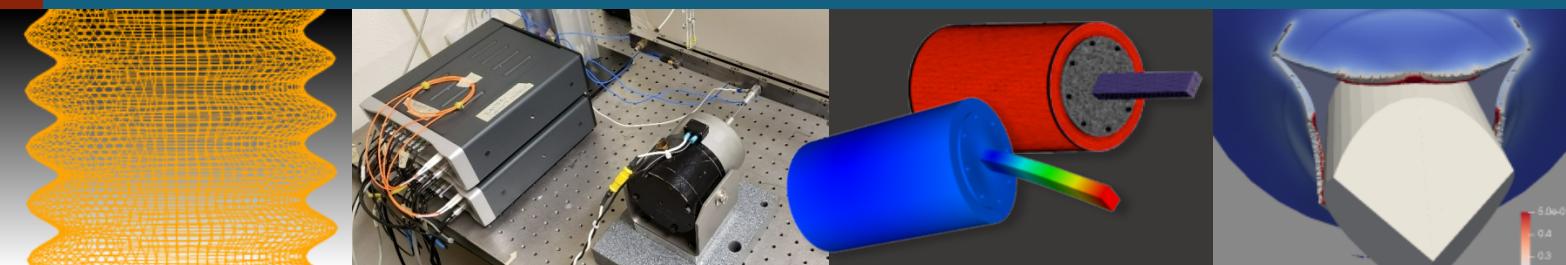




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# Experimental and computational investigation of nonlinear dynamics of a simplified bearing-and-shaft assembly



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# Background



## Low-amplitude vibrations

1. Long-duration random vibration
2. Linear responses produced
3. Classical modal analysis applicable

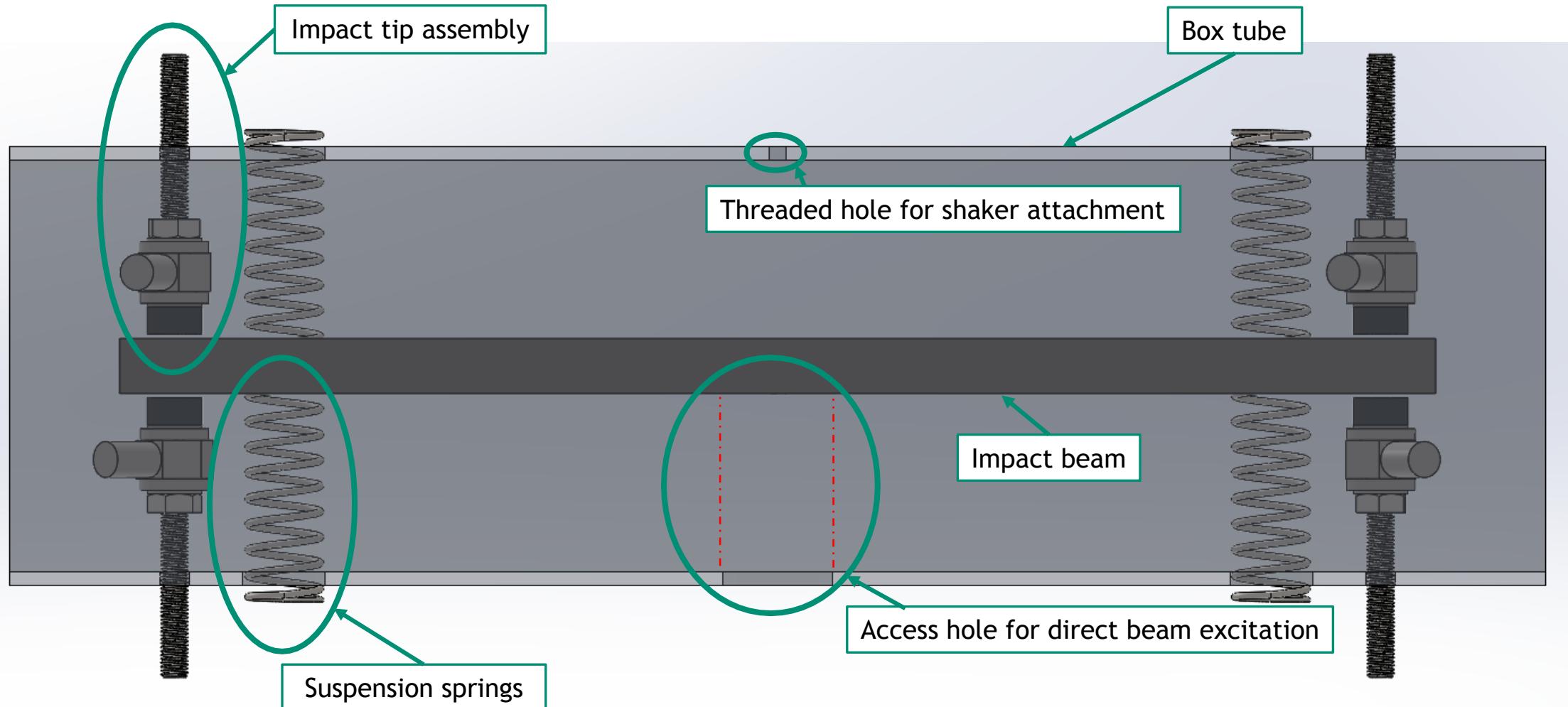
## High-amplitude vibrations

1. Short-duration mechanical shock
2. Nonlinear responses produced
3. Classical modal analysis not applicable\*

- Many electromechanical assemblies have sources of nonlinearity stemming from contact impacts
  - Gears, roller bearings, clearance within small mechanical envelopes, etc..
- Nonlinearities limit or invalidate applicability of linear modal analysis techniques



