

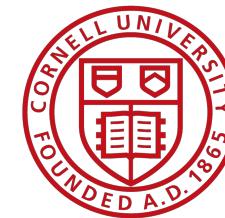
# Predictions of novel features in x-ray scattering spectra for thermometry

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Stephanie Hansen<sup>1</sup>, and Andrew Baczewski<sup>1</sup>



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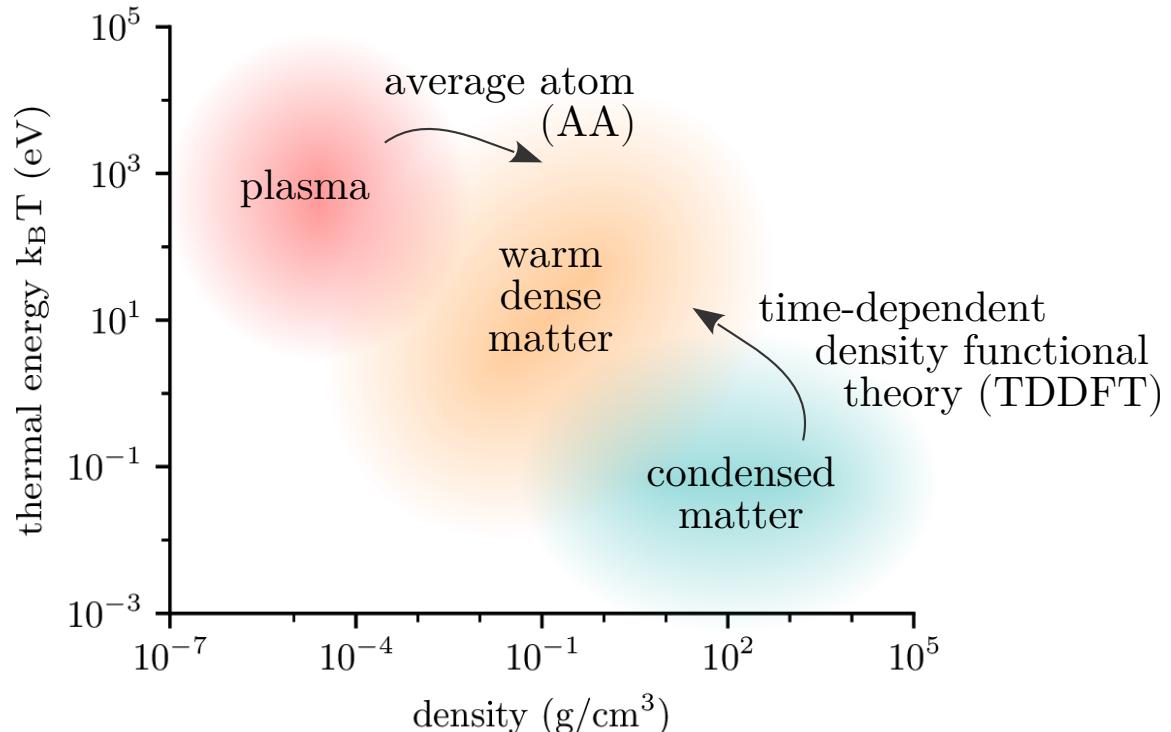
<sup>2</sup> Cornell University



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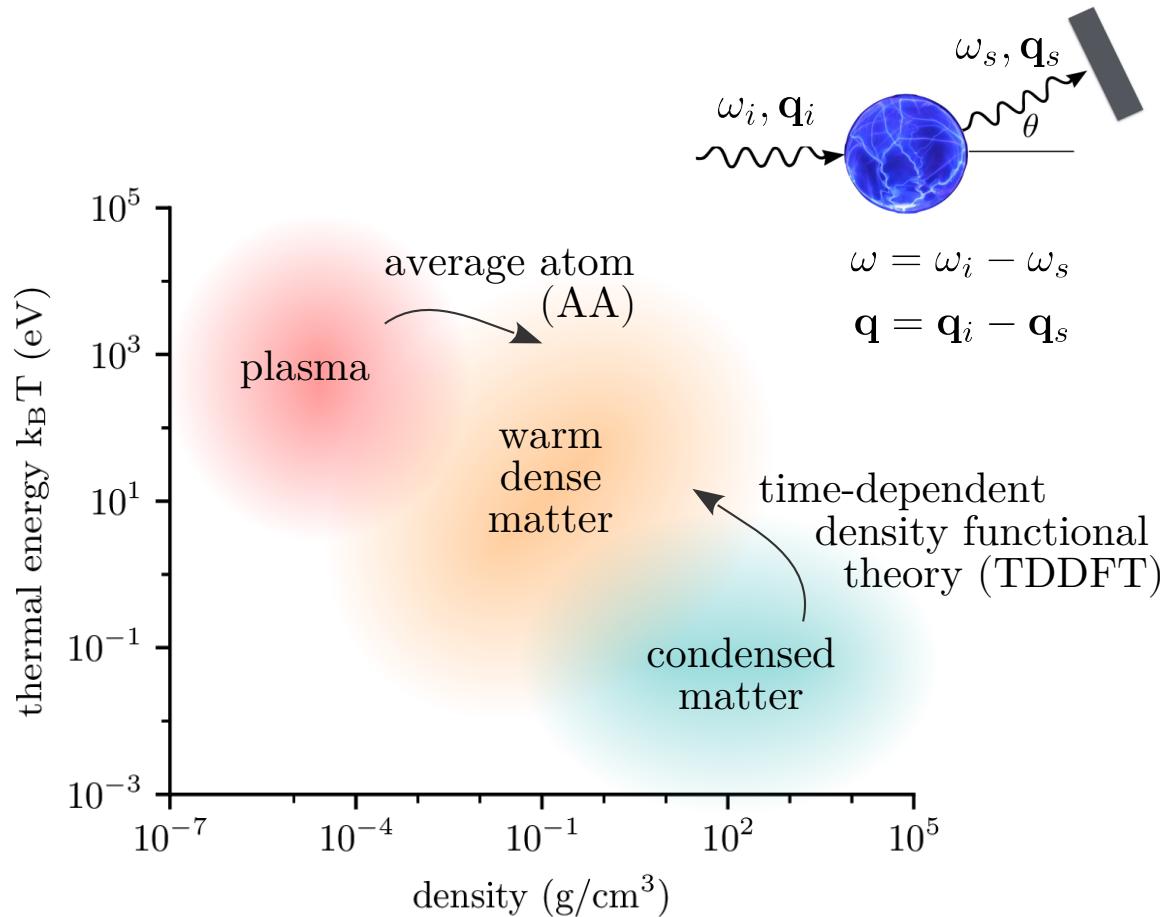
# Bridging the Gap between Plasma and Condensed Matter

