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SAND2023-12312C

On the Harmonic Balance Method Augmented with Non-Smooth Basis Functions for Contact/Impact Problems

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IMAC-XLI (41), submission #14374

February 13-16, 2023

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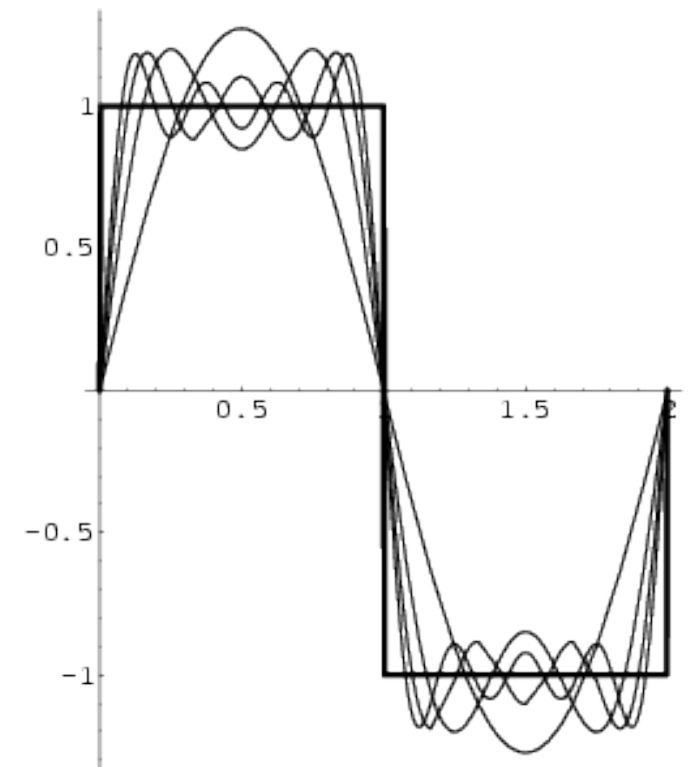
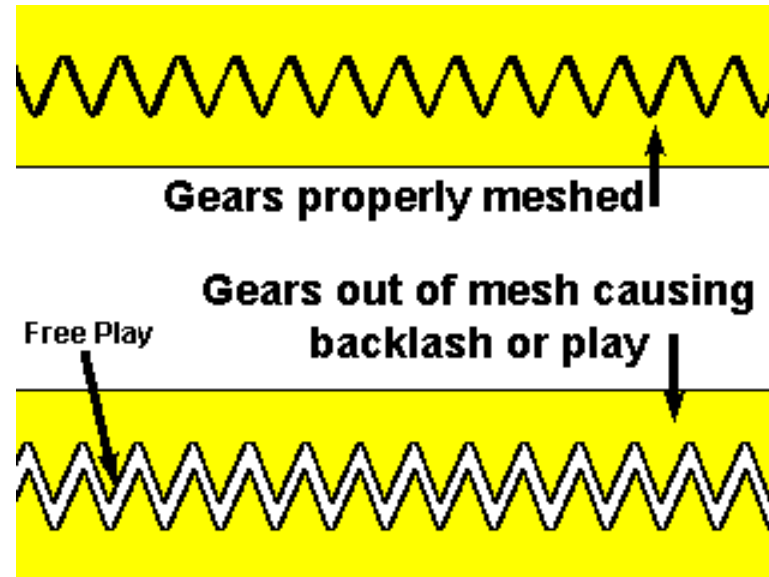
Presentation outline

- ☐ **Introduction and motivation**
- ☐ **Non-smooth Galerkin formulation**
- ☐ **Least-squares regression analysis**
- ☐ **Additional studies**
- ☐ **Conclusions**

Introduction and motivation

□ Recent developments have efficiently applied the harmonic balance method to strongly nonlinear systems

- Systems include aircraft, spacecraft, gear drives, bladed disks, etc.
- Applications include continuation procedures, bifurcation detection and tracking, nonlinear modal analysis, etc.
- Advantages include reduced computational costs and capture of unstable solutions
- Difficulties include chaotic/aperiodic responses and non-smooth or discontinuous solutions (Gibbs phenomenon)
- How can we simulate “stiff” contact or friction systems efficiently with HBM?



Introduction and motivation

❑ Review of the literature—simulation approaches include:

- Lanczos filtering to improve the values of the Fourier coefficients
- Append additional, non-smooth or discontinuous terms to a system's solution
- Replace some or all the terms in classical Fourier series with non-smooth terms
- Utilize event-driven schemes to find and integrate between the state transition times to compute nonlinear forces
- Non-smooth temporal and spatial transformations

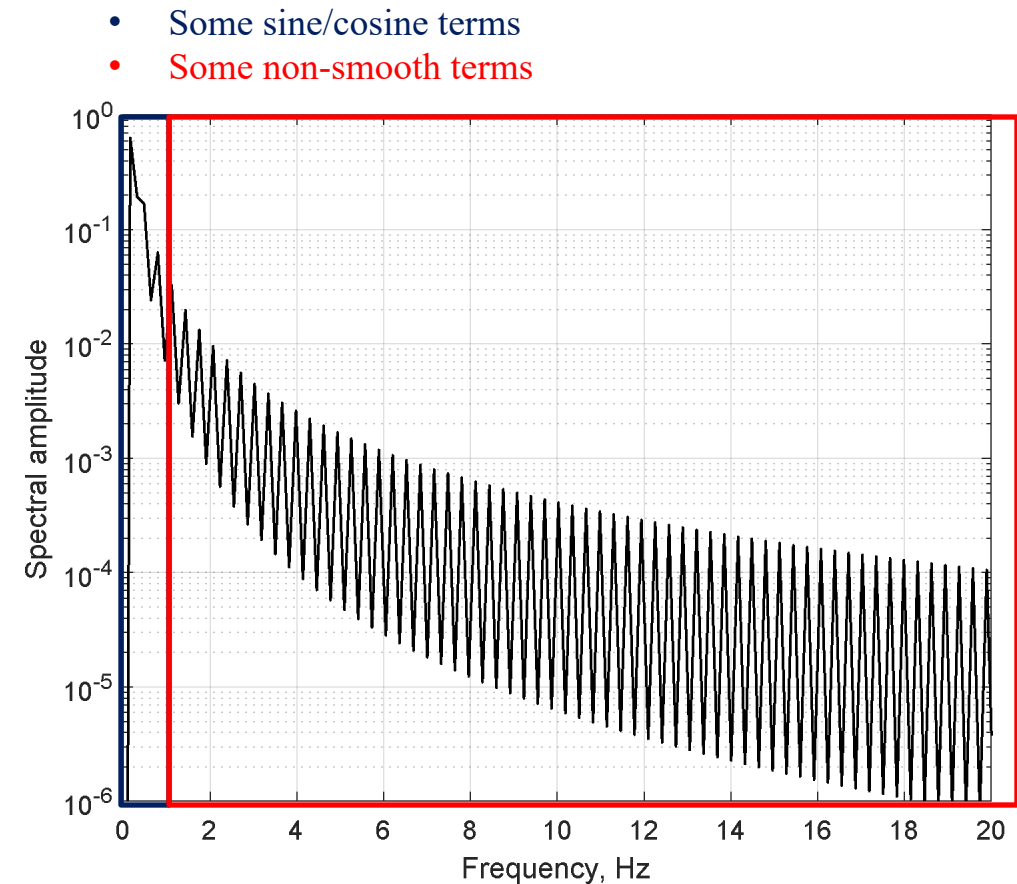
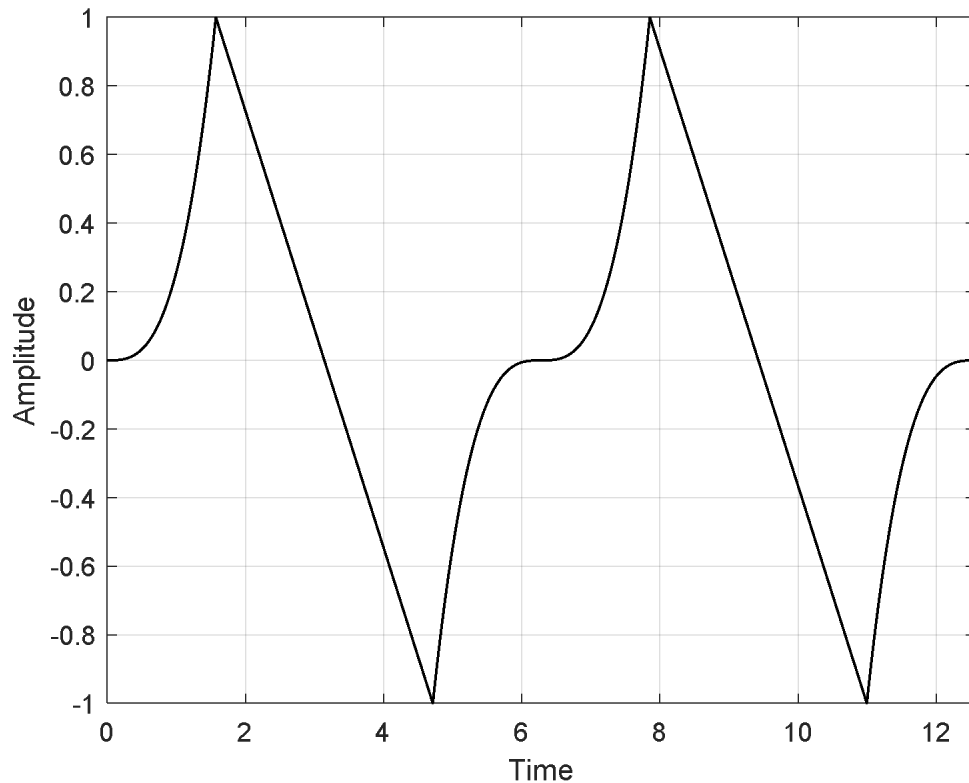
❑ Difficulties:

- Gibbs phenomenon: slow convergence (polynomial) compared to smooth systems (exponential)
- A priori knowledge of the state transition times may be required

Introduction and motivation

❑ Motivating idea:

- Non-smooth periodic motions can be represented by infinite Fourier series
- Can we append the Fourier series representation with entire non-smooth basis functions?
- This approach may allow us to capture a large set of harmonics with a small number of non-smooth functions



