



Bright-Spot Contributions to Hard-Photon Continuum K-Shell Yield from Argon and Stainless-Steel Implosions on Z

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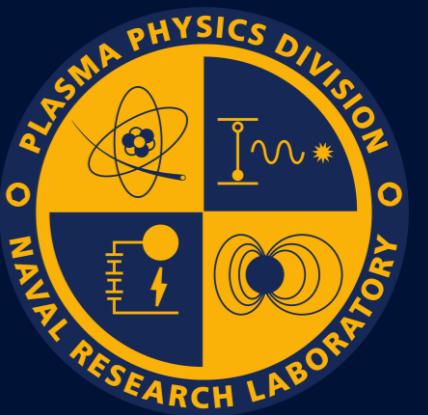
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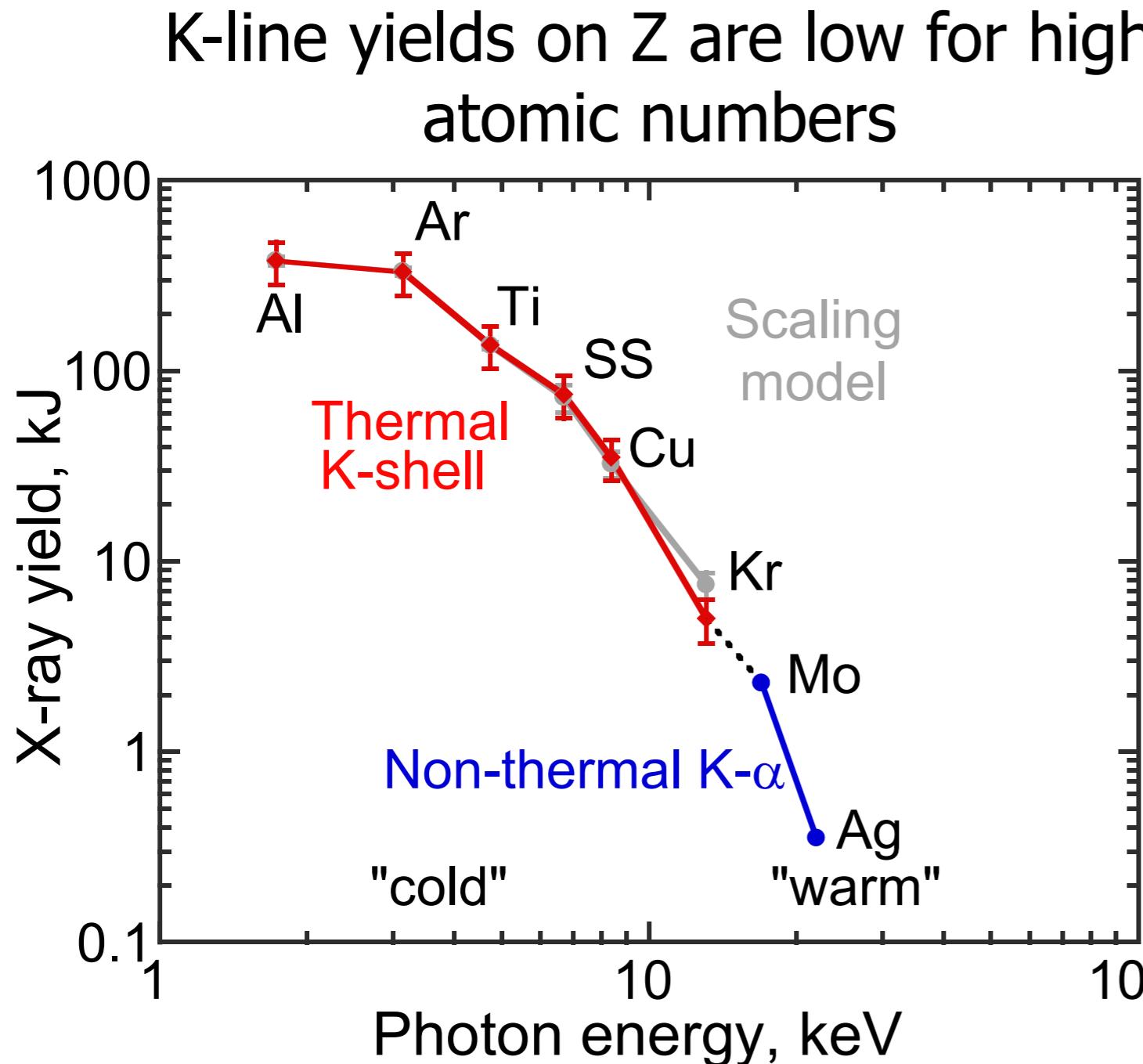
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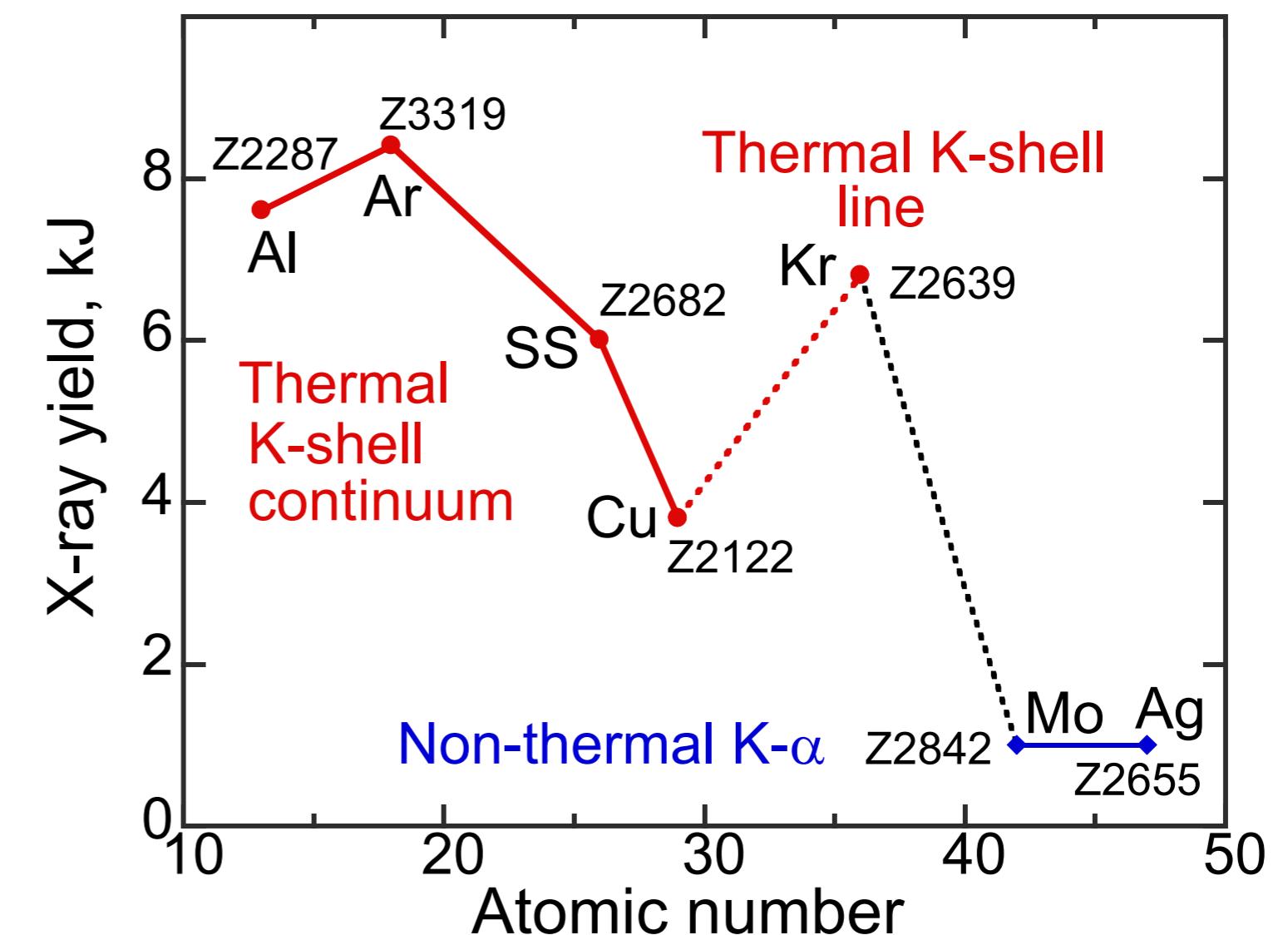


- Hard-photon direct recombination continuum alternative paths to “warm” x-ray production above 15 keV with pulsed power for next-gen pulsed power facility, NGPP
- A scaling model is needed to confidently predict the radiation yields from NGPP
 - Densities and temperatures of radiating plasmas at stagnation need to be confidently scaled up
- Scaling models are based on the conventional assumption of a uniform radiating plasma column, which is plausible for Ar but not for stainless steel
- Most likely, stainless steel emits K-shell from bright spots
- Continuum yield data from Z is consistent with the spectroscopic analysis indicating that the dense, hot bright spots containing a small fraction of the load mass radiate most of the K-shell x rays
 - Dimensions, densities, and temperatures of the bright spots can be inferred
 - The mechanism of their creation and energy replenishment is still obscure

Motivation



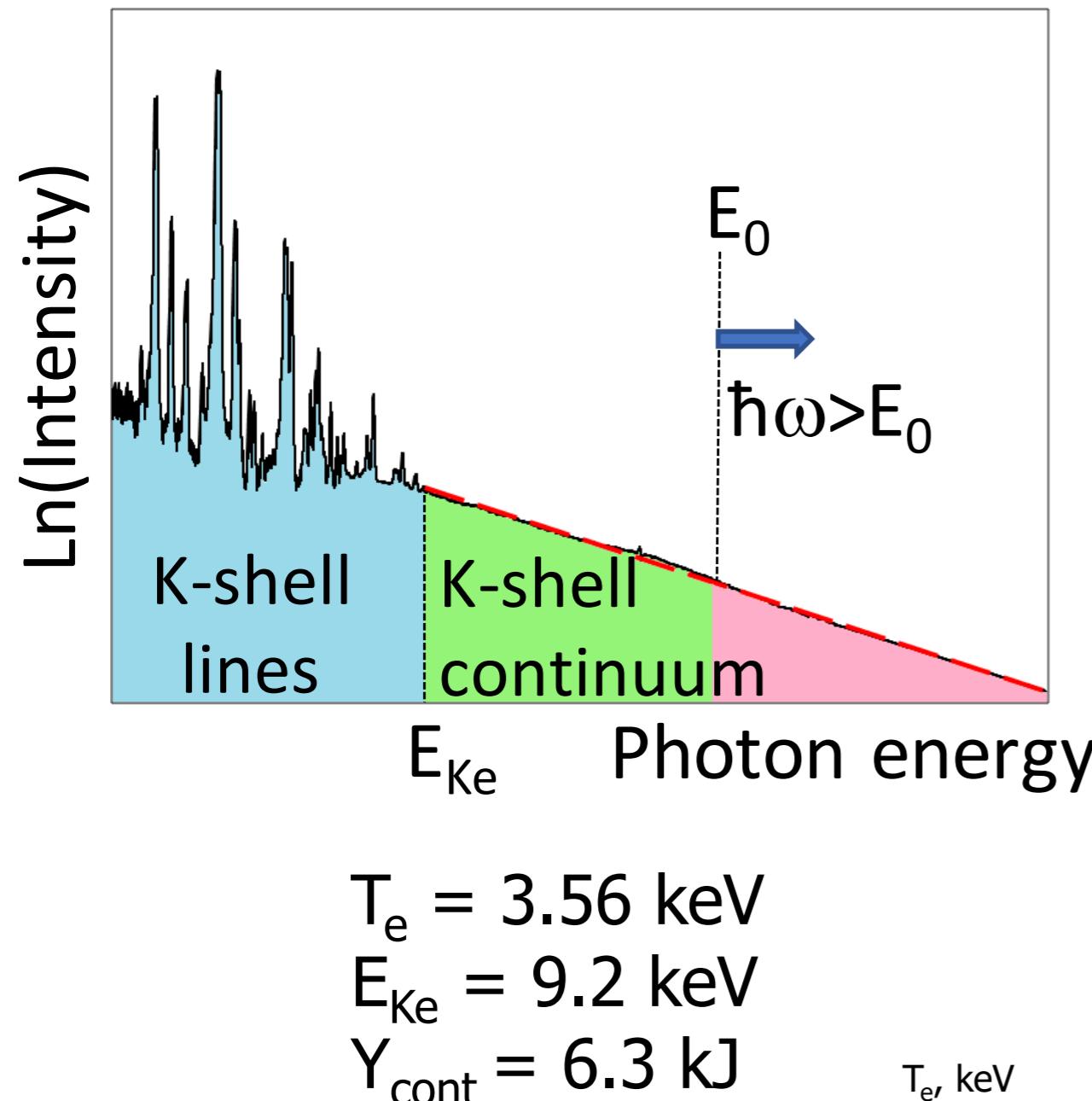
Continuum & line yields on Z in photon energies between 10 and 15 keV



