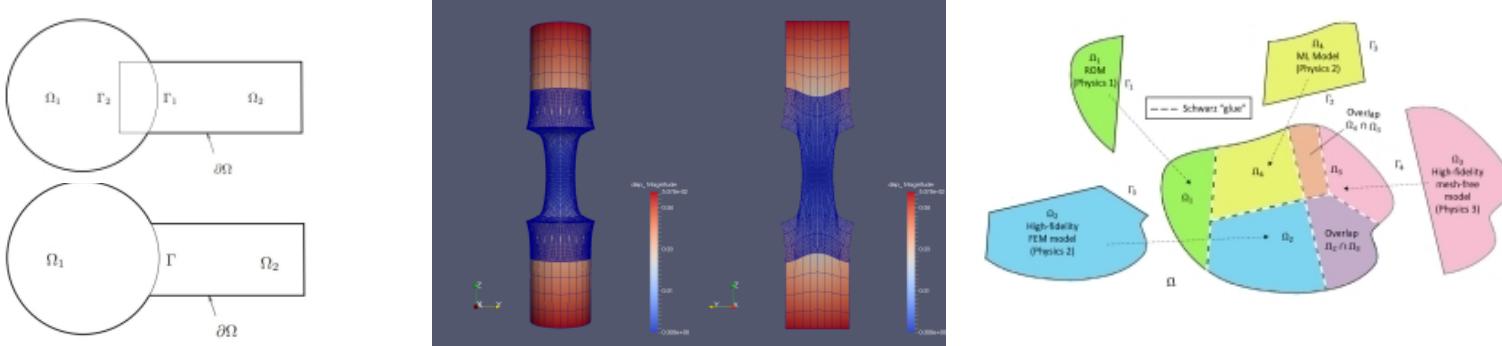




Accelerating mod/sim workflows through hybrid domain decomposition-based models and the Schwarz alternating method



Irina Tezaur¹, Alejandro Mota¹, Coleman Alleman¹, Greg Phlipot²,
 Chris Wentland¹, Francesco Rizzi³, Joshua Barnett⁴

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⁴Cadence Design Systems

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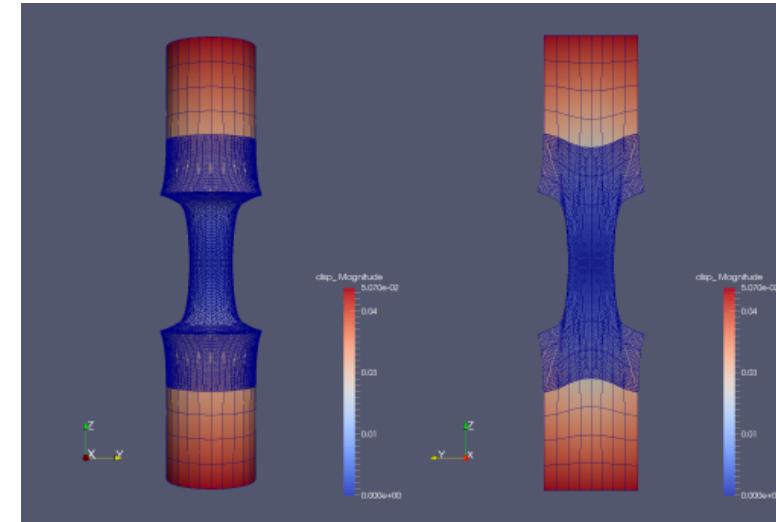
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Outline



1. Schwarz Alternating Method for Coupling of Full Order Models (FOMs) in Solid Mechanics

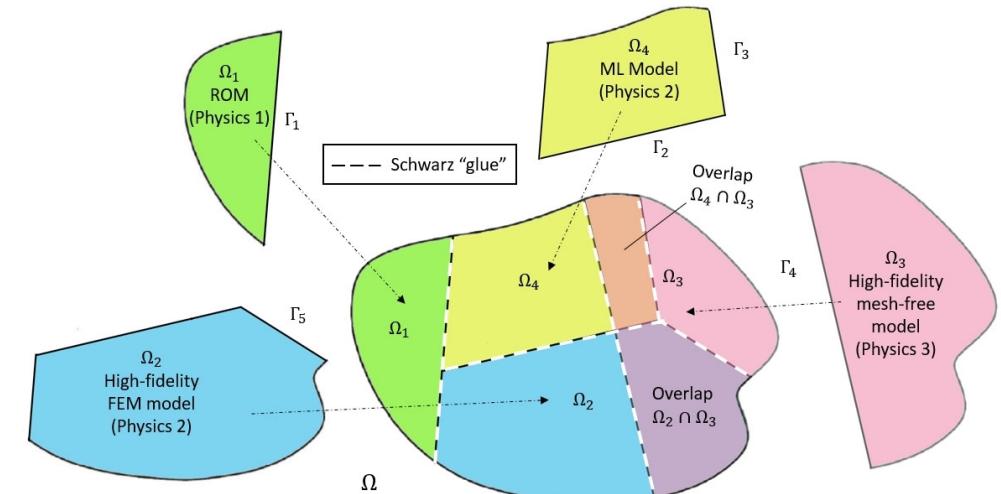
- Motivation & Background
- Quasistatic Formulation
 - Numerical Examples
- Extension to Dynamics
 - Numerical Examples



2. Schwarz Alternating Method for FOM-ROM* and ROM-ROM Coupling

- Motivation & Background
- Formulation
- Numerical Examples

3. Summary and Future Work

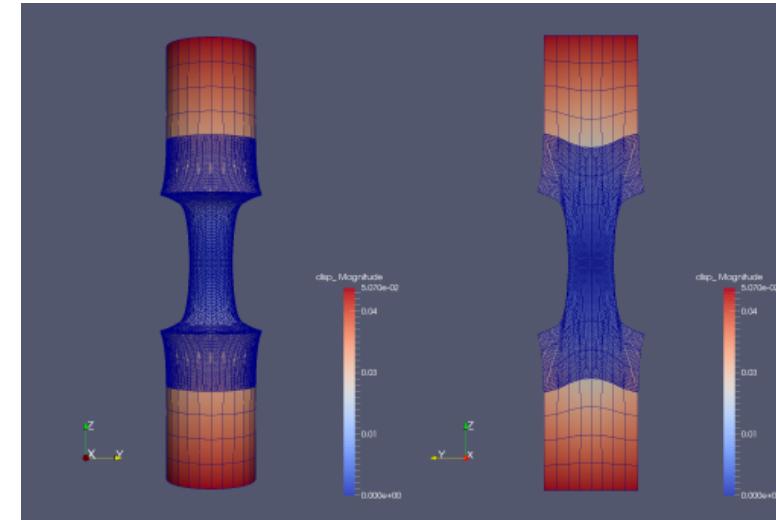


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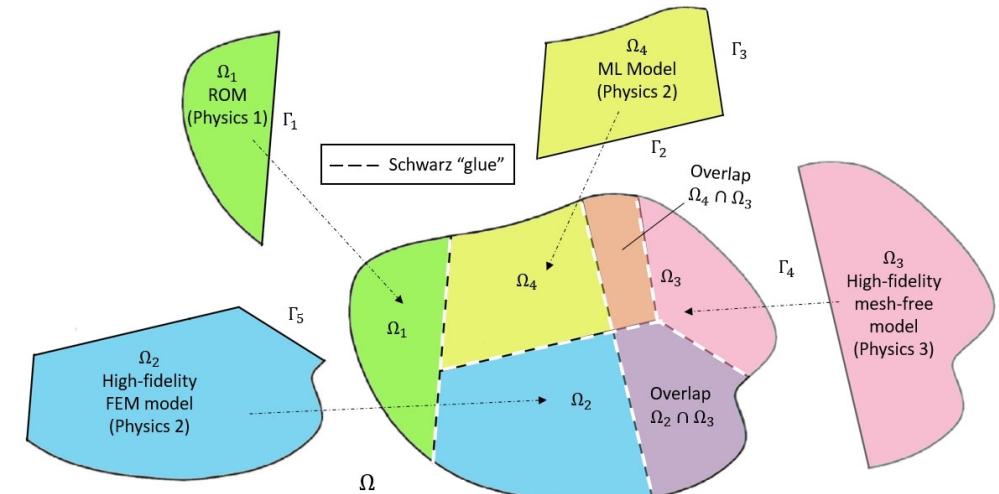
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Motivation for Coupling in Solid Mechanics

Concurrent multiscale coupling for predicting failure

- **Large scale** structural **failure** frequently originates from **small scale** phenomena such as defects, microcracks, inhomogeneities and more, which grow quickly in unstable manner
- Failure occurs due to **tightly coupled interaction** between small scale (stress concentrations, material instabilities, cracks, etc.) and large scale (vibration, impact, high loads and other perturbations)
- **Concurrent multiscale methods** are **essential** for understanding and prediction of behavior of engineering systems when a **small scale failure** determines the performance of the entire system

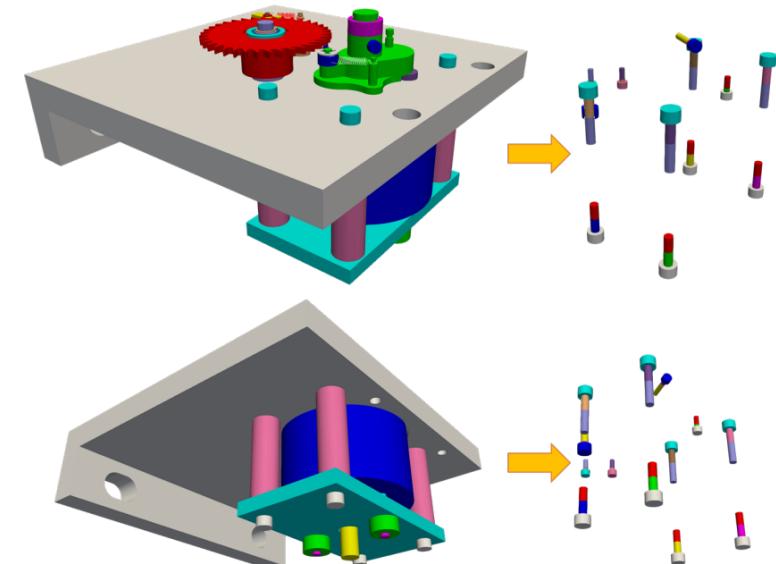
Simplification of mesh generation

- Creating a **high-quality mesh** for a **single component** can take **weeks**, making it “the single biggest bottleneck in analyses” [Sandia Lab News, 2020]!

Goal: develop a **concurrent multiscale coupling method** that is **minimally-intrusive** to implement into large HPC codes and can **simplify** the task of **meshing** complex geometries.



Roof failure of Boeing 737 aircraft due to fatigue cracks. From *imechanica.org*

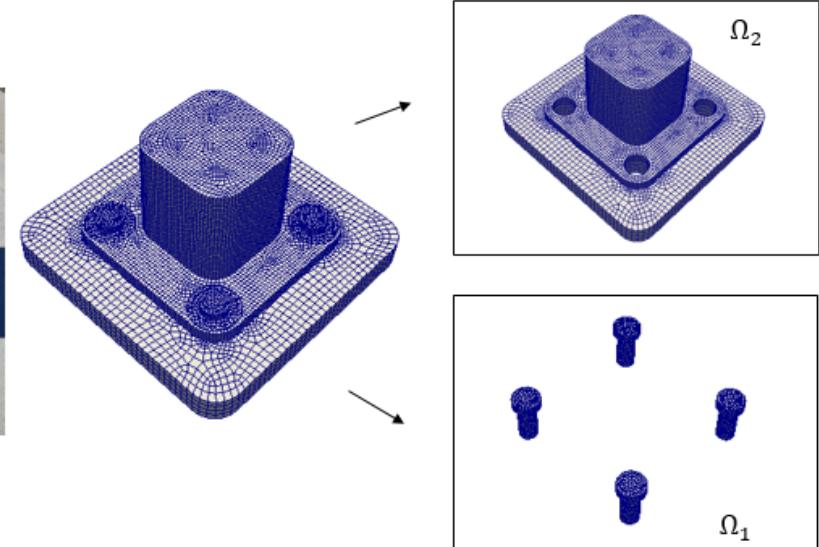


Schematic of difficult-to-mesh ratcheting mechanism with multiple threaded fasteners. From Parish *et al.*, 2024.

Requirements for Multiscale Coupling Method



- Coupling is ***concurrent*** (two-way)
- ***Ease of implementation*** into existing massively-parallel HPC codes
- “***Plug-and-play***” ***framework***: simplifies task of meshing complex geometries
 - Ability to couple regions with ***different non-conformal meshes, different element types and different levels of refinement***
 - Ability to use ***different solvers/time-integrators*** in different regions
- ***Scalable, fast, robust*** (we target ***real*** engineering problems, e.g., analyses involving failure of bolted components!)
- Coupling does not introduce ***nonphysical artifacts***
- ***Theoretical*** convergence properties/guarantees

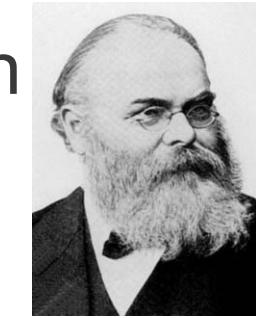


6 Schwarz Alternating Method for Domain Decomposition



- Proposed in 1870 by H. Schwarz for solving Laplace PDE on irregular domains.

Crux of Method: if the solution is known in regularly shaped domains, use those as pieces to iteratively build a solution for the more complex domain.



H. Schwarz (1843-1921)

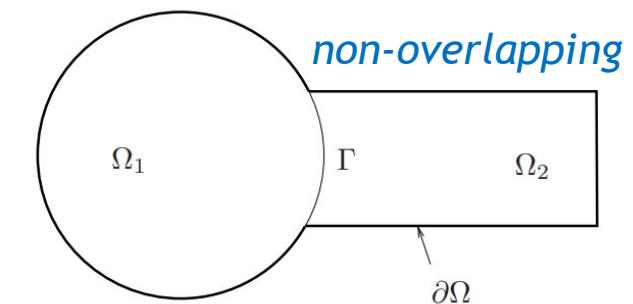
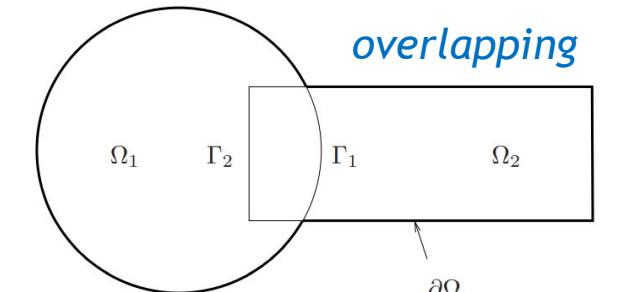
Basic Schwarz Algorithm

Initialize:

- Solve PDE by any method on Ω_1 w/ initial guess for transmission BCs on Γ_1 .

Iterate until convergence:

- Solve PDE by any method on Ω_2 w/ transmission BCs on Γ_2 based on values just obtained for Ω_1 .
- Solve PDE by any method on Ω_1 w/ transmission BCs on Γ_1 based on values just obtained for Ω_2 .



- Schwarz alternating method most commonly used as a *preconditioner* for Krylov iterative methods to solve linear algebraic equations.

Idea behind this work: using the Schwarz alternating method as a *discretization method* for solving multi-scale or multi-physics partial differential equations (PDEs).

How We Use the Schwarz Alternating Method



AS A **PRECONDITIONER**
FOR THE LINEARIZED
SYSTEM



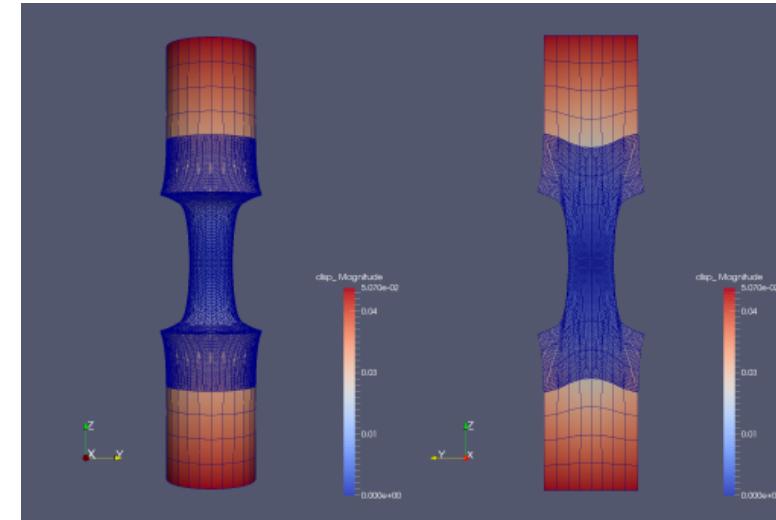
AS A **SOLVER** FOR THE
COUPLED
FULLY NONLINEAR
PROBLEM

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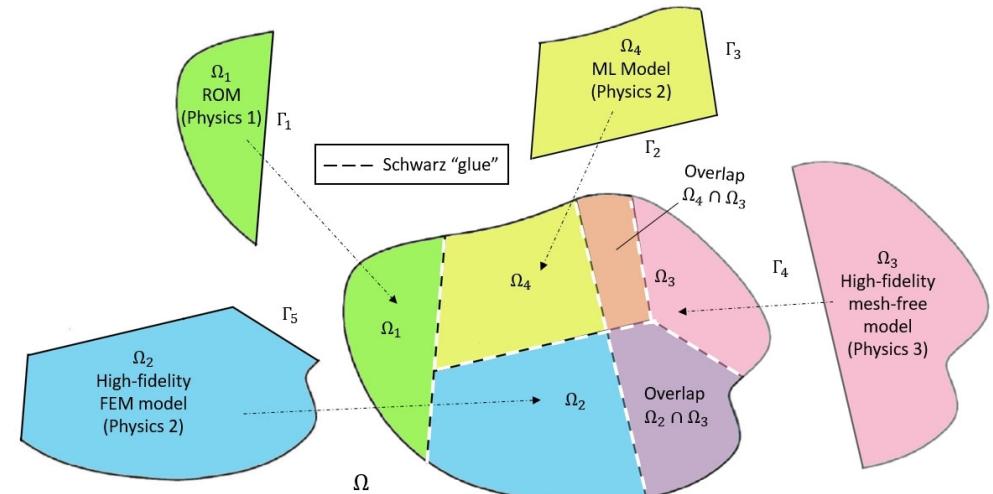
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- Energy functional defining weak form of the governing PDEs

$$\Phi[\boldsymbol{\varphi}] := \int_{\Omega} A(\mathbf{F}, \mathbf{Z}) dV - \int_{\Omega} \rho \mathbf{B} \cdot \boldsymbol{\varphi} dV$$

- $A(\mathbf{F}, \mathbf{Z})$: Helmholtz free-energy density
- $\mathbf{F} := \nabla \boldsymbol{\varphi}$: deformation gradient
- \mathbf{Z} : collection of internal variables (for plastic materials)
- ρ : density, \mathbf{B} : body force, $\mathbf{P} = \partial A / \partial \mathbf{F}$: Piola-Kirchhoff stress
- Euler-Lagrange equations, obtained by minimizing $\Phi[\boldsymbol{\varphi}]$: $\begin{cases} \operatorname{Div} \mathbf{P} + \rho \mathbf{B} = \mathbf{0}, & \text{in } \Omega \\ \boldsymbol{\varphi} = \boldsymbol{\chi}, & \text{on } \partial\Omega \end{cases}$
- Quasistatics solves **sequence of problems** in which loading (body force) \mathbf{B} is incremented quasistatically w.r.t. **pseudo time** t_i :

For $i = 1, \dots, n$

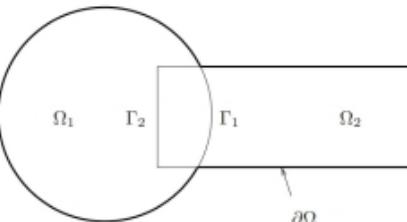
Solve $\operatorname{Div} \mathbf{P} + \rho \mathbf{B}(t_i) = \mathbf{0}$ with appropriate boundary conditions (BCs)
 Increment pseudo time t_i to obtain t_{i+1}

Spatial Coupling via (Multiplicative) Alternating Schwarz



Overlapping Domain Decomposition

$$\begin{cases} \operatorname{Div} \mathbf{P}_1^{(n+1)} + \rho \mathbf{B}(t_i) = \mathbf{0}, & \text{in } \Omega_1 \\ \boldsymbol{\varphi}_1^{(n+1)} = \chi, & \text{on } \partial\Omega_1 \setminus \Gamma_1 \\ \boldsymbol{\varphi}_1^{(n+1)} = \boldsymbol{\varphi}_2^{(n)} & \text{on } \Gamma_2 \\ \operatorname{Div} \mathbf{P}_2^{(n+1)} + \rho \mathbf{B}(t_i) = \mathbf{0}, & \text{in } \Omega_2 \\ \boldsymbol{\varphi}_2^{(n+1)} = \chi, & \text{on } \partial\Omega_2 \setminus \Gamma_2 \\ \boldsymbol{\varphi}_2^{(n+1)} = \boldsymbol{\varphi}_1^{(n+1)} & \text{on } \Gamma_2 \end{cases}$$



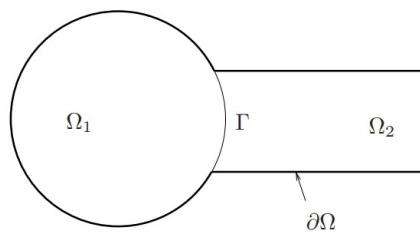
Model PDE:

$$\begin{cases} \operatorname{Div} \mathbf{P} + \rho \mathbf{B} = \mathbf{0}, & \text{in } \Omega \\ \boldsymbol{\varphi} = \chi, & \text{on } \partial\Omega \end{cases}$$

- Dirichlet-Dirichlet transmission BCs [Schwarz, 1870; Lions, 1988]

Non-overlapping Domain Decomposition

$$\begin{cases} \operatorname{Div} \mathbf{P}_1^{(n+1)} + \rho \mathbf{B}(t_i) = \mathbf{0}, & \text{in } \Omega_1 \\ \boldsymbol{\varphi}_1^{(n+1)} = \chi, & \text{on } \partial\Omega_1 \setminus \Gamma \\ \boldsymbol{\varphi}_1^{(n+1)} = \lambda_{n+1} & \text{on } \Gamma \\ \operatorname{Div} \mathbf{P}_2^{(n+1)} + \rho \mathbf{B}(t_i) = \mathbf{0}, & \text{in } \Omega_2 \\ \boldsymbol{\varphi}_2^{(n+1)} = \chi, & \text{on } \partial\Omega_2 \setminus \Gamma \\ \mathbf{P}_2^{(n+1)} \mathbf{n} = \mathbf{P}_2^{(n+1)} \mathbf{n}, & \text{on } \Gamma \\ \lambda_{n+1} = \theta \boldsymbol{\varphi}_2^{(n)} + (1 - \theta) \lambda_n, & \text{on } \Gamma, \text{ for } n \geq 1 \end{cases}$$



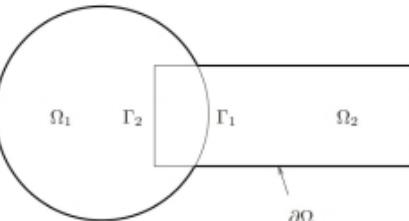
- Relevant for multi-material and multi-physics coupling
- Alternating Dirichlet-Neumann transmission BCs [Zanolli et al., 1987]
- Robin-Robin transmission BCs also lead to convergence [Lions, 1990]
- $\theta \in [0,1]$: relaxation parameter (can help convergence)

Spatial Coupling via (Multiplicative) Alternating Schwarz



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Model PDE:

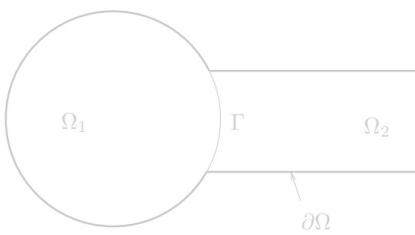
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- Dirichlet-Dirichlet transmission BCs [Schwarz, 1870; Lions, 1988]

Part 1 of talk

Non-overlapping Domain Decomposition

$$\begin{cases} \operatorname{Div} \mathbf{P}_1^{(n+1)} + \rho \mathbf{B}(t_i) = \mathbf{0}, & \text{in } \Omega_1 \\ \boldsymbol{\varphi}_1^{(n+1)} = \chi, & \text{on } \partial\Omega_1 \setminus \Gamma \\ \boldsymbol{\varphi}_1^{(n+1)} = \lambda_{n+1} & \text{on } \Gamma \\ \operatorname{Div} \mathbf{P}_2^{(n+1)} + \rho \mathbf{B}(t_i) = \mathbf{0}, & \text{in } \Omega_2 \\ \boldsymbol{\varphi}_2^{(n+1)} = \chi, & \text{on } \partial\Omega_2 \setminus \Gamma \\ \mathbf{P}_2^{(n+1)} \mathbf{n} = \mathbf{P}_2^{(n+1)} \mathbf{n}, & \text{on } \Gamma \\ \lambda_{n+1} = \theta \boldsymbol{\varphi}_2^{(n)} + (1 - \theta) \lambda_n, & \text{on } \Gamma, \text{ for } n \geq 1 \end{cases}$$



Part 2 of talk

- Relevant for multi-material and multi-physics coupling
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Additional Parallelism via Additive Schwarz



Multiplicative Overlapping Schwarz

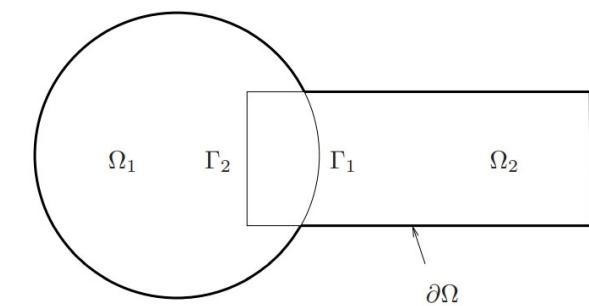
$$\begin{cases} \operatorname{Div} \mathbf{P}_1^{(n+1)} + \rho \mathbf{B}(t_i) = \mathbf{0}, & \text{in } \Omega_1 \\ \boldsymbol{\varphi}_1^{(n+1)} = \chi, & \text{on } \partial\Omega_1 \setminus \Gamma_1 \\ \boldsymbol{\varphi}_1^{(n+1)} = \boldsymbol{\varphi}_2^{(n)} & \text{on } \Gamma_2 \\ \\ \operatorname{Div} \mathbf{P}_2^{(n+1)} + \rho \mathbf{B}(t_i) = \mathbf{0}, & \text{in } \Omega_2 \\ \boldsymbol{\varphi}_2^{(n+1)} = \chi, & \text{on } \partial\Omega_2 \setminus \Gamma_2 \\ \boldsymbol{\varphi}_2^{(n+1)} = \boldsymbol{\varphi}_1^{(n+1)} & \text{on } \Gamma_2 \end{cases}$$

Additive Overlapping Schwarz

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Model PDE:

$$\begin{cases} \operatorname{Div} \mathbf{P} + \rho \mathbf{B} = \mathbf{0}, & \text{in } \Omega \\ \boldsymbol{\varphi} = \chi, & \text{on } \partial\Omega \end{cases}$$



- **Multiplicative Schwarz:** solves subdomain problems **sequentially** (in serial)
- **Additive Schwarz:** advance subdomains in **parallel**, communicate boundary condition data later
 - Typically requires a few more **Schwarz iterations**, but does not degrade **accuracy**
 - **Parallelism** helps balance additional **cost** due to Schwarz iterations
 - Applicable to both **overlapping** and **non-overlapping** Schwarz

Additional Parallelism via Additive Schwarz



Part 1 of talk

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Part 2 of talk

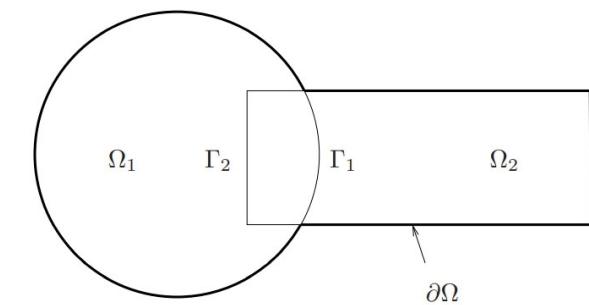
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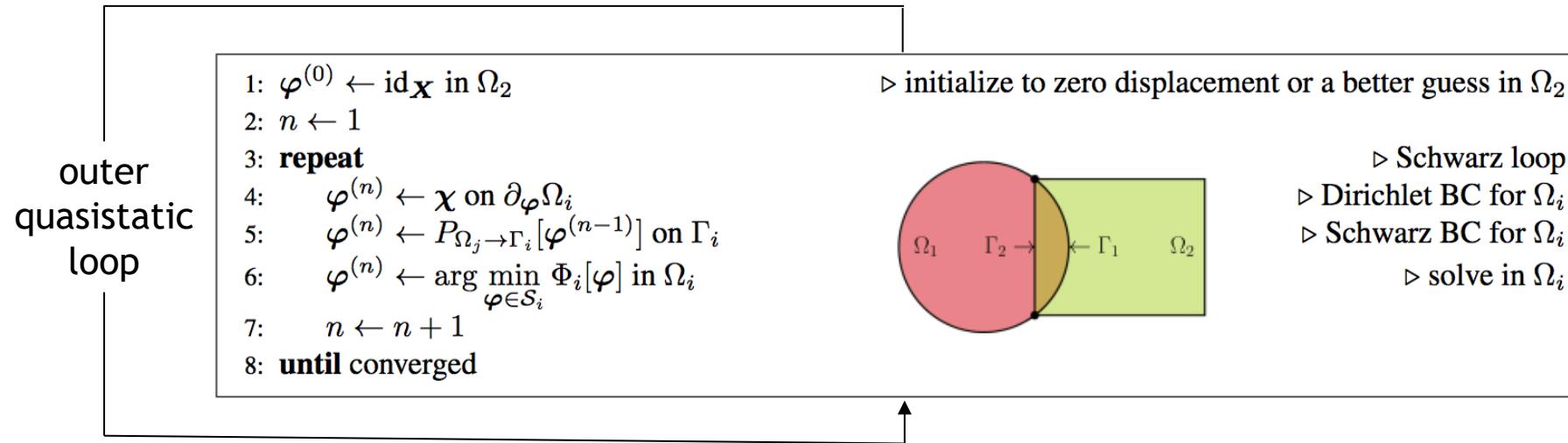
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 - Applicable to both **overlapping** and **non-overlapping** Schwarz

Overlapping Schwarz Coupling in Quasistatics



Advantages:

- Conceptually very *simple*.
- Allows the coupling of regions with *different non-conforming meshes, different element types, and different levels of refinement*.
- Information is exchanged among two or more regions, making coupling *concurrent*.
- *Different solvers* can be used for the different regions.
- *Different material models* can be coupled if they are compatible in the overlap region.
- Simplifies the task of *meshing complex geometries* for the different scales.

Convergence Proof*



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2 Formulation of the Schwarz Alternating Method

We start by defining the standard linear Galerkin variational formulation to establish notation before presenting the formulation of the coupling method.

2.1 Variational Formulation on a Single Domain

Consider the unit square $\Omega = [0, 1]^2$ with a traction boundary along the top edge $\Gamma = \{y=1\}$, where \mathbf{u}^0 is a displacement boundary, $(\mathbf{u}_0, \mathbf{p}_0)$ is a trial function. The traction boundary displacement is $\mathbf{u}^0 = \mathbf{u}_0 + \mathbf{p}_0 \mathbf{n}$ and the traction boundary conditions are $\mathbf{u}^0 \cdot \mathbf{n} = \mathbf{p}_0 \mathbf{n} \cdot \mathbf{n} = 0$. Let \mathbf{f} be the displacement gradient, $\mathbf{f} = \mathbf{f}^0$, and $\mathbf{f}^0 \cdot \mathbf{n} = \mathbf{f}^0 \mathbf{n} \cdot \mathbf{n} = 0$ be the body force, with \mathbf{f}^0 the traction in the reference configuration. Furthermore, introduce the energy functional

$$E(\mathbf{u}, \mathbf{p}) = \int_{\Omega} (\mathbf{f}, \mathbf{u}) \, d\Omega - \int_{\partial\Omega} \mathbf{u}^0 \cdot \mathbf{n} \, d\Gamma - \int_{\partial\Omega} \mathbf{p} \, d\Gamma, \quad (1)$$

in which $A(\mathbf{f}, \mathbf{u})$ is the Helmholtz free-energy density and \mathbf{f} is a collection of internal loads. The weak form of the energy functional is given by $\delta E(\mathbf{u}, \mathbf{p}) = 0$, where $\delta E(\mathbf{u}, \mathbf{p}) = \int_{\Omega} \delta \mathbf{u}^0 \cdot \mathbf{f} \, d\Omega + \int_{\partial\Omega} \delta \mathbf{p} \, d\Gamma$ is the variation of the energy functional that is composed of all the terms that are linearly independent and have a single sign that determines. Define $\mathcal{S} = \{\mathbf{u} \in H_0^1(\Omega) : \mathbf{u} = 0 \text{ on } \partial\Omega\}$.
(2)

and

$$V = \{\mathbf{u} \in H^1(\Omega) : \mathbf{u} = 0 \text{ on } \partial\Omega\}, \quad (3)$$

where $\mathcal{S} \subset V$ is a true subspace. The potential energy is minimized if and only if $\mathbf{u}^0 \in \mathcal{S}$ and $\mathbf{p} \in V$ for all $\mathbf{u} \in \mathcal{S}$ and $\mathbf{p} \in V$. It is also shown that the minimization of E is equivalent to that satisfied

$$\partial E(\mathbf{u}, \mathbf{p}) / \partial \mathbf{u} = \int_{\Omega} \mathbf{f} \, d\Omega - \int_{\partial\Omega} \mathbf{u}^0 \cdot \mathbf{n} \, d\Gamma - \int_{\partial\Omega} \mathbf{p} \, d\Gamma = 0, \quad (4)$$

where $\mathbf{f}^0 = \mathbf{f}^0 / \|\mathbf{f}^0\|$ represents the first Poincaré's state. The Euler-Lagrange equation corresponding to the variational statement (4) is

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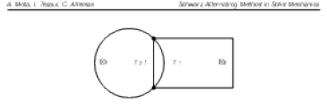


Figure 1: Two subdomains Ω_1 and Ω_2 and their corresponding boundary conditions 1 and 2 used for the Schwarz alternating method.