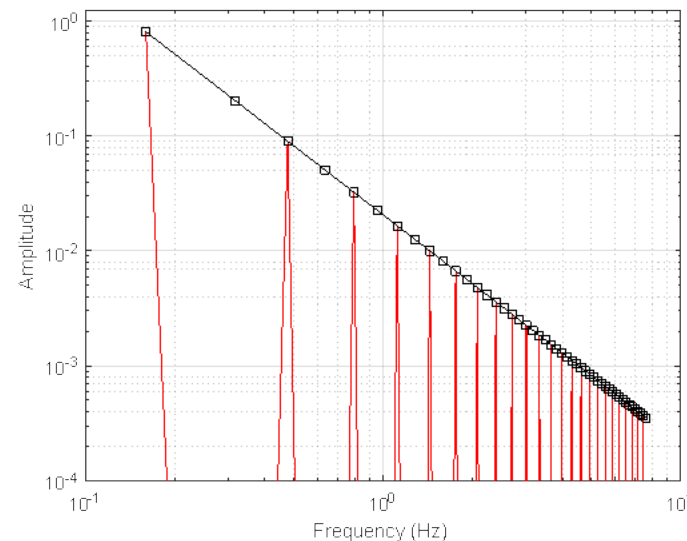
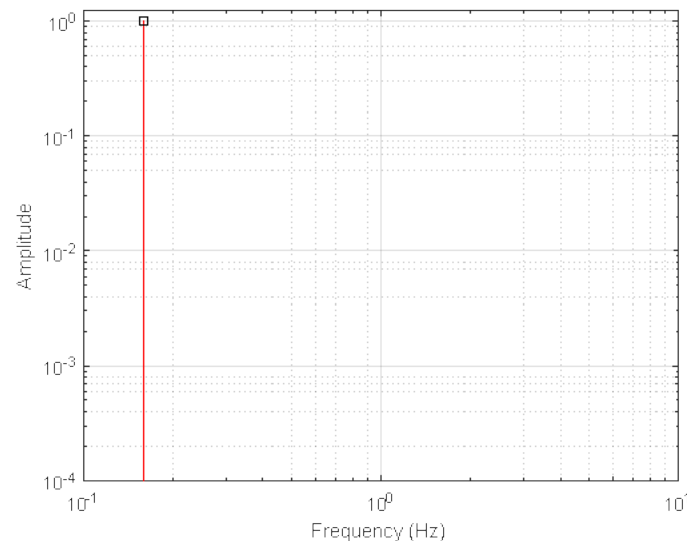
Introduction and motivation

❑ Desirable basis function traits:

- Easy to mathematically represent
- Convenient mathematical properties (Fourier series, derivatives, etc.)
- Intuitive, non-smooth counterparts to sine and cosine functions

❑ Goals of this work:

- Select functions with C^0 smoothness for capturing a non-smooth functional representation of a solution
- Use goodness-of-fit metrics to evaluate the classical Fourier series and the non-smooth basis functions
- Develop a framework that can later be implemented into harmonic balance formulations



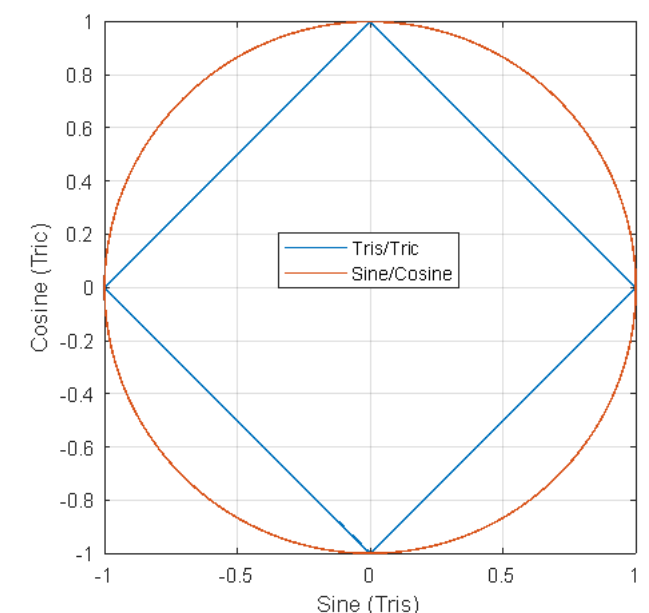
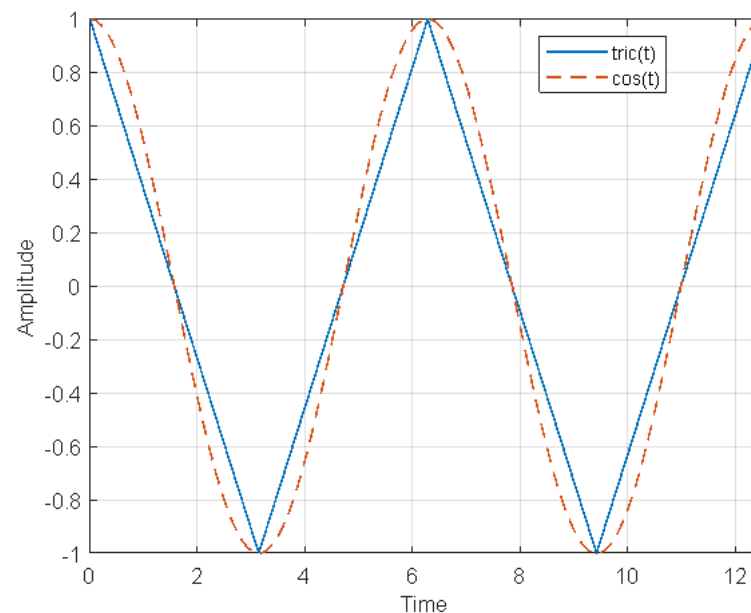
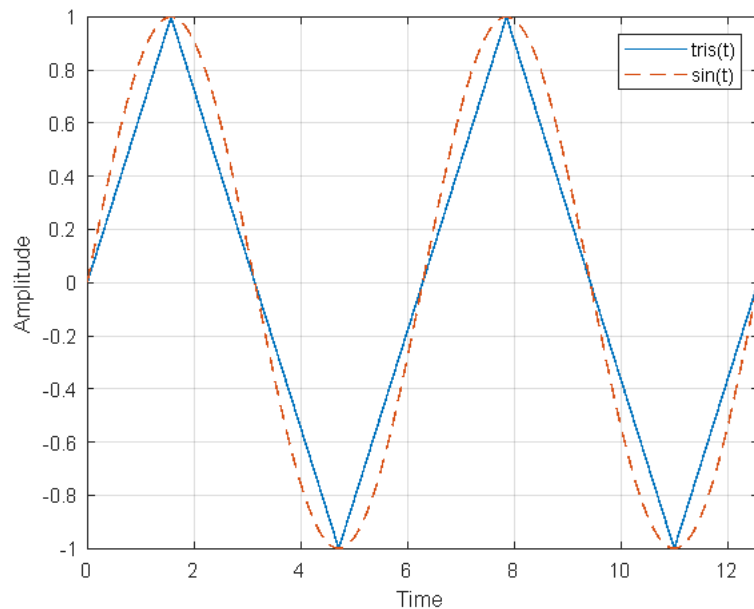
Non-smooth Galerkin formulation

□ Non-smooth triangle waves are chosen for study

- “Triangle sine” and “triangle cosine”
- The same periodicity, maxima, minima, and roots as sine and cosine waves

$$tris(\omega t) = \begin{cases} \frac{4}{T} t_m, & t_m < \frac{T}{4} \\ -\frac{4}{T} t_m + 2, & \frac{T}{4} \leq t_m \leq \frac{3T}{4} \\ \frac{4}{T} t_m - 4, & t_m > \frac{3T}{4} \end{cases}, \quad tric(\omega t) = \begin{cases} -\frac{4}{T} t_m + 1, & t_m \leq \frac{T}{2} \\ \frac{4}{T} t_m - 3, & t_m > \frac{T}{2} \end{cases},$$

$$t_m = \omega t (\text{mod } T), \\ T = 2\pi$$



Non-smooth Galerkin formulation

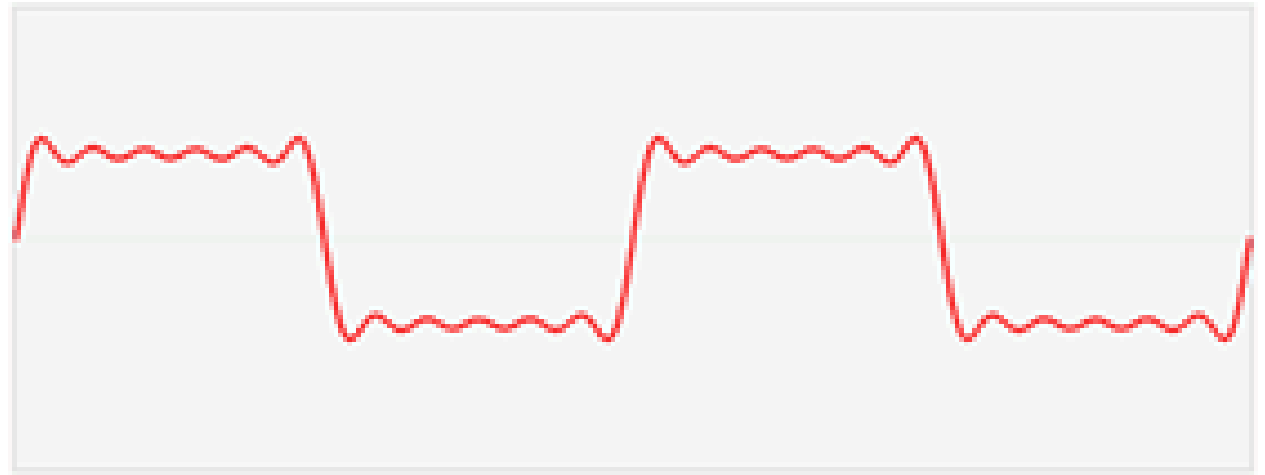
- ❑ **Infinite series representation:**

$$f(t) = a_0 + \sum_{n=1}^{\infty} a_n \text{tric}(n\omega t) + b_n \text{tris}(n\omega t)$$

- ❑ **A different form for numerical convenience**

$$f(t) = a_0 + \sum_{n=1}^{\infty} a_n \text{tris}(n\omega t + \theta_n)$$

- ❑ **Now we can study these functions for desirable properties and advantages**



Mathematical properties

❑ The non-smooth functions have similar properties to sine/cosine:

➤ Reflections:

θ reflected in $\alpha = 0$	θ reflected in $\alpha = \frac{\pi}{4}$	θ reflected in $\alpha = \frac{\pi}{2}$	θ reflected in $\alpha = \frac{3\pi}{4}$	θ reflected in $\alpha = \pi$
$\text{tris}(-\theta) = -\text{tris}(\theta)$	$\text{tris}\left(\frac{\pi}{2} - \theta\right) = \text{tric}(\theta)$	$\text{tris}(\pi - \theta) = +\text{tris}(\theta)$	$\text{tris}\left(\frac{3\pi}{2} - \theta\right) = -\text{tric}(\theta)$	$\text{tris}(2\pi - \theta) = -\text{tris}(\theta)$
$\text{tric}(-\theta) = +\text{tric}(\theta)$	$\text{tric}\left(\frac{\pi}{2} - \theta\right) = \text{tris}(\theta)$	$\text{tric}(\pi - \theta) = -\text{tric}(\theta)$	$\text{tric}\left(\frac{3\pi}{2} - \theta\right) = -\text{tris}(\theta)$	$\text{tric}(2\pi - \theta) = +\text{tric}(\theta)$

➤ Shifts and periodicity:

Shift by one quarter period	Shift by one half period	Shift by three quarter periods	Shift by full periods
$\text{tris}\left(\theta \pm \frac{\pi}{2}\right) = \pm \text{tric}(\theta)$	$\text{tris}(\theta \pm \pi) = -\text{tris}(\theta)$	$\text{tris}\left(\theta \pm \frac{3\pi}{2}\right) = \mp \text{tric}(\theta)$	$\text{tris}(\theta \pm 2\pi k) = +\text{tris}(\theta)$
$\text{tric}\left(\theta \pm \frac{\pi}{2}\right) = \mp \text{tris}(\theta)$	$\text{tric}(\theta \pm \pi) = -\text{tric}(\theta)$	$\text{tric}\left(\theta \pm \frac{3\pi}{2}\right) = \pm \text{tris}(\theta)$	$\text{tric}(\theta \pm 2\pi k) = +\text{tric}(\theta)$

Miscellaneous:

- $|\text{tris}(\theta)| + |\text{tric}(\theta)| = 1$
- Angle sums/phase shifts: not straightforward

Mathematical properties

□ Fourier series representations (ω integer):

$$tris(\omega t) = \sum_{n=1}^{\infty} b_n \sin(nt), \quad b_n = \frac{4\omega}{n^2\pi^2} \sum_{k=1}^{2\omega} (-1)^{k+1} \sin\left(\frac{(2k-1)n\pi}{2\omega}\right)$$

$$tric(\omega t) = a_0 + \sum_{n=1}^{\infty} a_n \cos(nt), \quad a_0 = 0, \quad a_n = \frac{4\omega}{n^2\pi^2} \sum_{k=1}^{2\omega-1} (-1)^k \cos\left(\frac{kn\pi}{\omega}\right)$$

➤ Examples:

$$tris(t) = \sum_{n=1}^{\infty} b_n \sin(nt), \quad b_n = \frac{8}{n^2\pi^2} \sin\left(\frac{n\pi}{2}\right) = \frac{8}{\pi^2}, 0, \frac{-8}{9\pi^2}, 0, \frac{8}{25\pi^2}, 0, \frac{-8}{49\pi^2}, \dots$$

$$tric(t) = \sum_{n=1}^{\infty} a_n \cos(nt), \quad a_n = \frac{4}{n^2\pi^2} [1 - \cos(n\pi)] = \frac{8}{\pi^2}, 0, \frac{8}{9\pi^2}, 0, \frac{8}{25\pi^2}, 0, \frac{8}{49\pi^2}, \dots$$

➤ First derivatives: square waves with coefficients $\propto \frac{8\omega}{n\pi^2}$

➤ Second derivatives: Dirac combs with coefficients $\propto \frac{8\omega}{\pi^2}$

Mathematical properties

□ The functions numerically satisfy the following *orthogonality* relationships:

$$\int_0^{2\pi} \text{tris}(mx) dx = 0,$$

$$\int_0^{2\pi} \text{tric}(mx) dx = 0,$$

$$\int_0^{2\pi} \text{tric}(mx) \text{tric}(nx) dx = c_m \delta_{mn},$$

$$\int_0^{2\pi} \text{tris}(mx) \text{tris}(nx) dx = d_m \delta_{mn},$$

$$\int_0^{2\pi} \text{tric}(mx) \text{tris}(nx) dx = 0,$$

➤ For integers m, n and constants c_m, d_m

➤ δ_{mn} denotes Kronecker delta

➤ Result: biorthogonal system, just like sine and cosine

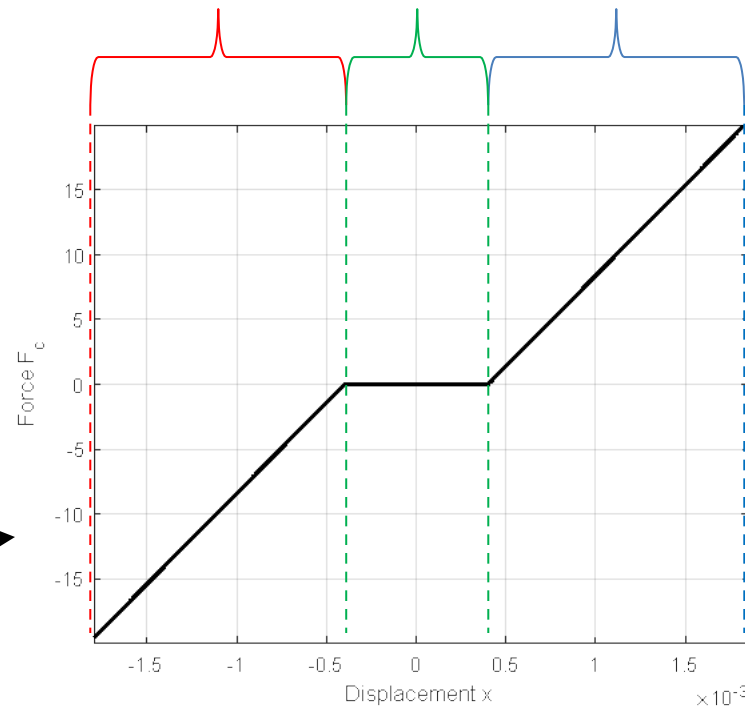
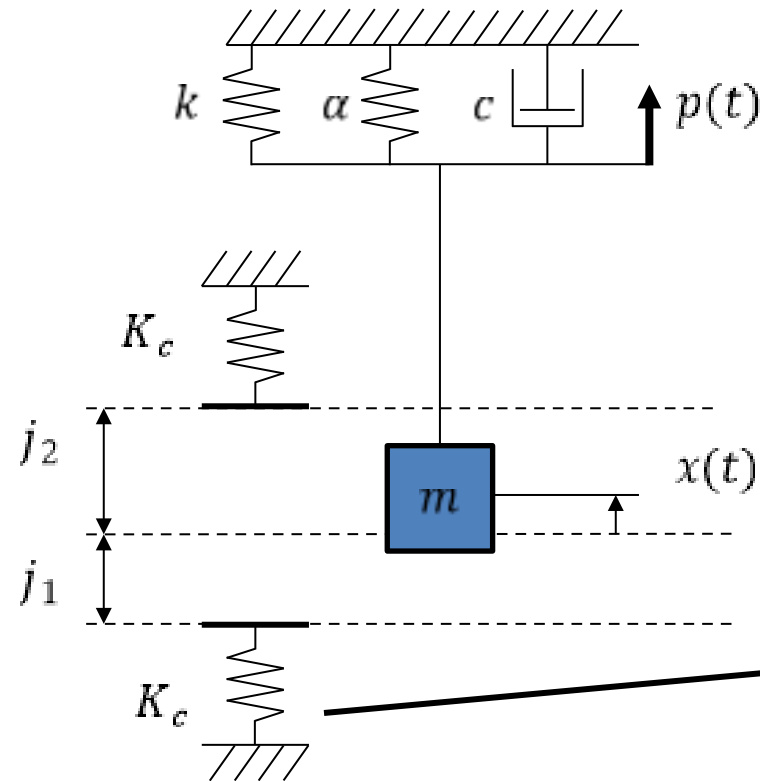
➤ A functional infinite series representation, similar to Fourier series, is possible

➤ Reiterate: no rigorous proof of this

Least-squares regression analysis

□ The system: a forced Duffing oscillator with freeplay

$$\ddot{x} + 2\omega_n\zeta\dot{x} + \omega_n^2x + \frac{\alpha}{m}x^3 + \frac{F_c(x)}{m} = \frac{p}{m}\cos(\omega t),$$



➤ Two different contact laws—contact penalty stiffness, and elastic impact:

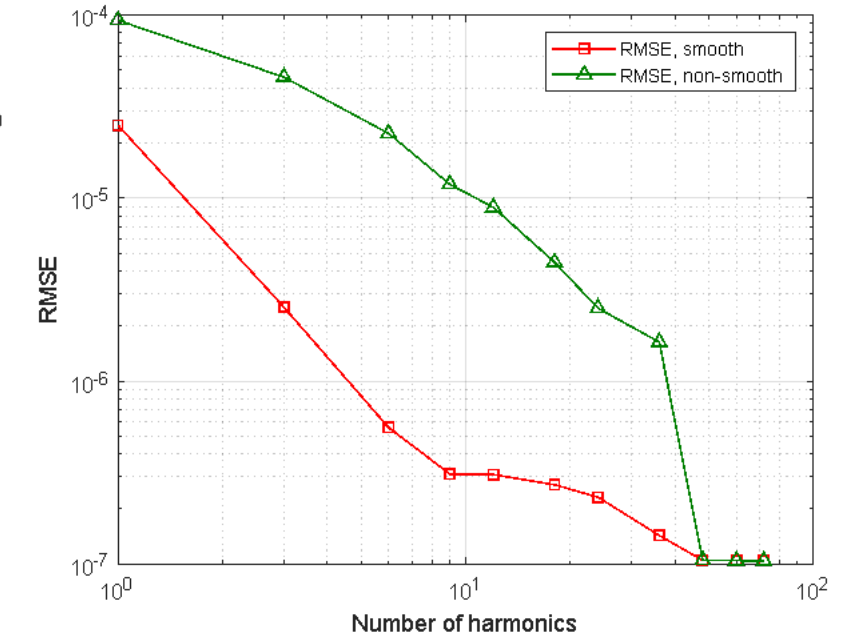
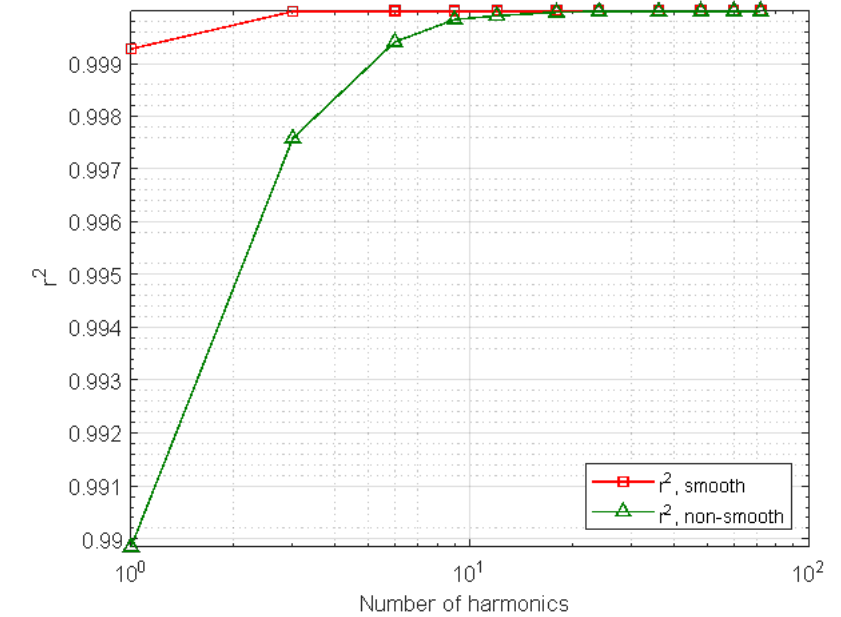
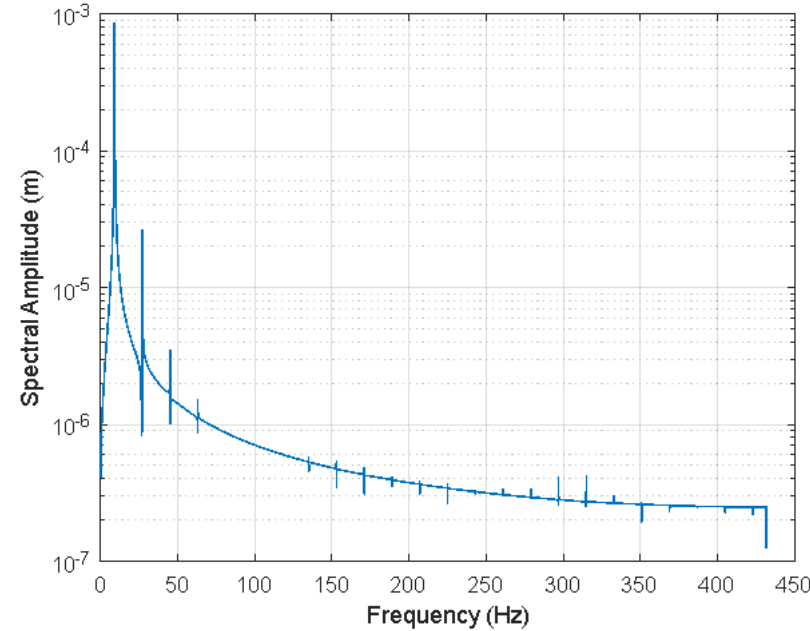
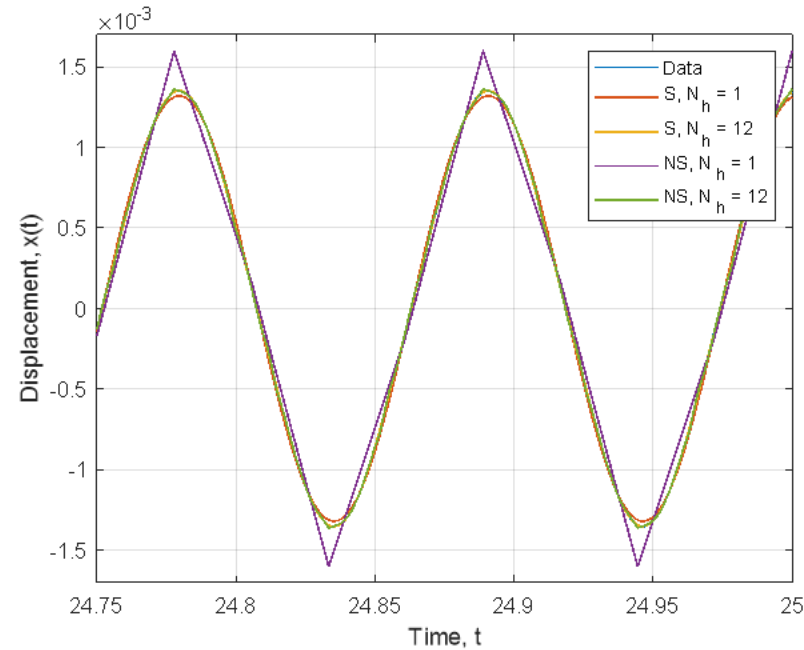
$$F_c = \begin{cases} K_c(x + j_1), & x < -j_1 \\ 0, & -j_1 \leq x \leq j_2 \\ K_c(x - j_2), & x > j_2 \end{cases}$$

vs.

$$-j_1 \leq x(t) \leq j_2, \\ x^+ = -rx^-, r = 1$$

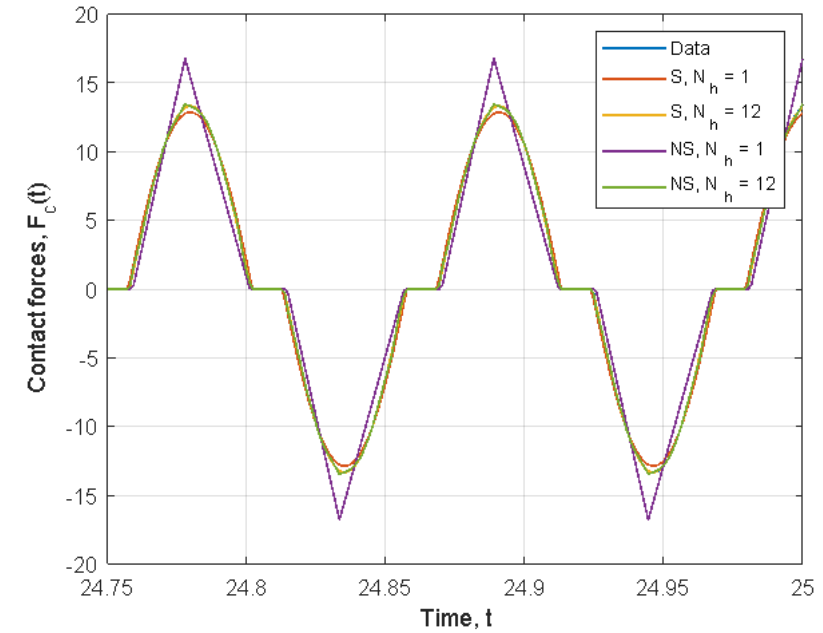
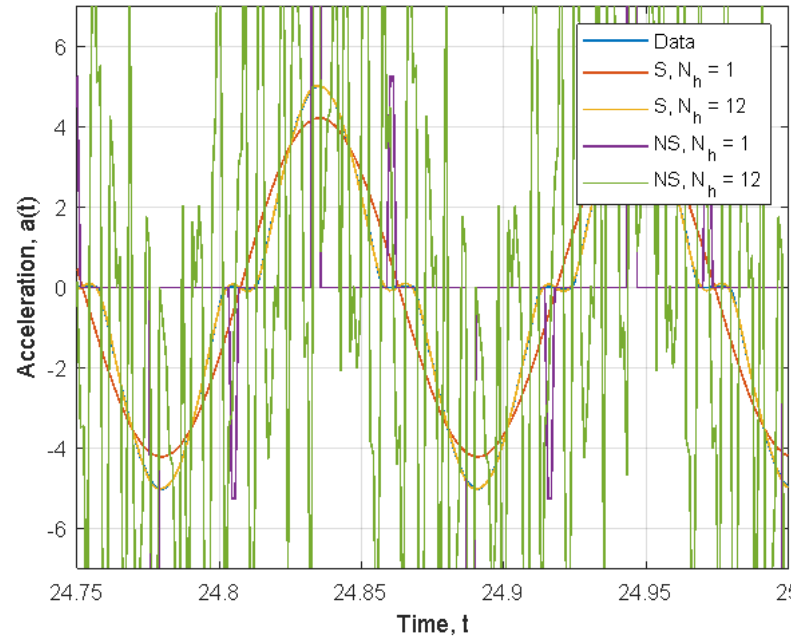
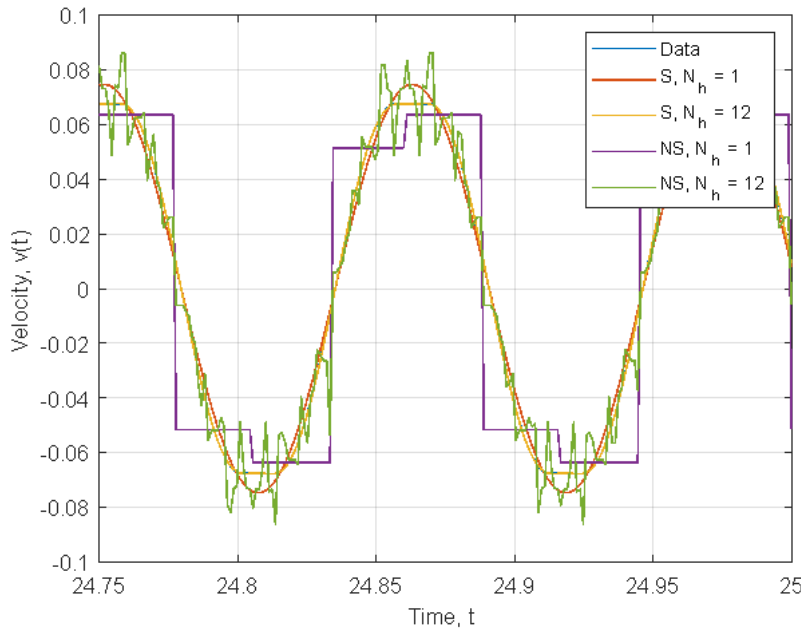
Least-squares regression analysis

Medium stiffness, $K_c = 1.4 * 10^4 \frac{N}{m}$, $\omega = 9 \text{ Hz}$



- Smooth system response
- Classical Fourier series converges much faster than non-smooth Fourier series

