

APS DPP  
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# Characterizing charge state distributions in photoionized plasma to test high density effects in astrophysical code XSTAR

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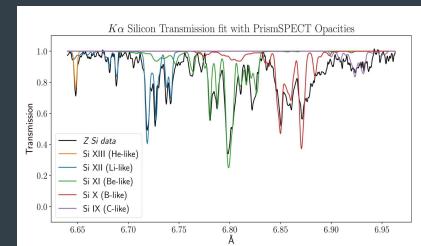
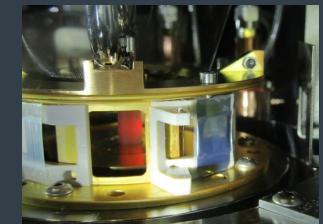
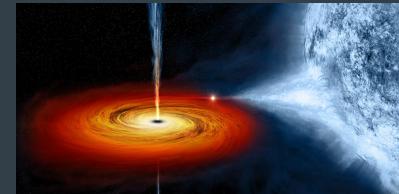
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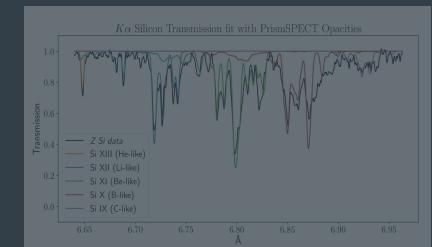
# Summary: Independent charge state distribution (CSD) analysis for Si agrees with astrophysical model predictions

1. There is limited scrutiny of CSDs predicted by astrophysical photoionized plasma models against laboratory data.
2. The expanding foil photoionized plasma experiment on Z can help interrogate the models.
3. We can use theoretical opacities to measure CSDs in a way that is not sensitive to plasma parameters.
4. Our analysis of Si data shows good agreement with XSTAR.
5. The Fe spectrum is much more complicated and work is ongoing.



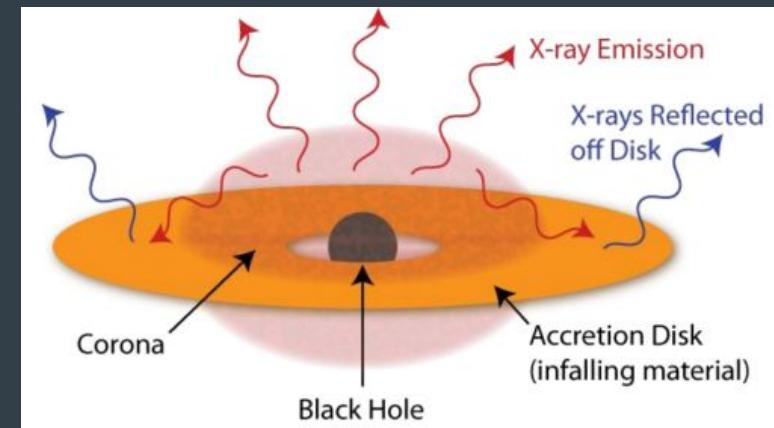
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# Fits to reflection spectra are used to determine properties of X-ray binaries and AGN

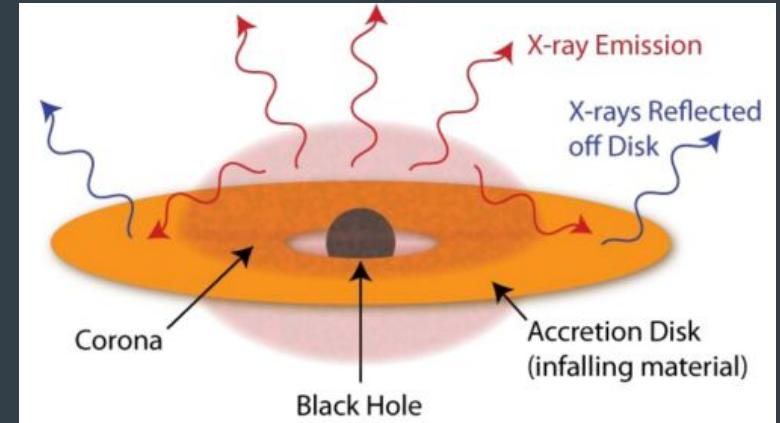
- Radiation emitted by the corona is “reflected” by the accretion disk.
- Fits to reflection spectra can be used to infer:
  - Composition
  - Density
  - BH spin
  - BH mass



**XSTAR is one widely used photoionized plasma model.**

# Model predictions for transmission spectra also need to be scrutinized

- Ion populations are established by photoionization.
- Ion populations are a key quantity predicted by photoionized plasma models.
- We see emission from these systems, but absorption is also occurring in the disk plasma.
- Both spectra are important to scrutinize.



