

A Next-Generation Transport Simulator for the Waste Isolation Pilot Plant (WIPP) Performance Assessment

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INTRODUCTION

Waste Isolation Pilot Plant (WIPP) performance assessment (PA) calculations estimate the probability and consequence of potential radionuclide releases from the repository to the accessible environment for a regulatory period of 10,000 years after facility closure. The proposed replacement of waste panels in the WIPP challenges the modeling assumptions inherent in the two-dimensional (2-D) flared grid used in PA calculations of flow and transport in and around the repository. Therefore, development of a new three-dimensional (3-D) model for use in PA was warranted. At the core of this task is the development of a single, efficient, 3-D flow and transport simulator (PFLOTRAN [1]) that incorporates WIPP-specific process models to eventually replace a portion of the suite of software (i.e., BRAGFLO [2] and NUTS [3]) that is currently utilized for flow and transport in WIPP PA. The current work presents the development of the new transport mode in PFLOTRAN, named the Nuclear Waste Transport (NWT) Mode, and describes the improvements ~~over the currently utilized NUTS transport simulator for WIPP PA~~. Most notably, the NWT Mode is formulated in terms of the total bulk concentration, rather than the more common aqueous (dissolved) mass concentration, allowing it to accurately accommodate completely dry conditions without numerical difficulty.

GOVERNING EQUATIONS FOR TRANSPORT WITH EQUILIBRIUM CHEMISTRY

The Nuclear Waste Transport Mode in PFLOTRAN models the transport of a species as described by the governing equations for conservation of mass. NWT includes species advection (without diffusion), assuming equilibrium chemistry for precipitation/dissolution, and radioactive decay processes. The NWT Mode process model is executed sequentially at each time step after the full flow field solution is calculated. The primary dependent variable is the total bulk mass of species j , in phase α , denoted by M_j^α . The governing equation is given by,

$$\frac{\partial}{\partial t} \sum_{\alpha} (M_j^\alpha) + \nabla \cdot \sum_{\alpha} (\mathbf{u}^\alpha M_j^\alpha) = \sum_{\alpha} (\dot{Q}_j^\alpha) + \sum_{\alpha} (\dot{R}_j^\alpha) \quad (1)$$

where the total bulk mass of species j , in phase α , is given by,

$$M_j^\alpha = M_j^A + M_j^P \quad (2)$$

and is a sum of the aqueous (A) and precipitated (P) phases which are assumed to be in chemical equilibrium. The units of M_j^α are $\text{mol} \cdot \text{m}_{\text{bulk}}^{-3}$. Furthermore, \mathbf{u}^α is the phase velocity in units of $\text{m} \cdot \text{sec}^{-1}$. On the right-hand side, \dot{Q}_j^α is a generic source/sink term of species j in each phase (in units of $\text{mol} \cdot \text{m}_{\text{bulk}}^{-3} \cdot \text{sec}^{-1}$), and \dot{R}_j^α represents a source or sink of species j mass in each phase due to radioactive decay or ingrowth (in units of $\text{mol} \cdot \text{m}_{\text{bulk}}^{-3} \cdot \text{sec}^{-1}$).

Transport is assumed to occur only in the liquid fluid phase and not in the gas phase. Furthermore, the aqueous component of species j is the only mobile phase, while the precipitated phase is considered immobile (e.g., $\mathbf{u}^P = 0$). With these assumptions, the governing equation simplifies to,

$$\frac{\partial}{\partial t} \sum_{\alpha} (M_j^\alpha) + \nabla \cdot (\mathbf{u}^A M_j^A) = \sum_{\alpha} (\dot{Q}_j^\alpha) + \sum_{\alpha} (\dot{R}_j^\alpha) \quad (3)$$

The aqueous and precipitated phases of the total mass for each species j are defined as,

$$M_j^A = \emptyset S_{liq} C_j^A, \text{ and} \quad (4)$$

$$M_j^P = \emptyset S_{ppt} C_j^P = f(M_j^A) \quad (5)$$

where S_{liq} and S_{ppt} are the liquid and precipitant saturations (in units of $\text{m}_{liq}^3 \cdot \text{m}_{void}^{-3}$ and $\text{m}_{ppt}^3 \cdot \text{m}_{void}^{-3}$, respectively) within the pore space \emptyset (in units of $\text{m}_{void}^3 \cdot \text{m}_{bulk}^{-3}$), and C_j^A and C_j^P are the aqueous and precipitated phase concentrations, respectively, of each species j (in units of $\text{mol}_j \cdot \text{m}_{liq}^{-3}$ and $\text{mol}_j \cdot \text{m}_{ppt}^{-3}$, respectively).

