



A hybrid mesh-based/element-free method using fine-scale triangulations for the solution of PDEs on geometrically complex domains without defeaturing

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Outline

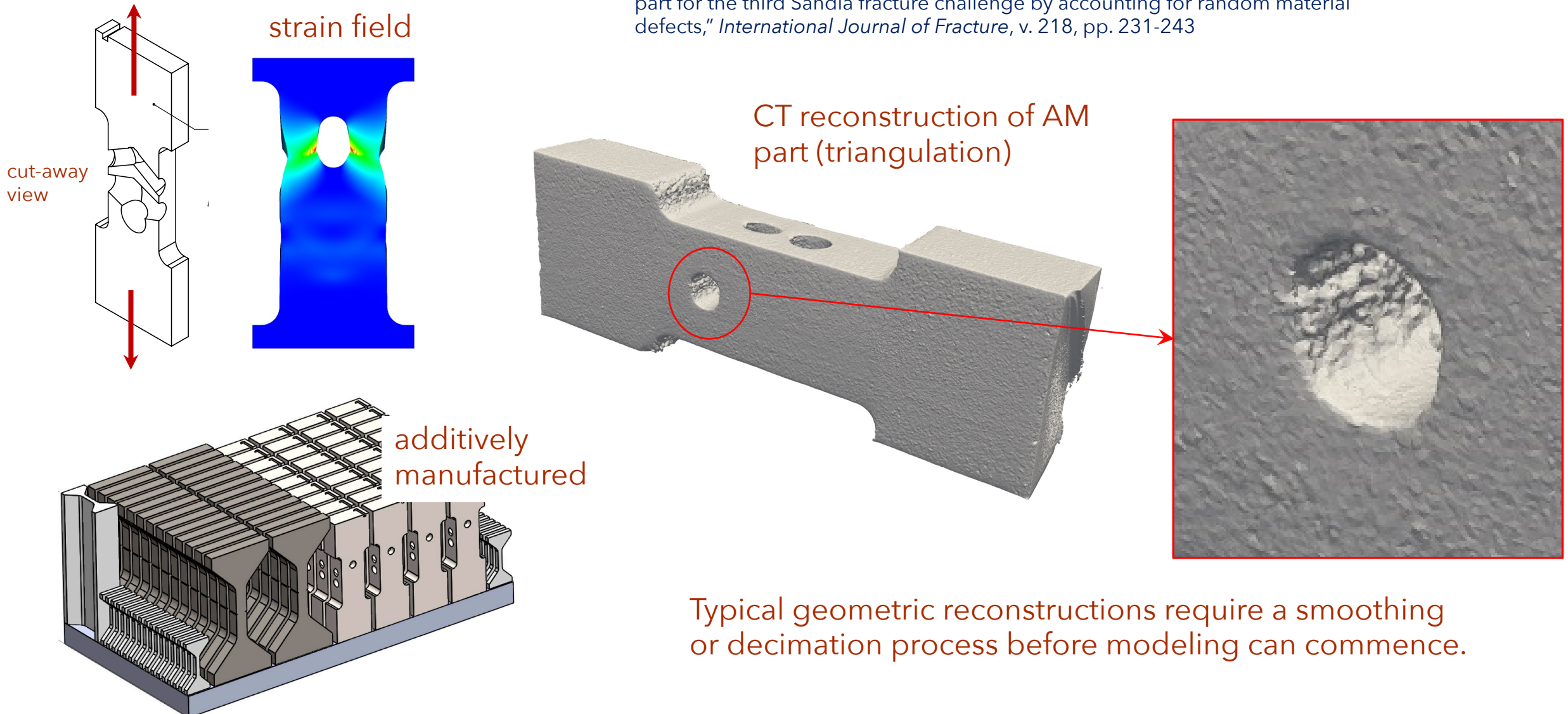


1. motivation
2. hybrid mesh/element-free discretizations
3. weight functions using manifold geodesics
4. shape functions using moving least squares
5. quadrature via secondary basis functions
6. strain projections for polynomial consistency
7. verification example, linear elasticity
8. summary

Motivation: Image-based analysis

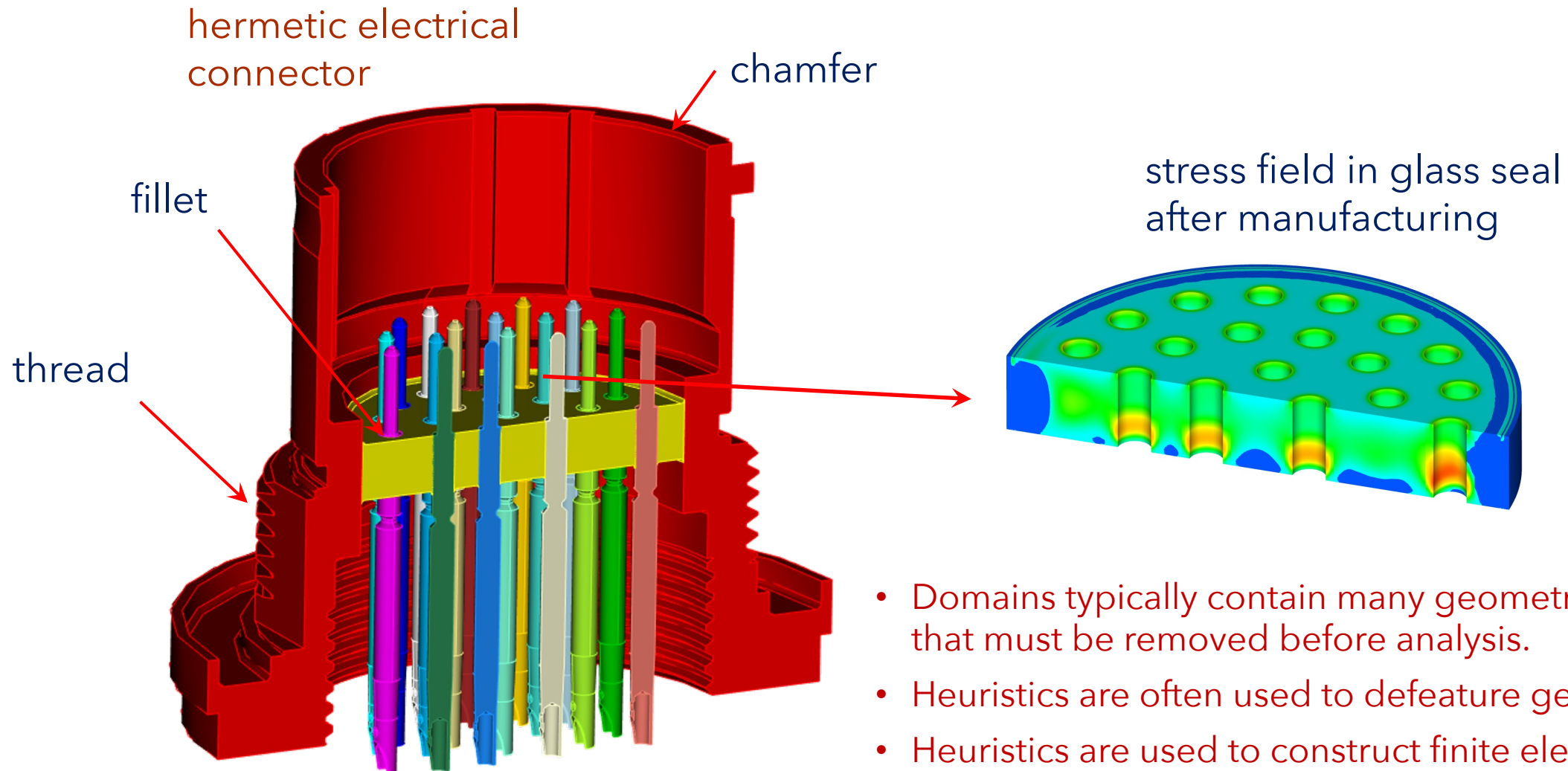


Johnson, et al., 2019, "Predicting the reliability of an additively-manufactured metal part for the third Sandia fracture challenge by accounting for random material defects," *International Journal of Fracture*, v. 218, pp. 231-243

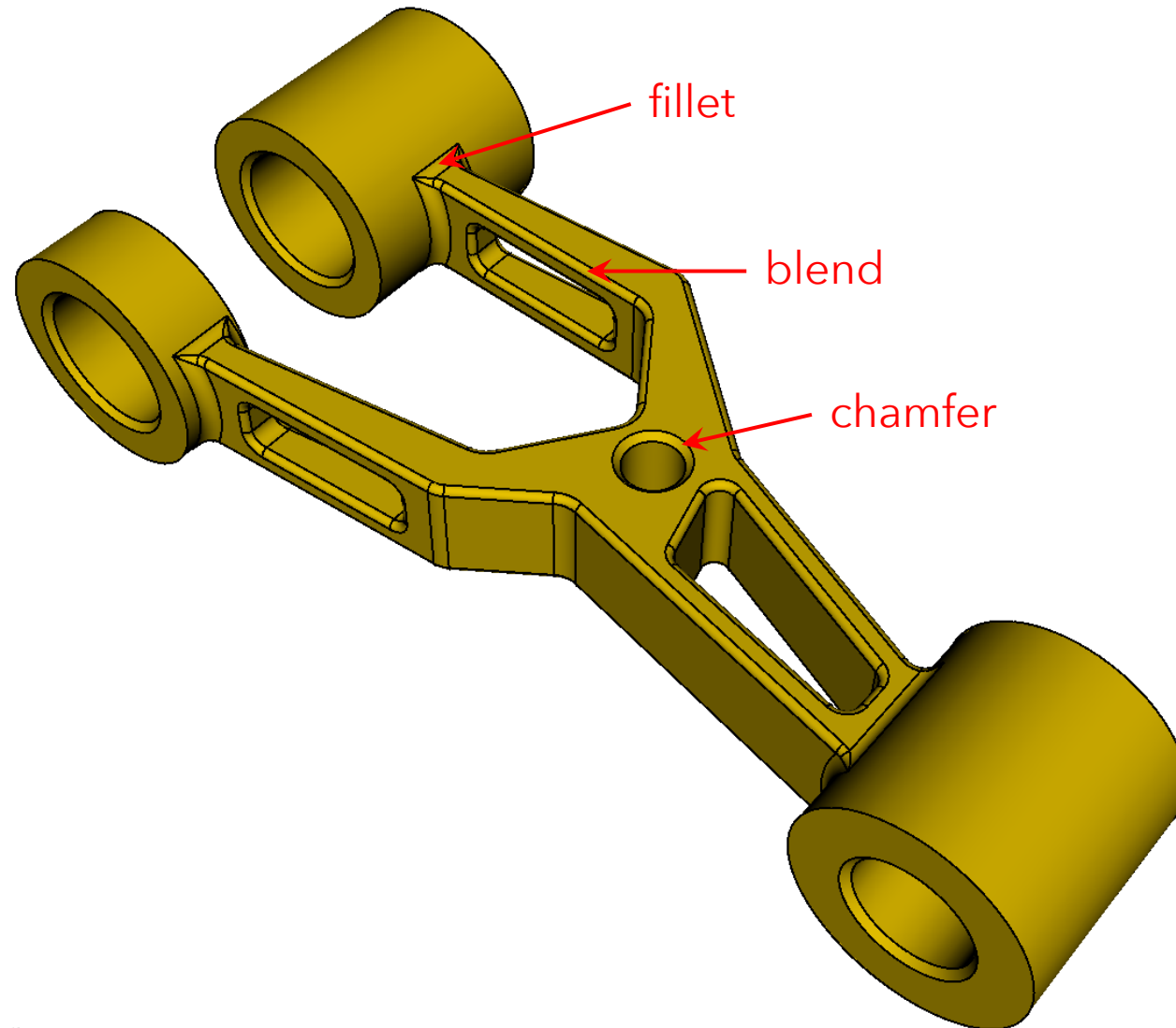


Typical geometric reconstructions require a smoothing or decimation process before modeling can commence.

Motivation: Agile simulation of complex assemblies



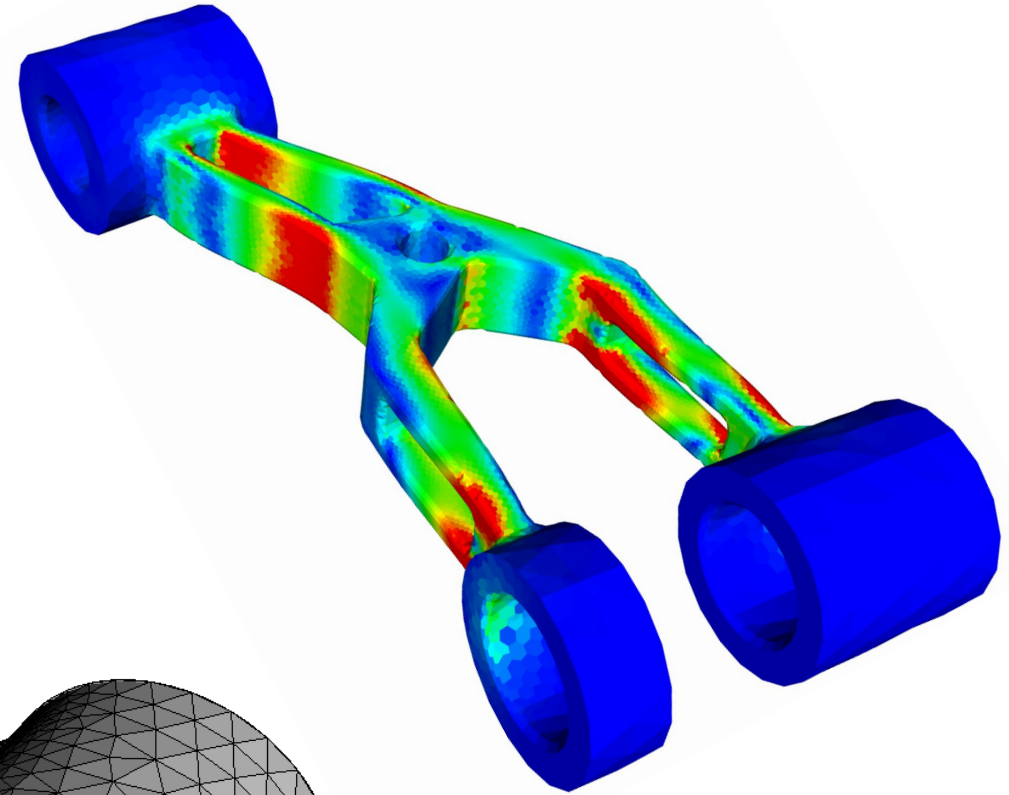
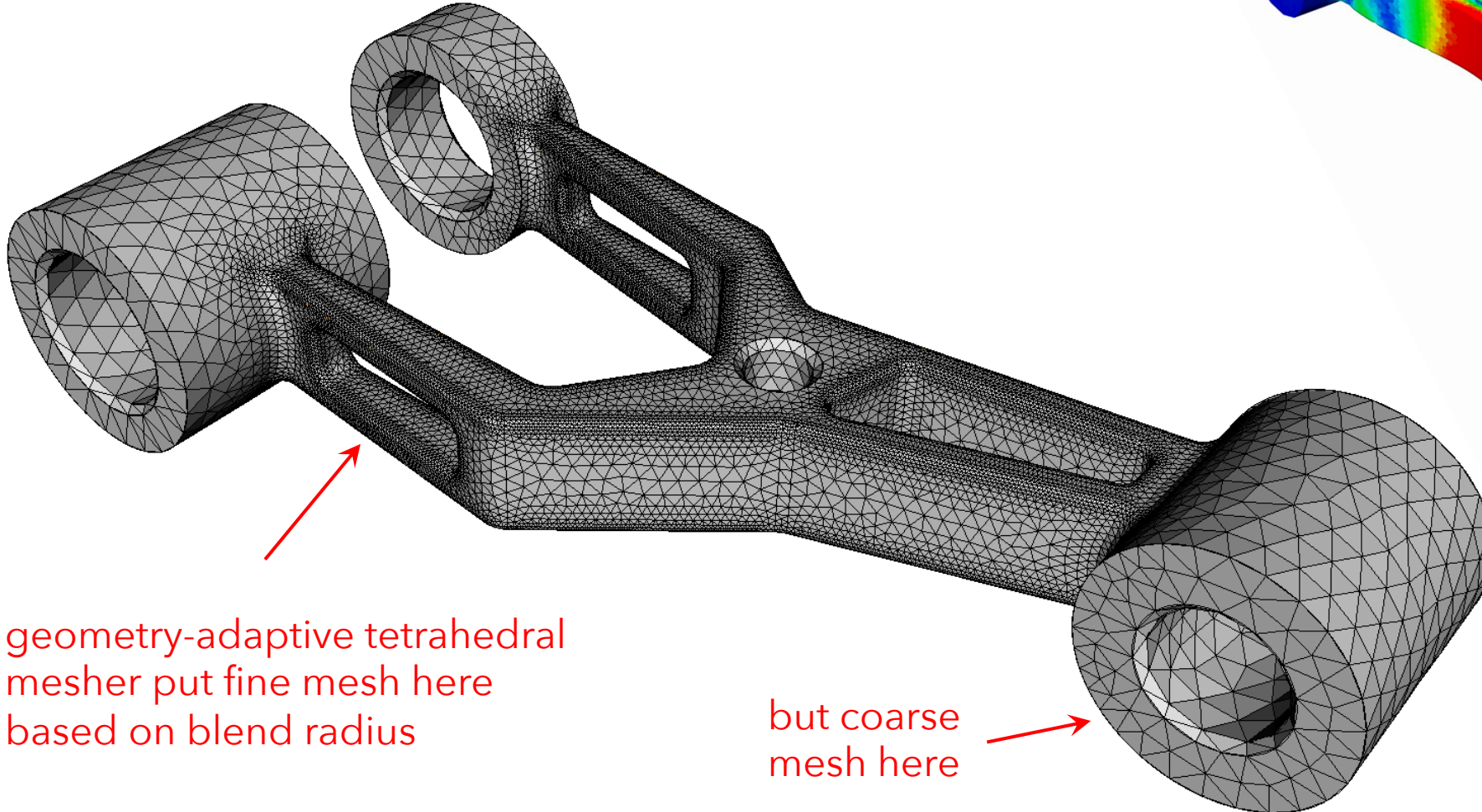
Motivation: Agile simulation



