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# On the Harmonic Balance Method Augmented with Non-Smooth Basis Functions for Contact/Impact Problems

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# **Presentation outline**

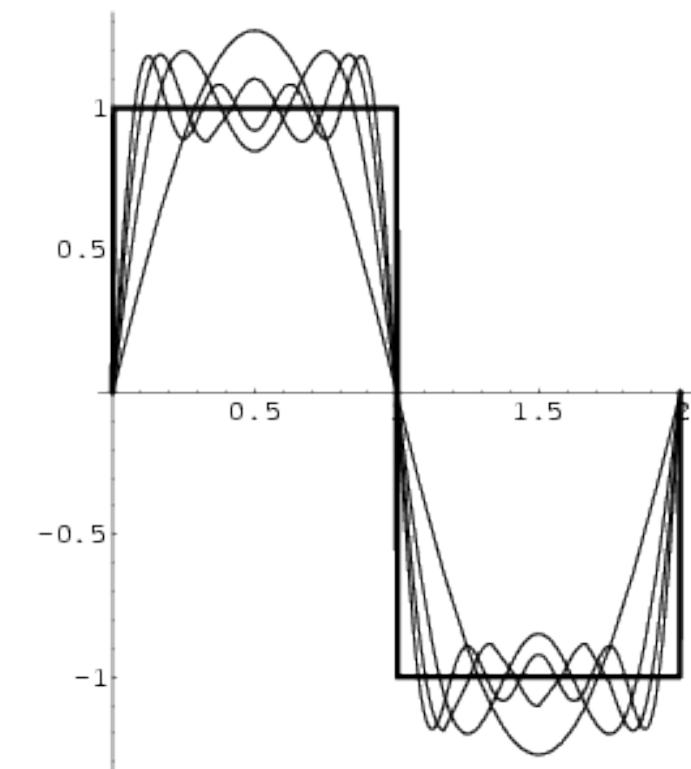
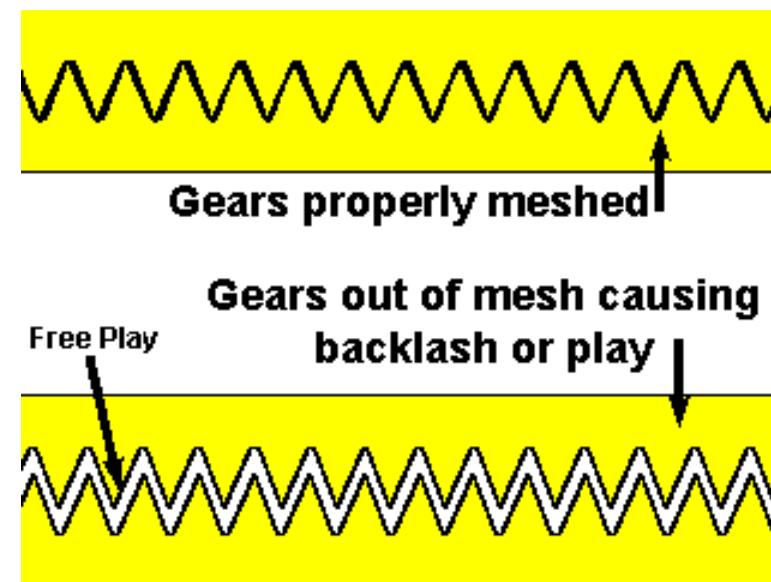
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- Introduction and motivation**
- Non-smooth Galerkin formulation**
- Least-squares regression analysis**
- Additional studies**
- Conclusions**

# Introduction and motivation

## Recent developments have efficiently applied the harmonic balance method to strongly nonlinear systems

- Systems include aircraft, spacecraft, gear drives, bladed disks, etc.
- Applications include continuation procedures, bifurcation detection and tracking, nonlinear modal analysis, etc.
- Advantages include reduced computational costs and capture of unstable solutions
- Difficulties include chaotic/aperiodic responses and non-smooth or discontinuous solutions (Gibbs phenomenon)
- How can we simulate “stiff” contact or friction systems efficiently with HBM?



# Introduction and motivation

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## □ Review of the literature—simulation approaches include:

- Lanczos filtering to improve the values of the Fourier coefficients
- Append additional, non-smooth or discontinuous terms to a system's solution
- Replace some or all the terms in classical Fourier series with non-smooth terms
- Utilize event-driven schemes to find and integrate between the state transition times to compute nonlinear forces
- Non-smooth temporal and spatial transformations

## □ Difficulties:

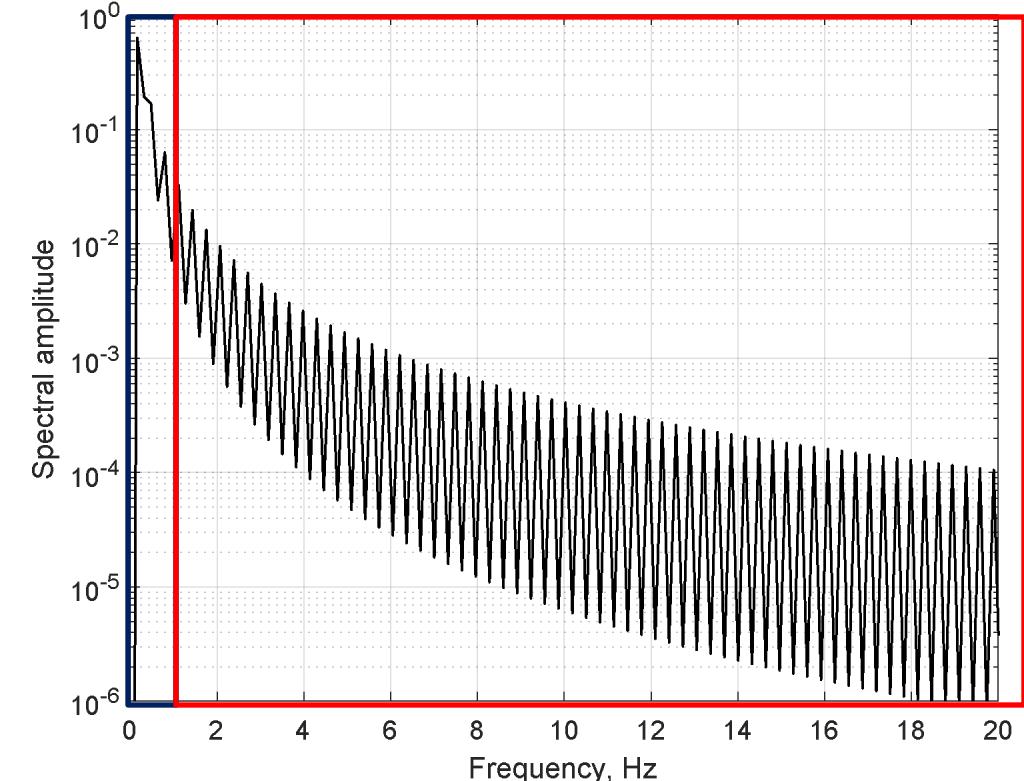
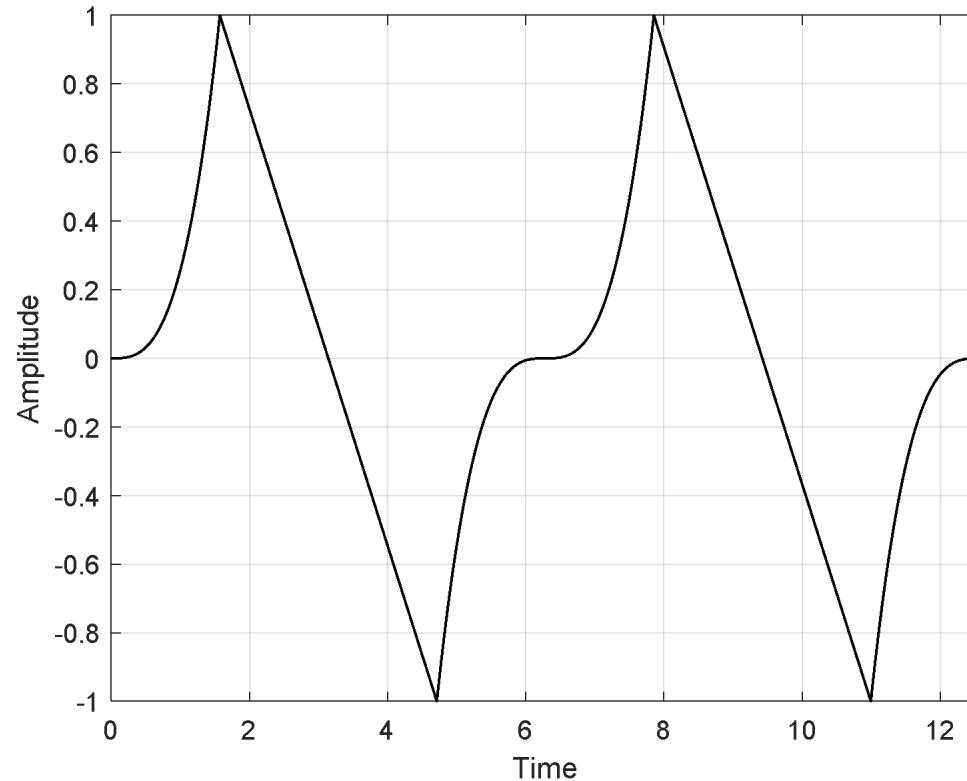
- Gibbs phenomenon: slow convergence (polynomial) compared to smooth systems (exponential)
- A priori knowledge of the state transition times may be required

# Introduction and motivation

## □ Motivating idea:

- Non-smooth periodic motions can be represented by infinite Fourier series
- Can we append the Fourier series representation with entire non-smooth basis functions?
- This approach may allow us to capture a large set of harmonics with a small number of non-smooth functions

