of Nuclear Data for Safeguards*", INDC (NDS) - 248, IAEA, 1991

peaks from each isotope. By taking intensity ratios on gamma peaks very close to the same energy, the detector efficiency and gamma attenuation differences in the sample will be small and to first order cancel.

When referring to the ^{238}U peaks in the following discussion the assumption is made that the gamma spectrum is in equilibrium with daughters ^{234}Th (24.1 d) , ^{234}Pa (6.70 hr.) , ^{234}mPa (1.17 min.) , but not ^{234}U (2.457E5 yr.) and its daughters. Similarly, the ^{235}U spectrum is assumed to be in equilibrium with its daughter ^{231}Th (25.52 hr.) , but not ^{231}Pa (3.276E4 yr.) and its daughters.

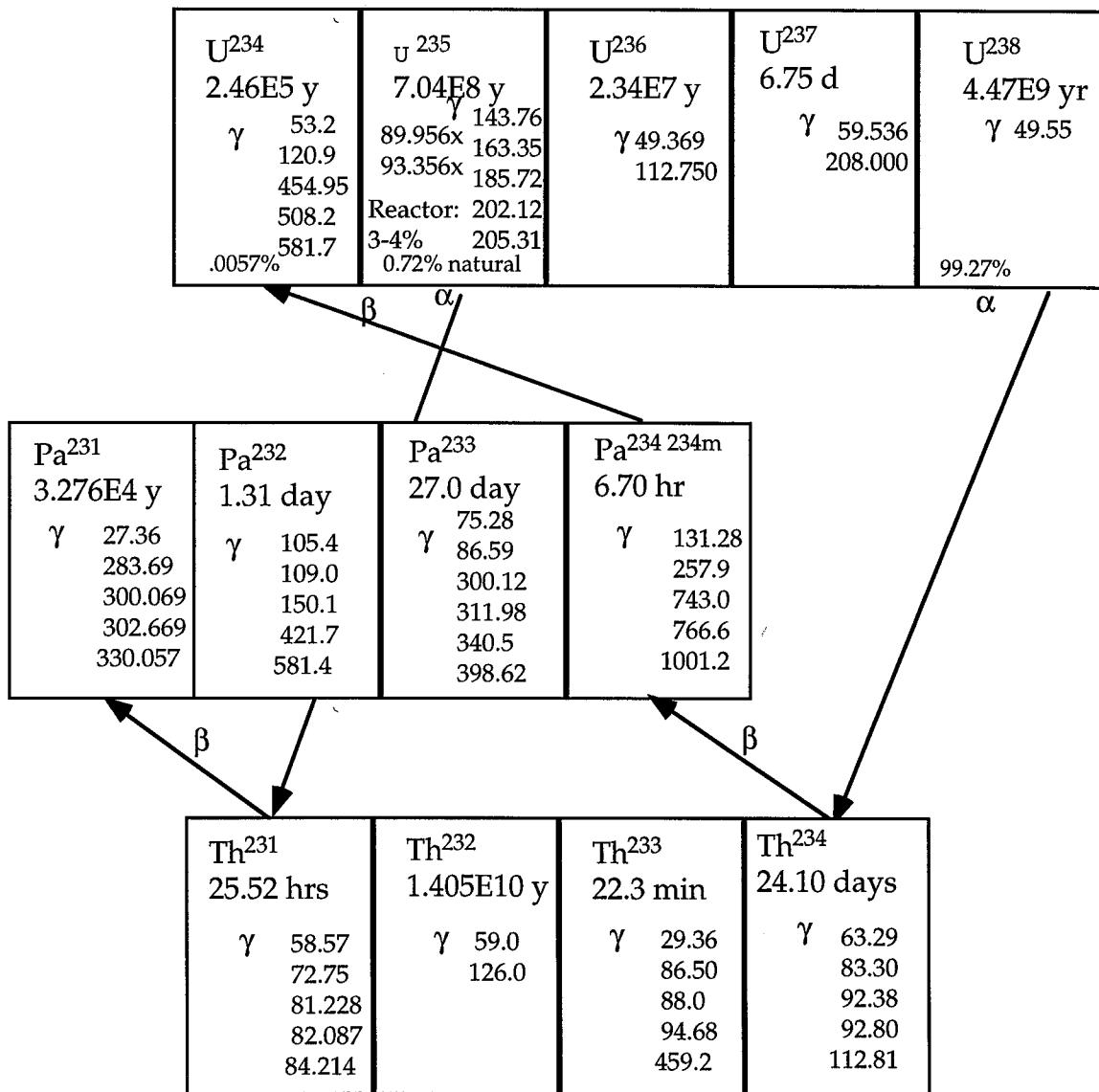


Figure 1. ^{235}U and ^{238}U decay scheme showing their principal daughters, their half-lives and their principal gamma rays.

The isotopic abundance is related to the observed peak intensities by the following relation:

$$I_1 = \lambda_1 A_1 B_1 \Omega_1 \epsilon_1 \tau_1 \text{ (counts/sec)}$$

Where:

I_1 = measured peak intensity of isotope one

$\lambda_1 = .6932/T_{1/2}$ = decay constant of isotope one

$T_{1/2}$ = material half-life (in seconds) of isotope one

A_1 = Number of atoms of isotope one

B_1 = branching ratio of isotope one

Ω_1 = Fractional solid angle of detector

ϵ_1 = Gamma counting efficiency of isotope one

τ_1 = Gamma transmission to detector

The isotopic ratio is given by the following equation:

$$A_1/A_2 = I_1 \lambda_2 B_2 \epsilon_2 \tau_2 / I_2 \lambda_1 B_1 \epsilon_1 \tau_1$$

Where:

A_1/A_2 = isotopic ratio

$\lambda_1 = .6932/T_1$ = decay constant

T_1 = material one half-life (in seconds)

I_1 and I_2 = measured peak intensities from isotope 1 and 2 respectively

T_1 and T_2 = half lives, in the same time units, of isotope 1 and 2 respectively

ϵ_1 and ϵ_2 = gamma counting efficiencies of isotope 1 and 2 respectively

B_1 and B_2 = branching ratios for characteristic gamma rays of isotope 1 and 2 respectively

Analysis is greatly simplified by the following observations:

$\epsilon_2 \tau_2 / \epsilon_1 \tau_1 \approx 1$ if the two gammas are close to the same energy

$\Omega_1 = \Omega_2$ The fractional solid angle of detector is the same for both gammas and cancels out.

λ_1 and λ_2 are known from the previously measured half lives.

B_1 and B_2 are known from the previously measured branching ratios

I_1 and I_2 have to be determined extremely accurately to get precise isotopic ratios.

The analysis proceeds on the assumptions that the solid angle terms cancel out and the half-lives and branching ratios of the respective gamma rays and x-rays can be determined. The efficiencies for detecting gamma rays are harder to determine; involving the intrinsic detector efficiencies and the overall detector and counting geometry used to obtain the data. Gamma and x-ray transmissions are nearly equal for energies close to each other. Fortunately, for gammas and x-rays close in energy the ratio of these terms, $\epsilon_2\tau_2 / \epsilon_1\tau_1 \approx 1$. Approximate detector efficiencies and gamma transmission corrections are used to make first order corrections to this ratio. The accuracy of determining the isotopic ratio, A_1/A_2 , is largely determined by the accuracy of determining the respective peak intensities, I_1/I_2 . The U235 code determines these peak intensities as accurately as possible and then applies the second order corrections for efficiency and transmission differences between the ratioed peaks to get as accurate an answer as possible.

A literature survey of ^{235}U and ^{238}U gammas and their daughter products found a number of peak groupings that could potentially provide information on isotopic ratios. These are listed in Appendix A. The survey also found that the state of information on the branching ratios and energies of ^{235}U and ^{238}U and their daughter products gammas (and x-rays) were of mixed quality.

The ^{235}U Analysis Code—U235

The ^{235}U Analysis Code—U235 was written to accomplish three main goals: first develop a tool that can accurately determine uranium isotopic ratios; second extend the applicability of the code to very low ^{235}U concentrations (depleted sources) and third to very high ^{235}U concentration (enriched) sources. Presently the code works for uranium samples that are 0.05% ^{235}U to 95% ^{235}U . Code algorithms have been found that very precisely subtract the "background" signal and fit the observed peak shapes. X-rays were fit with a Voight profile, the shape resulting from the Lorenzian profile emitted by the x-rays and the gaussian detector response. Gammas were fitted with a gaussian profile and a low energy exponential tail. (see Figure 5 and 6 for examples of these profiles). Techniques were developed to unfold the complex peak multiplets observed in the spectra using mathematical descriptions of the peak shapes and utilizing Taylor series minimization to fit the observed data as accurately as possible.

There are several potential energy regions in the uranium gamma ray spectra that can be used to calculate isotopic abundance ratios. In this report only gammas (and x-rays) less than 300 keV are considered. This energy region is measured by a "typical" low energy Ge detector set with a gain of .075 keV/channel and 4096 channels of data. A number of $^{235}\text{U} + ^{238}\text{U}$ spectra were available for this detector arrangement, making it a "natural" first place to start analysis. The only serious limitation this energy range imposes is the relative few ^{238}U (and daughters) peaks less than 300 keV. Fortunately there are two relatively strong $^{238}\text{U}/^{234}\text{Th}$ lines at 92.365 and 92.790 keV and a relatively strong IC x-ray at 93.356 keV $^{235}\text{U}/\text{Th-}\kappa\alpha 1$. (See Appendix A) Higher energy analysis suffers from the opposite problem, there are very few useful ^{235}U peaks above 205 keV. One of the disadvantages of using gammas in the 80 to 300 keV range is their limited transmission through "thick" material. This restricts the applicability of the analysis procedures to homogenous sources or "thin" heterogeneous uranium sources.

The 20--80 keV Energy Region

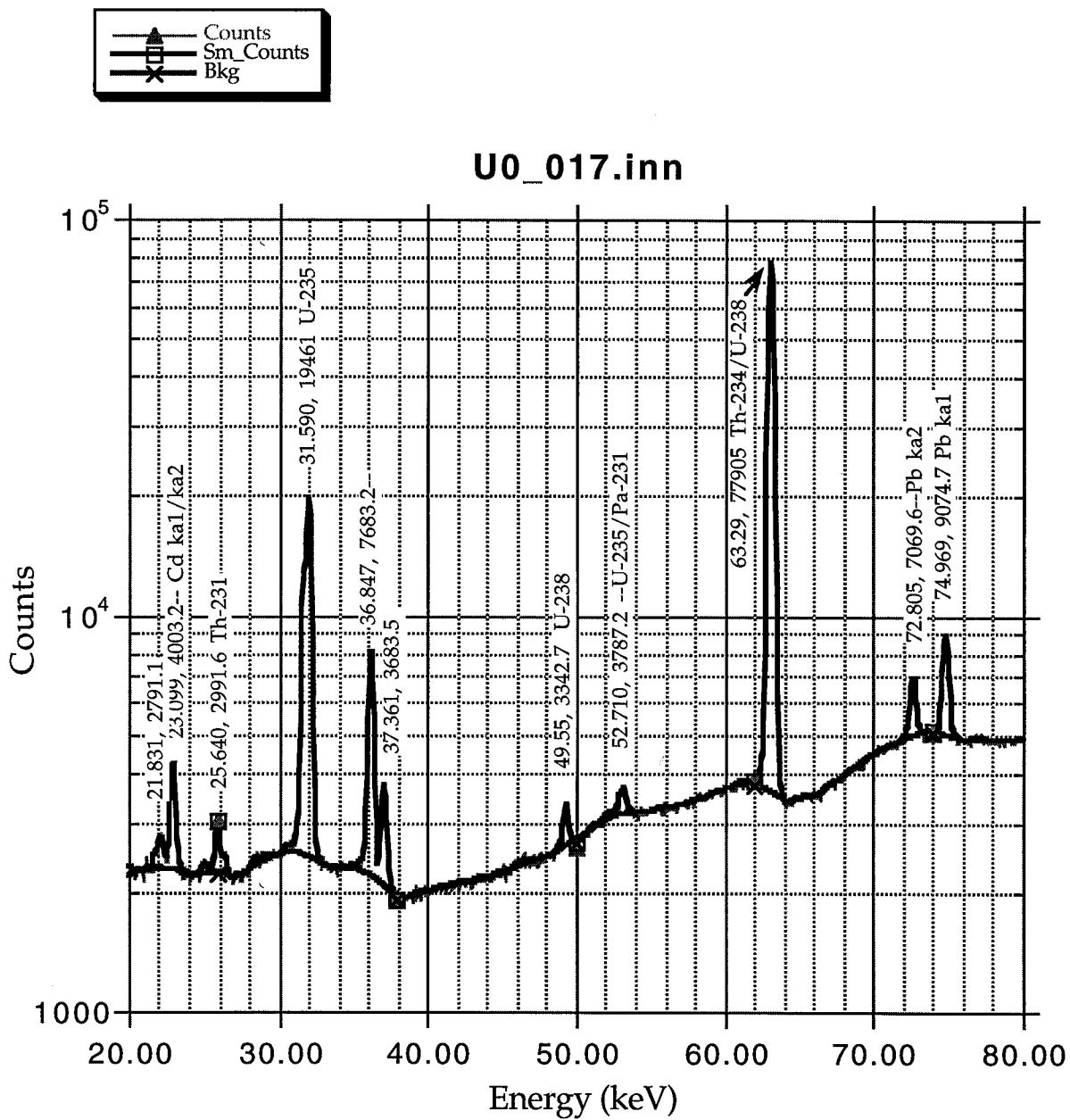


Figure 2. shows a plot of a 99.983% ^{238}U spectrum from 20 -80 keV. Clearly shown is the only gamma directly associated with the ^{238}U decay--the 49.55 keV peak. The strongest line in this region is the $^{238}\text{U} \rightarrow ^{234}\text{Th}$ daughter line at 63.29 keV. The first number on the peak labels above shows their energy, the second number their approximate peak counts and the third their origin.

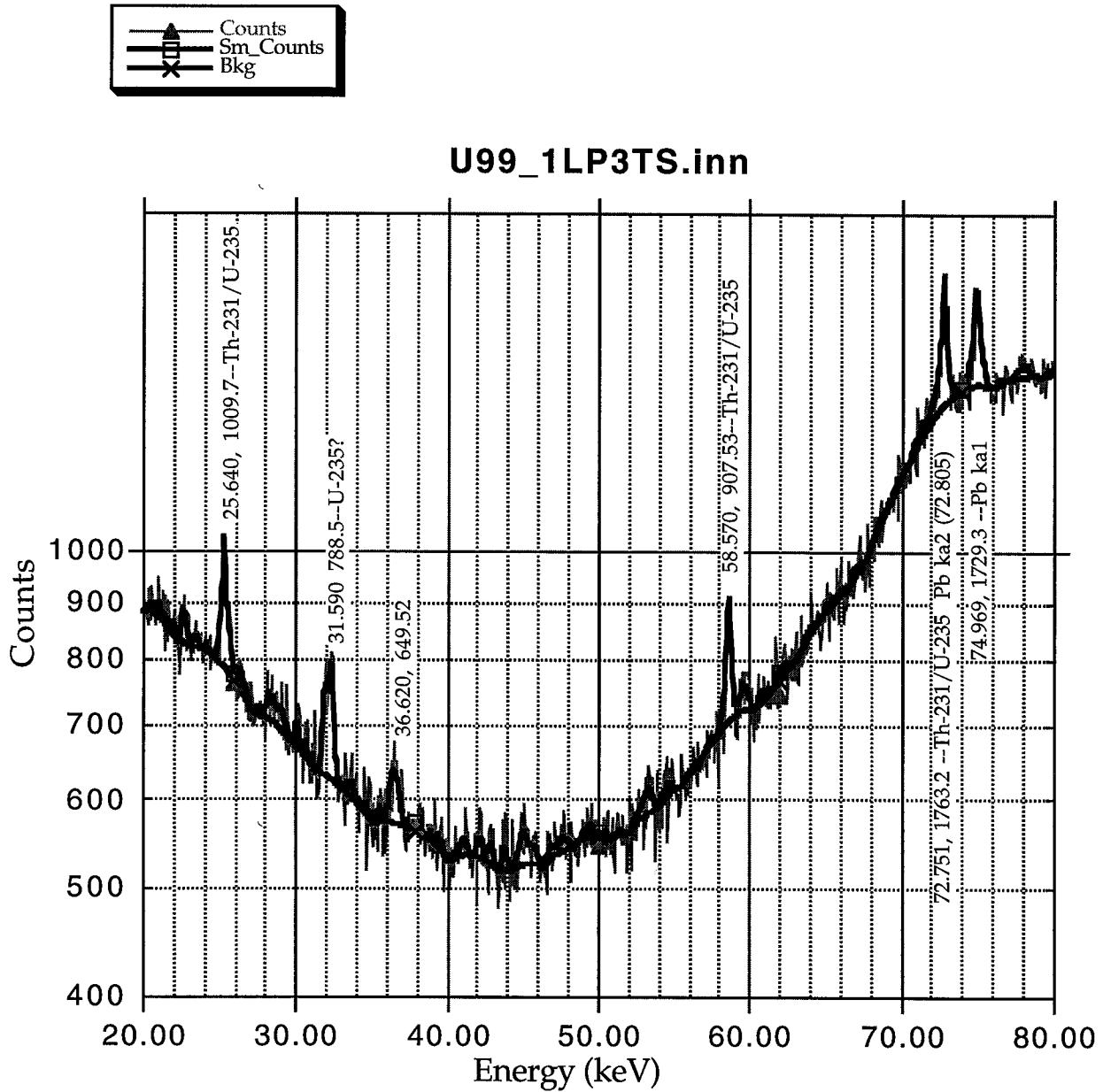


Figure 3. Shows the spectrum from a 99.1% sample of ^{235}U . It is considerably different than a ^{238}U spectrum shown above. There are no strong lines from ^{235}U or its daughters in this region. The lead ka1 and ka2 x-ray lines are a “typical” spectral contaminant resulting from fluorescent x-rays in the collimation etc.

