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Quantifying uncertainty for multi-megabar ramp compression of Pt

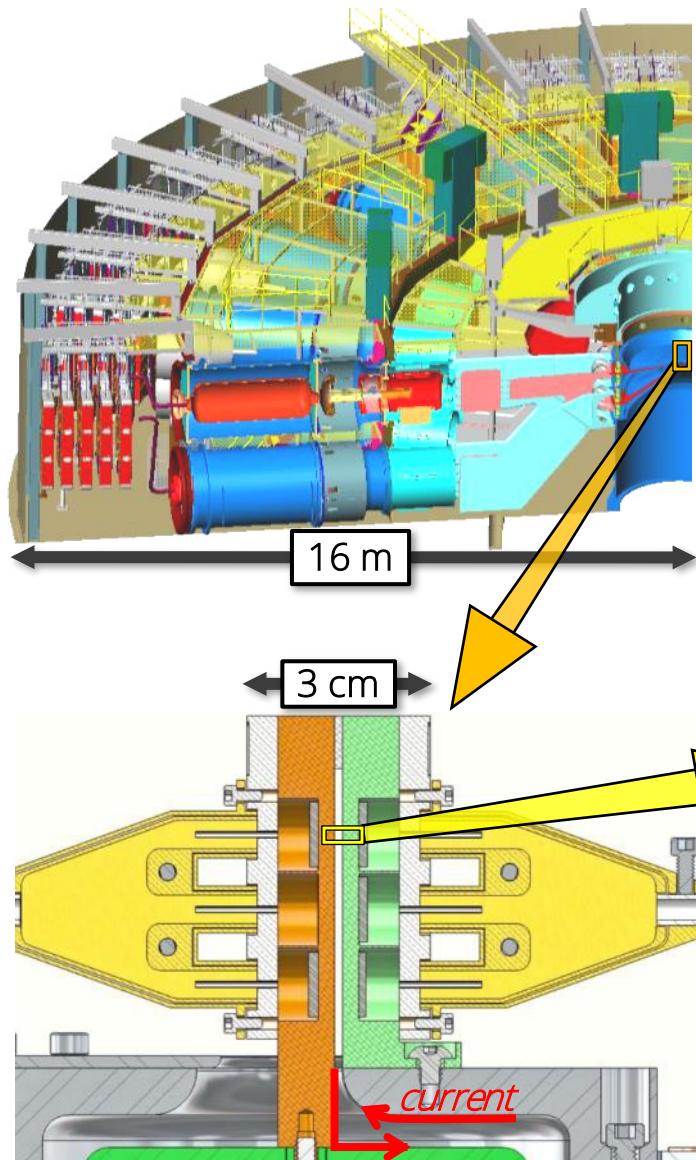
Jean-Paul Davis and Justin L. Brown

23rd Biennial Conference of the APS Topical Group on Shock
Compression of Condensed Matter

19-23 June 2023 in Chicago, IL, USA

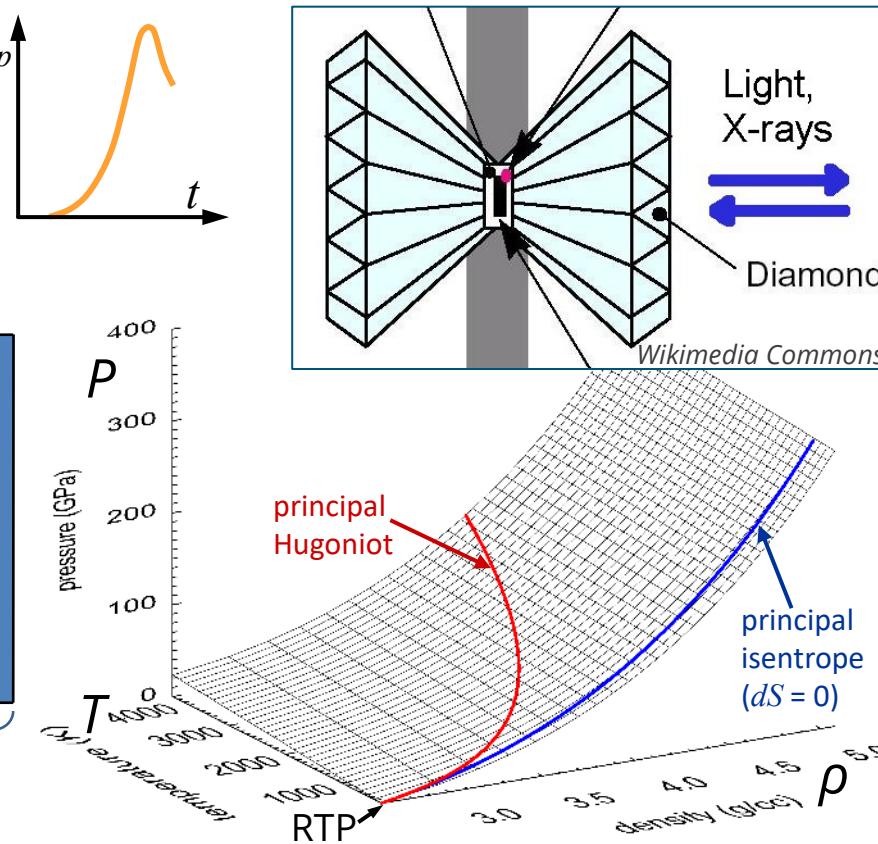
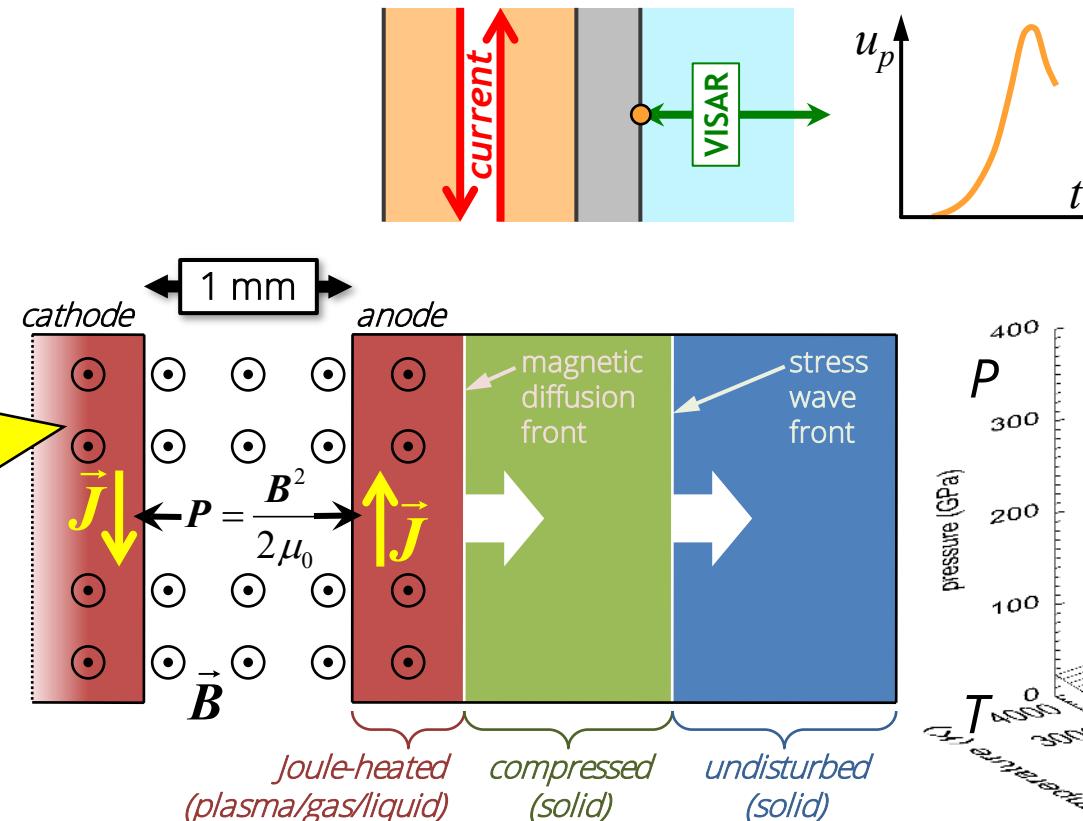
Session K01: Uncertainty Quantification and Propagation for EOS

Stripline short-circuit loads on the Z pulsed-power machine can produce planar shockless compression of solids to 400+ GPa