Now if $n = 1$ and $i = 2$ (i is odd), and $i = 2$ and $j = 1$ (i is even), introduce the following definition for each $\mathbf{u} \in V$:

- Closure: $\mathbf{u} \in \Omega_1 \cap \Omega_2$,
- Dirichlet boundary: $\mathbf{u} \in \Omega_2 \cap \partial\Omega$,
- Neumann boundary: $\mathbf{u} \in \Omega_2 \cap \Omega_1$,
- Relaxed boundary: $\mathbf{u} \in \Omega_1 \cap \partial\Omega$.

Now that with these definitions we guarantee that $\mathbf{u} \in \Omega_2 \cap \Omega_1 \Rightarrow \mathbf{u} \in \Omega_2 \cap \Omega_1 \cap \Omega_2$, and $\mathbf{u} \in \Omega_1 \cap \Omega_2 \Rightarrow \mathbf{u} \in \Omega_1 \cap \Omega_2 \cap \Omega_1$.

Since $\mathbf{u} \in \Omega_2 \cap \Omega_1 \cap \Omega_2$ we have $\mathbf{u} \in \Omega_2$ and $\mathbf{u} \in \Omega_1$, and $\mathbf{u} \in \Omega_1 \cap \Omega_2 \cap \Omega_1$ we have $\mathbf{u} \in \Omega_1$ and $\mathbf{u} \in \Omega_2$.
(7)

and

$$\mathbf{u} = \mathbf{u}_1 + \mathbf{u}_2 \quad \text{in } \Omega_2 \cap \Omega_1 \cap \Omega_2, \quad (8)$$

where $\mathbf{u}_1 = \mathbf{u} \in \Omega_1 \cap \Omega_2 \cap \Omega_1$ and $\mathbf{u}_2 = \mathbf{u} \in \Omega_2 \cap \Omega_1 \cap \Omega_2$.

where the operator $\mathcal{P}_{\Omega_2}(\mathbf{u}) = \int_{\Omega_2} \mathbf{f} \, d\Omega - \int_{\partial\Omega_2} \mathbf{u}^0 \cdot \mathbf{n} \, d\Gamma - \int_{\partial\Omega_2} \mathbf{p} \, d\Gamma$. The operator \mathcal{P}_{Ω_2} plays a prominent role in the Schwarz alternating method. The term \mathcal{P}_{Ω_2} is introduced to ensure that the operator \mathcal{P}_{Ω_2} is able to decompose the subdomain Ω_2 into two subdomains Ω_1 and Ω_2 . For the inversion of the operator \mathcal{P}_{Ω_2} it is necessary to assume that \mathcal{P}_{Ω_2} is invertible. The inverse of \mathcal{P}_{Ω_2} is given by $\mathcal{P}_{\Omega_2}^{-1}(\mathbf{u}) = \mathbf{u}_1 + \mathbf{u}_2$.

The scheme of averaging combined relations in a sequence of relations on Ω_2 and Ω_1 . The relations (7) provided

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$$\begin{aligned} 1. & \mathbf{u}_1^0 = \mathbf{u}_1^0 + \mathbf{u}_2^0 \quad \text{in } \Omega_2 \cap \Omega_1 \cap \Omega_2, & \text{it satisfies for } \mathbf{u}_1^0 \\ 2. & \mathbf{u}_2^0 = \mathbf{u}_2^0 + \mathbf{u}_1^0 \quad \text{in } \Omega_2 \cap \Omega_1 \cap \Omega_2, & \text{it satisfies for } \mathbf{u}_2^0 \\ 3. & \mathbf{u}_1^0 = \mathbf{u}_1^0 + \mathbf{u}_2^0 = \mathbf{u}_1^0 + \mathbf{u}_1^0 = \mathbf{u}_1^0, & \text{it is linear relation} \\ 4. & \mathbf{u}_2^0 = \mathbf{u}_2^0 + \mathbf{u}_1^0 = \mathbf{u}_2^0 + \mathbf{u}_2^0 = \mathbf{u}_2^0, & \text{it is linear relation} \\ 5. & \mathbf{u}_1^0 = \mathbf{u}_2^0 = \mathbf{u}_1^0 + \mathbf{u}_2^0 = \mathbf{u}_1^0, & \text{it is linear relation} \\ 6. & \mathbf{u}_2^0 = \mathbf{u}_1^0 = \mathbf{u}_2^0 + \mathbf{u}_1^0 = \mathbf{u}_2^0, & \text{it is linear relation} \\ 7. & \text{and } \left[\left(\mathbf{u}_1^0 \right)^2 + \left(\mathbf{u}_2^0 \right)^2 \right] = \left[\left(\mathbf{u}_1^0 + \mathbf{u}_2^0 \right)^2 \right]. & \text{it is linear relation} \end{aligned}$$

Appendix 1: Schwarz Alternating Method

Remark that $\|\mathbf{u}\|_1^2 = \|\mathbf{u}_1\|_1^2 + \|\mathbf{u}_2\|_1^2$, i.e. $\|\mathbf{u}_1\|_1^2 + \|\mathbf{u}_2\|_1^2 \geq \|\mathbf{u}\|_1^2$ in $\Omega_2 \cap \Omega_1 \cap \Omega_2$.

Theorem 1. Assume that the energy functional $\Phi[\varphi]$ satisfies properties 1–5 above. Consider the Schwarz alternating method of Section 2 defined by (9)–(13) and its equivalent form (39). Then

(a) $\Phi[\tilde{\varphi}^{(0)}] \geq \Phi[\tilde{\varphi}^{(1)}] \geq \dots \geq \Phi[\tilde{\varphi}^{(n-1)}] \geq \Phi[\tilde{\varphi}^{(n)}] \geq \dots \geq \Phi[\varphi]$, where φ is the minimizer of $\Phi[\varphi]$ over \mathcal{S} .

(b) The sequence $\{\tilde{\varphi}^{(n)}\}$ defined in (39) converges to the minimizer φ of $\Phi[\varphi]$ in \mathcal{S} .

(c) The Schwarz minimum values $\Phi[\tilde{\varphi}^{(n)}]$ converge monotonically to the minimum value $\Phi[\varphi]$ in \mathcal{S} starting from any initial guess $\tilde{\varphi}^{(0)}$.

A. Mota, I. Tezaur, C. Alleman. Schwarz Alternating Method in Solid Mechanics

Remark that $\|\mathbf{u}\|_1^2 = \|\mathbf{u}_1\|_1^2 + \|\mathbf{u}_2\|_1^2$, i.e. $\|\mathbf{u}_1\|_1^2 + \|\mathbf{u}_2\|_1^2 \geq \|\mathbf{u}\|_1^2$.

Theorem 1. Assume that the energy functional $\Phi[\varphi]$ satisfies properties 1–5 above. Consider the Schwarz alternating method of Section 2 defined by (9)–(13) and its equivalent form (39). Then

(a) $\Phi[\tilde{\varphi}^{(0)}] \geq \Phi[\tilde{\varphi}^{(1)}] \geq \dots \geq \Phi[\tilde{\varphi}^{(n-1)}] \geq \Phi[\tilde{\varphi}^{(n)}] \geq \dots \geq \Phi[\varphi]$, where φ is the minimizer of $\Phi[\varphi]$ over \mathcal{S} .

(b) The sequence $\{\tilde{\varphi}^{(n)}\}$ defined in (39) converges to the minimizer φ of $\Phi[\varphi]$ in \mathcal{S} .

(c) If $\Phi[\varphi]$ is Lipschitz continuous in a neighborhood of φ , then the sequence $\{\tilde{\varphi}^{(n)}\}$ converges monotonically to the minimizer φ .

Proof. Appendix A.

Finally, while most of proofs above present the analysis for the specific case of this subdomain, the proof of Theorem 1 is valid for the general case.

The case of multiple subdomains is considered in Section 3.

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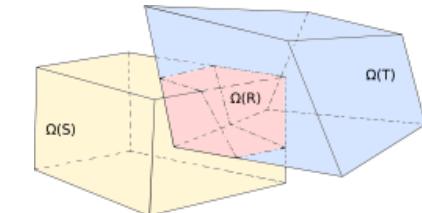
Implementation in *Albany-LCM* and *Sierra/SM* HPC Codes



The overlapping Schwarz alternating method has been implemented in two Sandia HPC codes: *Albany-LCM* and *Sierra/SM*

*Albany-LCM*¹

- *Open-source* parallel, C++, *multi-physics*, finite element code that relies heavily on Trilinos² libraries
- Parallel implementation of Schwarz alternating method uses the *Data Transfer Kit (DTK)*³



Data Transfer Kit (DTK)

Sierra/Solid Mechanics (Sierra/SM)

- Sandia proprietary production *Lagrangian 3D code* for finite element analysis of solids & structures
- Schwarz alternating method was “*implemented*” in *Sierra/SM* using *Arpeggio* loose coupling framework



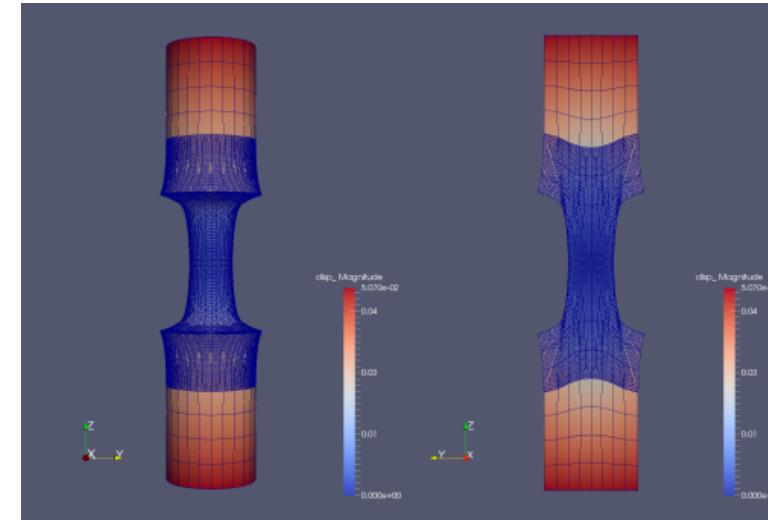
*We did not have to write any code in *Sierra/SM* to implement Schwarz!*

Outline



1. Schwarz Alternating Method for Coupling of Full Order Models (FOMs) in Solid Mechanics

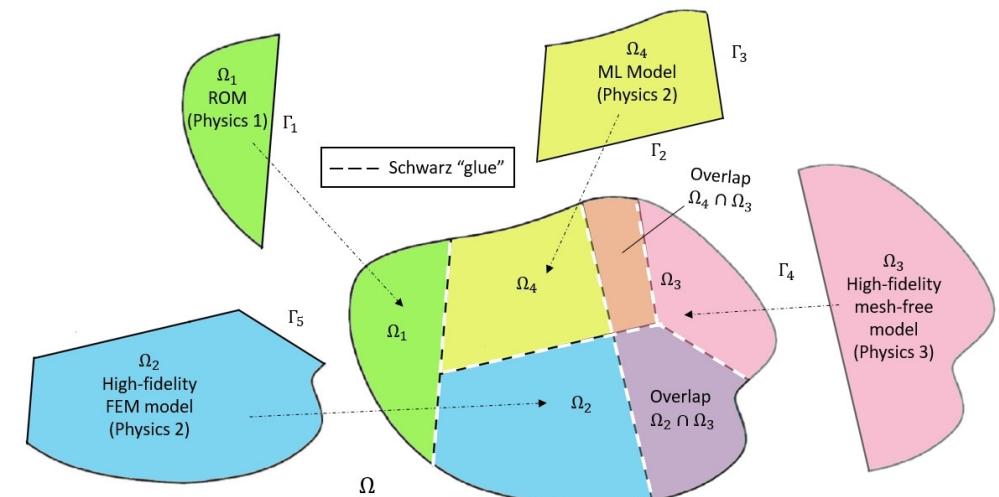
- Motivation & Background
- Quasistatic Formulation
 - Numerical Examples
- Extension to Dynamics
 - Numerical Examples



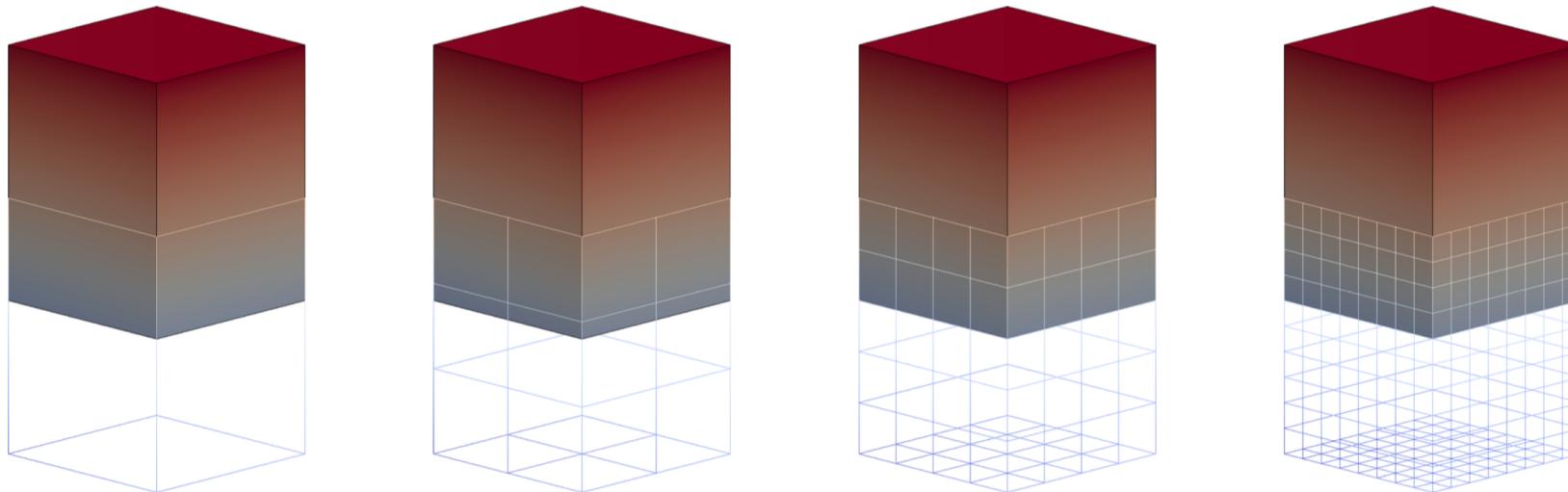
2. Schwarz Alternating Method for FOM-ROM* and ROM-ROM Coupling

- Motivation & Background
- Formulation
- Numerical Examples

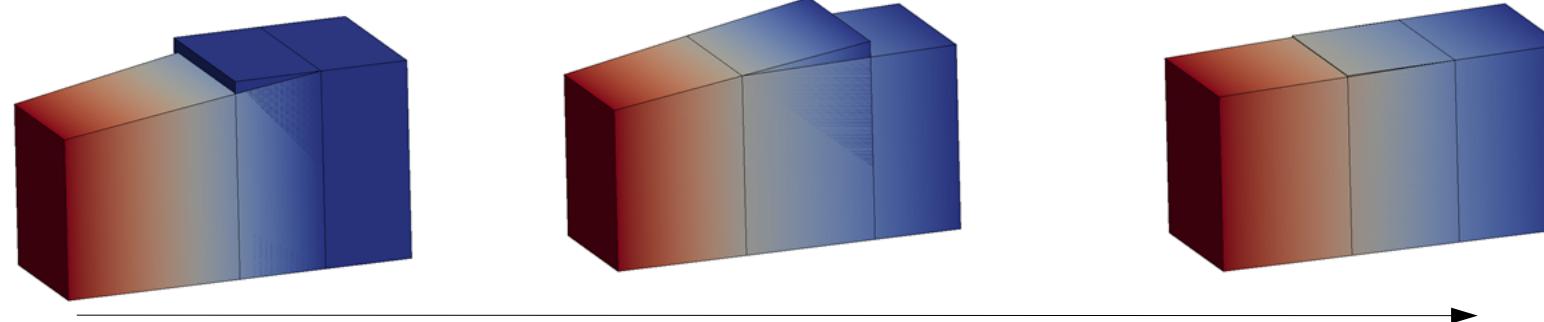
3. Summary and Future Work



Cuboid Problem



- Coupling of *two cuboids* with square base (above).
- *Neoookean*-type material model.



Schwarz Iteration

Cuboid Problem: Convergence and Accuracy

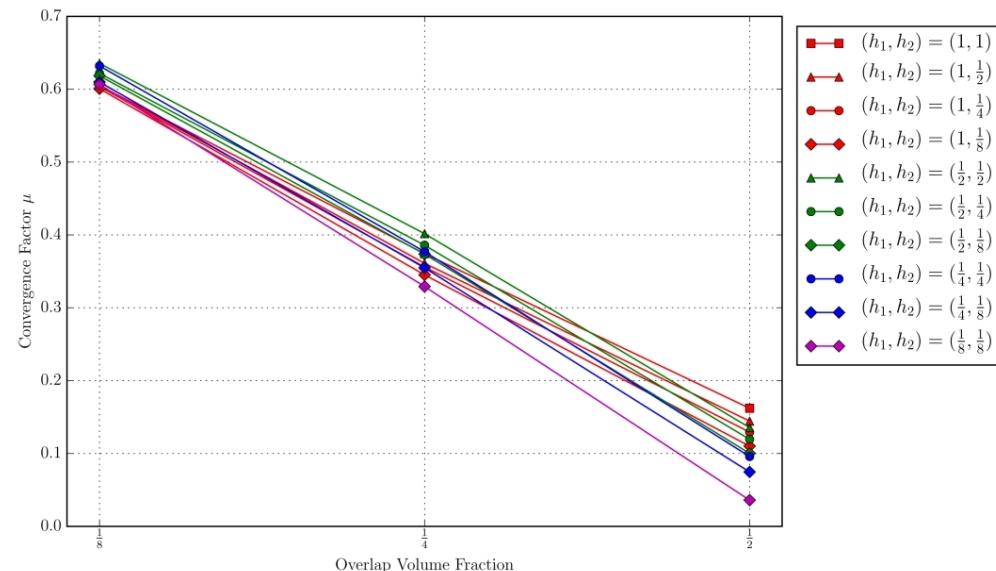
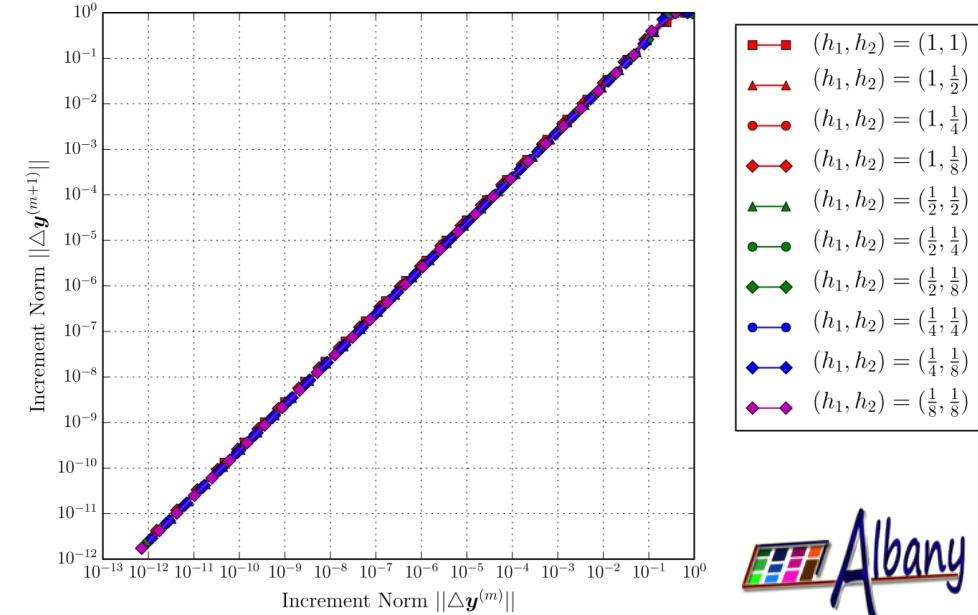


- **Top right:** convergence of the cuboid problem for *different mesh sizes* and *fixed overlap volume fraction*. The Schwarz alternating method converges *linearly*.
- **Bottom right:** convergence factor μ as a function of overlap volume and different mesh. There is *faster linear convergence* with increasing *overlap volume fraction*.

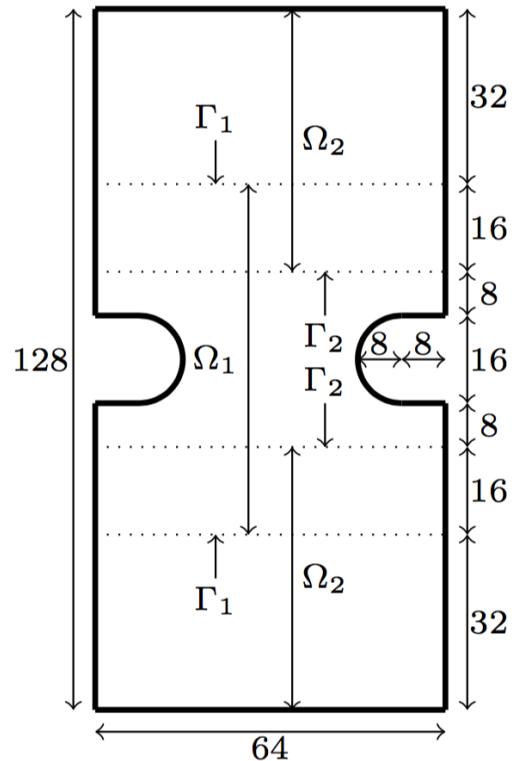
$$\Delta y^{(m+1)} \leq \mu \Delta y^{(m)}$$

- **Below:** *relative errors* in displacement and stress w.r.t. single-domain reference solution. Errors are on the order of *machine precision*.

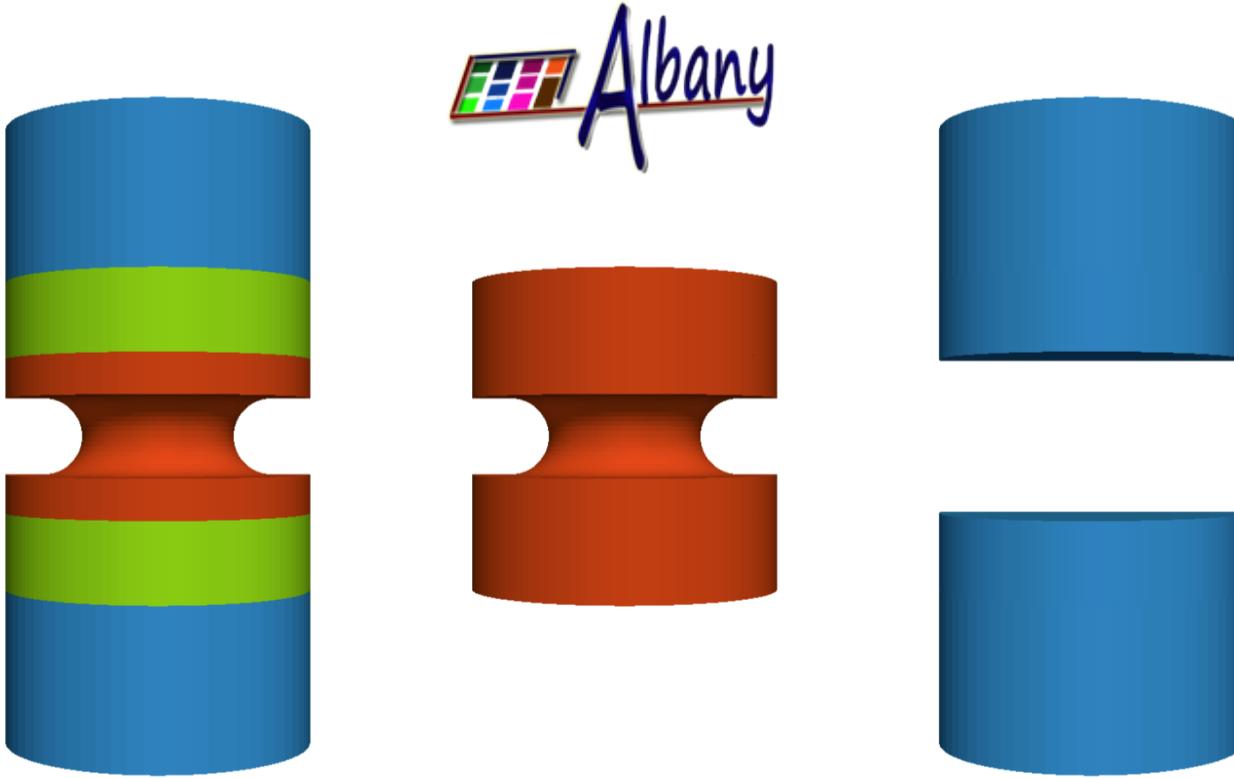
Subdomain	u_3 relative error	σ_{33} relative error
Ω_1	1.24×10^{-14}	2.31×10^{-13}
Ω_2	7.30×10^{-15}	3.06×10^{-13}



Notched Cylinder



(a) Schematic

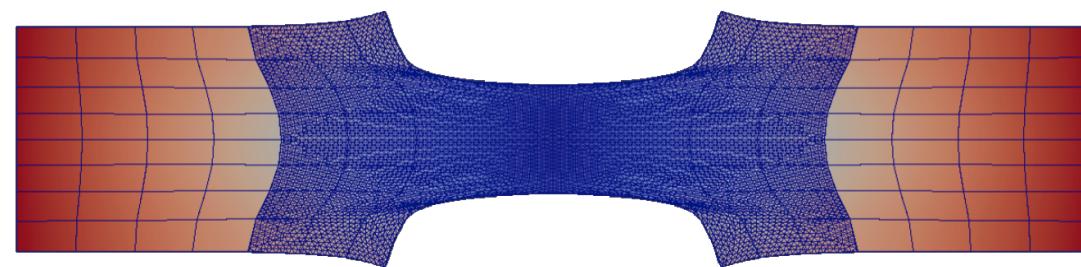
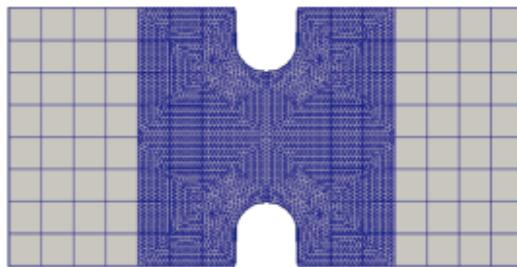
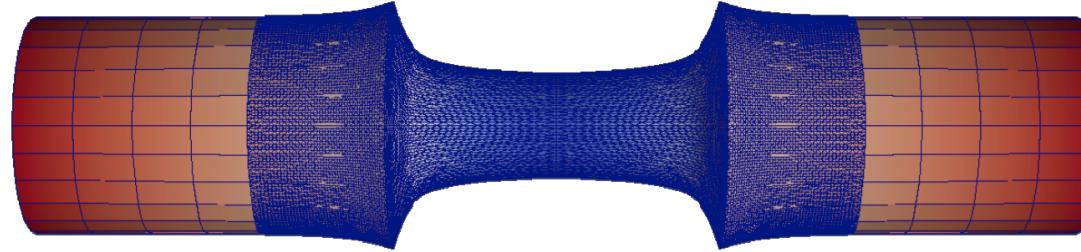
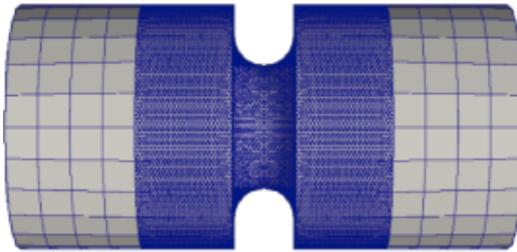
(b) Entire Domain Ω (c) Fine Region Ω_1 (d) Coarse Region Ω_2

- **Notched cylinder** that is stretched along its axial direction.
- Domain decomposed into **two subdomains**.
- **Neo-hookean**-type material model.

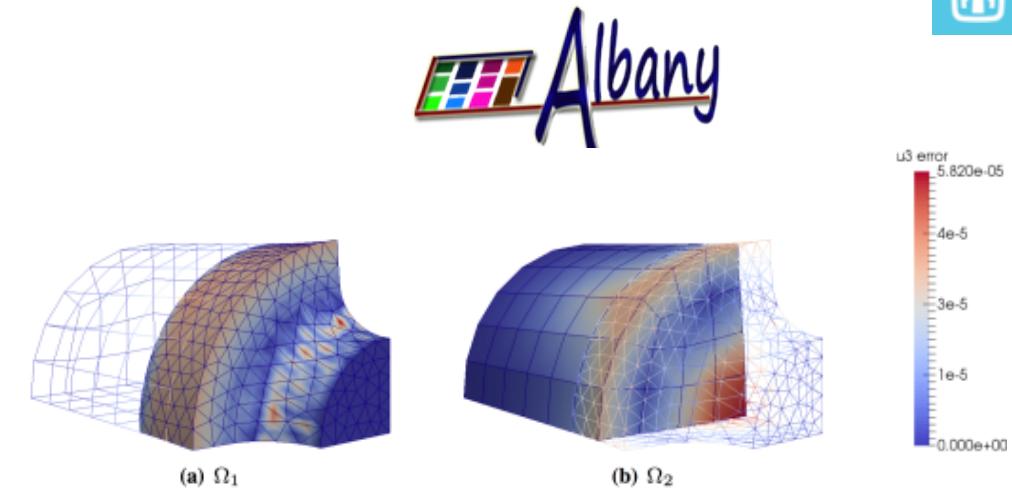
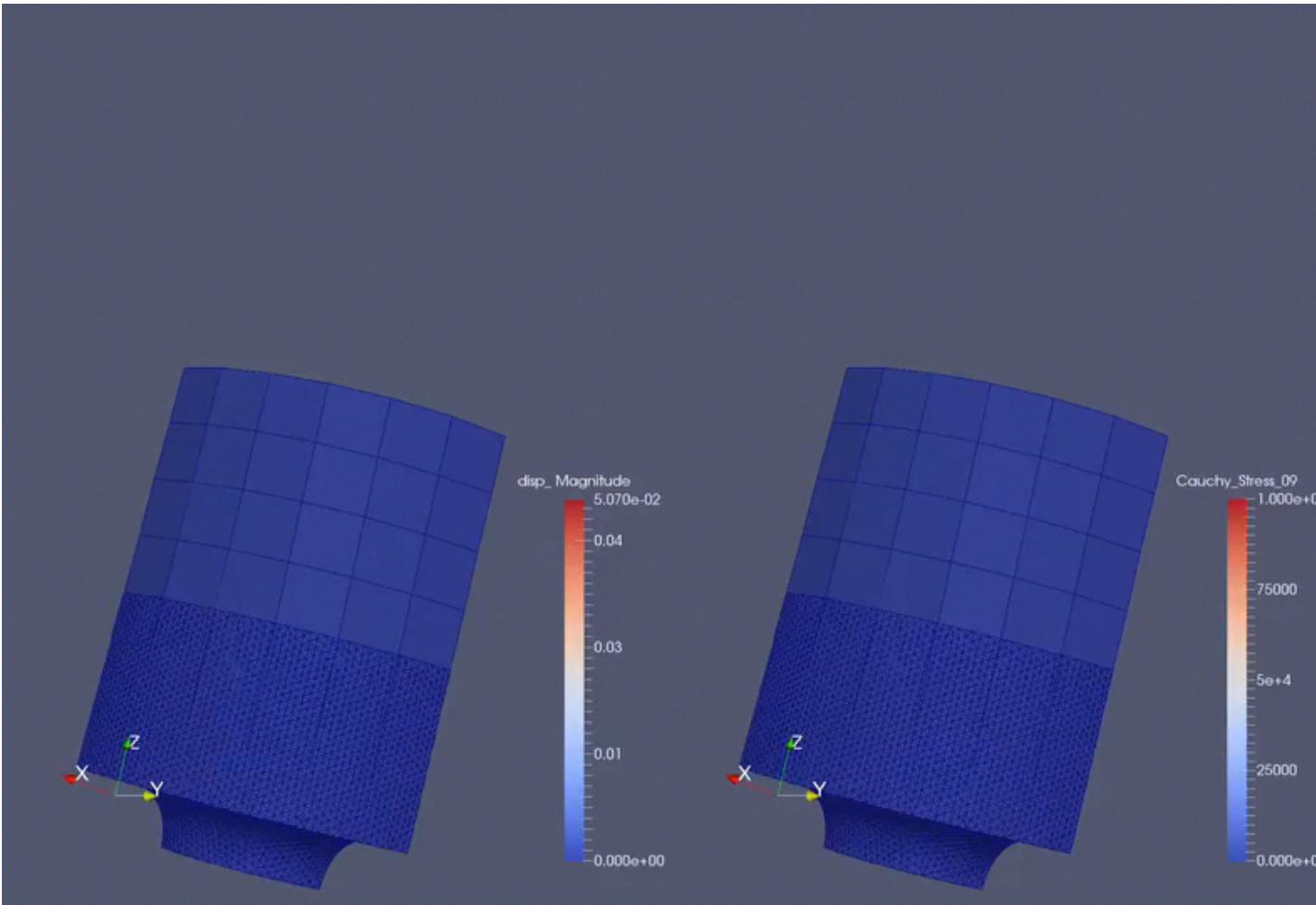
Notched Cylinder: TET - HEX Coupling



- The Schwarz alternating method is capable of coupling ***different mesh topologies***.
- The notched region, where stress concentrations are expected, is ***finely*** meshed with ***tetrahedral*** elements.
- The top and bottom regions, presumably of less interest, are meshed with ***coarser hexahedral*** elements.



