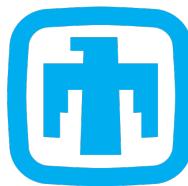


Core contributions to stopping powers in warm dense matter

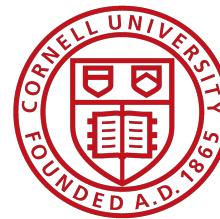
Alina Kononov¹, Thomas Hentschel²,
Stephanie B. Hansen¹, and Andrew D. Baczewski¹



**Sandia
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¹ Sandia National Laboratories

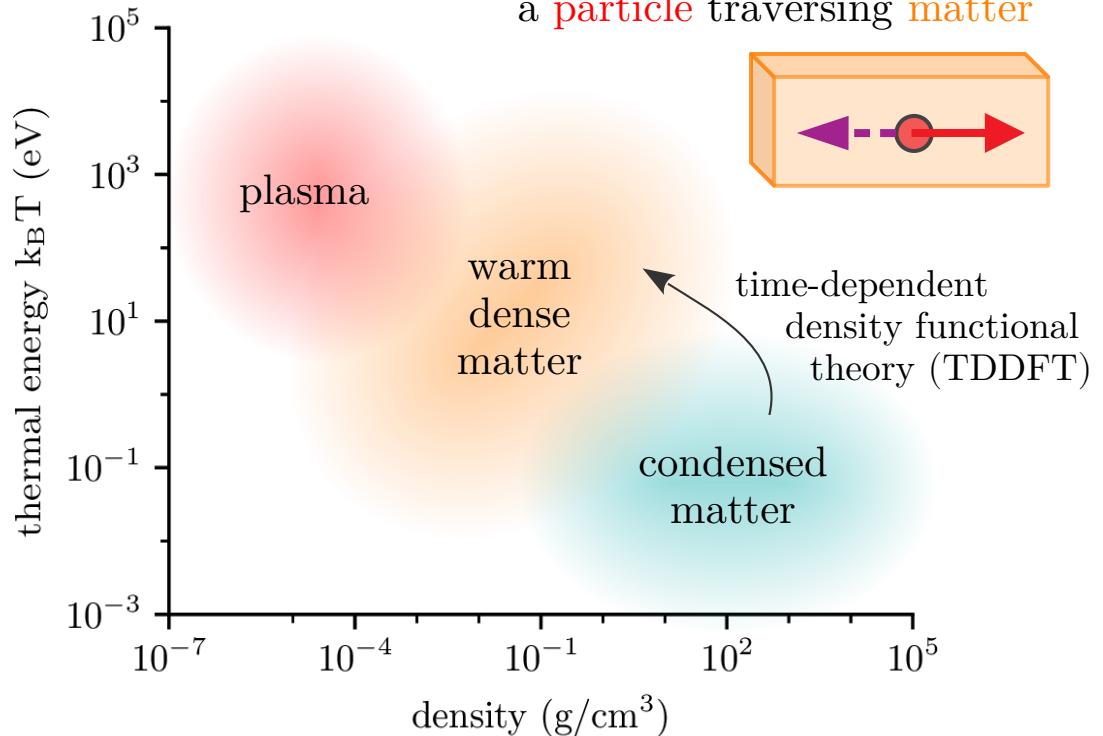
² Cornell University



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Bridging the gap between plasma and condensed matter

- Fusion fuel passes through poorly understood WDM regime on the way to ignition
 - competing physics challenge models
 - scarce experimental data, uncertain conditions limit validation
- Self-heating through electronic stopping influences ignition requirements
- This work: first-principles benchmark data using TDDFT
- Also studying other transport properties
 - XRTS: arXiv:2109.09576
 - conductivity: Robinson, B61.07



Electronic stopping powers from TDDFT

- Mean-field model of quantum electron dynamics
- Initial condition: equilibrium state from Mermin-DFT
- Evolve electron density $n(\mathbf{r}, t)$ over time
- Stopping power $\sim \mathbf{F}_\alpha[n](t)$

$$i \frac{\partial}{\partial t} \phi_j(\mathbf{r}, t) = \hat{H}[n(\mathbf{r}, t)] \phi_j(\mathbf{r}, t)$$

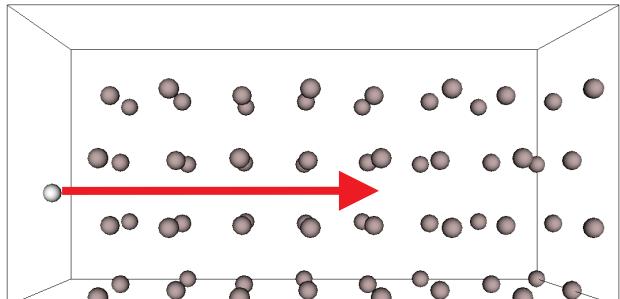
$$n(\mathbf{r}, t) = \sum_j f_j(T) |\phi_j(\mathbf{r}, t)|^2$$

$$\hat{H}[n(\mathbf{r}, t)](t) = -\frac{\nabla^2}{2} + V_{\text{ext}}(t) + \int \frac{n(\mathbf{r}', t)}{|\mathbf{r} - \mathbf{r}'|} d\mathbf{r}'^3 + V_{\text{xc}}[n(\mathbf{r}, t)]$$

projectile trajectory
is an approximation!