- Some sine/cosine terms
- Some non-smooth terms



# Introduction and motivation

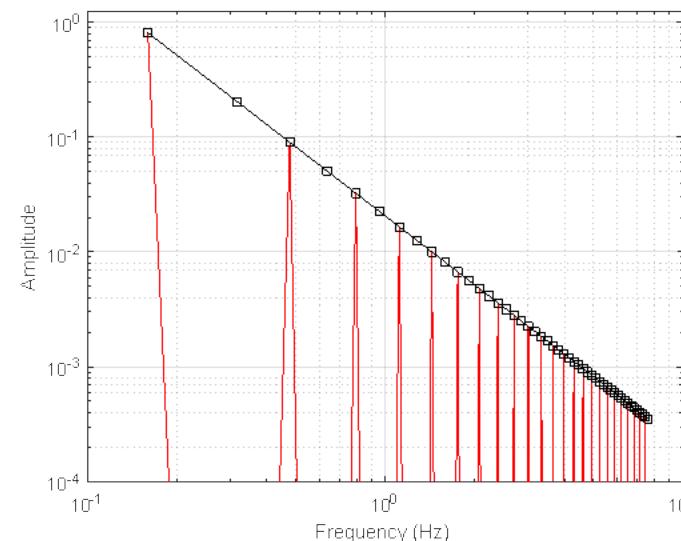
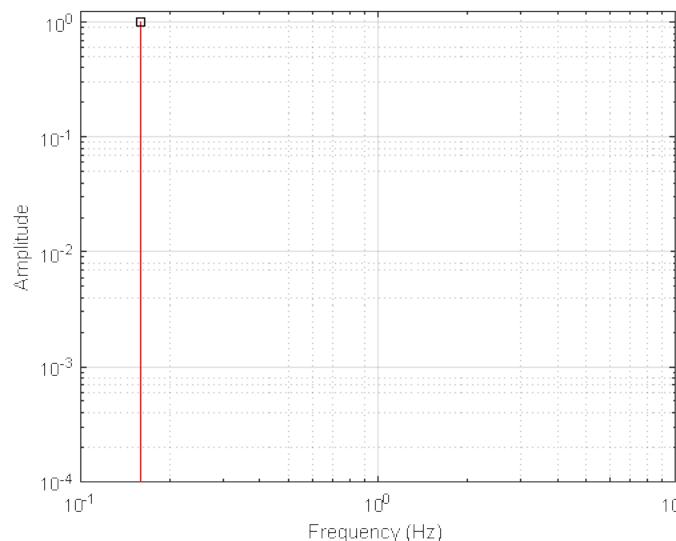
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## □ Desirable basis function traits:

- Easy to mathematically represent
- Convenient mathematical properties (Fourier series, derivatives, etc.)
- Intuitive, non-smooth counterparts to sine and cosine functions

## □ Goals of this work:

- Select functions with  $C^0$  smoothness for capturing a non-smooth functional representation of a solution
- Use goodness-of-fit metrics to evaluate the classical Fourier series and the non-smooth basis functions
- Develop a framework that can later be implemented into harmonic balance formulations



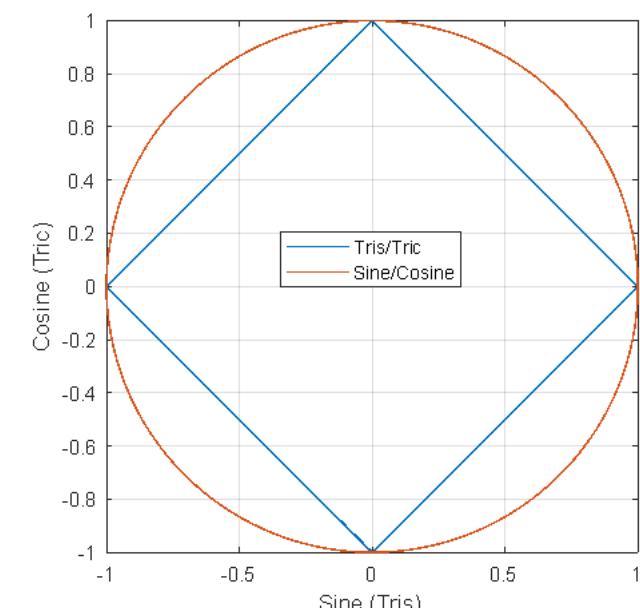
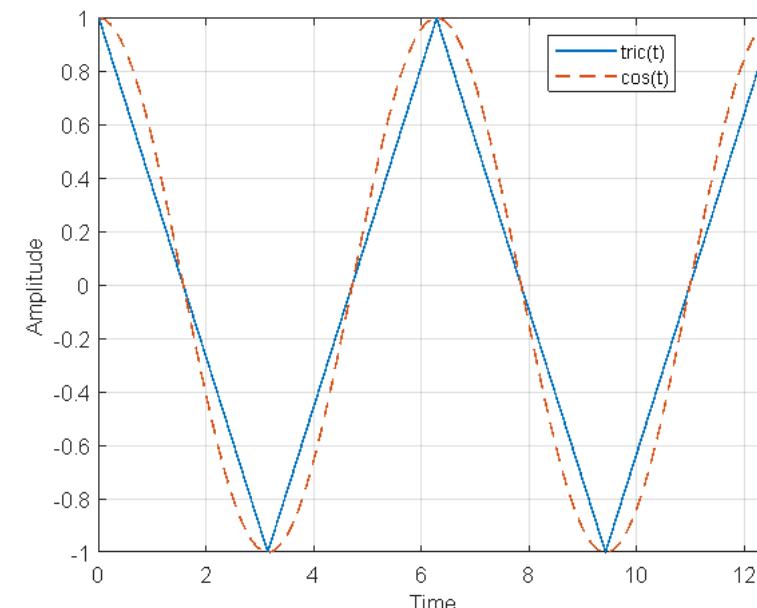
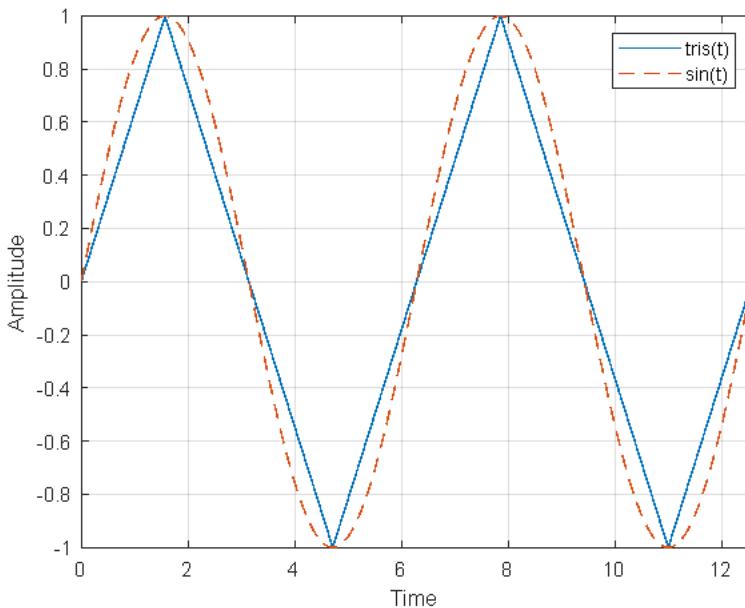
# Non-smooth Galerkin formulation

## □ Non-smooth triangle waves are chosen for study

- “Triangle sine” and “triangle cosine”
- The same periodicity, maxima, minima, and roots as sine and cosine waves

$$tris(\omega t) = \begin{cases} \frac{4}{T}t_m, & t_m < \frac{T}{4} \\ -\frac{4}{T}t_m + 2, & \frac{T}{4} \leq t_m \leq \frac{3T}{4}, \\ \frac{4}{T}t_m - 4, & t_m > \frac{3T}{4} \end{cases} \quad tric(\omega t) = \begin{cases} -\frac{4}{T}t_m + 1, & t_m \leq \frac{T}{2} \\ \frac{4}{T}t_m - 3, & t_m > \frac{T}{2} \end{cases}$$

$$t_m = \omega t(\bmod T), \quad T = 2\pi$$



# Non-smooth Galerkin formulation

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- **Infinite series representation:**

$$f(t) = a_0 + \sum_{n=1}^{\infty} a_n \text{tric}(n\omega t) + b_n \text{tris}(n\omega t)$$

- **A different form for numerical convenience**

$$f(t) = a_0 + \sum_{n=1}^{\infty} a_n \text{tris}(n\omega t + \theta_n)$$

- **Now we can study these functions for desirable properties and advantages**



# Mathematical properties

## □ The non-smooth functions have similar properties to sine/cosine:

➤ Reflections:

| $\theta$ reflected in $\alpha = 0$            | $\theta$ reflected in $\alpha = \frac{\pi}{4}$                         | $\theta$ reflected in $\alpha = \frac{\pi}{2}$     | $\theta$ reflected in $\alpha = \frac{3\pi}{4}$                          | $\theta$ reflected in $\alpha = \pi$                |
|---|--|--|--|---|
| $\text{tris}(-\theta) = -\text{tris}(\theta)$ | $\text{tris}\left(\frac{\pi}{2} - \theta\right) = \text{tric}(\theta)$ | $\text{tris}(\pi - \theta) = +\text{tris}(\theta)$ | $\text{tris}\left(\frac{3\pi}{2} - \theta\right) = -\text{tric}(\theta)$ | $\text{tris}(2\pi - \theta) = -\text{tris}(\theta)$ |
| $\text{tric}(-\theta) = +\text{tric}(\theta)$ | $\text{tric}\left(\frac{\pi}{2} - \theta\right) = \text{tris}(\theta)$ | $\text{tric}(\pi - \theta) = -\text{tric}(\theta)$ | $\text{tric}\left(\frac{3\pi}{2} - \theta\right) = -\text{tris}(\theta)$ | $\text{tric}(2\pi - \theta) = +\text{tric}(\theta)$ |

➤ Shifts and periodicity:

| Shift by one quarter period  | Shift by one half period                             | Shift by three quarter periods  | Shift by full periods                                   |
|--|--|---|---|
| $\text{tris}\left(\theta \pm \frac{\pi}{2}\right) = \pm \text{tric}(\theta)$ | $\text{tris}(\theta \pm \pi) = -\text{tris}(\theta)$ | $\text{tris}\left(\theta \pm \frac{3\pi}{2}\right) = \mp \text{tric}(\theta)$ | $\text{tris}(\theta \pm 2\pi k) = +\text{tris}(\theta)$ |
| $\text{tric}\left(\theta \pm \frac{\pi}{2}\right) = \mp \text{tris}(\theta)$ | $\text{tric}(\theta \pm \pi) = -\text{tric}(\theta)$ | $\text{tric}\left(\theta \pm \frac{3\pi}{2}\right) = \pm \text{tris}(\theta)$ | $\text{tric}(\theta \pm 2\pi k) = +\text{tric}(\theta)$ |

