

# MITL ELECTRODE SURFACE PLASMA DENSITY MEASUREMENTS

*at Sandia's 1 MA accelerator Mykonos*

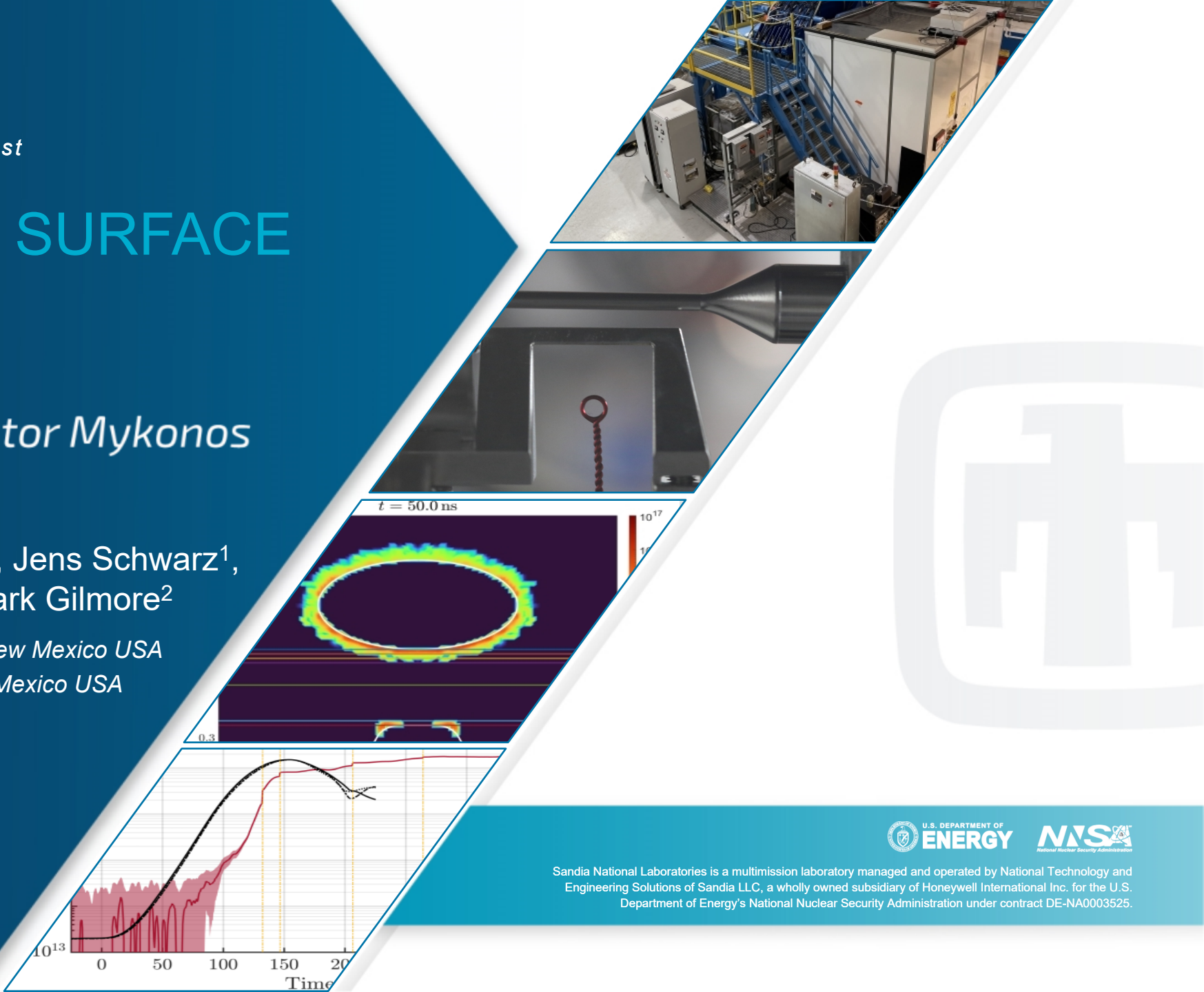
Nathan Hines<sup>1,2,a)</sup>, Derek Lamppa<sup>1</sup>, Jens Schwarz<sup>1</sup>,  
Michael Cuneo<sup>1</sup>, Thomas Awe<sup>1</sup>, Mark Gilmore<sup>2</sup>

<sup>1</sup> Sandia National Laboratories, Albuquerque, New Mexico USA

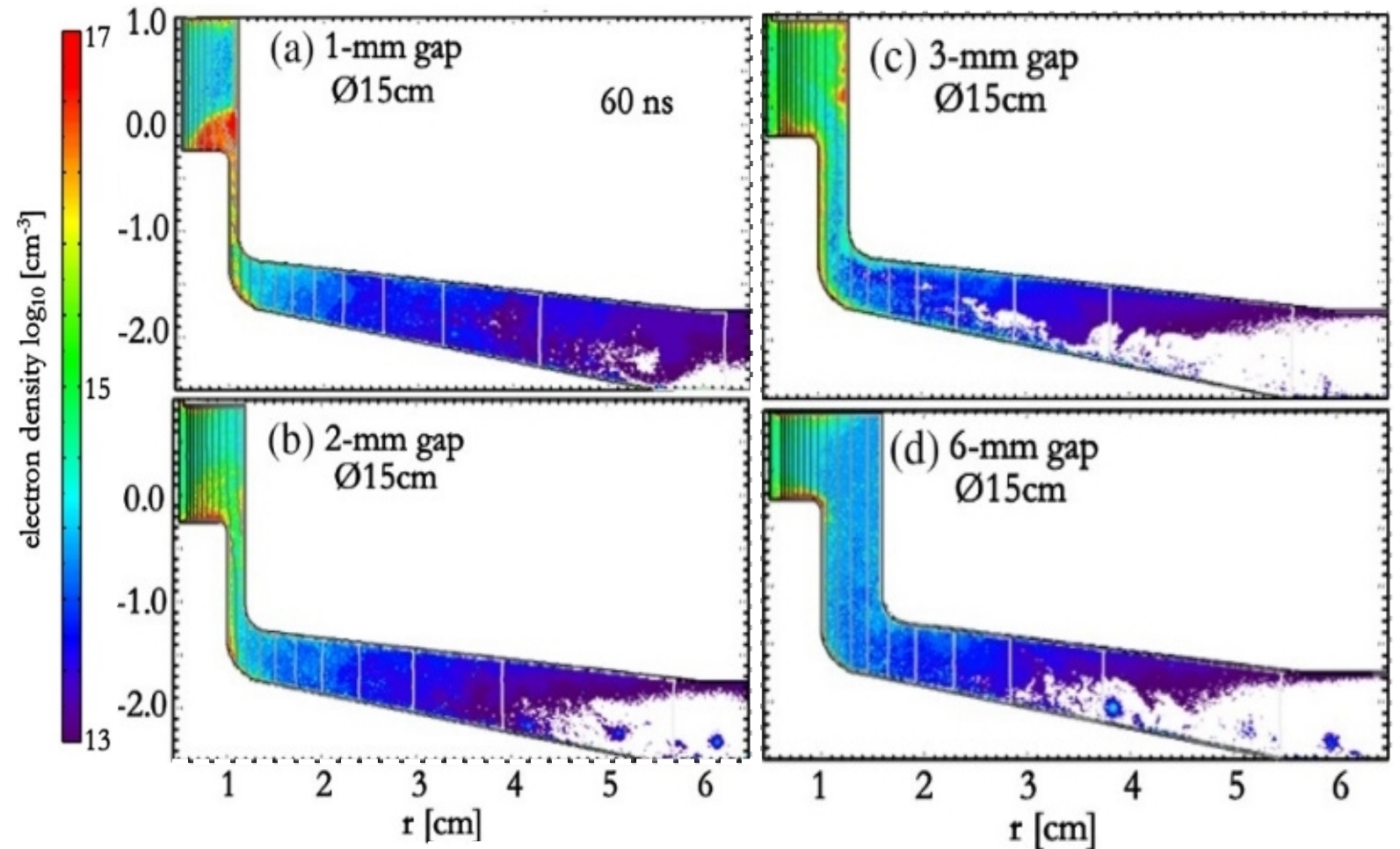
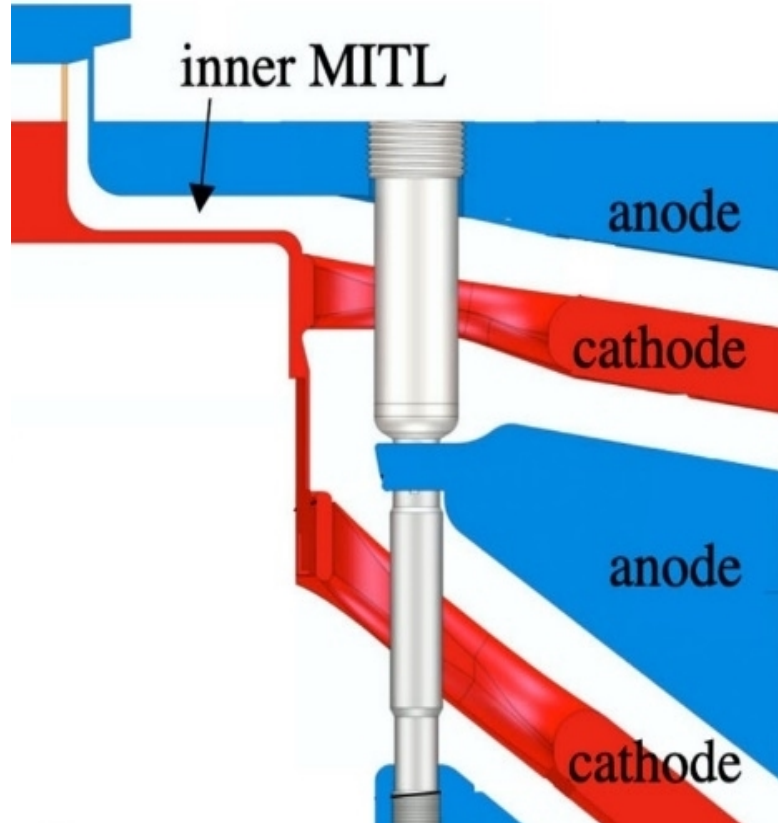
<sup>2</sup> University of New Mexico, Albuquerque, New Mexico USA