Least-squares regression analysis



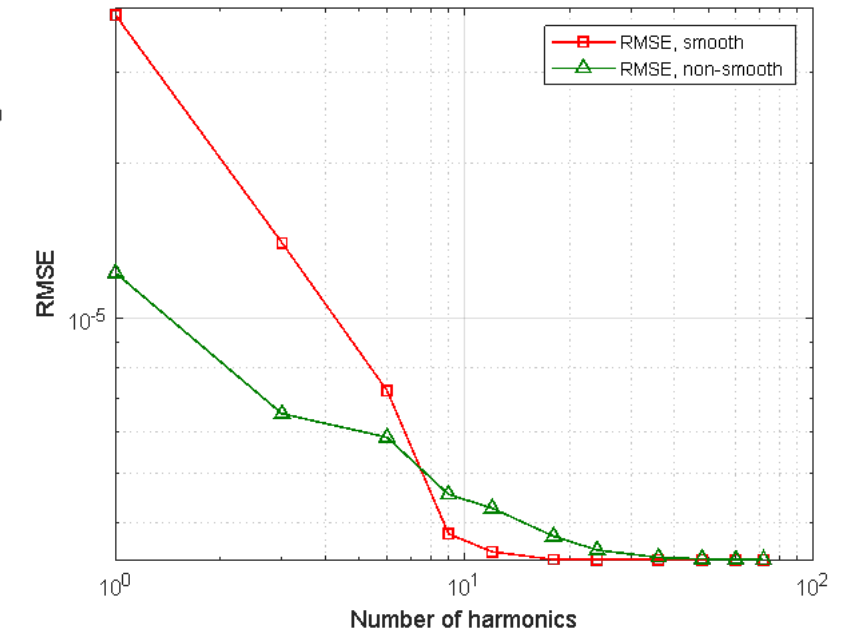
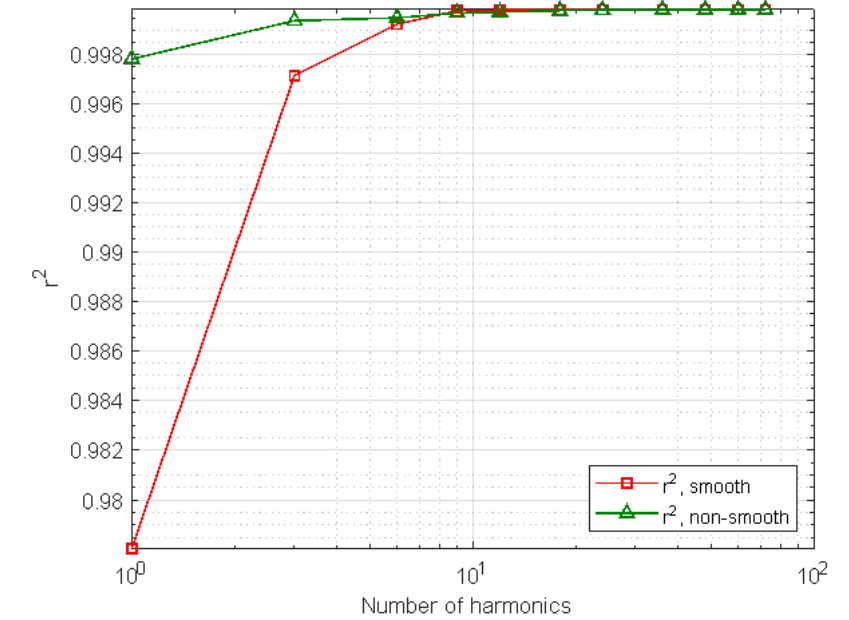
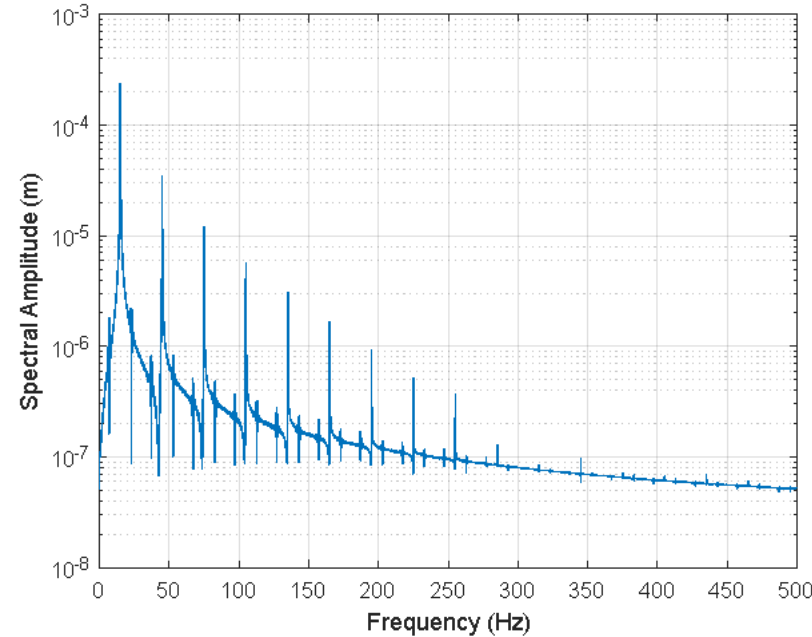
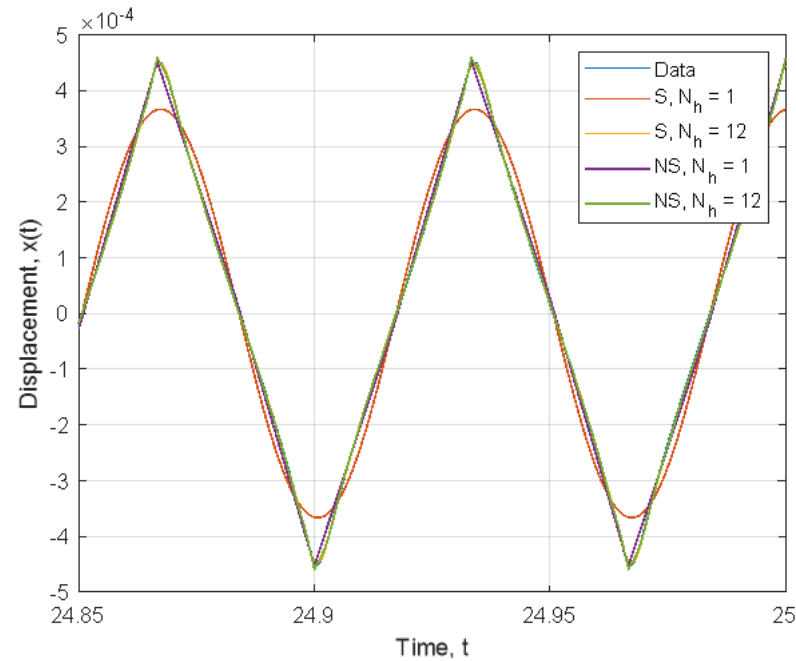
- Now let's look at the velocity, acceleration, and contact forces obtained using low- and high-quality fits
- Both smooth and non-smooth velocities use the coefficients obtained from curve-fitting the displacement:

$$x(t) = a_0 + \sum_{n=1}^{\infty} a_n \text{tric}(n\omega t) + b_n \text{tris}(n\omega t) \Rightarrow v(t) = \sum_{n=1}^{\infty} a_n n\omega \text{squc}(n\omega t) + b_n n\omega \text{squs}(n\omega t)$$

- Smooth acceleration is done similarly; non-smooth is computed using Matlab *gradient*
- Velocities and contact forces are captured well
- Accelerations show a limitation in the non-smooth formulation
- Dirac combs everywhere!

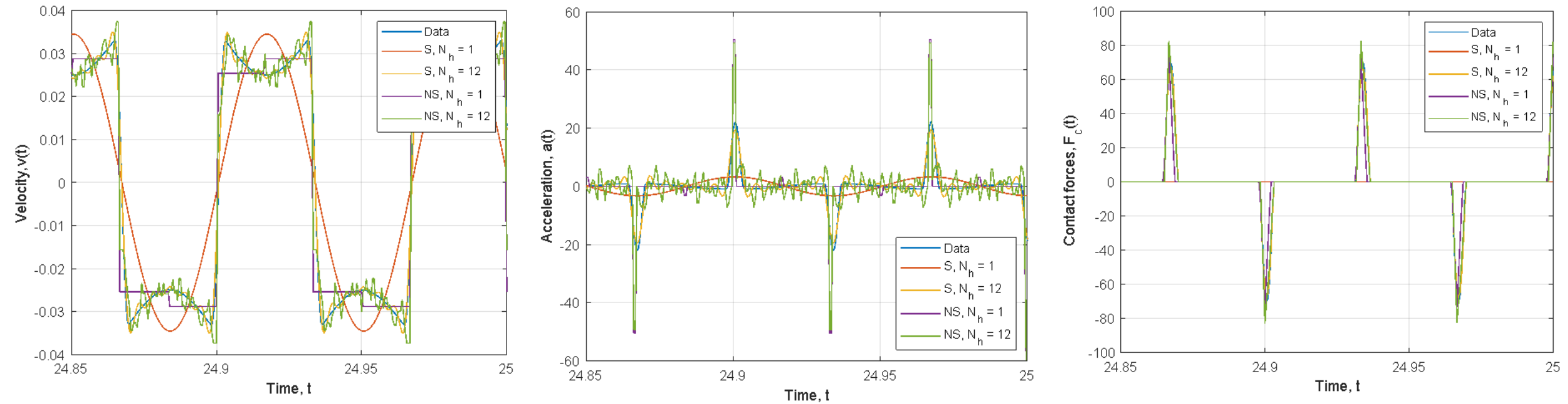
Least-squares regression analysis

Hard stiffness, $K_c = 1.4 \times 10^6 \frac{\text{N}}{\text{m}}$, $\omega = 15 \text{ Hz}$



- Less-smooth of a system response
- The non-smooth series converges faster until $N_h = 8$
- For more harmonics, the smooth curve-fit becomes better

