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Evidence of Alfvénic Poynting flux as the primary driver of auroral motion during a geomagnetic substorm

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Key Points:

- We report the one-to-one correlation between a discrete arc and enhanced Poynting flux
- Alfvénic Poynting flux directly powers auroral acceleration
- Dipolarization may be related to Alfvén wave generation

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Abstract

Geomagnetic substorms are major energy transfer events where energy stored in the Earth's magnetotail is released into the ionosphere. Substorm phenomena, including auroral activities, earthward Poynting flux, magnetic field dipolarization, etc, have been extensively studied. However, the complex interplay among them is not fully understood. In a fortuitous event on June 07, 2013, the twin Van Allen Probes (separated by 0.4 hour in local time) observed bursts of earthward Alfvénic Poynting flux in the vicinity of the plasma sheet boundary layer (PSBL). The Poynting flux bursts correlate with enhancements of auroral brightness around the footpoints of both spacecraft. This indicates a temporal and spatial correlation between the auroral brightening and Poynting flux bursts, and that the auroral motion is directly linked to the perpendicular expansion of the Alfvén wave. These observations suggest that the Alfvénic Poynting flux is a primary driver for the auroral electron acceleration. Around the time of auroral brightening, a dipolarization was seen to propagate more than 4 hours in local time during a 20 min period. The azimuthal phase speed of this dipolarization (2 deg/min) is too small to explain the azimuthal motion of the aurora (13.6 deg/min), but the dipolarization could be related to the generation of the Alfvénic Poynting flux through phase mixing at strong density gradients like those in the PSBL.

1 Introduction

Auroral substorms involve the sudden brightening and spatial expansion of auroral activity in the nightside ionosphere (Akasofu, 1964). Low-altitude in-situ observations have shown that the auroral activity is directly connected to the magnetic flux tubes of the upward current region (Chapter 4 in Paschmann et al., 2003, and references therein). Within these flux tubes, electrons are accelerated toward the ionosphere in the auroral acceleration region (about 2-4 Re geocentric distances). This occurs primarily through two major mechanisms: quasi-static potential drops (McFadden et al., 1999) and kinetic Alfvén waves (Chaston et al., 2004). Although strong wave-particle interaction occurs in the auroral acceleration region, it is clear that the free-energy source for the energy conversion lies at higher altitude.

In-situ measurements in the magnetotail have shown that the auroral substorm is one aspect of the more general concept of geomagnetic substorm, which involves sudden release of the magnetic energy stored in the tail. There are a variety of interrelated phenomena resulting from the energy release that can be observed including magnetic field dipolarization (Runov et al., 2009), enhanced earthward Poynting flux (Ergun et al., 2015), and plasma flow channels (Angelopoulos et al., 1992). As the magnetic flux tubes connect the ionosphere to the plasma sheet and its boundary layer in the magnetotail, the coupling between the ionosphere and magnetotail during substorms has been a focus of research for decades.

Conjunction studies between high-altitude in-situ observations and ionospheric or low-altitude measurements have been a powerful tool to study the coupling between the auroral ionosphere and magnetotail. The Polar spacecraft routinely sampled the plasma sheet boundary layer (PSBL) at 4-6 Re, providing in-situ measurements in magnetic conjunction with the auroral zone. Polar conjunctions have highlighted enhanced earthward Poynting fluxes as an important energy source for the auroral acceleration (Keiling et al., 2003). The Poynting fluxes are carried in the form of Alfvén waves and kinetic Alfvén waves (Wygant et al., 2000, 2002). However, one-to-one correlation has not been reported due to limited spatial and temporal resolution of auroral measurements. In this study, with improved auroral measurements in conjunction with two azimuthally separated Van Allen Probes, we show a spatial and temporal correlation holds between Alfvénic Poynting flux bursts and the brightening of a discrete arc in motion.

67 The azimuthal motion of the aurora has previously been linked to the azimuthal
 68 propagation of dipolarizations (Angelopoulos et al., 2008; Ogasawara et al., 2011).
 69 The azimuthal phase speed of the auroral motion is typically on the order of 10-20
 70 deg/min (Paschmann et al., 2003; Angelopoulos et al., 2008; Ogasawara et al., 2011).
 71 This speed corresponds to dipolarizations propagate at 150 km/s at geosynchronous
 72 orbit along the azimuthal direction. Although such “fast” dipolarizations are reported
 73 in many events (Angelopoulos et al., 2008; Ogasawara et al., 2011), statistical studies
 74 show that the typical value of the azimuthal phase speed of the dipolarizations is only
 75 30-50 km/s (Nagai, 1982; Ohtani et al., 2018) at geosynchronous orbit. In Section
 76 4, we discuss the nature of the “fast” dipolarizations, the apparent speed difference
 77 between the “fast” and “slow” dipolarizations, and their relation to the auroral motion
 78 and Alfvénic Poynting flux.

