

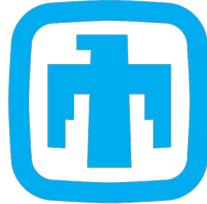
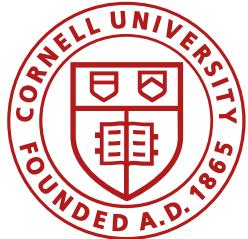
Effects of electron-ion collisions on stopping powers in warm dense matter

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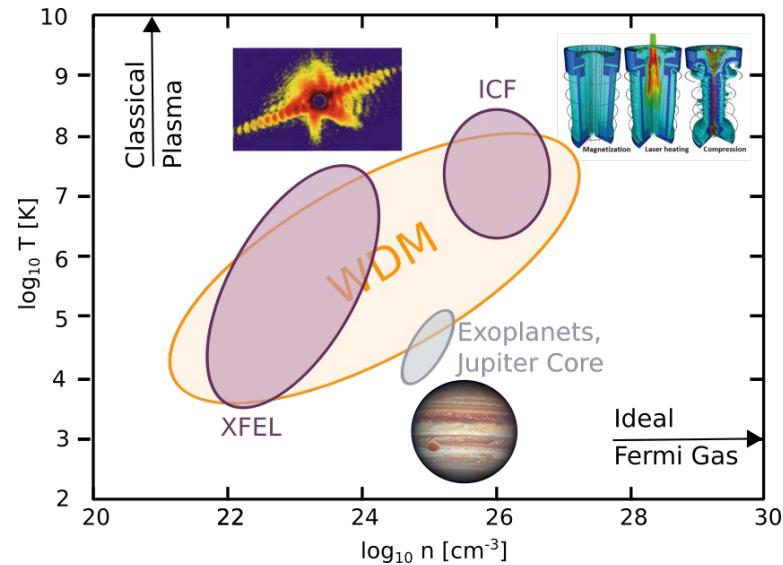
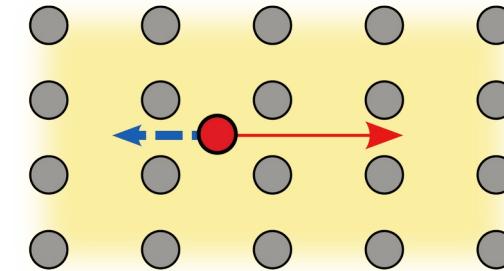
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Stopping powers are important in fusion experiments

- The stopping power is the energy lost as a charged particle moves through a plasma
- Stopping powers describe α -particle **self-heating** in inertial confinement fusion (ICF) experiments, which go through the **warm, dense matter** regime



T. Dornheim, S. Groth, M. Bonitz *Phys. Reports* **744** (2018)

P. F. Knapp et al. *Physics of Plasmas* **22** (2015)

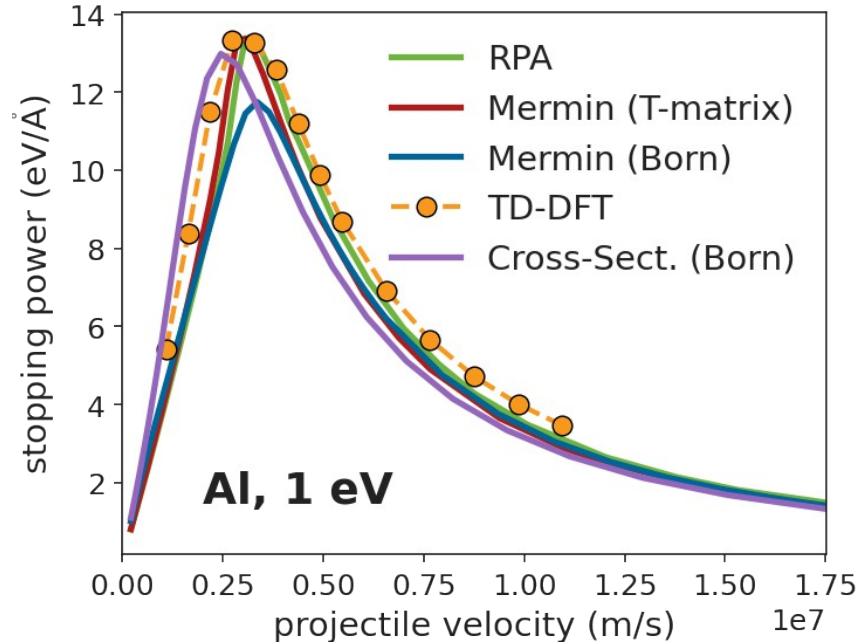
Various microscopic models for stopping exist

