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P. Asoka-Kumar

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Chemical identity of atoms using core electron annihilations

P. Asoka-Kumar

Lawrence Livermore National Laboratory, Livermore, CA 94550

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Abstract

Positron annihilation spectroscopy is a sensitive probe for studying the electronic structure of defects in solids. The high momentum part of the Doppler-broadened annihilation spectra can be used to distinguish different elements. This is achieved by using a new two-detector coincidence system and by imposing appropriate kinematic cuts to exclude background events. The new setup improves the peak to background ratio in the annihilation spectrum to $\sim 10^5$. As a result, the line shape variations arising from different core electrons can be studied. The new approach adds elemental specificity to the Doppler broadening technique, and is useful in studying elemental variations around a defect site. Results from several case studies are reviewed.

Introduction

The Doppler broadening of the annihilation gamma rays has been used extensively to study vacancies and vacancy clusters[1]. Despite the remarkable success of this observable in rapidly identifying dilute quantities of defects, the Doppler broadening parameters (S, W, etc.) have been difficult to calculate from the electron momentum distribution. Due to this difficulty, many in the positron community have shied away from the Doppler broadening measurements and have concentrated on theoretically verifiable observables (lifetime and angular correlation). In this paper, I show the recent improvements in Doppler broadening measurements that is making it an equally attractive positron tool.

Nearly two decades ago, Lynn et al. showed how the Doppler broadening measurements can yield the momentum profiles of core electrons[2,3]. This is achieved by replacing the standard single Ge(Li) detector with a coincident two Ge(Li) detector setup and by imposing appropriate kinematic cuts to the resulting spectrum[4]. Because the core electrons retain their atomic character even when atoms form a solid, these results can be directly compared to simple theoretical predictions. The importance of this work went largely unnoticed until recently when the Helsinki group started applying them to study defects in compound semiconductors[5,6]. A systematic study of the line shape variations originating from the high momentum electrons of different elements demonstrated the elemental specificity of this new approach[7].

The conventional Doppler broadening spectra obtained by recording one of the photons from the annihilation process is distorted by background events (high energy gamma rays, pulse pileup, incomplete charge collection, and three-photon decay of positronium). This problem is particularly acute in the energy region (tail regions of the annihilation spectrum) where the contributions from high momentum core electrons are dominant. As discussed by Saarinen here, a single Ge detector in coincidence with a timing detector (NaI or BGO) will reduce the background on the high-energy side of the annihilation peak[8]. The two-photon coincidence measurement provides a more informative kinematic picture of the annihilation process.

Two-photon Doppler system

The principle of the two-photon Doppler system is described elsewhere[4,7]. With the simultaneous recording of the energies of both gamma rays (E_1 & E_2), events meeting the total energy criteria of $E_T = E_1 + E_2 = 2m_0c^2 - E_b$ can be selected to suppress the background events, where E_b is the

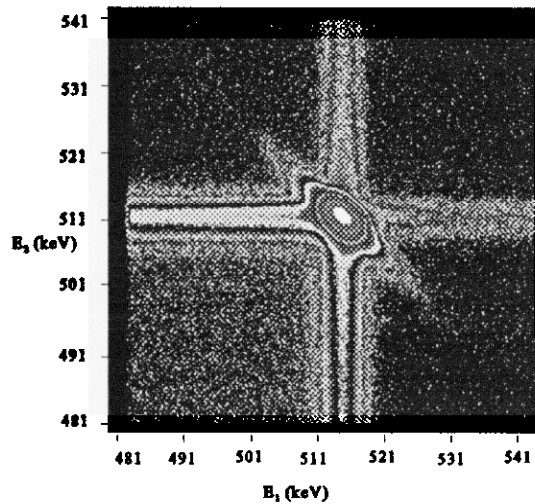


Figure 1 A two-dimensional display of coincidence spectrum

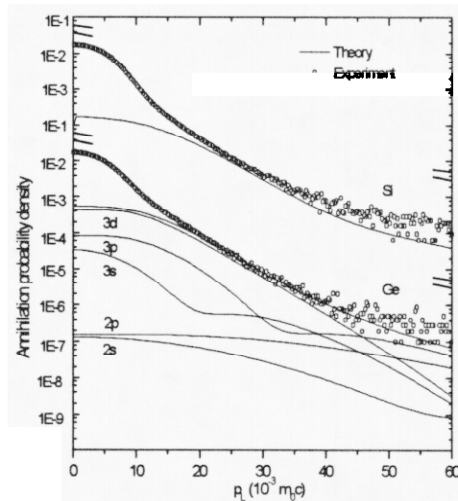


Figure 2 Annihilation probability density for Si and Ge as a function of the longitudinal momentum, p_L .

