



1-D MHD simulation of an argon gas puff implosion with time-dependent non-LTE kinetics*

N. D. Quart, A. Dasgupta, A. L. Velikovich, J. L. Giuliani[†], and V. Tangri

Plasma Physics Division, U.S. Naval Research Laboratory, Washington, DC 20375 USA

[†]NRL Voluntary Emeritus Program

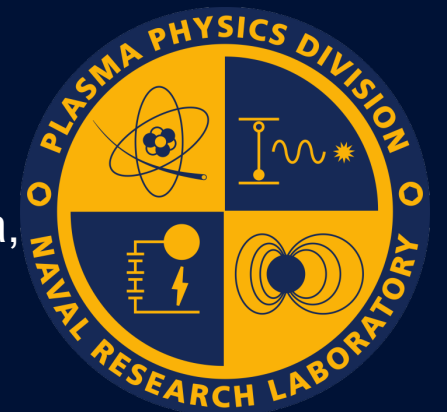
M.-A. Schaeuble, J. Schwarz, D. Ampleford, R. A. Vesey, C. Jennings, and B. Jones

Sandia National Laboratories[†], Albuquerque, NM 87185 USA

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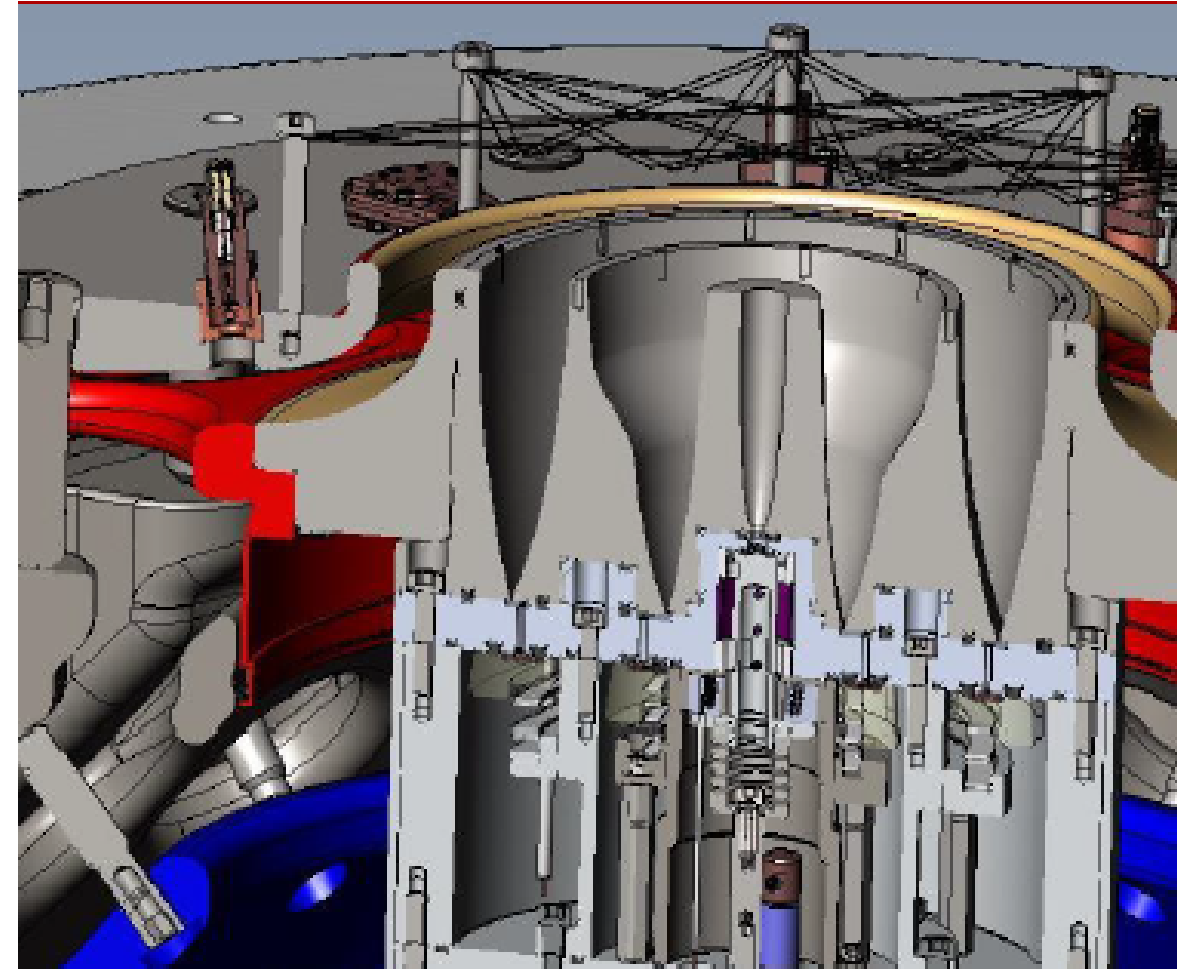
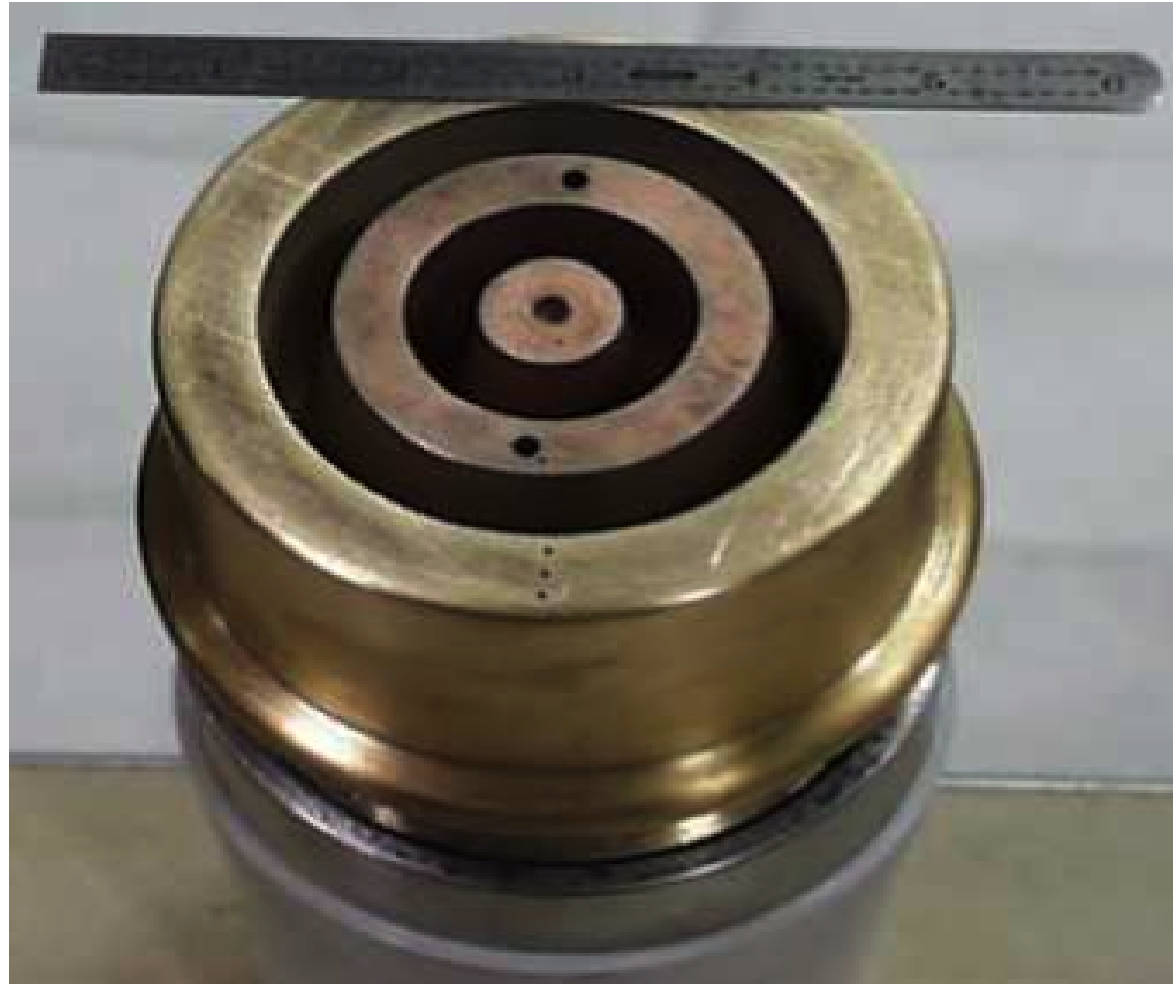


