

Nanoconfinement of carbon dioxide within interfacial aqueous/ionic liquid systems

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Abstract

Nanoporous, gas-selective membranes have shown encouraging results for the removal of CO₂ from flue gas, yet the optimal design for such membranes is often unknown. Therefore, we used molecular dynamics simulations to elucidate the behavior of CO₂ within aqueous and ionic liquid (IL) systems ([EMIM][TFSI] and [OMIM][TFSI]), both confined individually and as an interfacial aqueous/IL system. We found that, within aqueous systems, the mobility of CO₂ is reduced due to interactions between

the CO₂ oxygens and hydroxyl groups on the pore surface. Within the IL systems, we found that confinement has a greater effect on the [EMIM][TFSI] system as opposed to the [OMIM][TFSI] system. Paradoxically, the larger and more asymmetrical [OMIM]⁺ molecule undergoes less efficient packing, resulting in fewer confinement effects. Free energy surfaces of the nanoconfined aqueous/IL interface demonstrate that CO₂ will transfer spontaneously from the aqueous to IL phase.

Introduction

Ionic liquids (ILs) have garnered much attention for their unique and advantageous properties.^{1,2} Notably, ILs have negligible volatility, high chemical stability, and a large range of useful operating temperatures.^{3–5} These properties, as well as the fact that many ILs possess high carbon dioxide (CO₂) solubility, make ILs promising candidates to replace aqueous amines in post combustion capture of CO₂.^{6,7} For this reason, extensive research has been dedicated to understanding and enhancing CO₂ solubility and transport in ILs.^{8–13} One drawback of ILs is that their high cost relative to traditional solvents has led some technoeconomic analyses to suggest that it might not be viable to use large quantities of physically absorbing IL for CO₂ capture.^{14,15} However, this drawback can be remedied by the development of novel systems that use considerably less of the expensive IL component, such as thin film¹⁶ and membrane systems.¹⁷

Initial work integrating ILs into gas-selective polymer membranes by Noble^{18,19} and others^{20,21} achieved excellent performance, with CO₂/N₂ selectivity as high as 61.²² Maginn’s group performed molecular simulations of graphite slit pore IL membranes, finding that nanoconfinement may increase the permselectivity for CO₂ over CH₄.²³ Yet, further optimization is difficult as a trade-off between selectivity and permeability typically exists.²⁴ Fu et al. achieved unprecedentedly high performance, including a CO₂/N₂ selectivity of 788 and a CO₂ flux of 2600 GPU, using an ultrathin water (H₂O) membrane to stabilize the enzyme carbonic anhydrase (CA) at concentrations greater than those attainable in bulk water.²⁵

The gas separation membrane proposed by Fu et al. could be redesigned to include an IL layer directly following the ultra-thin CA-catalyzed aqueous layer. This new system would enable the CO₂ transporting through the membrane to be concentrated into a condensed phase for easier conversion into value-added products.

Previously, we used both classical molecular dynamics (MD) and laboratory experiments to investigate the behavior of CO₂ in bulk, interfacial aqueous-[EMIM][TFSI] and aqueous-[OMIM][TFSI] systems,²⁶ where [EMIM]⁺ = 1-ethyl-3-methylimidazolium, [OMIM]⁺ = 1-octyl-3-methylimidazolium, and [TFSI]⁻ = bis(trifluoromethanesulfonyl)imide (**Figure 1a**). We observed that CO₂ will transport spontaneously from the aqueous phase to the IL phase and that the diffusion of CO₂ in the ILs does not follow the conventional Stokes-Einstein relation.²⁶ To expand on this work, here we used MD simulations to investigate CO₂ dissolved in aqueous and IL systems when confined in cylindrical, silica nanopores (3, 5, and 8 nm diameters) with hydrophilic and hydrophobic pore surface termination (**Figure 1b-c**). Moreover, we designed an interfacial system that transitions between hydrophobic to hydrophilic surface chemistry, allowing us to examine the IL/aqueous interface while under the effects of nanoconfinement. The more complex systems modeled here are designed to match the experimental conditions of the ultra-thin enzymatic membrane better.

Methods

Molecular dynamics parameters

Classical molecular dynamics (MD) simulations were performed using GROMACS (2021).^{27,28} The leap-frog algorithm was used with a 1 fs timestep. The velocity-rescaling thermostat with a coupling constant of 1 ps and a reference temperature of 298 K was used.²⁹ During the pore loading procedure (*vide infra*), pressure was maintained using the Berendsen barostat and a 1 ps coupling constant.³⁰ The LINCS algorithm was used to constrain bond lengths for atoms bonded with hydrogen.³¹ Short-range, non-bonded interactions were cutoff at 13