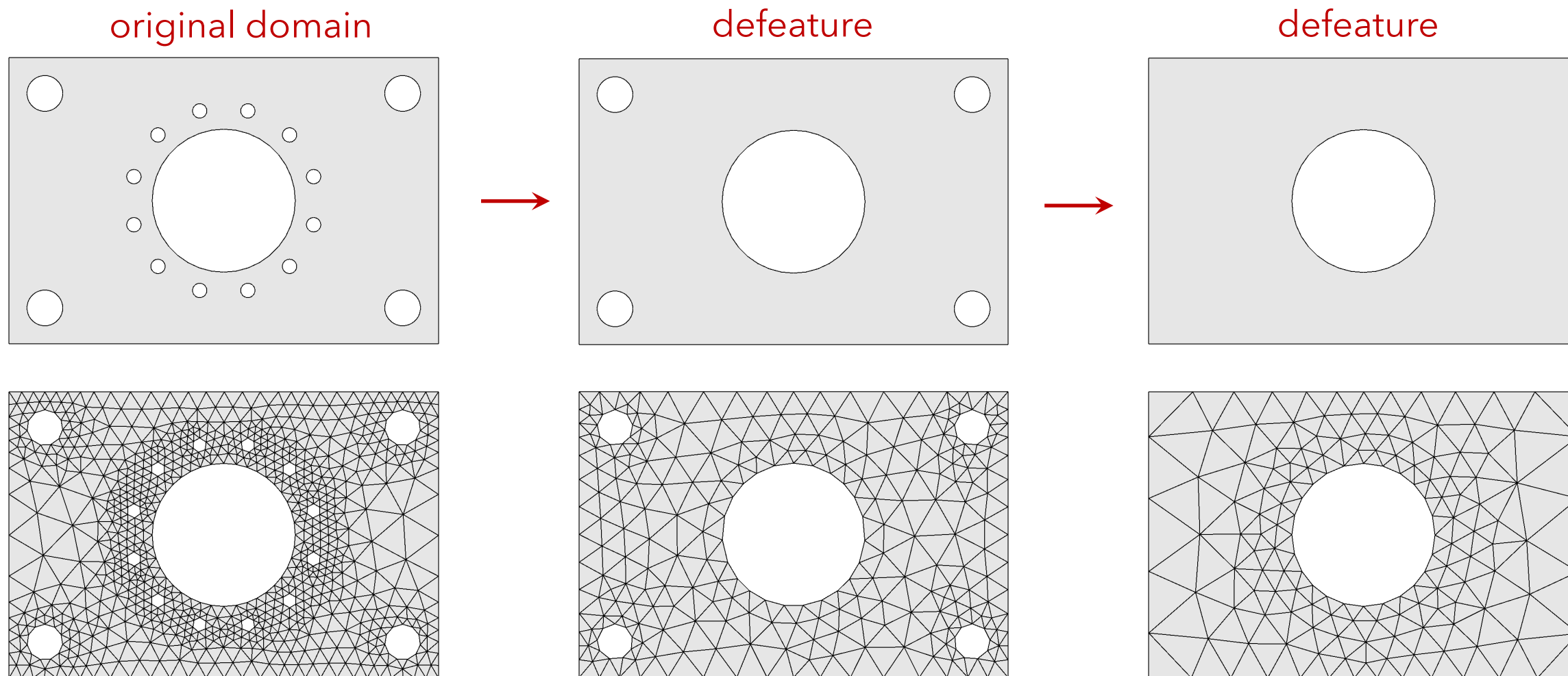
Motivation: Agile simulation



FEA discretization is intimately tied to domain geometry.



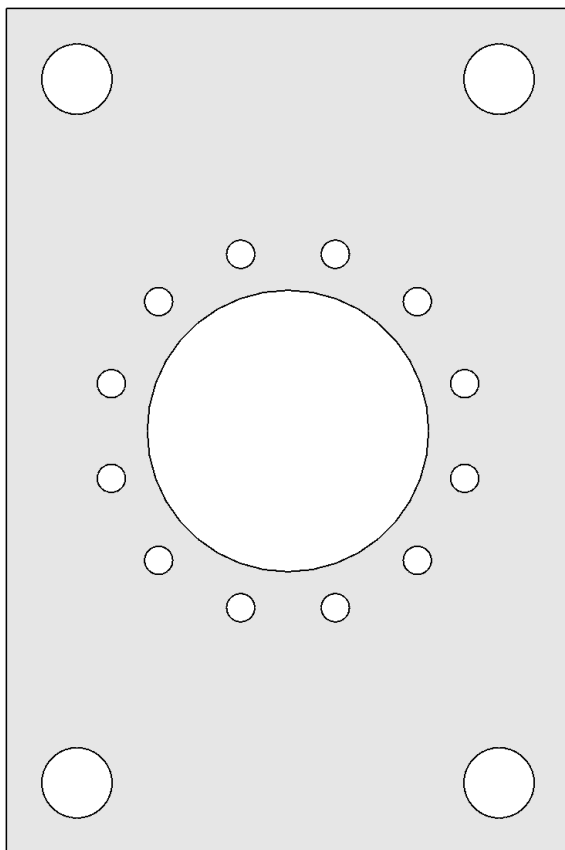
Motivation: Separate domain discretization from solution discretization



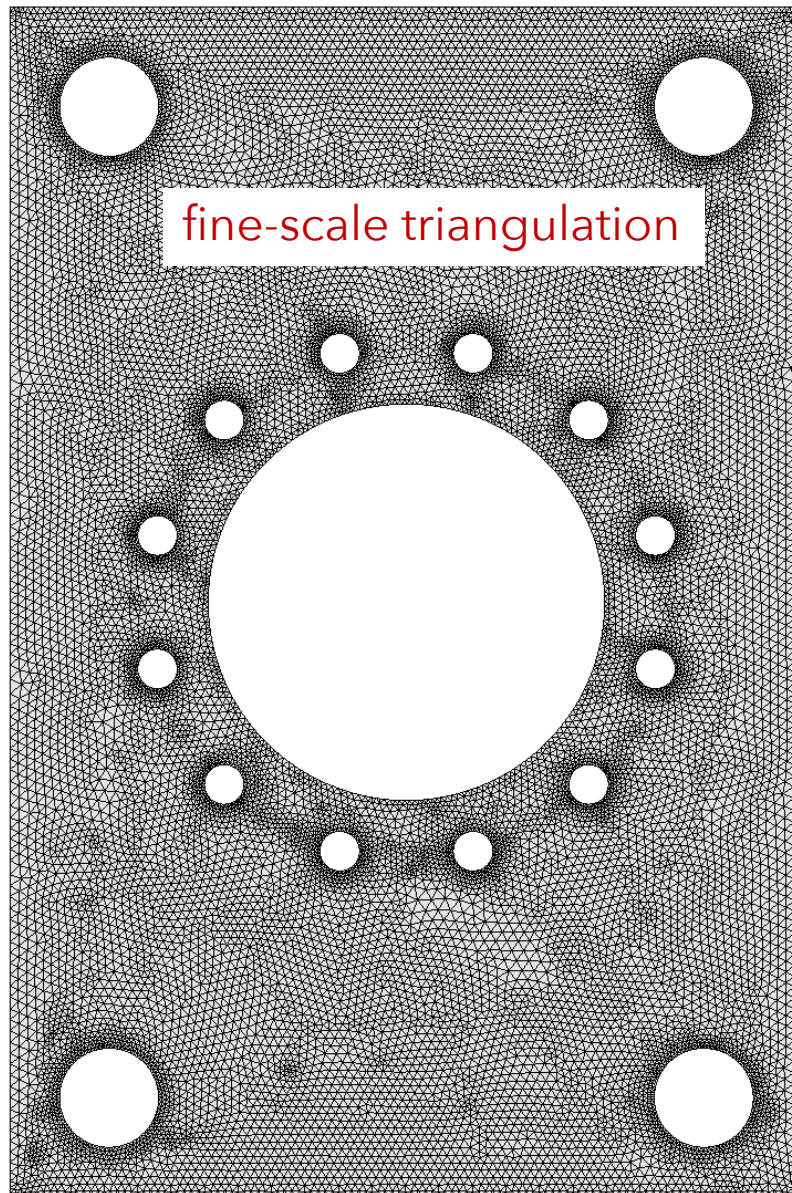
Hybrid approach: fine-scale triangulation



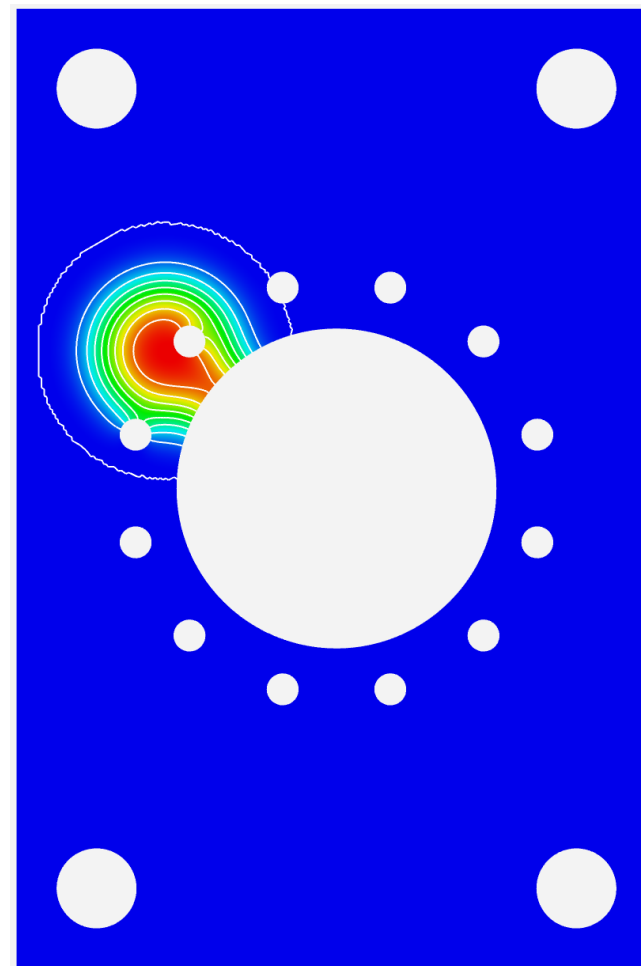
original domain



fine-scale triangulation



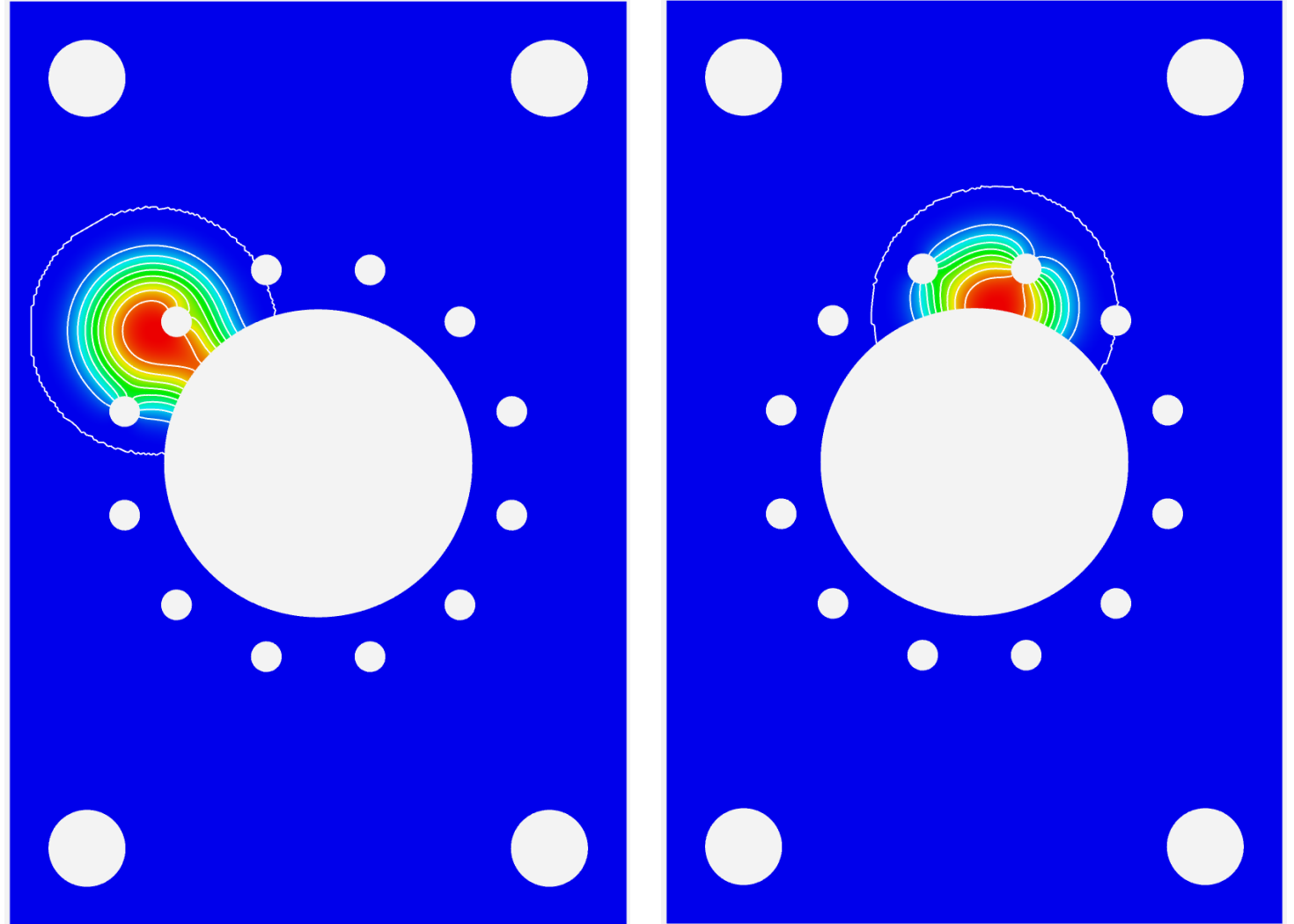
create an element-free basis using triangulation.



Element-free basis functions



- Element-free basis functions automatically include geometric features at all scales.
- Solution discretization is separate from domain discretization.
- No need to defeature domain.
- PDE solution is insensitive to quality of fine-scale triangulation.



Hybrid approach: fine-scale triangulation



Tet-meshing methods

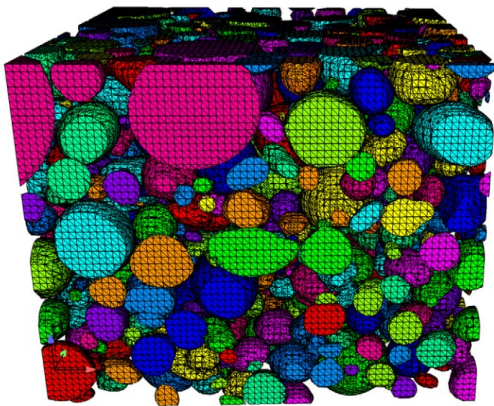
- Delaunay
- advancing front
- background grid
- envelope

TetWild

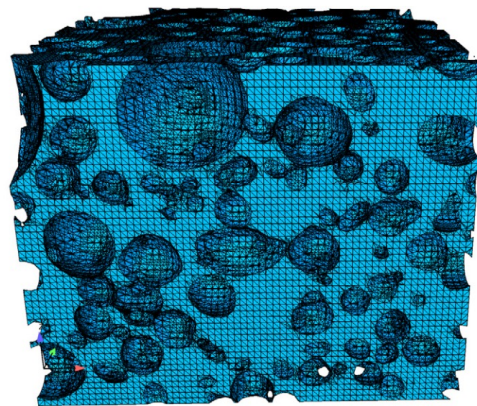
CDFEM

A verified conformal decomposition finite element method for implicit, many-material geometries

Scott A. Roberts*, Hector Mendoza, Victor E. Brunini, David R. Noble



(a) Particles



(b) Electrolyte

Fast Tetrahedral Meshing in the Wild

ACM Trans. Graph. 2020
vol. 39 Issue 4 Article 117

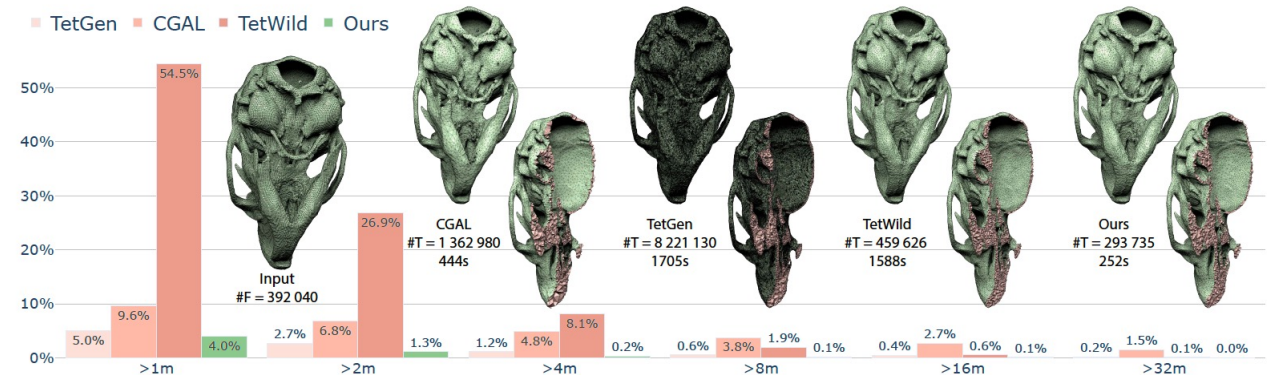
YIXIN HU, New York University, USA

TESEO SCHNEIDER, New York University, USA

BOLUN WANG, Beihang University, China and New York University, USA

DENIS ZORIN, New York University, USA

DANIELE PANOZZO, New York University, USA



Moving Least Squares (Reproducing Kernel)



