

# Spectral Gaps via Imaginary Time

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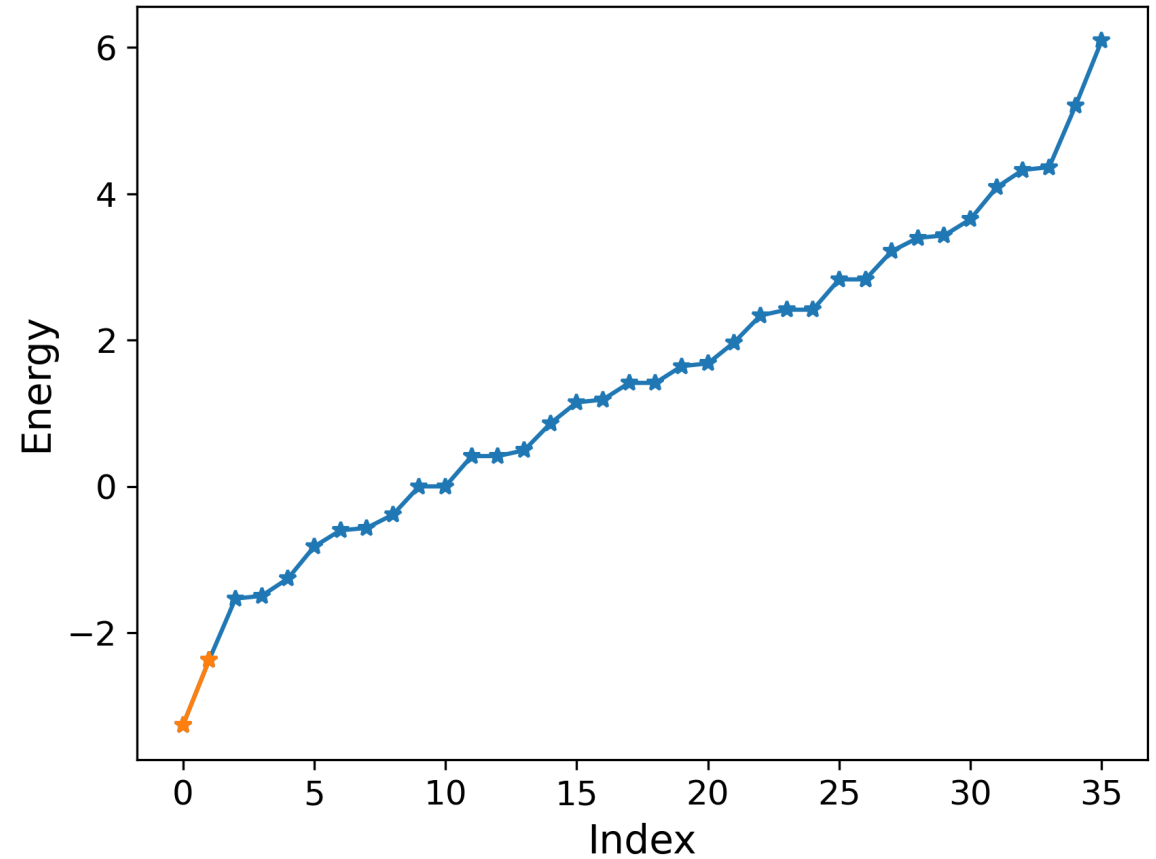
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<https://arxiv.org/abs/2303.02124>

# Spectral Gap

- Defined as:  $\Delta E = E_1 - E_0$



# Spectral Gap

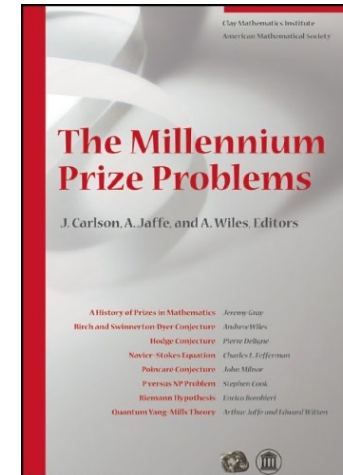
- Topological spin liquid phase
  - Artificial gauge fields
  - Superconductivity
  - Photosynthesis
  - Solar cells
- Fluorescence / Phosphorescence
  - Transfer of quantum states
  - Quark confinement
  - and more!