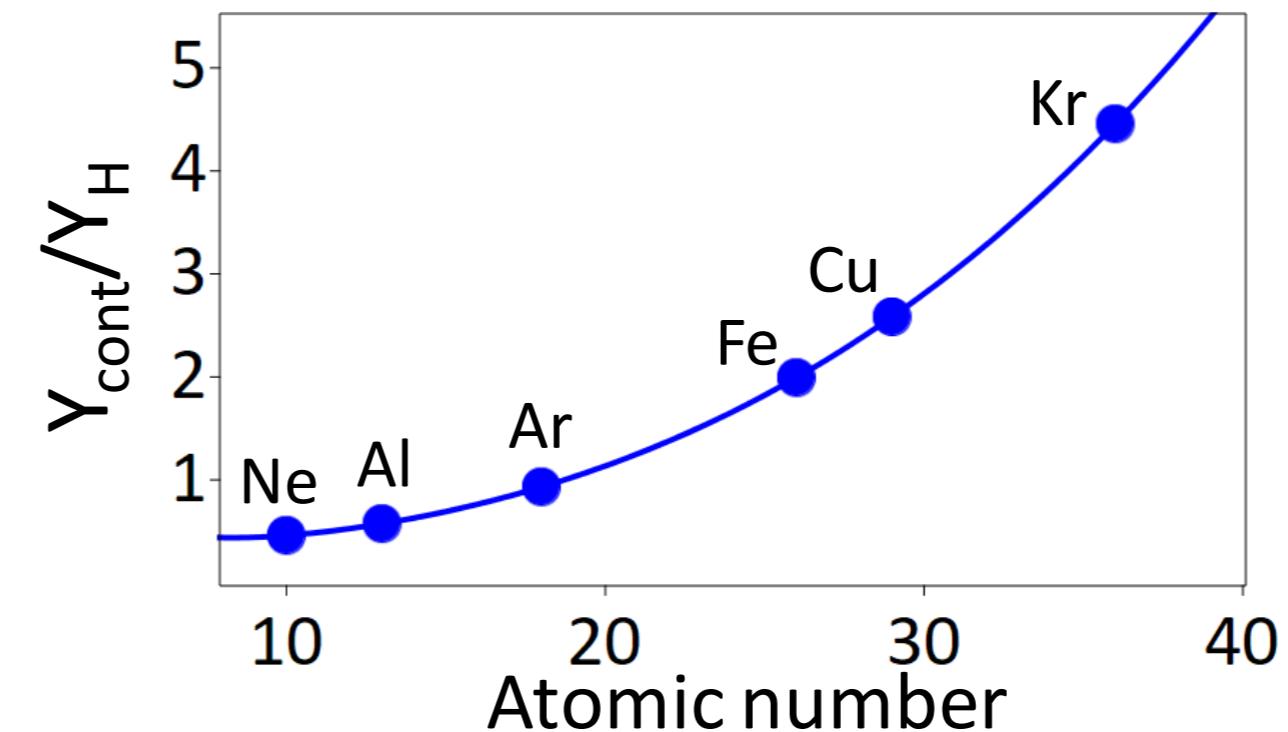
Scaling: J. Schwarz et al., "A model for K-shell x-ray yield from magnetic implosions at Sandia's Z machine," submitted to Phys. Plasmas 2022.

Background

Example: Z3151 stainless wire array



$$Y(\hbar\omega > E_0) = Y_{\text{cont}} \exp\left(-\frac{E_0 - \text{Ryd} \cdot Z^2}{T_e}\right)$$
$$\frac{Y_{\text{cont}}}{Y_H} = \frac{2\alpha^3}{3\pi\phi_H} \left(Z_A - \frac{1}{2}\right)^4 \exp\left[\frac{\text{Ryd}}{4T_e} (3Z_A - 1)(Z_A - 1)\right]$$



@ K-shell
production
optimal
temperature

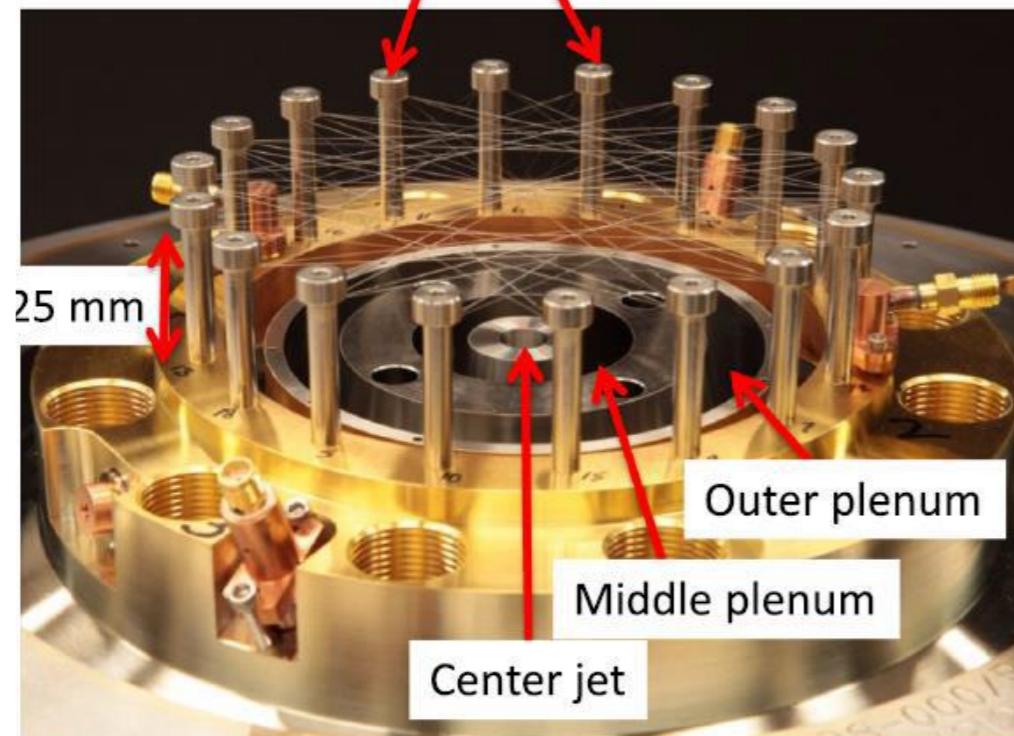
$$T_e = 0.3 \times Z_A^{2.9} \text{ eV}$$

Argon and stainless steel PRS load designs optimized for K-shell emission



Argon, 25 shots

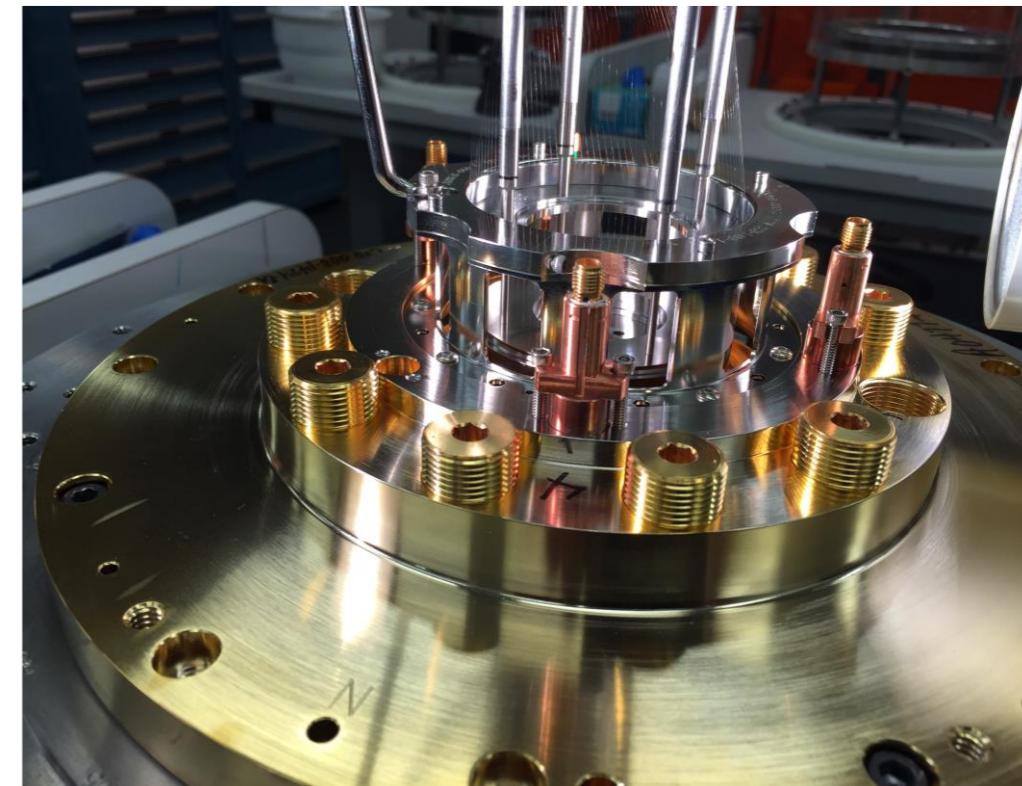
Return posts with anode mesh



- Double-shell $\varnothing 80$ mm gas puff
- Most shots w/o central jets
- Line masses $\sim 385/616$ $\mu\text{g}/\text{cm}$
- Height 3 cm

B. Jones *et al.*, Phys. Plasmas **22**, 020706 (2015); A. J. Harvey-Thompson *et al.*, Phys. Plasmas **23**, 101203 (2016); D. J. Ampleford, *et al.*, Phys. Plasmas **21**, 056708 (2014).