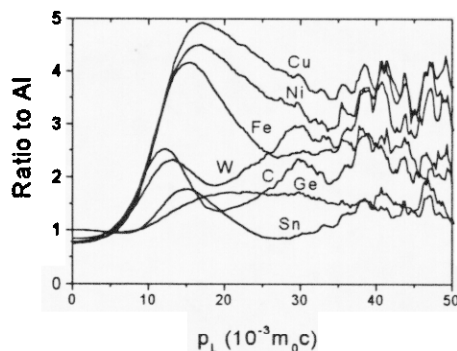


Figure 3 Ratio of the annihilation probability densities of different elements to Al.

electron binding energy. The energy difference, $E_1 - E_2$, is equal to $p_L c$, where p_L is the momentum of the electron-positron system along the direction of the gamma ray emission. Figure 1 shows an example of a two-dimensional spectrum. The gray tone indicates the number of counts. The central peak at $E_1 = E_2 \approx m_0 c^2$ corresponds to annihilations with small Doppler shifts. The horizontal and vertical bands extending from the central peak are produced by coincidences of a background photon with a 511 keV photon. The elliptical region extending diagonally with $E_1 + E_2 \sim 2m_0 c^2 = 1022$ keV originates from the high momentum electrons, and this region is nearly background free. Therefore, a cross section along the diagonal can be analyzed to study the core-electron momentum distributions.

Elemental Specificity

Figure 2 shows the shape variations arising from core annihilations of Si and Ge. The experimental curves are obtained by folding the Doppler-broadening curve and normalizing to obtain the annihilation probability density as a function of the longitudinal momentum. The details of the normalization procedure can be found in Ref. 7. The samples used in the present study are high-purity single crystals, and are "defect-free" as seen by positrons. The spectra were collected at an incident positron beam energy of 30 keV to exclude positron diffusion to the surface and subsequent annihilations from surface states. The larger amplitude of the Ge curve stems from the fact that the outermost core shell in the case of Ge is 3d, which is highly occupied and spatially extended. The theoretical curves are obtained using the momentum distribution for each core electron state within the independent particle model, using free atom wave functions[9]. The contributions from the individual shells of Ge are also included. Since the core electrons retain their atomic character even when atoms form a solid, atomic wave functions provide an accurate description of the core electron states. Atomic wave functions, however, are not an accurate description of the valence electrons and are not added to the total.

Since the curves in Fig. 2 span several orders of magnitude, the shape difference between them can be seen more easily after normalizing the measured spectrum to a reference spectrum. Fig. 3 shows the ratio curves for several elements and demonstrates the elemental specificity of the new method.

Vacancies and impurity-vacancy complexes

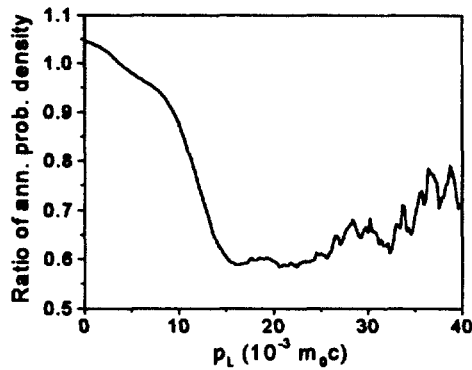


Figure 4 Ratio of annihilation probability density for thermally generated vacancies in Al.

