



# Quantum-dot-coupled Sn qubits in silicon: DFT analysis of coherent electronic shuttling

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## Introduction

- Quantum computing platforms leveraging homogeneous classical-quantum interfaces offer extraordinary engineering advantages.
- Kane's proposed architecture [1] based on donor nuclear spins in silicon was realized and it exhibits coherence times of  $\mathcal{O}(\text{seconds})$  [2].
- Single-qubit gates can be precisely controlled using NMR, but Kane's exchange-coupled two-qubit gates are difficult to fabricate.
- Electron shuttling is a promising alternative, wherein two-qubit entanglement is established through the hyperfine interaction (HFI).
- Here an ancilla electron, initially entangled with one nuclear spin, is coherently transported to interact with a second nuclear spin.
- This is a feasible route toward nuclear-nuclear entanglement and coherence transfer fidelities of 99.4% were recently achieved [3].
- Donor defect sites over-bind electrons, causing electron shuttling to be intractable, so here we consider other candidate defects.
- Isoelectronic defects bind electrons loosely and their HFIs may be sufficiently strong for shuttling, as was recently shown for  $^{29}\text{Si}$  [4].
- As the HFI of  $^{29}\text{Si}$  is still quite small, in this work we explore the prospect of using other isoelectronic defects as nuclear spin qubits.

## Theory

- A *hyperfine interaction* (HFI) occurs when an unpaired electronic spin encounters a nucleus possessing a non-vanishing magnetic moment.
- The hyperfine Hamiltonian is written as  $H_{\text{hf}} = \mathbf{I} \cdot \mathbf{A} \cdot \mathbf{S}$ , with  $\mathbf{I}$  and  $\mathbf{S}$  the nuclear and electronic spin operators and  $\mathbf{A}$  the coupling tensor.
- When the electron is moved through electrostatic controls to maximize the HFI,  $\mathbf{A}$  is dominated by the *Fermi contact interaction* (FCI):

$$A \approx A_{\text{FCI}} = \frac{8\pi}{3} \gamma_e \mu_e \gamma_I \mu_I |\Psi(R_I)|^2$$

- The *bunching factor*, or the electronic density enhancement at nuclear site  $I$ , was defined by Shulman and Wyluda [5] as  $\eta = |\Psi(R_I)|^2 / \langle \Psi^2 \rangle_{\text{Av}}$ .
- Van de Waale et al. [6] later formulated  $\eta$  using spin densities ( $\rho_{\text{spin}} = \rho_{\uparrow} - \rho_{\downarrow}$ ) while Assali et al. [7] outlined a DFT-based procedure to compute it for  $^{29}\text{Si}$  as

$$\eta = \frac{\rho_{\text{spin}}(R_I)}{[\rho_{\text{spin}}]_{\text{Av}}}$$

- This may be computed using conventional Kohn-Sham (KS) DFT in 3 steps:
  - 1) Introduce the defect and perform a structural optimization to relieve strain.
  - 2) Converge the spin-polarized, spin-orbit SCF cycles for the neutral system.
  - 3) Add an electron with a positive jellium background and solve KS equation.

