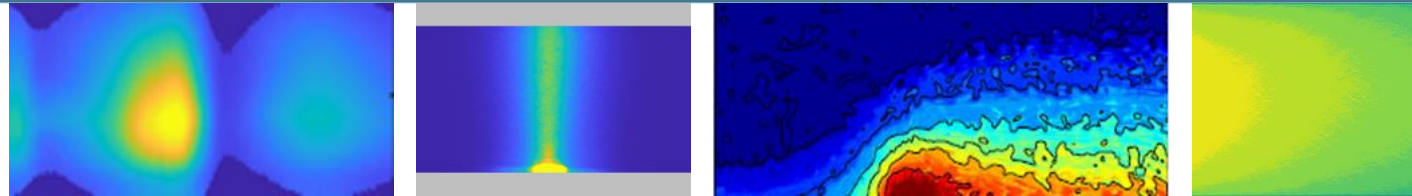




# SPRF: Photoemission Induced Plasma Breakdown in Argon



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Yang Zhou<sup>2</sup>, and Peng Zhang<sup>2</sup>

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# Photoelectron Emission

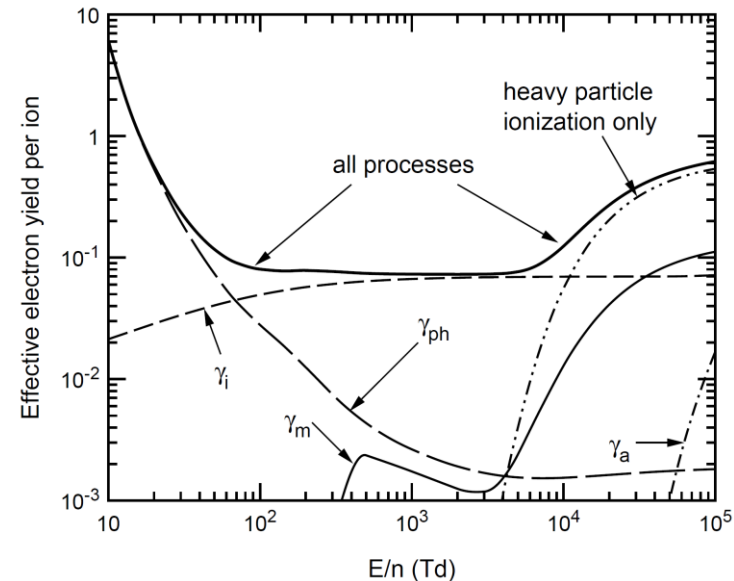
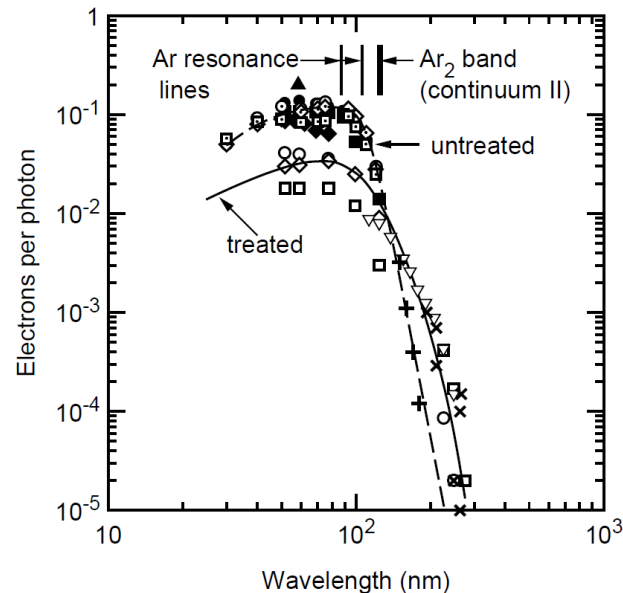