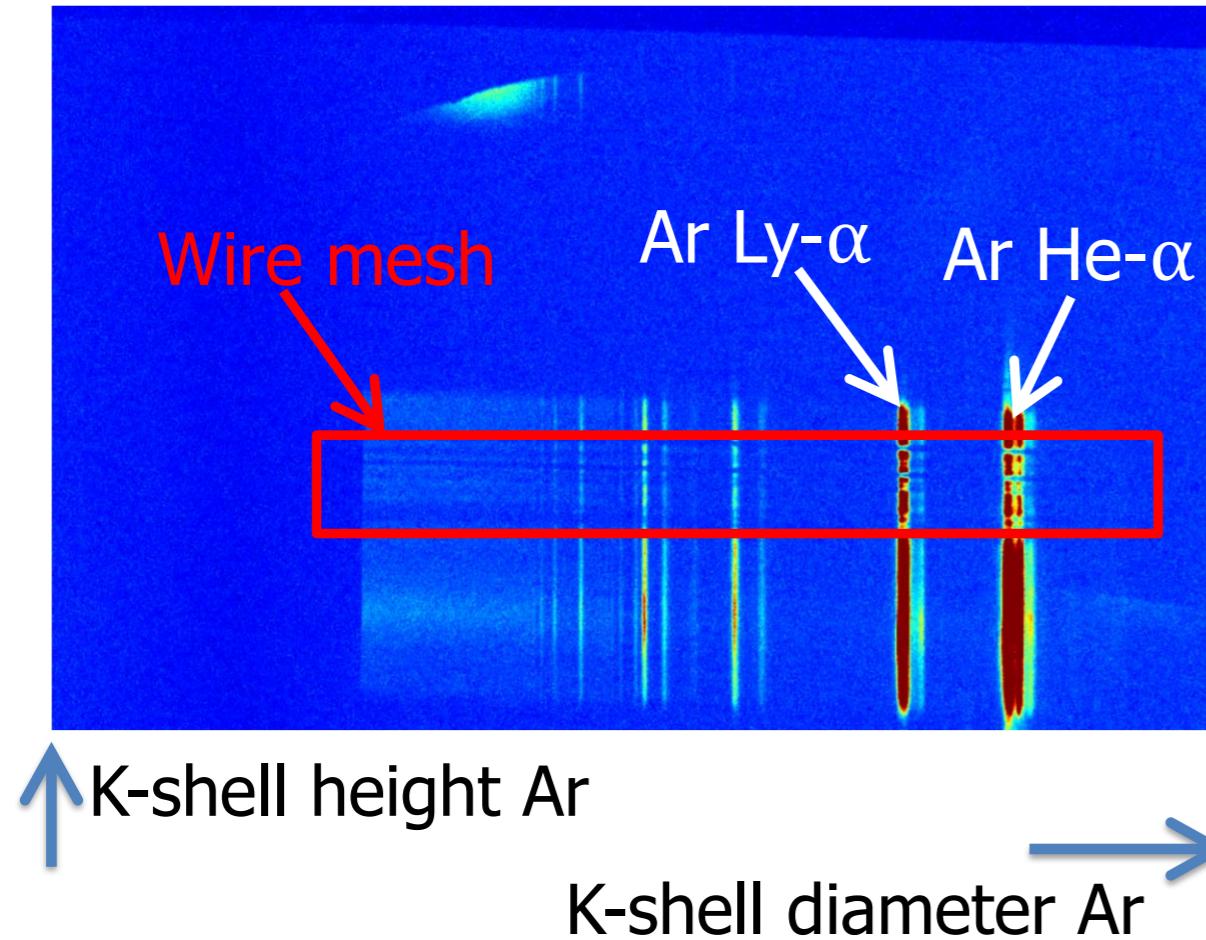
Stainless steel, 8 shots



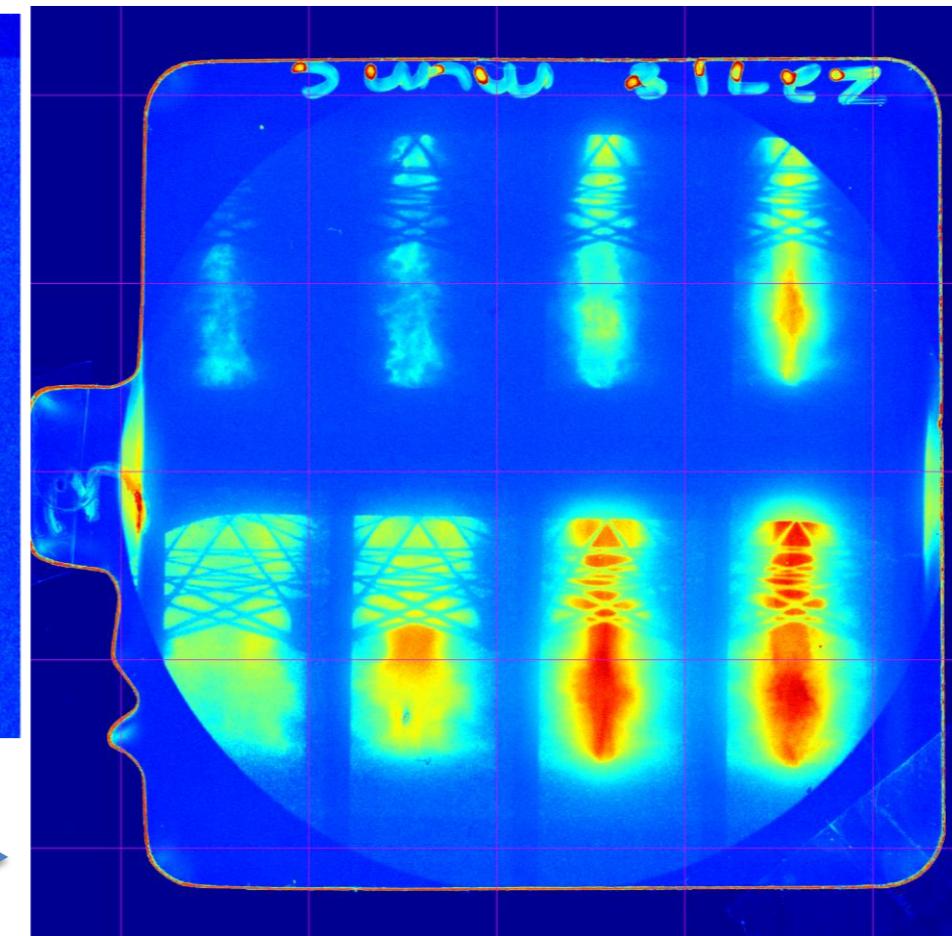
- Nested wire array, $\varnothing 70/35$ mm, 112/56 wires
- Wire diameter 8.56 μm
- Total mass ~ 1.533 mg
- Height 2 cm

Dimensions and of the K-shell emitting plasma are confidently measured

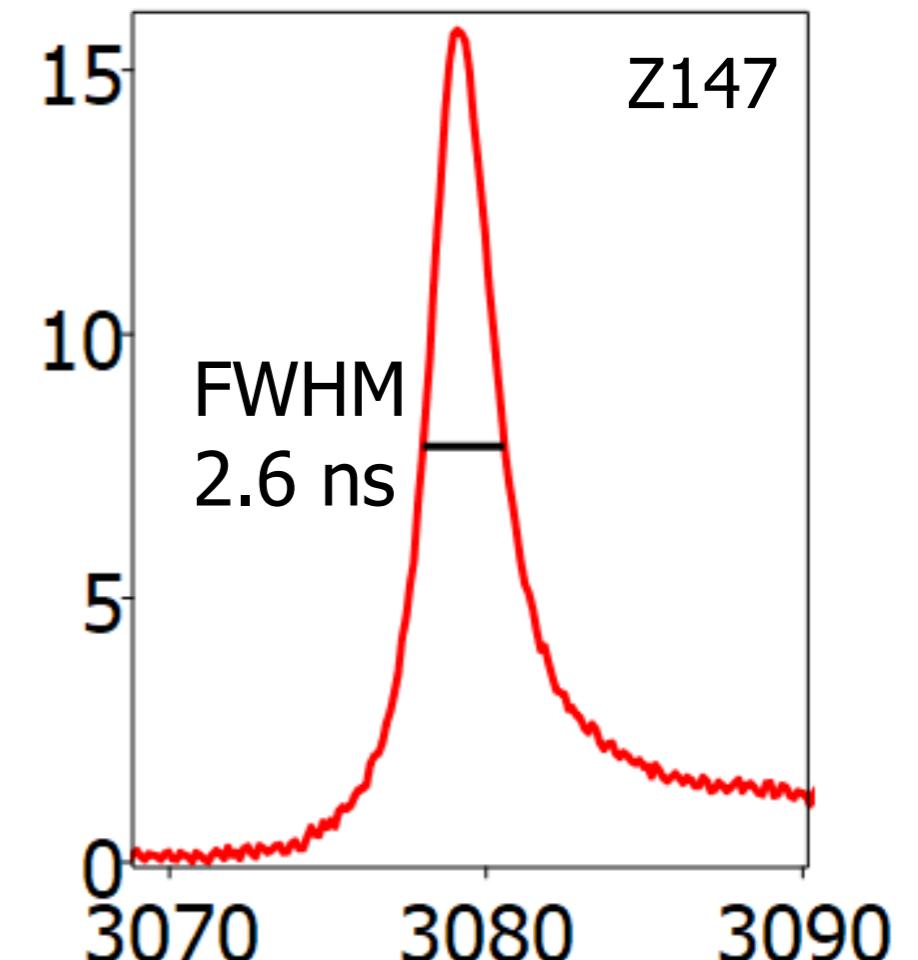
Time-integrated, axially resolved spectroscopy