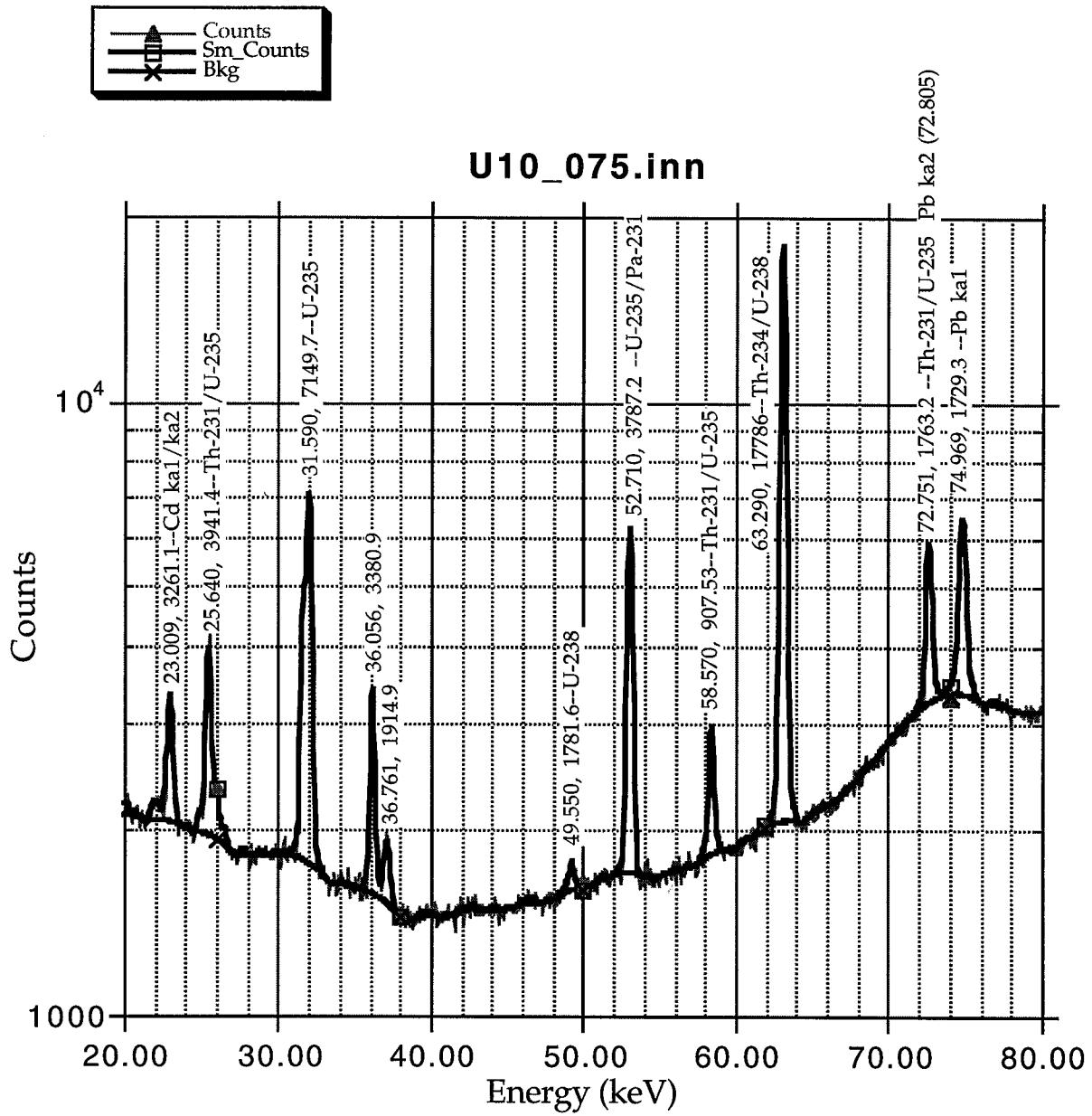


Figure 4.. Plot shows a 10.075% ^{235}U , 90.0 % ^{238}U spectrum from 20-80 keV. In this region is shown the only pure ^{238}U peak at 49.55 keV.

The 80-85 keV Energy Region

The lowest energy range of “practical” use is the 81-85 keV region. It contains peaks due to ^{235}U (81.228, 82.087, and 84.214 keV) as well as a 83.300 keV peak due to $^{238}\text{U} / ^{234}\text{Th}$ decay. Figure 5. shows the spectrum of a 10.075% ^{235}U sample. Even though the lead x-rays are weak they are a typical “contaminant” to spectra in this region and have to be accounted for accurate peak intensity determinations. The 83.30 $^{238}\text{U} / ^{234}\text{Th}$ peak is quite weak making its accurate determination difficult.

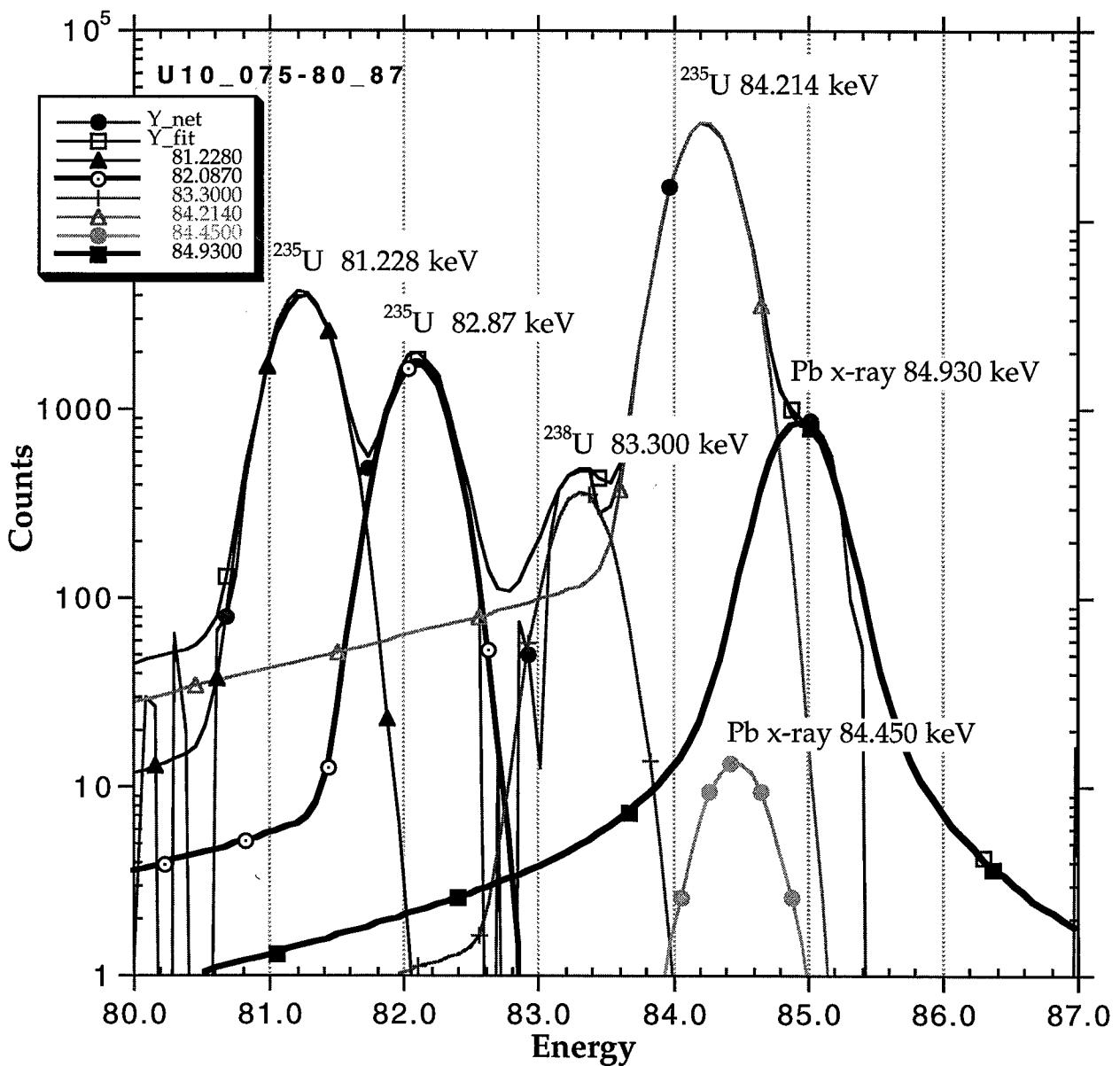


Figure 5. The net (minus background) uranium spectrum (10.075% ^{235}U) from 80 to 87 keV. As can be seen the $^{238}\text{U} / ^{234}\text{Th}$ 83.300 peak is quite weak, making good peak intensity measurements difficult for this sample and/or lower concentrations of ^{238}U .

The 87-100 keV Energy Region

This region has three peaks due to ^{238}U , a number of ^{235}U peaks and the two strong U $\text{k}\alpha 1$ and $\text{k}\alpha 2$ x-ray peaks. The tight clustering of peaks requires careful peak fitting and analysis. For most concentrations this is the region of primary interest since the 92.365 and 92.790 ^{238}U peaks are very near the 93.356 Th- $\text{k}\alpha 1$ / ^{235}U peak. The Th $\text{k}\alpha 1$ and Th $\text{k}\alpha 2$ x-ray peaks, due to ^{235}U decay, bracket the ^{238}U doublet. The ^{238}U 95.85 peak is so weak and has so much interference from the Pa $\text{k}\alpha 1$ 95.89 keV peak as to be virtually useless as a diagnostic tool. The main limitations on using this energy range are that at both high ^{235}U and low ^{235}U concentrations the signals of either the ^{238}U peaks or the ^{235}U peaks are too small, respectively, to accurately determine.

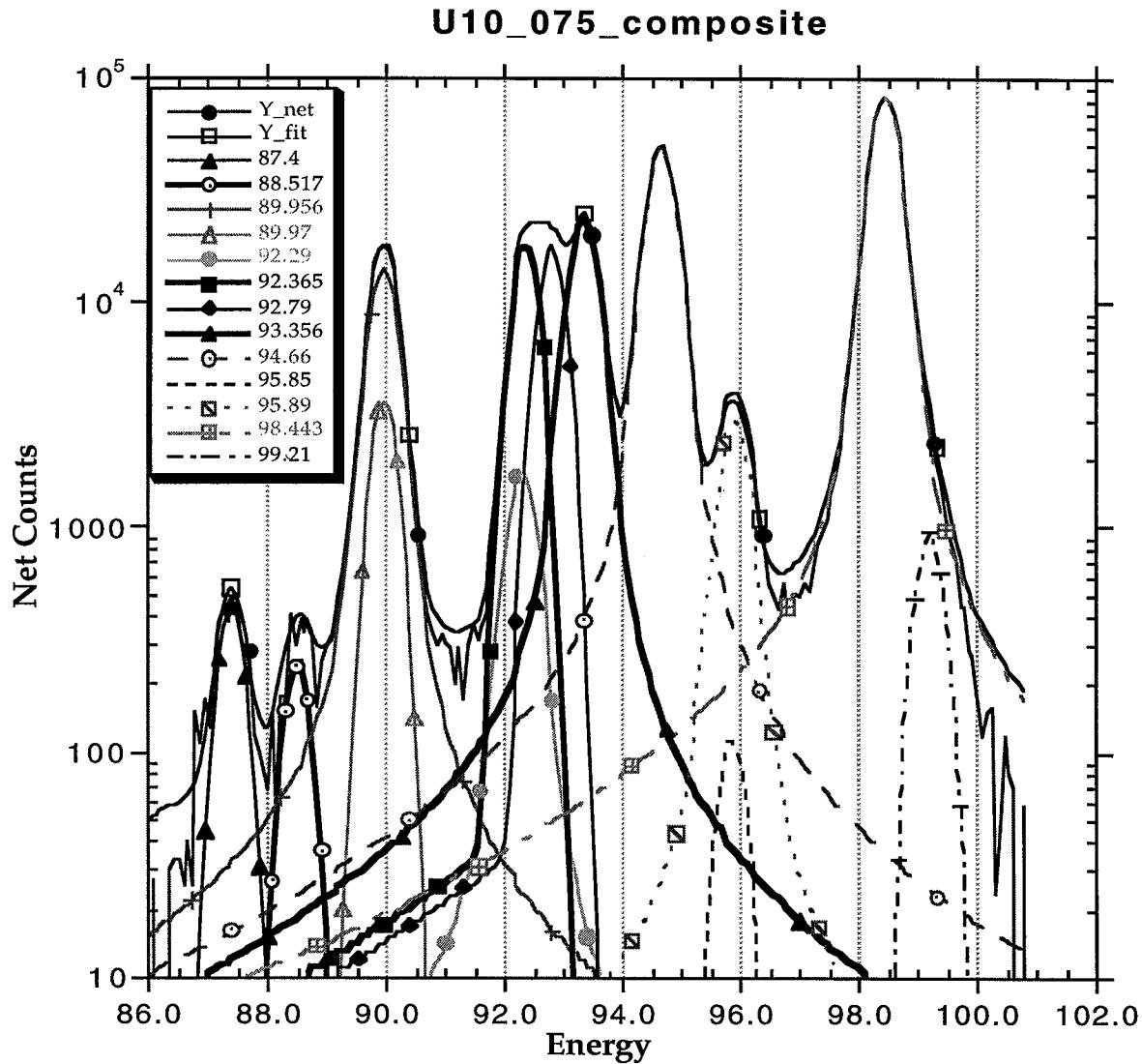


Figure 6a shows all 13 peaks used in fitting the data in the 86-102 keV region. Appendix 1 gives the identification of each of the energies and where they come from. Clearly seen is the gamma ray profile of the ^{238}U peaks and the much broader Voight x-ray profile of the ^{235}U /Th x-ray daughter peaks and the uranium x-rays.