The MLS shape functions $\phi_I(\mathbf{X})$ are defined as a spatial modulation of the nodal weight functions.

$$\phi_I(\mathbf{X}) = c_I(\mathbf{X})w_I(\mathbf{X})$$

where the modulation function $c_I(\mathbf{X})$ is found through a least square minimization process resulting in

$$c_I(\mathbf{X}) = \mathbf{g}^T(\mathbf{X})\mathbf{A}^{-1}(\mathbf{X})\mathbf{g}(\mathbf{X}_I)$$

where

$$\mathbf{A}(\mathbf{X}) = \sum_{I \in \mathcal{N}} w_I(\mathbf{X})\mathbf{g}(\mathbf{X}_I)\mathbf{g}^T(\mathbf{X}_I) \quad (\text{sum over neighbors})$$

$$\mathbf{g}^T(\mathbf{X}) = \{1 \ X_1 \ X_2\} \quad (\text{linear reproducibility})$$

$$\begin{aligned} \sum_I \phi_I(\mathbf{X}) &= 1 \\ \sum_I \mathbf{X}_I \phi_I(\mathbf{X}) &= \mathbf{X} \end{aligned}$$

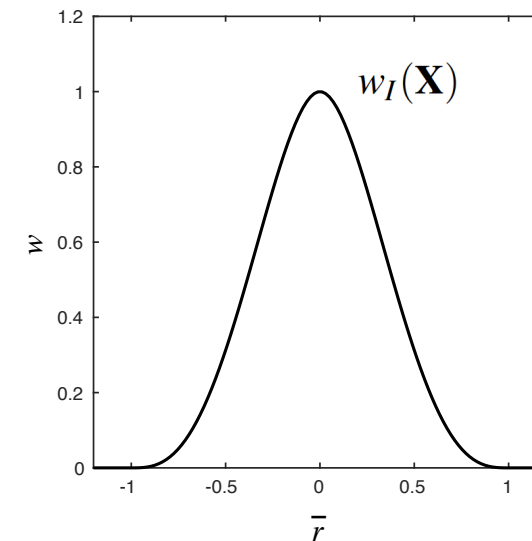
INTERNATIONAL JOURNAL FOR NUMERICAL METHODS IN ENGINEERING, VOL. 37, 229–256 (1994)

ELEMENT-FREE GALERKIN METHODS

T. BELYTSCHKO, Y. Y. LU AND L. GU

*Department of Civil Engineering, Robert R. McCormick School of Engineering and Applied Science,
The Technological Institute, Northwestern University, Evanston IL 60208-3109, U.S.A.*

nodal weight function

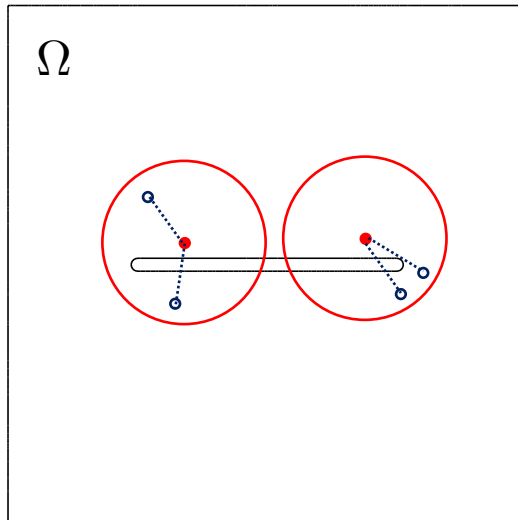


circular or rectangular support

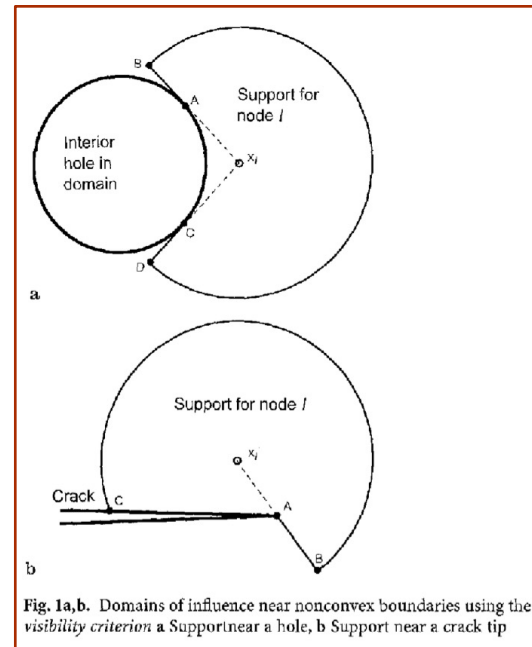
Continuous meshless approximations for nonconvex bodies by diffraction and transparency

D. Organ, M. Fleming, T. Terry, T. Belytschko

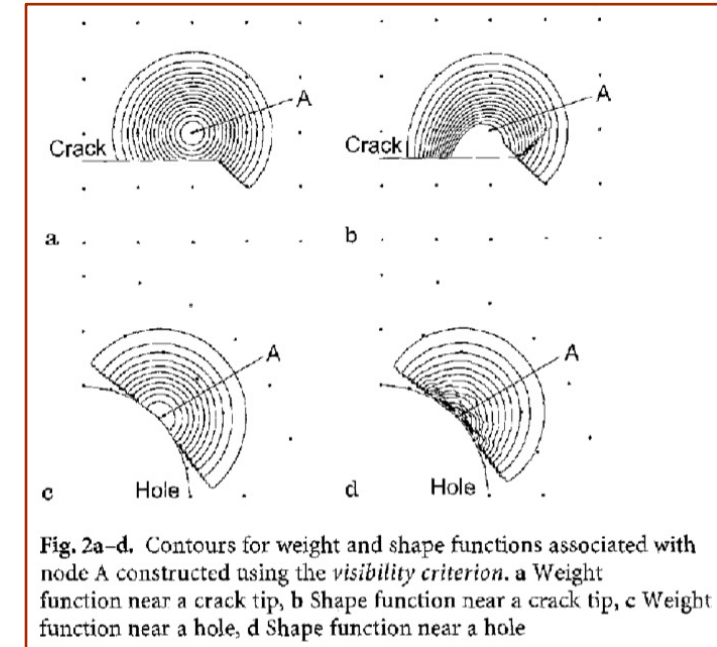
domain
with slot



visibility criterion



visibility criterion

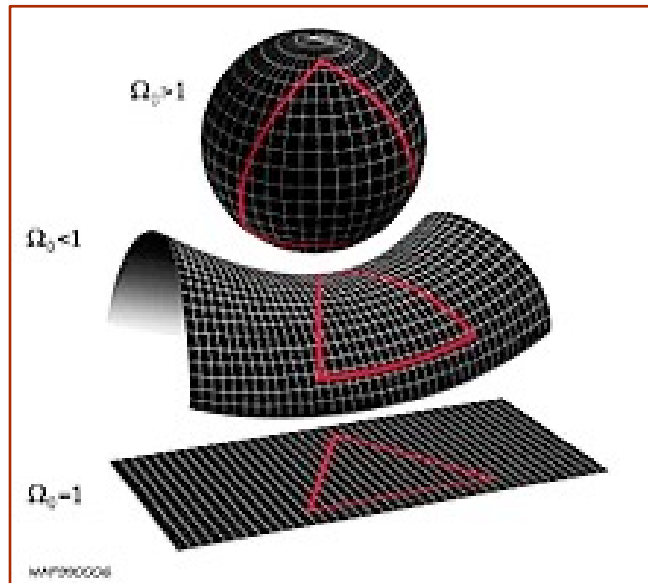


All these methods (visibility, transparency, diffraction) require use of computational geometry.

Manifold geodesic

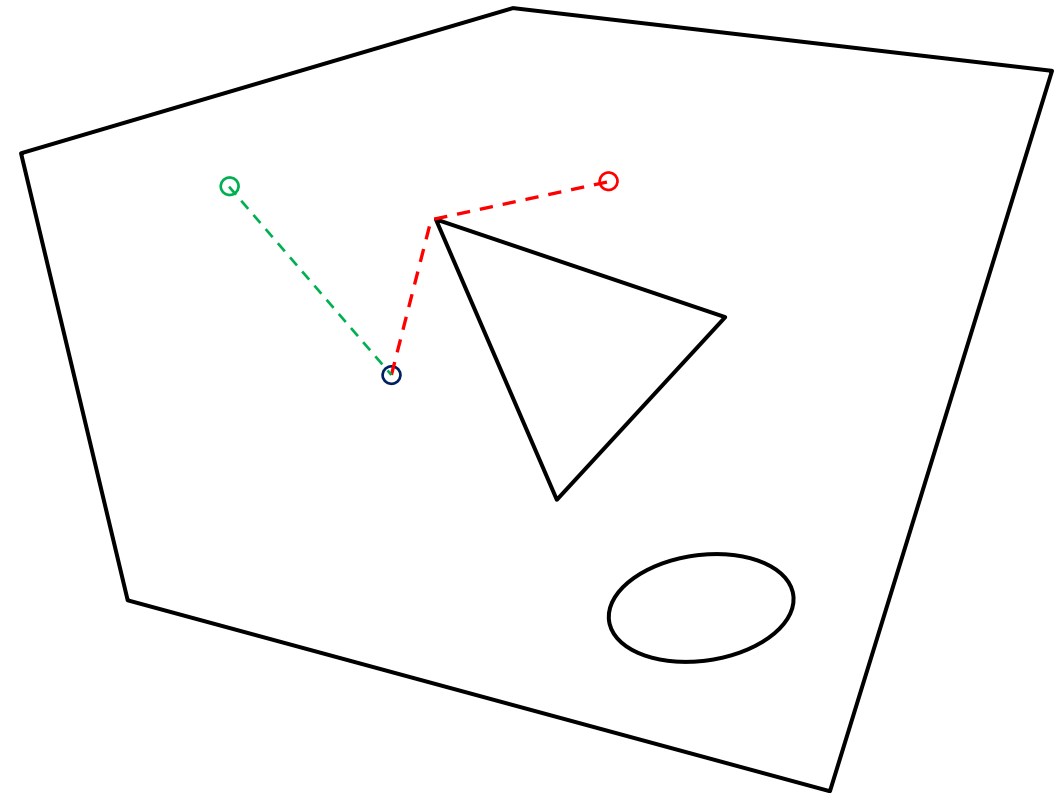


Geodesic: path that provides the shortest distance along a manifold



<https://en.wikipedia.org/wiki/Geodesic>

Euclidean manifold with boundary



Geodesics in Heat: A New Approach to Computing Distance Based on Heat Flow

KEENAN CRANE

Caltech

and

CLARISSE WEISCHEDEL and MAX WARDETZKY,

University of Göttingen

ACM Trans. Graph. 2013 Vol. 32 Issue 5 Pages Article 152

ALGORITHM 1: The Heat Method

- I. Integrate the heat flow $\dot{u} = \Delta u$ for some fixed time t .
 - II. Evaluate the vector field $X = -\nabla u / |\nabla u|$.
 - III. Solve the Poisson equation $\Delta \phi = \nabla \cdot X$.
-

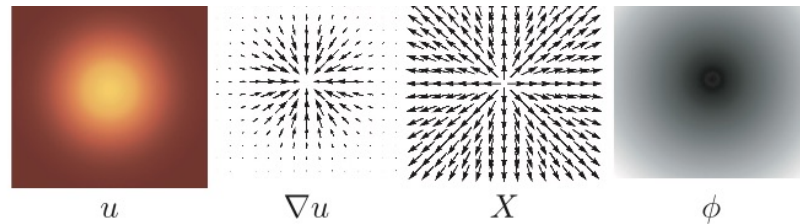


Fig. 5. Outline of the heat method. (I) Heat u is allowed to diffuse for a brief period of time (*left*). (II) The temperature gradient ∇u (*center left*) is normalized and negated to get a unit vector field X (*center right*) pointing along geodesics. (III) A function ϕ whose gradient follows X recovers the final distance (*right*).

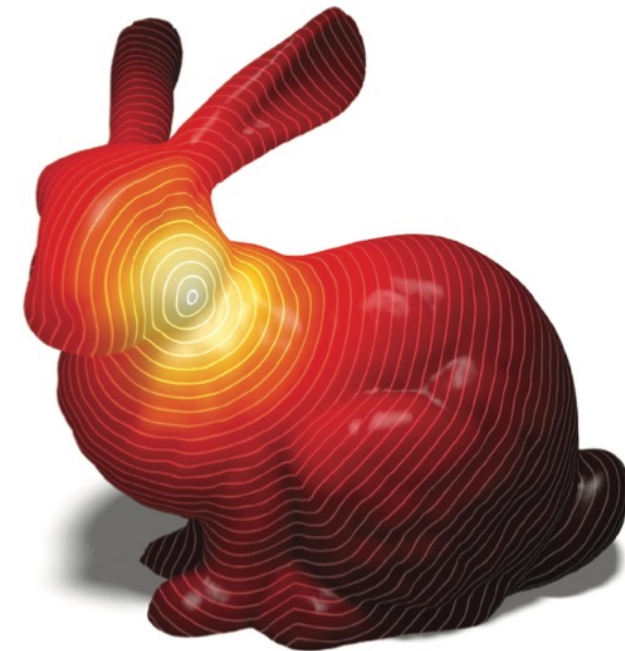
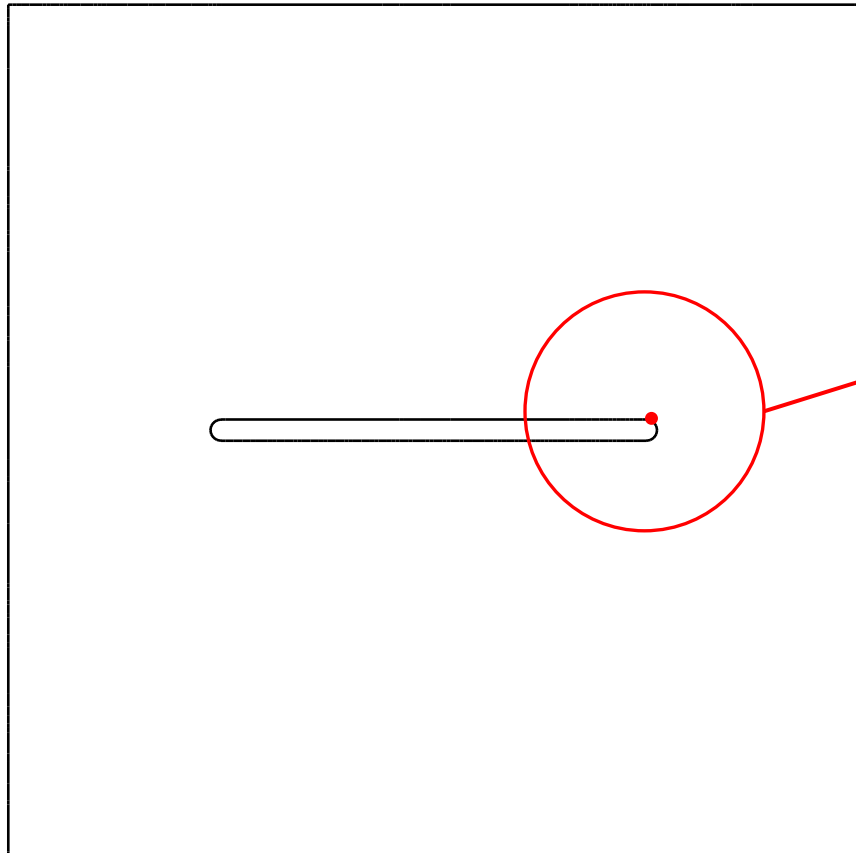
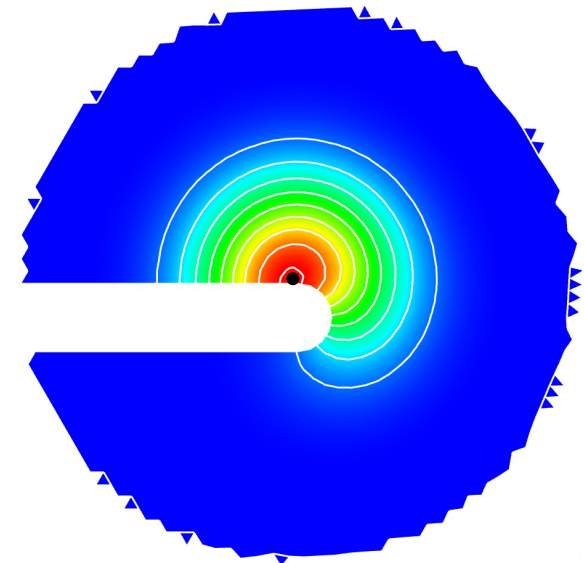
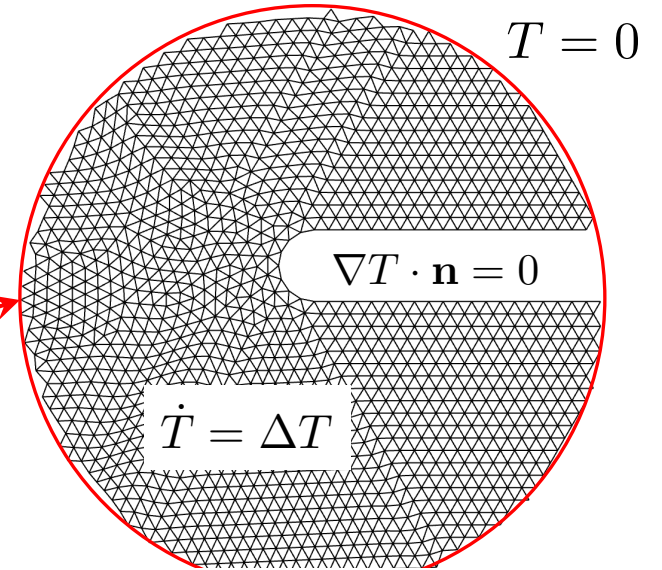


Fig. 1. Geodesic distance from a single point on a surface. The heat method allows distance to be rapidly updated for new source points or curves.

