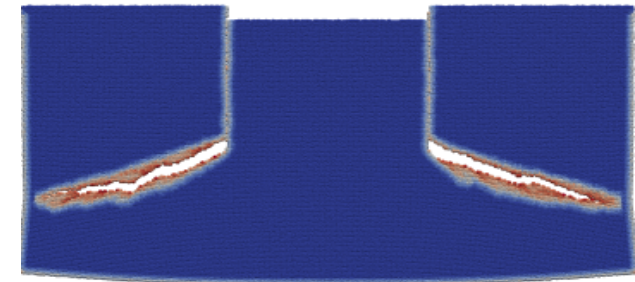
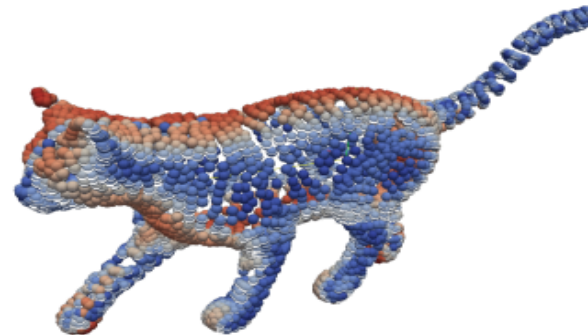
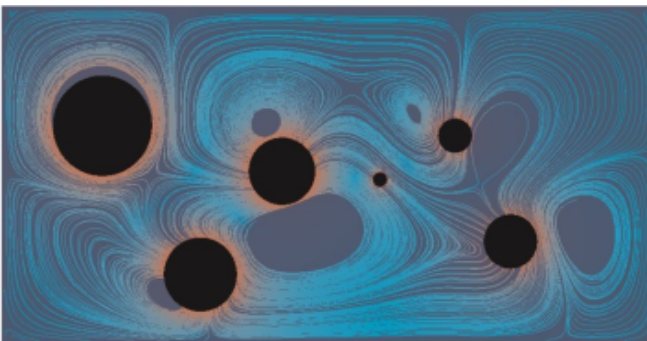


Exceptional service in the national interest



A data-driven exterior calculus for model discovery



Nat Trask
Center for Computing Research
Sandia National Laboratories



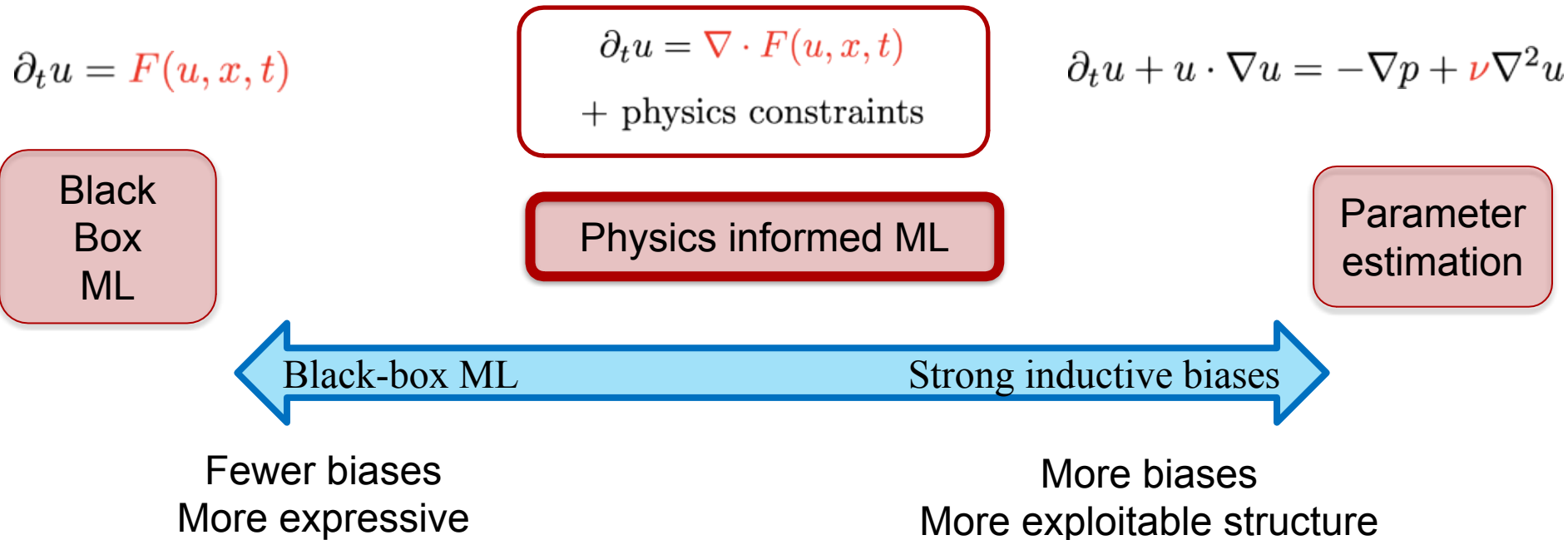
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Physical biases in data-driven modeling

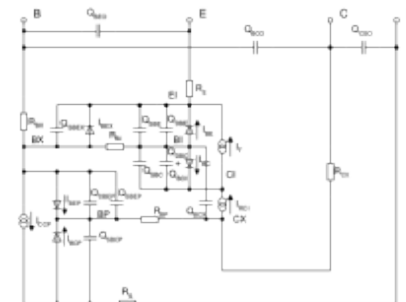
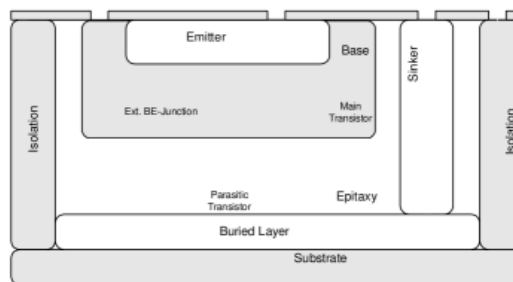
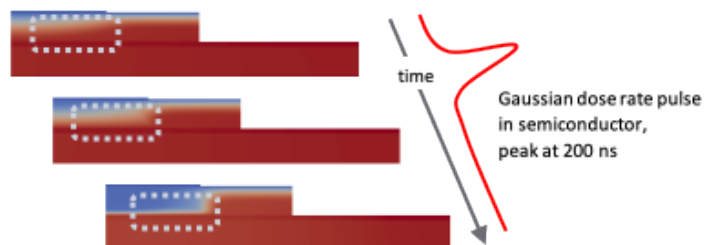
We extract models from data where first principles derivation is intractable, **while guaranteeing well-posed models in small data limits**

Ex:

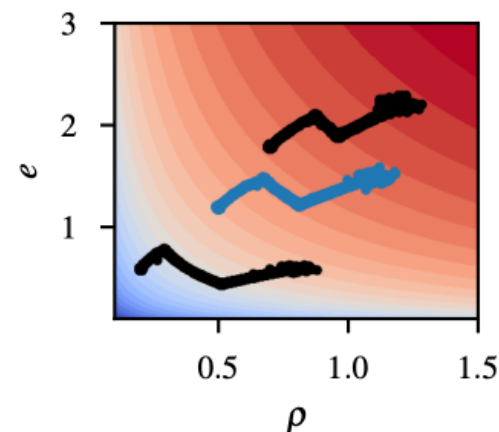
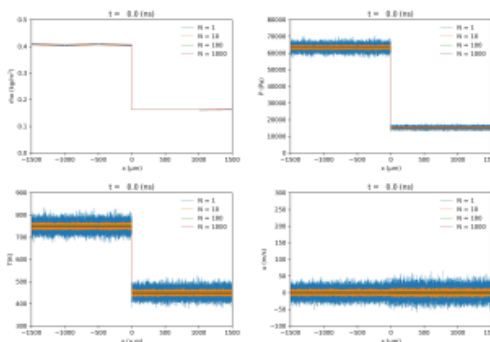
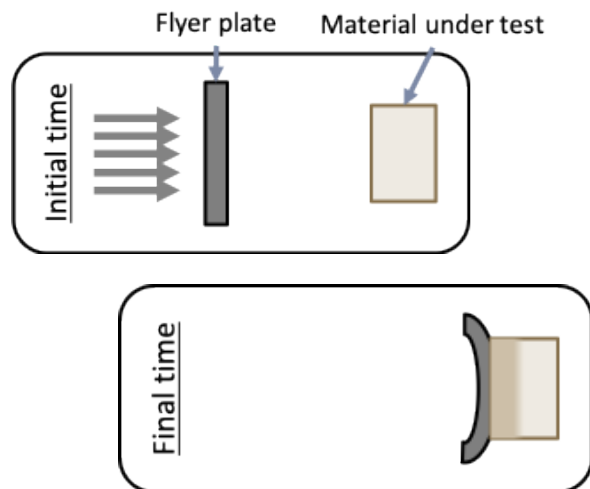
- Turbulence models
- Multiscale closures
- Equations of state
- Noneq. chemistry/kinetics



Data-driven modeling at SNL



Can high-fidelity E&M simulations for semiconductors be encoded as efficient circuit models?



Can a data-driven EOS be extracted from observations of Riemann problems at high energy states

Exact physics treatment and solvability guarantees are critical

The state-of-the-art in physics-informed ML

Make list of desired features and penalize them after the fact:
PDE structure, BC, IC, conservation, etc.

$$\mathbf{L} = \mathbf{L}_{data} + \epsilon \mathbf{L}_{physics}$$

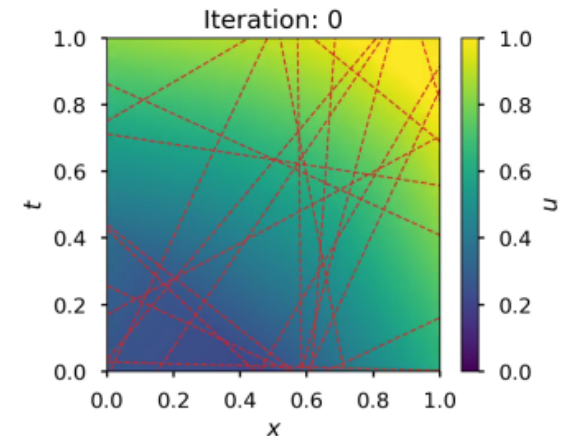
$$\mathbf{L} = ||u_{data} - \mathcal{NN}||_{\ell_2}^2 + \epsilon ||\mathcal{L}[u_{data}] - \mathcal{L}[\mathcal{NN}]||_{\ell_2}^2$$

Key Challenges:

Our applications need physics to hold exactly – not just by penalty
e.g. electromagnetics, fluid mechanics

What happens when the governing PDE and material parameters are unknown – is it still possible to impose physical constraints?

If physics can be imposed *exactly* then a source of uncertainty is eliminated



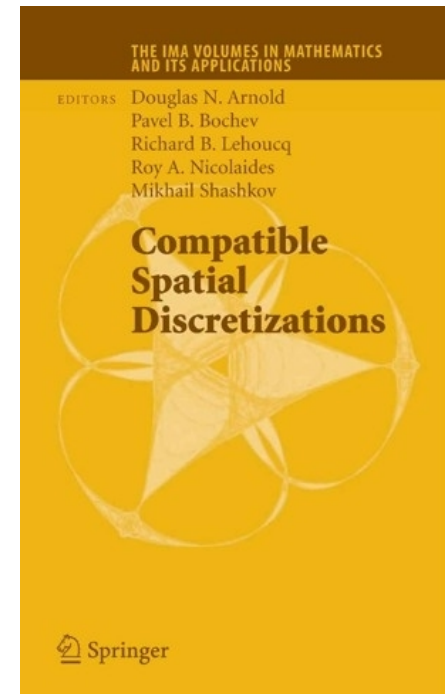
What are physics compatible discretizations for PDEs?

Methods for solving PDEs which:

Use generalized Stokes theorems to approximate differential operators

Preserve topological structure in governing equations

Mimic properties of continuum operators
(thus sometimes called **mimetic discretizations**)



Arnold, D. N., Bochev, P. B.,
Lehoucq, R. B., Nicolaides, R. A.,
& Shashkov, M. (Eds.). (2007).
Compatible spatial discretizations
(Vol. 142). Springer Science &
Business Media.

Two key ingredients:

1: A topological structure

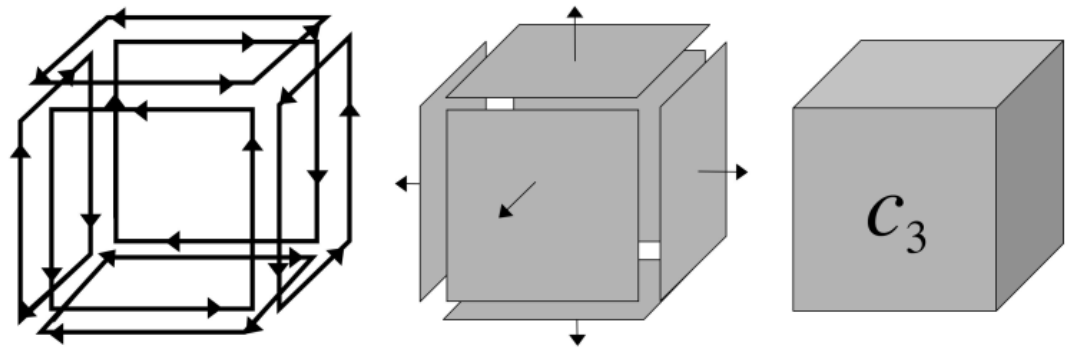
In PDE discretization this is a mesh, with boundary operators linking cells, faces, edges, and nodes

We will use a graph as an inexpensive low-dimensional mesh surrogate

2: Metric information

Measures associated with mesh entities, ensuring discrete exterior derivatives converge to div/grad/curl

Graphs are purely topological with no natural metric, we will use ML to extract metric information from data