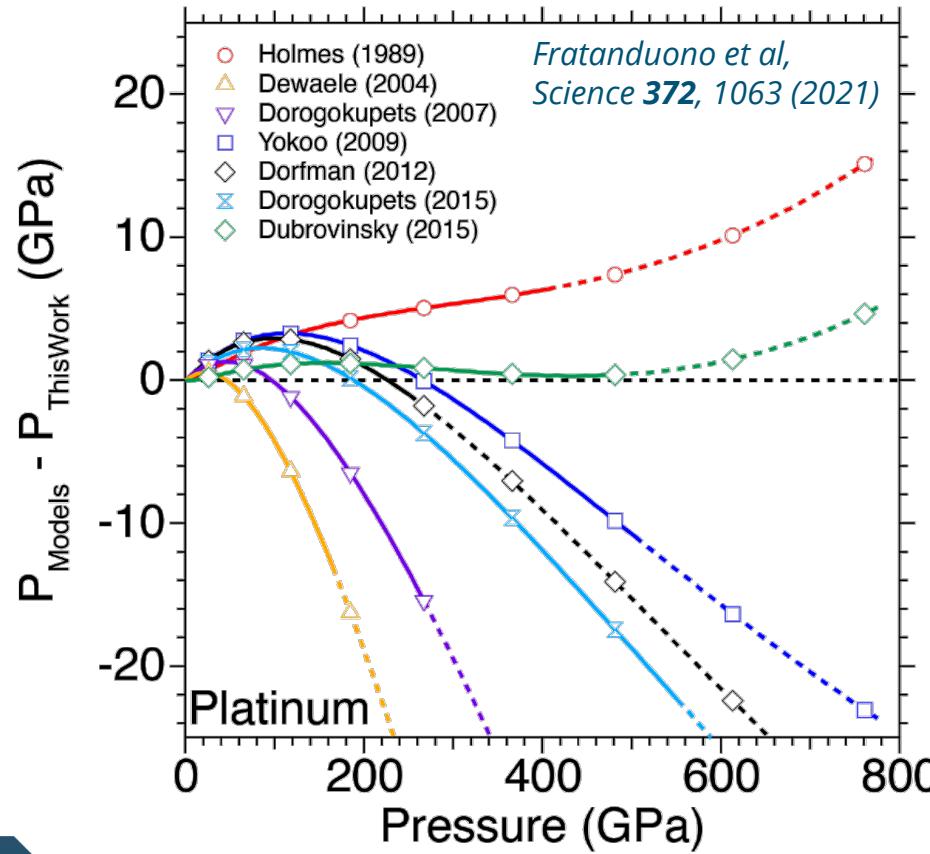
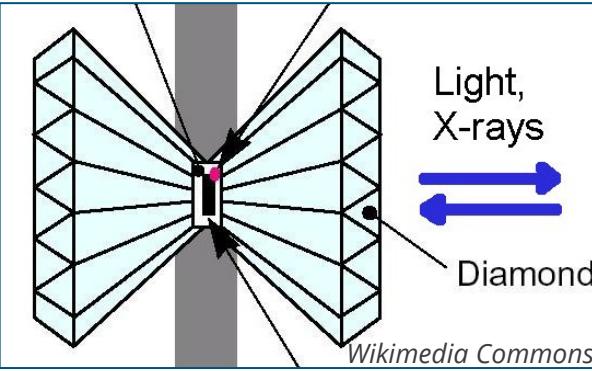


Magnetic drive propagates ramped stress wave into ambient material
Material's compressibility deduced from velocimetry measurements

- Shockless compression along quasi-isentrope close to isotherm
- Accurate absolute measurement of pressure standards (e.g., for DAC)