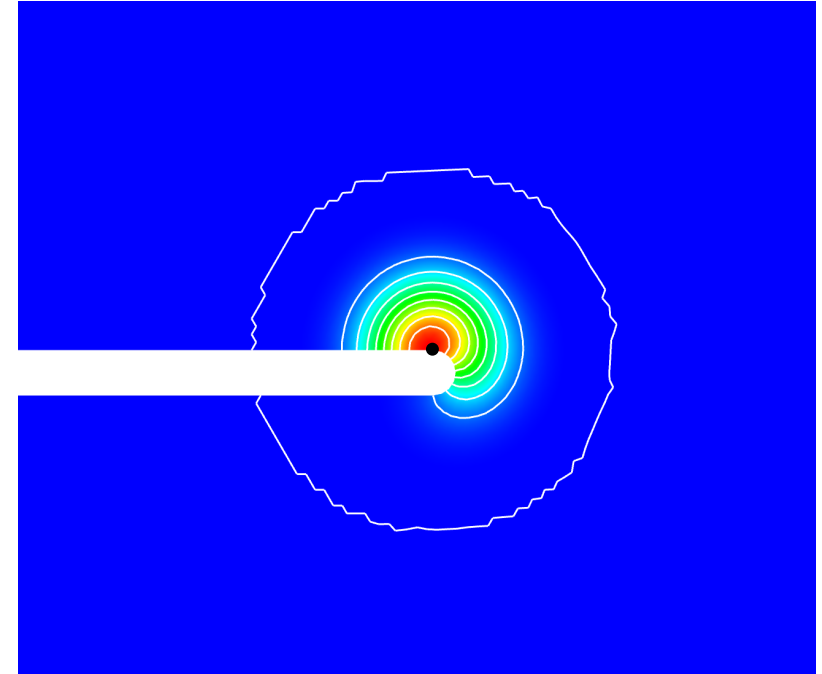
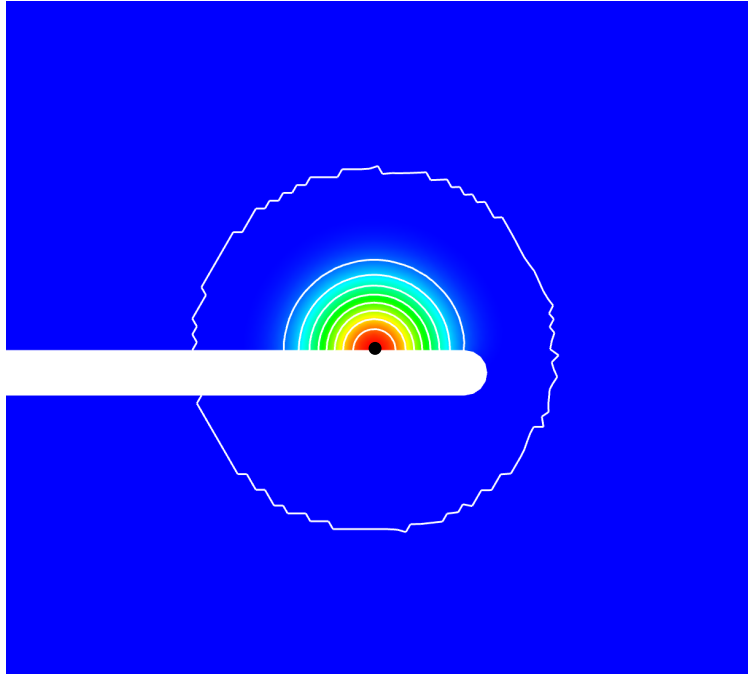
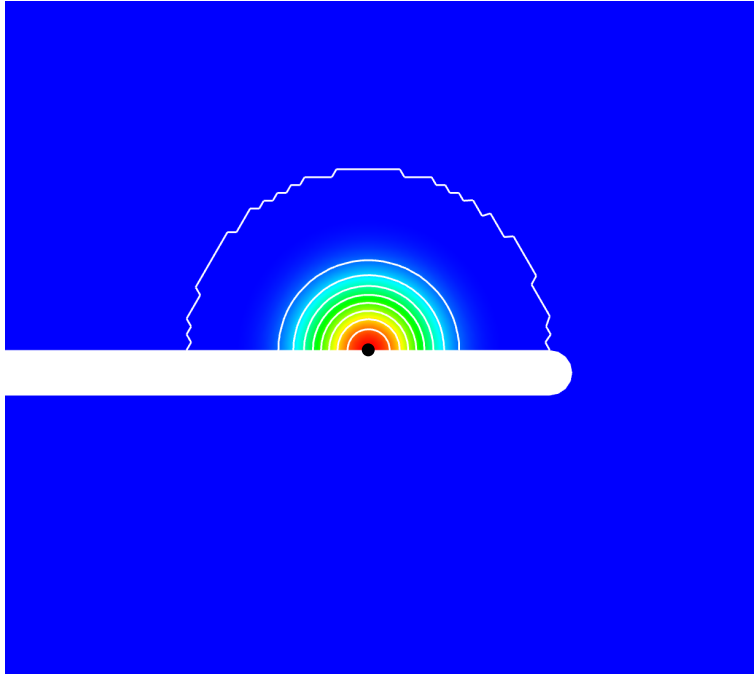
Weight functions using heat flow



- Solve local transient heat conduction problem with certain I.C. and B.C.
- Uses local tri-mesh within support radius.

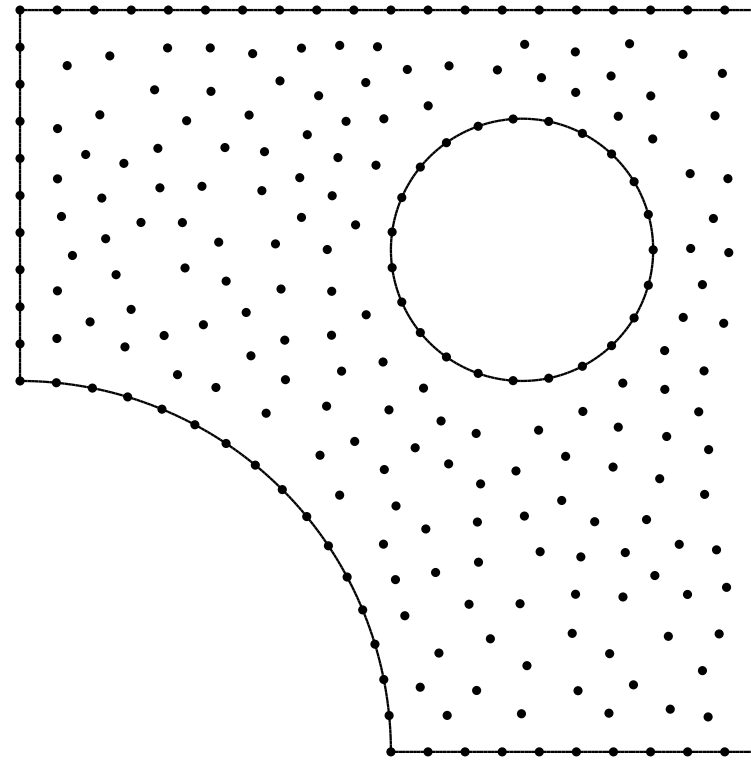
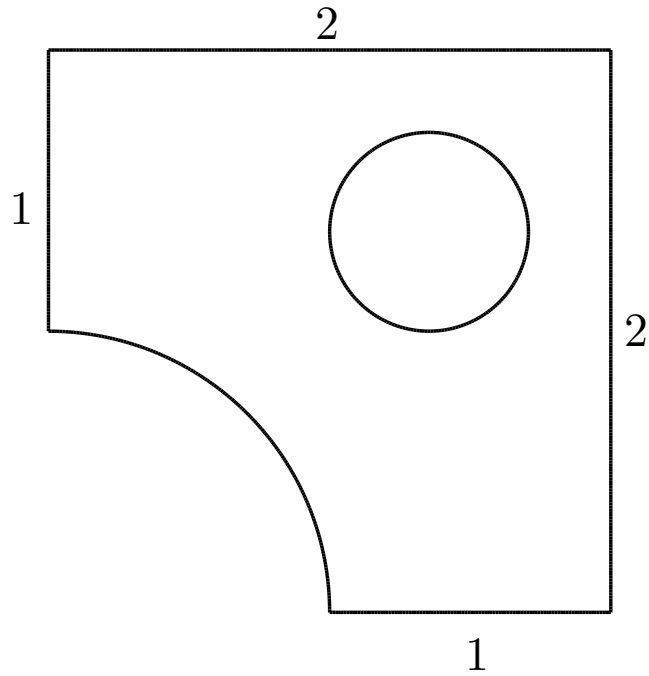


Weight functions using heat flow



Node placement

- uniform on boundary
- random close packing on interior (maximal Poisson sampling)

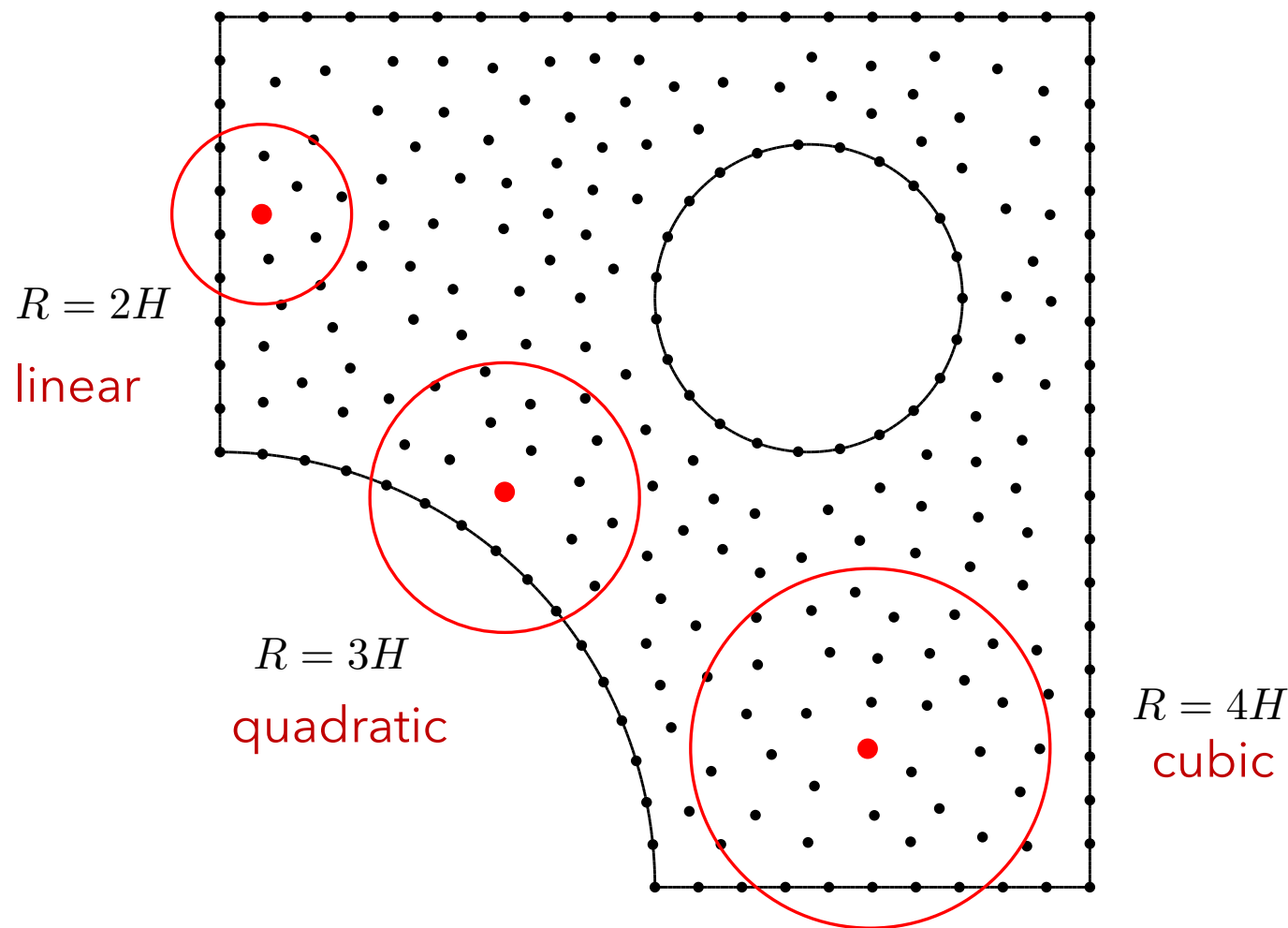


packing size:

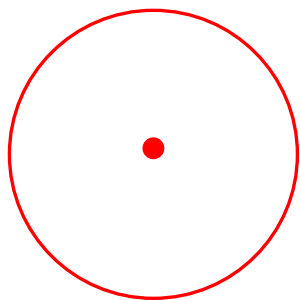
$$H = 0.1$$

Weight function support size

packing size: $H = 0.1$

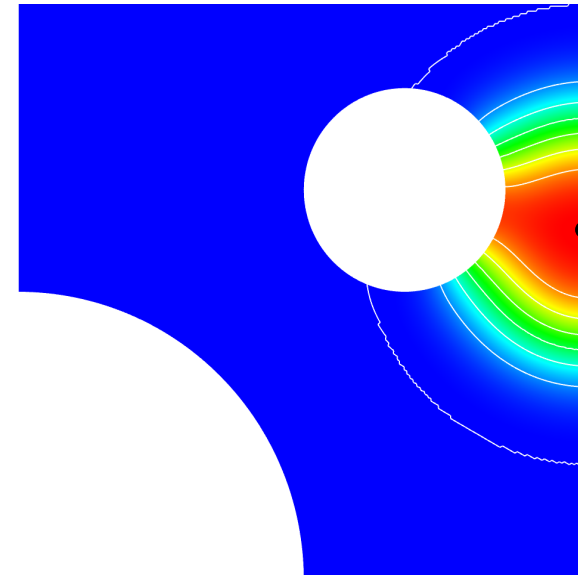
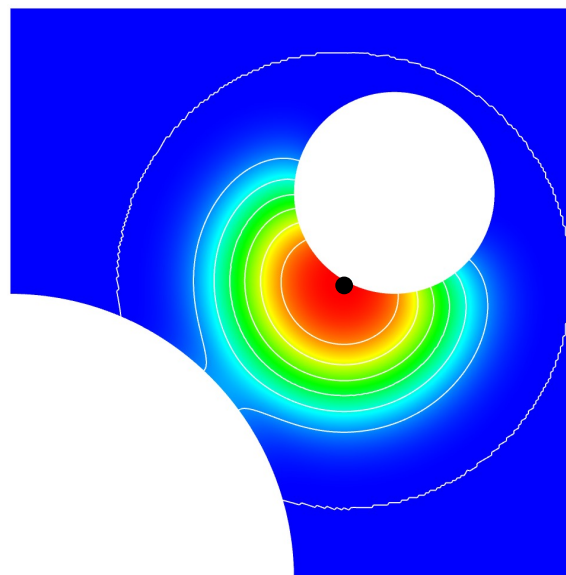
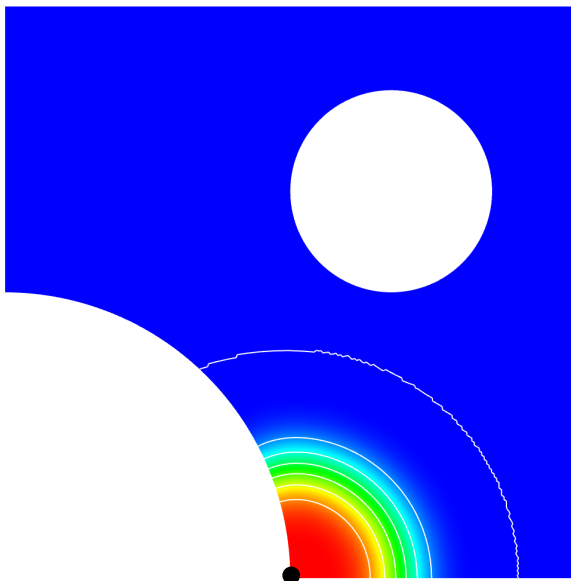


support size

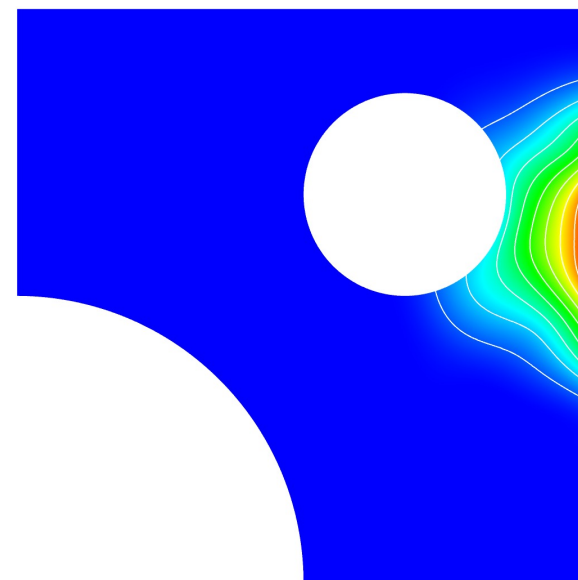
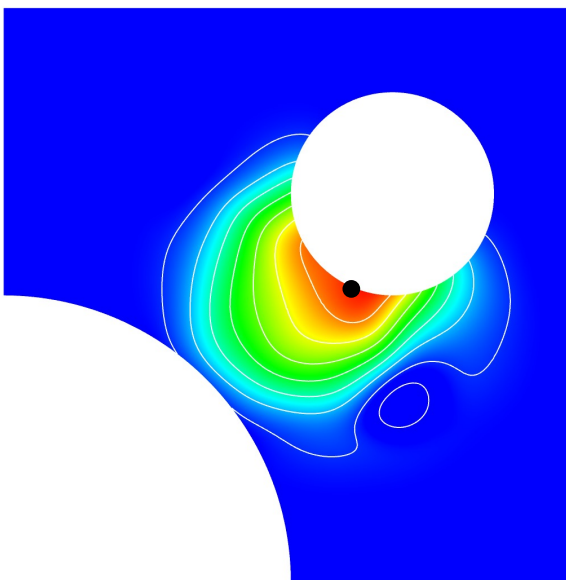
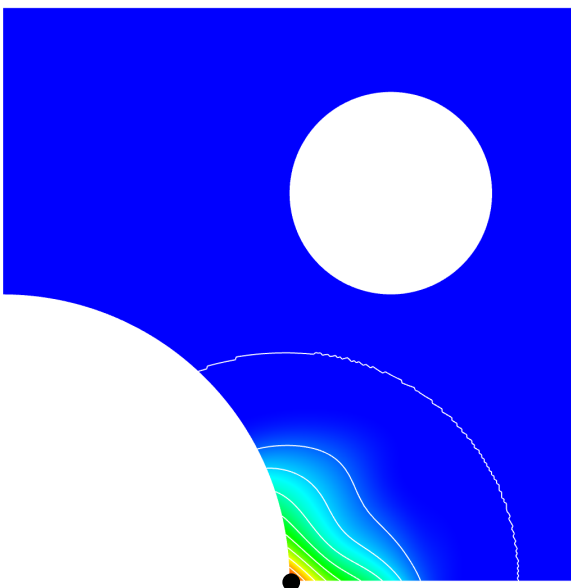


encloses underlying tri mesh

weight
functions



shape
functions
(basis)