1. Expensive **first-principles** methods

- real-time time-dependent density functional theory (TD-DFT)

2. Inexpensive *reduced-physics* methods

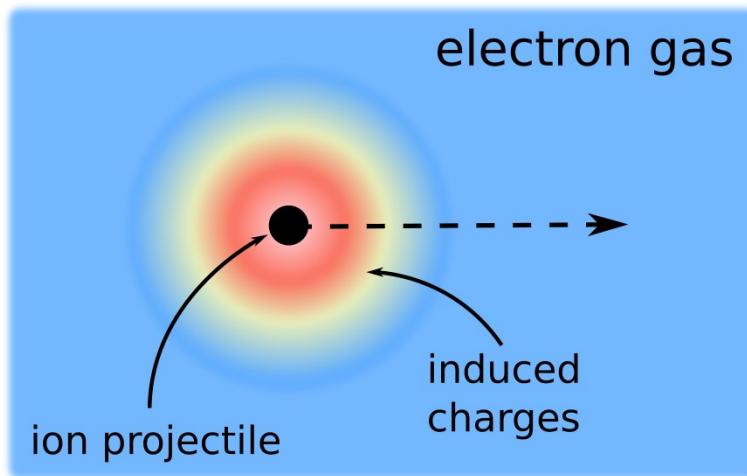
- Dielectric formalism
- Scattering theory [1,2]



[1] C. F. Clauser, N. R. Arista, *Phys. Rev. E* **97** (2018)

[2] N. R. Arista, P. Sigmund, *Phys. Rev. A* **76** (2007)

The dielectric function describes the dynamic electron screening



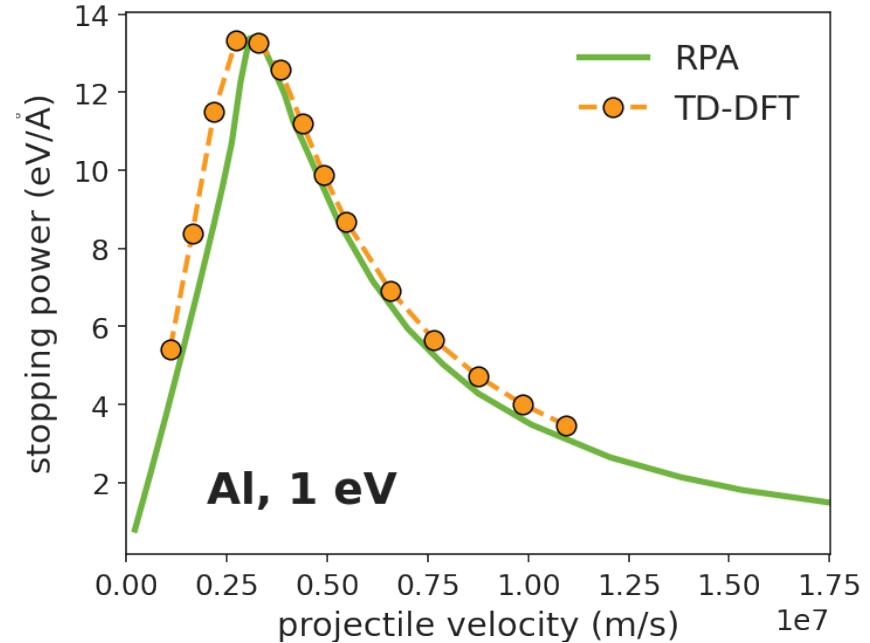
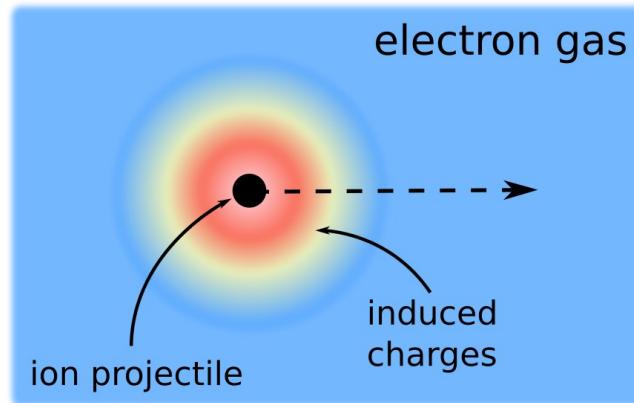
$$\phi^{\text{tot}}(\mathbf{q}, \omega) = \phi^{\text{ind}}(\mathbf{q}, \omega) + \phi^{\text{ext}}(\mathbf{q}, \omega)$$

