

## Interplay Between Kondo and Magnetic Interactions in $\text{Pr}_{0.75}\text{Gd}_{0.25}\text{ScGeH}$

Tyler Del Rose<sup>a,b,\*</sup>, Renu Choudhary<sup>b</sup>, Yaroslav Mudryk<sup>b</sup>, Daniel Haskel<sup>c</sup>, Arjun K. Pathak<sup>d</sup>, Gourab Bhaskar<sup>e</sup>, Julia V. Zaikina<sup>e</sup>, Duane D. Johnson<sup>a,b</sup>, and Vitalij K. Pecharsky<sup>a,b</sup>

a Department of Materials Science and Engineering, Iowa State University, Ames, Iowa 50011, USA

b Ames National Laboratory, U.S. Department of Energy, Iowa State University, Ames, Iowa 50011, USA

c Advanced Photon Source, Argonne National Laboratory, Argonne, IL 60439, USA

d Department of Physics, SUNY Buffalo State University, Buffalo, New York 14222, USA

e Department of Chemistry, Iowa State University, Ames, Iowa 50011, USA

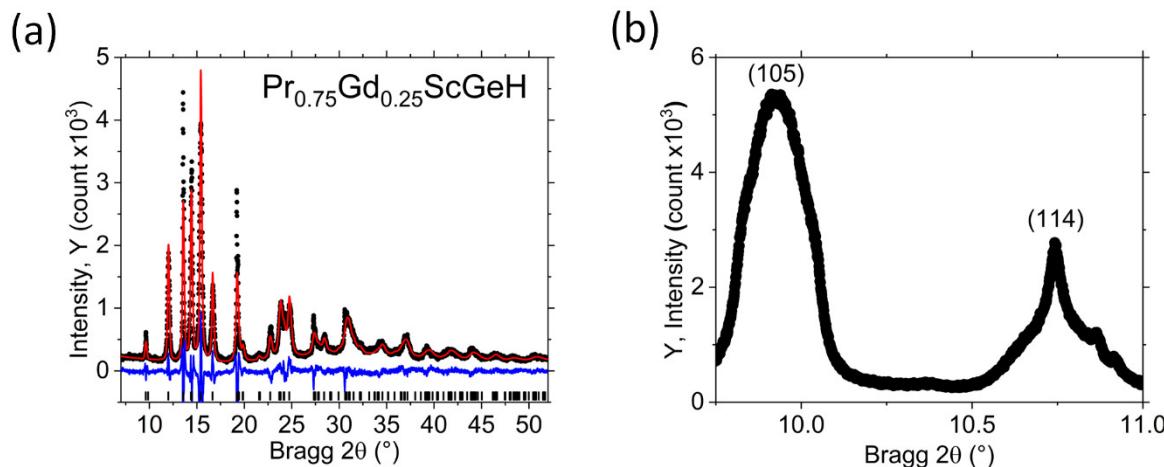


Figure S1. (a) Rietveld refinement and (b) the shapes of (105) and (114) Bragg peaks of  $\text{Pr}_{0.75}\text{Gd}_{0.25}\text{ScGeH}$  hydrogenated at 150 bar. The data presented in (a) are from a laboratory powder diffractometer using Mo-K $\alpha$  radiation, while the data in (b) are from APS, ANL using  $\lambda = 0.457897$   $\text{\AA}$ . The noticeable mismatch of intensities in (a) is due to deteriorated crystallinity and non-analytical peak shapes after hydrogenation, exemplified in (b).

Figure S2 depicts the heat capacity of  $\text{Pr}_{0.75}\text{Gd}_{0.25}\text{ScGeH}$  and  $\text{LaScGeH}$ . Contributions from magnetic and nuclear effects (exemplified by the low-temperature differences between  $\text{LaScGeH}$  and  $\text{Pr}_{0.75}\text{Gd}_{0.25}\text{ScGeH}$ ) makes low-temperature fitting to determine the electronic specific heat,  $\gamma$ , impractical. Assuming that isostructural  $\text{LaScGeH}$  has the same Debye

temperature,  $\theta_D$ , as  $\text{Pr}_{0.75}\text{Gd}_{0.25}\text{ScGeH}$ , it is still possible to estimate  $\gamma$  for the latter as the constant difference between their  $C_p/T$  in the region ( $\sim 30 - 80$  K) where both the nuclear and magnetic contributions become negligible, see Eq. S1 and Figure S2.