79 In the rest of the paper, we present a case study of a substorm where the Alfvénic
 80 Poynting flux is shown to be the driver of auroral motion. The instrumentation used
 81 is outlined in Section 2. The observations are presented in Section 3. The results are
 82 discussed in Section 4. In Section 5, we offer some conclusions based on these results.

83 2 Instrumentation

84 Although the orbits of the Van Allen Probes (RBSP) tend to maximize time
 85 around the equatorial plane, they can be in conjunction with the auroral zone for
 86 many hours per orbit due to the wobble of the earth’s dipole. Therefore, these probes
 87 can be surprisingly suitable to investigations of auroral physics. The spacecraft provide
 88 DC measurements of the electric and magnetic fields with the EFW (Wygant et al.,
 89 2013) and EMFISIS (Kletzing et al., 2013) instruments. The DC electric (magnetic)
 90 field data are available at 16 (32) samples/sec. Magnetic field data are thus down
 91 sampled to 16 samples/sec to calculate the 3D Poynting flux. Within the calculation,
 92 the DC electric field along the spin axis is not well measured and thus is calculated
 93 using the $\vec{E} \cdot \vec{B} = 0$ assumption. Note that for calculating the parallel Poynting flux,
 94 the spin axis electric fields are not important because the spin axis is primarily along
 95 the background magnetic field during nightside events like the one in this paper. In
 96 addition, the HOPE (Funsten et al., 2013; Spence et al., 2013) instrument onboard
 97 RBSP provides measurements of the electrons and H⁺, He⁺, O⁺ from several eV
 98 to 50 keV per charge at an alternating electron ion ~ 11 cadence, ~ 22 sec between
 99 subsequent ion measurements). The data are collected in 5 look directions, 16 spin
 100 angles, and 72 log-spaced energy bins. In conjunction with the RBSP spacecraft, the
 101 THEMIS all-sky imager (ASI) (Mende et al., 2008) at Pinawa, Manitoba, CA (PINA)
 102 records auroral images every 3 sec at a typical spatial resolution of 1 deg. This high
 103 temporal and spatial resolution enables us to resolve discrete arcs and their motions.

104 3 Observations

105 Figure 1a shows a substorm that occurred from around 04:47 to 04:57 UT on
 106 June 07, 2013. Panel a shows the Dst (black) and AE (red) indices from 04:30 to
 107 05:30 UT, with Dst at ~ -80 nT and AE above 1000 nT around 04:50 UT. This
 108 occurs in the early recovery phase of a moderate geomagnetic storm. Panel b shows
 109 the detrended tilt angle of the measured magnetic field as a function of UT and MLT,
 110 with positive (more dipolar) angles showing as red, as measured on the GOES-13 and
 111 -15 and RBSP-A and -B spacecraft. During the substorm, a dipolarization was seen
 112 from midnight to at least 20 MLT. The dipolarization is characterized by the sudden
 113 increase of the detrended tilt angle of the measured magnetic field $B_{x,y,z}$. The original
 114 tilt angle $\alpha = \arcsin B_z/|B|$ is first detrended using the T89 model (N. A. Tsyganenko,
 115 1989) and then smoothed over a 60 min window to remove the remaining offset. The
 116 detrended tilt angle is color-coded along the spacecraft tracks in panel b and plotted