Time-gated multi-layer-mirror (MLM) x-ray camera



X-ray power signal, arb. units



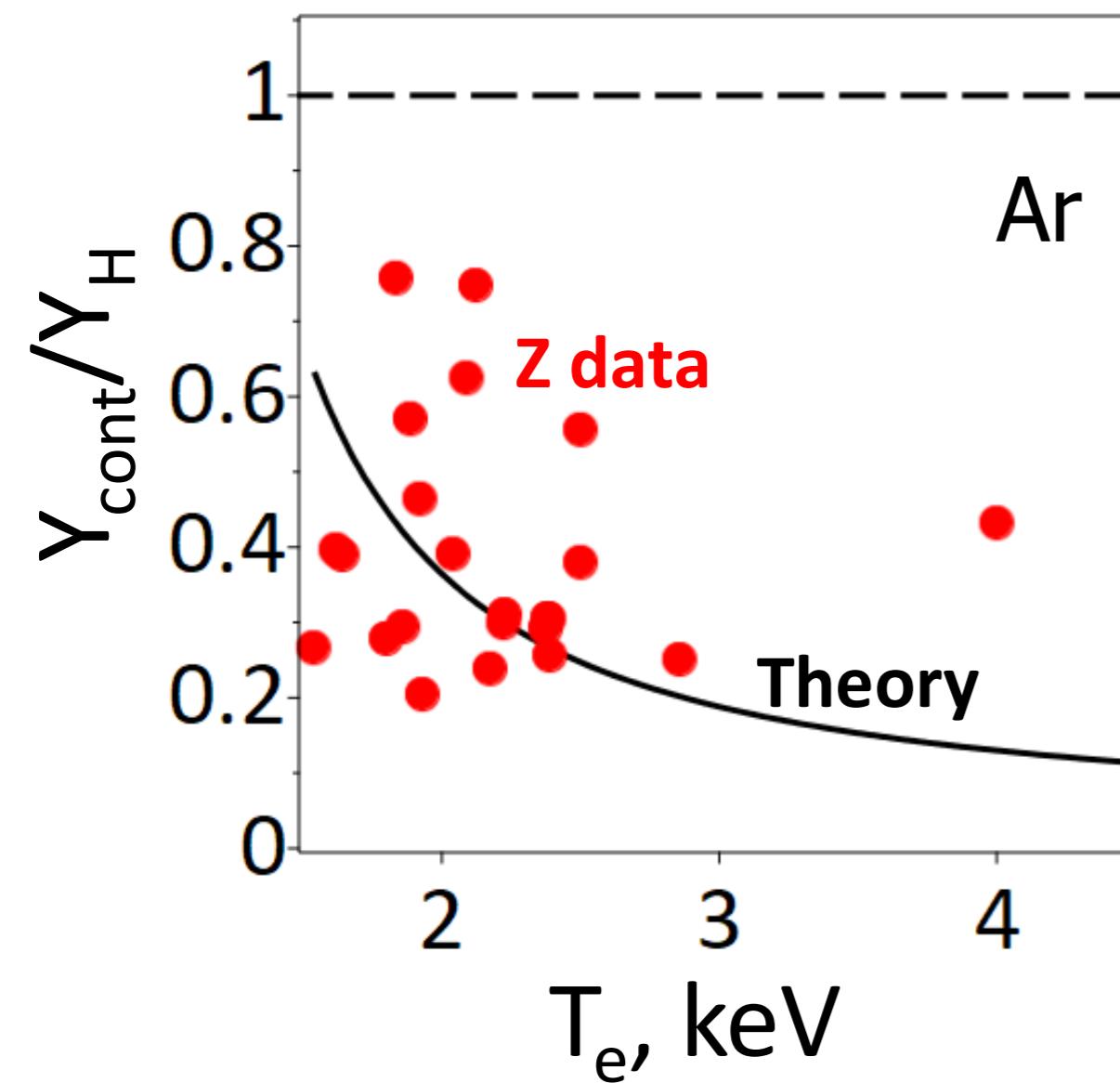
	Height, mm	Diameter, mm	Volume, cm ³	FWHM, ns
Argon	2.5 to 3.5	1 to 3	0.05 to 0.36	4 to 12
Stainless	2	0.8	0.01	2.5 to 3.5

↑
K-shell FWHM SS

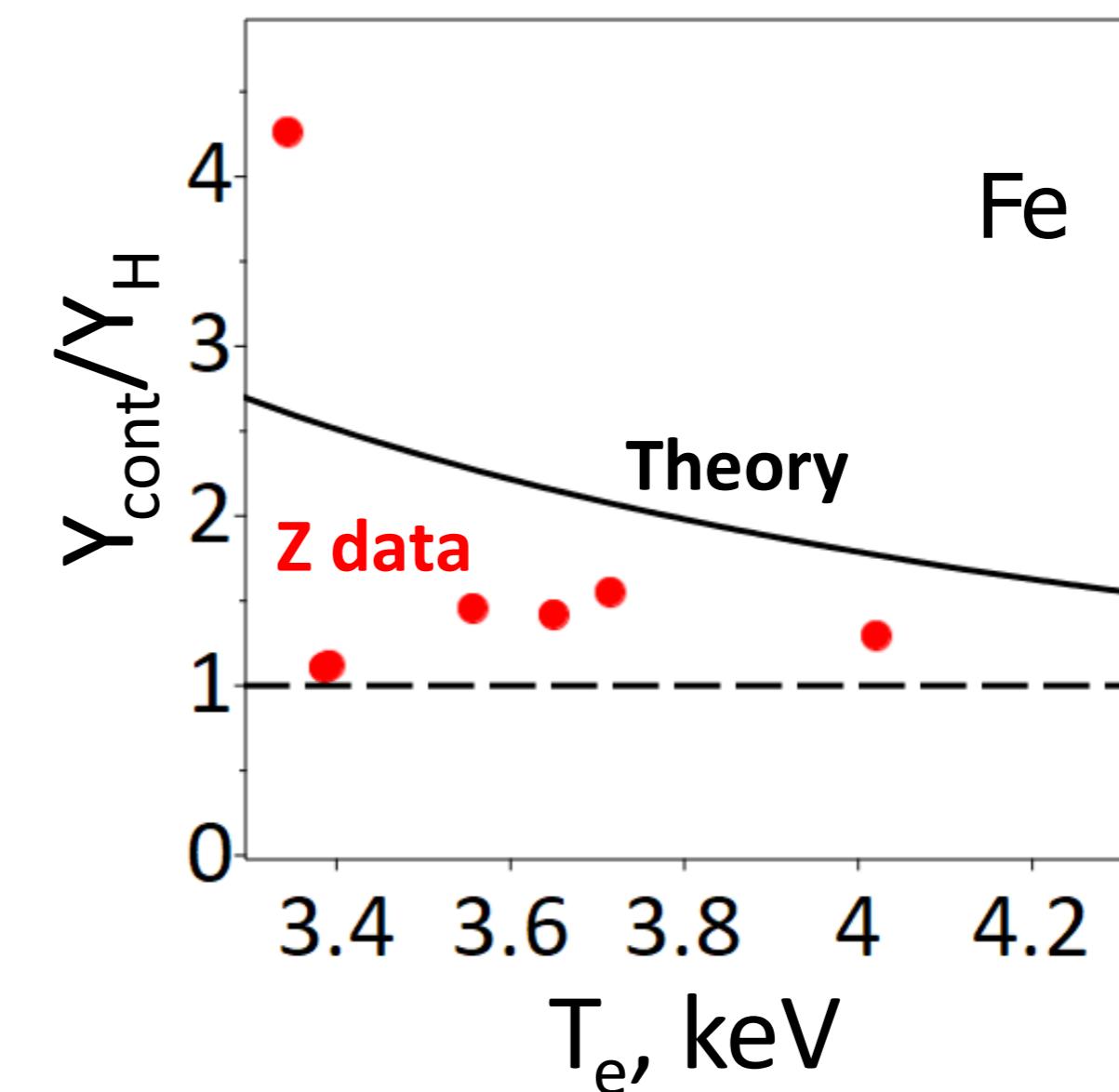
Higher-atomic-number materials radiate relatively more in the continuum



Argon: the yield in continuum is less than in Ly lines, even though bare ions contribute



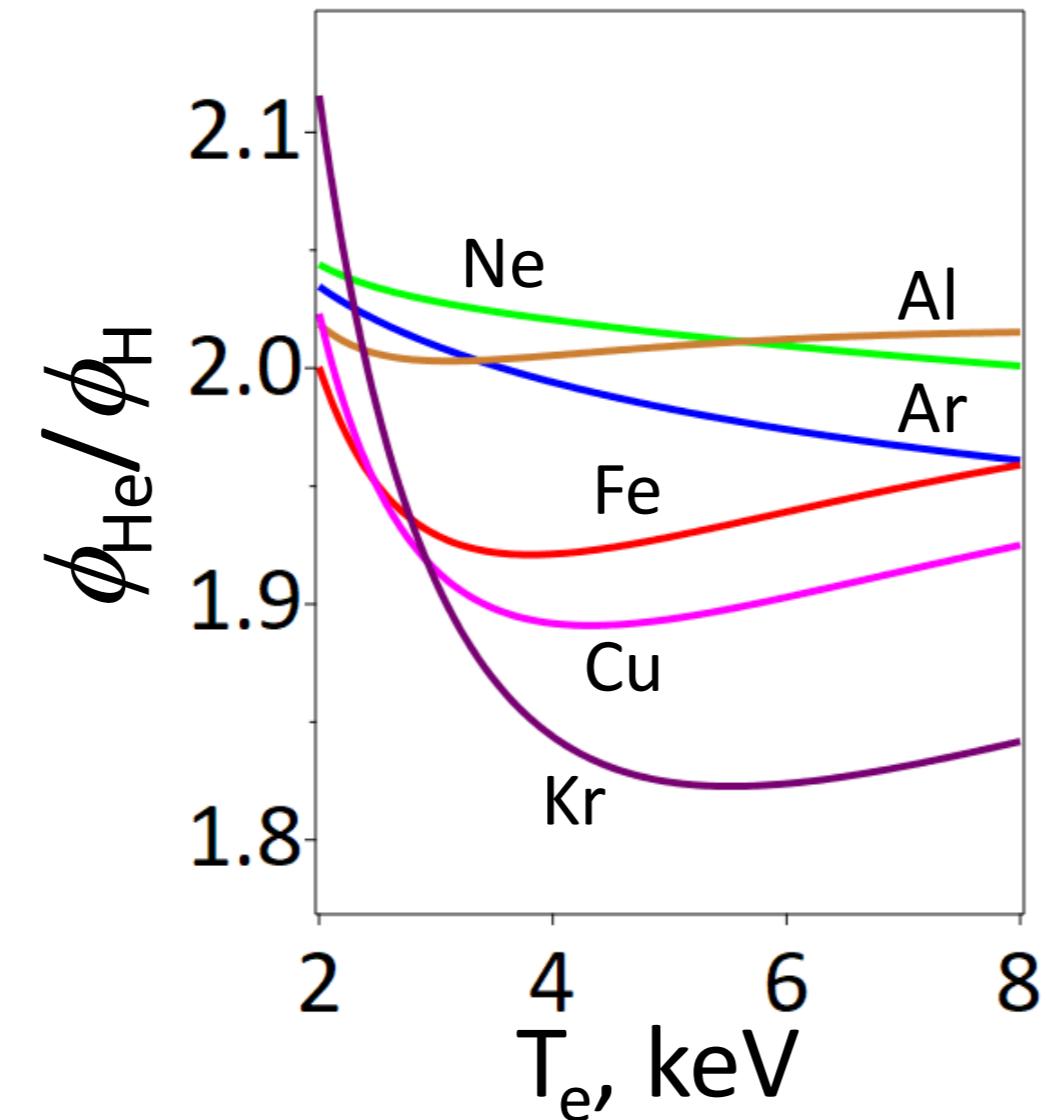
Iron: the yield in continuum is greater than in Ly lines; no contribution from bare ions



