

Interplay Between Kondo and Magnetic Interactions in $\text{Pr}_{0.75}\text{Gd}_{0.25}\text{ScGeH}$

Tyler Del Rose^{a,b,*}, Renu Choudhary^b, Yaroslav Mudryk^b, Daniel Haskel^c, Arjun K. Pathak^d, Gourab Bhaskar^e, Julia V. Zaikina^e, Duane D. Johnson^{a,b}, and Vitalij K. Pecharsky^{a,b}

^a Department of Materials Science and Engineering, Iowa State University, Ames, Iowa 50011, USA

^b Ames National Laboratory, U.S. Department of Energy, Iowa State University, Ames, Iowa 50011, USA

^c Advanced Photon Source, Argonne National Laboratory, Argonne, IL 60439, USA

^d Department of Physics, SUNY Buffalo State University, Buffalo, New York 14222, USA

^e Department of Chemistry, Iowa State University, Ames, Iowa 50011, USA

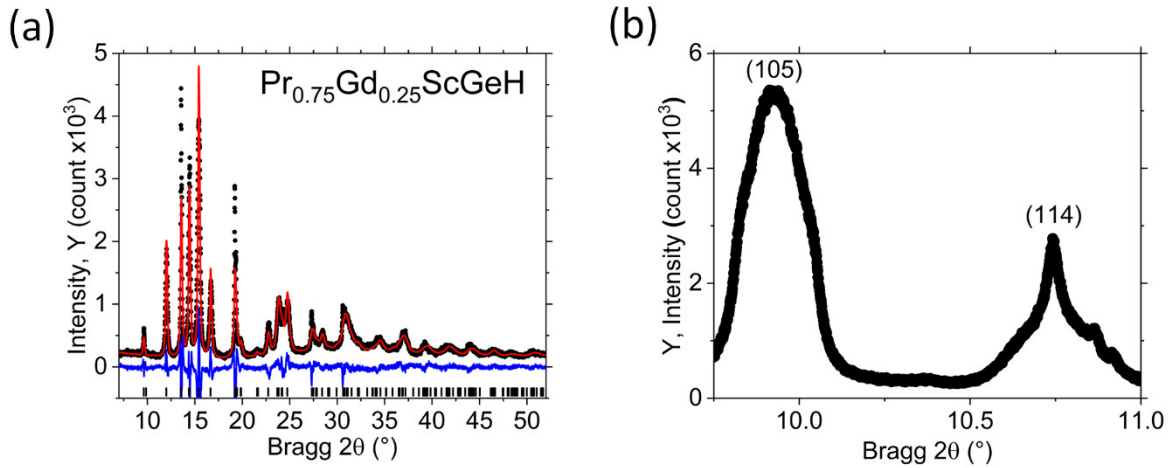


Figure S1. (a) Rietveld refinement and (b) the shapes of (105) and (114) Bragg peaks of $\text{Pr}_{0.75}\text{Gd}_{0.25}\text{ScGeH}$ hydrogenated at 150 bar. The data presented in (a) are from a laboratory powder diffractometer using Mo-K α radiation, while the data in (b) are from APS, ANL using $\lambda = 0.457897$ Å. The noticeable mismatch of intensities in (a) is due to deteriorated crystallinity and non-analytical peak shapes after hydrogenation, exemplified in (b).

Figure S2 depicts the heat capacity of $\text{Pr}_{0.75}\text{Gd}_{0.25}\text{ScGeH}$ and LaScGeH . Contributions from magnetic and nuclear effects (exemplified by the low-temperature differences between LaScGeH and $\text{Pr}_{0.75}\text{Gd}_{0.25}\text{ScGeH}$) makes low-temperature fitting to determine the electronic specific heat, γ , impractical. Assuming that isostructural LaScGeH has the same Debye

temperature, θ_D , as $\text{Pr}_{0.75}\text{Gd}_{0.25}\text{ScGeH}$, it is still possible to estimate γ for the latter as the constant difference between their C_P/T in the region ($\sim 30 - 80$ K) where both the nuclear and magnetic contributions become negligible, see Eq. S1 and Figure S2.

$$\gamma_{\text{Pr}_{0.75}\text{Gd}_{0.25}\text{ScGeH}} = \gamma_{\text{LaScGeH}} + \left(\left(\frac{C_P}{T} \right)_{\text{Pr}_{0.75}\text{Gd}_{0.25}\text{ScGeH}} - \left(\frac{C_P}{T} \right)_{\text{LaScGeH}} \right) \quad (\text{Eq. S1})$$