# Experimental Apparatus – Simplified Bearing-and-Shaft Assembly

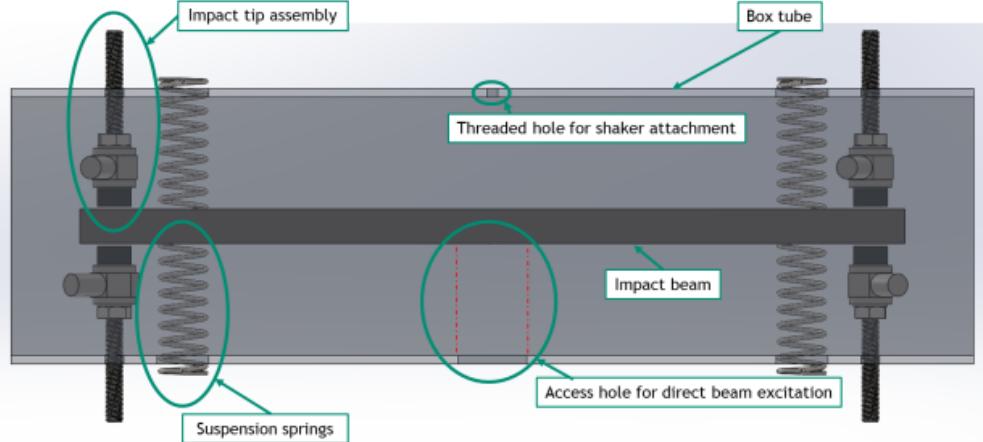


# Hierarchy of Models



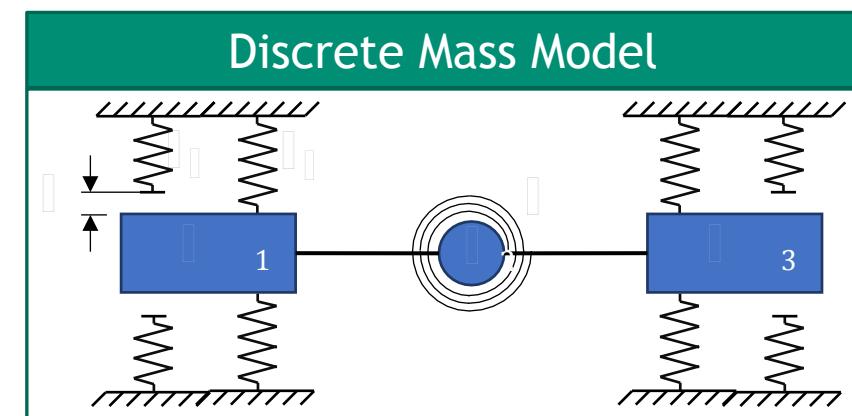
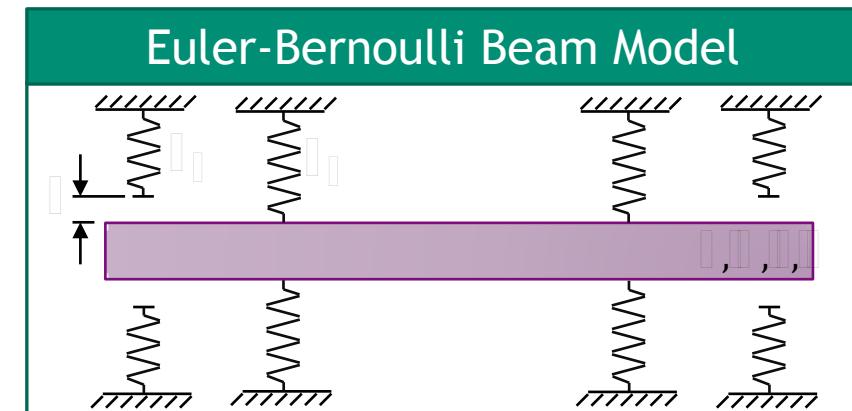
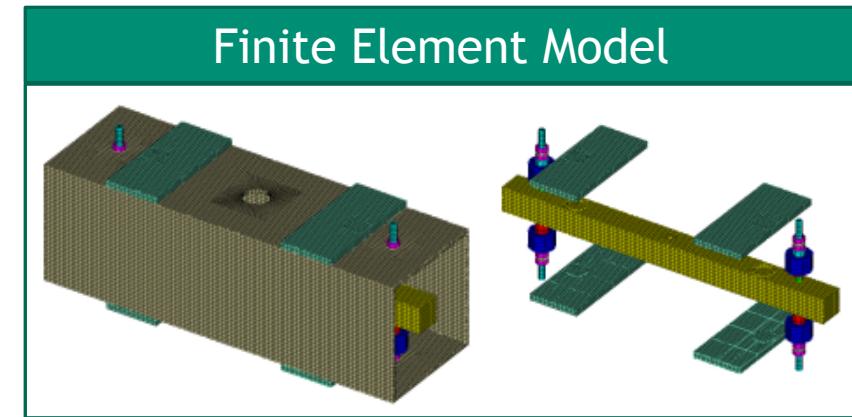
Shared characteristics:

- Symmetric pseudo-rigid body mode
- Antisymmetric pseudo-rigid body mode
- Symmetric bending mode