Olmstead, N59.05

arXiv preprint coming soon



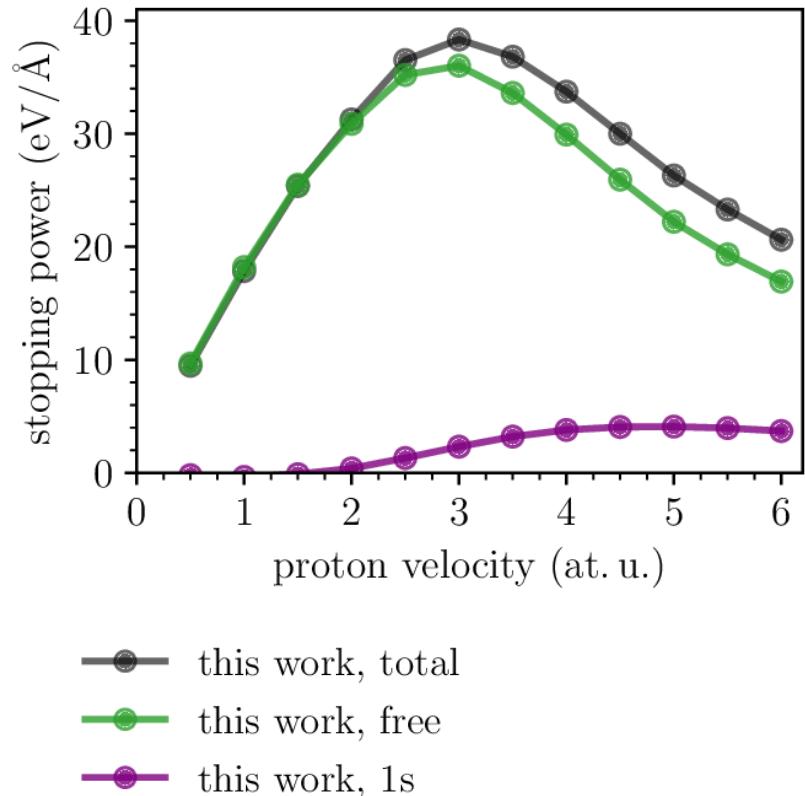
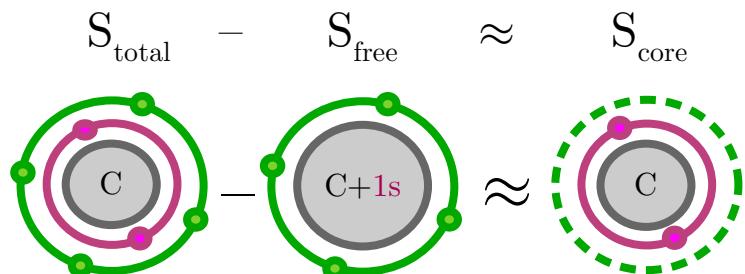
kinetic energy
external potential due to ions
Coulombic e-e interaction
exchange-correlation

pseudopotential
approximation
allows detailed
insights

adiabatic
local density
approximation

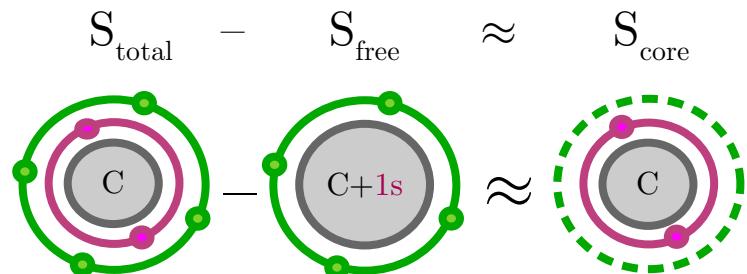
Orbital-resolved contributions through pseudization

- C at 10 g/cc, 10 eV



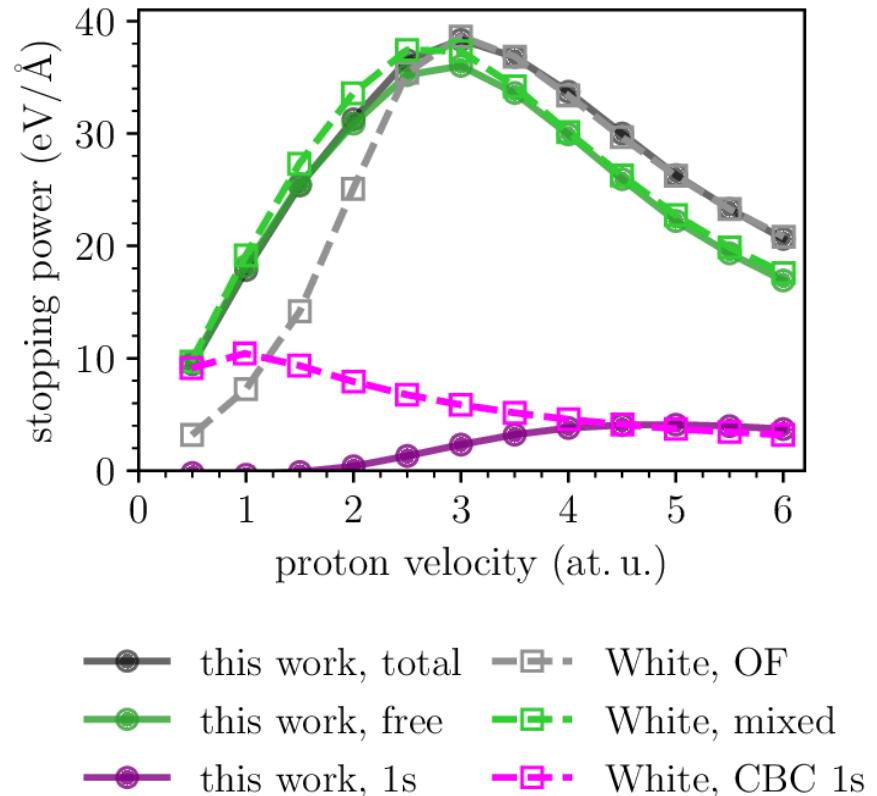
Simplified models of core contributions lose accuracy near/below Bragg peak

- C at 10 g/cc, 10 eV
- Reasonable agreement across TDDFT flavors
- Opposite trends in TDDFT and CBC 1s contributions
- Core contributions warrant further scrutiny