Governing equations for solid mechanics (Lagrangian)

strong form

$$\frac{\partial \mathbf{P}}{\partial \mathbf{X}} : \mathbf{I} = \rho_0 \ddot{\mathbf{u}}$$

$$\mathbf{u} = \bar{\mathbf{u}} \quad \text{on} \quad \Gamma_0^u \quad \text{and} \quad \mathbf{P} \cdot \mathbf{N} = \mathbf{t}_0 \quad \text{on} \quad \Gamma_0^t$$

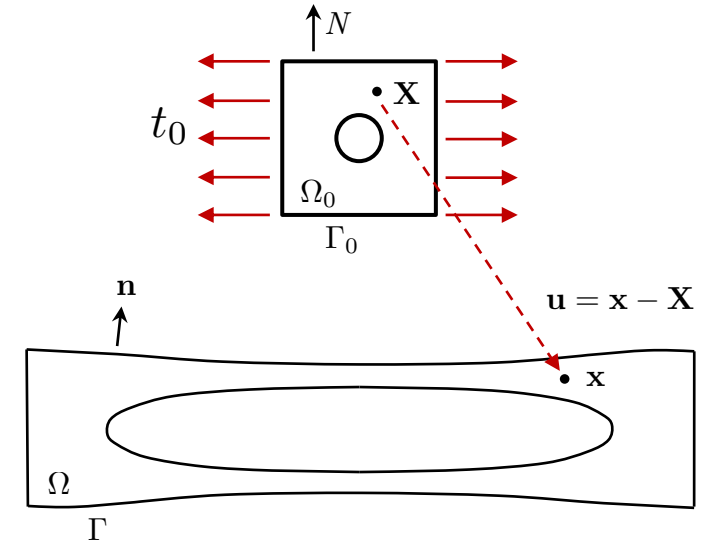
\mathbf{P} is first Piola-Kirchhoff stress tensor

weak form

find the trial functions $\mathbf{u} \in \mathbf{H}^1(\Omega_0)$ such that

$$\int_{\Gamma_0^t} \mathbf{t}_0 \cdot \mathbf{v} \, dS - \int_{\Omega_0} \mathbf{P} : (\partial \mathbf{v} / \partial \mathbf{X}) \, d\mathbf{X} = \int_{\Omega_0} \rho_0 \ddot{\mathbf{u}} \cdot \mathbf{v} \, d\mathbf{X}$$

for all test functions $\mathbf{v} \in \mathbf{H}_0^1(\Omega_0)$



Governing equations for linear elasticity

strong form $\frac{\partial \boldsymbol{\sigma}}{\partial \mathbf{x}} : \mathbf{I} + \mathbf{f} = \mathbf{0} \quad \mathbf{u} = \bar{\mathbf{u}} \text{ on } \Gamma_u \text{ and } \boldsymbol{\sigma} \mathbf{n} = \mathbf{t} \text{ on } \Gamma_t$

$\boldsymbol{\sigma} = \mathbb{C} \boldsymbol{\epsilon}$, where $\boldsymbol{\epsilon} := \text{sym}(\nabla \mathbf{u})$ *(linear elastic)*

$\exists \alpha_l, \alpha_u > 0$ such that $\alpha_l \boldsymbol{\epsilon} : \boldsymbol{\epsilon} \leq \boldsymbol{\epsilon} : (\mathbb{C}(\mathbf{x}) \boldsymbol{\epsilon}) \leq \alpha_u \boldsymbol{\epsilon} : \boldsymbol{\epsilon} \quad \forall \boldsymbol{\epsilon}$ *(uniform ellipticity)*

weak form find the trial functions $\mathbf{u} \in \mathbf{H}^1(\Omega_0)$ such that

$$\int_{\Omega} \boldsymbol{\sigma} : (\partial \mathbf{v} / \partial \mathbf{x}) d\Omega = \int_{\Omega} \mathbf{f} \cdot \mathbf{v} d\Omega + \int_{\Gamma_t} \mathbf{t} \cdot \mathbf{v} d\Gamma$$

for all test functions $\mathbf{v} \in \mathbf{H}_0^1(\Omega_0)$

abstract variational problem $a(u, v) = b(v)$ with bilinear form $a(u, v) = \int_{\Omega_e} \nabla u : \mathbb{C} \nabla v d\Omega$

How to do quadrature?



Too expensive to use fine-scale triangulation for quadrature!

Observe that partition-of-unity property allows us to approximate any continuous function arbitrarily closely using only point evaluations as long as basis functions have local support.

Given $\sum_K \phi_K(\mathbf{x}) = 1$ then it follows that $f(\mathbf{x}) \approx \sum_K f(\mathbf{x}_K) \phi_K(\mathbf{x}) := f_h(\mathbf{x})$ *non-interpolatory approximation*

Theorem: For every $\varepsilon > 0$ and $\mathbf{x} \in \overline{\Omega}$, there exists $h(\varepsilon) > 0$ such that $|f_h(\mathbf{x}) - f(\mathbf{x})| < \varepsilon$.

It follows that $\int f(\mathbf{x}) d\Omega \approx \int f_h(\mathbf{x}) d\Omega$

with $\left| \int f_h(\mathbf{x}) d\Omega - \int f(\mathbf{x}) d\Omega \right| \leq \int |f_h(\mathbf{x}) - f(\mathbf{x})| d\Omega \leq \int \varepsilon d\Omega = V \cdot \varepsilon$

Can obtain rates of convergence using Taylor's theorem.

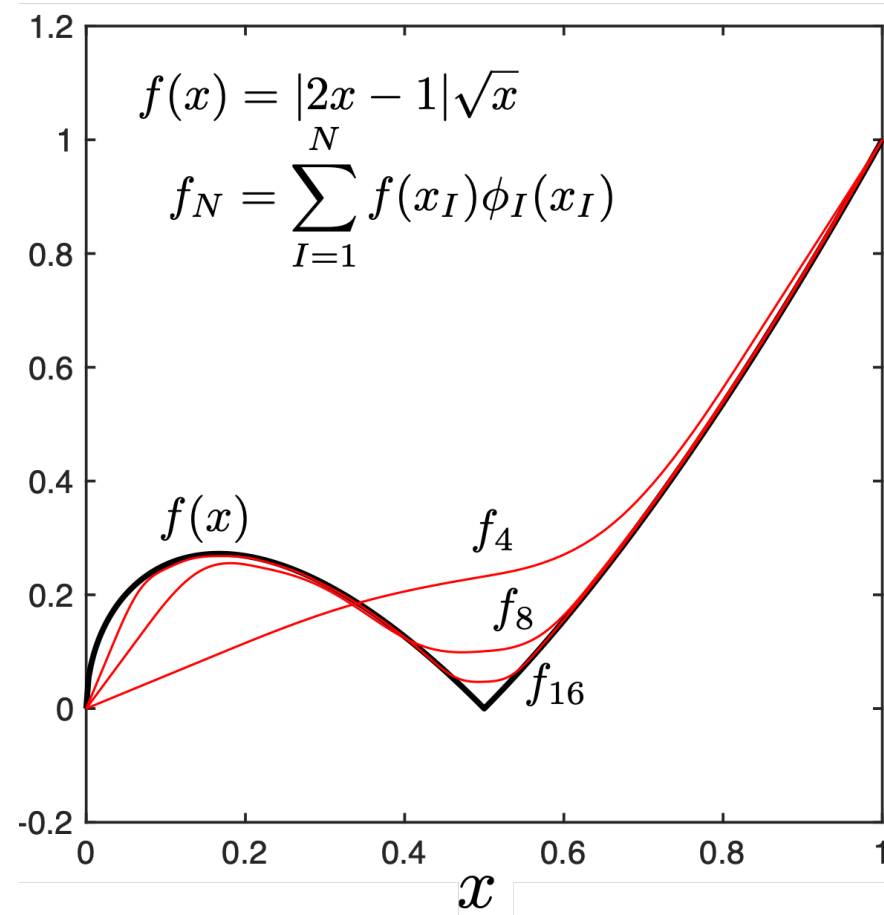
Approximation property



let $f(x) = |2x - 1|\sqrt{x}$

$$f_N = \sum_{K=1}^N f(x_K)\phi_K(x)$$

function approximation



$$\int f(\mathbf{x}) d\Omega \approx \int f_h(\mathbf{x}) d\Omega = \int \sum_K f(\mathbf{x}_K) \phi_K(\mathbf{x}) d\Omega = \sum_K f(\mathbf{x}_K) \int \phi_K(\mathbf{x}) d\Omega$$

Define quadrature weight as

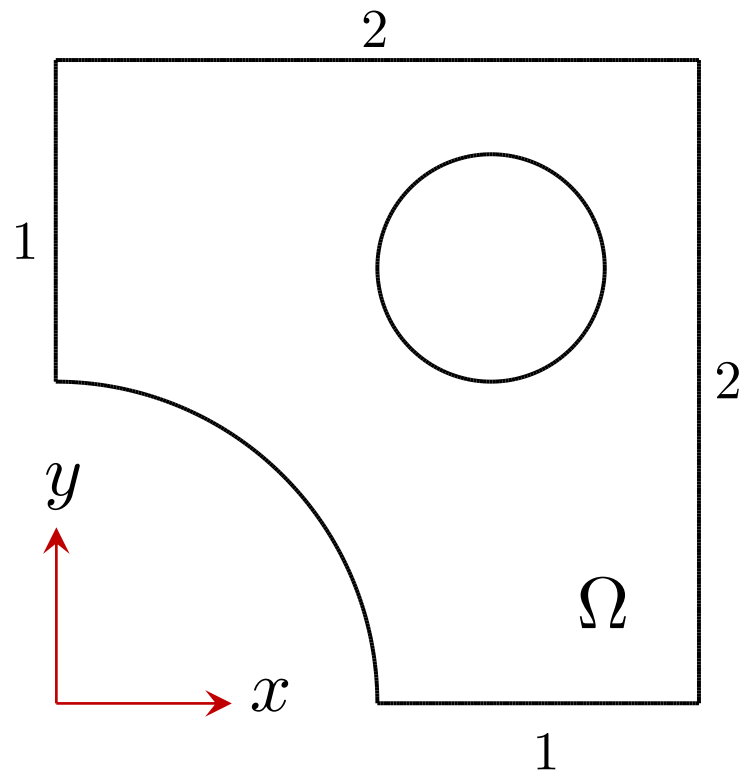
$$w_K = \int_{\Omega} \phi_K(\mathbf{x}) d\Omega$$

$$\int f(\mathbf{x}) d\Omega \approx \sum_K w_K f(\mathbf{x}_K)$$

Can show that $\sum_K w_K = V$ and $\sum_K w_K \mathbf{x}_K = \int_{\Omega} \mathbf{x} d\Omega$

Now have a second-order integration scheme that can integrate linear functions exactly.

Quadrature example



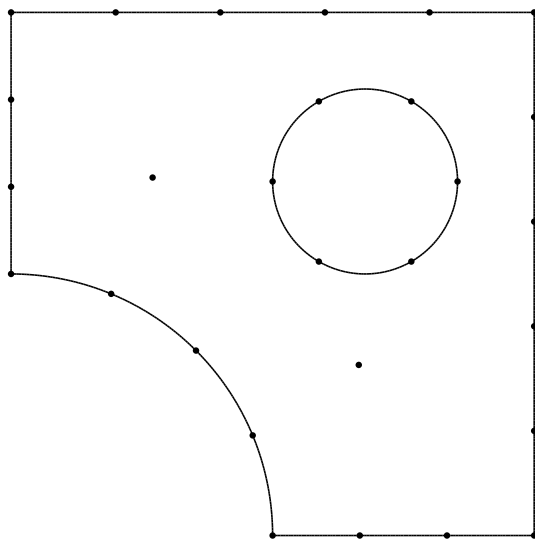
$$f(x, y) = \sin(\pi x/2) \sin(\pi y)$$

$$\text{error} := \left| \sum_K w_K f(\mathbf{x}_K) - \int_{\Omega} f d\Omega \right|$$

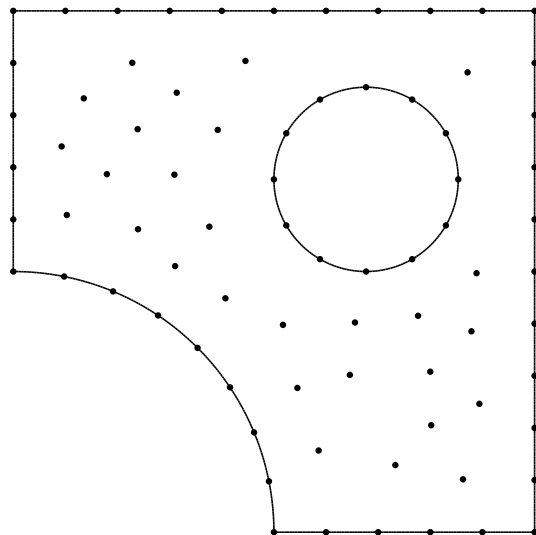
Quadrature example



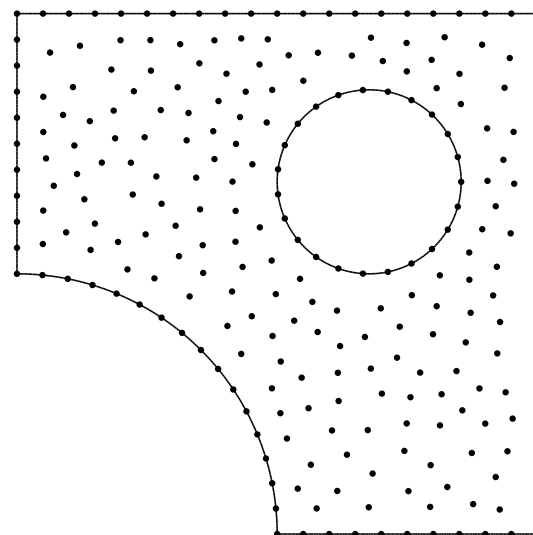
$H = 0.4$



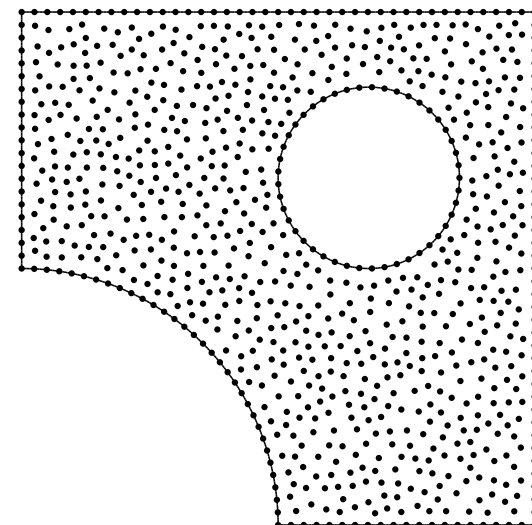
$H = 0.2$



$H = 0.1$



$H = 0.05$

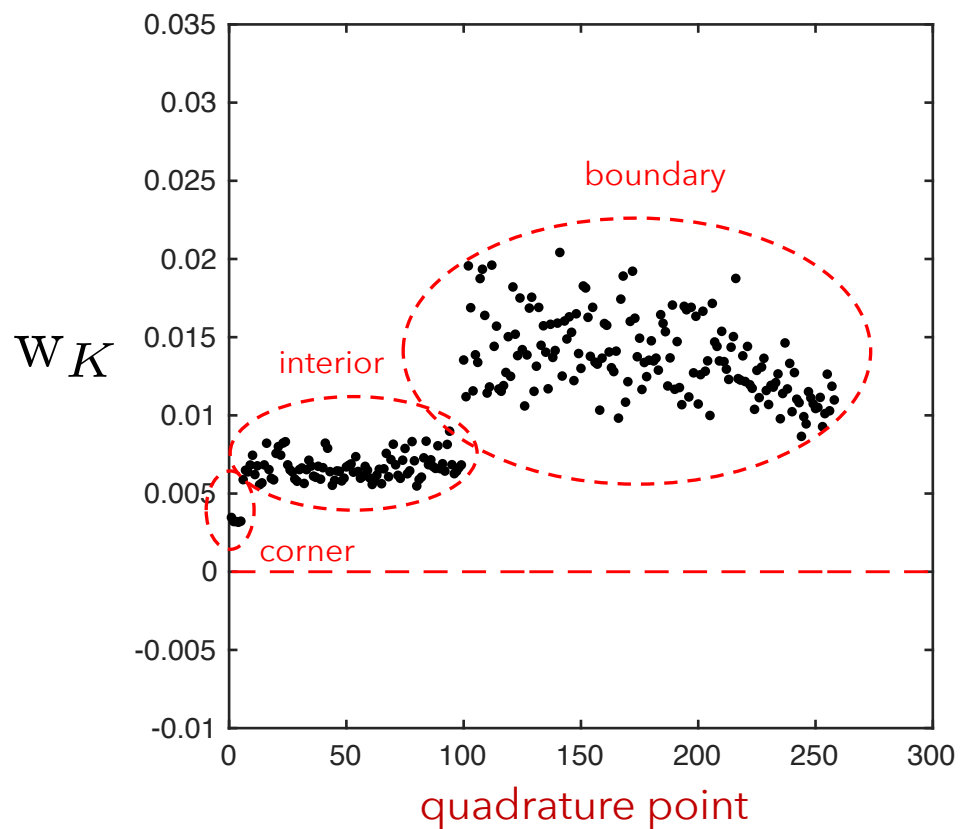


Evaluate error for 10 realizations.