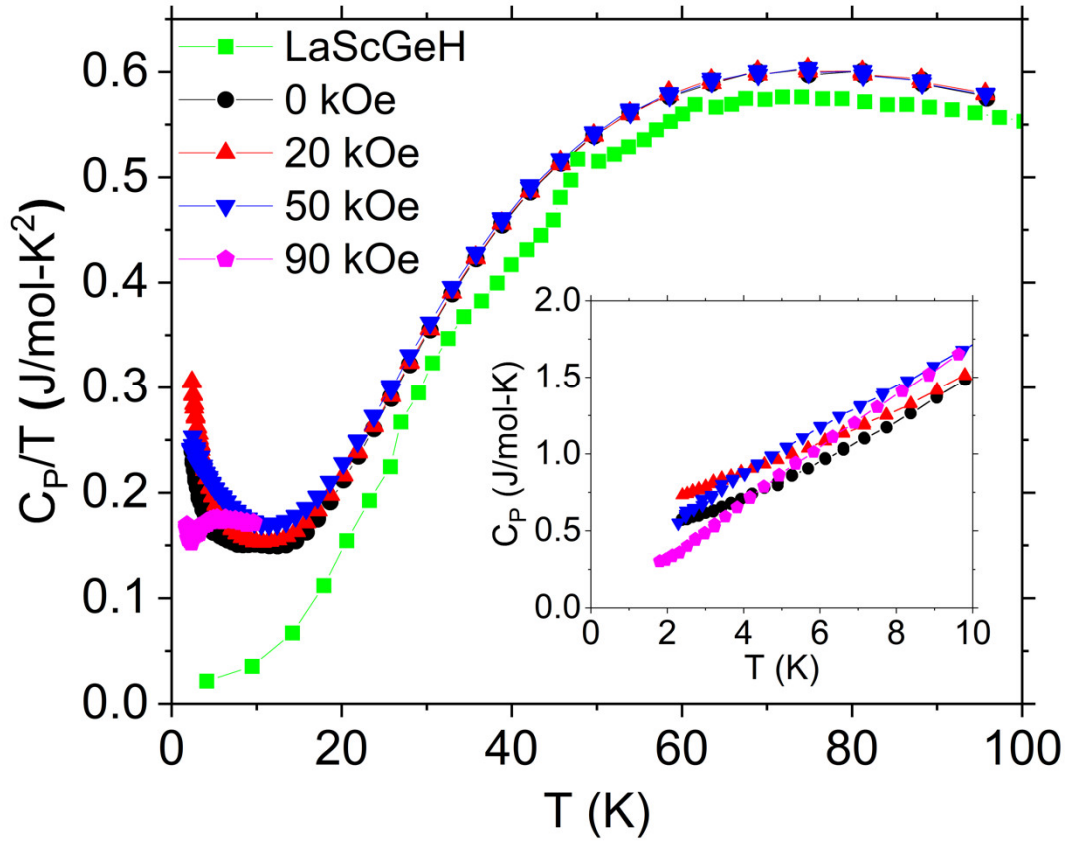


Figure S2. Heat capacity (C_P) of $\text{Pr}_{0.75}\text{Gd}_{0.25}\text{ScGeH}$ measured in zero and applied magnetic fields of 20, 50, and 90 kOe shown as C_P/T vs. temperature (T). LaScGeH data are taken from Ref [1]. Inset shows low-temperature details as C_P vs T .