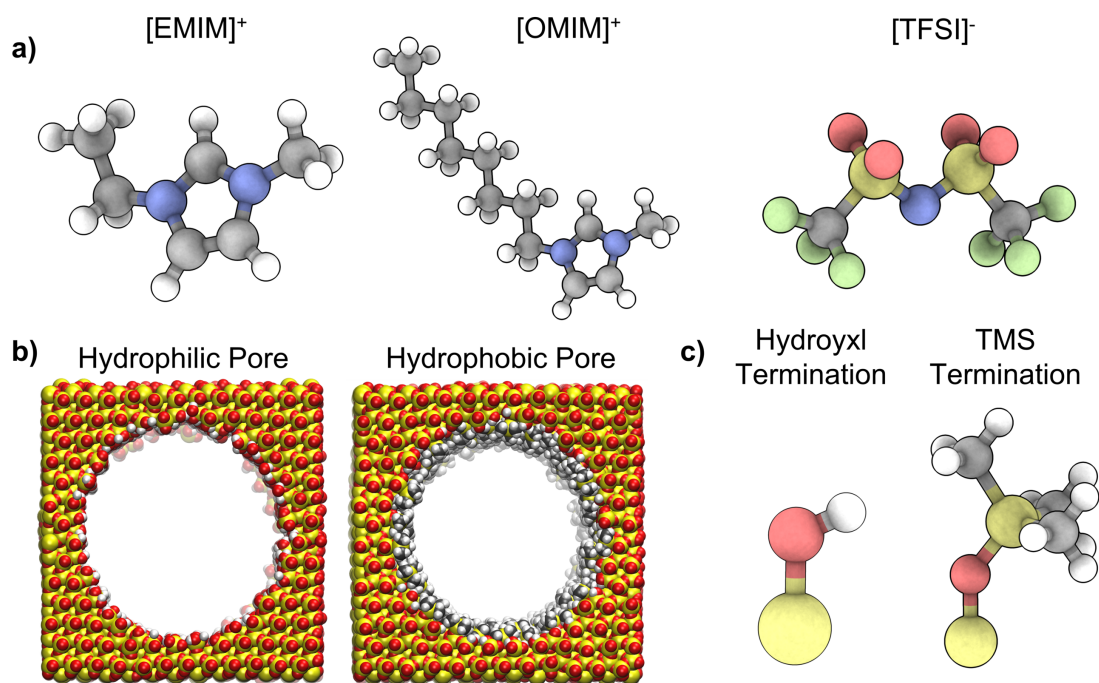


Figure 1: a) Molecular structures of ILs $[EMIM]^+$, $[OMIM]^+$, and $[TFSI]^-$. b) Molecular structure of the 5 nm diameter hydrophilic and hydrophobic pores. c) Hydrogen (-H) that makes surface hydroxyls and trimethylsilyl (TMS, $-Si(CH_3)_3$) groups terminate the hydrophilic and hydrophobic silica (SiO_2) pores, respectively. Grey = Carbon, Blue = Nitrogen, Red = Oxygen, Yellow = Sulfur in $[TFSI]^-$, else Silicon, Green = Fluorine, White = Hydrogen.

Angstroms. The particle mesh Ewald (PME) method was used for long-range, non-bonded interactions.³² Periodic boundary conditions were applied in all directions.

The ± 0.8 charge ionic liquid (IL) force field parameters reported by Doherty et al. were used for [EMIM]⁺, [OMIM]⁺ and [TFSI]⁻.³³ SPC/E parameters were used for water.³⁴ The parameters for CO₂ were taken from Cygan et al.³⁵ The Emami et al. silica surface parameters were used for the silica pore³⁶ and the LigParGen web server was used to generate OPLS-AA compatible parameters for the trimethylsilyl (TMS) groups terminating the hydrophobic silica pore.³⁷

Structure generation

The PYTHON module PoreMS (0.2.0) was used to generate hydrophilic and hydrophobic silica pore initial structures of approximately 3, 5, and 8 nm pore diameter and a length of approximately 7 nm.³⁸ Hydrophilic pores were generated with 100% hydroxyl group coverage within the pore. Hydrophobic pores were generated by replacing 40% of the H in the hydroxyl groups with TMS groups. Due to the larger size of the TMS groups compared with the hydroxyl groups that were replaced, 40% coverage is near the maximum amount of TMS groups that can fit without collisions. While real silica surfaces typically have ~ 9 -18% hydroxyl group ionization at neutral pH, we did not include surface group ionization in our models. Previously, no major effects were observed by this degree of ionization (e.g., identical water density profiles in Emami et al. Figure 8).³⁶ The empty pore structures underwent an energy minimization with the steepest descent algorithm for 5000 steps, followed by a 500 ps NPT simulation at 1 atm and 298 K. Afterwards, the silica molecules were frozen in place for the remainder of the study; however, the hydroxyl and TMS functional groups attached to the silica pore interior remained mobile.