- Fusion fuel passes through poorly understood WDM regime on the way to ignition
- Competing physics challenges models
- Scarce experimental data, uncertain conditions limit model validation
- Validate modified AA against more accurate but expensive TDDFT
- Predict potential improvements for x-ray diagnostic techniques



# Bridging the Gap between Plasma and Condensed Matter

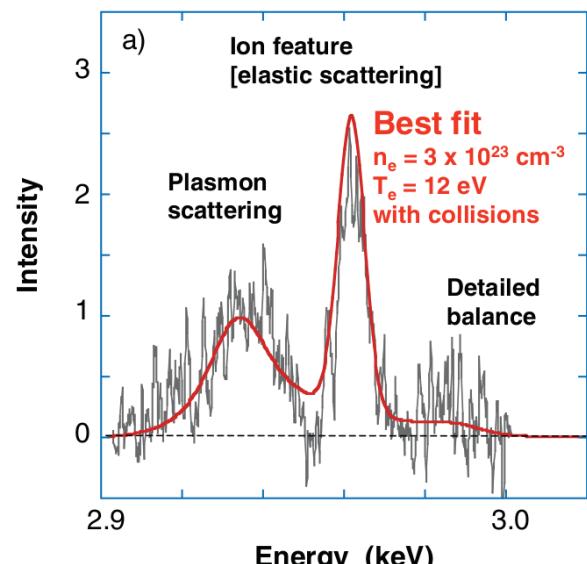
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# Diagnosing Temperatures with X-ray Scattering

Electron temperature:

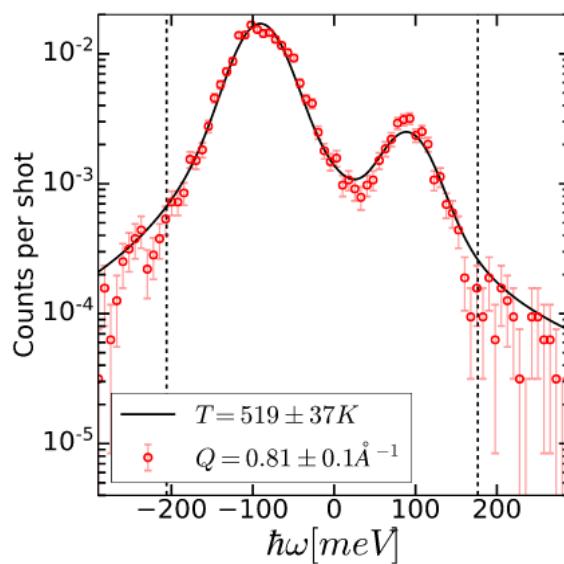
- detailed balance between  $\sim 10$  eV red- and blue-shifted plasmon features
- sensitivity limited to  $T \sim 10$  eV
- alternative: bound-bound features



$$\frac{S(\mathbf{q}, \omega)}{S(\mathbf{q}, -\omega)} = e^{\hbar\omega/(k_B T)}$$

Ion temperature:

- detailed balance between  $\sim 100$  meV red- and blue-shifted phonon features
- detector limitations challenging
- complement: band structure effects



# Simulating X-ray Scattering with Time-Dependent DFT

- Initial condition: equilibrium state from Mermin-DFT
- Evolve response to probe in real time

$$i \frac{\partial}{\partial t} \phi_j(\mathbf{r}, t) = \hat{H}[n(\mathbf{r}, t)] \phi_j(\mathbf{r}, t)$$

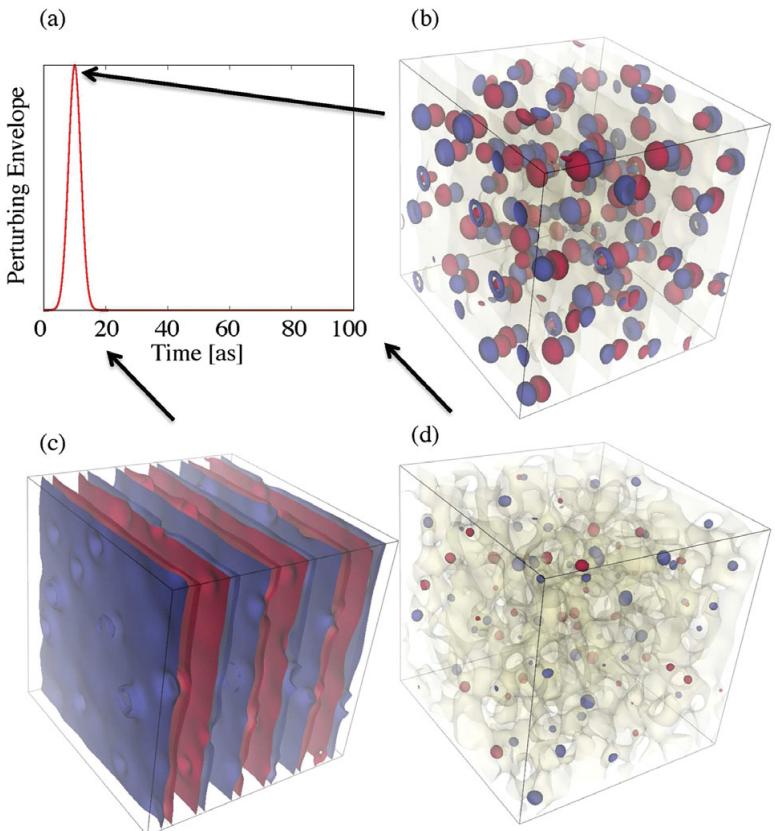
$$n(\mathbf{r}, t) = \sum_j f_j(T) |\phi_j(\mathbf{r}, t)|^2$$

$$V_{\text{probe}}(\mathbf{r}, t) = V_0 e^{i\mathbf{q} \cdot \mathbf{r}} f(t)$$

- Dynamic structure factor (DSF) related to Fourier transform of density response

$$\tilde{\chi}(\mathbf{q}, -\mathbf{q}, \omega) = \delta \tilde{n}(\mathbf{q}, \omega) / (V_0 \tilde{f}(\omega))$$

$$S(\mathbf{q}, \omega) = -\frac{1}{\pi} \frac{\text{Im} [\tilde{\chi}(\mathbf{q}, -\mathbf{q}, \omega)]}{1 - e^{-\omega/(k_B T)}}$$