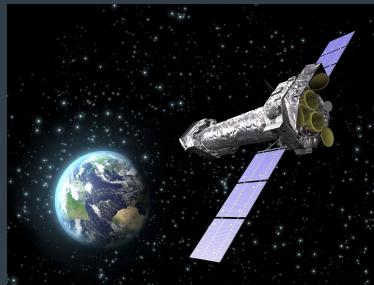
# The x-ray astrophysical community has called for independent measurements of CSDs



Chandra  
- NASA

*“The two primary needs are precise and accurate determination of the wavelengths for common transitions as well as the [CSDs] of [L-shell ions of Mg, Si, Fe...] under a range of thermodynamic and radiative conditions.”*

XMM-Newton  
- ESA

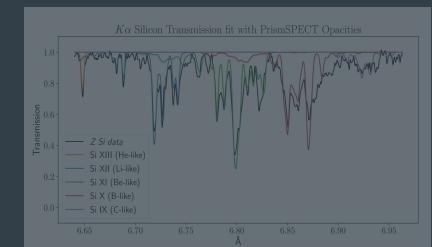
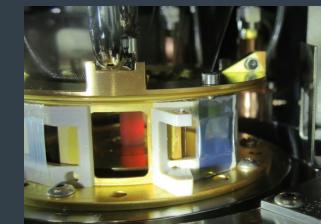
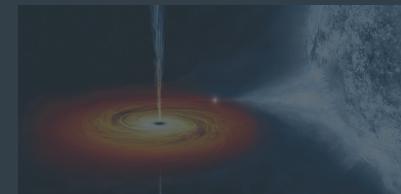


- Smith et al., Laboratory Astrophysics Needs for X-ray Grating Spectrometers, Astro 2020 Decadal Review

**Laboratory photoionized plasma data is sparse!**

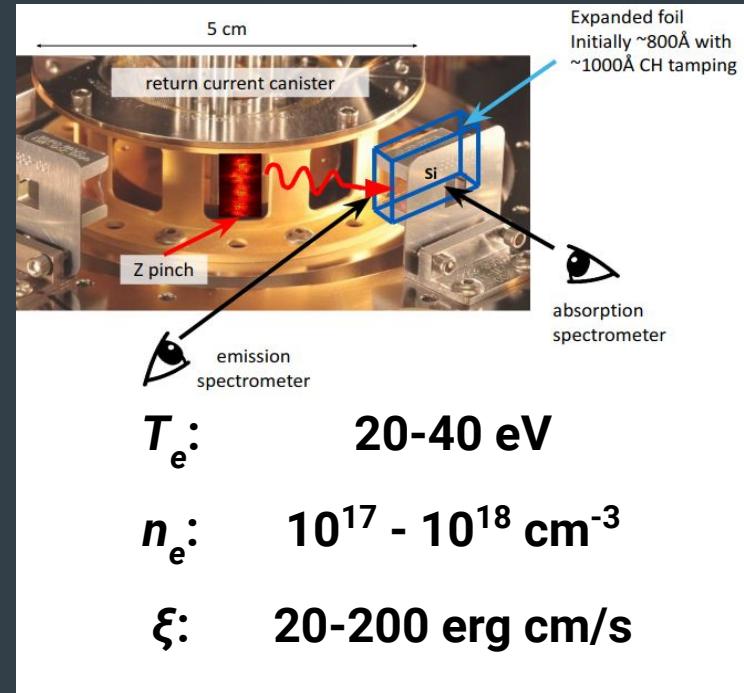
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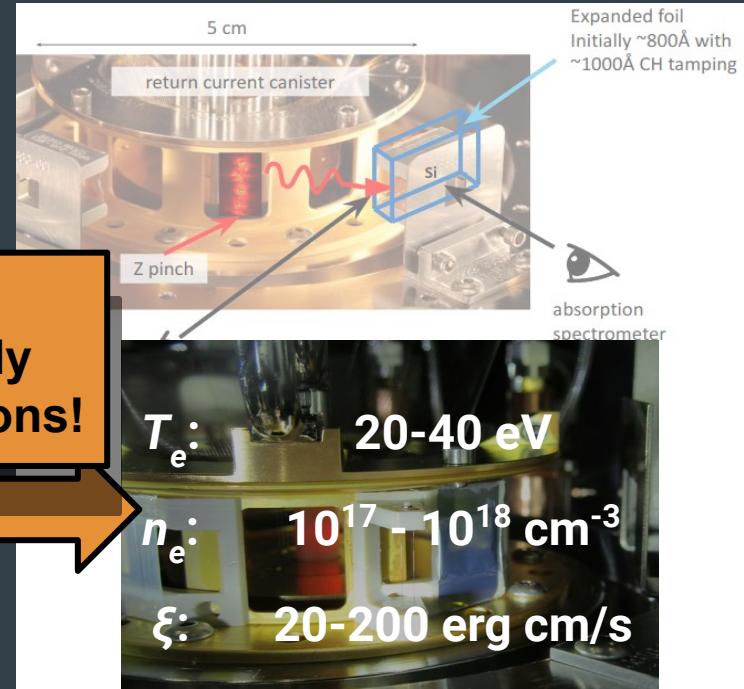
# The foil photoionized plasma experiment on Z reaches regimes relevant to accreting systems

- A Si or Fe foil sample irradiated by the Z-pinch.
- We observe the plasma in both emission and absorption.