Miscellaneous:

- $|\text{tris}(\theta)| + |\text{tric}(\theta)| = 1$
- Angle sums/phase shifts: not straightforward

# Mathematical properties

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## □ Fourier series representations ( $\omega$ integer):

$$tris(\omega t) = \sum_{n=1}^{\infty} b_n \sin(nt), \quad b_n = \frac{4\omega}{n^2\pi^2} \sum_{k=1}^{2\omega} (-1)^{k+1} \sin\left(\frac{(2k-1)n\pi}{2\omega}\right)$$

$$tric(\omega t) = a_0 + \sum_{n=1}^{\infty} a_n \cos(nt), \quad a_0 = 0, \quad a_n = \frac{4\omega}{n^2\pi^2} \sum_{k=1}^{2\omega-1} (-1)^k \cos\left(\frac{kn\pi}{\omega}\right)$$

➤ Examples:

$$tris(t) = \sum_{n=1}^{\infty} b_n \sin(nt), \quad b_n = \frac{8}{n^2\pi^2} \sin\left(\frac{n\pi}{2}\right) = \frac{8}{\pi^2}, 0, \frac{-8}{9\pi^2}, 0, \frac{8}{25\pi^2}, 0, \frac{-8}{49\pi^2}, \dots$$

$$tric(t) = \sum_{n=1}^{\infty} a_n \cos(nt), \quad a_n = \frac{4}{n^2\pi^2} [1 - \cos(n\pi)] = \frac{8}{\pi^2}, 0, \frac{8}{9\pi^2}, 0, \frac{8}{25\pi^2}, 0, \frac{8}{49\pi^2} \dots$$

➤ First derivatives: square waves with coefficients  $\propto \frac{8\omega}{n\pi^2}$

➤ Second derivatives: Dirac combs with coefficients  $\propto \frac{8\omega}{\pi^2}$

- The functions numerically satisfy the following *orthogonality* relationships:

$$\int_0^{2\pi} tris(mx)dx = 0,$$

$$\int_0^{2\pi} tric(mx)dx = 0,$$

$$\int_0^{2\pi} tric(mx)tric(nx)dx = c_m \delta_{mn},$$

$$\int_0^{2\pi} tris(mx)tris(nx)dx = d_m \delta_{mn},$$

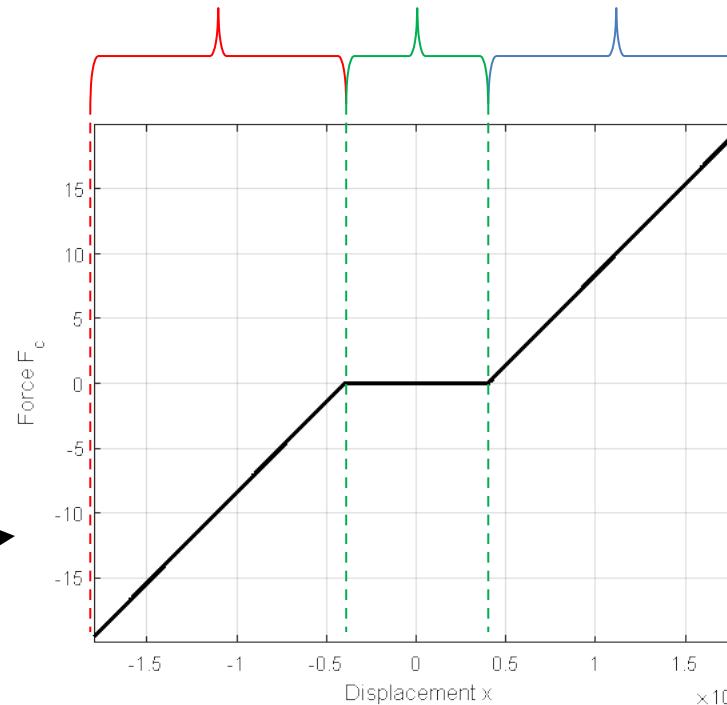
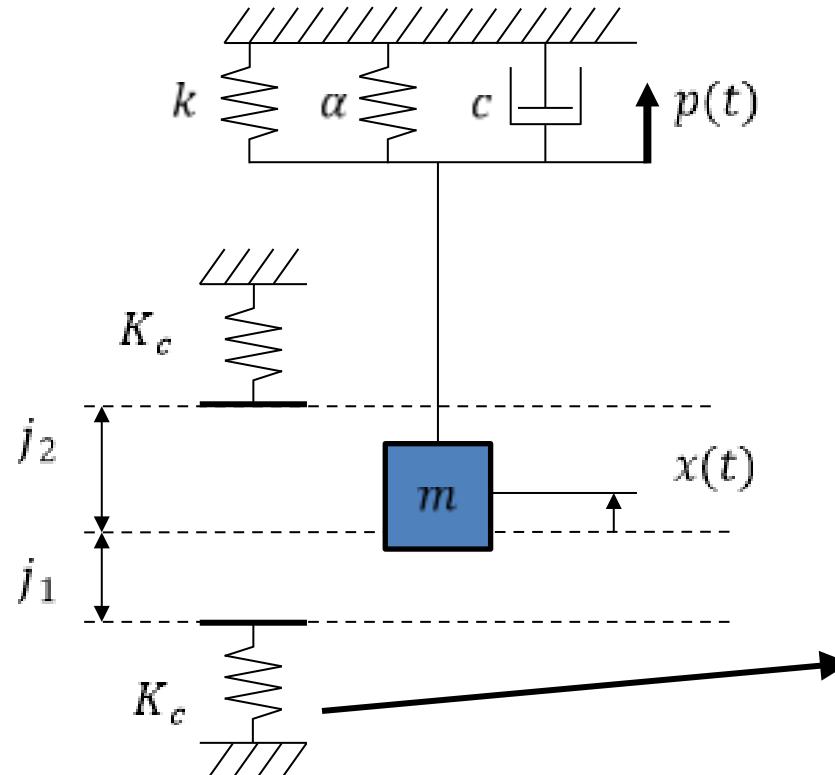
$$\int_0^{2\pi} tric(mx)tris(nx)dx = 0,$$

- For integers  $m, n$  and constants  $c_m, d_m$
- $\delta_{mn}$  denotes Kronecker delta
- Result: biorthogonal system, just like sine and cosine
- A functional infinite series representation, similar to Fourier series, is possible
- Reiterate: no rigorous proof of this

# Least-squares regression analysis

## □ The system: a forced Duffing oscillator with freeplay

$$\ddot{x} + 2\omega_n \zeta \dot{x} + \omega_n^2 x + \frac{\alpha}{m} x^3 + \frac{F_c(x)}{m} = \frac{p}{m} \cos(\omega t),$$



- Two different contact laws—contact penalty stiffness, and elastic impact:

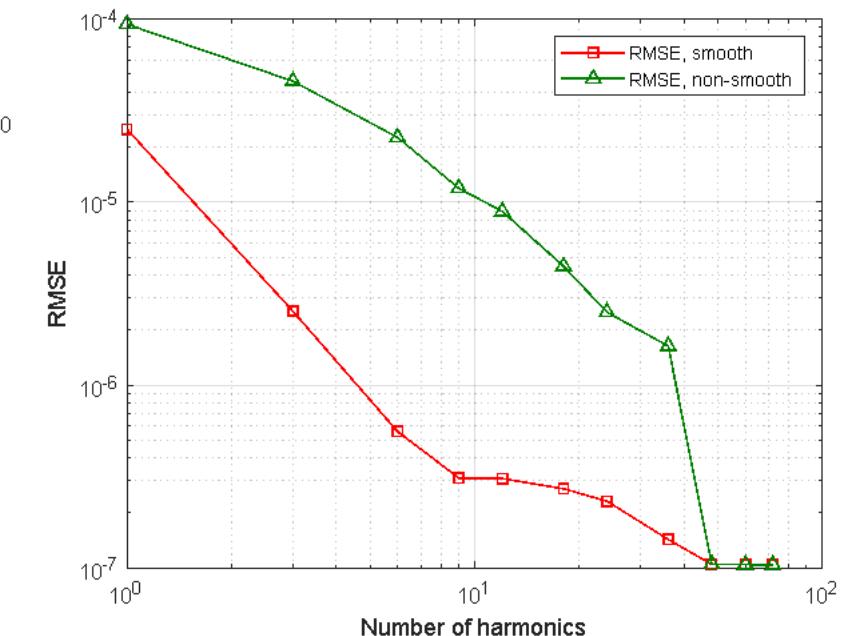
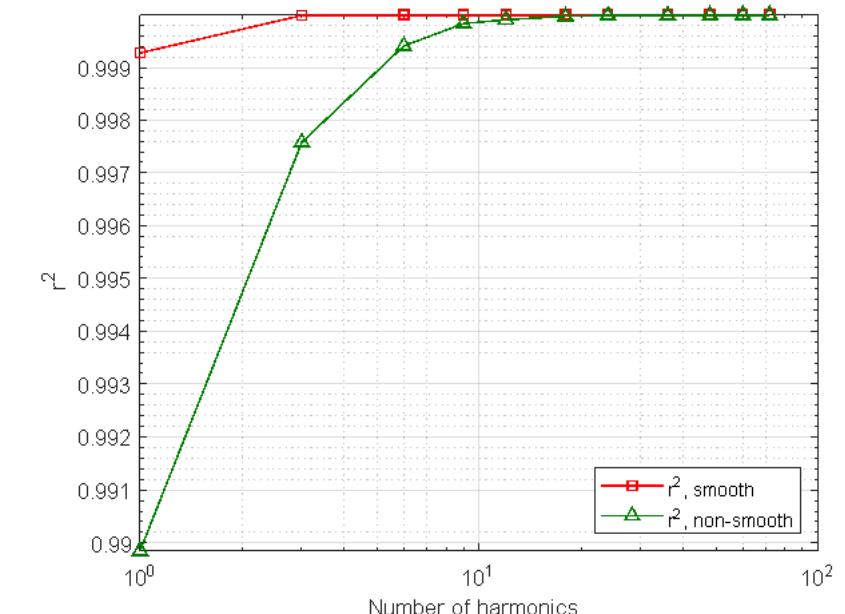
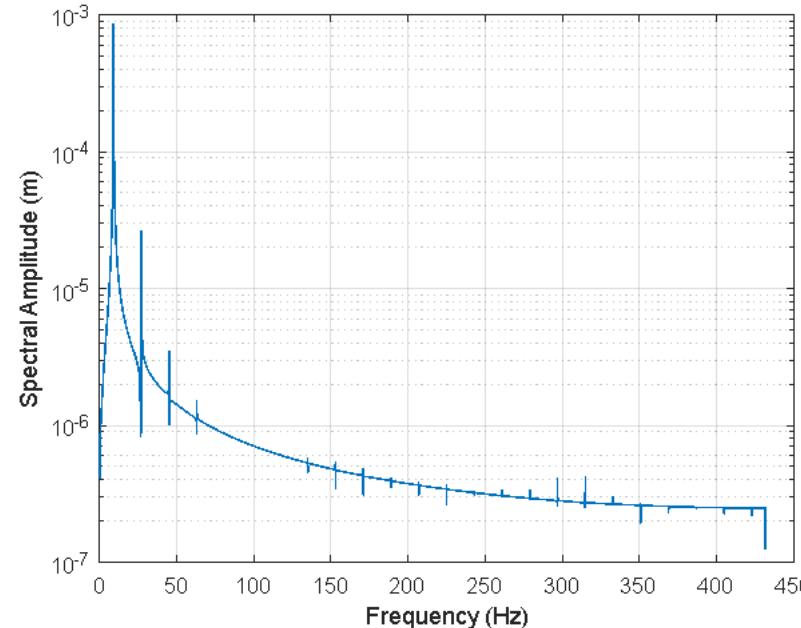
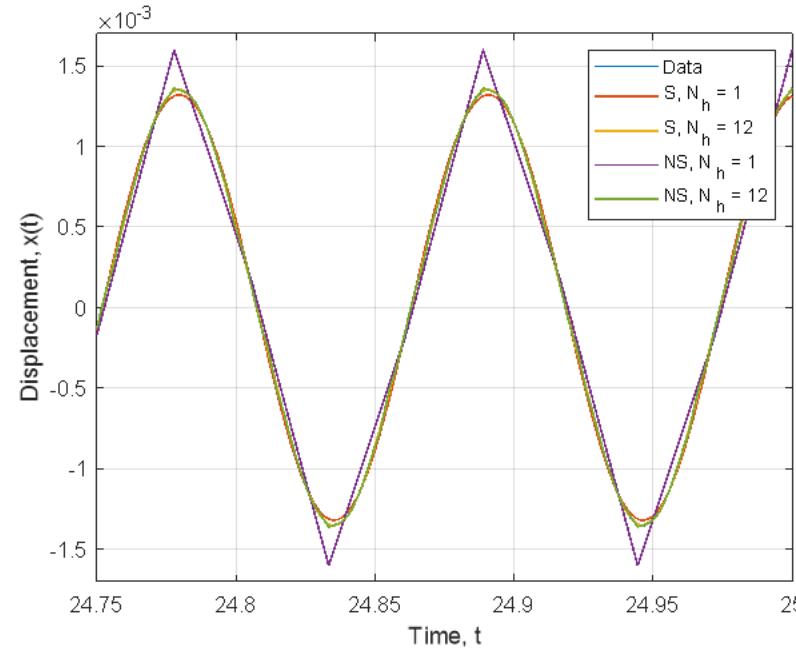
$$F_c = \begin{cases} K_c(x + j_1), & x < -j_1 \\ 0, & -j_1 \leq x \leq j_2 \\ K_c(x - j_2), & x > j_2 \end{cases}$$