Photons supply energy to release electrons from surfaces

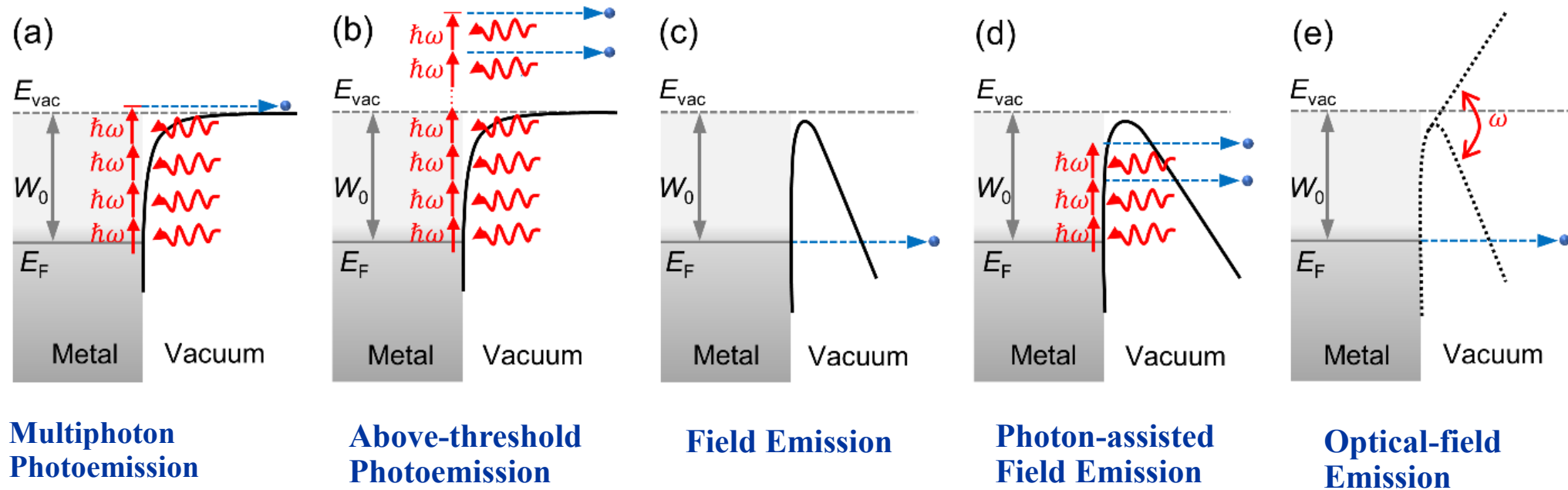
- Photoemission influences or sustains plasma (at low  $E/N$ )

Laser-induced photoemission not as well understood

- Important for laser-triggered breakdown



# Laser-Induced Photoemission



Emission process depends on laser wavelength, intensity, DC bias, surface properties

Pursued experiments to validate photoemission quantum model [1]

- Better understanding may provide more control of plasma breakdown and plasma properties