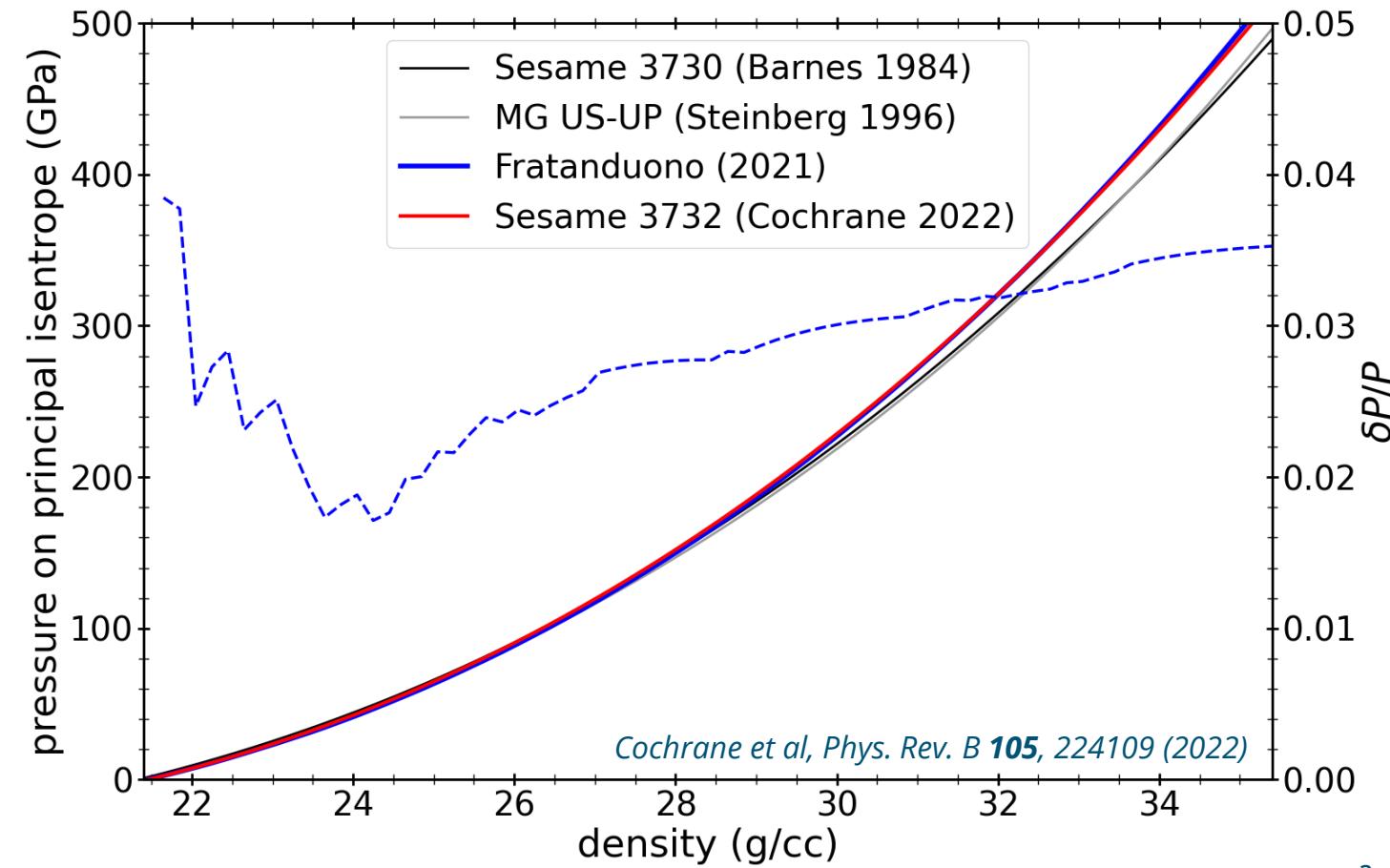


Platinum is a widely used pressure calibrant in DAC work

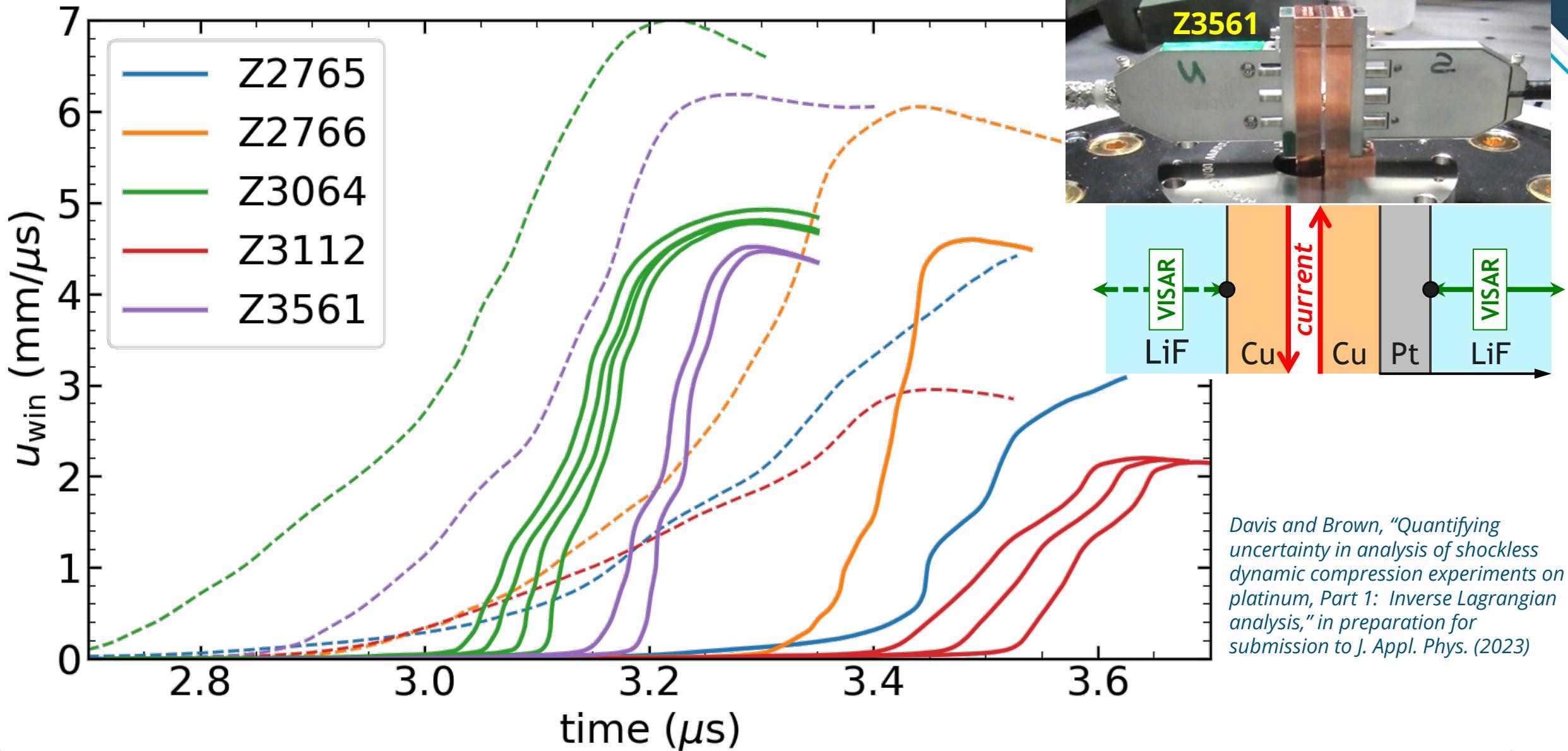


New DAC techniques to $P > 300$ GPa require precise calibration
Extrapolated models differ 5-10%

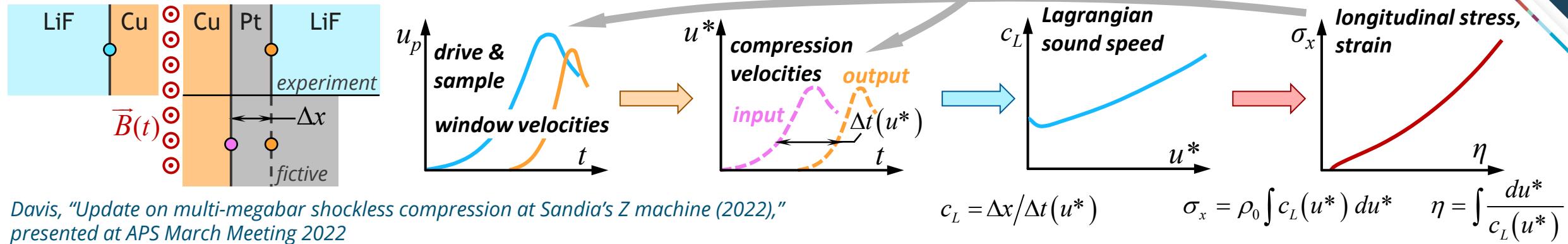
- recent NIF/Z ramp-compression result with uncertainty 3.5%
- included 3 Z measurements to ~ 400 GPa (NIF data to ~ 800 GPa)



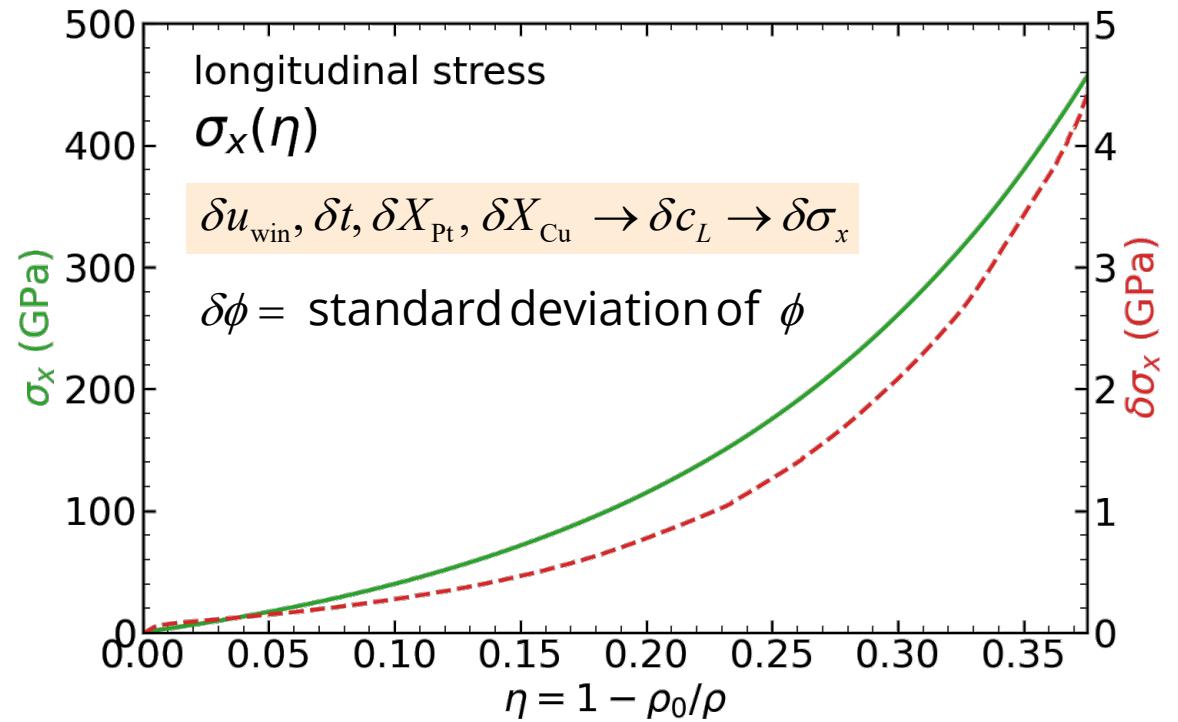
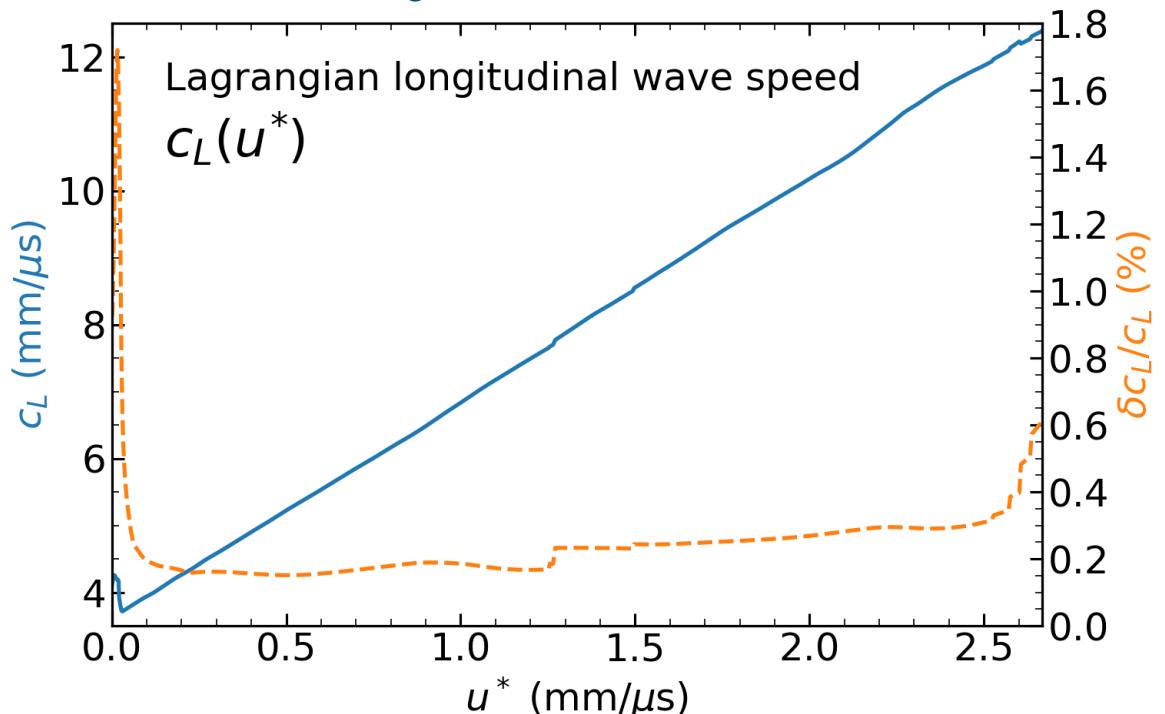
High-precision data are available from 11 single-sample measurements on Pt to peak stresses in the range 160-580 GPa