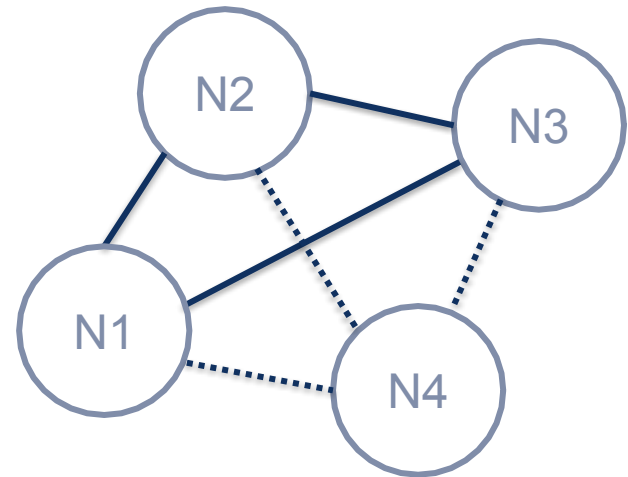
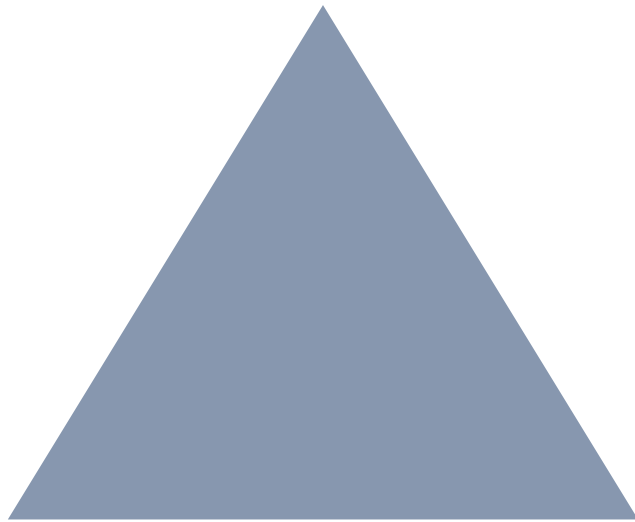


$$0 \leftarrow \partial \partial c_3 \xleftarrow{\partial} \partial c_3 \xleftarrow{\partial} c_3$$

$$\nabla \cdot \mathbf{u} = \frac{1}{\mu(C)} \sum_{f \in \partial C} \int_f \mathbf{u} \cdot d\mathbf{A}$$

Exterior calculus preliminaries: chain complex

$$C_0 \xleftarrow{\partial_0} C_1 \xleftarrow{\partial_1} C_2 \xleftarrow{\partial_2} C_3$$



Compat. PDE

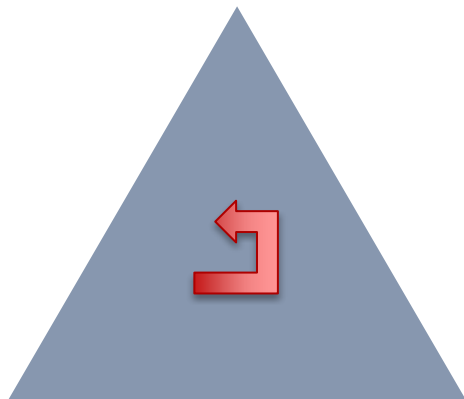
Mesh entities

Comb. Hodge

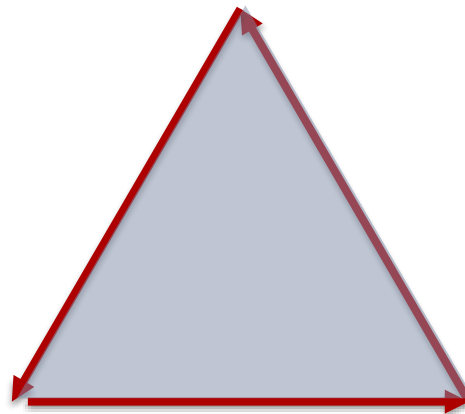
K-cliques

Exterior calculus preliminaries: chain complex

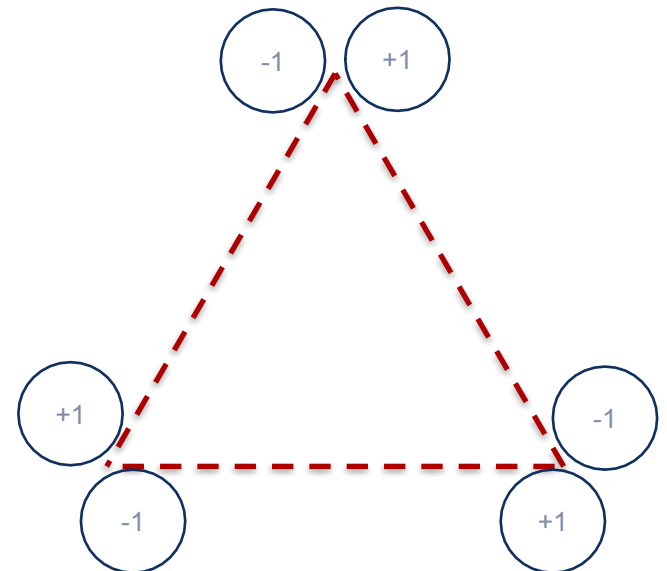
$$C_0 \xleftarrow{\partial_0} C_1 \xleftarrow{\partial_1} C_2 \xleftarrow{\partial_2} C_3$$



$$f \in C_2$$



$$\partial_2 f \in C_1$$



$$\partial_1 \partial_2 f \in C_0$$

Exact sequence property: $\forall k, \partial_k \partial_{k+1} = 0$

Exterior calculus preliminaries: cochain complex

$$C_0 \xleftarrow{\partial_0} C_1 \xleftarrow{\partial_1} C_2 \xleftarrow{\partial_2} C_3$$

$$C^0 \xrightarrow{d_0} C^1 \xrightarrow{d_1} C^2 \xrightarrow{d_2} C^3$$

Coboundary operators define maps $d_k : C^k \rightarrow C^{k+1}$ satisfying $d_{k+1}d_k = 0$

Boundary and coboundary operators satisfy the *generalized Stokes theorem*

$$\int_{\omega} du = \int_{\partial\omega} u$$

Comb. Hodge	Compat. PDE
$\text{grad}[s](i, j) = \int_{e_{ij}} \nabla s \cdot d\mathbf{l} = s_j - s_i$ $\text{curl}[X](F) = \int_F \nabla \times X \cdot d\mathbf{A} = \sum_{e \in \partial F} \int_e X \cdot d\mathbf{l}$	$\text{grad}[s](i, j) = s_j - s_i$ $\text{curl}[X](i, j, k) = X_{ij} + X_{jk} + X_{ki}$

$$C_0 \xleftarrow{\partial_1} C_1 \xleftarrow{\partial_2} C_2 \xleftarrow{\partial_3} C_3$$