$$\gamma_{\text{Pr}_{0.75}\text{Gd}_{0.25}\text{ScGeH}} = \gamma_{\text{LaScGeH}} + \left( \left( \frac{C_p}{T} \right)_{\text{Pr}_{0.75}\text{Gd}_{0.25}\text{ScGeH}} - \left( \frac{C_p}{T} \right)_{\text{LaScGeH}} \right) \quad (\text{Eq. S1})$$

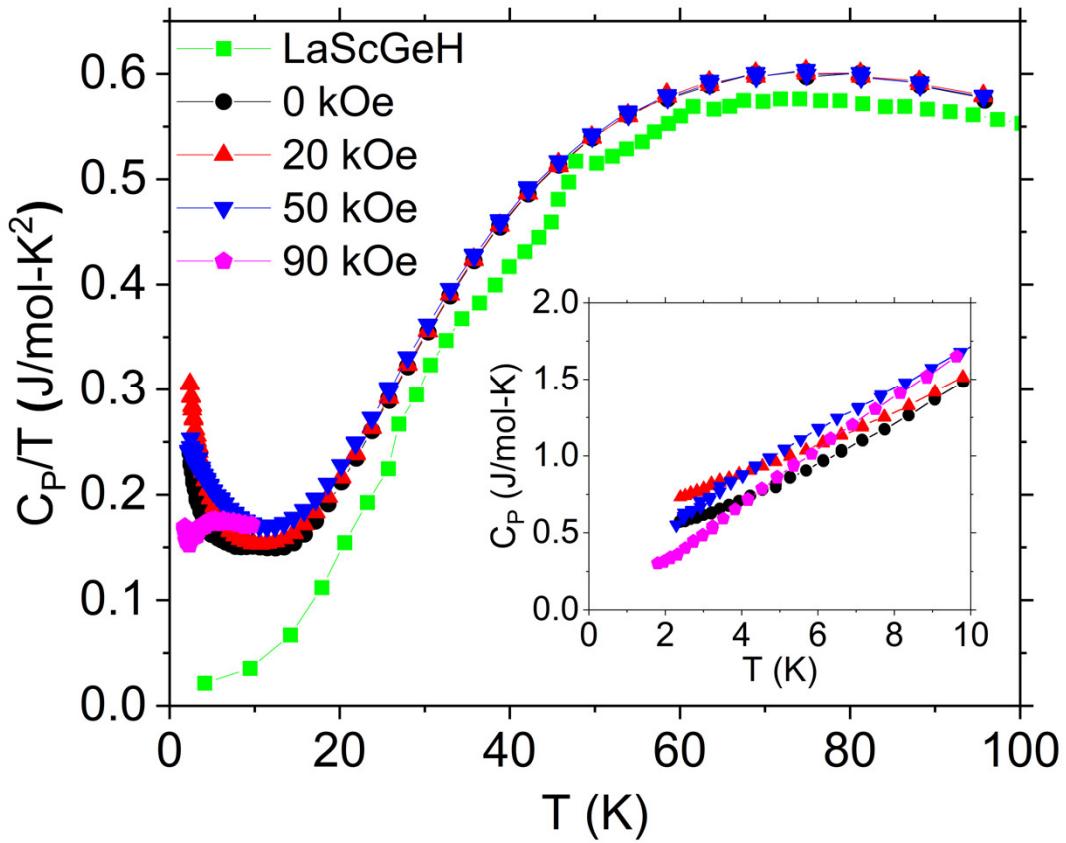


Figure S2. Heat capacity ( $C_p$ ) of  $\text{Pr}_{0.75}\text{Gd}_{0.25}\text{ScGeH}$  measured in zero and applied magnetic fields of 20, 50, and 90 kOe shown as  $C_p/T$  vs. temperature ( $T$ ). LaScGeH data are taken from Ref [1]. Inset shows low-temperature details as  $C_p$  vs  $T$ .