## Spin liquids in frustrated magnets

Leon Balents<sup>1</sup>

## Fractionalized excitations in the spin-liquid state of a kagome-lattice antiferromagnet

Tian-Heng Han<sup>1</sup>, Joel S. Helton<sup>2</sup>, Shaoyan Chu<sup>3</sup>, Daniel G. Nocera<sup>4</sup>, Jose A. Rodriguez-Rivera<sup>2,5</sup>, Collin Broholm<sup>2,6</sup> & Young S. Lee<sup>1</sup>



**Excitation and entanglement transfer  
versus spectral gap**

**M J Hartmann<sup>1,2,3</sup>, M E Reuter<sup>1,2</sup>, and M B Plenio<sup>1,2</sup>**

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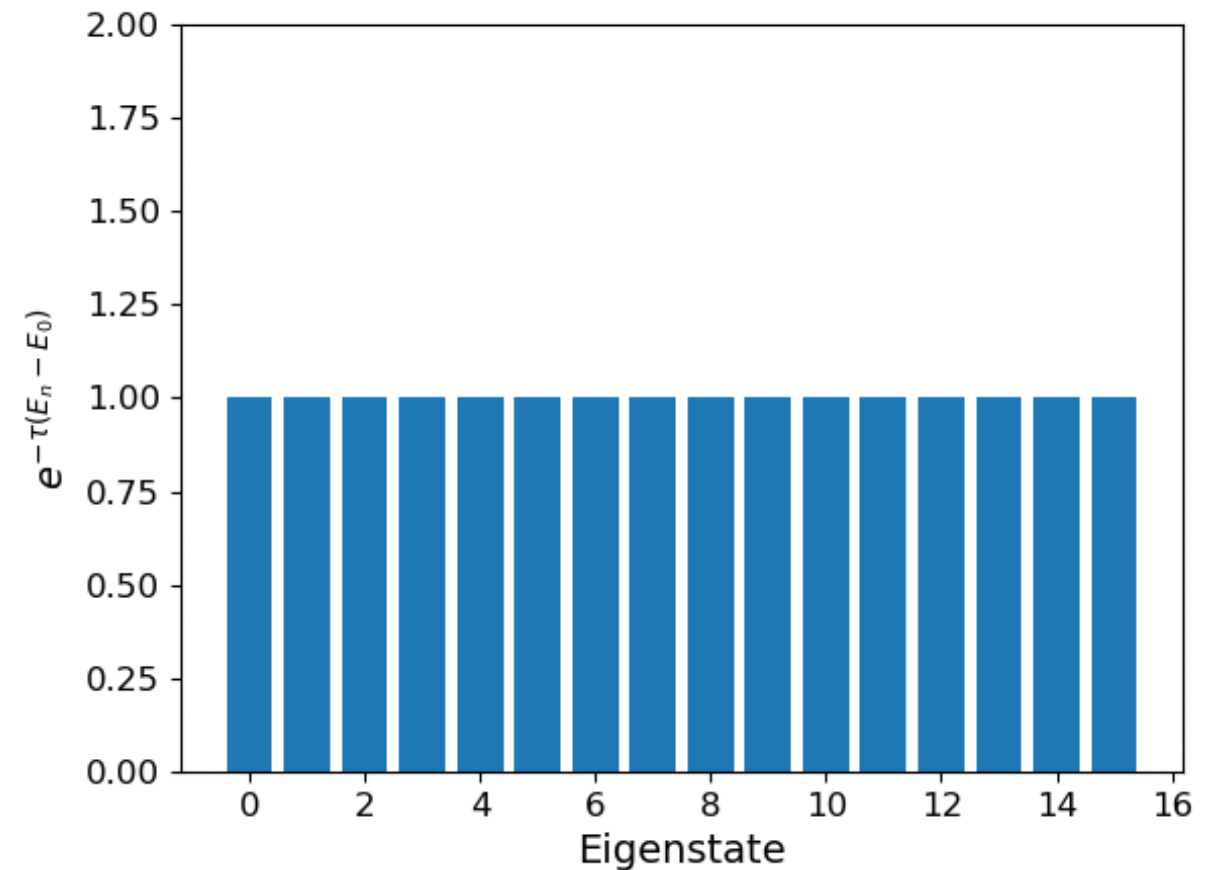
# Imaginary Time Evolution