Presented at the 2024 APS DPP Meeting  
October 7-11, Atlanta, Georgia USA

<sup>a)</sup> Corresponding author email: [nhines@sandia.gov](mailto:nhines@sandia.gov)



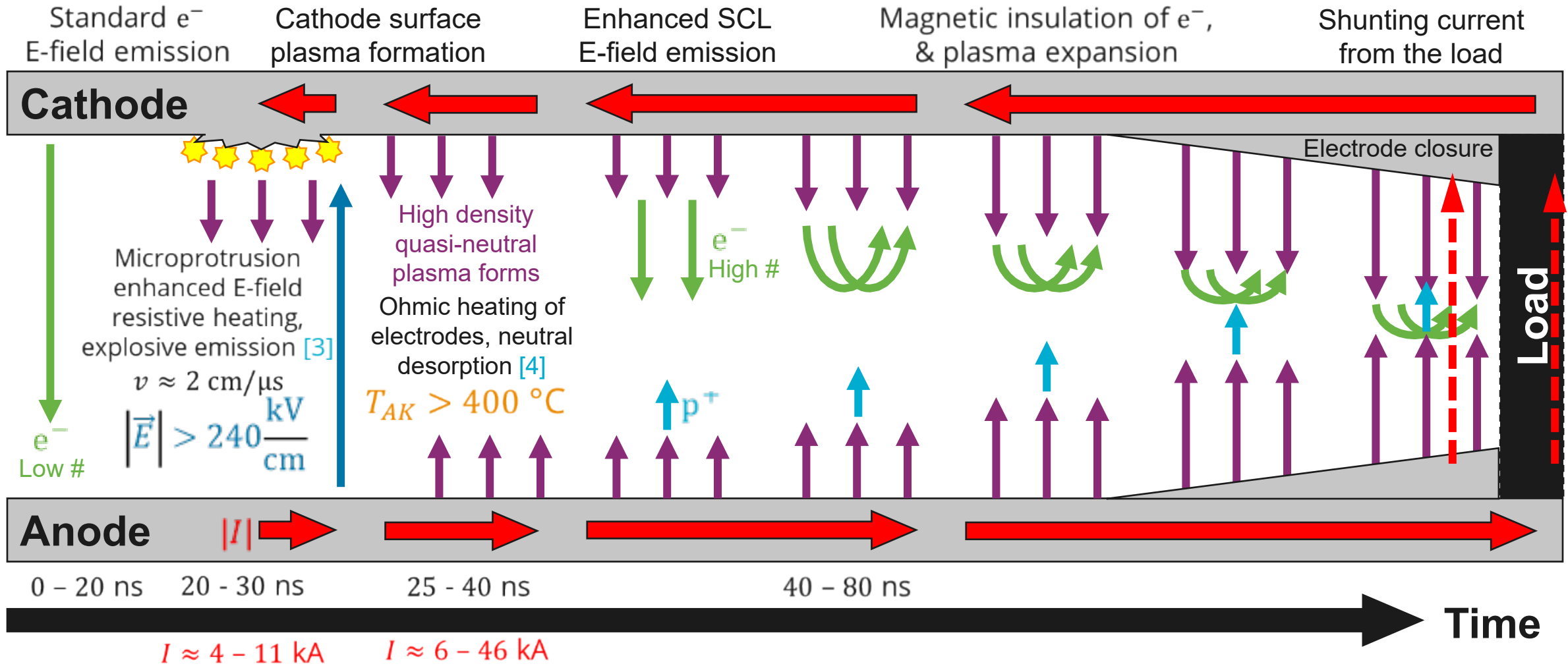
Z's inner MITL experiences current loss from charged particle cross-gap flow of expected  $e^-$  densities of  $10^{13} - 10^{17} \text{ cm}^{-3}$



[1] W. A. Stygar et al., "55-TW magnetically insulated transmission-line system: Design, simulations, and performance," Phys. Rev. Accel. Beams, vol. 12, no. 12, p. 120401, 12/07/2009.

[2] N. Bennett, D. R. Welch, G. Laity, D. V. Rose, and M. E. Cuneo, "Magnetized particle transport in multi-MA accelerators," Physical Review Accelerators and Beams, vol. 24, no. 6, p. 060401, 06/23/2021.

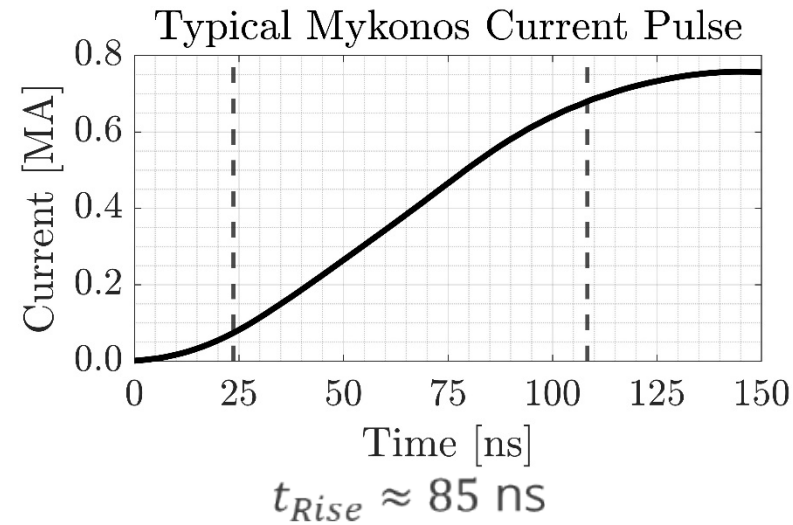
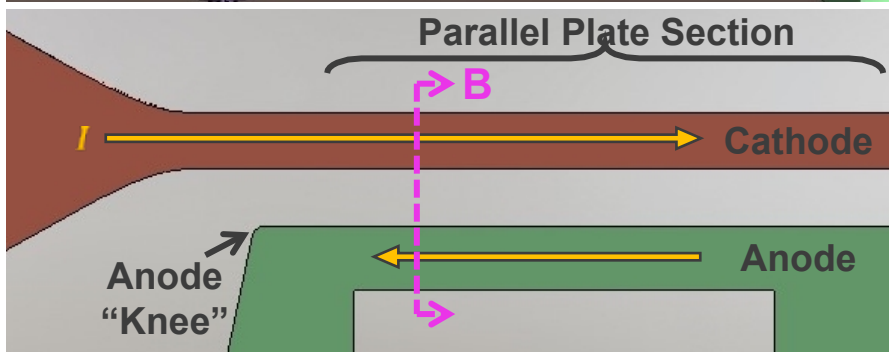
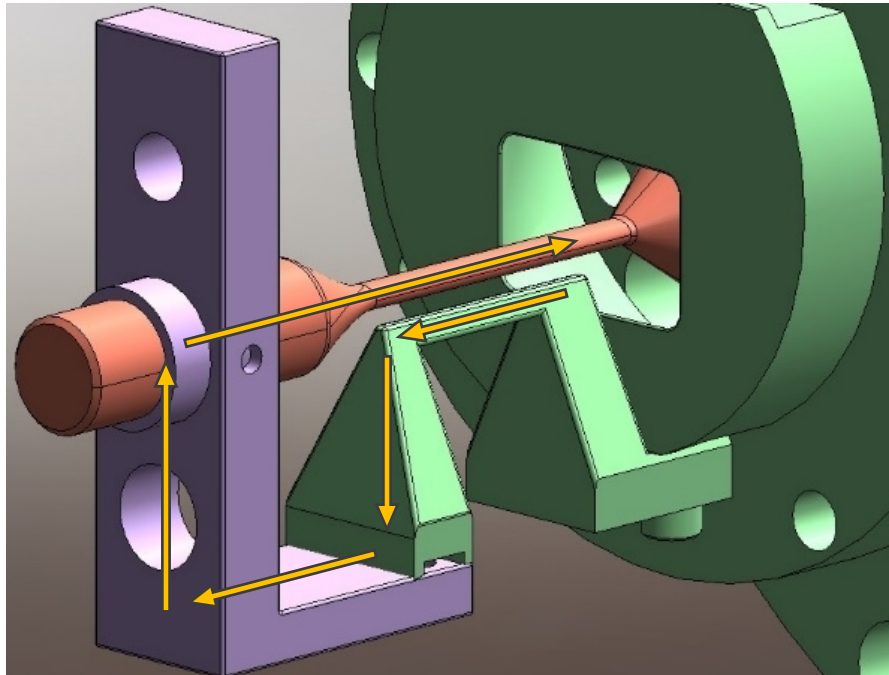
# Power Flow Physics of a MITL (Specifically P<sup>3</sup> at Mykonos)



[3] S. P. Bugaev et al., Sov. Phys. Usp. 18 51 (1975)

[4] T. W. L. Sanford et al., J. Appl. Phys. 1 July (1989)

An existing platform on Mykonos provides diagnostically accessible A-K gap geometry scaled to match Z's inner MITL field strengths