Bulk liquid IL/CO₂ simulation boxes were taken from our previous work²⁶ and are available on GitHub.³⁹ These simulation boxes contain 512 IL ion pairs along with 100 CO₂ molecules. This mole fraction of CO₂ corresponds to that expected for [EMIM][TFSI] in

equilibrium with a CO_2 partial pressure of ~ 0.35 MPa.⁴⁰ This value is within the CO_2 partial pressures for either integrated gasification combined cycle (IGCC) synthesis gas after gasification or integrated reforming combined cycles (IRCC) synthesis gas after reforming.^{41,42} The bulk $\text{H}_2\text{O}/\text{CO}_2$ simulation box was initially created in Packmol⁴³ with 149 CO_2 molecules and $\sim 10,000$ H_2O molecules, chosen to match the CO_2 per nm^3 number density found in the IL systems. No N_2 was included in these systems as Fu et al. showed that an ultra-thin enzymatic layer can serve as an N_2 barrier.²⁵

Simulation procedure

As the density of a liquid within nanoconfinement differs from that of the bulk liquid,^{44,45} a pore loading procedure was developed instead of using the *gmx solvate* command. A simulation box was generated by placing an equilibrated bulk liquid box (i.e., IL/ CO_2 or water/ CO_2) next to the empty pore (**Figure S1a**). The system underwent an energy minimization followed by an NPT simulation at elevated pressure to ensure rapid and complete loading of the liquid into the pore. The system then underwent an NPT simulation at 1 atm and 298 K (**Figure S1b**). This process of allowing the pore liquid to equilibrate with a liquid reservoir allows the liquid density within the pore to be dictated by the liquid-liquid and liquid-pore intermolecular interactions, rather than simply constraining the pore liquid density to be the same as bulk liquid density. Finally, the reservoir liquid not located within the pore was removed, leaving only the pore as a periodic system (**Figure S1c**). The production runs were conducted in the NVT ensemble at 298 K for a minimum of 80 ns and 1000 ns for the water/ CO_2 systems and IL/ CO_2 systems, respectively. A similar procedure was used to generate interfacial pores with IL and aqueous layers, which is described in detail in the Supporting Information.

Simulation analysis

Radial distribution functions (RDFs), density profiles, spatial distribution functions (SDFs), and continuous dimer existence autocorrelation functions (DACFs) were computed

with TRAVIS.^{46,47} Free energy surfaces (FES) were computed by the histogram re-weighting methods using PLUMED.^{48,49} The built-in GROMACS command *gmx msd* was used to calculate diffusion coefficients.⁵⁰ In the pore systems, the diffusion coefficient was only computed in the direction parallel with the pore walls. Molecular renderings in **Figures** 1a and c, 4, 6 and S6 were created using Speck⁵¹ while the renderings in **Figures** 1b, 3, 7a, S1, S3, S4, and S5 were created using Visual Molecular Dynamics (VMD).⁵²

Results and Discussions

Aqueous pore systems

We began by observing the hydrophilic pore systems, containing only water and CO₂ molecules. In **Figure 2**, we show the diffusion coefficients and the diffusion ratios, defined as the pore diffusion coefficient divided by the bulk diffusion coefficient, for water and CO₂ as a function of pore diameter. The diffusion coefficient was calculated using a linear regression fit of the mean-squared displacement (MSD) and the error bars are estimated by taking the difference of diffusion coefficients from fits over the two halves of the fit interval.⁵⁰ Here, we see that the diffusion coefficients of both water and CO₂ are reduced by confinement, yet the magnitude of this reduction is considerably greater for CO₂. The diffusion ratios obtained for confined water are in excellent agreement with experimental results. For example, Takahara et al. found the diffusion ratio of water confined in mesoporous silica (2.84 nm diameter) to be 0.63 using neutron scattering experiments⁵³ while the diffusion ratio we found for our 3 nm pore is 0.64 ± 0.02 . In the largest hydrophilic pore, we find the absolute value of the CO₂ diffusion coefficient to be $\sim 1.86 * 10^{-5} \text{cm}^2/\text{s}$, which is more than 1 magnitude greater than the diffusion coefficient in an equivalently sized hydrophobic, IL filled pore (*vide infra*).

To explain the difference in confinement effects on the transport of water and CO₂, we turn to the liquid structure. The number density of CO₂ and H₂O molecules, starting at the center of the pore and proceeding radially, is shown in **Figure 3** for the various pore diameters. For all pore sizes, the density of water molecules remains relatively constant

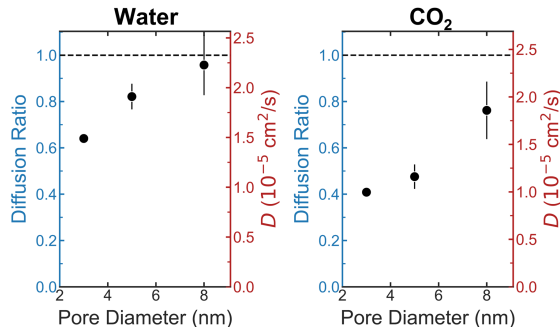


Figure 2: The diffusion coefficients of water (left) and CO₂ (right) in 3, 5, and 8 nm hydrophilic pores filled with water and CO₂. The left y-axis shows the diffusion coefficient compared against the value for an equivalent bulk solvent system while the right y-axis displays the diffusion coefficient absolute value. The dashed horizontal line denotes the bulk diffusion coefficient.

within the pores. In contrast, the CO₂ density is greatly increased near the pore surface to values as high as 4 times the CO₂ density found near the center of the pore. This partitioning likely results from interactions between the surface hydroxyl group that can act as hydrogen bond donors toward the CO₂ oxygens. Previously, researchers showed CO₂/H₂O partitioning is much different near surfaces without hydroxyl groups (e.g., calcite⁵⁴) where the surface can only act as a hydrogen bond acceptor, rather than a donor. The result of the partitioning seen here, with increased CO₂ near the pore walls, is that a much larger fraction of the CO₂ molecules are interacting with the mostly immobile pore surface groups. In fact, **Figure 4** shows a two-dimensional combined radial/radial distribution function (RDF), where the distance between each oxygen in CO₂ and different hydroxyl hydrogens on the pore surface are plotted. The most common orientation for CO₂ is where the molecule is aligned with the pore wall and both oxygens are only 3.5 Å from hydroxyl hydrogens. These CO₂-OH interactions result in CO₂ having a smaller diffusion ratio compared with H₂O.