Baczewski et al., Phys. Rev. Lett. 116 (2016)

# Simulating X-ray Scattering with Average Atom

- Chihara decomposition:  $S(q, \omega) = S_{ii}(q, \omega) + S_{FF}(q, \omega) + S_{BF}(q, \omega) + S_{BB}(q, \omega)$

- Modified  $S_{FF}$  captures non-Drude behavior
  - Mermin dielectric function
  - non-ideal, quantum DOS
  - T-matrix elastic + inelastic collisions

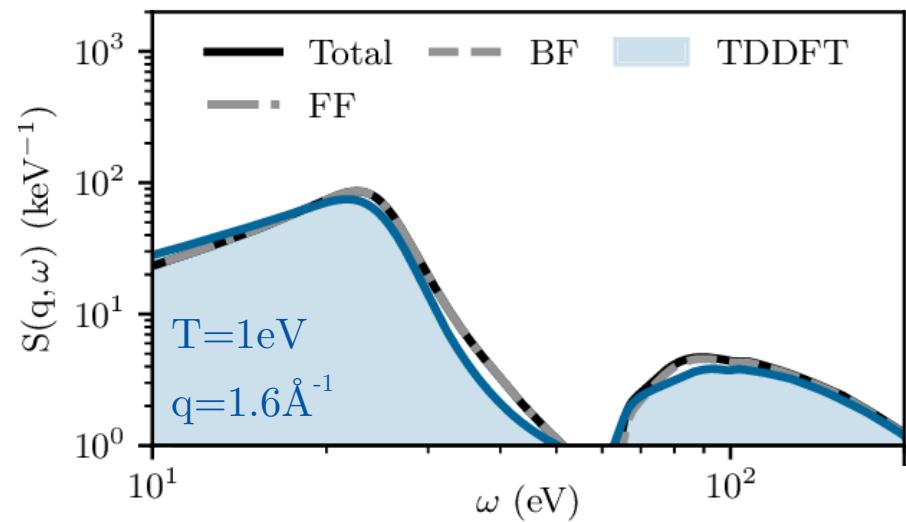
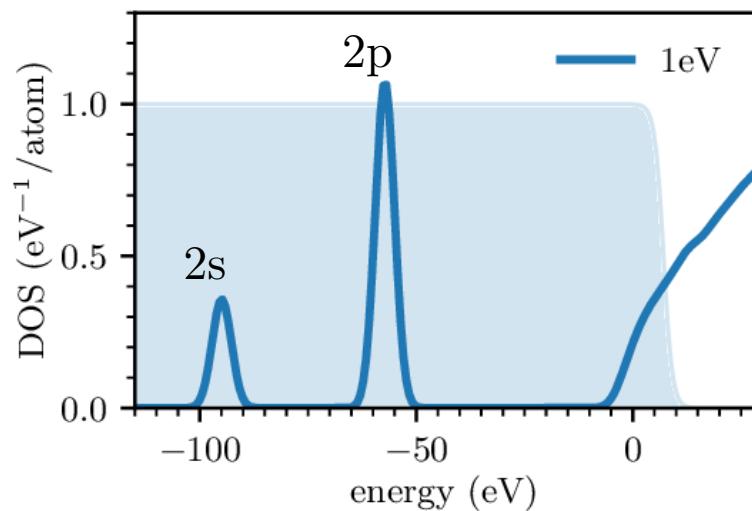
$$S_{FF}(q, \omega) = -\frac{1}{1 - e^{-\omega/(k_B T)}} \frac{q^2}{4\pi n_e} \text{Im} \left[ -\frac{1}{\varepsilon(q, \omega)} \right]$$

$$\epsilon_{\text{MA}}(k, \omega) = 1 + \frac{(\omega + i\nu)[\epsilon_{\text{RPA}}(k, \omega + i\nu) - 1]}{\omega + i\nu \frac{\epsilon_{\text{RPA}}(k, \omega + i\nu) - 1}{\epsilon_{\text{RPA}}(k, 0) - 1}}$$

- Introduce bound-bound  $S_{BB}$  to achieve continuity under ionization of bound states

$$S_{BF}(q, \omega) = \sum_i g_i f(\varepsilon_i) (1 - f(\varepsilon_i + \omega)) \mathcal{M}_i^{BF}(q, \omega) \quad \text{sum over initial and final bound states}$$

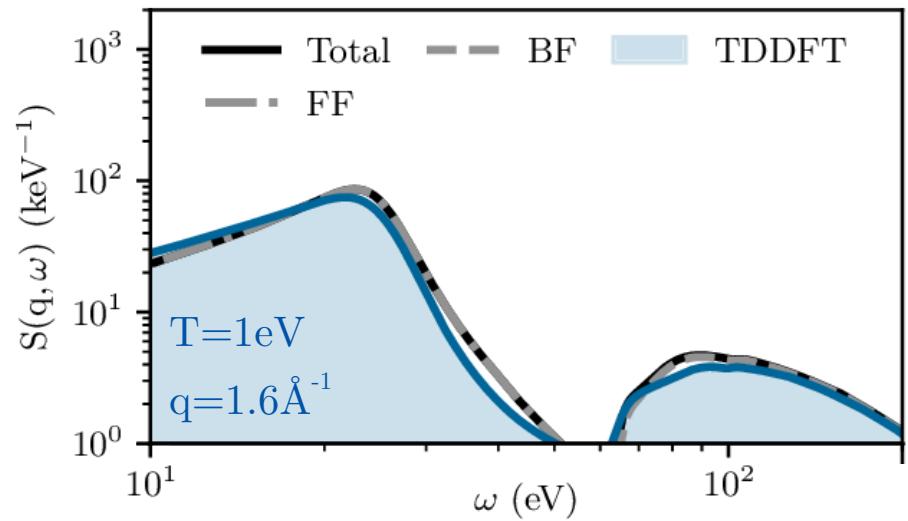
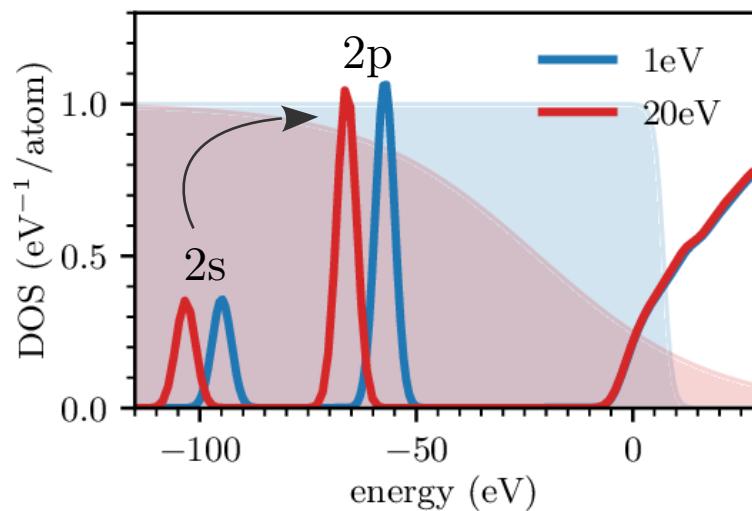
$$S_{BB}(q, \omega) = \sum_{i,f} g_i f(\varepsilon_i) (1 - f(\varepsilon_f)) \phi(\omega) \mathcal{M}_{if}^{BB}(q, \omega) \quad \text{line shape}$$



