

# **Group-V acceptor ionization energies and compensation centers in CdTe revisited**

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# In-situ antimony doped polycrystalline CdTe films for simplified cell processing and maximized energy

US DOE - Solar Energy Technologies Office: Award # DE-EE0009344

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# Polycrystalline CdTe-based thin-film solar cells

Record power conversion efficiency (PCE) > 22%, First Solar (2016-2017)

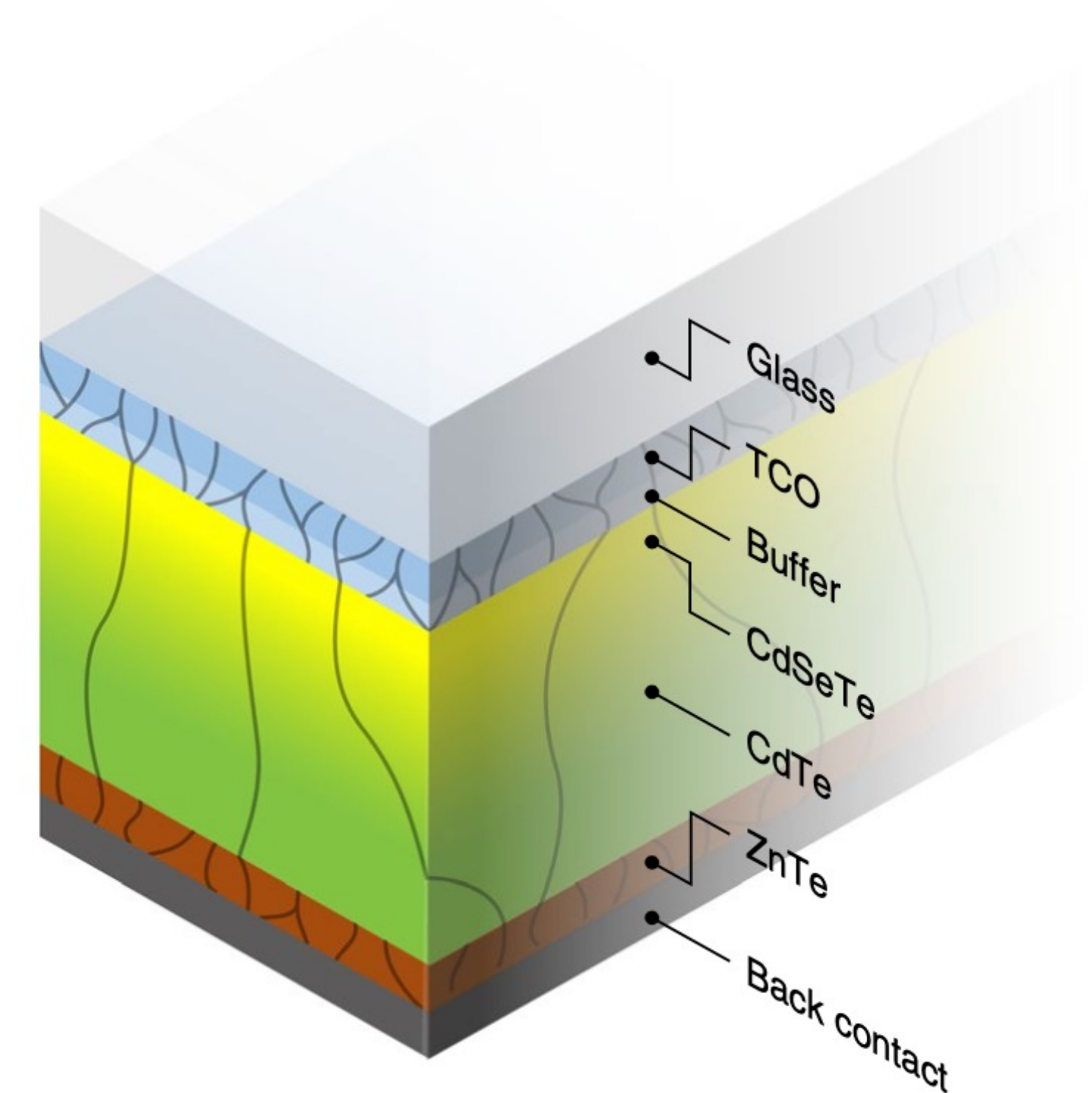
> 18% for modules

- limited by low  $V_{OC} < 0.9$  V
- low carrier concentration ( $\sim 2 \times 10^{14} \text{ cm}^{-3}$ )
- poor bulk and surface passivation

Device modeling indicates that efficiency can be increased from 22% to 25% by increasing hole density to  $> 2 \times 10^{16} \text{ cm}^{-3}$ , provided that long carrier lifetimes and low interface recombination velocities can be attained

Single-crystal CdTe with hole density  $p > 10^{16} \text{ cm}^{-3}$  demonstrated,  
 $\Rightarrow$  maintaining  $\sim 20$ -nanosecond carrier lifetimes can increase the  $V_{OC}$  to  $> 1$

**20.8% efficiency without antireflection coatings demonstrated for polycrystalline film, with hole density of  $10^{16}$ - $10^{17} \text{ cm}^{-3}$  without compromising the lifetime**  
 $\Rightarrow$  resolving interfacial and potential fluctuation issues may enable further improvement



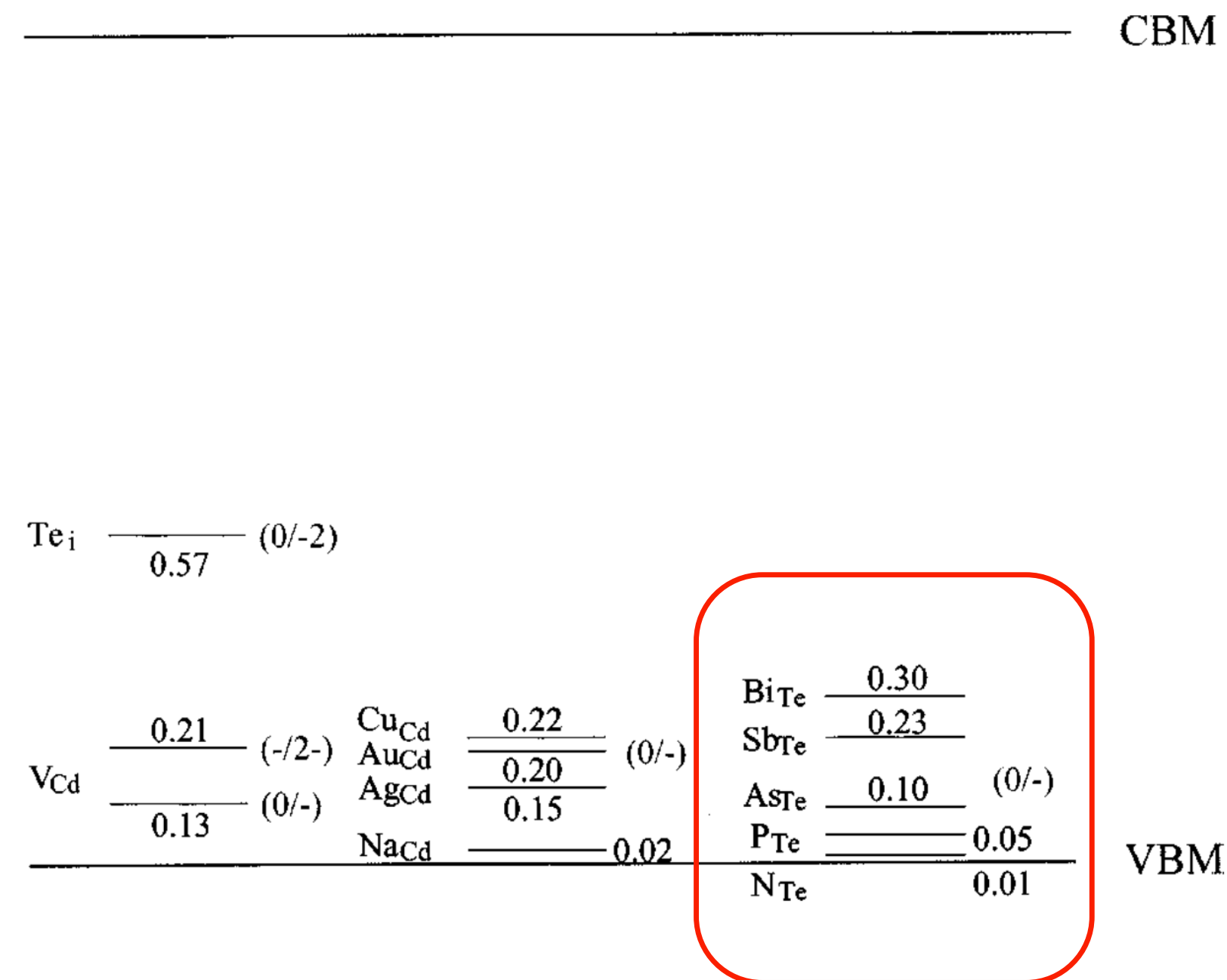
W. K. Metzger *et al.*, Nature Energy **4**, 837 (2019)

Green *et al.*, Solar cell efficiency tables (version 56), *Prog. Photovolt. Res. Appl.* **28**, 629 (2020)

A. Romeo and E. Arregiani, Energies **14**, 1684 (2021)

M. Gloeckler, FSLR, presented at 44th IEEE PVSC, Washington (2017)

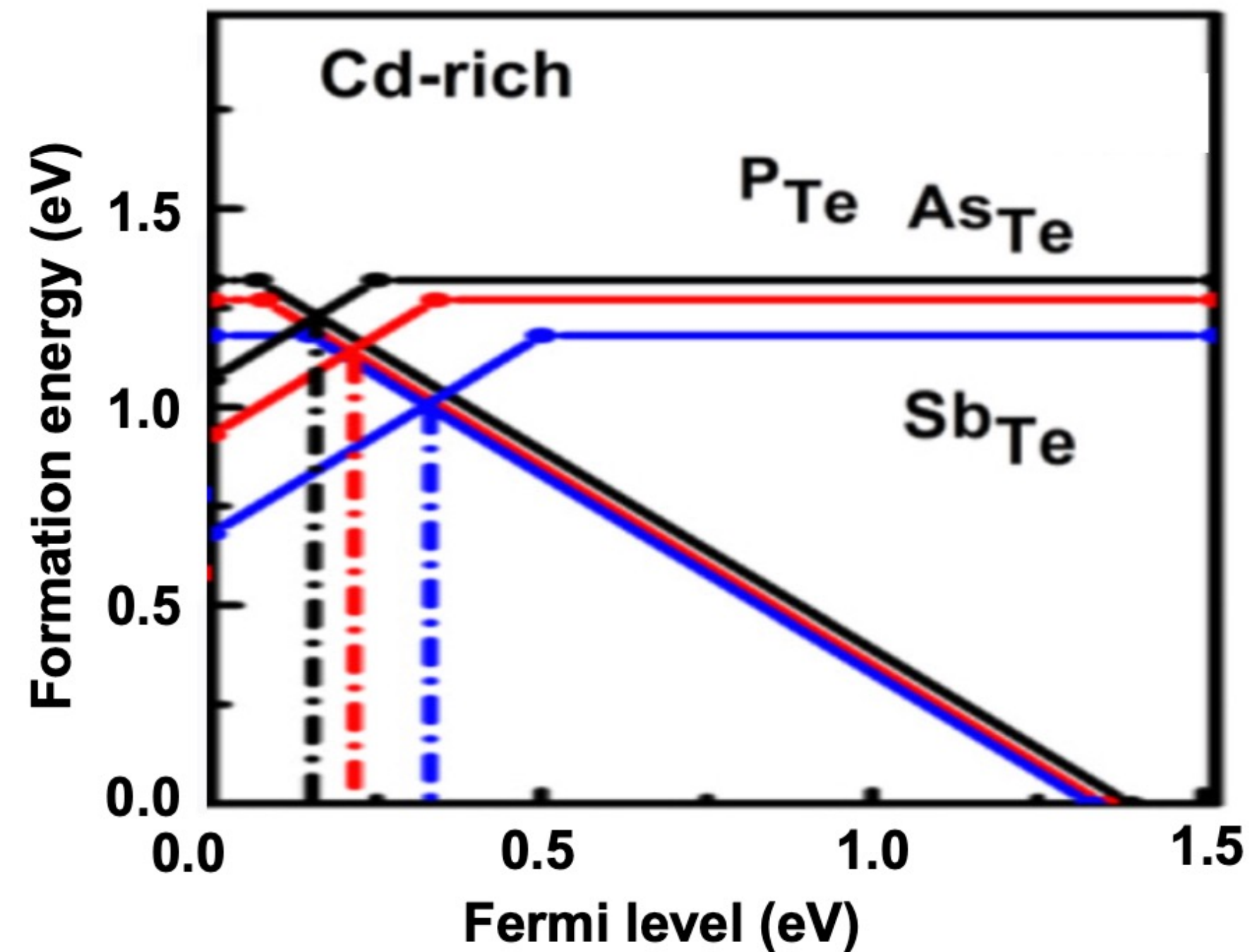
# Are Sb, As, and P shallow acceptors in CdTe? How shallow? Do they suffer from self compensation (i.e., by AX formation)?



S.-H. Wei and S. B. Zhang,  
phys. stat. sol. (b) **229**, 305 (2002);  
Phys. Rev. B **66**, 155211 (2002)

**DFT-LDA, supercell of 32 atoms,  
no spin-orbit coupling (SOC)**

- Very large ionization energy for Sb



B. Dou, Q. Sun, and S.-H. Wei,  
Phys. Rev. Appl. **15**, 054045 (2021)

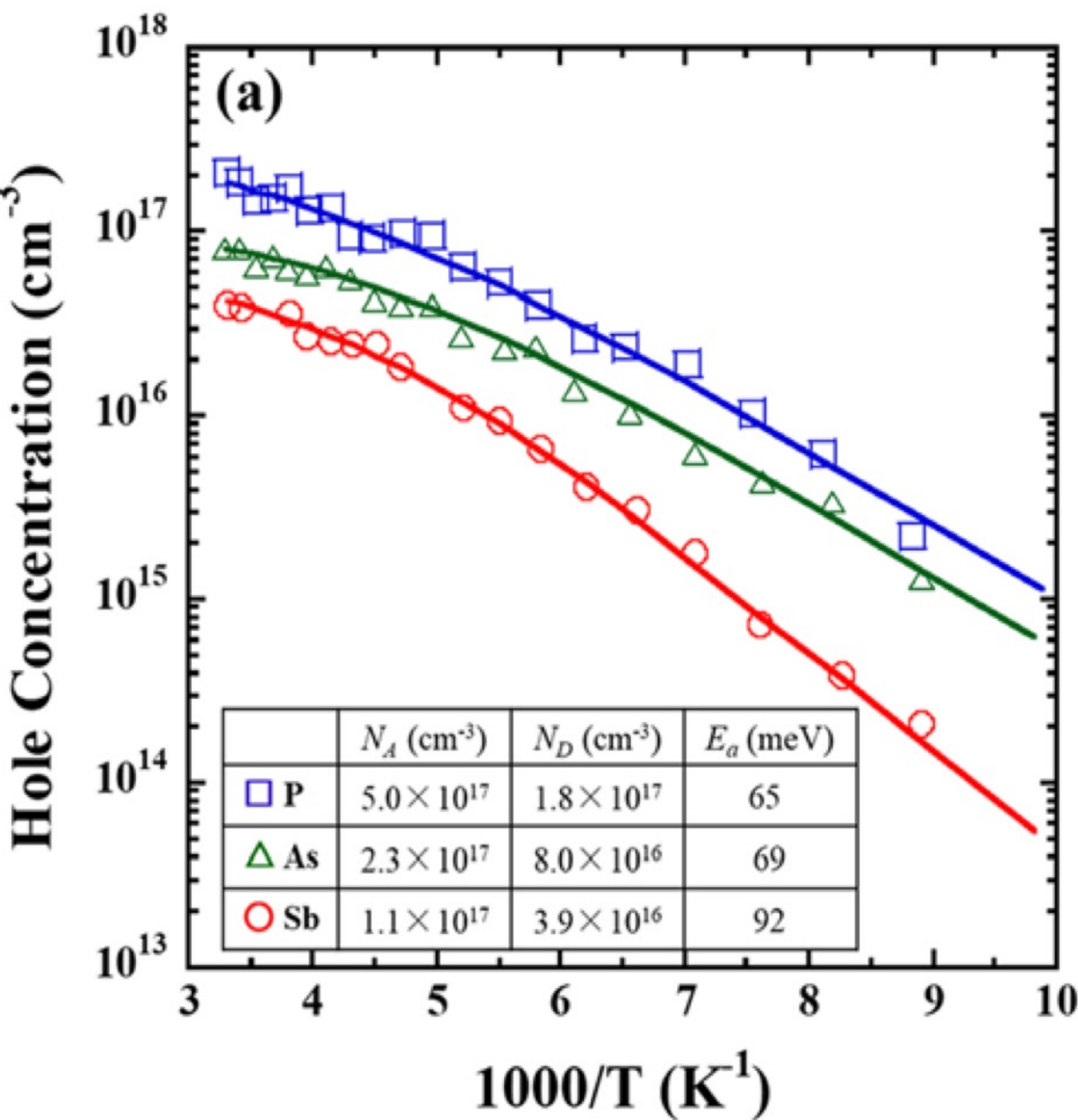
**DFT- hybrid functional,  
supercell of 216 atoms, no SOC**