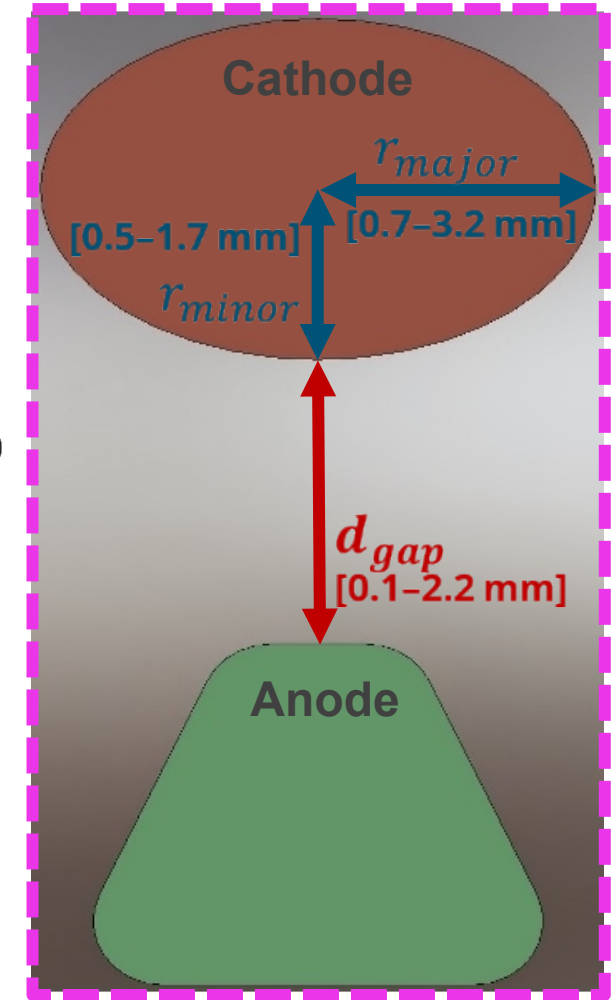
$$1/d_{\text{gap}} \propto |\vec{E}| \text{ (electric field)}$$

$$|\vec{E}| \sim 0.5\text{-}5 \text{ MV/cm}$$

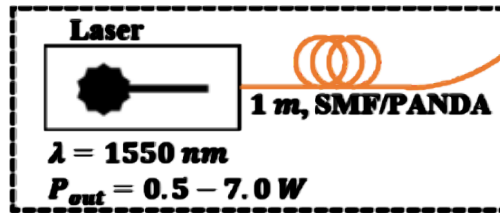
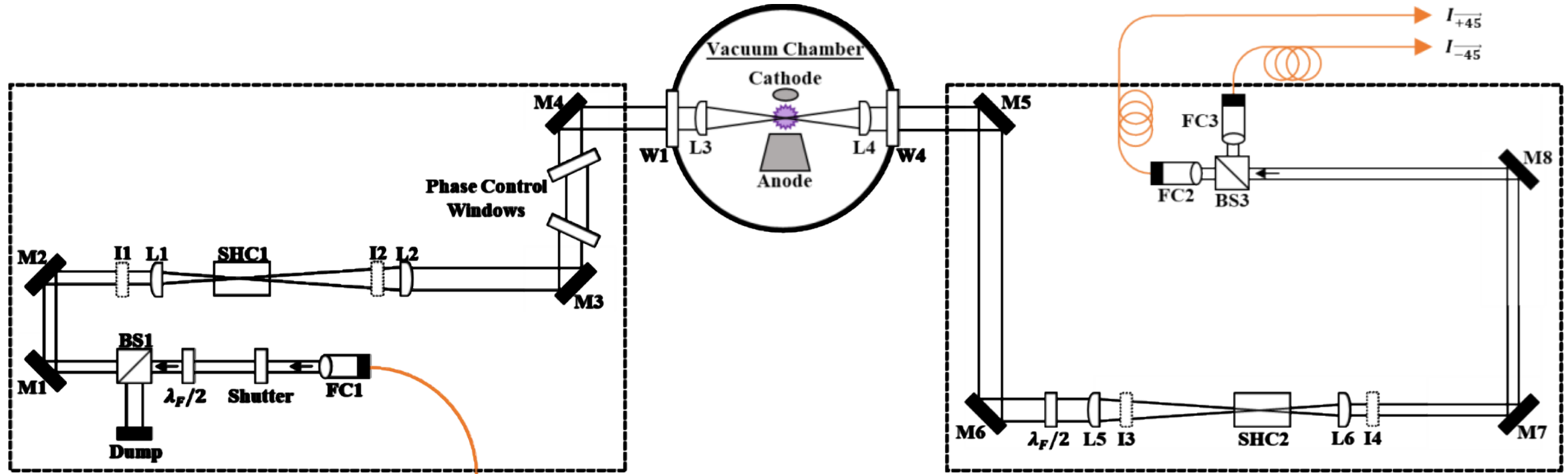
$$1/r_{\text{major}}, 1/r_{\text{minor}} \propto |\vec{J}| \propto |\vec{B}|$$

$$|\vec{B}| \sim 50\text{-}500 \text{ T}$$

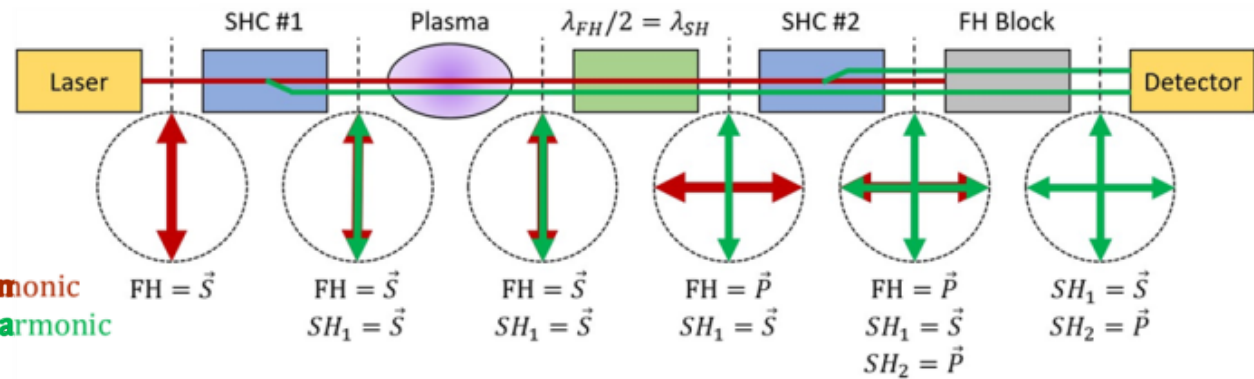
Cross Section: B



# Second-Harmonic Orthogonally Polarized Dispersion Interferometer (SHOP-DI) diagnostic design for Mykonos



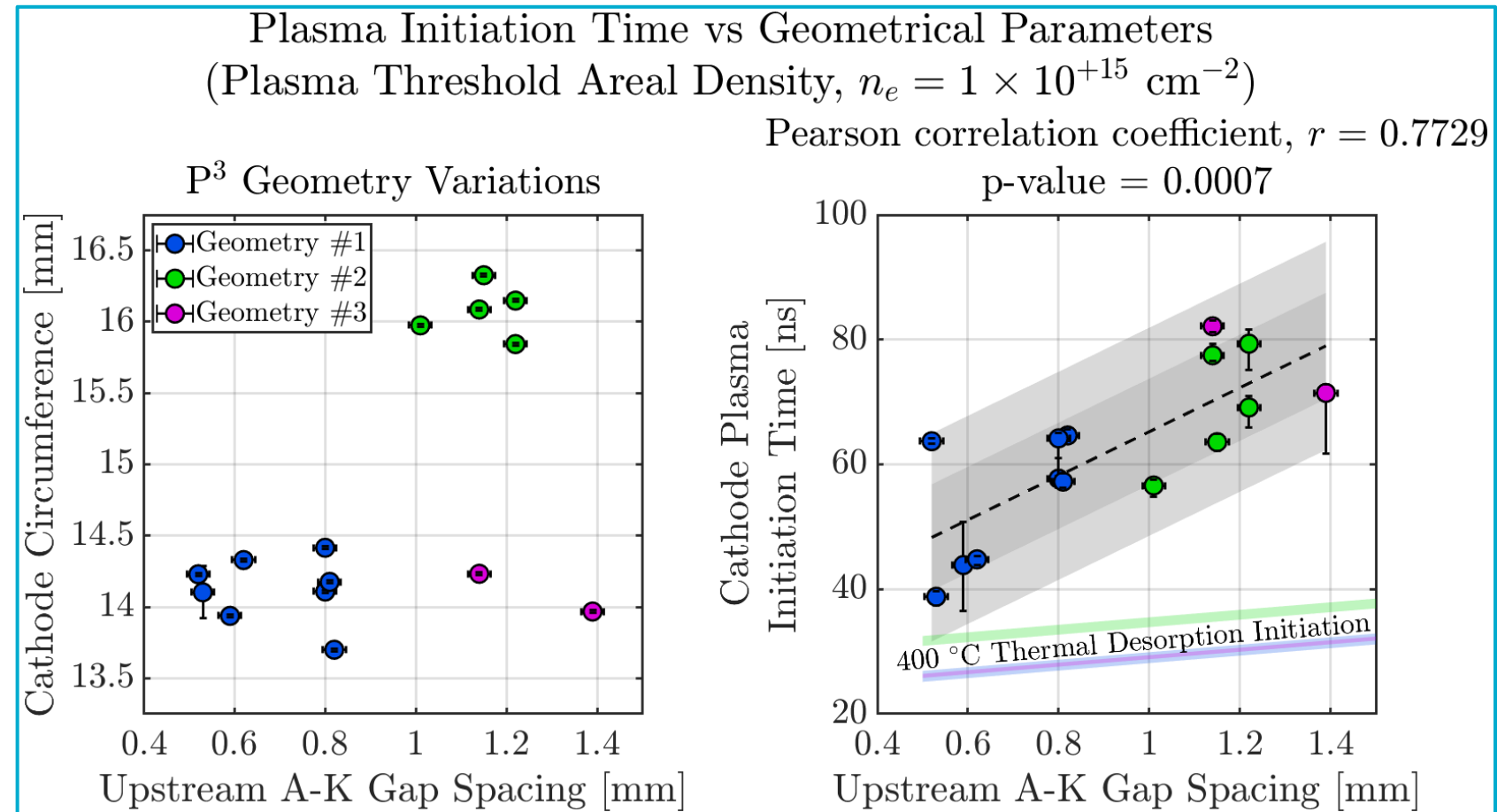
## Basic Design:



# P<sup>3</sup> Hardware Geometric Effects on the Plasma Initiation Time



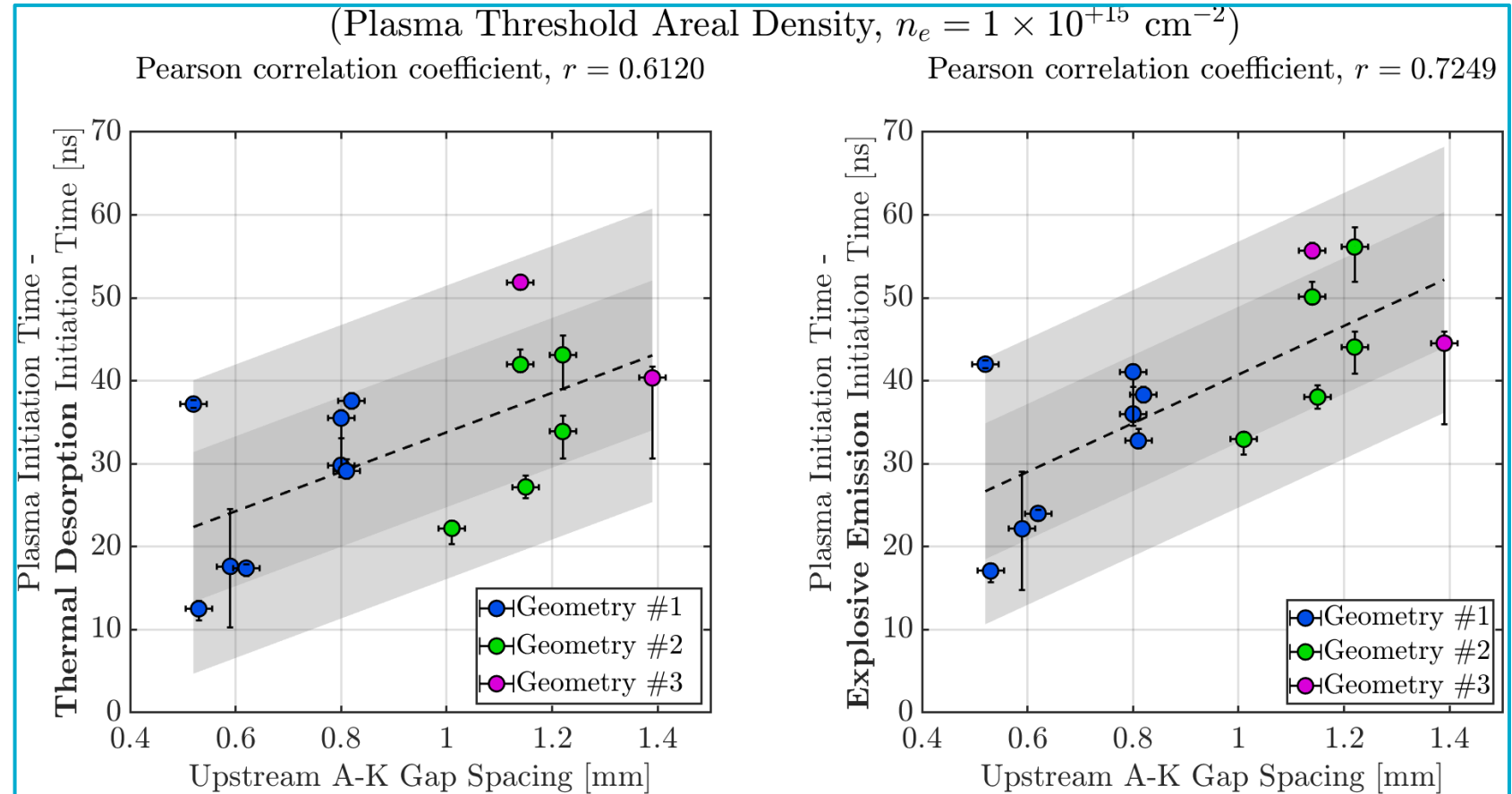
- Several P<sup>3</sup> geometries were utilized, varying the Cathode circumference and A-K gap spacing:
  - Larger Cathode circumferences reduce the current density and thus reduce Ohmic heating.
    - ❖ Leads to an expected variance in 400 °C thermal desorption initiation of 7 ns.
  - Larger A-K gap spacings reduce both the E-field and current density, thus delaying field emitted particles and reducing ohmic heating.
    - ❖ Leads to an expected variance in 400 °C thermal desorption initiation of 6 ns.
    - ❖ Leads to a variance in the 240 kV/cm field emission initiation of 7 ns.



# Standard (400 °C) Thermal Emission vs. Enhanced E-Field (240 kV/cm) Explosive Emission



- We can account for the variation in geometric effects by subtracting the expected thermal desorption initiation time.
  - A moderate correlation is shown between the  $10^{15} \text{ cm}^{-2}$  Cathode plasma detection time as a function of expected 400 °C thermal desorption initiation and the P<sup>3</sup> upstream A-K gap.
- However, the physics sourcing neutrals into the gap (that then ionize to source free electrons), may be dominated by field enhanced explosive emission.
  - A stronger correlation is made when factoring instead the earlier 240 kV/cm field enhanced explosive emission initiation.



# Why is There an A-K Gap Dependent Delay Between Plasma Initiation and Plasma Detection?

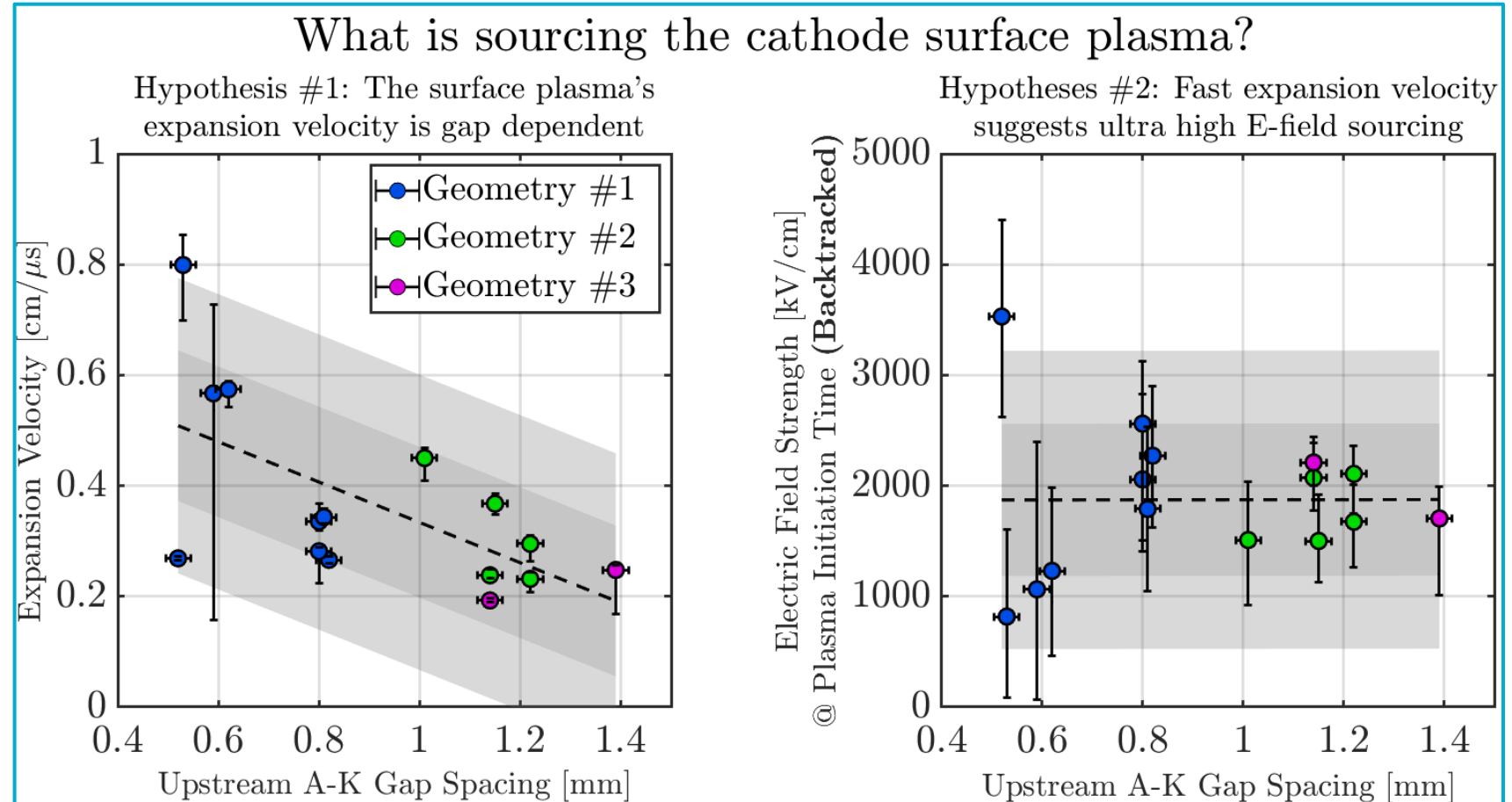


■ Perhaps the plasma cross-gap expansion velocity is reduced as the Electric field strength is reduced.

