



Bright-Spot Contributions to Hard-Photon Continuum K-Shell Yield from Argon and Stainless-Steel Implosions on Z

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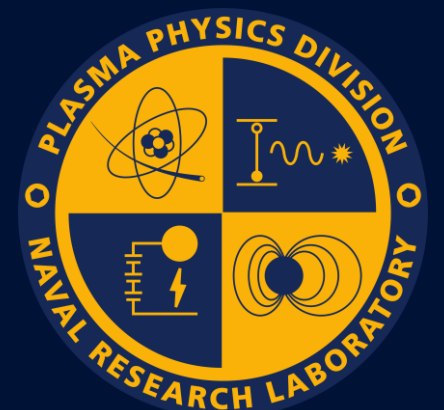
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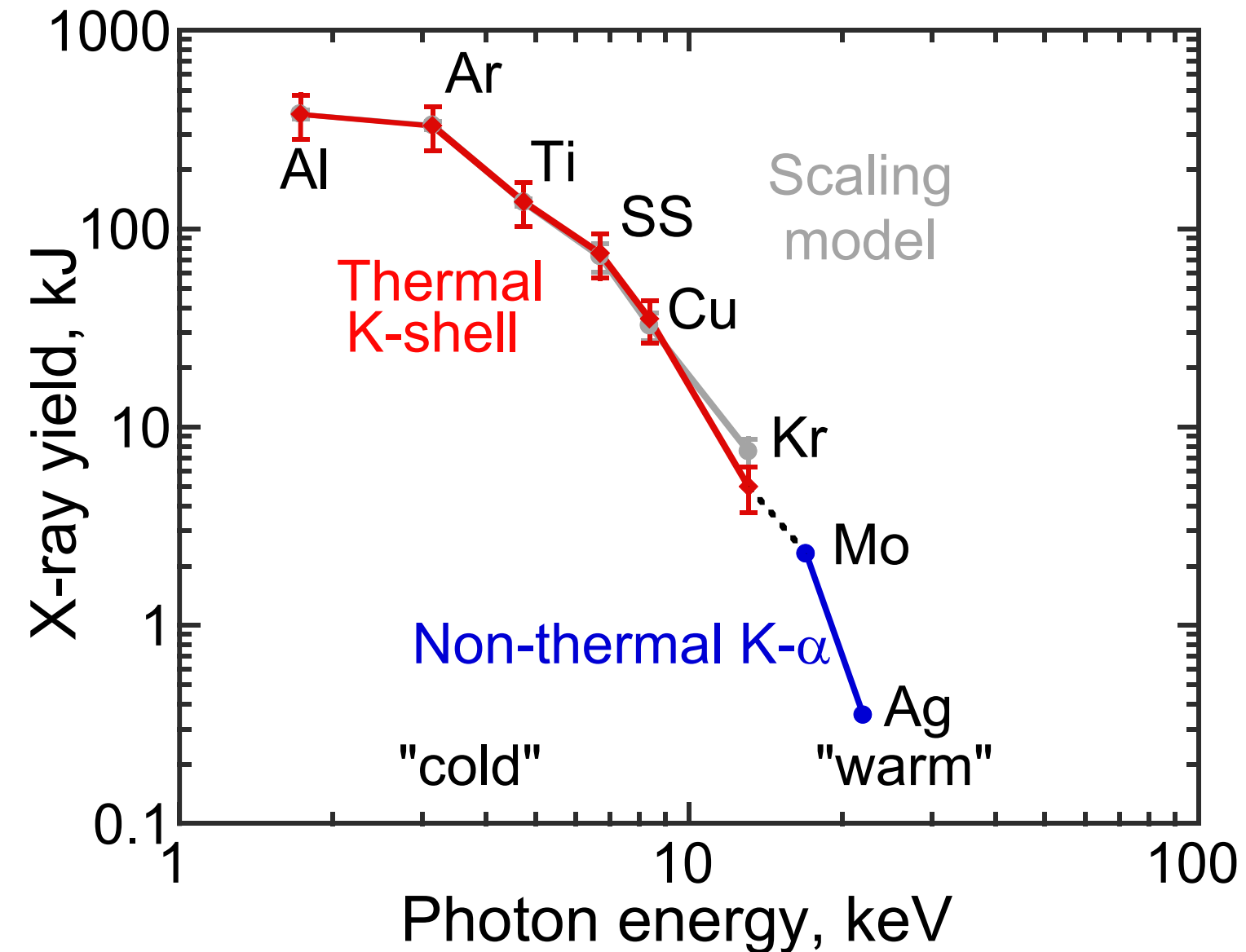


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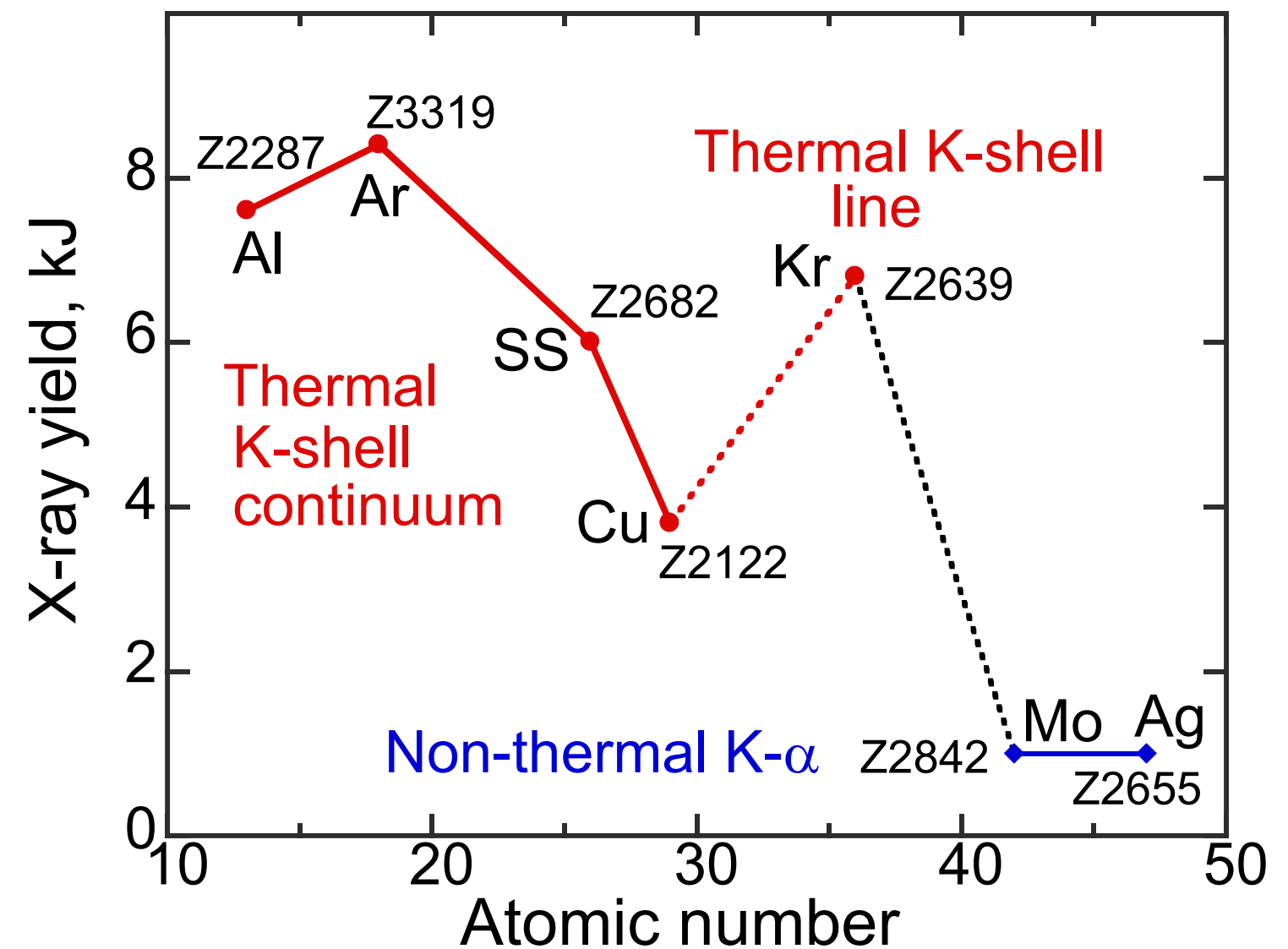
- Hard-photon direct recombination continuum alternative paths to “warm” x-ray production above 15 keV with pulsed power for next-gen pulsed power facility, NGPP
- A scaling model is needed to confidently predict the radiation yields from NGPP
 - Densities and temperatures of radiating plasmas at stagnation need to be confidently scaled up
- Scaling models are based on the conventional assumption of a uniform radiating plasma column, which is plausible for Ar but not for stainless steel
- Most likely, stainless steel emits K-shell from bright spots
- Continuum yield data from Z is consistent with the spectroscopic analysis indicating that the dense, hot bright spots containing a small fraction of the load mass radiate most of the K-shell x rays
 - Dimensions, densities, and temperatures of the bright spots can be inferred
 - The mechanism of their creation and energy replenishment is still obscure