- Still large ionization energy for Sb (150 meV)
- P, As, and Sb will form AX, killing hole conductivity

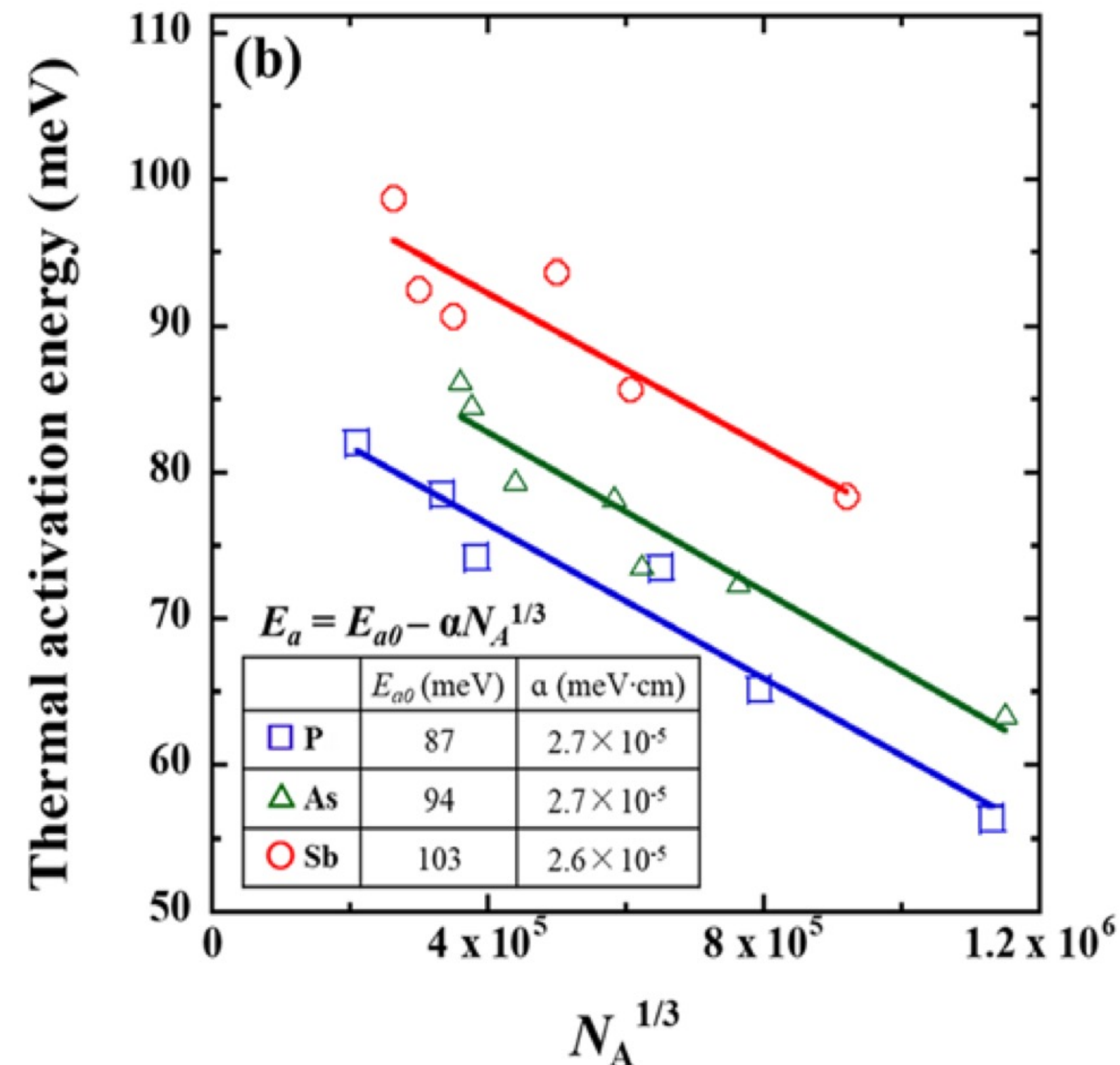


# Sb, As and P doping of CdTe single crystals

Recent experiments indicate that Sb, As and P are shallow acceptors with ionization energies of ~100 meV



Temperature dependence of hole concentration for group V-doped samples

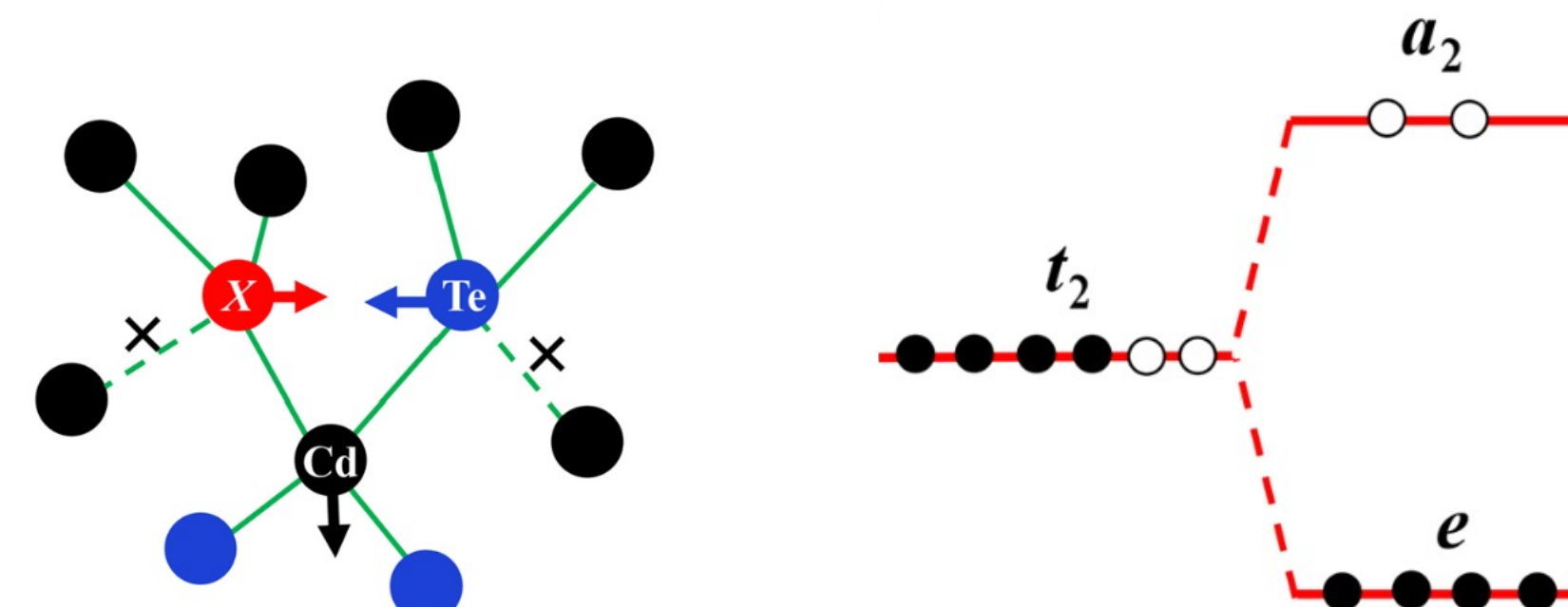
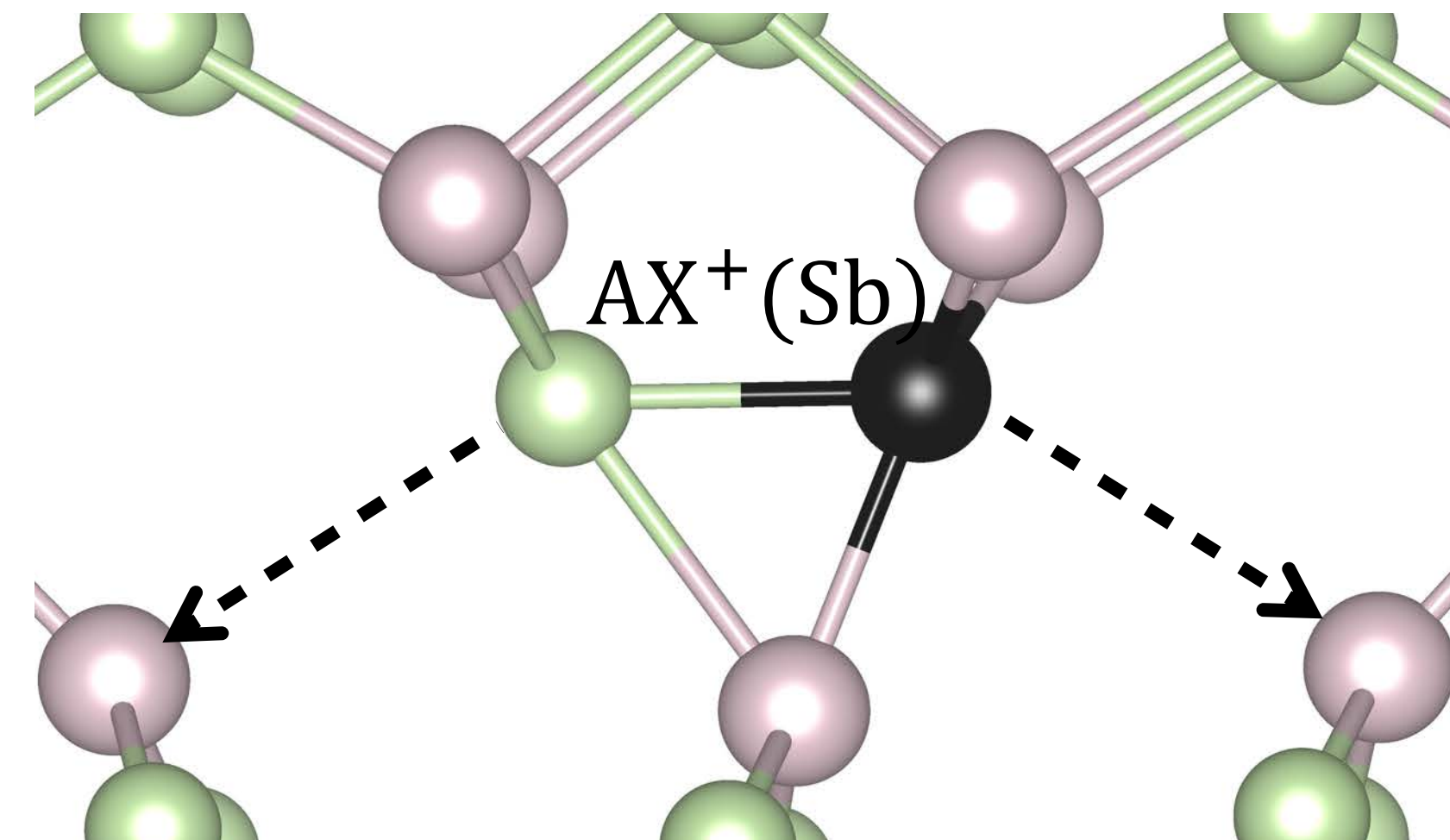
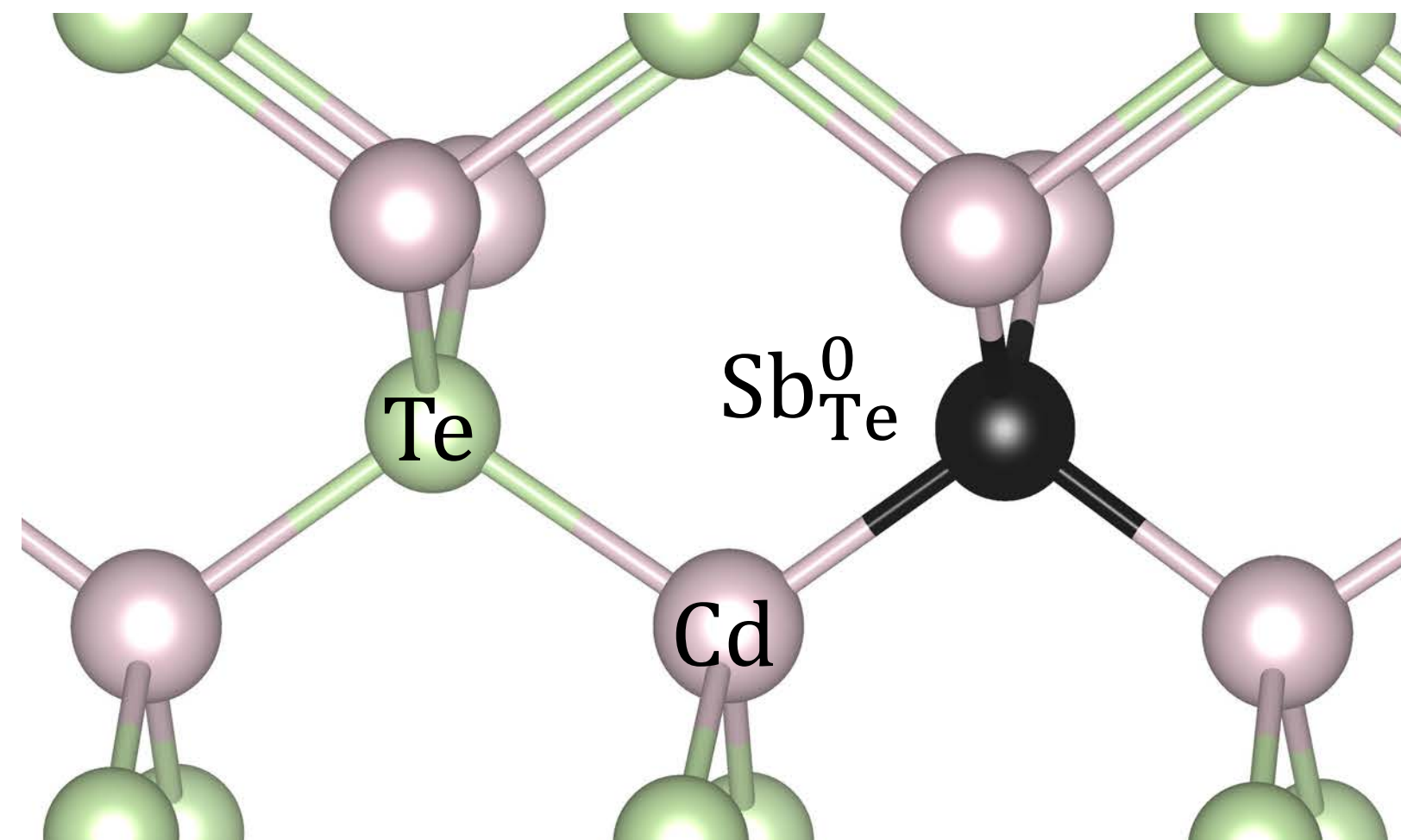