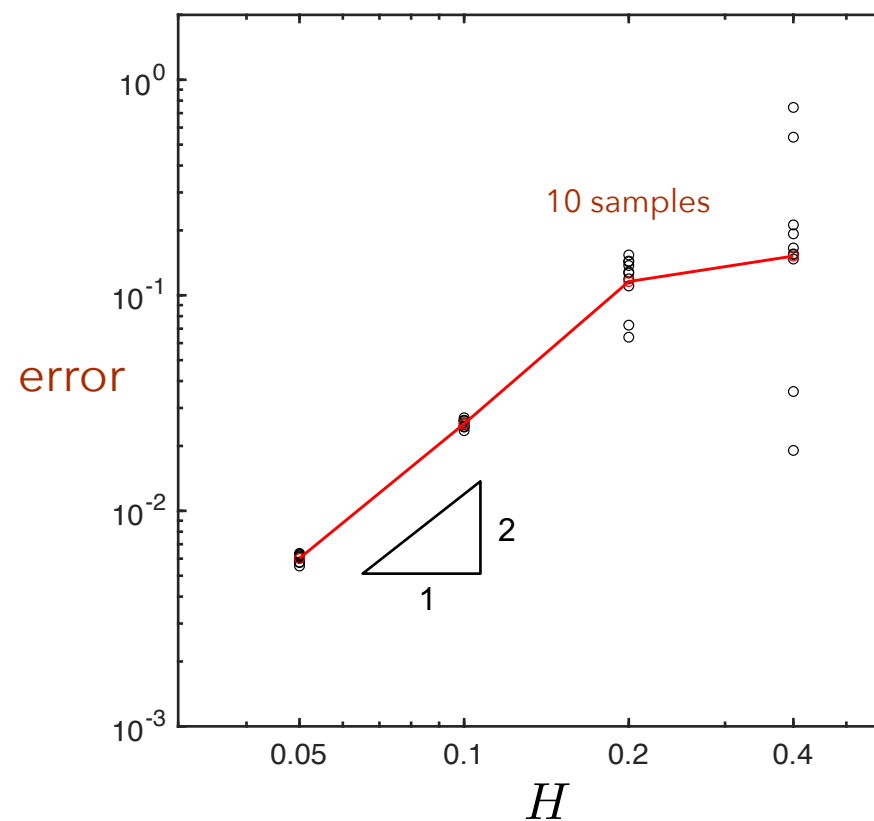
Quadrature weights



$$w_K = \int_{\Omega} \phi_K(\mathbf{x})$$



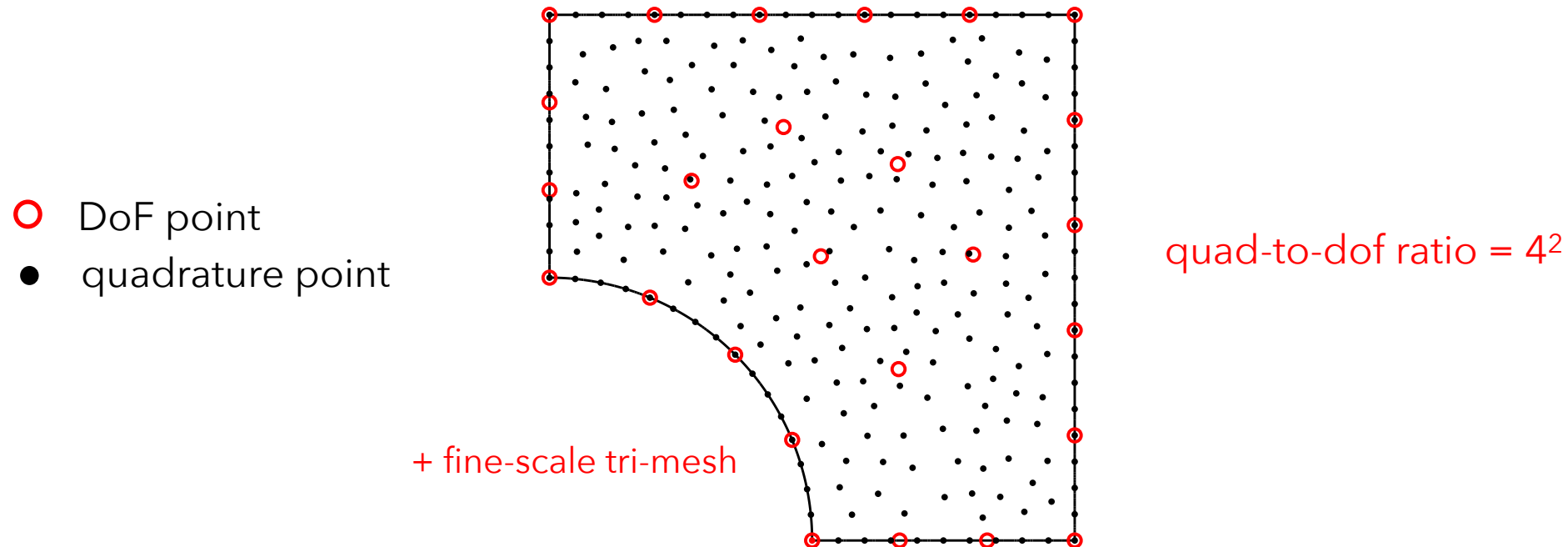
$$\text{error} := \left| \sum_K w_K f(\mathbf{x}_K) - \int_{\Omega} f d\Omega \right|$$



Element-free approach to solve BVPs

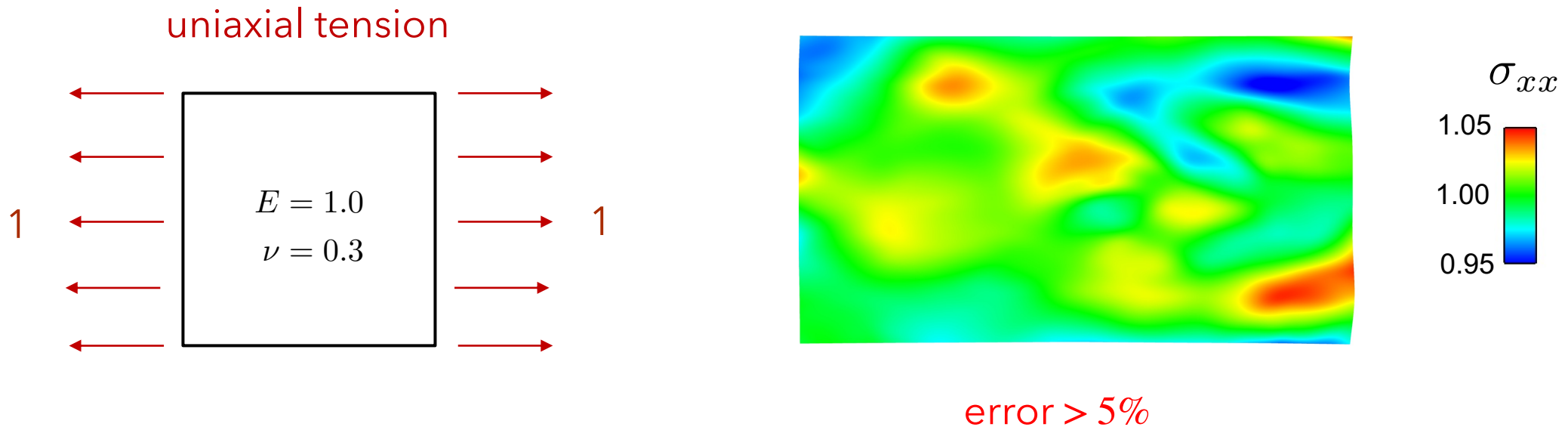


Use two meshfree clouds: one for solution discretization (DoF) and one for quadrature.



What ratio of quad points to dof points is needed for stability (coercivity of bilinear form)?

Patch test (linear consistency)



Consistency of discrete form (integration)

- For convergence of discrete approximation, need to ensure consistency of discrete and continuous bilinear forms.
- Requires polynomial consistency of shape-function gradients (including quadrature).
- To obtain quadrature consistency, project the DoF shape function gradients to the subspace of quadrature shape functions.
- Only performed once in a pre-processing step.

$\{\phi_I, I = 1, \dots, N\}$ *DoF basis (shape functions)*

$\{\Phi_K, K = 1, \dots, M\}$ *Quadrature basis (shape functions)*

$$\bar{\nabla} \phi_I := \arg \min \int_{\Omega} \left(\nabla \phi_I - \sum_{K=1}^M a^K \Phi_K \right)^2 d\Omega \quad (L_2 \text{ projection})$$

The projection can be written in terms of the dual or conjugate basis $\{\Phi^J\}$

$$(\Phi_K, \Phi^J) = \delta_K^J \quad \text{bi-orthogonal}$$

$$\bar{\nabla} \phi_I = \sum_K \underbrace{(\nabla \phi_I, \Phi_K)}_{\text{covariant components}} \Phi^K = \sum_K \underbrace{(\nabla \phi_I, \Phi^K)}_{\text{contravariant components}} \Phi_K$$

Can prove polynomial consistency up to the order of the precision of $\{\Phi_K\}$

Theorem: $\int_{\Omega} \mathbf{p} \bar{\nabla} \phi_I d\Omega = \int_{\Omega} \mathbf{p} \nabla \phi_I d\Omega \quad \text{for all } \mathbf{p} \in \mathbb{P}_k(\Omega)$

This ensures satisfaction of the patch test.

Replace the original bilinear form $a(u, v) = \int_{\Omega} \nabla u : \mathbb{C} \nabla v \, d\Omega$

with this modified bilinear form $\bar{a}(u, v) = \int_{\Omega} \bar{\nabla} u : \mathbb{C} \bar{\nabla} v \, d\Omega$ *Note: This modified bilinear form is still symmetric (Bubnov-Galerkin).*

$$\bar{a}(u, v) = \int_{\Omega} \left[\sum_I (\nabla u, \Phi_I) \Phi^I \right] \mathbb{C} \left[\sum_J (\nabla v, \Phi_J) \Phi^J \right] d\Omega$$

$$\bar{a}(u, v) = \sum_{I, J} (\nabla u, \Phi_I) \mathbb{C} (\nabla v, \Phi_J) \underbrace{\int_{\Omega_e} \Phi^I \Phi^J \, d\Omega}_{G^{IJ}} \quad \text{inverse of the Gram matrix}$$

Can show that $G^{IJ} = (G_{IJ})^{-1}$

where $G_{IJ} = \int_{\Omega_e} \Phi_I \Phi_J d\Omega$ is the Gram matrix for the basis $\{\Phi_K\}$

Can show that $\Phi^I = G^{IJ} \Phi_J$ and $\Phi_I = G_{IJ} \Phi^J$ *"raising" and "lowering" of indices*

$$\bar{a}(u, v) = \sum_{I, J} G^{IJ} (\nabla u, \Phi_I) \mathbb{C} (\nabla v, \Phi_J) = \sum_K (\nabla u, \Phi^K) \mathbb{C} (\nabla v, \Phi_K)$$

Looks like a sum over quadrature points.

Since $G^{IJ} = (G_{IJ})^{-1}$ is dense:

Replace G_{IJ} with row-sum lumped version: $G_{IJ}^L := \sum_J G_{IJ} = \text{diag}\{w_K\}$

where recall $w_K = \int_{\Omega} \phi_K(\mathbf{x}) d\Omega$

Then $\bar{a}(u, v) \rightarrow \bar{a}^L(u, v) = \sum_K \frac{1}{w_K} (\nabla u, \Phi_K) \mathbb{C} (\nabla v, \Phi_K)$ where $(G_{IJ}^L)^{-1} = \text{diag}\left\{\frac{1}{w_K}\right\}$

Can write $\bar{a}^L(u, v)$ as $\bar{a}^L(u, v) = \sum_K w_K (\bar{\nabla} u)_K : \mathbb{C} (\bar{\nabla} v)_K$

where $(\bar{\nabla} u)_K := \frac{1}{w_K} \int_{\Omega} (\nabla u) \Phi_K d\Omega$

which has the form of a discrete derivative at a quadrature point K .

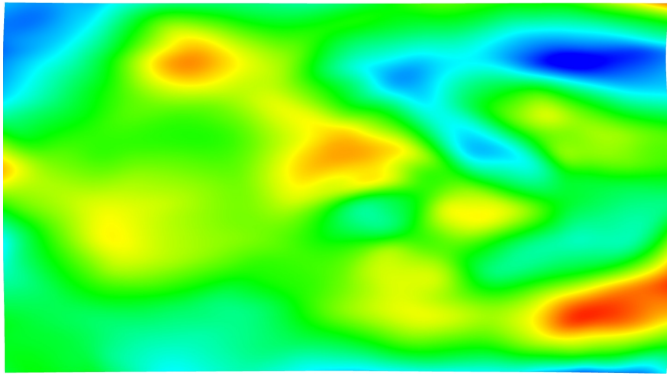
Our discrete bilinear form is now "sparse."

Patch test (linear consistency)



uniaxial tension

no projection



error > 5%

σ_{xx}

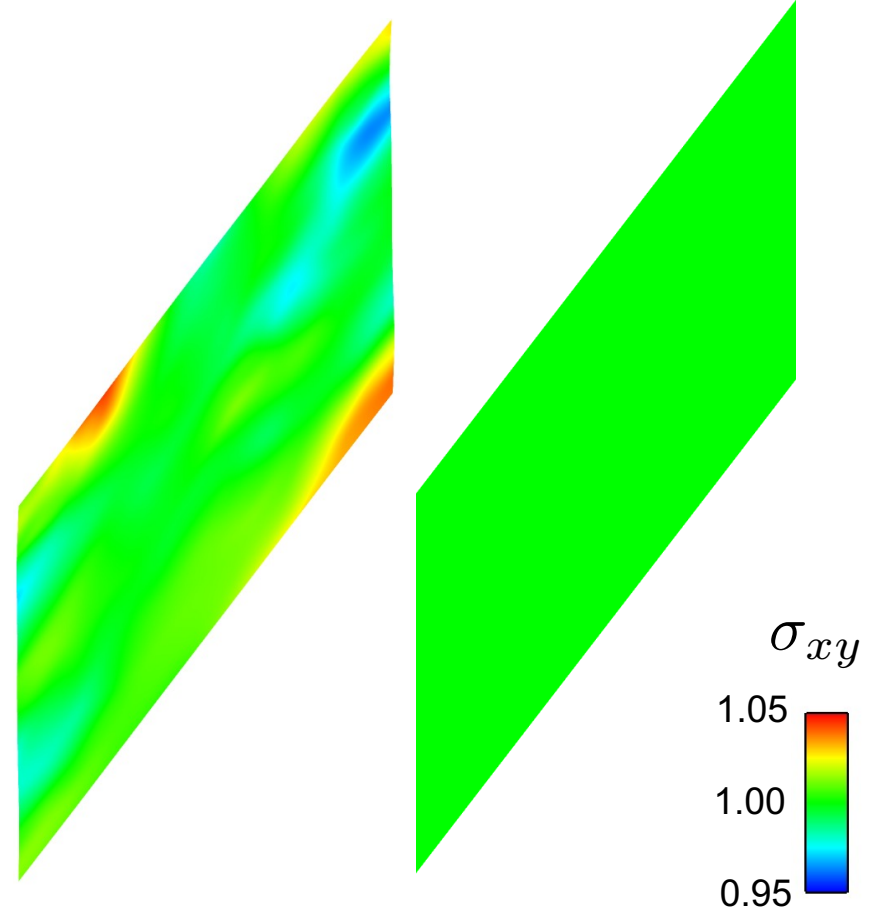
1.05
1.00
0.95

with projection



error < 10^{-13}

pure shear

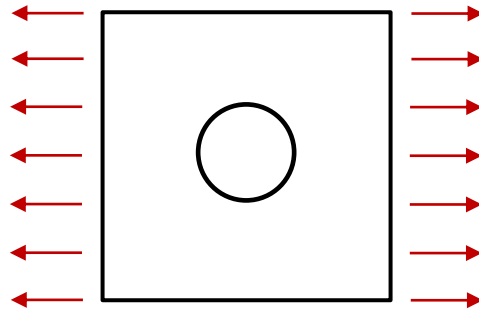


σ_{xy}

1.05
1.00
0.95

Example: plate with hole

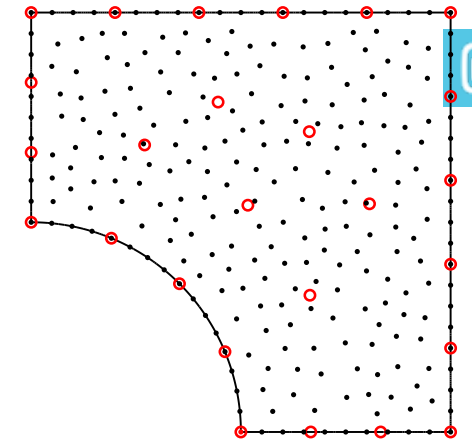
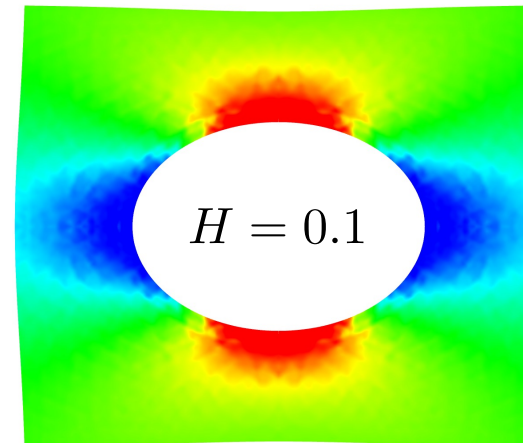
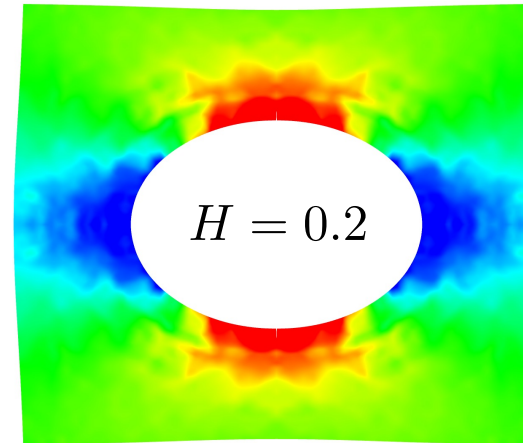
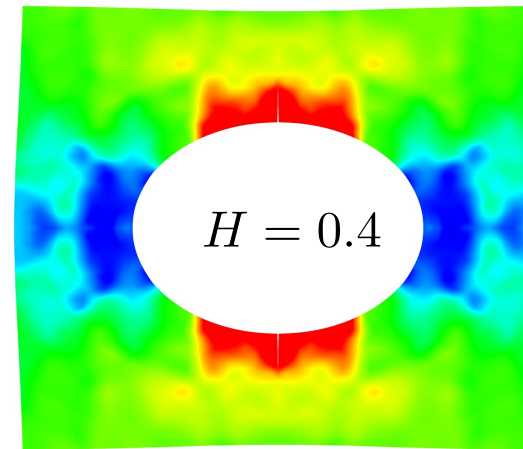
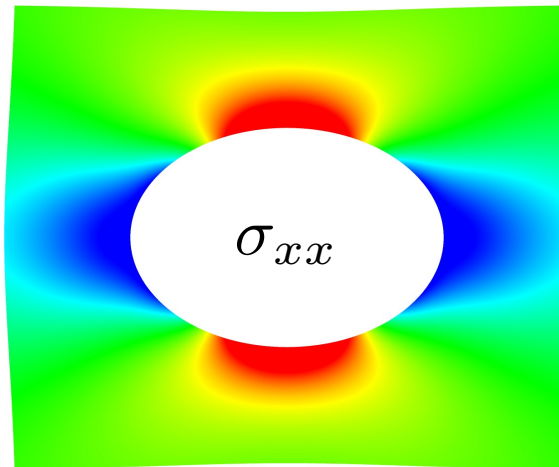
uniaxial tension