- The dielectric function generalizes the dielectric constant for dynamic/non-uniform fields
- The **induced charge density** will create an **electric field** that acts against the ion projectile
- This induced field depends on $\epsilon(\mathbf{q}, \omega)$ and the external charge

$$\phi^{\text{tot}}(\mathbf{q}, \omega) = \frac{1}{\epsilon(\mathbf{q}, \omega)} \phi^{\text{ext}}(\mathbf{q}, \omega)$$

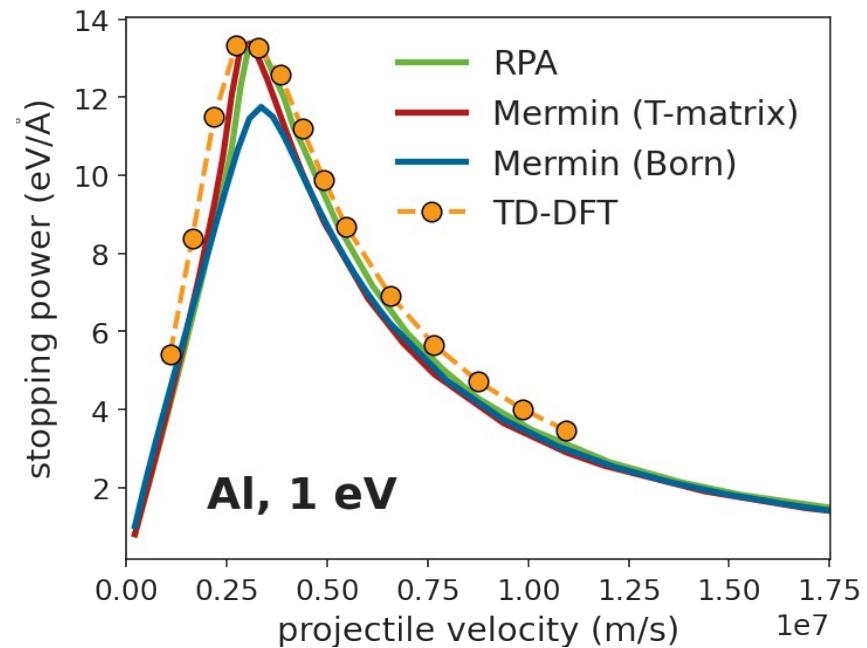
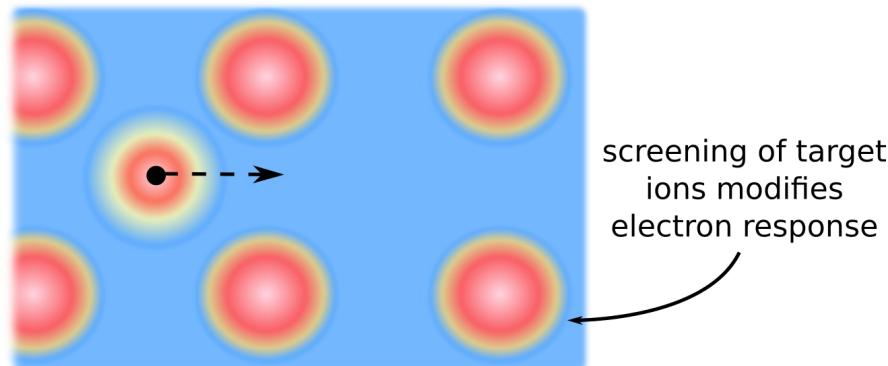
We rely on approximations to calculate the dielectric function

- Our dielectric model is based on a uniform electron gas (UEG) description of the free electrons – **random phase approximation (RPA)**