Thermal activation energy vs acceptor concentration

$$p = -A + \sqrt{A^2 + \frac{N_V}{2} (N_A - N_D) \exp\left(-\frac{E_a}{k_B T}\right)},$$

$$A = \frac{1}{2} \left[ N_D + \frac{N_V}{2} \exp\left(-\frac{E_a}{k_B T}\right) \right],$$

# Substitutional acceptor vs. AX donor



S.-H. Wei and S. B. Zhang, Phys. Rev. B **66**, 155211 (2002)  
 J.-H. Yang et al., Semicond. Sci. Technol. **31**, 083002 (2016)  
 B. Dou et al., Phys. Rev. Appl. **15**, 054045 (2021)  
 C. H. Park and D. J. Chadi, Phys. Rev. Lett. **75**, 1134 (1995)

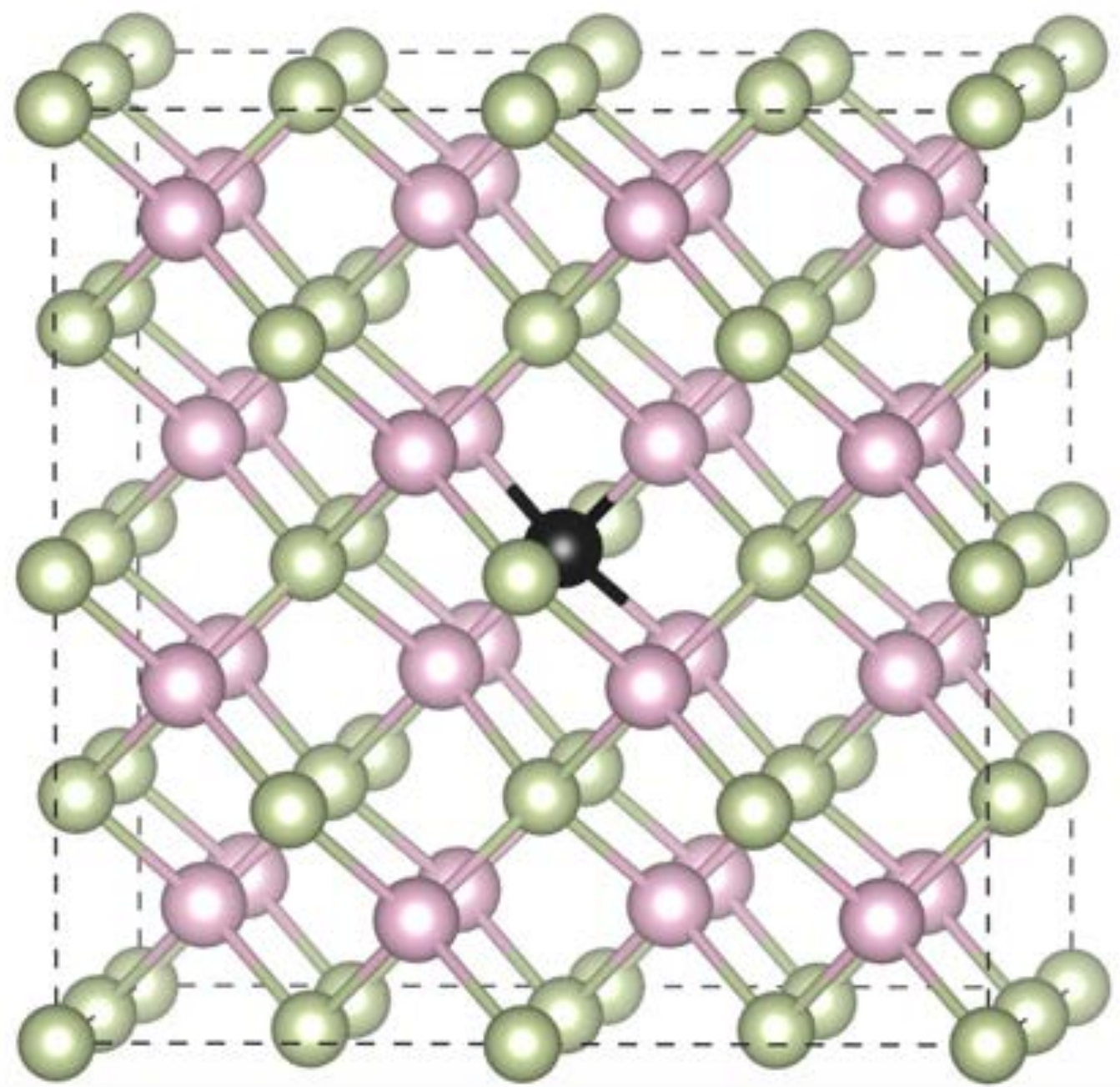
The acceptor is displaced along the  $[110]$  direction, forming a bond with a second nearest neighbor Te becoming a donor



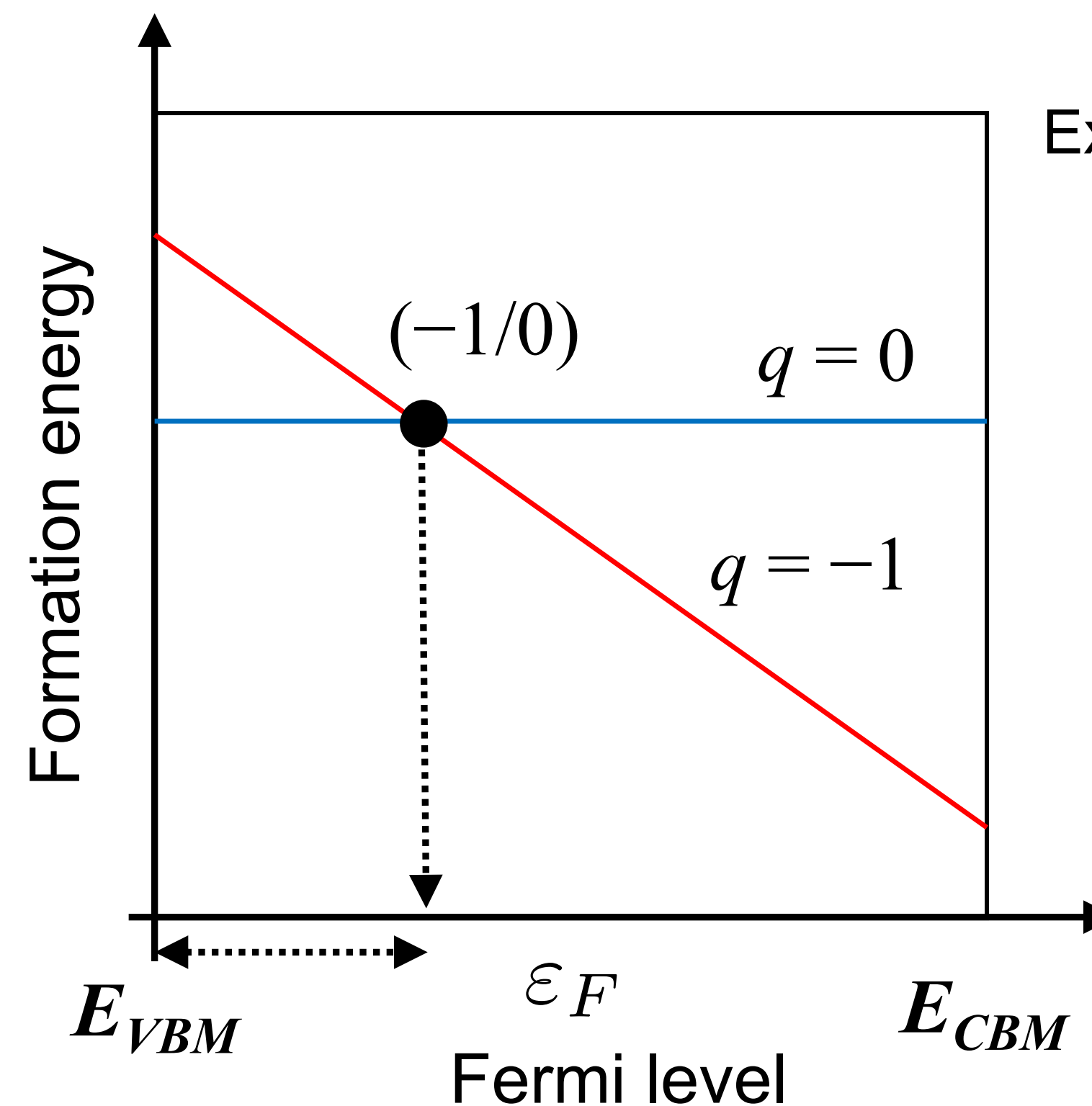
# How to calculate defect formation energies

$$E^f(X^q) = E_{tot}(X^q) - E_{tot}(\text{bulk}) + \sum_i [E_{tot}(X_i) + n_i \mu_i] + q(\epsilon_F + E_{VBM}) + \Delta^q$$

64-atoms, zinc blende



Supercell with a defect  
periodically repeated in 3D  
All atoms in the supercell are allowed to relax  
(minimizing forces and total energy)



Ex: acceptor defect

$(-1/0)$ : acceptor transition level  
or ionization energy

For shallow acceptors,  
 $(-1/0) \sim \text{few } k_B T \text{ at RT}$

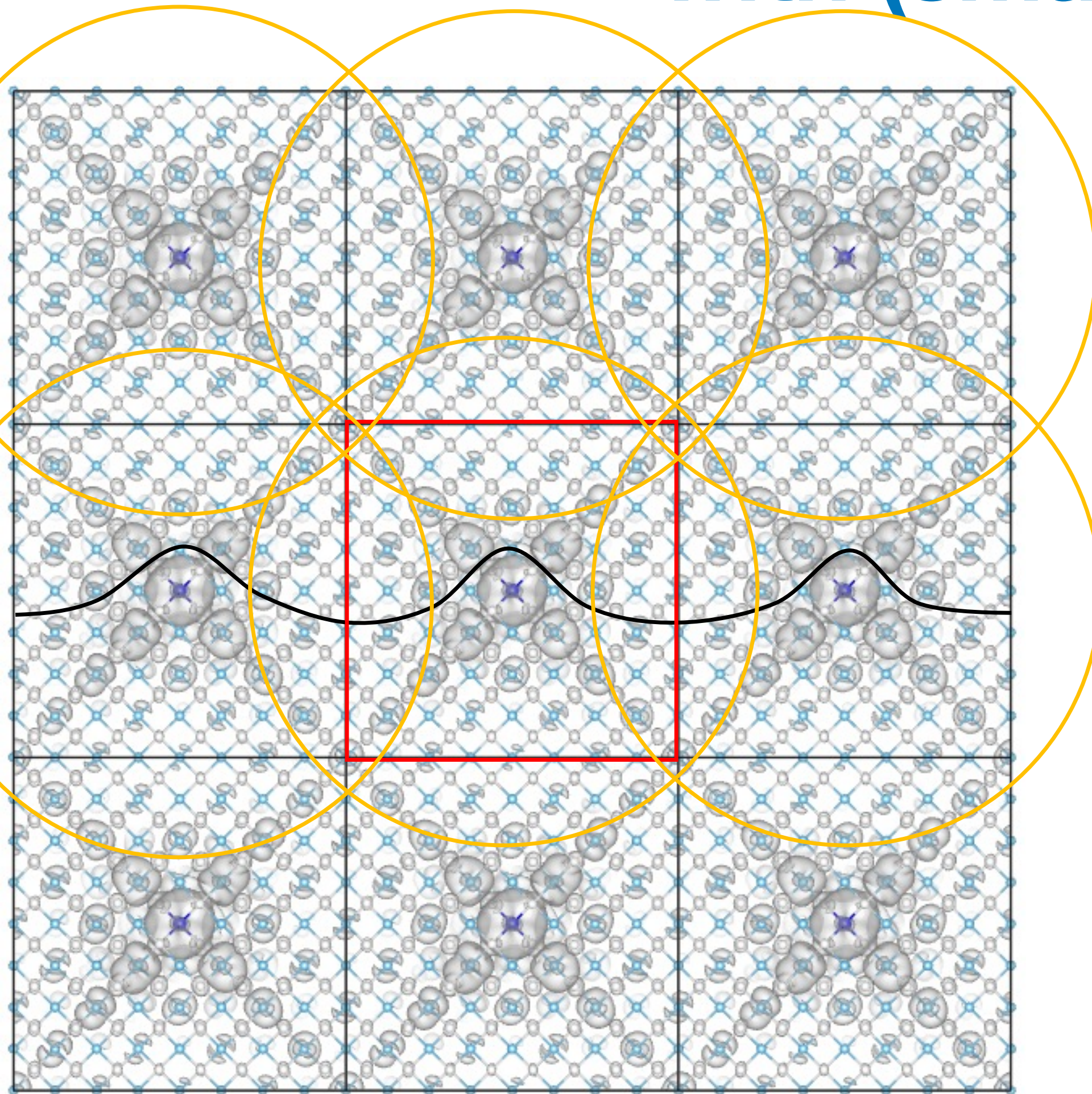
Dependence on atom chemical potential

$$\mu_{\text{Cd}} + \mu_{\text{Te}} = \Delta H_f(\text{CdTe})$$

Cd-rich or Te-rich just shift up/down  
the two curves together



# Problems with describing shallow acceptors or donors with (small) finite supercell sizes



$\text{Sb}_{\text{Te}}$  related state at the top of VBM

supercell of 216 atoms, periodically repeated

Overlap of “hydrogenic” wavefunctions  
leads to artificially increased ionization energies

- 1) DFT within LDA/GGA severely underestimate band gaps  
⇒ large errors in formation energies and transition levels ( $\gg 0.2$  eV)
- 2) Need to correct band structure (both band gap and ionization potential)  
⇒ Hybrid functionals, have to include spin-orbit coupling  
Still, errors in transition levels are typically  $\sim 0.1$  eV

In the case of shallow donors/acceptor,  
the overlap of the hydrogenic wavefunction between the impurity  
and its images using typical supercell sizes lead to overestimation of  
ionization energies

“central-cell” effects

M.W. Swift *et al.*, Npj Comput. Mater. **6**, 181(2020)