Inverse Lagrangian analysis (ILA) of these data gives weighted-mean longitudinal stress as function of strain with uncertainty



Davis, "Update on multi-megabar shockless compression at Sandia's Z machine (2022)," presented at APS March Meeting 2022



Davis and Brown, "Quantifying uncertainty in analysis of shockless dynamic compression experiments on platinum, Part 1: Inverse Lagrangian analysis," in preparation for submission to J. Appl. Phys. (2023)

Yield stress and thermal EOS models allow reduction of long. stress to pressures on quasi-isentrope, isentrope, and isotherm

Standard application of Von Mises yield criterion and Mie-Grüneisen EOS

Mie-Grüneisen EOS

$$P = P_{\text{ref}} + \frac{\rho_0 \Gamma (E - E_{\text{ref}})}{1 - \eta}$$

power-law Grüneisen parameter

$$\Gamma = \Gamma_0 (1 - \eta)^q$$

Assume constant c_V and $\beta = 0.9$

Need functions $Y(P)$, $G(P)$, and $\Gamma(\rho)$

Iterate to self-consistent $P_Q \leftrightarrow Y$

Integrate ordinary diff. eqn. to get P_T

quasi-isentrope $P_Q = \sigma_x - \sigma' = \sigma_x - \frac{2}{3}Y$

isentrope $P_S = P_Q - \frac{\rho_0 \Gamma}{1 - \eta} \beta \int dW_p$

isotherm $P_T = P_S - \frac{\Gamma}{1 - \eta} \int \left(P_S - P_T - \frac{\rho_0^2 \Gamma c_V T}{1 - \eta} \right) d\eta$

σ' deviatoric stress

Y yield stress

Γ Grüneisen parameter

β integral Taylor-Quinney coefficient

$$dW_p = \frac{2}{3} \frac{Y}{\rho_0} \left(\frac{1}{1 - \eta} - \frac{1}{2G} \frac{dY}{d\eta} \right) d\eta \quad \text{incremental plastic work}$$

G shear modulus

c_V isochoric specific heat

Simply subtracting deviatoric stress to obtain quasi-isentrope pressure can introduce elastic-region errors, limiting accuracy



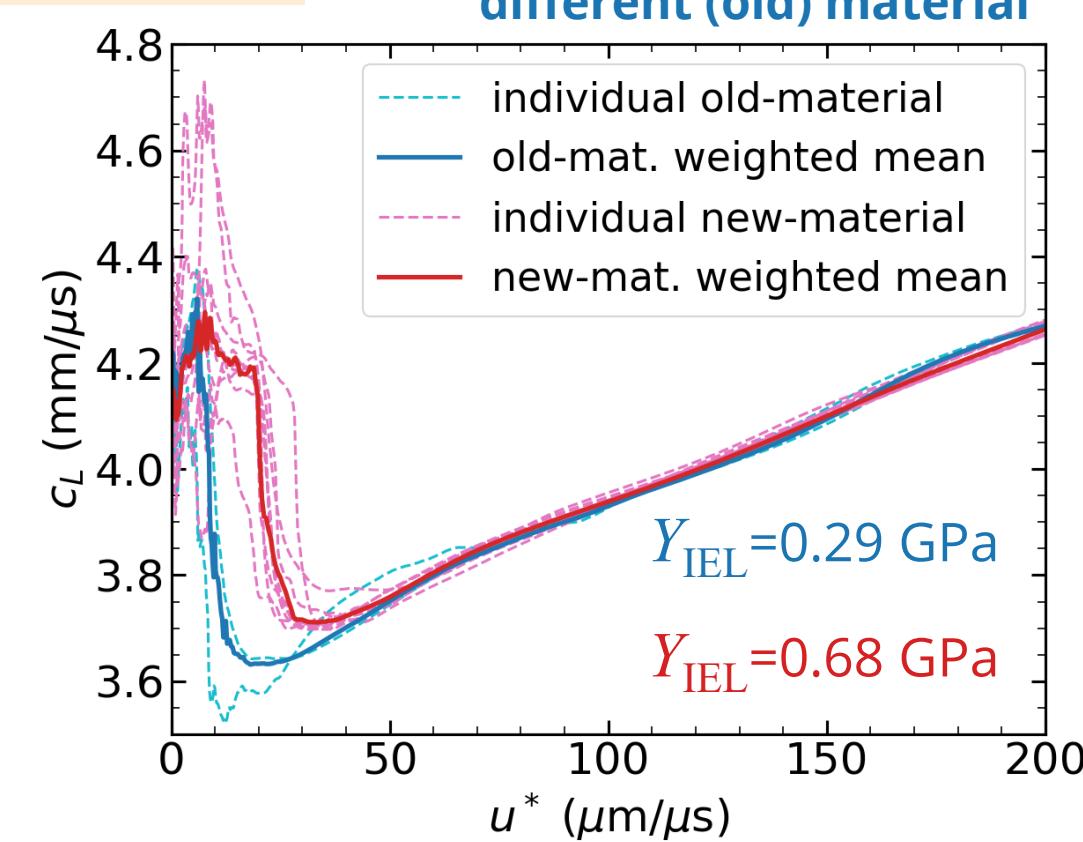
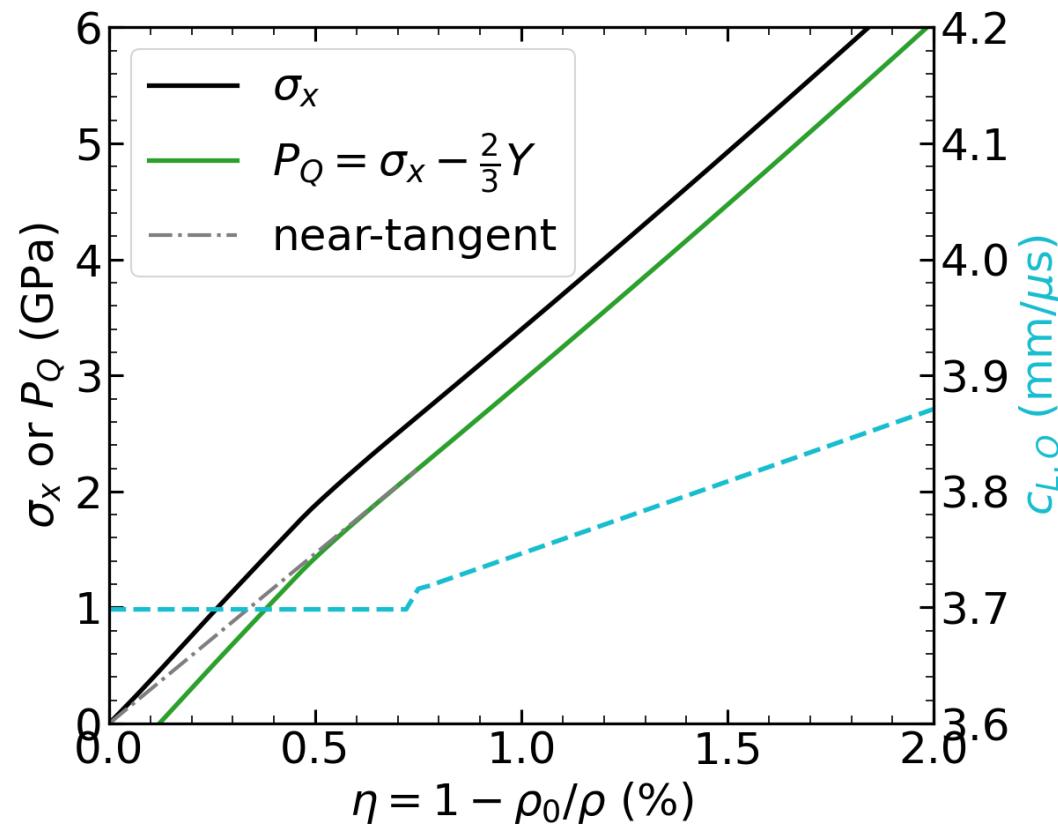
A straight line in u , ρ , or η from $P_Q=0$...