CAD of experimental apparatus

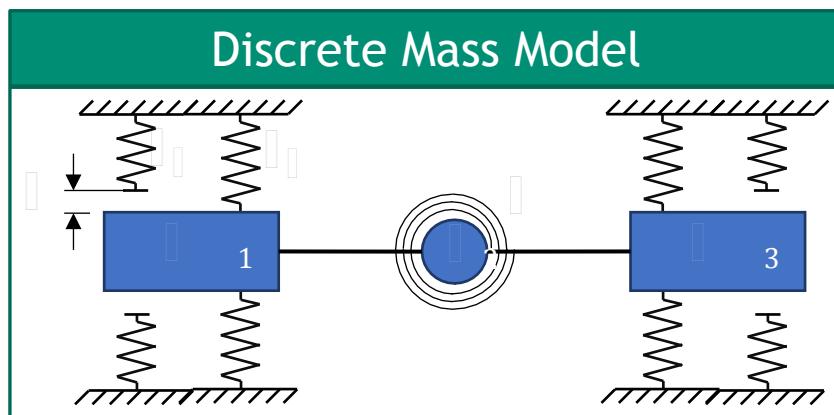
Simplicity



# Objective



- Investigate the nonlinear dynamics in response to high-amplitude broadband excitation
- Identify the spectral content of the ring-down response and correlate to the underlying nonlinear normal modes of the system
  - Key Question – how to treat boundary conditions between impacts? Linearize about open/closed state? Or model as nonlinear displacement dependent boundary condition?
- Observe how the low frequency pseudo-rigid body modes evolve into higher-frequency elastic modes of the structure



Low-energy vibrations:

- No impacts occur
- Linear response

High-energy vibrations:

- Impacts make response nonlinear
- Frequency and energy are related

# Nonlinear Normal Mode Theory

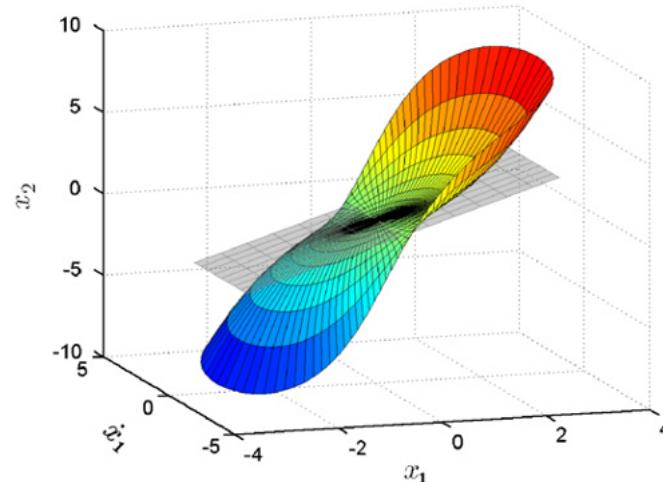


Many definitions exist for either damped or undamped systems [1-4]

For a conservative (undamped) system, a nonlinear normal mode (NNM) is defined as a ***not necessarily synchronous periodic response of the undamped nonlinear system***

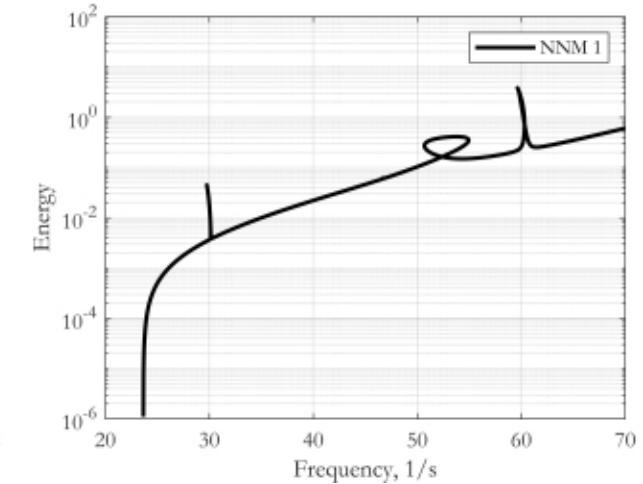
$$\mathbf{M}\ddot{\mathbf{x}}(t) + \mathbf{K}\mathbf{x}(t) + \mathbf{f}_{nl}(\mathbf{x}(t)) = 0$$

For an MDOF system, there exists  $N$  NNM solution branches that are extensions of linear normal modes at low energy [2]



Mode Shapes  
(Manifolds from [2])

## Nonlinear Normal Modes



Frequency-Energy Plots

- [1] Rosenberg, R.M., *The normal modes of nonlinear n-degree-of-freedom systems*. Journal of Applied Mechanics, 1962.
- [2] Kerschen, G., et al., *Nonlinear normal modes. Part I. A useful framework for the structural dynamicist*. Mechanical Systems and Signal Processing, 2009.
- [3] Shaw, S.W., Pierre, C., *Non-linear normal modes and invariant manifolds*. Journal of Sound and Vibration, 1991.
- [4] Haller, G., Ponsioen, S., *Nonlinear normal modes and spectral submanifolds: Existence, uniqueness, and use in model reduction*. Nonlinear Dynamics, 2016.