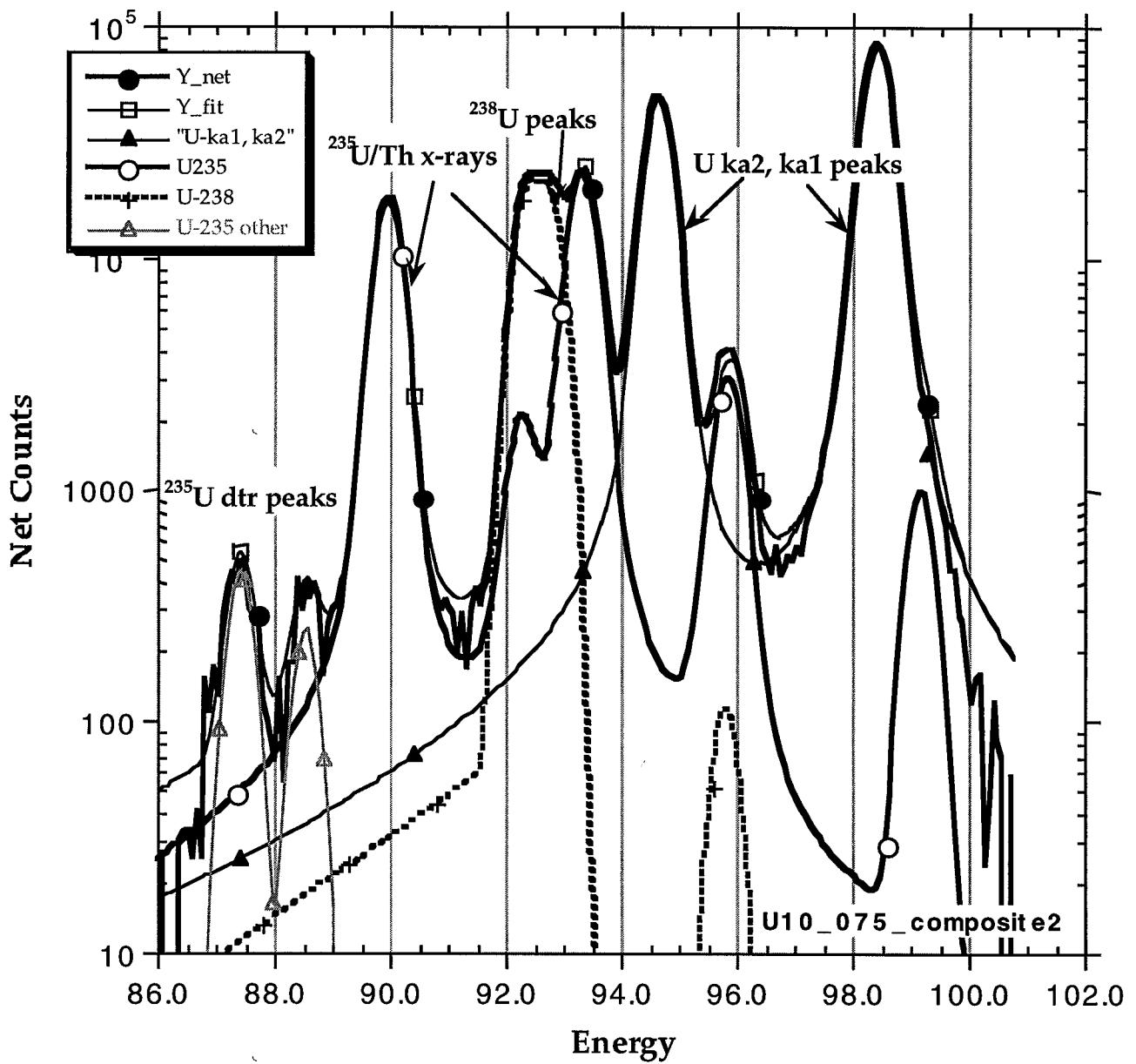


Figure 6b. The net count spectrum from 86 to 102 keV of a 10.075% ^{235}U sample with the peaks grouped into their respective components. At this ^{235}U concentration the ^{235}U and ^{238}U peaks are approximately equal. The fitting process uses both the protactinium and thorium x-rays from the ^{235}U daughters to find the best fit to the combined ^{235}U and ^{238}U spectrum.

The 100-118 keV Region

This region is very complex with 21 peaks containing all the $k\beta$ x-rays of U, Th and Pa plus a 109.2 keV gamma from ^{235}U and a 112.82 keV peak from ^{238}U . The large number of peaks and the overlap of peaks due to the wide Voight profile of the x-ray signals makes extracting useful peak ratios difficult. The Th and Pa x-ray peaks are tied to the ^{235}U decay and might give a useful data if good branching ratios were available—unfortunately they are not. Figure 7. shows the different x-ray multiplets in this region (each the sum of six x-ray peaks) and the two gamma rays. This energy region is not used in the analysis due to the difficult nature of the signals and the "poor" information available on branching ratios.

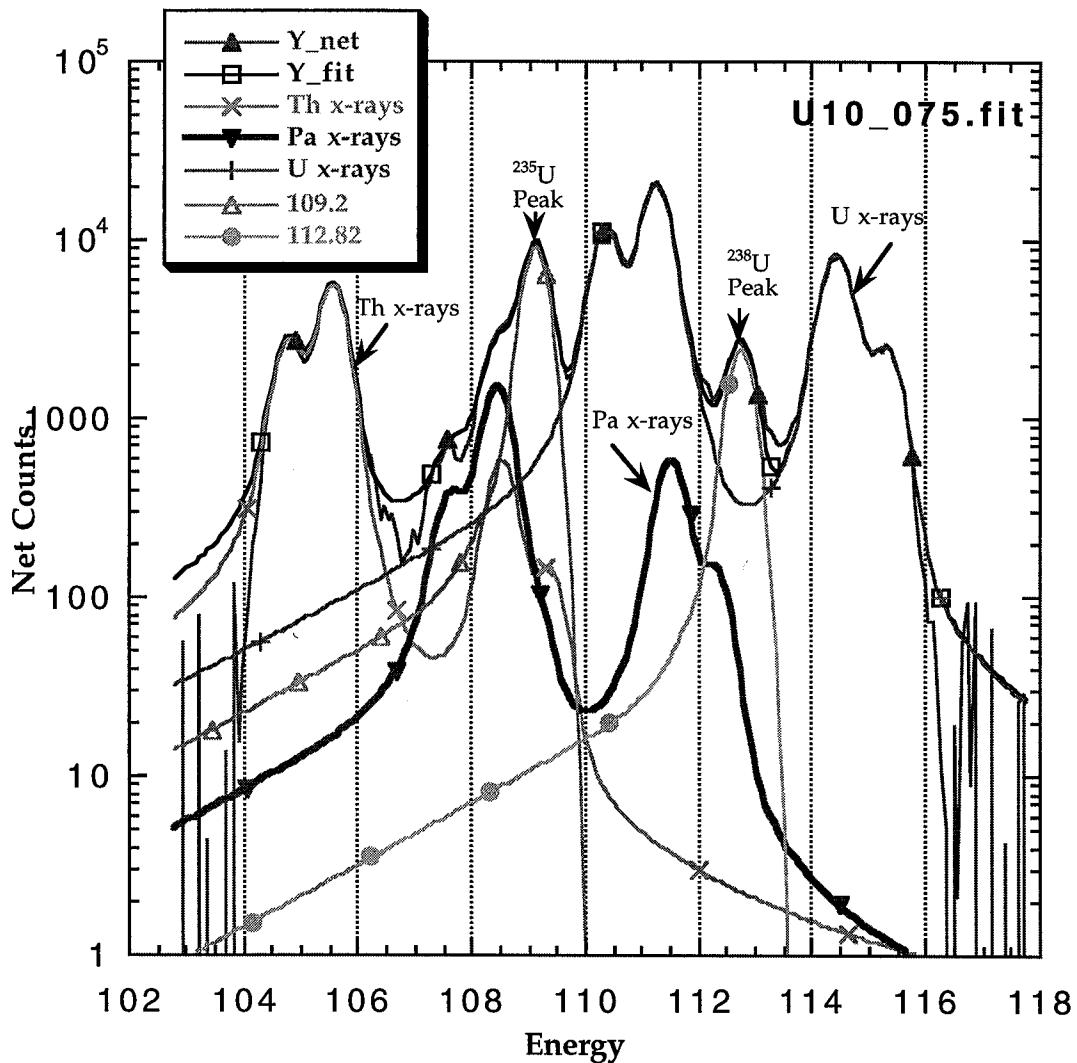


Figure 7. The net count spectrum from 102 to 118 keV of a 10.075% ^{235}U sample. In this energy range are 21 peaks, mostly thorium, protactinium and uranium $k\beta$ x-rays.

The 118-180 keV Region

The 118-180 keV region has relatively few peaks. The usually clean 120.90 keV peak of ^{234}U is useful for obtaining an estimate of that isotope. This peak suffers in that it is usually weak (sometimes too weak to analyze), giving poor statistic answers and there are no near by peaks to ratio it to. For good accuracy the 120.90 keV peak intensity needs to be corrected for efficiency and gamma transmission.

There are usually no ^{238}U /daughter peaks of sufficient intensity to be of interest in this region. The only exception is for depleted uranium spectra where the normally weak 131.300 keV $^{238}\text{U}/^{234}\text{Pa}$ peak is enhanced and the 143.760 keV ^{235}U peak is one of the cleanest ^{235}U peak available (see Figure 8b). Trying to dig the 93.35 keV $^{235}\text{U}/\text{Th}$ x-ray peak out of the 82-102 spectral regions is very inaccurate at very low ^{235}U concentrations and these isolated ^{235}U and ^{238}U /daughter peaks in the 118-180 keV region can be more accurately analyzed. The 143.76 and 163.33 keV ^{235}U peaks are also potentially useful for establishing the average material thickness in the sample by analyzing their relative intensities. Both of these techniques are utilized in the U235 code for low ^{235}U concentrations and transmission corrections.

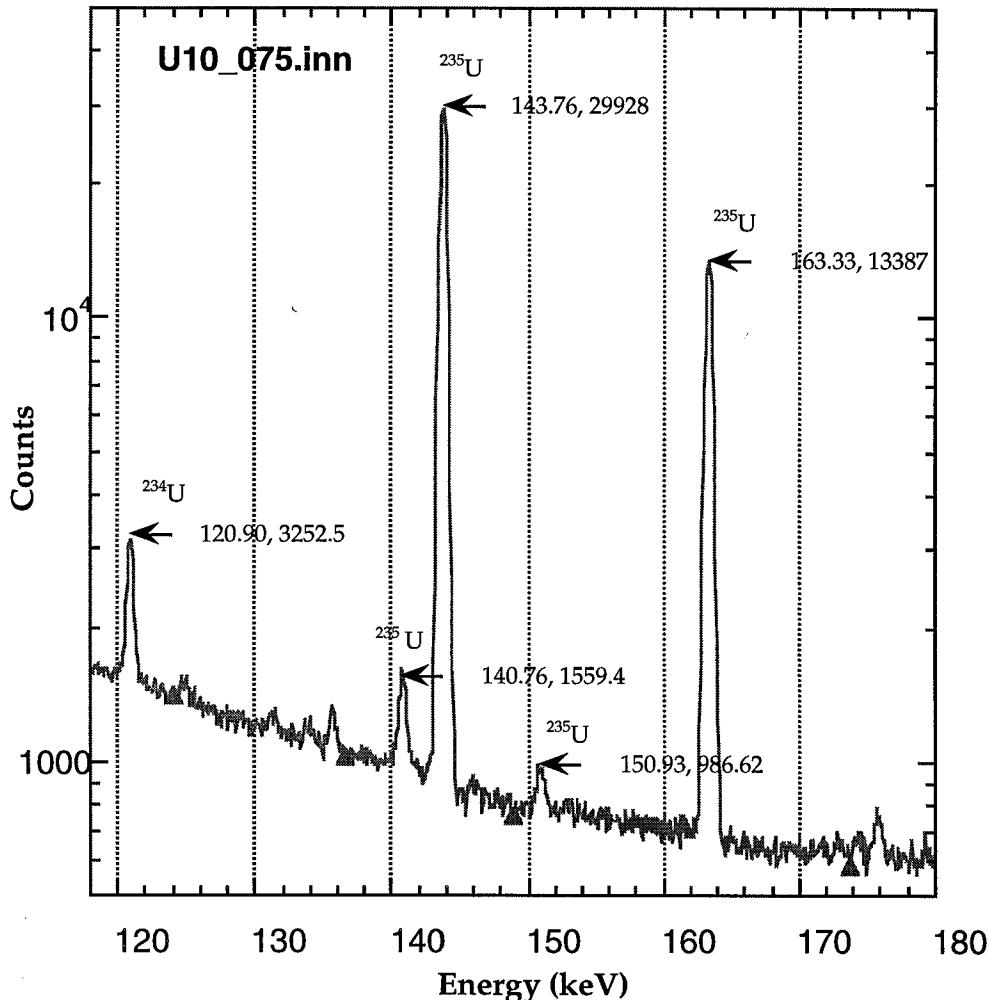


Figure 8a. The gamma spectrum from 118 to 180 keV of a 10.075% ^{235}U sample. Peaks are rather sparse in this region with usable ^{238}U peaks mainly showing up at low ^{235}U concentrations. The

^{234}U peak at 120.90 keV is usually quite weak but can usually be analyzed because of its isolation. Lack of good statistics on this peak may limit its accuracy.

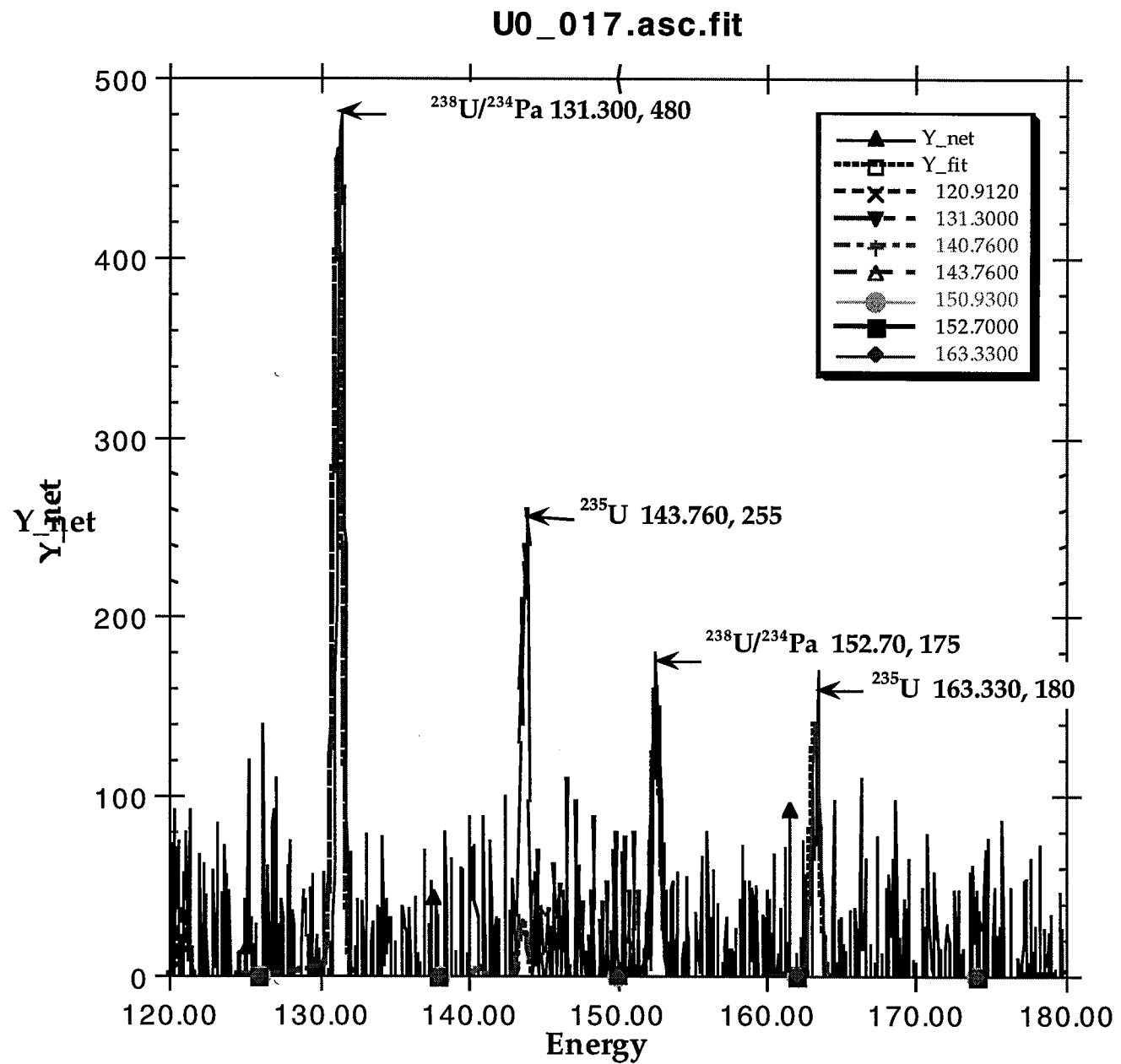


Figure 8b. showing the same 120-180 keV range for a 0.017 percent ^{235}U sample. this shows the ^{238}U peaks are enhanced; but the low count rate makes getting decent statistics for analysis very time consuming. The ^{234}U 120.90 keV peak is normally too weak to analyze at low ^{235}U concentrations.