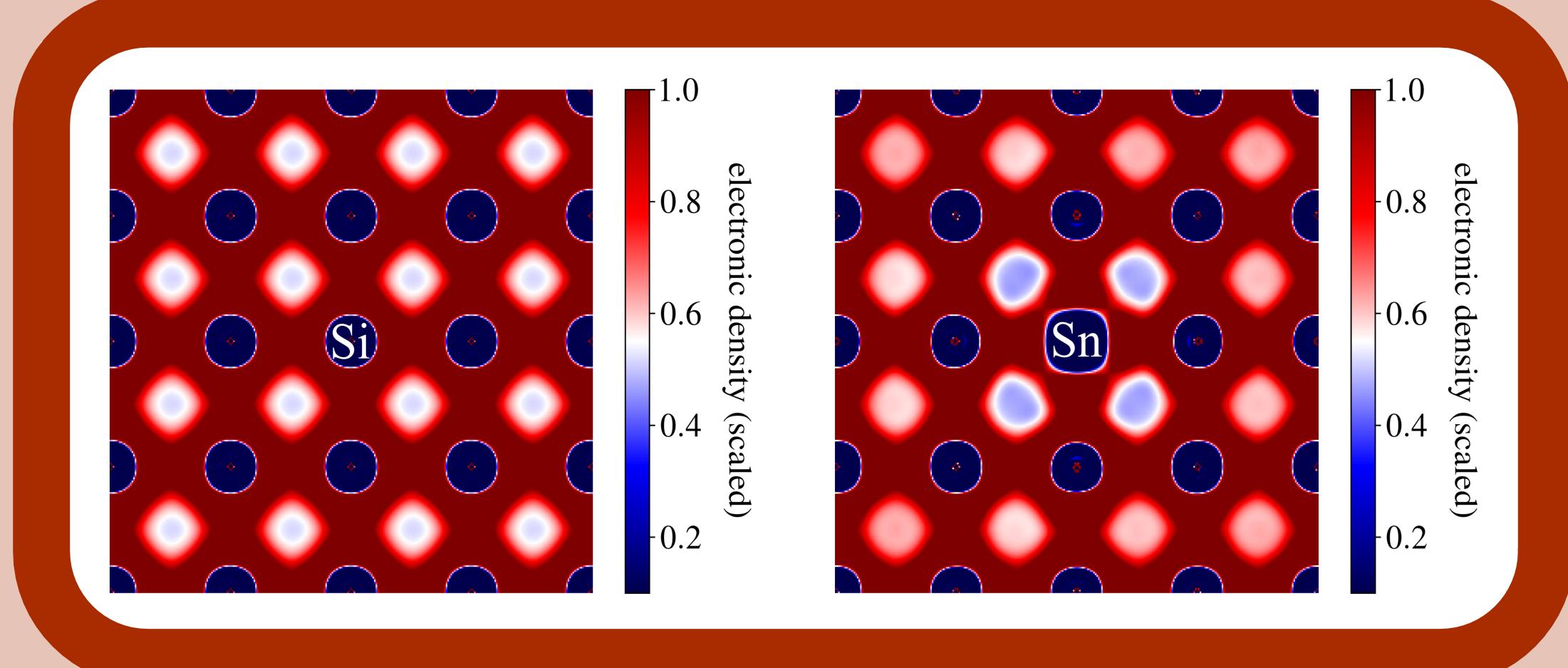
## Objective

- Simulation can facilitate a high-throughput analysis of designer defects.
- We must first develop a method capable of predicting HFI strengths and then apply it to select the optimal isoelectronic defect in silicon.

## Methods

- We performed full-potential spin-polarized DFT-PBE calculations using WIEN2k within a basis of linearized augmented plane waves plus local orbitals (LAPW+lo), while treating spin-orbit coupling in a separate variational optimization step.