$$|\psi(t)\rangle = \sum_{n=0}^N e^{-itE_n} c_n |n\rangle \quad \xrightarrow{t = -i\tau} \quad |\psi(\tau)\rangle = \sum_{n=0}^N e^{-\tau E_n} c_n |n\rangle$$

$$|\psi(\tau)\rangle = c_0 |0\rangle + \sum_{n=1}^N e^{-\tau(E_n - E_0)} c_n |n\rangle$$

# Imaginary Time Evolution

$$|\psi(\tau)\rangle = c_0 |0\rangle + \sum_{n=1}^N e^{-\tau(E_n - E_0)} c_n |n\rangle$$



# Our Method

Definitions:

$$\hat{H} = \sum_{n=0}^N E_n \hat{\Pi}_n$$

$$|\psi(\tau)\rangle = \sum_{n=0}^N e^{-\tau E_n} \hat{\Pi}_n |\phi_0\rangle$$

$$[\hat{H}, \hat{O}]_M = \underbrace{[\hat{H}, [\hat{H}, \dots, [\hat{H}, \hat{O}]] \dots]}_{M\text{-times}}$$

For 3x3 H with *non-degenerate* spectrum:

$$\hat{H} = E_0 \hat{\Pi}_0 + E_1 \hat{\Pi}_1 + E_2 \hat{\Pi}_2$$

$$\hat{\Pi}_n = |n\rangle\langle n| \quad \delta_{nm} = \langle n|m\rangle$$

For 3x3 H with *degenerate* spectrum:

$$\hat{H} = E_0 \hat{\Pi}_0 + E_1 \hat{\Pi}_1$$

$$\hat{\Pi}_0 = |0\rangle\langle 0| + |1\rangle\langle 1|$$

$$\hat{\Pi}_1 = |2\rangle\langle 2|$$

# Our Method

Definitions:

$$\hat{H} = \sum_{n=0}^N E_n \hat{\Pi}_n \quad \hat{\Pi}_n = |n\rangle\langle n|$$

$$|\psi(\tau)\rangle = \sum_{n=0}^N e^{-\tau E_n} \hat{\Pi}_n |\phi_0\rangle$$

$$[\hat{H}, \hat{O}]_M = [\hat{H}, [\hat{H}, \dots, [\hat{H}, \hat{O}] \dots]]$$

Conditions:

$$[\hat{H}, \hat{O}]_M \neq 0$$

$$[\hat{H}, \hat{O}]_{M+2} \neq 0$$

$$\langle \phi_0 | \hat{\Pi}_0 \hat{O} \hat{\Pi}_1 | \phi_0 \rangle \neq 0$$

$$\langle \psi(\tau) | [\hat{H}, \hat{O}]_M | \psi(\tau) \rangle = \sum_{l=0}^N \sum_{k=0}^N e^{-\tau(E_l + E_k)} (E_l - E_k)^M \langle \phi_0 | \hat{\Pi}_l \hat{O} \hat{\Pi}_k | \phi_0 \rangle$$

# Our Method

- Leading order terms:

$$\langle \psi(\tau) | [\hat{H}, \hat{O}]_M | \psi(\tau) \rangle = e^{-\tau(E_0 + E_1)} (E_0 - E_1)^M \langle \phi_0 | \hat{\Pi}_0 \hat{O} \hat{\Pi}_1 + (-1)^M h.c. | \phi_0 \rangle + \mathcal{O}(e^{-\tau(E_0 + E_2)})$$