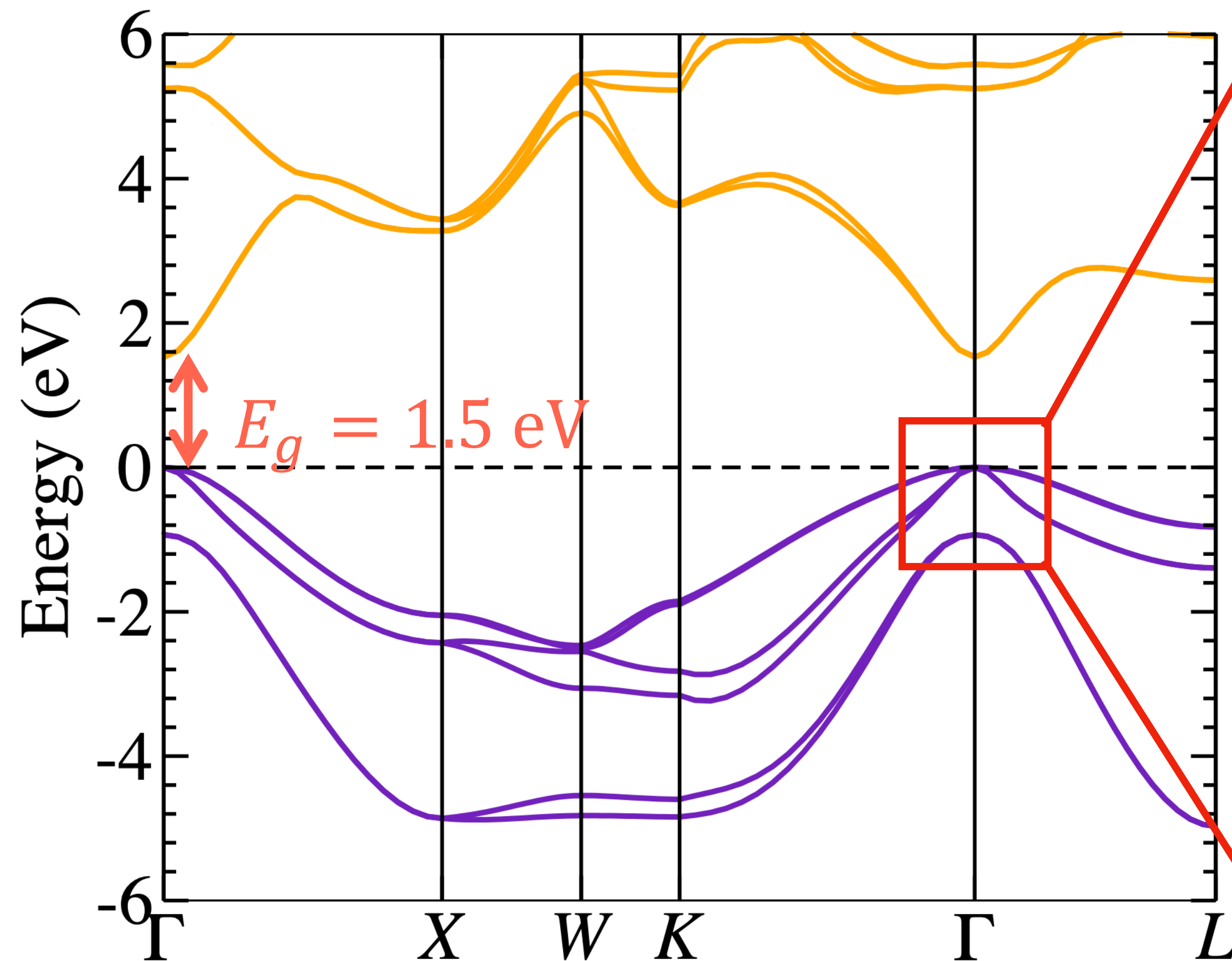
L.-W. Wang, J. Appl. Phys. **105**, 123712 (2009)

**Solution** ⇒ calculate transition levels as function of supercell size  
and extrapolate to the dilute limit

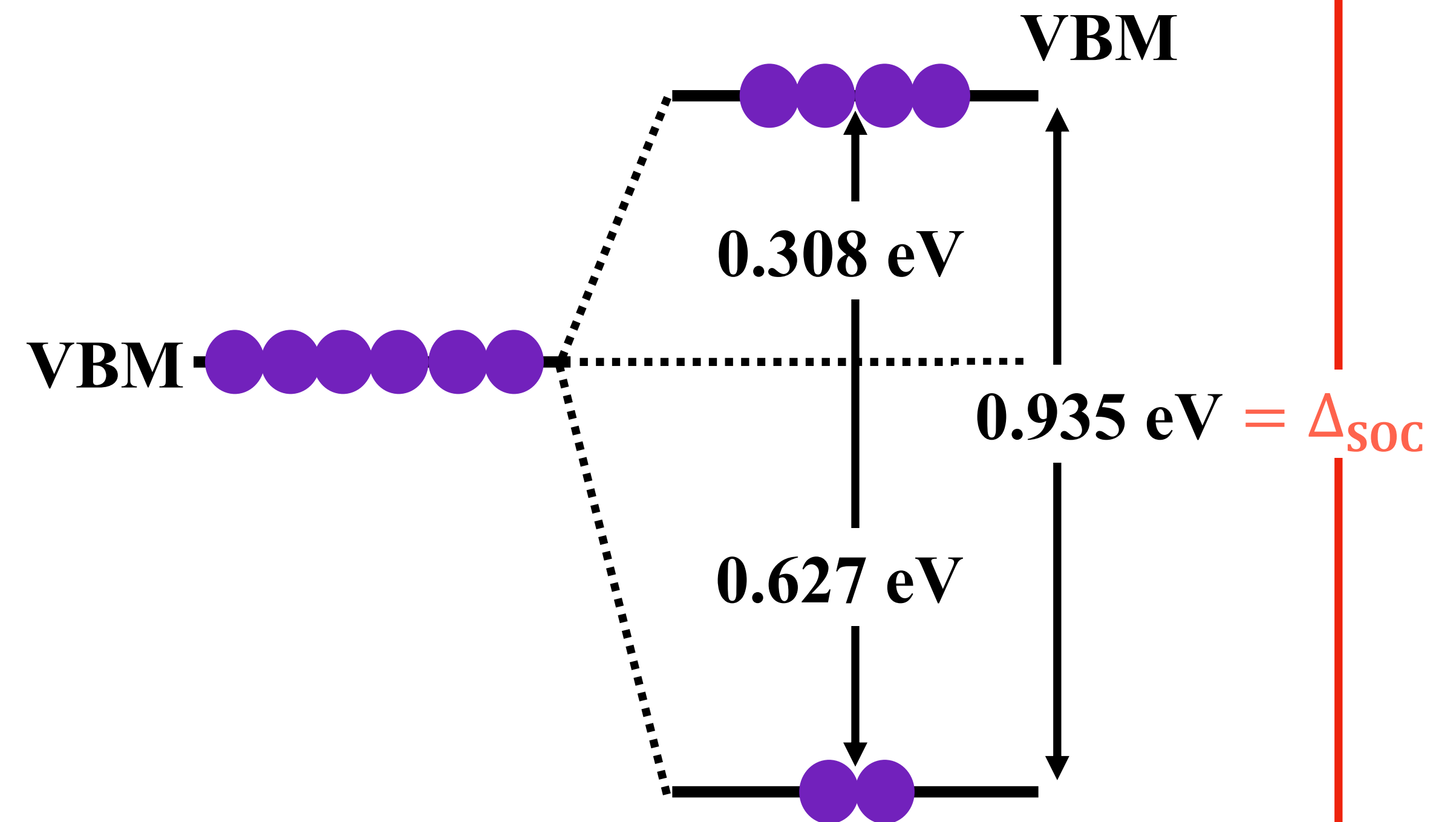


# Band structure of CdTe, effects of spin-orbit coupling

HSE  $\alpha = 0.33$  with SOC



Effect of SOC

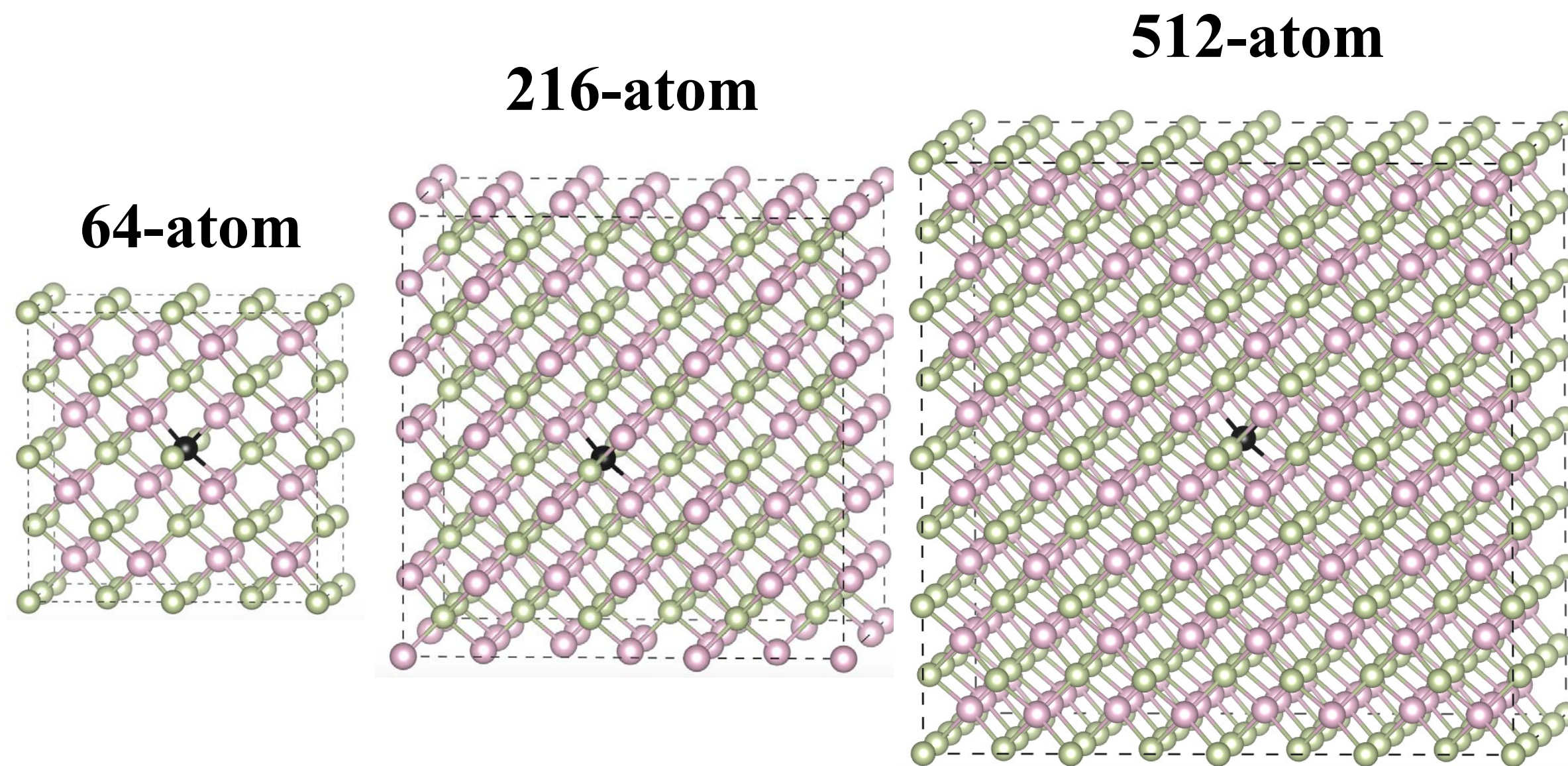


Ionization potential in much better agreement with exp. data compared to DFT-LDA/GGA

See also J. Pan, *et al.*, Phys. Rev. B **98**, 054108 (2018)

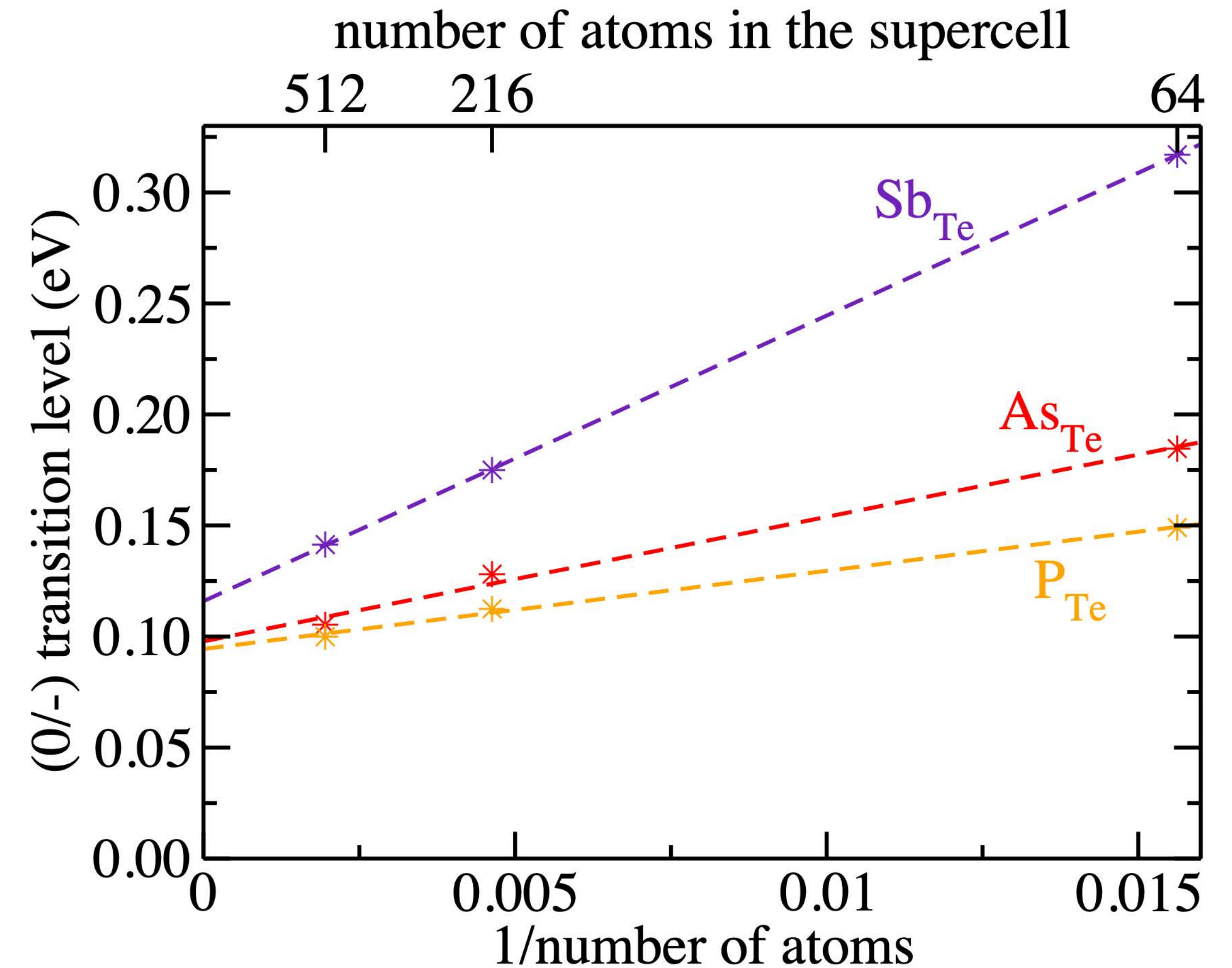


# Extrapolation of the acceptor level to the dilute limit



- At dilute limit,  $\text{Sb}(0/-)=116 \text{ meV}$ ,  $\text{As}(0/-)=99 \text{ meV}$  and  $\text{P}(0/-)=93 \text{ meV}$  (close to hydrogenic model  $\sim 100 \text{ meV}$ )
- From experiment,  $\text{Sb}(0/-)=0.103 \text{ eV}$ ,  $\text{As}(0/-)=0.094 \text{ eV}$  and  $\text{P}(0/-)=0.087 \text{ eV}$