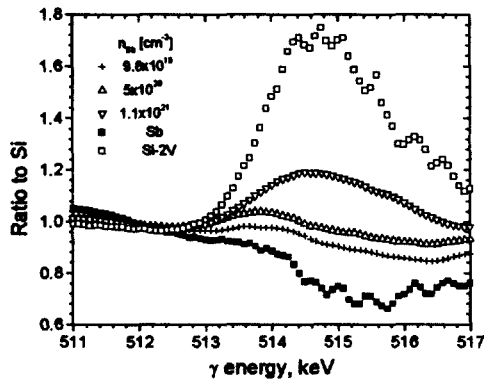


Figure 5 Ratio curves for Sb doped Si as a function of gamma ray energies. Results for Si divacancies are also shown (after ref. 10).

about $N_{\text{Sb}} = 7 \times 10^{20} \text{ cm}^{-3}$, and further increase in dopant density does not show a corresponding increase in the carrier concentration (for details see Ref. 10).

The core annihilation line shape clearly shows the presence of Sb for concentrations beyond the critical dopant density. When Sb occupies a lattice site in Si and is activated, the Sb ion is positively charged. Therefore, positrons cannot sense them. Figure 5 shows the ratio curves at three Sb concentrations along with the results from bulk Si. The Sb signal is clearly evident at the highest dopant concentration. Since positrons are not sensitive to ionized Sb atoms, the presence of Sb signal at the high dopant concentration can be understood in terms of a Sb-vacancy complex. Modeling of the shape variations show that the data points are consistent with a linear superposition of Si vacancy signal and Sb signal [10]. The fraction of the Sb signal required is consistent with the observed reduction in carrier concentration. Since Sb-vacancy complex is an acceptor, it can compensate the substitutional Sb donors.

Ion implantation is routinely used in device fabrication to introduce controlled quantities of dopant atoms in isolated regions of a semiconductor wafer, thus transcending the limits set by thermal diffusion. The ion implantation, besides introducing foreign atoms, creates many lattice defects in the solid. Several impurity-vacancy related problems are present in ion implantation. When the implantation dose exceeds solid solubility, the precipitation can become severe. The implantation damage can act as gettering centers. The oxygen-vacancy complex can be formed in Czochralski Si. The contamination of the primary beam from the residual gas scattering can occur in high energy implantations using multiply

When positrons are localized at an open-volume defect site, the positron overlap with core electrons is reduced. As a result, the Doppler-broadening curves become narrow. The reduction in core overlap reduces counts at high momentum values as seen in Fig. 4, for thermally generated vacancies in Al. The ratio curve reduces by $\sim 40\%$ at $p_L \sim 15 \times 10^{-3} m_0 c$. In comparison, the traditional single detector measurement of S and W parameters show an increase of only $\sim 2\text{-}5\%$. Thus, the new setup increases the sensitivity of the positron technique to vacancy-like defects.

The technique can also be used to study the elemental variations around a defect site. When the momentum distribution of the core electrons of an impurity atom is clearly different from the host atoms, the technique can be used to identify the foreign atoms decorating the vacancy sites. These studies are relevant to understand the structure and the kinetics of defects in several technologically important systems. Below, I list some of them.

The low temperature MBE growth of Si allows to produce sharp dopant profiles at high concentrations. However, at high doping values the carrier concentration fails to follow the unity activation line. While many studies have attempted to determine the defect responsible for the compensation at high concentration, no unambiguous evidence exists for a specific defect that limits the dopant activity in Si. In Sb-doped Si, the carrier concentration increases linearly with increasing defect concentration to

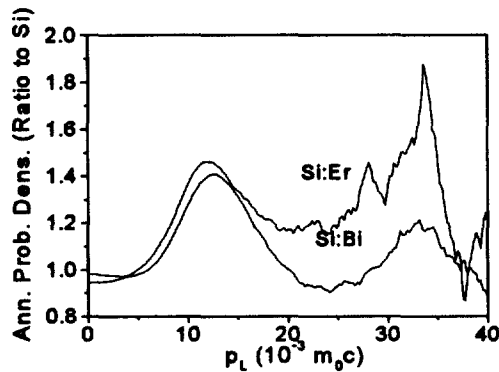


Figure 6 Annihilation probability density for Er and Bi implanted in Si, after normalizing to bulk Si.