# Bound-Bound Transitions in Aluminum

arXiv:2109.09576

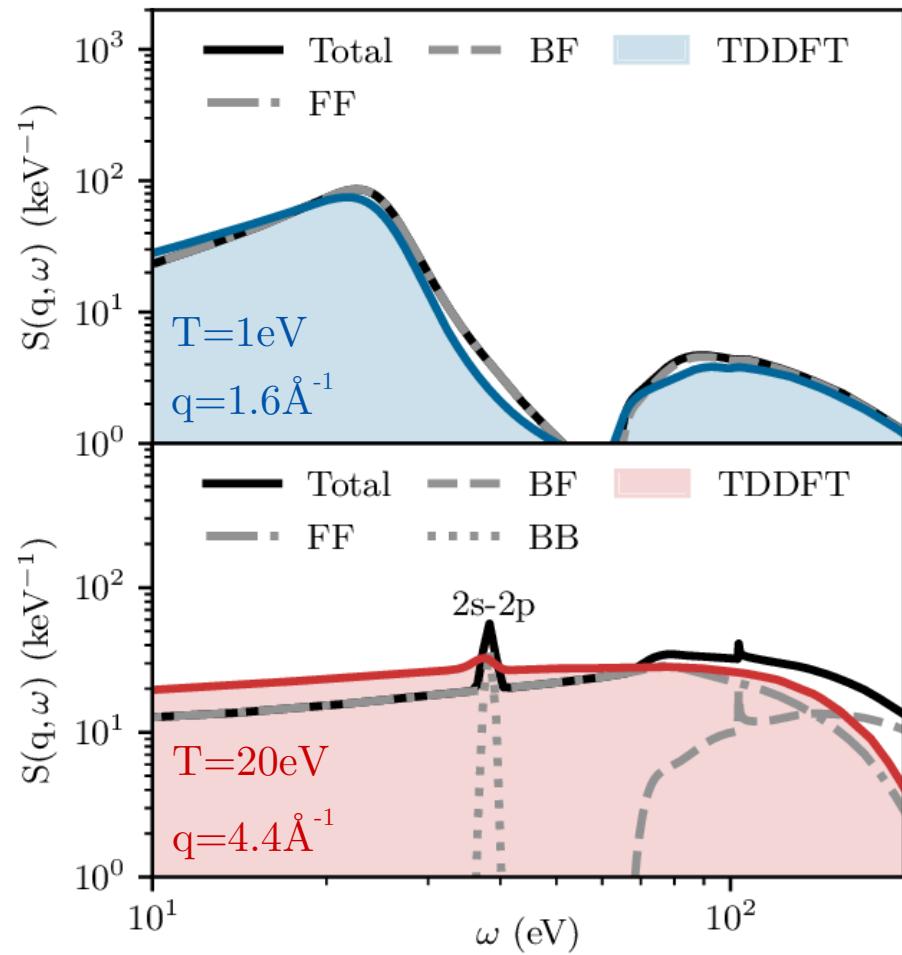
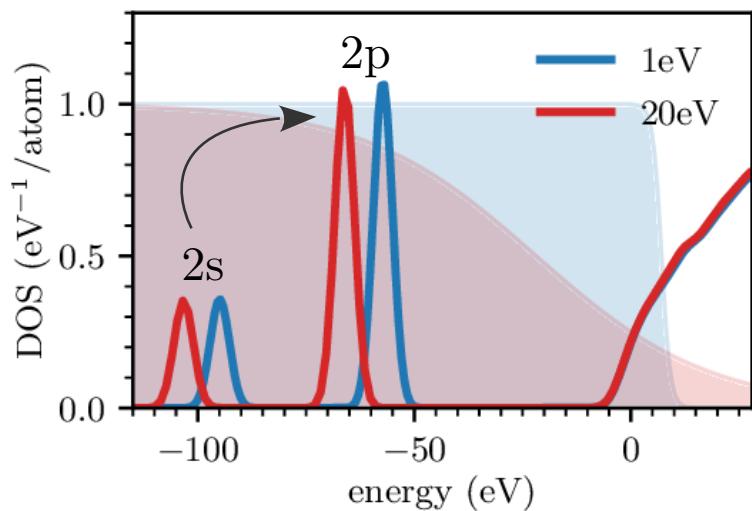
- Thermal depletion of 2p states for  $T \gtrsim 20\text{eV}$