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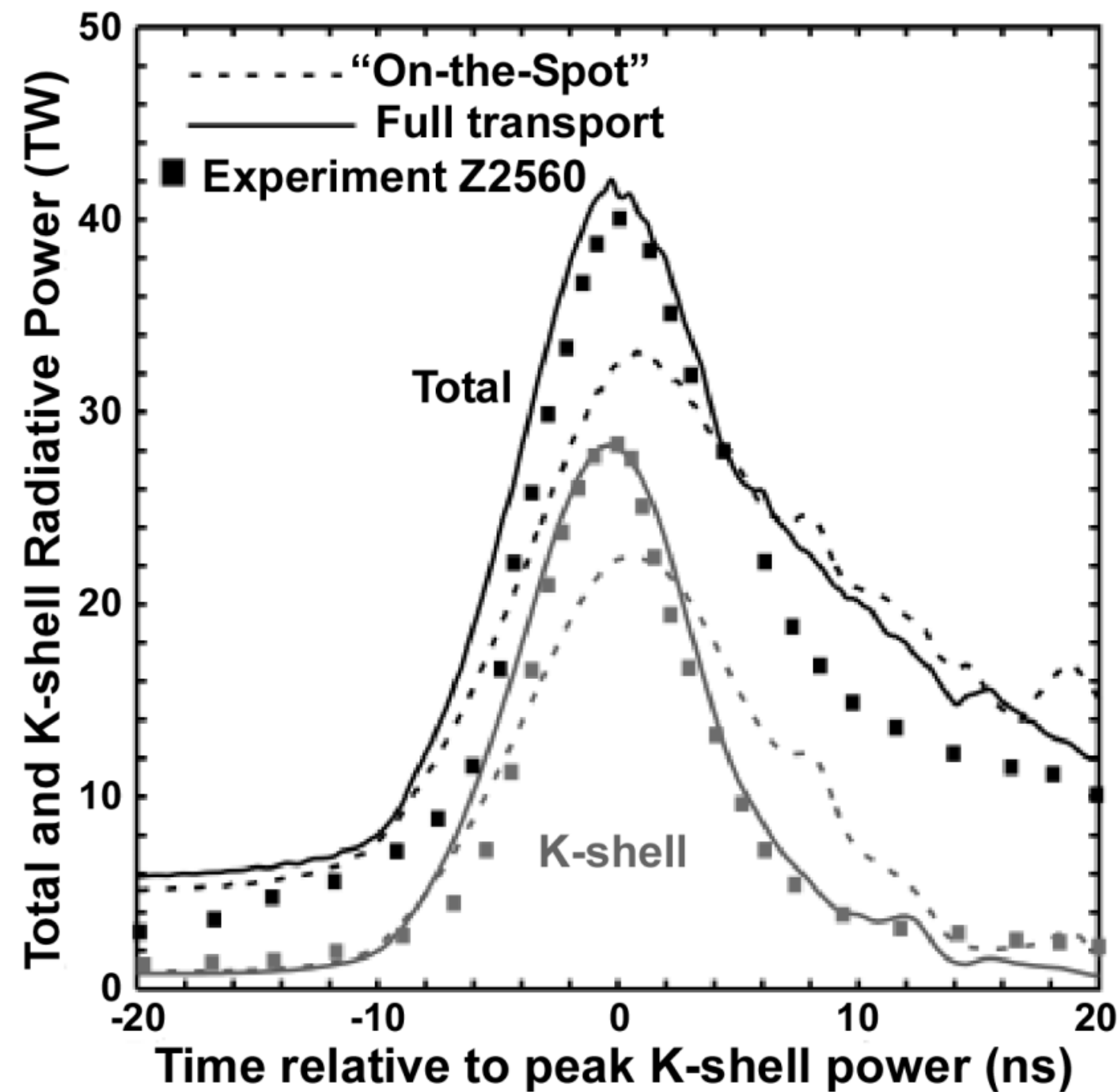
- Argon gas puff experiments on the Z machine have produced $330 \text{ kJ} \pm 9\%$ above 3 keV photon energies
- Simulations using Mach2-TCRE were able to reproduce the radiative powers and yield from shot Z2560
 - The Mach2-TCRE simulation's Ly- α /He- α +IC line ratio had good agreement after peak power, but was higher prior to peak power. This was attributed to 3-D effects and the implicit assumption of steady-state populations in the TCRE table.
- The 1-D NRL DZAPP code was used to assess the effects of time-dependent ground state level populations compared to steady-state for the Ly- α /He- α +IC line ratio
 - There was a decrease in the line ratio in the case of the time-dependent ground state level populations
 - However, the line ratio was still larger compared with experiment.
 - Time-dependent level kinetics and multi-dimensional effects may both be important.
- Future work will assess the effects of time-dependent level kinetics for NGPP loads