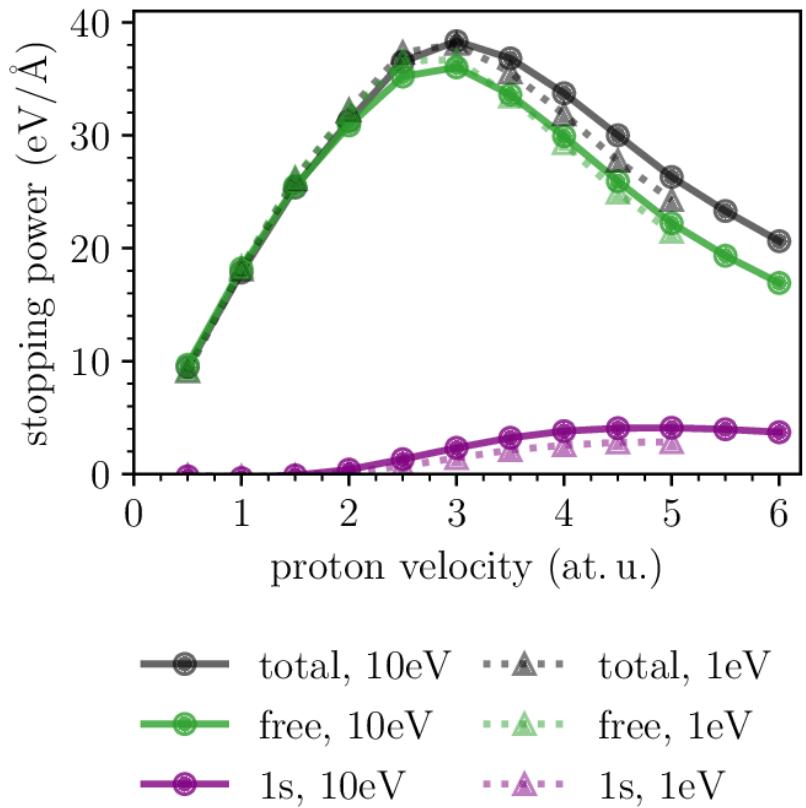
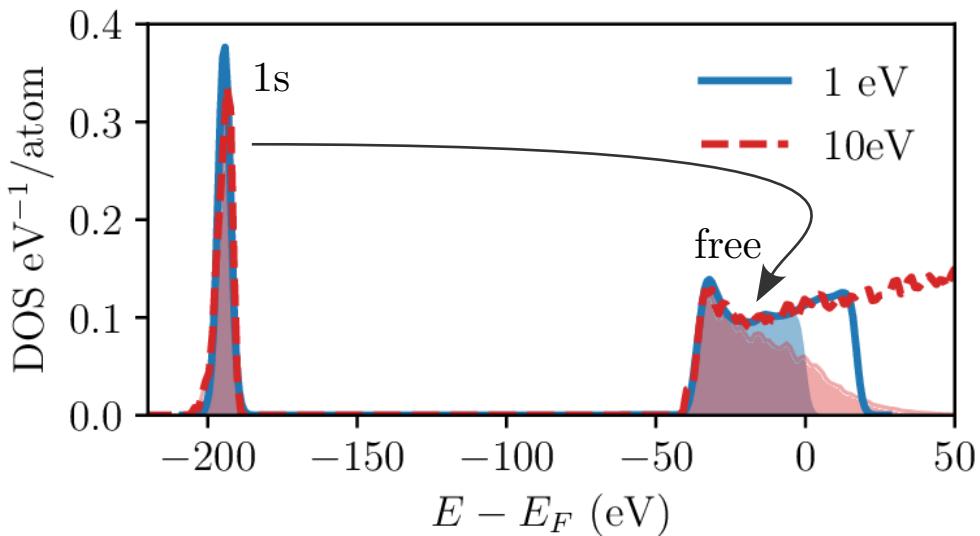
White, J. Phys. Condens. Matter 34 (2022)

Barriga-Carrasco and Casas, Laser Part. Beams 31 (2013)



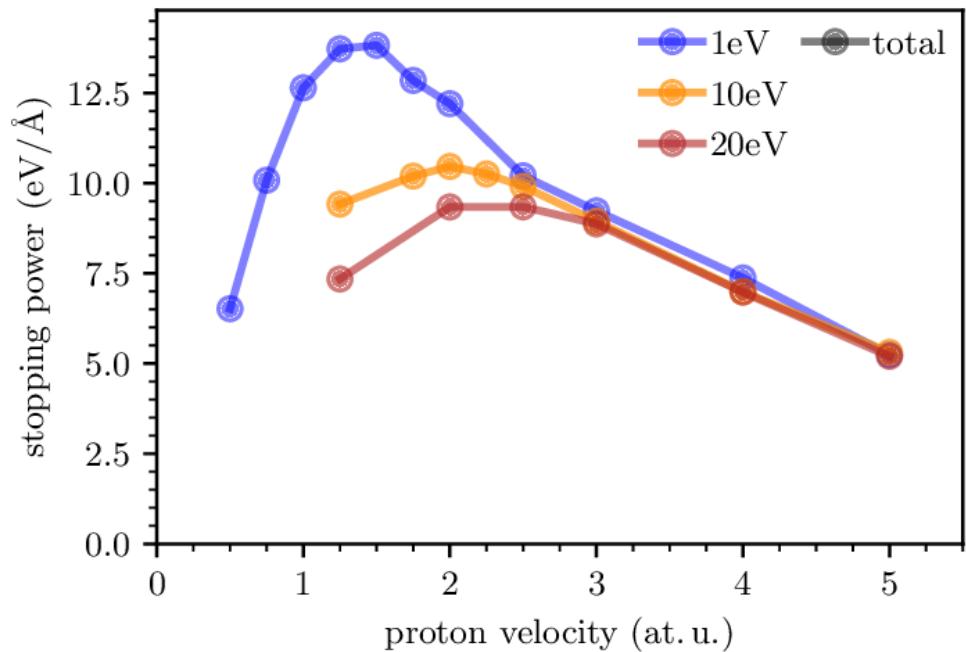
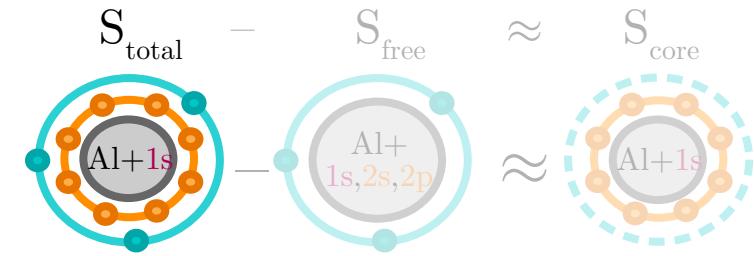
Electron temperature enhances core contributions in C