117 over a shorter time range in panel c and d. In panel b, a linear fit on the times and
 118 MLTs of the major ramps of increasing tilt angle (where blue turns to red) determines
 119 the angular speed as 2 deg/min, which scales to ~ 20 km/s at 5 Re (arrow in panel e).
 120 Panels c and d zoom in to a shorter time range and show the tilt angle (red) and density
 121 (gray and black) observed on RBSP-A (panel c) and -B (panel d). At both spacecraft,
 122 the dipolarization was accompanied by density fluctuations. The dipolarization and
 123 density fluctuation are clearly time-lagged, with peaks on RBSP-B occurring ~ 2 min
 124 after the peaks are observed on RBSP-A. The density is consistently determined by the
 125 particle instrument (gray, low resolution) and by the spacecraft potential (black, high
 126 resolution). The time lag between the density structure is 147.2 sec, determined by
 127 cross-correlating the density from the spacecraft potential. Based on this time lag, the
 128 local angular speed of the density structure is 2.4 deg/min around 22 MLT, which is
 129 consistent with the global value (2 deg/min) of the dipolarization speed calculated from
 130 the data shown in panel b. In addition, the MHD plasma velocity along the azimuthal
 131 direction was calculated to be 20(26) km/s westward at RBSP-A(-B). The speed match
 132 along the azimuthal direction among the dipolarization, density structure, and plasma
 133 motion suggests that the dipolarization propagated along with plasma motion. This
 134 speed will be further discussed and compared to previous studies (Angelopoulos et
 135 al., 2008; Ogasawara et al., 2011; Nagai, 1982; Ohtani et al., 2018) in Section 4.
 136 RBSP-A and -B were in the PSBL before around 05:05 UT, when both spacecraft
 137 saw increases in the electron temperature from several keV to several 10s of keV and
 138 electron density from several 0.1s to above 1 cm^{-3} (not shown). These are typical
 139 features when a spacecraft enters the central plasma sheet from PSBL.

140 Figure 2 shows the Poynting flux (panel 4) and the frequency spectra of the
 141 parallel component of the Poynting flux (panel 5), calculated from the wave electric
 142 (panel 2) and magnetic (panel 3) fields observed on RBSP-A (Figure 2a) and -B (Figure
 143 2b). The electric field data along the spin-axis is calculated using the perpendicular
 144 components based on the $\vec{E} \cdot \vec{B} = 0$ assumption. This calculation does not affect the
 145 parallel Poynting flux much, representing a 10% correction in this event, because the
 146 spin-axis was approximately along the background magnetic field during the substorm.
 147 The Poynting flux was calculated using a Morlet wavelet in the frequency range of 4
 148 mHz to 10 Hz, or from 0.1 sec to 250 sec in terms of wave period. Panels a-4 and
 149 b-4 show that the Poynting fluxes were predominantly earthward (positive parallel
 150 Poynting flux, red curve) at both spacecraft, $>10\text{ mW/m}^2$ after normalized to 100
 151 km altitude to account for the converging flux tube. Panels a-5 and b-5 show that
 152 the main power is between 0.01 to 0.1 Hz. The E/B ratio in this frequency range
 153 was comparable to the local Alfvén speed (panels a-8 and b-8), suggesting that the
 154 earthward Poynting flux was due to a shear Alfvén wave (shear mode as no significant
 155 $|B|$ is seen). The power spectra of the magnetic field follows the Kolmogorov law
 156 (panels a-7 and b-7), suggesting a turbulent environment which may lead to cascades
 157 from shear Alfvén waves to smaller scale waves like kinetic Alfvén waves.

158 Figure 3 shows several snapshots of auroral activity observed every 15 seconds
 159 from 04:55 to 04:56 UT (panels a-e), and at 04:58 UT (panel f) at the PINA THEMIS
 160 All-Sky Imager (the Movie S1 of the supporting information includes all images ob-
 161 served, every 3 seconds, during this time interval). Aurora first brightened around
 162 04:55 UT, then traveled westward (panels a and b), expanded poleward (panels b to
 163 d), and finally faded (panels e and f). Because the in-situ measurements are primarily
 164 azimuthal, we only focus on the westward motion of the aurora. Figure 4 shows four
 165 snapshots capturing this east/west motion (panels a-d), as well as an EWogram (pho-
 166 ton count within an MLat bin as a function of MLon and time) of the aurora within
 167 MLat of 63-64.5 deg to further illustrate and quantify this westward motion. The time
 168 of each of the snapshots is indicated by red arrows to the left of this panel, the location
 169 of the brightest aurora shown by red boxes, and the line with the best fit to these boxes
 170 is over plotted in red as well. The slope of this line gives us the westward motion of

171 the aurora, determined to be 13.6 deg/min through this linear fit. This speed scales
 172 to ~ 120 km/s at 5 Re for comparisons to the in-situ measurements. This speed is ~ 6
 173 times larger than the calculated speed of the dipolarization. The relationship between
 174 the dipolarization and aurora, which were co-observed during the substorm, will be
 175 further discussed in Section 4.