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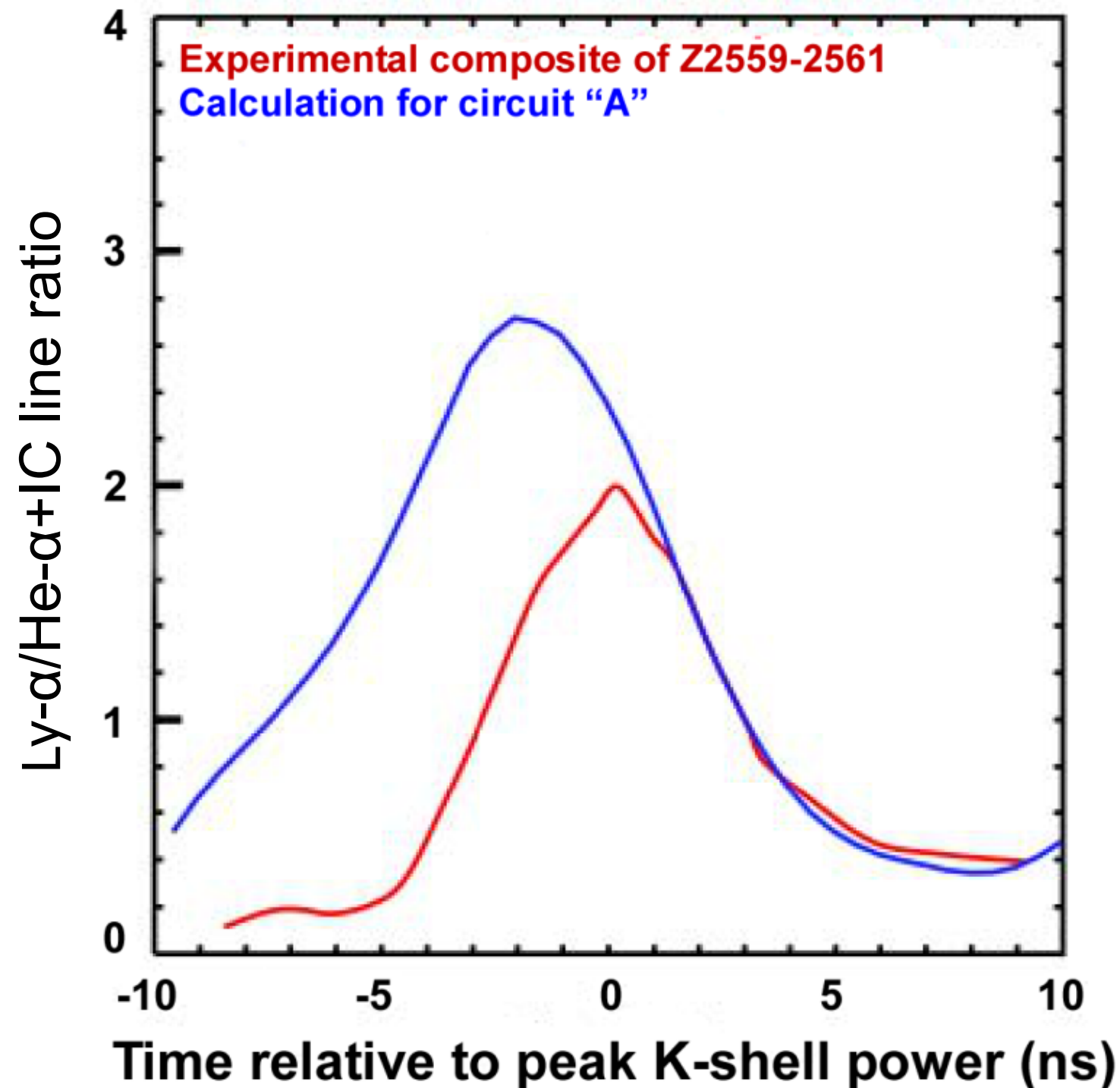


An 8 cm diameter double-shell gas nozzle was used.
The mass of the argon gas was 1 mg/cm.
The inner- to outer-shell mass ratio was 1.6:1.
The length was 2.5 cm.

Simulations using Mach2-TCRE were able to reproduce the radiative powers and yield from shot Z2560



The simulation's Ly- α /He- α +IC line ratio had good agreement after peak power, but was higher prior to peak power.



Label	Upper	Lower	Energy
IC	1s2p 3P_1	1s 2 1S_0	3.124 keV
He- α	1s2p 1P_1	1s 2 1S_0	3.140 keV
Ly- α_2	2p $^2P_{1/2}$	1s $^2S_{1/2}$	3.318 keV
Ly- α_1	2p $^2P_{3/2}$	1s $^2S_{1/2}$	3.323 keV

Thornhill *et al.* attributed this discrepancy prior to peak power to:

1. 3-D effects
2. Implicit assumption of steady-state populations in the TCRE table.

This presentation will explore the effect of time-dependent level population kinetics on the line ratio.

The NRL DZAPP code was used to assess the effect of time-dependent and steady-state atomic level kinetics



DZAPP is a coupled 1-D MHD, detailed non-LTE atomic physics with radiation transport, incorporating a transmission line circuit for the driving generator.

The non-LTE population dynamics and radiation transport are time-split from the MHD.

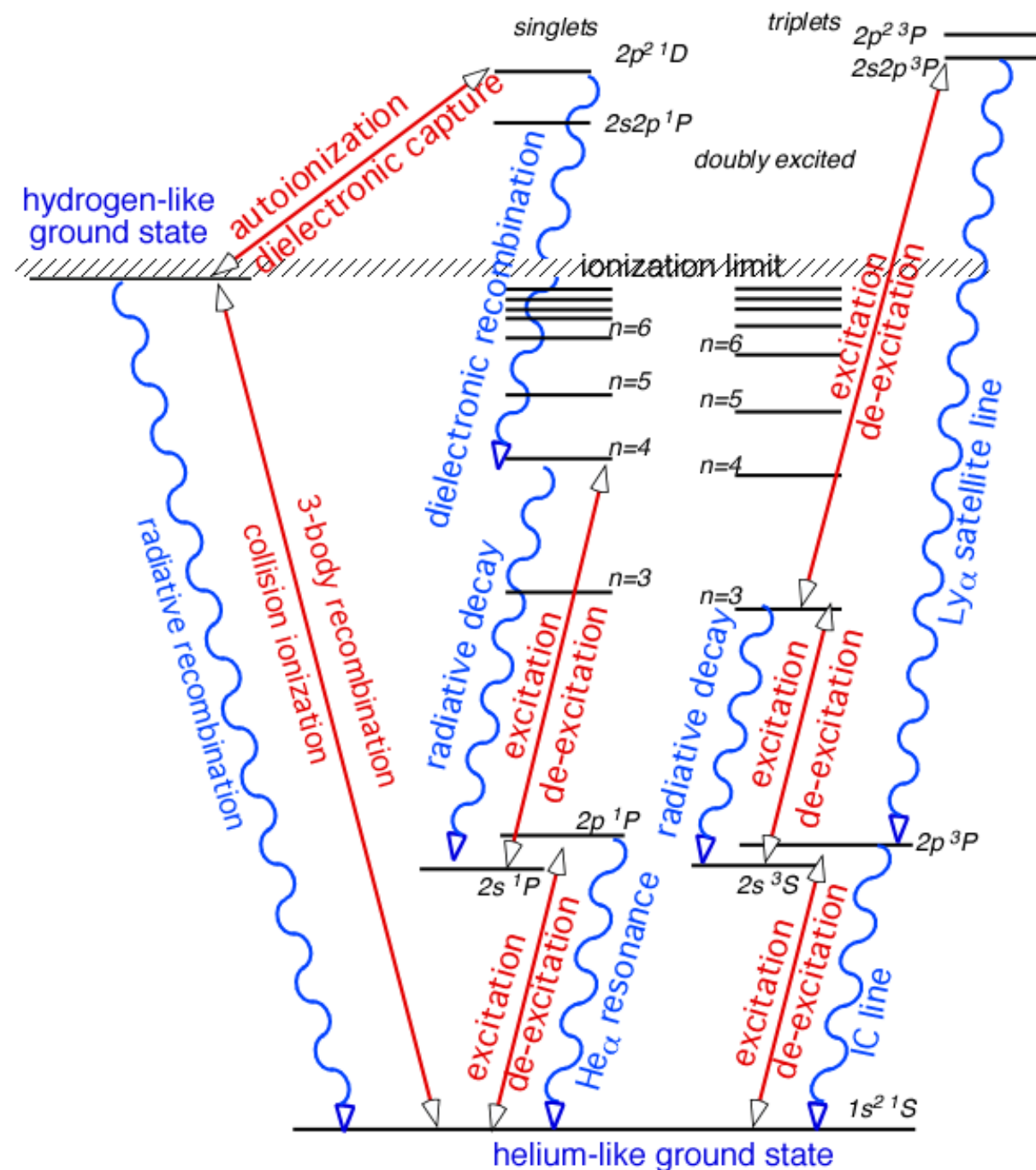
For steady-state populations the time derivative is set to zero