- ~50% higher 1s contribution at $T_e = 10\text{eV}$ vs 1eV
- Thermal vacancies facilitate $1s \rightarrow \text{free}$ excitations



Competing mechanisms affect core contributions in Al

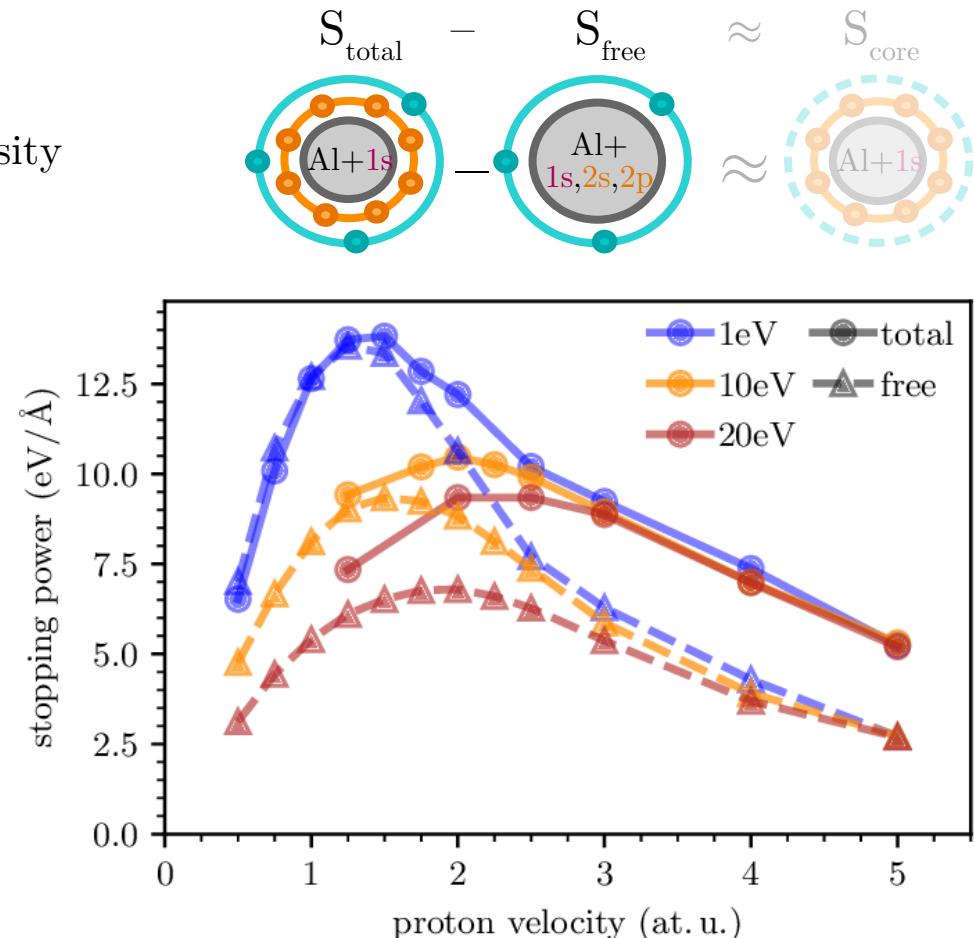
- At high T_e , Bragg peak lowers and shifts to higher velocities



Competing mechanisms affect core contributions in Al

- At high T_e , Bragg peak lowers and shifts to higher velocities
- Thermal excitations increase free-electron density
 - 3e PP underestimates S_{free} at $T_e=20\text{eV}$