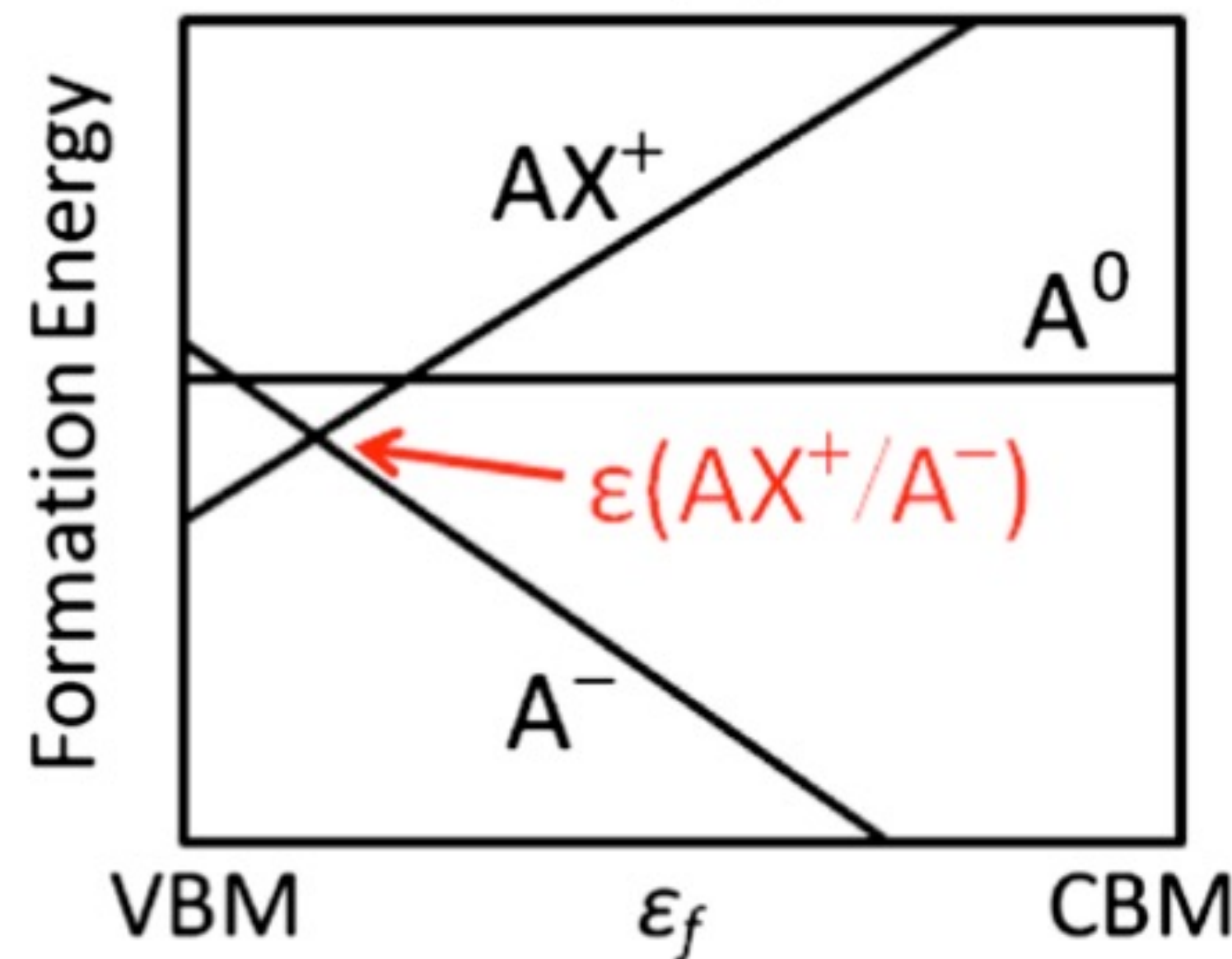
A. Nagaoka *et al.*, Appl. Phys. Lett. **116**, 132102 (2020)



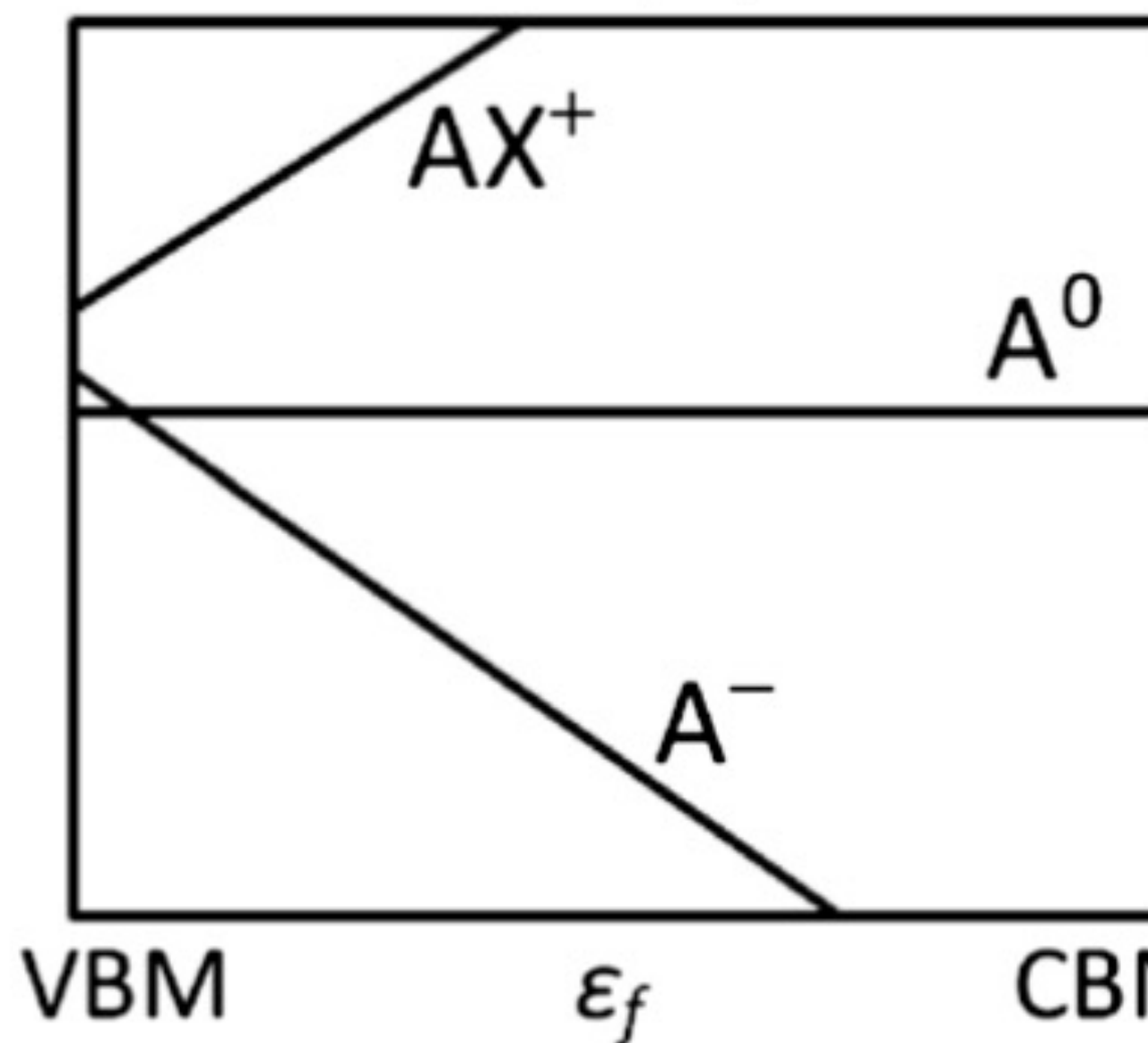


# Formation of AX centers: Three possible scenarios

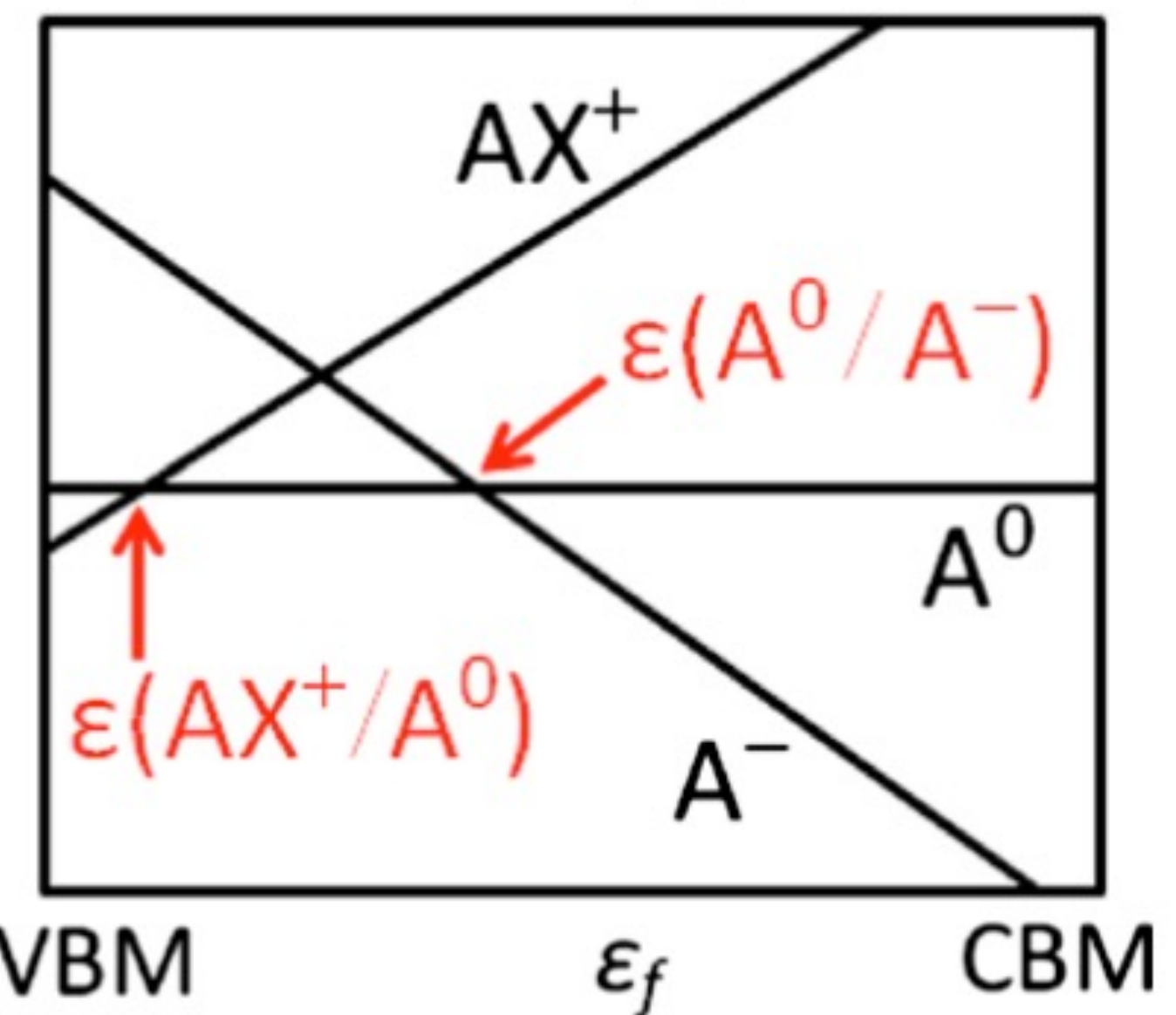
AX is stable,  
kills hole conductivity



AX is unstable,  
hole conductivity limited by ( $A^0/A^-$ )

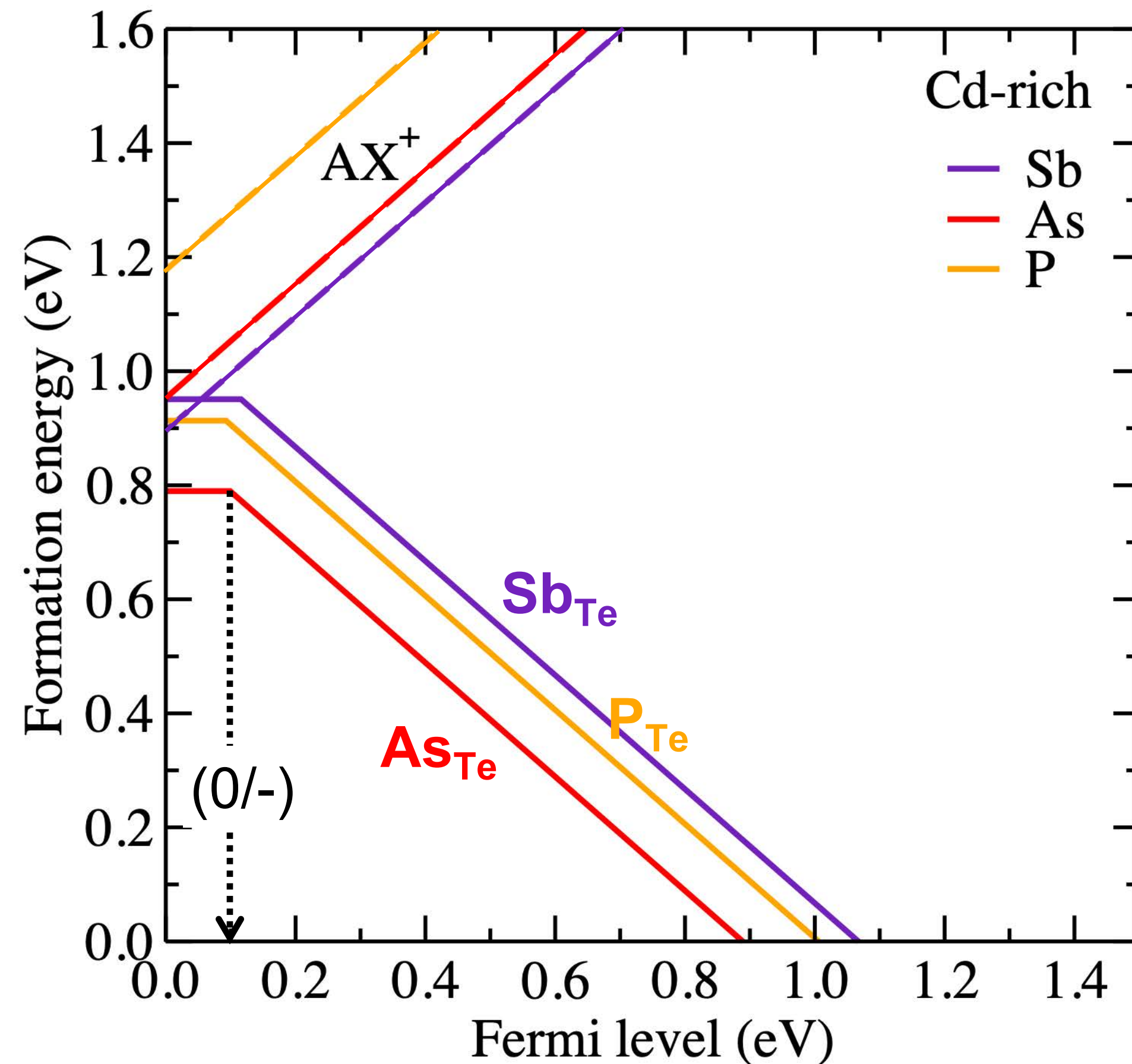


AX is stable,  
hole conductivity may be limited  
by ( $A^0/A^-$ ) and AX





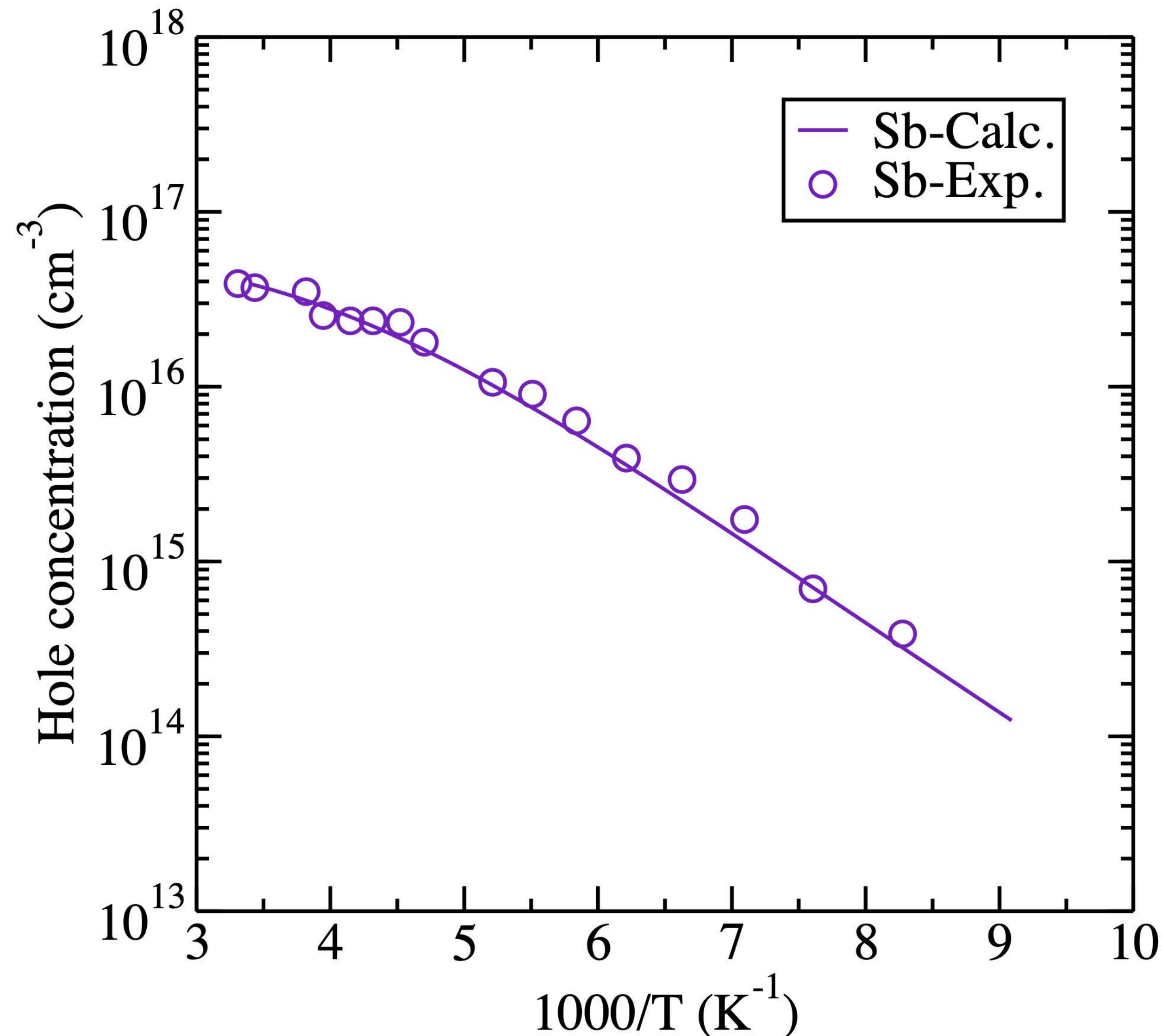
# Group-V acceptors and their corresponding AX centers