# Equations of Motion



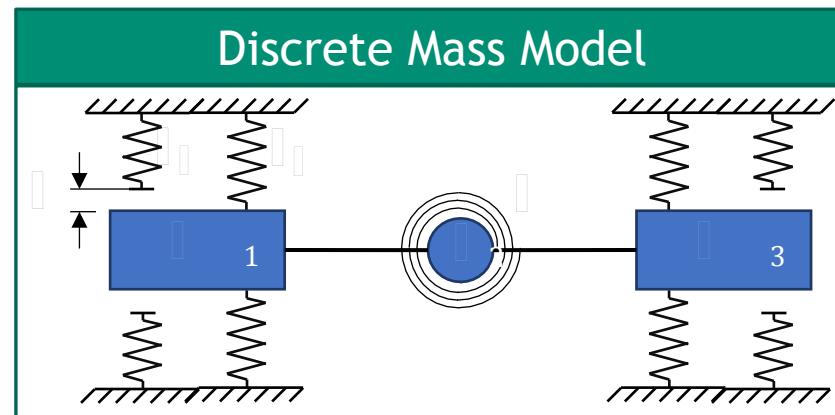
Assuming small rotations in torsional spring (undamped):

$$\begin{bmatrix} m_1 & 0 & 0 \\ 0 & m_2 & 0 \\ 0 & 0 & m_3 \end{bmatrix} \begin{bmatrix} \ddot{x}_1 \\ \ddot{x}_2 \\ \ddot{x}_3 \end{bmatrix} + \begin{bmatrix} 2k_s + \kappa/L^2 & -2\kappa/L^2 & \kappa/L^2 \\ -2\kappa/L^2 & 4\kappa/L^2 & -2\kappa/L^2 \\ \kappa/L^2 & -2\kappa/L^2 & 2k_s + \kappa/L^2 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \\ x_3 \end{bmatrix} + \begin{bmatrix} f(x_1) \\ 0 \\ f(x_3) \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

where

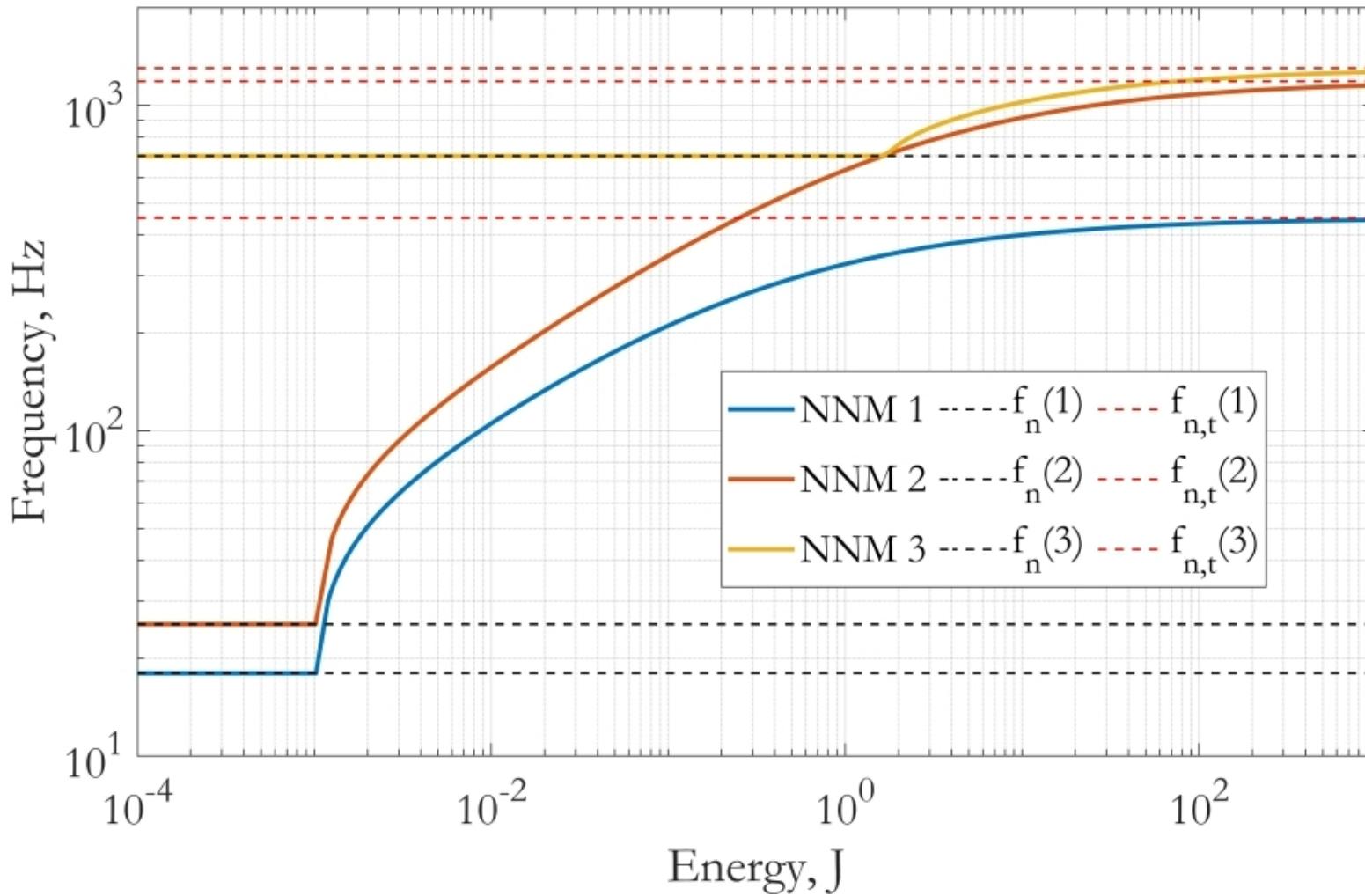
$$f(x) = \begin{cases} k_g(x - g), & x > g \\ 0, & -g \leq x \leq g \\ k_g(x + g), & x < -g \end{cases}.$$

For transient simulations, we include proportional damping matrix  $cI$ .