The 180-210 keV Region

The 180-210 keV region has several prominent ^{235}U peaks including the most intense ^{235}U peak at 185.715 keV. This peak, in conjunction with the 98.443 keV uranium x-ray peak, is used to determine a more accurate gain and zero for the spectrum and to verify that ^{235}U is present in the spectrum. There are no easily observable ^{238}U peaks in this region. The only major uncertainty here is the 185.712 keV peak which has several other weak peaks around it that must be corrected for to get a good 185.715 keV peak intensity. One of the significant variations observed in this region is the 185.715 keV peak height to 188 keV background ratio. This ratio is found to vary from about 1000 at 90% enrichment to 1 at .02% enrichment. This change is attributed to the high energy gammas in ^{238}U decay and the contribution they make to the Compton continuum in this energy region. Further investigation is needed to see if these ratios are a function of sample, and detector type—or a universal characteristic of uranium spectra. This ratio is of potential usefulness in establishing a quick estimate of the ^{235}U enrichment. A spectrum with a high 185.712 to 188 keV ratio has almost certainly a "high" ^{235}U enrichment. Conversely a weak 185.715 peak with a high Compton continuum has a low ^{235}U enrichment.

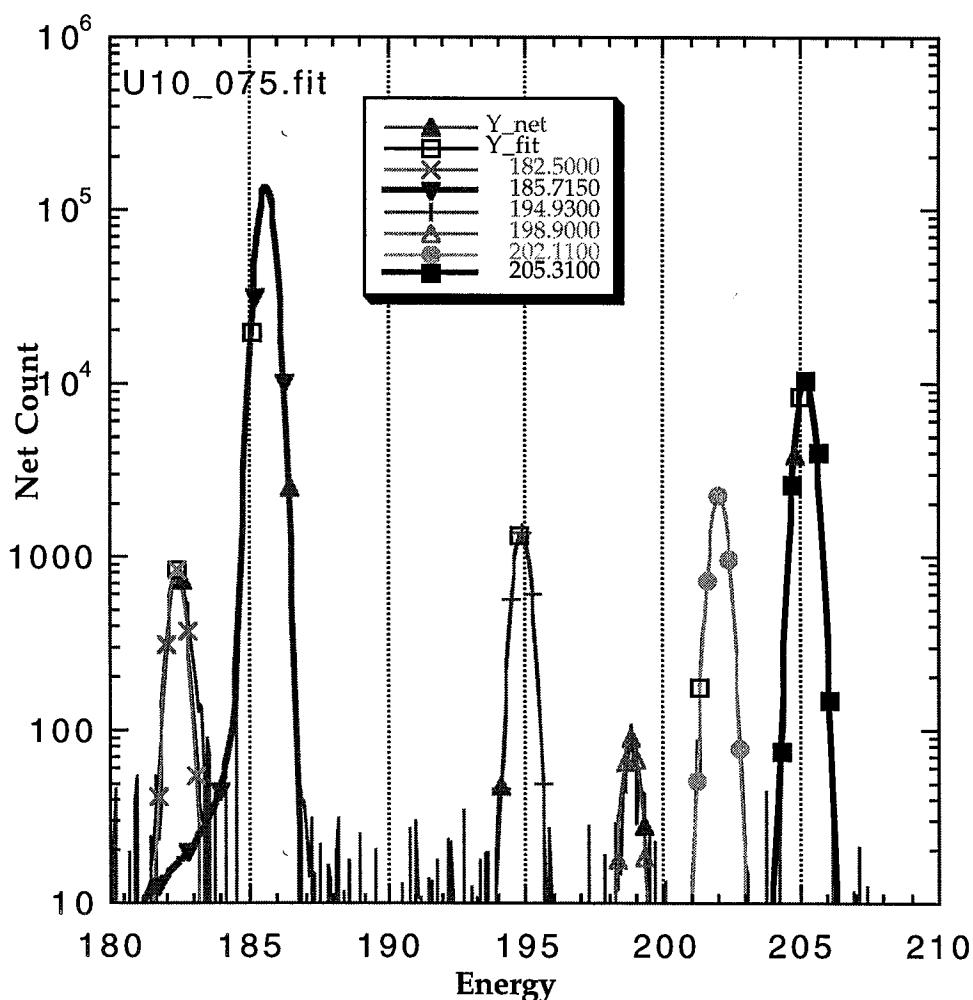


Figure 9. The net count spectrum from 180 to 210 keV of a 10.075% ^{235}U sample. In this region there are typically no ^{238}U peaks intense enough for any peak analysis.

The 210--300 keV Region

The 210-300 keV region only has one strong $^{238}\text{U}/^{234\text{m}}\text{Pa}$ peak at 258.2 keV. This peak is too weak to be of any great interest. The overall low intensity of this region lowers its utility in analyzing isotopic ratios.

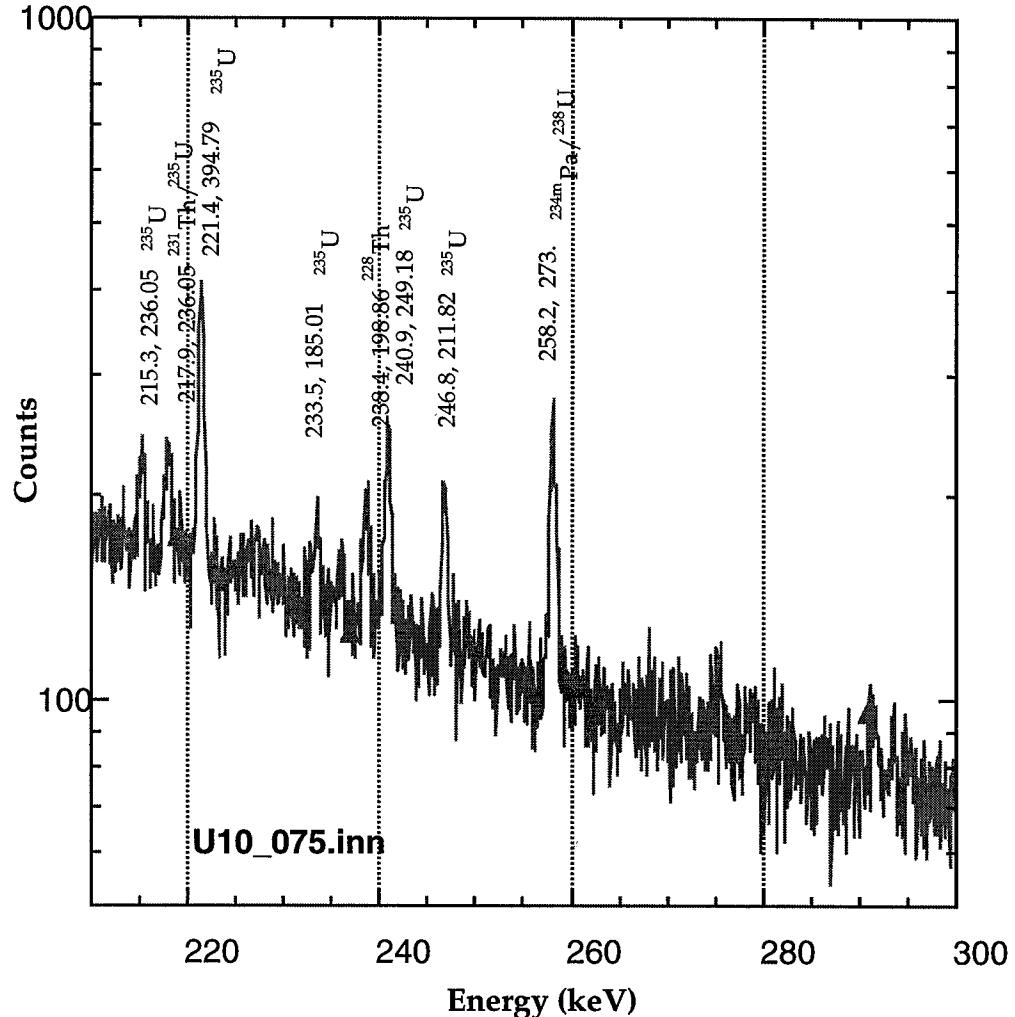


Figure 10.. The net count spectrum from 210 to 300 keV of a 10.075% ^{235}U sample. In this region there is only one ^{238}U peak of interest and a few ^{235}U peaks.

U235 Calculation Results

Table 1 shows the results of the U235 code calculations using the ^{238}U peaks at 92.29 and 92.7 keV, ratioed with the 93.356 keV $\text{Th k}\alpha 1/^{235}\text{U}$ peak. The first column shows the calculated enrichment on a number of standard (known) sources with the isotopic composition shown in column 2-4. The data in the last column shows the percent difference between the measured and calculated enrichment. As expected the samples that are the most difficult to analyze are where there is either very small ^{235}U or ^{238}U signals.

The code has been modified at low ^{235}U concentrations to use the 131.3 keV $^{238}\text{U}/^{234}\text{Pa}$ line and the ^{235}U lines at 143.76 and 163.33 keV to obtain a more accurate estimate of the $^{235}\text{U}/^{238}\text{U}$ ratio for low ^{235}U concentrations. Normally the 131.3 keV $^{238}\text{U}/^{234}\text{Pa}$ line is too weak to analyze; but at low concentrations the long counting time required to accumulate enough counts combined with the low ^{235}U emission rate make this line analyzable. It has the additional advantage that there are no competing lines near it; hence even though weak, it is unambiguous to analyze.

Table 1
U235 Code Preliminary Calculation Results

| Calc. U235 | Standard % ^{235}U | Standard % ^{234}U | Standard % ^{238}U | Std-Calc % diff. |
|----------------------|--|--|--|---------------------|
| 0.0169 | 0.017 | 0.001 | 99.982 | -0.59% |
| 0.5133 | 0.483 | 0.005 | 99.512 | 6.27% |
| 1.0126 | 0.991 | 0.008 | 99.001 | 2.18% |
| 2.0302 | 2.013 | 0.016 | 97.971 | 0.85% |
| 2.9991 | 3.009 | 0.031 | 96.960 | -0.33% |
| 5.0127 | 4.949 | 0.041 | 95.010 | 1.29% |
| 10.0806 | 10.075 | 0.077 | 89.848 | 0.06% |
| 49.1959 | 49.380 | 0.672 | 49.948 | -0.37% |
| 75.4537 | 75.130 | 0.584 | 24.286 | 0.43% |
| 92.982 | 93.076 | 1.452 | 5.472 | -0.10% |

Summary:

Algorithms are being developed to better subtract backgrounds and improve the high concentration ^{235}U results. Future plans include expanding the coverage to higher than 300 keV for more complicated spectra and expanding the code to handle mixed Pu and U spectra better.

Appendix A
Gamma and X-ray Decay of ^{235}U and ^{238}U and their Daughters
 (data from 49-300 keV)

The gammas energies listed in bold are used with these branching ratios listed in column three to determine ^{235}U / ^{238}U / ^{234}U ratios by the U235 code. All x-rays listed as IC-decay are internally converted in the isotope and decay with the isotopes decay characteristics—half-life and isotopic composition. All x-rays labeled as fluorescence are caused (nearly completely) by photoelectric absorption in the material and subsequent L-K shell electron decay. These x-rays are characteristic of the physical properties of the material and not its isotopic composition. Gammas and x-rays are listed by energy. This should allow quicker identification of observed spectra and may help pinpoint potential interference in a given measurement. Only the "strongest" lines are listed, many other gammas in this range are normally too weak to observe. These lines will occur with different intensities depending upon the isotopic concentration being observed. The branching ratios listed in column three are the ones presently used at LLNL.

| | | Branch Ratio — E(keV) | g=0 BR x 100 | x=-1 | Source | Parent | URADOS ^c Branch Ratio - BR x 100 | IAEA ^a Branch Ratio x100 |
|----------------|---------------|-----------------------------|-----------------|------|----------------|----------------|---|---|
| Group-1 | | | | | | | | |
| 1 | 49.550 | 0.064 | 0 | | U-238 | U-238 | | 0.064 ± 0.008 |
| 2 | 53.200 | 0.123 | 0 | | U-234 | | | 0.123 ± 0.002 |
| 3 | 58.570 | 0.500 | 0 | | Th-231 | U-235 | 0.46 ± 0.060 | 0.5 ± 0.05 |
| 4 | 63.290 | 4.470 | 0 | | Th-234 | U-238 | 3.94 ± 0.010 | 4.47 ± 0.88 |
| 5 | 72.751 | 0.260 | 0 | | Th-231 | U-235 | | 0.26 ± 0.02 |
| 6 | 72.804 | 27.700 | -1 | | Pb- κ 2 | fluorescence | | |
| 7 | 73.920 | 0.202 | 0 | | Pa-234m | U-238 | | |
| 8 | 74.000 | 0.036 | 0 | | Pa-234 | U-238 | | |
| 9 | 74.910 | 0.510 | 0 | | U-235 | U-235 | | |
| 10 | 74.969 | 46.200 | -1 | | Pb- κ 1 | fluorescence | | |
| 11 | 81.228 | 0.850 | 0 | | Th-231 | U-235 | | 0.85 ± 0.03 |
| 12 | 82.087 | 0.370 | 0 | | Th-231 | U-235 | | 0.37 ± 0.02 |
| 13 | 83.300 | 0.073 | 0 | | Th-234 | U-238 | 0.064 ± 0.10 | 0.073 |
| 14 | 84.214 | 6.710 | 0 | | Th-231 | U-235 | | 6.71 ± 0.1 |
| 15 | 84.450 | 5.580 | -1 | | Pb- κ 3 | fluorescence | | |
| 16 | 84.930 | 10.700 | -1 | | Pb- κ 1 | fluorescence | | |
| 17 | 87.300 | 3.910 | -1 | | Pb- κ 2 | fluorescence | | |
| Group-2 | | | | | | | | |
| | E(keV) | BR x 100 | g=0,x=1 | | Source | Parent | URADOS | IAEA |
| 1 | 87.700 | 0.050 | 0 | | Pa-231 | Th-231/U-235 | | |
| 2 | 88.500 | 0.030 | 0 | | Th-227 | U-235 | | |
| 3 | 89.956 | 3.360 | -1 | | Th κ 2 | U-235 IC-decay | 3.17 ± 0.08 | 3.4 ± 0.8 |
| 4 | 89.970 | 0.742 | 0 | | Th-231 | U-235 | 0.97 ± 0.05 | |
| 5 | 92.290 | 0.470 | -1 | | Pa κ 2 | U-235 IC-decay | 0.451 ± 0.036 | 0.39 ± 0.03 |
| 6 | 92.365 | 2.600 | 0 | | Th-234 | U-238 | 2.52 ± 0.06 | 2.60 ± 0.53 |
| 7 | 92.790 | 2.560 | 0 | | Th-234 | U-238 | 2.50 ± 0.06 | 2.56 ± 0.52 |
| 8 | 93.356 | 5.500 | -1 | | Th κ 1 | U-235 IC-decay | 5.22 ± 0.14 | 5.6 ± 1.3 |
| 9 | 94.660 | 9.161†† | -1 | | U κ 2 | fluorescence | 61.2 (norm) | 28.2 ± 0.6 |

| | | | | | | | |
|----|--------|----------|----|----------------|----------------|---------------|-------------|
| 10 | 94.700 | 0.0321 | | Pa-234 | U-238 | | |
| 11 | 95.850 | 0.0024 | 0 | Th-234 | U-238 | | |
| 12 | 95.860 | 0.880 | -1 | Pa $k\alpha 1$ | U-235 IC-decay | 0.776 ± 0.043 | 0.63 ± 0.05 |
| 13 | 98.443 | 14.800†† | -1 | U $k\alpha 1$ | fluorescence | 100.0 (norm) | 45.1 ± 0.9 |
| 14 | 99.270 | 0.400 | 0 | Th-231 | U-235 | 0.14 ± .03 | |