- Cannot avoid discontinuity in wave speed above elastic limit
- Does not match known elastic response at ambient

isentropic elastic limit Y_{IEL}

Initial yield stress depends strongly on microstructure

- Isentrope pressure P_S should **not** depend on yield stress



Z2765 & Z2766 used different (old) material

To avoid these issues, compute quasi-isentrope & isentrope reductions in wave speed instead of stress/pressure

Lagrangian wave speeds $c_L = \sqrt{\frac{1}{\rho_0} \frac{d\sigma_x}{d\eta}}$, $c_{L,Q} = \sqrt{\frac{1}{\rho_0} \frac{dP_Q}{d\eta}}$, $c_{L,S} = \sqrt{\frac{1}{\rho_0} \frac{dP_S}{d\eta}} = \frac{\rho}{\rho_0} \sqrt{\left(\frac{\partial P}{\partial \rho}\right)_S}$ (sound speed)

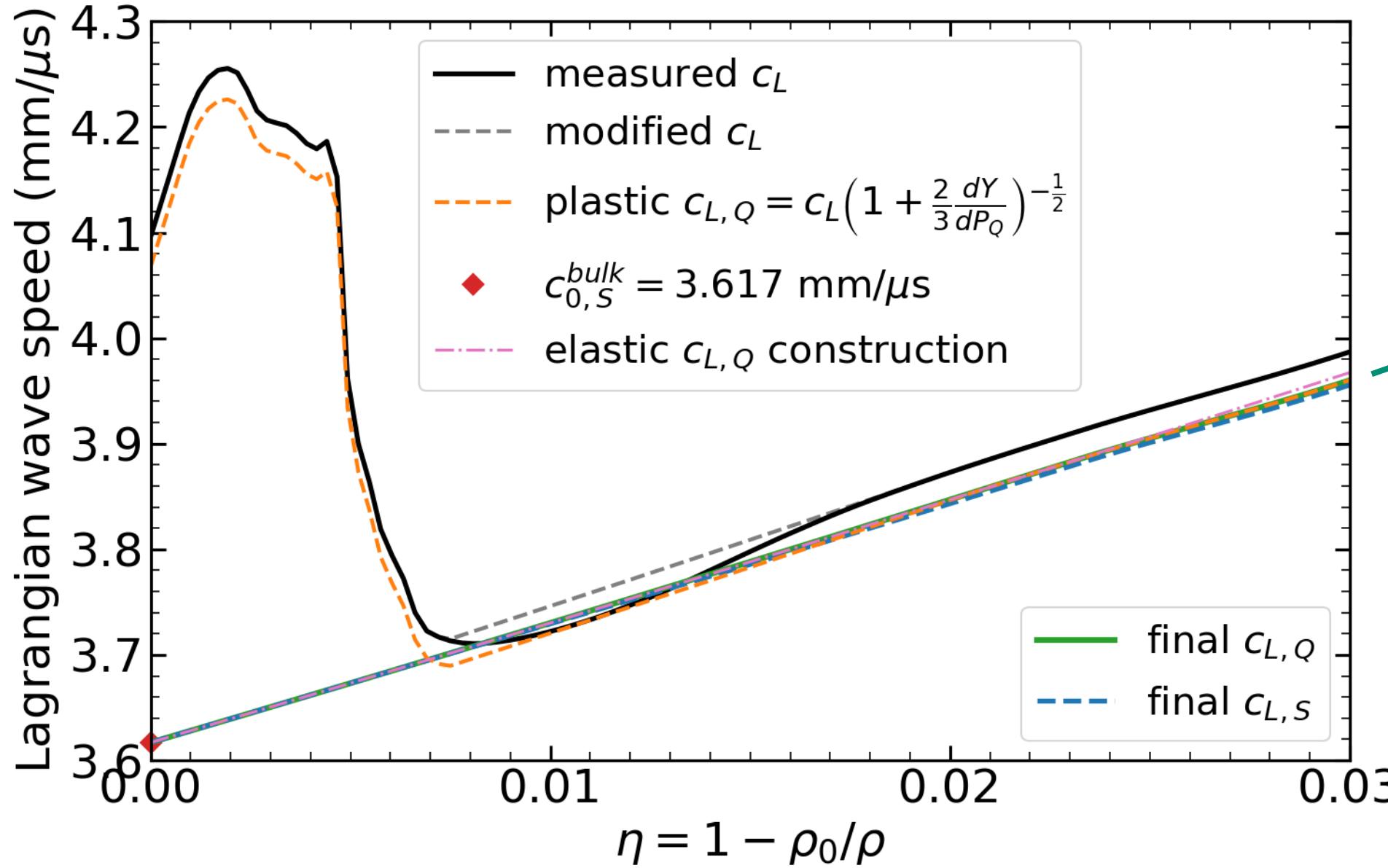
$$\sigma_x = P_Q + \frac{2}{3}Y \rightarrow \frac{d\sigma_x}{du^*} = \frac{dP_Q}{d\eta} \frac{d\eta}{du^*} \left(1 + \frac{2}{3} \frac{dY}{dP_Q} \right) \quad u^* = \int c_L d\eta \neq u_Q \neq u_S$$

$$c_{L,Q}^2 = c_L^2 \left(1 + \frac{2}{3} \frac{dY}{dP_Q} \right)^{-1} \quad u_Q = \int c_{L,Q} d\eta \quad P_Q = \rho_0 \int c_{L,Q} du_Q = \rho_0 \int c_{L,Q}^2 d\eta$$

$$c_{L,S}^2 = c_{L,Q}^2 - \beta \Gamma_0 (1-\eta)^{q-1} \left(\frac{dW_p}{d\eta} - \frac{q-1}{1-\eta} W_p \right) \quad u_S = \int c_{L,S} d\eta \quad P_S = \rho_0 \int c_{L,S} du_S = \rho_0 \int c_{L,S}^2 d\eta$$

$$dW_p = \frac{2}{3} \frac{Y}{\rho_0} \left(\frac{1}{1-\eta} - \frac{1}{2G} \frac{dY}{d\eta} \right) d\eta$$

Interpolating bulk wave speed between known ambient value and plastic region gives smooth pressure curve

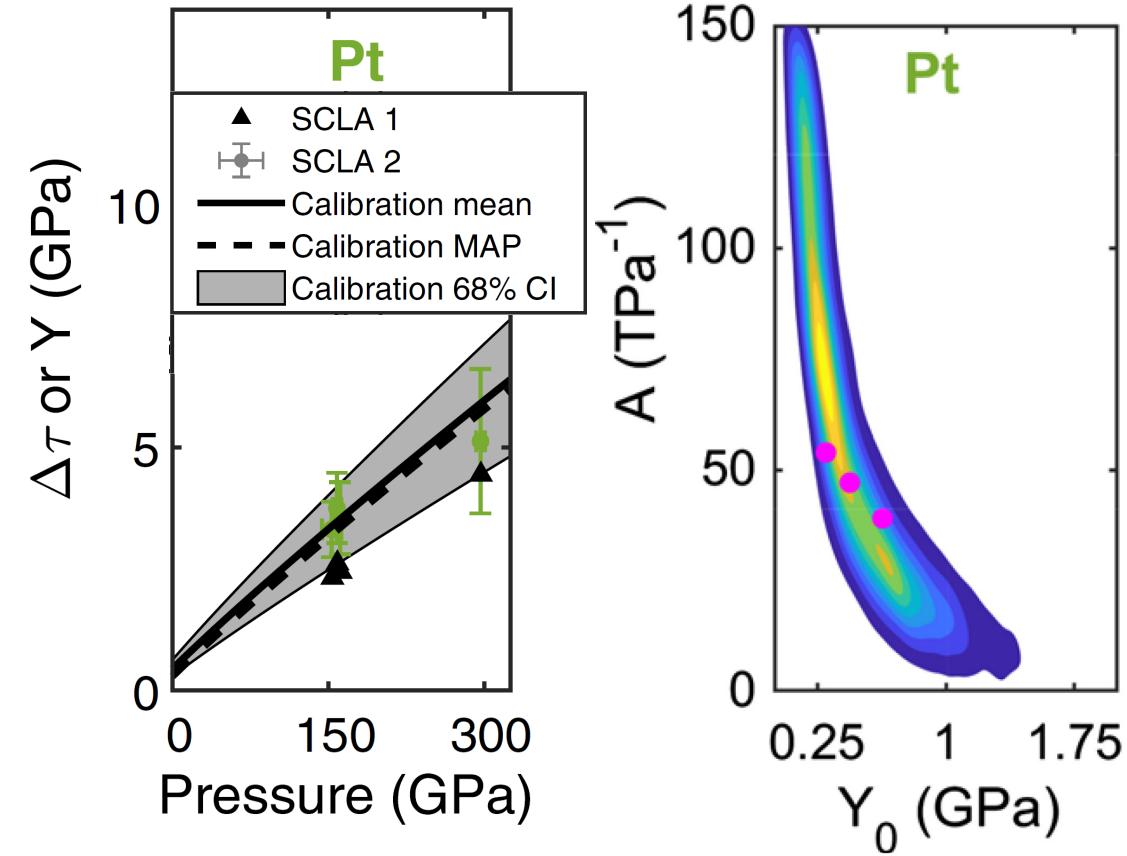


Monte-Carlo UQ analysis of the reduction procedure requires statistical models for 10 input parameters