- A-K gaps expansion velocities would then vary between  $0.1 < v_p < 0.5 \text{ cm}/\mu\text{s}$ .

■ Maybe ultra high E-field sources ions that bombard the Cathode, creating the cathode plasma that expands at faster velocities.

- Backtrack formation time assuming a faster cathode plasma expansion velocity and known probing location.
- Find that plasma initiation consistently potentially occurred at Electric fields of  $\sim 1.87 \pm 0.66 \text{ MV/cm}$ .
- This matches the expected ultra high field strengths from [6]



[6] D. J. Johnson et al., IEEE Transactions on Dielectrics and Electrical Insulation, vol. 13, no. 1, pp. 52-64, Feb. (2006)

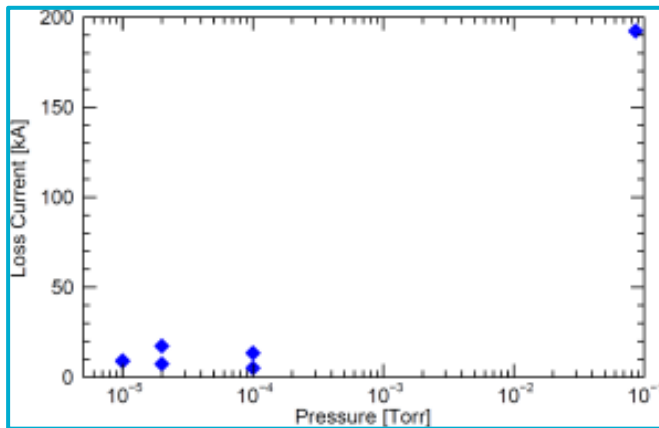
# The P<sup>3</sup> Vacuum Pressure Does Not Effect the Plasma Initiation



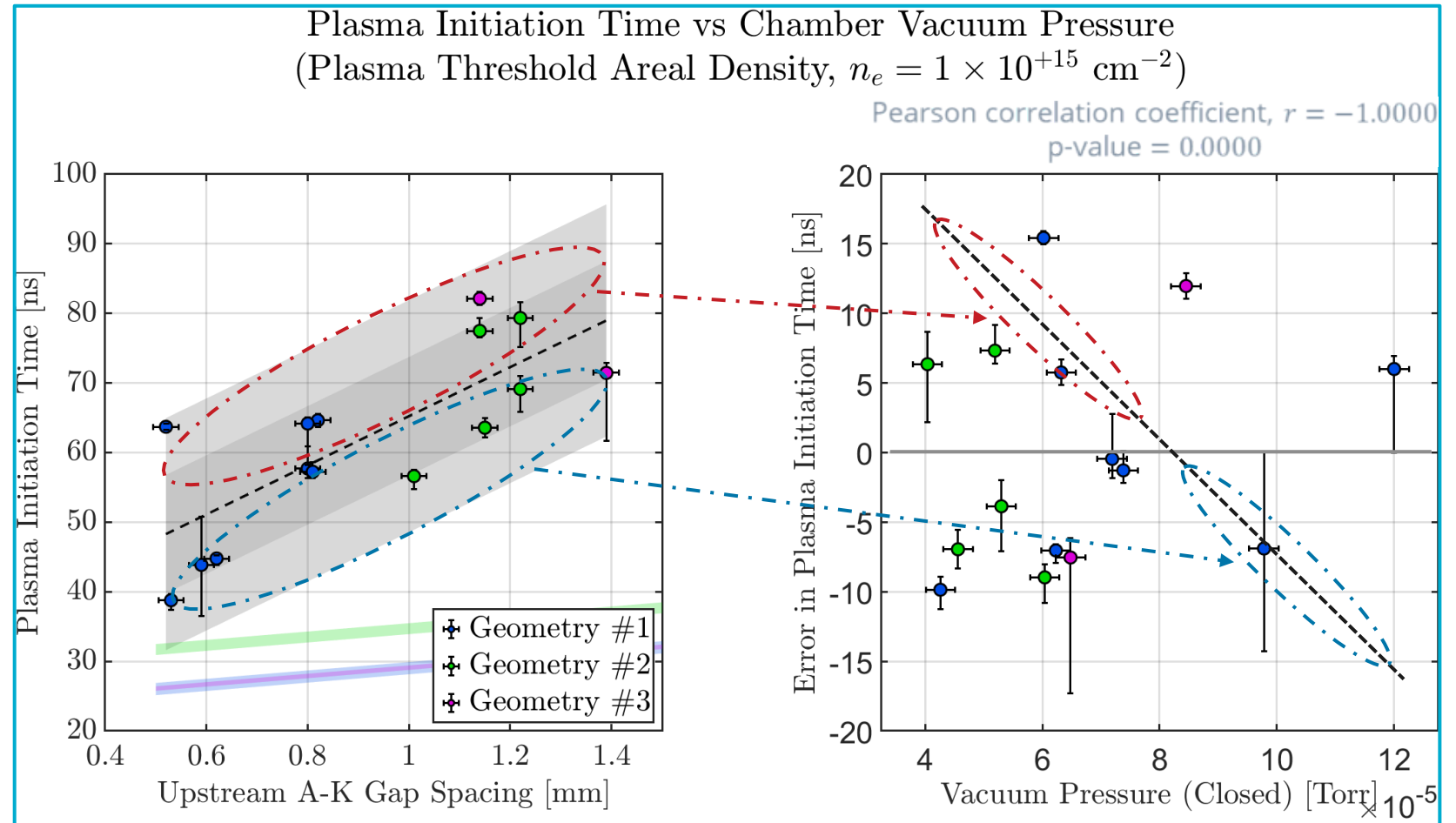
- Are the later than expected plasma formations caused due to less surface contaminants via a longer pump down duration (i.e. a lower downline vacuum pressure)?



Within the mid- to high- $10^{-5}$  Torr vacuum pressure region, there is no significant correlation to the plasma initiation times for these P<sup>3</sup> hardware sets.



- There is evidence of more extreme vacuum pressure differences effecting current loss.



[7] B. Hutsel et al., "Millimeter-Gap Magnetically Insulated Transmission Line Power Flow Experiments," Sandia National Lab.(SNL-NM), Albuquerque, NM (United States), 2014.

# Conclusions



- There is a need to better understand Power Flow physics and Current Loss contributors.
- The SHOP-DI diagnostic at SNL can measure electrode plasma free electron effective areal densities that are expected to form in the inner MITL and convolute regions of TW-class accelerators like the Z machine.
  - $\langle n_e L \rangle_{\min} = 6.3 \times 10^{13} \text{ [cm}^{-2}]$  (4.11 [mrad] sensitivity)
  - $\Delta \langle n_e L \rangle_{\max} = 4.8 \times 10^{16} \text{ [cm}^{-2}/\text{ns}]$  (2 [GHz] bandwidth)
- Preliminary experimental studies of Cathode plasma formation were conducted.
  - Variety of Ohmic heating rates and E-field strengths.
    - ❖ Perhaps the cathode plasma expansion velocity is lower than thought and is inversely related to A-K gap.
    - ❖ Or maybe plasma expansion velocity is higher and ultra high E-field cathode plasma initiation is occurring.
  - Mid- to high-  $10^{-5}$  Torr vacuum pressure variances do not play a significant role.