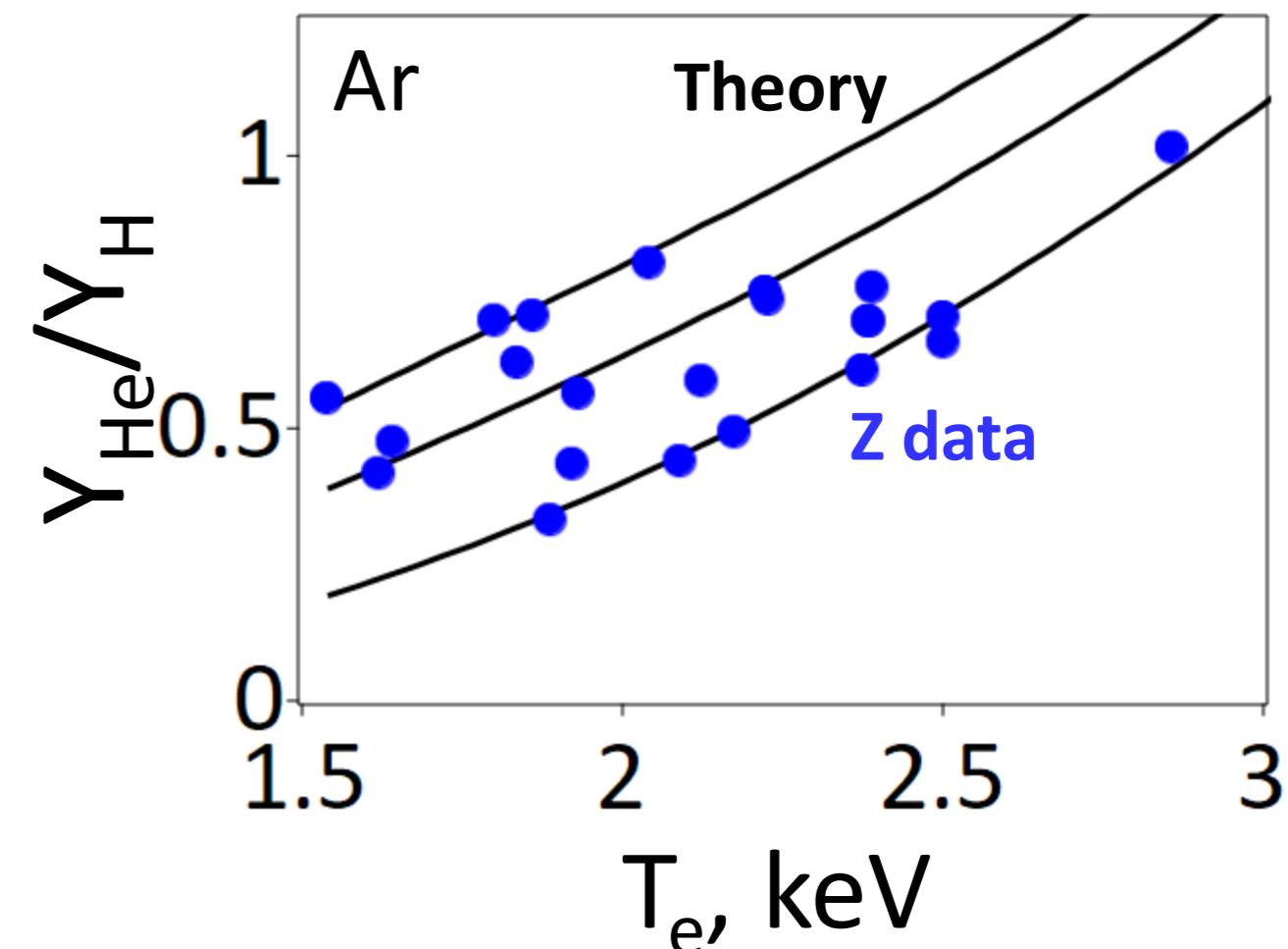
Ratio of H-like
to He-like yields

$$\frac{Y_H}{Y_{He}} = \frac{n_{i,H}}{n_{i,He}} \cdot \frac{\phi_H}{\phi_{He}}$$

Seaton's formula
fit with FAC
atomic data



Argon on Z: more H-like than He-like ions,
Yield in He lines slightly higher than in Ly lines

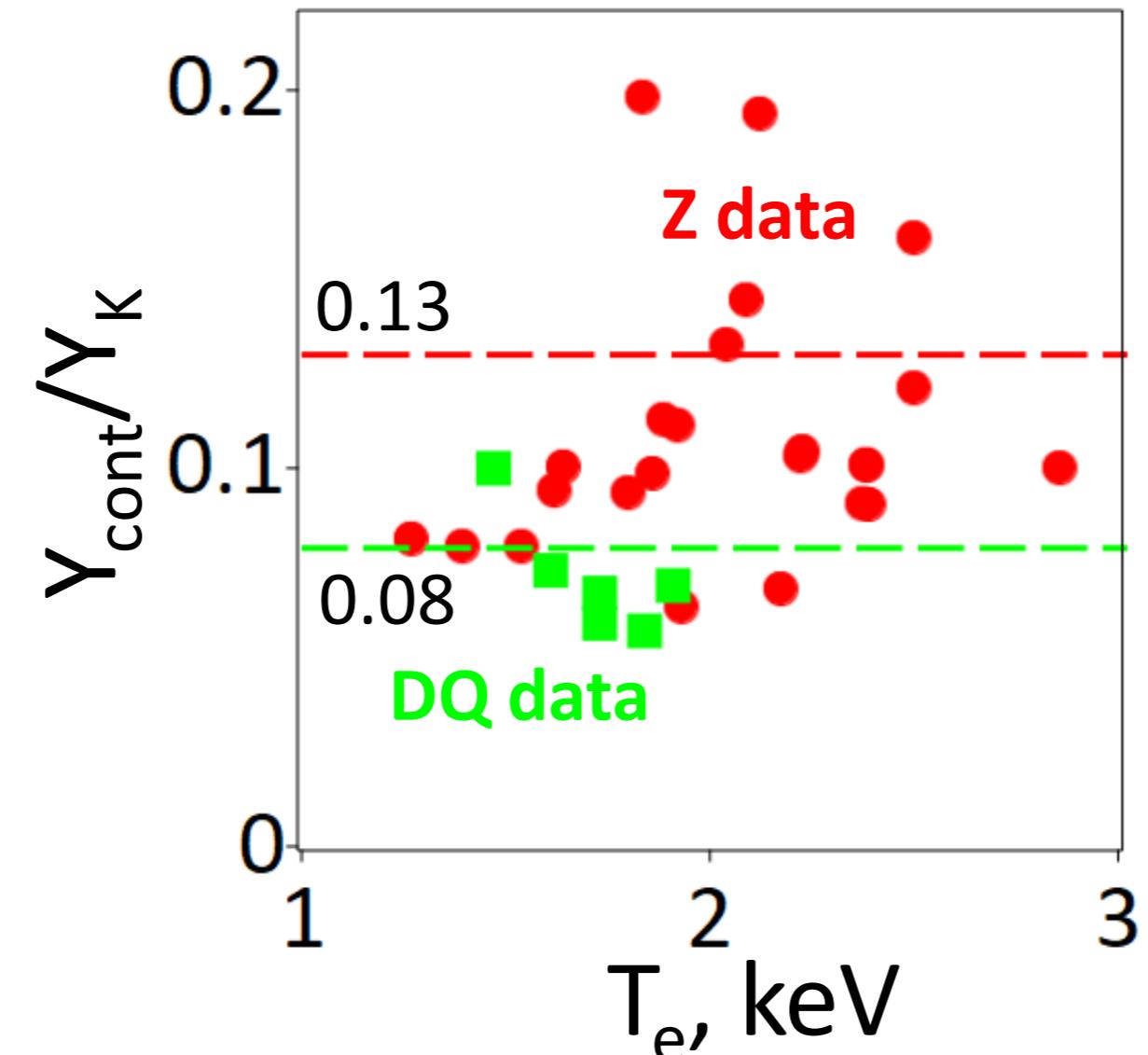


The highest hard-photon continuum fractions we can reasonably expect

Moderate- Z_A argon, accessible now

$$\begin{aligned} Z &= \\ Y_H/Y_{He} &= 0.6 \\ Y_{cont}/Y_H &= 0.4 \\ \Rightarrow & \\ Y_{cont}/Y_K &\sim 0.13 \end{aligned}$$

$$\begin{aligned} DQ &= \\ Y_H/Y_{He} &= 0.4 \\ Y_{cont}/Y_H &= 0.3 \\ \Rightarrow & \\ Y_{cont}/Y_K &\sim 0.08 \end{aligned}$$



DQ data from F. C. Young *et al.*, IEEE Trans. Plasma Sci. **34**, 2312 (2006).

High- Z_A iron or copper on NGPP

$$\begin{aligned} \text{If } Y_H/Y_{He} &= 0.5 \\ \text{Then } Y_{cont}/Y_H &> 1 \end{aligned}$$

$$\Rightarrow$$

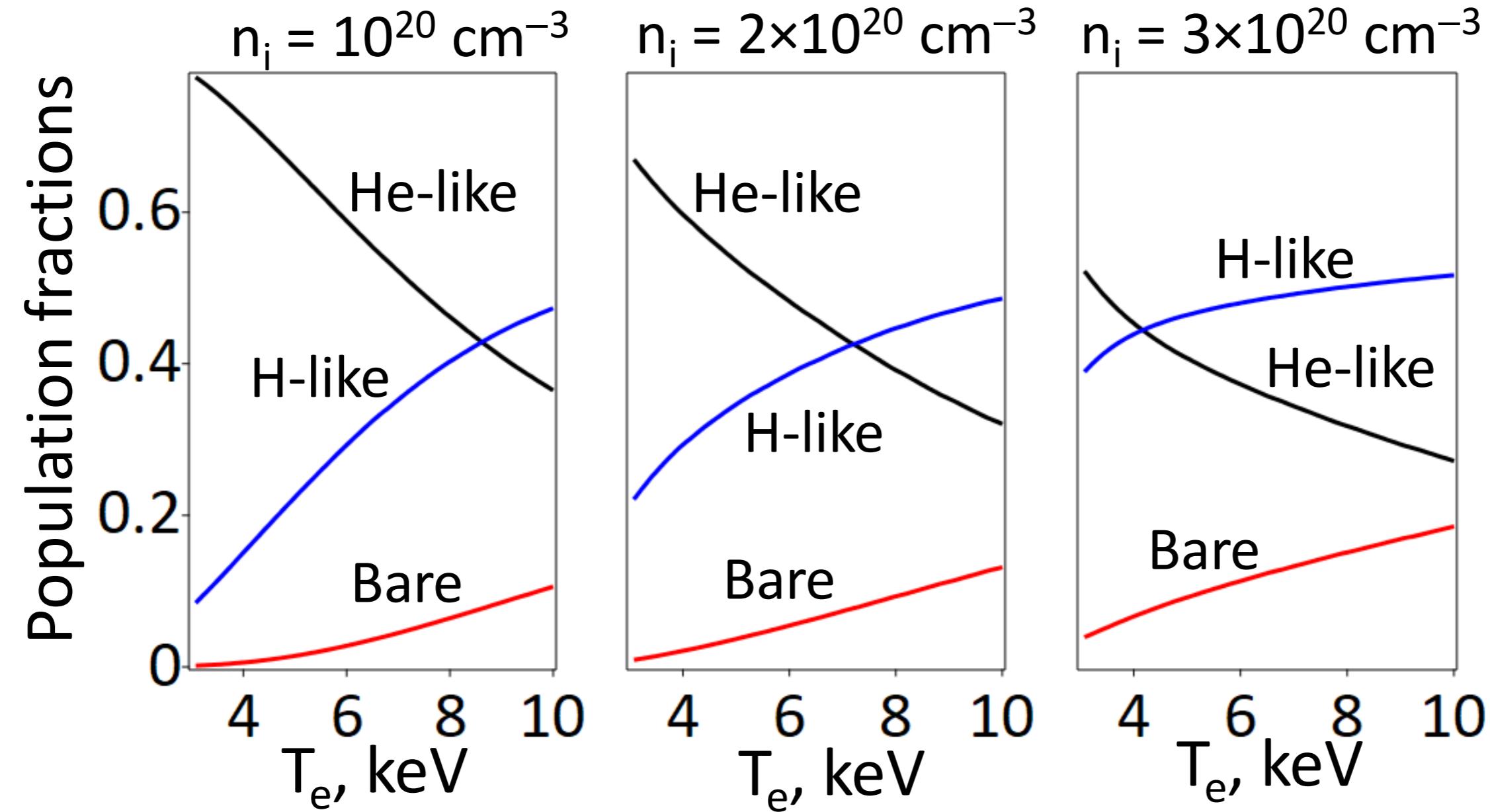
$$Y_{cont}/Y_K > 0.25$$

A 25% continuum efficiency at high Z_A is possible if we can produce an H-like population no less than He-like

