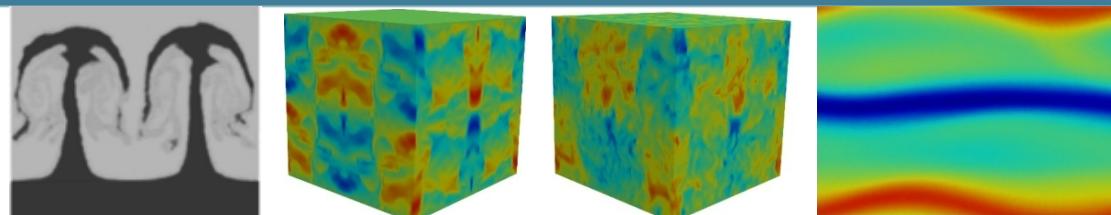
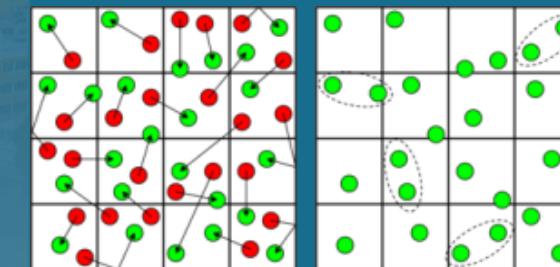




Sandia  
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# The Navier-Stokes Equations Do Not Describe the Smallest Scales of Turbulence



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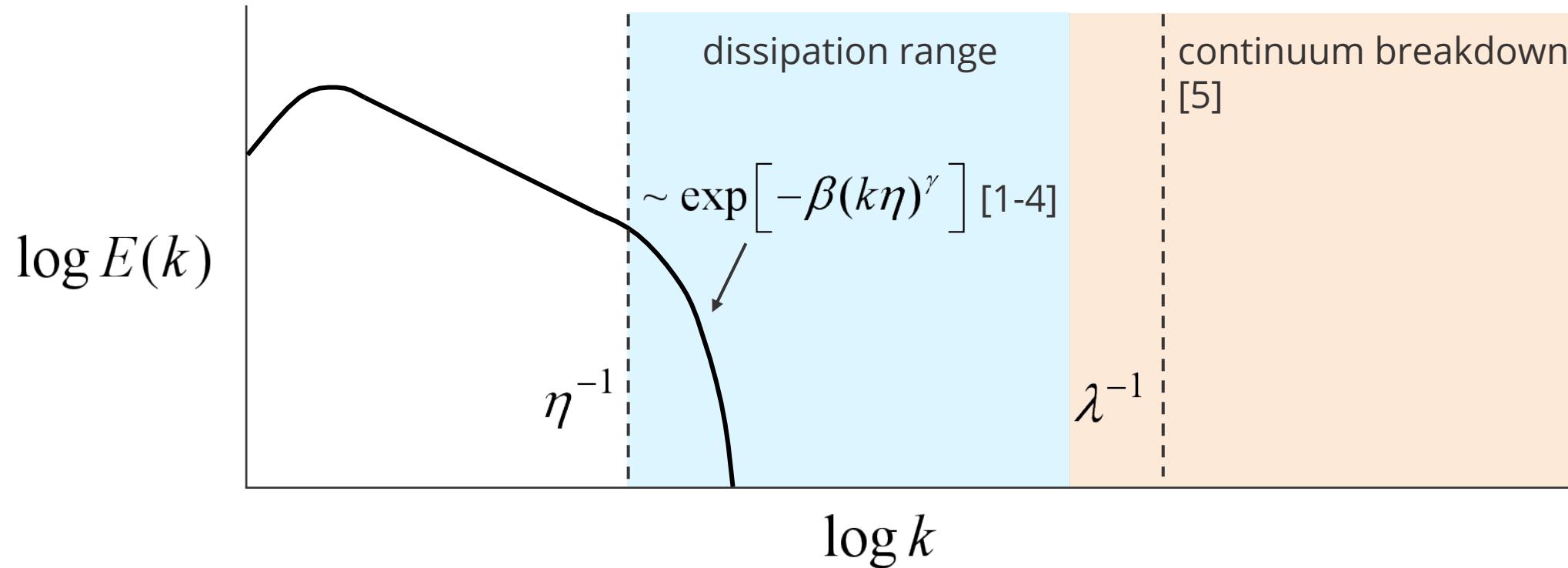


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# The turbulent energy spectrum according to Navier-Stokes



[1] Chen et al., Phys. Rev. Lett. (1993)

[2] Sirovich et al., Phys. Rev. Lett. (1994)

[3] Khurshid et al., Phys. Rev. Fluids (2018)

[4] Buaria & Sreenivasan, Phys. Rev. Fluids (2020)

[5] Bird, Clarendon Press (1994)

# What about thermal fluctuations?



All finite systems exhibit random molecular fluctuations

Center-of-mass velocity  $u_{\text{th}}$  in volume  $l^3$ :

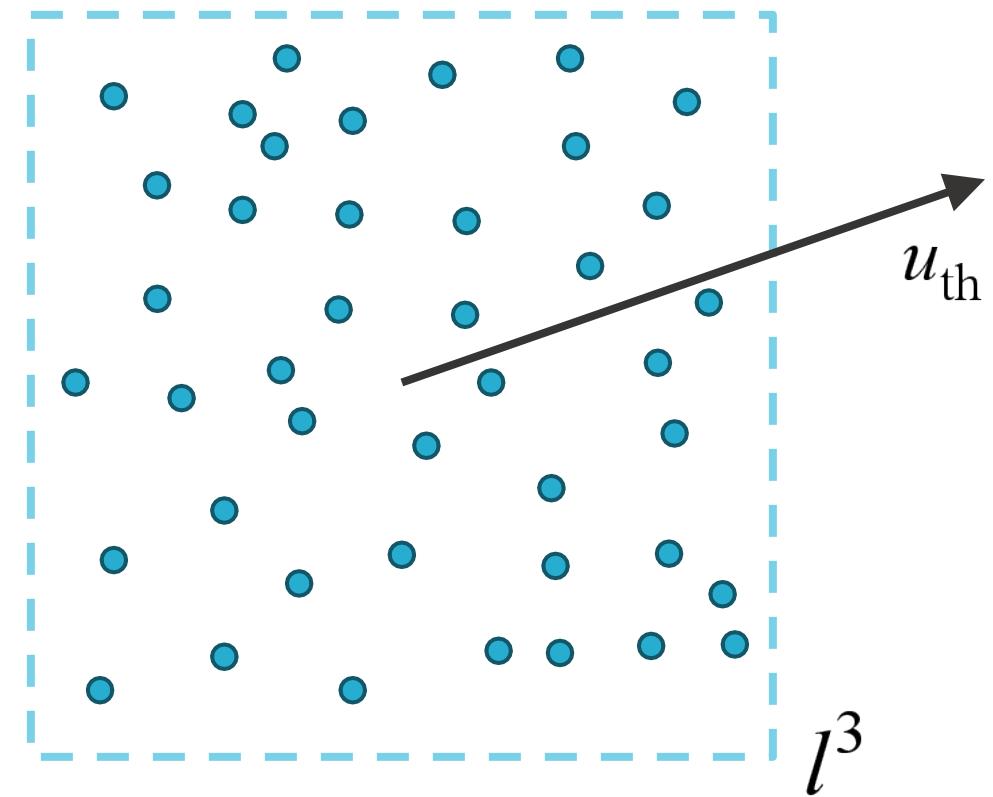
$$\langle u_{\text{th}} \rangle = 0, \quad \langle u_{\text{th}}^2 \rangle \sim \frac{k_B T}{\rho l^3}$$

Compare to kinetic energy in Kolmogorov-scale eddy [1]

$$\theta_\eta \equiv \frac{k_B T}{\rho u_\eta^2 \eta^3} \sim 10^{-9} - 10^{-6}$$

But  $E(k)$  decays exponentially fast for  $k\eta > 1$ ...