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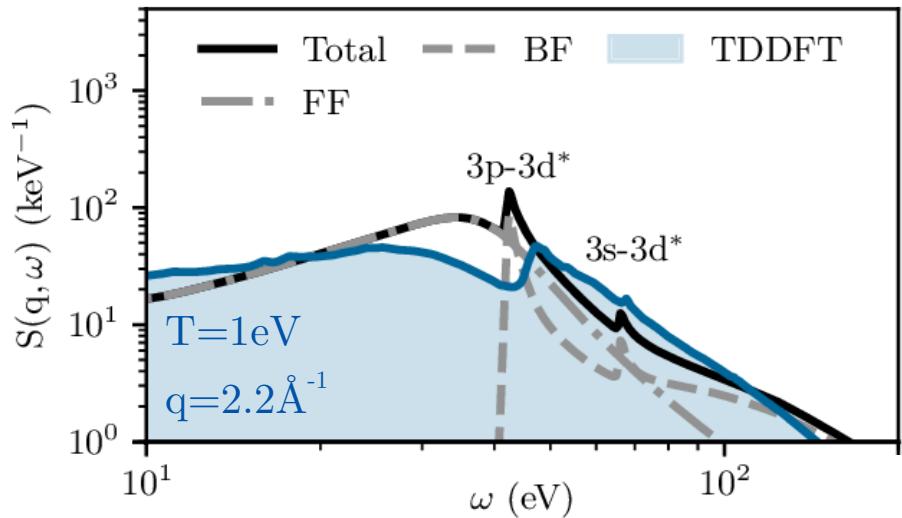
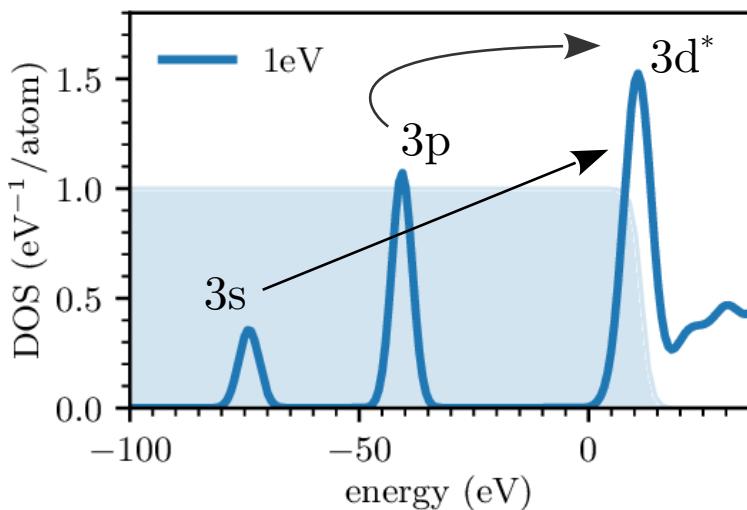
- Thermal depletion of 2p states for  $T \gtrsim 20\text{eV}$
- Prominent  $2s \rightarrow 2p$  feature at  $\sim 40\text{eV}$  in DSF
  - particularly towards higher angles
  - very sensitive to  $T$  via 2p occupation



# Bound-Bound Transitions in Iron

arXiv:2109.09576

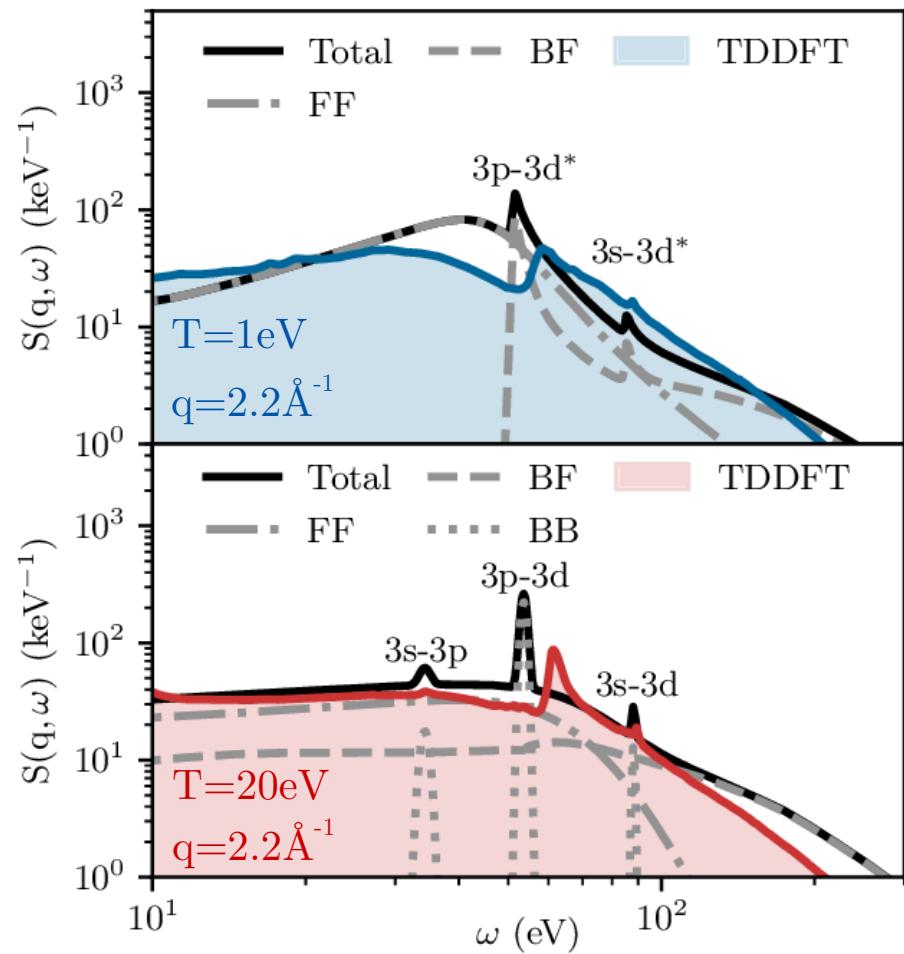
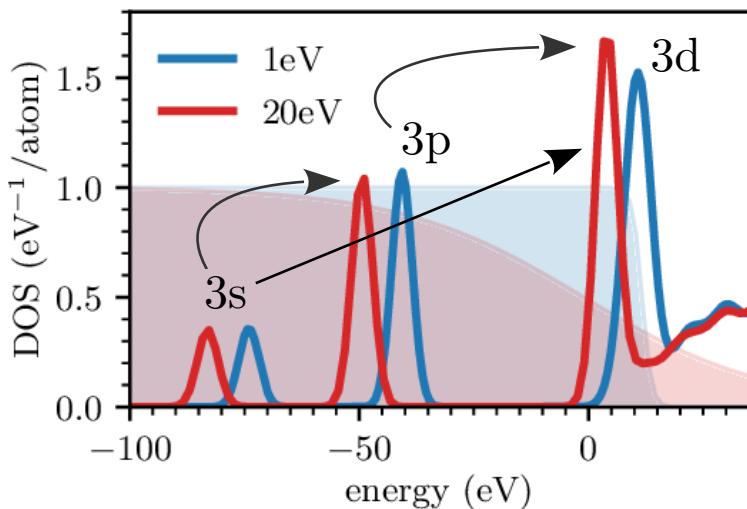
- Richer set of transitions:
  - $3p \rightarrow 3d$  at  $\sim 55\text{eV}$
  - $3s \rightarrow 3d$  at  $\sim 85\text{eV}$