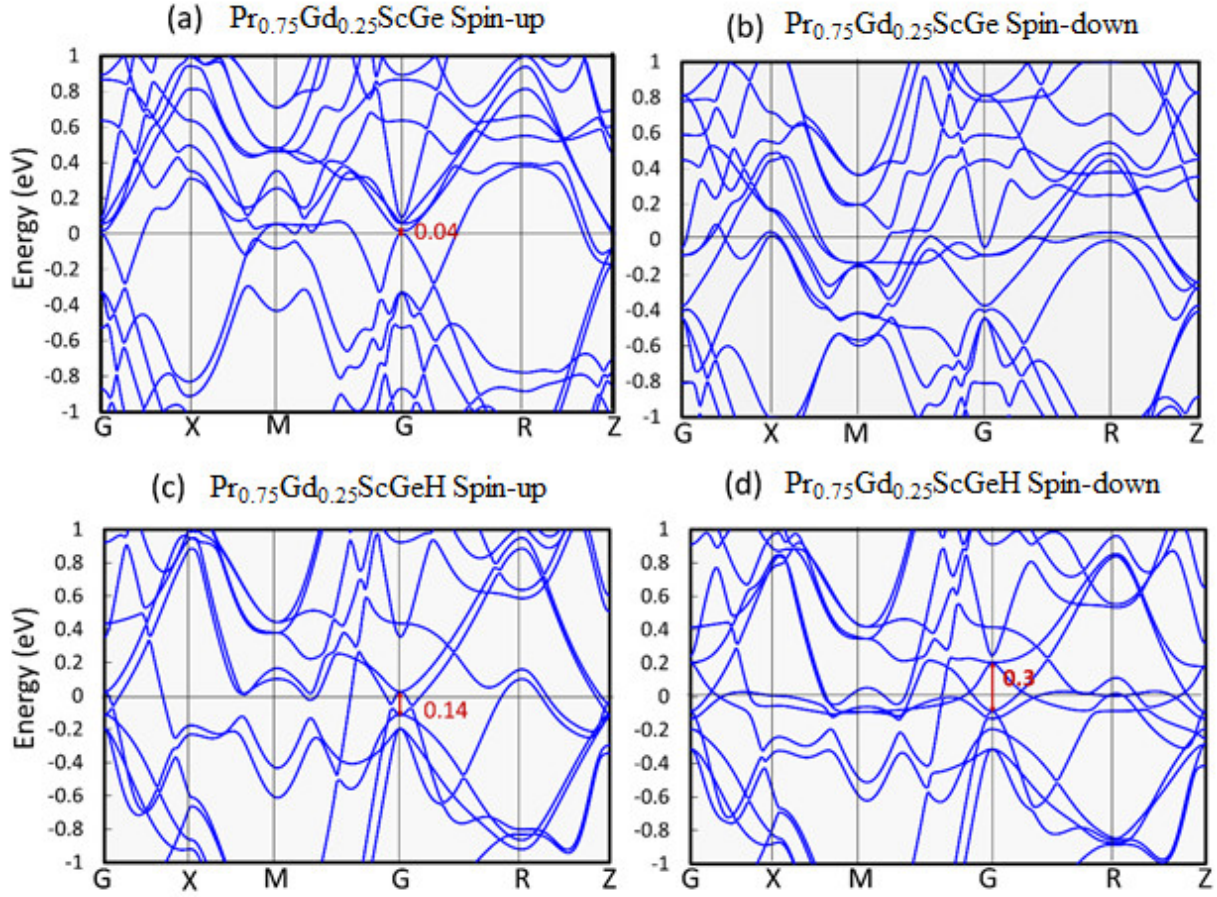


Figure S3: Spin-polarized DFT band structure of (a, b) $\text{Pr}_{0.75}\text{Gd}_{0.25}\text{ScGe}$ and (c, d) its hydride, showing both spin-up (a, c) and spin-down (b, d) electron manifolds.

The band structures, plotted in Figure S3, indicate a sharp peak of the $4f$ states close to the E_F . Additionally, the Gamma point has a very narrow gap at E_F in the majority-spin channel (Figure S3a), whereas the minority-spin channel is metallic in $\text{Pr}_{0.75}\text{Gd}_{0.25}\text{ScGe}$ (Figure S3b). As the density of states for a system is inversely proportional to the derivative of energy dispersion E_k , the sharp minority-spin $4f$ -peak in PDOS (see Figure 4b) of hydrogenated $\text{Pr}_{0.75}\text{Gd}_{0.25}\text{ScGe}$ around E_F is flat in the band structure (Figure S3b). Flat bands mean larger effective electron mass or possible heavy-fermion behavior however, they are not exactly at the E_F . Therefore, the flat $4f$ band leads to an enhanced effective electron mass without necessarily reflecting heavy-fermion-