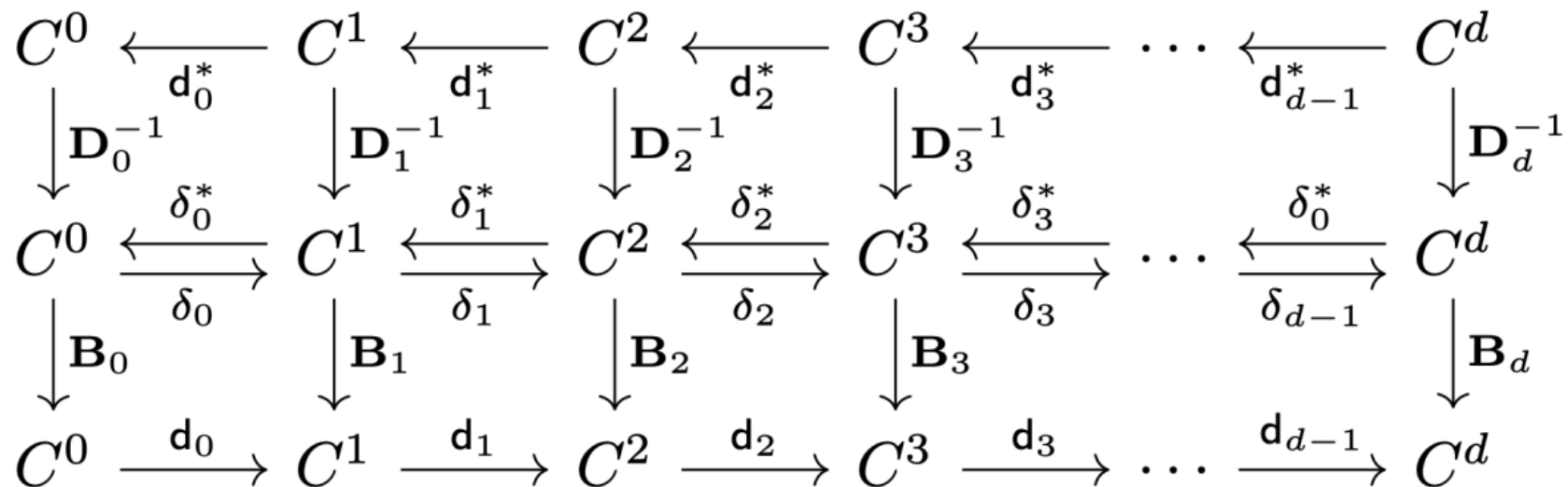
$$C^0 \begin{matrix} \xrightarrow{d_0} \\ \xleftarrow{d_0^*} \end{matrix} C^1 \begin{matrix} \xrightarrow{d_1} \\ \xleftarrow{d_1^*} \end{matrix} C^2 \begin{matrix} \xrightarrow{d_2} \\ \xleftarrow{d_2^*} \end{matrix} C^3$$

Introducing inner products $(\cdot, \cdot)_k$, we define the *codifferential* operator $d_k^* : C^{k+1} \rightarrow C^k$ as

$$(v, d_k^* u)_k = (d_k v, u)_{k+1}$$

$$\text{Again, } d_{k-1}^* * d_k^* = 0$$

A data-driven exterior calculus (DDEC)



Idea: Take graph calculus and introduce learnable inner products

$$(x, y)_{\mathbf{B}_k} = x^\top \mathbf{B}_k y$$

$$(x, y)_{\mathbf{D}_k} = x^\top \mathbf{D}_k y$$

to find data-driven exterior calculus operators that inherit the structure of graph exterior calculus

What does all this give you?

$$\begin{array}{ccccccc}
 C^0 & \xleftarrow{d_0^*} & C^1 & \xleftarrow{d_1^*} & C^2 & \xleftarrow{d_2^*} & C^3 & \xleftarrow{d_3^*} & \dots & \xleftarrow{d_{d-1}^*} & C^d \\
 \downarrow \mathbf{D}_0^{-1} & & \downarrow \mathbf{D}_1^{-1} & & \downarrow \mathbf{D}_2^{-1} & & \downarrow \mathbf{D}_3^{-1} & & & & \downarrow \mathbf{D}_d^{-1} \\
 C^0 & \xrightleftharpoons[\delta_0]{\delta_0^*} & C^1 & \xrightleftharpoons[\delta_1]{\delta_1^*} & C^2 & \xrightleftharpoons[\delta_2]{\delta_2^*} & C^3 & \xrightleftharpoons[\delta_3]{\delta_3^*} & \dots & \xrightleftharpoons[\delta_{d-1}]{\delta_{d-1}^*} & C^d \\
 \downarrow \mathbf{B}_0 & & \downarrow \mathbf{B}_1 & & \downarrow \mathbf{B}_2 & & \downarrow \mathbf{B}_3 & & & & \downarrow \mathbf{B}_d \\
 C^0 & \xrightarrow{d_0} & C^1 & \xrightarrow{d_1} & C^2 & \xrightarrow{d_2} & C^3 & \xrightarrow{d_3} & \dots & \xrightarrow{d_{d-1}} & C^d
 \end{array}$$

- Differential operators which locally and globally conserve fluxes, circulations, potentials
- Invertible Hodge Laplacians $\Delta_k = d_{k+1}^* d_{k+1} + d_k d_{k+1}^*$
- Exact sequence properties $d_{k+1} d_k = d_k^* d_{k+1}^* = 0$
- Hodge decomposition $u = d^* \alpha + d \beta + \gamma$
- Corollary: treatment of nontrivial null-spaces in electromagnetism

Theorem 3.1. The discrete derivatives \mathbf{d}_k in (11) form an exact sequence if the simplicial complex is exact, and in particular $\mathbf{d}_{k+1} \circ \mathbf{d}_k = 0$. In \mathbb{R}^3 , we have $CURL_h \circ GRAD_h = DIV_h \circ CURL_h = 0$.

Theorem 3.2. The discrete derivatives \mathbf{d}_k^* in (11) form an exact sequence of the simplicial complex is exact, and in particular $\mathbf{d}_k^* \circ \mathbf{d}_{k+1}^* = 0$. In \mathbb{R}^3 , $DIV_h^* \circ CURL_h^* = CURL_h^* \circ GRAD_h^* = 0$.

Theorem 3.3 (Hodge Decomposition). For C^k , the following decomposition holds

$$C^k = \text{im}(\mathbf{d}_{k-1}) \bigoplus_k \ker(\Delta_k) \bigoplus_k \text{im}(\mathbf{d}_k^*), \quad (17)$$

where \bigoplus_k means the orthogonality with respect to the $(\cdot, \cdot)_{\mathbf{D}_k \mathbf{B}_k^{-1}}$ -inner product.

Theorem 3.4 (Poincaré inequality). For each k , there exists a constant $c_{P,k}$ such that

$$\|\mathbf{z}_k\|_{\mathbf{D}_k \mathbf{B}_k^{-1}} \leq c_{P,k} \|\mathbf{d}_k \mathbf{z}_k\|_{\mathbf{D}_{k+1} \mathbf{B}_{k+1}^{-1}}, \quad \mathbf{z}_k \in \text{im}(\mathbf{d}_k^*),$$

and another constant $c_{P,k}^*$ such that

$$\|\mathbf{z}_k\|_{\mathbf{D}_k \mathbf{B}_k^{-1}} \leq c_{P,k}^* \|\mathbf{d}_{k-1}^* \mathbf{z}_k\|_{\mathbf{D}_{k-1} \mathbf{B}_{k-1}^{-1}}, \quad \mathbf{z}_k \in \text{im}(\mathbf{d}_{k-1}).$$

Thus, for $\mathbf{u}_k \in C^k$, we have

$$\inf_{\mathbf{h}_k \in \ker(\Delta_k)} \|\mathbf{u}_k - \mathbf{h}_k\|_{\mathbf{D}_k \mathbf{B}_k^{-1}} \leq C \left(\|\mathbf{d}_k \mathbf{u}_k\|_{\mathbf{D}_{k+1} \mathbf{B}_{k+1}^{-1}} + \|\mathbf{d}_{k-1}^* \mathbf{u}_k\|_{\mathbf{D}_{k-1} \mathbf{B}_{k-1}^{-1}} \right),$$

where constant $C > 0$ only depends on $c_{P,k}$ and $c_{P,k}^*$.