Thermal fluctuations may dominate dissipation range,  
even when  $k \ll \lambda^{-1}$  [1,2]



[1] Eyink et al., Phys. Rev. E (2022)

[2] Betchov, J. Fluid Mech. (1957)

# Molecular-level simulations of turbulence



Dissipation range is extremely difficult to measure experimentally

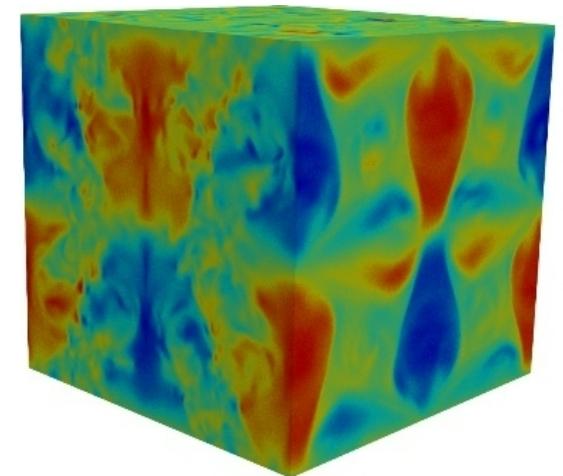
But can investigate with molecular-level simulations

Simulate Taylor-Green (TG) vortex flow [1] using DSMC

- $Re = \rho_0 VL / \mu_0 = 500, 1000, 1500$
- $Ma = V/a_0 = 0.3, 0.6, 0.9$
- $Kn = \lambda/L = 9.7 \times 10^{-4}$

Compare with compressible Navier-Stokes (NS) simulations

- Use Sandia's finite-volume code SPARC [2]



[1] Taylor & Green, Proc. R. Soc. Lond. A (1937)

[2] Howard et al., 23<sup>rd</sup> AIAA CFD (2017)

# Direct simulation Monte Carlo (DSMC)



No PDEs solved – tracks very large numbers ( $\sim 10^{12}$ ) of computational “molecules”

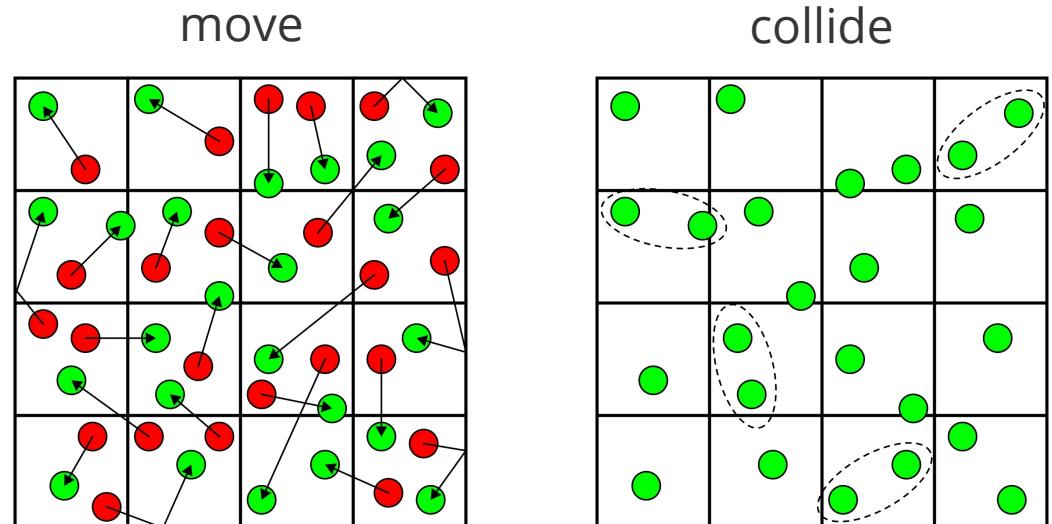
- Each represents  $F \gg 1$  actual molecules
- Move ballistically, collide & reflect stochastically
- Flow quantities from averages over molecules in each cell

Inherently includes physics usually not in NS

- Thermal and chemical nonequilibrium
- Pressure and heat-flux tensor anisotropy
- Thermal fluctuations

Simulates gas flows very accurately

- Converges to solutions of the Boltzmann Equation [1]
- Reproduces Chapman-Enskog distribution [2]



$$\frac{d\mathbf{x}_i}{dt} = \mathbf{v}_i$$

$$\frac{d(m_i \mathbf{v}_i)}{dt} = \mathbf{F}(\mathbf{x}_i) + \mathbf{C}(\mathbf{v}_i)$$

[1] Wagner, J. Stat. Phys. (1992)

[2] Gallis et al., Phys. Rev. E (2004)

# SPARTA: An exascale DSMC code



SPARTA: Stochastic PArallel Rarefied-gas Time-accurate Analyzer

Implementation is similar to Molecular Dynamics (MD)

- Single-processor to massively-parallel platforms
- Load balancing, in-situ visualization, on-the-fly FFTs, adaptive grid

Extremely complicated geometries can be treated

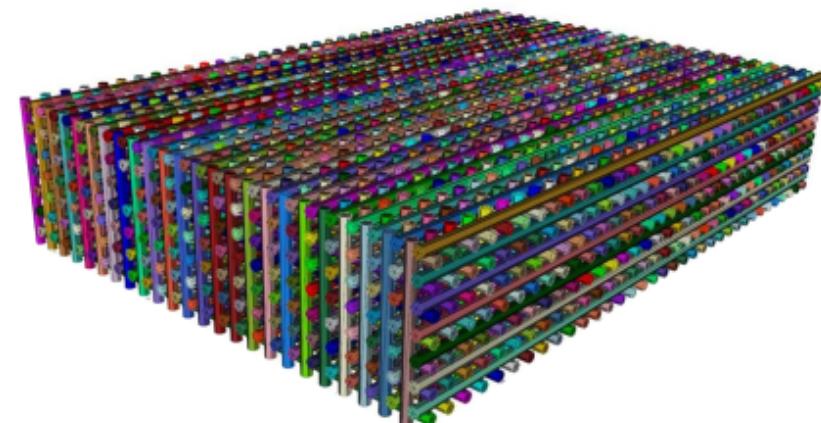
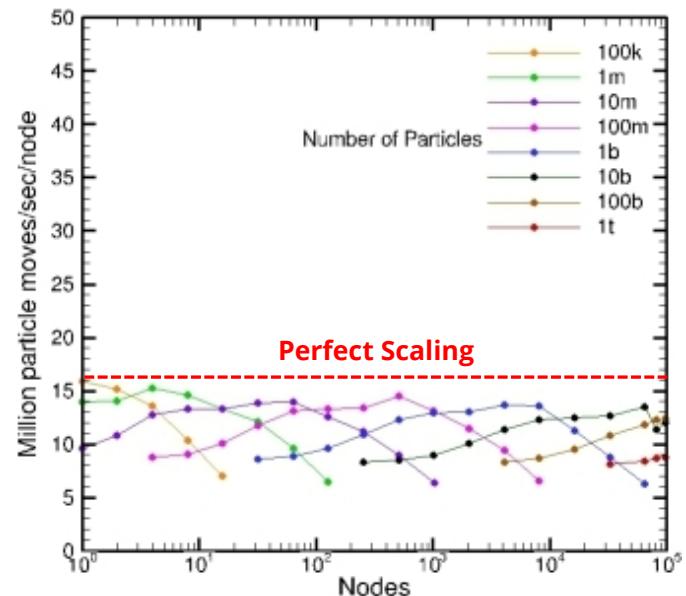
- Domain can be 2D, axisymmetric, 3D
- Gas molecules use hierarchical Cartesian “ijk” grid
- Body surfaces represented by triangular elements which cut gas grid cells