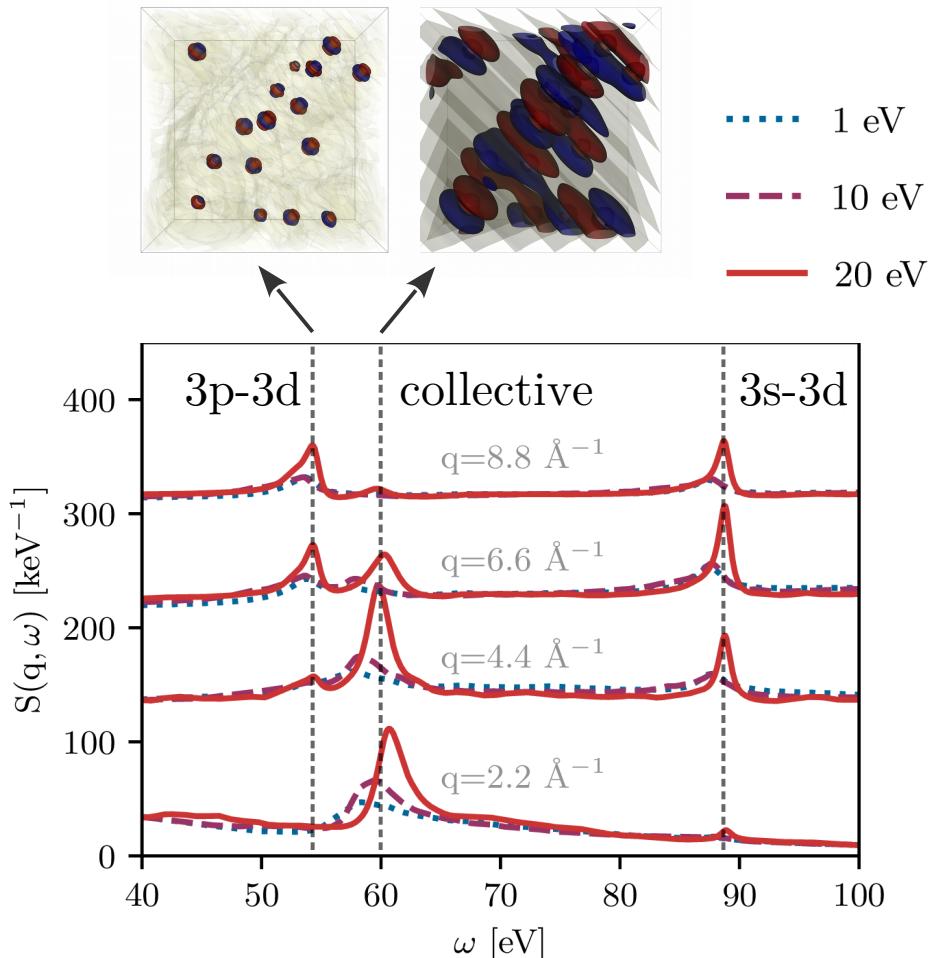
## Bound-Bound Transitions in Iron

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- Richer set of transitions:
  - $3p \rightarrow 3d$  at  $\sim 55\text{eV}$
  - $3s \rightarrow 3d$  at  $\sim 85\text{eV}$
  - $3s \rightarrow 3p$  at  $\sim 35\text{eV}$
- Apparent discrepancy in  $3p$ - $3d$  energy!
  - $\sim 5\text{eV}$  higher in TDDFT than AA



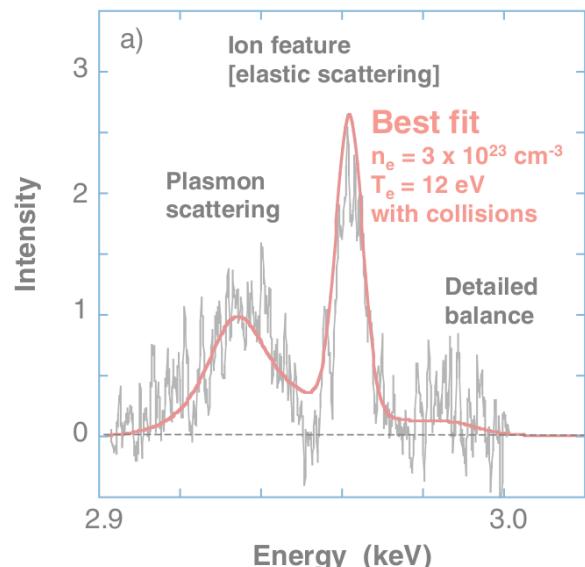
- For high  $q$ :
  - 3p-3d feature appears at expected energy
  - density response localized
- For low  $q$ :
  - density response delocalized
  - higher-energy feature arises from new collective mechanism!