Motivation

K-line yields on Z are low for high atomic numbers



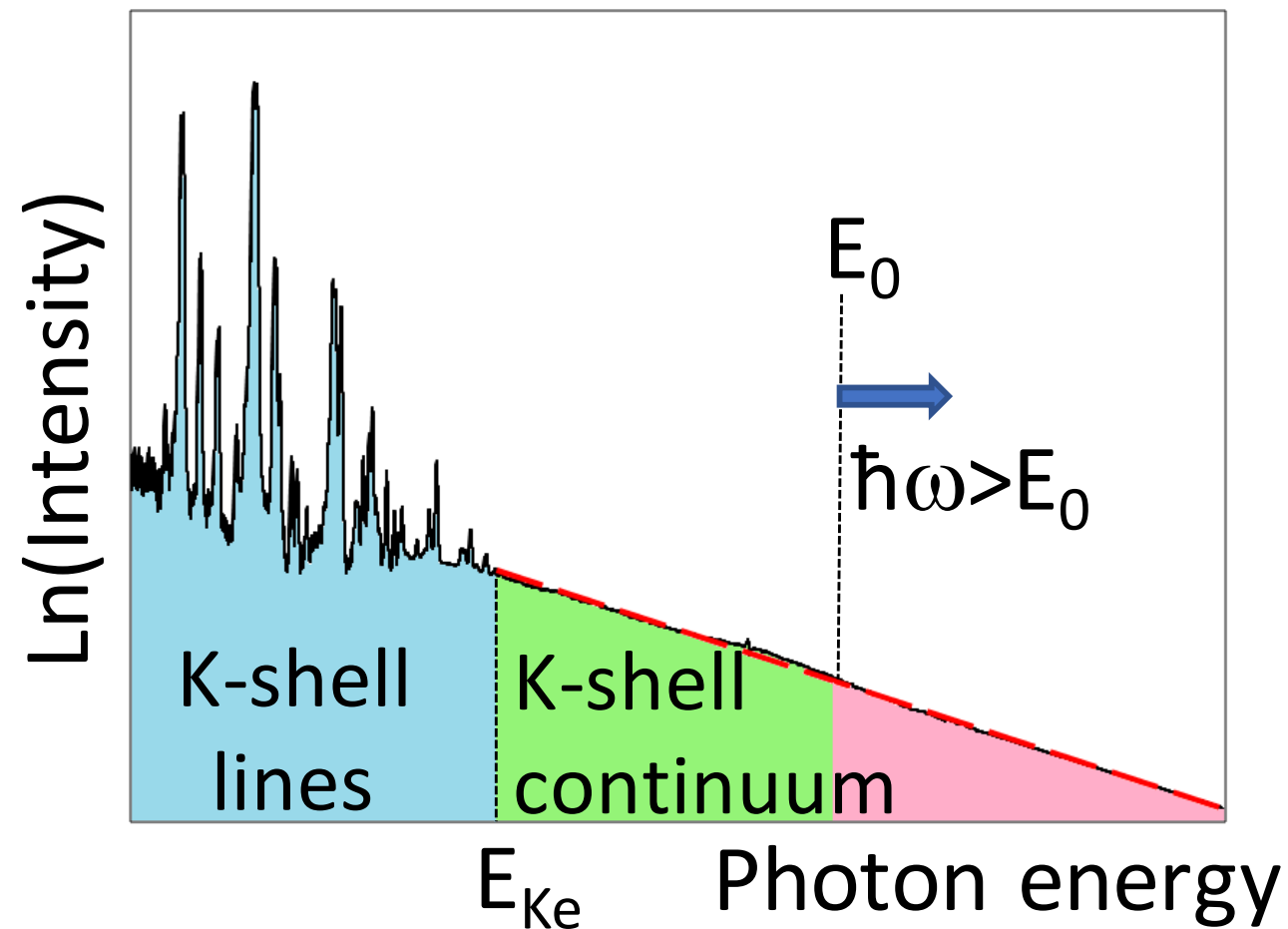
Continuum & line yields on Z in photon energies between 10 and 15 keV



Scaling: J. Schwarz et al., "A model for K-shell x-ray yield from magnetic implosions at Sandia's Z machine," submitted to Phys. Plasmas 2022.

Background

Example: Z3151 stainless wire array



$$T_e = 3.56 \text{ keV}$$

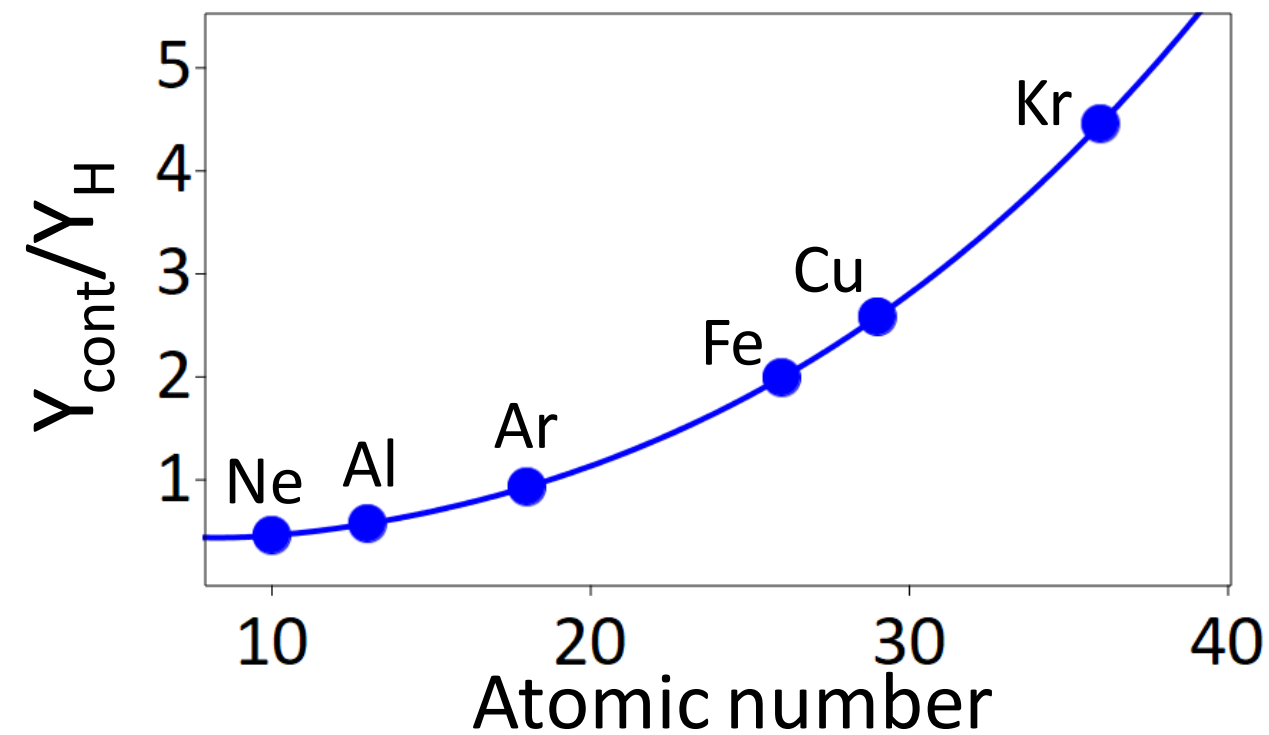
$$E_{Ke} = 9.2 \text{ keV}$$

$$Y_{cont} = 6.3 \text{ kJ}$$

T_e , keV

$$Y(\hbar\omega > E_0) = Y_{cont} \exp\left(-\frac{E_0 - \text{Ryd} \cdot Z^2}{T_e}\right)$$

$$\frac{Y_{cont}}{Y_H} = \frac{2\alpha^3}{3\pi\phi_H} \left(Z_A - \frac{1}{2}\right)^4 \exp\left[\frac{\text{Ryd}}{4T_e} (3Z_A - 1)(Z_A - 1)\right]$$

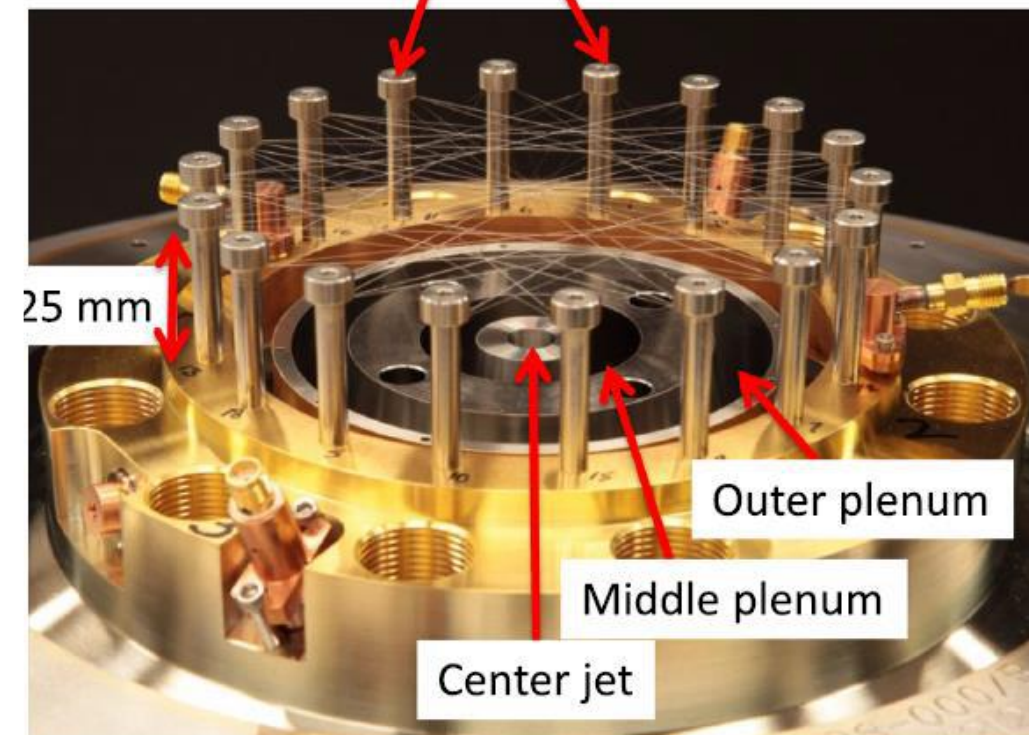


@ K-shell
production
optimal
temperature
 $T_e = 0.3 \times Z_A^{2.9} \text{ eV}$

Argon and stainless steel PRS load designs optimized for K-shell emission

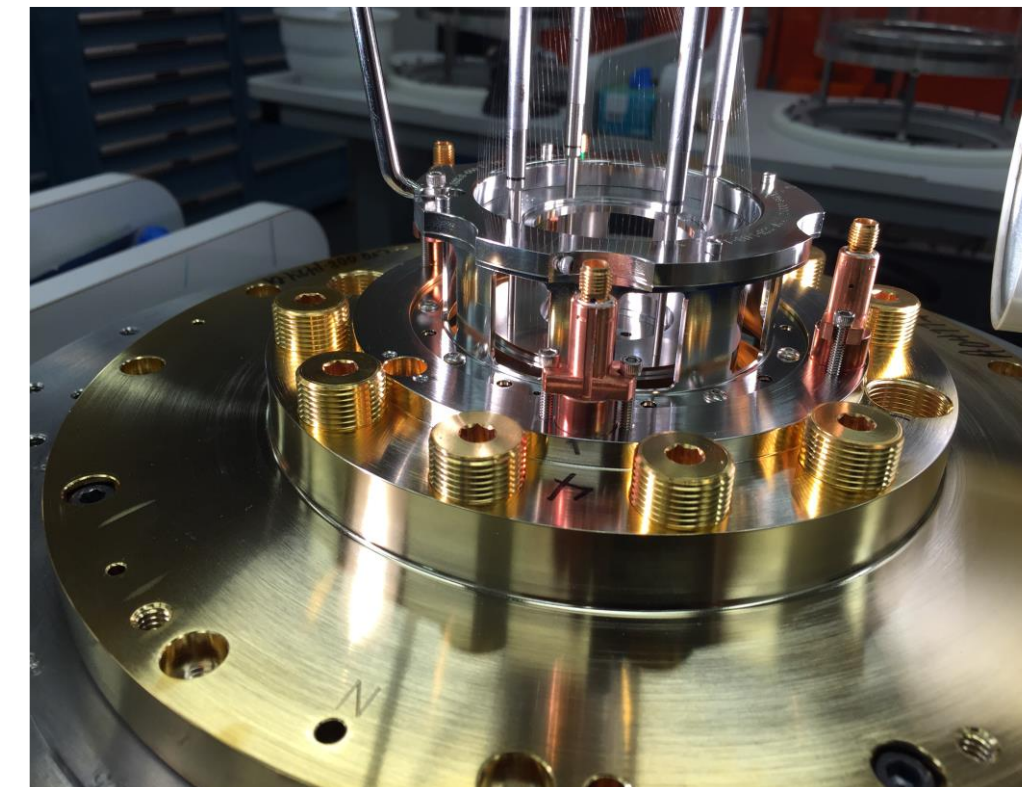
Argon, 25 shots

Return posts with anode mesh



- Double-shell $\varnothing 80$ mm gas puff
- Most shots w/o central jets
- Line masses $\sim 385/616 \mu\text{g/cm}$
- Height 3 cm