We need high temperatures and densities to ionize Fe to H-like



CRE equilibrium populations calculated for a $\varnothing 3$ mm plasma cylinder with DRACHMA-II



Ion populations
are very
sensitive to the
ion densities
above
 $n_i \sim 2 \times 10^{20} \text{ cm}^{-3}$

Ion number density can be inferred from the hard-photon continuum power



$$W_{cont} = F(T)m_K n_i (2f_b + f_H)$$

Radiation power is proportional to the product of the radiating mass and ion density

$$n_i = \frac{W_{cont}}{F(T)m_K (2f_b + f_H)}$$

where the ion populations for our parameter range are estimated as

$$2f_b + f_H \approx \begin{cases} 1 \text{ for Ar} \\ f_{He} \cdot 2Y_H / Y_{He}, \quad f_{He} \approx 0.7 \text{ for SS} \end{cases}$$

$$m_K = R_K \sqrt{\frac{W_{cont} \pi m_i}{F(2f_b + f_H)}}$$

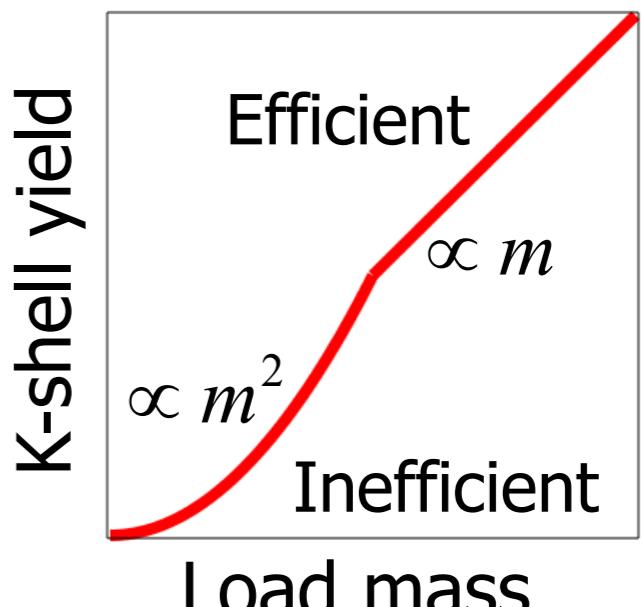
Self-consistent value of the continuum K-shell emitting mass, obtained by assuming the radiating plasma to be uniform

Assumptions for scaling line and continuum K-shell yields to higher currents



$$Y_K \propto m \cdot (\text{K-shell mass participation}) \cdot n_i$$

- Matched load mass increases as I^2 . Some options for self-similar scaling:¹
 - Mass participation and ion density stay constant²
 - Efficient K-shell emission, $Y_K \propto I^2$, radiating volume increases $\propto m$
 - Works well for argon above 5 MA
 - Mass participation and radiating volume stay constant²
 - Inefficient K-shell emission, $Y_K \propto I^4$
 - Works well for argon³ below 5 MA and for stainless⁴ up to 20 MA
 - Uncertain about stainless above 20 MA – what stays constant, what changes?



¹P. F. Schmit and D. E. Ruiz, Phys. Plasmas **27**, 062707 (2020).

²J. W. Thornhill *et al.*, IEEE Trans. Plasma Sci. **34**, 2377 (2006).

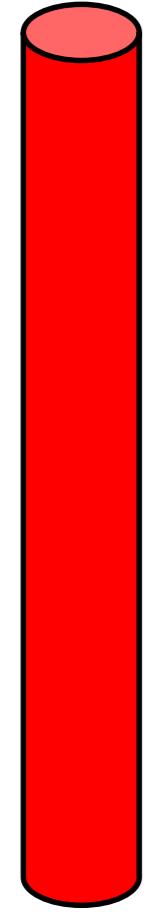
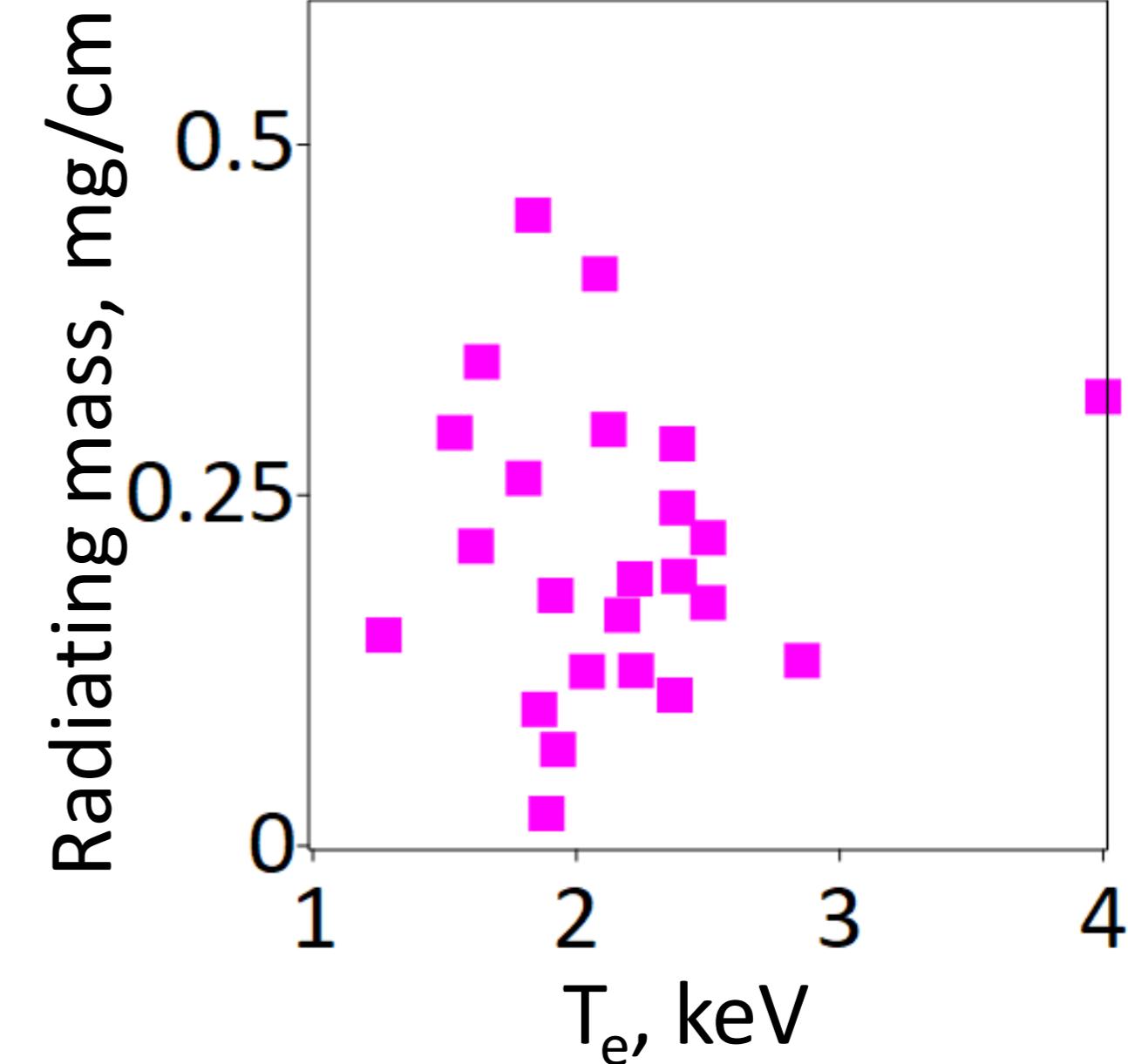
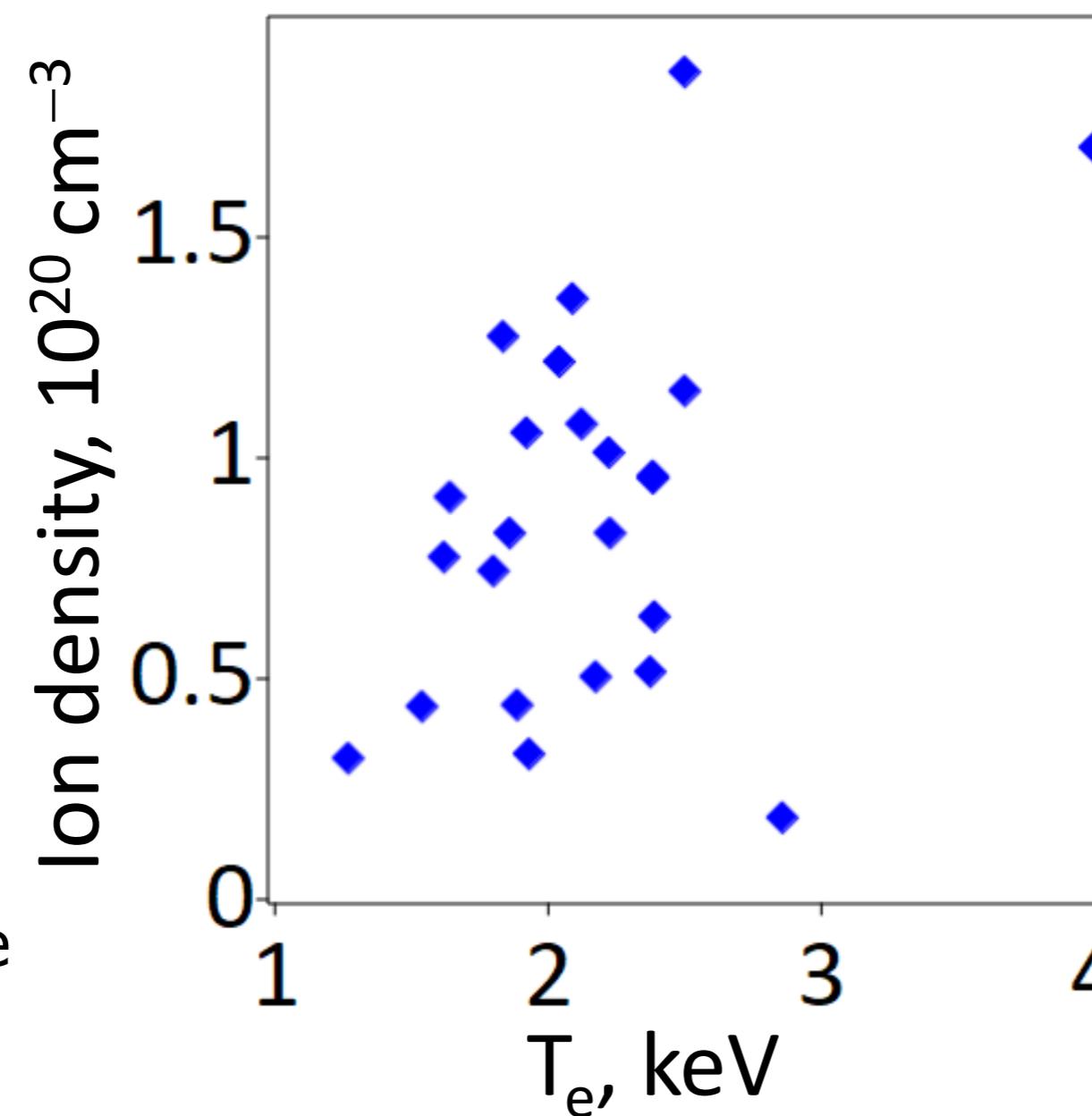
³This conference: V. Tangri *et al.*, PO 4.33 “Scaling of efficient Ar K-shell emission from fast gas-puff Z-pinch in 10 to 100 MA current range”; A. Esaulov *et al.*, PO 4.37 “Progress in the refining of the K-shell yield scaling model for Z-pinch plasma radiation sources.”