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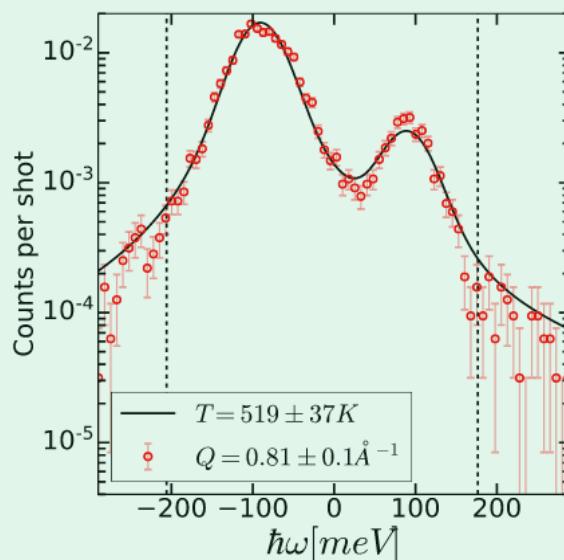
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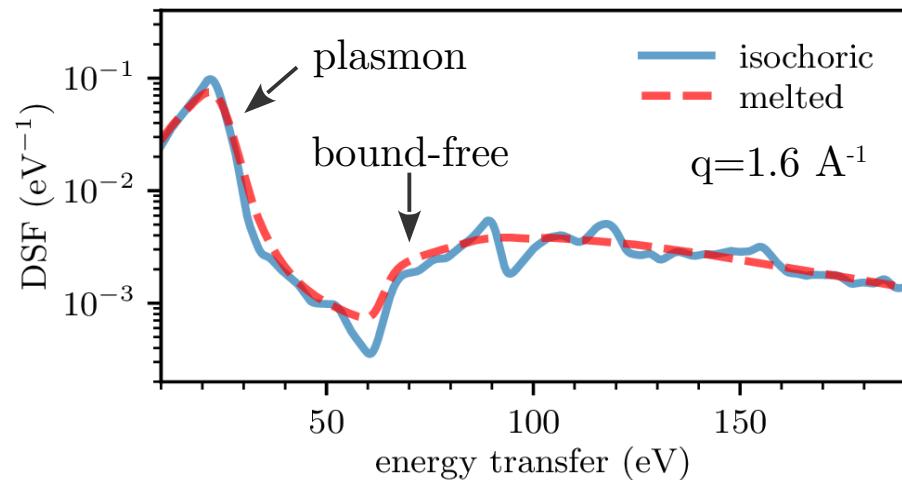
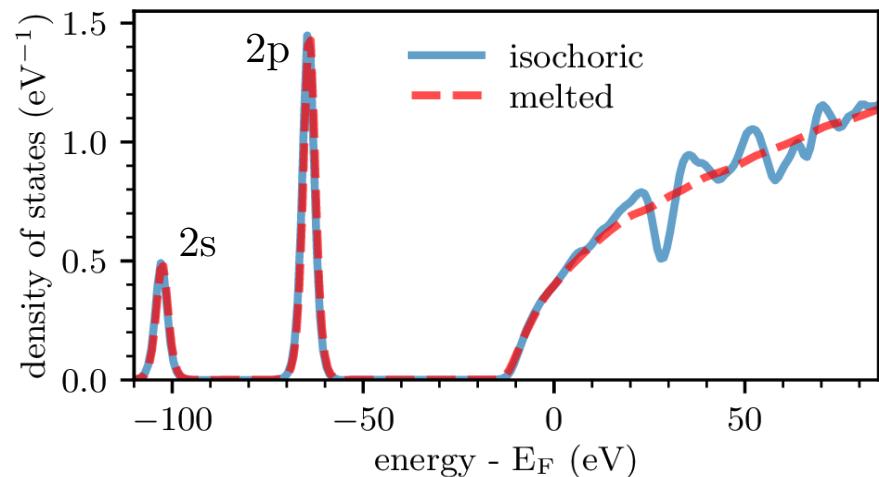
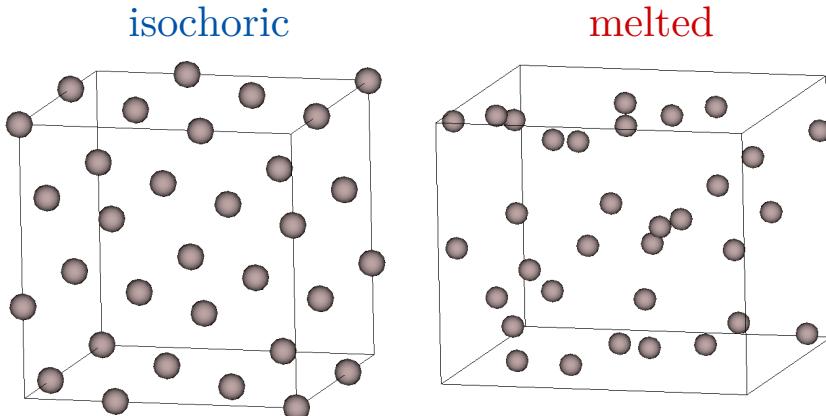
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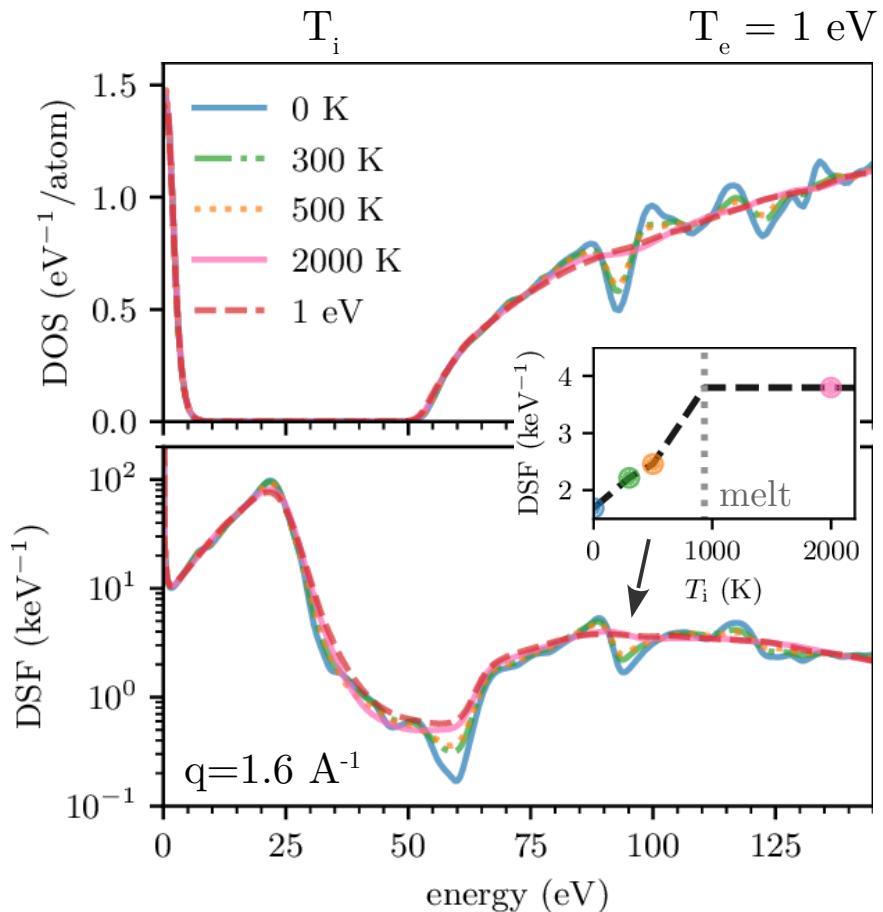
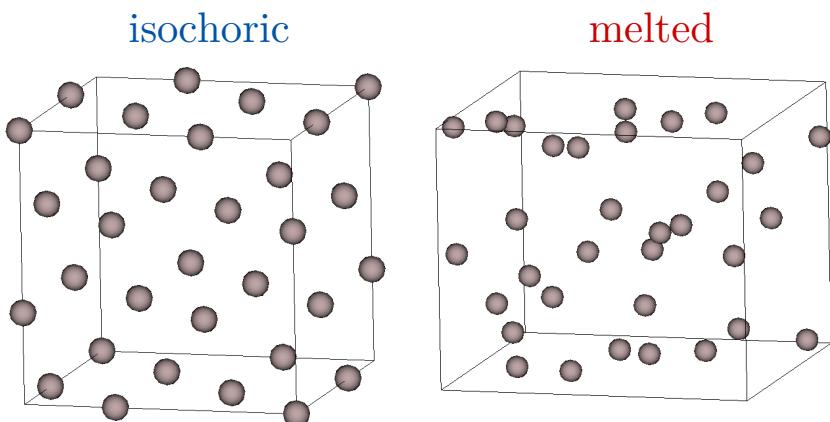
# Band Structure Effects out of Equilibrium