Population of atomic level i in spatial zone k

$$\frac{df_{ik}}{dt} = \sum_j (W_{jik}f_{jk} - W_{ijk}f_{ik})$$

Net rate describing a transitions from atomic level j to atomic level i

All physical processes in our collisional-radiative modeling are driven by non-LTE atomic kinetics



Spontaneous decay/Resonant photoabsorption

$$X_j^{Z*} \Leftrightarrow X_i^Z + h\nu$$

Electron impact ionization/3-body recombination

$$X_i^Z + e \Leftrightarrow X_j^{Z+1} + e' + e''$$

Electron impact excitation/deexcitation

$$X_i^Z + e \Leftrightarrow X_j^{Z*} + e'$$

Photoionization/radiative recombination

$$X_i^Z + h\nu \Leftrightarrow X_j^{Z+1} + e$$

Autoionization/resonant capture

$$X_i^{Z+1} + e \rightarrow X_j^{Z**}$$

Bremsstrahlung/inverse bremsstrahlung

$$X_i^Z + e \rightarrow X_i^Z + e + h\nu$$

Dielectronic recombination

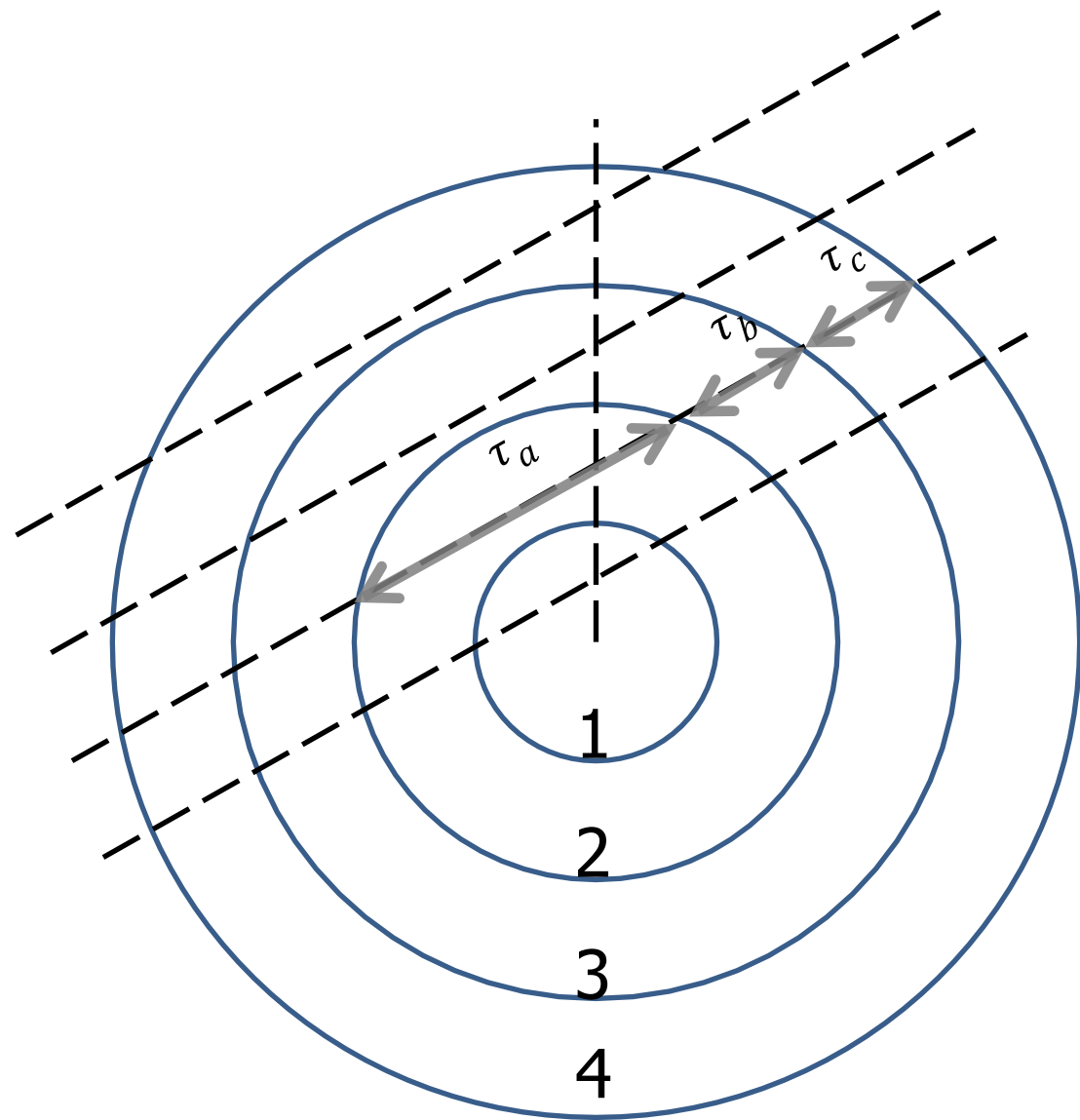
$$X_i^{Z+1} + e \rightarrow X_j^{Z**}$$

$$X_j^{Z**} \rightarrow X_i^Z + h\nu \quad \text{and/or} \quad X_j^{Z**} \rightarrow X_i^{Z+1} + e$$

The populations depend on the rates of these atomic processes and can be determined by:

$$\frac{df_{ik}}{dt} = \sum_j (w_{jik}f_{jk} - w_{ijk}f_{ik}) + \text{photo-pumping by non-local radiation field}$$

Multi-zone radiation transport via the coupling constant formalism provides good overall energetics at a reasonable cost in computer time



Multi-zone radiation transport is necessary because photons are more likely to escape near the edge

$$C_{24} = \frac{1}{\tau_a} \int_0^{\tau_a} [P_e(\tau_b + \tau) - P_e(\tau_b + \tau_c + \tau)] d\tau$$

Reciprocity relation:

$$C_{42} = C_{24} \frac{N_2}{N_4}$$

Ratio of number
of absorbers

The photo-excitation rate for transition $i'-i$ in spatial zone k :

$$W_{ii'k} = 4\pi \sum_{k' \neq k} C_{k'k}^{ii'} j_{k'}^{ii'} \frac{V_{k'}}{N_k f_{ik} E_{ii'} V_k}$$