Stainless steel, 8 shots

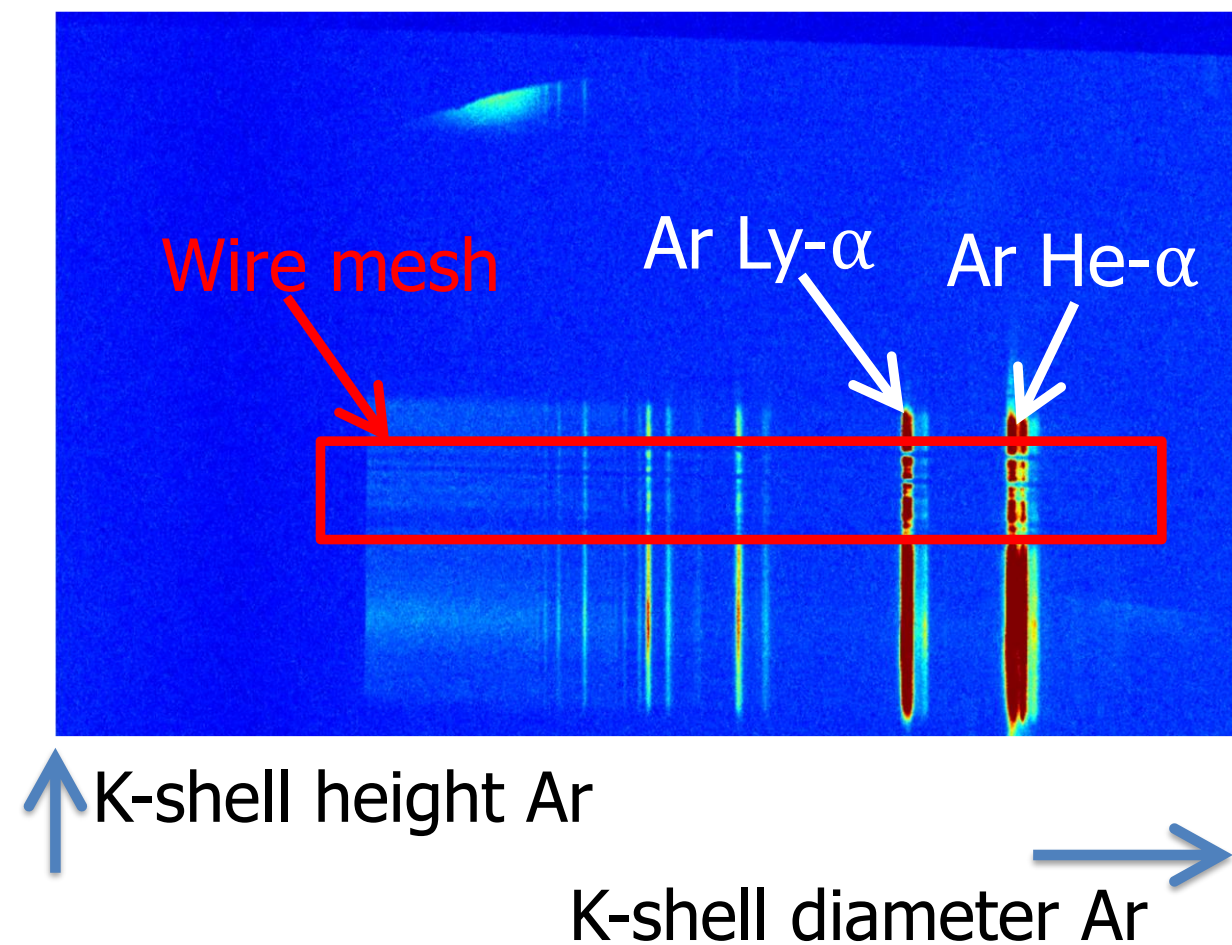


- Nested wire array, $\varnothing 70/35$ mm, 112/56 wires
- Wire diameter $8.56 \mu\text{m}$
- Total mass ~ 1.533 mg
- Height 2 cm

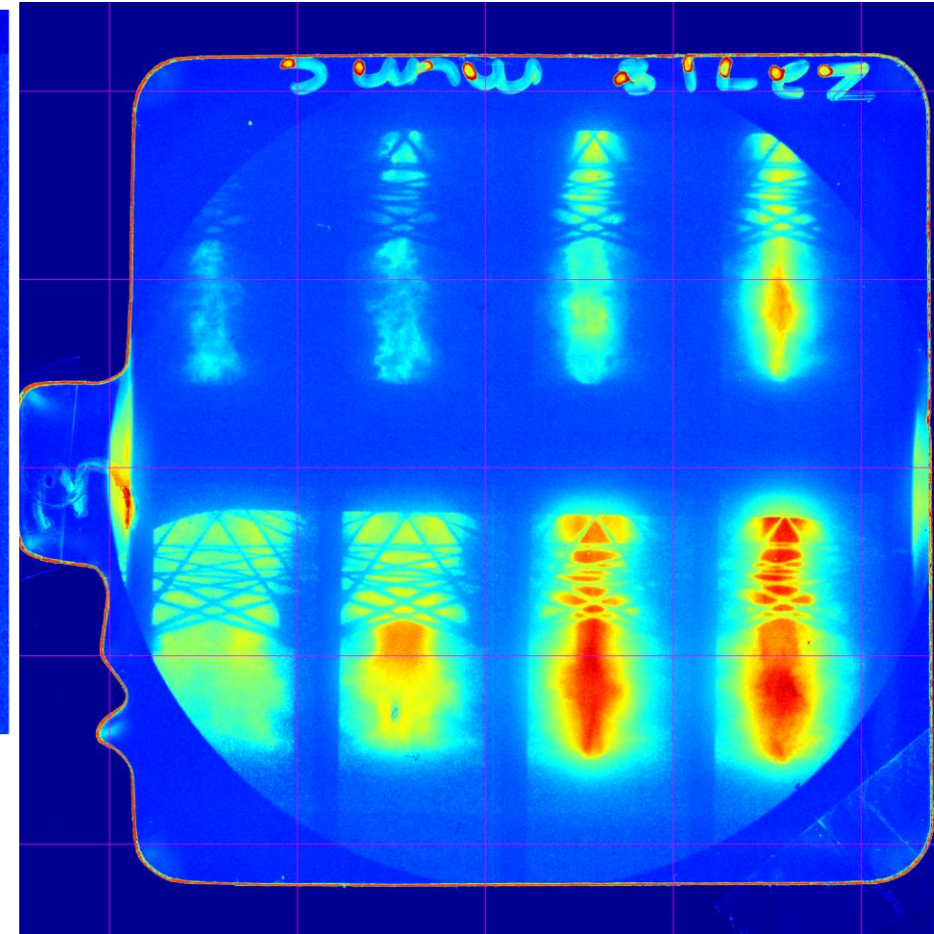
B. Jones *et al.*, Phys. Plasmas **22**, 020706 (2015); A. J. Harvey-Thompson *et al.*, Phys. Plasmas **23**, 101203 (2016); D. J. Ampleford, *et al.*, Phys. Plasmas **21**, 056708 (2014).

Dimensions and of the K-shell emitting plasma are confidently measured

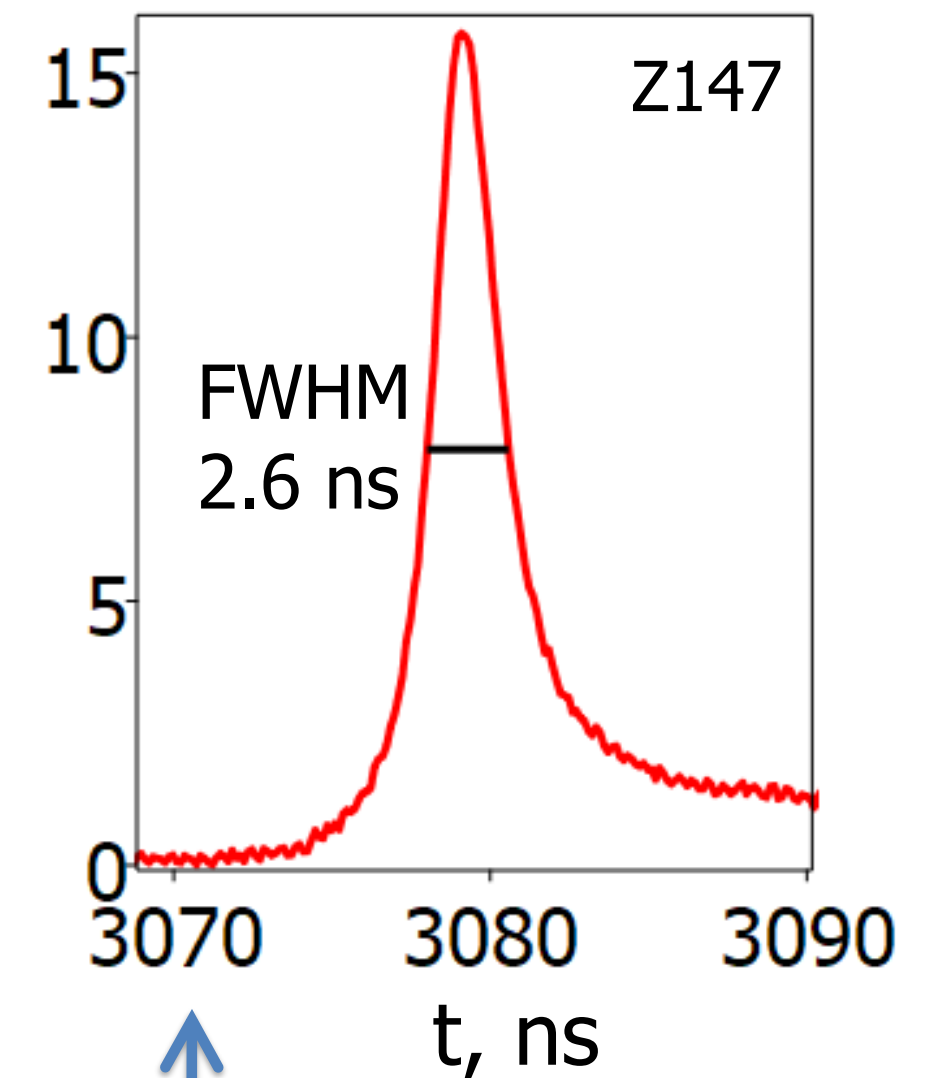
Time-integrated, axially resolved spectroscopy