- AX is unstable in the case of As and P  
i.e., does not play any role in the doping efficiency  
⇒ problems with As doping must be caused by something else (or some other center)
- In the case of Sb, AX will have a limited effect since the Fermi level typically does not reach (+/0) level near the VBM



# Temperature-dependent hole concentration for Sb doping of CdTe



Use:

- Calculated  $E_a = 116$  meV
- Calculated (+/-) level ( $AX+/Sb_{Te}-$ )

Solve partially compensated semiconductor equation to obtain  $N_a$  that best fits experimental data

$$\frac{p(p + N_d)}{N_a - N_d - p} = \frac{N_V}{\beta} e^{-E_a/k_B T}$$

J. S. Blakemore, Semiconductor Statistics (Pergamon, New York, 2013)

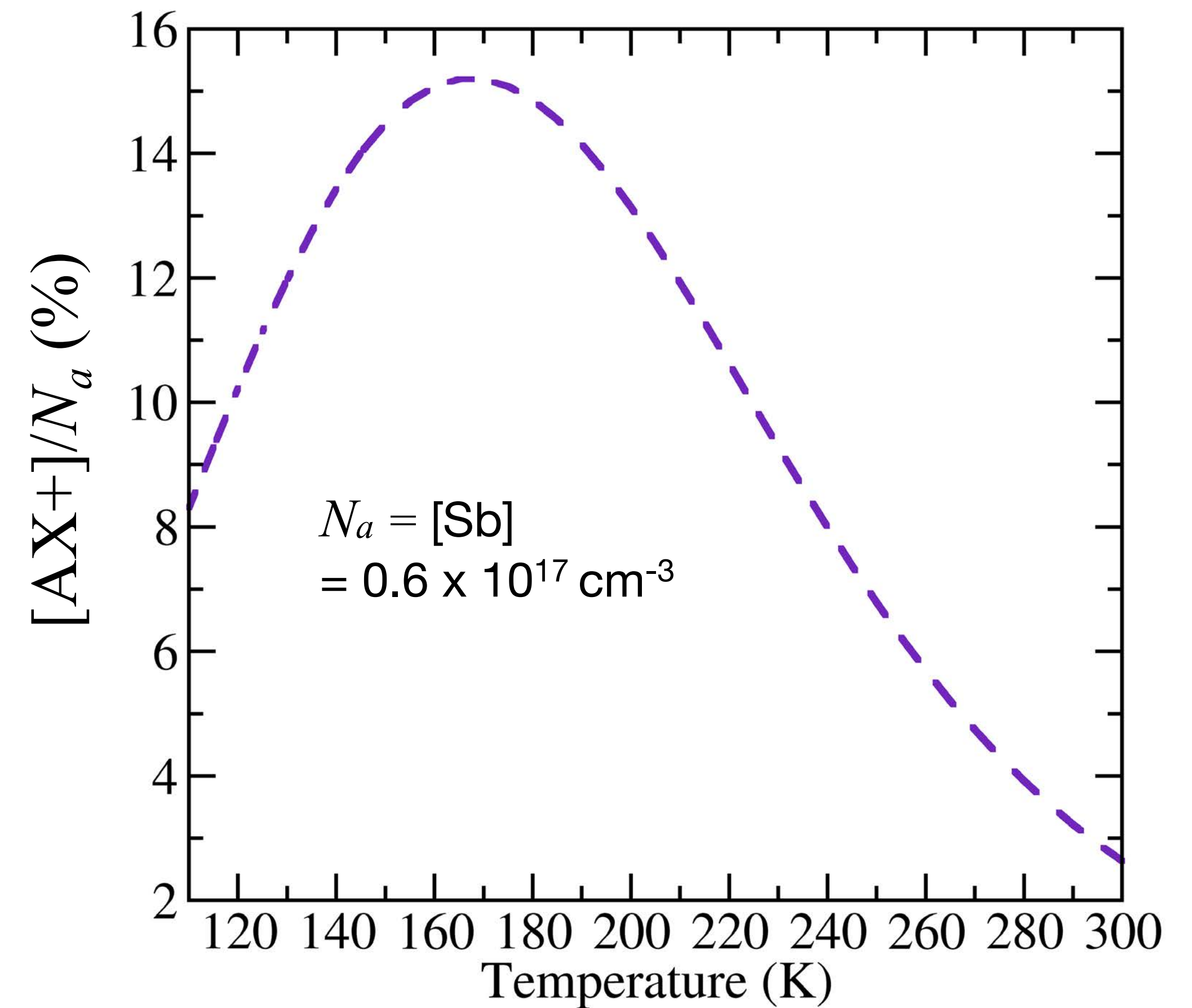
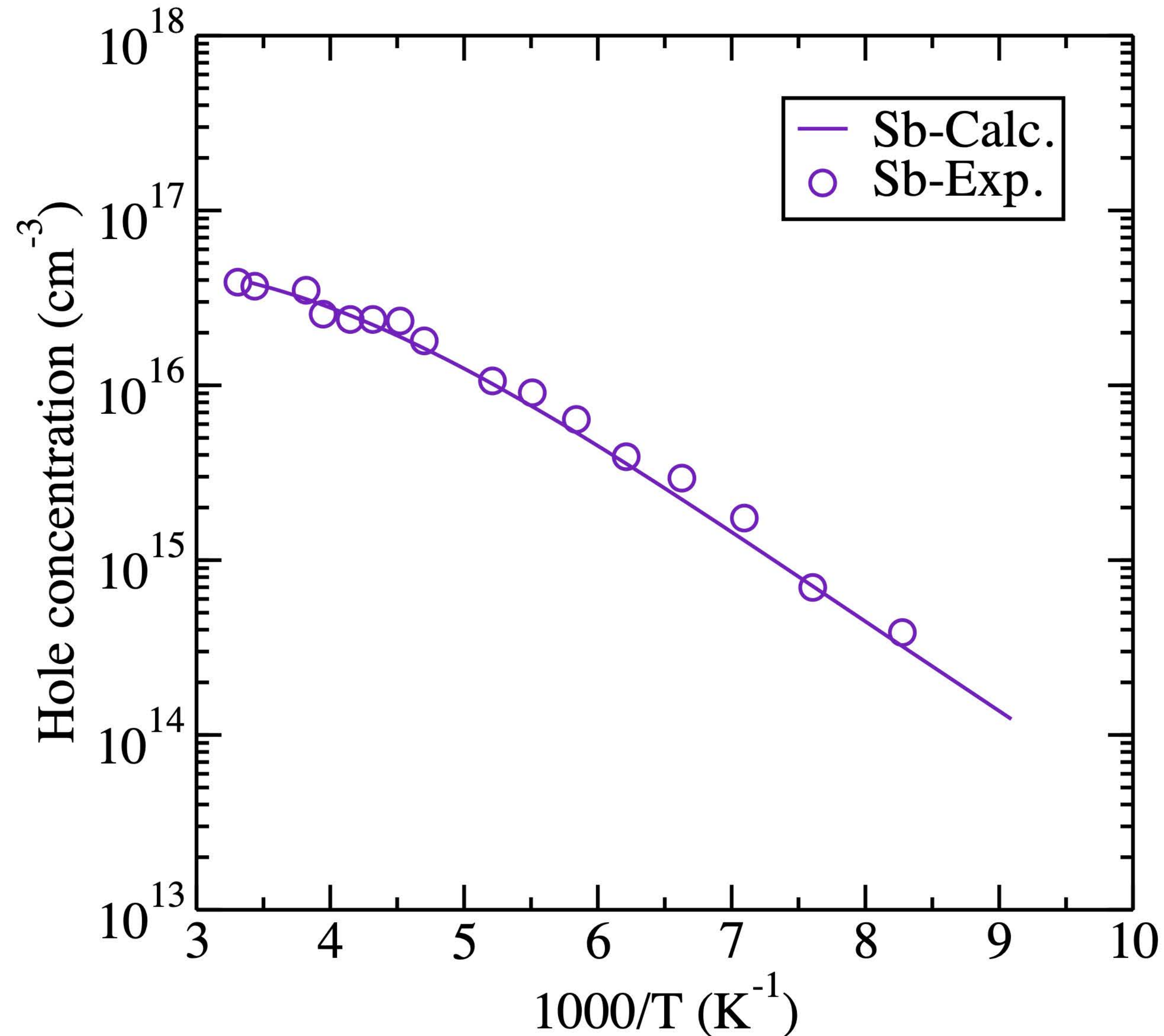
Best fit:  $N_a = [Sb] = 0.6 \times 10^{17} \text{ cm}^{-3}$

Compared to  $1.1 \times 10^{17} \text{ cm}^{-3}$  from exp.

Exp. data from  
A. Nagaoka *et al.*, Appl. Phys. Lett. **116**, 132102 (2020)



# Temperature-dependent hole concentration for Sb doping of CdTe

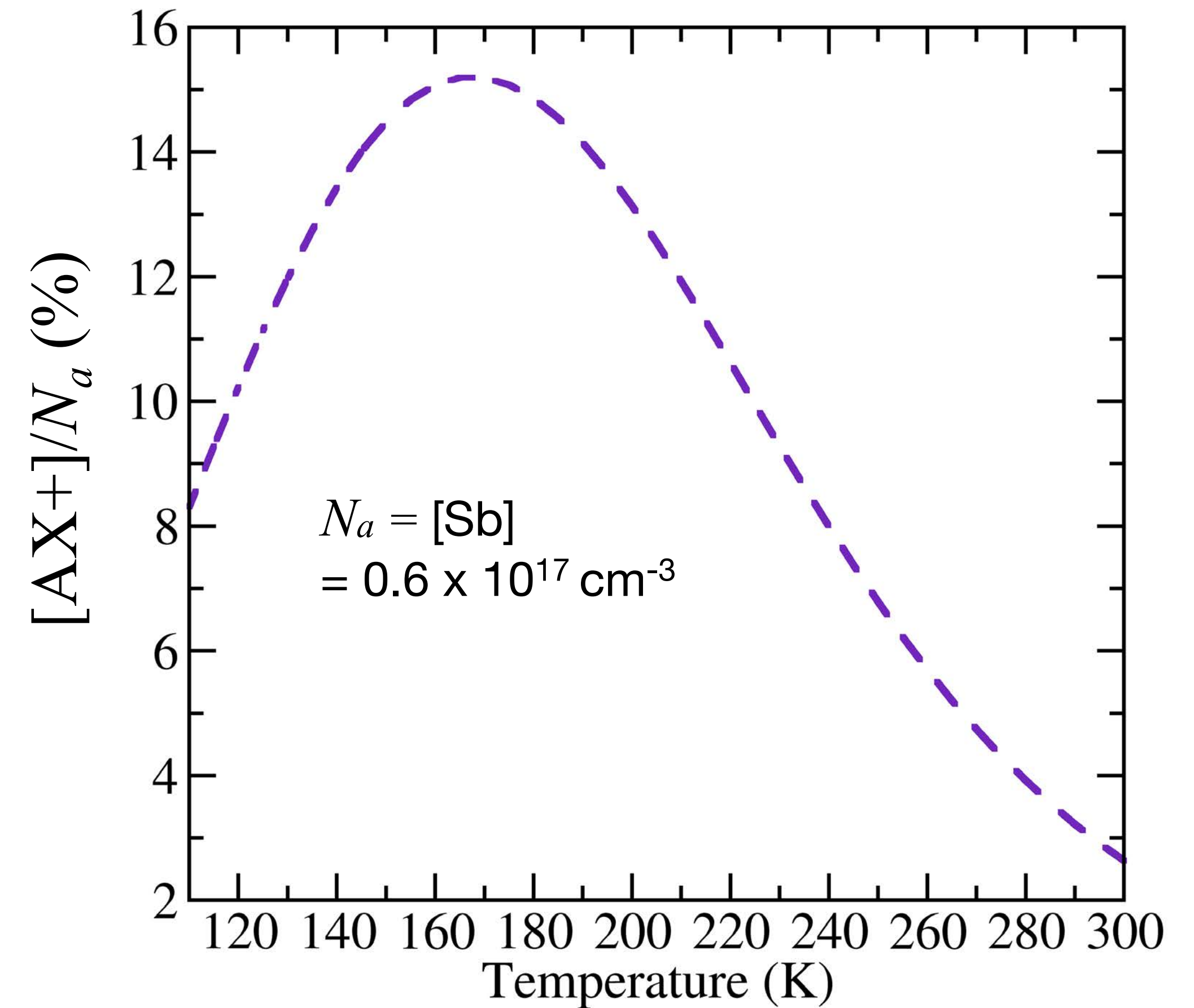
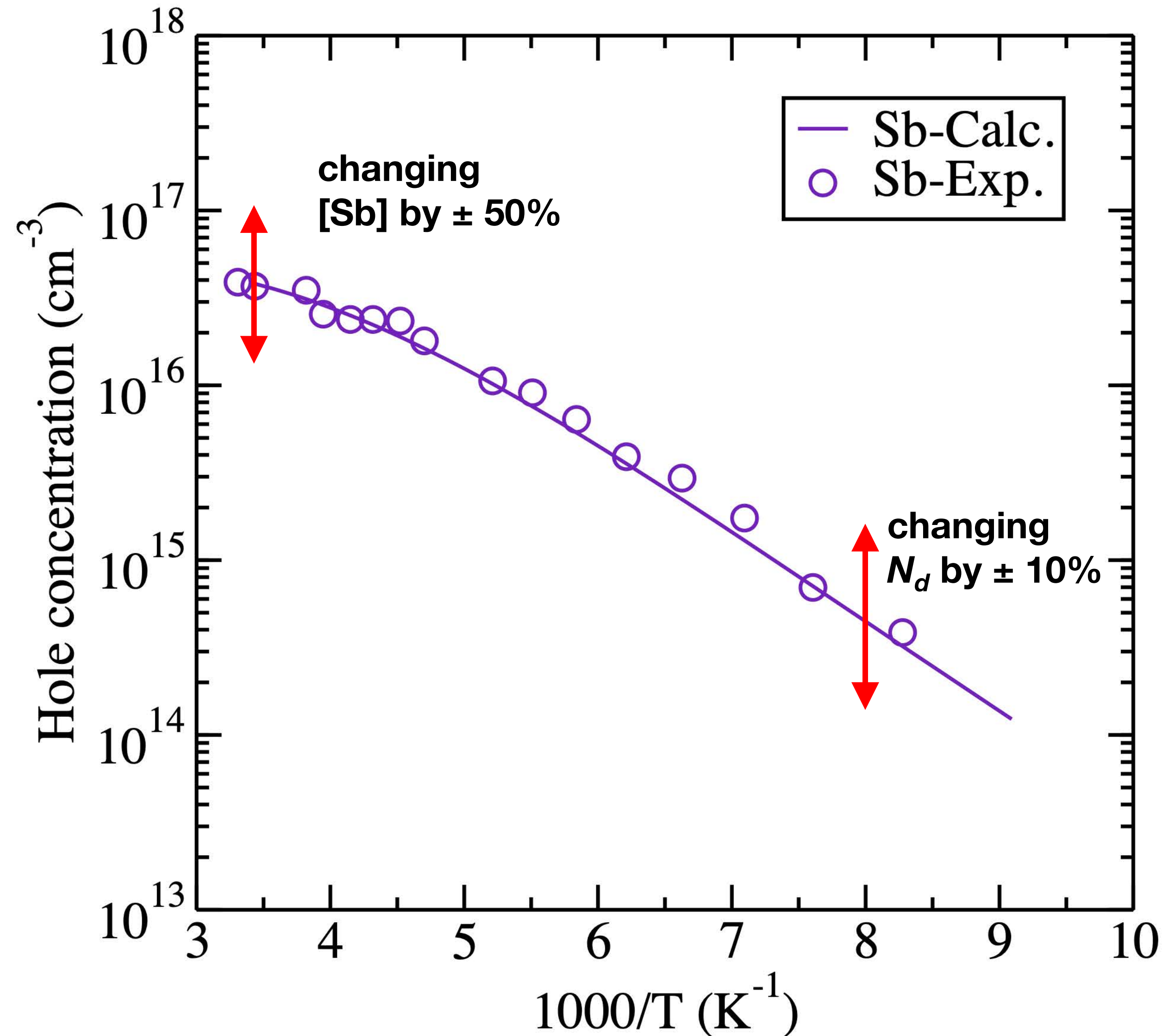