$$E = 1.0$$

$$\nu = 0.3$$

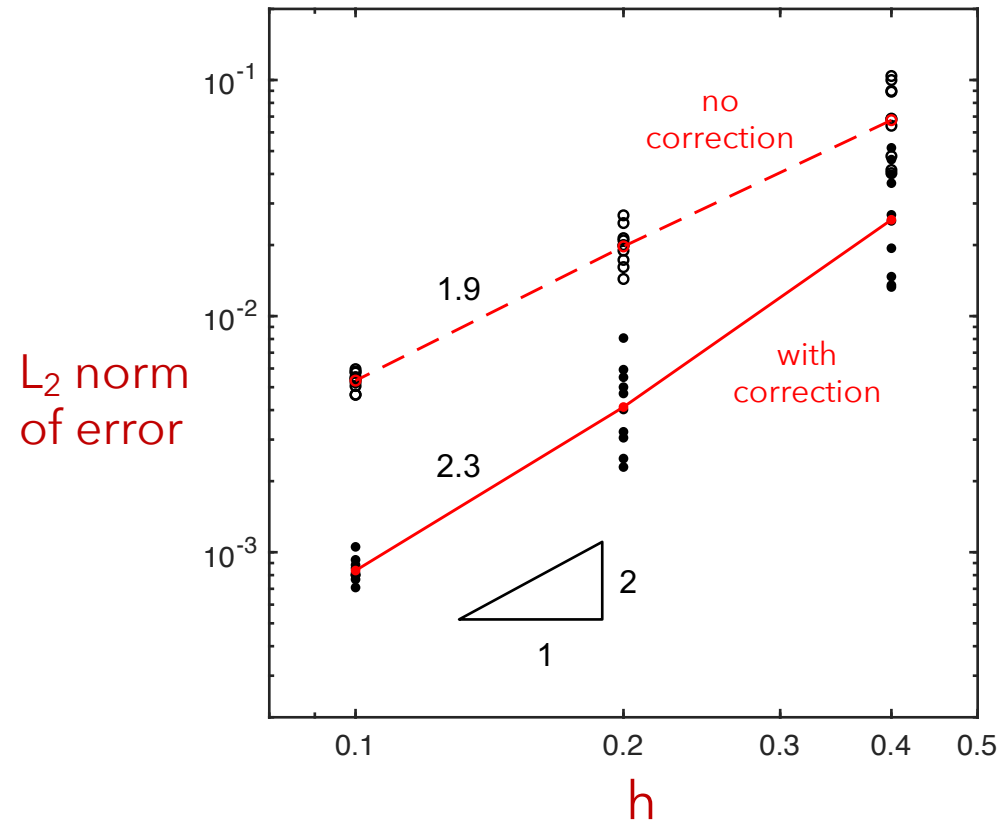
exact



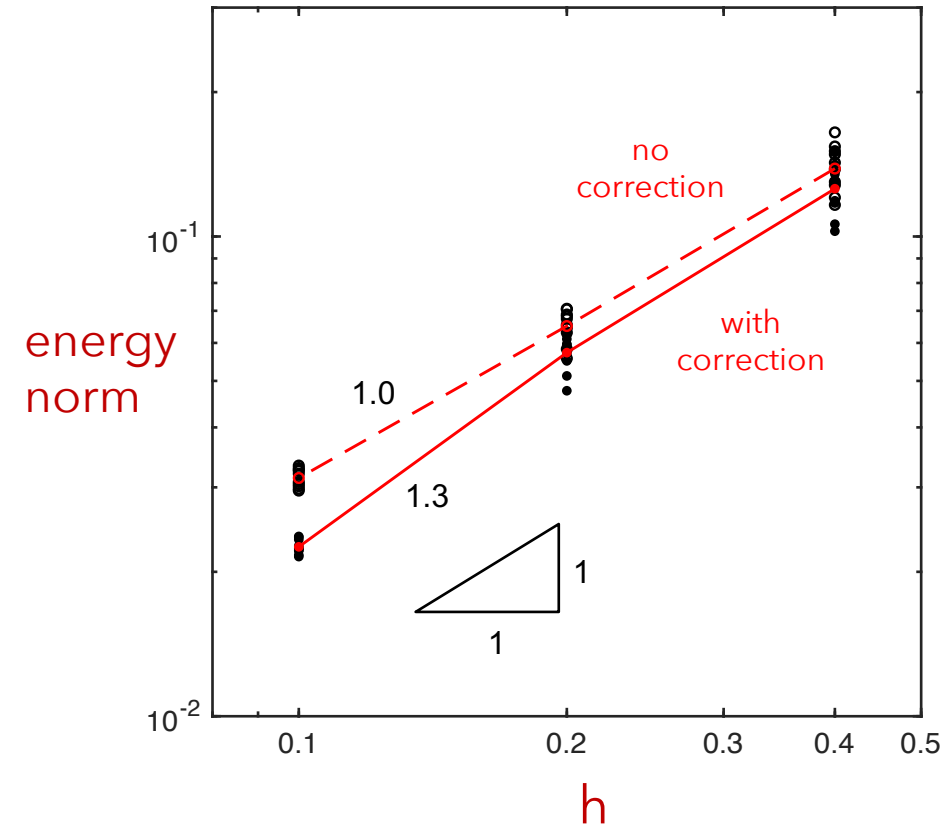
Example: plate with hole



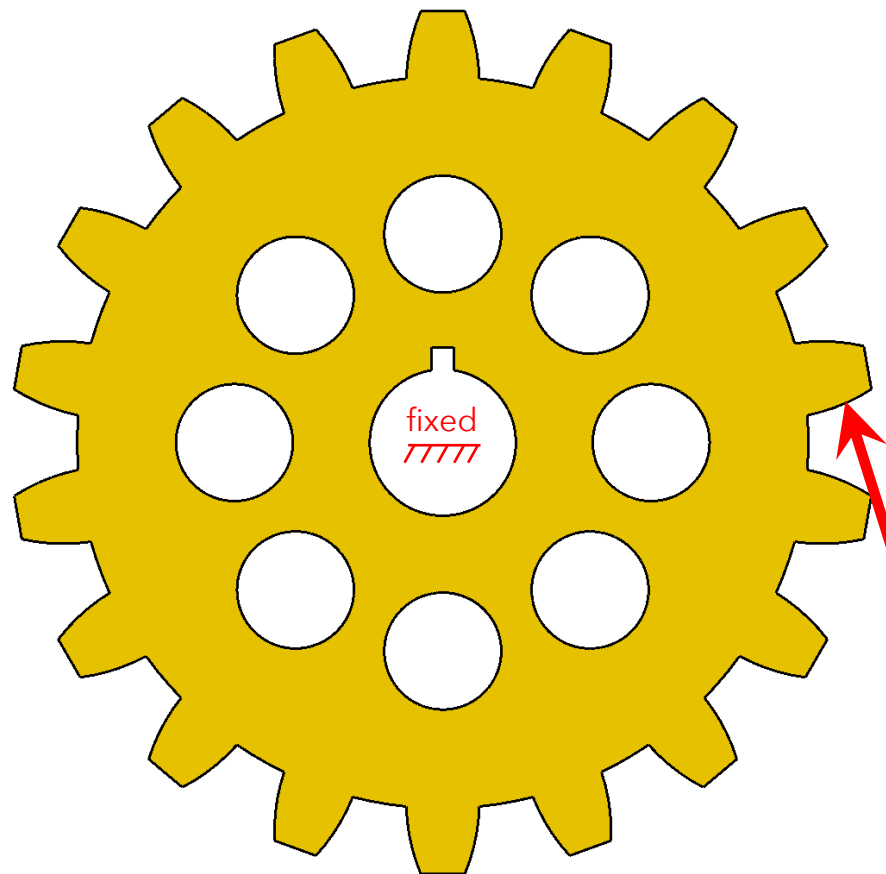
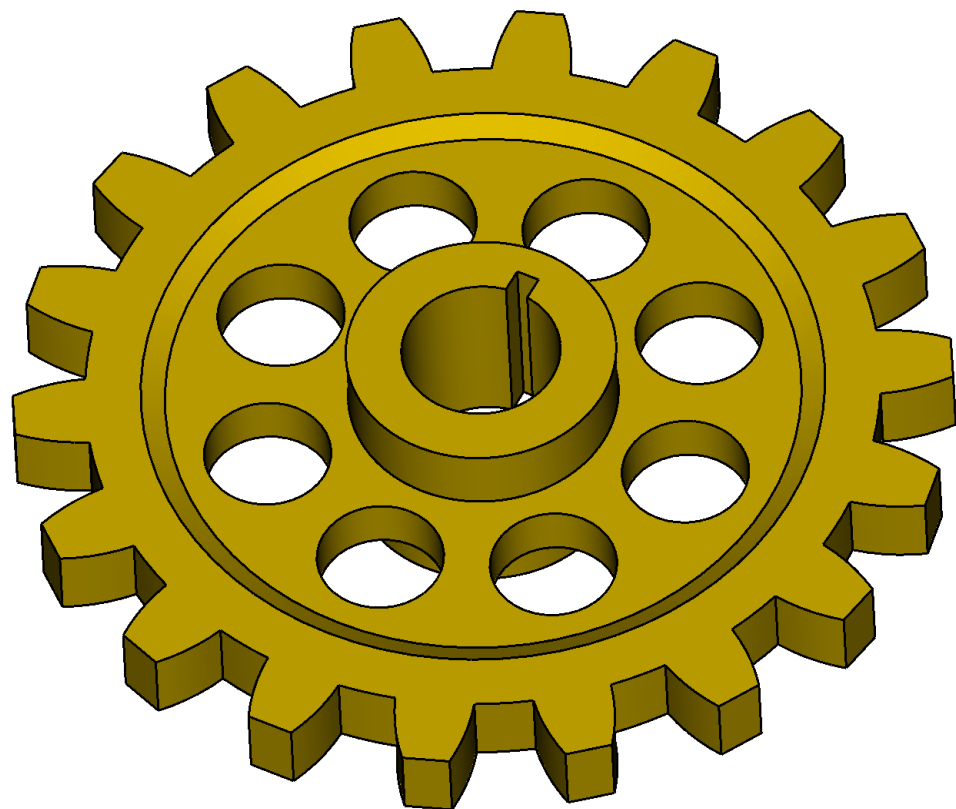
L_2 norm

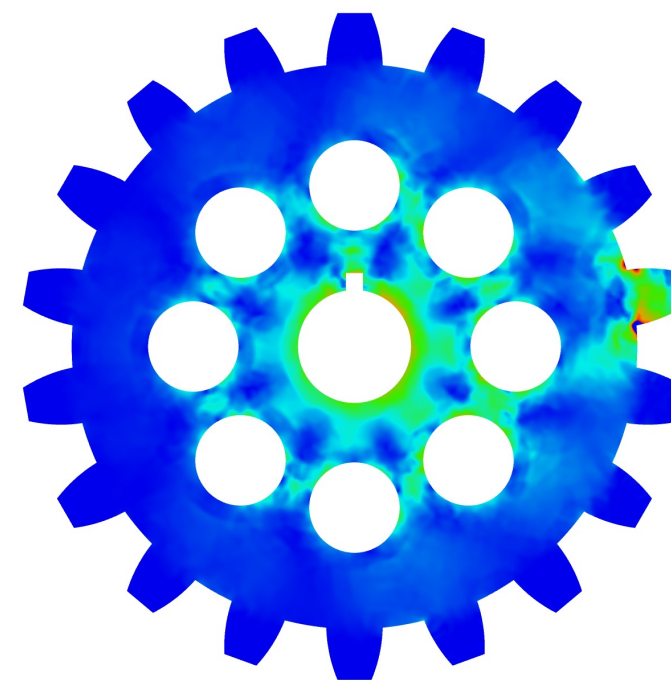
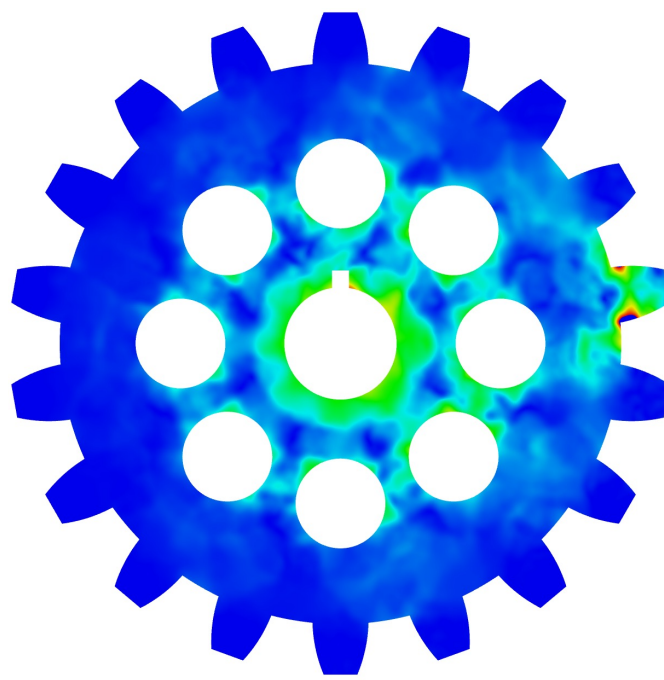
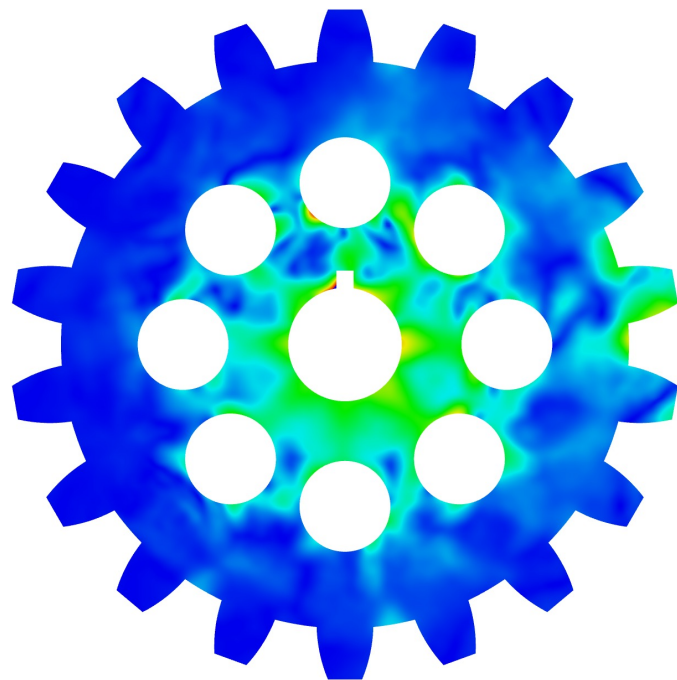
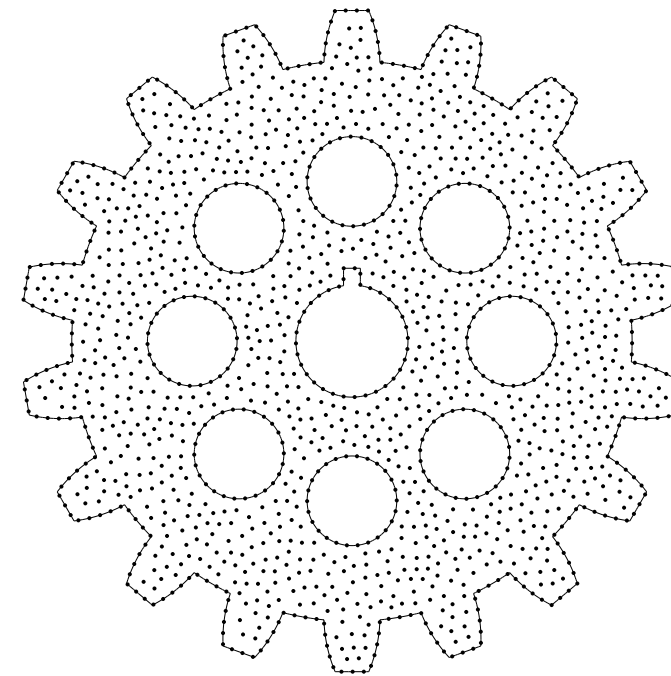
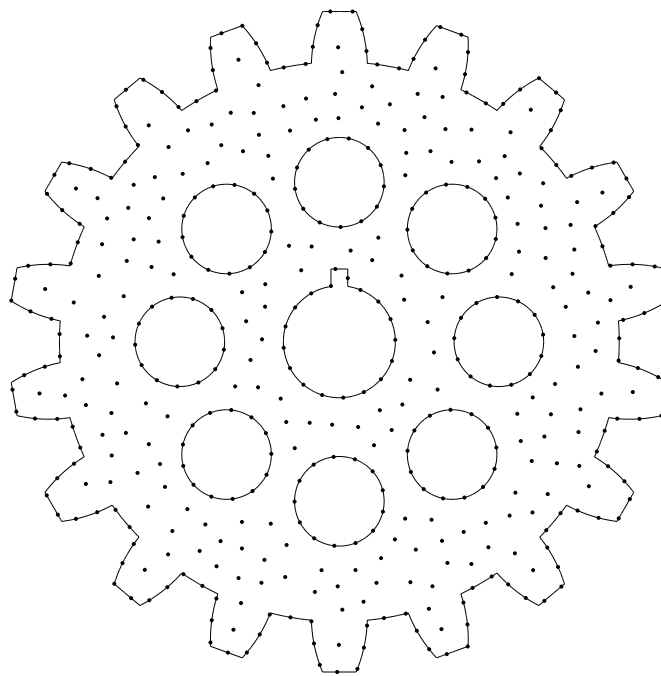
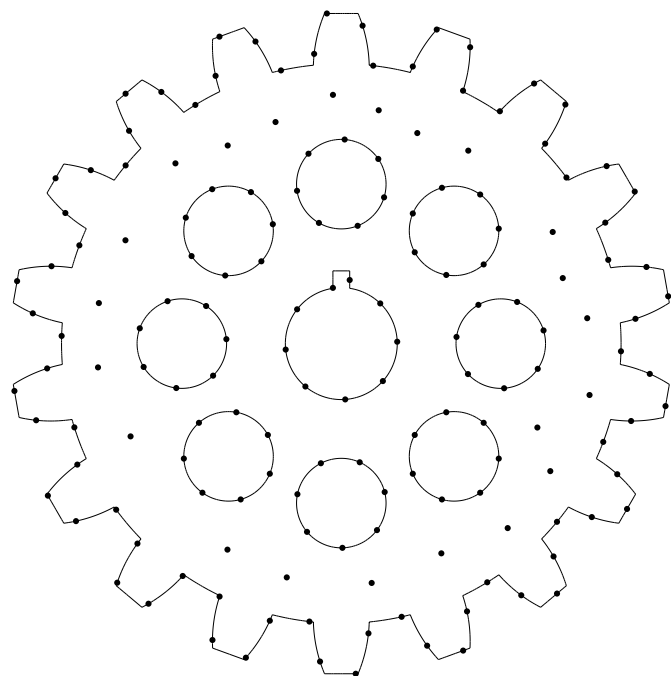