The aqueous concentration for each species is calculated by first disbursing the total species mass within the pore space for each grid cell. Immediately after, the resulting aqueous concentration for each species is compared to the species solubility limit. If it is below the solubility limit, the aqueous concentration calculated remains unchanged. If instead it is above the solubility limit, the aqueous concentration is set to the species solubility limit and the

amount of species mass in excess of the solubility limit is assigned as the precipitated mass.

With these definitions, the governing equation expands to,

$$\frac{\partial}{\partial t} \Sigma_{\alpha} (M_j^{\alpha}) + \nabla \cdot (\mathbf{q}^A S_{liq} C_j^A) = \Sigma_{\alpha} (\dot{Q}_j^{\alpha}) + \Sigma_{\alpha} (\dot{R}_j^{\alpha}) \quad (6)$$

where \mathbf{q}^A is the liquid Darcy flux (in units of $m_{liq}^3 \cdot m_{bulk}^{-2} \cdot sec^{-1}$) and has the relationship $\mathbf{q}^A = \mathbf{u}^A \phi$.

The generic source/sink term, \dot{Q}_j^{α} , is defined by,

$$\dot{Q}_j^{\alpha} = \frac{U_j^{\alpha}}{\forall} M_j^{\alpha} = \frac{U_j^A}{\forall} M_j^A = \frac{Q_j^A}{\forall} S_{liq} C_j^A \quad (7)$$

where U_j^{α} is a volumetric flow of species j in each phase (in units of $m_{bulk}^3 \cdot sec^{-1}$). Because the only mobile phase is the aqueous phase, the volumetric flow of each species is defined entirely by U_j^A , the volumetric flow of each species in the aqueous phase. On the right-hand side of Eq. (7), we define the generic source/sink term using Q_j^A , the volumetric Darcy flux (in units of $m_{liq}^3 \cdot sec^{-1}$), and \forall , which is a volume. In this formulation, if the term in a source, C_j^A represents the aqueous phase concentration of species j in the fluid source. On the other hand, if the term is a sink, C_j^A represents the aqueous phase concentration of species j in the domain at the location of the sink.

The reaction term, \dot{R}_j^{α} , consists of decay and ingrowth of radioactive isotope species. Radioactive decay and ingrowth is defined by the Bateman equations,

$$\dot{R}_j^{\alpha} = \frac{\partial M_j^{\alpha}}{\partial t} = -\lambda_j M_j^{\alpha} + \lambda_p M_p^{\alpha} \quad (8)$$

where the subscript j represents the radioactive isotope species, and the subscript p represents the radioactive parent of species j . The radioactive decay rate (in units of sec^{-1}) is given by λ . Decay and ingrowth is calculated for the total bulk mass, which is then distributed across the aqueous and precipitated phases after the Bateman equations are solved.

NUMERICAL METHODS FOR THE SOLUTION OF THE TRANSPORT GOVERNING EQUATIONS

The governing equation described by Eq. (1) is a set of equations in space and time. The set of equations is solved by first discretizing them using the finite volume method. The finite volume method uses a volume integral formulation of the set of equations with a finite partitioning set of volumes to discretize the equations in space and time. The expression in Eq. (6) becomes,

$$\int_V \left[\int_t \left[\frac{\partial}{\partial t} \Sigma_{\alpha} (M_j^{\alpha}) + \nabla \cdot (\mathbf{q}^A S_{liq} C_j^A) \right] dt \right] dV =$$

$$\int_V \left[\int_t \left[\Sigma_{\alpha} (\dot{Q}_j^{\alpha}) + \Sigma_{\alpha} (\dot{R}_j^{\alpha}) \right] dt \right] dV \quad (9)$$

which integrates the governing equation in space and time. Upon taking these integrals and applying the divergence theorem (noting several mathematical details omitted for brevity), the discretized equation becomes,

$$\begin{aligned} \frac{\partial \Sigma_{\alpha} (M_j^{\alpha})}{\partial t} \forall + \Delta t \Sigma_F (\mathbf{q}^A S_{liq} C_j^A \cdot \mathbf{n} A_F) = \\ \Delta t \forall \left[\Sigma_{\alpha} (\dot{Q}_j^{\alpha}) + \Sigma_{\alpha} (\dot{R}_j^{\alpha}) \right] \end{aligned} \quad (10)$$

where Δt is the time step duration (in units of sec). The summation term, $\Sigma_F (\mathbf{q}^A S_{liq} C_j^A \cdot \mathbf{n} A_F)$, results from a surface integration on the discretized volume, \forall . This surface integral is discretized as a sum over a finite number of faces, where \mathbf{n} is the unit normal vector for each face and A_F is the area of each face. In the case of a 3-D structured rectilinear grid, each volume represents a grid cell that has 6 faces.