Notched Cylinder: TET - HEX Coupling



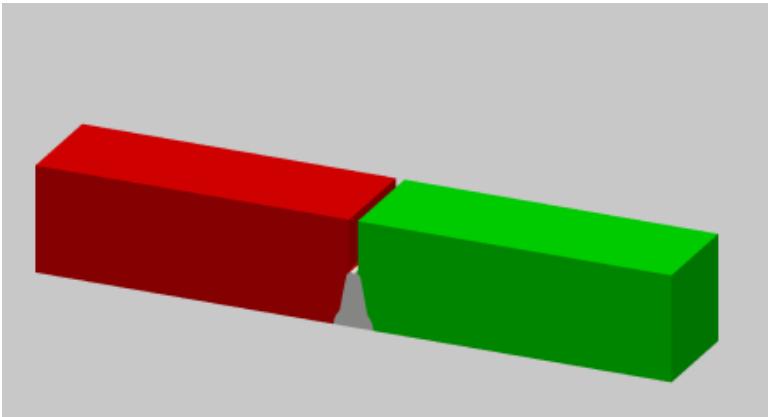
Absolute residual tolerance	u_3 relative error	
	Ω_1	Ω_2
1.0×10^{-14}	9.27×10^{-3}	3.70×10^{-3}

- Relative errors in displacement w.r.t. single-domain reference solution are dominated by **geometric** (rather than coupling) **error**.

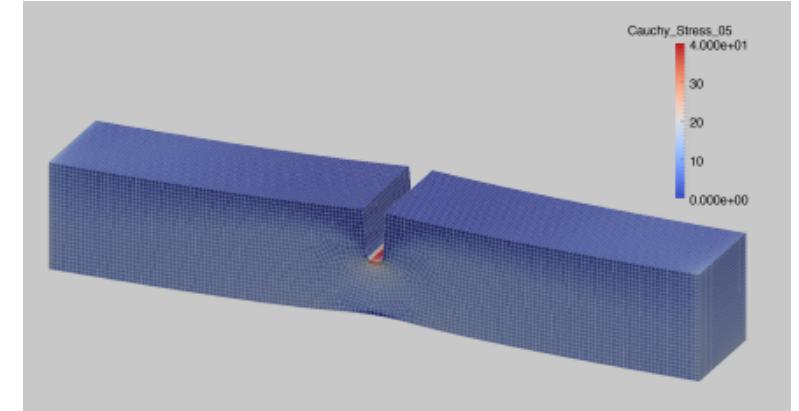
Laser Weld #1



Laser weld specimen

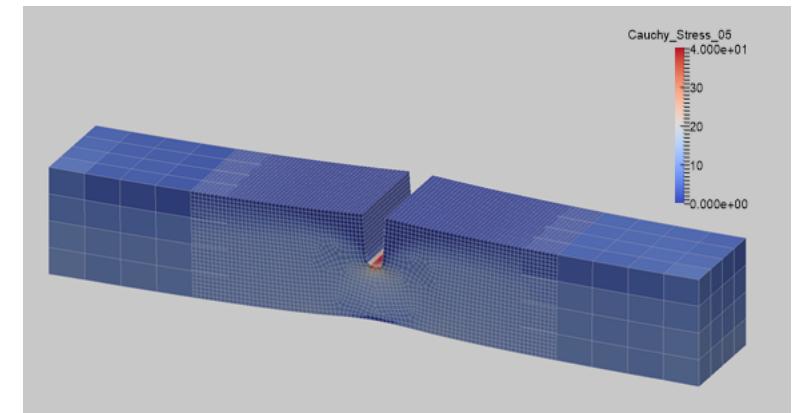


Single domain discretization

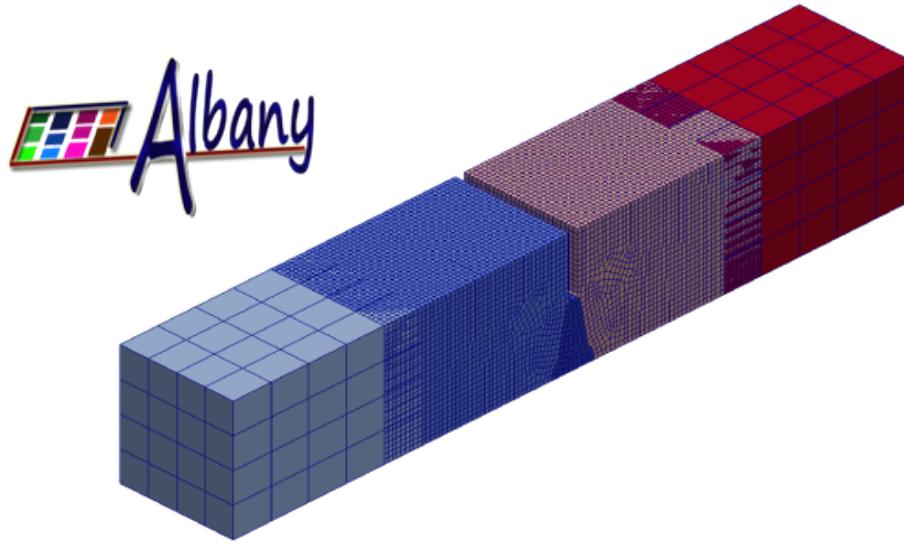


Coupled Schwarz discretization
(50% reduction in model size)

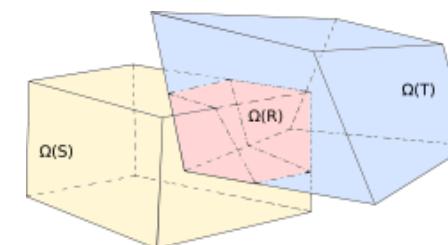
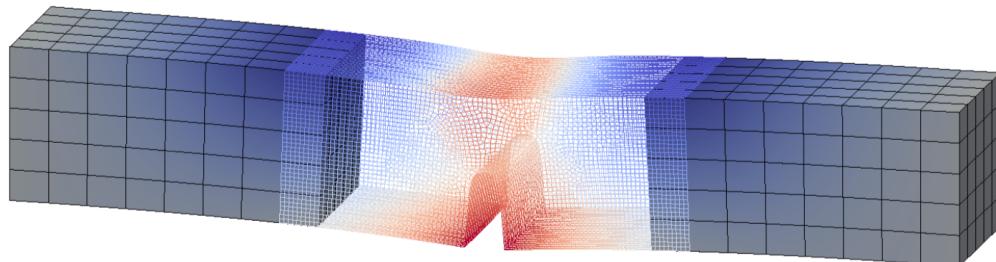
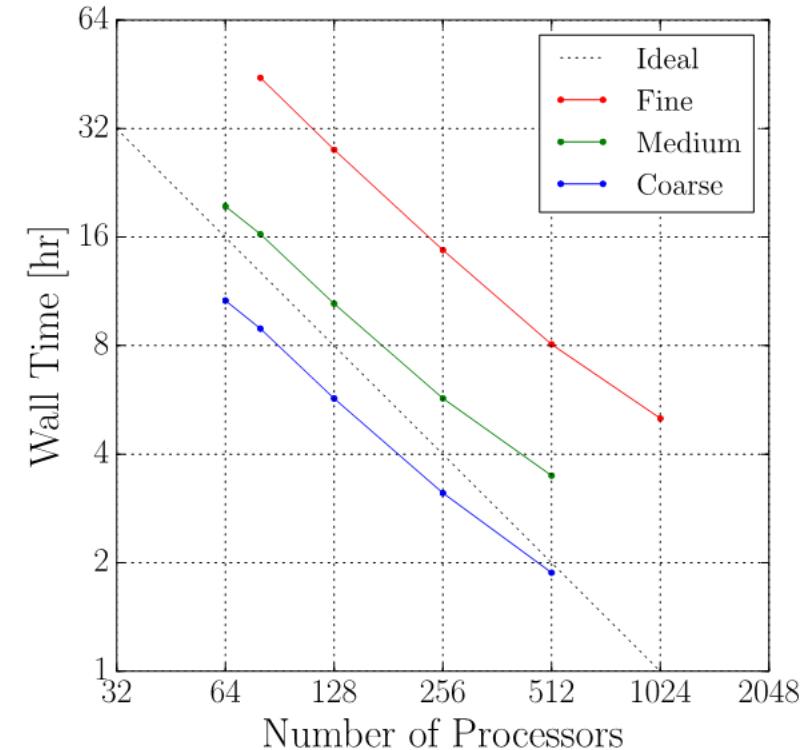
- Problem of *practical scale*.
- *Isotropic elasticity* and *J_2 plasticity* with linear isotropic hardening.
- *Identical parameters* for weld and base materials for proof of concept, to become independent models.



Laser Weld # 1 : Strong Scalability of Parallel Schwarz with DTK

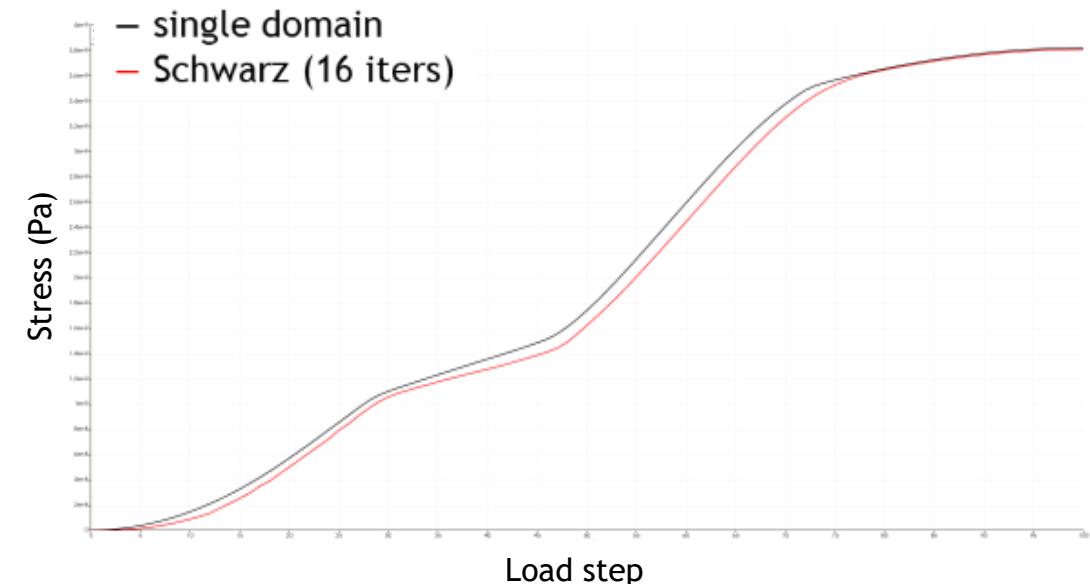
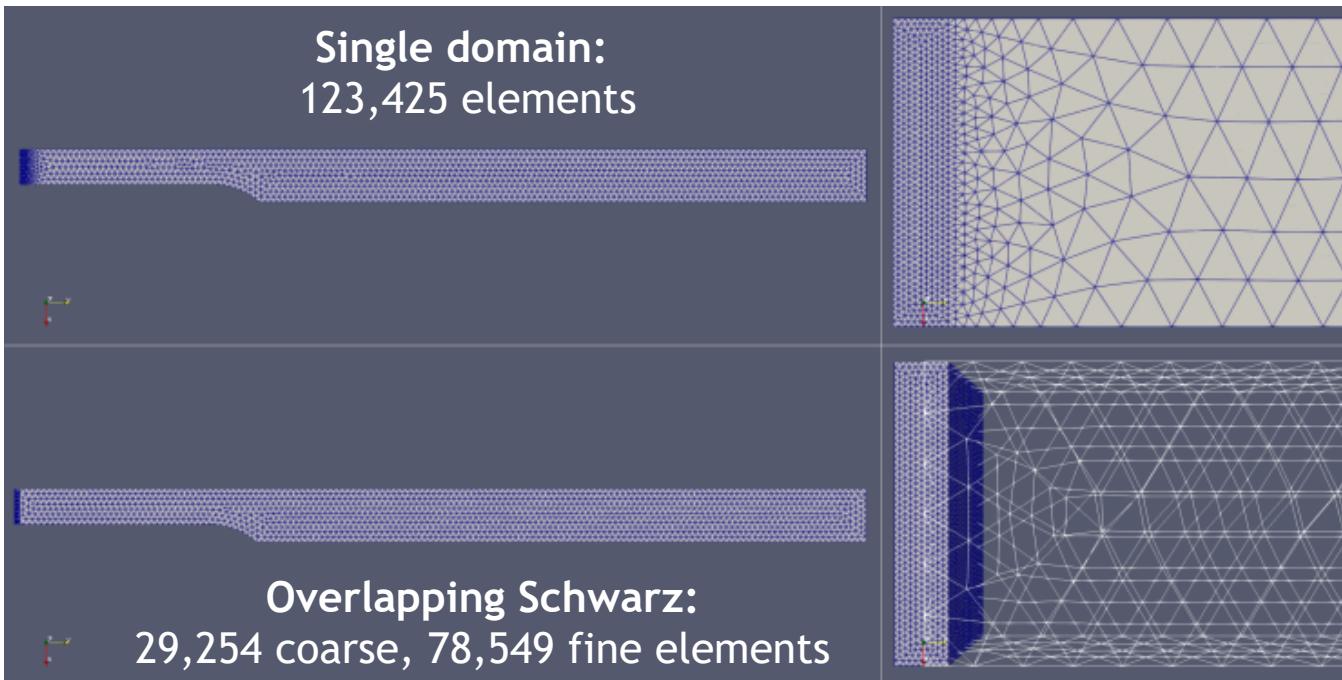


- *Near-ideal linear speedup (64-1024 cores).*



Data Transfer Kit (DTK)

Laser Weld #2: Uniaxial Tension-Test Models



- The domains for Schwarz coupling are **meshed independently**
- This provides the ability to try **different meshing schemes** for each subdomain
- **No need to re-mesh** entire domain
- Schwarz gives **accurate prediction** of stress states if tight enough Schwarz tolerance is used
 - **Tight Schwarz tolerance** needed due to **large disparity** between **element sizes**
- For now, Schwarz is **slower** on this problem, but we are optimizing this

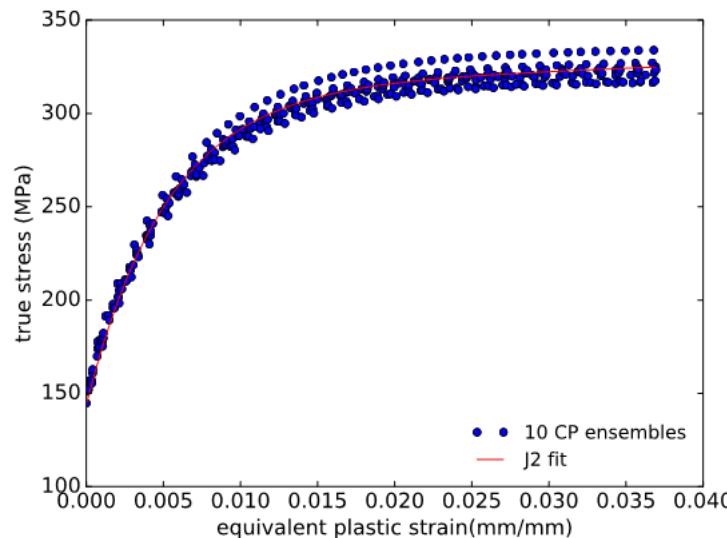


Tensile Bar

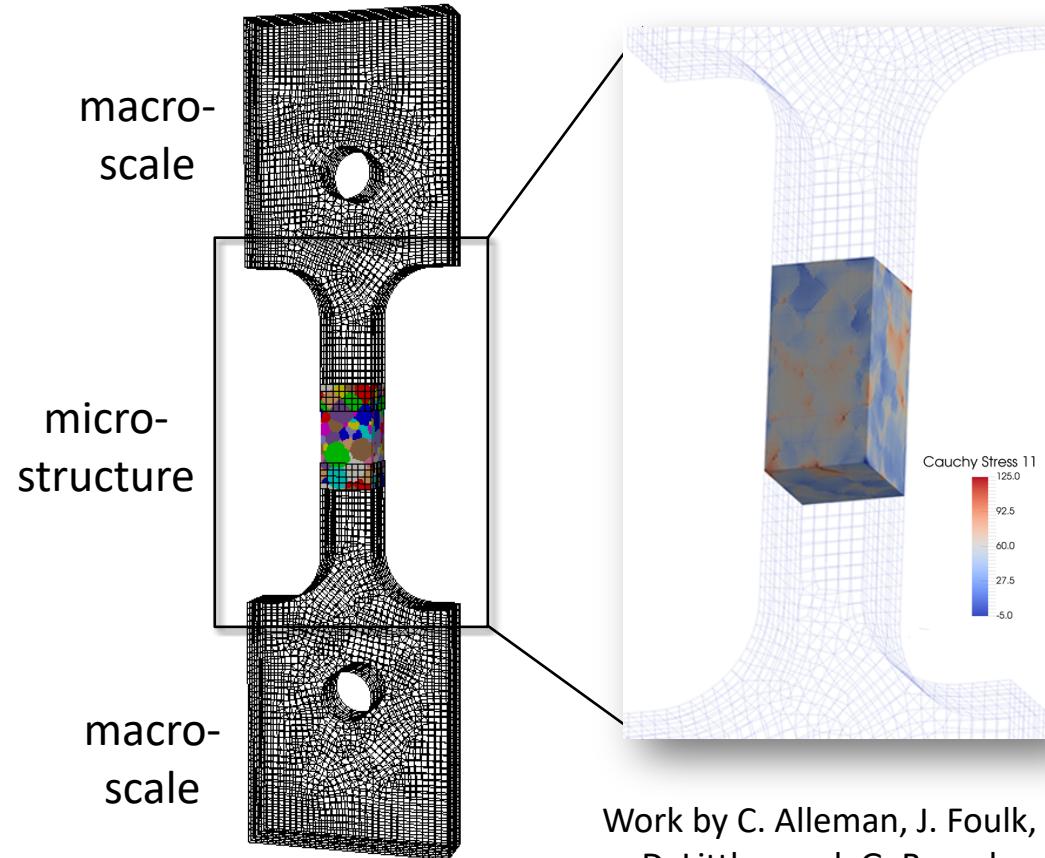


The alternating Schwarz method can be used as part of a *homogenization* (upscaling) process to bridge gap b/w *microscopic* and *macroscopic* regions

- **Microstructure** embedded in ASTM tensile geometry (right).
- Fix microstructure, investigate *ensemble* of uniaxial loads.
- Fit flow curves with a *macroscale* J_2 plasticity model (below).



Goal: study strain localization in microstructure.

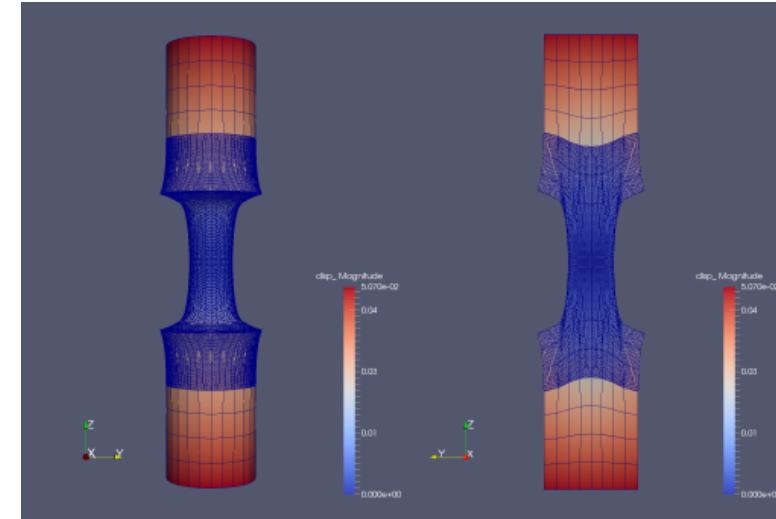


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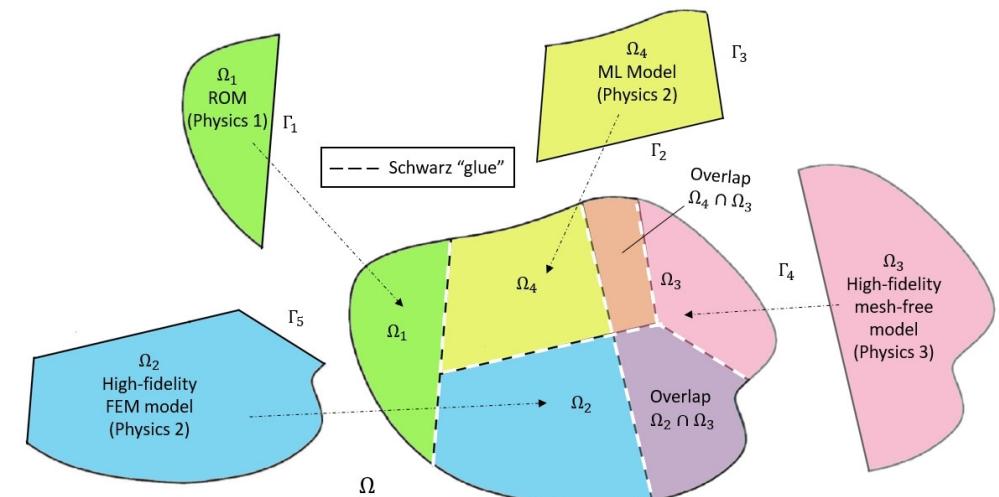
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3. Summary and Future Work



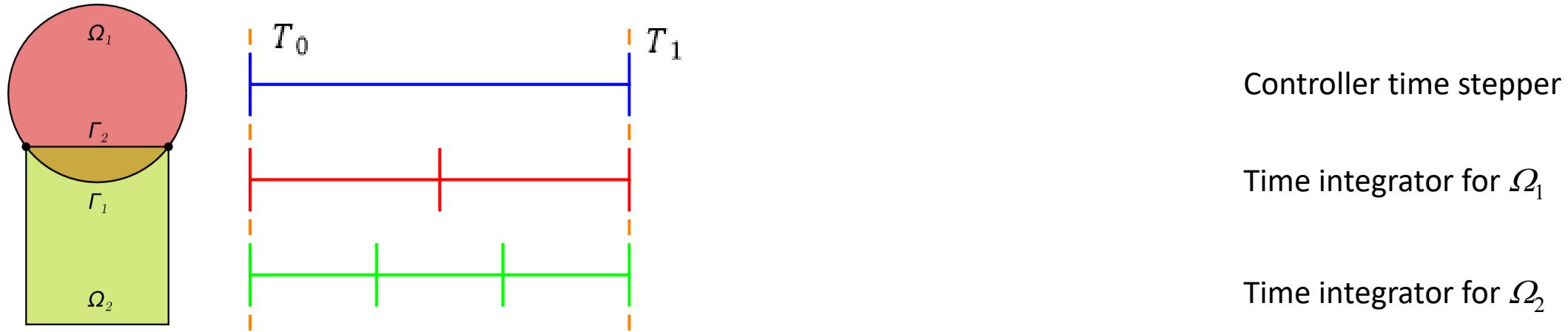
Solid Dynamics Formulation



- Kinetic energy:
$$T(\dot{\boldsymbol{\varphi}}) := \frac{1}{2} \int_{\Omega} \rho \dot{\boldsymbol{\varphi}} \cdot \dot{\boldsymbol{\varphi}} \, dV$$
- Potential energy:
$$V(\boldsymbol{\varphi}) := \int_{\Omega} A(\mathbf{F}, \mathbf{Z}) \, dV - \int_{\Omega} \rho \mathbf{B} \cdot \boldsymbol{\varphi} \, dV$$
- Lagrangian:
$$L(\boldsymbol{\varphi}, \dot{\boldsymbol{\varphi}}) := T(\dot{\boldsymbol{\varphi}}) - V(\boldsymbol{\varphi})$$
- Action functional:
$$S[\boldsymbol{\varphi}] := \int_I L(\boldsymbol{\varphi}, \dot{\boldsymbol{\varphi}}) \, dt$$
- Euler-Lagrange equations:
$$\begin{cases} \text{Div } \mathbf{P} + \rho \mathbf{B} = \rho \ddot{\boldsymbol{\varphi}}, & \text{in } \Omega \times I \\ \boldsymbol{\varphi}(\mathbf{X}, t_0) = \mathbf{x}_0, & \text{in } \Omega \\ \dot{\boldsymbol{\varphi}}(\mathbf{X}, t_0) = \mathbf{v}_0, & \text{in } \Omega \\ \boldsymbol{\varphi}(\mathbf{X}, t) = \boldsymbol{\chi}, & \text{on } \partial\Omega \times I \end{cases}$$
- Semi-discrete problem following FEM discretization in space:

$$\mathbf{M} \ddot{\mathbf{u}} + \mathbf{f}_{\text{int}}(\mathbf{u}, \dot{\mathbf{u}}) = \mathbf{f}_{\text{ext}}$$

Time-Advancement Within the Schwarz Framework

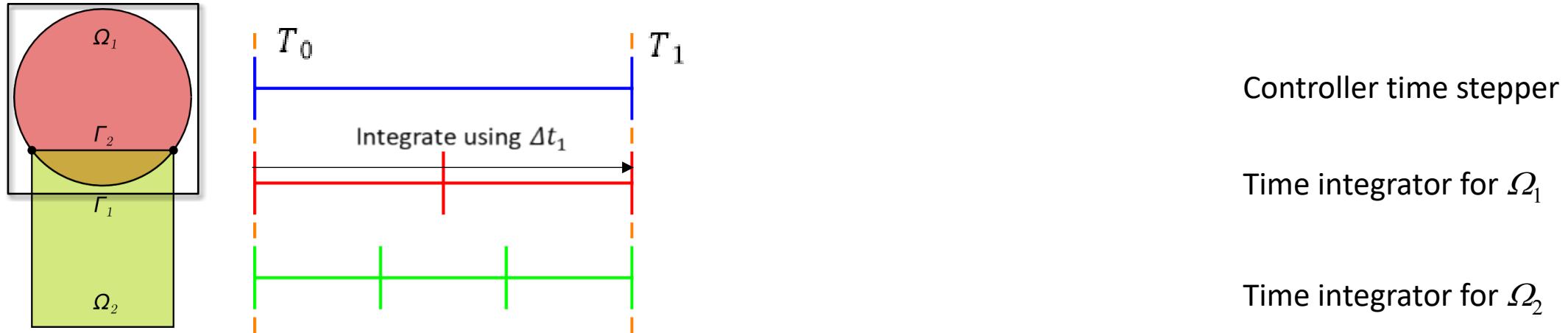


Step 0: Initialize $i = 0$ (controller time index).

Model PDE:

$$\begin{cases} M\ddot{u} + f_{\text{int}}(u, \dot{u}) = f_{\text{ext}} \\ u(x, 0) = u_0 \end{cases}$$

Time-Advancement Within the Schwarz Framework

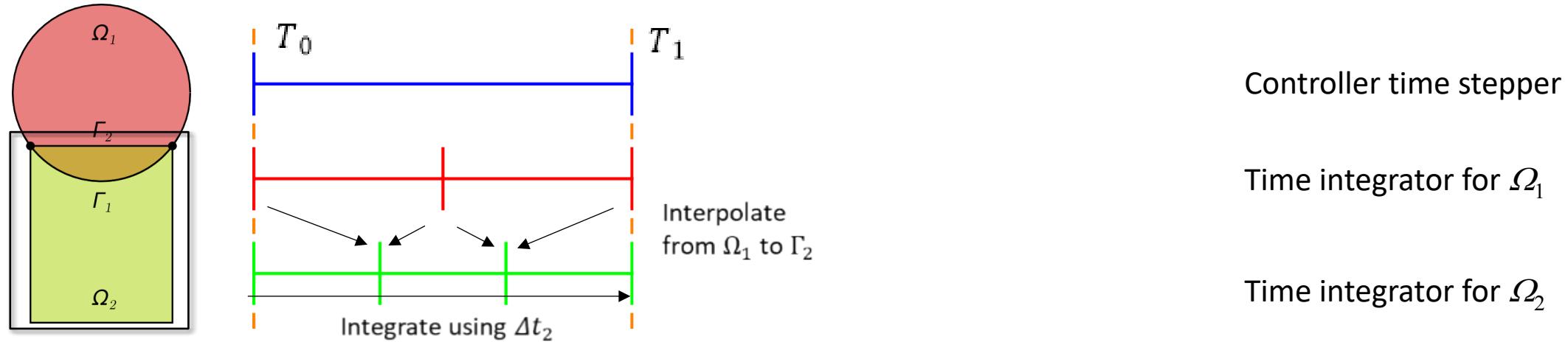


Step 0: Initialize $i = 0$ (controller time index).

Step 1: Advance Ω_1 solution from time T_i to time T_{i+1} using time-stepper in Ω_1 with time-step Δt_1 , using solution in Ω_2 interpolated to Γ_1 at times $T_i + n\Delta t_1$.

Model PDE:
$$\begin{cases} M\ddot{u} + f_{\text{int}}(u, \dot{u}) = f_{\text{ext}} \\ u(x, 0) = u_0 \end{cases}$$

Time-Advancement Within the Schwarz Framework

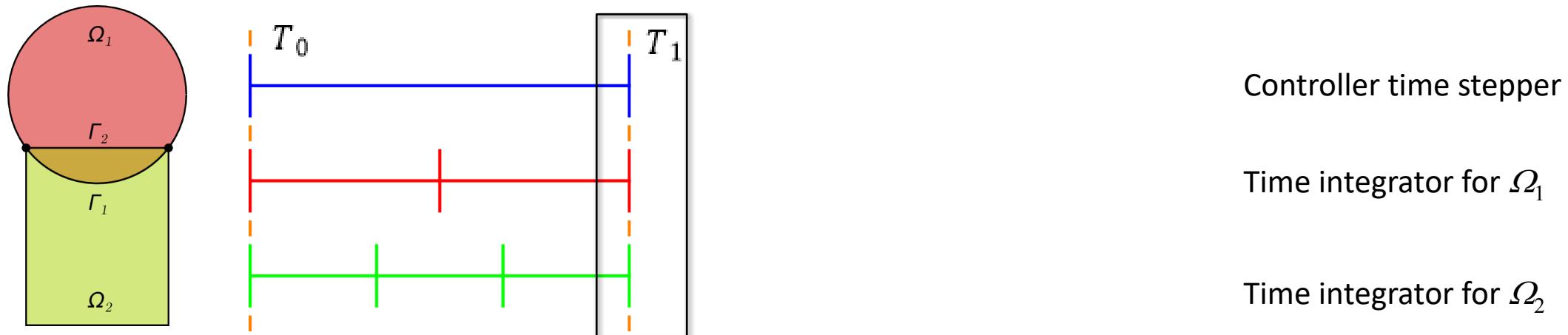


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Step 2: Advance Ω_2 solution from time T_i to time T_{i+1} using time-stepper in Ω_2 with time-step Δt_2 , using solution in Ω_1 interpolated to Γ_2 at times $T_i + n\Delta t_2$.

Model PDE:
$$\begin{cases} M\ddot{u} + f_{\text{int}}(u, \dot{u}) = f_{\text{ext}} \\ u(x, 0) = u_0 \end{cases}$$



Step 0: Initialize $i = 0$ (controller time index).

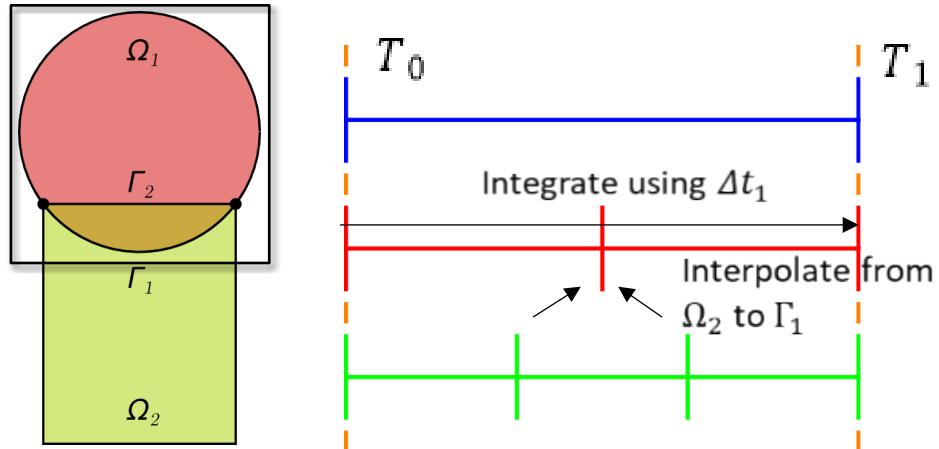
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Step 2: Advance Ω_2 solution from time T_i to time T_{i+1} using time-stepper in Ω_2 with time-step Δt_2 , using solution in Ω_1 interpolated to Γ_2 at times $T_i + n\Delta t_2$.