| Group-3 | E(keV) | BR x 100 | g=0,x=1 | Source | Parent | URADOS | IAEA |
|---------|---------|----------|---------|---------------|--------------|--------|---------------|
| 1 | 102.270 | 0.400 | 0 | Th-231 | U-235 | | 0.40 ± 0.02 |
| 2 | 104.819 | 0.137 | -1 | Th $k\beta 3$ | IC-decay | | |
| 3 | 105.604 | 0.262 | -1 | Th $k\beta 1$ | IC-decay | | |
| 4 | 106.239 | 0.009 | -1 | Th $k\beta 5$ | IC-decay | | |
| 5 | 107.595 | 0.022 | -1 | Pa $k\beta 3$ | IC-decay | | |
| 6 | 108.422 | 0.042 | -1 | Pa $k\beta 1$ | IC-decay | | |
| 7 | 108.582 | 0.100 | -1 | Th $k\beta 2$ | IC-decay | | |
| 8 | 108.955 | 0.003 | -1 | Th $k\beta 4$ | IC-decay | | |
| 9 | 109.072 | 0.002 | -1 | Pa $k\beta 5$ | IC-decay | | |
| 10 | 109.160 | 1.540 | 0 | U-235 | U-235 | | 1.54 ± 0.05 |
| 11 | 109.442 | 0.022 | -1 | Th KO2_3 | IC-decay | | |
| 12 | 110.480 | 0.555 | -1 | U $k\beta 3$ | fluorescence | | |
| 13 | 110.500 | .0043 | -1 | U-238 | U-238 | | |
| 14 | 111.350 | 1.000 | -1 | U $k\beta 1$ | fluorescence | | |
| 15 | 111.486 | 0.017 | -1 | Pa $k\beta 2$ | IC-decay | | |
| 16 | 111.870 | 0.001 | -1 | Pa $k\beta 4$ | IC-decay | | |
| 17 | 111.964 | 0.037 | -1 | U $k\beta 5$ | fluorescence | | |
| 18 | 112.380 | 0.004 | -1 | Pa KO2_3 | IC-decay | | |
| 19 | 112.820 | 0.040 | 0 | Th-234 | U-238 | | 0.256 ± 0.054 |
| 20 | 114.540 | 0.388 | -1 | U $k\beta 2$ | fluorescence | | |
| 21 | 114.844 | 0.011 | -1 | U $k\beta 4$ | fluorescence | | |
| 22 | 114.900 | 0.0064 | 0 | Pa-234m | U-238 | | |
| 23 | 115.377 | 0.089 | -1 | U KO23 | fluorescence | | |

| Group-4 | E(keV) | BR x100 | g=0,x=1 | Source | Parent | URADOS | IAEA |
|---------|---------|---------|---------|--------|--------|---------------|-------------------------|
| 1 | 120.900 | 0.0342 | 0 | U-234 | | 0.041 ± 0.006 | 0.0342 ± 0.0005 |
| 2 | 124.914 | 0.0600 | 0 | Th-231 | U-235 | | 0.06 ± 0.003 |
| 3 | 131.300 | 0.0286 | 0 | Pa-234 | U-238 | | |
| 4 | 134.030 | 0.0250 | 0 | Th-231 | U-235 | | 0.025 ± 0.005 |
| 5 | 135.664 | 0.0840 | 0 | Th-231 | U-235 | | 0.084 ± 0.007 |
| 6 | 140.760 | 0.2200 | 0 | U-235 | | | 0.22 ± 0.02 |
| 7 | 143.760 | 10.9600 | 0 | U-235 | | 10.95 ± 0.15 | 10.96 ± 0.08 |
| 8 | 150.930 | 0.0800 | 0 | U-235 | | | .08 ± 0.02 ^d |
| 9 | 152.700 | 0.0083 | | Pa-234 | U-238 | | |
| 10 | 163.330 | 5.0800 | 0 | U-235 | | 5.11 ± 0.05 | 5.08 ± 0.04 |

| Group-5 | E(keV) | BR x100 | g=0,x=1 | Source | Parent | URADOS | IAEA |
|---------|---------|---------|---------|--------|--------|-------------|-------------|
| 1 | 182.610 | 0.3400 | 0 | U-235 | | 0.37 ± 0.02 | 0.34 ± 0.02 |
| 2 | 183.500 | 0.0329 | 0 | U-235 | | | .0329 |
| 3 | 184.800 | 0.2200 | 0 | Th-234 | U-238 | | |
| 4 | 185.715 | 57.2000 | 0 | U-235 | | 57.2 ± 0.02 | 57.2 ± 0.5 |

| | | | | | | | |
|-----------------|---------------|----------------|----------------|---------------|---------------|---------------|---------------|
| 5 | 185.900 | 0.0039 | 0 | Pa-234 | U-238 | | 3.89E-3 |
| 6 | 194.940 | 0.6300 | 0 | U-235 | | | 0.630 ± 0.01 |
| 7 | 198.900 | 0.0420 | 0 | U-235 | | | 0.042 |
| 8 | 202.110 | 1.0800 | 0 | U-235 | | | 1.080 ± 0.02 |
| 9 | 205.311 | 5.0100 | 0 | U-235 | | | 5.010 ± 0.05 |
| Group -6 | E(keV) | BR x100 | g=0,x=1 | Source | Parent | URADOS | IAEA |
| 1 | 215.30 | 0.0288 | 0 | U-235 | | | |
| 2 | 217.94 | 0.0370 | 0 | Th-231 | U-235 | | 0.037 ± 0.001 |
| 3 | 221.38 | 0.1200 | 0 | U-235 | | | 0.12 ± 0.01 |
| 4 | 226.63 | 0.0059 | 0 | Pa-234 | U-238 | | |
| 5 | 227.17 | 0.0055 | 0 | Pa-234 | U-238 | | |
| 6 | 233.50 | 0.0290 | 0 | U-235 | | | |
| 7 | 238.50 | 0.0092 | 0 | Th-231 | U-235 | | |
| 8 | 240.85 | 0.0540 | 0 | U-235 | | | |
| 9 | 246.84 | 0.0530 | 0 | U-235 | | | |
| 10 | 258.20 | 0.0730 | 0 | Pa-234m | U-238 | | |
| 11 | 291.63 | 0.0180 | 0 | U-235 | | | |
| 12 | 293.90 | 0.0039 | 0 | Pa-234 | U-238 | | |

a. "Handbook of Nuclear Data for Safeguards", INDC (NDS) - 248, IAEA, 1991

b. GAMGEN code LLNL

c. Presented in CEA meeting by DAMPRI-LPRI May 1996

d. "Decay Data of the Transactinium Nuclides", Report # 261, IAEA, 1986

238U and Daughters— 234Pa and 234Th

GAMGEN Calculation showing gammas/sec/gram of ^{238}U (g/s/gm) at 5 years since separation and the implied Branching Ratio.(Branching Ratio normalized to 2.60% at 92.3 keV)

| E(keV) | g/s/gm | g=0,x=1 | BR x100 | Source | Parent1 | Emiter2 | Parent2 |
|--------|--------|---------|---------|---------|---------|---------|---------|
| 62.9 | 2.36 | 0 | 0.0182 | Th-234 | U-238 | | |
| 73.9 | 1.36 | 0 | 0.0105 | Pa-234m | U-238 | | |
| 74.0 | 5.30 | 0 | 0.0408 | Th-234 | U-238 | | |
| 83.3 | 8.71 | 0 | 0.0670 | Th-234 | U-238 | | |
| 92.3 | 338.0 | 0 | 2.6000 | Th-234 | U-238 | | |
| 92.8 | 335.0 | 0 | 2.5769 | Th-234 | U-238 | | |
| 94.7 | 21.6 | 0 | 0.1662 | Pa-234m | U-238 | Pa-234 | U-238 |
| 95.9 | 1.62 | 0 | 0.0125 | Th-234 | U-238 | | |
| 110.5 | 2.98 | 0 | 0.0229 | U-238 | U-238 | | |
| 114.9 | 4.32 | 0 | 0.0332 | Pa-234m | U-238 | Pa-234 | U-238 |
| 131.3 | 3.23 | 0 | 0.0248 | Pa-234 | U-238 | | |
| 152.7 | 1.08 | 0 | 0.0083 | Pa-234 | U-238 | | |
| 184.8 | 1.49 | 0 | 0.0115 | Th-234 | U-238 | | |
| 258.2 | 9.02 | 0 | 0.0694 | Pa-234m | U-238 | | |

Appendix B

Uranium and Daughter X-rays*

| Th X-rays | | Pa X-rays | | U X-rays | |
|-----------|--------|-----------|-------|----------|-------|
| E(keV) | % | E(keV) | % | E(keV) | % |
| 93.350 | 45.400 | 95.863 | 45.3 | 98.434 | 45.1 |
| 89.957 | 28.100 | 92.282 | 28.1 | 94.654 | 28.2 |
| 105.604 | 10.700 | 108.422 | 10.7 | 111.298 | 10.7 |
| 108.582 | 4.100 | 111.486 | 4.163 | 114.445 | 4.15 |
| 104.819 | 5.610 | 107.595 | 5.64 | 110.421 | 5.65 |
| 108.955 | 0.110 | 111.870 | 0.11 | 114.844 | 0.12 |
| 106.239 | 0.380 | 109.072 | 0.389 | 111.964 | 0.397 |
| 109.442 | 0.900 | 112.380 | 0.93 | 115.377 | 0.95 |
| | | | | | k02_3 |

% refers to percent decay per 100 k-shell vacancies

* Browne, E. and Firestone, R., "Table of Radioactive Isotopes", LBL, 2986, pg. C-23

X-rays associated with Uranium decay—sorted by energy.

Uranium x-ray fluorescence intensity is set to 1.00 for U-k α 1 in this comparison. The other branching ratios are derived from the observed decay of ^{235}U . The Pa branching ratios determined from the protactinium fluorescent decay ratios normalized to a 0.042 branching ratio for the Pa k β 1 line at 108.422 keV. The thorium branching ratios determined from the Th fluorescent decay ratios normalized to a .503 branching ratio for the Th k α 1 line at 93.350 keV. [All data preliminary.]

| E(keV) | % | Branch Ratio | Type | x-ray |
|---------|--------|--------------|------|--------------|
| 89.957 | 28.100 | 0.3113 | Th | k α 2 |
| 92.282 | 28.100 | 0.1103 | Pa | k α 2 |
| 93.350 | 45.400 | 0.5030 | Th | k α 1 |
| 94.654 | 28.200 | 0.6253 | U | k α 2 |
| 95.863 | 45.300 | 0.1778 | Pa | k α 1 |
| 98.434 | 45.100 | 1.0000 | U | k α 1 |
| 104.819 | 5.610 | 0.0622 | Th | k β 3 |
| 105.604 | 10.700 | 0.1185 | Th | k β 1 |
| 106.239 | 0.380 | 0.0042 | Th | k β 5 |
| 107.595 | 5.640 | 0.0221 | Pa | k β 3 |
| 108.422 | 10.700 | 0.0420 | Pa | k β 1 |
| 108.582 | 4.100 | 0.0454 | Th | k β 2 |
| 108.955 | 0.110 | 0.0012 | Th | k β 4 |
| 109.072 | 0.389 | 0.0015 | Pa | k β 5 |
| 109.442 | 0.900 | 0.0100 | Th | k02_3 |
| 110.421 | 5.650 | 0.1253 | U | k β 3 |
| 111.298 | 10.700 | 0.2373 | U | k β 1 |
| 111.486 | 4.163 | 0.0163 | Pa | k β 2 |
| 111.870 | 0.110 | 0.0004 | Pa | k β 4 |

| | | | | | |
|---------|-------|--------|----|------------|--------------|
| 111.964 | 0.397 | 0.0088 | U | $k\beta 5$ | fluorescence |
| 112.380 | 0.930 | 0.0037 | Pa | $k02_3$ | IC-decay |
| 114.445 | 4.150 | 0.0920 | U | $k\beta 2$ | fluorescence |
| 114.844 | 0.120 | 0.0027 | U | $k\beta 4$ | fluorescence |
| 115.377 | 0.950 | 0.0211 | U | $k02_3$ | fluorescence |

| E(keV) | Th X-rays | | Norm. meas. | Intensity % | Scofield Calculation | Th X-rays | | Calc-Meas. % diff |
|------------------|-----------|-------|-------------|-------------|----------------------|-----------|-------------|-------------------|
| | measured* | meas. | | | | | % | |
| 93.350 | 93.348 | 100 | | 45.4 | 100 | 45.40 | $k\alpha 1$ | 0 |
| 89.957 | 89.957 | 61 | | 27.694 | 61.9 | 28.10 | $k\alpha 2$ | 0 |
| 105.604 | 105.606 | 19 | | 8.626 | 22.35 | 10.70 | $k\beta 1$ | 2.47 |
| 108.582 | 108.471 | 10 | | 4.54 | 8.5601 | 4.10 | $k\beta 2$ | 2.5 |
| 104.819 | 104.822 | | | | 11.466 | 5.61 | $k\beta 3$ | 3.53 |
| 108.955 | | | | | | 0.11 | $k\beta 4$ | |
| 106.239 | | | | | 0.8247 | 0.38 | $k\beta 5$ | 0.68 |
| 109.442 | | | | | | 0.90 | $k02_3$ | |
| Pa X-rays | | | | | | | | |
| E(keV) | | | | | | % | | |
| 95.863 | 95.867 | 100 | | 45.3 | 100 | 45.300 | $k\alpha 1$ | 0 |
| 92.282 | 92.284 | 62 | | 28.086 | 62.2 | 28.100 | $k\alpha 2$ | -0.1 |
| 108.422 | 108.418 | 24 | | 10.872 | 22.45 | 10.700 | $k\beta 1$ | 2.36 |
| 111.486 | | | | 0 | 8.6882 | 4.163 | $k\beta 2$ | 2.62 |
| 107.595 | 107.586 | 11 | | 4.983 | 11.472 | 5.640 | $k\beta 3$ | 3.86 |
| 111.870 | | | | | | 0.110 | $k\beta 4$ | |
| 109.072 | | | | | 0.8441 | 0.389 | $k\beta 5$ | 0.78 |
| 112.380 | | | | | | 0.930 | $k02_3$ | |
| U X-rays | | | | | | | | |
| E(keV) | | | | | | % | | |
| 98.434 | 98.435 | | | | 100 | 45.100 | $k\alpha 1$ | 0 |
| 94.654 | 94.656 | | | | 62.5 | 28.200 | $k\alpha 2$ | 0.02 |
| 111.298 | 111.300 | | | | 22.6 | 10.700 | $k\beta 1$ | 2.25 |
| 114.445 | | | | | 8.7462 | 4.150 | $k\beta 2$ | 2.35 |
| 110.421 | 110.416 | | | | 11.549 | 5.650 | $k\beta 3$ | 3.82 |
| 114.844 | | | | | | 0.120 | $k\beta 4$ | |
| 111.964 | 112.043 | | | | 0.8656 | 0.397 | $k\beta 5$ | 0.77 |
| 115.377 | | | | | | 0.950 | $k02_3$ | |

| | $k\alpha 2/k\alpha 1$ | $k\beta 1/k\alpha 1$ | $k\beta 3/k\beta 1$ | $k\beta 5/k\beta 1$ |
|----|-----------------------|----------------------|---------------------|---------------------|
| Th | 0.619 | 0.224 | 0.513 | 0.0369 |
| Pa | 0.622 | 0.225 | 0.512 | 0.0376 |
| U | 0.625 | 0.226 | 0.511 | 0.0383 |

Scofield, J. D. "Relativistic Hartree-Slater Values of the K and L X-ray Emission"
 Atomic Data and Nuclear Data Tables 14, 121-137(1974)

*Barreau, G. et. al., "Z. Phys. A. Atoms and Nuclei" 308, 209-213 (1982)