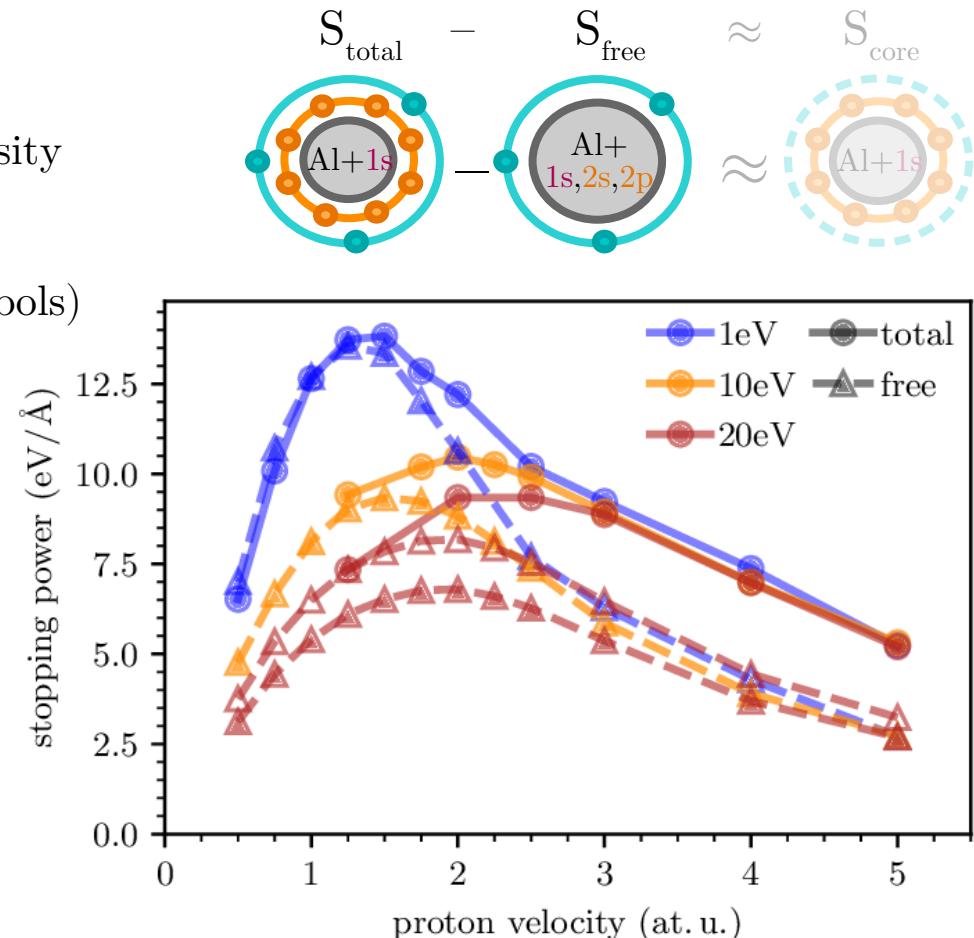
T_e	1eV	10eV	20eV
free e^- per atom	3.00	3.02	3.61



Competing mechanisms affect core contributions in Al

- At high T_e , Bragg peak lowers and shifts to higher velocities
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 - correct with 3.61/3 scale factor (open symbols)

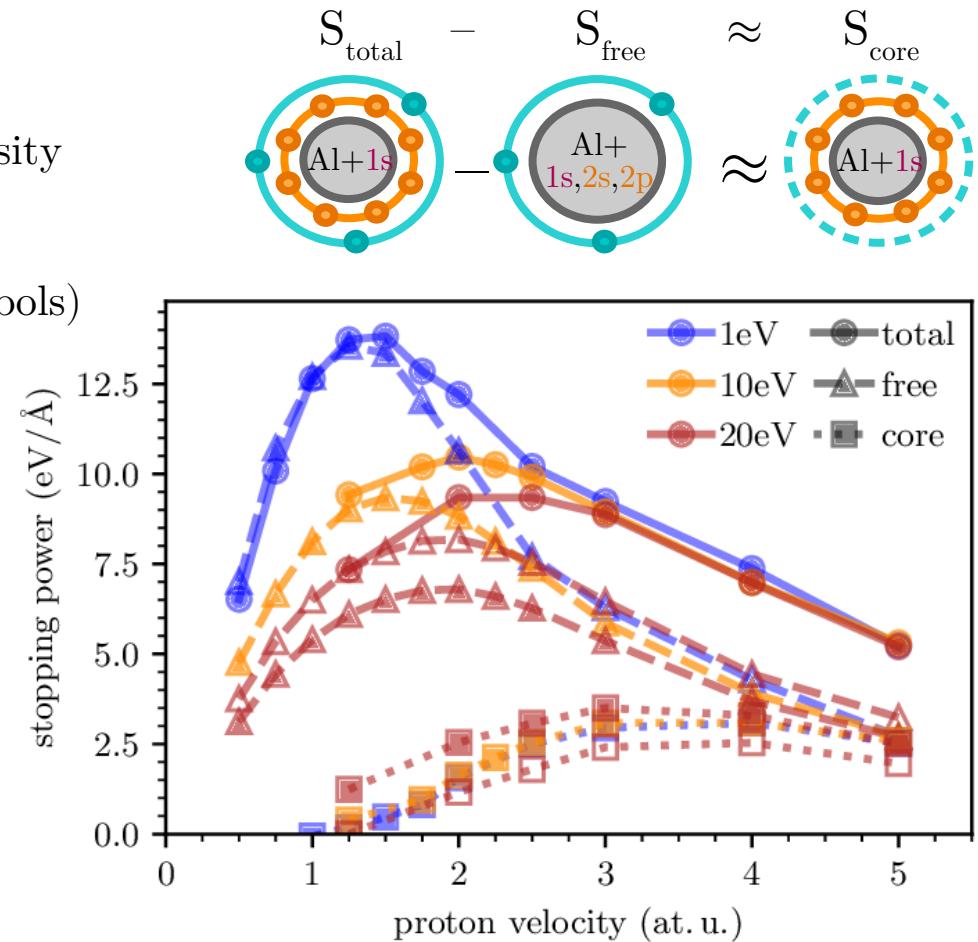
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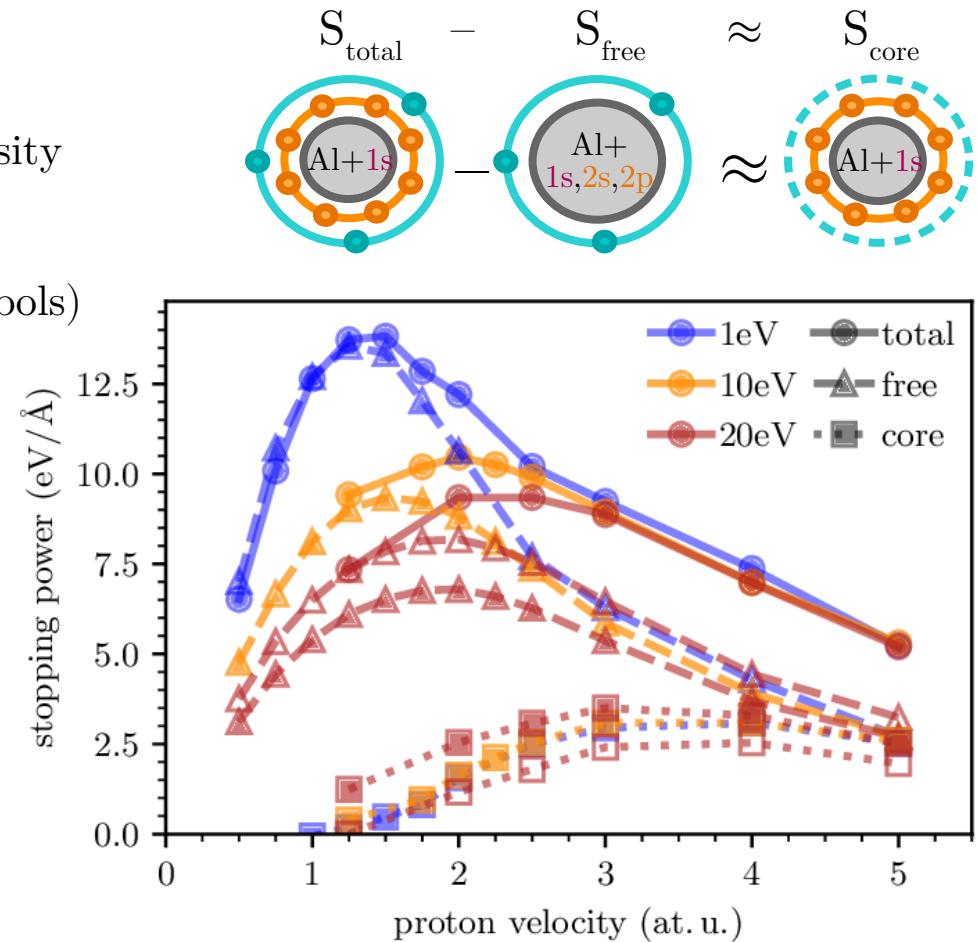
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- Thermal depletion of low-energy free states and deeper 2p binding alter $2\text{p} \rightarrow \text{free } \Delta E$