Time-gated multi-layer-mirror (MLM) x-ray camera



X-ray power signal, arb. units

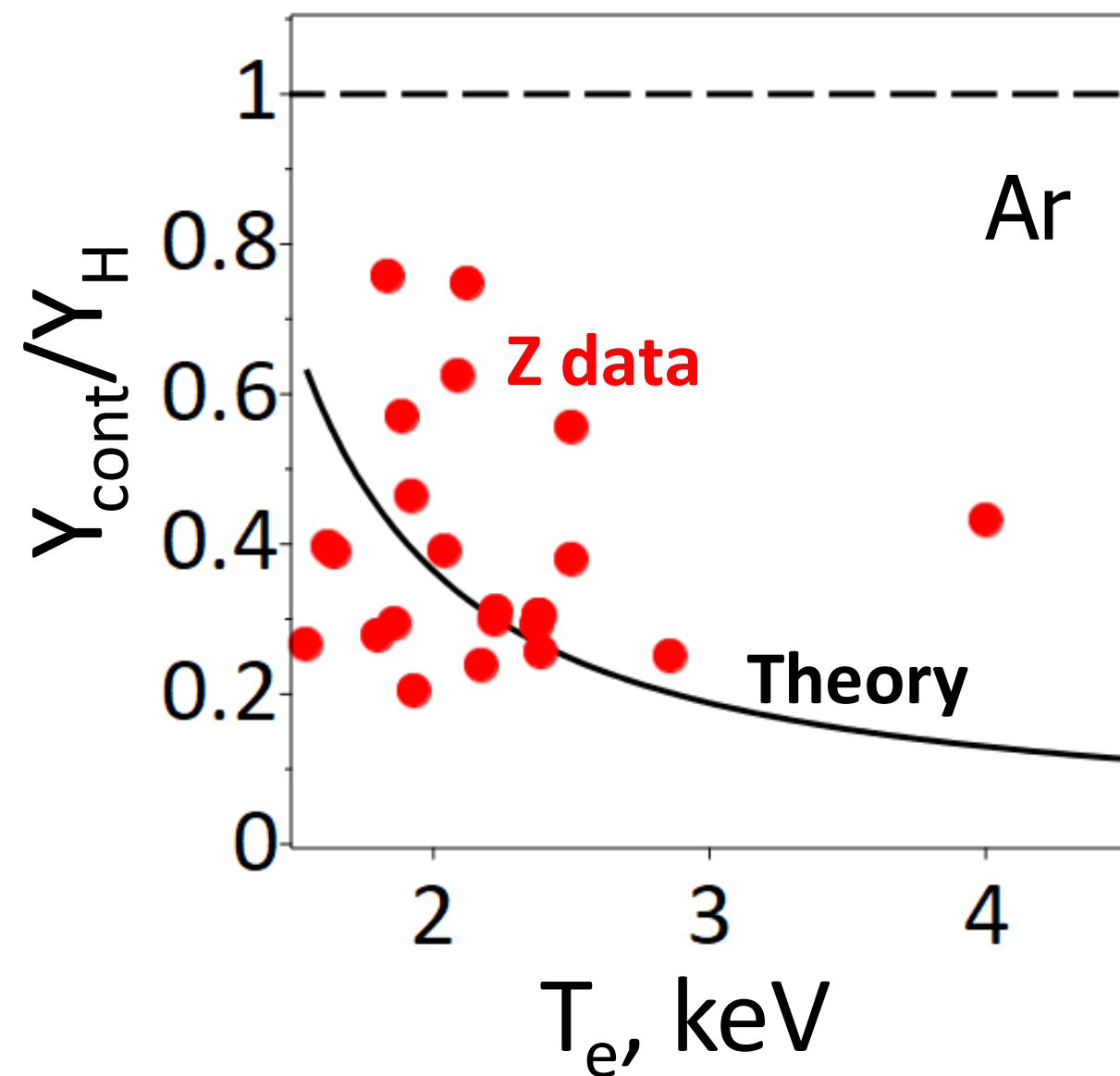


K-shell FWHM SS

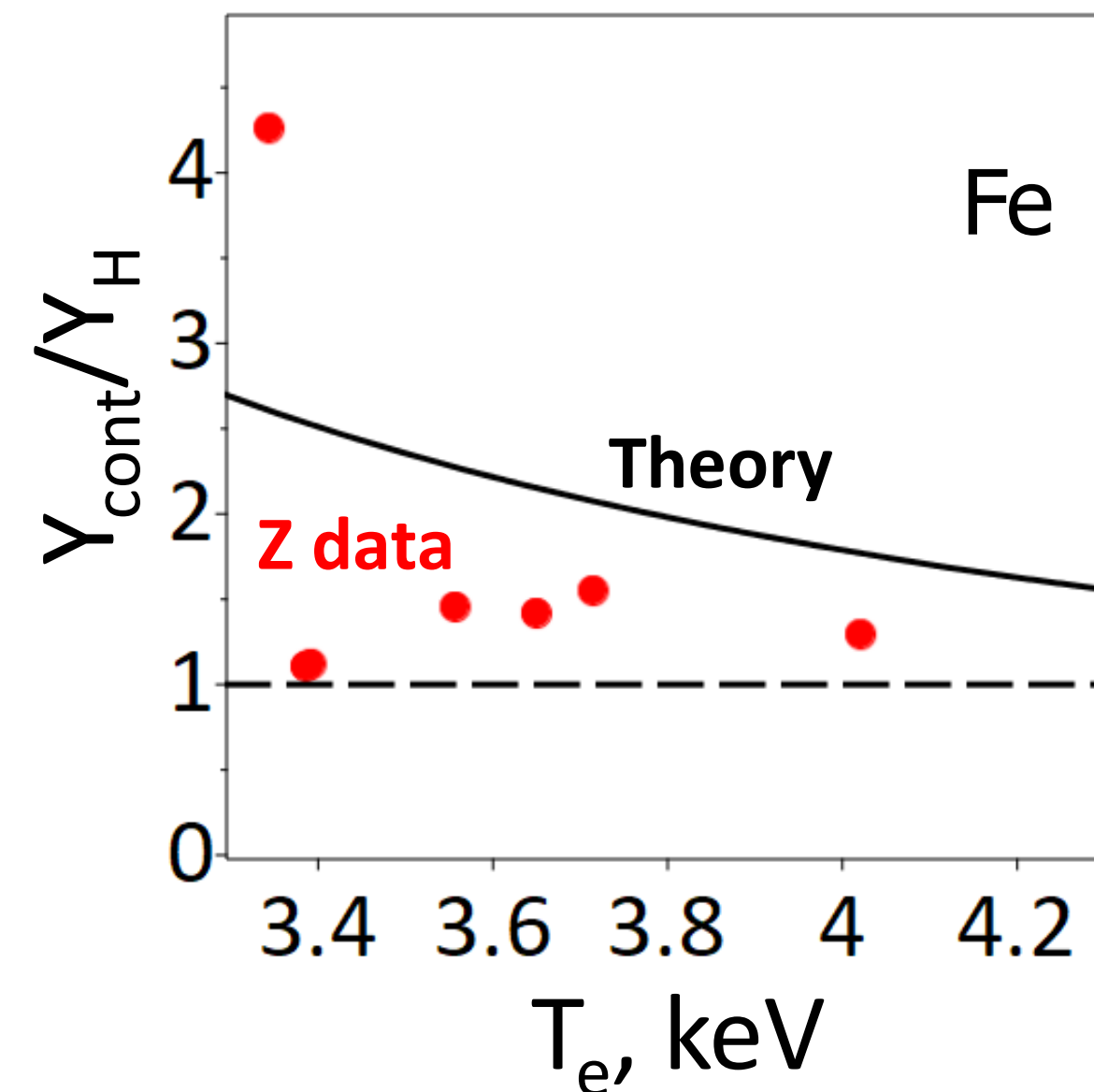
	Height, mm	Diameter, mm	Volume, cm ³	FWHM, ns
Argon	2.5 to 3.5	1 to 3	0.05 to 0.36	4 to 12
Stainless	2	0.8	0.01	2.5 to 3.5

Higher-atomic-number materials radiate relatively more in the continuum

Argon: the yield in continuum is less than in Ly lines, even though bare ions contribute



Iron: the yield in continuum is greater than in Ly lines; no contribution from bare ions