Appendix C
 $^{235}\text{U Daughter}$ — ^{231}Th Gammas, Pa x-rays, and Branching Ratios

| E(keV) | $\pm\Delta E$ | IAEA E(keV) | x=- 1 | * Rel. Int ^b | $\pm\Delta I$ | Notes: | Branch Ratio | Imp. Branch |
|---------|---------------|----------------|----------|----------------------------|---------------|------------------------|--------------|--------------------|
| | | | | | | | (IAEA) | Ratio ^a |
| 26.560 | | 25.640 | | 202 | 20 | | 0.146 | 0.135542 |
| 44.100 | 0.3 | | | 0.06 | 0.04 | | | 4.03E-05 |
| 58.470 | 0.05 | 58.570 | | 7.2 | 0.7 | | 0.005 | 0.004831 |
| 63.700 | 0.2 | | | 0.68 | 0.14 | | | 0.000456 |
| 72.660 | 0.06 | 72.751 | | 4 | 0.4 | | 0.0026 | 0.002684 |
| 73.000 | 0.1 | | | 0.1 | 0.04 | | | 6.71E-05 |
| 81.180 | 0.05 | 81.228 | | 14.2 | 1.4 | | 0.0085 | 0.009528 |
| 82.020 | 0.06 | 82.087 | | 7.2 | 0.7 | | 0.0037 | 0.004831 |
| 84.170 | | 84.214 | | 100 | | Reference ^a | 0.0671 | 0.067100 |
| 89.940 | 0.05 | | | 15.3 | 1.5 | | | 0.010266 |
| 92.230 | 0.05 | | -1 | 6 | 0.6 | Pa $\kappa\alpha 2$ | | 0.004026 |
| 93.300 | 0.1 | | | 0.5 | 0.05 | | | 0.000336 |
| 95.870 | 0.05 | | -1 | 10.3 | 1 | Pa $\kappa\alpha 1$ | | 0.006911 |
| 99.300 | 0.05 | | | 2.1 | 0.2 | | | 0.001409 |
| 102.300 | 0.05 | 102.270 | | 6.7 | 0.7 | | 0.004 | 0.004496 |
| 105.730 | 0.1 | | | 0.14 | 0.02 | | | 9.39E-05 |
| 106.580 | 0.1 | | | 0.34 | 0.04 | | | 0.000228 |
| 107.620 | 0.1 | | -1 | 1.29 | 0.14 | Pa $\kappa\beta 3$ | | 0.000866 |
| 108.490 | 0.1 | | -1 | 2.43 | 0.24 | Pa $\kappa\beta 1+5$ | | 0.001631 |
| 111.590 | 0.1 | | -1 | 0.9 | 0.1 | Pa $\kappa\beta 2$ | | 0.000604 |
| 112.460 | 0.1 | | -1 | 0.34 | 0.04 | Pa κ_0 | | 0.000228 |
| 115.500 | 0.2 | | | 0.04 | 0.01 | | | 2.68E-05 |
| 116.910 | 0.05 | | | 0.39 | 0.04 | | | 0.000262 |
| 125.100 | 0.05 | 124.914 | | 0.95 | 0.09 | | 0.0006 | 0.000637 |
| 134.140 | 0.08 | 134.030 | | 0.42 | 0.05 | | 0.00025 | 0.000282 |
| 135.770 | 0.06 | 135.664 | | 1.3 | 0.1 | | 0.00084 | 0.000872 |
| 136.780 | 0.2 | | | 0.09 | 0.03 | | | 6.04E-05 |
| 145.150 | 0.3 | | | 0.12 | 0.03 | | | 8.05E-05 |
| 146.000 | 0.07 | | | 0.58 | 0.06 | | | 0.000389 |
| 163.160 | 0.06 | | | 2.6 | 0.03 | | | 0.001745 |
| 164.940 | 0.1 | | | 0.06 | 0.03 | | | 4.03E-05 |
| 169.580 | 0.1 | | | 0.03 | 0.01 | | | 2.01E-05 |
| 174.190 | 0.08 | | | 0.31 | 0.03 | | | 0.000208 |
| 183.470 | 0.07 | | | 0.57 | 0.06 | | | 0.000382 |
| 188.770 | 0.2 | | | 0.08 | 0.01 | | | 5.37E-05 |
| 218.000 | 0.07 | 217.94 | | 0.67 | 0.07 | | 0.00037 | 0.00045 |
| 236.170 | 0.07 | | | 0.18 | 0.02 | | | 0.000121 |
| 240.400 | 0.2 | | | 0.005 | 0.0005 | | | 3.36E-06 |
| 242.600 | 0.1 | | | 0.013 | 0.0006 | | | 8.72E-06 |
| 249.800 | 0.3 | | | 0.01 | 0.002 | | | 6.71E-06 |
| 250.500 | 0.3 | | | 0.011 | 0.002 | | | 7.38E-06 |
| 267.800 | 0.07 | | | 0.023 | 0.0006 | | | 1.54E-05 |
| 308.900 | 0.3 | | | 0.008 | 0.001 | | | 5.37E-06 |
| 311.000 | 0.1 | | | 0.054 | 0.005 | | | 3.62E-05 |
| 318.000 | 0.4 | | | 0.002 | 0.0002 | | | 1.34E-06 |
| 320.200 | 0.3 | | | 0.004 | 0.0003 | | | 2.35E-06 |

a. Normalized to 0.0671 for 84.17 keV transition.

b. Browne, E and Asaro F. (Phys Rev C V7n6 p2545) find the 84.17 keV transition branching ratio = 0.070 \pm 0.003

235U and Daughters Gammas

| keV | Branch ratio* | g/s/gm | Emiter1 | Parent1 | Emiter2 | Parent2 |
|---------|---------------|----------|---------|---------|---------|---------|
| 11.400 | 0.03050 | 2.40E+03 | Th-231 | U-235 | | |
| 13.000 | 0.22367 | 1.76E+04 | U-235 | U-235 | Ac-227 | U-235 |
| 13.700 | 0.49817 | 3.92E+04 | Th-231 | U-235 | | |
| 14.500 | 0.00224 | 1.76E+02 | U-235 | U-235 | Ac-227 | U-235 |
| 15.000 | 0.00407 | 3.20E+02 | Th-231 | U-235 | | |
| 16.100 | 0.15250 | 1.20E+04 | U-235 | U-235 | | |
| 16.600 | 0.37617 | 2.96E+04 | Th-231 | U-235 | | |
| 17.200 | 0.00224 | 1.76E+02 | Th-231 | U-235 | | |
| 19.100 | 0.02643 | 2.08E+03 | U-235 | U-235 | | |
| 19.800 | 0.07523 | 5.92E+03 | Th-231 | U-235 | | |
| 25.600 | 0.14869 | 1.17E+04 | Th-231 | U-235 | | |
| 42.000 | 0.00061 | 4.80E+01 | U-235 | U-235 | | |
| 42.800 | 0.00059 | 4.64E+01 | Th-231 | U-235 | | |
| 58.600 | 0.00488 | 3.84E+02 | Th-231 | U-235 | | |
| 72.700 | 0.00112 | 8.80E+01 | U-235 | U-235 | | |
| 72.800 | 0.00255 | 2.01E+02 | Th-231 | U-235 | | |
| 74.800 | 0.00061 | 4.80E+01 | U-235 | U-235 | | |
| 81.200 | 0.00915 | 7.20E+02 | Th-231 | U-235 | | |
| 84.200 | 0.06710 | 5.28E+03 | Th-231 | U-235 | | |
| 90.000 | 0.03419 | 2.69E+03 | U-235 | U-235 | | |
| 90.000 | 0.00956 | 7.52E+02 | Th-231 | U-235 | | |
| 92.300 | 0.00397 | 3.12E+02 | Th-231 | U-235 | | |
| 93.400 | 0.05592 | 4.40E+03 | U-235 | U-235 | | |
| 95.900 | 0.00641 | 5.04E+02 | Th-231 | U-235 | | |
| 96.200 | 0.00087 | 6.88E+01 | U-235 | U-235 | | |
| 99.300 | 0.00122 | 9.60E+01 | Th-231 | U-235 | | |
| 102.300 | 0.00417 | 3.28E+02 | Th-231 | U-235 | | |
| 105.400 | 0.02008 | 1.58E+03 | U-235 | U-235 | | |
| 108.200 | 0.00231 | 1.82E+02 | Th-231 | U-235 | | |
| 109.000 | 0.00671 | 5.28E+02 | U-235 | U-235 | | |
| 109.200 | 0.01563 | 1.23E+03 | U-235 | U-235 | | |
| 111.900 | 0.00077 | 6.08E+01 | Th-231 | U-235 | | |
| 116.100 | 0.00071 | 5.60E+01 | U-235 | U-235 | | |
| 124.900 | 0.00057 | 4.48E+01 | Th-231 | U-235 | | |
| 135.700 | 0.00079 | 6.24E+01 | Th-231 | U-235 | | |
| 140.800 | 0.00224 | 1.76E+02 | U-235 | U-235 | | |
| 143.800 | 0.11133 | 8.76E+03 | U-235 | U-235 | | |
| 150.900 | 0.00081 | 6.40E+01 | U-235 | U-235 | | |
| 163.100 | 0.00158 | 1.24E+02 | Th-231 | U-235 | | |
| 163.300 | 0.05160 | 4.06E+03 | U-235 | U-235 | | |
| 182.600 | 0.00346 | 2.72E+02 | U-235 | U-235 | | |
| 185.700 | 0.58077 | 4.57E+04 | U-235 | U-235 | | |
| 194.900 | 0.00641 | 5.04E+02 | U-235 | U-235 | | |
| 198.900 | 0.00427 | 3.36E+02 | U-235 | U-235 | | |
| 202.100 | 0.01098 | 8.64E+02 | U-235 | U-235 | | |
| 205.300 | 0.05096 | 4.01E+03 | U-235 | U-235 | | |
| 221.400 | 0.00122 | 9.60E+01 | U-235 | U-235 | | |
| 240.900 | 0.00055 | 4.32E+01 | U-235 | U-235 | | |

* Normalized to .0671 at 84.214 keV

^{238}U and Daughters Gammas

| E(keV) * | Branch Ratio | uncert. | g/s/gm | Emiter1 | Parent1 |
|----------|--------------|---------|----------|---------|---------|
| | % * | % * | | | |
| 63.24 | 3.6000 | 3 | 4.73E+02 | Th-234 | U-238 |
| 131.31 | 0.0286 | 1.4 | 3.23E+00 | Pa-234 | U-238 |
| 152.76 | 0.0083 | 3.7 | 1.08E+00 | Pa-234 | U-238 |
| 203.12 | 0.0027 | 8 | 3.37E-01 | Pa-234 | U-238 |
| 226.85 | 0.0167 | 1.3 | 9.54E-01 | Pa-234 | U-238 |
| 249.21 | 0.0035 | 4.7 | 4.53E-01 | Pa-234 | U-238 |
| 258.26 | 0.0730 | 0.46 | 9.02E+00 | Pa-234m | U-238 |
| 272.20 | 0.0018 | 9.1 | 1.62E-01 | Pa-234 | U-238 |
| 293.74 | 0.0049 | 3.1 | 6.31E-01 | Pa-234 | U-238 |
| 369.52 | 0.0044 | 3.5 | 4.69E-01 | Pa-234 | U-238 |
| 372.02 | 0.0023 | 6.9 | 2.10E-01 | Pa-234 | U-238 |
| 450.96 | 0.0030 | 5.2 | 3.36E-01 | Pa-234m | U-238 |
| 453.58 | 0.0019 | 8.4 | 2.71E-01 | Pa-234m | U-238 |
| 458.63 | 0.0020 | 8 | 2.43E-01 | Pa-234 | U-238 |
| 468.44 | 0.0023 | 6.8 | 2.63E-01 | Pa-234m | U-238 |
| 475.75 | 0.0023 | 6.5 | 3.18E-01 | Pa-234m | U-238 |
| 506.70 | 0.0035 | 5.5 | 2.59E-01 | Pa-234 | U-238 |
| 543.98 | 0.0036 | 4.7 | 4.14E-01 | Pa-234m | U-238 |
| 569.30 | 0.0203 | 1.3 | 1.73E+00 | Pa-234 | U-238 |
| 654.37 | 0.0022 | 7.6 | 9.70E-02 | Pa-234 | U-238 |
| 666.42 | 0.0015 | 9.8 | 2.59E-01 | Pa-234 | U-238 |
| 669.64 | 0.0017 | 8.9 | 2.26E-01 | Pa-234 | U-238 |
| 691.08 | 0.0090 | 2.1 | 8.75E-01 | Pa-234m | U-238 |
| 699.02 | 0.0059 | 2.6 | 7.44E-01 | Pa-234 | U-238 |
| 702.05 | 0.0071 | 2.4 | 8.59E-01 | Pa-234m | U-238 |
| 705.90 | 0.0065 | 2.4 | 9.16E-01 | Pa-234 | U-238 |
| 733.38 | 0.0115 | 1.5 | 1.39E+00 | Pa-234 | U-238 |
| 737.88 | 0.0021 | 8.3 | 1.62E-01 | Pa-234 | U-238 |
| 739.95 | 0.0118 | 2.1 | 1.13E+00 | Pa-234m | U-238 |
| 742.77 | 0.0946 | 0.7 | 9.38E+00 | Pa-234m | U-238 |
| 755.00 | 0.0021 | 8.1 | 1.62E-01 | Pa-234 | U-238 |
| 766.37 | 0.3220 | 0.65 | 3.29E+01 | Pa-234m | U-238 |
| 781.73 | 0.0078 | 2.2 | 8.43E-01 | Pa-234m | U-238 |
| 786.25 | 0.0554 | 0.93 | 5.67E+00 | Pa-234m | U-238 |
| 796.42 | 0.0054 | 4.3 | 6.14E-01 | Pa-234 | U-238 |
| 805.74 | 0.0088 | 1.8 | 1.04E+00 | Pa-234 | U-238 |
| 808.20 | 0.0026 | 10 | 3.60E-01 | Pa-234m | U-238 |
| 819.21 | 0.0037 | 3.9 | 4.20E-01 | Pa-234 | U-238 |
| 824.94 | 0.0068 | 2.6 | 6.47E-01 | Pa-234 | U-238 |
| 831.39 | 0.0078 | 1.9 | 8.89E-01 | Pa-234 | U-238 |
| 851.57 | 0.0070 | 2 | 6.88E-01 | Pa-234m | U-238 |
| 875.94 | 0.0042 | 3 | 6.47E-01 | Pa-234 | U-238 |
| 880.45 | 0.0212 | 0.9 | 1.46E+00 | Pa-234 | U-238 |
| 883.22 | 0.0211 | 0.9 | 2.13E+00 | Pa-234 | U-238 |
| 887.28 | 0.0071 | 1.8 | 8.27E-01 | Pa-234m | U-238 |
| 898.52 | 0.0059 | 2.2 | 6.63E-01 | Pa-234 | U-238 |
| 921.70 | 0.0127 | 1.1 | 1.32E+00 | Pa-234m | U-238 |
| 924.98 | 0.0142 | 1.2 | 1.78E+00 | Pa-234 | U-238 |
| 926.61 | 0.0192 | 1.1 | 1.60E+00 | Pa-234 | U-238 |

| | | | | | |
|----------------|--------------|-------------|-----------------|---------|-------|
| 941.94 | 0.0025 | 4.2 | 3.45E-01 | Pa-234m | U-238 |
| 945.90 | 0.0335 | 0.86 | 2.44E+00 | Pa-234 | U-238 |
| 947.43 | 0.0031 | 4.4 | 1.29E+00 | Pa-234 | U-238 |
| 980.42 | 0.0045 | 3 | 4.85E-01 | Pa-234 | U-238 |
| 984.09 | 0.0030 | 4.2 | 3.07E-01 | Pa-234 | U-238 |
| 994.93 | 0.0057 | 2.1 | 4.61E-01 | Pa-234m | U-238 |
| 1000.99 | 0.839 | 0.56 | 1.03E+02 | Pa-234m | U-238 |
| 1041.70 | 0.0012 | 8 | 1.54E-01 | Pa-234m | U-238 |
| 1061.89 | 0.0023 | 5.2 | 2.23E-01 | Pa-234m | U-238 |
| 1084.25 | 0.0012 | 7.5 | 2.22E-01 | Pa-234 | U-238 |
| 1124.93 | 0.0042 | 3.1 | 3.34E-01 | Pa-234m | U-238 |
| 1193.69 | 0.0135 | 0.96 | 1.43E+00 | Pa-234m | U-238 |
| 1220.37 | 0.0009 | 10.2 | 1.11E-01 | Pa-234m | U-238 |
| 1237.24 | 0.0053 | 1.8 | 5.73E-02 | Pa-234m | U-238 |
| 1292.66 | 0.0009 | 11.2 | 9.70E-02 | Pa-234 | U-238 |
| 1352.80 | 0.0019 | 4.1 | 2.75E-01 | Pa-234 | U-238 |
| 1393.57 | 0.0039 | 2.5 | 4.85E-01 | Pa-234 | U-238 |
| 1413.88 | 0.0023 | 4.2 | 2.39E-01 | Pa-234m | U-238 |
| 1434.13 | 0.0097 | 1.3 | 9.07E-01 | Pa-234m | U-238 |
| 1452.63 | 0.0012 | 7.3 | 1.62E-01 | Pa-234 | U-238 |
| 1510.20 | 0.0129 | 1.2 | 1.45E+00 | Pa-234m | U-238 |
| 1527.27 | 0.0024 | 3.7 | 2.47E-01 | Pa-234m | U-238 |
| 1548.12 | 0.0014 | 5.9 | 2.07E-01 | Pa-234m | U-238 |
| 1553.74 | 0.0081 | 1.6 | 1.00E+00 | Pa-234m | U-238 |
| 1570.67 | 0.0011 | 7.8 | 1.21E-01 | Pa-234m | U-238 |
| 1591.65 | 0.0019 | 5.2 | 4.30E-01 | Pa-234m | U-238 |
| 1593.88 | 0.0027 | 3.6 | 9.70E-02 | Pa-234 | U-238 |
| 1668.44 | 0.0012 | 6.2 | 1.94E-01 | Pa-234 | U-238 |
| 1694.08 | 0.0013 | 5.9 | 1.94E-01 | Pa-234 | U-238 |
| 1737.73 | 0.0212 | 1.1 | 2.26E+00 | Pa-234m | U-238 |
| 1759.81 | 0.0014 | 4.4 | 2.55E-01 | Pa-234m | U-238 |
| 1765.44 | 0.0087 | 1.4 | 9.71E-01 | Pa-234m | U-238 |
| 1809.04 | 0.0037 | 2.1 | 4.77E-01 | Pa-234m | U-238 |
| 1819.69 | 0.0009 | 7.3 | 1.32E-01 | Pa-234m | U-238 |
| 1831.36 | 0.0172 | 1.3 | 1.78E+00 | Pa-234m | U-238 |
| 1863.09 | 0.0012 | 4.3 | 1.35E-01 | Pa-234m | U-238 |
| 1867.68 | 0.0092 | 1.4 | 8.43E-01 | Pa-234m | U-238 |
| 1874.85 | 0.0082 | 1.5 | 8.75E-01 | Pa-234m | U-238 |
| 1877.21 | 0.00165 | 3.4 | 3.07E-02 | Pa-234 | U-238 |
| 1893.50 | 0.00219 | 2.9 | 2.39E-01 | Pa-234m | U-238 |
| 1911.17 | 0.0063 | 1.6 | 5.89E-01 | Pa-234m | U-238 |
| 1925.42 | 0.0005 | 10.1 | 8.08E-02 | Pa-234 | U-238 |
| 1937.01 | 0.0029 | 2.3 | 3.34E-01 | Pa-234m | U-238 |