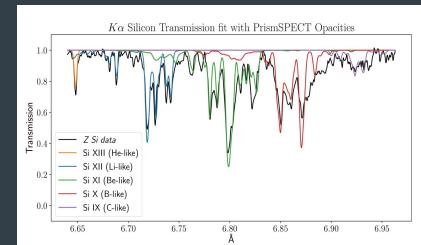
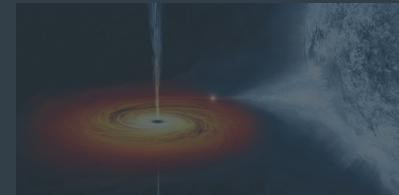
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# We extract areal densities from a transmission spectrum by fitting opacities

- $X_\nu$  [cm<sup>2</sup>] is an opacity “cross-section”: it is \*almost\* independent of plasma parameters except for the line shape  $\phi_\nu$ .
- PrismSPECT + ATBASE to get  $X_\nu$  per ion.
- The product of areal density and opacity cross section across all wavelengths gives a transmission spectrum.
- We can fit areal densities for several ions to a transmission spectrum, giving a CSD measurement.

$$T_\nu = \exp \left[ - \sum_i X_{\nu,i}^{\text{prism}} \times \frac{n l_i}{\text{cm}} \right]$$

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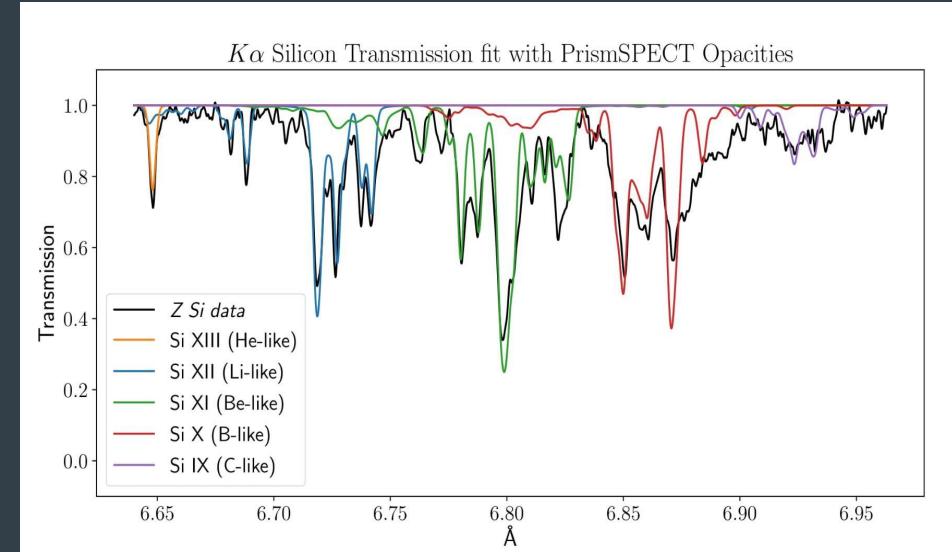
Steps:

- Get  $X_{\nu,i}$ .
- Convert to transmission.
- Convolve with instrumental resolution.
- Levenberg-Marquardt minimization to get best fit  $nl_i$ .

$$T_\nu = \exp \left[ - \sum_i X_{\nu,i}^{\text{prism}} \times \underline{nl}_i \right]$$

# Results with Si show the method is effective

- Spectrum from Loisel et al. (2017) was obtained by averaging several Si transmission spectra collected across multiple shots.
- Fits to the ion areal densities return a total areal density that agrees with independent Rutherford back-scattering measurements.