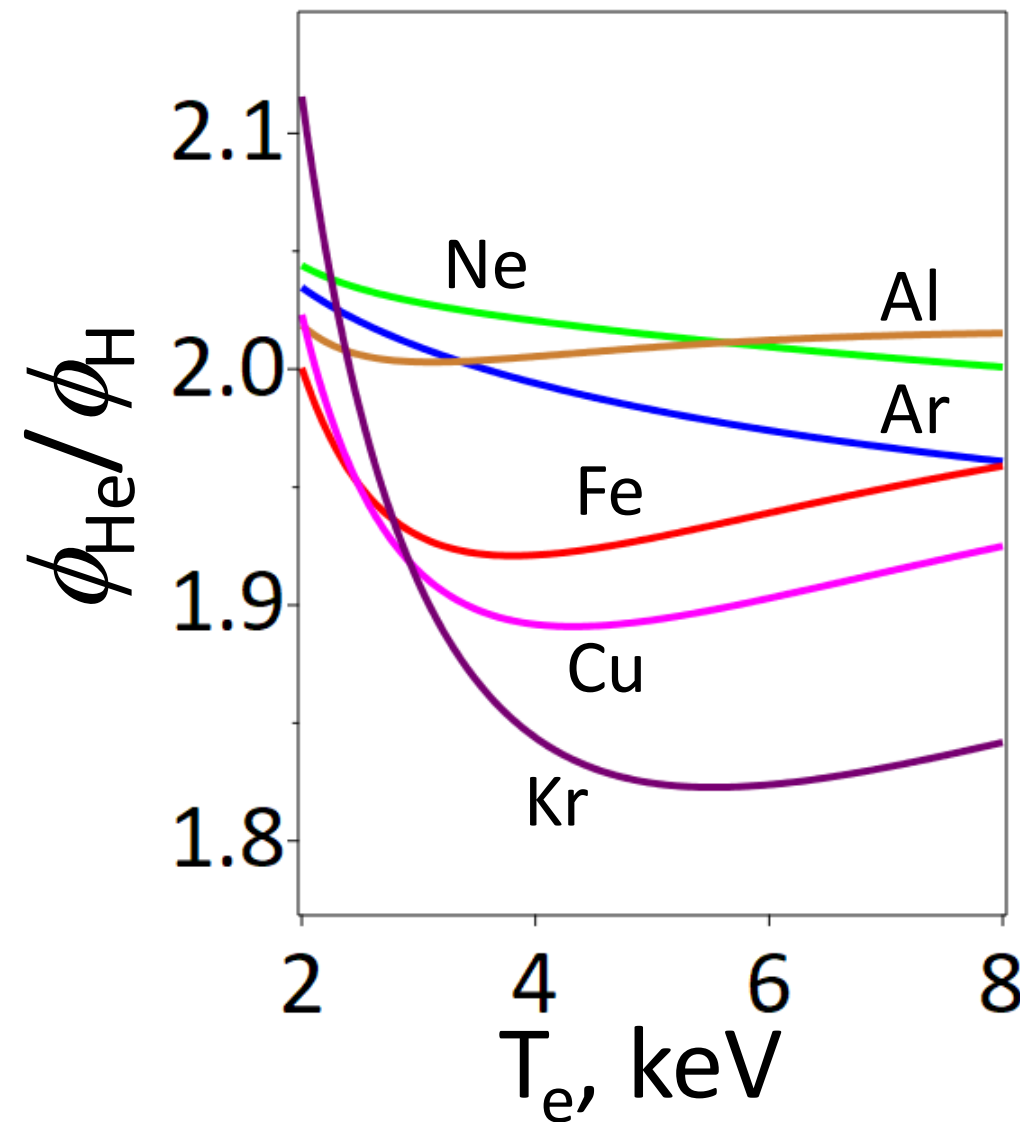


He-like ions are more efficient K-line radiators than H-like by a factor of ~ 2

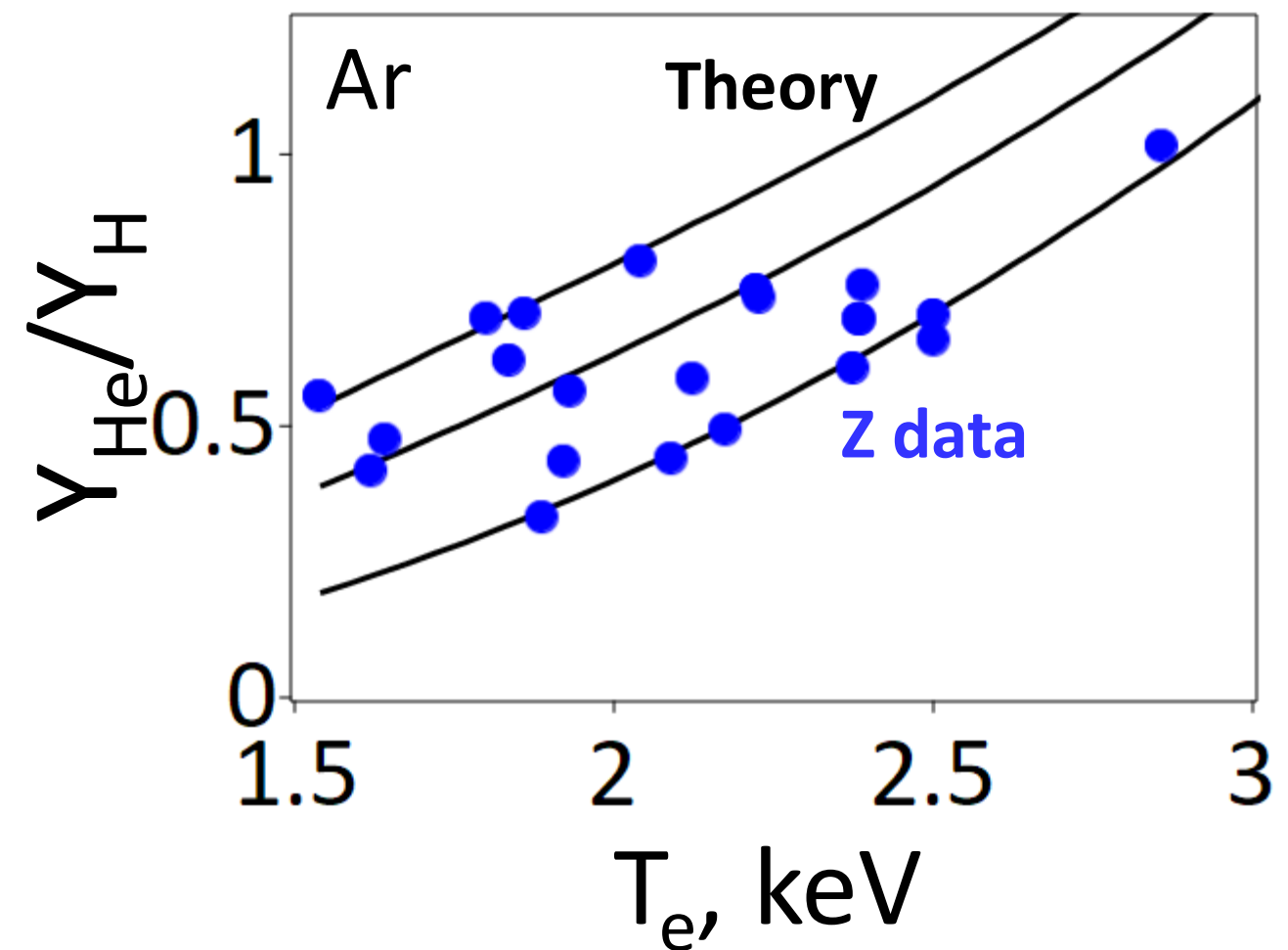
Ratio of H-like
to He-like yields

$$\frac{Y_H}{Y_{He}} = \frac{n_{i,H}}{n_{i,He}} \cdot \frac{\phi_H}{\phi_{He}}$$

Seaton's formula
fit with FAC
atomic data



Argon on Z: more H-like than He-like ions,
Yield in He lines slightly higher than in Ly lines

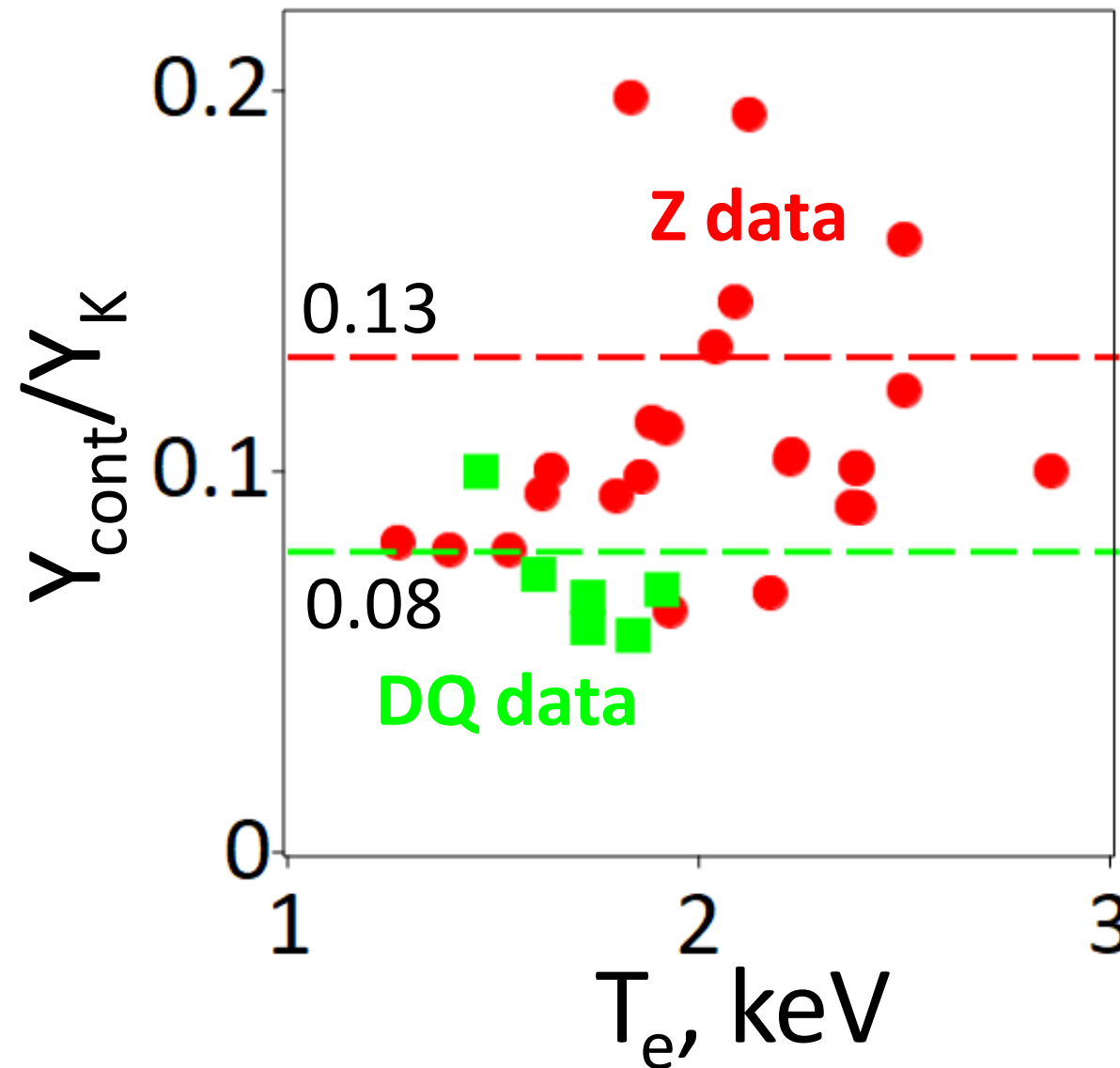


The highest hard-photon continuum fractions we can reasonably expect

Moderate- Z_A argon, accessible now

$$\begin{aligned} &Z \\ &Y_H/Y_{He} = 0.6 \\ &Y_{cont}/Y_H = 0.4 \\ &\Rightarrow \\ &Y_{cont}/Y_K \sim 0.13 \\ \\ &DQ \\ &Y_H/Y_{He} = 0.4 \\ &Y_{cont}/Y_H = 0.3 \\ &\Rightarrow \\ &Y_{cont}/Y_K \sim 0.08 \end{aligned}$$

DQ data from F. C. Young *et al.*, IEEE Trans. Plasma Sci. **34**, 2312 (2006).