	distribution	$\mu, s^{1/2}$	units
Y_0	$p([Y_0, A] D)$	0.29, N/A	GPa
A	$p([Y_0, A] D)$	72, N/A	TPa ⁻¹
Γ_0	$N([\mu_{\ln \Gamma_0}, \mu_q], \Sigma)$	0.971, 0.118	
q	$N([\mu_{\ln \Gamma_0}, \mu_q], \Sigma)$	0.94, 0.411	
β	$B(6.4, 1.6)$	0.90; [0.66, 0.93]	
ρ_0	$N(\mu_{\rho_0}, s_{\rho_0})$	21.421, 0.043	g/cm ³
$c_{0,S}$	$N(\mu_{c_{0,S}}, s_{c_{0,S}})$	3.617, 0.055	mm/μs
G_0	$N(\mu_{G_0}, s_{G_0})$	63.7, 3.6	GPa
c_V	$N(\mu_{c_V}, s_{c_V})$	130.2, 0.4	J/(kg·K)
T_0	$N(\mu_{T_0}, s_{T_0})$	300, 3	K

simplified Steinberg-Guinan yield

$$Y = Y_0 \left[1 + AP_Q (1 - \eta)^{\frac{1}{3}} \right], \quad G = G_0 \frac{Y}{Y_0}$$



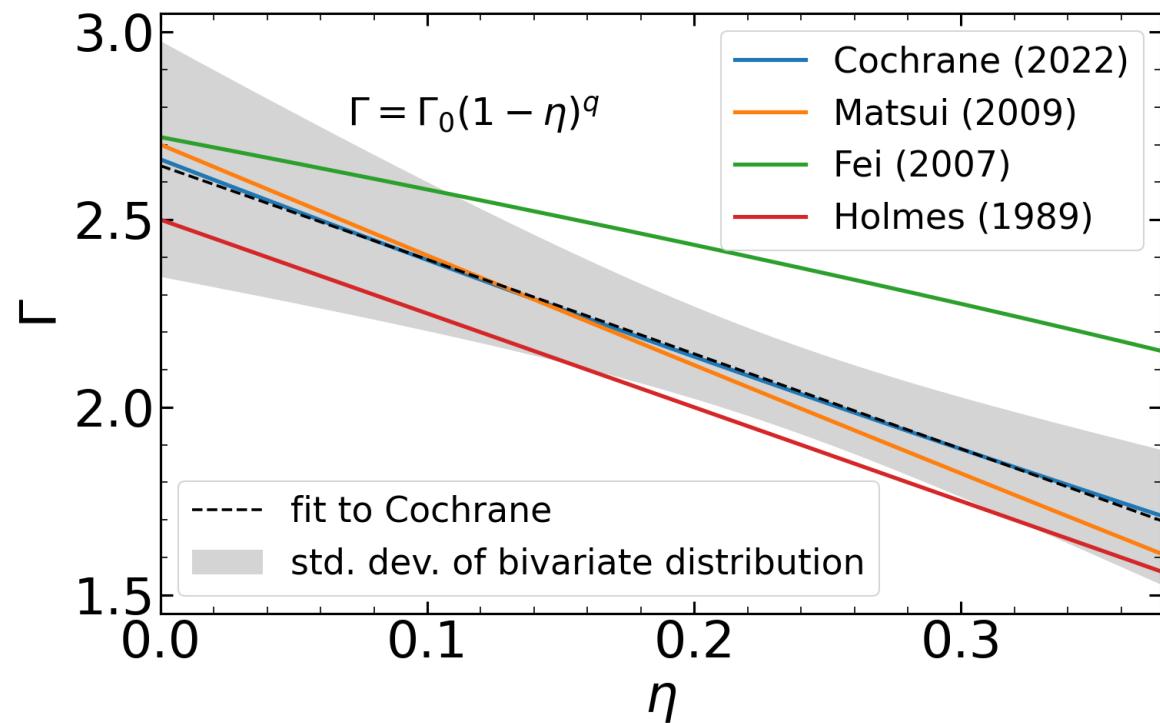
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mean: fit to Cochrane et al (2022)

std. dev: fit $\Delta \ln \Gamma = \Delta \ln \Gamma_0 + \Delta q \ln(1-\eta)$

where $\Delta \Gamma = \Gamma_{\text{Matsui, Fei, or Holmes}} - \Gamma_{\text{Cochrane}}$



Cochrane et al, *Phys. Rev. B* **105**, 224109 (2022)

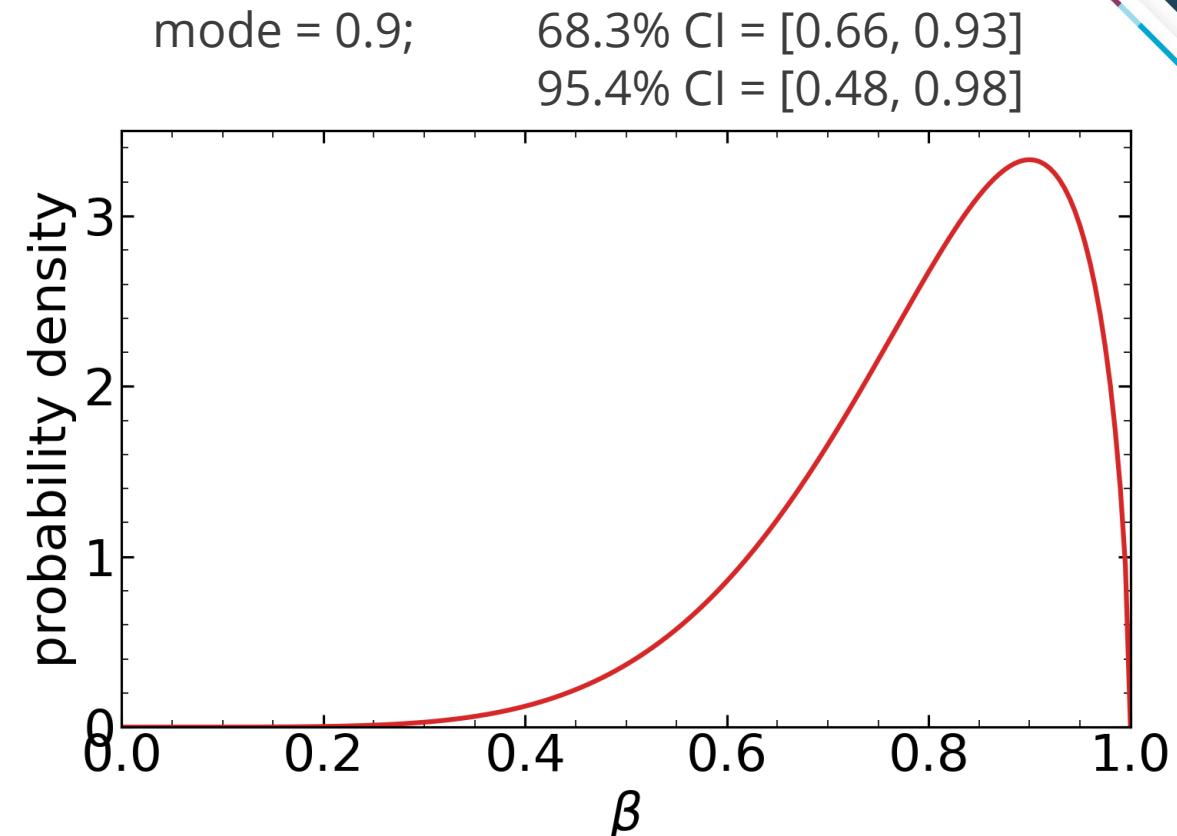
Matsui et al, *J. Appl. Phys.* **105**, 013505 (2009)