Developed with next-generation architectures in mind

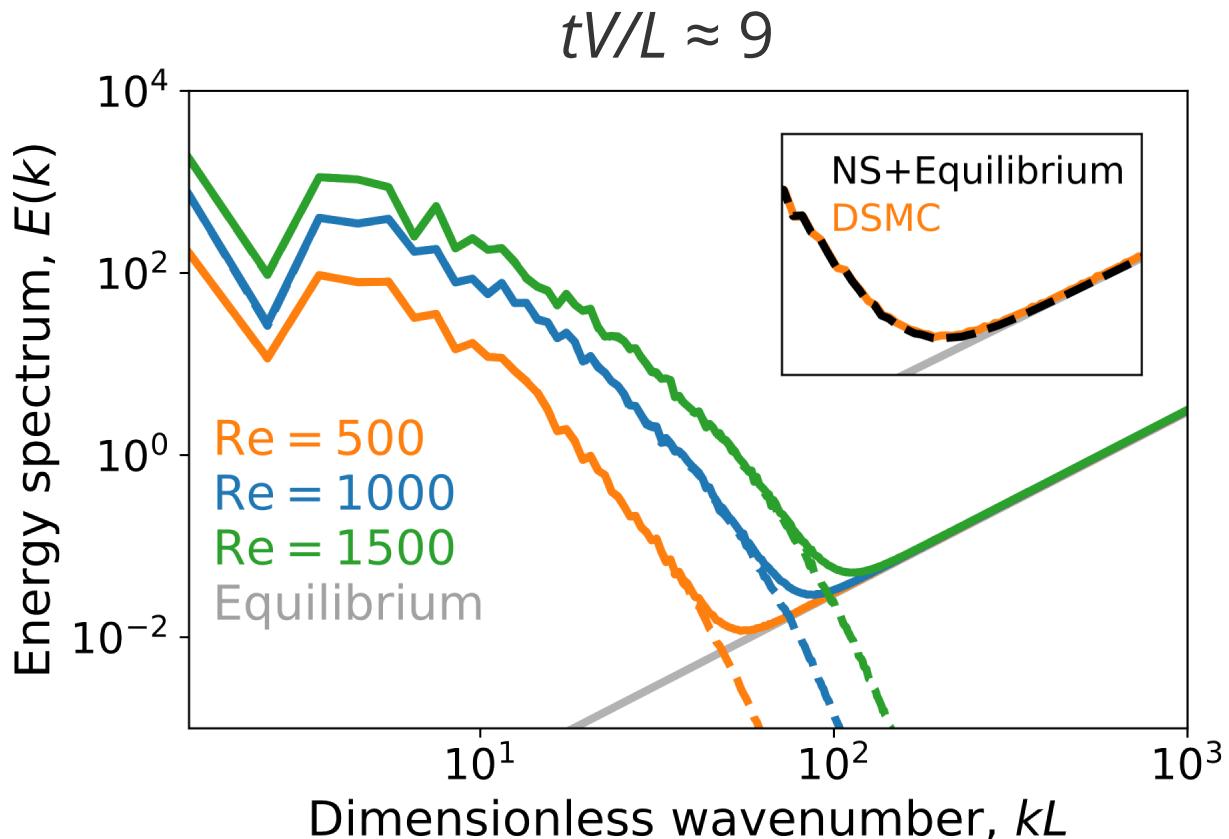
- Write application kernels only once
- Efficient on many platforms: GPU, manycore, heterogeneous, ...

Open-source code available: <http://sparta.sandia.gov>

- 10,000+ downloads, 100+ verified users worldwide
- Collaborators: ORNL, LANL, ANL, LBNL, NASA, Purdue, UIUC, ESA



# DSMC and NS spectra comparison



Excellent agreement for low  $k$

DSMC spectra show large- $k$  departure from NS due to thermal fluctuations

NS equations are inaccurate for  $k > k_c$

Fluctuation variance overestimated in DSMC when  $F > 1$

- These simulations use  $F = 16,154$
- Want to determine the  $k_c$  that would be observed in a physical gas

$$\text{Equilibrium spectrum: } E(k) = F \frac{3k_B T}{4\pi^2 \rho} k^2$$

# Equilibrium DSMC spectra



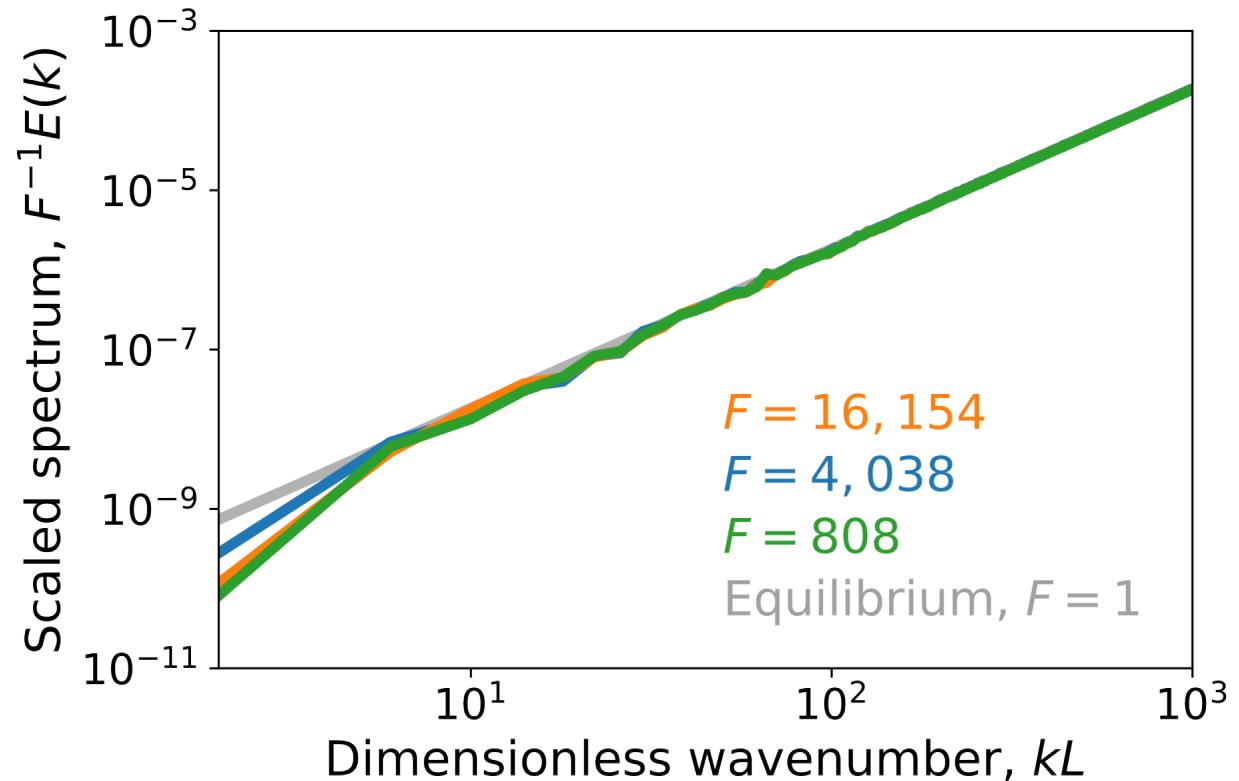
Equilibrium spectrum:

$$E(k) = F \frac{3k_B T}{4\pi^2 \rho} k^2$$

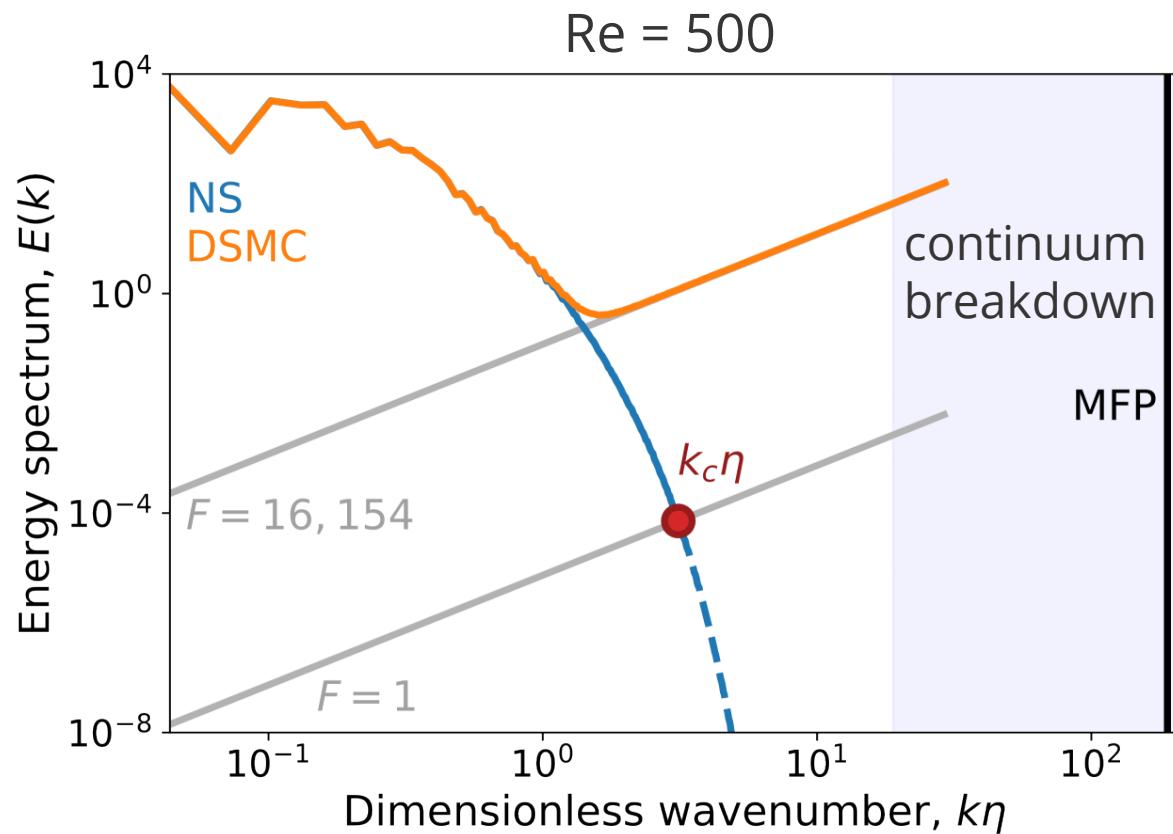
DSMC spectra obey simulation ratio scaling

$F = 1$  corresponds to physical gas

- Use this to determine  $k_c$



# Crossover wavenumber



$$k_c\eta \approx 3.1$$

Thermal fluctuations dominate almost the entire dissipation range

- Similar for other Re
- Agrees with fluctuating hydrodynamics predictions for liquids [1]

$$l_c / \lambda \approx 61$$

Crossover scale is in a regime where NS equations widely believed to be valid [2]

[1] Bell et al., J. Fluid Mech. (2022)

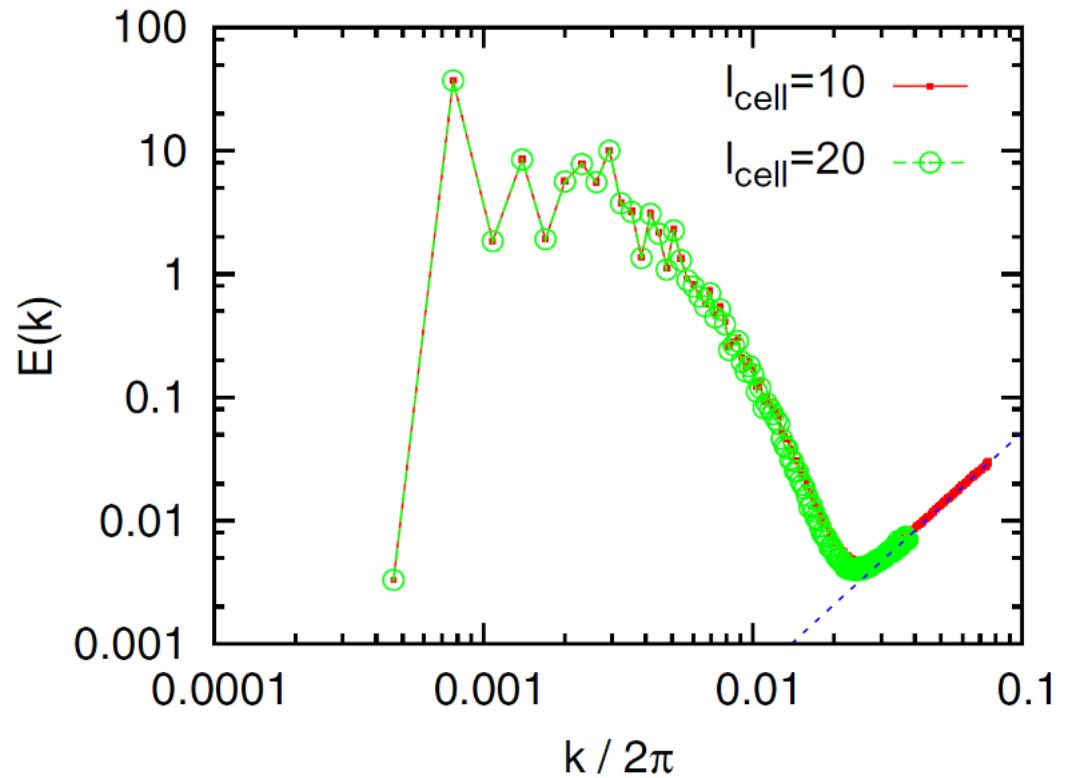
[2] Bird, Clarendon Press (1994)