# Experiments in Low Space Charge Regime ( $<100 \mu\text{A}$ )

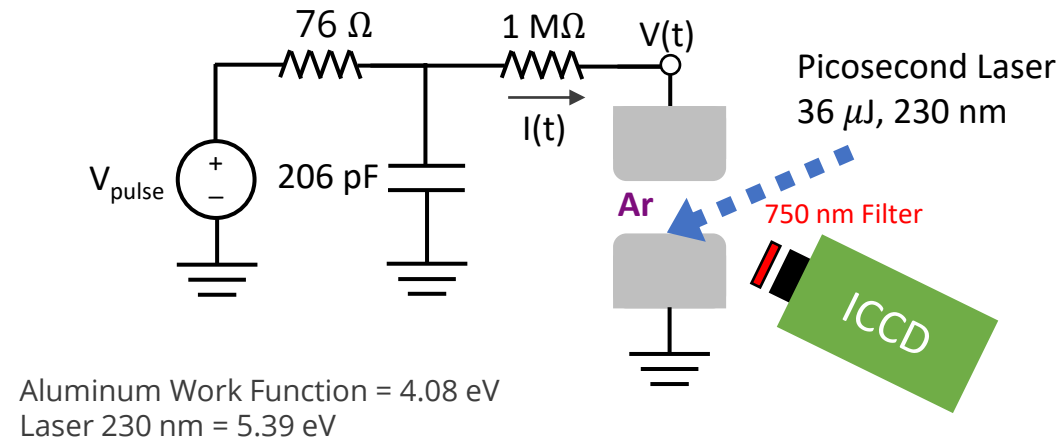


## De-couple electrode phenomena from space charge effects

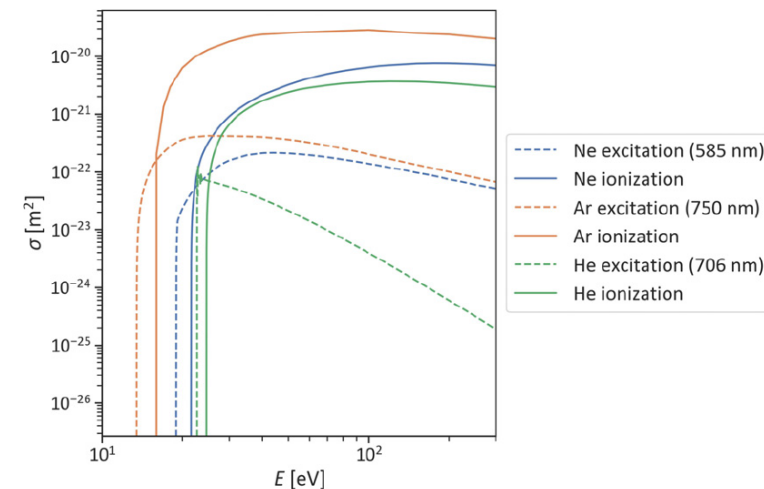
- Effect of photoemission becomes observable

## Townsend breakdown is a good approximation

- Provides global model to combine with photoemission quantum model



Cross sections from ground state into excited states [1]