vs.

$$-j_1 \leq x(t) \leq j_2, \\ x^+ = -rx^-, r = 1$$

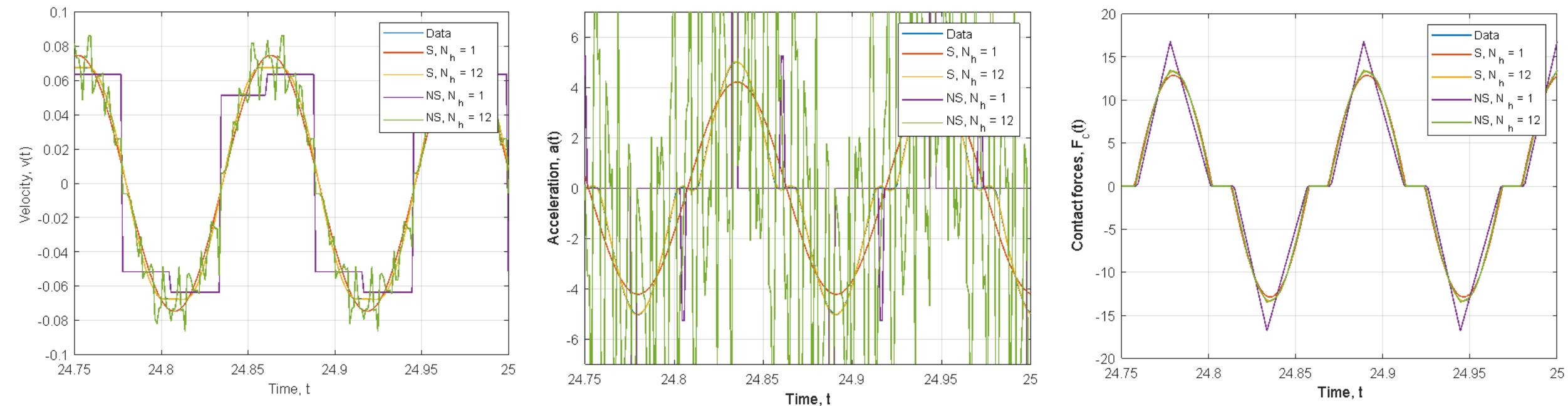
# Least-squares regression analysis

Medium stiffness,  $K_c = 1.4 * 10^4 \frac{N}{m}$ ,  $\omega = 9 \text{ Hz}$



- Smooth system response
- Classical Fourier series converges much faster than non-smooth Fourier series

# Least-squares regression analysis



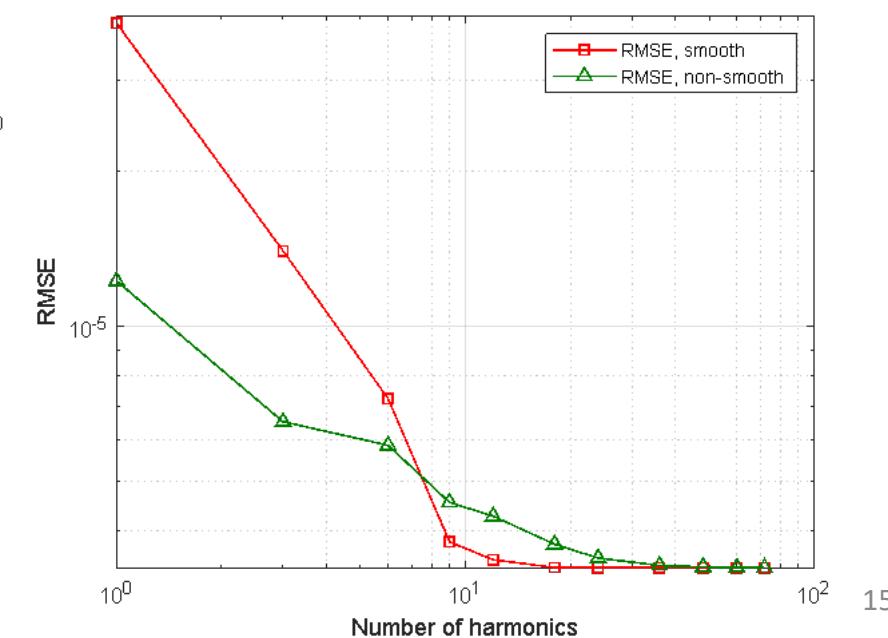
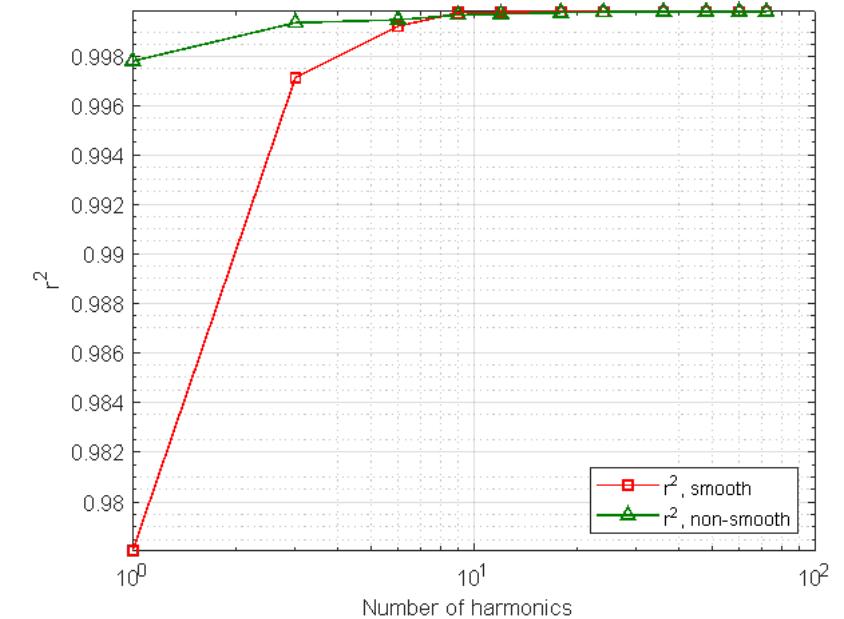
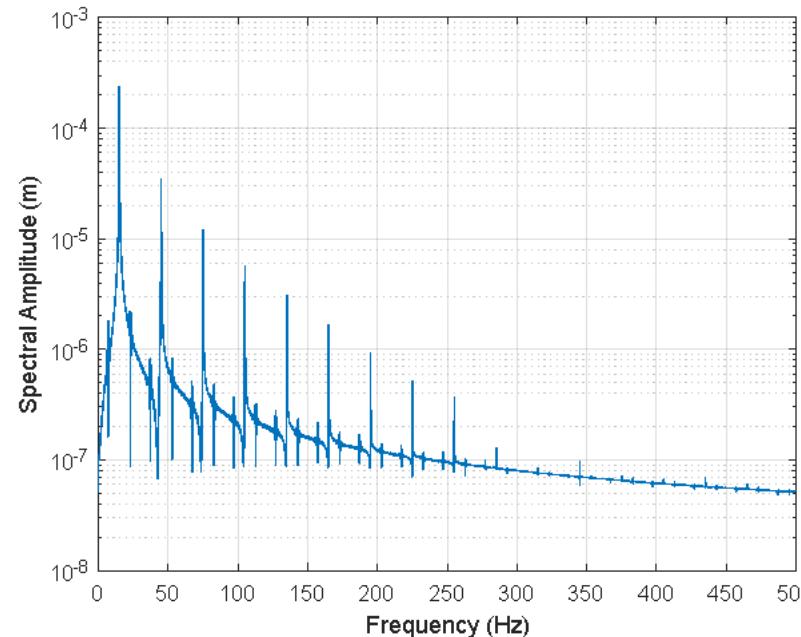
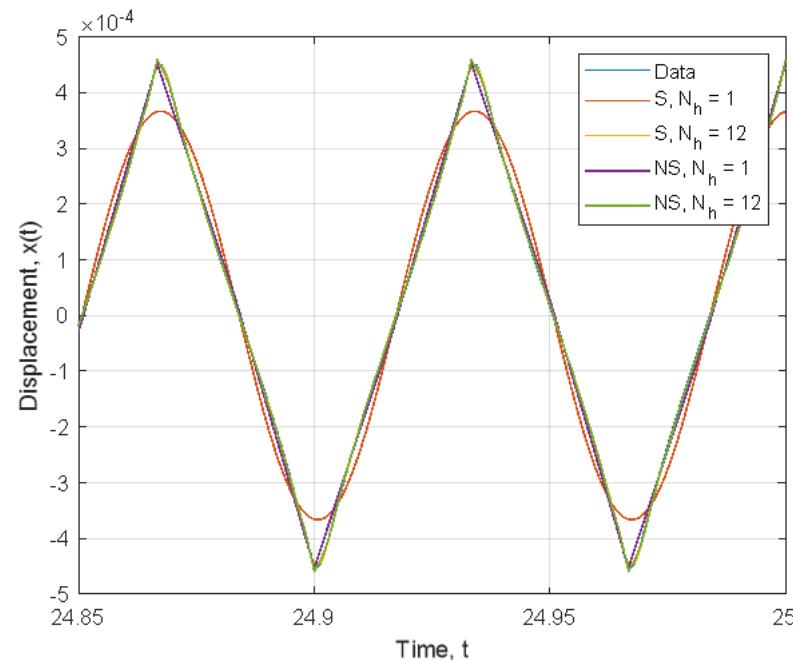
- Now let's look at the velocity, acceleration, and contact forces obtained using low- and high-quality fits
- Both smooth and non-smooth velocities use the coefficients obtained from curve-fitting the displacement:

$$x(t) = a_0 + \sum_{n=1}^{\infty} a_n \text{tric}(n\omega t) + b_n \text{tris}(n\omega t) \Rightarrow v(t) = \sum_{n=1}^{\infty} a_n n\omega \text{squc}(n\omega t) + b_n n\omega \text{squs}(n\omega t)$$

- Smooth acceleration is done similarly; non-smooth is computed using Matlab *gradient*
- Velocities and contact forces are captured well
- Accelerations show a limitation in the non-smooth formulation
- Dirac combs everywhere!

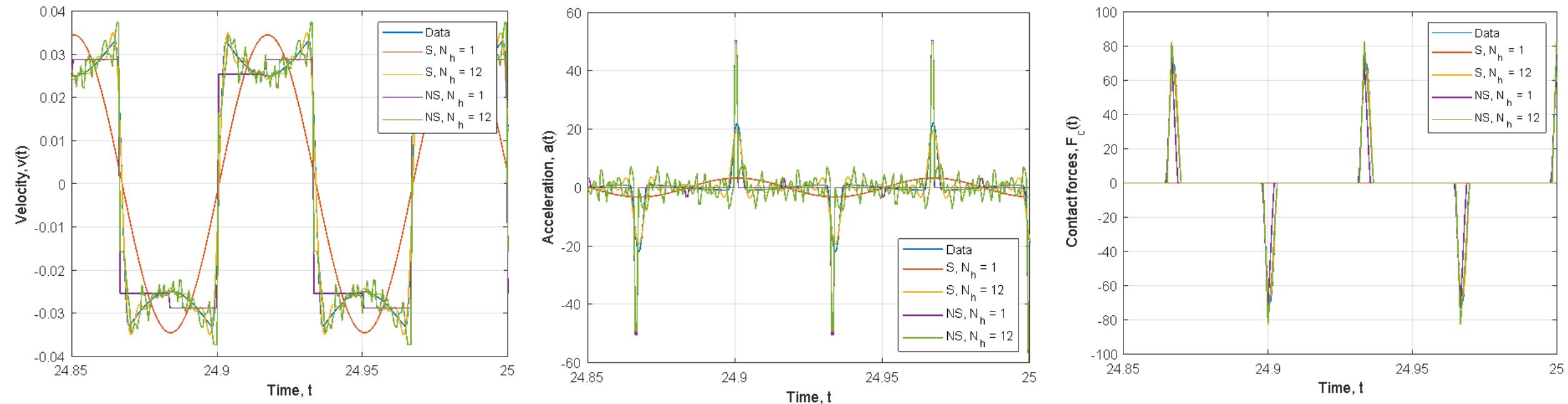
# Least-squares regression analysis

Hard stiffness,  $K_c = 1.4 * 10^6 \frac{N}{m}$ ,  $\omega = 15 \text{ Hz}$



- Less-smooth of a system response
- The non-smooth series converges faster until  $N_h = 8$
- For more harmonics, the smooth curve-fit becomes better

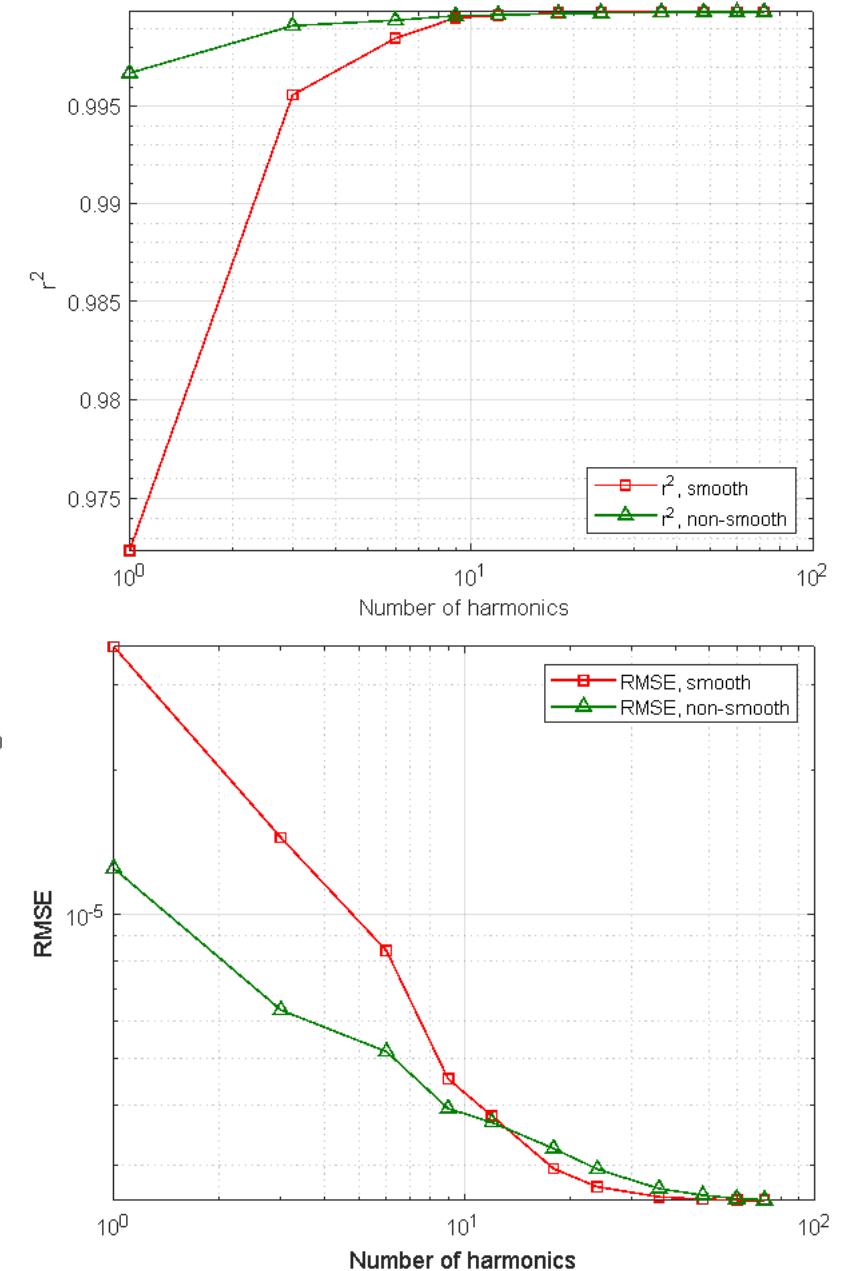
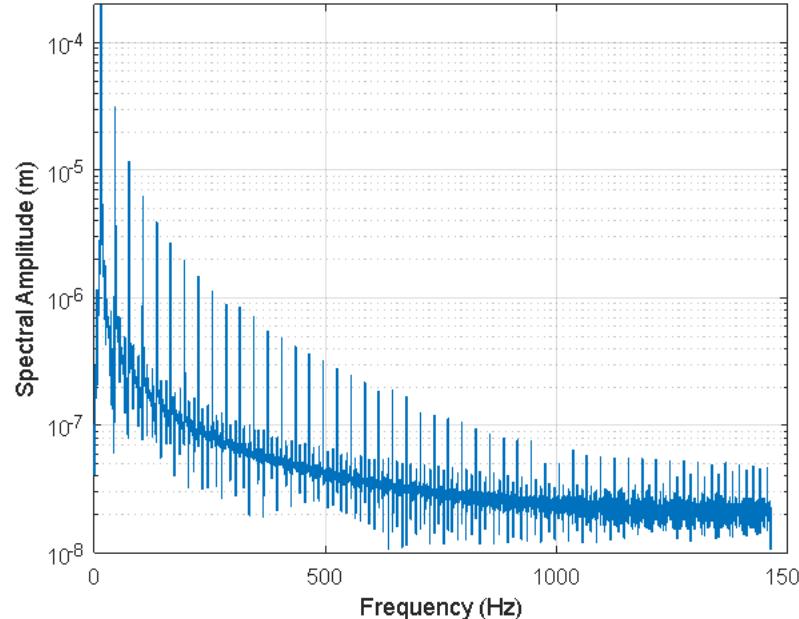
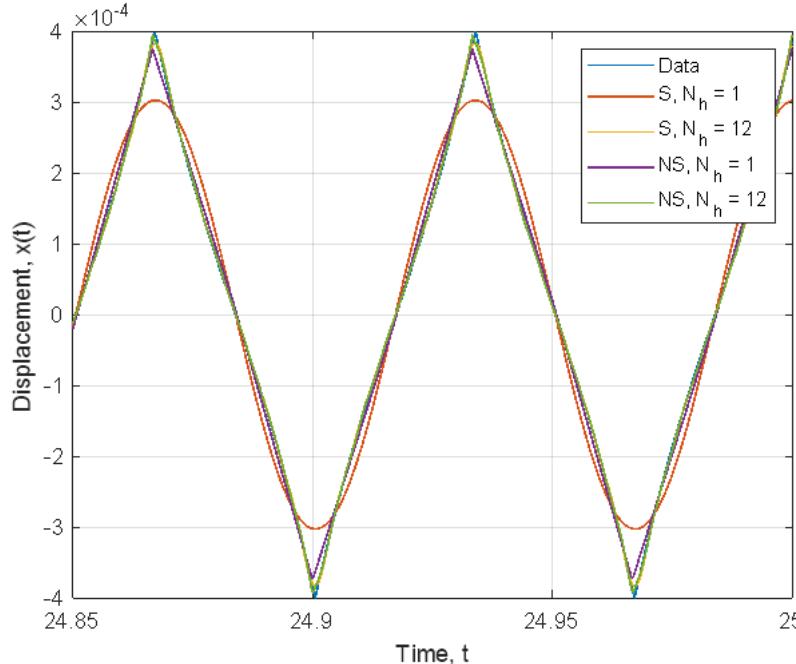
# Least-squares regression analysis