As the liquid structure dictates the diffusion ratio in these systems, we would expect that the diffusion is a function of CO₂ concentration. At larger concentrations of CO₂, we would expect more CO₂ near the center of the pore, as much of the pore wall is already occupied by CO₂ molecules. To test this hypothesis, we modeled two additional systems at $\sim 3\times$ and $\sim 6\times$ the concentration of CO₂. At triple the CO₂ concentration, the diffusion coefficient

of CO_2 is modestly increased by $\sim 8\%$, yet in the 6x system the diffusion coefficient is dramatically reduced (**Figure S2**). **Figure S3** shows snapshots comparing the initial system with the 6x CO_2 system, where it is clear that the CO_2 phase separates from the water in the high concentration system, but not in the initial system. In phase-separated systems, the boundary between the CO_2 and H_2O serves as a barrier to diffusion, resulting in slow CO_2 diffusion along the length of the pore. Nevertheless, the typical CO_2 partial pressure in flue gas is not large enough to reach the CO_2 concentration modeled in this system (> 0.14 CO_2 mole fraction). Within the expected concentration range for CO_2 , we expect greater concentrations of CO_2 to result in a small increase in the diffusion coefficient of CO_2 due to an increase in CO_2 partitioning near the center of the pore.

[EMIM]⁺ pore systems

Next, we observed the hydrophobic pore systems, constructed by replacing $\sim 40\%$ of the pore hydrogens (-H) from the hydroxylated silica surfaces with trimethylsilyl groups (TMS, $-\text{Si}(\text{CH}_3)_3$), and loaded with either [EMIM][TFSI]/ CO_2 or [OMIM][TFSI]/ CO_2 . The diffusion coefficients for [EMIM]⁺, [TFSI]⁻, and CO_2 as a function of pore diameter in the [EMIM][TFSI]/ CO_2 systems are shown in **Figure 5a**. The IL components show a significant decrease in diffusion coefficient, even in the largest diameter pore. At the smallest pore size, the [EMIM]⁺ and [TFSI]⁻ components are nearly frozen in place, with diffusion ratios of only 0.031 and 0.015, respectively.

The diffusion coefficient of CO_2 in the 8 nm hydrophobic pore is not significantly different than in bulk IL. This result contrasts to the equivalently-sized aqueous pore system, where the diffusion coefficient of CO_2 is noticeably reduced relative to within bulk water. One explanation for this difference is the reduction of surface hydroxyl (-OH) groups available to interact with the CO_2 molecules in the IL pore systems. While the hydrophobic pores still have some hydroxyl groups on the pore surface, these are obscured by the much larger TMS groups. This explanation is further supported by the density profiles of the [EMIM]⁺,

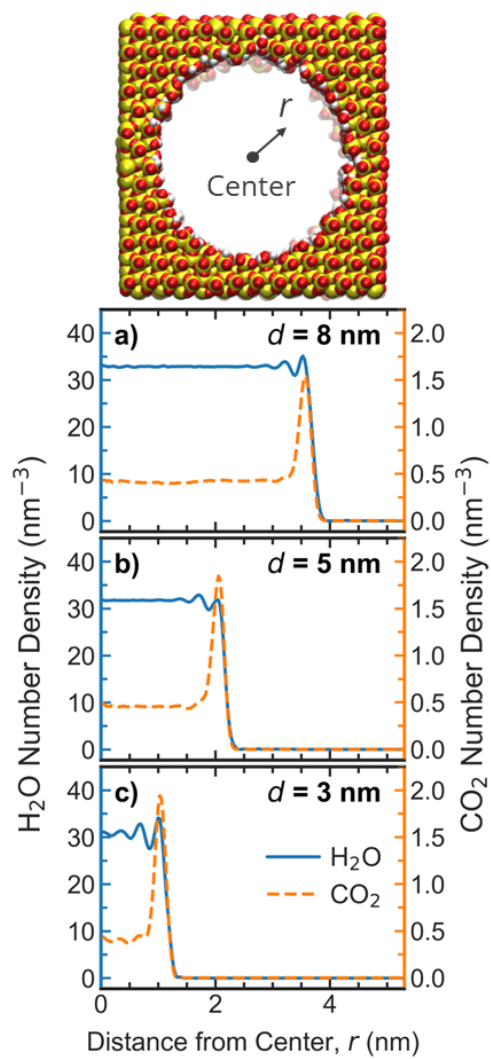


Figure 3: The number density of H₂O and CO₂ molecules shown as a function of distance from the center of the pore in the a) 8 nm, b) 5 nm, and c) 3 nm hydrophilic pore systems filled with water and CO₂.