176 This observed auroral activity can be directly compared to in-situ observations of
 177 Alfvénic Poynting flux in the plasma sheet boundary layer along the same magnetic
 178 field lines. Figure 5 shows the Poynting flux observed on RBSP and the auroral
 179 brightness around the ionospheric footpoint of the satellites. The footpoint determined
 180 from the Tsyganenko models, including T89, T96, T01, and T04 (N. A. Tsyganenko,
 181 1989; N. Tsyganenko, 1996; N. A. Tsyganenko, 2002; N. A. Tsyganenko & Sitnov,
 182 2005), is consistent in MLon but uncertain in MLat. Because the mapping is accurate
 183 in MLon, the KEOgram (photon count within an MLon bin as a function of MLat and
 184 time) is used to track the auroral motion, and is shown in panels a-1 and b-1. Also in
 185 these panels, the MLat footpoint location given by the four models is shown and can
 186 be seen to vary from ~ 61.8 to 65.9 degrees. To eliminate this uncertainty, two methods
 187 are used to pin down the MLat of the footpoint. First, we compare the measured in-situ
 188 magnetic field to the models. The T01 model is found to provide the best prediction.
 189 Secondly, we cross-correlate the auroral brightness at all possible MLat to the Poynting
 190 flux. The results of this cross-correlation are shown as a function of MLat in panels
 191 a-2 and b-2, with the peak of the correlation and the full width-half max (FWHM)
 192 shown as well. At both spacecraft, the cross-correlation peaks around the T01 MLat
 193 within the FWHM error bar. Therefore, the T01 model is used for mapping in this
 194 event. Panels a-3 and b-3 plot the auroral brightness at the T01 footpoint against
 195 the parallel Poynting flux. The comparison shows a clear correlation between the two
 196 quantities at each satellite. More importantly, the fact that this correlation holds at
 197 two spacecraft separated along the direction of the auroral westward motion suggests
 198 a dynamical picture where the auroral motion is directly linked to the perpendicular
 199 expansion of the Alfvénic Poynting flux. The spatial and temporal correlation provide
 200 strong evidence that the Poynting flux, transmitting significant electromagnetic energy
 201 earthward, is an important (if not the only) energy source that powers the auroral
 202 acceleration.

203 4 Discussion

204 We have used a close magnetic conjunction between the Van Allen Probes and
 205 the PINA ASI during a geomagnetic substorm on June 07, 2013, to investigate the
 206 coupling between the auroral ionosphere and the near-Earth (5-6 Re) plasma sheet
 207 and its boundary layer. Significant bursts of earthward Alfvénic Poynting flux (>50
 208 mW/m^2 when normalized to 100 km altitude) were observed during the brightening
 209 and expansion of a discrete auroral arc. Although previous studies have highlighted
 210 the correlation between Poynting flux and aurora (Wygant et al., 2000; Keiling et al.,
 211 2003), this is the first time this correlation has been resolved down to auroral structures
 212 like a discrete arc, which can now be spatially resolved due to the improved auroral
 213 observations. In addition, the fact that both RBSP-A and -B were in conjunction with
 214 the discrete arc enables us to study the relation between the Poynting flux and aurora
 215 in a dynamical picture. We find that the Poynting flux correlated with the auroral
 216 brightness at the footpoint of spacecraft at both Van Allen Probes. The correlation
 217 at each footpoint shows that the auroral activity was driven by the Alfvénic Poynting
 218 flux, i.e., the electromagnetic energy transmitted in the form of Alfvén waves along the
 219 magnetic field line. The correlation at both footpoints shows that the auroral motion
 220 was directly linked to the perpendicular expansion of the Alfvén waves/Poynting flux.
 221 The perpendicular expansion of the Alfvén wave refers to either the perpendicular

222 propagation of the Alfvén wave or the perpendicular expansion of the source region
223 for generating the Alfvén wave.

224 The auroral motion was 13.6 deg/min in the azimuthal direction. This is con-
225 sistent with the typical value obtained in previous studies (Paschmann et al., 2003;
226 Angelopoulos et al., 2008; Ogasawara et al., 2011). In addition, we observed a dipol-
227 arization with the azimuthal phase speed of 2 deg/min, which is consistent with the
228 typical azimuthal speed of 30-50 km/s at geosynchronous orbit (Nagai, 1982; Ohtani
229 et al., 2018). Therefore, in our event, the auroral motion is not directly linked to the
230 dipolarization propagation. Instead, the auroral motion is directly linked to the per-
231 pendicular expansion of Alfvén waves/Poynting flux, based on the correlation between
232 the auroral brightness and Poynting flux at multiple locations during the auroral mo-
233 tion. Given the correlation between the aurora and Poynting flux, we infer that the
234 perpendicular expansion of the Alfvén wave/Poynting flux is also not directly linked
235 to the dipolarization.