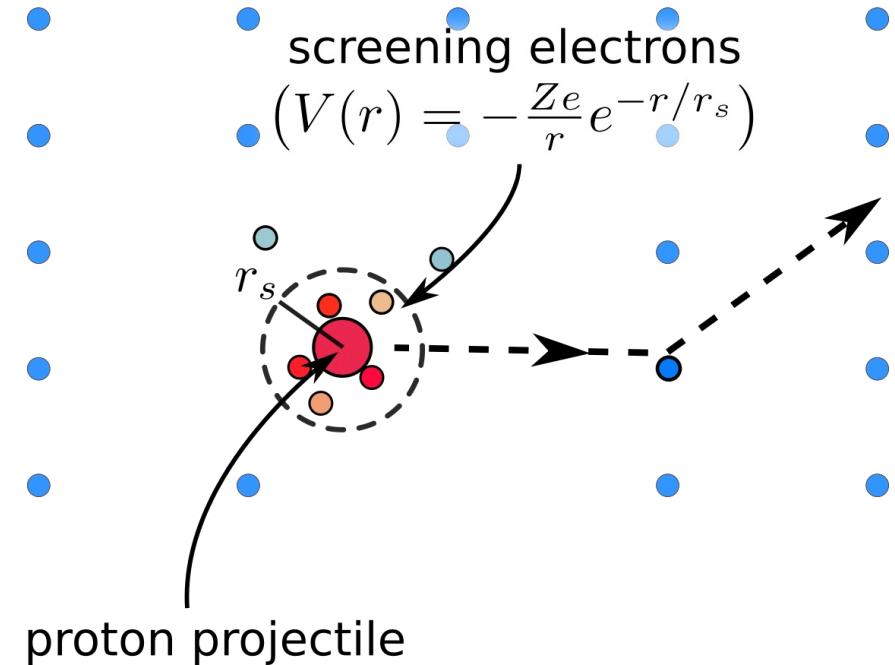
Going beyond the RPA picture with collisions frequencies

- We modify the RPA by including dynamic **electron-ion collision frequencies** – **Mermin dielectric**
- These are calculated with **Born** or **transition-matrix** cross-sections using an **average-atom** model – **depend on the target material**



We can also compute stopping powers directly from scattering cross-sections [1,2]

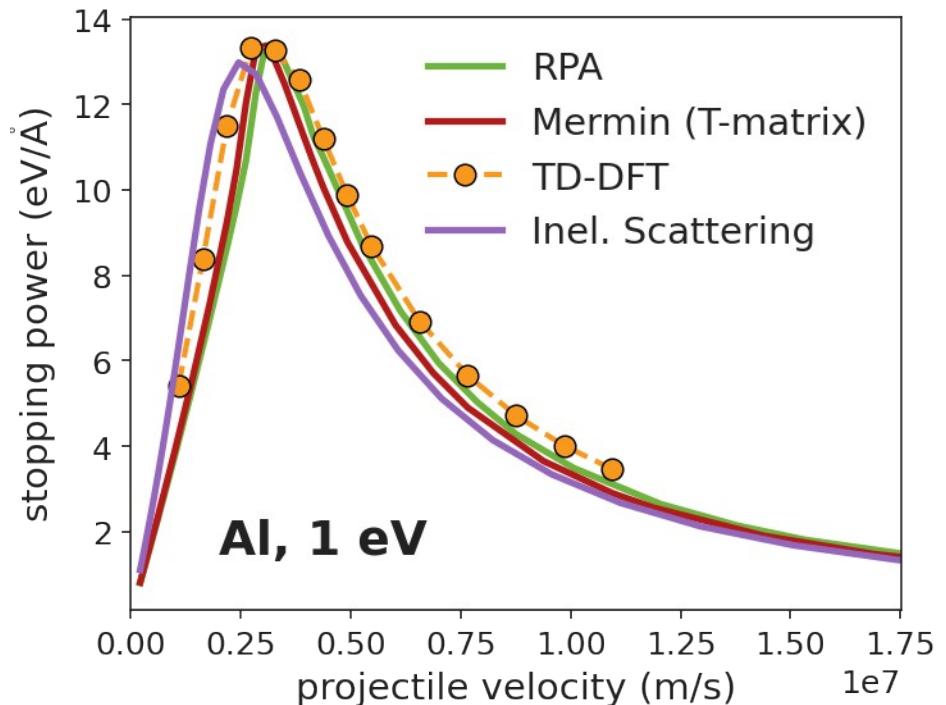
- Projectile traveling through a UEG
- Energy loss as a series of **inelastic scattering** events between projectile and electrons
- Yukawa potential for screening