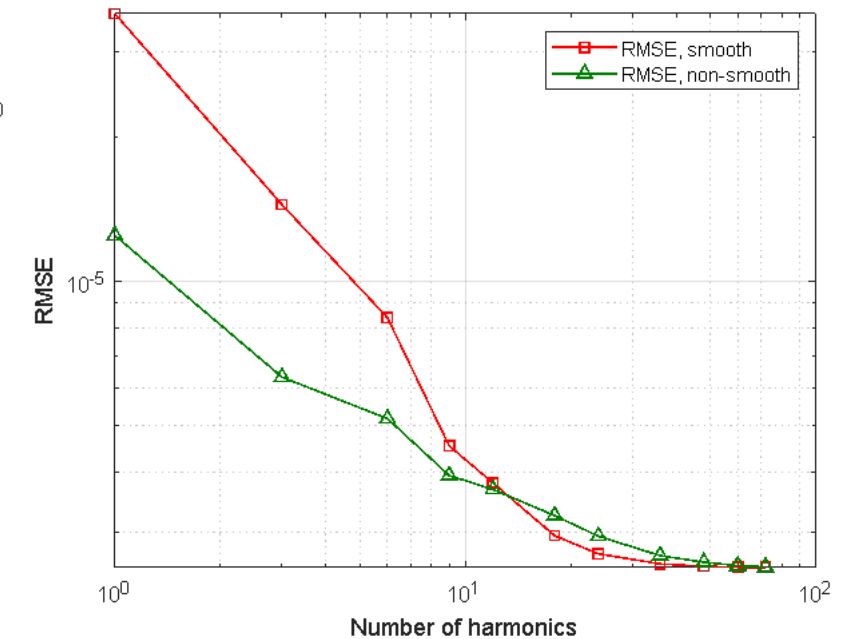
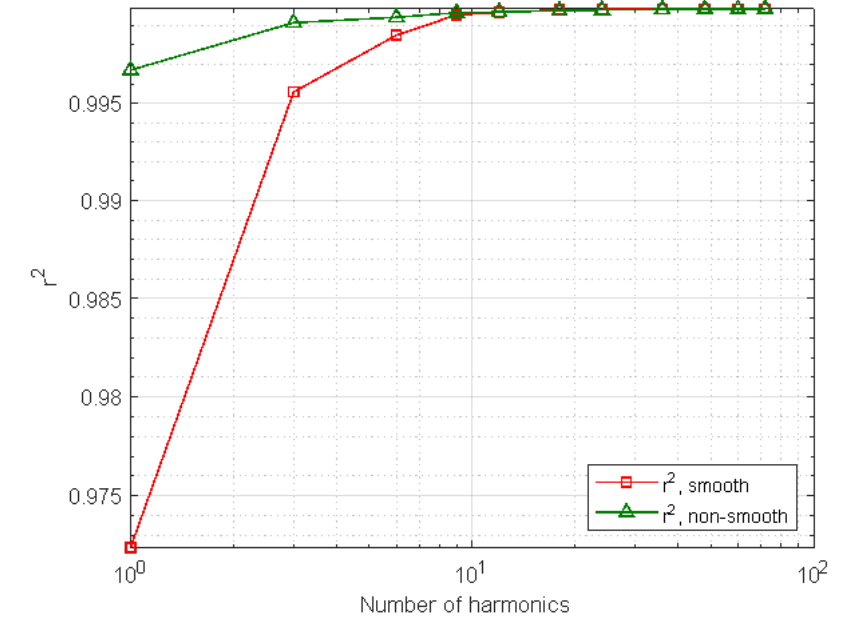
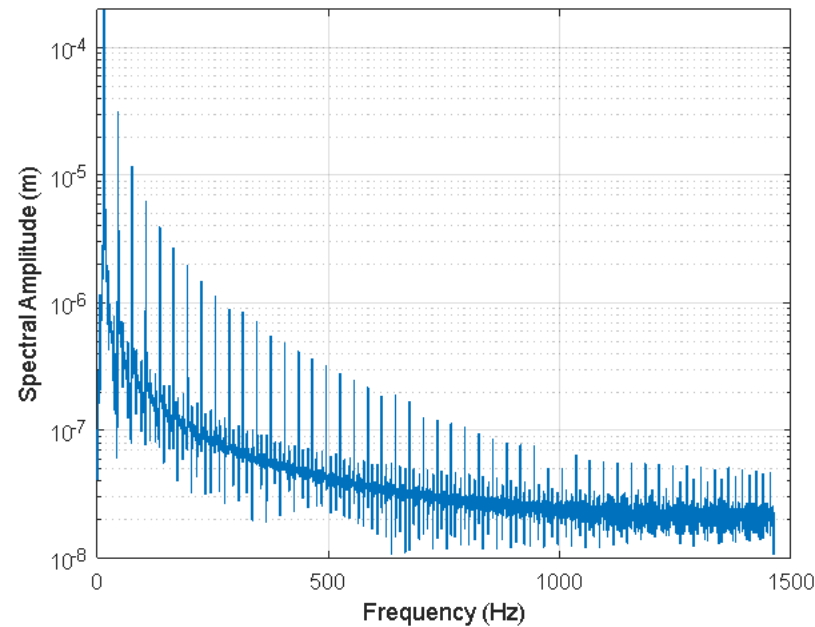
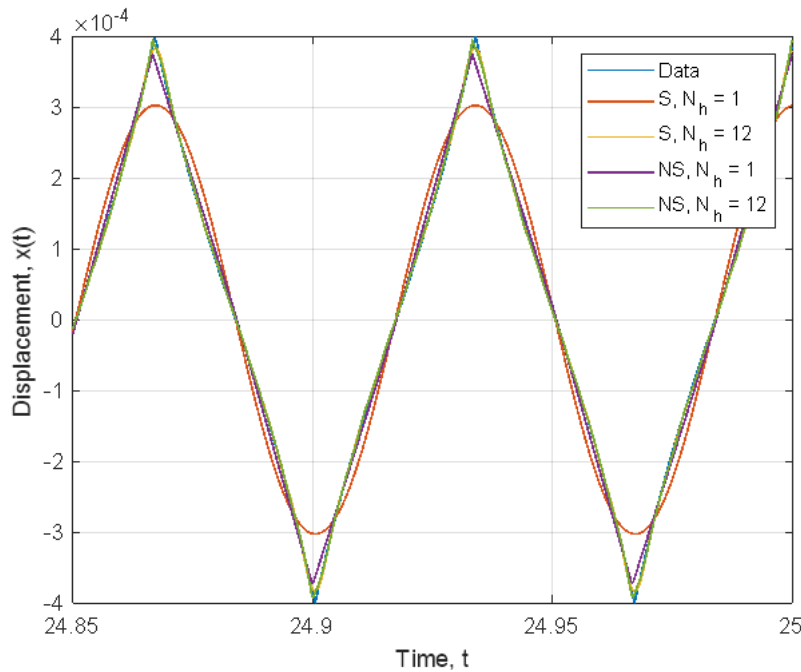
Least-squares regression analysis



- Both smooth and non-smooth velocities agree well and match closely to the data
- Smooth and non-smooth accelerations agree well away from the points of contact
- Non-smooth acceleration overshoots significantly
- Contact forces agree very well

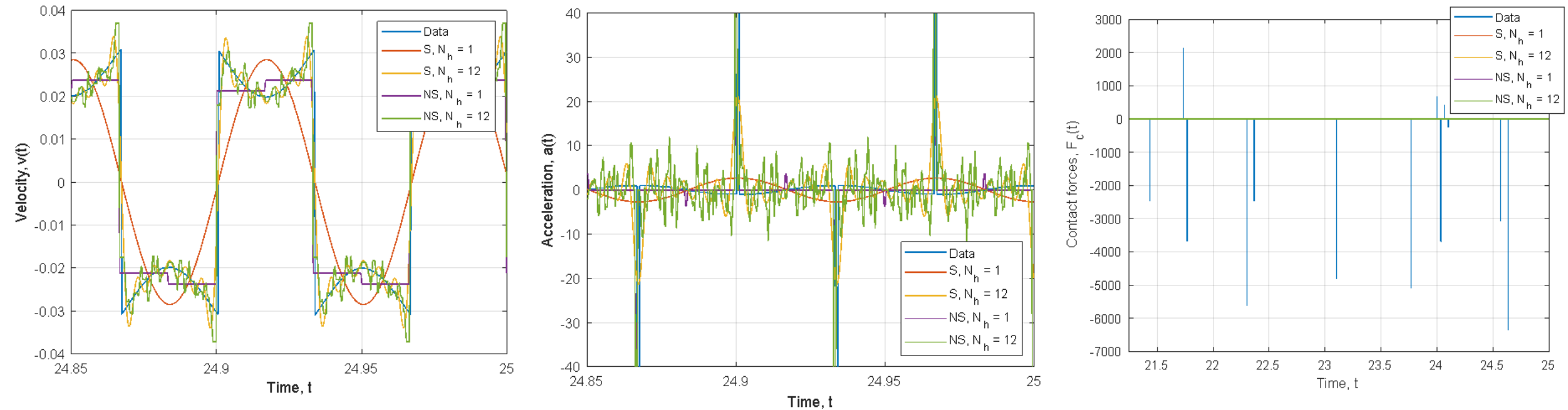
Least-squares regression analysis

Very hard stiffness, $K_c = 1.4 * 10^{10} \frac{N}{m}$, $\omega = 15 \text{ Hz}$



- Strongly non-smooth response
- The non-smooth series converges faster until $N_h = 12$
- For more harmonics, the smooth curve-fit becomes better

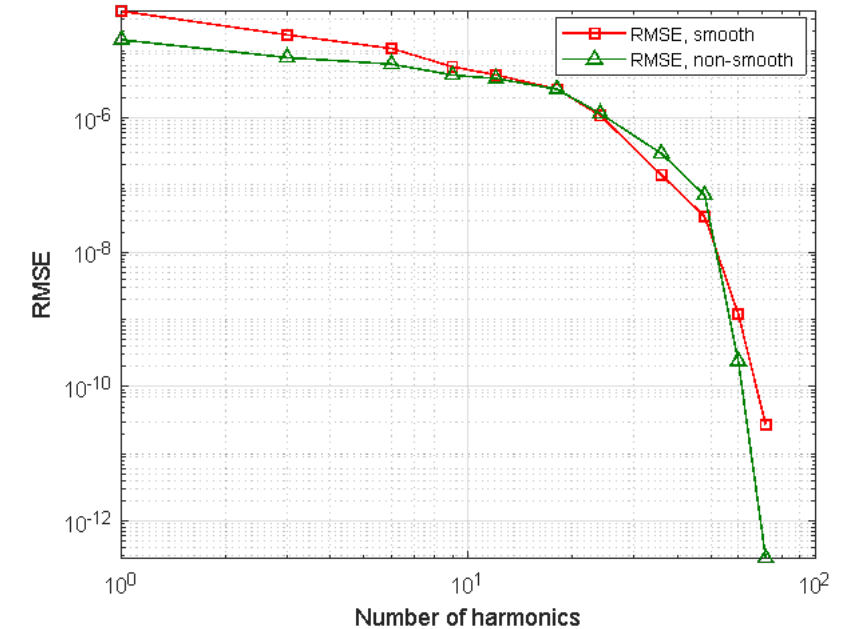
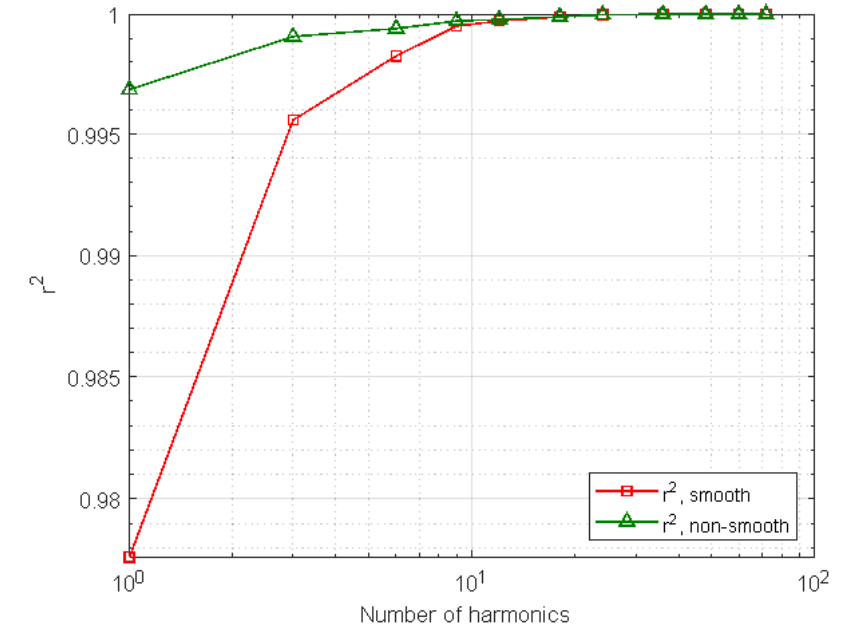
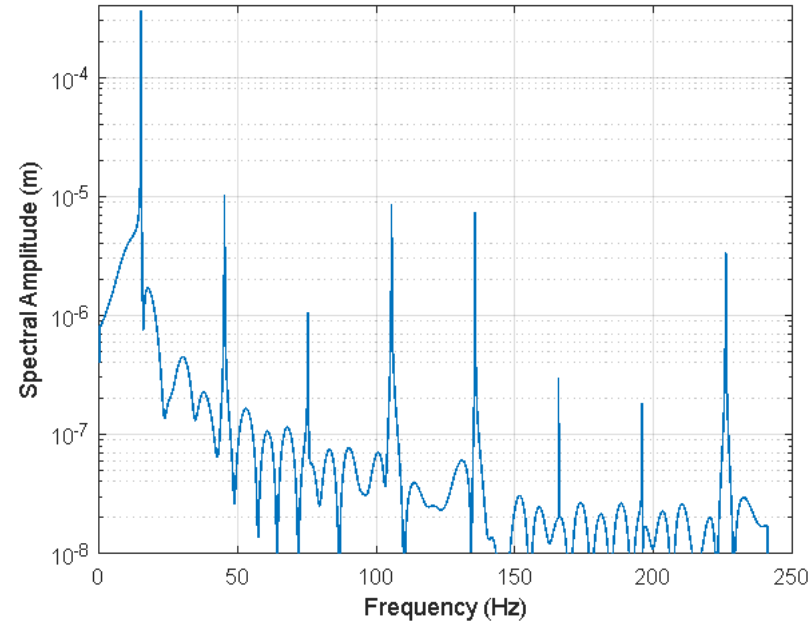
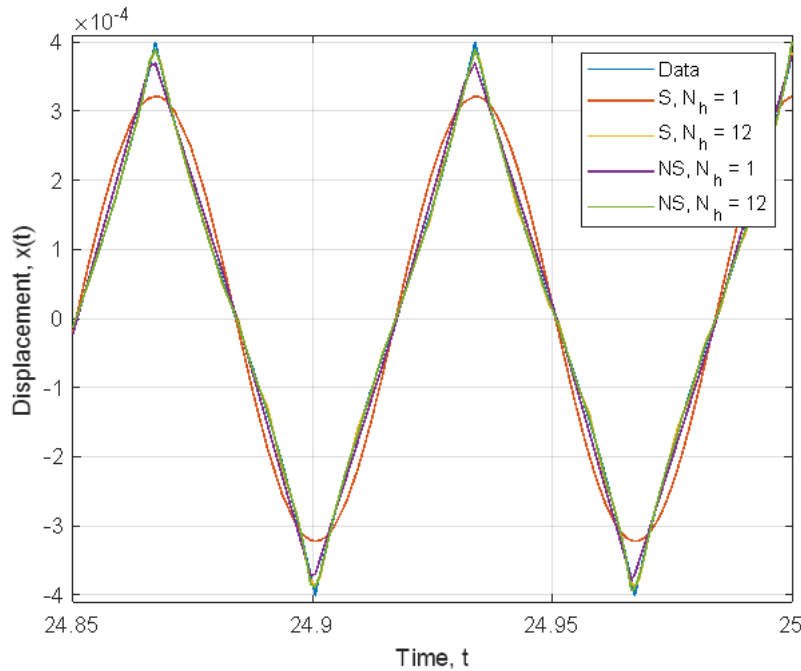
Least-squares regression analysis