236 The azimuthal motion of the aurora has previously been linked to the azimuthal
237 propagation of dipolarization (Angelopoulos et al., 2008; Ogasawara et al., 2011).
238 These dipolarizations, which we refer to as the “fast” dipolarizations, are identified as
239 B_z increase or B_x decrease on the order of several 10s sec. The “fast” dipolarizations
240 would have a typical azimuthal speed around 150 km/s at geosynchronous orbit, which
241 is much faster than the typical value of 30-50 km/s found in previous statistical studies
242 (Nagai, 1982; Ohtani et al., 2018). Further investigation on all of the reported “fast”
243 dipolarizations when both electric and magnetic field data are available shows that
244 they are associated with Alfvénic Poynting fluxes (not shown). The period of the peak
245 in the power of the Poynting flux ranges from ~ 10 -100 sec, equivalent to a frequency
246 of 10-100 mHz in the spacecraft frame. The E/B ratio in this frequency range is on
247 the order of local Alfvén speed. All these features resemble the Poynting flux events
248 in this and previous studies (Wygant et al., 2000; Dombeck et al., 2005; Keiling et al.,
249 2003). Therefore, we propose that the signatures in B_x and B_z used for timing the
250 previously reported “fast” dipolarizations can be explained as the signatures of Alfvén
251 waves in B_x and B_z . Consequently, the speed match between the auroral motion and
252 “fast” dipolarizations may effectively be explained as a direct correlation between the
253 auroral brightening and Alfvénic Poynting flux during auroral motion, just as we have
254 shown in this study.

255 Based on our observations, the dipolarization at 2 deg/min cannot be directly
256 linked to the much faster auroral motion and the perpendicular expansion of Alfvén
257 waves. However, an indirect link may exist because dipolarizations are often observed
258 around the time of the Poynting flux bursts, as in the event shown here and other events
259 (Ergun et al., 2015). One potential explanation is the following coupling mechanism.
260 The dipolarization and the associated flow may provide the free-energy source for
261 generating Alfvén waves. The Alfvén wave growth rates are the largest around density
262 gradients (e.g. the plasma sheet boundary layer) because velocity shear and phase
263 shear are large (Lysak et al., 2009). In this scenario, during the slow (2 deg/min) and
264 global (> 4 hours in MLT) azimuthal propagation of the dipolarization, there could
265 be local generation and fast (> 10 deg/min) perpendicular expansion of Alfvén waves
266 near density gradients in the plasma sheet and its boundary layer. These Alfvén waves
267 are associated with earthward Poynting flux bursts which directly drive the auroral
268 brightening and auroral motion.

269 5 Conclusions

270 Figure 6 summarizes the observations in the June 07, 2013 event. As shown in
271 Section 3, a dipolarization was seen to propagate from 0 MLT to 20 MLT during 25
272 min. The westward phase speed of the dipolarization was 2 deg/min, which translates

273 to about 20 km/s at 5 Re. Around the time when the dipolarization passed the RBSP
274 satellites, intense earthward Alfvénic Poynting flux was observed at both satellites and
275 westward auroral motion of 13.6 deg/min was observed in the ionosphere. A temporal
276 and spatial correlation was found between the Poynting flux and the auroral brightness
277 at the spacecraft footpoint, suggesting that the Alfvénic Poynting flux powers the
278 auroral acceleration and that the Alfvén waves couple the auroral acceleration region
279 to the near-earth plasma sheet boundary layer (PSBL).

280 These observations offer strong evidence for Alfvén waves as the primary driver of
281 auroral acceleration during a substorm. The results also suggest that the dipolarization
282 was not directly involved in the coupling between the auroral and PSBL. However,
283 dipolarizations have long been known to play a key role during geomagnetic substorms
284 (e.g. Fu et al., 2020, and references therein). For example, our study suggests that the
285 dipolarization may be related to the generation of the observed Alfvénic Poynting flux.
286 The relationship between dipolarizations and Alfvén waves/Poynting flux is important
287 for understanding the major energy conversions during geomagnetic substorms, and
288 further study is necessary to acquire a complete picture of the substorm driving process.