⁴J. Schwarz *et al.*, “A model for K-shell x-ray yield from magnetic implosions at Sandia’s Z machine,” submitted to Phys. Plasmas 2022.

For argon, the assumption of uniformity makes sense

For 1 mg/cm, $\emptyset 3$ mm Ar load, the assumption of a $\sim 20\%$ average mass participation in the K-shell, leads to reasonable estimates of the ion density, 10^{20} cm $^{-3}$ and line mass, 0.2 mg/cm.

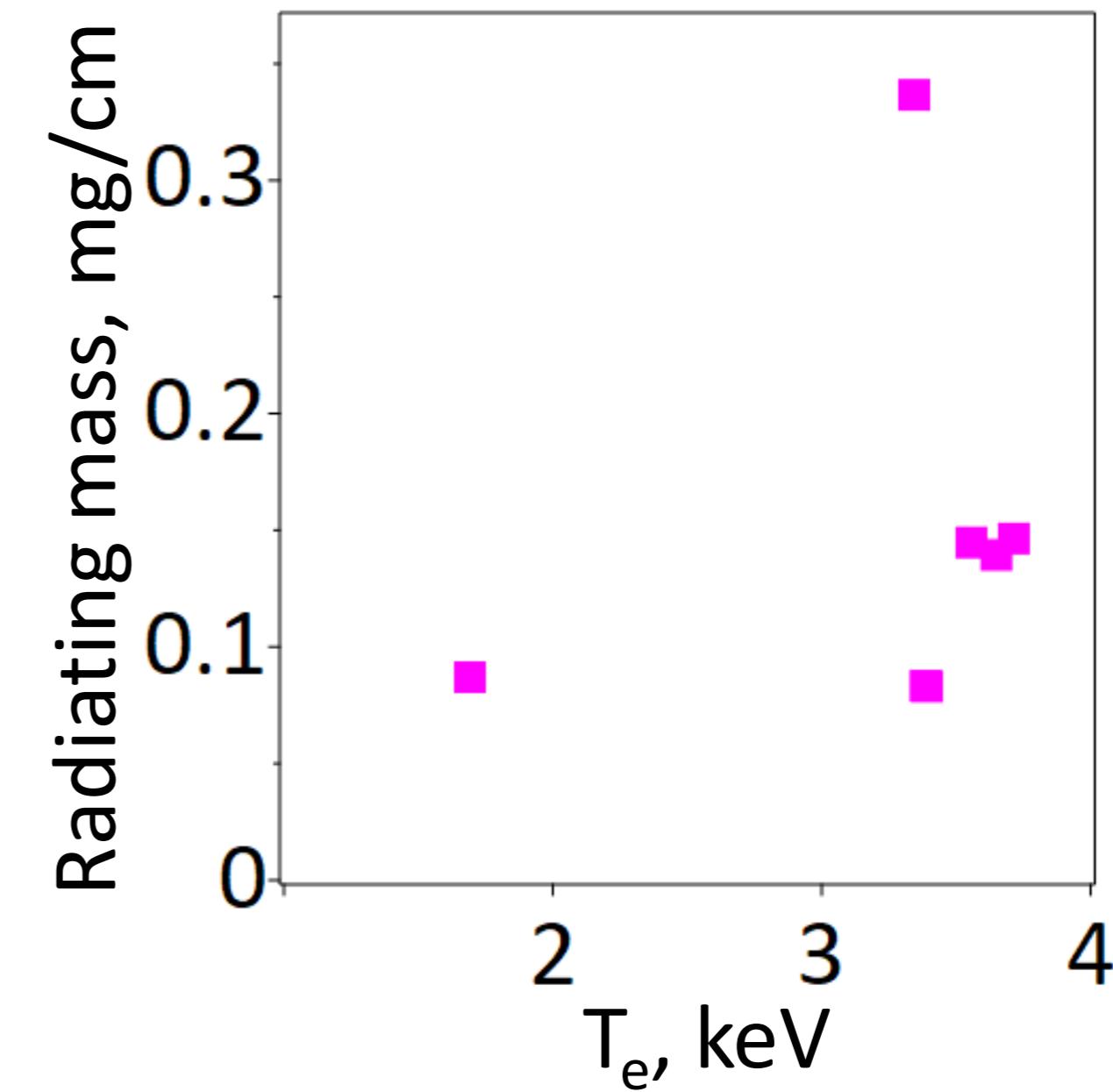
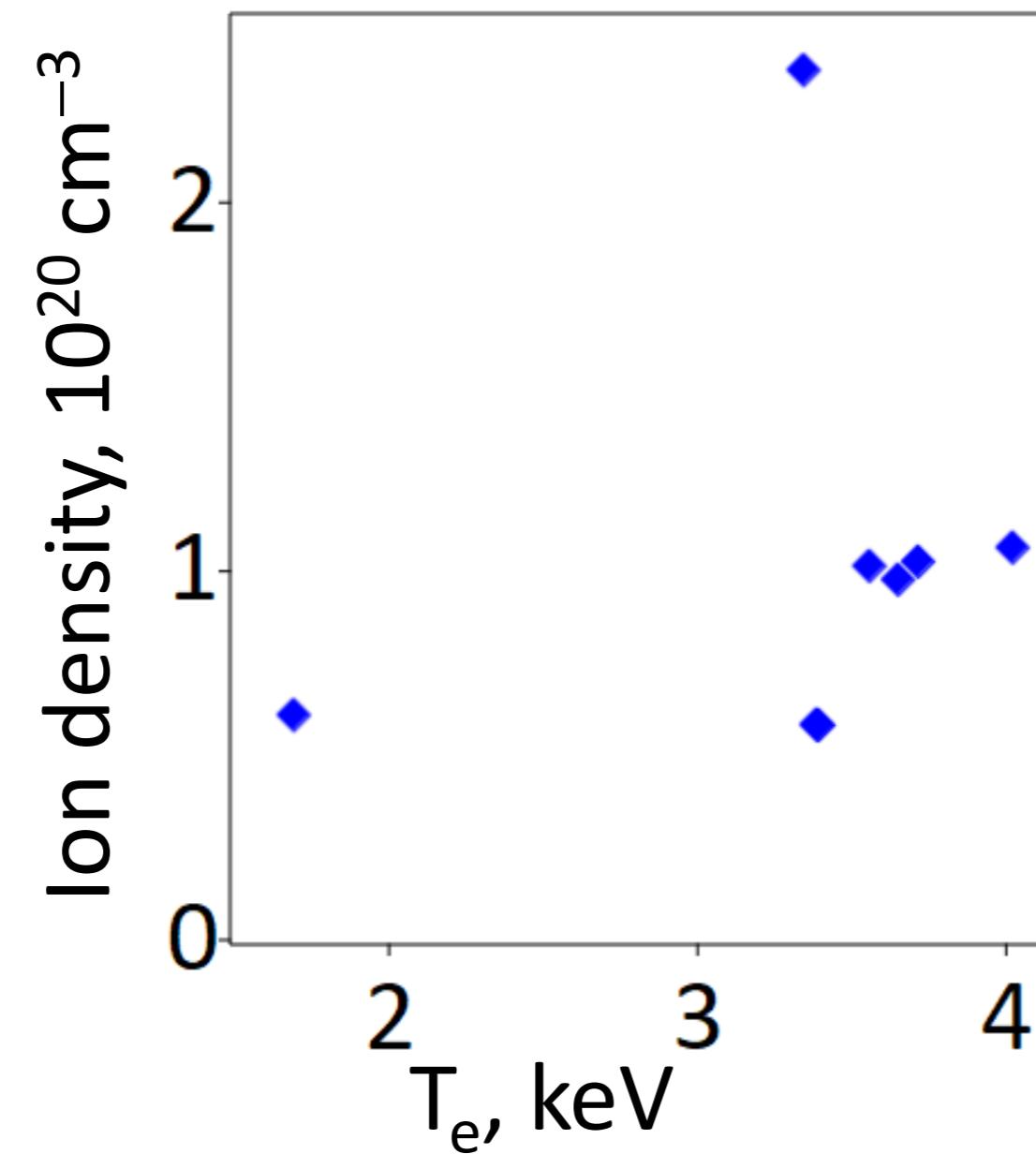
Uniform hot core plasma emits K-shell photons

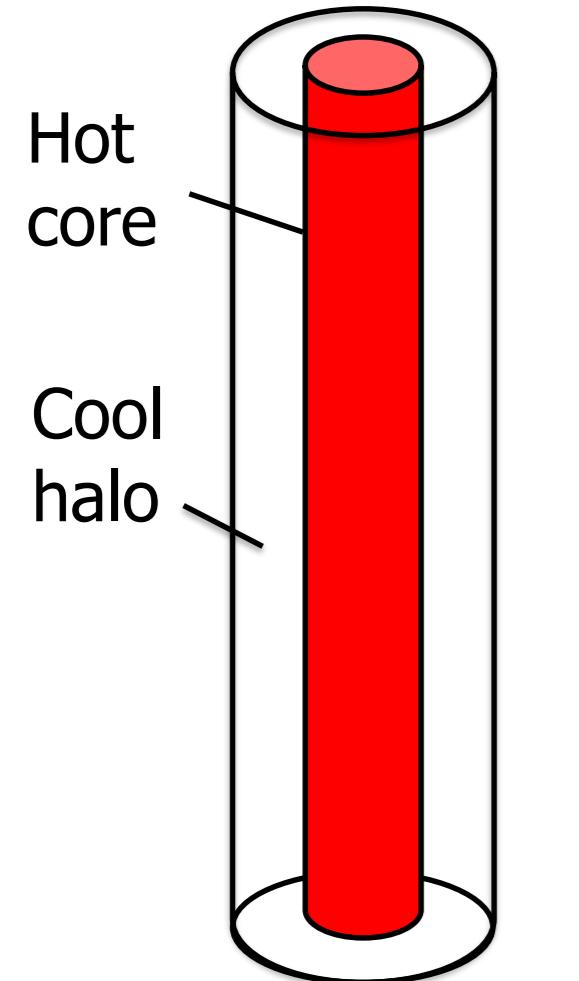
For stainless, the uniformity implies high mass participation, same as in Ar

For 0.75 mg/cm, $\emptyset 1.4$ mm Ar load, the assumption of a $\sim 18\%$ average mass participation in the K-shell, leads to estimates of the ion density, $\sim 10^{20}$ cm $^{-3}$ and line mass, 0.14 mg/cm.