Fei et al, *Proc. Nation. Acad. Sci.* **104**, 9182-9186 (2007)

Holmes et al, *J. Appl. Phys.* **66**, 2962-2967 (1989)

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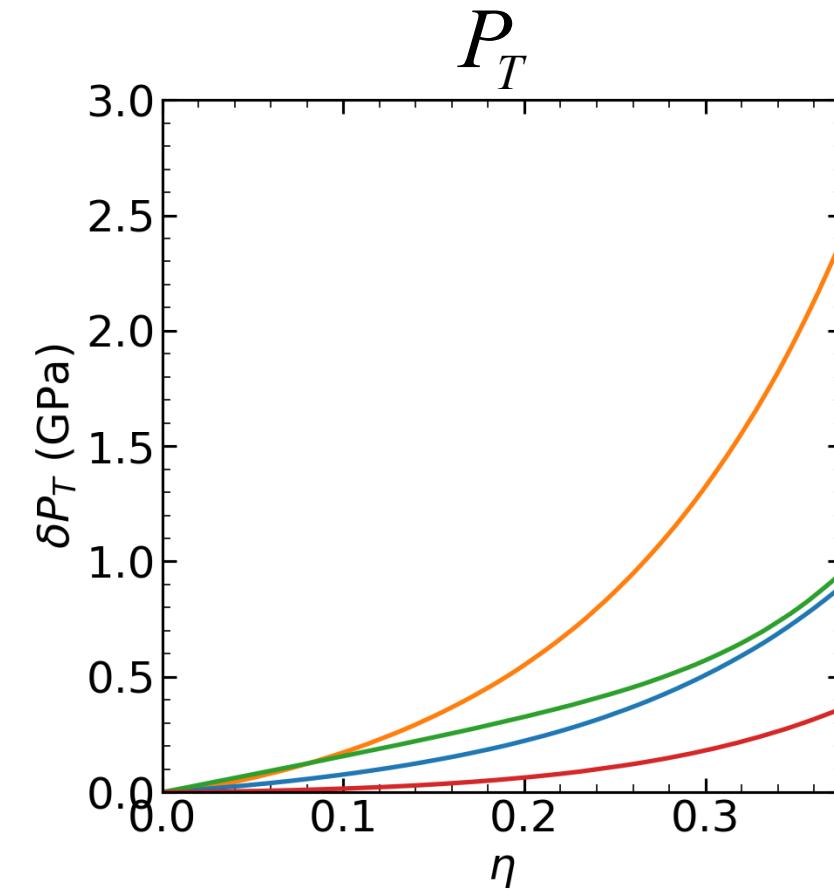
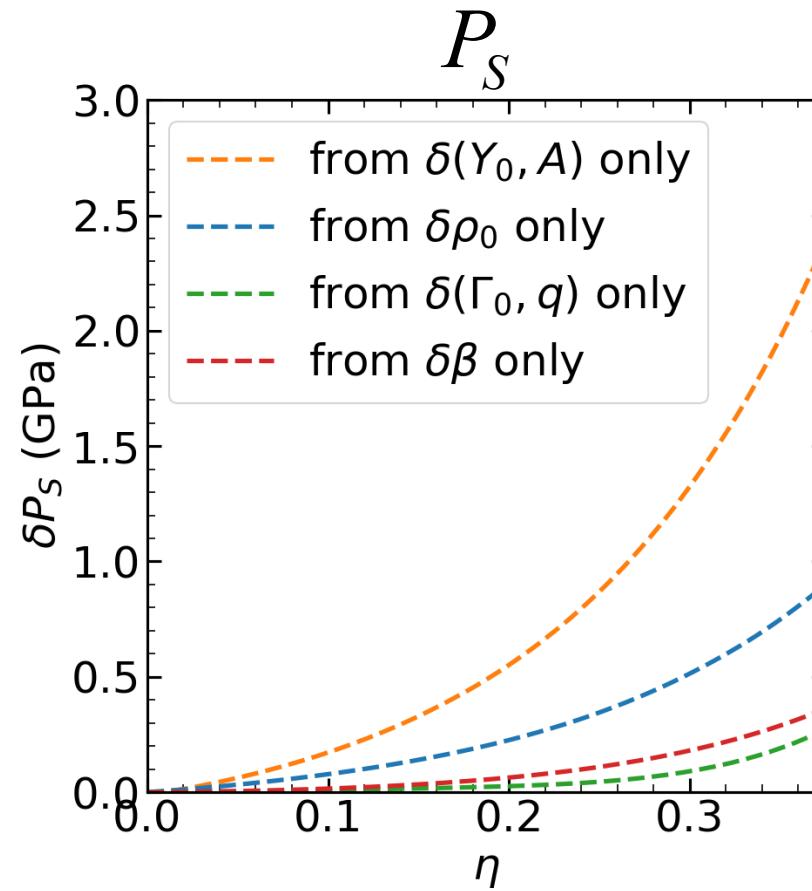
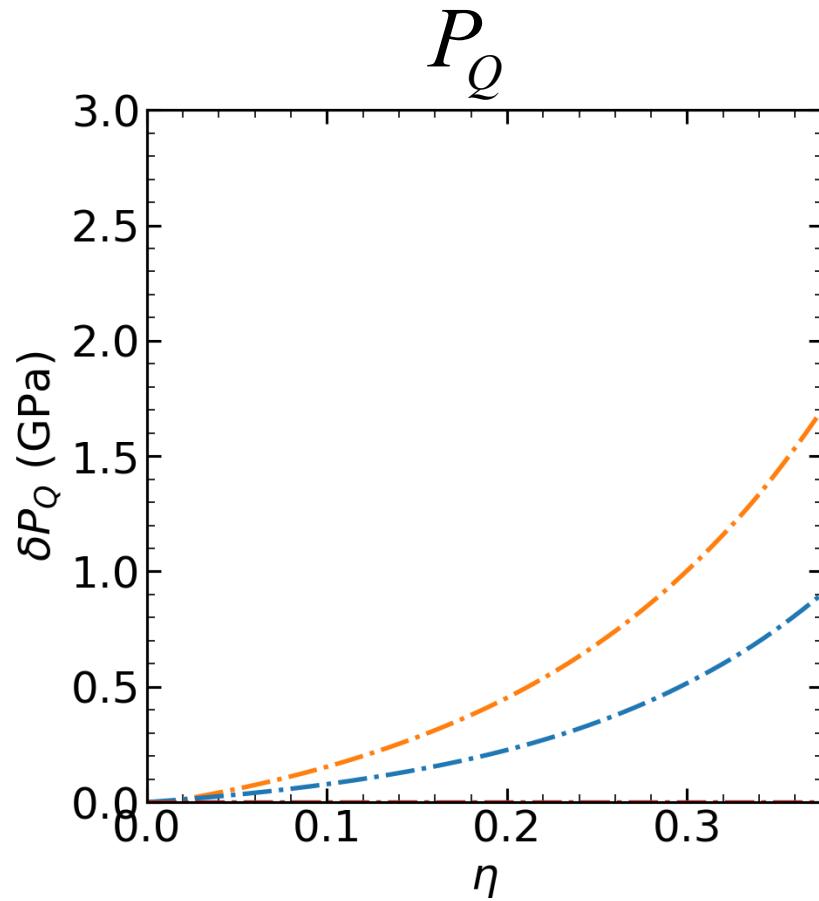
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Remaining 5 parameters = independent univariate normal distributions

Obtain rough estimate of sensitivities by “screening” for effect of one univariate or bivariate parameter distribution at a time

- Screening Monte Carlo (MC) converged by 100,000 samples
- Non-normal distributions described by average median-centered 68.3% confidence interval (CI)