* Scott, H. L. and Marlow, K. W., NIM A286(1990) 549-55

^{238}U and Daughters— ^{234}Pa and ^{234}Th

| E(keV) | g/s/gm | Emitter | Parent1 | Emiter2 | Parent2 |
|--------|--------|---------|---------|---------|---------|
| 62.9 | 2.36 | Th-234 | U-238 | | |
| 73.9 | 1.36 | Pa-234m | U-238 | | |
| 74.0 | 5.30 | Th-234 | U-238 | | |
| 83.3 | 8.71 | Th-234 | U-238 | | |

| | | | | | |
|-------|-------|---------|-------|--------|-------|
| 92.3 | 338.0 | Th-234 | U-238 | | |
| 92.8 | 335.0 | Th-234 | U-238 | | |
| 94.7 | 21.6 | Pa-234m | U-238 | Pa-234 | U-238 |
| 95.9 | 1.62 | Th-234 | U-238 | | |
| 110.5 | 2.98 | U-238 | U-238 | | |
| 114.9 | 4.32 | Pa-234m | U-238 | Pa-234 | U-238 |
| 131.3 | 3.23 | Pa-234 | U-238 | | |
| 152.7 | 1.08 | Pa-234 | U-238 | | |
| 184.8 | 1.49 | Th-234 | U-238 | | |
| 258.2 | 9.02 | Pa-234m | U-238 | | |

^{238}U Daughter— Protactinium: Gammas and Branching Ratios

$^{234}\text{m}\text{Pa}$ IC-decay*

| E(keV) | IC-decay prob % | | |
|---------|-----------------|---------------|--------|
| | X1000 | uncert. \pm | |
| 257.90 | 57.000 | 0.230 | Pa234m |
| 691.00 | 5.500 | 0.500 | Pa234m |
| 701.60 | 5.400 | 0.500 | Pa234m |
| 740.10 | 7.100 | 0.700 | Pa234m |
| 743.00 | 56.600 | 0.230 | Pa234m |
| 766.60 | 207.800 | 0.800 | Pa234m |
| 782.30 | 5.300 | 0.500 | Pa234m |
| 786.40 | 34.200 | 0.130 | Pa234m |
| 887.50 | 5.200 | 0.500 | Pa234m |
| 922.30 | 8.300 | 0.800 | Pa234m |
| 946.30 | 7.000 | 0.700 | Pa234m |
| 1001.20 | 590.000 | | Pa234m |
| 1738.20 | 14.200 | 0.600 | Pa234m |
| 1831.50 | 11.200 | 0.400 | Pa234m |
| 1868.20 | 5.300 | 0.500 | Pa234m |
| 1911.80 | 3.700 | 0.400 | Pa234m |
| 1937.70 | 2.100 | 0.200 | Pa234m |

NORMALIZED 1001. = 590 (0.59% \times 1000)

* Ardisson G. and Marsol c., Nuovo Chimie 11v28A 155(75)

^{234}Pa IC-decay **

| E(keV) | IC-decay prob % | |
|--------|-----------------|--------|
| | | |
| 63.0 | 4.10 | Pa 234 |
| 131.3 | 20.00 | Pa 234 |
| 153.0 | 6.60 | Pa 234 |
| 226.9 | 11.50 | Pa 234 |
| 569.3 | 13.80 | Pa 234 |
| 699.1 | 4.75 | Pa 234 |
| 805.5 | 3.10 | Pa 234 |
| 824.7 | 3.70 | Pa 234 |
| 831.1 | 5.10 | Pa 234 |
| 926.7 | 16.80 | Pa 234 |
| 945.8 | 18.40 | Pa 234 |
| 980.5 | 3.90 | Pa 234 |
| 1394.1 | 2.40 | Pa 234 |

** Radiochem. Radioanal Lett. 221 357(75)

Appendix D

Uranium Standard Measurements

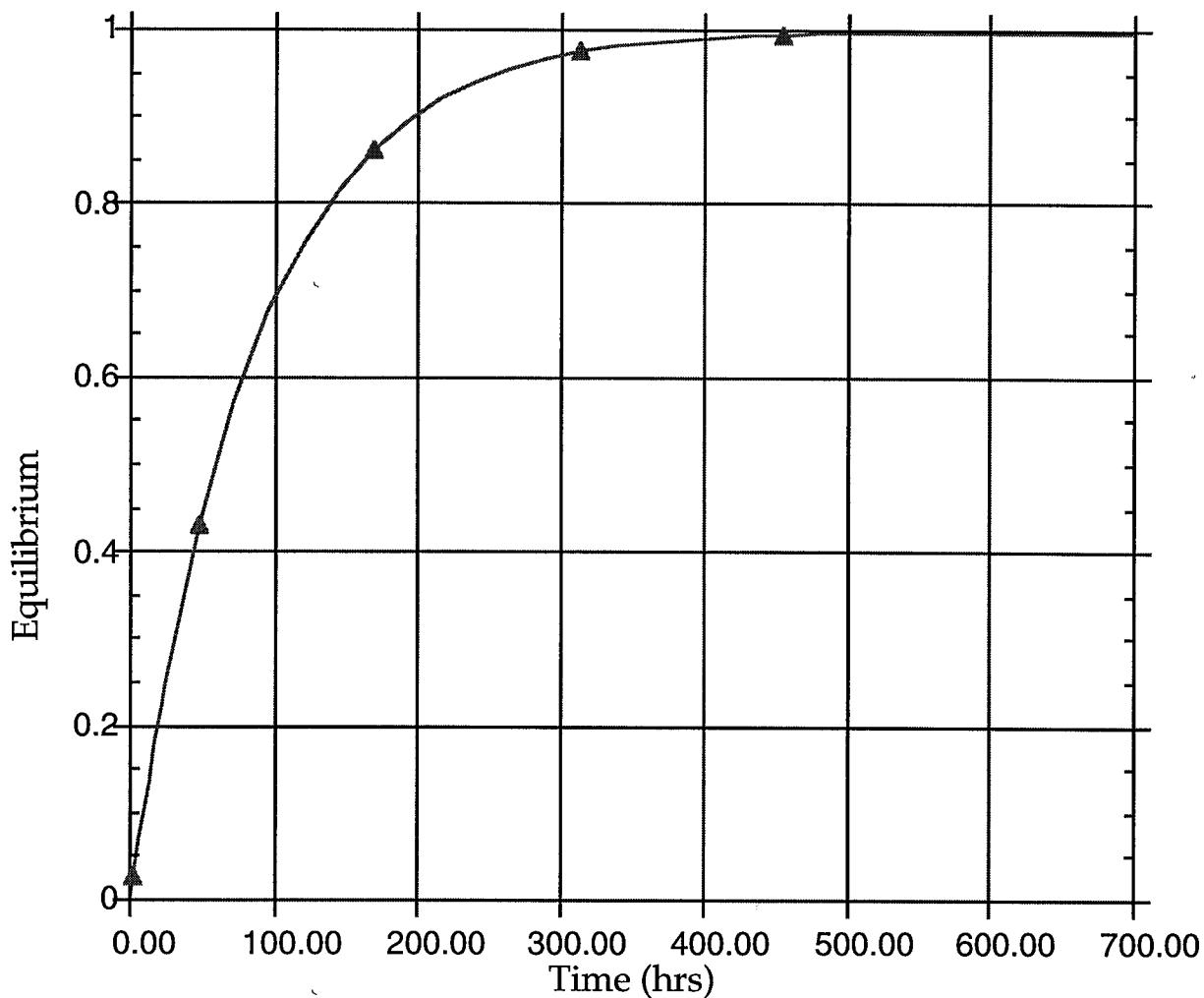
LANL Uranium Standards: mass spec. results

| <u>U238</u> | <u>U235</u> | <u>U-234</u> | <u>U-236</u> | <u>u234/5</u> | <u>U236/5</u> | <u>U234/6</u> |
|--------------------|--------------------|---------------------|---------------------|----------------------|----------------------|----------------------|
| <u>%</u> | <u>%</u> | <u>%</u> | <u>%</u> | <u>ratio</u> | <u>ratio</u> | <u>ratio</u> |
| 99.275 | 0.7194 | 0.0050 | 0.0010 | 0.0070 | 0.0014 | 5.0000 |
| 99.272 | 0.7225 | 0.0050 | 0.0010 | 0.0069 | 0.0014 | 5.0000 |
| 99.267 | 0.7257 | 0.0050 | 0.0020 | 0.0069 | 0.0028 | 2.5000 |
| 99.267 | 0.7257 | 0.0050 | 0.0020 | 0.0069 | 0.0028 | 2.5000 |
| 98.024 | 1.9608 | 0.0160 | 0.0090 | 0.0082 | 0.0046 | 1.7778 |
| 98.022 | 1.9632 | 0.0160 | 0.0090 | 0.0082 | 0.0046 | 1.7778 |
| 96.895 | 3.0637 | 0.0250 | 0.0160 | 0.0082 | 0.0052 | 1.5625 |
| 96.894 | 3.0637 | 0.0250 | 0.0170 | 0.0082 | 0.0056 | 1.4706 |
| 89.664 | 10.1923 | 0.0530 | 0.0910 | 0.0052 | 0.0089 | 0.5824 |
| 89.662 | 10.1961 | 0.0520 | 0.0910 | 0.0051 | 0.0089 | 0.5714 |
| 89.627 | 10.2273 | 0.0540 | 0.0920 | 0.0053 | 0.0090 | 0.5870 |
| 87.883 | 11.9281 | 0.0730 | 0.1160 | 0.0061 | 0.0097 | 0.6293 |
| 86.730 | 13.0841 | 0.0840 | 0.1020 | 0.0064 | 0.0078 | 0.8235 |
| 82.278 | 17.4129 | 0.1400 | 0.1690 | 0.0080 | 0.0097 | 0.8284 |
| 72.499 | 26.9932 | 0.2370 | 0.2710 | 0.0088 | 0.0100 | 0.8745 |
| 62.677 | 37.8417 | 0.2630 | 0.2180 | 0.0070 | 0.0058 | 1.2064 |
| 46.749 | 52.4462 | 0.5360 | 0.2690 | 0.0102 | 0.0051 | 1.9926 |
| 32.833 | 66.3254 | 0.5830 | 0.2590 | 0.0088 | 0.0039 | 2.2510 |
| 7.342 | 91.4086 | 0.9150 | 0.3340 | 0.0100 | 0.0037 | 2.7395 |

| <u>Standard</u> | <u>U234</u> | <u>U235</u> | <u>U236</u> | <u>U238</u> |
|------------------------|-----------------------|-----------------------|-----------------------|-----------------------|
| <u>ID#</u> | <u>(at. %)</u> | <u>(at. %)</u> | <u>(at. %)</u> | <u>(at. %)</u> |
| u0_017 | 0.0010 | 0.0170 | 0.0000 | 99.9710 |
| CRM U002 | 0.0002 | 0.0176 | 0.0000 | 99.9823 |
| u0_483 | 0.0050 | 0.4830 | 0.0008 | 99.4970 |
| NBS U-005a | 0.0034 | 0.5064 | 0.0012 | 99.4890 |
| AI-409 | 0.0050 | 0.7190 | 0.0010 | 99.3000 |
| AI-408-2 | 0.0050 | 0.7226 | 0.0010 | 99.3000 |
| AI-1127-2 | 0.0050 | 0.7256 | 0.0020 | 99.3000 |
| AI-1127-1 | 0.0050 | 0.7258 | 0.0020 | 99.3000 |
| u0_991 | 0.0080 | 0.9910 | 0.0019 | 98.9930 |
| NBS U-010 | 0.0054 | 1.0037 | 0.0068 | 98.9840 |
| AI-1125-1 | 0.0160 | 1.9610 | 0.0090 | 98.0000 |
| AI-1125-2 | 0.0160 | 1.9620 | 0.0090 | 98.0000 |
| u2_013 | 0.0160 | 2.0130 | 0.0044 | 98.0070 |
| u3_009 | 0.0310 | 3.0090 | 0.0070 | 96.9830 |
| CRM U030-A | 0.0278 | 3.0404 | 0.0006 | 96.9312 |
| AI-1126-1 | 0.0250 | 3.0630 | 0.0160 | 96.9000 |
| AI-1126-2 | 0.0250 | 3.0650 | 0.0170 | 96.9000 |
| u4_949 | 0.0410 | 4.9490 | 0.0128 | 94.9630 |
| NBS U-050 | 0.0279 | 5.0100 | 0.0480 | 94.9150 |
| U10_075 | 0.0770 | 10.0750 | 0.0299 | 89.8880 |
| CRM U100 | 0.0676 | 10.1900 | 0.0379 | 89.7040 |
| AI-324-2 | 0.0530 | 10.2000 | 0.0910 | 89.7000 |
| AI-324-1 | 0.0520 | 10.2020 | 0.0910 | 89.7000 |

| | | | | |
|----------|--------|---------|--------|---------|
| AI-323-1 | 0.0540 | 10.2200 | 0.0920 | 89.6000 |
| UISO-12 | 0.0730 | 11.9300 | 0.1160 | 87.8000 |
| UISO-13 | 0.0840 | 13.0890 | 0.1020 | 86.7000 |
| UISO-17 | 0.1400 | 17.4200 | 0.1690 | 82.2000 |
| UISO-27 | 0.2370 | 27.0000 | 0.2710 | 72.7000 |
| UISO-38 | 0.2630 | 37.8480 | 0.2180 | 61.6000 |
| u49_38 | 0.6720 | 49.3800 | 0.2009 | 50.0740 |
| UISO-52 | 0.5360 | 52.4260 | 0.2690 | 46.2000 |
| UISO-66 | 0.5830 | 66.3170 | 0.2590 | 32.8000 |
| U75_13 | 0.5840 | 75.1300 | 0.3320 | 23.5120 |
| CRM U750 | 0.5923 | 75.3570 | 0.2499 | 23.8010 |
| UISO-91 | 0.9150 | 91.4190 | 0.3340 | 7.3300 |
| u93_076 | 1.4520 | 93.0760 | 0.4291 | 5.4720 |

Appendix E
The approach to daughter product equilibrium.