Estimated SS plasma parameters are close to argon's, but:
 The radius is twice smaller and the temperature is $\sim 40\%$ higher



Time-resolved spectroscopy agrees with this early in the K-shell pulse



Cool halo absorbs some K-shell photons, affecting observed spectra

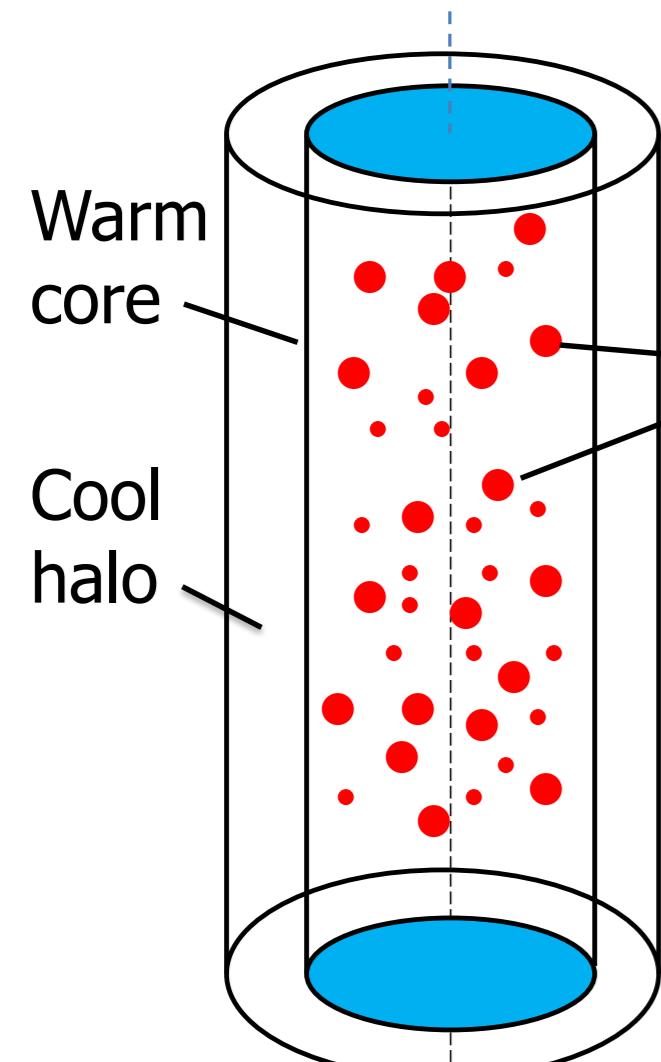
- Analysis of time-resolved spectroscopy from Z stainless nested wire-array shots done at Weizmann in 2014-2016
- The ratios of the FWHMs of resonant lines of Fe, Cr, Ni, and Mn, that are of different abundances, are analyzed. It also requires the determination of Doppler broadening for each element, and the correction for splitting effects.
- At early time, the spectra are consistent with a hot core surrounded with a cool halo

t, ns	Cylinder diameter, mm	n_i, cm^{-3}	$T_{e'}$ keV	Mass participation, %
-2.7	1.8	3×10^{19}	2	12
-1.8	1.6	5×10^{19}	3	17

Near the K-shell peak, the K-line shapes are consistent with coming from a large number of small hot spots: higher density and higher temperature



The spot parameters are constrained by the low opacity inferred and the requirement to generate the entire K power



Bright spots emit all K-shell photons, warm core background emits softer photons

$t = -0.8 \text{ ns}, T_e = 4 \text{ keV}, T_i(\text{hydro}) = 40 \text{ keV}$

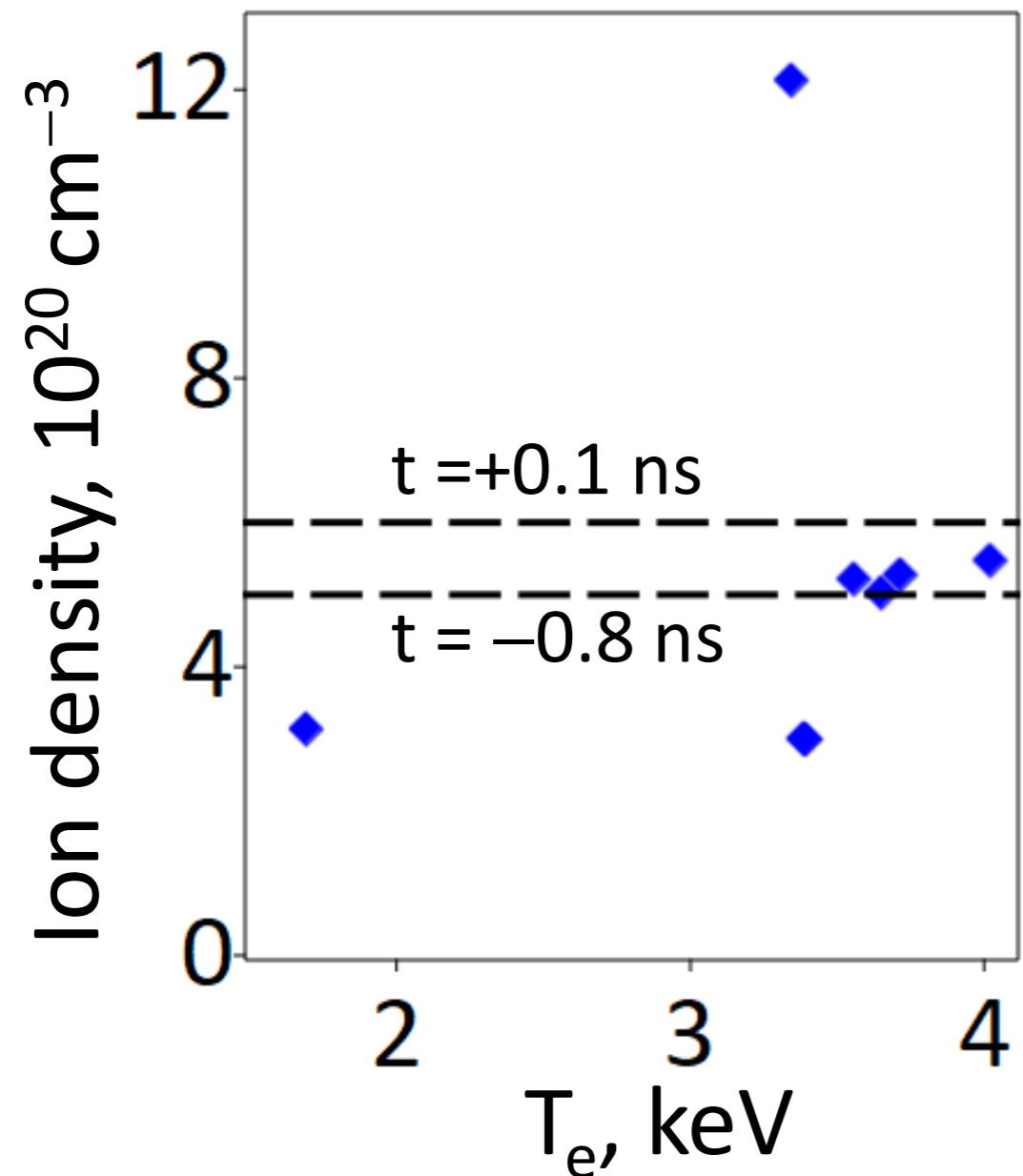
D, μm	No. of spots	n_i, cm^{-3}	Mass participation, %	Photon mfp, mm
50	8900	5×10^{20}	3.1	1.3
100	4200	2.5×10^{20}	5.9	0.68
150	2930	1.67×10^{20}	9.2	0.43

$t = +0.1 \text{ ns}, T_e = 6 \text{ keV}, T_i(\text{hydro}) = 30 \text{ keV}$

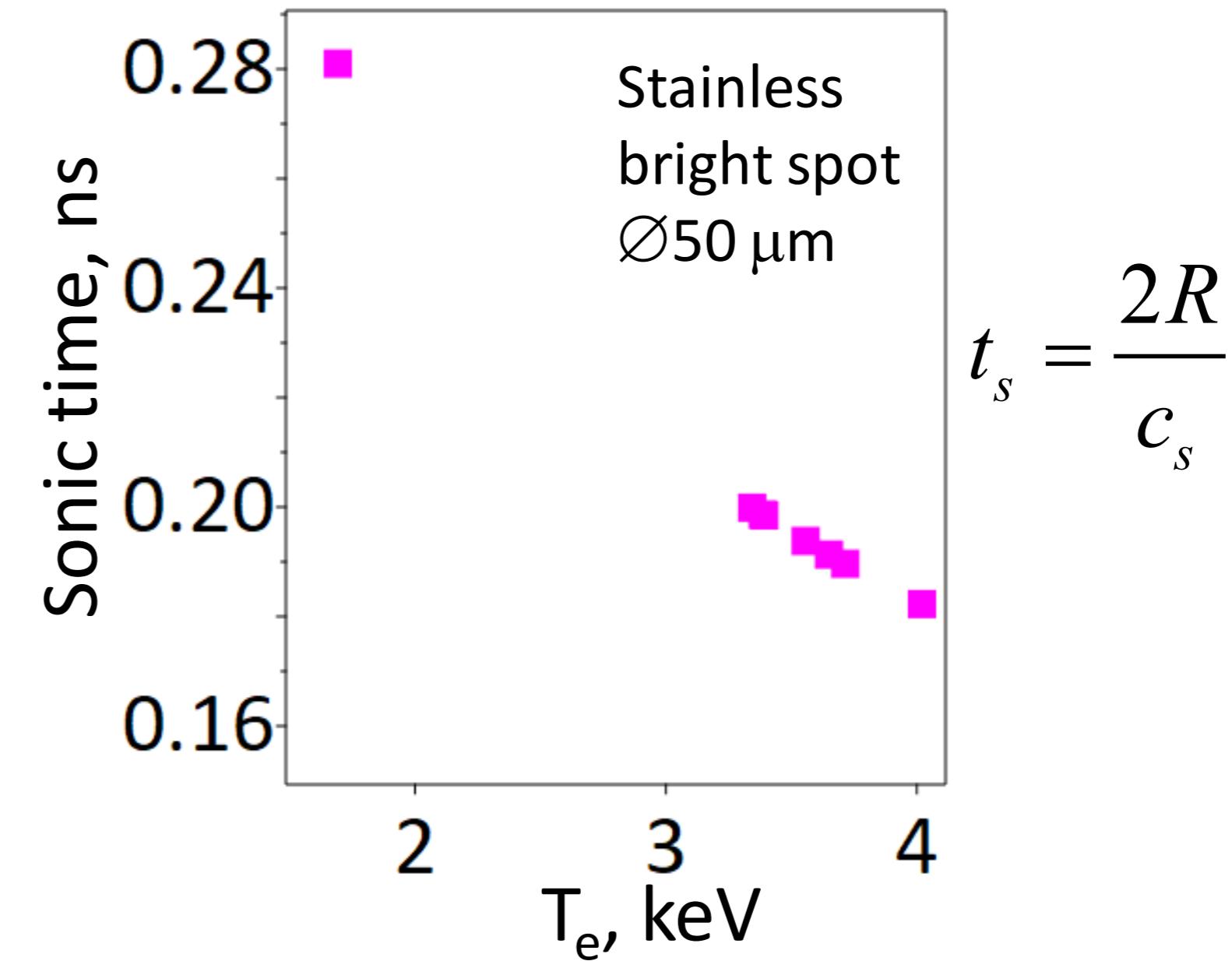
D, μm	No. of spots	n_i, cm^{-3}	Mass participation, %	Photon mfp, mm
50	8200	6×10^{20}	3.4	1.4
100	3800	3×10^{20}	6.4	0.76
150	2350	2×10^{20}	9.2	0.54

Continuum yield are also consistent with this assumption

Stainless mass participation 3.5%



Life time of spots is less than 1 ns. They must be replaced or re-energized.



Summary/Conclusions

- Hard-photon direct recombination continuum alternative paths to “warm” x-ray production above 15 keV with pulsed power for next-gen pulsed power facility, NGPP
- A scaling model is needed to confidently predict the radiation yields from NGPP
 - Densities and temperatures of radiating plasmas at stagnation need to be confidently scaled up
- Scaling models are based on the conventional assumption of a uniform radiating plasma column, which is plausible for Ar but not for stainless steel
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