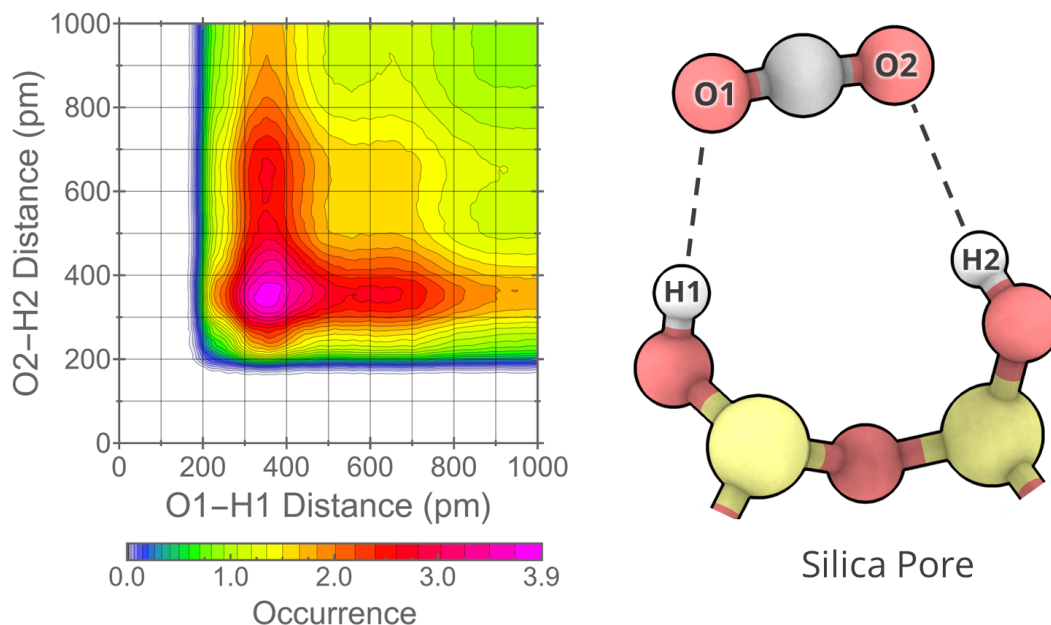


Figure 4: The combined radial/radial distribution function for the two oxygen atoms in CO_2 and different hydroxyl group hydrogens on the pore surface illustrated by the schematic on the right. Grey = Carbon, Red = Oxygen, Yellow = Silicon, White = Hydrogen.

$[\text{TFSI}]^-$, and CO_2 within the hydrophobic pores, which show that the increase in CO_2 density near the pore surface is considerably smaller than the increase observed in the hydrophilic water/ CO_2 -filled pores (**Figure S4**). Nevertheless, in the smallest pore size, the CO_2 diffusion ratio drops to ~ 0.21 . This significant drop is likely due to interactions between the CO_2 and the IL components, which are mostly immobile at this degree of confinement.

$[\text{OMIM}]^+$ pore systems

The diffusion coefficients for $[\text{OMIM}]^+$, $[\text{TFSI}]^-$, and CO_2 as a function of pore diameter in the $[\text{OMIM}][\text{TFSI}]/\text{CO}_2$ system are shown in **Figure 5b**. In the 8 nm diameter pore, the $[\text{OMIM}]^+$ and $[\text{TFSI}]^-$ undergo a larger reduction in diffusion coefficient compared with $[\text{EMIM}][\text{TFSI}]$. This reduction occurs because $[\text{OMIM}]^+$ is a larger molecule than $[\text{EMIM}]^+$ and thus the effects of confinement are comparatively greater. Counterintuitively, we do not see the same trend in the 5 and 3 nm diameter pores. This result arises because of important properties of ILs that make them unique and useful materials – their irregular