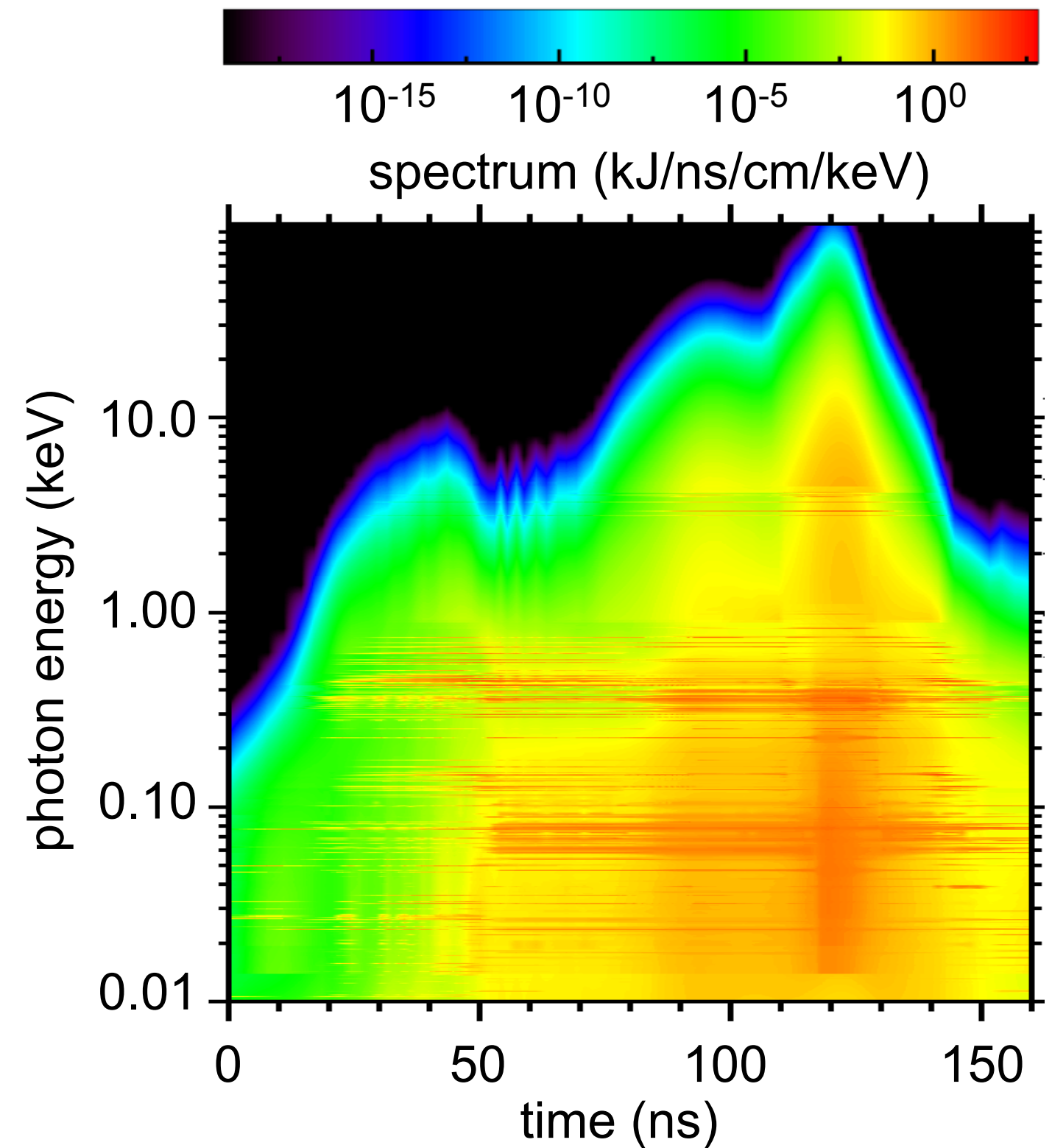
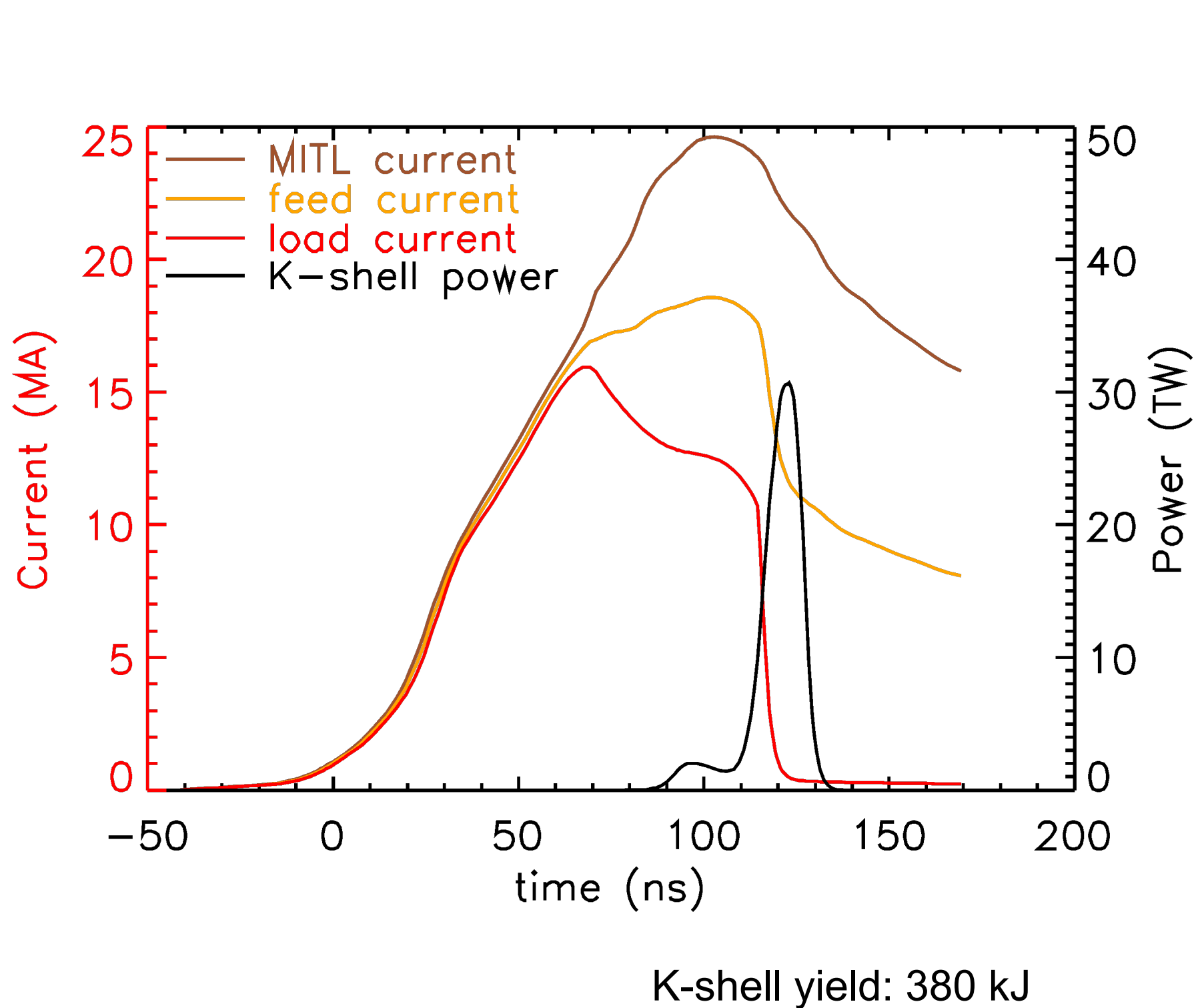
J.P. Apruzese *et al.* JQSRT 23, pp. 479-487 (1980)

J.P. Apruzese. JQSRT 25, pp. 419-425 (1981)