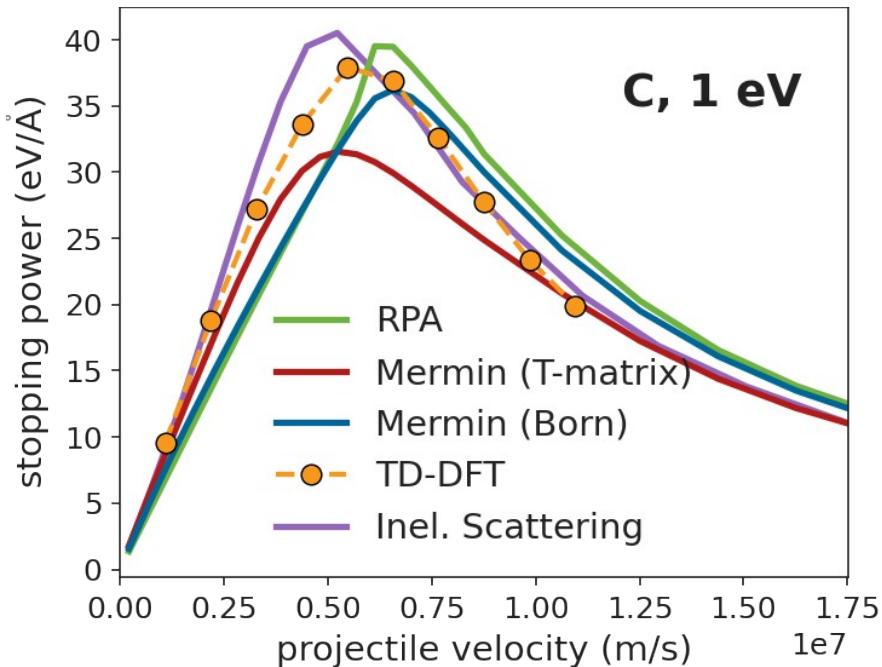
[1] C. F. Clauser, N. R. Arista, *Phys. Rev. E* **97** (2018)
[2] N. R. Arista, P. Sigmund, *Phys. Rev. A* **76** (2007)

Comparing reduced-physics models to TD-DFT for aluminum

- Proton projectile with a warm, dense aluminum target ($k_B T = 1 \text{ eV}$, 2.7 g/cm^3)
- **Free-electron stopping** (ignoring bound electron contribution)
- Low-velocity shift in the inelastic scattering Bragg peak



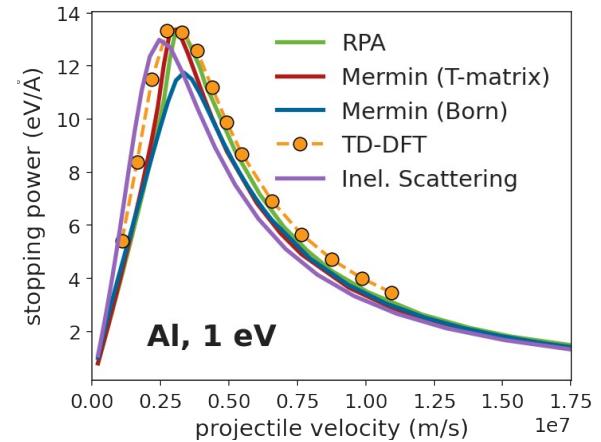
Stopping results in warm, dense carbon



- Proton projectile with a carbon target ($k_B T = 1 \text{ eV}$, 10 g/cm^3)
- (Free-electron stopping)
- Mermin dielectric with T-matrix collisions and inelastic scattering theory stopping models agree best with TD-DFT's **low velocity behavior**

Reduced-physics models are computationally inexpensive

- Overall, **good agreement** with our TD-DFT data for aluminum at 1 eV; for carbon, fair agreement with Mermin (T-matrix) and scattering theory model
- **Dielectric formalism:** Born failed to describe collisions between free-electrons and (average-) **atoms** in our target



Future Work:

- Consider non-ideal density of states calculated from AA in scattering theory stopping calculations
- Within the dielectric formalism, can we approximate non-ideal stopping with a modified e-i collisions?