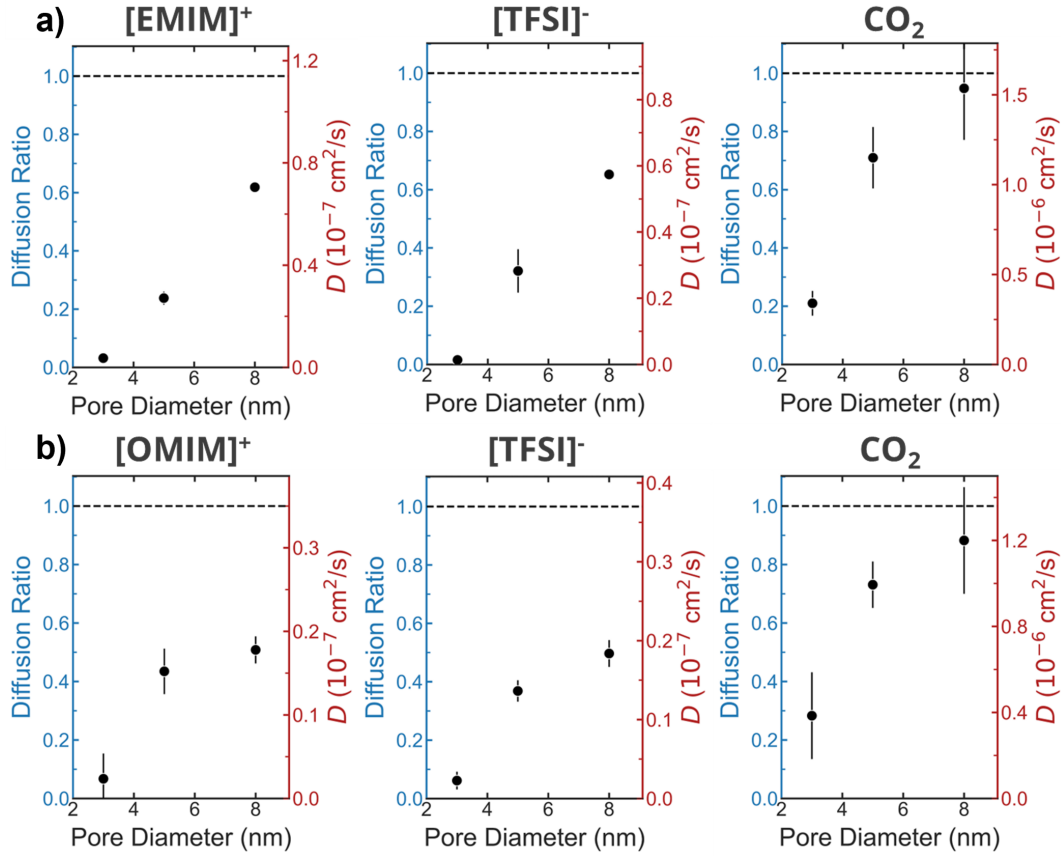


Figure 5: a) The diffusion coefficients of [EMIM]⁺ (left), [TFSI]⁻ (middle), and CO₂ (right) as a function of hydrophobic pore size for the [EMIM][TFSI] system. b) The diffusion coefficients of [OMIM]⁺ (left), [TFSI]⁻ (middle), and CO₂ (right) as a function of hydrophobic pore size for the [OMIM][TFSI] system. The left y-axis shows the diffusion coefficient compared against the value for an equivalent bulk liquid system while the right y-axis displays the diffusion coefficient absolute value. The dashed horizontal line denotes the bulk diffusion coefficient.

and asymmetrical shapes. In bulk liquid, the asymmetrical shape of ILs prevent efficient packing and result in ILs remaining liquid at room temperature despite strong intermolecular forces. Under nanoconfinement in the 3 nm pores, we see a similar effect where $[\text{OMIM}]^+$, having a much larger alkyl group and more asymmetric shape than $[\text{EMIM}]^+$, cannot undergo efficient packing, maintains a larger diffusion ratio, and more liquid character when compared with $[\text{EMIM}]^+$. This explanation can be corroborated by comparing the density profile for the $[\text{OMIM}][\text{TFSI}]/\text{CO}_2$ system (**Figure S5c**) to the $[\text{EMIM}][\text{TFSI}]/\text{CO}_2$ system (**Figure S4c**). Here, the density profiles can be interpreted similarly to RDFs, where particularly large peak values along with small valley values indicate highly structured or immobilized fluids.⁵⁵ When looking at the prominent features in the density profiles, the maximum and minimum closest to the pore wall, we find a higher degree of structure for the $[\text{EMIM}]^+$ system, with an average maximum value/minimum value ratio of 6.45 for the three liquid components, compared with the $[\text{OMIM}]^+$ system ratio of only 3.78. This result confirms that the long $[\text{OMIM}]^+$ alkyl tail disrupts the packing of all liquid components and results in larger diffusion ratios in the 3 nm pore when compared with the $[\text{EMIM}]^+$ counterpart system.

The diffusion coefficient of CO_2 in the $[\text{OMIM}][\text{TFSI}]/\text{CO}_2$ system is similar to the IL components. At the largest pore size, the diffusion ratio for CO_2 in the $[\text{OMIM}]^+$ system is smaller than for CO_2 in the $[\text{EMIM}]^+$ system, yet this trend is once again broken at the smaller pore diameters. Interestingly, in the 3 nm pore system, the absolute diffusion coefficient of CO_2 is larger in the $[\text{OMIM}]^+$ system compared to the $[\text{EMIM}]^+$ system (albeit within the error of the MD simulations). This finding is surprising as bulk $[\text{OMIM}][\text{TFSI}]$ is $\sim 3\times$ more viscous than bulk $[\text{EMIM}][\text{TFSI}]$ ⁵⁶ and we had previously found CO_2 diffuses faster in bulk $[\text{EMIM}][\text{TFSI}]$ relative to bulk $[\text{OMIM}][\text{TFSI}]$;²⁶ however, the nanoconfinement effects of the 3 nm pore disrupts the liquid structure of $[\text{OMIM}][\text{TFSI}]/\text{CO}_2$ comparatively more than $[\text{EMIM}][\text{TFSI}]/\text{CO}_2$, enough to reverse this trend. Moreover, CO_2 - $[\text{EMIM}]^+$ and CO_2 - $[\text{OMIM}]^+$ RDFs for bulk liquid and all three pore sizes are shown in **Figure S6**, allowing

us to see the effect of pore size on CO_2 /cation structure directly. For both ILs, the RDFs for the bulk, 8 nm pore, and 5 nm pore systems are similar. Nevertheless, for the 3 nm pore systems, we observe an increase in the first maximum value for some of the CO_2 -[EMIM] $^+$ RDFs, while we observe a decrease in first maximum values for all of the CO_2 -[OMIM] $^+$ RDFs. This data indicates that, in the 3 nm pores, the [EMIM] $^+$ system is more structured than the [OMIM] $^+$ system, consistent with the density profiles and the diffusion coefficients.