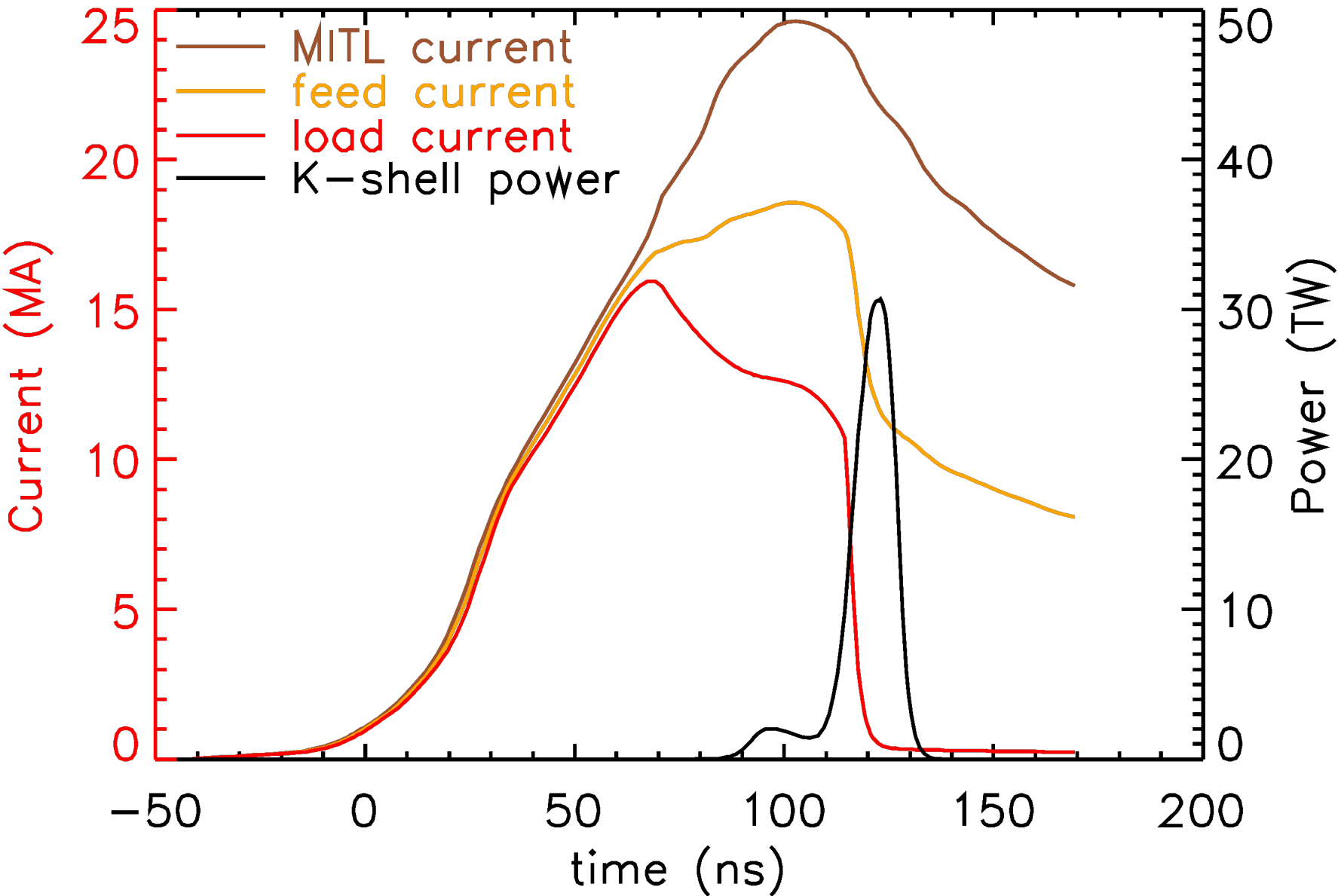
J.P. Apruzese. JQSRT 34, pp. 447-452 (1985)

R.W. Clark *et al.* JQSRT **53**, pp. 307-320 (1995)

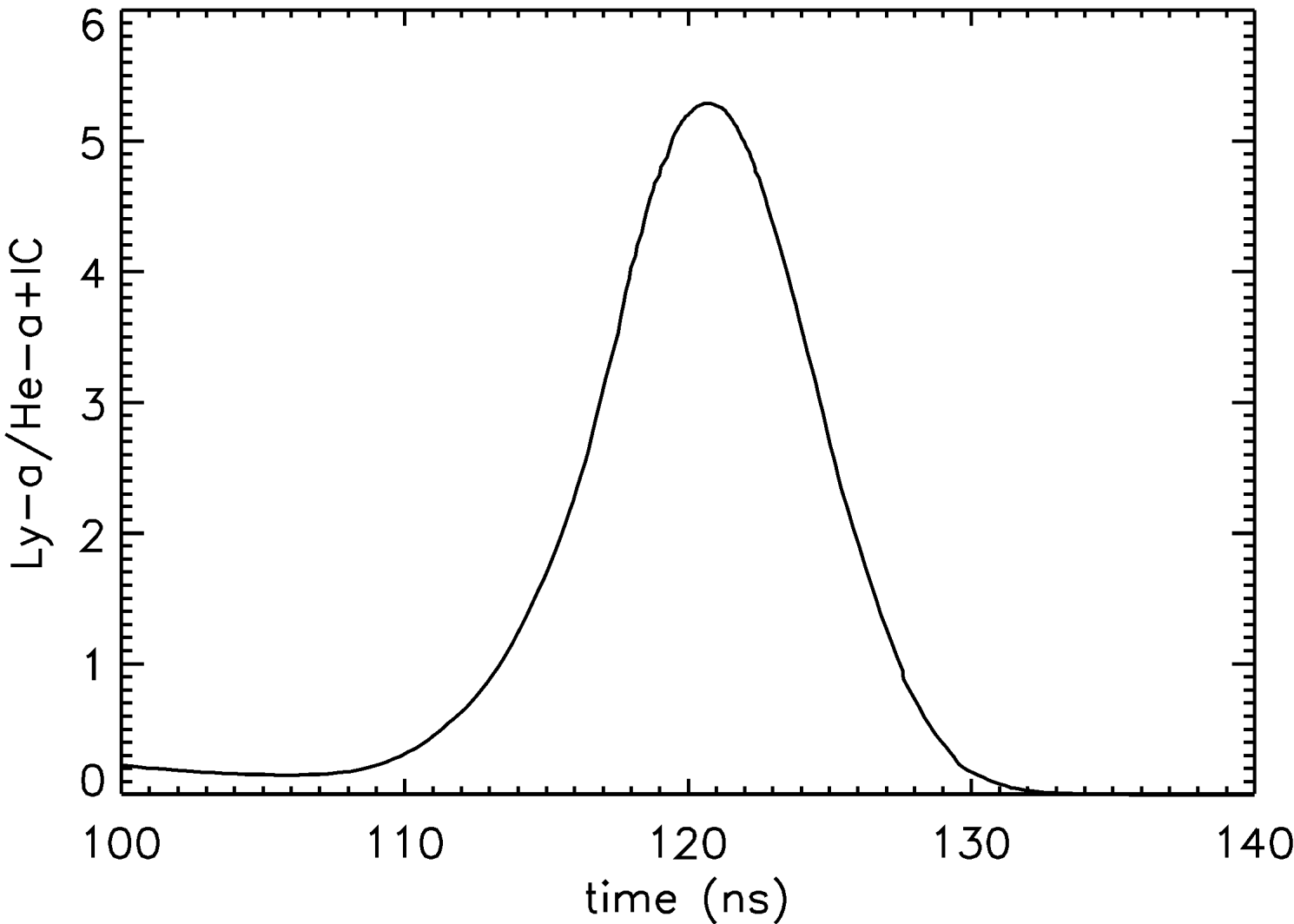
Current waveforms, x-ray powers, and spectra assuming collisional-radiative steady-state level populations



The plasma is too hot as determined from the Ly- α /He- α +IC line ratio
(steady-state level populations)