# Molecular-dynamics simulations



Komatsu et al. [1] observed thermal fluctuations in MD simulations of TG flow

- Crossover at  $k\eta \approx 1$
- But  $\eta \approx 7$  molecular diameters ( $\sim \lambda$ )
  - $\theta_\eta \approx 0.5$
  - Present DSMC:  $\theta_\eta \sim 10^{-6}$
- Concluded observable only due to microscopic system size



# Thermal-fluctuation effects on other statistics



Equilibrium spectrum for  $k > k_c$  suggests simple model for velocity

$$u_i = u_i^{\text{NS}} + u_i^{\text{th}},$$

$$u_i^{\text{th}} \sim N \left( 0, Fk_B T / \rho l^3 \right)$$

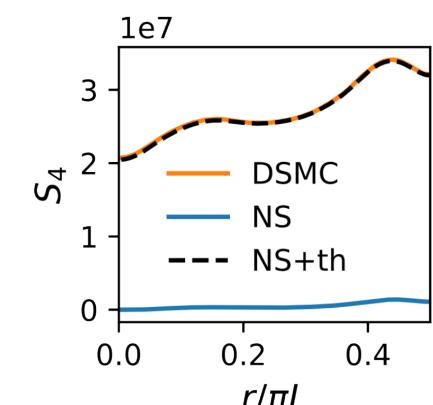
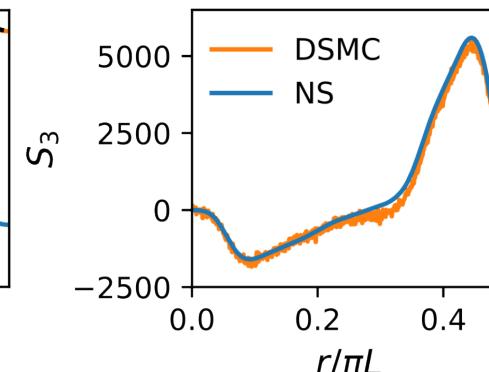
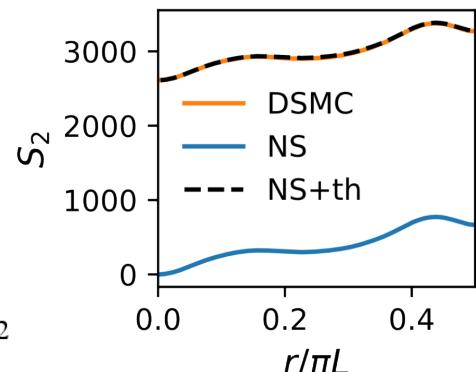
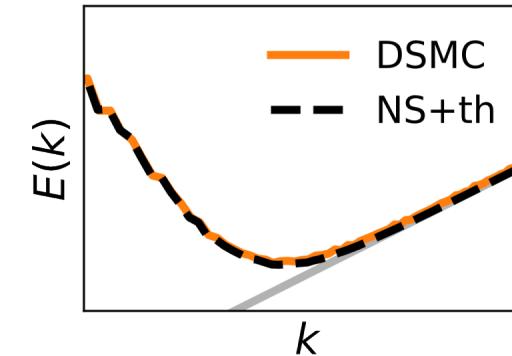
Use to look at other statistics, e.g. longitudinal structure functions

$$S_n(r) = \langle [u(x+r) - u(x)]^n \rangle$$

$$S_2(r) = S_2^{\text{NS}}(r) + 2 \frac{Fk_B T}{\rho l^3},$$

$$S_3(r) = S_3^{\text{NS}}(r),$$

$$S_4(r) = S_4^{\text{NS}}(r) + 12 \frac{Fk_B T}{\rho l^3} S_2^{\text{NS}}(r) + 12 \left( \frac{Fk_B T}{\rho l^3} \right)^2$$



# Structure function predictions for $F = 1$



Separate into NS and thermal correction:

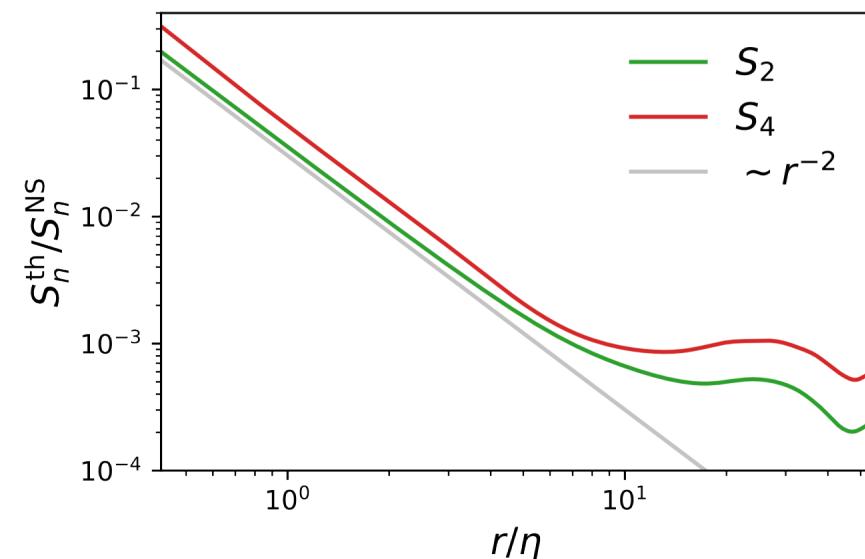
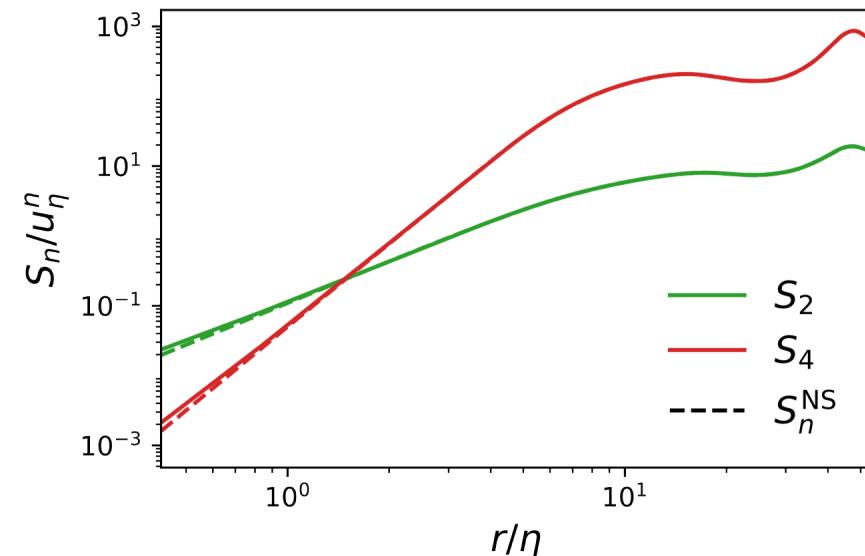
$$S_2(r) = S_2^{\text{NS}}(r) + 2 \underbrace{\frac{k_B T}{\rho l^3}}_{S_2^{\text{th}}}$$

$$S_4(r) = S_4^{\text{NS}}(r) + 12 \underbrace{\frac{k_B T}{\rho l^3} S_2^{\text{NS}}(r) + 12 \left( \frac{k_B T}{\rho l^3} \right)^2}_{S_4^{\text{th}}}$$