The accumulation term (first term on the left-hand side of Eq. (10)) is discretized in time using a forward difference, and all other terms are taken at the new time level, which is described as an implicit backward Euler discretization method (as opposed to an explicit standard Euler discretization where all other terms are taken at the current time level). The implicit time discretization method has the advantage of numerical stability for any time step taken but requires more computational effort to solve than an explicit method.

The Newton-Krylov iteration method is used to solve the discretized governing equations in space and time. First, Eq. (10) is transformed into a residual equation by putting all terms on one side of the equation and setting their sum to zero, as expressed by,

$$R(M_j^{\alpha, t+1}) = 0 \quad (11)$$

This method states that a generic system of non-linear equations, $f_i(x_1, x_2, \dots, x_n) = 0$, can be expressed as a new system of linear equations, described by,

$$\begin{aligned} f_i(x_1, x_2, \dots, x_n) = \\ f_i(x_1^k, x_2^k, \dots, x_n^k) + \sum_{j=1}^n (x_j^{k+1} - x_j^k) \frac{\partial f_i(x_1^k, x_2^k, \dots, x_n^k)}{\partial x_j} = 0 \end{aligned} \quad (12)$$

of the form $Ax = b$. The matrix A is called the Jacobian and is the set of partial derivatives of the discretized residual equation (i.e., Eq. (10)) with respect to the primary variable, M_j^{α} . In the NWT Mode, analytical derivatives are defined and computed. The right-hand side, b , is the residual at the k^{th} iteration. Finally, x is the solution vector update, thus giving,

$$\sum_{j=1}^n \left(\frac{\partial R_i(M_j^{\alpha, t+1})}{\partial M_j^{\alpha}} (\delta M_j^{\alpha, t+1}) \right) = -R(M_j^{\alpha, t+1})^k \quad (13)$$

187
188 The solution is then updated after each iteration
189 according to,
190

$$191 \quad (M_j^{\alpha,t+1})^{k+1} = (M_j^{\alpha,t+1})^k + (\delta M_j^{\alpha,t+1}) \quad (14)$$

192
193 and the updated solution is used to evaluate the residual
194 equation. This process is repeated until the residual value
195 zero. However, because the residual will never truly be zero
196 convergence is declared according to a variety of tolerances
197 on $\delta M_j^{\alpha,t+1}$ and the residual value.

198 The specific convergence criteria used in NWT Mode
199 include the infinity norm of the absolute residual value, the
200 infinity norm of the scaled residual value, and the infinity
201 norm of the relative solution update. The scaled residual
202 value is the proportion of the residual relative to the
203 accumulation term (i.e., first term on the left-hand side of Eq.
204 (10)), described by
205

$$206 \quad \frac{R(M_j^{\alpha,t+1})}{(\sum_{\alpha} (M_j^{\alpha,t}) \forall)}. \quad (15)$$

207
208 The relative solution update is the proportion of the solution
209 update relative to the previous solution, described by
210

$$211 \quad \delta M_j^{\alpha,t+1} / M_j^{\alpha,t}. \quad (16)$$

212
213 Convergence is defined as meeting the absolute OR scaled
214 residual criteria, AND the relative solution update criteria
215 each grid cell (e.g., infinity norm).

216 An uncommon and very advantageous feature in the
217 NWT Mode of PFLOTTRAN is the ability to define
218 convergence criteria for each individual species. This
219 useful when there are large differences (e.g., several orders
220 of magnitude) in the amount of mass between species, which
221 may result from differences in the initial inventories
222 differences in solubility limits, or differences in the
223 radioactive decay constants between species. Whether
224 result of any of these reasons, when large differences in the
225 amount of mass between species exists, it may not be
226 appropriate to assign a single value for the convergence
227 criteria to all species. Having the capability to define species
228 specific convergence criteria values allows the numerical
229 solution to converge with less error.

230 231 **NWT MODE IMPROVEMENTS OVER PREVIOUS** 232 **SOFTWARE FOR WIPP PA CALCULATIONS**

233
234 Several NWT Mode features represent improvements for
235 WIPP PA calculations, over the previous software, NUTS.

236 First, the NWT Mode is formulated in terms of the total
237 bulk concentration (moles per bulk cubic meter), allowing
238 to accurately accommodate completely dry conditions
239 without numerical difficulty. Formulations in terms of the

aqueous concentration, ~~(a more common formulation as is the~~
~~ease with NUTS), cannot naturally handle completely dry~~
~~conditions because the aqueous concentration becomes~~
~~undefined, and the system of equations cannot be solved for~~
~~the aqueous concentration. In such cases, modification to the~~
~~governing equations is done to avoid mathematical~~
~~singularity at dry-out conditions.~~

Second, NWT Mode uses finite volume discretization
and solves the discretized equations with a fully implicit,
backward Euler approach based on Newton-Krylov iteration.
The finite volume discretization conserves mass, by
definition, ~~as opposed to the explicit, finite difference method~~
~~that is used by NUTS.~~ Furthermore, using a backward Euler
approach based on Newton-Krylov iteration allows
numerical stability for even large time steps, and can handle
non-linearities in the problem formulation. The uncommon
capability allowing species-specific assignment of
convergence criteria in NWT Mode reduces numerical error
when multiple species with disparate inventories are
simulated and is a state-of-the-art feature that is lacking in
many other transport simulators.