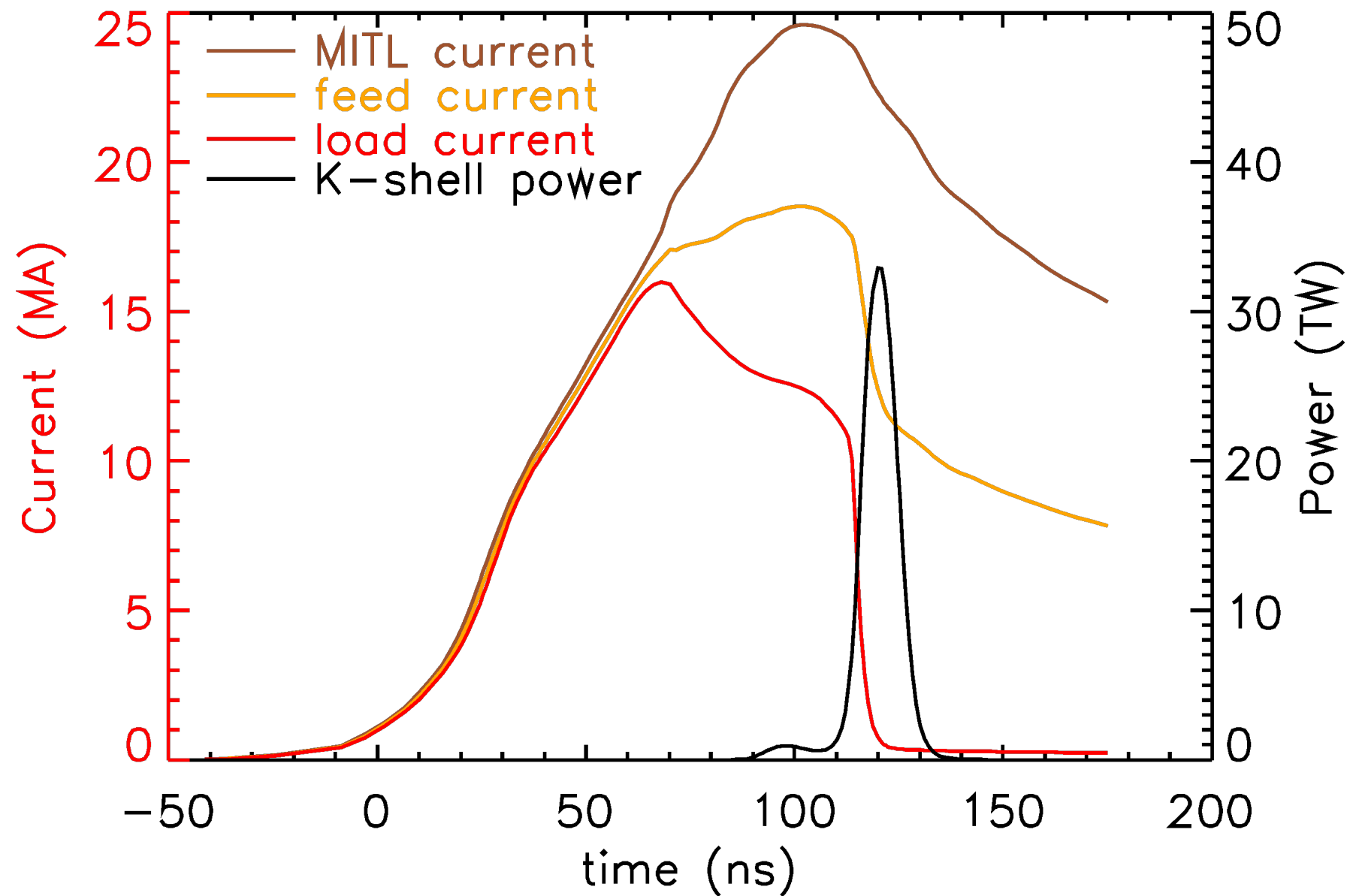


K-shell yield: 380 kJ

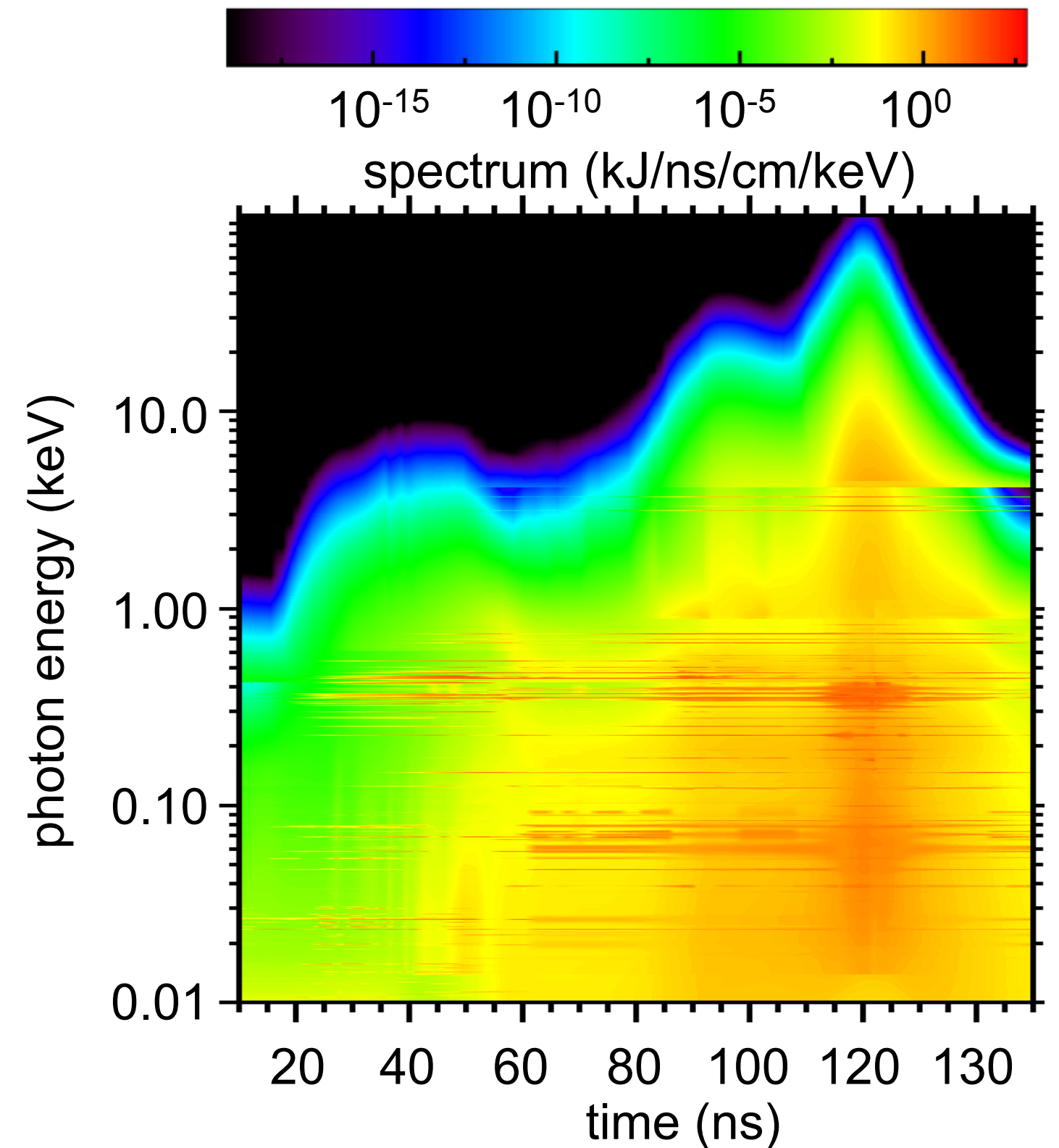


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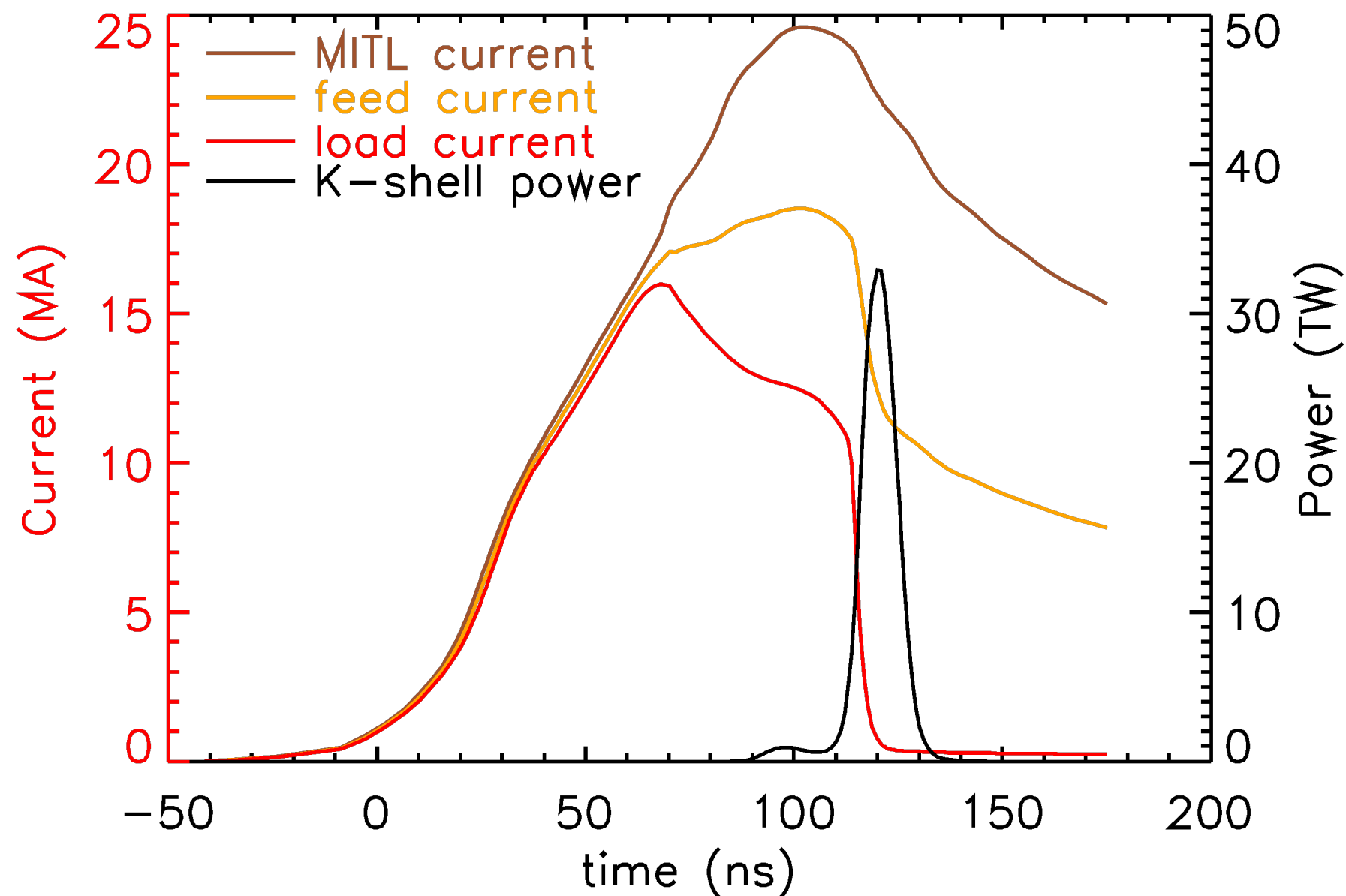
Current waveforms, x-ray powers, and spectra assuming