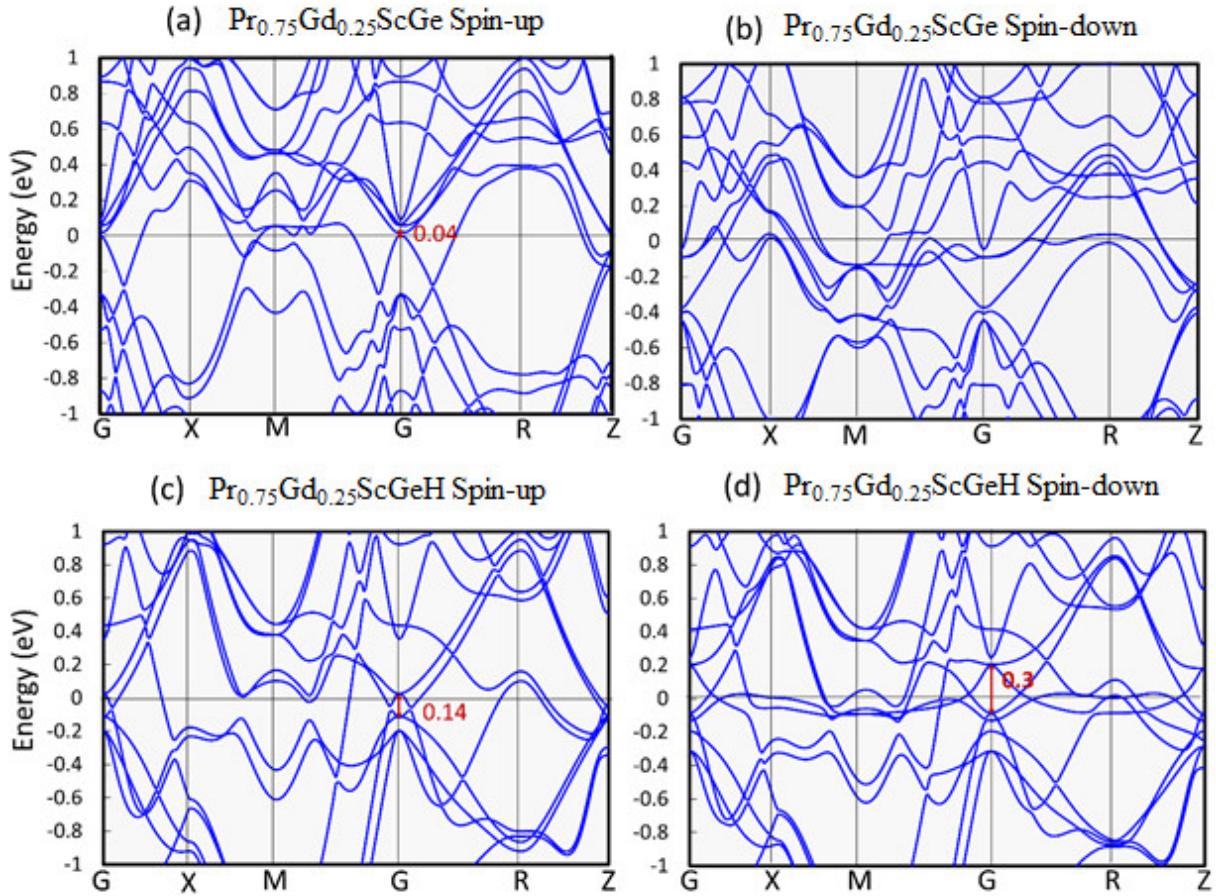


Figure S3: Spin-polarized DFT band structure of (a, b)  $\text{Pr}_{0.75}\text{Gd}_{0.25}\text{ScGe}$  and (c, d) its hydride, showing both spin-up (a, c) and spin-down (b, d) electron manifolds.

The band structures, plotted in Figure S3, indicate a sharp peak of the  $4f$  states close to the  $E_F$ . Additionally, the Gamma point has a very narrow gap at  $E_F$  in the majority-spin channel (Figure S3a), whereas the minority-spin channel is metallic in  $\text{Pr}_{0.75}\text{Gd}_{0.25}\text{ScGe}$  (Figure S3b). As the density of states for a system is inversely proportional to the derivative of energy dispersion  $E_k$ , the sharp minority-spin  $4f$ -peak in PDOS (see Figure 4b) of hydrogenated  $\text{Pr}_{0.75}\text{Gd}_{0.25}\text{ScGe}$  around  $E_F$  is flat in the band structure (Figure S3b). Flat bands mean larger effective electron mass or possible heavy-fermion behavior however, they are not exactly at the  $E_F$ . Therefore, the flat  $4f$  band leads to an enhanced effective electron mass without necessarily reflecting heavy-fermion-