- Melted Al ( $T_i = T_e = 1\text{eV}$ )
  - $\sim$ ideal free-electron gas
- Isochorically heated Al ( $T_i = 0$ ,  $T_e = 1\text{eV}$ )
  - gaps among high-energy bands
- DSF reflects details of band structure
- May help simultaneously infer  $T_e$  and  $T_i < T_{\text{melt}}$  !



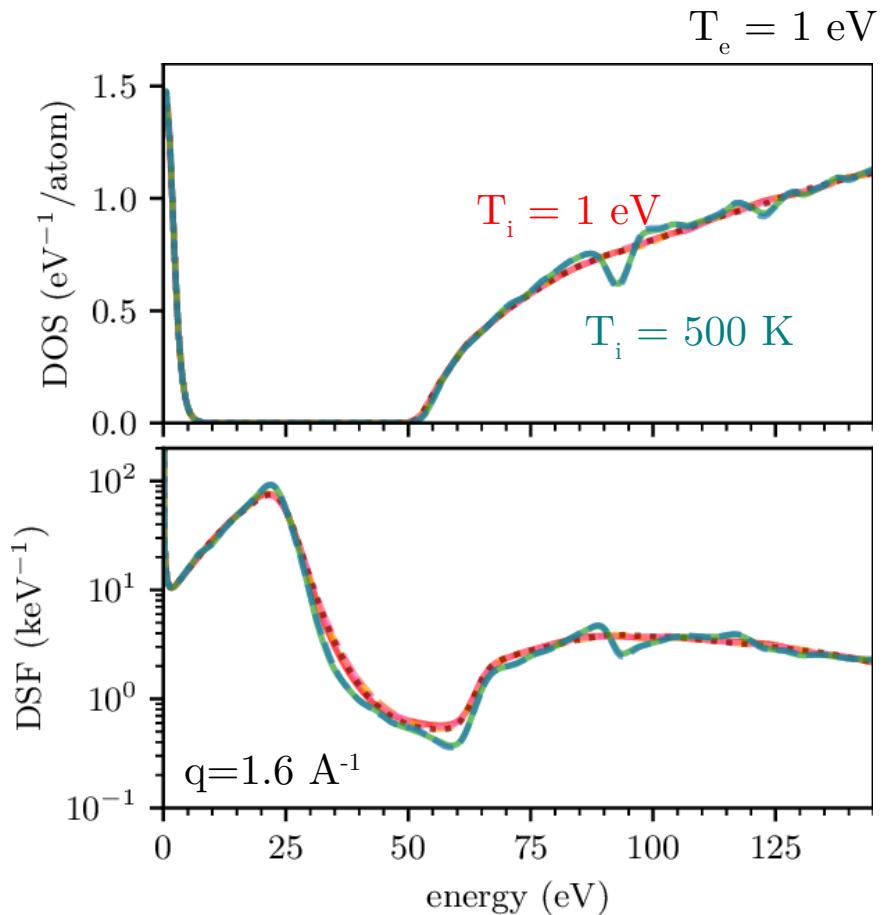
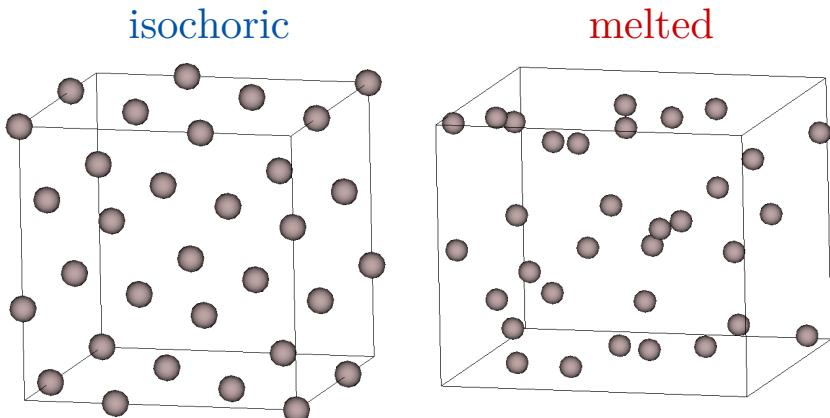
## Band Structure Effects

- Lower amplitude features with increasing  $T_i$
- $T$ -sensitivity appears limited
  - may improve at higher scattering angles



# Band Structure Effects

- Lower amplitude features with increasing  $T_i$
- T-sensitivity appears limited
  - may improve at higher scattering angles
- Results insensitive to atomic configuration at given T



## Conclusions

- Predict prominent bound-bound features in DSFs at high T
  - promising alternative to plasmon-based thermometry for  $k_B T_e > \hbar\omega_p$
  - novel collective behavior in iron's 3p – 3d transition
- Predict subtle band structure features in DSFs of isochorically heated materials
  - promising alternative to phonon-based thermometry for  $T_i < T_{melt}$
  - possibilities for simultaneously constraining  $T_e$  and  $T_i$
  - promising for studies of ultrafast melting