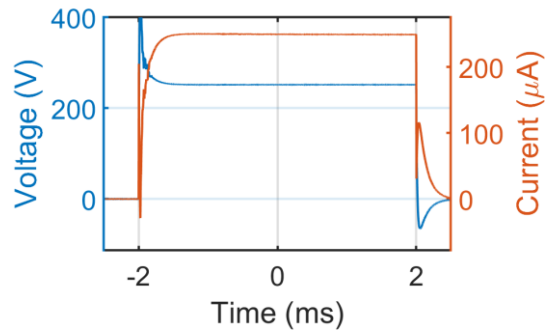




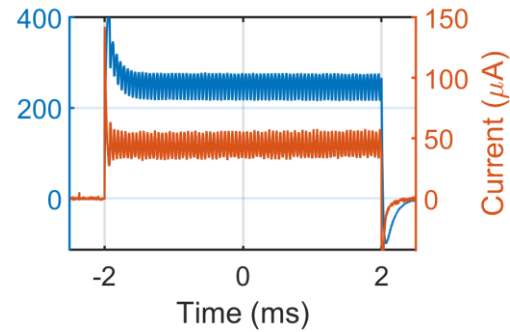
Decreasing Current

No Laser

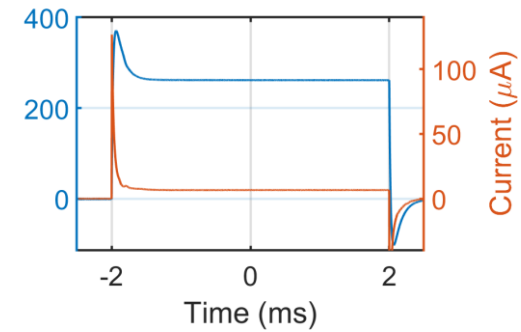
Constricted  
Stable



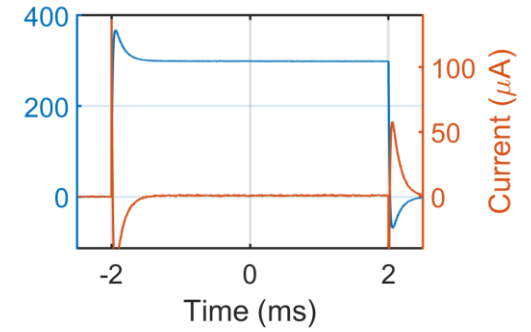
Constricted  
Oscillatory



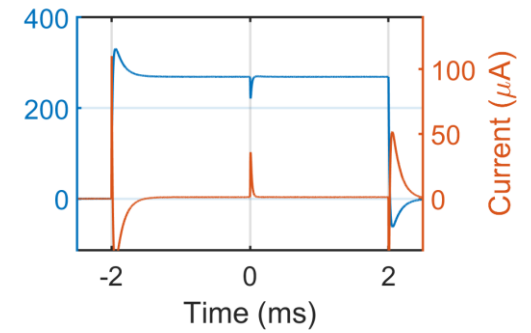
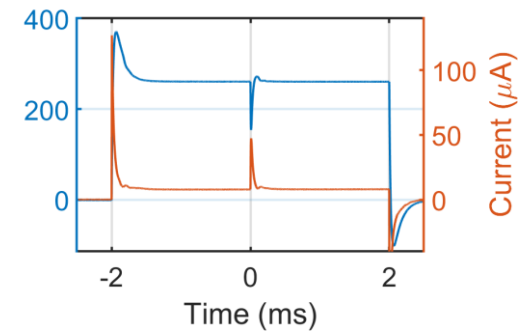
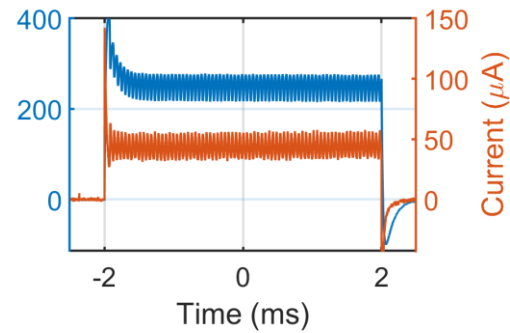
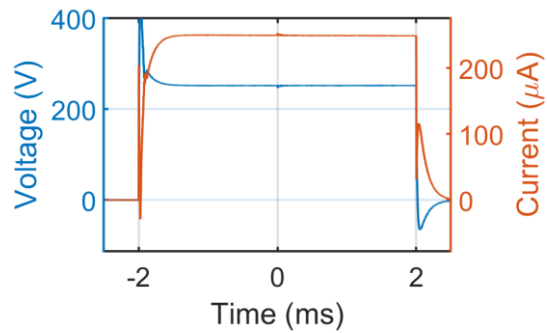
Diffuse