Third, in NWT Mode, the transport process model is
solved sequentially after fluid flow at each time step. ~~This is~~
~~an improvement over NUTS in PA calculations, which solves~~
~~for transport after 55 years of the flow solution has been~~
~~solved.~~ Solving for the transport solution at each flow time
step, ~~as done by NWT Mode,~~ reduces operator-splitting
error, producing a more accurate transport solution.
Moreover, in NWT Mode, the transport process model is
allowed to sub-step the flow solution in time in order to
satisfy convergence criteria and tolerances. This additional
feature reduces time truncation error in the numerical
solution and is automatically activated when the transport
solution cannot be accurately solved at as large of a time step
duration as the flow solution was (due to complex flow
regimes, for example).

~~Fourth, the NWT Mode can be sequentially coupled to a~~
~~new wellbore model (currently under development by the~~
~~authors) to additionally simulate transport in a wellbore~~
~~emplaced within a larger reservoir. In the previous WIPP PA~~
~~software architecture, BRAGFLO explicitly meshes the~~
~~wellbore in the computational domain and uses a flared grid~~
~~geometry to represent 3-D space on a 2-D grid. In the~~
~~development of a new 3-D model for use in PA, explicitly~~
~~meshing the borehole in the computational domain quickly~~
~~becomes computational intractable due to the large number~~
~~of grid cells required to do this with good grid quality. Using~~
~~a wellbore model rather than explicitly meshing the wellbore,~~
~~therefore, significantly reduces the computational expense.~~
~~The flexibility of the NWT Mode to plug into updated~~
~~features of the sequentially coupled flow mode, such the new~~
~~wellbore model, significantly streamlines the software~~
~~development process and allows PA calculations to remain~~
~~practical in 3-D.~~

Finally, because it is part of PFLOTTRAN, the NWT
Mode inherits its massively parallel computing capability,

which greatly speeds up 3-D simulations. In the previous PAFLOTRAN, which NUTS was limited to serial-2-D computations, serial computations were adequate. However, serial computations must move towards parallelization practically inhibiting its use for the new 3-D PA model.

CONCLUSIONS AND FUTURE WORK

Presented here is the development of the new 3-D transport mode in PFLOTRAN, named the Nuclear Waste Transport (NWT) Mode. Its development was undertaken as a result of the proposed replacement of waste panels in the WIPP, which challenges the modeling assumptions inherent in the 2-D flared grid used in PA calculations of flow and transport in and around the repository. The NWT Mode is part of a single, efficient, 3-D flow and transport simulator (PFLOTRAN) that incorporates WIPP-specific processes and models that will eventually replace a portion of the suite of software (i.e., BRAGFLO and NUTS) that is currently utilized for flow and transport in WIPP PA.

Besides its role in the updated WIPP PA calculations, the new NWT Mode in PFLOTRAN will also be broadly applicable for the PA community outside of WIPP and has already shown applicability in transport calculations which have been traditionally difficult to solve numerically, due to large non-linearities in the system of equations or complex partially fluid-saturated regimes.

NOMENCLATURE

α = phase index (aqueous or precipitated)
 j = species index
 p = species' parent index
 M_j^α = total bulk mass of species j , in phase α
 M_j^A = total bulk mass of species j , in aqueous phase
 M_j^P = total bulk mass of species j , in precipitated phase
 q^A = liquid Darcy flux
 u^α = phase velocity
 u^A = aqueous phase velocity
 u^P = precipitated phase velocity, assumed zero
 U_j^α = volumetric flow of species j , in phase α
 U_j^A = volumetric flow of species j , in aqueous phase
 C_j^A = aqueous phase concentration of species j
 C_j^P = precipitated phase concentration of species j
 \dot{Q}_j^α = generic source/sink of species j , in phase α
 \dot{R}_j^α = radioactive source/sink of species j , in phase α
 \emptyset = pore space
 S_{liq} = liquid saturation
 S_{ppt} = precipitant saturation
 \forall = discretized volume

A_F = area of each discretized volume face
 \mathbf{n} = normal vector of each discretized volume face
 $t, \Delta t$ = time, time step
 m = meter
 mol = mole
 sec = second

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