- Velocities are again in agreement with each other and the data
- The non-smooth response best captures the peaks in acceleration
- Poor agreement everywhere else, however, for both smooth and non-smooth
- Contact forces are essentially Dirac impulses by now
- *None* of the curve-fits capture the contact force
- Why? Because the penalty stiffness force definition

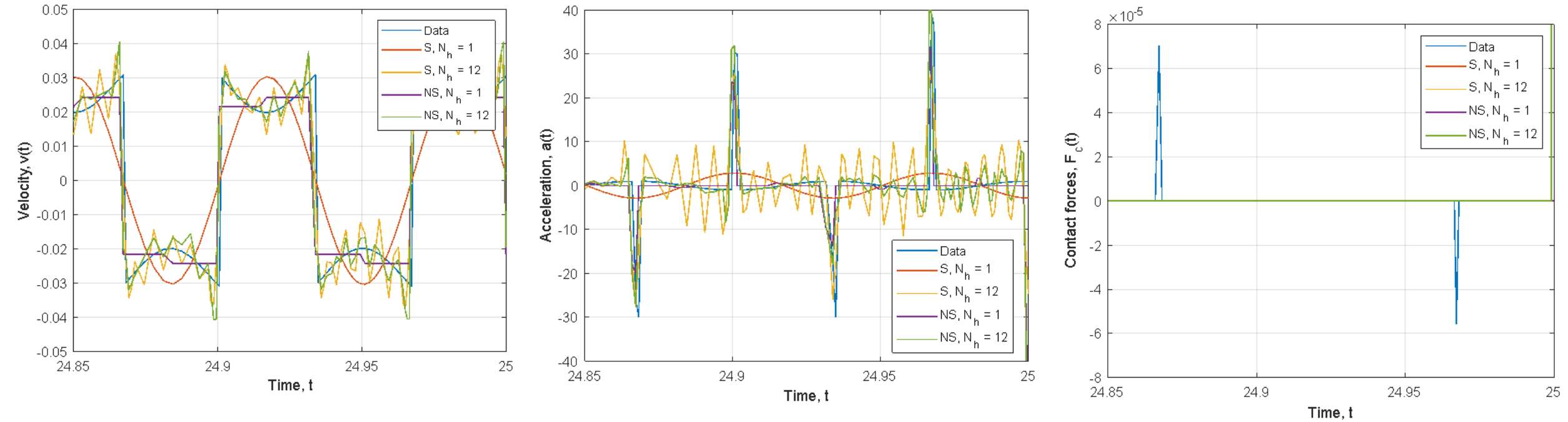
Least-squares regression analysis

Elastic impact, $COR = 1, \omega = 15 \text{ Hz}$



- Explicit hard impact using coefficient of restitution instead of penalty stiffness, aka the limit as $K_c \rightarrow \infty$
- The non-smooth fit converges faster until $N_h = 18$
- Then the smooth fit until $N_h \approx 54$
- Then the non-smooth fit again

Least-squares regression analysis



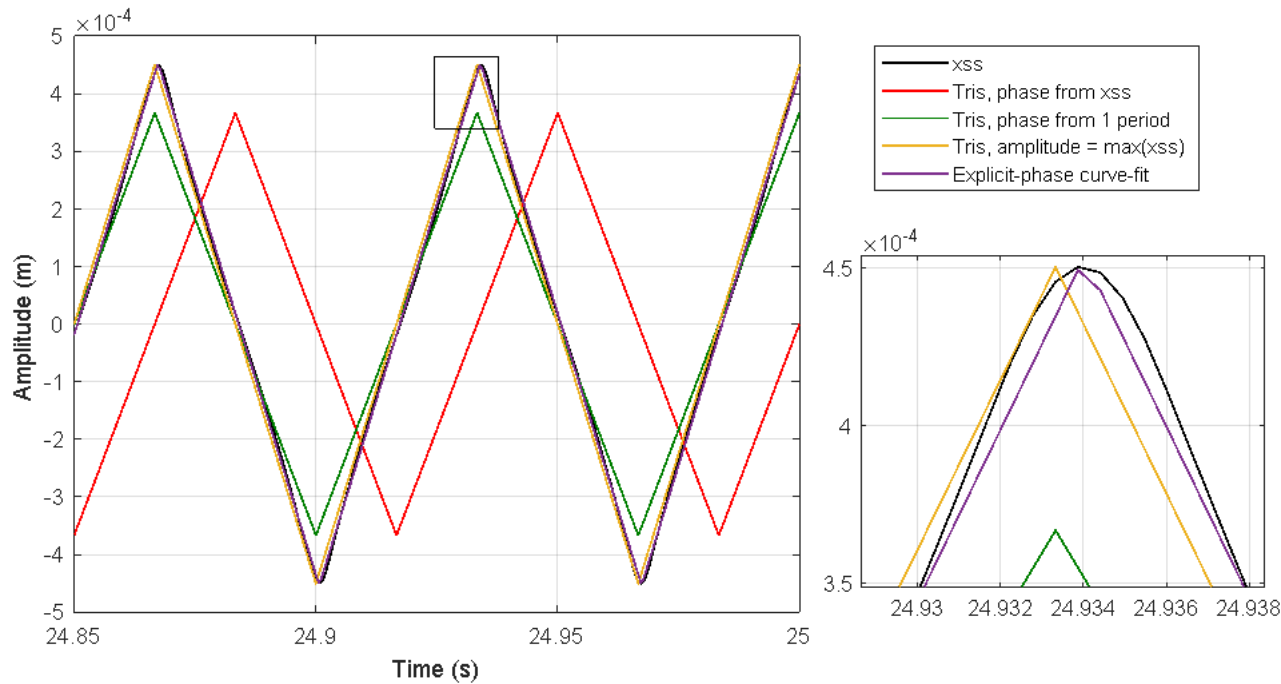
- Velocities are still in good agreement
- Both smooth and non-smooth capture the peaks in acceleration
- Poor agreement everywhere else, still
- Contact forces are Dirac impulses

Additional studies

❑ Consider a different curve fit of the following form:

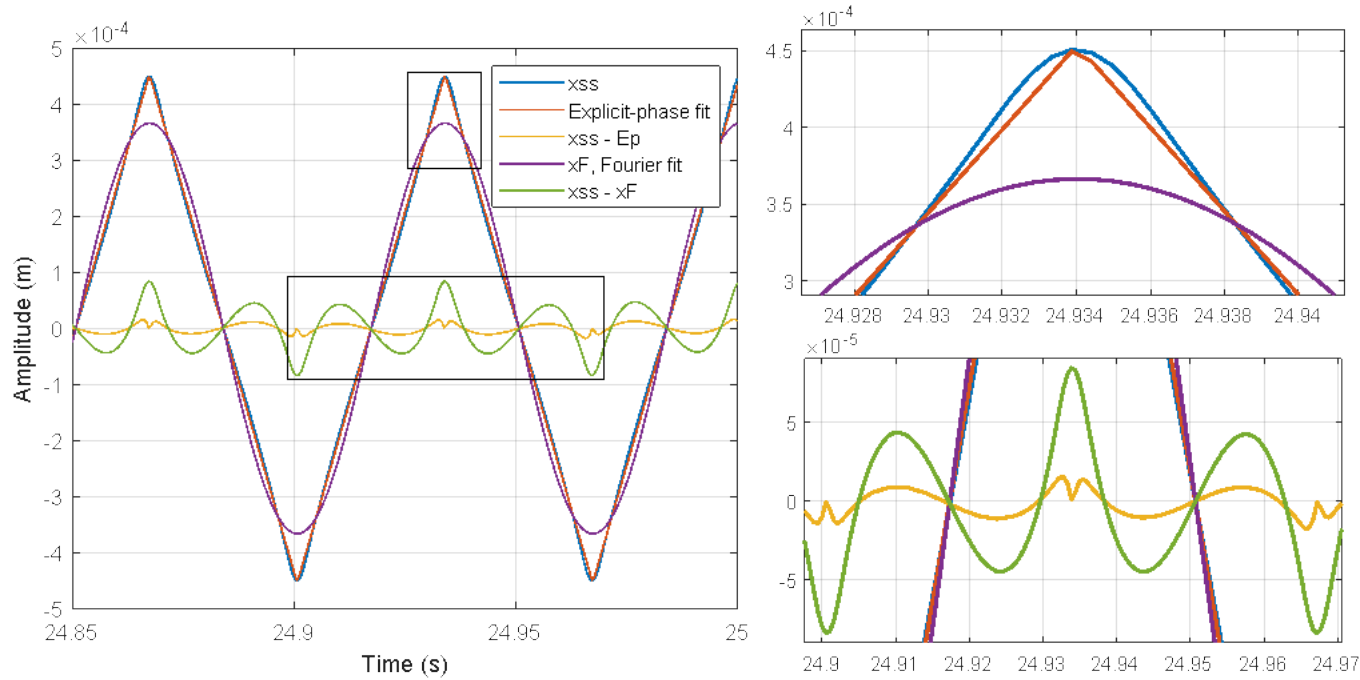
$$f(t) = a_0 + \sum_{n=1}^{\infty} a_n \text{tris}(n\omega t + \theta_n)$$