No  
Breakdown



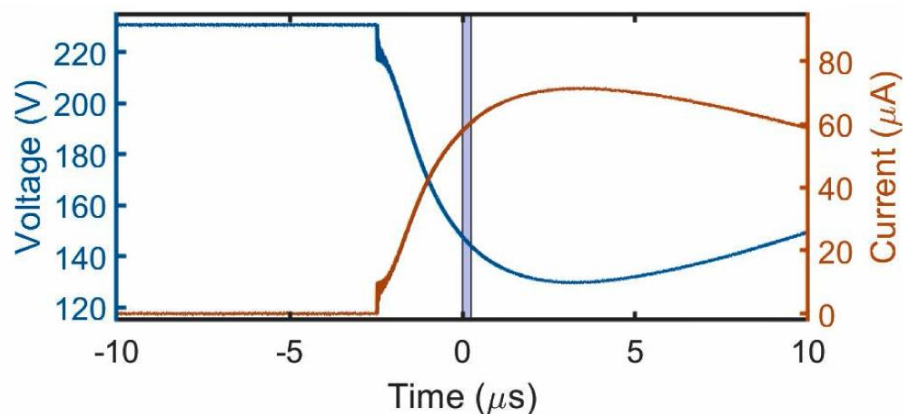
With Laser  
at  $t = 0$  ms



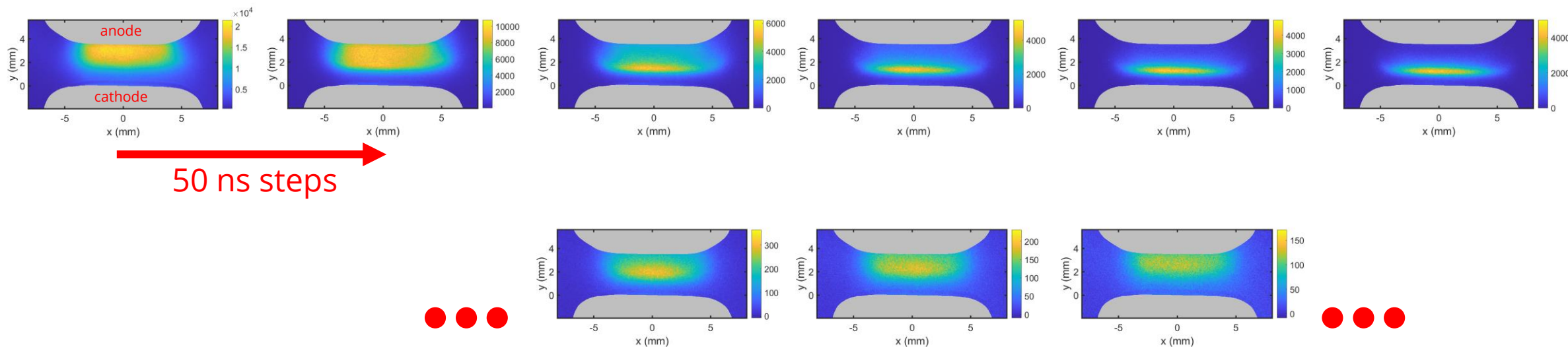
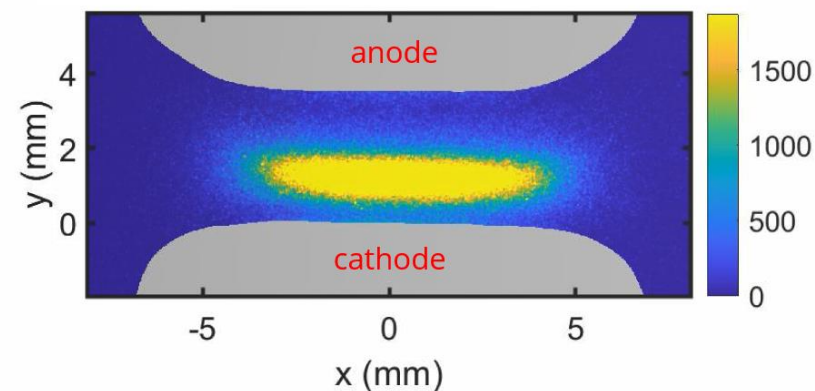
# Laser-Induced Breakdown at $pd = 1$ (Torr-cm)



Waveforms (with Laser)