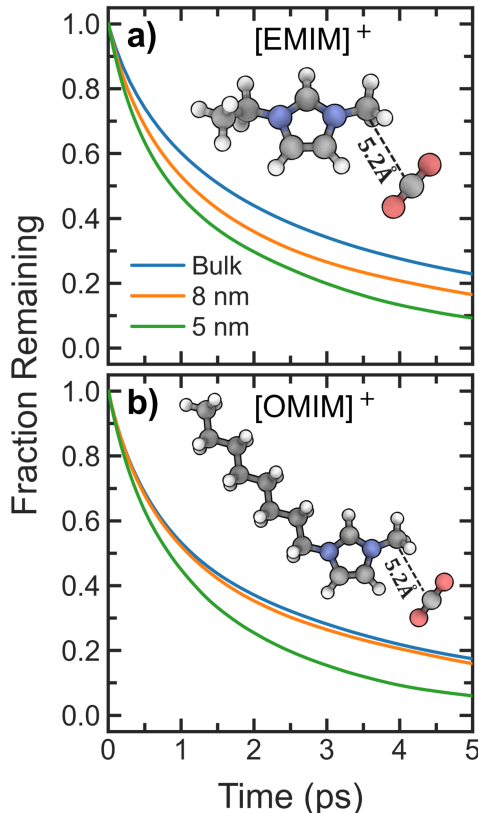


Figure 6: Dimer existence autocorrelation functions for a) [EMIM] $^+$ - CO_2 and b) [OMIM] $^+$ - CO_2 dimer pairs in bulk and 8 nm and 5 nm pores.

CO_2 /cation interactions

The RDFs shown here and those reported in the literature show strong associations between various imidazolium cation atoms and CO_2 atoms.^{26,57} This observation poses a question as to how, under nanoconfinement, the CO_2 diffusion ratio is larger than that of the

imidazolium cations they appear to be strongly interacting with. To elucidate this result, we examined the continuous dimer existence autocorrelation functions (DACFs) for the methyl carbon on the imidazolium cation and carbon in CO_2 shown in **Figure 6**. The methyl carbon on the cation was chosen for this analysis as it shows the strongest interaction with CO_2 , and a dimer cutoff value of 5.2 Å was chosen based on the distance where the first solvation shell ends in the RDFs (**Figure S6**). Here, we see that nanoconfinement destabilizes the CO_2 solvation of the imidazolium cations, resulting in shorter dimer lifetimes. Nevertheless, the RDFs in **Figure S6** for the bulk, 8 nm, and 5 nm systems show little differences under confinement. Therefore, under confinement, the amount of time a CO_2 molecule interacts with an individual imidazolium cation decreases, but the likelihood that a CO_2 is interacting with *any* imidazolium cation is comparable to that in bulk. As such, the diffusion ratio of CO_2 exceeds that of the imidazolium cation component because the CO_2 molecules are jumping between imidazolium cations more frequently.

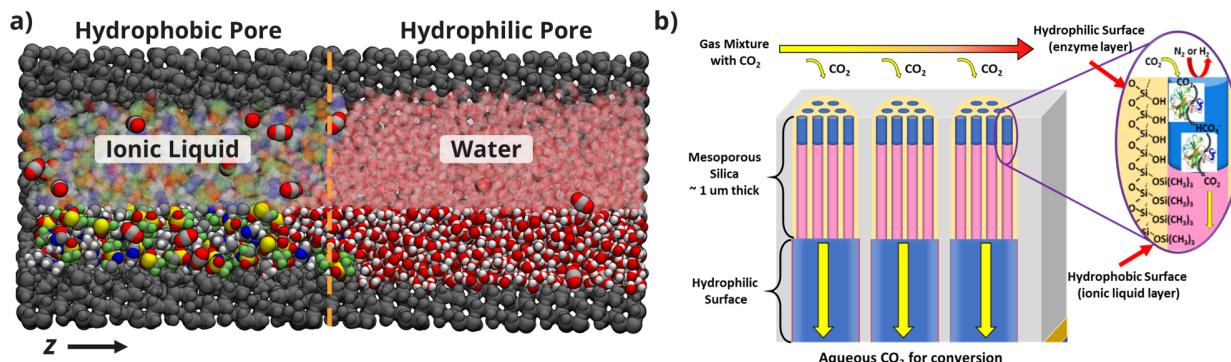


Figure 7: a) Snapshot of [EMIM][TFSI]/H₂O sliced open with the hydrophobic and hydrophilic surface termination to the left and right of the dashed orange line, respectively. CO_2 emphasized to show presence located predominately in the ionic liquid phase. Grey = Carbon, Blue = Nitrogen, Red = Oxygen, Yellow = Sulfur, Green = Fluorine, White = Hydrogen, Dark Grey = Pore. b) Schematic of nanoporous CO_2 capture device with aqueous layers (blue) and ionic liquid layer (pink).