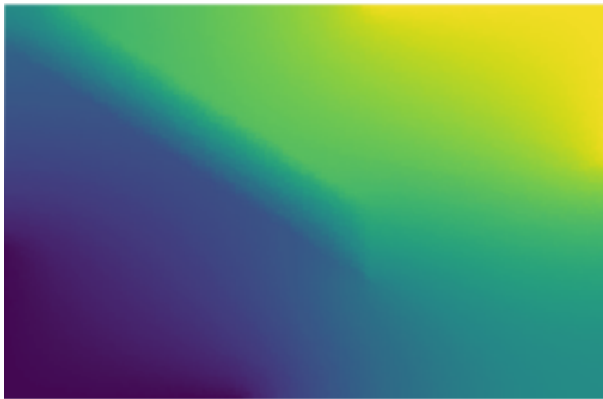
Theorem 3.5 (Invertibility of Hodge Laplacian). The k^{th} -order Hodge Laplacian Δ_k is positive-semidefinite, with the dimension of its null-space equal to the dimension of the corresponding homology $H^k = \ker(\mathbf{d}_k) / \text{im}(\mathbf{d}_{k-1})$.

$$\nabla \cdot \mathbf{F} = f$$

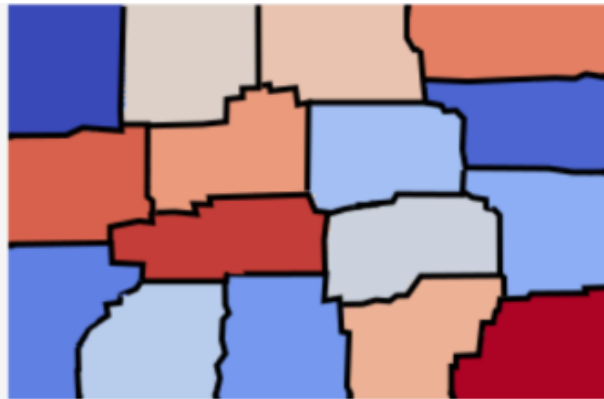
$$d_0^* \mathbf{F} = f$$

$$\mathbf{F} + \kappa \nabla \phi = 0$$

$$\mathbf{F} + \xi d_0 \phi + \mathcal{N}_\eta(\phi) = 0$$



High-fidelity PDE
solution



Apply graph-cut to
coarse-grain
chain complex



Average over
partitions to obtain
training data

General optimization problem

Fluxes: $\mathbf{w}_{k+1} = \mathbf{d}_k \mathbf{u}_k + \epsilon \mathcal{N}\mathcal{N}(\mathbf{d}_k \mathbf{u}_k; \xi),$

Conservation: $\mathbf{d}_{k-1} \mathbf{d}_{k-1}^* \mathbf{u}_k + \mathbf{d}_k^* \mathbf{w}_{k+1} = \mathbf{f}_k.$

➔ $a(\mathbf{v}, \mathbf{u}; \mathbf{B}, \mathbf{D}) + N_{\mathbf{v}}[\mathbf{u}; \xi] = b(\mathbf{v})$

Invertible bilinear
form

Nonlinear
perturbation

If we can fit the model to data while imposing equality constraint, then during training we restrict to manifold of solvable models preserving physics

$$\operatorname{argmin}_{\mathbf{B}, \mathbf{D}, \xi} ||\mathbf{w} - \mathbf{w}_{\text{data}}||^2$$

such that $\mathcal{L}[\mathbf{w}, \mathbf{u}; \mathbf{B}, \mathbf{D}, \xi] = 0$

Is PDE constraint well posed?

$$a(\mathbf{v}, \mathbf{u}; \mathbf{B}, \mathbf{D}) + N_{\mathbf{v}}[\mathbf{u}; \xi] = b(\mathbf{v})$$

Theorem 3.6. *The equation (24) has at least one solution $\mathbf{u}_k \in \mathbb{V}$ satisfies*

$$\|\mathbf{u}_k\| \leq \frac{\|\mathbf{f}\|}{(C_p - C_N)}. \quad (26)$$

Theorem 3.7. *If $\frac{C_{\nabla N} \|\mathbf{f}\|}{C_p(C_p - C_N)} < 1$, then the equation (24) has at most one solution in \mathbb{V} .*

A unique solution exists if the Hodge-Laplacian is sufficiently large relative to the nonlinear part, following standard elliptic PDE arguments

- Poincare constant easily estimated from matrix eigenvalues
- Lipschitz constant on nonlinearity straightforward for DNNs

Solvability constraint could be enforced during training if desired

“PDE”-constrained optimization

$$\mathbf{L}_{\mathbf{u},\lambda,\mathbf{B},\mathbf{D},\xi} = ||\mathbf{w} - \mathbf{w}_{\text{data}}||^2 + \lambda^\top \mathcal{L}[\mathbf{w}, \mathbf{u}; \mathbf{B}, \mathbf{D}, \xi]$$

$$\mathcal{L}[\mathbf{w}, \mathbf{u}; \mathbf{B}, \mathbf{D}, \xi] = 0$$

- Solve forward problem with current model parameters

$$\mathbf{w}, \mathbf{u} \leftarrow \nabla_{\lambda} \mathbf{L}_{\mathbf{u},\lambda,\mathbf{B},\mathbf{D},\xi} = 0$$

- Solve adjoint problem with current forward solution

$$\lambda \leftarrow \nabla_{\mathbf{u}} \mathbf{L}_{\mathbf{u},\lambda,\mathbf{B},\mathbf{D},\xi} = 0$$

- Apply gradient descent to update model

$$\mathbf{B}, \mathbf{D}, \xi \leftarrow \nabla_{\mathbf{B},\mathbf{D},\xi} \mathbf{L}_{\mathbf{u},\lambda,\mathbf{B},\mathbf{D},\xi} = 0$$

An iterative algorithm
guaranteeing exact
enforcement of physics
at each iteration:

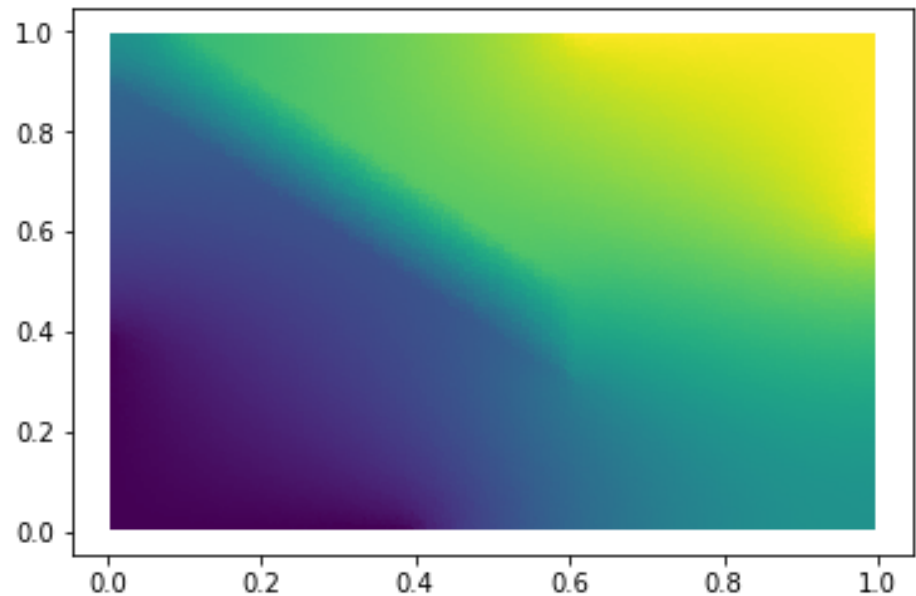
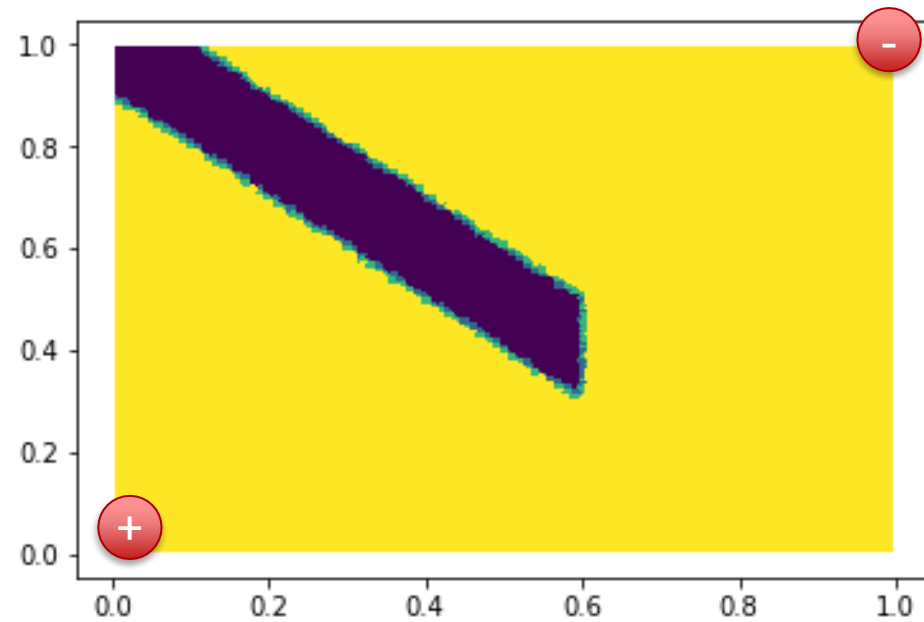
Back to Darcy...

$$\nabla \cdot \mathbf{F} = f$$

$$\mathbf{F} + \kappa \nabla \phi = 0$$

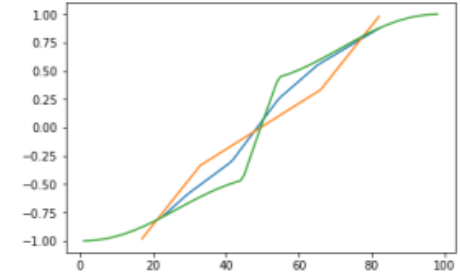
$$d_0^* \mathbf{F} = f$$

$$\mathbf{F} + \xi d_0 \phi + \cancel{\mathcal{N}_\eta(\phi)} = 0$$

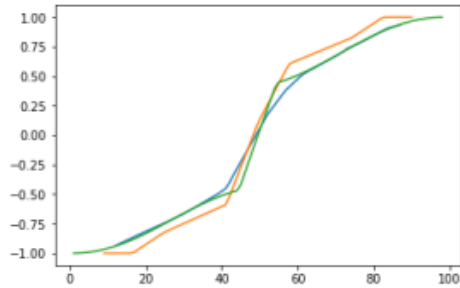
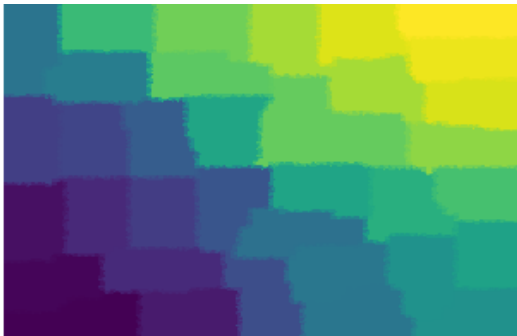


Comparison to traditional covolume: improved accuracy at low resolution

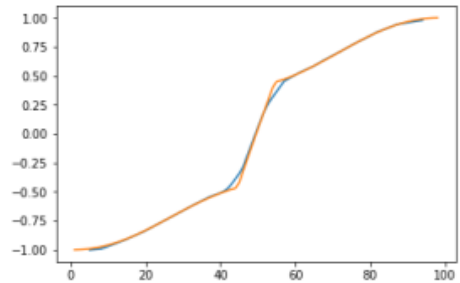
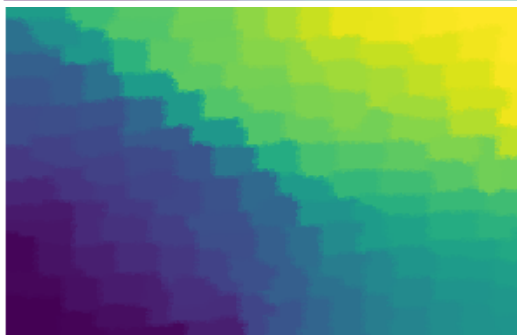
$N = 2^2$