time-dependent ground state level populations



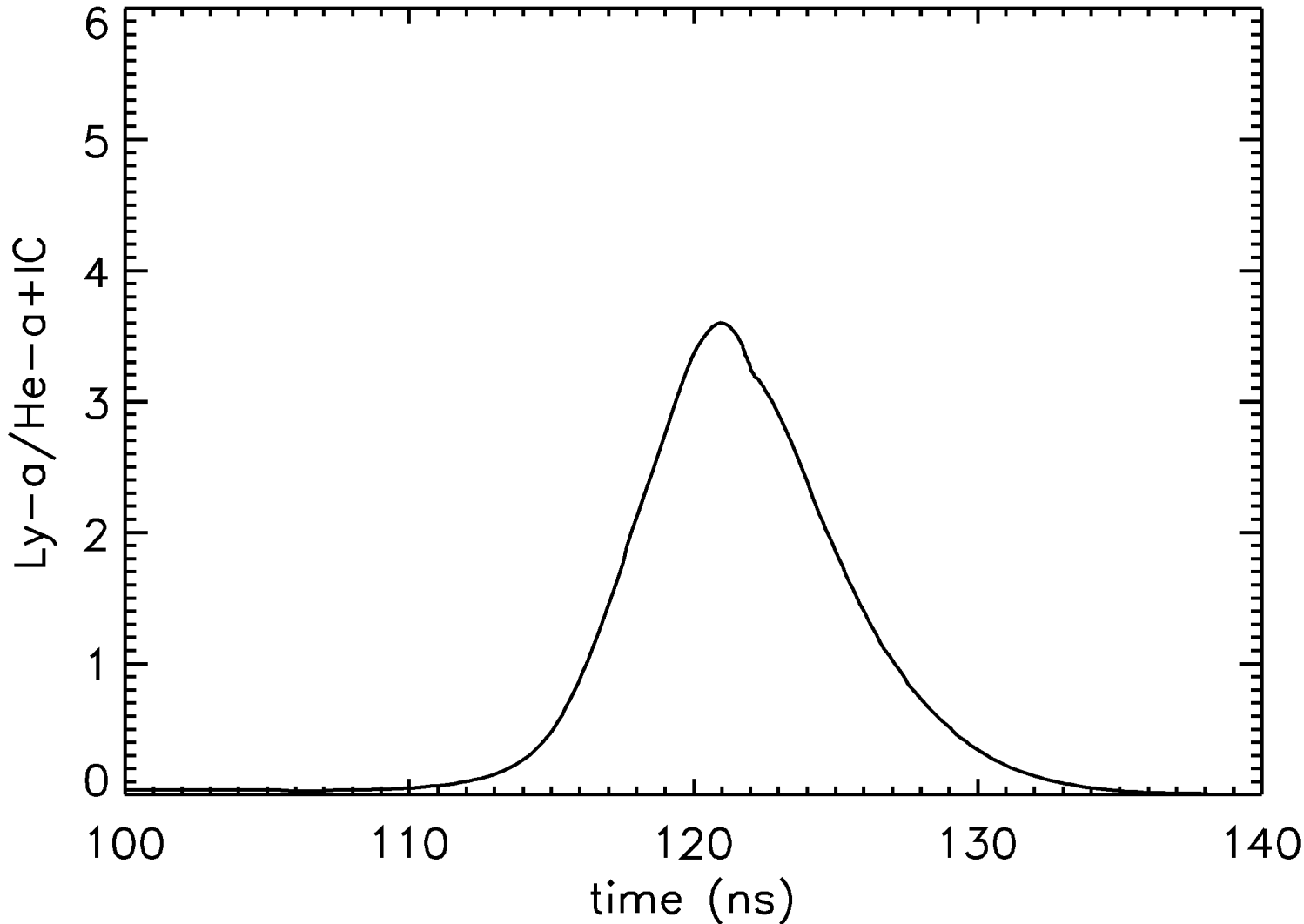
K-shell yield: 358 kJ



The Ly- α /He- α +IC line ratio decreased when using time-dependent ground state level populations



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