Interfacial aqueous/ionic liquid systems

Finally, we constructed interfacial systems with a pore diameter of 5 nm and which

contain regions of hydrophilic and hydrophobic surface groups, with water/CO₂ loaded into the hydrophilic region and IL/CO₂ loaded into the hydrophobic region. A snapshot of the [EMIM][TFSI]/CO₂/H₂O system is shown in **Figure 7a** and the schematic for a CO₂ capture device with aqueous and ionic liquid layers, that incorporates an enzymatic capture layer, is shown in **Figure 7b**. Free energy profiles for the CO₂ and H₂O components in the [EMIM]⁺ and [OMIM]⁺ systems are shown in **Figure 8a** and **b**, respectively. In both systems, CO₂ has a negative free energy (ΔG) in the IL phase, indicating favorable transport from the water to IL phase. The magnitude of ΔG is comparable for the [EMIM]⁺ ($\Delta G \sim -4.62$ kJ/mol) and [OMIM]⁺ ($\Delta G \sim -3.97$ kJ/mol) solutions. These values are slightly less than the values found in bulk (-5.86 kJ/mol and -5.66 kJ/mol for the [EMIM]⁺ and [OMIM]⁺ systems).²⁶ This reduction could be a result of the hydroxyl groups in the hydrophilic region making the water phase more favorable and, by comparison, reducing the ΔG when transporting from the aqueous to IL phases. Conversely, the ΔG for H₂O is positive and large (14.39 kJ/mol and 17.28 kJ/mol for the [EMIM]⁺ and [OMIM]⁺ systems), indicating that H₂O molecules prefer maintaining a separate phase and do not mix with the IL components.

Conclusion

CO₂ separation membranes are a critical technology to reduce emissions. Our MD results reveal numerous insights that can help understand and improve CO₂ separation membrane technologies. In confined aqueous systems, a large density of hydroxyl groups result in increased CO₂ density near the pore walls and an overall greater reduction in diffusion in the axial direction of the pore, even at the relatively large pore diameter of 8 nm. Conversely, in [EMIM][TFSI] and [OMIM][TFSI], CO₂ maintains diffusion rates comparable to bulk in 8 nm hydrophobic pores. Counterintuitively, at greater degrees of nanoconfinement (e.g., 3 nm diameter), bulkier and more asymmetric ILs (for example, [OMIM]⁺ as opposed to [EMIM]⁺) may undergo less efficient packing, allowing the system to maintain higher liquid character and faster diffusion rates. We expect this phenomenon, where asymmetric ILs maintain more

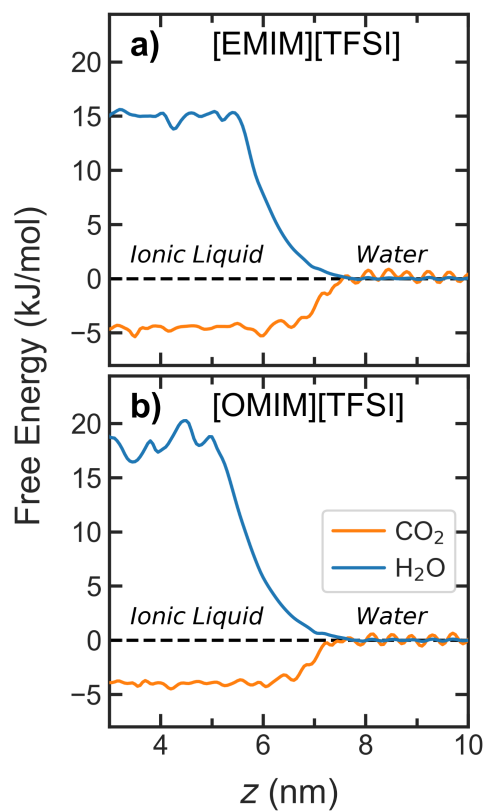


Figure 8: Free energy profiles for CO₂ and H₂O in the a) [EMIM][TFSI]/H₂O and b) [OMIM][TFSI]/H₂O interfacial, 5 nm diameter pores. The interface between the ionic liquid and water occurs at $z \sim 7$ nm.

liquid character under confinement, to be generalizable to other nanoconfined IL membrane systems. The nanoconfined [EMIM][TFSI]/H₂O and [OMIM][TFSI]/H₂O systems modeled here maintain separate IL and aqueous phases, where CO₂ is shown to transport spontaneously from the aqueous to IL phase. Future experimental studies, such as with pulsed-field gradient nuclear magnetic resonance spectroscopy and neutron scattering, could look to verify the structure and diffusion rates of nanoconfined imidazolium ILs and CO₂.

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Supporting Information Available

Supporting Figures S1-S6 showing ionic liquid and carbon dioxide density profiles, cation/carbon dioxide radial distribution functions, a schematic illustration of the pore loading procedure, snapshots from aqueous systems with an increased concentration of CO₂, and diffusion coefficients for aqueous systems with an increased concentration of CO₂

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TOC Graphic