Step 3: Check for convergence at time T_{i+1} .

Model PDE:
$$\begin{cases} M\ddot{u} + f_{\text{int}}(u, \dot{u}) = f_{\text{ext}} \\ u(x, 0) = u_0 \end{cases}$$

Time-Advancement Within the Schwarz Framework



Controller time stepper

Time integrator for Ω_1

Time integrator for Ω_2

Step 0: Initialize $i = 0$ (controller time index).

Step 1: Advance Ω_1 solution from time T_i to time T_{i+1} using time-stepper in Ω_1 with time-step Δt_1 , using solution in Ω_2 interpolated to Γ_1 at times $T_i + n\Delta t_1$.

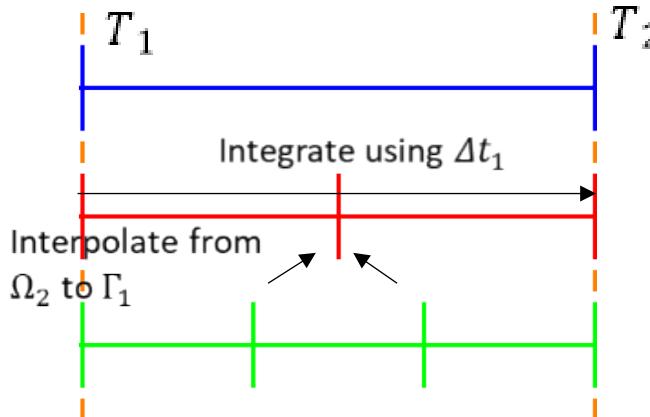
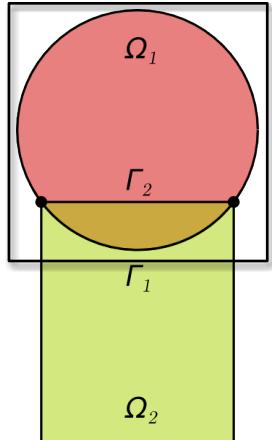
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Step 3: Check for convergence at time T_{i+1} .

➤ If unconverged, return to Step 1.

Model PDE:
$$\begin{cases} M\ddot{u} + f_{\text{int}}(u, \dot{u}) = f_{\text{ext}} \\ u(x, 0) = u_0 \end{cases}$$

Time-Advancement Within the Schwarz Framework



Controller time stepper

Time integrator for Ω_1

Time integrator for Ω_2

Can use *different integrators* with *different time steps* within each domain!

Step 0: Initialize $i = 0$ (controller time index).

Step 1: Advance Ω_1 solution from time T_i to time T_{i+1} using time-stepper in Ω_1 with time-step Δt_1 , using solution in Ω_2 interpolated to Γ_1 at times $T_i + n\Delta t_1$.

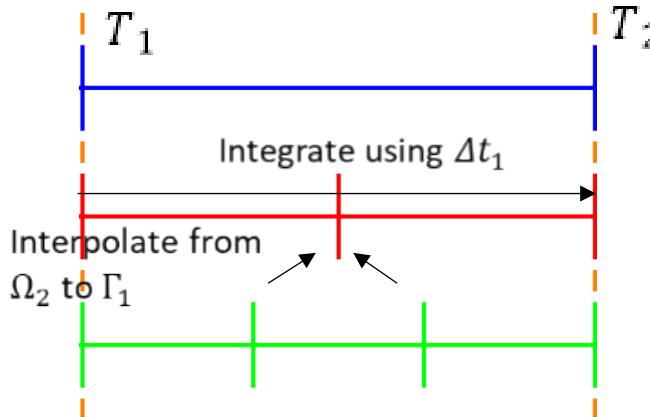
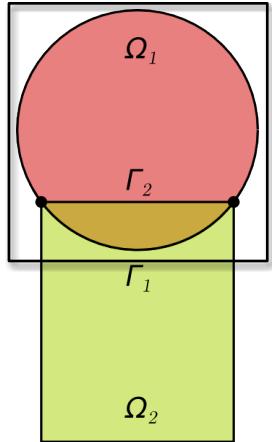
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Model PDE:
$$\begin{cases} M\ddot{u} + f_{\text{int}}(u, \dot{u}) = f_{\text{ext}} \\ u(x, 0) = u_0 \end{cases}$$

Time-Advancement Within the Schwarz Framework



Controller time stepper

Time integrator for Ω_1

Time integrator for Ω_2

Time-stepping procedure is equivalent to doing Schwarz on space-time domain [Mota *et al.* 2022].

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Model PDE:
$$\begin{cases} M\ddot{u} + f_{\text{int}}(u, \dot{u}) = f_{\text{ext}} \\ u(x, 0) = u_0 \end{cases}$$



- Like for quasistatics, dynamic alternating Schwarz method converges provided each single-domain problem is **well-posed** and **overlap region is non-empty**, under some **conditions** on Δt .
- **Well-posedness** for the dynamic problem requires that action functional $S[\boldsymbol{\varphi}] := \int_I \int_{\Omega} L(\boldsymbol{\varphi}, \dot{\boldsymbol{\varphi}}) dV dt$ be **strictly convex** or **strictly concave**, where $L(\boldsymbol{\varphi}, \dot{\boldsymbol{\varphi}}) := T(\dot{\boldsymbol{\varphi}}) + V(\boldsymbol{\varphi})$ is the Lagrangian.
 - This is studied by looking at its second variation $\delta^2 S[\boldsymbol{\varphi}_h]$
- We can show assuming a **Newmark** time-integration scheme that for the **fully-discrete** problem:

$$\delta^2 S[\boldsymbol{\varphi}_h] = \mathbf{x}^T \left[\frac{\gamma^2}{(\beta \Delta t)^2} \mathbf{M} - \mathbf{K} \right] \mathbf{x}$$

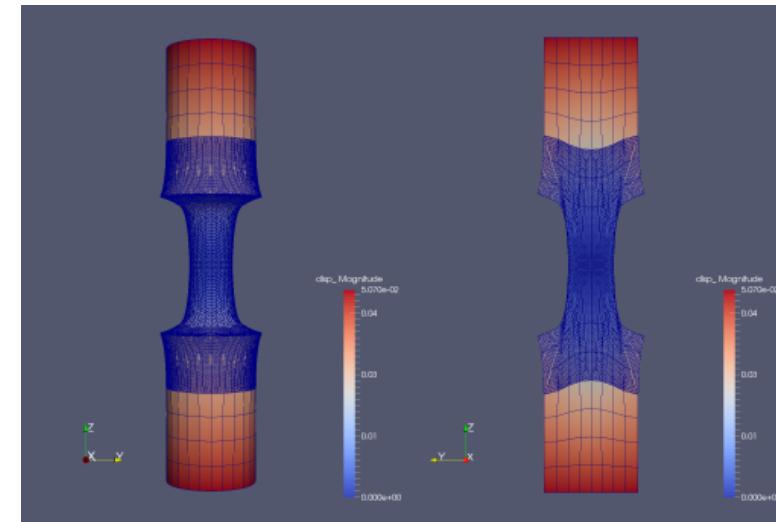
- $\delta^2 S[\boldsymbol{\varphi}_h]$ can always be made positive by choosing a **sufficiently small** Δt
- Numerical experiments reveal that Δt requirements for **stability/accuracy** typically lead to automatic satisfaction of this bound.

Outline



1. Schwarz Alternating Method for Coupling of Full Order Models (FOMs) in Solid Mechanics

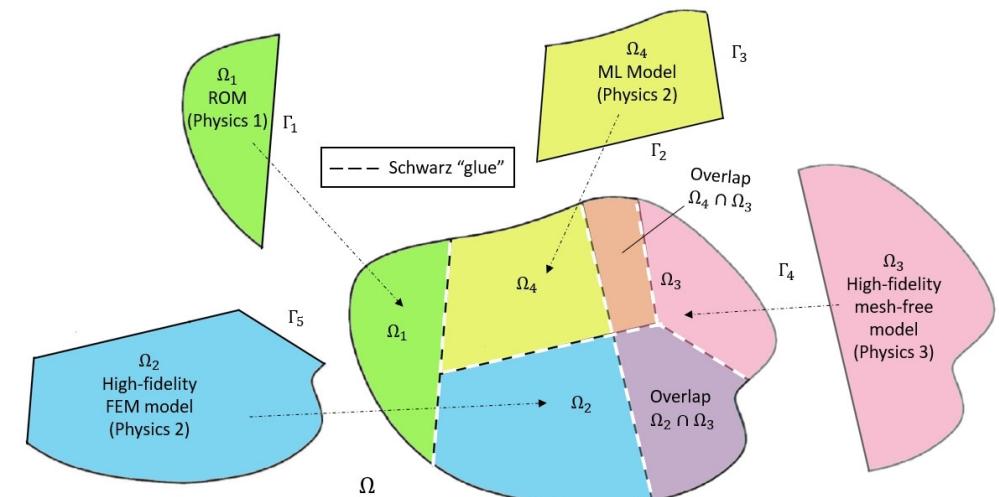
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- Quasistatic Formulation
 - Numerical Examples
- Extension to Dynamics
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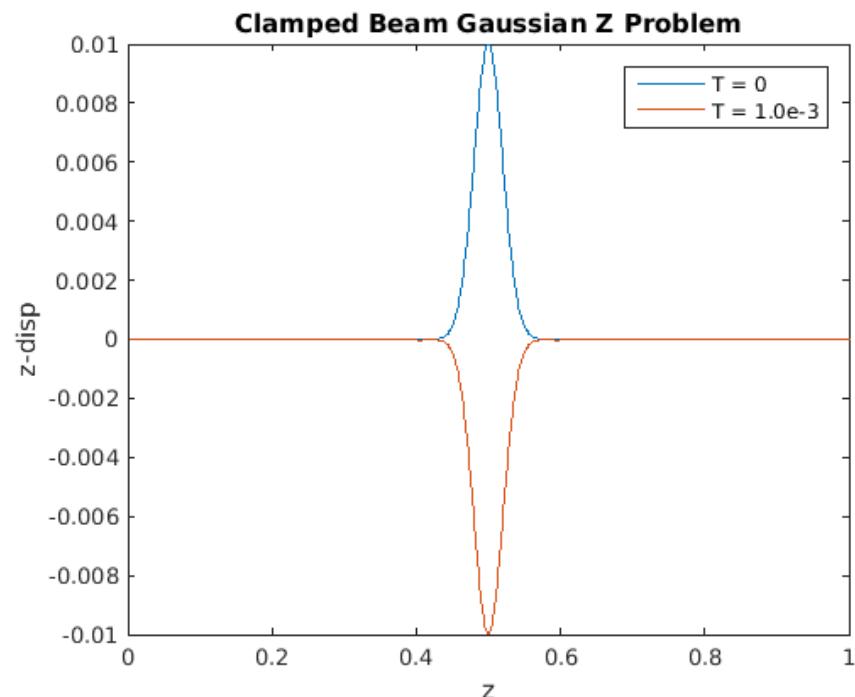
2. Schwarz Alternating Method for FOM-ROM* and ROM-ROM Coupling

- Motivation & Background
- Formulation
- Numerical Examples

3. Summary and Future Work



- Linear elastic **clamped beam** with Gaussian initial condition for the z -displacement.
- Simple problem with analytical exact solution but very **stringent test** for discretization methods.
- Test Schwarz with **2 subdomains**: $\Omega_0 = (0,0.001) \times (0,0.001) \times (0,0.75)$, $\Omega_1 = (0,0.001) \times (0,0.001) \times (0.25,1)$.



Left: Initial condition (blue) and final solution (red). Wave profile is negative of initial profile at time $T = 1.0e-3$.

Time-discretizations:
Newmark (implicit, explicit).
Meshes: HEX, TET

Elastic Wave: Different Integrators, Same Δt s

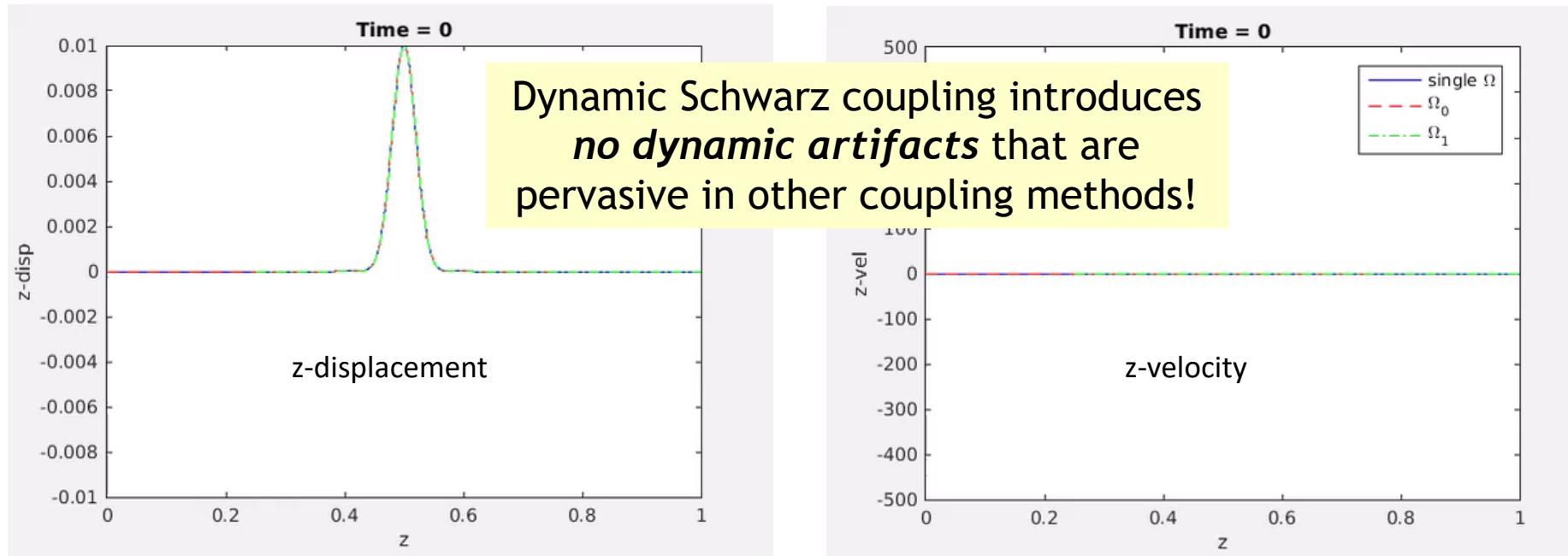
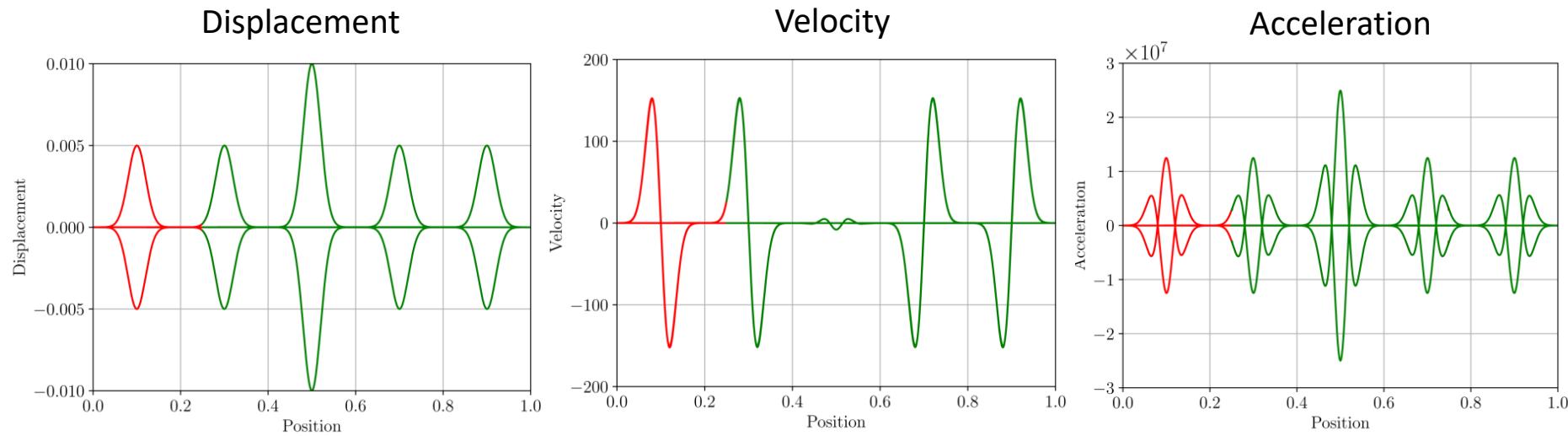


Table 1: Averaged (over times + domains) relative errors in **z-displacement (blue)** and **z-velocity (green)** for several different Schwarz couplings, 50% overlap volume fraction

	Implicit-Implicit		Explicit(CM)-Implicit		Explicit(LM)-Implicit	
Conformal HEX - HEX	2.79e-3	7.32e-3	3.53e-3	8.70e-3	4.72e-3	1.19e-2
Nonconformal HEX - HEX	2.90e-3	7.10e-3	2.82e-3	7.29e-3	2.84e-3	7.33e-3
TET - HEX	2.79e-3	7.58e-3	3.52e-3	8.92e-3	4.72e-3	1.19e-2

Elastic Wave: Different Integrators, Different Δt s



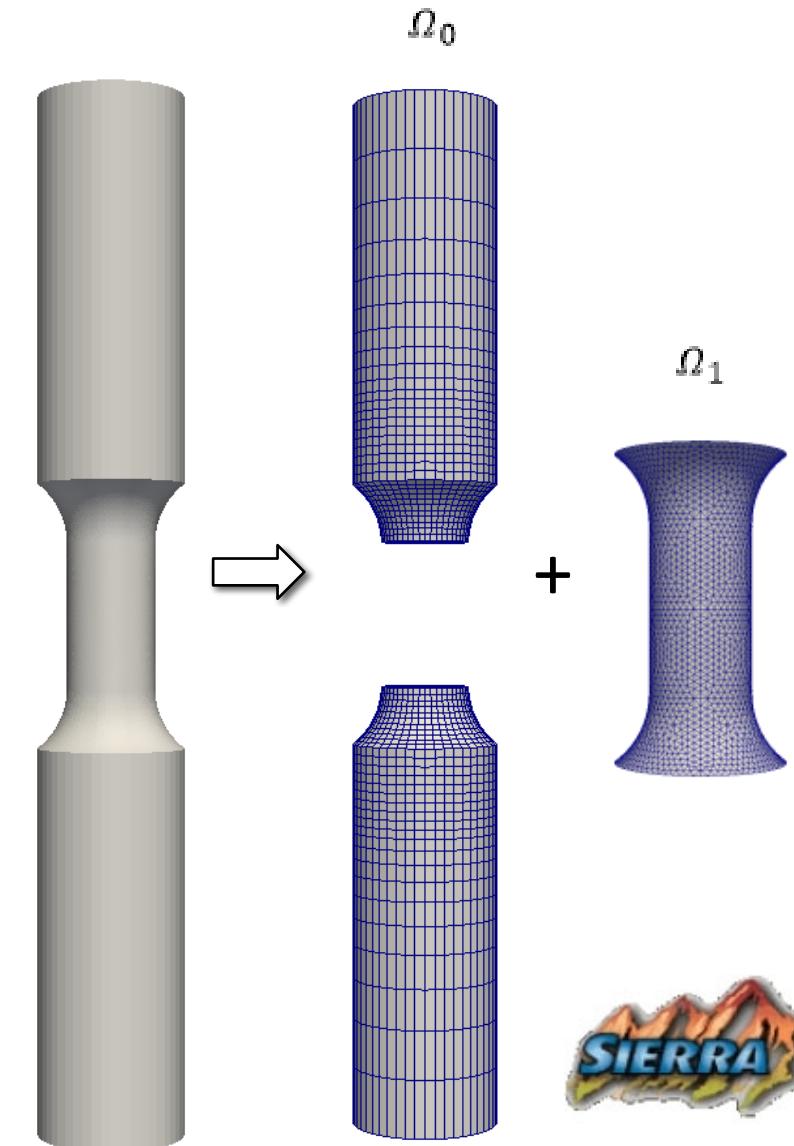
Figures above: Plots of displacement, velocity and acceleration for the elastic wave propagation problem using different time integrators (implicit and explicit) and different time steps ($1e-2s$ and $2e-7s$) for each subdomain, superimposed over the analytic single domain solution.

The analytic solution is *indistinguishable* from Schwarz solutions (hidden behind the solutions for Ω_0 (red) and Ω_1 (green))!

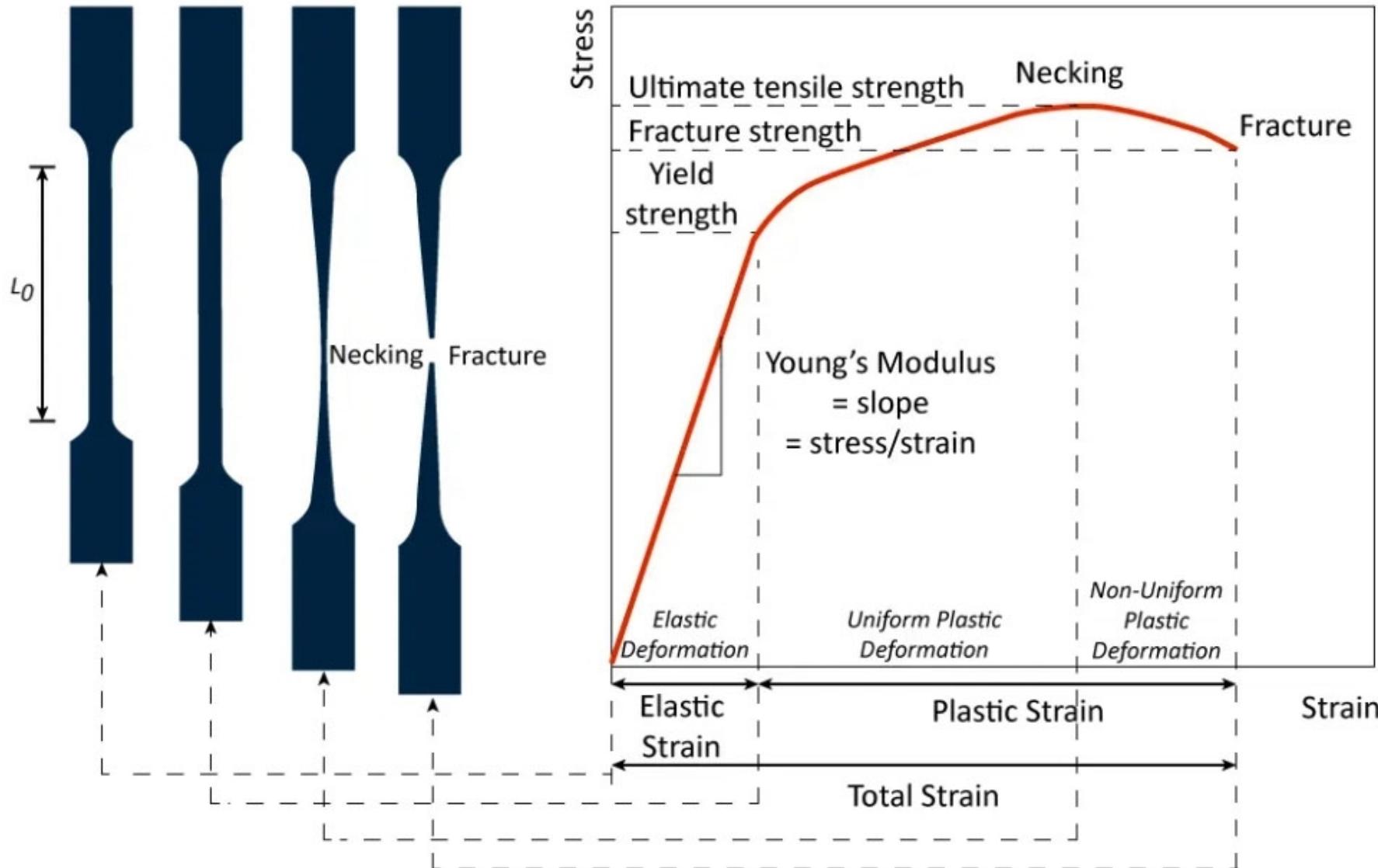
Tension Specimen



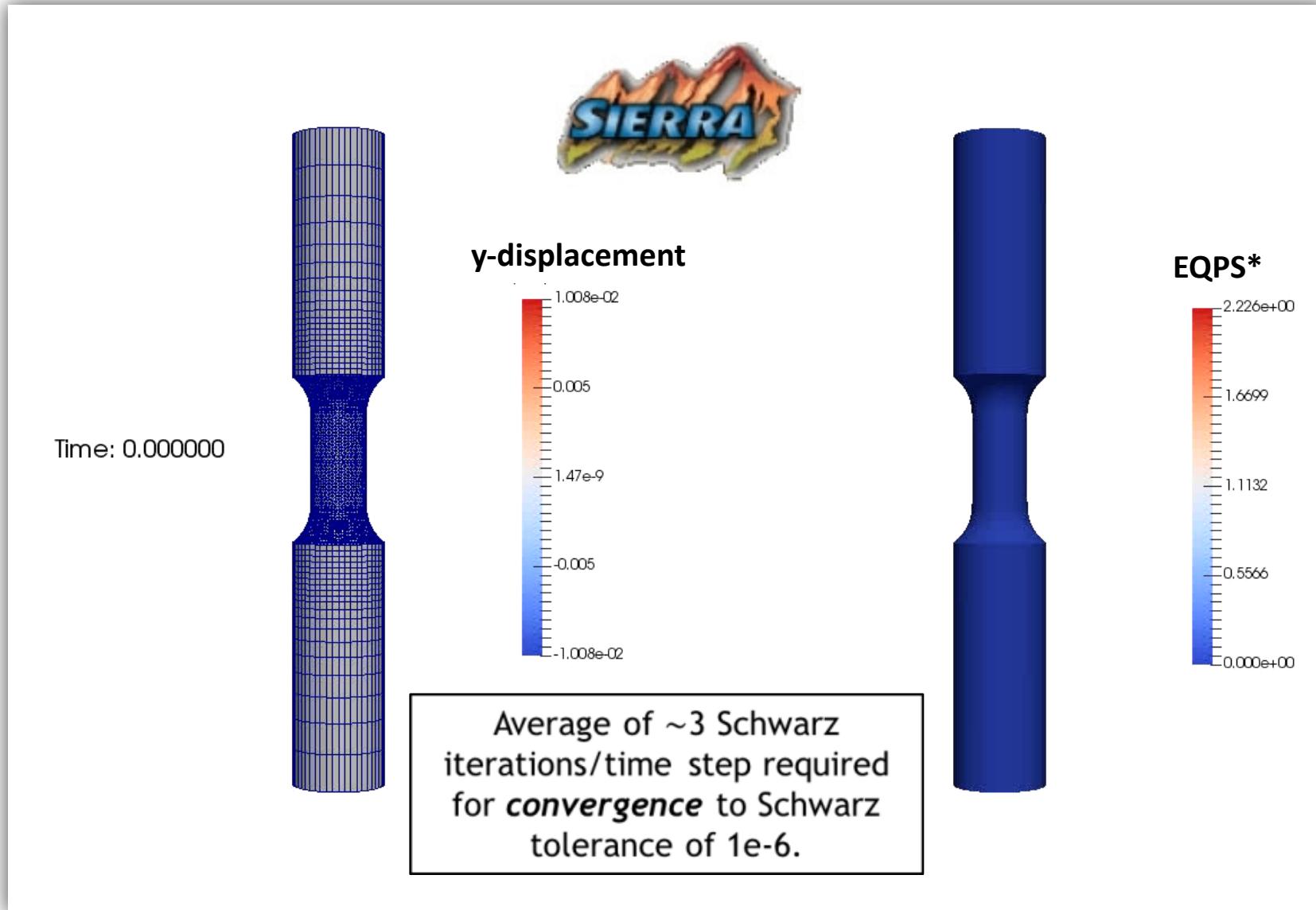
- Uniaxial aluminum cylindrical tensile specimen with *inelastic J₂ material model*.
- Domain decomposition into *two subdomains* (right): Ω_0 = ends, Ω_1 = gauge.
- *Nonconformal HEX + composite TET10* coupling via Schwarz.
- *Implicit* Newmark time-integration with *adaptive time-stepping* algorithm employed in both subdomains.
- Slight *imperfection* introduced at center of gauge to force *necking* upon pulling in vertical direction.



Tension Specimen: Expected Result



Tension Specimen: Displacement & EQPS*

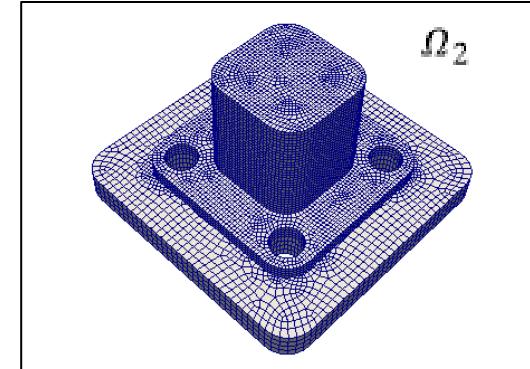
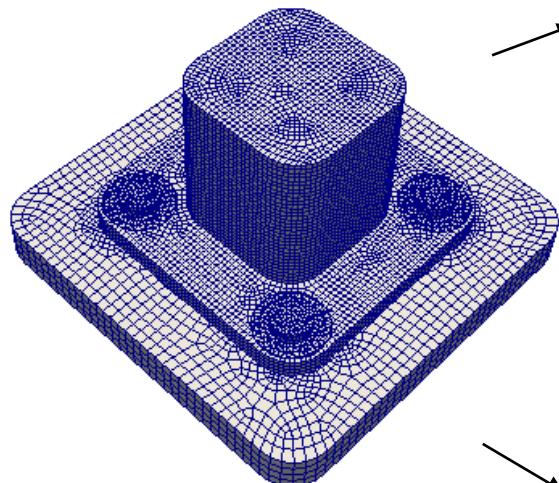


Bolted Joint Problem



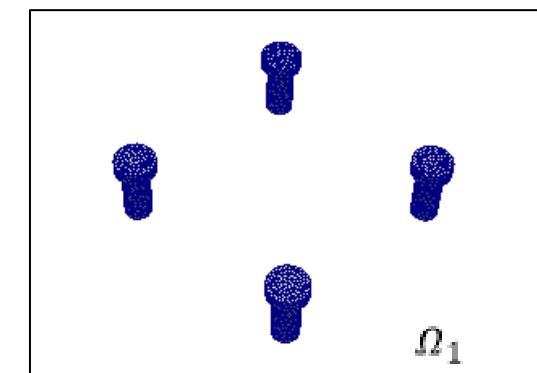
Problem of *practical scale*.

- Schwarz solution compared to single-domain solution on composite TET10 mesh.



- BC: $x\text{-disp} = 0.02$ at $T = 1.0\text{e-}3$ on top of parts.
- Run until $T = 5.0\text{e-}4$ w/ $dt = 1\text{e-}5$ + implicit Newmark with analytic mass matrix for composite tet 10s.