energy norm



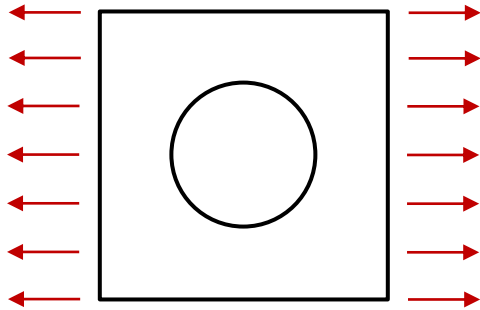
Example





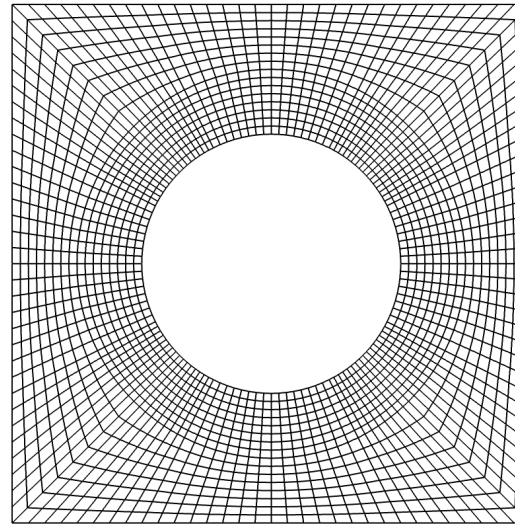
Example: hyperelastic, hole-in-plate

uniaxial extension

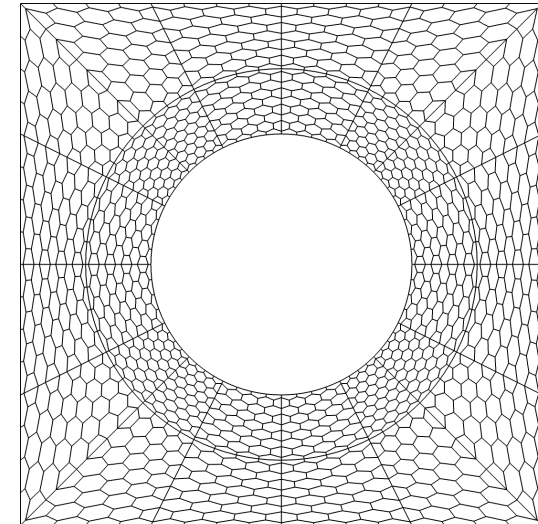


- plane strain
- quarter symmetry model used

quad mesh



mapped hexagon mesh

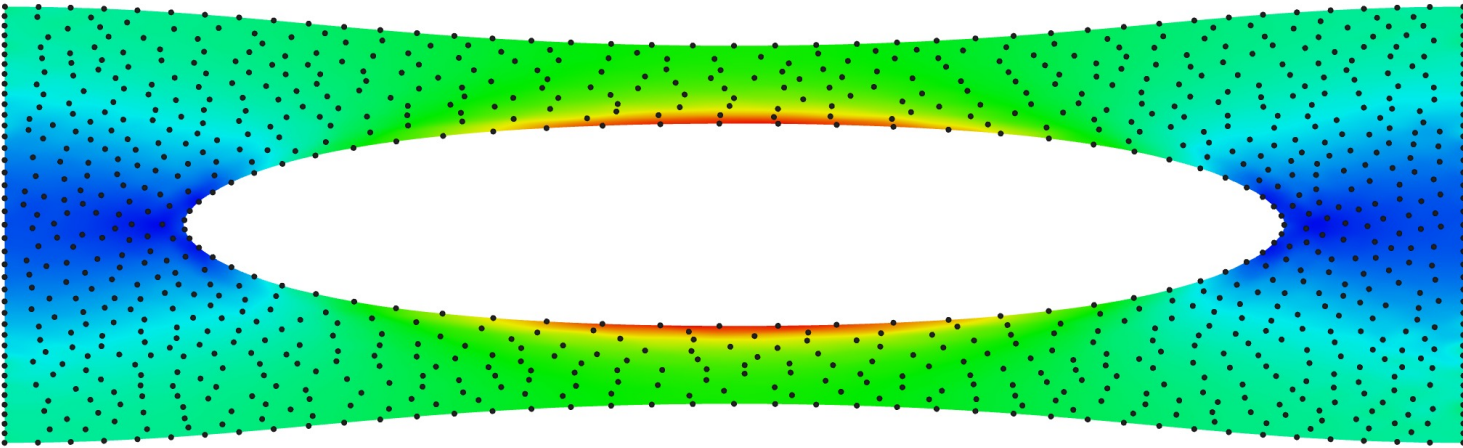
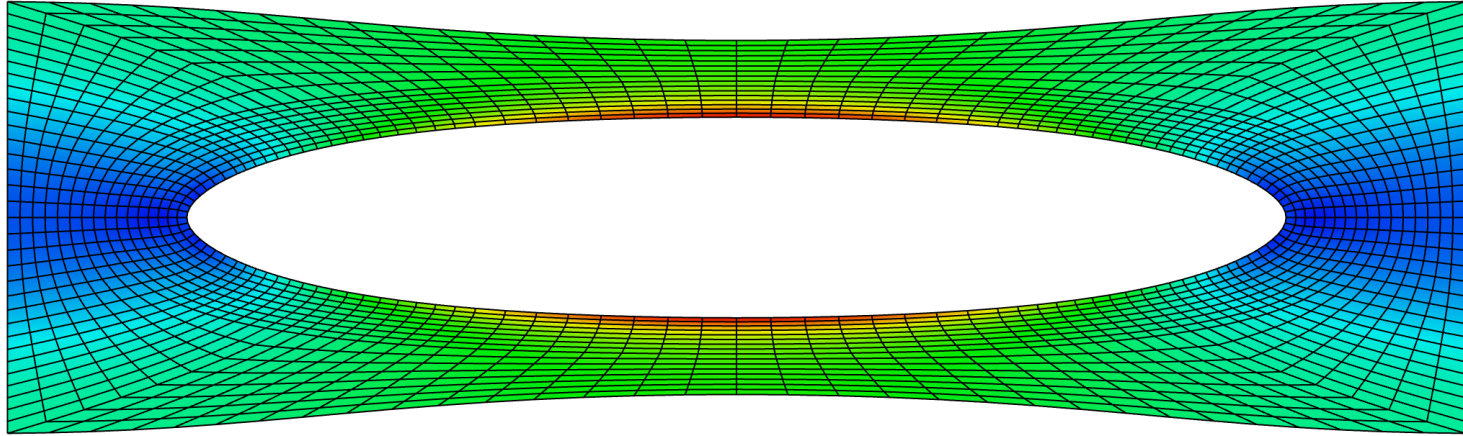


compressible neo-Hookean material

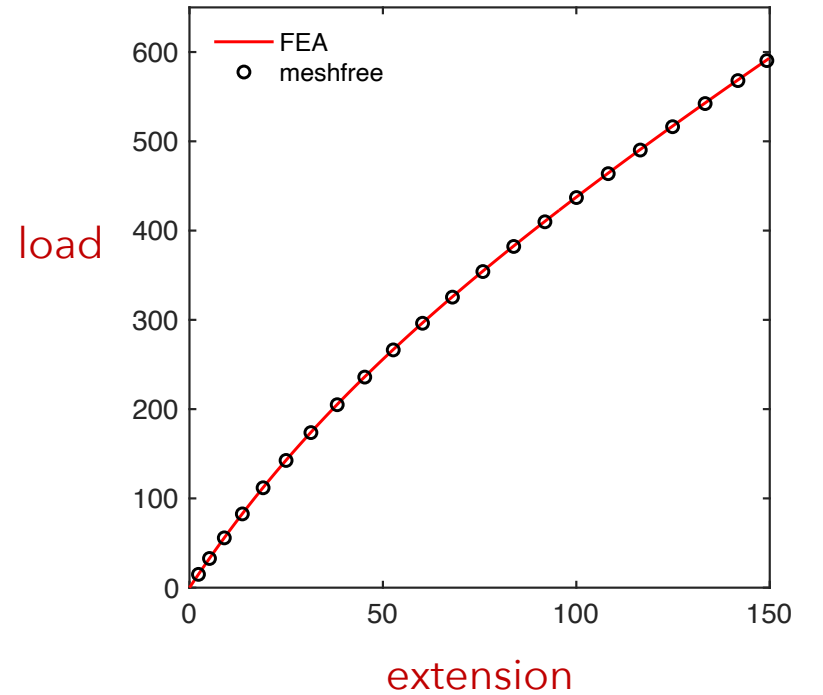
$$\boldsymbol{\sigma} = \frac{\mu}{J}(\mathbf{F}\mathbf{F}^T - \mathbf{I}) + \frac{\lambda}{\ln J}\mathbf{I}$$

$$J = \det \mathbf{F} \quad \mathbf{F} = \frac{\partial \mathbf{x}}{\partial \mathbf{X}}$$

Example: hyperelastic, hole-in-plate



load vs. extension



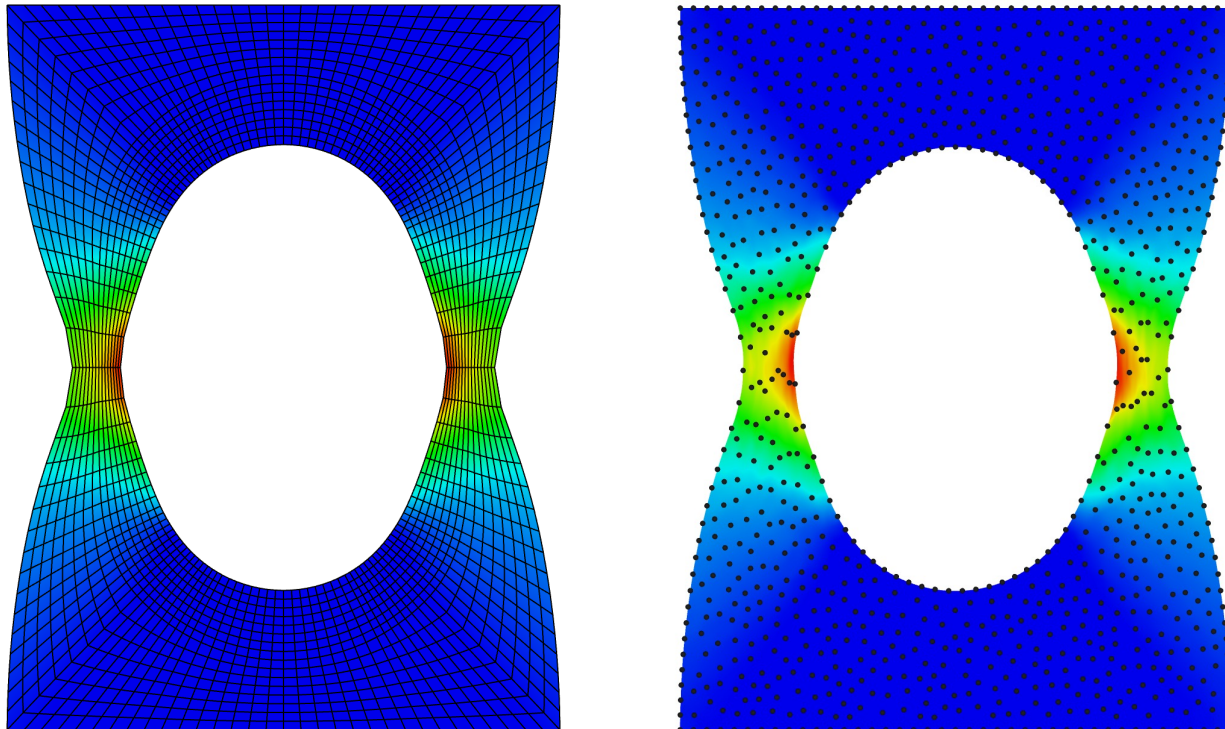
Example: elastic-plastic, hole-in-plate



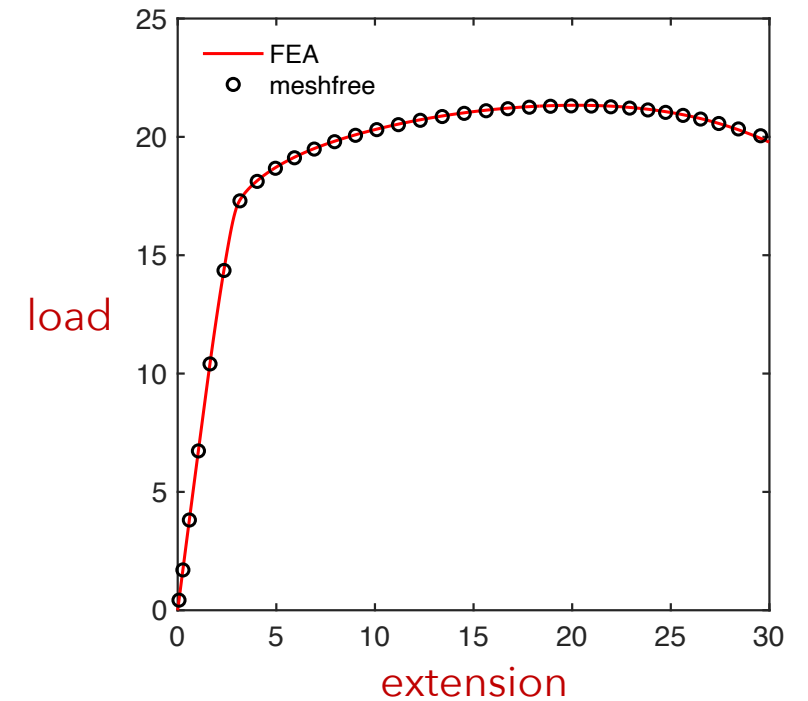
yield surface $f(\sigma, \bar{\epsilon}^p) = \phi(\sigma) - \sigma_y(\bar{\epsilon}^p) = 0$

$$\phi(\sigma) = \left\{ \frac{1}{2} (|\sigma_1 - \sigma_2|^2 + |\sigma_1 - \sigma_3|^2 + |\sigma_2 - \sigma_3|^2) \right\}^{1/2}$$

plastic strain field



load vs. extension



Summary



1. Separate domain discretization from solution discretization (fine-scale domain triangulation with coarse-scale solution discretization).
2. Example of discretization-based reduced order model.
3. Generation of meshfree weight functions using manifold geodesics.
4. New approach to quadrature for meshfree methods based on secondary basis.
5. Projected shape-function gradients using dual basis for polynomial consistency.
6. Observed optimal convergence rates for 2D elasticity.
7. Applicable to nonlinear solid mechanics too (plasticity).
8. Examples here were in H^1 , also can be extended to $H(\text{div})$ and $H(\text{curl})$.
9. Use VMS (LoD) for material interfaces and multiscale.
10. Adaptivity framework