- Both smooth and non-smooth velocities agree well and match closely to the data
- Smooth and non-smooth accelerations agree well away from the points of contact
- Non-smooth acceleration overshoots significantly
- Contact forces agree very well

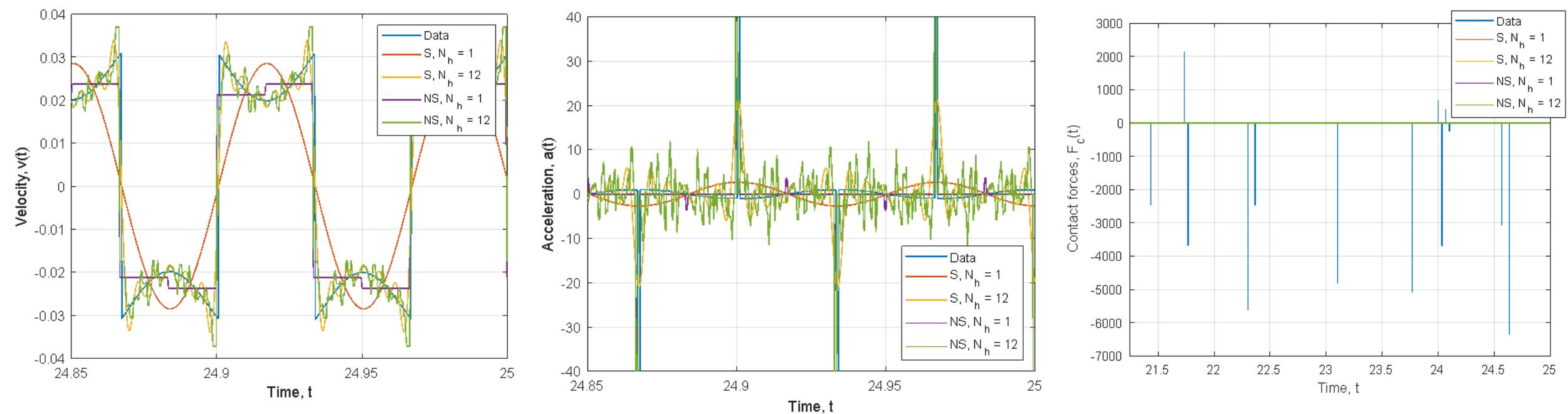
# Least-squares regression analysis

Very hard stiffness,  $K_c = 1.4 * 10^{10} \frac{N}{m}$ ,  $\omega = 15 \text{ Hz}$



- Strongly non-smooth response
- The non-smooth series converges faster until  $N_h = 12$
- For more harmonics, the smooth curve-fit becomes better

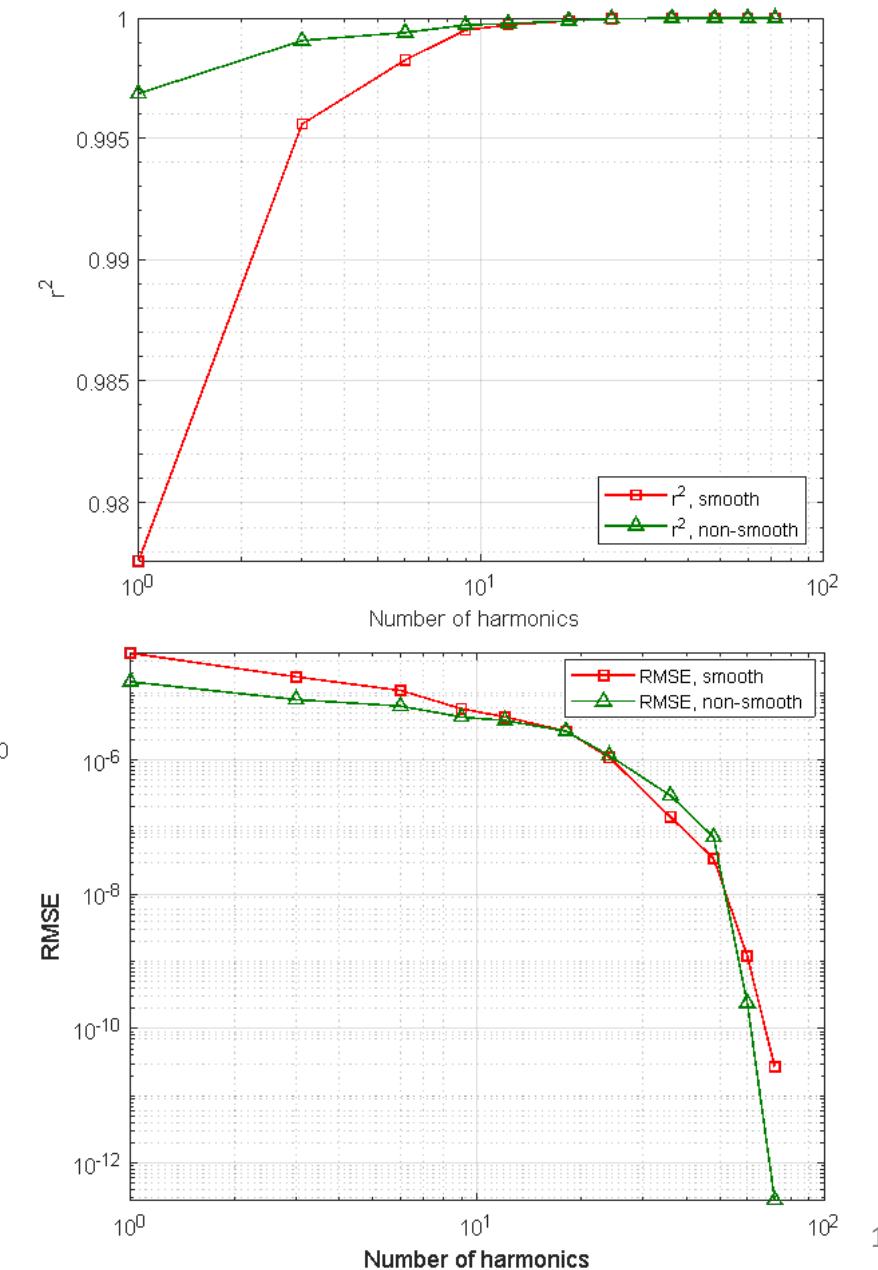
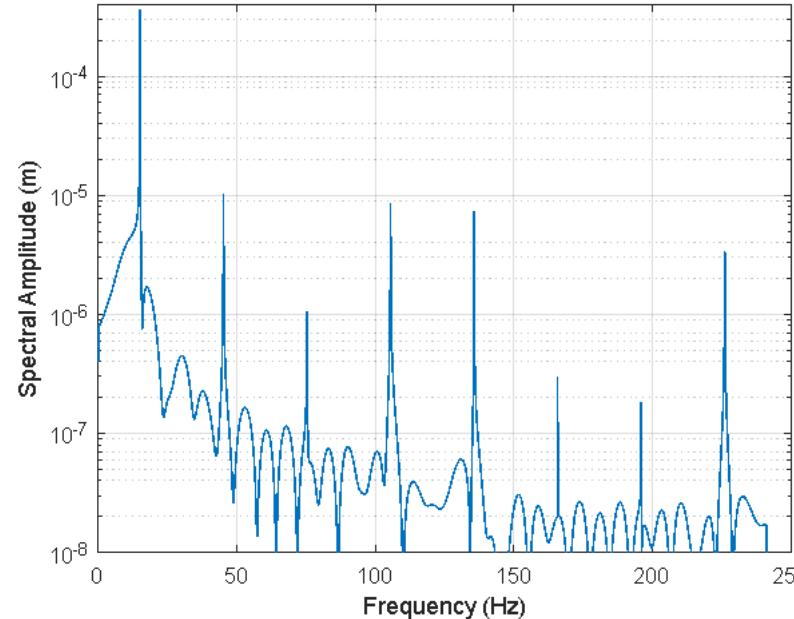
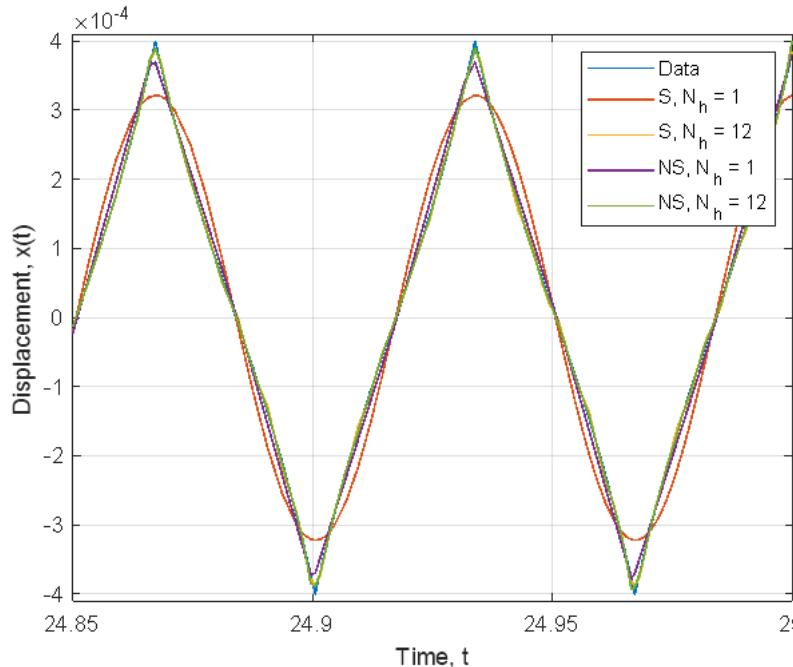
# Least-squares regression analysis