Description	Symbol	Value
Suspension spring stiffness	$k_s$	8.03e3 N/m
Torsional spring stiffness	$\kappa$	79161 N m/rad
Gap spring stiffness	$k_g$	3.502e7 N/m
Half-length	$L$	0.1614 m
Gap	$g$	2.54e-4 m
Left mass	$m_1$	0.629 kg
Middle mass	$m_2$	1.258 kg
Right mass	$m_3$	0.629 kg
Damping coefficient	$c$	4 N s/m

# Approximate Frequency-Energy Relations from NNM Theory



# Exact Frequency-Energy Relation for Antisymmetric Motions



When  $m_1 = m_3$ , antisymmetric initial conditions yield antisymmetric responses (i.e.,  $x_1(t) = x_3(t)$  and  $x_2(t) = 0$ )

EOM simplify to a single EOM:

$$m_1 \ddot{x}_1 + 2k_s x_1 + f(x_1) = 0.$$

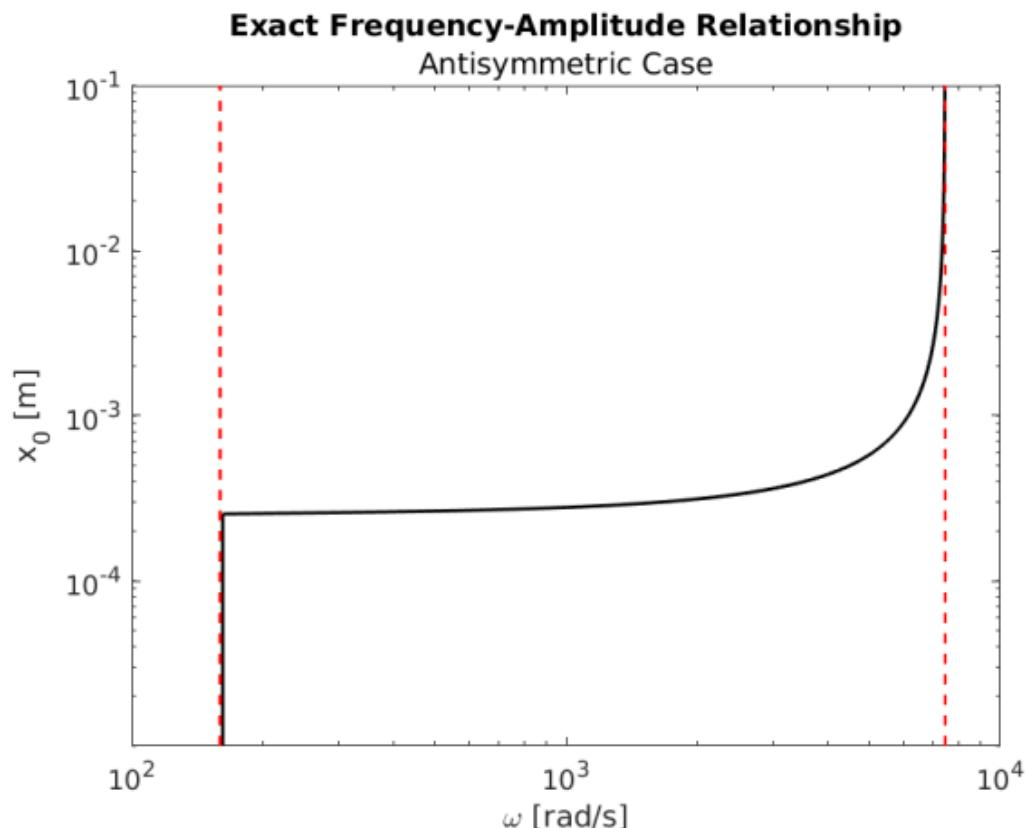
Piecewise linear, can be studied analytically.