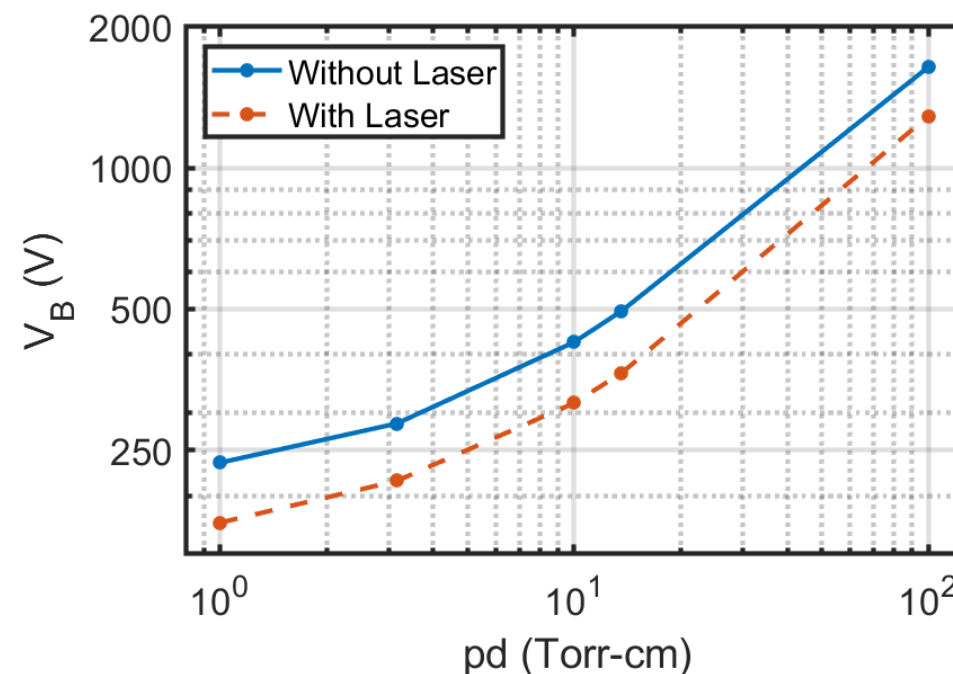
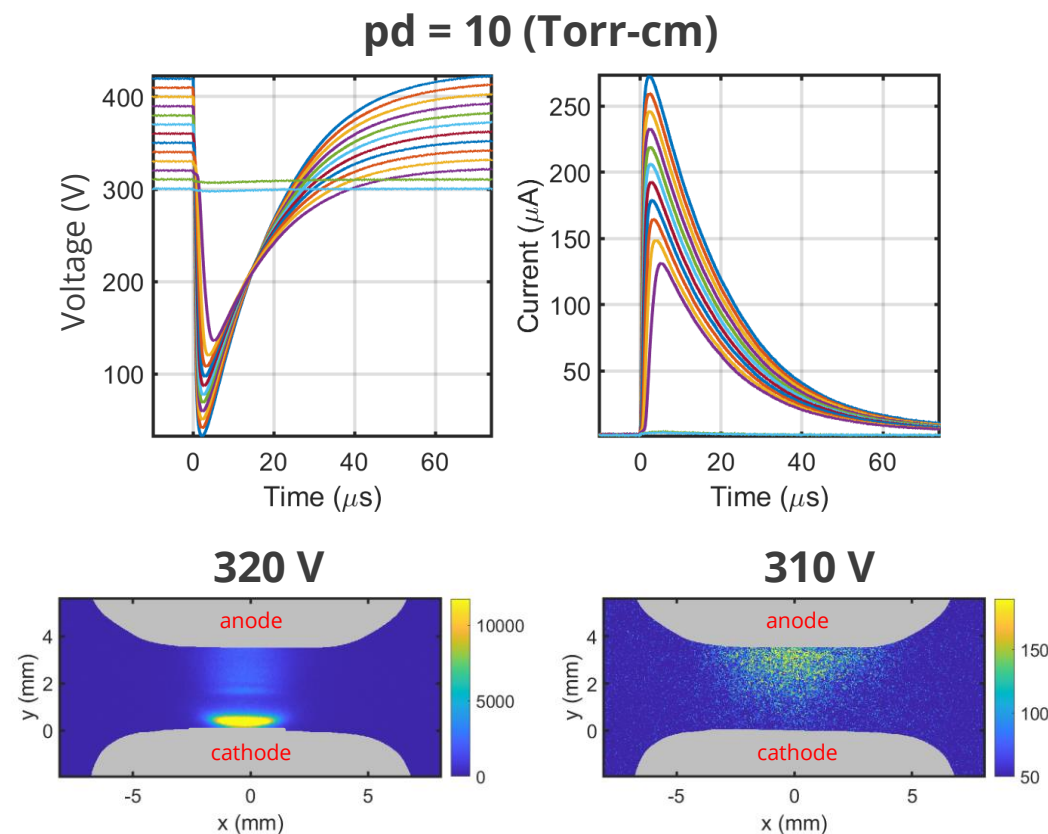
750 nm Emission (with Laser)