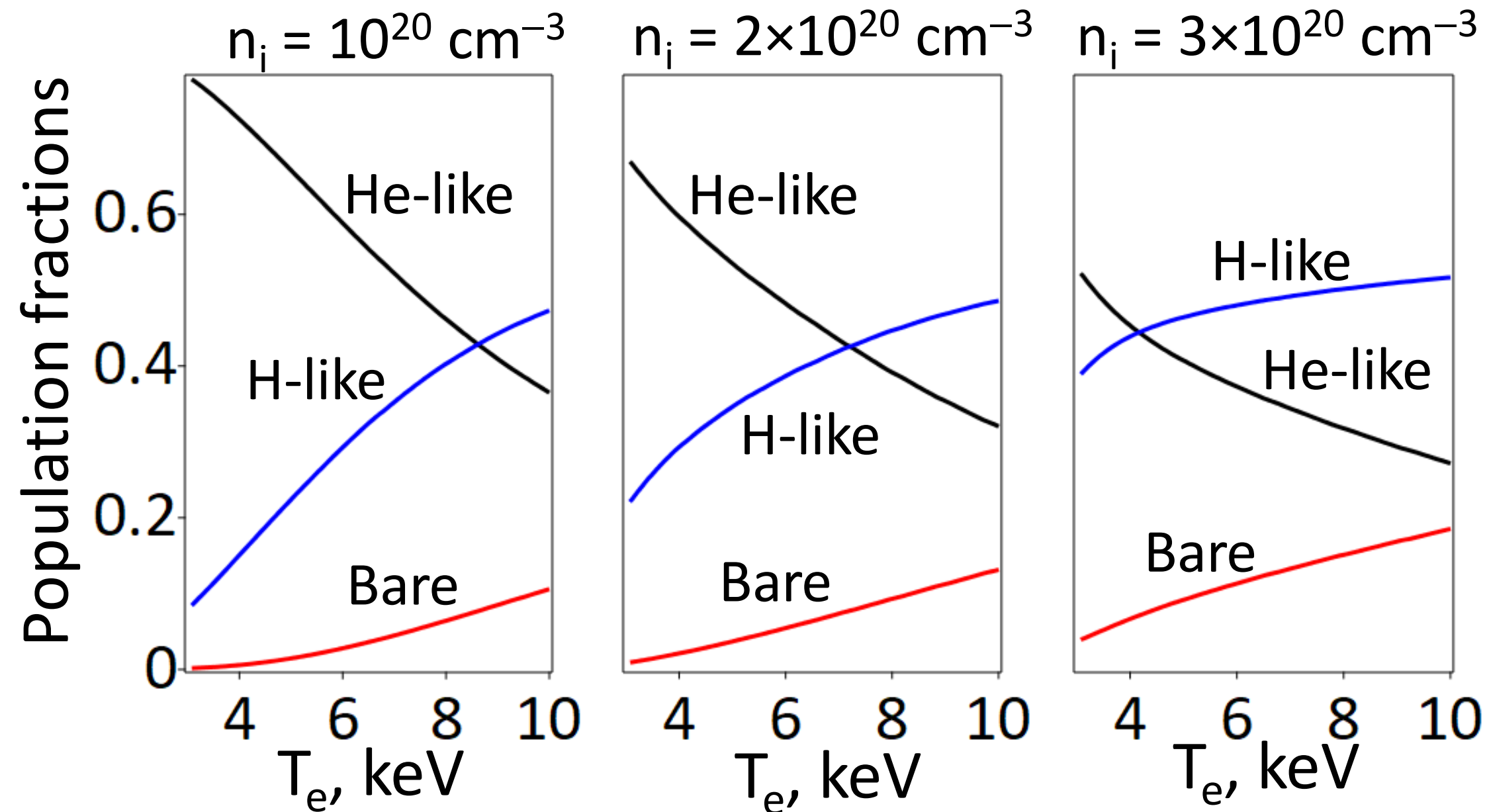
High- Z_A iron or copper on NGPP

$$\begin{aligned} &\text{If } Y_H/Y_{He} = 0.5 \\ &\text{Then } Y_{cont}/Y_H > 1 \\ &\Rightarrow \\ &Y_{cont}/Y_K > 0.25 \end{aligned}$$

A 25% continuum efficiency at high Z_A is possible if we can produce an H-like population no less than He-like

We need high temperatures and densities to ionize Fe to H-like

CRE equilibrium populations calculated for a $\varnothing 3$ mm plasma cylinder with DRACHMA-II



Ion populations are very sensitive to the ion densities above $n_i \sim 2 \times 10^{20} \text{ cm}^{-3}$

Ion number density can be inferred from the hard-photon continuum power



$$W_{cont} = F(T) m_K n_i (2f_b + f_H)$$

Radiation power is proportional to the product of the radiating mass and ion density

$$n_i = \frac{W_{cont}}{F(T) m_K (2f_b + f_H)}$$

where the ion populations for our parameter range are estimated as

$$2f_b + f_H \approx \begin{cases} 1 & \text{for Ar} \\ f_{He} \cdot 2Y_H / Y_{He}, & f_{He} \approx 0.7 \text{ for SS} \end{cases}$$

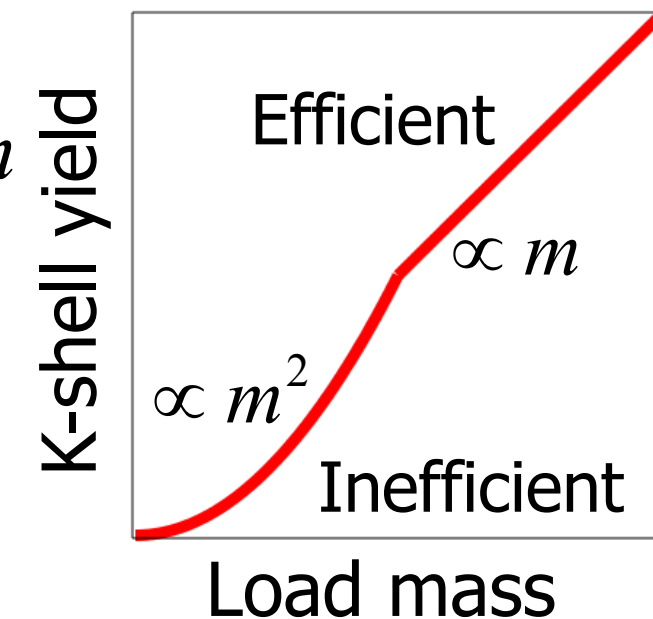
$$m_K = R_K \sqrt{\frac{W_{cont} \pi m_i}{F(2f_b + f_H)}}$$

Self-consistent value of the continuum K-shell emitting mass, obtained by assuming the radiating plasma to be uniform

Assumptions for scaling line and continuum K-shell yields to higher currents

$$Y_K \propto m \cdot (\text{K-shell mass participation}) \cdot n_i$$

- Matched load mass increases as I^2 . Some options for self-similar scaling:¹
 - Mass participation and ion density stay constant²
 - Efficient K-shell emission, $Y_K \propto I^2$, radiating volume increases $\propto m$
 - Works well for argon above 5 MA
 - Mass participation and radiating volume stay constant²
 - Inefficient K-shell emission, $Y_K \propto I^4$
 - Works well for argon³ below 5 MA and for stainless⁴ up to 20 MA
 - Uncertain about stainless above 20 MA – what stays constant, what changes?



¹P. F. Schmit and D. E. Ruiz, Phys. Plasmas **27**, 062707 (2020).

²J. W. Thornhill *et al.*, IEEE Trans. Plasma Sci. **34**, 2377 (2006).

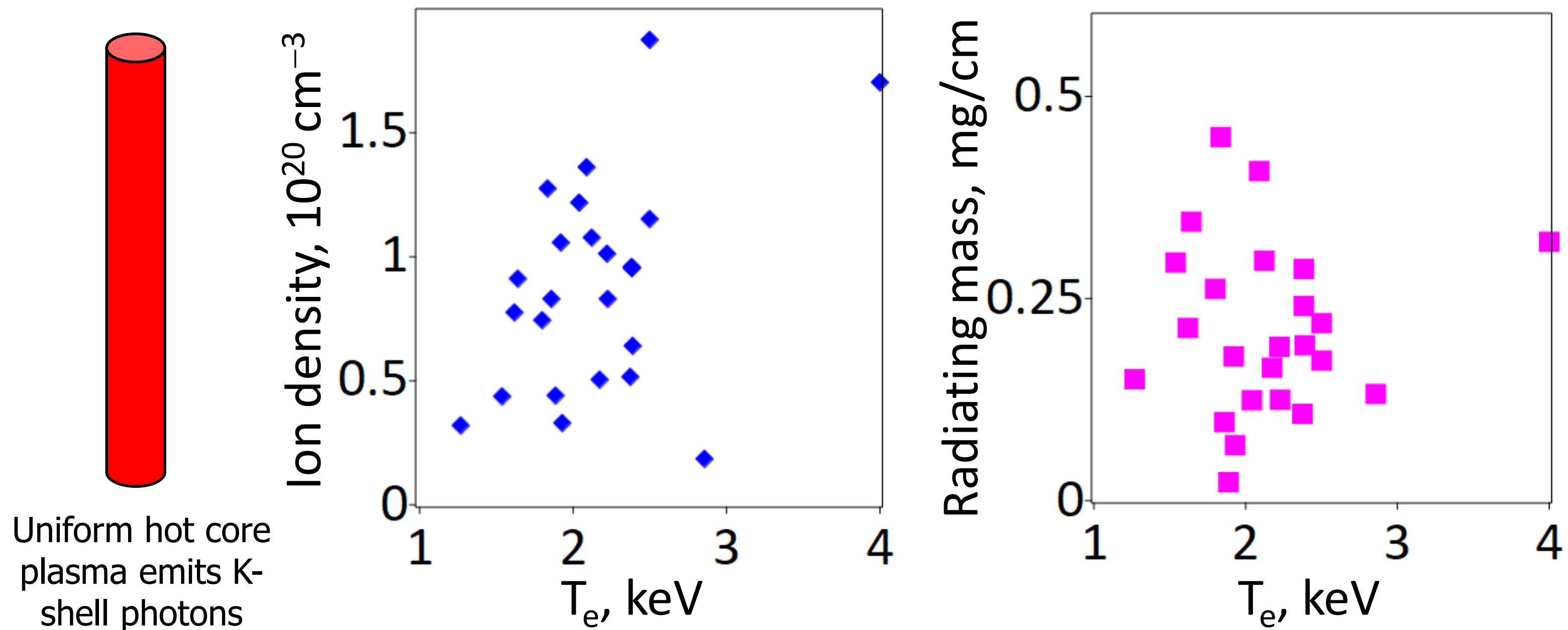
³This conference: V. Tangri *et al.*, PO 4.33 "Scaling of efficient Ar K-shell emission from fast gas-puff Z-pinch in 10 to 100 MA current range"; A. Esaulov *et al.*, PO 4.37 "Progress in the refining of the K-shell yield scaling model for Z-pinch plasma radiation sources."