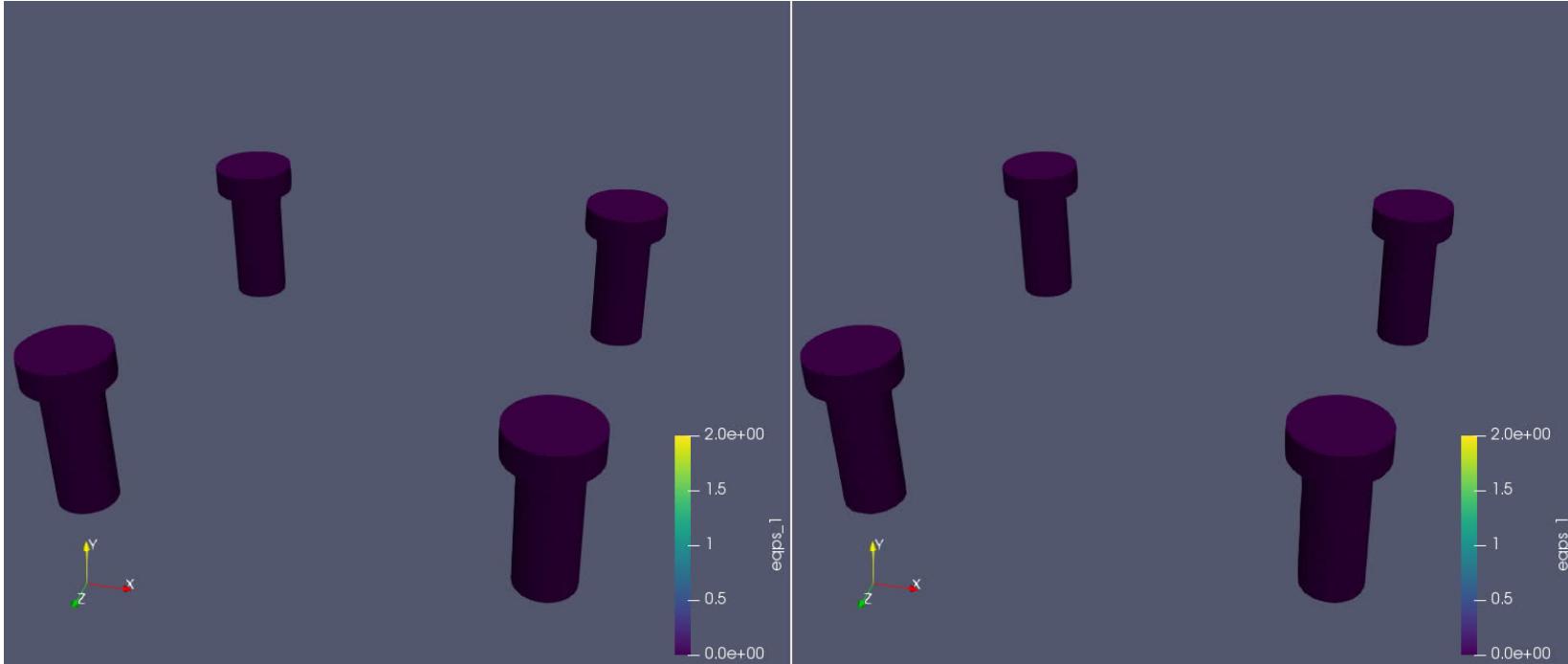
- Ω_1 = bolts (Composite TET10), Ω_2 = parts (HEX).
- *Inelastic J_2 material model* in both subdomains.
 - Ω_1 : steel
 - Ω_2 : steel component, aluminum (bottom) plate



Bolted Joint Problem: Displacement



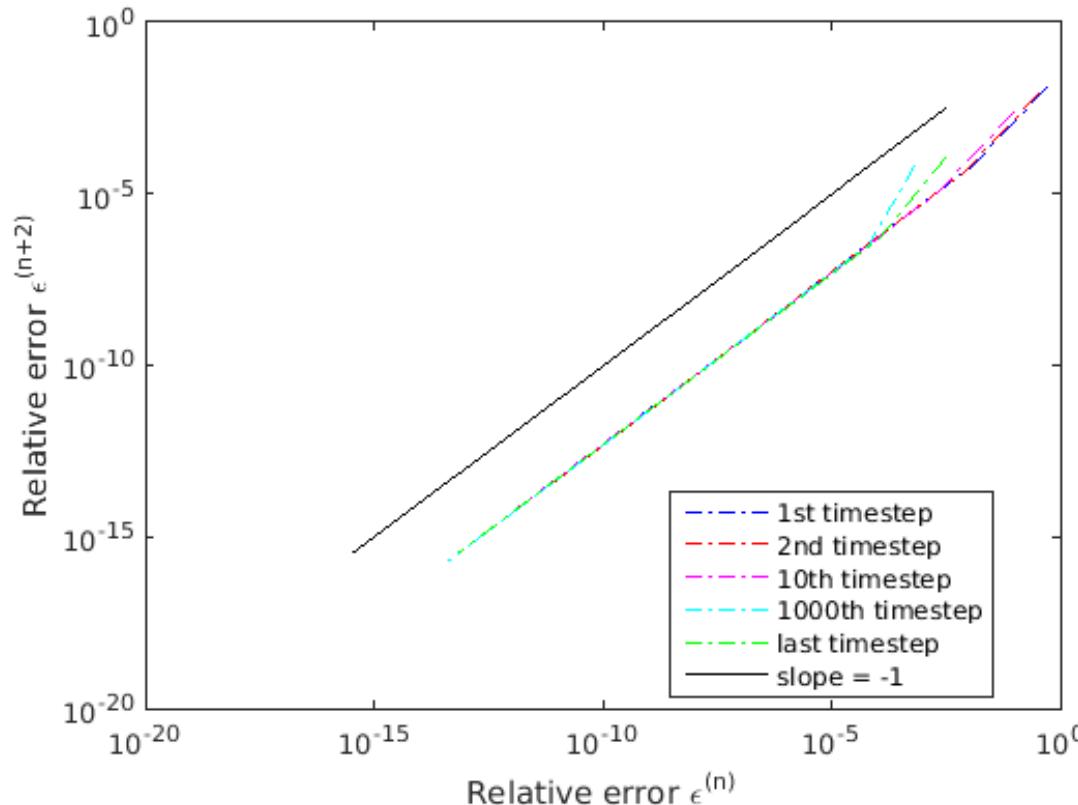
Bolted Joint Problem: Equivalent Plastic Strain (EQPS)



Single Ω

Schwarz

Bolted Joint Problem: Convergence Rate



Linear convergence rate is observed for **dynamic** Schwarz algorithm, as for the quasistatic Schwarz algorithm.



Figure above: Convergence behavior of the dynamic Schwarz algorithm for the bolted joint problem

Bolted Joint Problem: Performance



	CPU times (64 procs*)	Avg # Schwarz iters	Max # Schwarz iters
Single Domain	3h 34m	—	—
Schwarz	2h 42m	3.22	4
Single Domain (finer)	17h 00m	—	—
Schwarz (finer mesh of bolts)	29h 29m	3.28	4



Bolted Joint Problem: Performance



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Schwarz (finer mesh of bolts)	20h 29m	2.78	4



- Despite its iterative nature, Schwarz can actually be **faster** than single-domain run for discretizations having comparable # of elements in the bolts.

Bolted Joint Problem: Performance



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- Despite its iterative nature, Schwarz can actually be **faster** than single-domain run for discretizations having comparable # of elements in the bolts.
 - Even if the method is more computationally expensive for some resolutions, it may be preferred for its ability to **rapidly change** and **evaluate a variety of engineering designs** (our typical use case).

Bolted Joint Problem: Performance



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- Despite its run for days, the Schwarz method converges in 2-4 iterations.
- Even if the method is more computationally expensive for some resolutions, it may be preferred for its ability to *rapidly change* and *evaluate a variety of engineering designs* (our typical use case).
- Dynamic Schwarz converges in between **2-4 Schwarz iterations** per time-step despite the *overlap* region being *very small* for this problem.



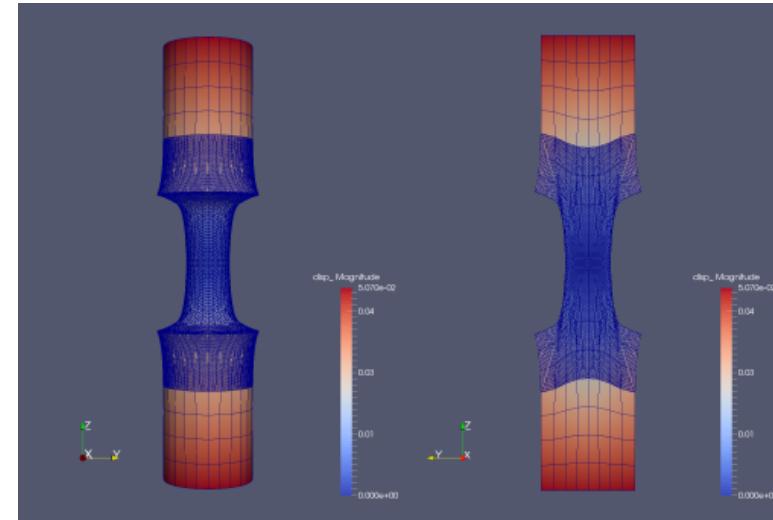
* On SNL asicgpu15, 16, 17 machines (Intel Skylake CPU processor), Schwarz tol = 1e-6.

Outline



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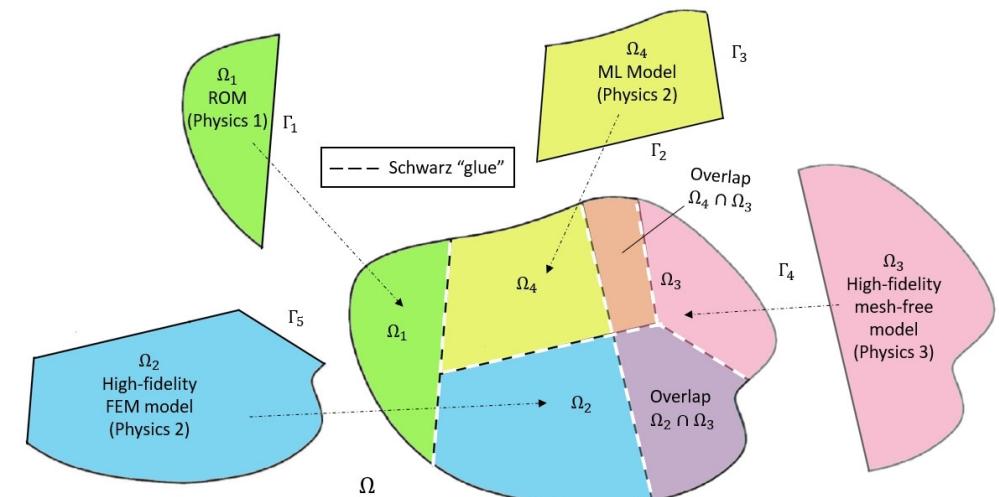
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2. Schwarz Alternating Method for FOM-ROM* and ROM-ROM Coupling

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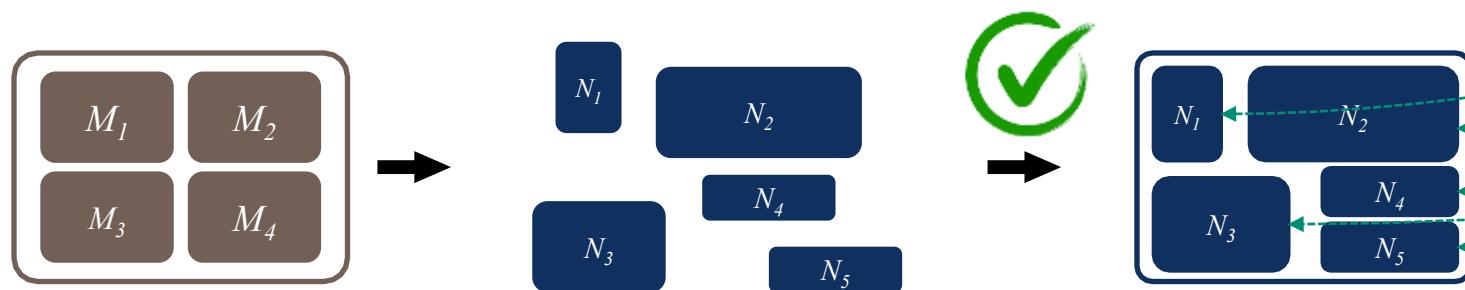


* Projection-based Reduced Order Model

Motivation

The past decades have seen tremendous investment in **simulation frameworks for coupled multi-scale and multi-physics problems**.

- Frameworks rely on **established mathematical theories** to couple physics components.
- Most existing coupling frameworks are based on **traditional discretization methods**.



Complex System Model

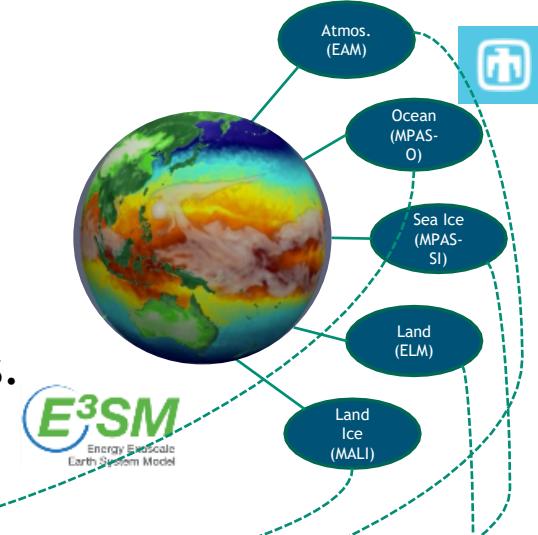
- PDEs, ODEs
- Nonlocal integral
- Classical DFT
- Atomistic, ...

Traditional Methods

- Mesh-based (FE, FV, FD)
- Meshless (SPH, MLS)
- Implicit, explicit
- Eulerian, Lagrangian...

Coupled Numerical Model

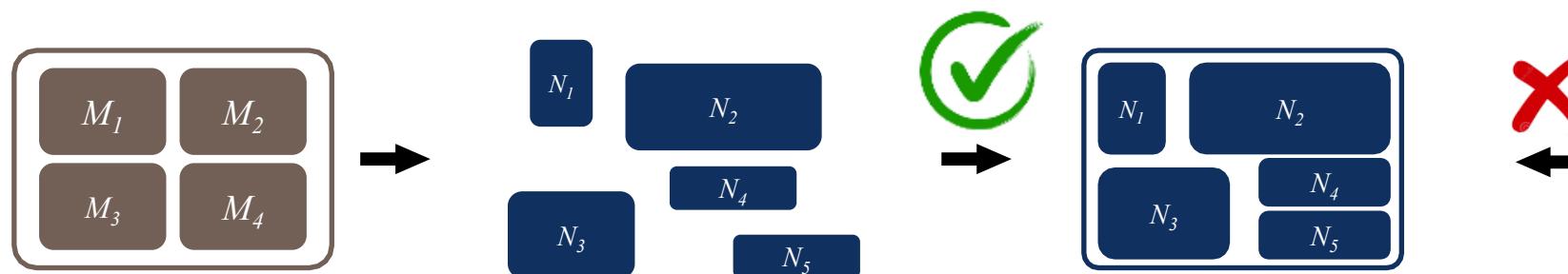
- Monolithic (Lagrange multipliers)
- Partitioned (loose) coupling
- Iterative (Schwarz, optimization)



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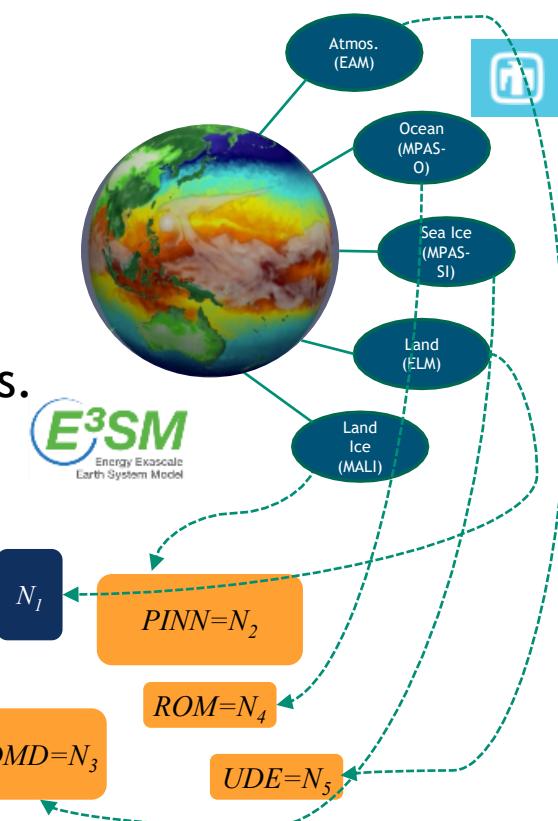
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Traditional + Data-Driven Methods

- PINNs
- Neural ODEs
- Projection-based ROMs, ...

- There is currently a big push to integrate **data-driven methods** into modeling & simulation toolchains.

Unfortunately, existing algorithmic and software infrastructures are **ill-equipped** to handle plug-and-play integration of **non-traditional, data-driven models**!



Principal research objective:

- Discover mathematical principles guiding the assembly of standard and data-driven numerical models in stable, accurate and physically consistent ways.

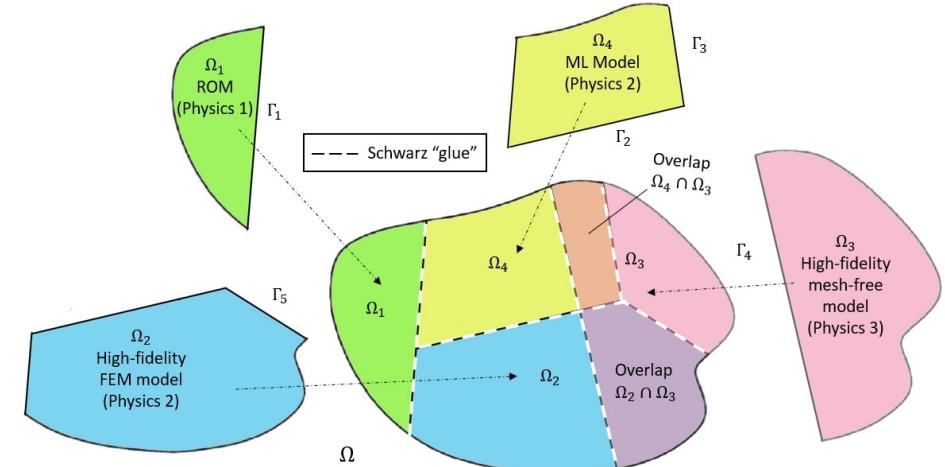
Principal research challenges:

we lack mathematical and algorithmic understanding of how to

- “Mix-and-match” standard and data-driven models from three-classes
 - *Class A*: projection-based reduced order models (ROMs) [This talk.](#)
 - *Class B*: machine-learned models, i.e., Physics-Informed Neural Networks (PINNs)
 - *Class C*: flow map approximation models, i.e., dynamic model decomposition (DMD) models
- Ensure well-posedness & physical consistency of resulting heterogeneous models.
- Solve such heterogeneous models efficiently.

Three coupling methods:

- Alternating Schwarz-based coupling [This talk.](#)
- Optimization-based coupling
- Coupling via generalized mortar methods



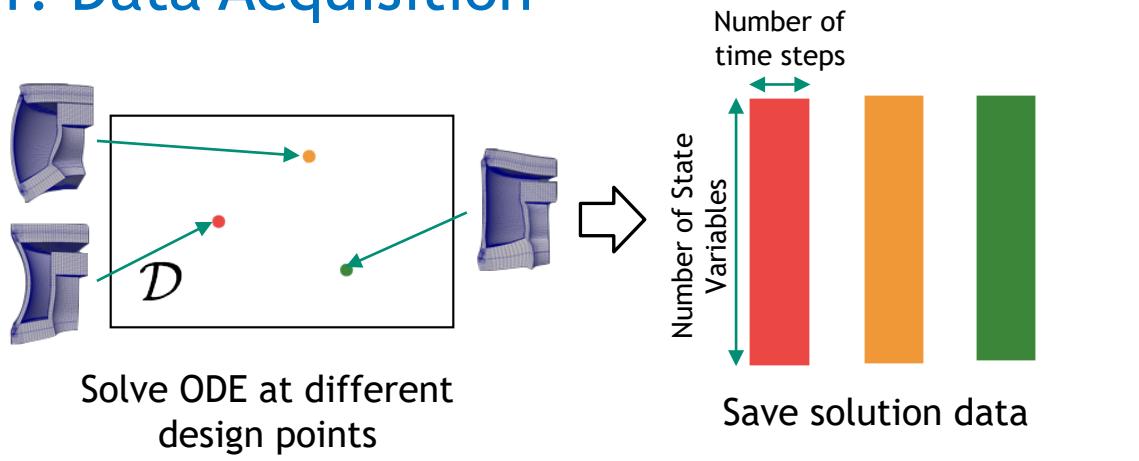
Projection-Based Model Order Reduction via the POD/LSPG*

Method

$$\text{Full Order Model (FOM): } \frac{\partial \mathbf{q}}{\partial t} = \mathbf{f}(\mathbf{q}, t; \boldsymbol{\mu})$$

*Least-Squares Petrov-Galerkin

1. Data Acquisition



2. Learning of Reduced Basis

Proper Orthogonal Decomposition (POD):

$$\mathbf{X} = \begin{array}{|c|c|c|} \hline \text{Red} & \text{Orange} & \text{Green} \\ \hline \end{array} = \begin{array}{|c|} \hline \text{Brown} \\ \hline \end{array} \begin{array}{|c|} \hline \text{Blue} \\ \hline \end{array} \begin{array}{|c|} \hline \Sigma \\ \hline \end{array} \begin{array}{|c|} \hline \text{Blue} \\ \hline \end{array}$$

ROM = projection-based Reduced Order Model

3. Projection-Based Reduction

Discretize FOM in time

$$\dot{\mathbf{q}} = \mathbf{f}(\mathbf{q}, t; \boldsymbol{\mu})$$

$$\mathbf{r}^n(\mathbf{q}^n; \boldsymbol{\mu}) = \mathbf{0}, \quad n = 1, \dots, T$$

Reduce the number of unknowns

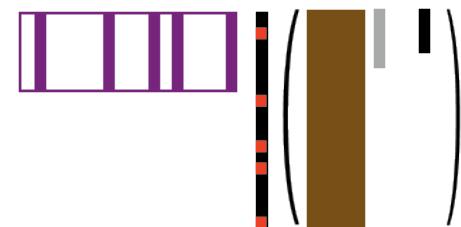
$$\mathbf{q}(t) \approx \tilde{\mathbf{q}}(t) = \Phi \hat{\mathbf{q}}(t)$$

Apply hyper-reduction and minimize residual



Hyper-reduction/sample mesh

$$\text{minimize}_{\boldsymbol{\vartheta}} \parallel \mathbf{A} \mathbf{r}^n(\Phi \boldsymbol{\vartheta}; \boldsymbol{\mu}) \parallel_2$$



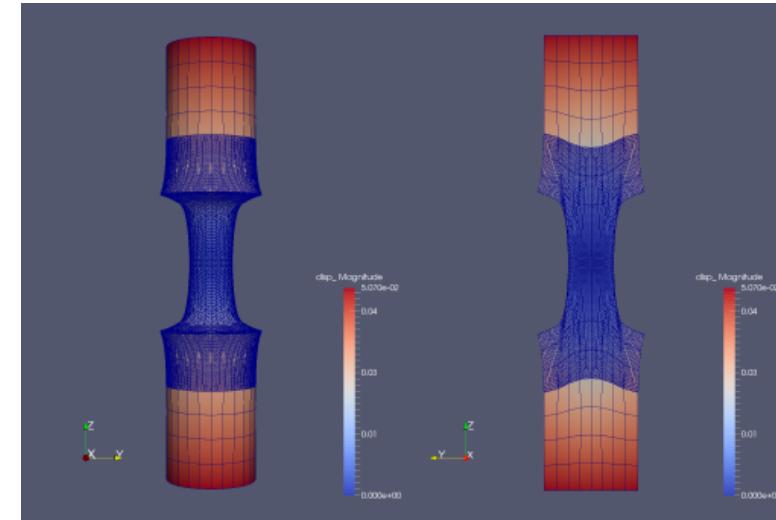
HROM = Hyper-reduced ROM

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1. Schwarz Alternating Method for Coupling of Full Order Models (FOMs) in Solid Mechanics

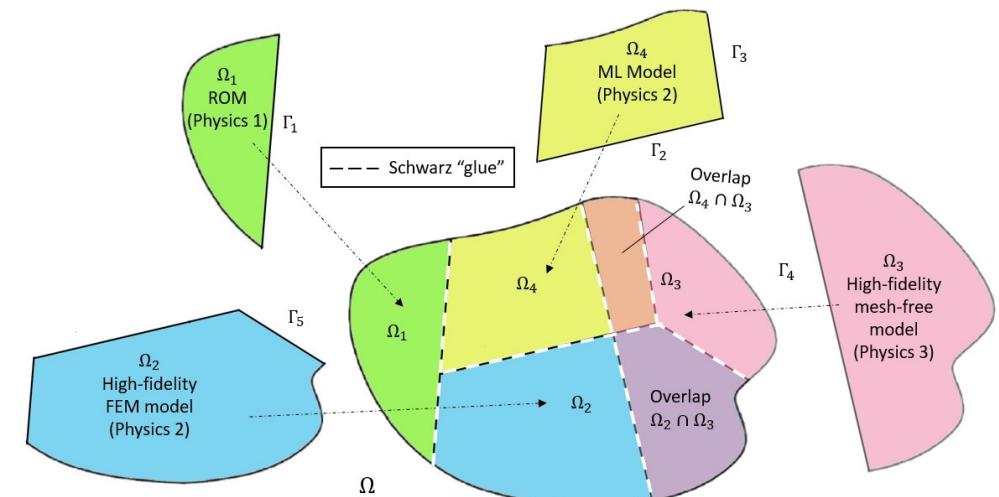
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* Projection-based Reduced Order Model

Schwarz Extensions to FOM-ROM and ROM-ROM Couplings



Choice of domain decomposition

- Overlapping vs. non-overlapping domain decomposition?
 - Non-overlapping more flexible but typically requires more Schwarz iterations
- FOM vs. ROM subdomain assignment?
 - Do not assign ROM to subdomains where they have no hope of approximating solution

Snapshot collection and reduced basis construction

- Are subdomains simulated independently in each subdomains or together?

Enforcement of boundary conditions (BCs) in ROM at Schwarz boundaries

- Strong vs. weak BC enforcement?
 - Strong BC enforcement difficult for some models (e.g., cell-centered finite volume, PINNs)
- Optimizing parameters in Schwarz BCs for non-overlapping Schwarz?

Choice of hyper-reduction

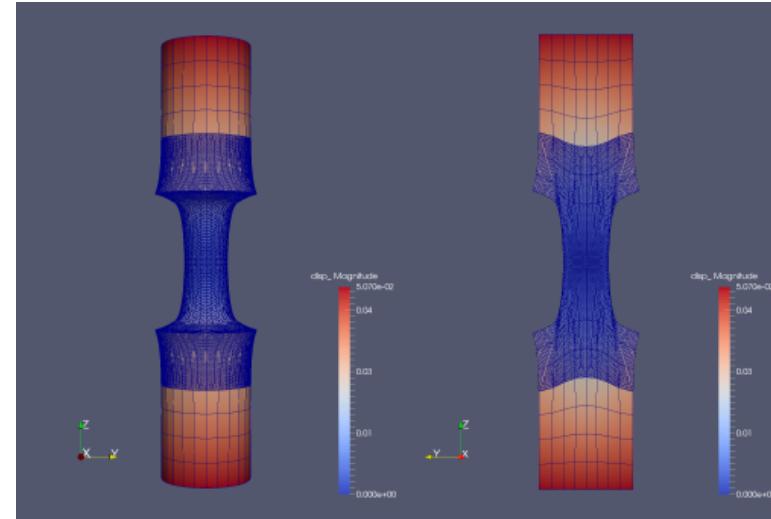
- What hyper-reduction method to use?
 - Application may require particular method (e.g., ECSW for solid mechanics problems)
- How to sample Schwarz boundaries in applying hyper-reduction?
 - Need to have enough sample mesh points at Schwarz boundaries to apply Schwarz

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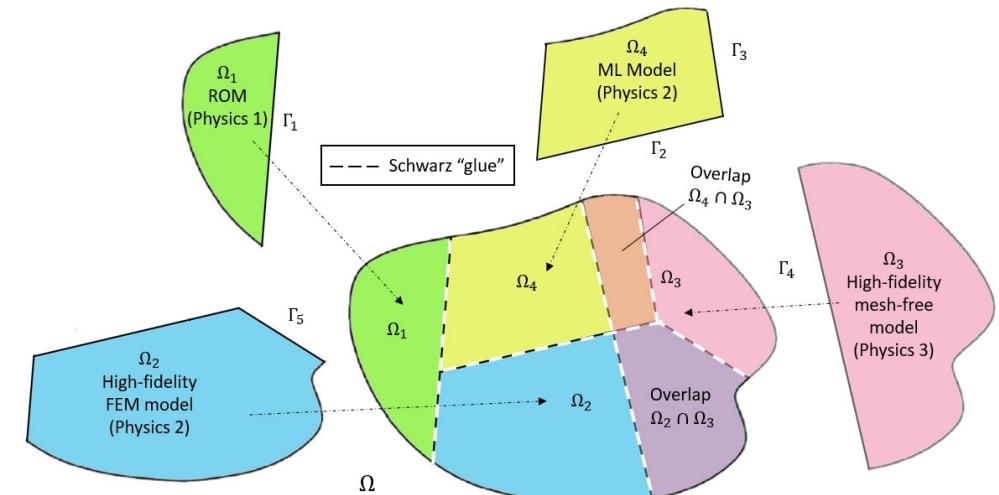
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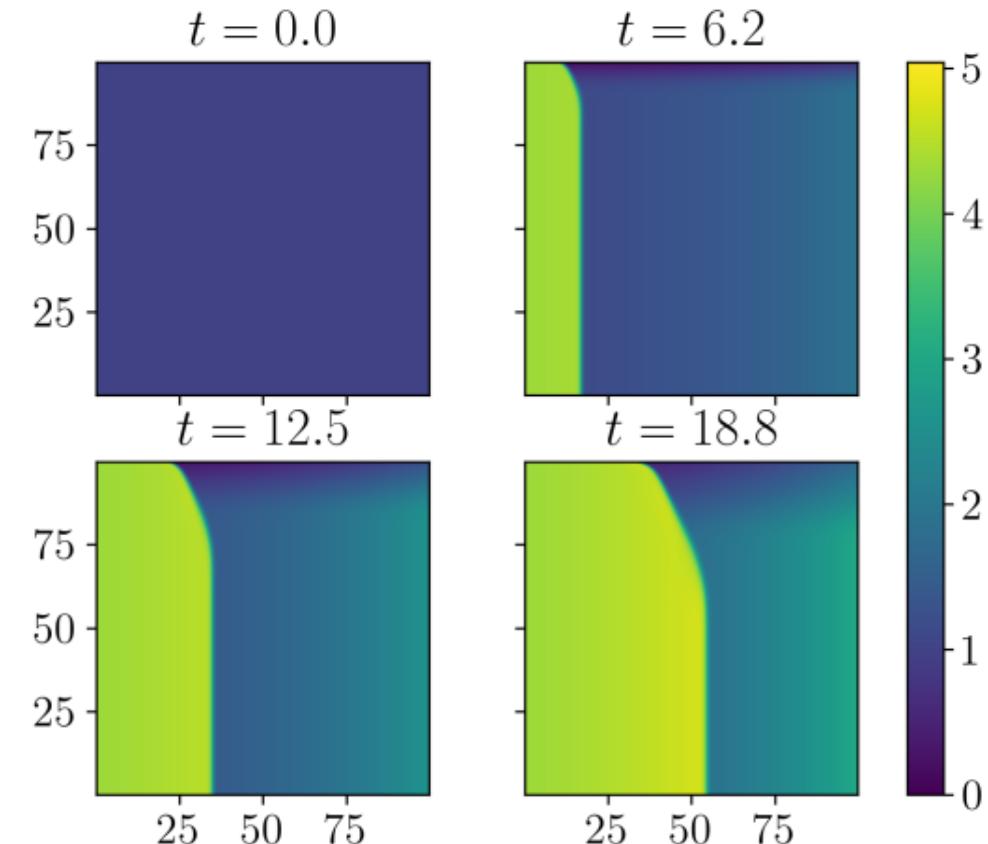
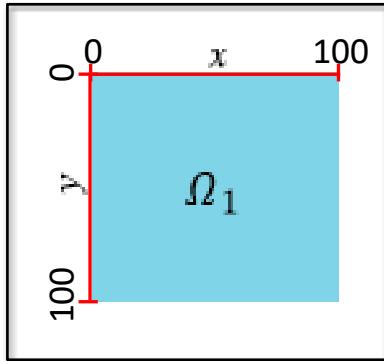


2D Inviscid Burgers Equation



Popular analog for fluid problems where **shocks** are possible, and particularly **difficult** for conventional projection-based ROMs

$$\begin{aligned}\frac{\partial u}{\partial t} + \frac{1}{2} \left(\frac{\partial (u^2)}{\partial x} + \frac{\partial (uv)}{\partial y} \right) &= 0.02 \exp(\mu_2 x) \\ \frac{\partial v}{\partial t} + \frac{1}{2} \left(\frac{\partial (vu)}{\partial x} + \frac{\partial (v^2)}{\partial y} \right) &= 0 \\ u(0, y, t; \mu) &= \mu_1 \\ u(x, y, 0) &= v(x, y, 0) = 1\end{aligned}$$



Problem setup:

- $\Omega = (0, 100)^2$, $t \in [0, 25]$
- Two parameters $\mu = (\mu_1, \mu_2)$. **Training:** uniform sampling of $=[4.25, 5.50] \times [0.015, 0.03]$ by a 3×3 grid. **Testing:** query unsampled point $\mu = [4.75, 0.02]$

FOM discretization:

- Spatial discretization given by a **Godunov-type scheme** with $N = 250$ elements in each dimension
- Implicit **trapezoidal method** with fixed $\Delta t = 0.05$

Figure above: solution of u component at various times

Schwarz Coupling Details



Choice of domain decomposition

- Overlapping DD of Ω into 4 subdomains coupled via multiplicative Schwarz
- Solution in Ω_1 is most difficult to capture by ROM