Significant deviations ( $\sim 5\%$ ) at  $r = \eta$

- Grow like  $r^{-2}$
- Thermal fluctuations dominate for separations smaller than

$$\frac{r^*}{\eta} \sim \theta_\eta^{1/2} \left( \frac{l}{\eta} \right)^{-3/2}$$



# Thermal-fluctuation effects on larger scales

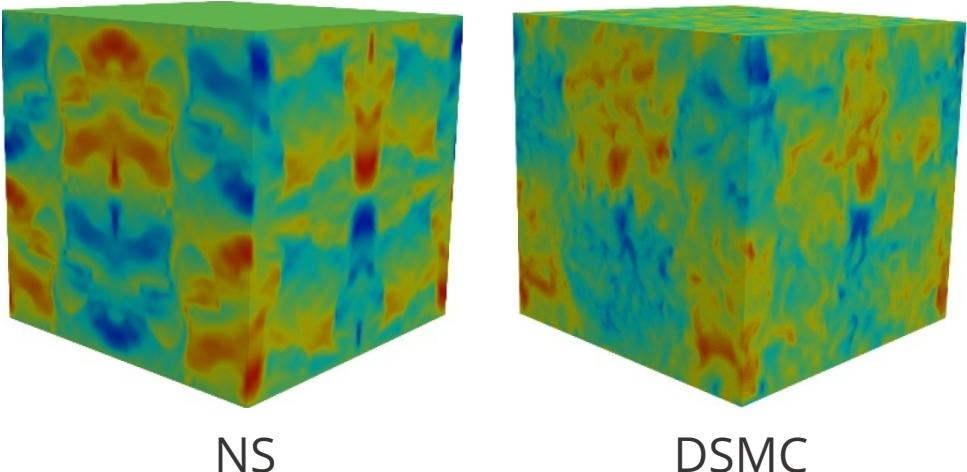


No observed effects on global or large-scale statistics

However, they do influence realizations of larger scales [1,2]

Implications for predictability?

- Maximal Lyapunov exponent? [1,3,4]
  - $\delta(t) \sim \exp(\text{Re}^\beta t)$
- Superfast amplification? [5]
  - $\delta(t) \sim \exp(C \text{Re}^{1/2} t^{1/2} + C_1 t)$
- Spontaneous stochasticity? [6-9]
  - “Intrinsic randomness”



NS

DSMC

[1] Ruelle, Phys. Lett. A (1979)

[2] Gallis et al. Phys. Rev. Fluids (2021)

[3] Boffetta & Musacchio, Phys. Rev. Lett. (2017)

[4] Berera & Ho, Phys. Rev. Lett. (2018)

[5] Li et al., J. Fluid Mech. (2020)

[6] Lorenz, Tellus (1969)

[7] Kupiainen, Ann. Henri Poincaré (2003)

[8] Eyink & Bandak, Phys. Rev. Res. (2020)

[9] Thalabard et al., Comms. Phys. (2020)

# Summary

R. M. McMullen, M. C. Krygier, J. R. Torczynski, and M. A. Gallis,  
 Phys. Rev. Lett. **128**, 114501 (2022).

DSMC shows that the Navier-Stokes equations are not accurate in the dissipation range for turbulence in gases

$$k_c \eta = O(1),$$

$$l_c / \lambda \approx 1$$

Agrees with previous estimates [1,2] fluctuating hydrodynamics simulations of liquids [3], and microscopic MD simulations [4]

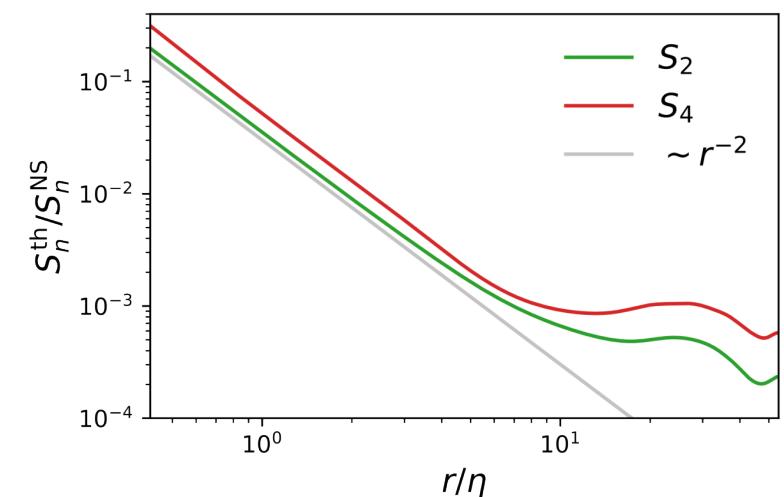
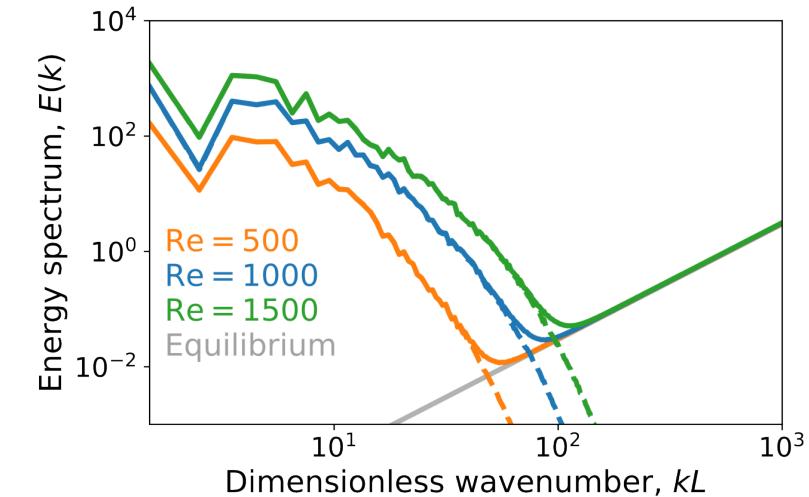
Simple model for velocity field reproduces DSMC structure functions

- Predicts significant modifications below Kolmogorov scale

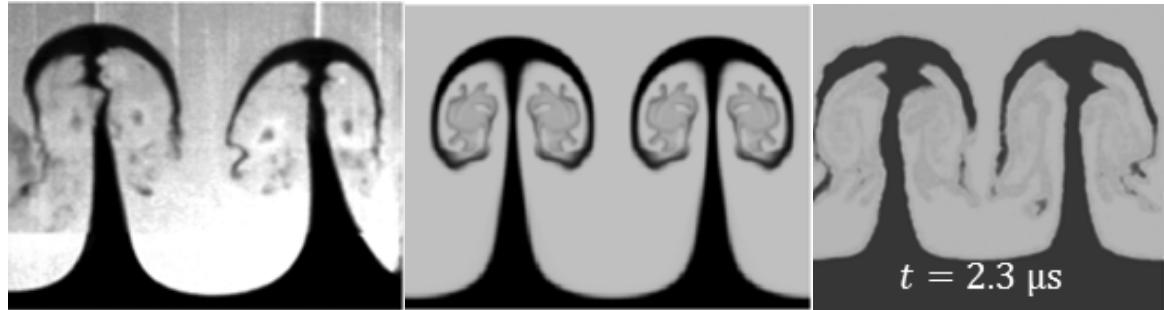
What role do thermal fluctuations play in predictability?