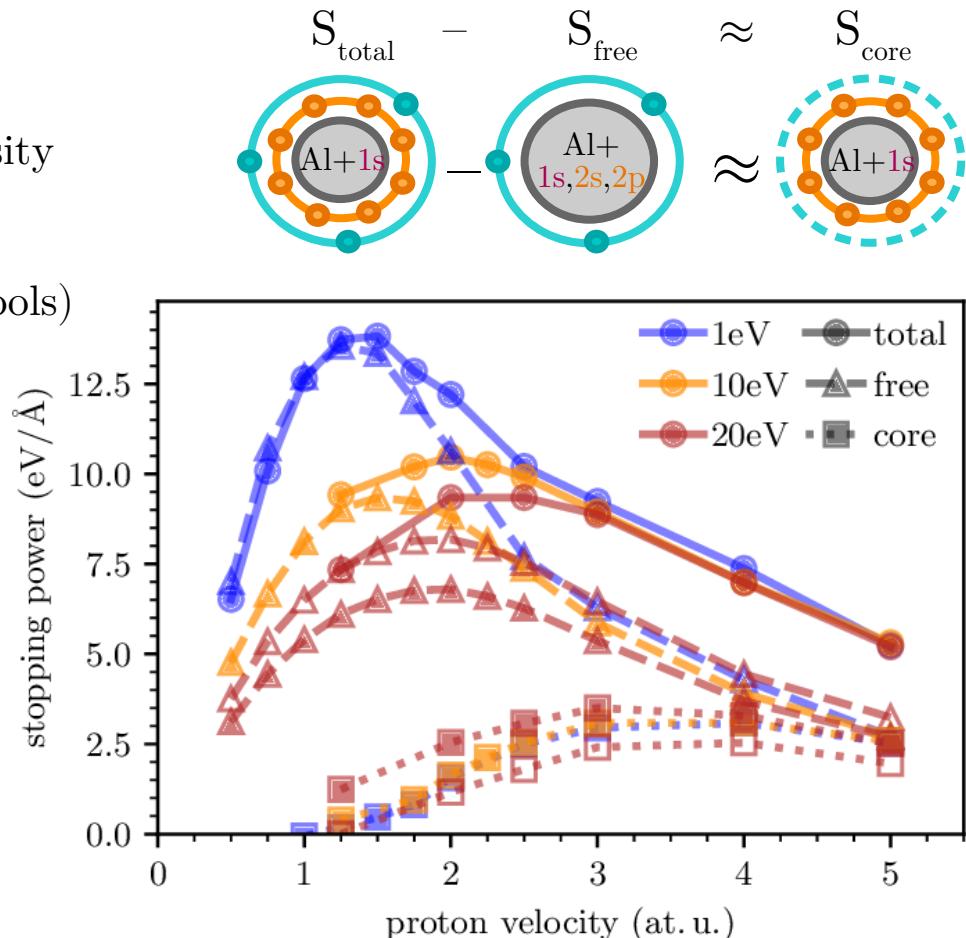
T_e	1eV	10eV	20eV
free e^- per atom	3.00	3.02	3.61
$E_{\text{free}} - E_{2\text{p}}$ (eV)	65	55	62.5