- [AX] < 16% of total Sb concentration
- Becomes less important at room temp.

Exp. data from  
A. Nagaoka *et al.*, Appl. Phys. Lett. **116**, 132102 (2020)



# Temperature-dependent hole concentration for Sb doping of CdTe

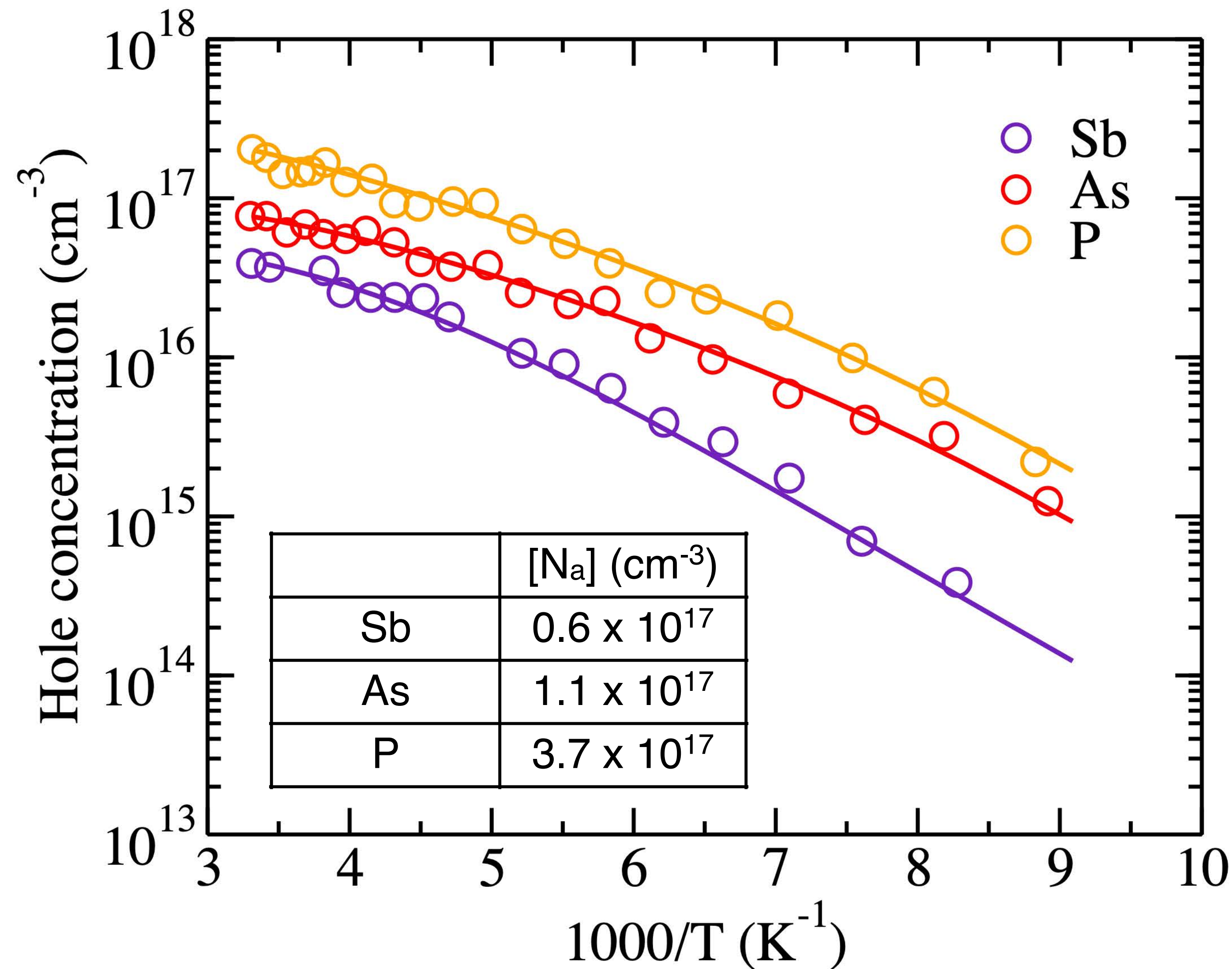


- $[AX] < 16\%$  of total Sb concentration
- Becomes less important at room temp.

Exp. data from  
A. Nagaoka *et al.*, Appl. Phys. Lett. **116**, 132102 (2020)



# Temp.-dependent hole concentration for Sb/As/P doping of CdTe



**Sb** →  $E_a = 116 \text{ meV}$

**As** →  $E_a = 99 \text{ meV}$   
**Compensation = 6%**

**P** →  $E_a = 93 \text{ meV}$   
**Compensation = 6%**

Have to assume ~6% compensation (not AX)  
in the case of As and P doping

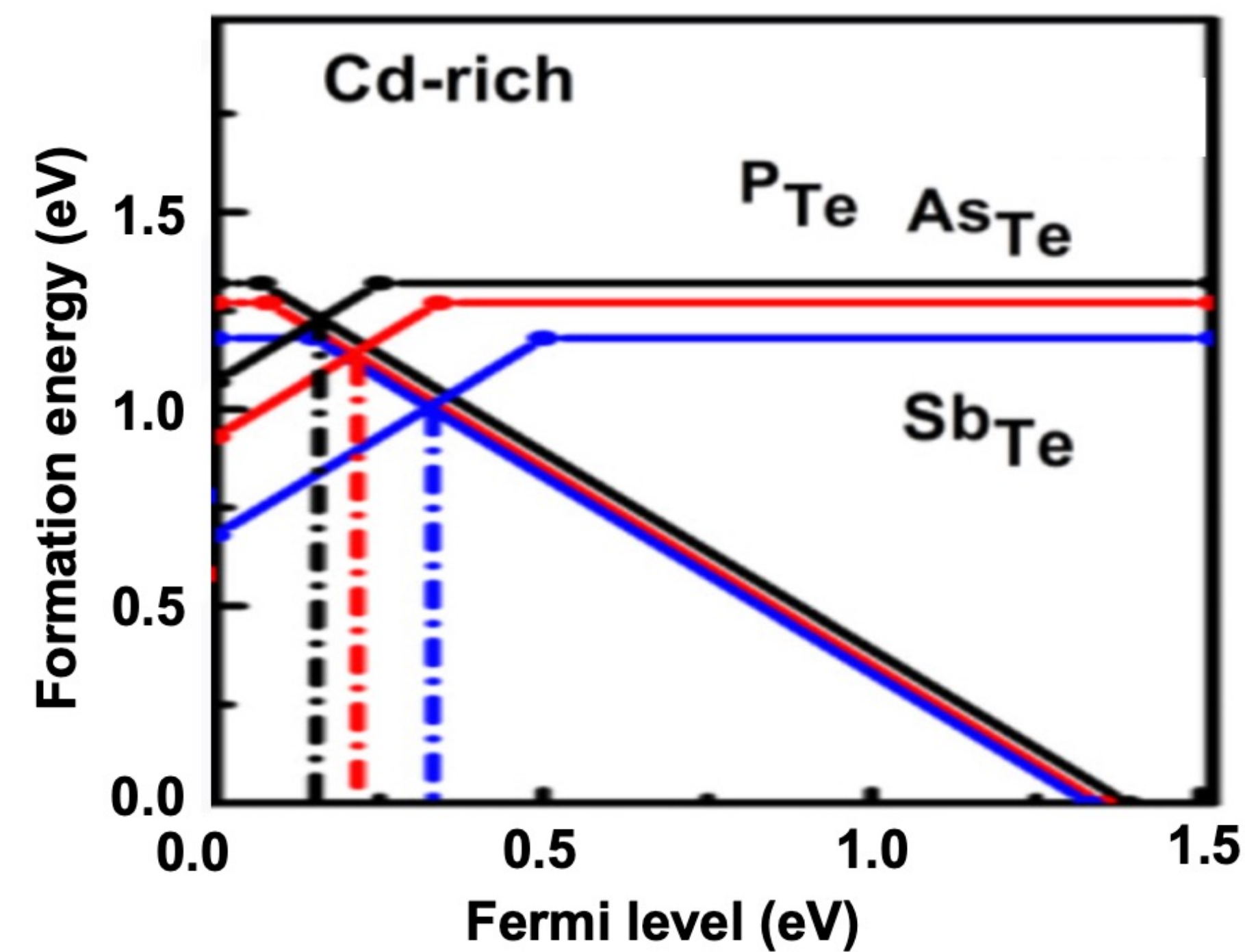
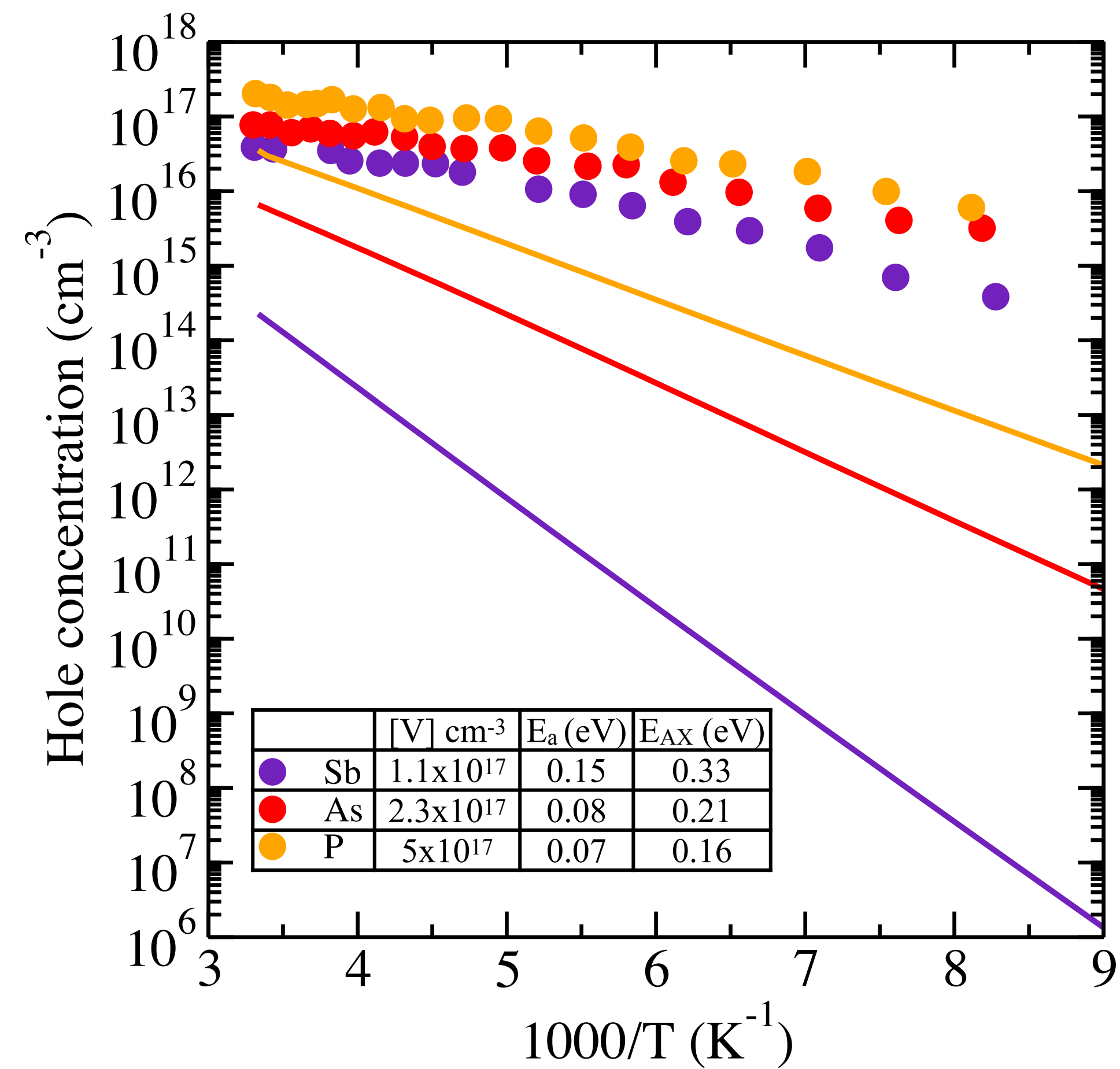
Similar compensation mechanisms for As  
and P? Point defects?