- Can obtain

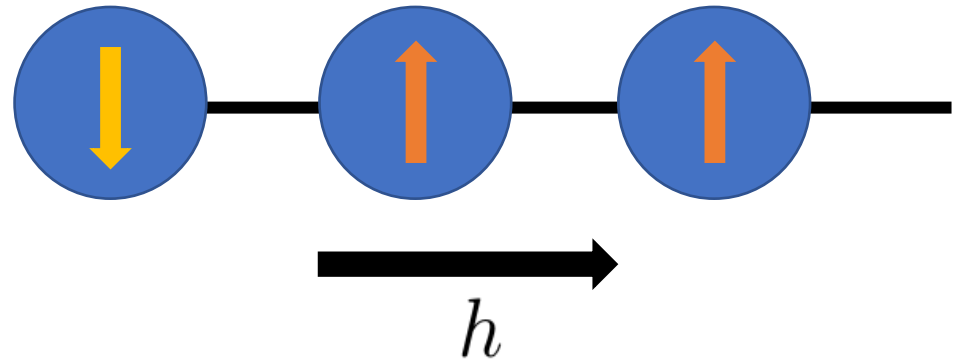
$$\frac{\langle \psi(\tau) | [\hat{H}, [\hat{H}, [\hat{H}, \hat{O}]]] | \psi(\tau) \rangle}{\langle \psi(\tau) | [\hat{H}, \hat{O}] | \psi(\tau) \rangle} = (E_0 - E_1)^2 + \mathcal{O}(e^{-\tau(E_2 - E_1)})$$