[1] Betchov, J. Fluid Mech. (1957)  
 [2] Eyink et al., Phys. Rev. E (2022)

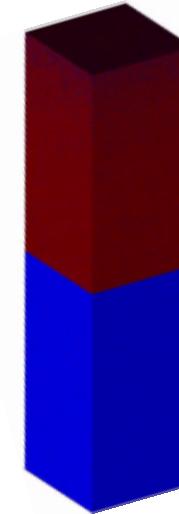
[3] Bell et al., J. Fluid Mech (2022)  
 [4] Komatsu et al., Int. J. Mod. Phys. C (2014)



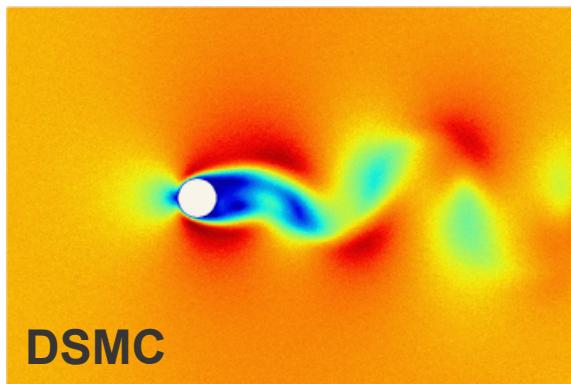
# DSMC simulations of other near-continuum flows