Snapshot collection and reduced basis construction

- Single-domain FOM on Ω used to generate snapshots/POD modes

Enforcement of boundary conditions (BCs) in ROM at Schwarz boundaries

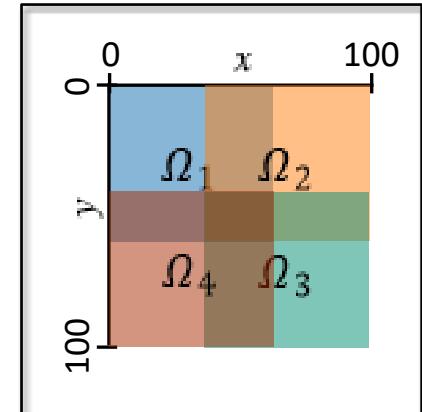
- BCs imposed strongly using Method 1 of [Gunzburger *et al.*, 2007] at indices i_{Dir}

$$\mathbf{q}(t) \approx \bar{\mathbf{q}} + \Phi \hat{\mathbf{q}}(t)$$

- POD modes made to satisfy homogeneous DBCs: $\Phi(i_{\text{Dir}}, :) = \mathbf{0}$
- BCs imposed by modifying $\bar{\mathbf{q}}$: $\bar{\mathbf{q}}(i_{\text{Dir}}) \leftarrow \chi_q$

Choice of hyper-reduction

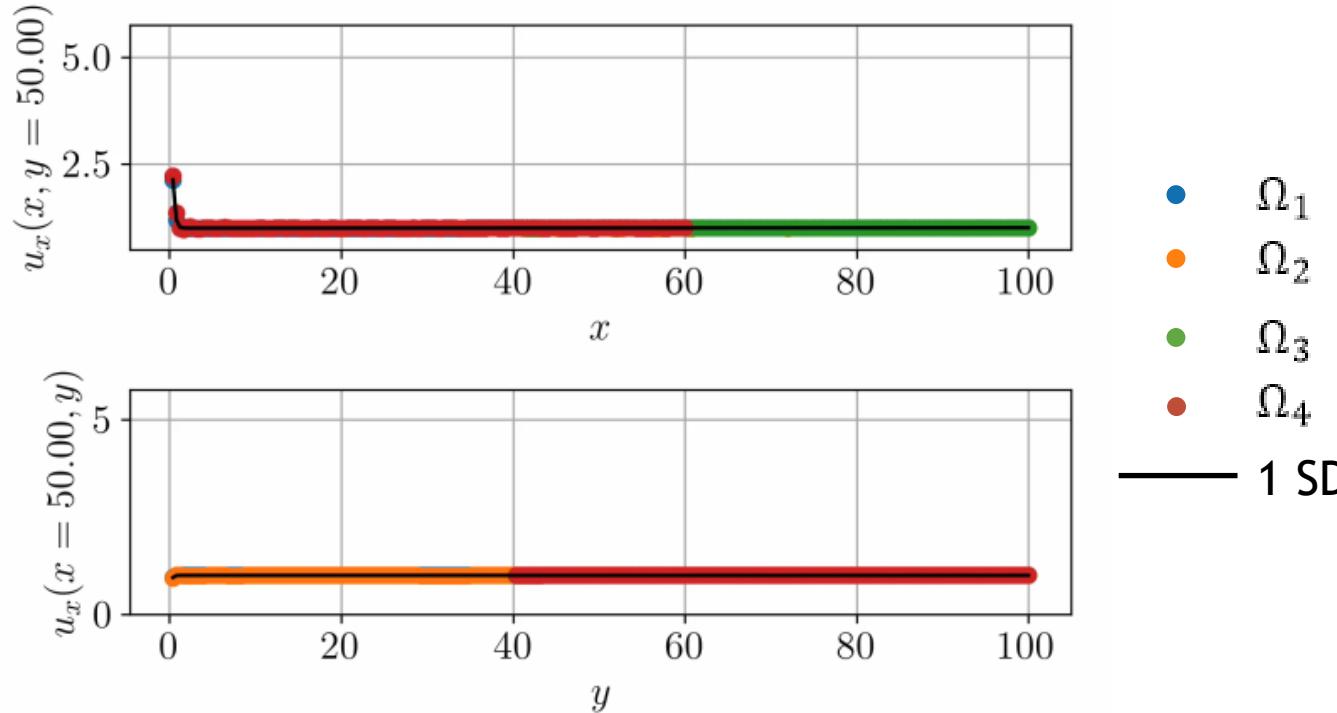
- Energy Conserving Sampling & Weighting (ECSW) method for hyper-reduction
- All points on Schwarz boundaries are included in the sample mesh



FOM-HROM-HROM-HROM Coupling



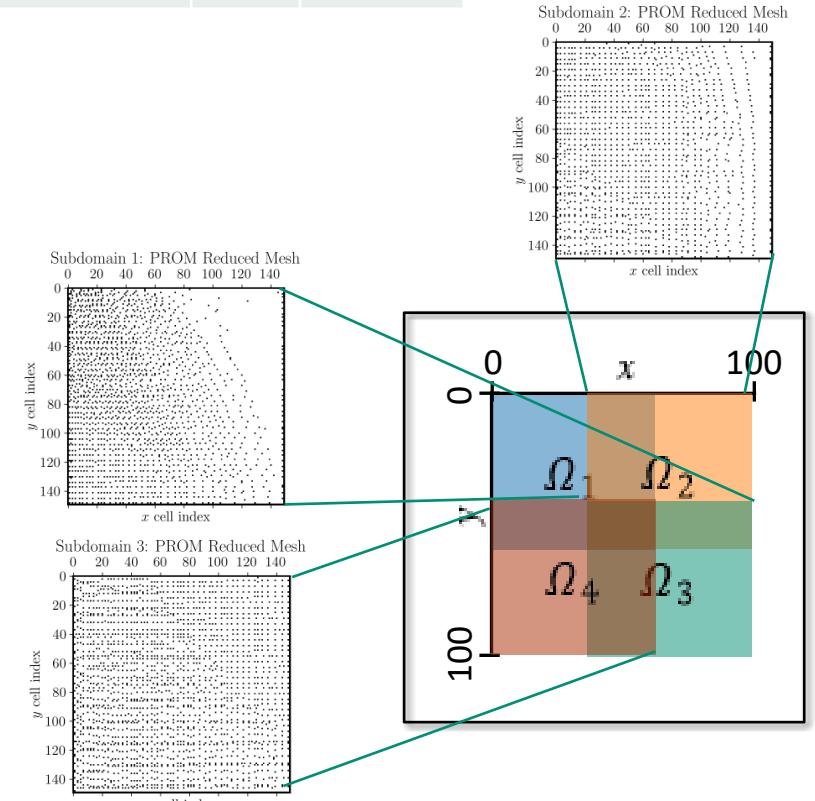
62



- **FOM in Ω_1** as this is “**hardest**” subdomain for ROM
- **HROMs in $\Omega_2, \Omega_3, \Omega_4$ capture 99% snapshot energy**
- Method converges in **3 Schwarz iterations** per controller time-step
- Errors $O(0.1\%)$ with 0 error in Ω_1
- **2.26× speedup** achieved over all-FOM coupling

Further speedups possible via **code optimizations**,
additive Schwarz and reduction of # sample mesh points.

Subdomains	99% SV Energy		
	M	MSE (%)	CPU time (s)
Ω_1	—	0.0	95
Ω_2	120	0.26	26
Ω_3	60	0.43	17
Ω_4	66	0.34	21
Total			159



2D Shallow Water Equations (SWE)



Hyperbolic PDEs modeling wave propagation below a pressure surface in a fluid (e.g., atmosphere, ocean).

$$\begin{aligned}\frac{\partial h}{\partial t} + \frac{\partial(hu)}{\partial x} + \frac{\partial(hv)}{\partial y} &= 0 \\ \frac{\partial(hu)}{\partial t} + \frac{\partial}{\partial x} \left(hu^2 + \frac{1}{2} gh^2 \right) + \frac{\partial}{\partial y} (huv) &= -\mu v \\ \frac{\partial(hv)}{\partial t} + \frac{\partial}{\partial x} (huv) + \frac{\partial}{\partial y} \left(hv^2 + \frac{1}{2} gh^2 \right) &= \mu u\end{aligned}$$

Problem setup:

- $\Omega = (-5, 5)^2$, $t \in [0, 10]$, Gaussian initial condition
- Coriolis parameter** $\mu \in \{-4, -3, -2, -1, 0\}$ for training, and $\mu \in \{-3.5, -2.5, -1.5, -0.5\}$ for testing

FOM discretization:

- Spatial discretization given by a first-order **cell-centered finite volume** discretization with $N = 300$ elements in each dimension
- Implicit first order temporal discretization: **backward Euler** with fixed $\Delta t = 0.01$
- Implemented in **Pressio** (<https://github.com/Pressio/pressio-demoapps>)

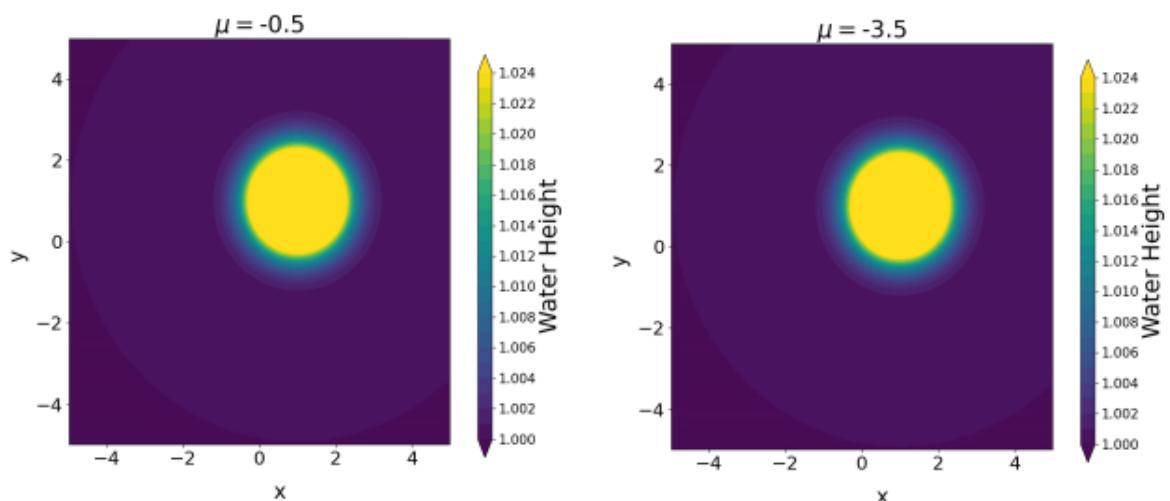


Figure above: FOM solutions to SWE for $\mu = -0.5$ (left) and $\mu = -3.5$ (right).

Schwarz Coupling Details

Choice of domain decomposition

- Non-overlapping DD of Ω into 4 subdomains coupled via additive Schwarz
 - OpenMP parallelism with 1 thread/subdomain
- All-ROM or All-HROM coupling via Pressio*

Snapshot collection and reduced basis construction

- Single-domain FOM on Ω used to generate snapshots/POD modes

Enforcement of boundary conditions (BCs) in ROM at Schwarz boundaries

- BCs are imposed approximately by fictitious ghost cell states
 - Implementing Neumann and Robin BCs is challenging
- Ghost cells introduce some overlap even with non-overlapping DD
 - \Rightarrow Dirichlet-Dirichlet non-overlapping Schwarz is stable/convergent!

Choice of hyper-reduction

- Collocation for hyper-reduction: min residual at small subset DOFs
- Assume fixed budget of sample mesh points at Schwarz boundaries

<https://github.com/Pressio/pressio-demoapps>

Green: different from Burgers' problem

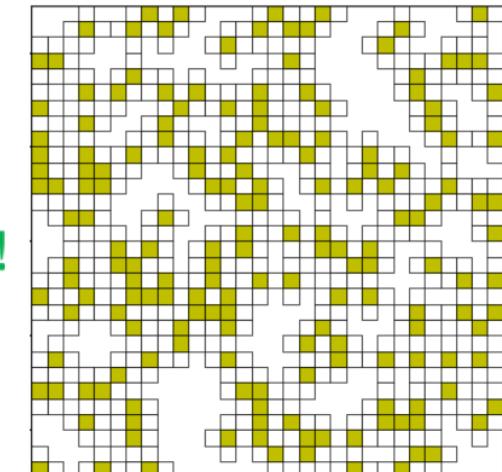
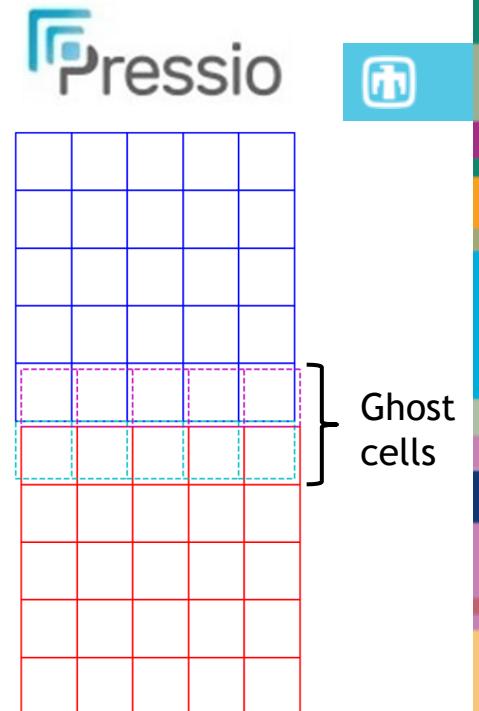
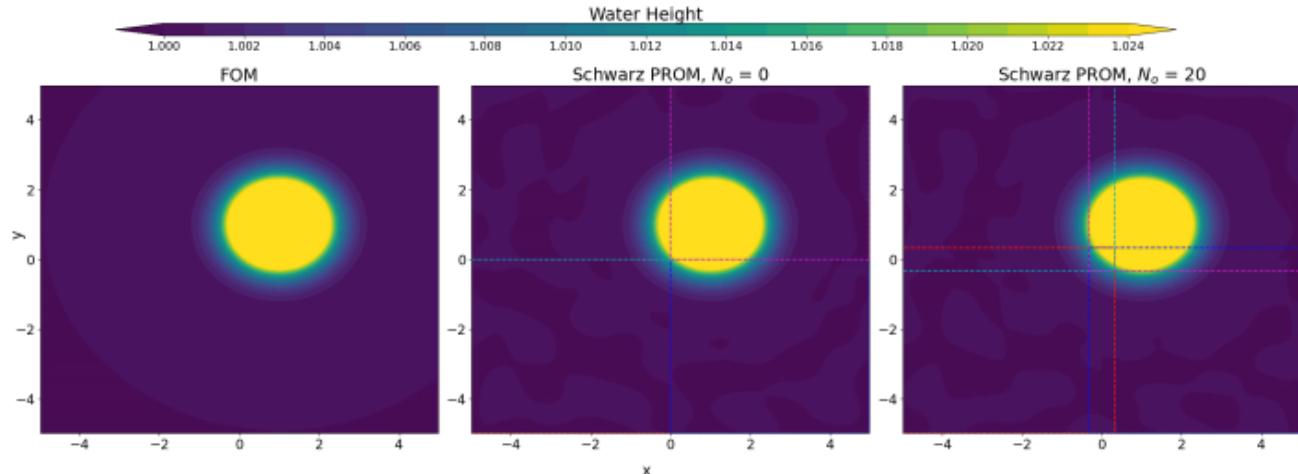


Figure above: sample mesh (yellow) and stencil (white) cells

Schwarz All-ROM Domain Overlap Study



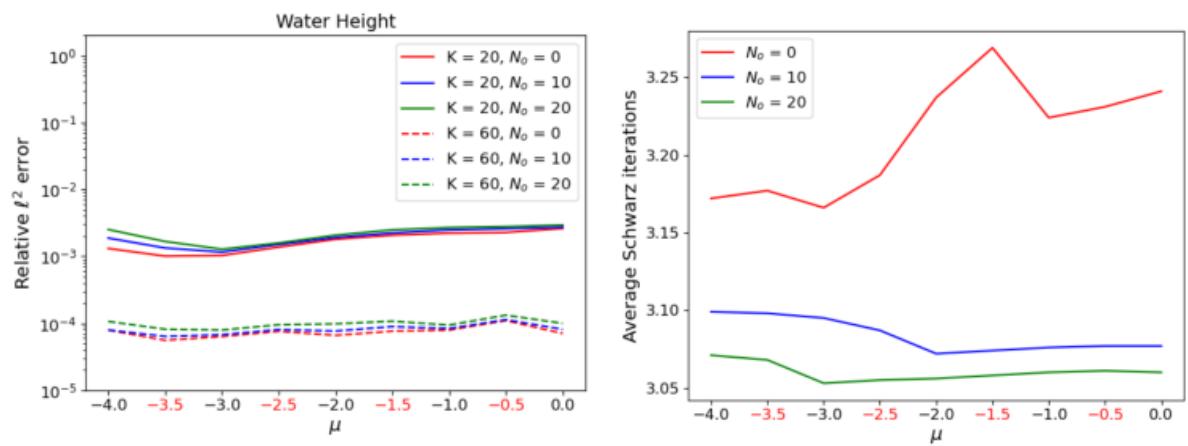
Study of Schwarz convergence for all-ROM coupling as a function of $N_o :=$ cell width of overlap region (not including ghost cells).



Movie above: FOM (left), 4 subdomain ROM coupled via non-overlapping Schwarz (middle), and 4 subdomain ROM coupled via overlapping Schwarz (right) for predictive SWE problem with $\mu = -0.5$. All ROMs have $K = 80$ POD modes.

- **Schwarz iterations** decrease (very roughly) with $N_o^{0.25}$ (figure, right) whereas evaluating $r(q)$ scales with N_o^2
 - \Rightarrow there is no reason not to do **non-overlapping coupling** for this problem

- Dirichlet-Dirichlet coupling with **no-overlap** ($N_o = 0$) performs well with **no convergence issues** (movie, left) and **errors comparable** to Dirichlet-Dirichlet coupling with overlap (figure below, left)



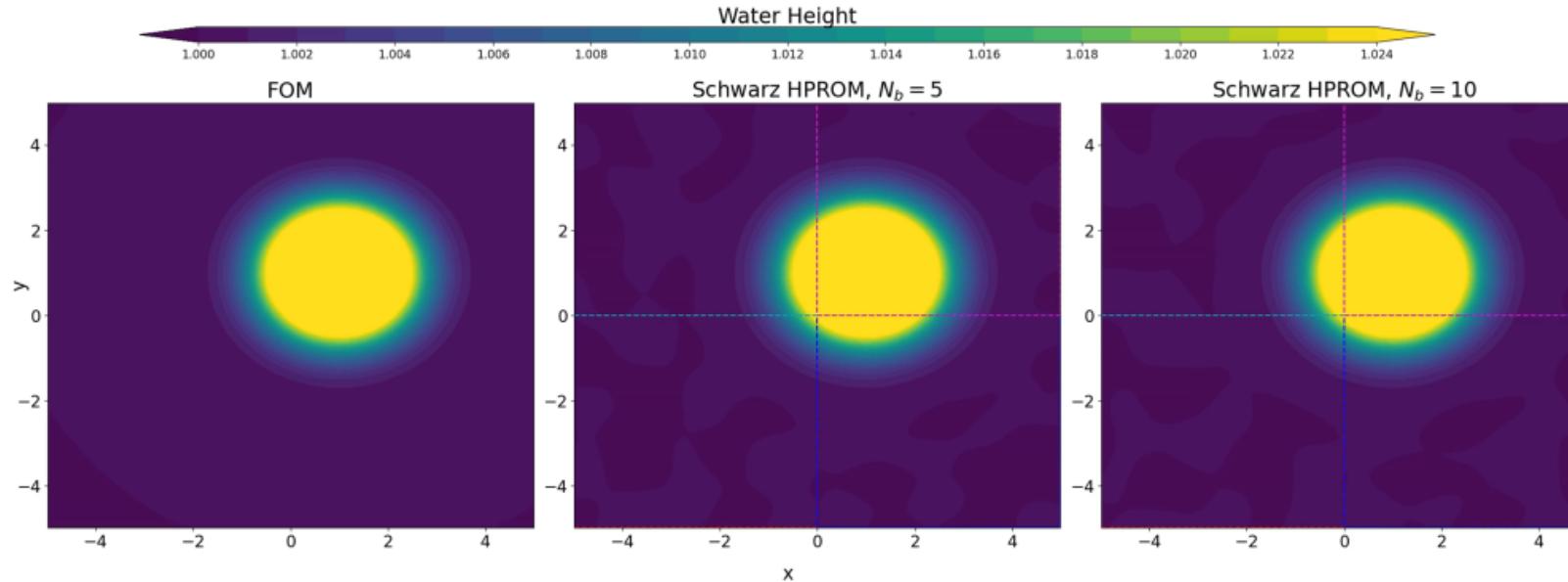
Figures above: relative error and average # Schwarz iterations as a function of μ and N_o . Black μ : training, red μ : testing.

Schwarz Boundary Sampling for All-HROM Coupling



Key question: how many Schwarz boundary points need to be included in sample mesh when performing HROM coupling?

- Naive/sparsely-sampled Schwarz boundary results in **failure** to transmit coupling information during Schwarz



Movie above: FOM (left), all HROM with $N_b = 5\%$ (middle) and all HROM with $N_b = 10\%$ (left). ROMs have $K = 100$ modes and $N_s = 0.5\%N$ sample mesh points.

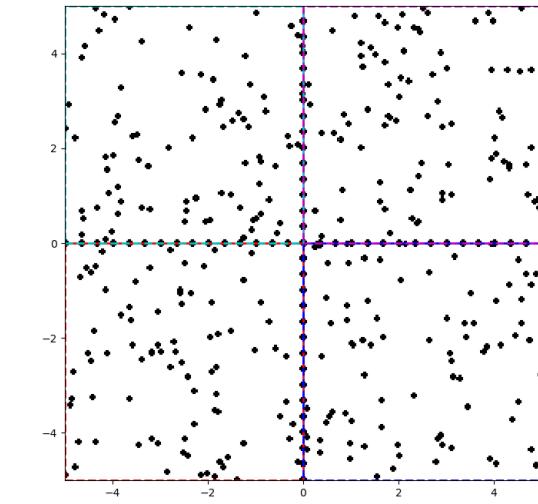
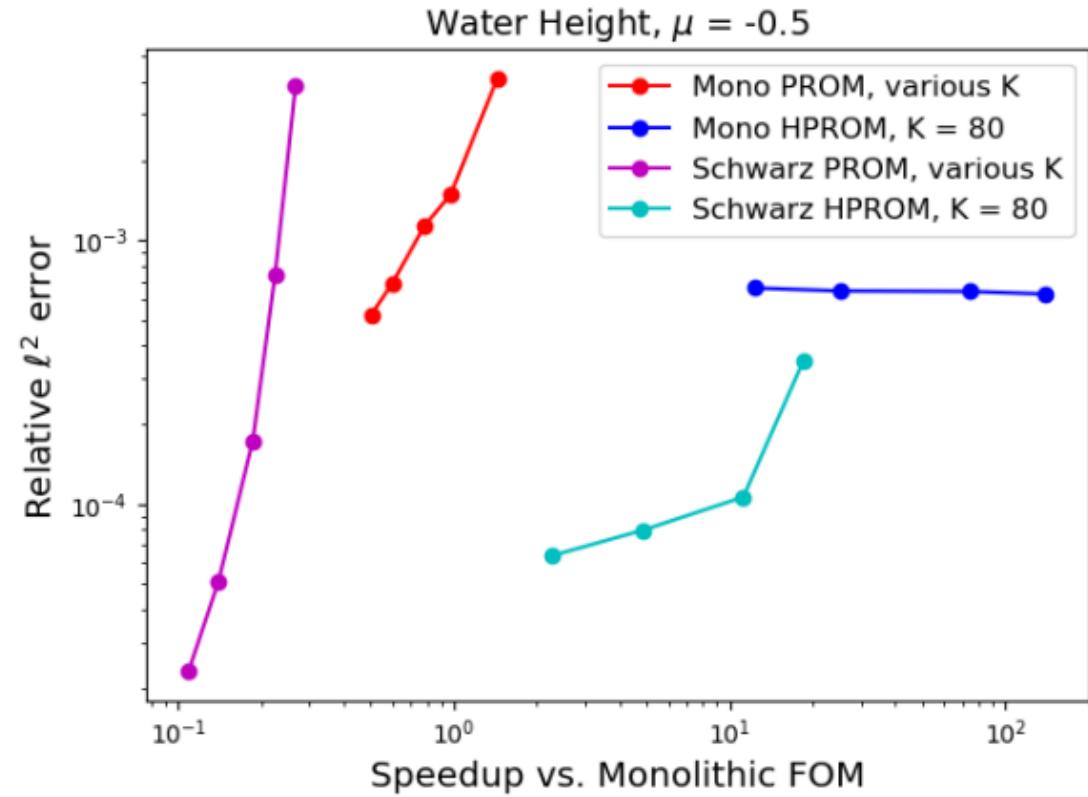
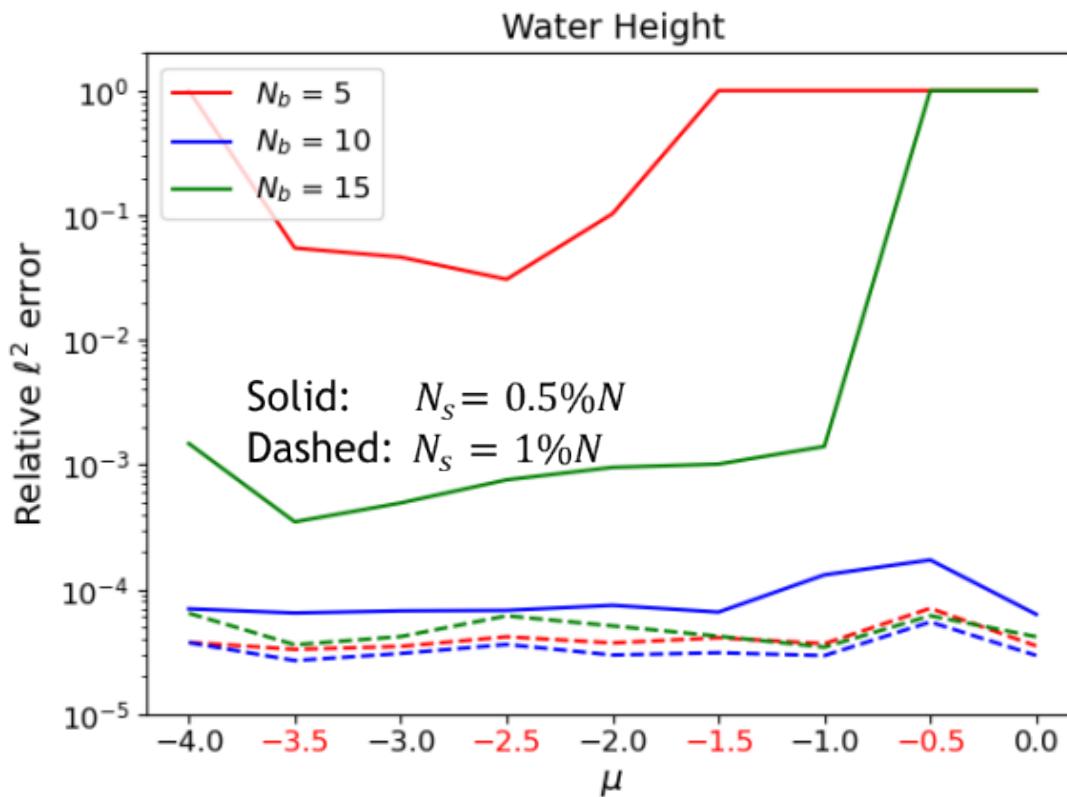


Figure above: example sample mesh with sampling rate $N_b = 10\%$

- Including **too many Schwarz boundary points** (N_b) in sample mesh given **fixed budget** of N_s sample mesh points may lead to **too few sample mesh points** in interior
- For SWE problem, we can get away with **~10% boundary sampling** (movie above, right-most frame)

Coupled HROM Performance



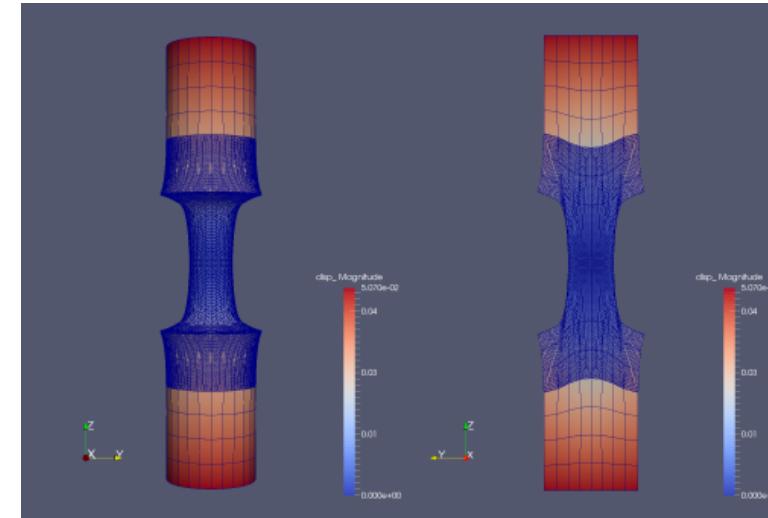
- For a fixed ROM dimension, Schwarz delivers **lower error and comparable cost!**
- There are noticeable **cost savings** relative to **monolithic FOM!**
- Accuracy similar for **predictive μ** (red) and **non-predictive μ** (black) cases.

Outline



1. Schwarz Alternating Method for Coupling of Full Order Models (FOMs) in Solid Mechanics

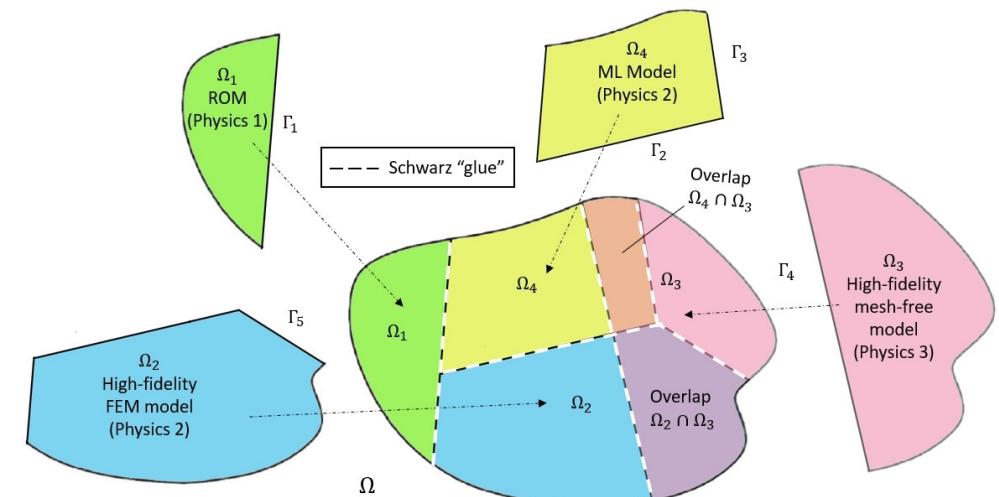
- Motivation & Background
- Quasistatic Formulation
 - Numerical Examples
- Extension to Dynamics
 - Numerical Examples



2. Schwarz Alternating Method for FOM-ROM* and ROM-ROM Coupling

- Motivation & Background
- Formulation
- Numerical Examples

3. Summary and Future Work



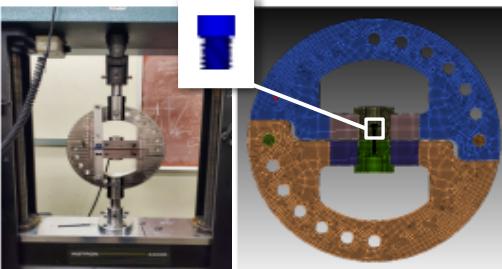


The Schwarz alternating method has been developed for concurrent multi-scale coupling of conventional and data-driven models.

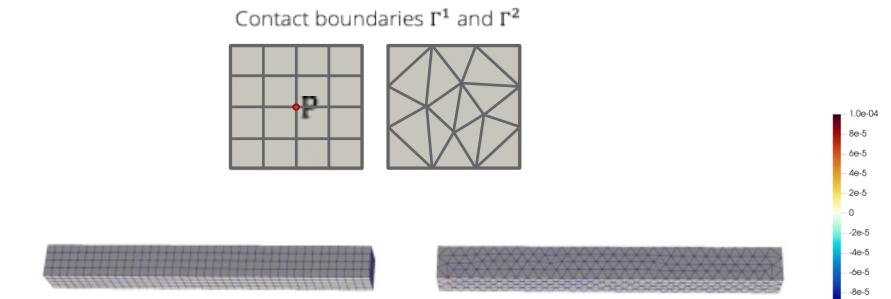
- 😊 Coupling is ***concurrent*** (two-way).
- 😊 ***Ease of implementation*** into existing massively-parallel HPC codes.
- 😊 “***Plug-and-play***” ***framework***: simplifies task of meshing complex geometries!
- 😊 Ability to couple regions with ***different non-conformal meshes***, ***different element types*** and ***different levels of refinement***.
- 😊 Ability to use ***different solvers (including ROM/FOM)*** and ***time-integrators*** in different regions.
- 😊 ***Scalable, fast, robust*** on ***real*** engineering problems
- 😊 Coupling does not introduce ***nonphysical artifacts***.
- 😊 ***Theoretical*** convergence properties/guarantees.