$N = 5^2$



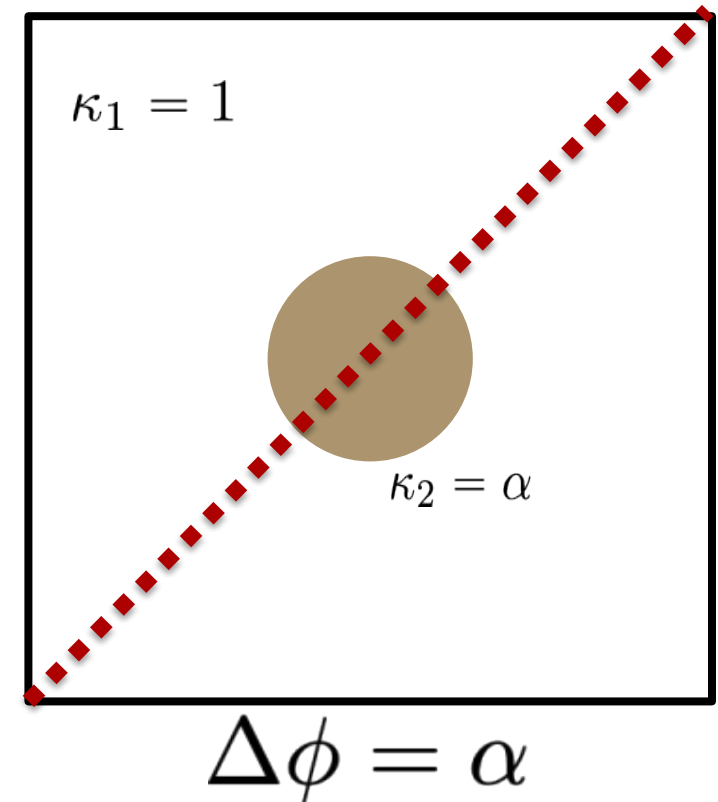
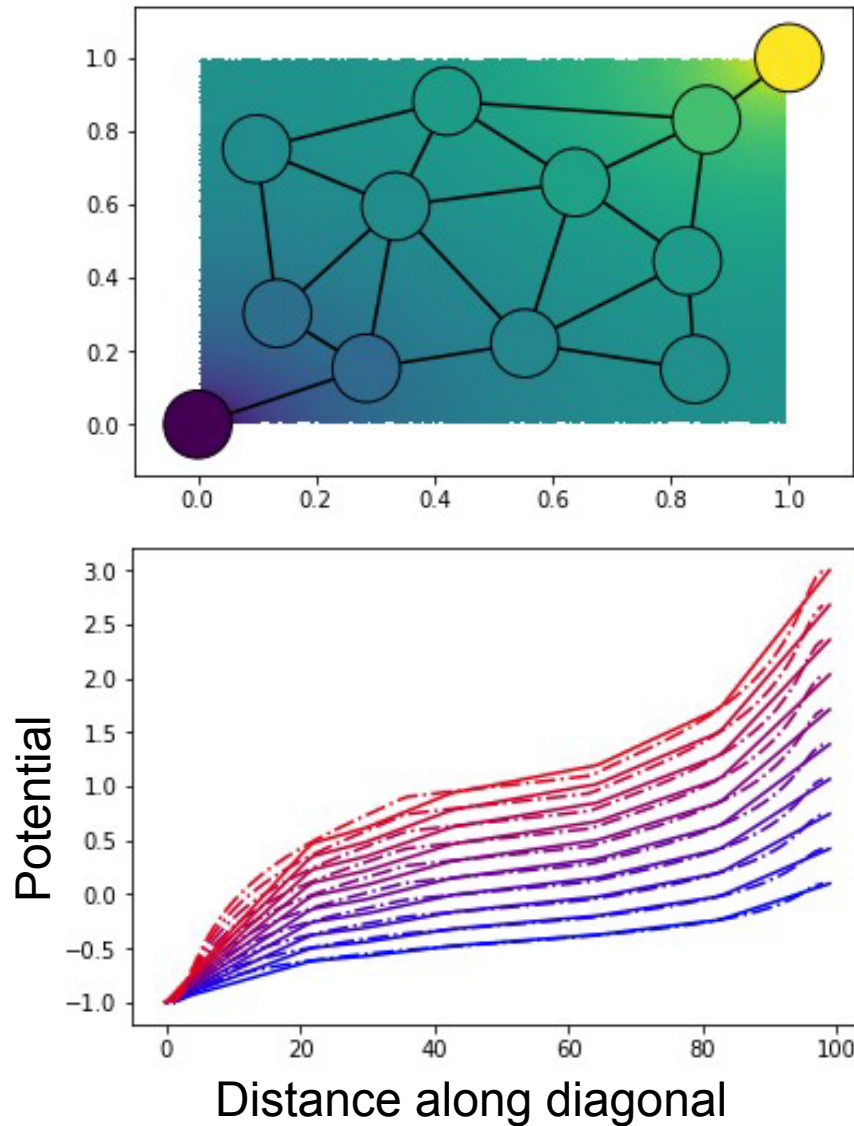
$N = 10^2$



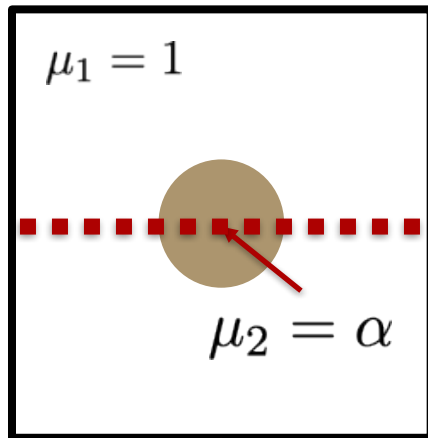
Comparison of pressure for same # DOF for FVM (left) and DDEC (center)

Right: profile along diagonal shows better fit to solution (green) by DDEC (blue) vs FVM (orange)

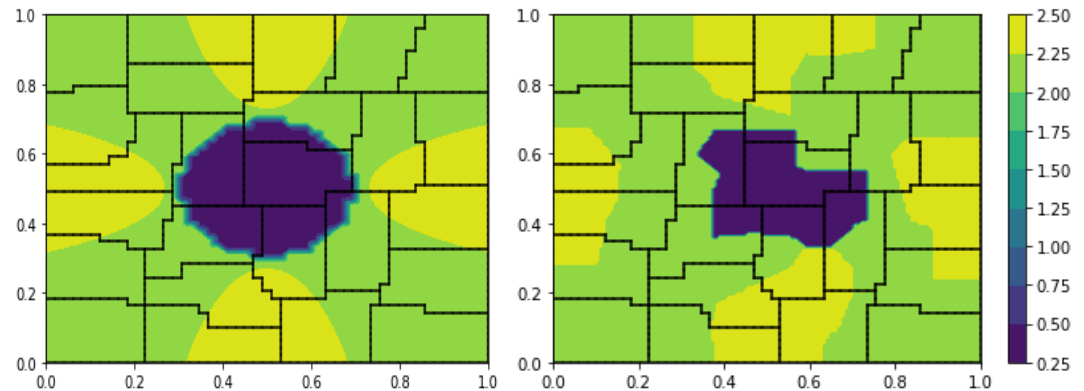
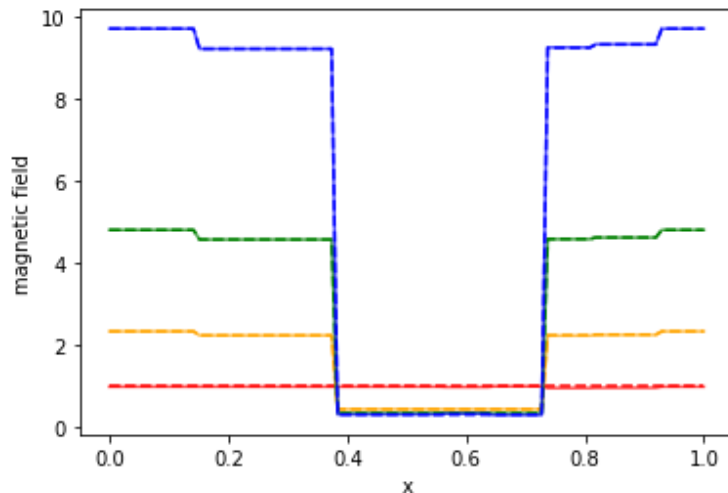
Nonlinear Darcy: potential profile across diagonal



Magnetostatics



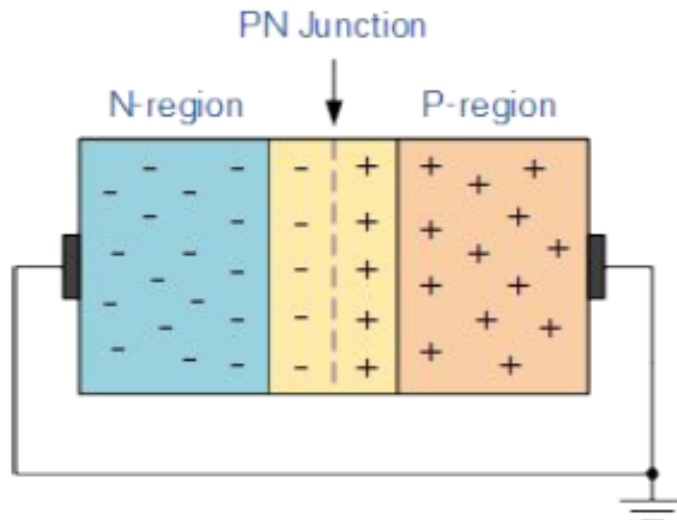
$$B|_{\partial\Omega} = \alpha$$



Extracted surrogate:
 Is exactly div free
 Provides sharp interfaces
 Conserves circulation
 Guaranteed solvable
 Generalizes to other BCs