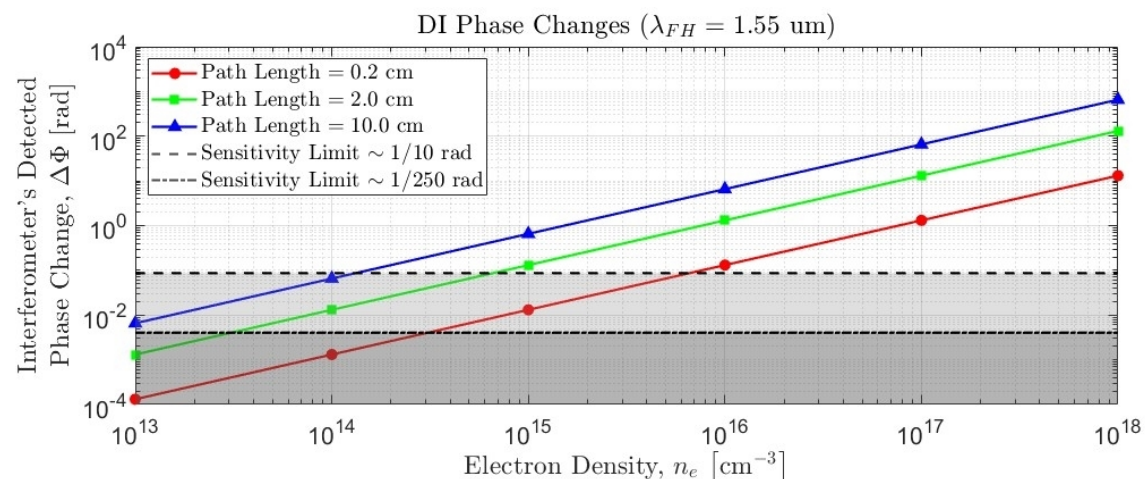
# SUPPLEMENTAR Y SLIDES

# Conclusion

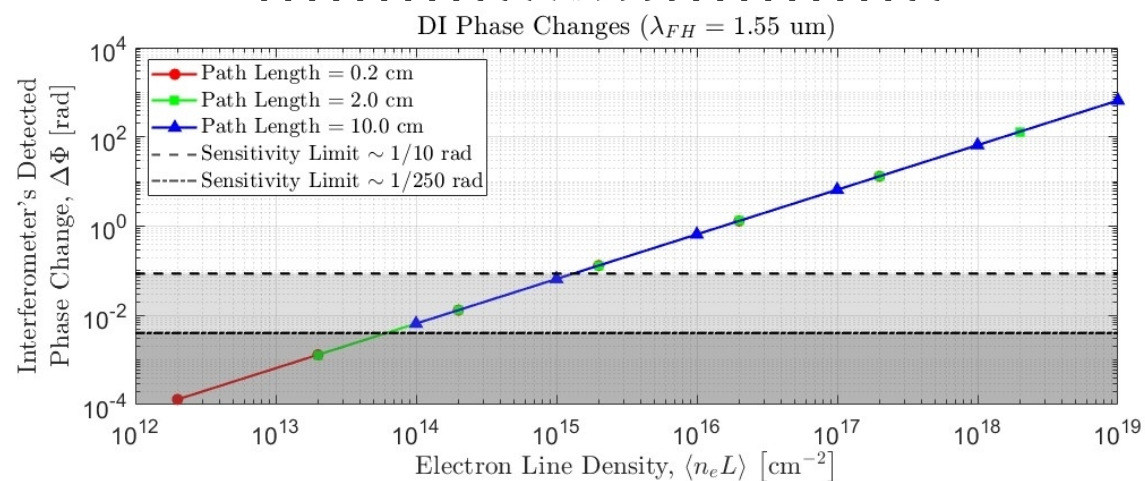


- The areal density is considered “free electron effective” since it’s calculated assuming the plasma’s **refractive index** is solely affected by free electrons.
  - There may in fact be substantial bound electrons that are affecting (positively) the detected phase change, thus affecting (negatively) the calculated electron areal density below reality. So, this is really a minimum free electron areal density measurement.

# SUPPLEMENTARY PLOTS OF PARAMETER RELATIONSHIPS



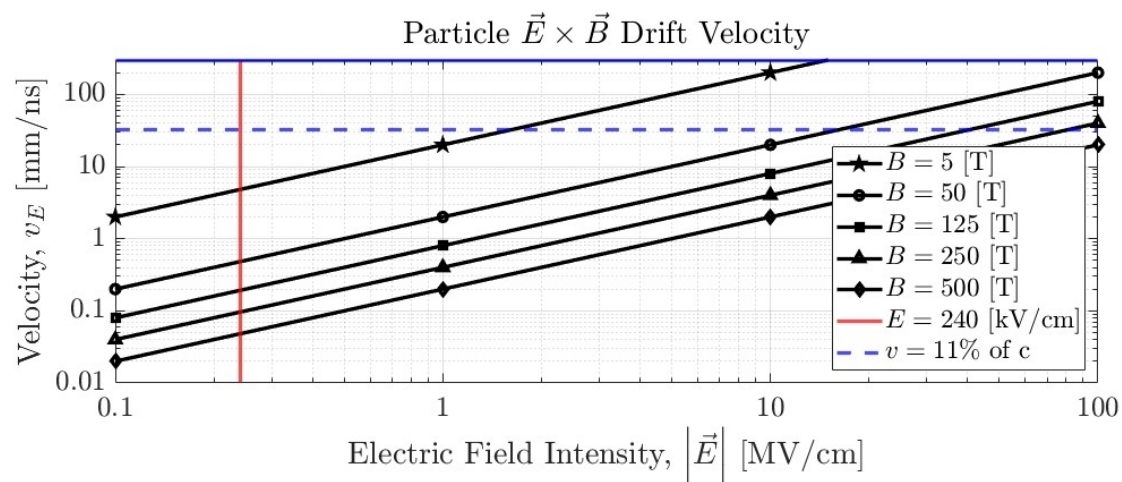
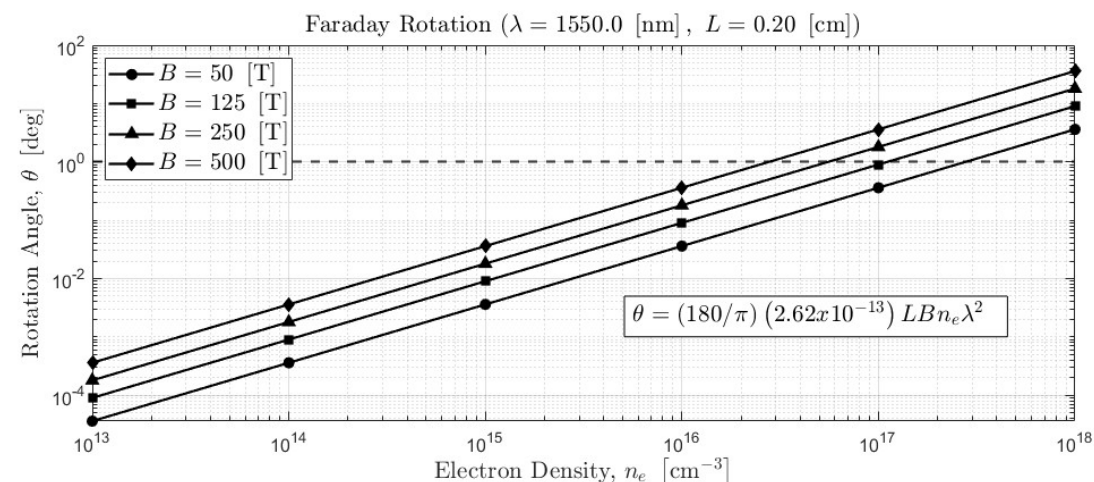
$$\langle n_e l \rangle [\text{cm}^{-2}] = (-1.53 \times 10^{16}) \Delta\Phi [\text{rad}]$$



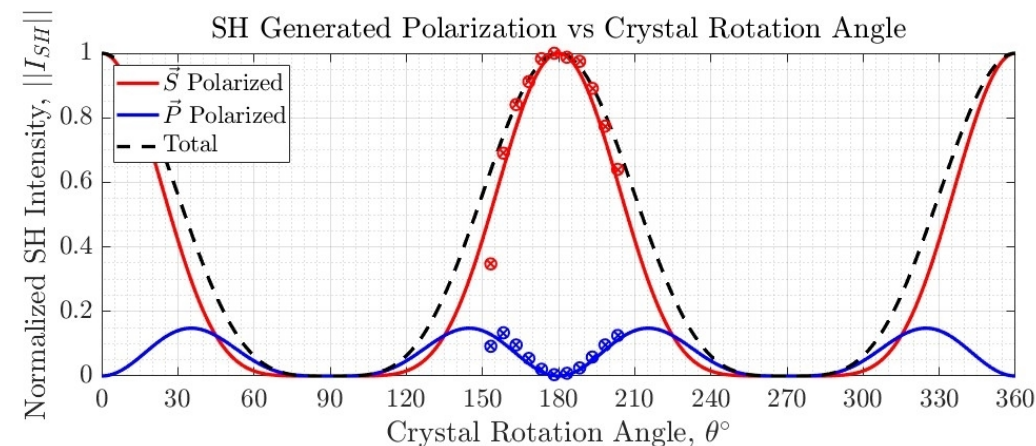
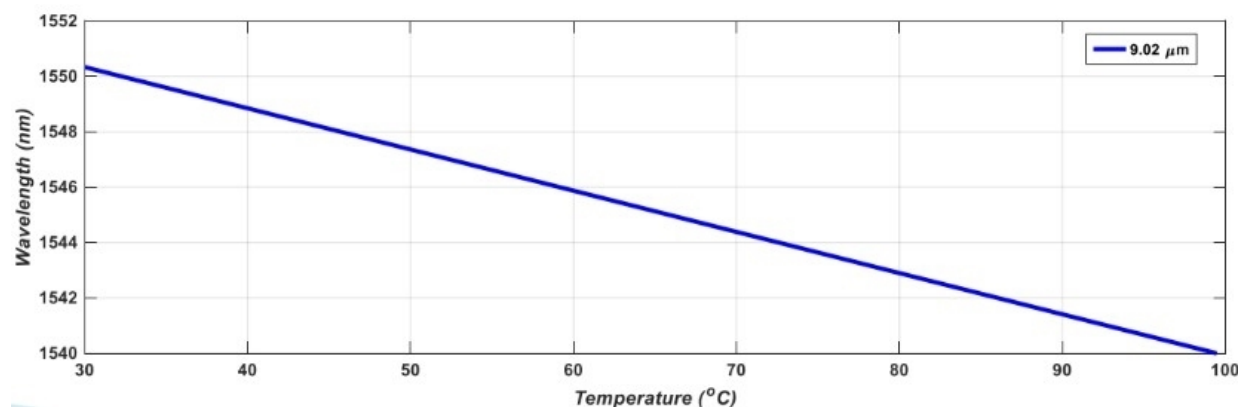
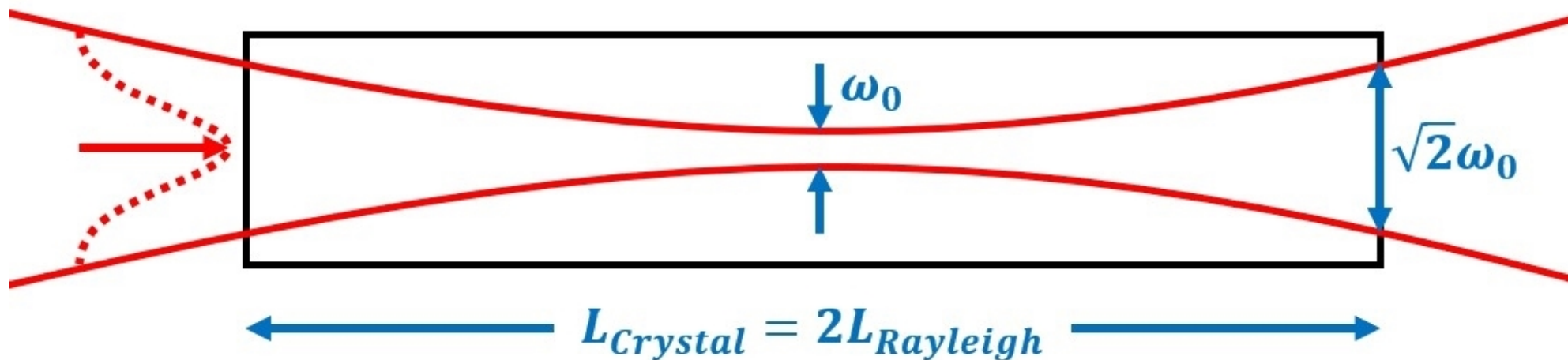
R-L Wave:  $\vec{k} \parallel \vec{B}$

$$\tilde{n}^2 = 1 - \frac{\omega_p^2/\omega^2}{1 - (\omega_c/\omega)} \quad (\text{R Wave})$$