- Velocities are again in agreement with each other and the data
- The non-smooth response best captures the peaks in acceleration
- Poor agreement everywhere else, however, for both smooth and non-smooth
- Contact forces are essentially Dirac impulses by now
- *None* of the curve-fits capture the contact force
- Why? Because the penalty stiffness force definition

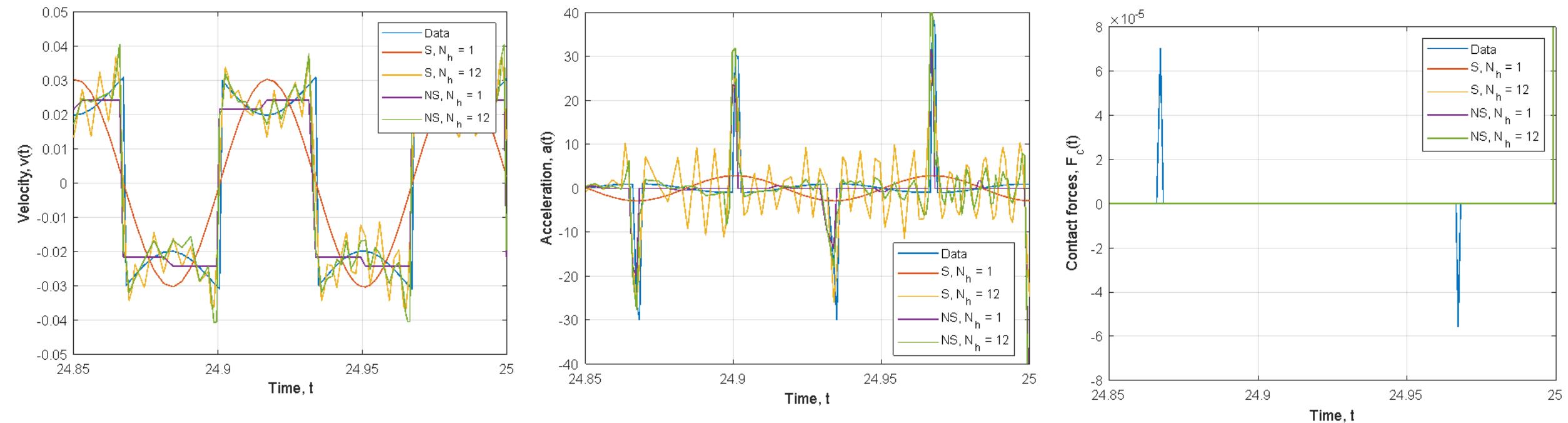
# Least-squares regression analysis

Elastic impact,  $COR = 1, \omega = 15 \text{ Hz}$



- Explicit hard impact using coefficient of restitution instead of penalty stiffness, aka the limit as  $K_c \rightarrow \infty$
- The non-smooth fit converges faster until  $N_h = 18$
- Then the smooth fit until  $N_h \approx 54$
- Then the non-smooth fit again

# Least-squares regression analysis

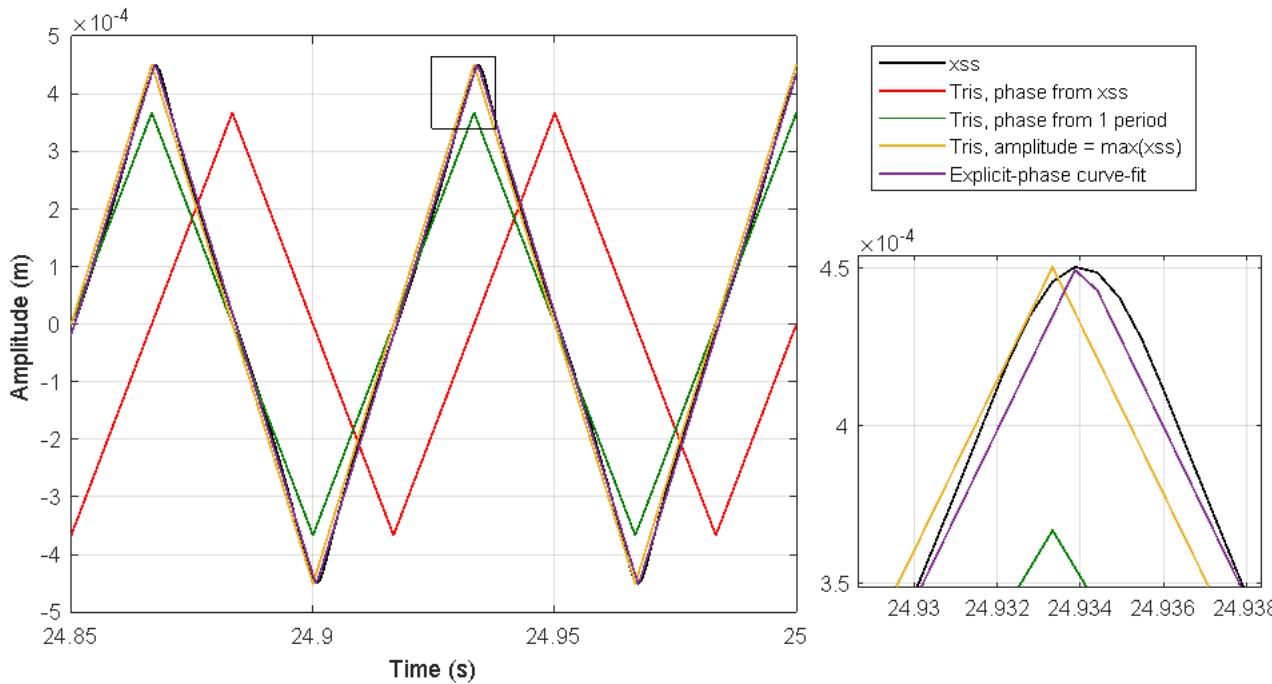


- Velocities are still in good agreement
- Both smooth and non-smooth capture the peaks in acceleration
- Poor agreement everywhere else, still
- Contact forces are Dirac impulses

□ Consider a different curve fit of the following form:

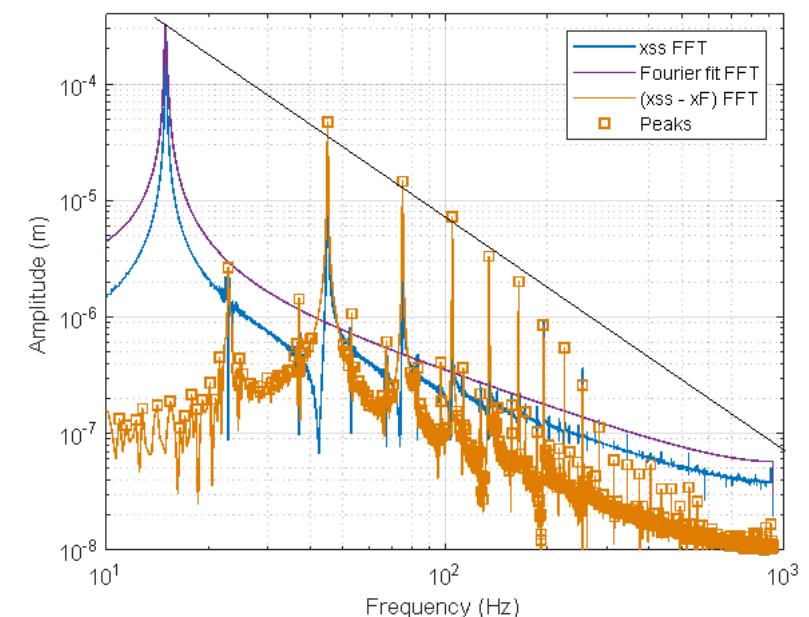
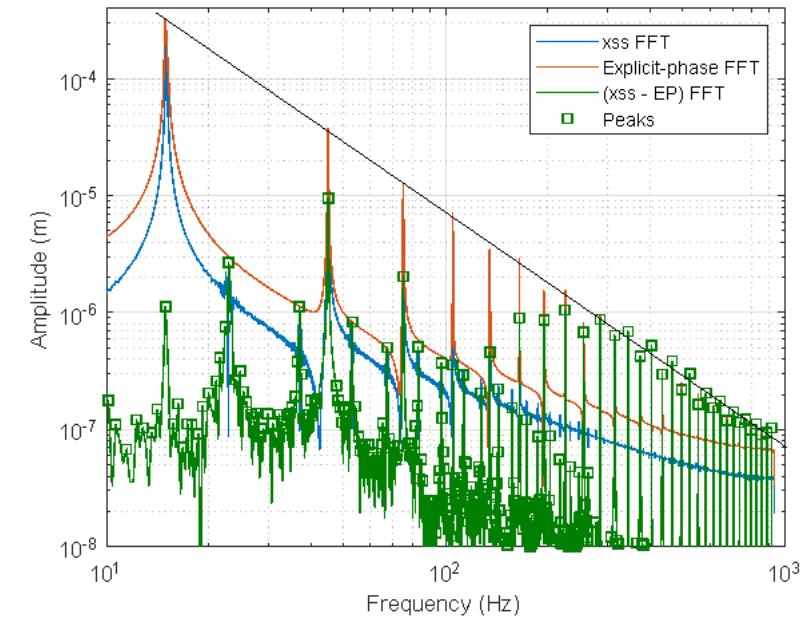
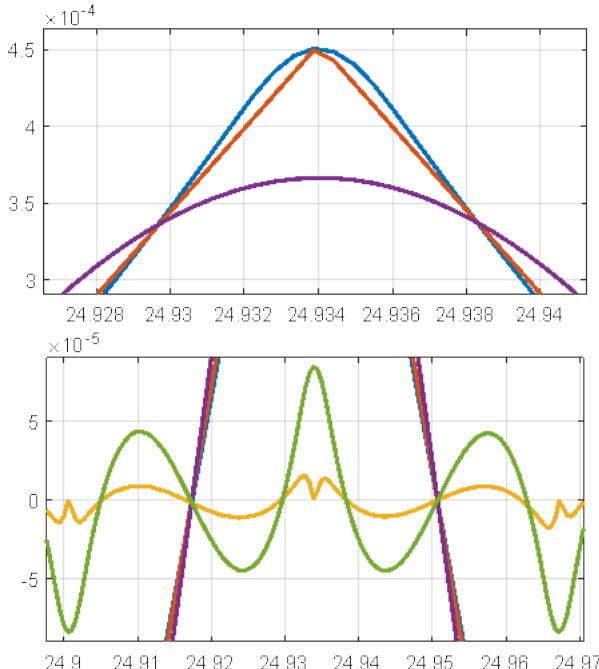
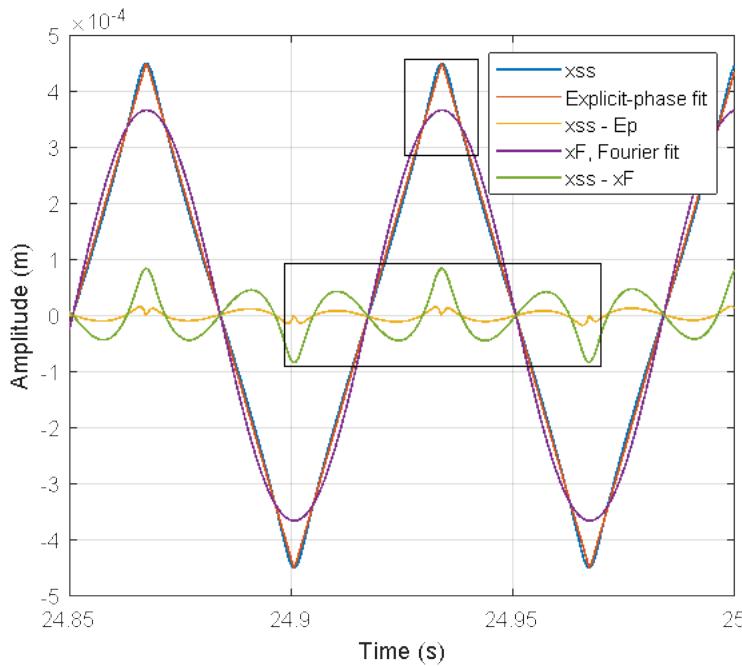
$$f(t) = a_0 + \sum_{n=1}^{\infty} a_n \text{tris}(n\omega t + \theta_n)$$

- Motivation: to better match up the locations of contacts/impacts with the phase-shifts that would produce them, hopefully reducing the number of terms needed
- Every additional non-smooth term means additional contacts and impulses/Dirac combs in the acceleration, which becomes unrealistic



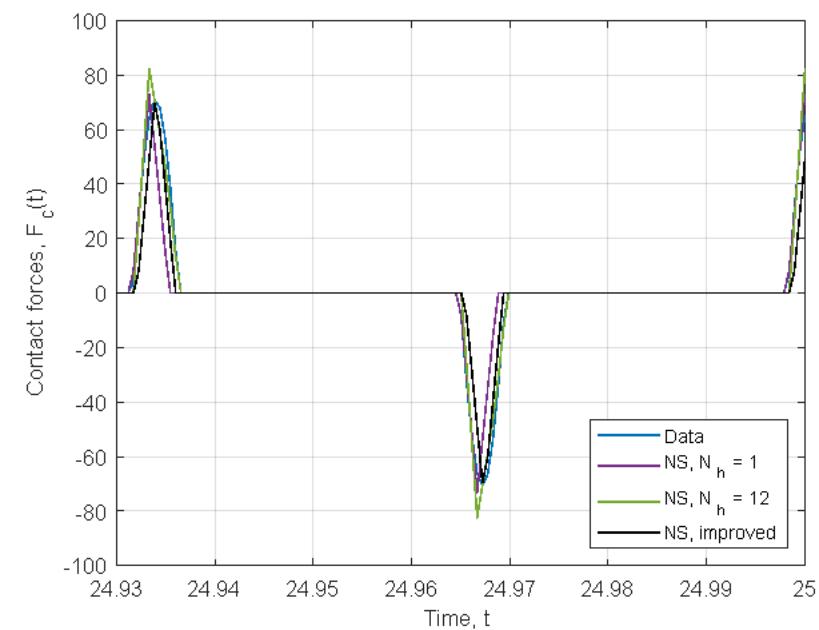
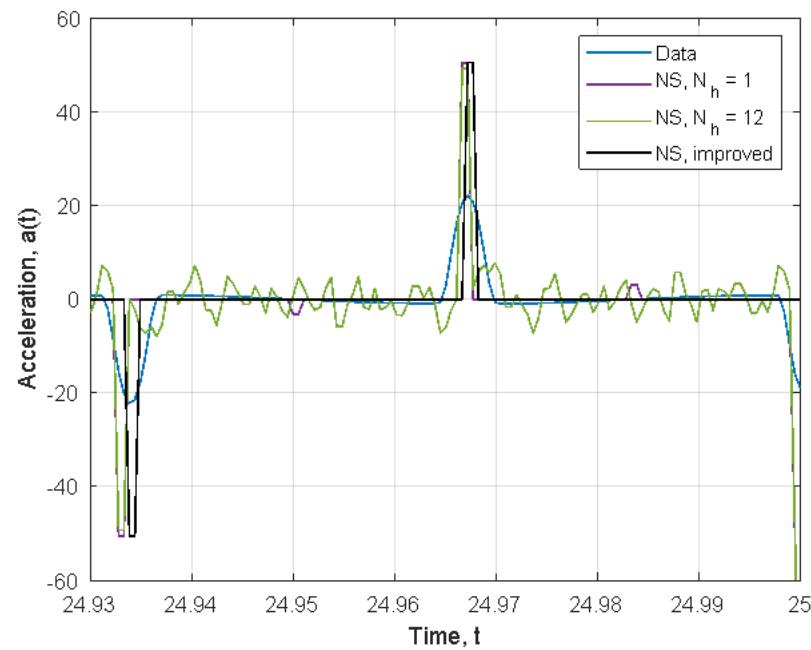
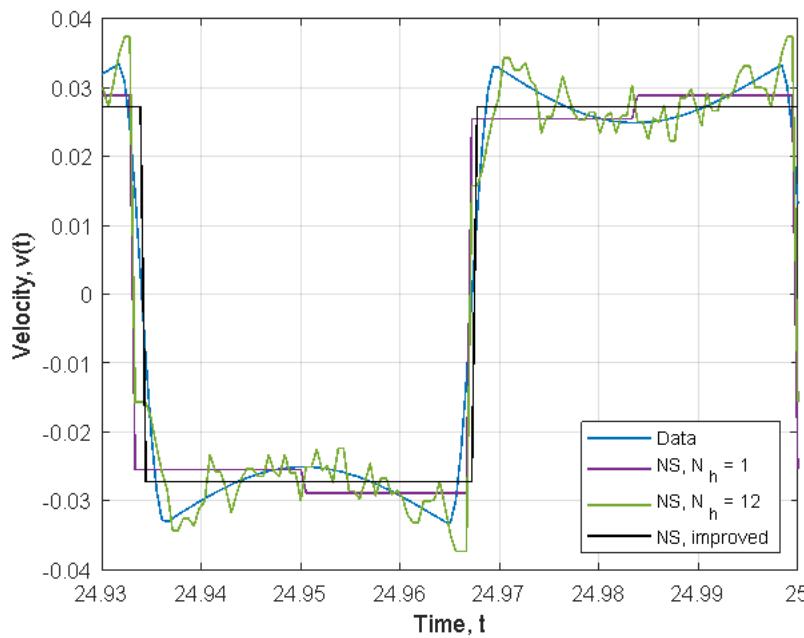
- Left: hard-contact case with different curve fits of a single triangle function
- 3 different manual curve fits
- 4<sup>th</sup> curve fit using Matlab nonlinear least squares
- First 3 tend to capture only amplitude or phase accurately, not both
- 4<sup>th</sup> one captures both well

# Additional studies



- Compare the improved curve fit to the Fourier fit for  $N_h = 1$
- Original tris:  $RMSE = 1.22 * 10^{-5}$
- Fourier:  $RMSE = 3.88 * 10^{-5}$
- New tris:  $RMSE = 8.55 * 10^{-6}$
- 1.5 times better now!

## Additional studies



- The improved non-smooth fit keeps good agreement with the velocity, acceleration, and contact force
- Further evidence that more terms does not necessarily improve the fit if performed naively

## Conclusions

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- Evaluated the usefulness of non-smooth basis functions for obtaining the response of a contact/impact system
- Non-smooth, triangular sine and cosine functions were defined
- Mathematical properties were highlighted
- Applied curve fits to time histories of a contact/impact system and studied for quality
- Results show Fourier series is superior for smooth responses, as expected
- The non-smooth series becomes superior for increasingly non-smooth responses
- Fourier series tends to become more accurate again when many harmonics are used
- A modified series form showed better results than the original naïve series form

## Future work

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- ❑ Continue studying non-smooth series representation and how to improve accuracy
- ❑ Optimal combinations of smooth and non-smooth terms based on when Fourier series regains highest accuracy
- ❑ Mathematical properties amenable to addition in a harmonic balance code

# Acknowledgements

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**Thank you for your attention!  
Please ask any questions**

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