# Laser-Induced Breakdown Voltage

- Transition shown in waveform and image data
- Implies applied voltage threshold for laser-induced breakdown



- Laser-induced photoemission provides sufficient electrons to reduce breakdown voltage
- Plasma could be sustained at a **lower breakdown voltage** or E/N

# Modeling Photoemission Induced Plasma Breakdown



A. V. Phelps, Z. Lj. Petrović, and B. M. Jelenković [1] → Transient current ( $I$ ) and voltage ( $V$ ) model

- A pulse of photoelectrons released from the cathode (corresponding to a photoelectric current pulse,  $I_p$ )
- A pulse of voltage ( $V_0$ ) applied to the discharge circuit

## Assumptions [1]

- Small space-charge distortion of the electric field in the observed range of currents
- The times for significant changes in the electric field and current are long compared to the electron and ion transient times
- Electrons are produced at the cathode only by ions → contributions from photons and metastables are small

Effective yield of electrons per ions arriving at the cathode,  $\gamma = \gamma_p + k_V V + k_I I$

$\gamma_p$  → "potential ejection" of electrons

$k_V$  → "kinetic ejection" of electrons

$k_I$  → first-order effects of space charge on the electric field



## Assumptions [1] (continued)

- Ionization by electrons only → ionization by heavy particles neglected
- The electron multiplication is an exponential function of position with a spatial ionization coefficient ( $\alpha$ ) [2]

$$\alpha = Cpe^{-D\sqrt{\frac{p}{E}}} \left\{ \begin{array}{l} p \rightarrow \text{pressure (Torr)} \\ E \rightarrow \text{electric field (V cm}^{-1}\text{)} \\ C, D \rightarrow \text{fitting coefficients for the specific gas} \end{array} \right.$$

- Round-trip electron number gain ( $g$ ) resulting from an electron released from the cathode → a unique function of the gas density ( $n$ ), electrodes separation ( $d$ ), and discharge voltage ( $V$ )

$$g = \gamma[(1 + \delta)e^{\alpha d}]$$

- $\delta$  → yield of ions produced by backscattered electrons per electron arriving at the anode
- Electron and ion currents → uniformly distributed over the surface of the electrodes

[1] A. V. Phelps, Z. Lj. Petrović, and B. M. Jelenković, *Phys. Rev. E* **47**, 2825, 1993.

[2] Fu et al., *Plasma Sources Sci. Technol.* **27**, 095014, 2018.

# Discharge Characteristics and Circuit Equations



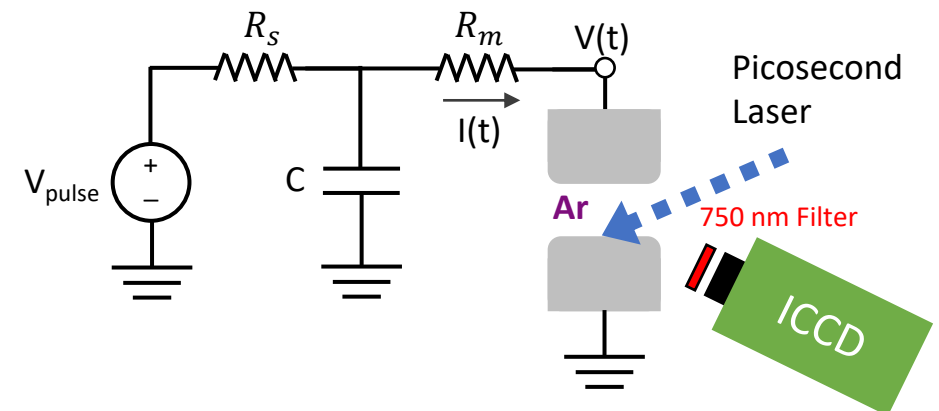
Discharge characteristics equation [1]:

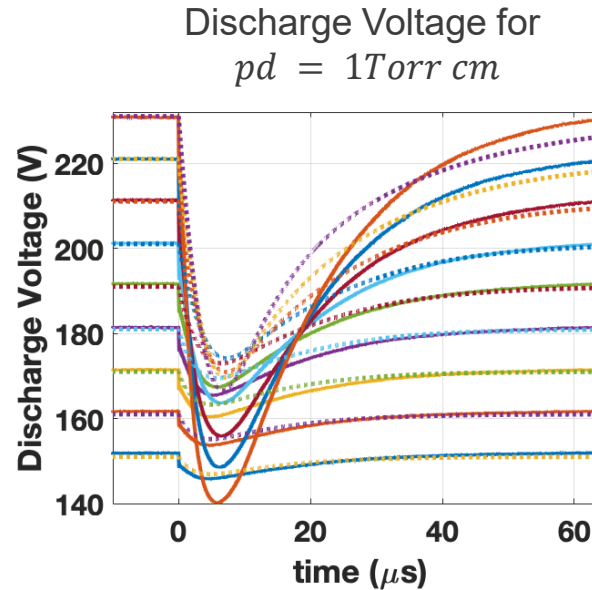
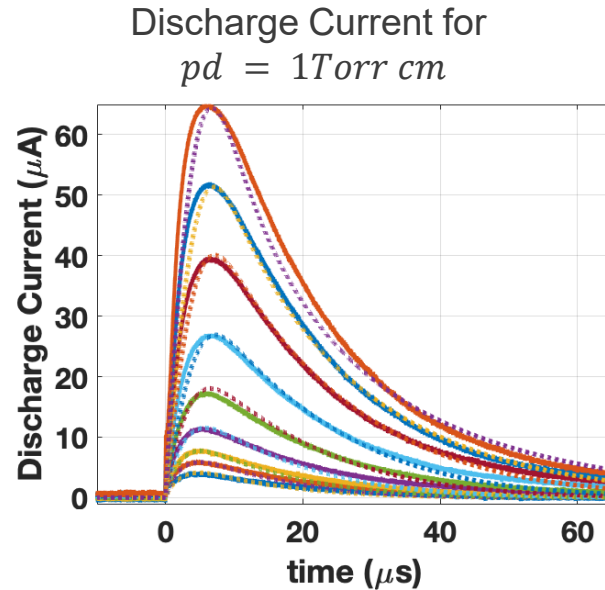
$$\frac{dI}{dt} = \left[ \frac{(1 + \gamma)I_p}{\gamma T} + \frac{(g - 1)I}{T} - \frac{Ik_V}{\gamma(1 + \gamma)} \frac{dV}{dt} \right] \times \left[ 1 + \frac{Ik_I}{\gamma(1 + \gamma)} \right]^{-1}$$

Ion transit time,  $T = d/v_d$  ;  $d \rightarrow$  Electrodes separation,  $v_d \rightarrow$  Ion drift velocity

Equation for the discharge voltage ( $V$ ) in terms of the discharge current ( $I$ ) and the circuit components [1]:

$$\frac{dV}{dt} = \frac{1}{R_s C} [V_0 - V - I(R_s + R_m)] - R_m \frac{dI}{dt}$$





**Dotted lines: Model**  
**Solid lines: Experiment**

Good match to experiment for currents  $< 10 \mu\text{A}$

Empirical fitting: **Similar observations were made by Petrović and Phelps [1].**

- Gaussian-shaped simulated photoemission pulse,  $I_p(t)$ , must have  $25 \mu\text{s}$  pulse width
- A small charge transfer ( $10^{-12} \text{C}$  to  $10^{-13} \text{C}$ ) due to photocurrent pulse is necessary to fit the observed current amplitude → **much smaller than the calculated emitted charge due to photoemission** from the cathode ( $2.1033 \times 10^{-10} \text{C}$ ).
- An improved discharge characteristics equation is needed to capture effects of photoemission

# Conclusions



Effects of laser-induced photoemission become observable at low currents/space charge

0D models provide an avenue to study surface effects like photoemission on breakdown

- Requires small space-charge distortion of the electric field in the observed range of currents

Higher order effects of photoemission are hard to capture, modifications to the discharge characteristics equation are needed

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