**Open-gap natural frequency:**  $\omega_0 = \sqrt{2k_s/m_1}$ .

**Closed-gap nat. freq.:**  $\bar{\omega}_0 = \sqrt{(2k_s + k_g)/m_1}$

**Frequency-amplitude relationship:**

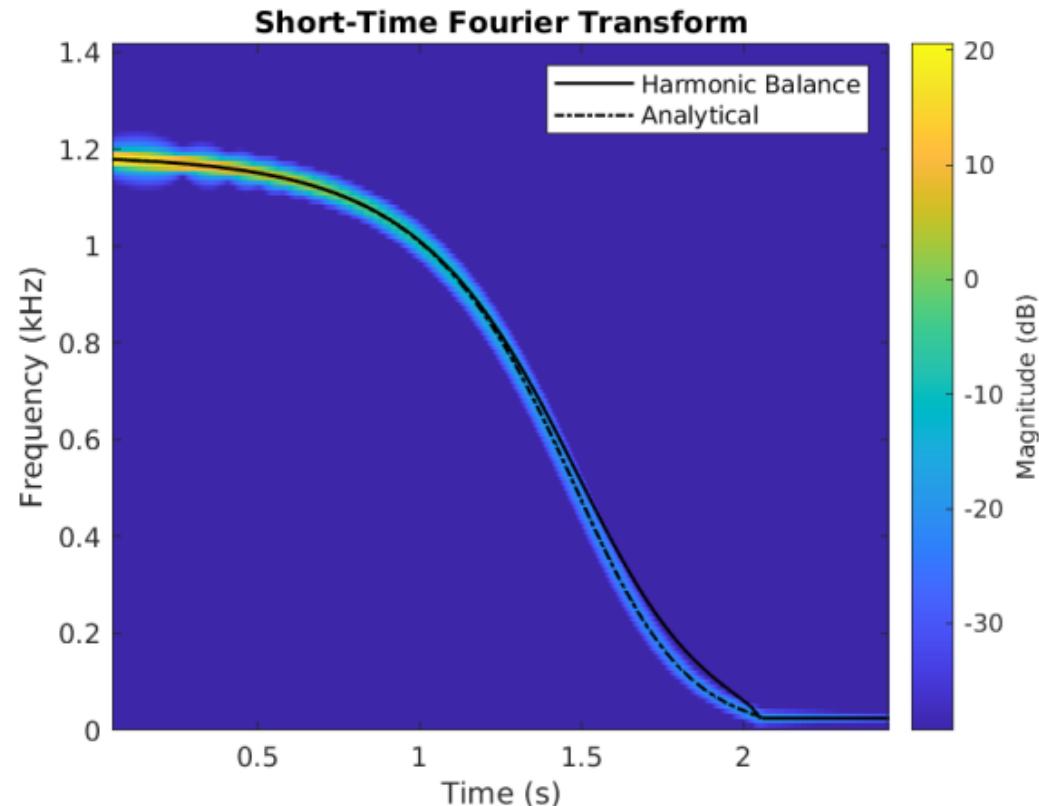
$$\hat{\omega}(x_0) = \begin{cases} 2\pi [T_g(x_0) + T_s(x_0)]^{-1} & x_0 > g \\ \omega_0 & x_0 < g \end{cases}.$$



# Lightly Damped Ringdown: Antisymmetric Initial Conditions



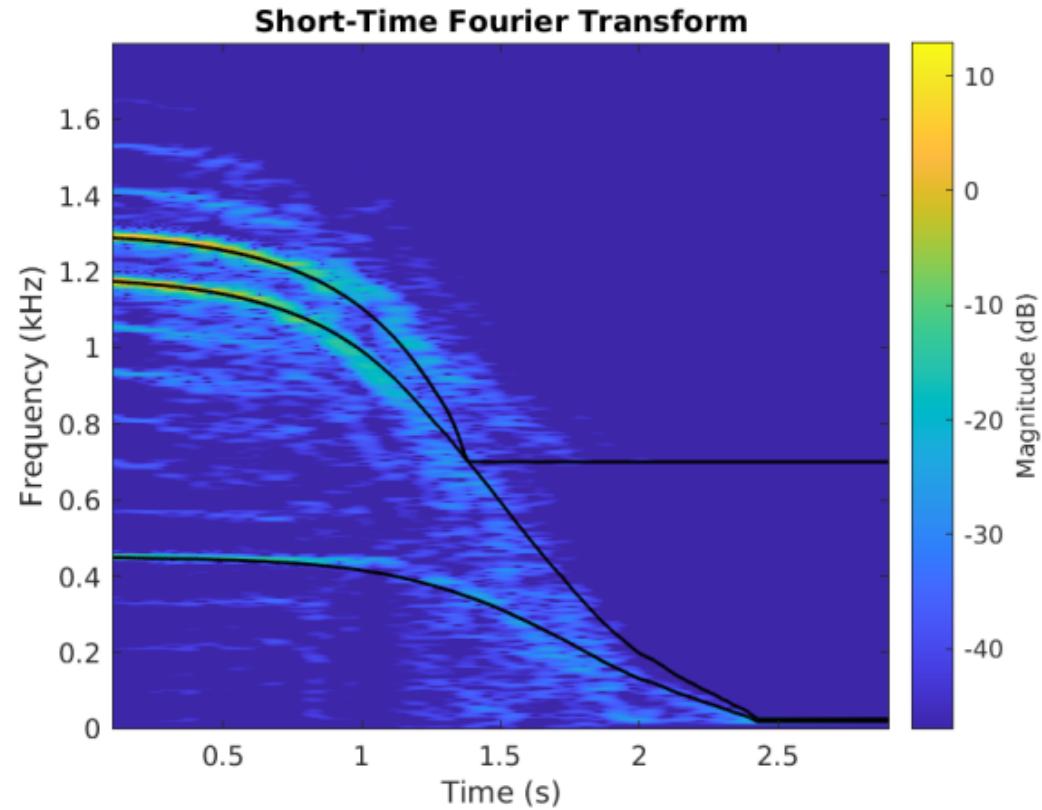
1. Solve (damped) antisymmetric EOM with  $x_1(0) = 100g$  and  $\dot{x}_1(0) = 0$ .
2. Create spectrogram of  $x_1(t)$ .
3. Compute total mechanical energy  $E(t)$ .
4. Use exact and approximate frequency-energy relations to overlay  $\hat{\omega}(E(t))/(2\pi)$  and  $\check{\omega}_1(E(t))/(2\pi)$ .