Plot of the equilibrium buildup of ^{231}Th ($T_{1/2} = 25.52$ hours) from the decay of ^{235}U ($T_{1/2} = 7.037 \times 10^8$ years). Full (99.98%) equilibrium is seen in 30 days. At equilibrium, $N_{\text{U}}\lambda_{\text{U}} = N_{\text{Th}}\lambda_{\text{Th}}$ where N_{U} is the number of uranium atoms and N_{Th} is the number of thorium atoms; λ_{U} is the decay constant of ^{235}U and λ_{Th} is the decay constant of ^{231}Th . At equilibrium $N_{\text{Th}}/N_{\text{U}} = 4.137 \times 10^{-12}$ ^a.

$$\frac{N_d}{N_p} = \frac{\lambda_p \cdot N_p}{\lambda_d - \lambda_p} \cdot \left(e^{-\lambda_p t} - e^{-\lambda_d t} \right), \quad N_d \approx \frac{\lambda_p}{\lambda_d - \lambda_p} \cdot N_p e^{-\lambda_p t}$$

At equilibrium

$$N_d \lambda_d = N_p \lambda_p$$

a. Reference: Glasstone, S., Sesonske, A., *Nuclear Reactor Engineering*, 3rd Ed. p42(1991)

Appendix F

Uranium General Properties

| | |
|-------------------------|----------|
| Atomic number: | 92 |
| Standard atomic weight: | 238.0289 |
| Specific gravity | 18.95 |
| Symbol: | U |

General information

Standard state: solid
 Color: metallic gray
 Discoverer: Martin Klaproth
 Date discovered: 1789
 Discovered at: Germany
 Meaning of name: Planet Uranus

Radii (pm)

| | |
|----------------|-------|
| Atomic: | 138.5 |
| van der Waals: | |
| Covalent: | 142 |
| Metallic: | 143 |

Isotopes

Uranium has sixteen isotopes, all of which are radioactive. Naturally occurring uranium nominally contains

| | |
|-----------------|------------------|
| 99.28305% by wt | ^{238}U |
| 0.7200% | ^{235}U |
| 0.0054% | ^{234}U |

X-ray properties

Data for U ; Z = 92

atomic weight = 238.070; density = 19.0500

K-edge at: 115.603 keV

L-edges at: 21.7560, 20.9470, 17.1670 keV

M-edge at: 5.54900 keV

K- α 1, K- β at: 98.4280 111.289 keV

L- α , L- β 1 at: 13.6130 17.2180 keV

K, L1, L2, L3 cross section jumps:

Inf NaN 1.41000 2.29200

Fluorescence yield for K, L1, L2, L3: 0.9720, 0.1760, 0.4670, 0.4890

Calculations are based on data compiled By W. H. McMaster et. al.

Fluorescence yield data by M. O. Krause, J. Phys. Chem. Ref. Data. 8, 307(1979).

Electron Configuration: -21-9-2
 heat Vapor: 110
 heat Fusion: 2.7

Electrical Conductivity: 0.034

Thermal Conductivity: 0.064

Specific Heat Capacity: 0.028

Temperatures °K

| | |
|-----------------|--------|
| melting: | 1405.5 |
| boiling: | 4028 |
| Debye: | 200 |
| superconducton: | 0.68 |

Uranium Isotope Properties

| Z | Abund | mass (AMU) | half-life |
|-----|--------|------------|------------|
| 232 | 0 | 69.8 | yr. |
| 233 | 0 | 1.592E5 | yr. |
| 234 | 0.0054 | 234.040946 | 2.457E5 |
| 235 | 0.720 | 235.043924 | 7.037E8 |
| 236 | 0 | 236.045562 | 2.342E7 |
| 237 | 0 | 237.04827 | 6.75 |
| 238 | 99.275 | 238.050784 | 4.468E9 |
| 239 | 0 | | 23.47 min. |
| 240 | 0 | | 14.1 hr |

92-uranium-235

(U235 is also called actinouranium)

U235(n,f) is standard for neutrons below 20 MeV.

General Properties

Atomic Mass:

$235.0439222 \pm 0.0000021$ AMU

Excess Mass: 40913.215 ± 2.002 keV

Binding Energy:

1783871.153 ± 2.028 keV

Beta IC-decay energy: β^-
 -123.716 ± 0.869 keV

Atomic Percent Abundance: 0.720%

Spin: 7/2-

Half life: $703.8\text{E}+6$ Y (0.0710 %)

Mode of IC-decay: Alpha to Th-231

IC-decay energy: 4.6787 MeV

Mode of IC-decay: Spontaneous fission

Branch ratio: 7.0E-9 %

Fission spectrum avg. = 89.07 mb
g-factor = 0.9898

Meta state at 0.0768 keV

Spin: 1/2+

Half life: 25 M

Mode of IC-decay: Internal Transition

Total Cross Section

Cross Section at 0.0253 eV = 698.2 b

Maxwell avg. at 0.0253 eV = 608.4 b

Cross Section at 14 MeV = 5.865 b

Fission spectrum avg. = 7.705 b

g-factor = 0.9833

Elastic scattering Cross Section

Cross Section at 0.0253 eV = 15.04 b

Maxwell avg. at 0.0253 eV = 14.95 b

Cross Section at 14 MeV = 2.871 b

Fission spectrum avg. = 4.566 b

g-factor = 1.1215

Total inelastic Cross Section

Cross Section at 14 MeV = 350.3 mb

Fission spectrum avg. = 1.804 b

(n,2n) Cross Section

Cross Section at 14 MeV = 542.9 mb

Fission spectrum avg. = 11.56 mb

(n,3n) Cross Section

Cross Section at 14 MeV = 41.79 mb

Fission spectrum avg. = 7.074 μ b

Total fission Cross Section

Cross Section at 0.0253 eV = 584.4 b

Maxwell avg. at 0.0253 eV = 506.8 b

Resonance integral = 278.1 b

Cross Section at 14 MeV = 2.056 b

Fission spectrum avg. = 1.235 b

g-factor = 0.9786

(n,4n) Cross Section

Fission spectrum avg. = 0.008408 μ b

Radiative capture Cross Section

Cross Section at 0.0253 eV = 98.81 b

Maxwell avg. at 0.0253 eV = 86.67 b

Resonance integral = 133.0 b

Cross Section at 14 MeV = 0.1607 μ barn

92-uranium-238

(^{238}U is also called uranium I)

General Properties

Atomic Mass:

$238.0507835 \pm 0.0000022$ AMU

Excess Mass:

47304.541 ± 2.024 keV

Binding Energy:

1801693.796 ± 2.050 keV

Beta IC-decay energy:

$\beta^- 145.345 \pm 1.351$ keV

Atomic Percent Abundance:

99.2745% 15

Spin: 0+

Half life: $4.468\text{E}+9$ Y (0.0671 %)

Mode of IC-decay: Alpha to Th-234

IC-decay energy: 4.270 MeV

Mode of IC-decay: Spontaneous fission

Branch ratio: 5.45×10^{-5} %

Mode of IC-decay: Beta

Branch ratio: 2.2×10^{-10} %

Cross Section at 14 MeV = 350.0 mb

Fission spectrum avg. = 61.90μ barn

Total fission Cross Section

Cross Section at 0.0253 eV = 11.77μ barn

Maxwell avg. at 0.0253 eV = 10.45μ barn

Resonance integral = 2.020 b

Cross Section at 14 MeV = 1.136 b

Fission spectrum avg. = 308.4 mb

g-factor = 1.0015

Radiative capture Cross Section

Cross Section at 0.0253 eV = 2.717 b

Maxwell avg. at 0.0253 eV = 2.414 b

Resonance integral = 278.1 b

Cross Section at 14 MeV = 1.943 mb

Fission spectrum avg. = 66.40 mb

g-factor = 1.0024

Total Cross Section

Cross Section at 0.0253 eV = 12.09 b

Maxwell avg. at 0.0253 eV = 11.77 b

Cross Section at 14 MeV = 5.805 b

Fission spectrum avg. = 7.786 b

g-factor = 1.0994

Elastic scattering Cross Section

Cross Section at 0.0253 eV = 9.360 b

Maxwell avg. at 0.0253 eV = 9.356 b

Cross Section at 14 MeV = 2.704 b

Fission spectrum avg. = 4.804 b

g-factor = 1.1279

Total inelastic Cross Section

Cross Section at 14 MeV = 698.3 mb

Fission spectrum avg. = 2.595 b

(n,2n) Cross Section

Cross Section at 14 MeV = 910.0 mb

Fission spectrum avg. = 12.07 mb

(n,3n) Cross Section

| Isotope[1] | | delta | | T1/2 or Abundance[4] | | Decay Mode[5] |
|------------|----|-------|------------|-------------------------|--------------|----------------------------|
| Z | El | A | Jpi[2] | MeV)[3] | | |
| 92 | U | 218 | 0+ | 21.88s | 1.5 ms +73-7 | A |
| | | 219 | | | 23.2s | 42 us +34-13 A |
| | | 220 | 0+ | | 23.0s | |
| | | 221 | | | 24.5s | |
| | | 222 | 0+ | | 24.3s | 1.0 us +10-4 A |
| | | 223 | | | 25.82 | 18 us +10-5 A |
| | | 224 | 0+ | 25.70 | 1.0 ms 4 | A |
| | | 225 | | | 27.37 | 95 ms 15 A |
| | | 226 | 0+ | 27.32 | 0.20 s 5 | A |
| | | 227 | (3/2+) | 29.00 | | 1.1 m 1 A |
| | | 228 | 0+ | 29.22 | | 9.1 m 2 A > 95%,EC < 5% |
| | | 229 | (3/2+) | 31.204 | 58 m 3 | EC ~ 80%,A ~ 20% |
| | | 230 | 0+ | 31.603 | 20.8 d | A |
| | | 231 | (5/2-) | 33.78 | | 4.2 d 1 EC |
| | | 231 | (3/2+,5/2+ | 33.78 | | 4.2 d 1 A ~ 4E-3% |
| | | 232 | 0+ | 34.601 | 68.9 y 4 | A,Ne 9E-10% |
| | | 233 | 5/2+ | 36.912 | 1.592E+5 y 2 | A,SF <6.0E-9% Ne 7E-11% |
| | | 234 | 0+ | 38.140 | 2.455E+5 y 6 | A,SF 1.7E-9%,Mg 1E-11% |
| | | | | | 0.0055% 5 | Ne 9E-12% |
| | | 235 | 7/2- | 40.913 | 703.8E+6 y 5 | A,SF 7.0E-9%,Ne 8E-10% |
| | | | | | 0.720% 1 | |
| | | 235M | 1/2+ | 40.913 | ~ 25 m | IT |
| | | 236 | 0+ | 42.440 | 2.342E7 y 3 | A,SF 9.6E-8% |
| | | 236M | | | 42.440 | 121 NS 2 SF 0.013% |
| | | 237 | 1/2+ | 45.385 | 6.75 d 1 | B- |
| | | 238 | 0+ | 47.305 | 4.468E+9 y 3 | A,SF 0.5E-4% |
| | | | | | 99.2745% 15 | |
| | | 238M | 0+ | 47.305 | 267 NS 3 | SF 0.015% |
| | | 239 | 5/2+ | 50.570 | 23.45 m 2 | B- |
| | | 240 | 0+ | 52.708 | 14.1 h 1 | B-,A |
| | | 242 | 0+ | 58.830C | 16.8 m 5 | B- |

Explanation of Table

• Column 1, Isotope (Z, El, A)

Nuclides are listed in order of increasing atomic number (Z), and are subordered by increasing mass number (A). All isotopic species are included as well as all isomers with half-life > or = 0.1 s, and some other isomers which decay by SF or alpha emissions. A nuclide is included even if only its mass estimate or its production cross section is available. For the latter nuclides $T_{1/2}$ limit is given[8].

Isomeric states are denoted by the symbol "m" after the mass number and are given in the order of increasing excitation energy.

The ^{235}U thermal fission products, with fractional cumulative yields > or = 10-6, are italicized in the table. The information on fission products is taken from the ENDF/B-VI fission products file[11].

The names for elements Z=104-109 are those adopted by the American Chemical Society Nomenclature Committee. The symbols Rf (Rutherfordium) and Ha (Hahnium) have, not been accepted internationally due to conflicting claims about the discovery of these elements.

• Column 2, J^{π}

Spin and parity assignments, without and with parentheses, are based upon strong and weak arguments, respectively. See the introductory pages of any January issue of Nuclear Data Sheets[2] for description of strong and weak arguments for J^{π} assignments.

• Column 3, Mass Excess, Delta

Mass excesses, M-A, are given in MeV with $\Delta(^{12}\text{C})=0$, by definition. For isomers the values are obtained by adding the excitation energy to the $\Delta(\text{g.s.})$ values. Wherever the excitation energy is not known, the mass excess for the next lower isomer (or g.s.) is given. The values are given to the accuracy determined by uncertainty in $\Delta(\text{g.s.})$ (maximum of three figures after the decimal). The uncertainty is > or = 9 in the last significant figure. An appended "s" denotes that the value is obtained from systematics.

• Column 4, $T_{1/2}$ or Abundance

The half-life and the abundance (in bold face) are shown followed by their units ("%" symbol in the case of abundance) which are followed by the uncertainty, in italics, in the last significant figure. For example, 8.1 s 10 means 8.1 ± 1.0 s. For some very short-lived nuclei, level widths rather than half-lives are given. There also, the width is followed by units (e.g., eV, keV, or MeV) which are followed by the uncertainty in italics, if known.

• Column 5, Decay Mode

Decay modes are given in decreasing strength from left to right, followed by the percentage branching, if known ("w" indicates a weak branch). The percentage branching is omitted where there is no competing mode of decay or no other mode has been observed.

The various modes of decay are given below:

| | |
|-------------------------|---|
| B- | beta- decay |
| E | epsilon (electron capture), or epsilon+beta+, or beta+ decay |
| IT | isomeric transition (through gamma or conversion-electron decay) |
| n, p, A, ... | neutron, proton, alpha, decay |
| SF | spontaneous fission |
| 2B-, 3A, ... | double beta- decay (beta-beta-), decay through emission of 3 alpha's, |
| B-N, B-P, | delayed n, p, alpha, ... B-A, ..emission following beta decay |
| EP, EA, ESF, ... | delayed p, alpha, SF, decay following epsilon or beta+ decay |

The appendices conform to the Fundamental Physical Constants[13]. For properties of the elementary particles and for the astrophysical constants see the Review of Particle Properties, Physical Review D50, 1173 (1994) and its subsequent biennial updates. See also the World Wide Web at URL: <http://pdg.lbl.gov/>

References:

1. **Evaluated Nuclear Structure Data File**- a computer file of evaluated experimental nuclear structure data maintained by the National Nuclear Data Center, Brookhaven National Laboratory (file as of June 1995).
2. **Nuclear Data Sheets** - Academic Press, San Diego. Evaluations published by mass number for $A = 45$ to 266. See page ii of any issue for the index to A-chains.
3. **Nuclear Physics** - North Holland Publishing Co., Amsterdam - Evaluations by F. Ajzenberg-Selove and by D. R. Tilley, H. R. Weller, C. M. Cheves, and R. M. Chasteler for $A = 3$ to 20.
4. **Energy Levels of $A = 21$ -44 Nuclei (VII)**, P. M. Endt, Nuclear Physics A521, 1 (1990).
5. **Nuclear Science Reference File**- a bibliographic computer file of nuclear science references continually updated and maintained by the National Nuclear Data Center, Brookhaven National Laboratory. Recent literature scanned by S. Ramavataram.
6. **Table of Isotopes** (1996), 8th edition, Editors:R. B. Firestone, et al., John Wiley, New York.
7. Spontaneous Fission, D. C. Hoffman, T. M. Hamilton, and M. R. Lane, Rept. LBL-33001 (1992).
8. **NUBASE: A Database of Nuclear and Decay Properties**, G. Audi, O. Bersillon, J. Blachot, and A. H. Wapstra, Intl. Symposium on Radionuclide Metrology and its Applications (1995).
9. **The 1993 Atomic Mass Evaluation**, G. Audi and A. H. Wapstra, computerized list of recommended values based on authors' publication Nuclear Physics A565, 1 (1993)
10. **Table of the Isotopes**, N. E. Holden, Rept BNL-61460 (1995) and private communication.
11. **Evaluation and Compilation of Fission Product Yields 1993**, T. R. England and B. F. Rider; Rept. LA-UR-94-3106 (1994). ENDF/B-VI evaluation; MAT #9228, Revision 1.
12. **Table of Isotopes** (1978), 7th edition, Editors: C. M. Lederer, V. S. Shirley, Authors: E. Browne, J. M. Dairiki, R. E. Doebler, A. A. Shihab-Eldin, J. Jardine, J. K. Tuli, and A. B. Buryn, John Wiley, New York.
13. **The Fundamental Physical Constants**, E. R. Taylor and B. N. Taylor, Physics Today BG9 (August, 1995).

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