Compact models for semiconductors: PN-diode

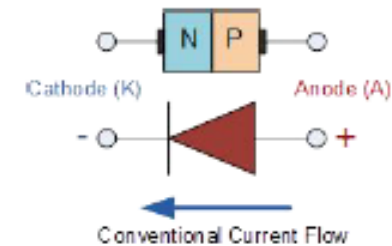
https://www.electronics-tutorials.ws/diode/diode_3.html



$$\nabla \cdot \epsilon \nabla \phi = -(p - n + N_D^+ - N_A^-)$$

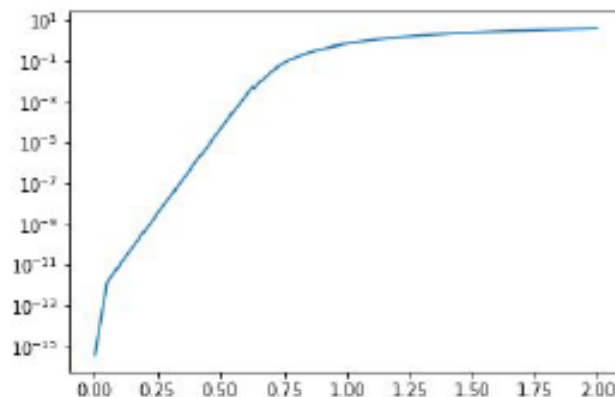
$$\frac{\partial n}{\partial t} = \frac{1}{q} \nabla \cdot (-\mu_n n E - D_n \nabla n) - R_n(n, p)$$

$$\frac{\partial p}{\partial t} = -\frac{1}{q} \nabla \cdot (\mu_p p E - D_p \nabla p) - R_p(n, p)$$

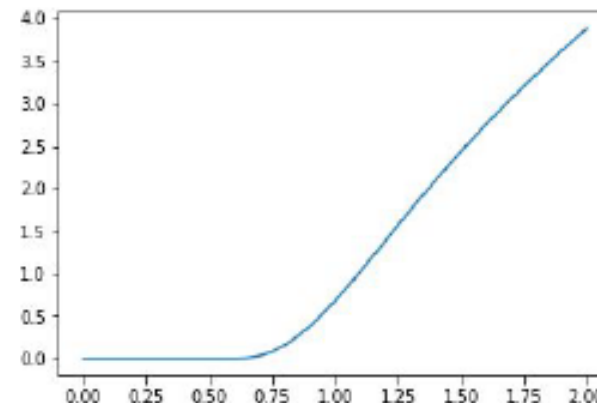


Traditional compact models fit ideal diode + resistor, and can be tuned to match either small or large voltage regimes

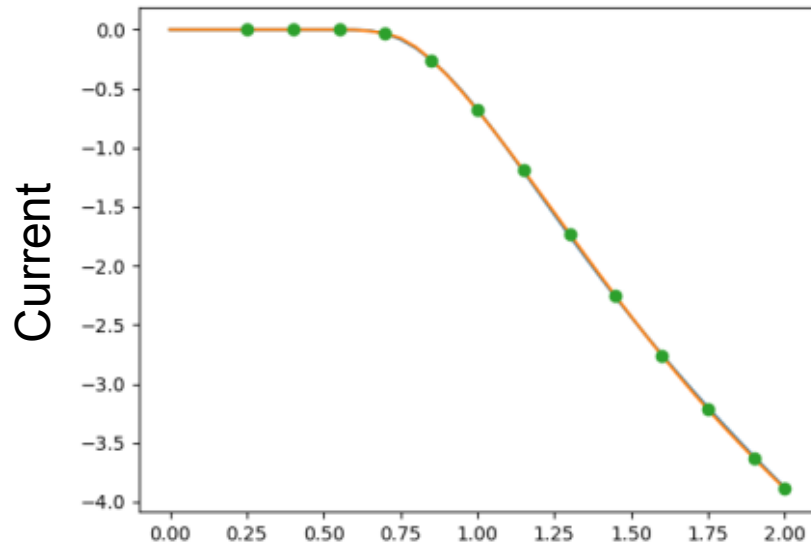
Locally exponential:
I.D. model



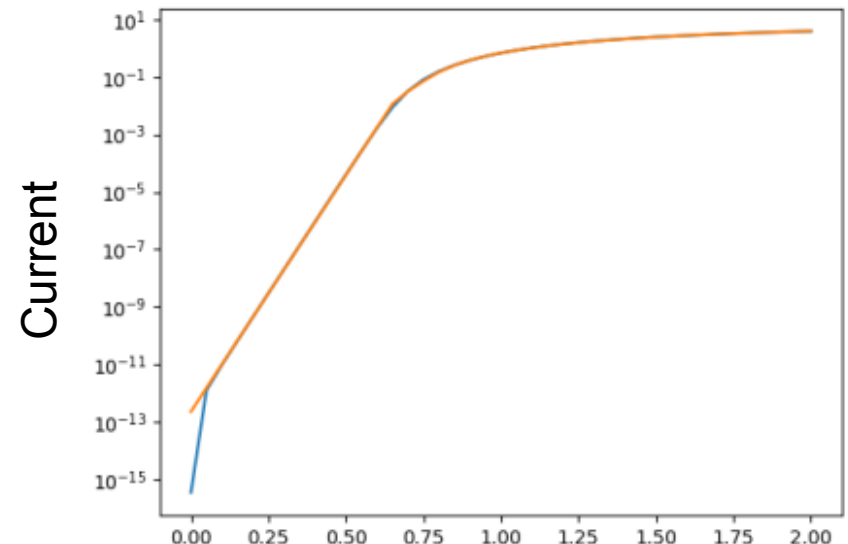
Locally linear:
resistor model



Matching IV-curve – linear scale



Voltage drop



Voltage drop

Extract a conservative surrogate accurate over
fifteen orders of magnitude

May be embedded in a circuit simulator (e.g. Xyce) to
couple coarse-grained high-fidelity PDE model in
multiscale model w/ millions of components

Acknowledgements

- **PHILMs – Physics Informed Learning Machines** for multiscale/multiphysics problems
 - ASCR MMICCs center at the intersection of machine learning and scientific computing
 - PI: George Karniadakis
 - SNL team: Mike Parks (PI), Pavel Bochev, Marta D’Elia, Mamikon Gulian, Ravi Patel, Mauro Perego, Nathaniel Trask
- **PIRAMID – Physics Informed Rapid and Automated ML** for compact model development
 - SNL LDRD to extract efficient compact circuit models from high-fidelity PDE simulation
 - Team: Andy Huang (PI), Xujiao Gao, Shahed Reza, Nathaniel Trask
- **DOE Early Career – Physics informed graph neural networks** for multiscale physics

Applications

Non-equilibrium closures for autoignition in turbulent combustion

Pulse shaping for pulsed power fusion applications on Z-machine

Development of surrogate models for radiation modeling of circuits

Fracture mechanics closures for ice sheet models

Multiscale modeling of lithium-ion batteries during failure

Multiscale closure for subsurface flow through fracture networks

Multiscale data-driven closures for kinetic effects and turbulence in plasmas

Several new projects – please contact for postdoc/collaboration opportunities

(natrask@sandia.gov)