# Lightly Damped Ringdown: Asymmetric Initial Conditions



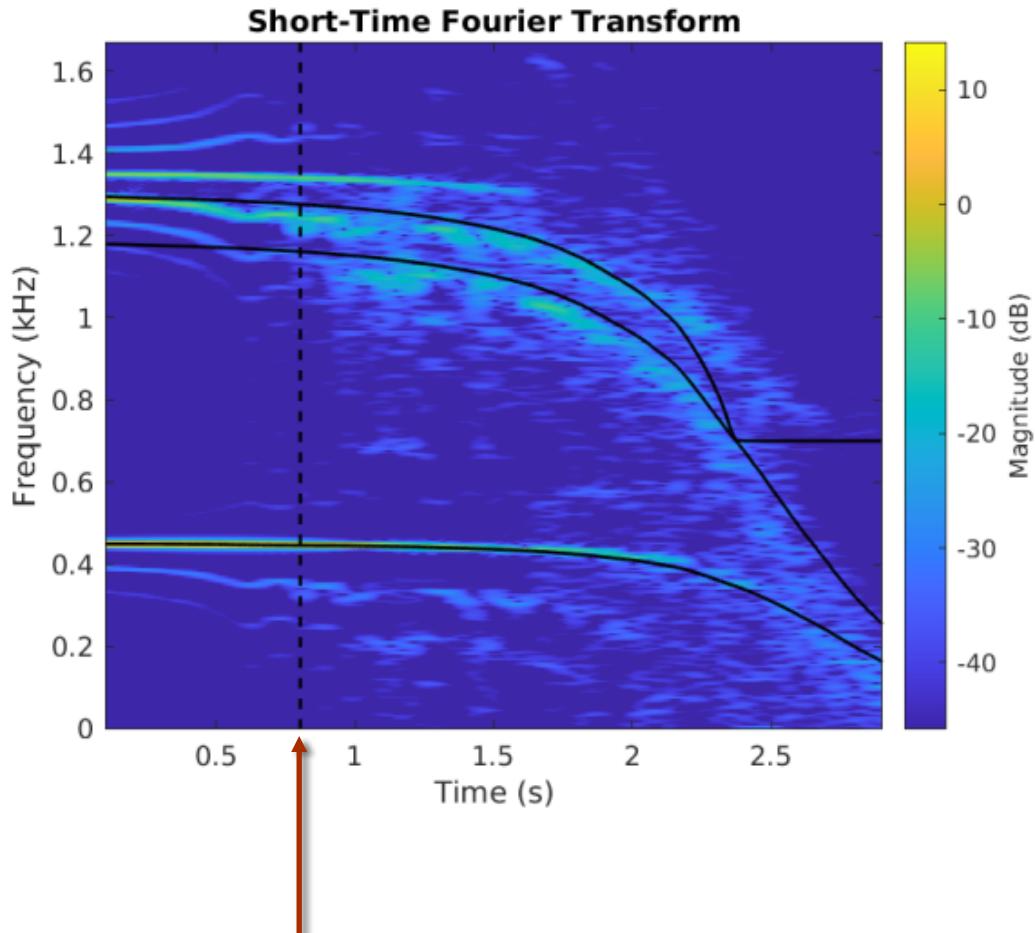
$$x(0) = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix} \text{ m}, \quad \dot{x}(0) = \begin{bmatrix} 0 \\ 0 \\ 200 \end{bmatrix} \text{ m/s}$$



# Lightly Damped Ringdown: Symmetric Initial Conditions

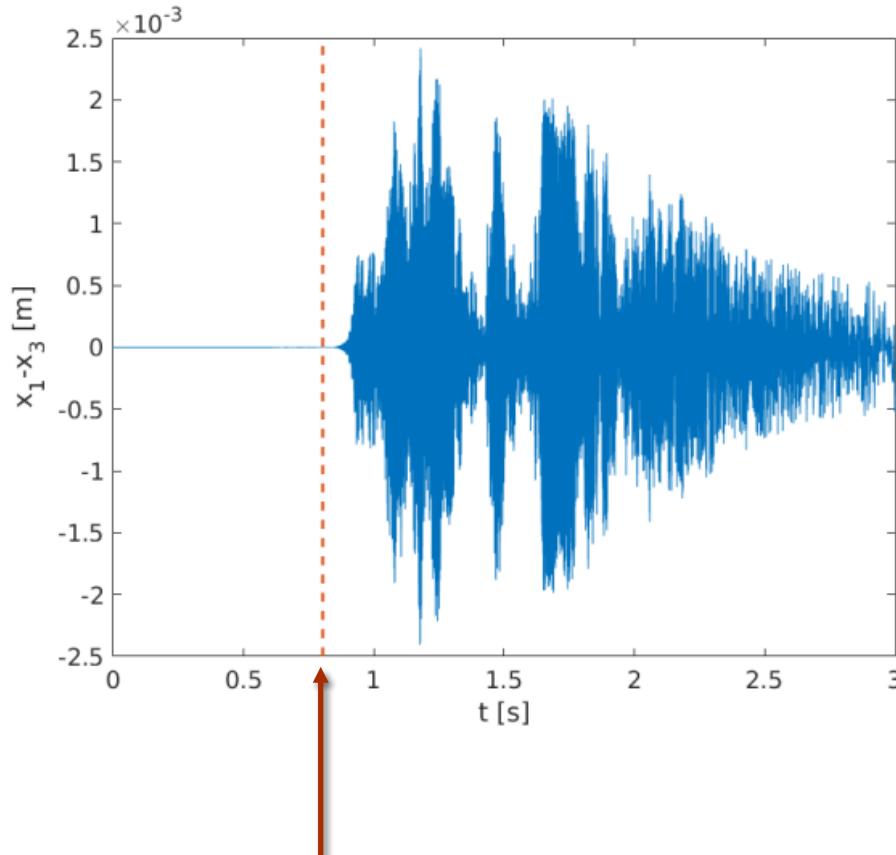


$$x(0) = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix} \text{ m}, \quad \dot{x}(0) = \begin{bmatrix} 0 \\ 200 \\ 0 \end{bmatrix} \text{ m/s}$$