like state. To confirm this, electronic specific heat (γ) has been calculated for both $\text{Pr}_{0.75}\text{Gd}_{0.25}\text{ScGe}$ and its hydride, which are 9.35 and 13.22 mJ/mol. K^2 , respectively. The theoretical value of γ for the hydride is smaller than experimental values (28.7 mJ/mol. K^2) by a factor of nearly two, whereas both theory and experiment agree for the $\text{Pr}_{0.75}\text{Gd}_{0.25}\text{ScGe}$ system (Table 1). Electron mass enhancement in rare-earth metals and the effect on specific heat were discussed by Fulde et al. [2]. Experimental electronic specific heat of Pr metal determined from heat capacity data between 1 and 6 K is larger than the DFT value by a factor of ~ 4 , and the difference is reduced by applying a magnetic field. The reason for the mass enhancement of conduction electrons in Pr-metal (and other lanthanides) at low temperature and low magnetic field is the interaction of conduction electrons with localized $4f$ moments, which is neglected in calculating γ by DFT methods. Hence, theoretically calculated values are commonly different from experimental ones. Here, the predicted interactions of conduction and $4f$ -electrons of Pr1 atoms around E_F leads to electron mass enhancements in hydrogenated $\text{Pr}_{0.75}\text{Gd}_{0.25}\text{ScGe}$ compared to the non-hydrogenated parent. Although γ is enhanced with H-insertion, the increase is far from sufficient to classify the hydride as a heavy-fermion system.

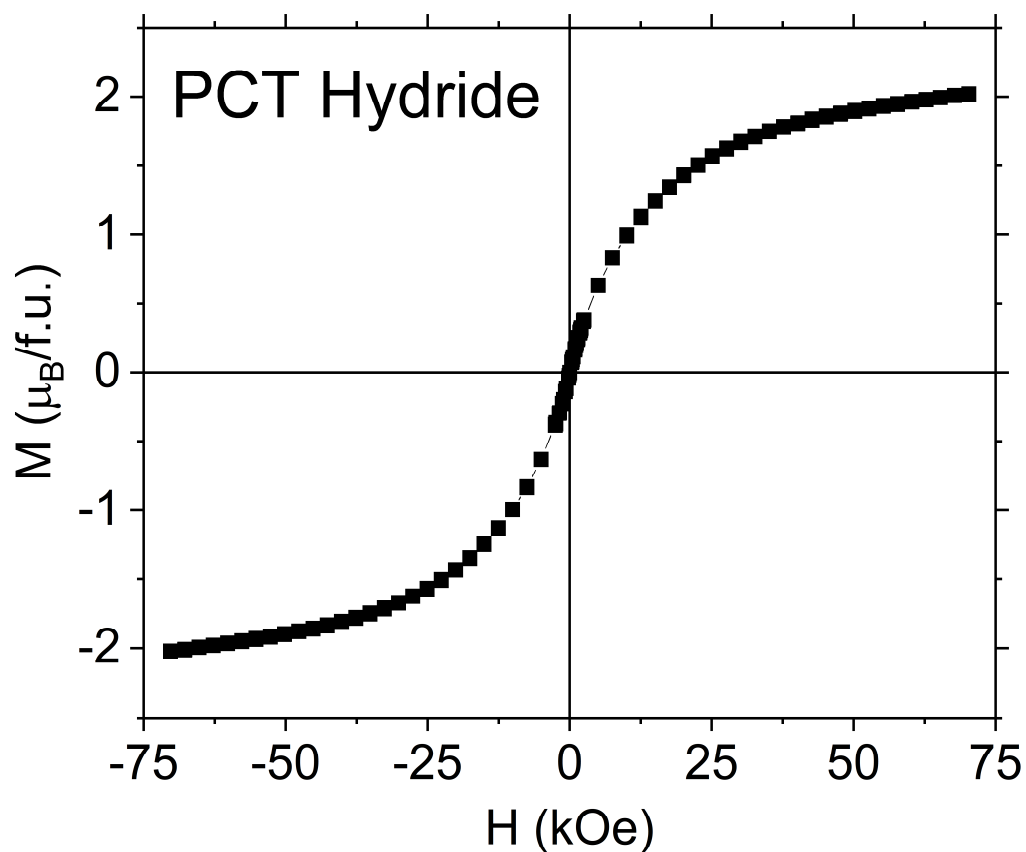


Figure S4: Magnetization of the PCTPro hydride measured as a function of applied magnetic field at $T = 2$ K.

References

- [1] Mahon, T. *et al.* "Hydrogen insertion in the intermetallic GdScGe: A drastic reduction of the dimensionality of the magnetic and transport properties." *Inorganic Chemistry* 57, 14230 (2018).
- [2] Fulde, P. & Jensen, J. "Electronic heat capacity of the rare-earth metals." *Physical Review B* 27, 4085 (1983).