Richtmyer-Meshkov

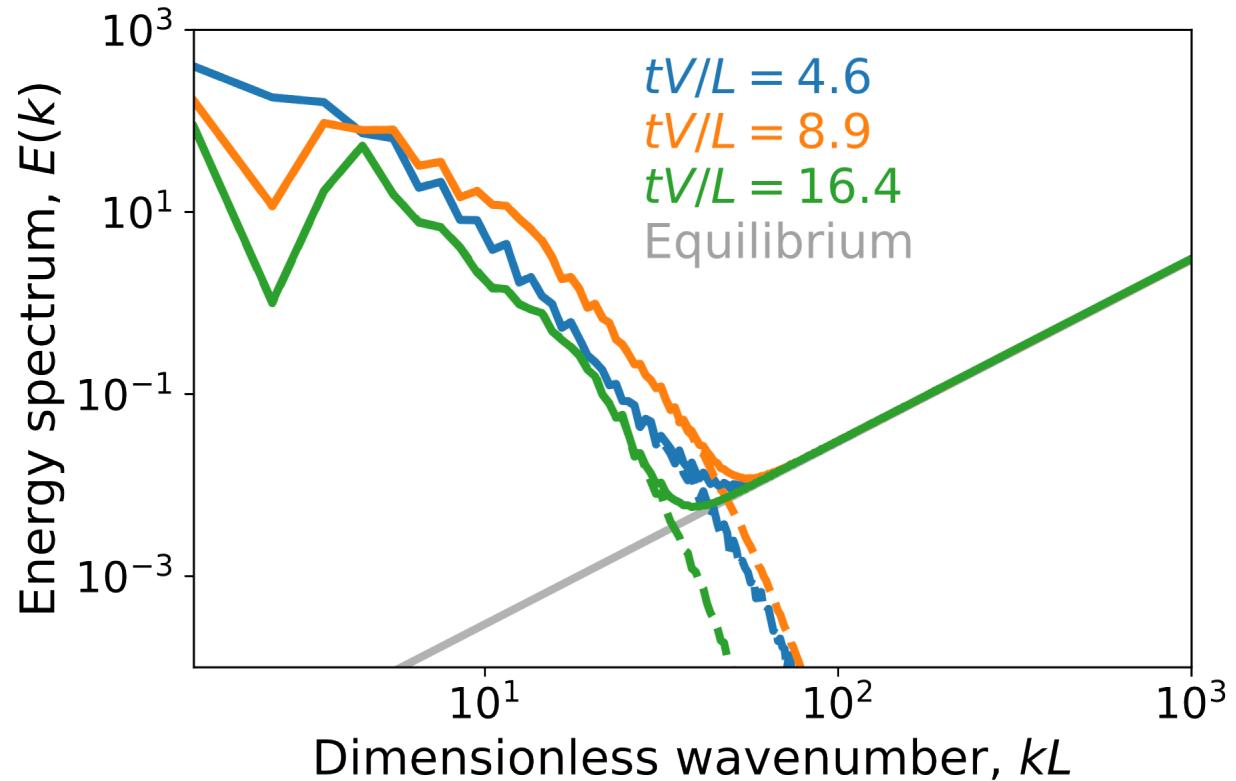


Rayleigh-Taylor



Vortex shedding with slip

# Re = 500 spectra at different times



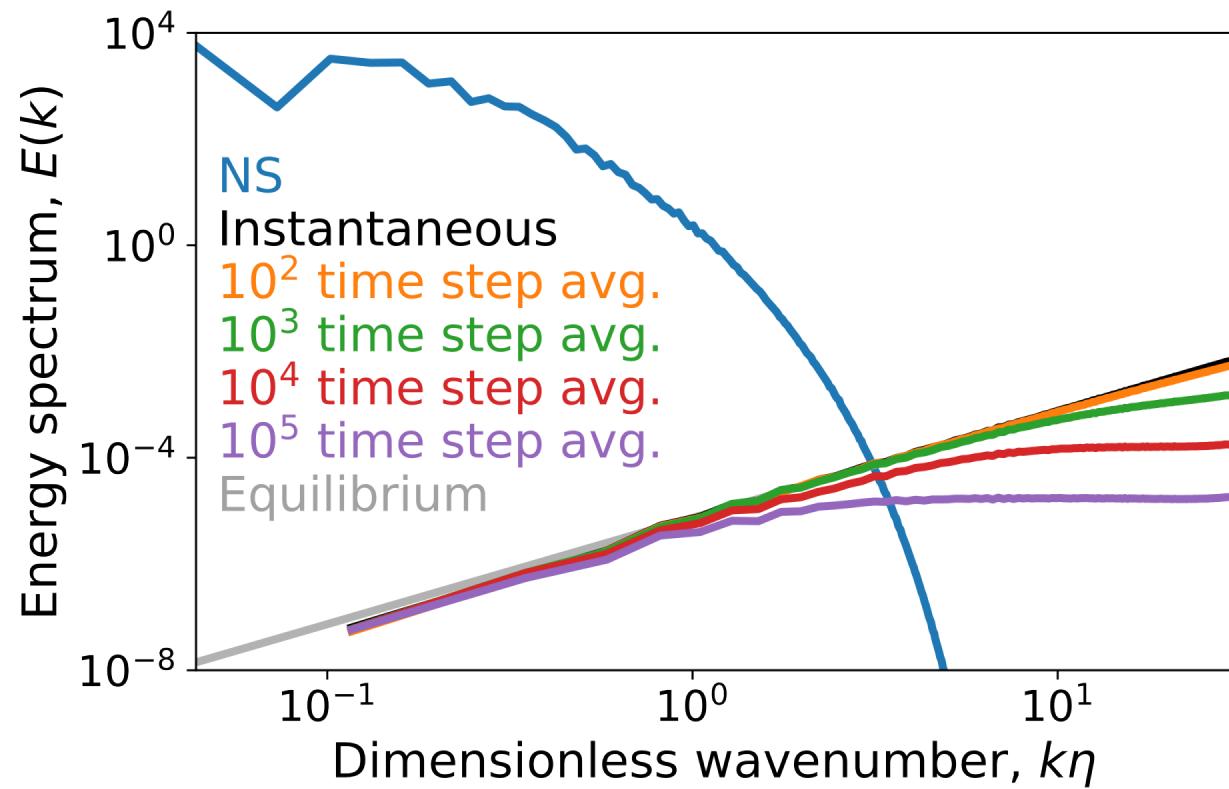
# Time averaging



Averaging is common practice to reduce statistical noise in DSMC

Here,  $10^5$  timesteps corresponds to  $\approx 0.5\tau_\eta$

Only changes crossover scale by  $\approx 10\%$

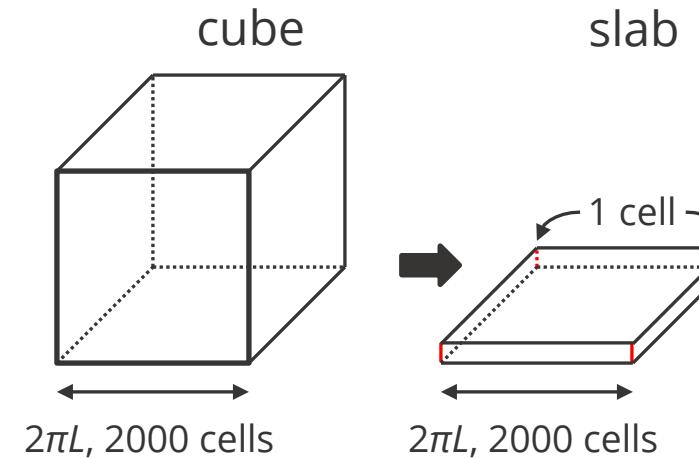


# Viscosity determination for DSMC



Cells are large, so transport is enhanced

- Viscosity is 36% larger than molecular value
- Near-neighbor collisions reduce mean collision separation

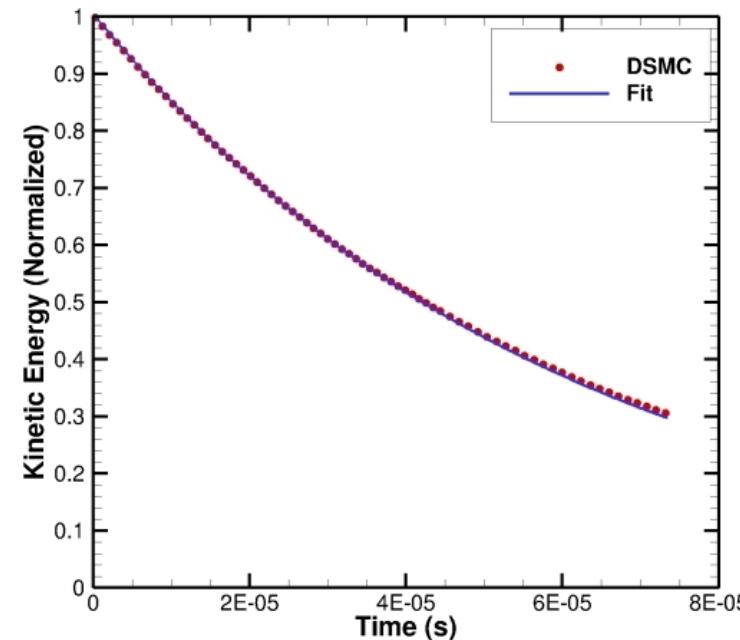


Simulate some other flow to find viscosity

- Use a similar but much easier flow
- 2D TG vortex energy decay:

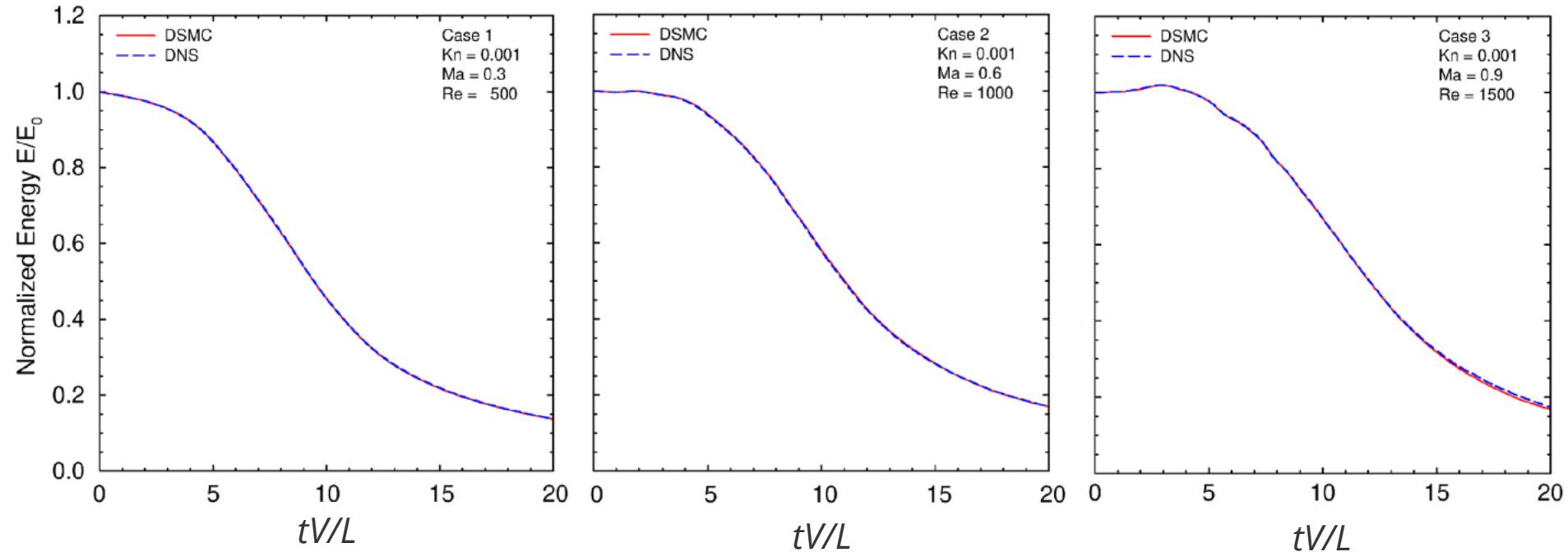
$$E = E_0 \exp(-4\mu_{\text{eff}} t / \rho_0 L^2)$$

Use effective viscosity in NS simulations for comparison



# Kinetic energy decay

Excellent agreement between MGD and NS! \*



\*DSMC data are time-averaged before computing energy