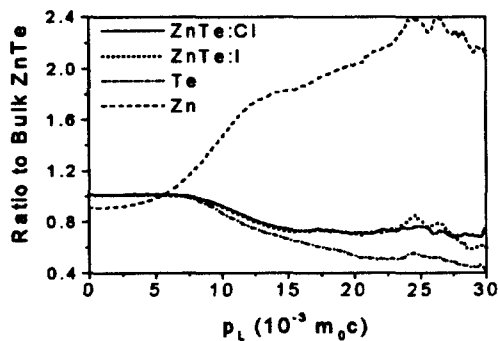


Figure 7 Annihilation probability density of ZnTe films doped with Cl and I. Results for bulk Zn and Te are also shown. The data is normalized to a ZnTe polycrystal showing a single lifetime value of 286 ps.

charged ions, where the contaminant has the same magnetic rigidity as the primary ion. Since the new method is elemental specific, it is ideally suited to study the creation and evolution of these defects. A library of defect signatures for different ionic species will enhance the attraction of the positron probe as a characterization method in Si industry. As an example, Figure 6 shows Bi and Er decorated vacancies in Si.

In compound materials, the positrons localized at vacancies in different sublattices will be surrounded by different chemical species. Hence, the high momentum components can be used to identify the vacancies in different sublattices. The Helsinki group has studied several compound semiconductor systems[6]. Here, as an example, I address the doping difficulties encountered in wide band gap II-VI semiconductors. The as-grown materials contain excess carriers of one type (electrons or holes) showing a native n-type or p-type conductivity. Attempts to change the conductivity type by adding dopant atoms have been difficult. For example, ZnTe crystals show native p-type behavior. Several theoretical calculations have been performed to explain them, and point defects are believed to be central to these difficulties[11]. We examined Cl and I doped (both n-type dopants) ZnTe using the annihilation line shapes of the inner-shell electrons to understand the role of vacancies in the compensation mechanism. Results from ZnTe:Cl and ZnTe:I show an enhancement of annihilations with Te electrons compared to undoped samples, and can be

understood as arising due to the first neighbor of a Zn vacancy (see Figure 7).

Novel studies using the coincidence method

Before concluding, I also wanted to point out few novel studies that one could undertake using the two-photon coincidence method. Positronium (Ps) formation can easily be detected using the two-dimensional spectra. Ps formation adds two new features to the spectrum shown in Fig. 1. The central peak narrows due to the two-photon decay of Ps. Also, there will be more counts in the lower-left quadrant ($E_1 \& E_2 < 511$ keV) from the three photon decay mode of Ps. A more systematic study is needed to take advantage of these new features.

The projection of the two-dimensional spectrum onto E_1 - E_2 axis provides the annihilation probability density as a function of the longitudinal momentum p_L . The projected spectrum includes contributions from all electrons. The high momentum part of the spectrum is dominated by the inner-shell electrons, and the approach described above can detect only the sum contributions of all inner-shell electrons. The annihilation of the inner-shell electrons with positrons will create a core hole, which will be filled with the emission of an Auger electron or X-ray. Therefore, the detection of the characteristic X-ray or Auger electron can be used to gate the two-dimensional spectrum to study the contribution of a specific atomic level. The partial annihilation rate and annihilation probability distribution can then be compared directly to understand the positron overlap with the inner-shell electrons. Positron lifetime

measurements and the corresponding theoretical values show remarkable agreement. However, the positron lifetime is dominated by outer electron contributions and is insensitive to core contributions. Hence, the measurement of the inner-shell annihilation probability distributions will be valuable in validating the theoretical calculations.

The annihilation of energetic positrons can be used to study the spin-dependent properties of a material. With a polarized positron beam (positron beams derived from a radioactive source retains some amount of polarization) the annihilation-in-flight cross section will depend on the spin alignment of the electrons. The annihilation-in-flight events will have a total energy greater than 1022 keV and will be clearly distinguishable in the upper right quadrant of the two-dimensional spectrum of Fig. 1.

Conclusion

The recent developments in Doppler broadening technique is exciting and is useful in distinguishing between different elements. Because the core electrons retain their atomic character even when atoms form a solid, their annihilation spectra reflect the atomic character and are elemental specific. The method can be used to study a wide variety of problems, like simple open-volume defects, open-volume defects decorated with foreign elements, and vacancies in different sublattices of a compound material.

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Technical Information Department • Lawrence Livermore National Laboratory
University of California • Livermore, California 94551