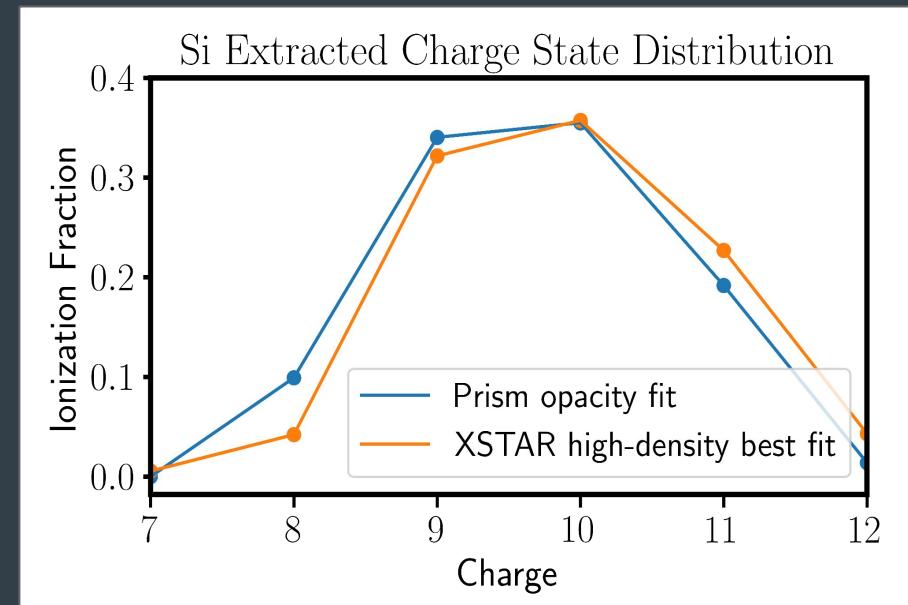


**RBS Areal Density:  $3.10 \times 10^{17} \text{ Si/cm}^2 \pm 5\%$**   
**Fitted Areal Density:  $2.97 \times 10^{17} \text{ Si/cm}^2$**

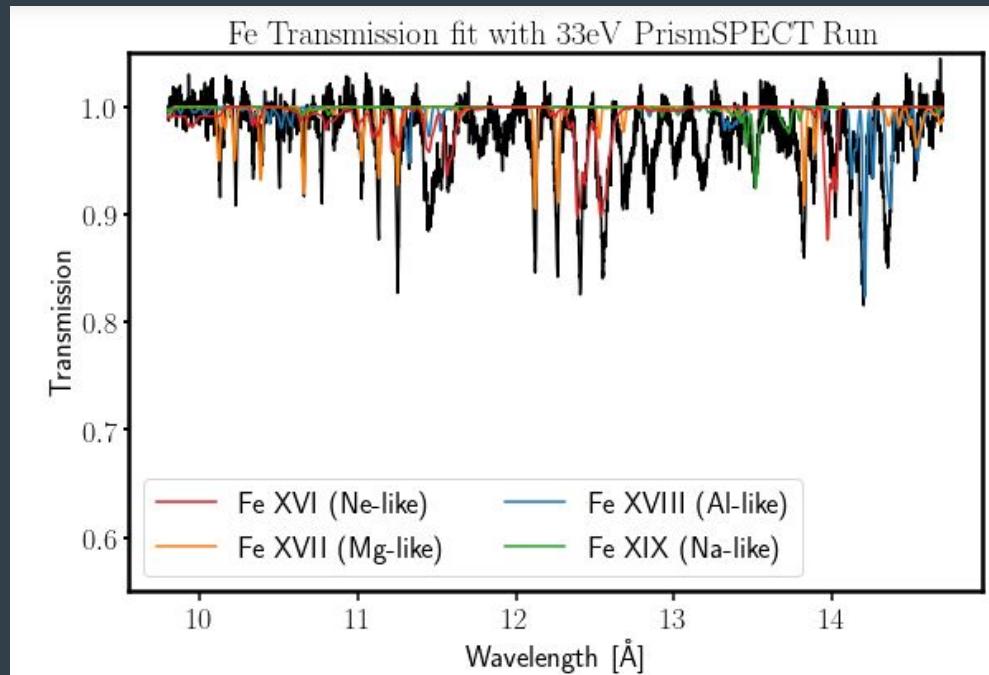
# XSTAR predicts similar CSD to our measurements

Next steps:

- Robust error analysis
- Investigate  $T_e$  differences



# CSD analysis for Fe is in progress



## Next Steps:

- Apply wavelength-dependent instrumental resolution
- Re-evaluate continuum in absorption spectrum

# Thanks!

# More on opacity cross sections and fitting routine

$$\frac{\kappa_\nu}{\rho} = \frac{\pi e^2}{mc} \frac{N_A}{M} f_l f_{l \rightarrow u} \phi_\nu \rightarrow X_\nu \equiv \frac{M}{N_A f_l} \left( \frac{\kappa_\nu}{\rho} \right) = \frac{\pi e^2}{mc} f_{l \rightarrow u} \phi_\nu$$

$$T_\nu = \exp \left[ - \sum_i X_{\nu,i}^{\text{prism}} \times \underline{n l_i} \right] \quad \frac{\lambda}{\delta\lambda} \sim 2400$$

Instrumental resolution  
for Si data