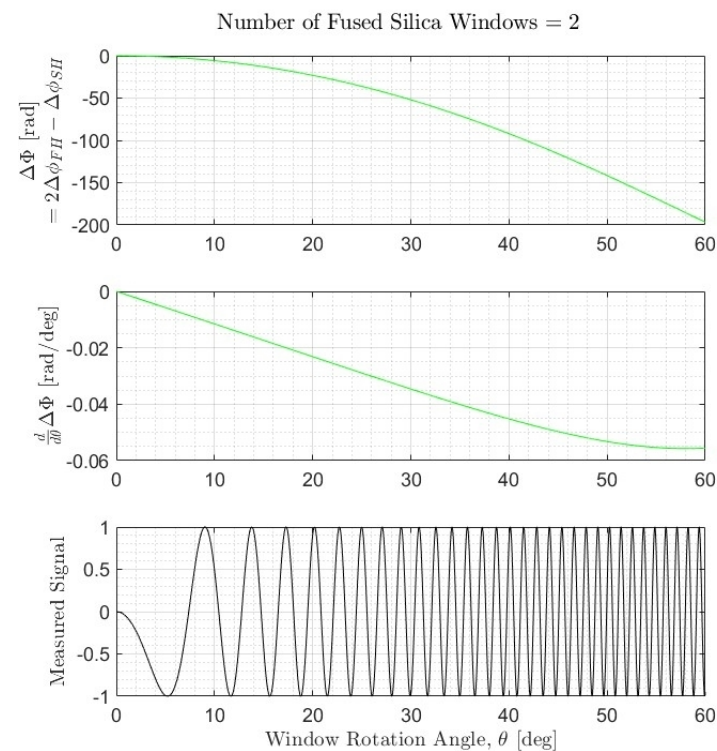
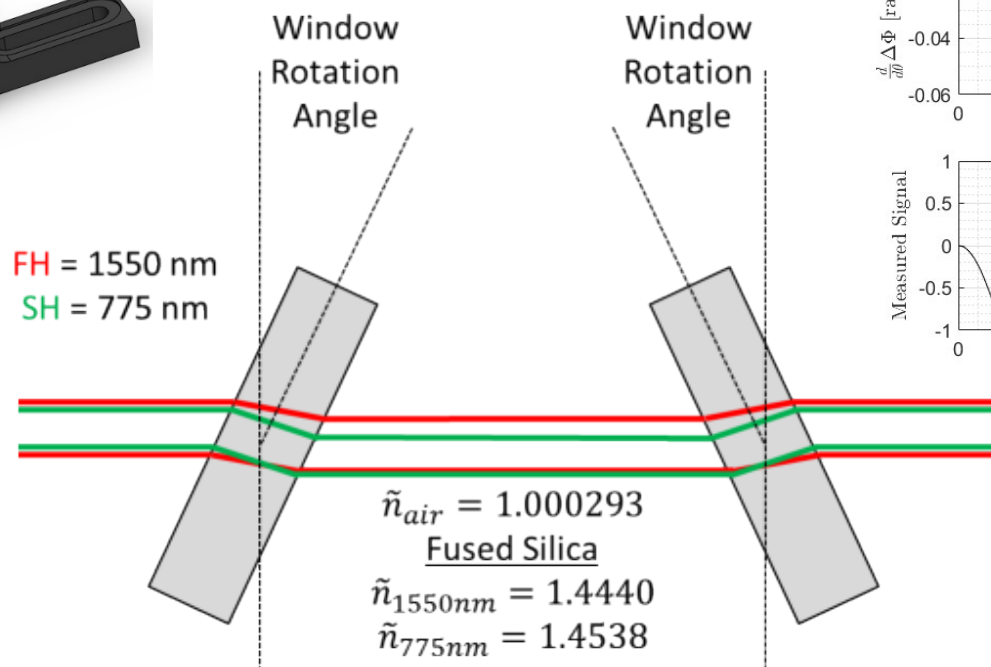
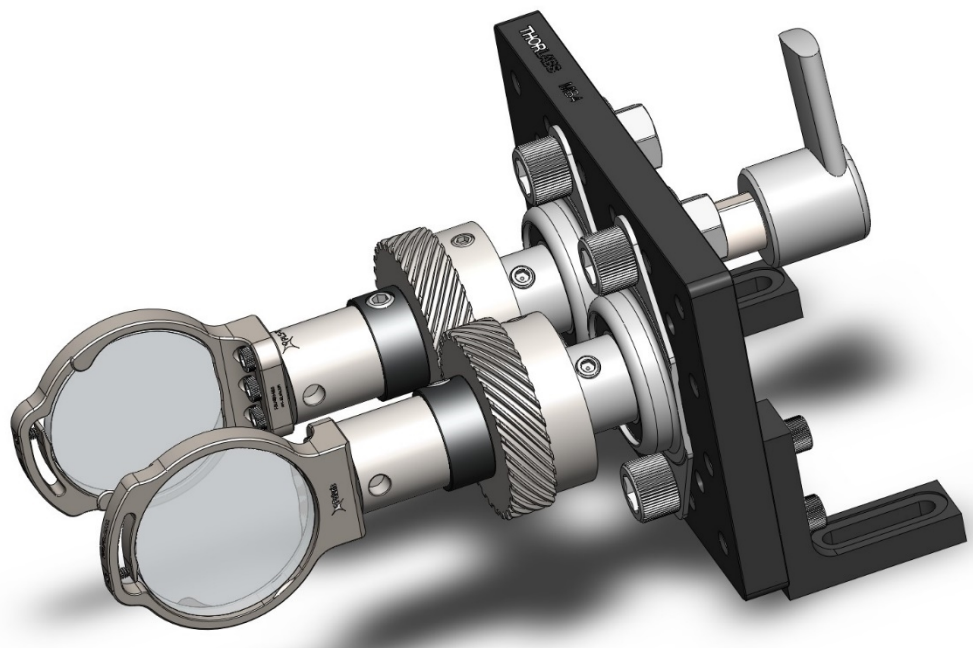
$$\tilde{n}^2 = 1 - \frac{\omega_p^2/\omega^2}{1 + (\omega_c/\omega)} \quad (\text{L Wave})$$