$[h]$  from Sb doping drops faster at low  
temperatures due to higher ionization  
energy



# Temp.-dependent hole concentration for Sb/As/P doping of CdTe

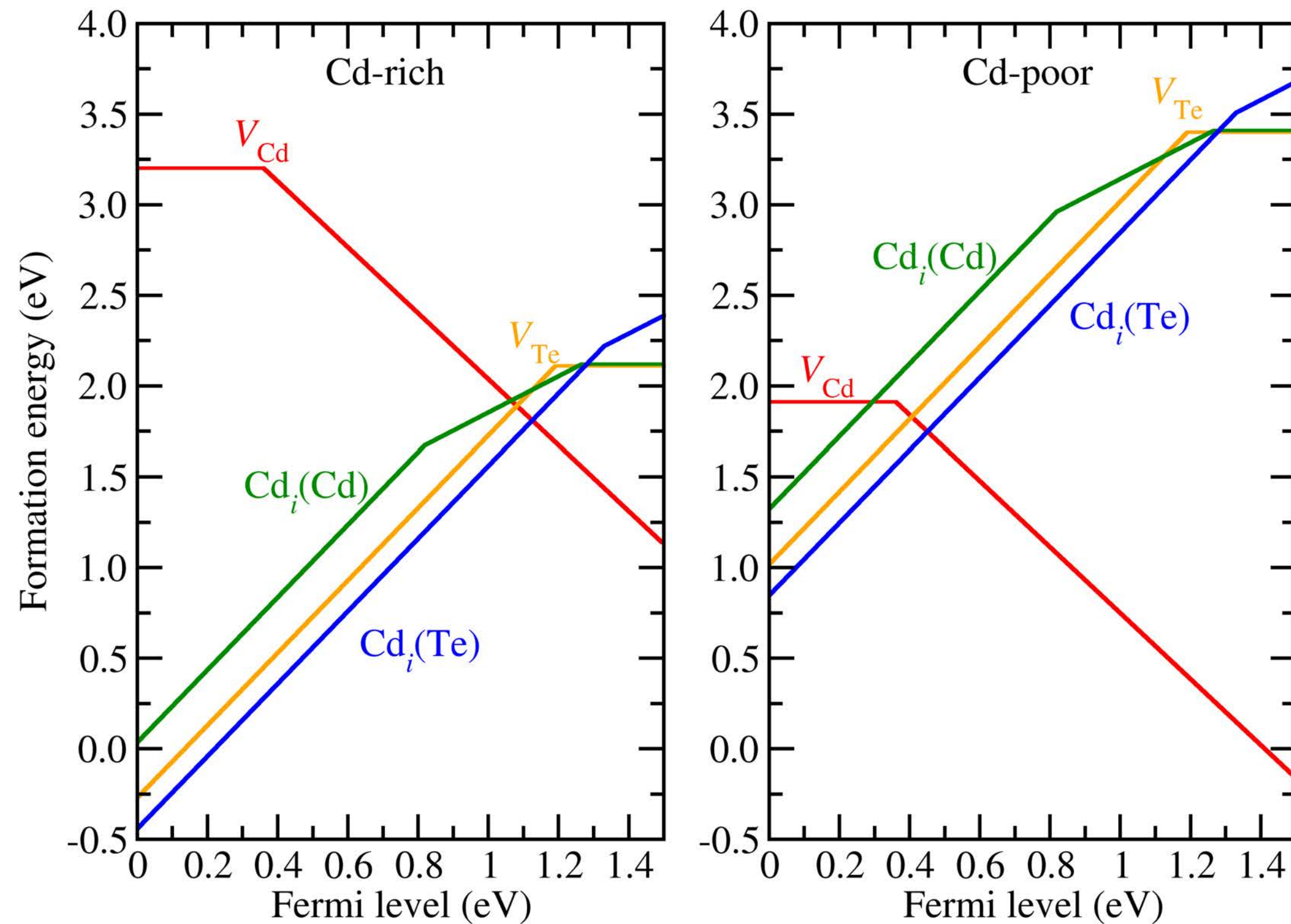
- According to previous theoretical prediction where AX centers are stable



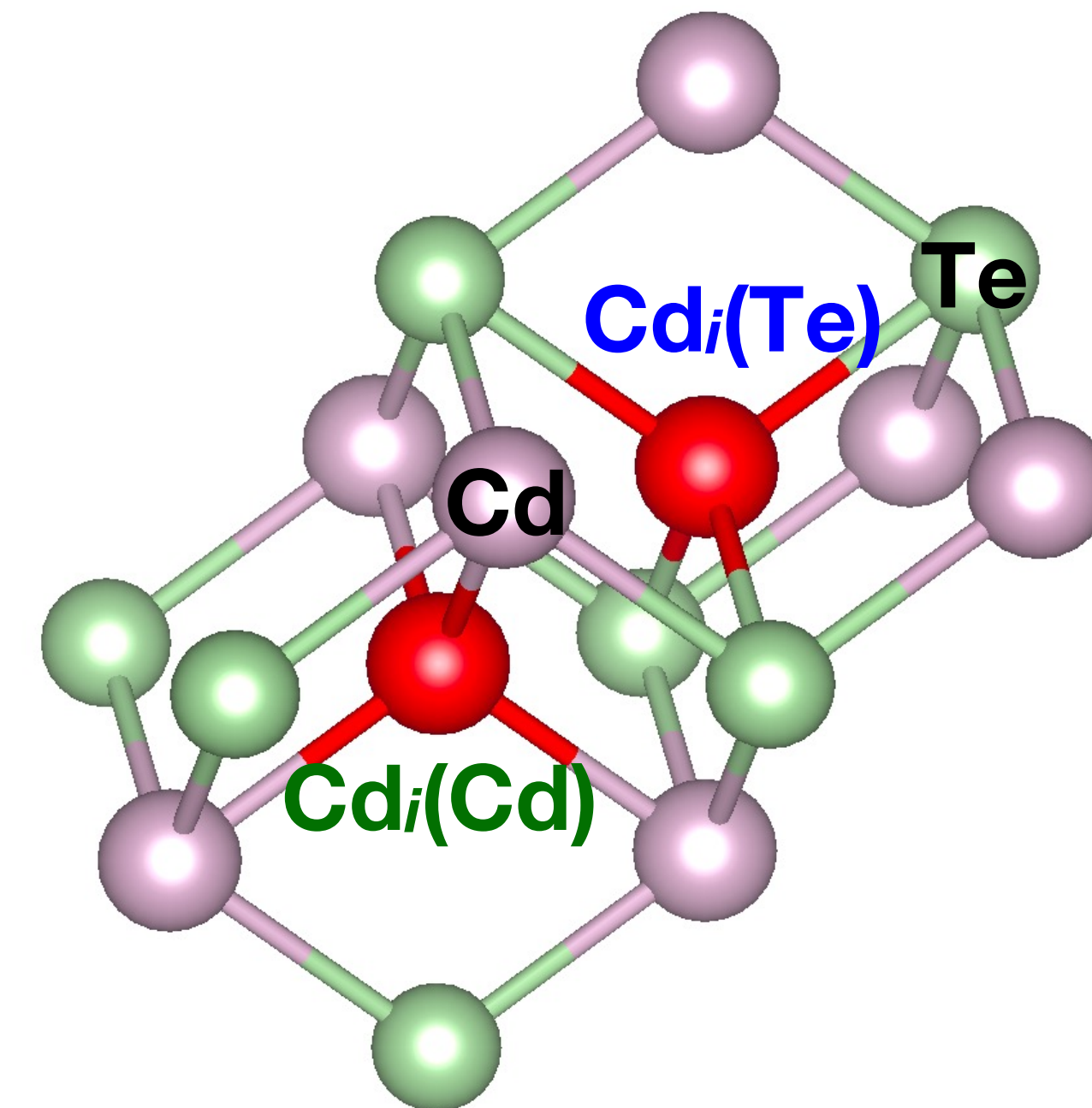
B. Dou, Q. Sun, and S.-H. Wei,  
Phys. Rev. Appl. **15**, 054045 (2021)

Exp. data from  
A. Nagaoka *et al.*, Appl. Phys. Lett. **116**, 132102 (2020)

# Possible intrinsic defects that act as compensation centers in p-type CdTe



- $Cd_i$  and  $V_{Te}$  are the lowest energy defects in p-type and in Cd-rich (Te-poor) limiting conditions



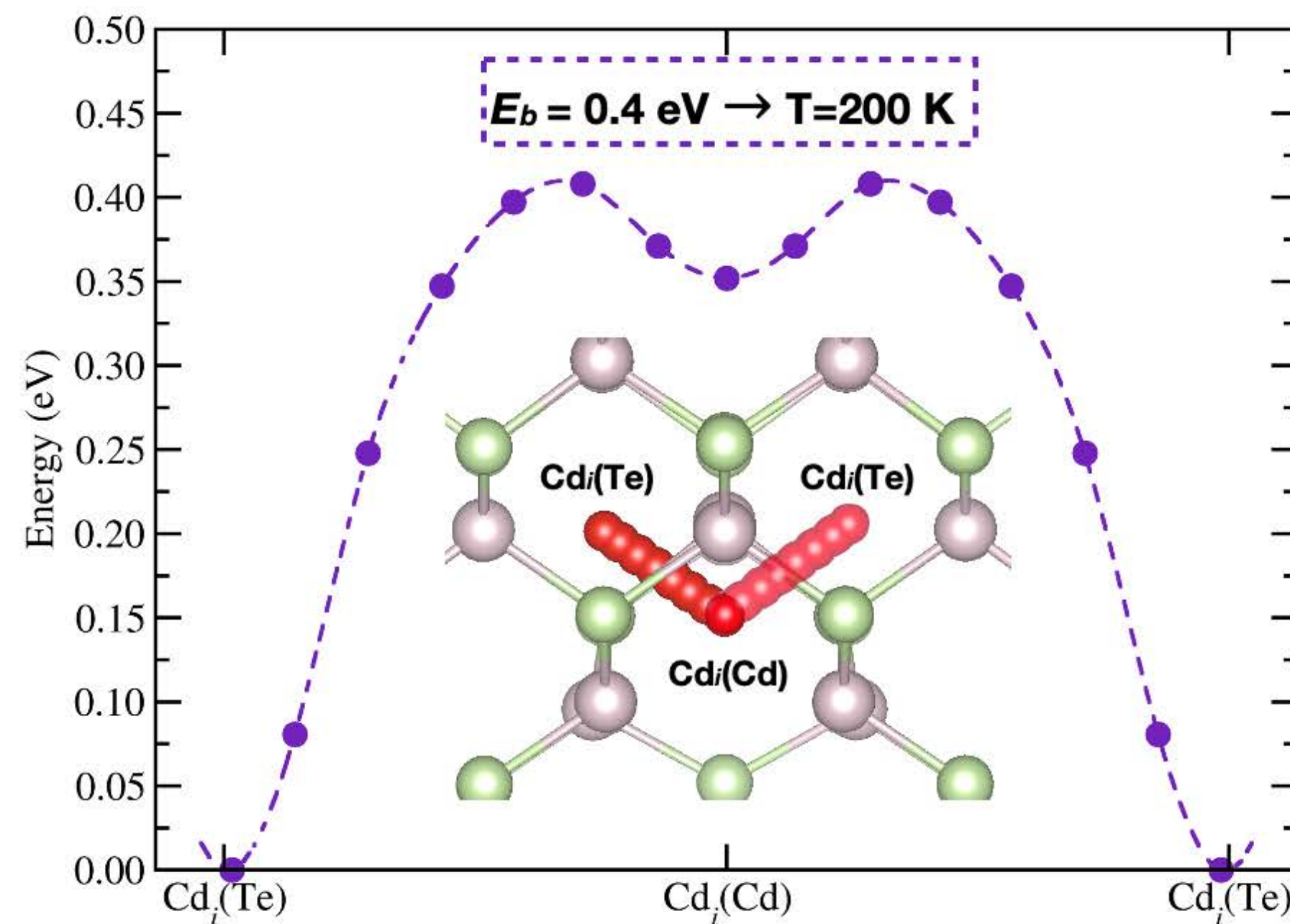
Similar to J. Pan, W. K. Metzger, and S. Lany,  
Phys. Rev. B **98**, 054108 (2018)



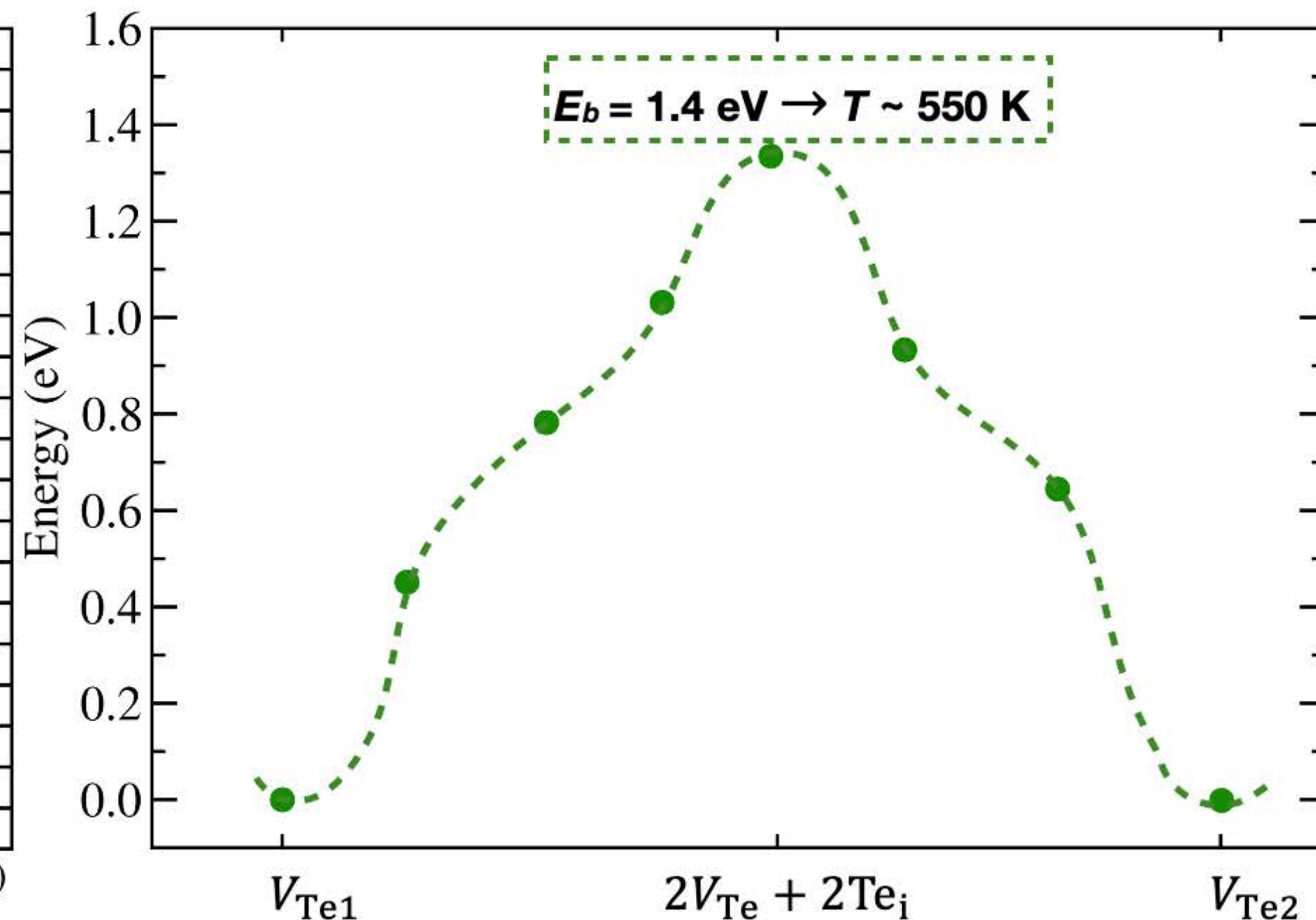
# Calculated migration barriers of $\text{Cd}_i$ and $V_{\text{Te}}$

- $\text{Cd}_i$  is unstable, with very low migration barrier, will move even at room temp.

$$\Gamma = \Gamma_0 e^{-E_b/k_B T}, \text{ assuming } \Gamma_0 \sim 5 \text{ THz}$$



Cd interstitial will be highly mobile at room temperature, making it unstable; will either move out or combine with other defects



Te vacancy more likely to be a compensating donor, stable up to 550 K

# Summary

- SOC is crucial to describe the band structure of CdTe and the properties of group-V acceptors in CdTe; explains the difference between present work and previous calc.
- P, As, and Sb are shallow acceptors in CdTe, with the ionization energies  $\sim 100$  meV, in agreement with recent Hall measurements in bulk crystals
- AX center is not the dominant compensation center in *p*-type CdTe; unstable in the case of As and P, barely stable in the case of Sb doping
- Intrinsic defects, such as  $V_{\text{Te}}$ , are potentially important compensation centers
- $\text{Cd}_i$  is unstable with low migration barrier

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