- Motivation: to better match up the locations of contacts/impacts with the phase-shifts that would produce them, hopefully reducing the number of terms needed
- Every additional non-smooth term means additional contacts and impulses/Dirac combs in the acceleration, which becomes unrealistic

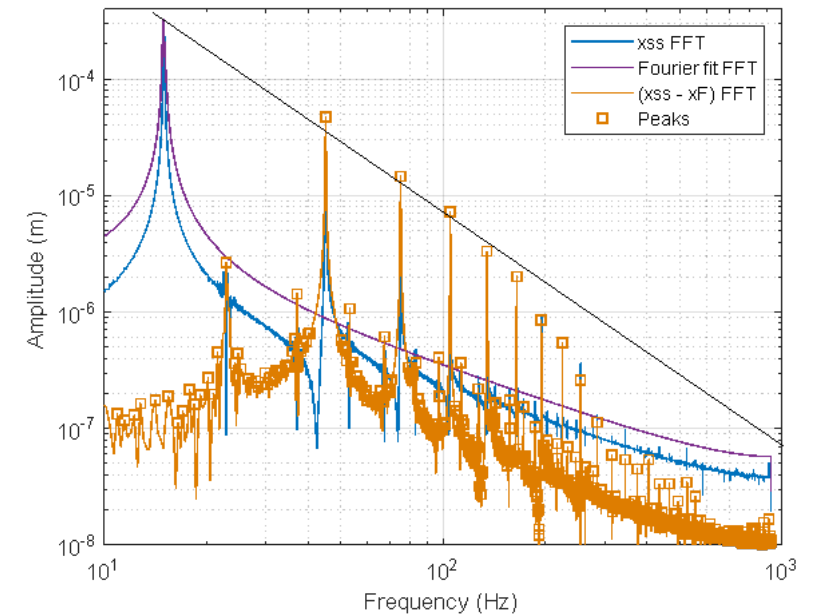
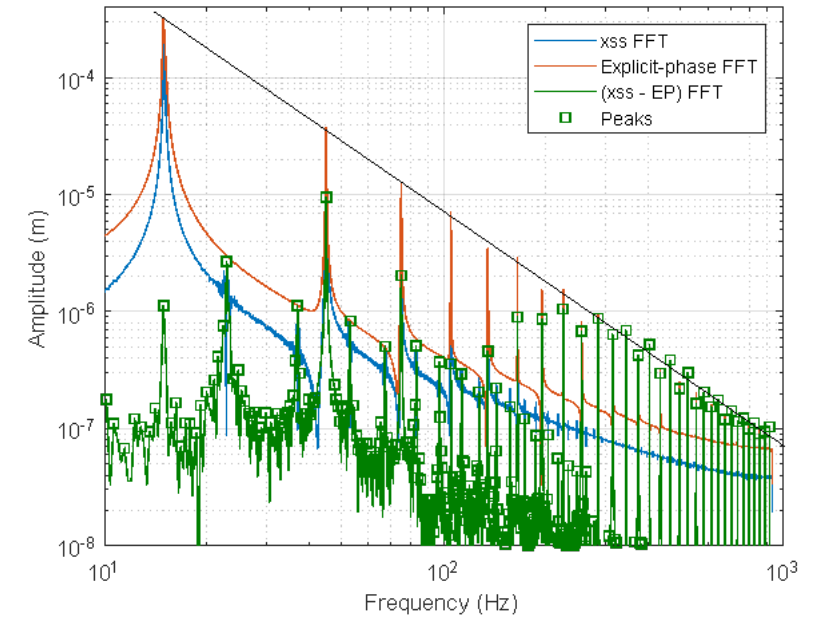


- Left: hard-contact case with different curve fits of a single triangle function
- 3 different manual curve fits
- 4th curve fit using Matlab nonlinear least squares
- First 3 tend to capture only amplitude or phase accurately, not both
- 4th one captures both well

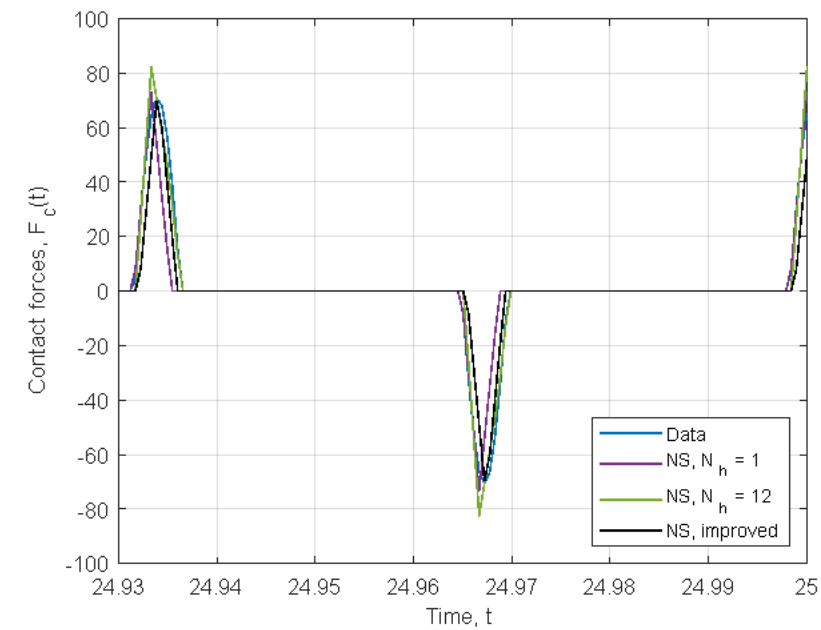
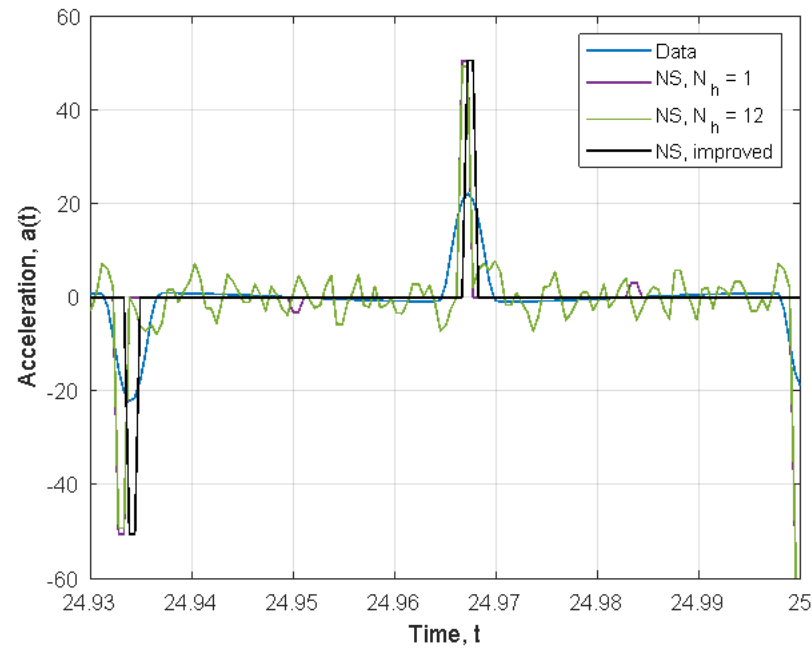
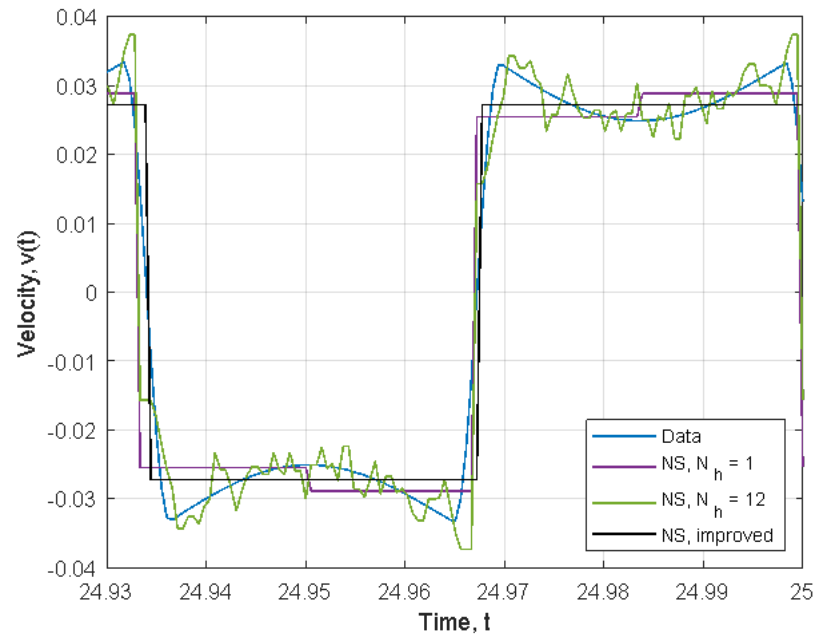
Additional studies



- Compare the improved curve fit to the Fourier fit for $N_h = 1$
- Original tris: $\text{RMSE} = 1.22 \times 10^{-5}$
- Fourier: $\text{RMSE} = 3.88 \times 10^{-5}$
- New tris: $\text{RMSE} = 8.55 \times 10^{-6}$
- 1.5 times better now!



Additional studies



- The improved non-smooth fit keeps good agreement with the velocity, acceleration, and contact force
- Further evidence that more terms does not necessarily improve the fit if performed naively

Conclusions

- ❑ Evaluated the usefulness of non-smooth basis functions for obtaining the response of a contact/impact system
- ❑ Non-smooth, triangular sine and cosine functions were defined
- ❑ Mathematical properties were highlighted
- ❑ Applied curve fits to time histories of a contact/impact system and studied for quality
- ❑ Results show Fourier series is superior for smooth responses, as expected
- ❑ The non-smooth series becomes superior for increasingly non-smooth responses
- ❑ Fourier series tends to become more accurate again when many harmonics are used
- ❑ A modified series form showed better results than the original naïve series form

Future work

- ❑ Continue studying non-smooth series representation and how to improve accuracy
- ❑ Optimal combinations of smooth and non-smooth terms based on when Fourier series regains highest accuracy
- ❑ Mathematical properties amenable to addition in a harmonic balance code

Acknowledgements

New Mexico State University

Sandia National Laboratories—Laboratory-Directed Research and Development (LDRD)

São Paulo State University (UNESP)



Thank you for your attention!
Please ask any questions

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