Competing mechanisms affect core contributions in Al

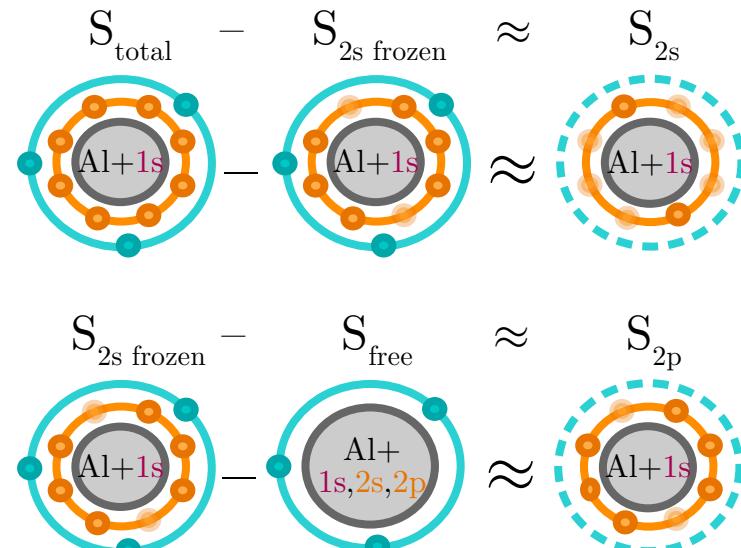
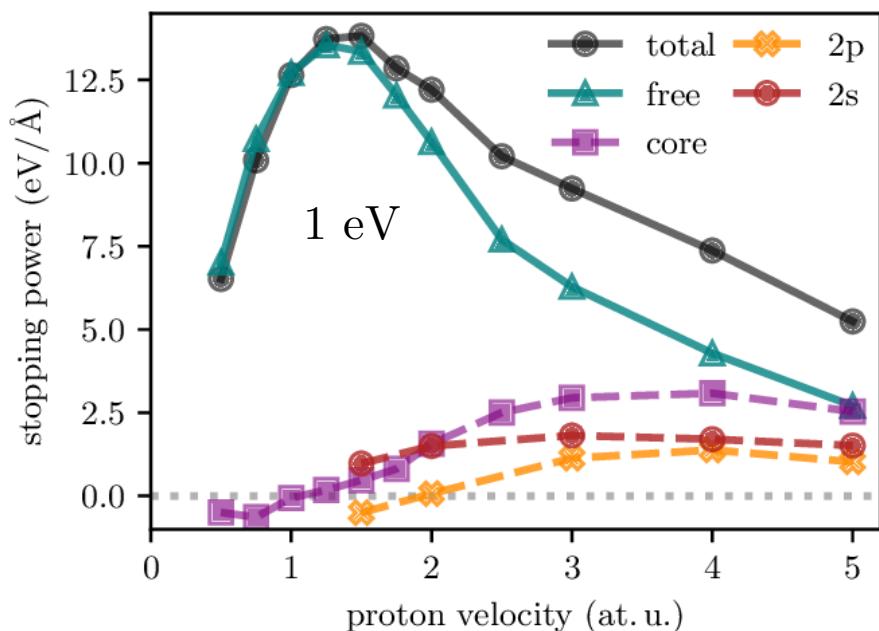
- At high T_e , Bragg peak lowers and shifts to higher velocities
- Thermal excitations increase free-electron density
 - 3e PP underestimates S_{free} at $T_e=20\text{eV}$
 - correct with $3.61/3$ scale factor (open symbols)
- Thermal depletion of low-energy free states and deeper 2p binding alter $2\text{p} \rightarrow \text{free } \Delta E$
- Thermal depletion of 2p allows $2\text{s} \rightarrow 2\text{p}$ at $T_e=20\text{eV}$

T_e	1eV	10eV	20eV
free e^- per atom	3.00	3.02	3.61
$E_{\text{free}} - E_{2\text{p}}$ (eV)	65	55	62.5
2p vacancy (%)	0	0.5	9.6



Surprising orbital-resolved core contributions in Al

- Further disentangle by freezing 2s
- 2s contribution surprisingly large: symmetry breaking of 2p excitations
- 2s onset surprisingly low v: inconsistent with threshold models
- Negative values: core excitations can inhibit free response

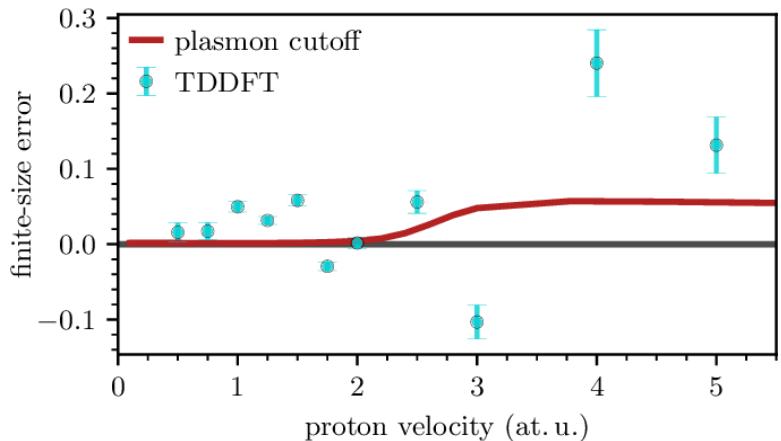


Finite-size errors beyond plasmonic interpretation

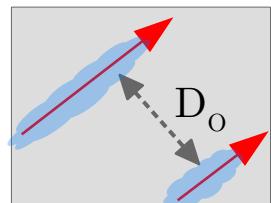
- Trajectory-dependent and non-monotonic finite-size effects inconsistent with plasmonic model

$$S(v) = \frac{2Z^2}{\pi v^2} \int_0^\infty \frac{dk}{k} \int_0^{kv} d\omega \omega \text{Im} \left[\frac{-1}{\epsilon(k, \omega)} \right]$$
$$k_{\text{cut}} = 2\pi/L$$