⁴J. Schwarz *et al.*, "A model for K-shell x-ray yield from magnetic implosions at Sandia's Z machine," submitted to Phys. Plasmas 2022.

$$Y_K \propto m \cdot (\text{K-shell mass participation}) \cdot n_i$$

For argon, the assumption of uniformity makes sense

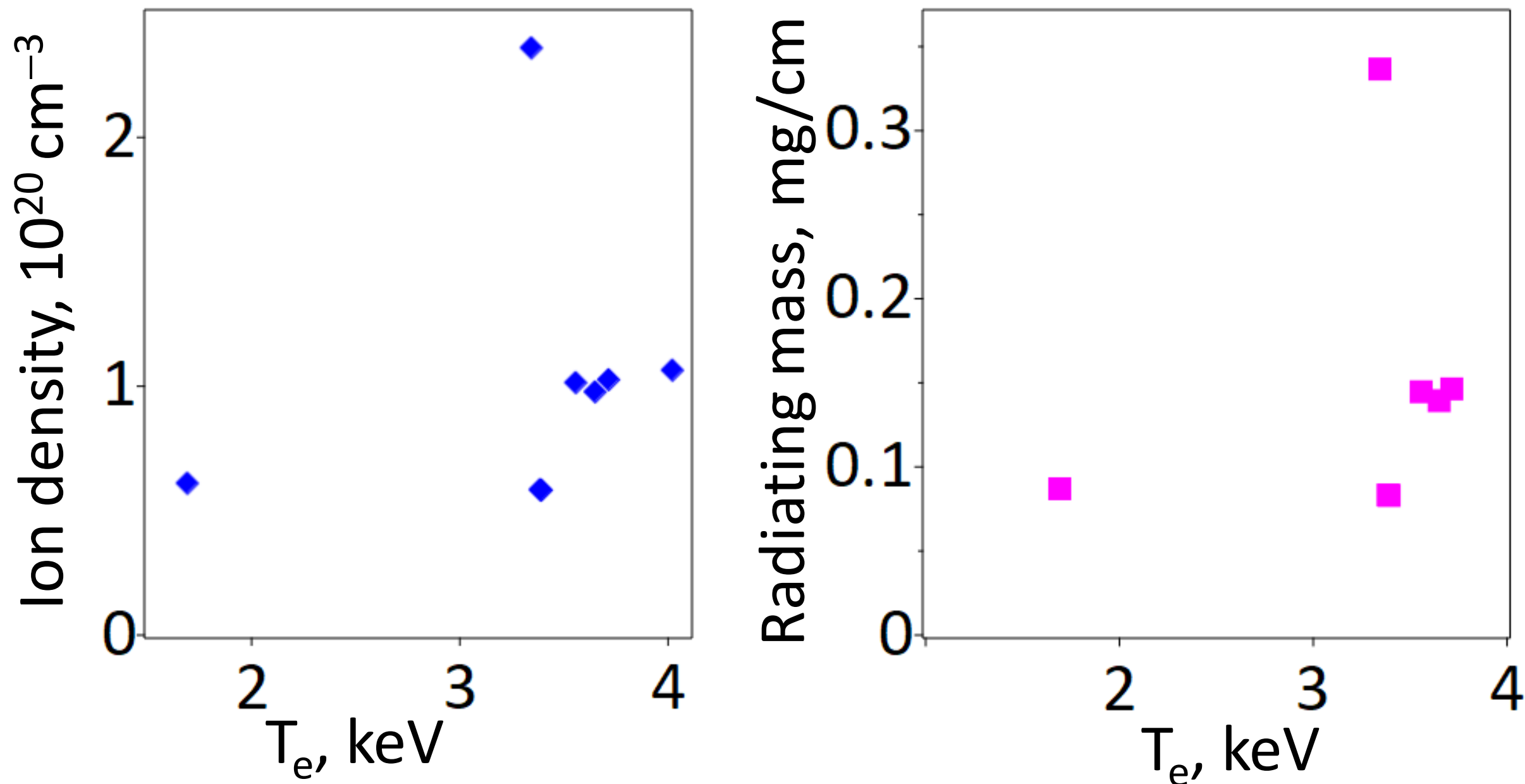
For 1 mg/cm, $\varnothing 3$ mm Ar load, the assumption of a $\sim 20\%$ average mass participation in the K-shell, leads to reasonable estimates of the ion density, 10^{20} cm^{-3} and line mass, 0.2 mg/cm.



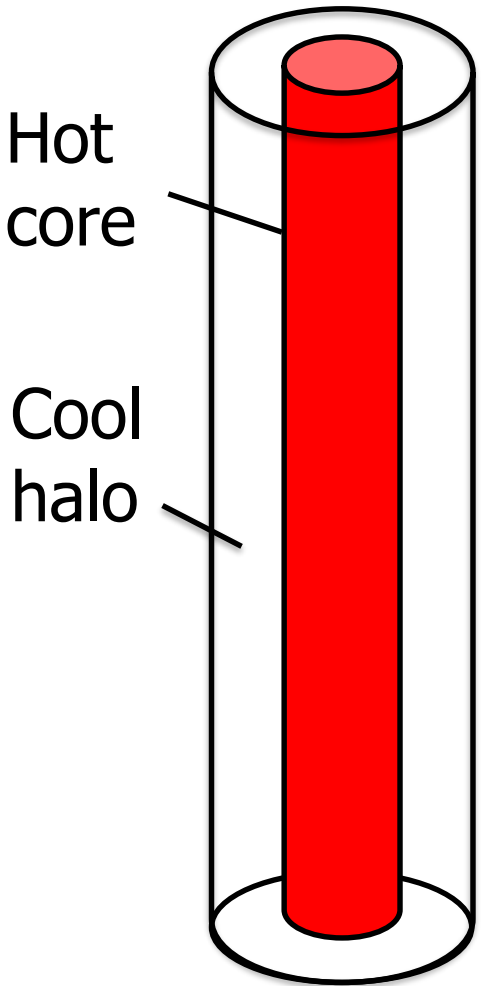
For stainless, the uniformity implies high mass participation, same as in Ar

For 0.75 mg/cm, $\varnothing 1.4$ mm Ar load, the assumption of a $\sim 18\%$ average mass participation in the K-shell, leads to estimates of the ion density, $\sim 10^{20} \text{ cm}^{-3}$ and line mass, 0.14 mg/cm.

Estimated SS plasma parameters are close to argon's, but:
The radius is twice smaller and the temperature is $\sim 40\%$ higher



Time-resolved spectroscopy agrees with this early in the K-shell pulse



Cool halo absorbs some K-shell photons, affecting observed spectra

- Analysis of time-resolved spectroscopy from Z stainless nested wire-array shots done at Weizmann in 2014-2016
- The ratios of the FWHMs of resonant lines of Fe, Cr, Ni, and Mn, that are of different abundances , are analyzed. It also requires the determination of Doppler broadening for each element, and the correction for splitting effects.
- At early time, the spectra are consistent with a hot core surrounded with a cool halo

t, ns	Cylinder diameter, mm	n_i , cm^{-3}	T_e , keV	Mass participation, %
-2.7	1.8	3×10^{19}	2	12
-1.8	1.6	5×10^{19}	3	17

Near the K-shell peak, the K-line shapes are consistent with coming from a large number of small hot spots: higher density and higher temperature