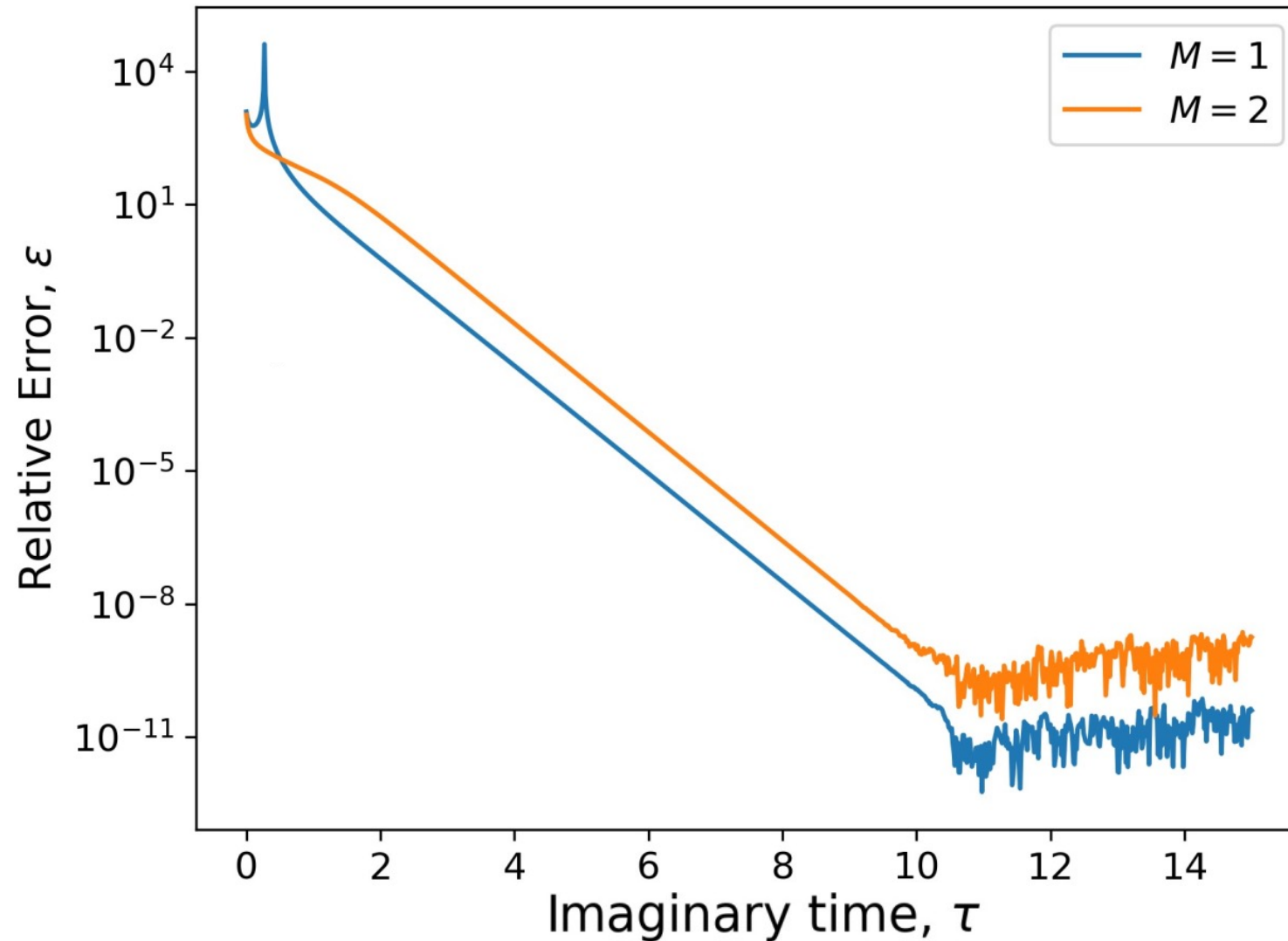
# Transverse Field Ising Model

$$\hat{H} = \sum_{k=0}^{K-1} -J\sigma_{k+1}^z\sigma_k^z - h\sigma_k^x$$

- Very simple model
- Nearest neighbor interactions governed by spins



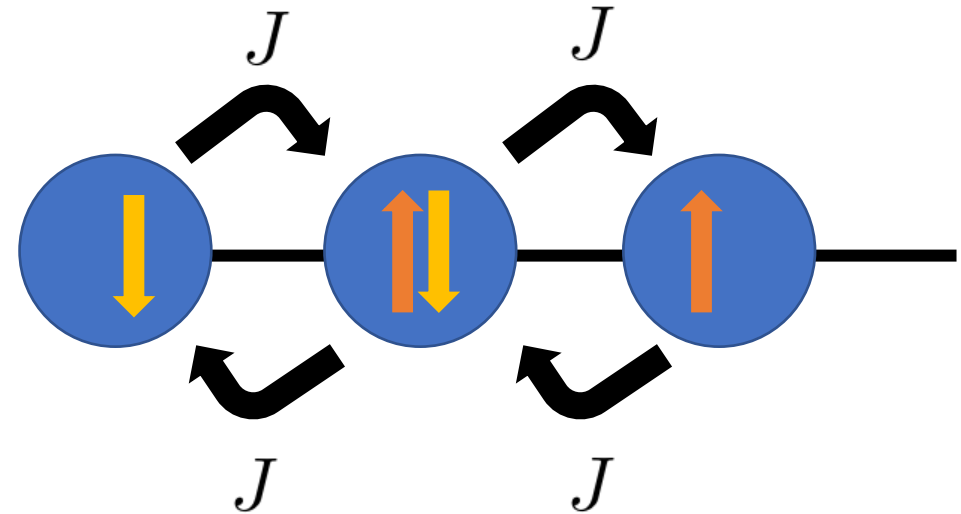
# Transverse Field Ising Model Results



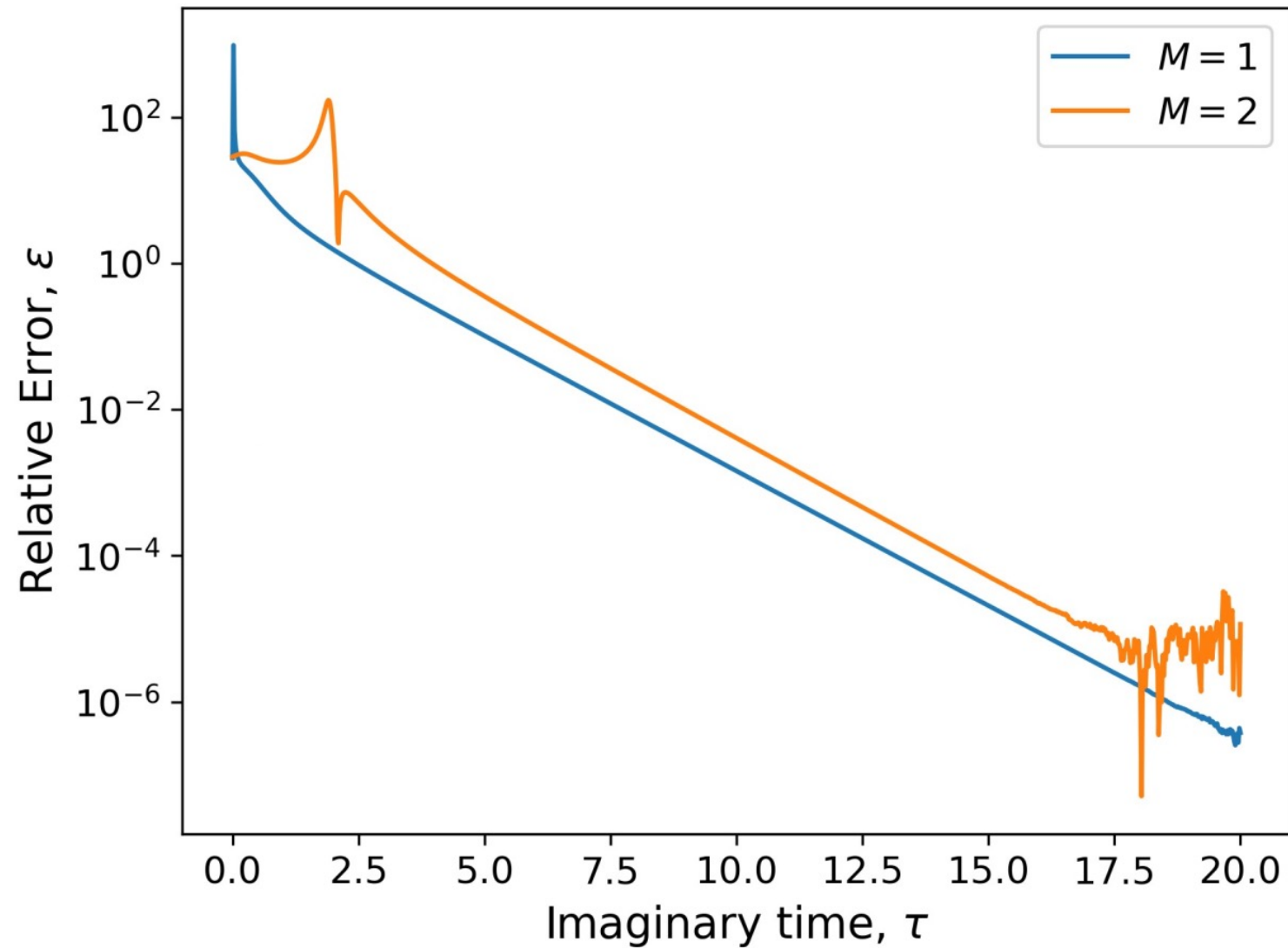
# Fermi-Hubbard Model

$$\hat{H} = -J \sum_{k=0, \sigma}^{K-1} (c_{k, \sigma}^{\dagger} c_{k+1, \sigma} - c_{k, \sigma} c_{k+1, \sigma}^{\dagger}) - \mu \sum_{k=0, \sigma}^{K-1} n_{k \sigma} + U \sum_{k=0}^{K-1} n_{k \uparrow} n_{k \downarrow}$$

- Simple model for correlated systems



# Fermi-Hubbard Model Results



# Conclusion

- Presented a method for calculating spectral gaps using imaginary time propagation to quench out higher energy states
- Demonstrated effectiveness using two paradigmatic models
- More details -> <https://arxiv.org/abs/2303.02124>