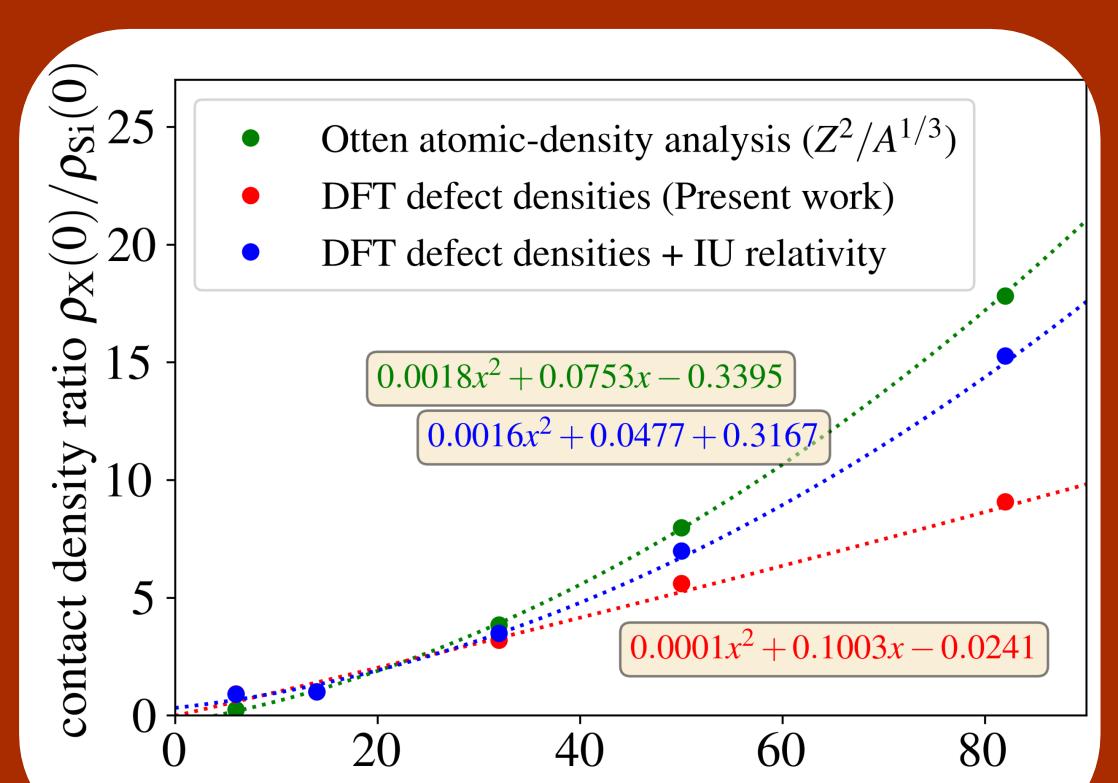
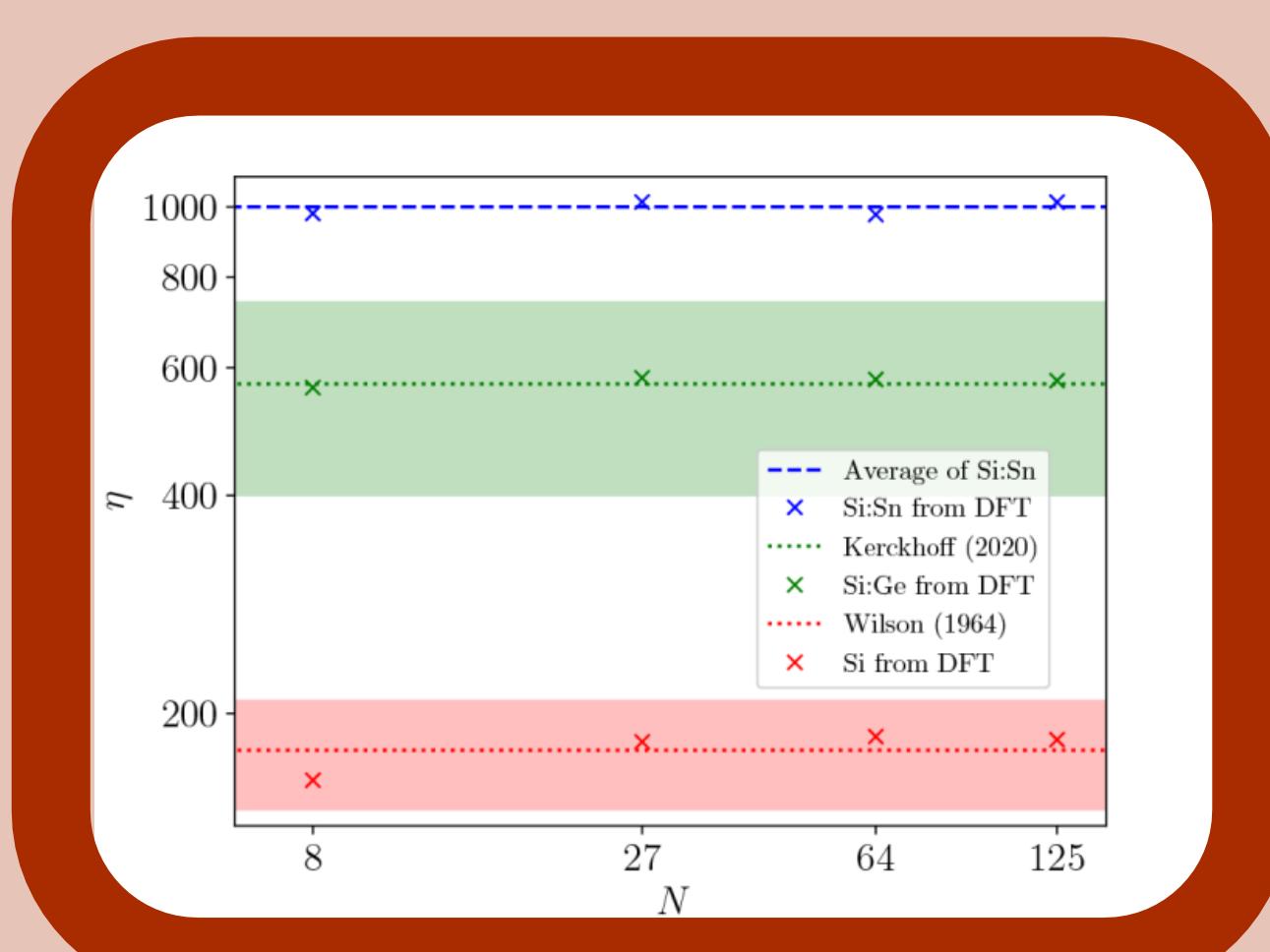
## Results and discussion



- DFT-level computation of electron spin densities shows the conduction electron will bunch slightly at a Sn site, but localization is still very weak compared to the P donor.
- Meanwhile comparing to a Si nucleus in pristine bulk, the density at the Sn site is at least five times larger, which will conveniently distinguish it from the surrounding Si atoms.
- Thus the character of extrinsic isoelectronic defects looks promising, but why do we favor Sn over C, Ge, or Pb?

## Future work

- Nuclear flip-flop dynamics simulations of Sn in silicon are currently being planned. These will utilize in-house codes which have been in development for several years.
- Experimental fabrication and validation is also underway.



- Computed values for the hyperfine enhancement ( $\eta$ ) of Si and Ge nuclei agree well with the experimental values of Wilson and Kerckhoff et al.
- $^{73}\text{Ge}$  nuclei are spin-9/2, so we opt for the simpler 1/2 spins of  $^{117}\text{Sn}$  or  $^{119}\text{Sn}$ .
- The bunching factor ratio of  $\eta_{\text{Sn}} / \eta_{\text{Si}} \sim 5$  translates to a hyperfine interaction enhancement for Sn of about 10 times with respect to the bath Si nuclei.
- Taking into account relativistic corrections to the Fermi contact interaction, the HFI enhancement for Sn grows to 13.7 times with respect to Si nuclei.

## References

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- A publication associated with this work is currently available on the arXiv at <https://doi.org/10.48550/arXiv.2206.06285>