# SECOND HARMONIC GENERATION EFFICIENCY IS FUNCTION OF TEMPERATURE, BEAM FOCUSING PARAMETERS, AND POLARIZATION



# Phase Control Windows

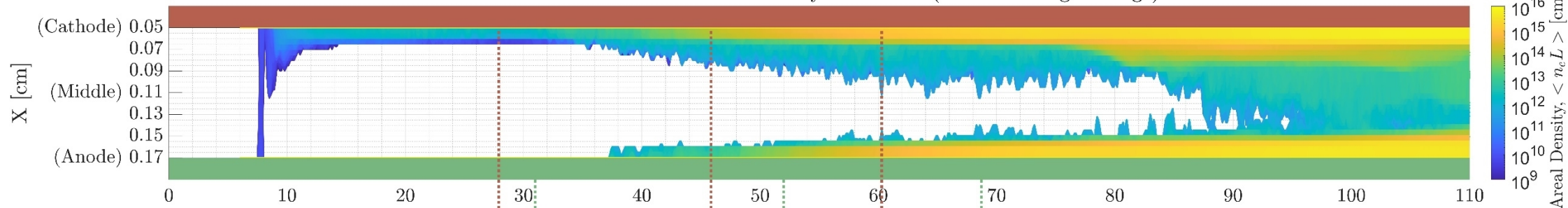


# CHICAGO & ALEGRA simulations

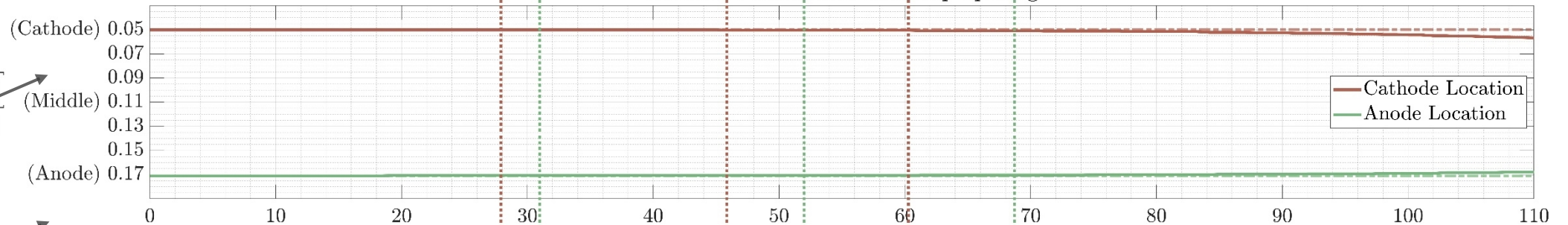


Cathode:  $1.50 \times 1.00$  mm, Gap: 1.21 mm

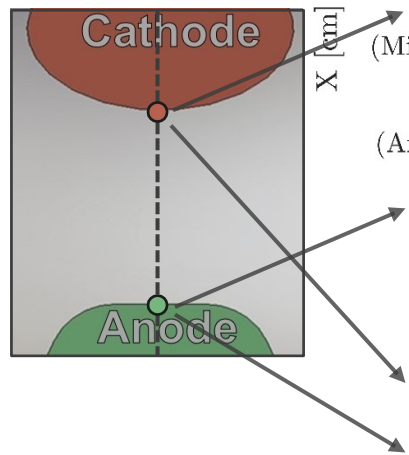
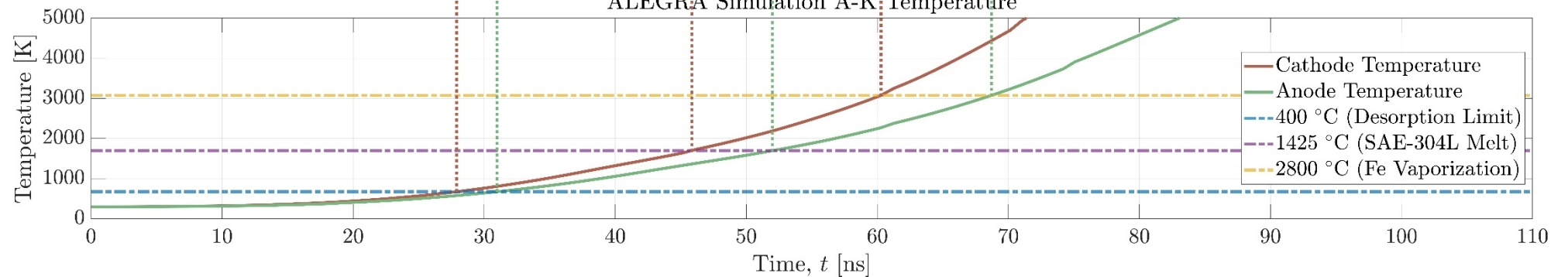
CHICAGO Simulation Areal Density Countours (0.2 ns Moving Average)



ALEGRA Simulation A-K Gap Spacing

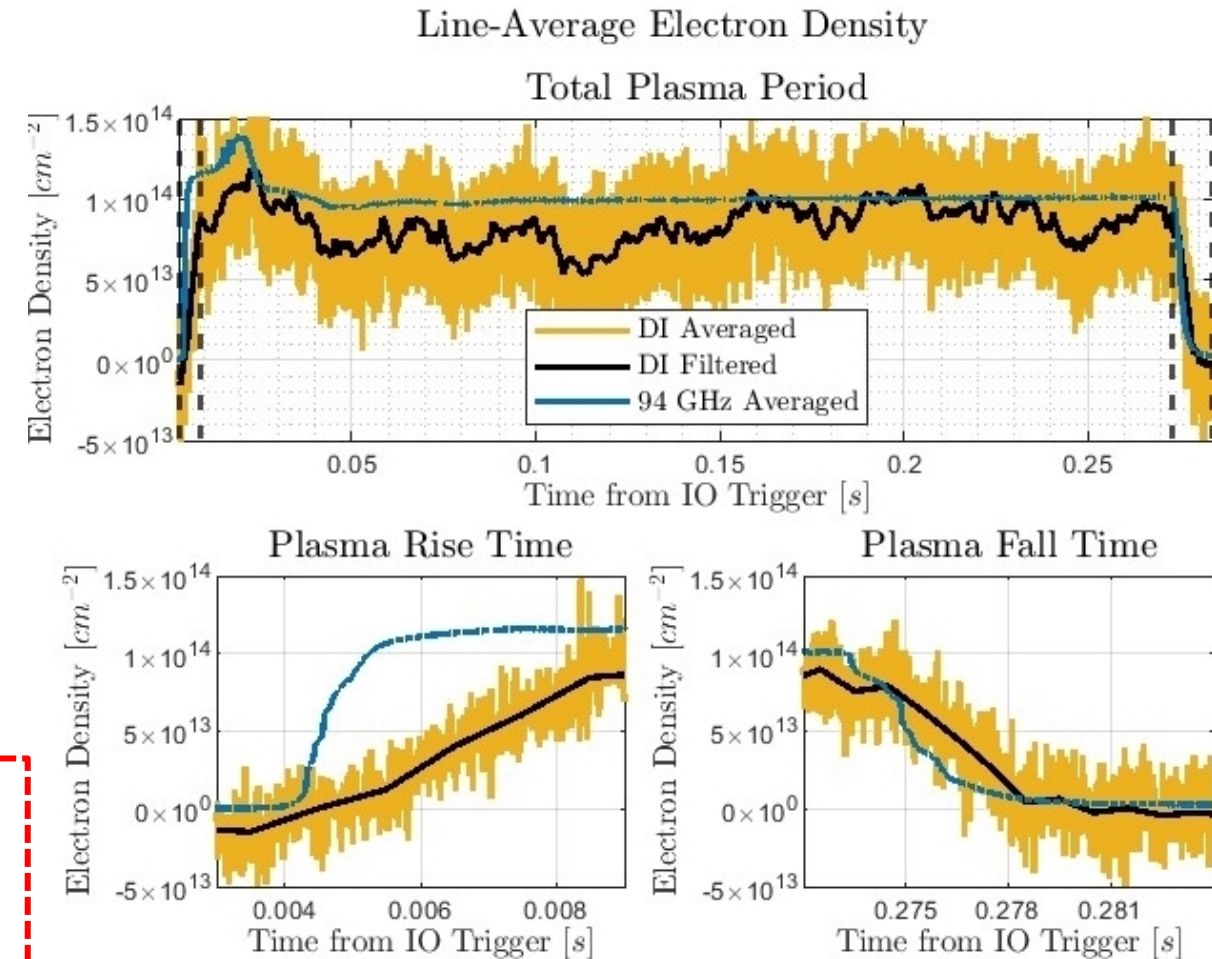
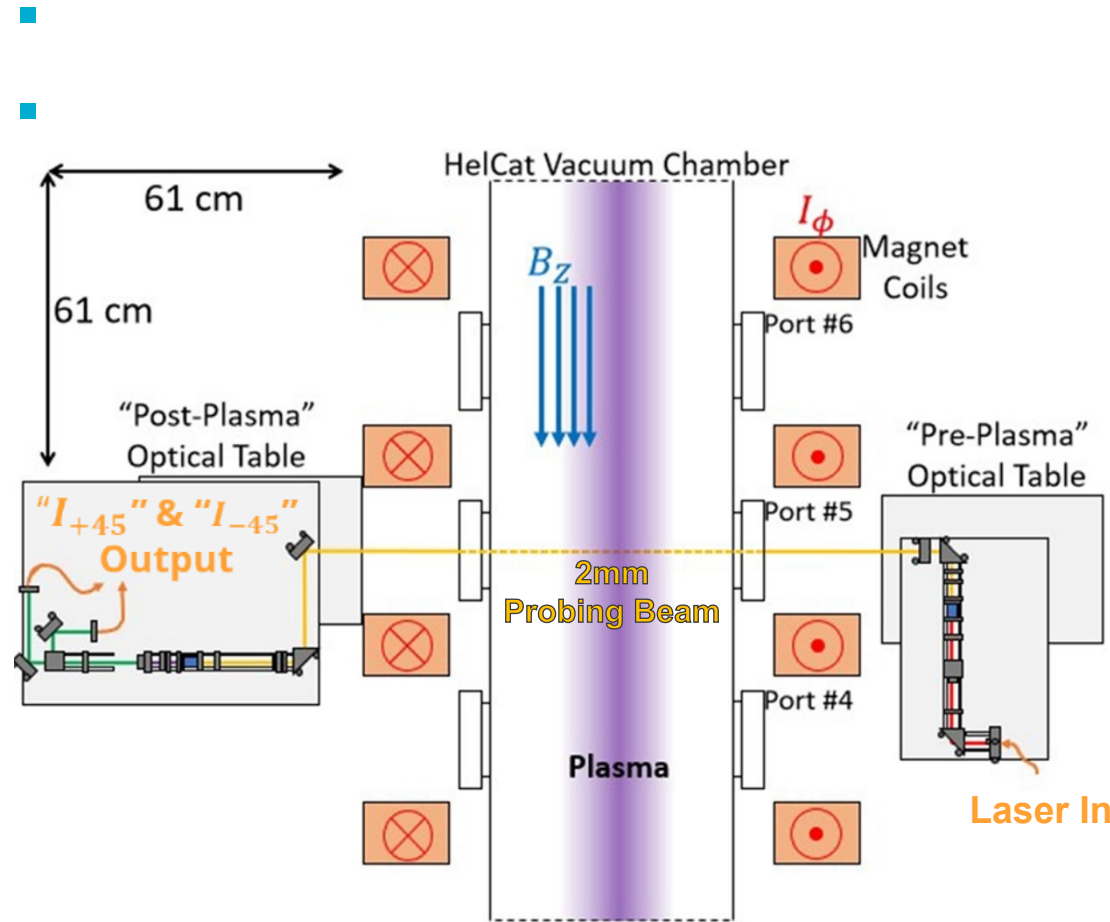


ALEGRA Simulation A-K Temperature



# SHOP Interferometer Install on the UNM HelCat Plasma Device

## Verified Functionality for Low Density Measurements



[5] N. R. Hines et al., "A fiber-coupled dispersion interferometer for density measurements of pulsed power transmission line electron sheaths on Sandia's Z machine," Review of Scientific Instruments, vol. 93, no. 11, p. 113505, November 2022.