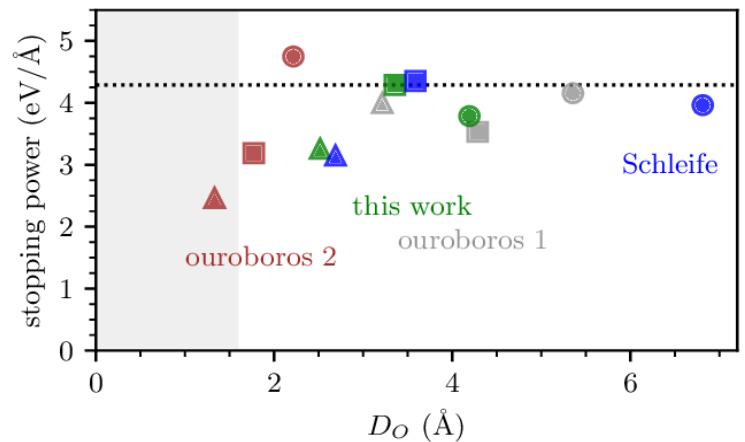
see Hentschel et al., arXiv:2301.09700



- Attribute to “ouroboros” effects
- Propose converging with increasing D_O



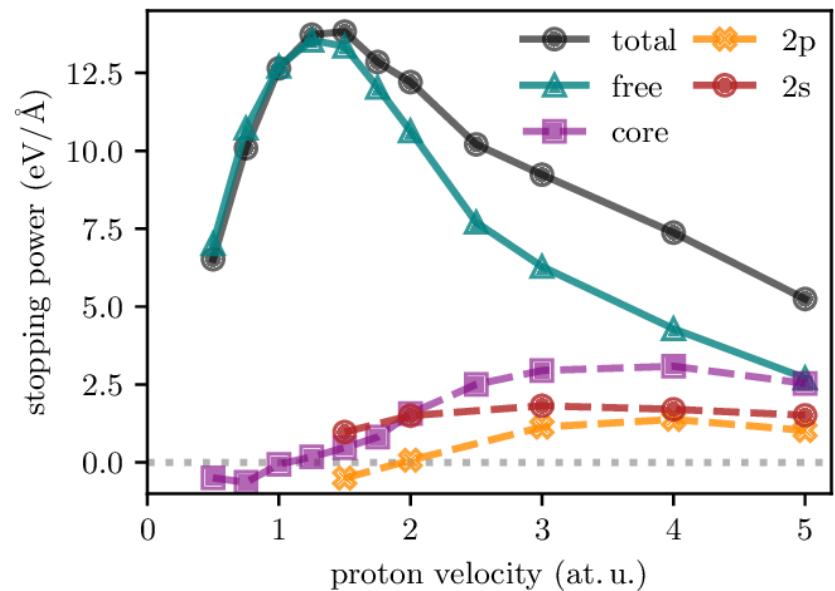
arXiv preprint
coming soon!



Summary and Outlook

- Simplified models incorrectly describe onset of core contributions to stopping power
- Isolating orbital-resolved contributions through pseudization in TDDFT
- New interpretation of finite-size errors arising from fictitious interaction with excitations
- In C, thermal vacancies in free electron states enhance 1s contribution
- In Al, competing effects modify 2s+2p contributions at high T_e
- Planning detailed investigation with cheaper systems, e.g., Li and Na

akonono@sandia.gov



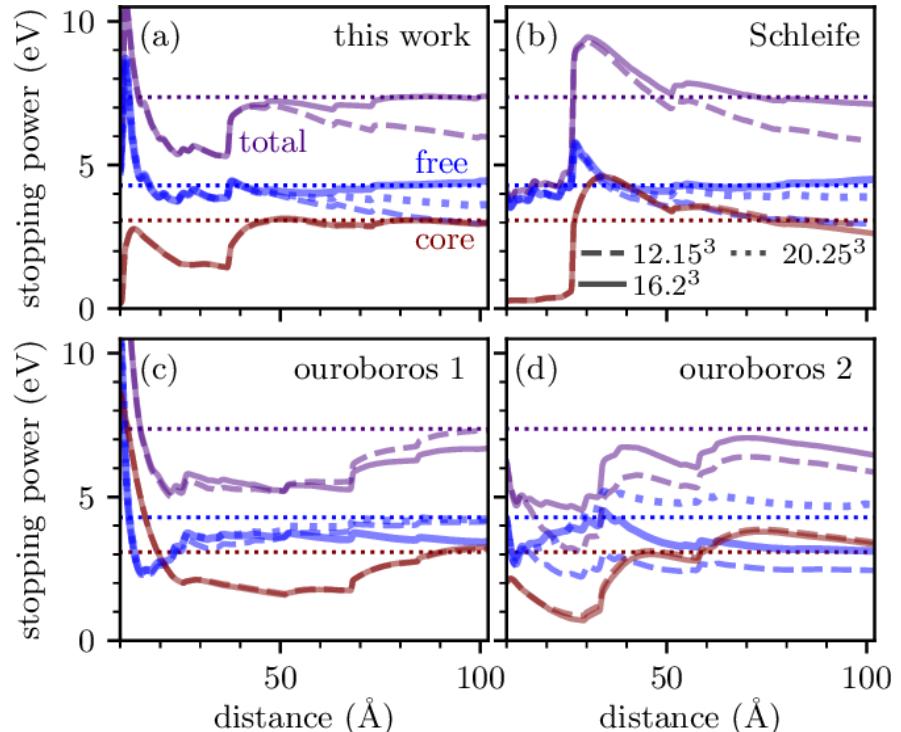
Partitioned finite-size effects allow cost reductions

- Finite-size effects depend on trajectory!
- Core contribution insensitive to supercell size
 - can use smaller supercell, reduce cost
- Behavior in free contribution inconsistent with plasmonic interpretation of finite-size effects

$$S(v) = \frac{2Z^2}{\pi v^2} \int_0^\infty \frac{dk}{k} \int_0^{kv} d\omega \omega \text{Im} \left[\frac{-1}{\epsilon(k, \omega)} \right]$$

$k_{\text{cut}} = 2\pi/L$

see Hentschel et al., arXiv:2301.09700



Kononov et al., in prep. (2023)
Schleife et al., PRB 91 (2015)