like state. To confirm this, electronic specific heat ( $\gamma$ ) has been calculated for both  $\text{Pr}_{0.75}\text{Gd}_{0.25}\text{ScGe}$  and its hydride, which are 9.35 and 13.22  $\text{mJ/mol. K}^2$ , respectively. The theoretical value of  $\gamma$  for the hydride is smaller than experimental values (28.7  $\text{mJ/mol. K}^2$ ) by a factor of nearly two, whereas both theory and experiment agree for the  $\text{Pr}_{0.75}\text{Gd}_{0.25}\text{ScGe}$  system (Table 1). Electron mass enhancement in rare-earth metals and the effect on specific heat were discussed by Fulde et al. [2]. Experimental electronic specific heat of Pr metal determined from heat capacity data between 1 and 6 K is larger than the DFT value by a factor of  $\sim 4$ , and the difference is reduced by applying a magnetic field. The reason for the mass enhancement of conduction electrons in Pr-metal (and other lanthanides) at low temperature and low magnetic field is the interaction of conduction electrons with localized  $4f$  moments, which is neglected in calculating  $\gamma$  by DFT methods. Hence, theoretically calculated values are commonly different from experimental ones. Here, the predicted interactions of conduction and  $4f$ -electrons of Pr1 atoms around  $E_F$  leads to electron mass enhancements in hydrogenated  $\text{Pr}_{0.75}\text{Gd}_{0.25}\text{ScGe}$  compared to the non-hydrogenated parent. Although  $\gamma$  is enhanced with H-insertion, the increase is far from sufficient to classify the hydride as a heavy-fermion system.

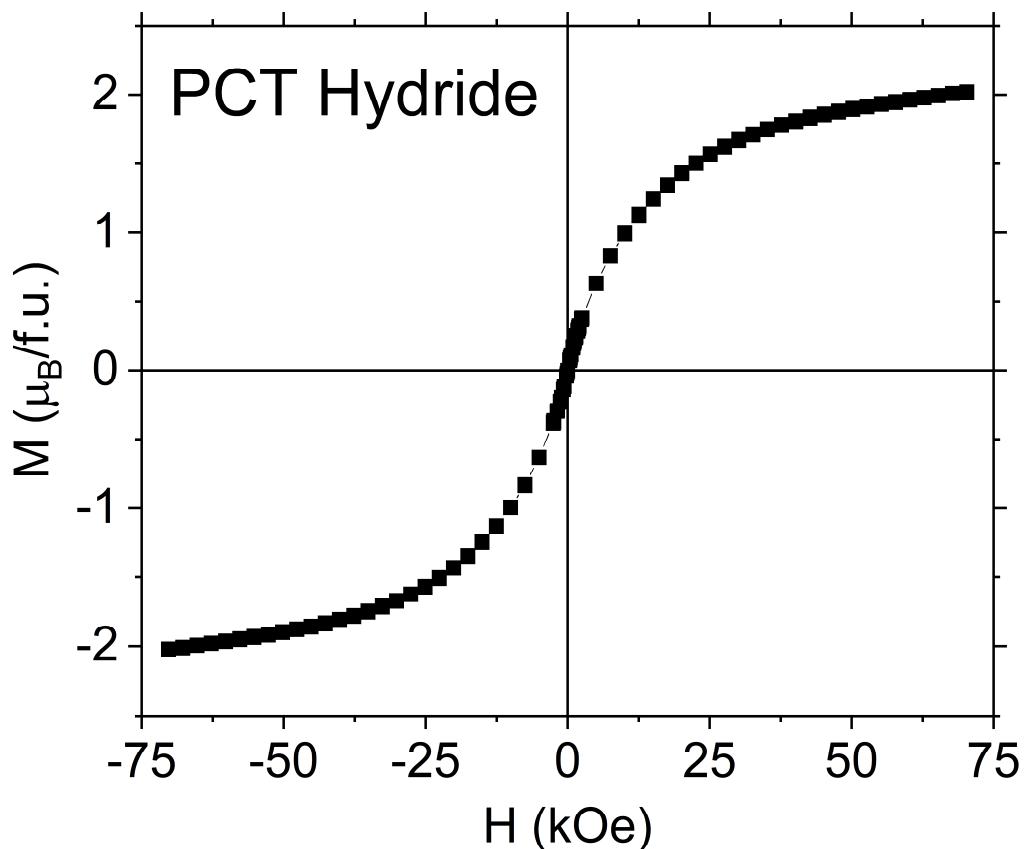


Figure S4: Magnetization of the PCTPro hydride measured as a function of applied magnetic field at  $T = 2$  K.

## References

- [1] Mahon, T. *et al.* “Hydrogen insertion in the intermetallic GdScGe: A drastic reduction of the dimensionality of the magnetic and transport properties.” *Inorganic Chemistry* 57, 14230 (2018).
- [2] Fulde, P. & Jensen, J. “Electronic heat capacity of the rare-earth metals.” *Physical Review B* 27, 4085 (1983).