Ongoing & Future Work

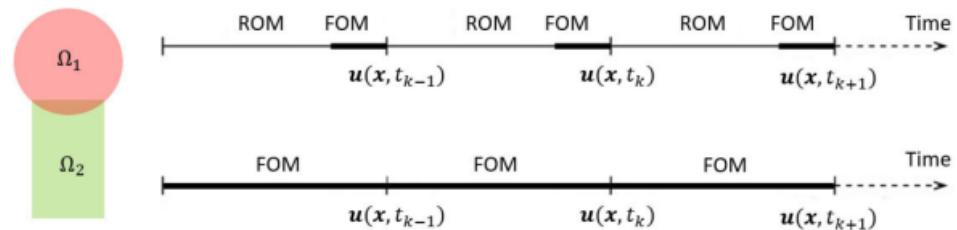
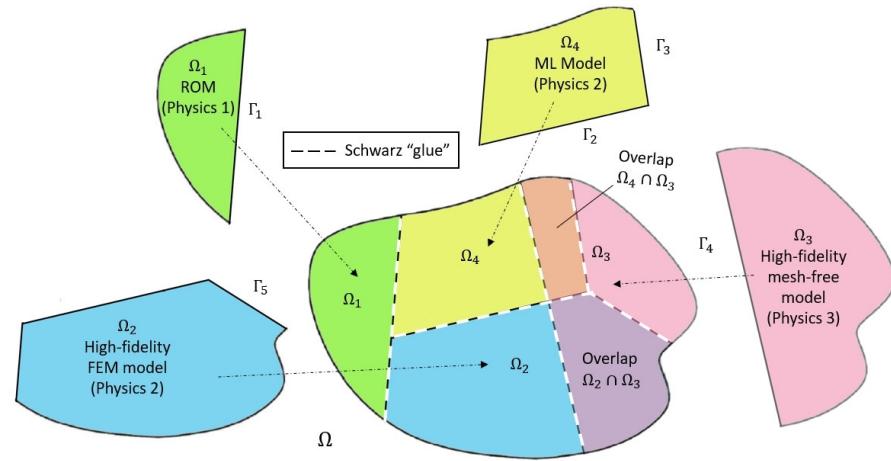
- Development fundamentally new approach for simulating multi-scale mechanical contact using the **Dirichlet-Neumann Schwarz alternating method**
 - Contact constraints are replaced with **boundary conditions** applied iteratively at contact boundaries
- Implementation of **non-overlapping** Schwarz in Sierra/SM
- Working with **analysts to apply** Schwarz to problems of interest to Sandia missions
 - **Laser welds**
 - **Fastener** modeling for joints
 - **Salt caverns for oil storage**
- **Rigorous analysis** of why Dirichlet-Dirichlet BC “work” when employing non-overlapping Schwarz with discretizations that employ ghost cells
- Extension to coupling of **non-intrusive ROMs** (dynamic mode decomposition, operator inference, neural networks)
- Development of **automated criteria** to determine appropriate use of less refined or reduced-order models w/o sacrificing accuracy, enabling **real-time transitions** between different model fidelities



From Murugesan *et al.*, 2020.



Impact of two 3D beams having different meshes with Schwarz contact method. From [Mota *et al.*, 2023].



Team & Acknowledgments



Irina Tezaur



Joshua Barnett



Alejandro Mota



Chris Wentland



Francesco Rizzi



Sandia
National
Laboratories



Coleman Alleman



Greg Phlipot





[1] A. Mota, **I. Tezaur**, C. Alleman. “The Schwarz Alternating Method in Solid Mechanics”, *Comput. Meth. Appl. Mech. Engng.* 319 (2017), 19-51.

[2] A. Mota, **I. Tezaur**, G. Phlipot. “The Schwarz Alternating Method for Dynamic Solid Mechanics”, *Comput. Meth. Appl. Mech. Engng.* 121 (21) (2022) 5036-5071.

[3] J. Barnett, **I. Tezaur**, A. Mota. “The Schwarz alternating method for the seamless coupling of nonlinear reduced order models and full order models”, ArXiv pre-print, 2022.

<https://arxiv.org/abs/2210.12551>

[4] W. Snyder, **I. Tezaur**, C. Wentland. “Domain decomposition-based coupling of physics-informed neural networks via the Schwarz alternating method”, ArXiv pre-print, 2023.

<https://arxiv.org/abs/2311.00224>

[5] A. Mota, D. Koliesnikova, **I. Tezaur**. “A Fundamentally New Coupled Approach to Contact Mechanics via the Dirichlet-Neumann Schwarz Alternating Method”, ArXiv pre-print, 2023.

<https://arxiv.org/abs/2311.05643>

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URL: www.sandia.gov/~ikalash

Start of Backup Slides

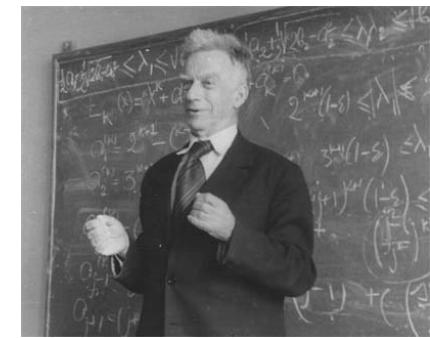
Theoretical Foundation

Using the Schwarz alternating as a *discretization method* for PDEs is natural idea with a sound *theoretical foundation*.

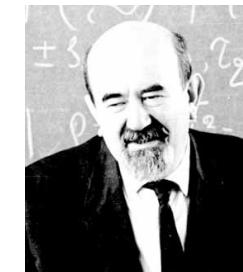
- [S.L. Sobolev \(1936\)](#): posed Schwarz method for *linear elasticity* in variational form and *proved method's convergence* by proposing a convergent sequence of energy functionals.
- [S.G. Mikhlin \(1951\)](#): *proved convergence* of Schwarz method for general linear elliptic PDEs.
- [P.-L. Lions \(1988\)](#): studied convergence of Schwarz for *nonlinear monotone elliptic problems* using max principle.
- [A. Mota, I. Tezaur, C. Alleman \(2017\)](#): proved *convergence* of the alternating Schwarz method for *finite deformation quasi-static nonlinear PDEs* (with energy functional $\Phi[\varphi]$) with a *geometric convergence rate*.

$$\Phi[\varphi] = \int_B A(\mathbf{F}, \mathbf{Z}) dV - \int_B \mathbf{B} \cdot \varphi dV$$

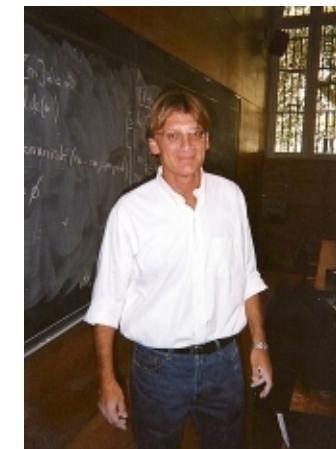
$$\nabla \cdot \mathbf{P} + \mathbf{B} = \mathbf{0}$$



S.L. Sobolev (1908 - 1989)



S.G. Mikhlin
(1908 - 1990)



P.- L. Lions (1956-)



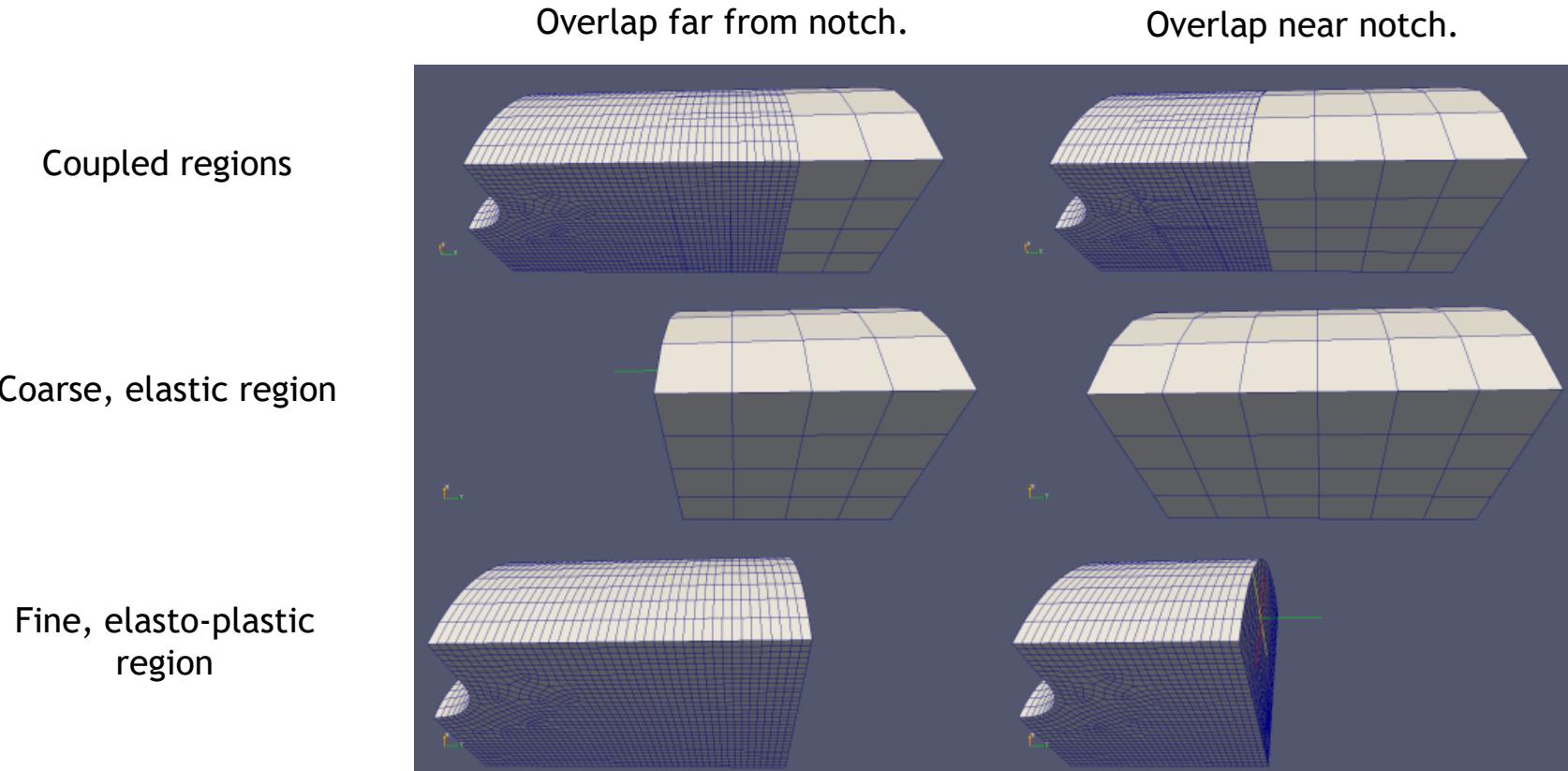
A. Mota, I. Tezaur, C. Alleman

Notched Cylinder: Coupling Different Materials



The Schwarz method is capable of coupling regions with *different material models*.

- Notched cylinder subjected to tensile load with an *elastic* and *J2 elasto-plastic* regions.
- Coarse* region is *elastic* and *fine* region is *elasto-plastic*.
- The *overlap region* in the first mesh is nearer the notch, where plastic behavior is expected.

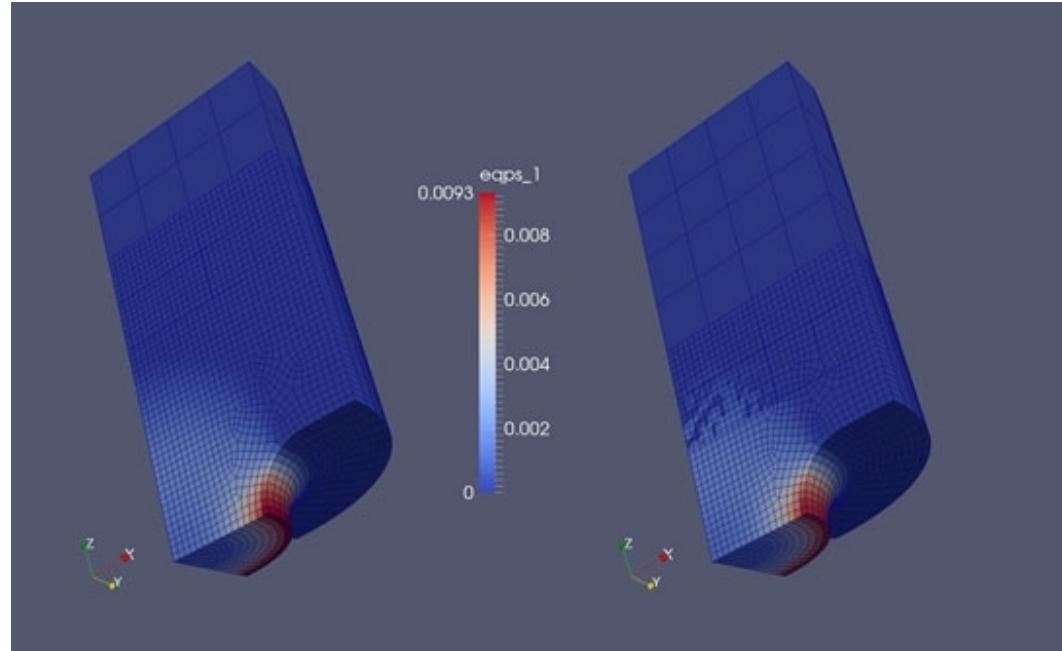


Notched Cylinder: Coupling Different Materials



Need to be careful to do domain decomposition so that material models are *consistent* in overlap region.

- When the *overlap* region is *far from the notch*, no plastic deformation exists in it: the coarse and fine regions predict the *same behavior*.
- When the *overlap* region is *near the notch*, plastic deformation spills onto it and the two models predict different behavior, affecting convergence *adversely*.



Overlap far from notch.

Overlap near notch.

Single Domain Predictive ROM



- Uniform sampling of $\mathcal{D} = [4.25, 5.50] \times [0.015, 0.03]$ by a 3×3 grid
 $\Rightarrow 9$ training parameters characterized by $\Delta\mu_1 = 0.625, \Delta\mu_2 = 0.0075$
 - > 200 POD modes required to capture 99% snapshot energy
- Queried but **unsampled parameter** point $\mu = [4.75, 0.02]$
- **Reduced mesh** resulting from solving non-negative least squares problem defining ECSW gives $n_e = 5,689$ elements (9.1% of $N_e = 62,500$ elements).

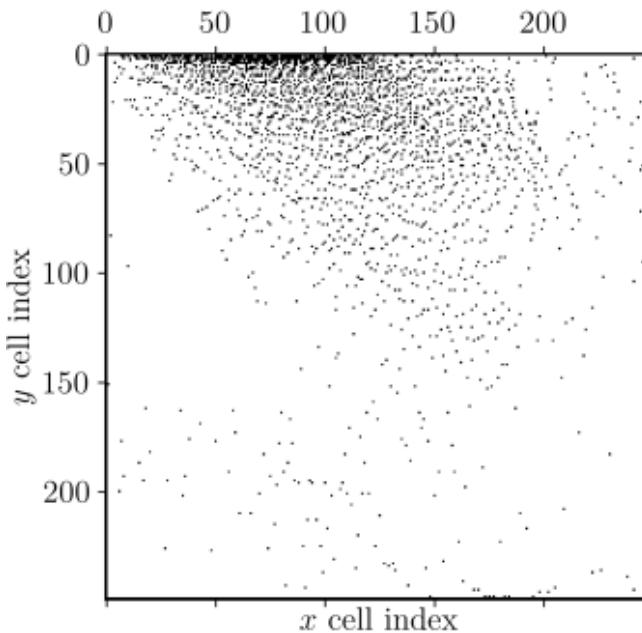


Figure above: Reduced mesh of single domain HROM

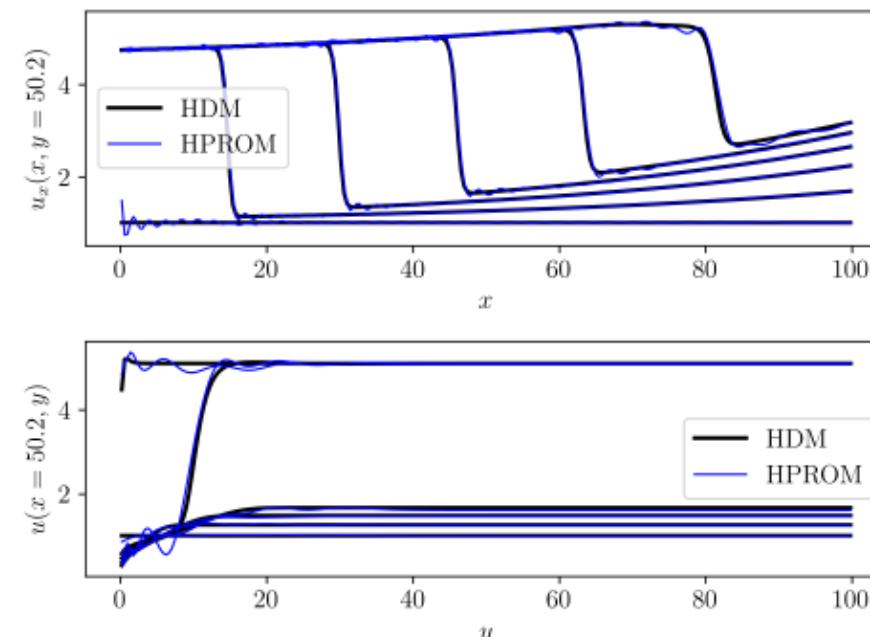
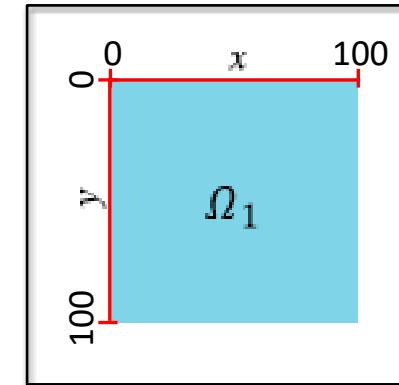
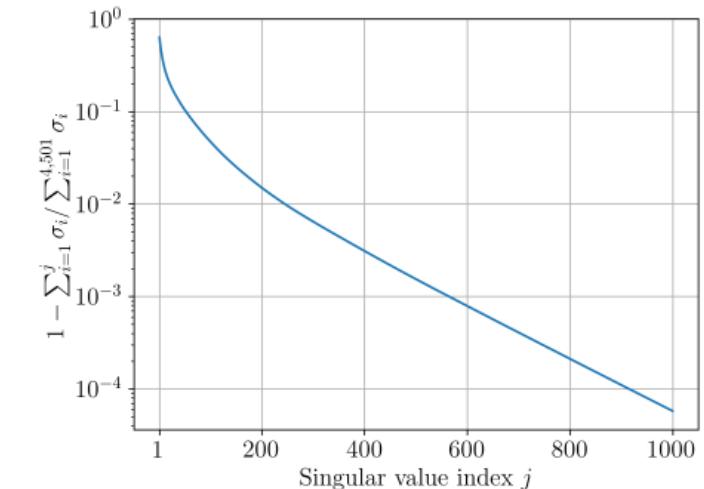


Figure above: HROM and FOM results at various time steps

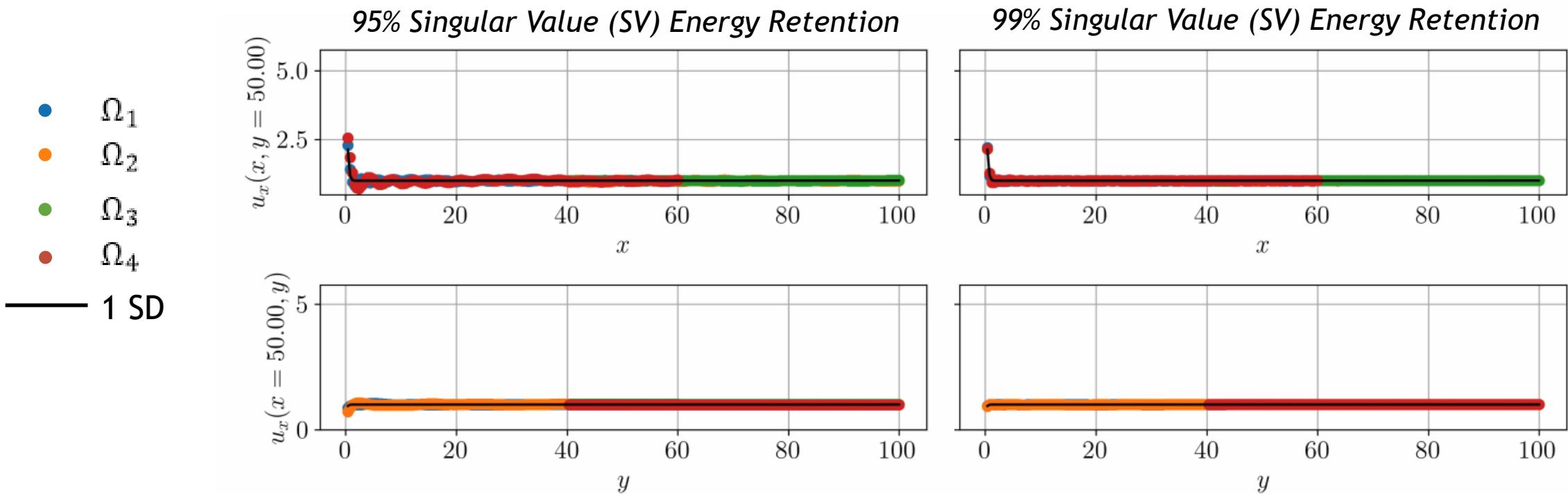


% SV Energy	M	MSE* (%)	CPU time* (s)
95	69	1.1	138
99	177	0.17	447

* Numbers in table are w/o hyper-reduction

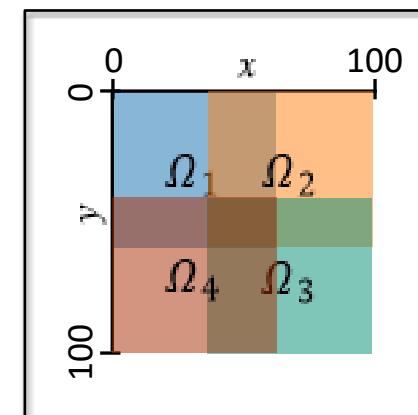


All-ROM Coupling



- Method converges in **only 3 Schwarz iterations** per controller time-step
- Errors $O(1\%)$ or less
- 1.47 \times speedup** over all-FOM coupling for 95% SV energy retention case

Subdomains	95% SV Energy		99% SV Energy		
	M	MSE (%)	CPU time (s)	M	MSE (%)
Ω_1	57	1.1	85	146	0.18
Ω_2	44	1.2	56	120	0.18
Ω_3	24	1.4	43	60	0.16
Ω_4	32	1.9	61	66	0.25
Total			245		700



Schwarz Boundary Sampling for All-HROM Coupling



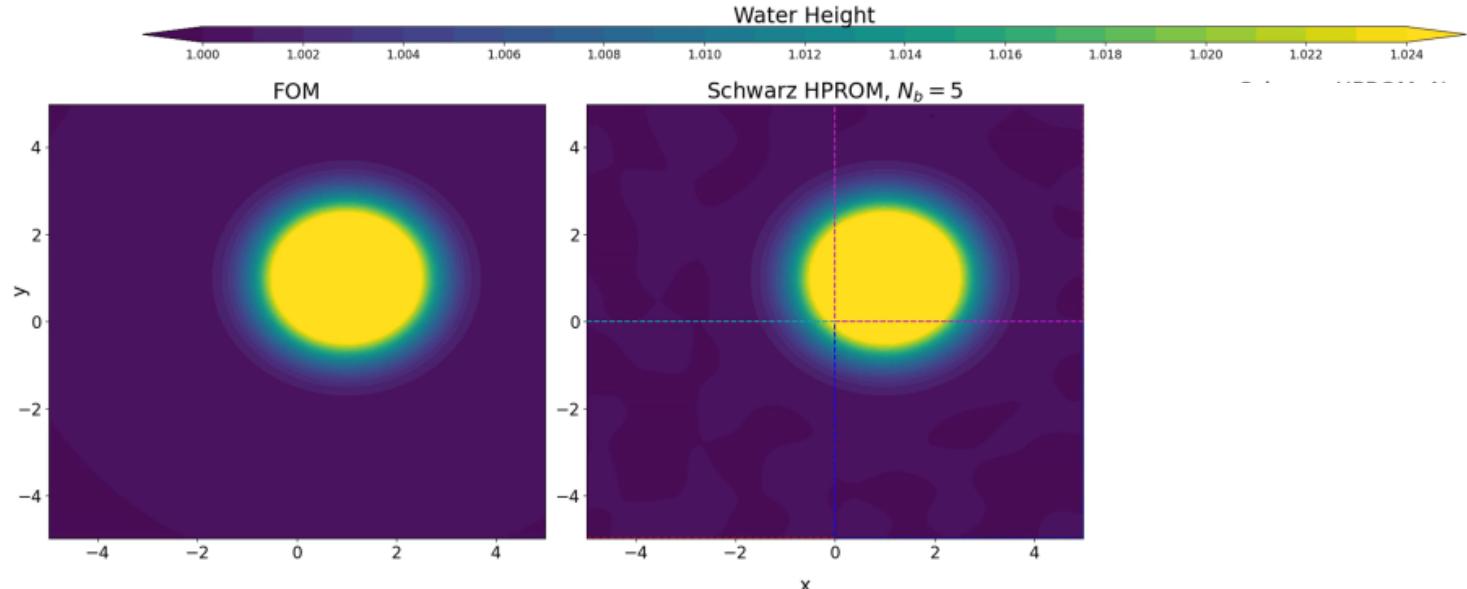
Key question: how many Schwarz boundary points need to be included in sample mesh when performing HROM coupling?

Schwarz Boundary Sampling for All-HROM Coupling



Key question: how many Schwarz boundary points need to be included in sample mesh when performing HROM coupling?

- Naive/sparsely-sampled Schwarz boundary results in **failure** to transmit coupling information during Schwarz



Movie above: FOM (left) and all HROM with $N_b = 5\%$ (right).
 ROMs have $K = 100$ modes and $N_s = 0.5\%N$ sample mesh points.

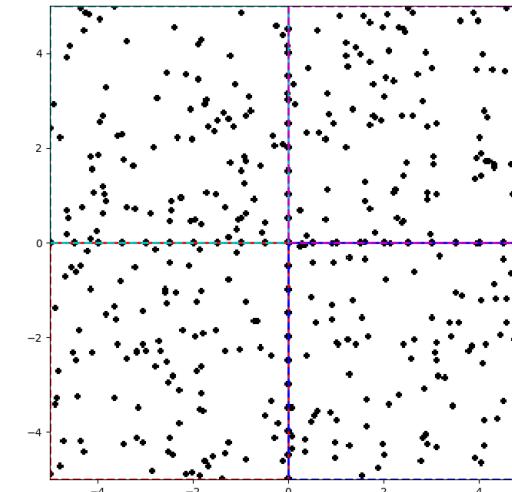


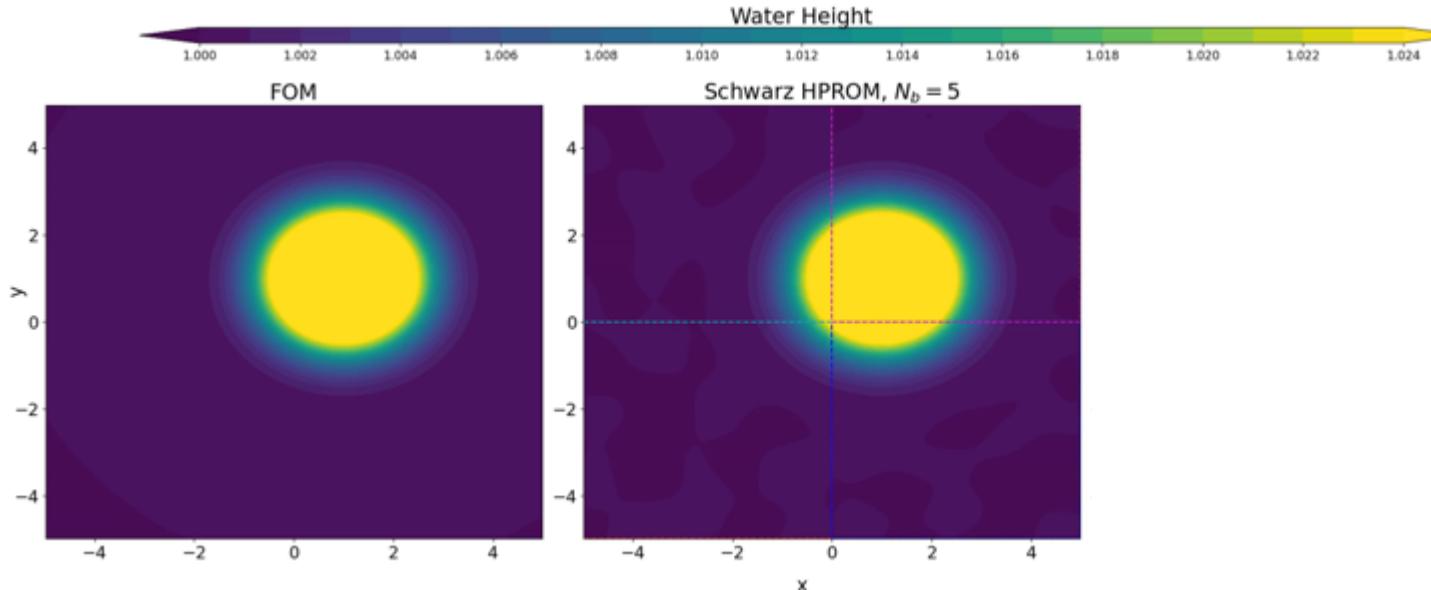
Figure above: example sample mesh with sampling rate $N_b = 5\%$.

Schwarz Boundary Sampling for All-HROM Coupling



Key question: how many Schwarz boundary points need to be included in sample mesh when performing HROM coupling?

- Naive/sparsely-sampled Schwarz boundary results in **failure** to transmit coupling information during Schwarz



Movie above: FOM (left) and all HROM with $N_b = 5\%$ (right).
 ROMs have $K = 100$ modes and $N_s = 0.5\%N$ sample mesh points.

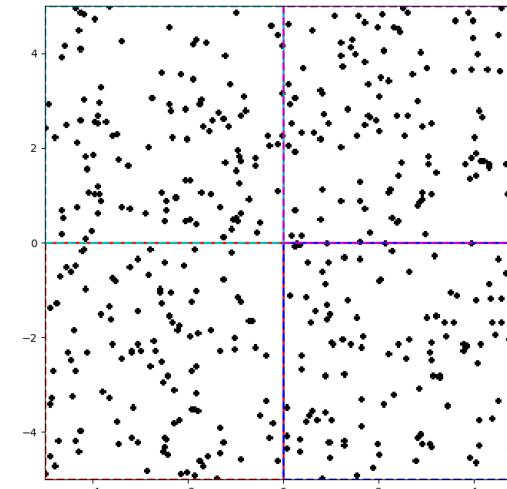


Figure above: example sample mesh with sampling rate $N_b = 0$.

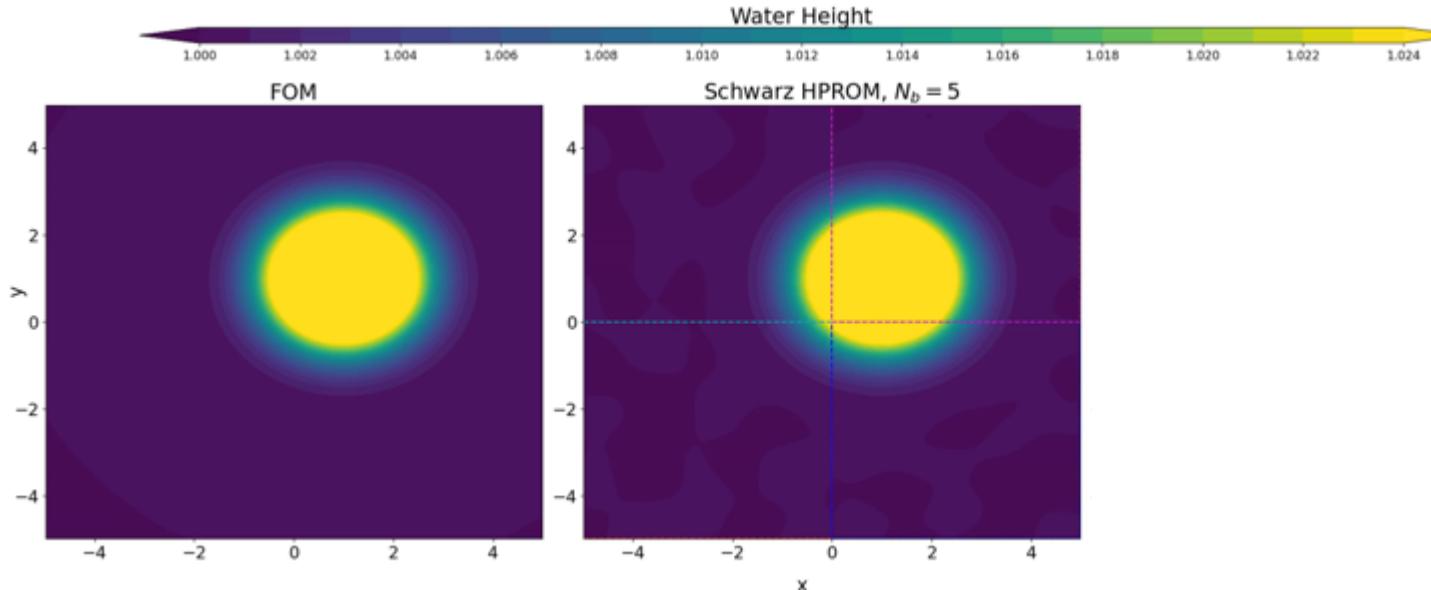
- Including **too many Schwarz boundary points** (N_b) in sample mesh given **fixed budget** of N_s sample mesh points may lead to **too few sample mesh points** in interior

Schwarz Boundary Sampling for All-HROM Coupling



Key question: how many Schwarz boundary points need to be included in sample mesh when performing HROM coupling?