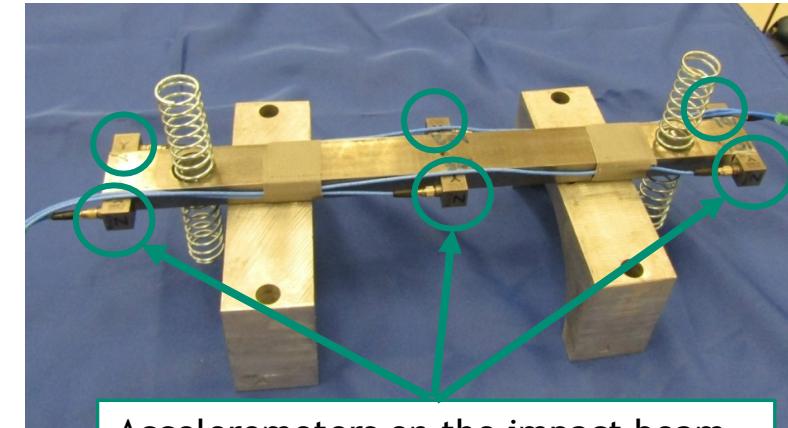
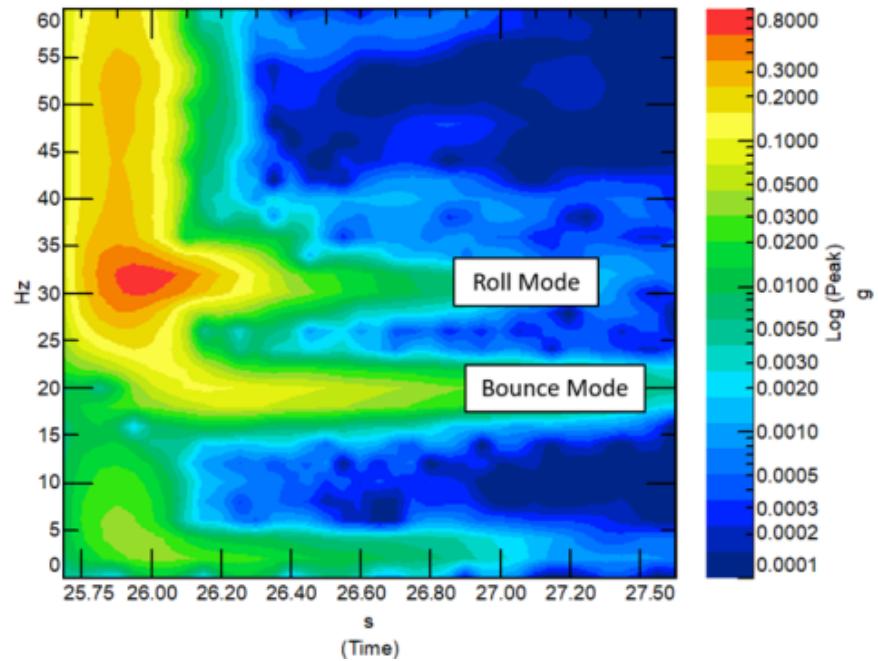
What is special about this time?

# Lightly Damped Ringdown: Symmetry Breaking

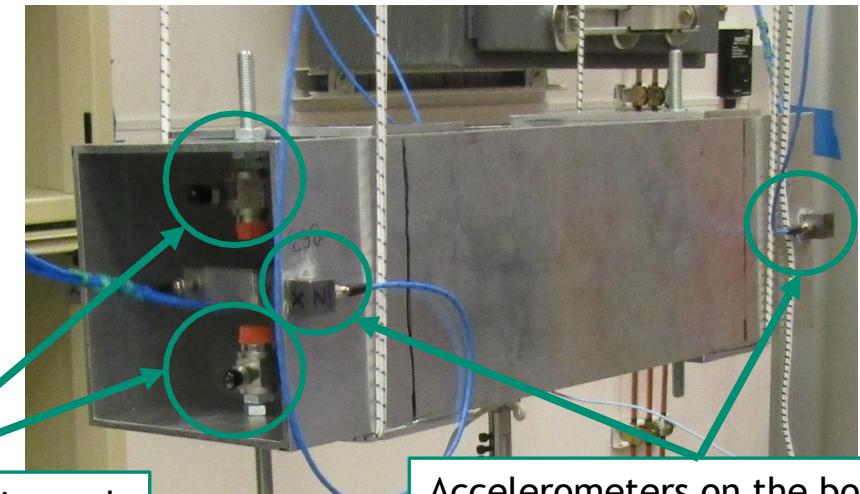


Left-right symmetry is spontaneously broken!

# Evidence of Ringdown After Impact



Accelerometers on the impact beam



Impact load cells (mirrored on the other end)

Accelerometers on the box tube (mirrored on the opposite side)

# Conclusions



- Developed a simple discrete mass system to study the nonlinear dynamics of an idealized bearing-and-shaft assembly
- Nonlinear normal modes reveal the evolution of the resonant behavior starting from pseudo-rigid body motion at low energy to elastic mode deformation at high energy
  - Help inform operating range for modeling assumptions (linear vs. nonlinear boundary conditions)
- Transient free-response simulations show how the NNM frequencies contribute to the response to high-amplitude broadband excitation
  - Observed resonant frequency can depend on energy of excitation
- Future work to further investigate the dynamics of the experimental apparatus to identify spectral content of transient ring-down

# Acknowledgements



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