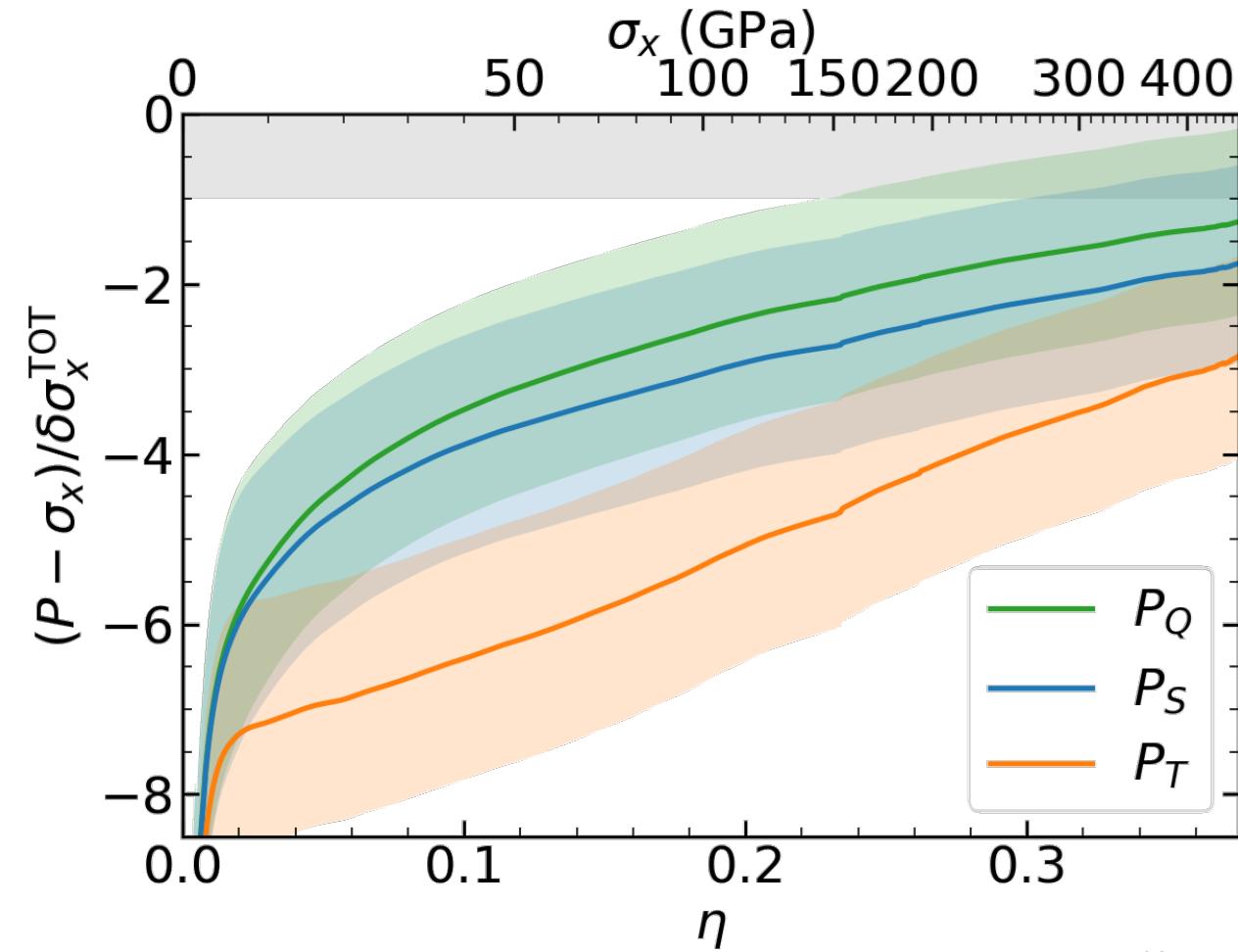
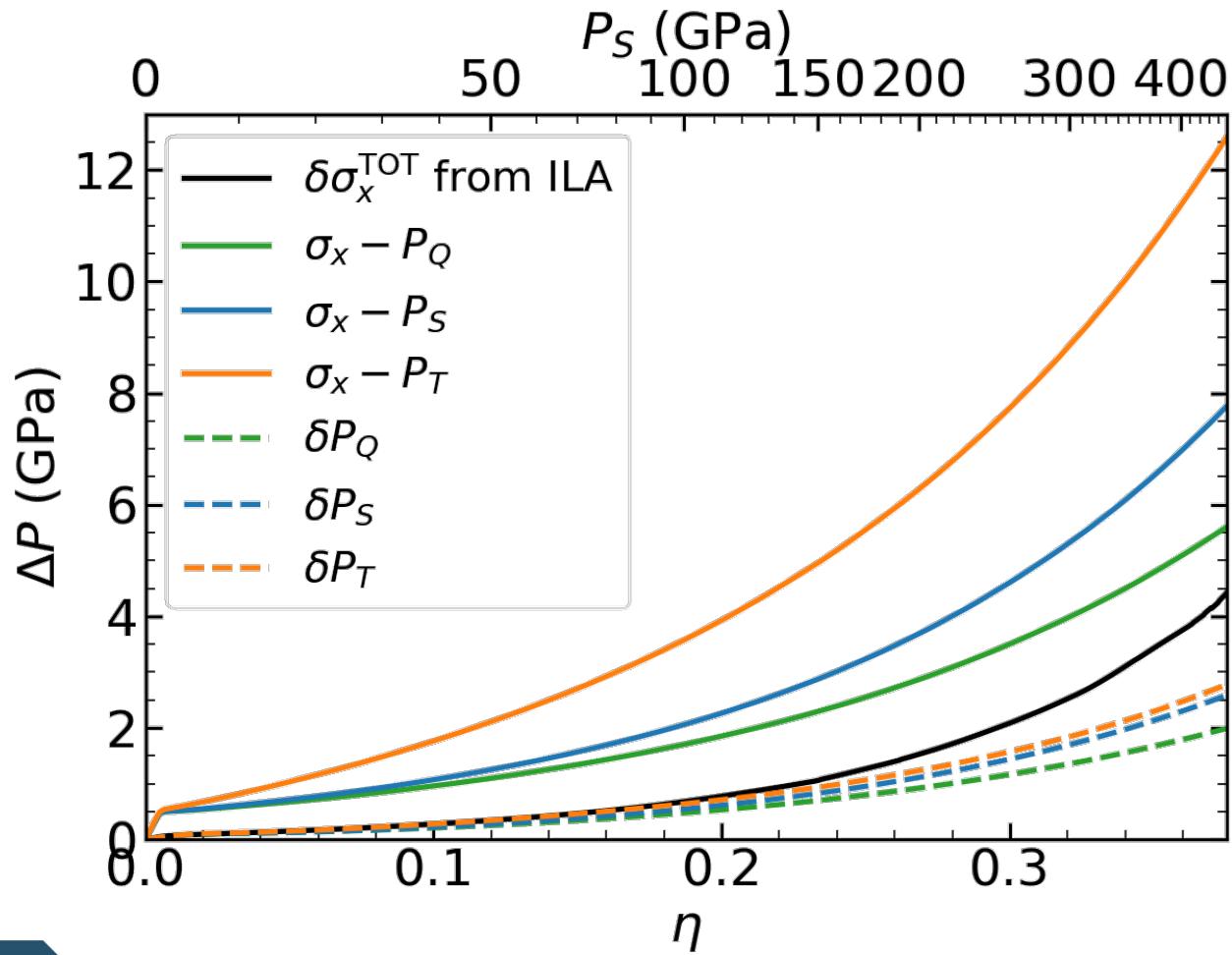


None of the remaining 4 parameters surpassed 0.1-GPa effect

Uncertainty contributed by reduction procedure quantified by 500,000-sample Monte Carlo added in quadrature to total $\delta\sigma_x$



- Non-normal distributions again described by average median-centered 68.3% confidence interval (CI)
- Two views of reduction corrections and CI: (left) absolute, before adding to total $\delta\sigma_x$ std. dev., and (right) in units of $\delta\sigma_x$, after adding to total $\delta\sigma_x$ std. dev.



Conclusion

Rigorous propagation of experimental uncertainties to total $\delta\sigma_x(\eta)$

- about 1% at ~ 450 GPa from weighted averaging of 11 high-precision Z measurements on Pt

Improved procedure for reduction to pressures on quasi-isentrope, isentrope, isotherm

- work in wave speeds instead of stress and pressure
- interpolate across elastic-plastic transition from known ambient bulk $c_{0,S}$ to plastic-flow $c_{L,Q}$
- P_Q depends only on yield-stress derivative dY/dP , not initial yield stress Y_{IEL}
- P_S depends also on absolute $Y(P)$ but with nearly inconsequential sensitivity

Monte Carlo analysis quantifies reduction procedure contribution to uncertainty

- in descending order of importance, parameters of measurable effect are those for strength, initial density, Grüneisen parameter, and Taylor-Quinney coefficient
- ambient bulk sound speed, shear modulus, heat capacity, and temperature are inconsequential
- non-normal distribution described well by averaged median-centered confidence intervals
- Reduction to isotherm P_T is correction of $\sim 3 \delta\sigma_x$ at ~ 450 GPa, resulting in δP_T around 1.2%

UQ does not yet include uncertainty in Cu and LiF material models used in ILA

Bayesian model calibration (BMC) suggests decreased sensitivity at highest pressures should inflate uncertainty there (1.2% from ILA might really be 1.5%)