- Naive/sparsely-sampled Schwarz boundary results in **failure** to transmit coupling information during Schwarz



Movie above: FOM (left) and all HROM with $N_b = 5\%$ (right).
ROMs have $K = 100$ modes and $N_s = 0.5\%N$ sample mesh points.

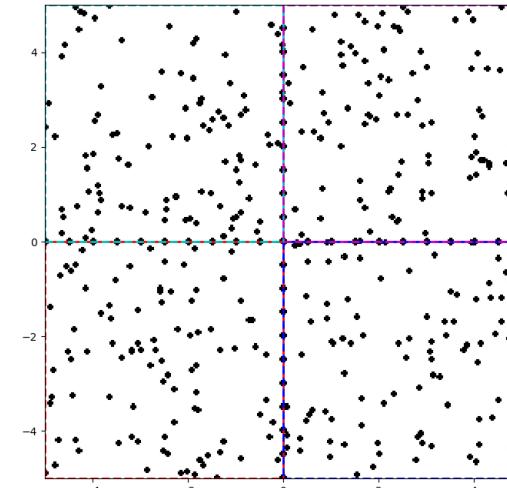


Figure above: example sample mesh with sampling rate $N_b = 5\%$.

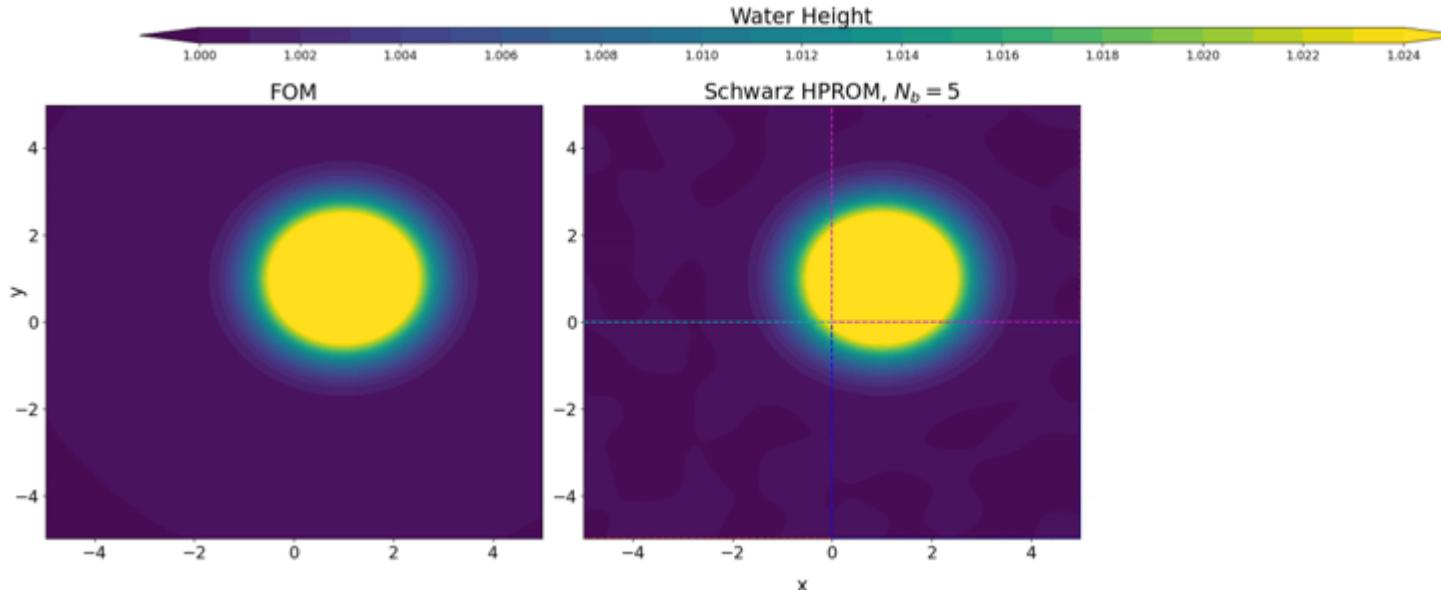
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Schwarz Boundary Sampling for All-HROM Coupling



Key question: how many Schwarz boundary points need to be included in sample mesh when performing HROM coupling?

- Naive/sparsely-sampled Schwarz boundary results in **failure** to transmit coupling information during Schwarz



Movie above: FOM (left) and all HROM with $N_b = 5\%$ (right).
ROMs have $K = 100$ modes and $N_s = 0.5\%N$ sample mesh points.

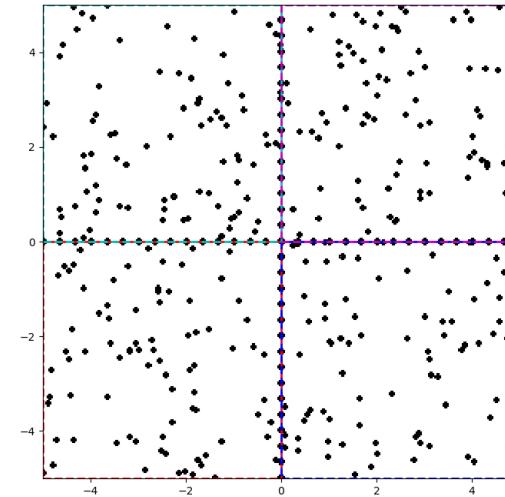


Figure above: example sample mesh with sampling rate $N_b = 10\%$.

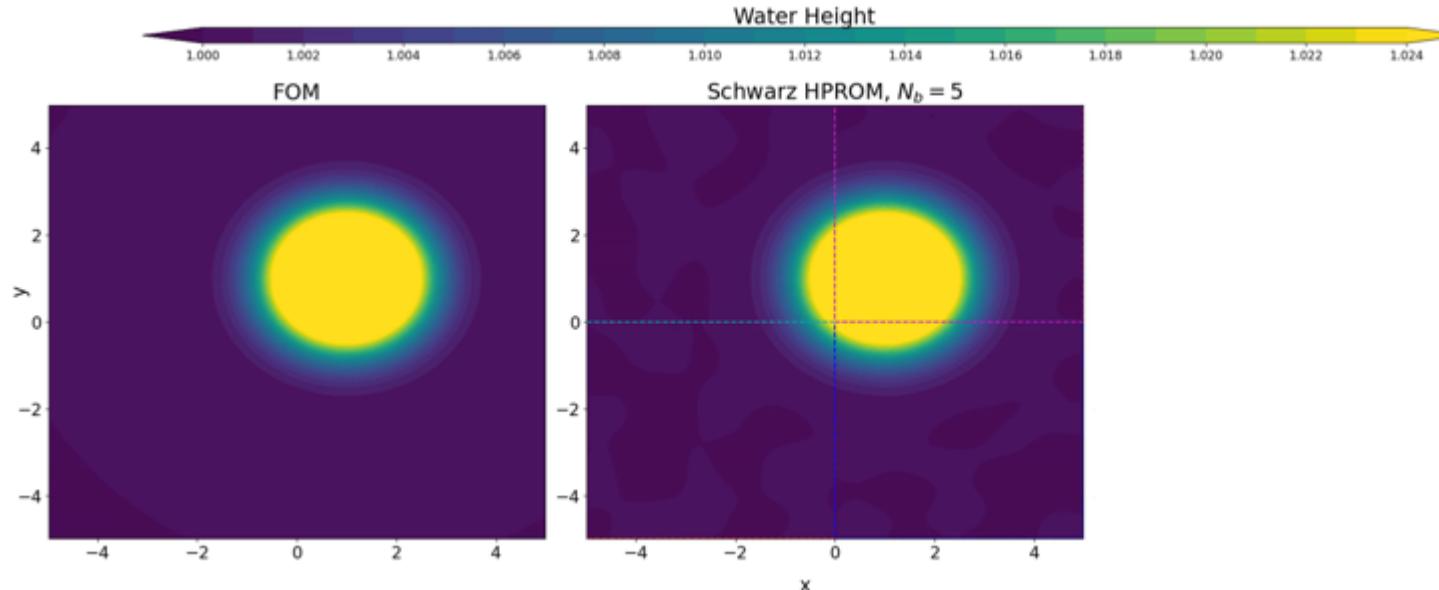
- Including **too many Schwarz boundary points** (N_b) in sample mesh given **fixed budget** of N_s sample mesh points may lead to **too few sample mesh points** in interior

Schwarz Boundary Sampling for All-HROM Coupling



Key question: how many Schwarz boundary points need to be included in sample mesh when performing HROM coupling?

- Naive/sparsely-sampled Schwarz boundary results in **failure** to transmit coupling information during Schwarz



Movie above: FOM (left) and all HROM with $N_b = 5\%$ (right).
ROMs have $K = 100$ modes and $N_s = 0.5\%N$ sample mesh points.

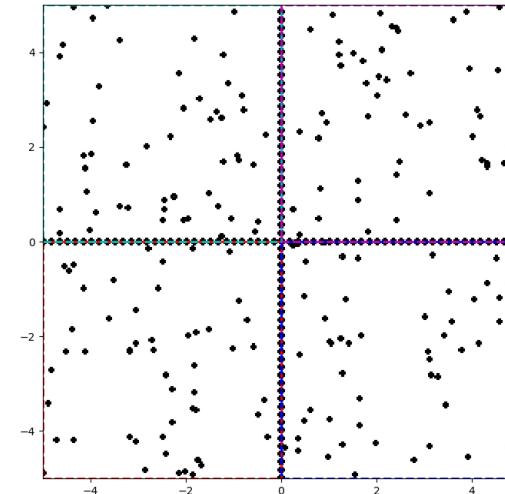


Figure above: example sample mesh with sampling rate $N_b = 15\%$.

- Including **too many Schwarz boundary points** (N_b) in sample mesh given **fixed budget** of N_s sample mesh points may lead to **too few sample mesh points** in interior

$$\frac{\partial}{\partial t} \begin{pmatrix} \rho \\ \rho u \\ \rho v \\ \rho E \end{pmatrix} + \frac{\partial}{\partial x} \begin{pmatrix} \rho u \\ \rho u^2 + p \\ \rho u v \\ (E + p)u \end{pmatrix} + \frac{\partial}{\partial y} \begin{pmatrix} \rho v \\ \rho u v \\ \rho v^2 + p \\ (E + p)v \end{pmatrix} = \mathbf{0}$$

$$p = (\gamma - 1) \left(\rho E - \frac{1}{2} \rho (u^2 + v^2) \right)$$

Problem setup:

- $\Omega = (0,1)^2$, $t \in [0, 0.8]$, homogeneous Neumann BCs
- Fix $\rho_1 = 1.5$, $u_1 = v_1 = 0$, $p_3 = 0.029$
- Vary p_1 ; IC from compatibility conditions*
 - Training: $p_1 \in [1.0, 1.25, 1.5, 1.75, 2.0]$
 - Testing: $p_1 \in [1.125, 1.375, 1.625, 1.875]$

FOM discretization:

- Spatial discretization given by a first-order **cell-centered finite volume** discretization with $N = 300$ or $N = 100$ elements in each dimension
- Implicit first order temporal discretization: **backward Euler** with fixed $\Delta t = 0.005$
- Implemented in **Pressio-demoapps** (<https://github.com/Pressio/pressio-demoapps>)

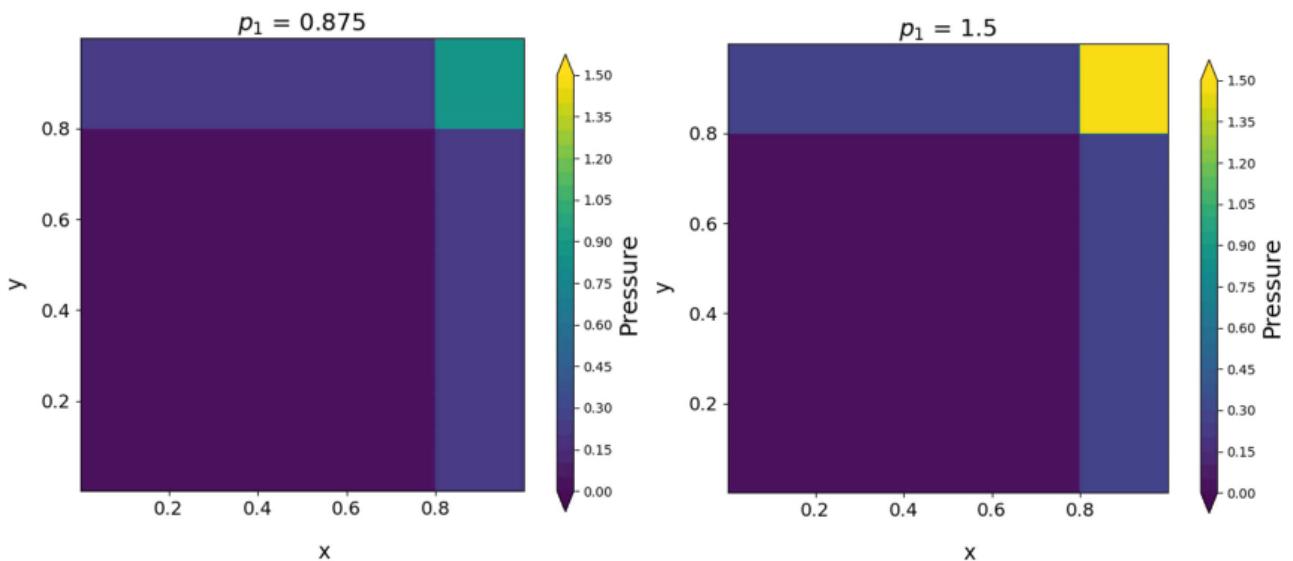


Figure above: FOM solutions to Euler Riemann problem for $p_1 = 0.875$ (left) and $p_1 = 1.5$ (right).

Preliminary results (WIP)

*Schulz-Rinne, 1993.

Schwarz Coupling Details

Choice of domain decomposition

- Overlapping and non-overlapping DD of Ω into 4 subdomains coupled via additive/multiplicative Schwarz
- All-ROM or All-HROM coupling via Pressio*



Snapshot collection and reduced basis construction

- Single-domain FOM on Ω used to generate snapshots/POD modes

Enforcement of boundary conditions (BCs) in ROM at Schwarz boundaries

- BCs are imposed approximately by fictitious ghost cell states
- Dirichlet-Dirichlet BCs for both overlapping and non-overlapping

Choice of hyper-reduction

- Collocation and gappy POD for hyper-reduction
- Assume fixed budget of sample mesh points at Schwarz boundaries

*<https://github.com/Pressio/pressio-demoapps>

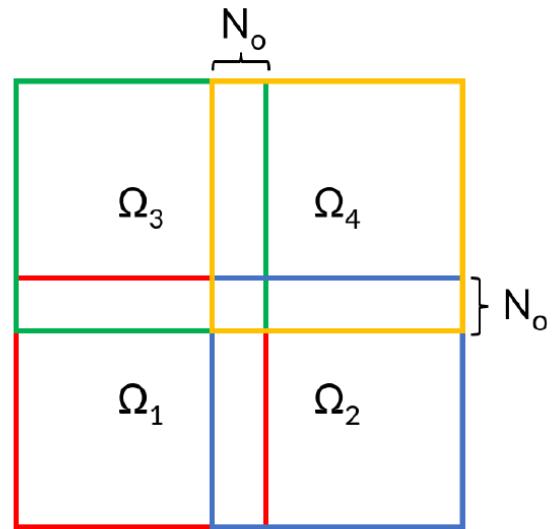


Figure above: DD of Ω into 4 subdomains

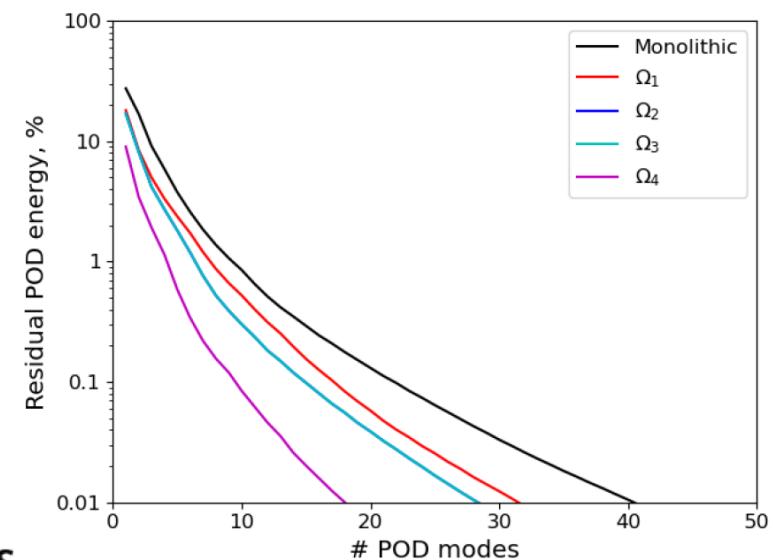
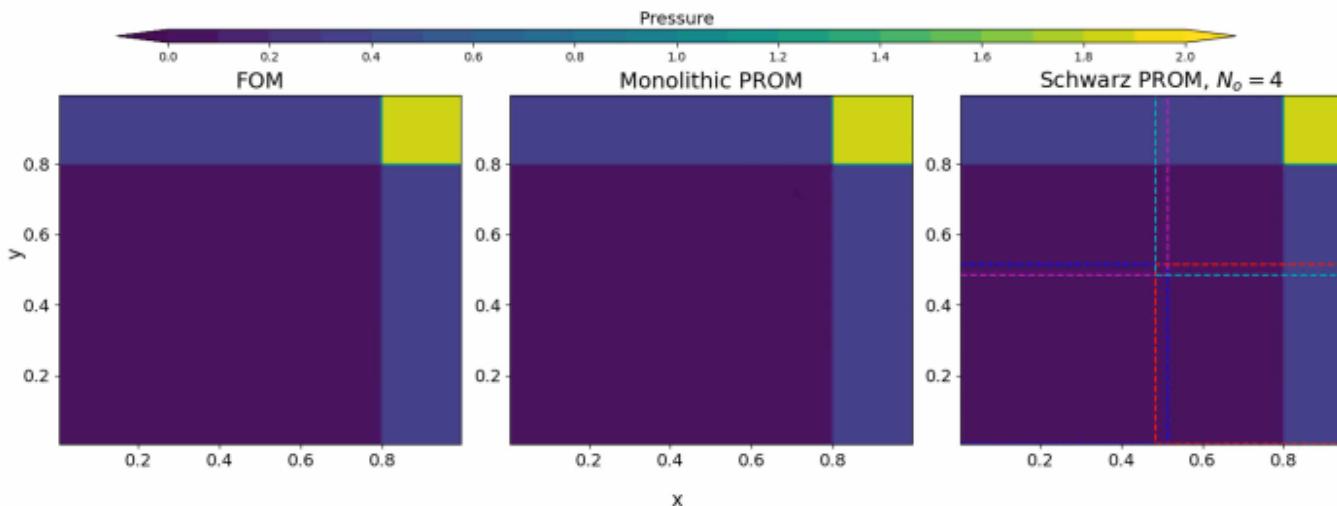


Figure above: Slow decay of POD energy for Euler problem

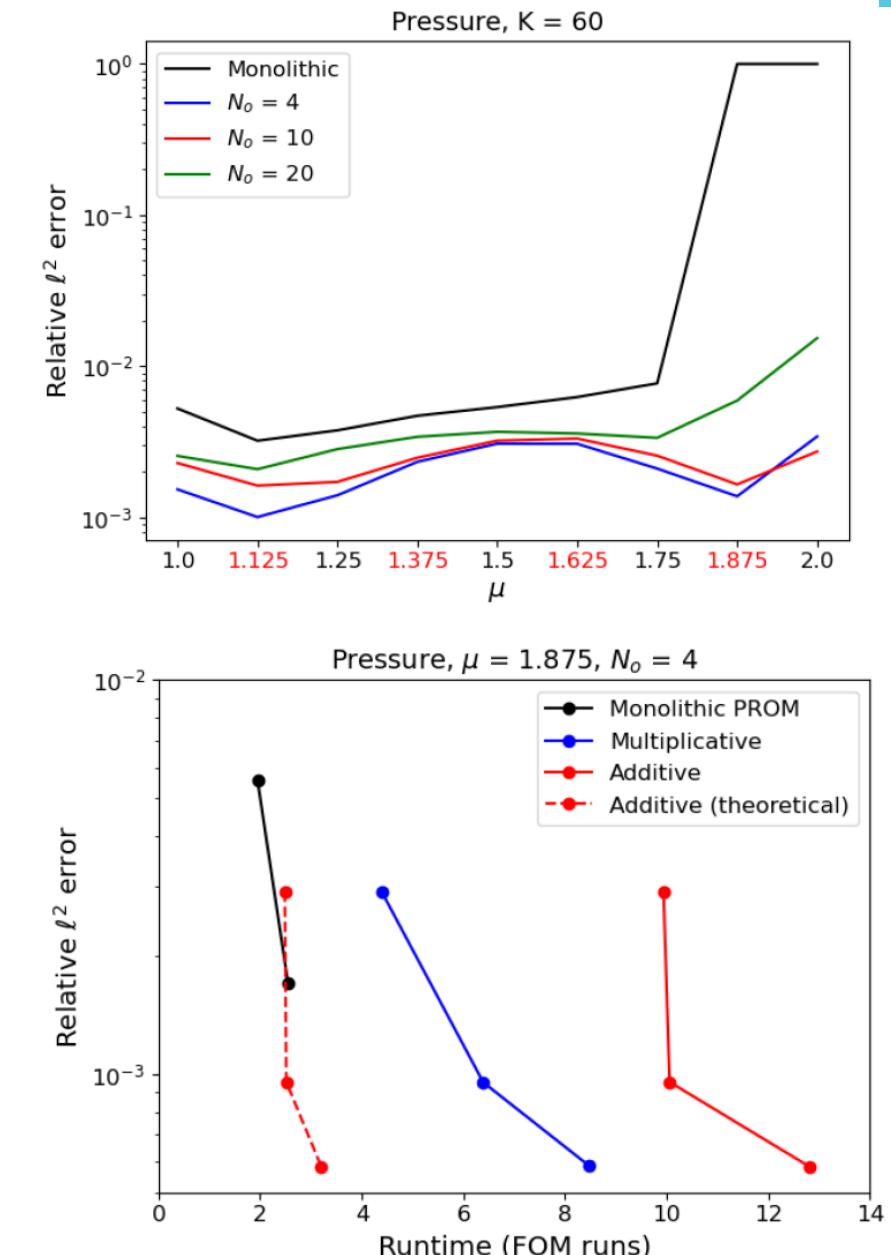
Model Problem 3: All-ROM Coupling + Overlapping Schwarz



- For smaller basis sizes and larger p_1 , monolithic ROM is **unstable** whereas **Schwarz ROM** gives **accurate solution!**
- Increased **overlap** degrades accuracy (top right)
- Shock transmission **error** significantly **increases with overlap**
- ~4.4 **average # Schwarz iterations** with additive Schwarz vs. ~3.6 for multiplicative Schwarz
- With **additive Schwarz**, can achieve **lower error** than monolithic ROM for **same CPU time** (bottom right)



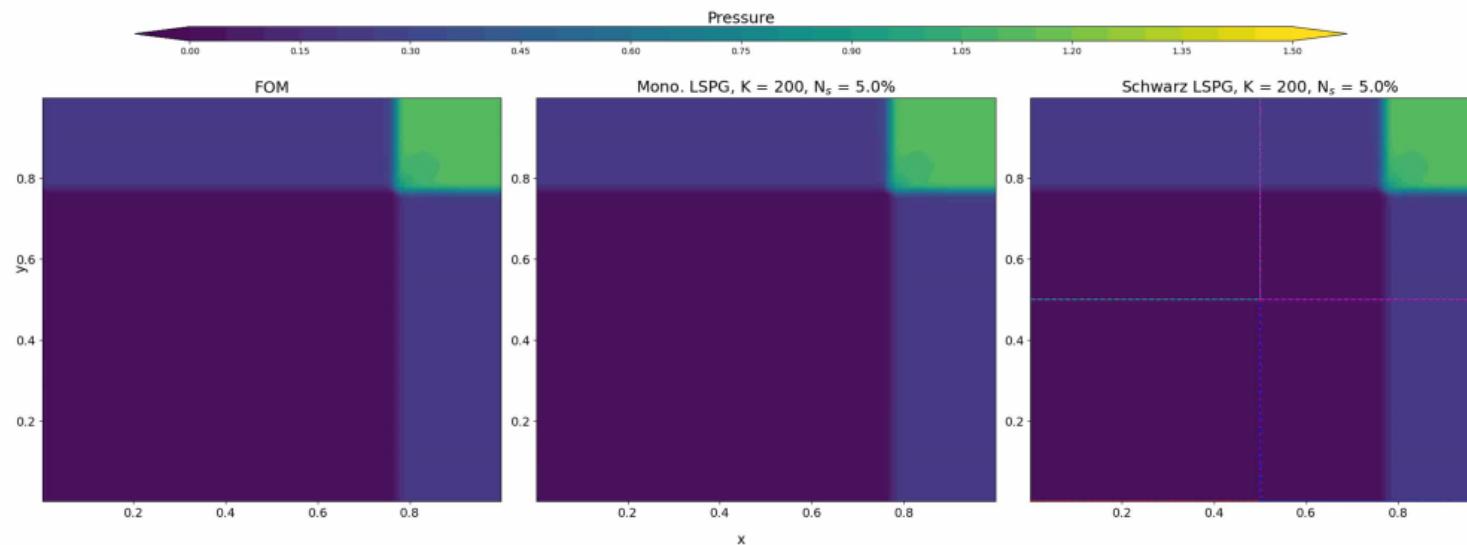
Movie above: FOM (left), $K = 50$ monolithic ROM (middle), and $K = 50$ overlapping Schwarz ROM with $N_o = 4$ (left) for $p_1 = 1.875$.



Model Problem 3: All-HROM Coupling + Non-Overlapping Schwarz



- Hyper-reduction via collocation works better than gappy POD
- Schwarz can give improved accuracy relative to monolithic ROM
- Achieving cost-savings w.r.t. monolithic FOM is WIP



Movie above: FOM (left), HROM (middle) and Schwarz All-HROM (right) solution. HROMs have 5% sampling rate and 200 POD modes.

Preliminary results (WIP)

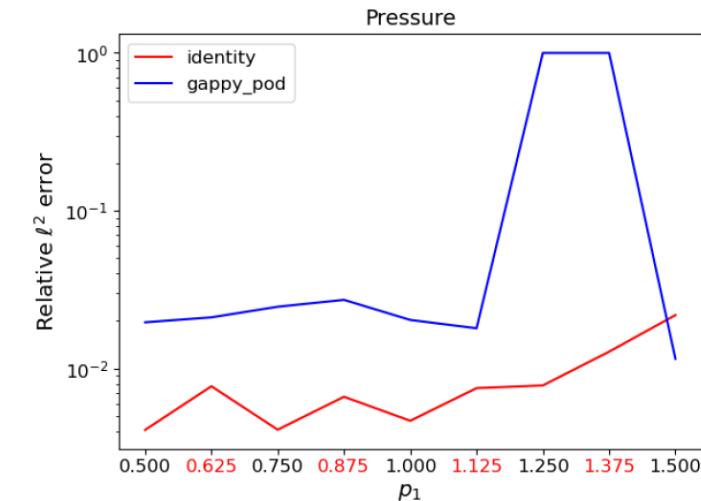


Figure above: collocation and gappy POD relative errors for K=200, 1% sampling rate.

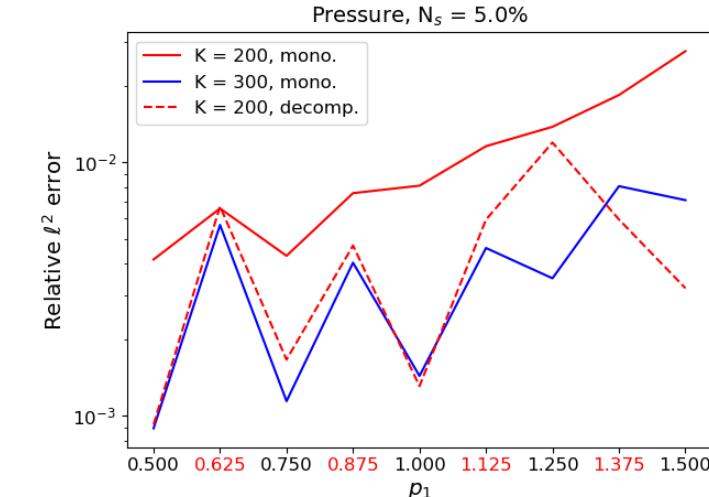


Figure above: monolithic vs. decomposed HROM errors with 5% sampling rate no overlap.

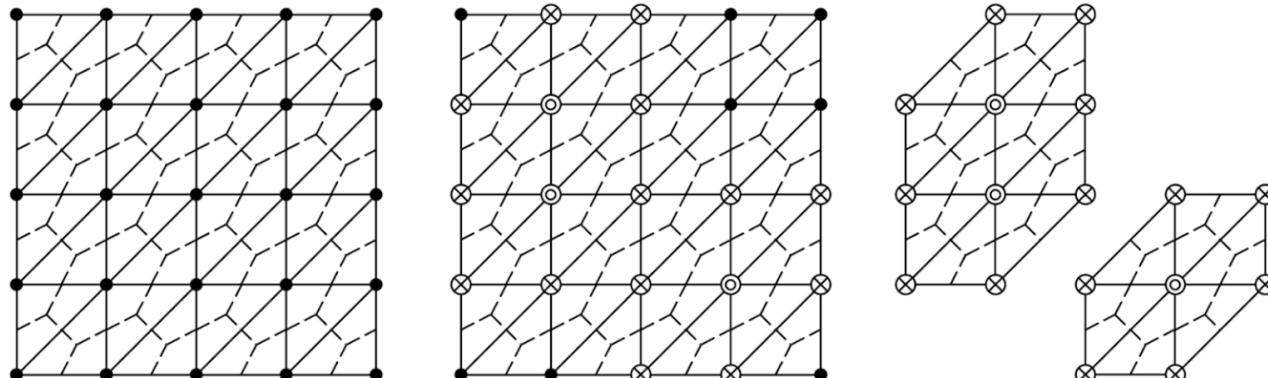
Energy-Conserving Sampling and Weighting (ECSW)



- Project-then-approximate paradigm (as opposed to approximate-then-project)

$$\begin{aligned}
 r_k(q_k, t) &= W^T r(\tilde{u}, t) \\
 &= \sum_{e \in \mathcal{E}} W^T L_e^T r_e(L_e + \tilde{u}, t)
 \end{aligned}$$

- $L_e \in \{0,1\}^{d_e \times N}$ where d_e is the **number of degrees of freedom** associated with each mesh element (this is in the context of meshes used in first-order hyperbolic problems where there are N_e mesh elements)
- $L_e^+ \in \{0,1\}^{d_e \times N}$ selects degrees of freedom necessary for **flux reconstruction**
- Equality can be **relaxed**



Augmented reduced mesh: \circledcirc represents a selected node attached to a selected element; and \otimes represents an added node to enable the full representation of the computational stencil at the selected node/element

ECSW: Generating the Reduced Mesh and Weights



- Using a subset of the same snapshots $u_i, i \in 1, \dots, n_h$ used to generate the **state basis** V , we can train the reduced mesh
- Snapshots are first **projected** onto their associated basis and then **reconstructed**

$$c_{se} = W^T L_e^T r_e \left(L_e^+ \left(u_{ref} + V V^T (u_s - u_{ref}) \right), t \right) \in \mathbb{R}^n$$

$$d_s = r_k(\tilde{u}, t) \in \mathbb{R}^n, \quad s = 1, \dots, n_h$$

- We can then form the **system**

$$\mathbf{C} = \begin{pmatrix} c_{11} & \dots & c_{1N_e} \\ \vdots & \ddots & \vdots \\ c_{n_h 1} & \dots & c_{n_h N_e} \end{pmatrix}, \quad \mathbf{d} = \begin{pmatrix} d_1 \\ \vdots \\ d_{n_h} \end{pmatrix}$$

- Where $\mathbf{C}\xi = \mathbf{d}, \xi \in \mathbb{R}^{N_e}$, $\xi = \mathbf{1}$ must be the solution
- Further relax the equality to yield **non-negative least-squares problem**:

$$\xi = \arg \min_{x \in \mathbb{R}^n} \|\mathbf{C}x - \mathbf{d}\|_2 \text{ subject to } x \geq \mathbf{0}$$

- Solve the above optimization problem using a **non-negative least squares solver** with an **early termination condition** to promote **sparsity** of the vector ξ