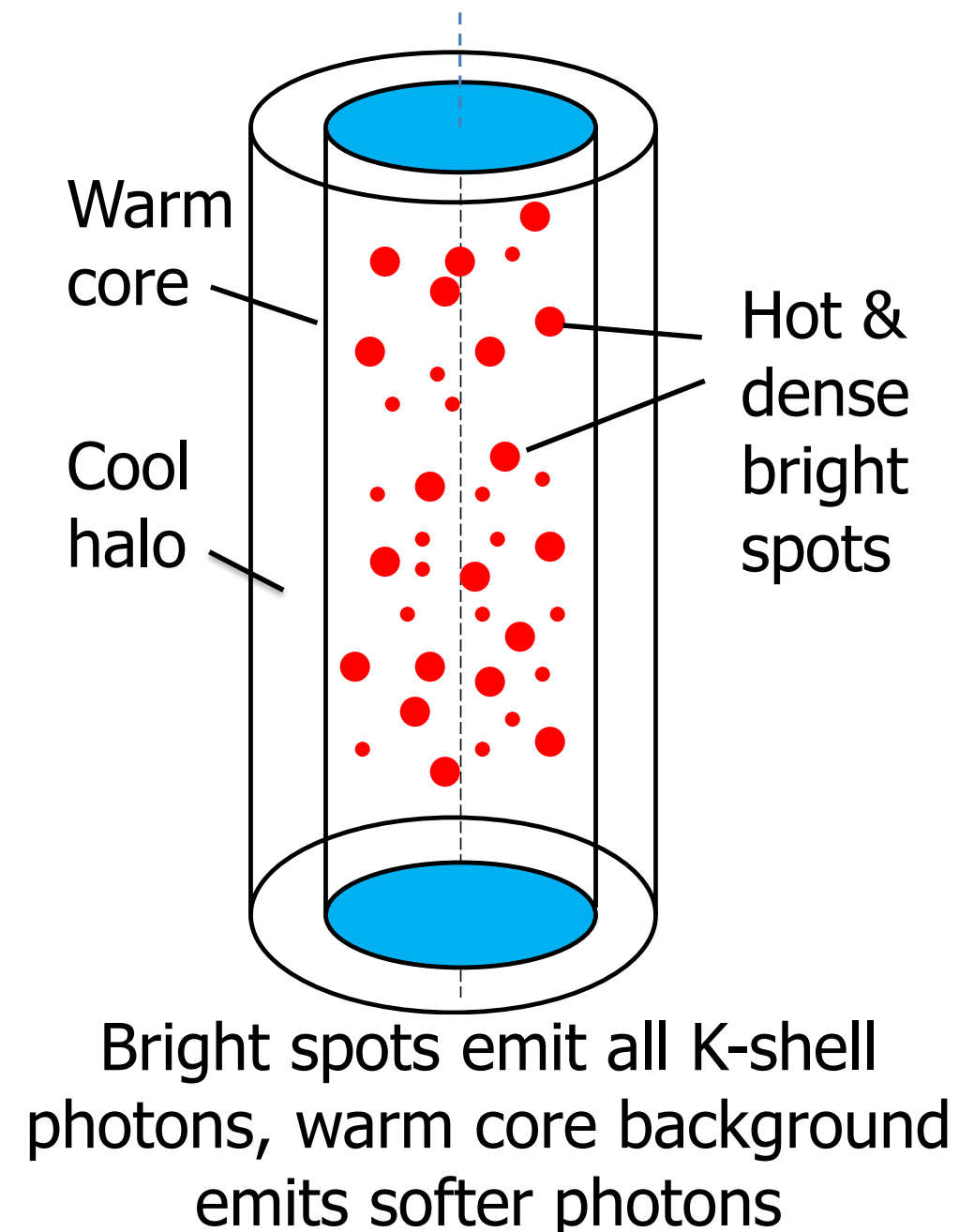
The spot parameters are constrained by the low opacity inferred and the requirement to generate the entire K power

$t = -0.8$ ns, $T_e = 4$ keV, T_i (hydro) = 40 keV

D, μm	No. of spots	n_i , cm^{-3}	Mass participation, %	Photon mfp, mm
50	8900	5×10^{20}	3.1	1.3
100	4200	2.5×10^{20}	5.9	0.68
150	2930	1.67×10^{20}	9.2	0.43

$t = +0.1$ ns, $T_e = 6$ keV, T_i (hydro) = 30 keV

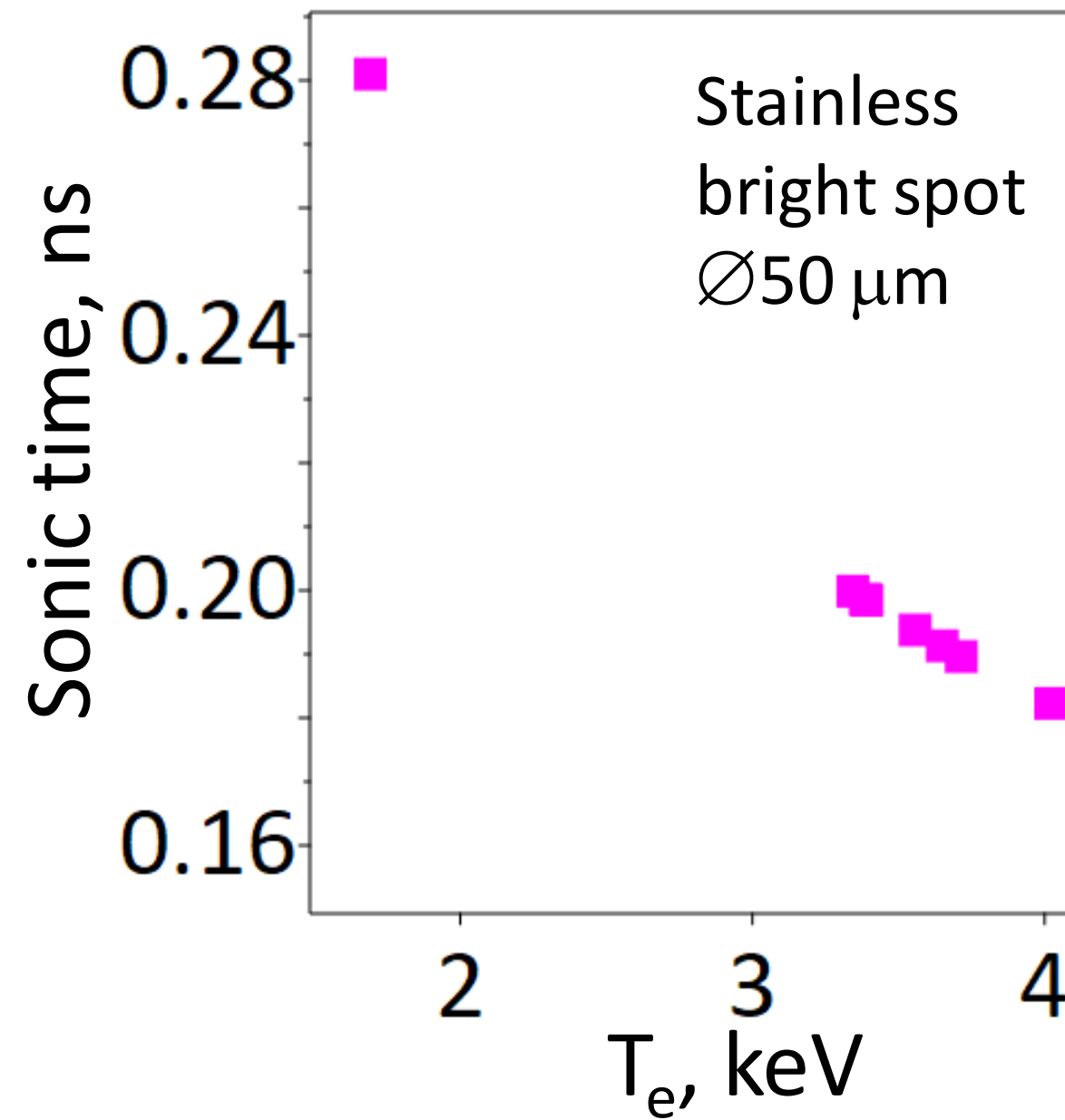
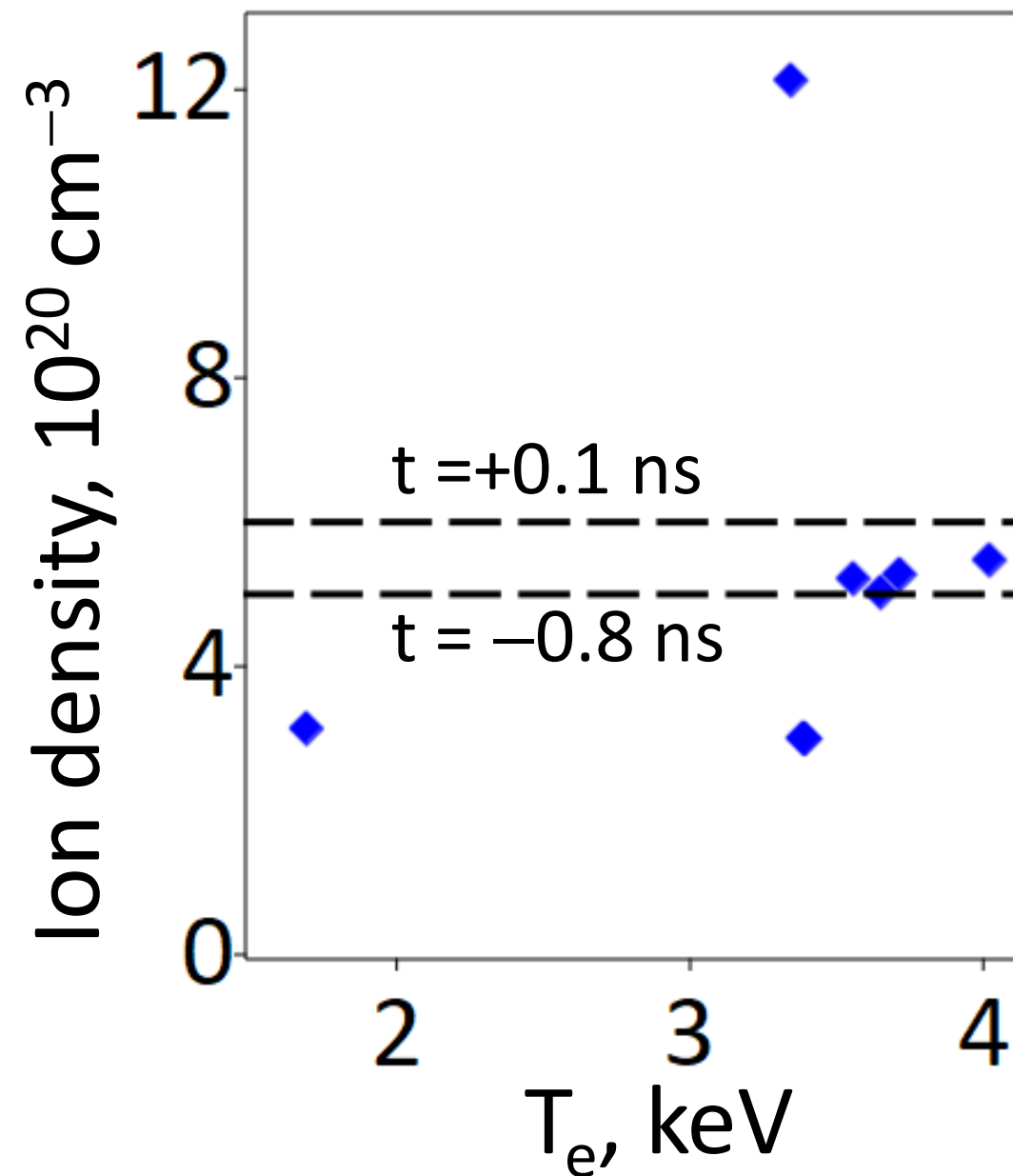
D, μm	No. of spots	n_i , cm^{-3}	Mass participation, %	Photon mfp, mm
50	8200	6×10^{20}	3.4	1.4
100	3800	3×10^{20}	6.4	0.76
150	2350	2×10^{20}	9.2	0.54



Continuum yield are also consistent with this assumption

Stainless mass participation 3.5%

Life time of spots is less than 1 ns. They must be replaced or re-energized.



$$t_s = \frac{2R}{c_s}$$

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