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294 rbsp-ect.newmexicoconsortium.org/rbsp_ect.php](https://rbsp-ect.newmexicoconsortium.org/rbsp_ect.php). EFW and EMFISIS data are
295 available at <https://cdaweb.gsfc.nasa.gov/pub/data/rbsp>. THEMIS all-sky im-
296 ager data are available at [http://themis.ssl.berkeley.edu/data/themis/thg/11/
297 asi/](http://themis.ssl.berkeley.edu/data/themis/thg/11/asi/).

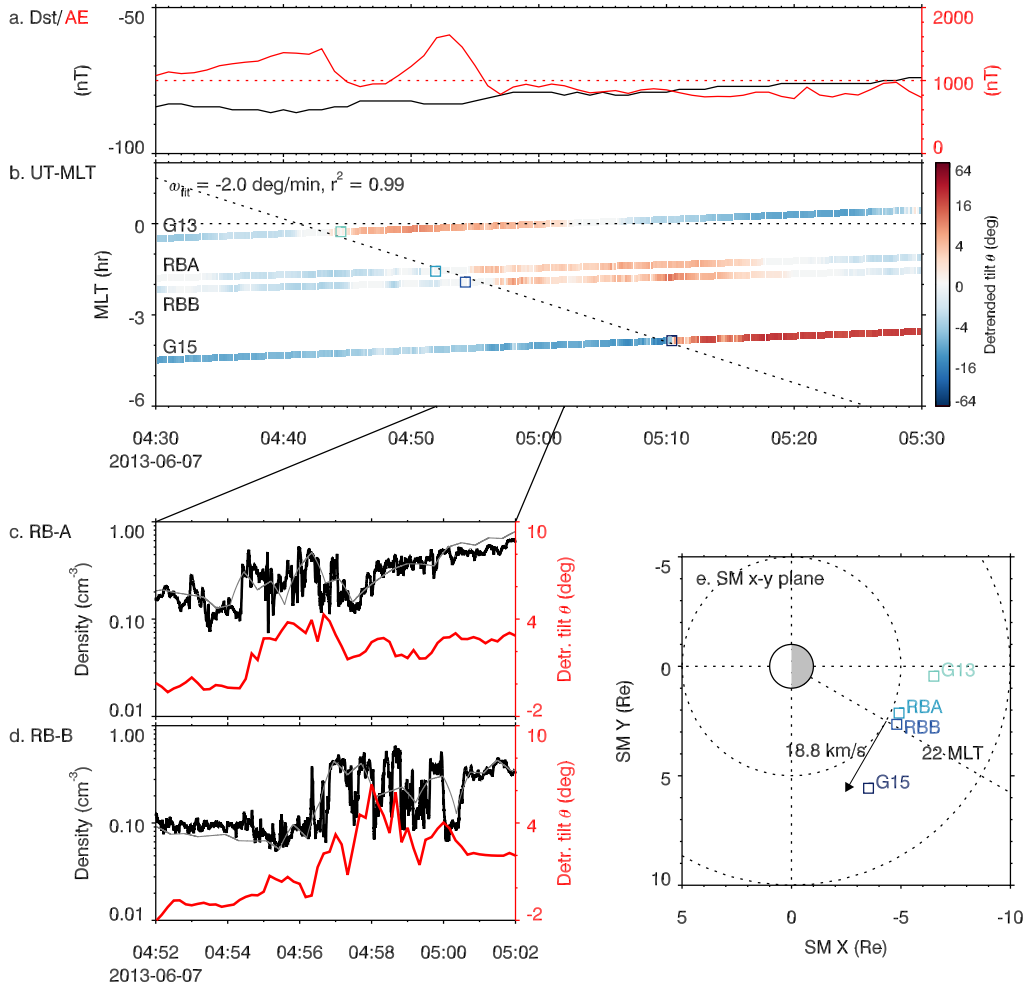


Figure 1. The dipolarization propagated westward with density structures during a substorm on June 07, 2013. Panel a shows the Dst and AE indices, showing a substorm around 04:47 to 04:57 UT during the recovery phase of a storm. Panel b shows the dipolarization as measured in tilt angle color-coded along tracks of 4 spacecraft (GOES-13 and -15 and RBSP-A and -B) in the pre-midnight section. At each satellite, the time of dipolarization is when the increasing tilt angle crosses 0 (where blue turns red). A linear fit on the times and MLTs of the dipolarization determines the drift of dipolarization in MLT to be 2 deg/min. Panel c and d show the density as determined from the particle instrument (low-resolution, gray) and spacecraft potential (high-resolution, black), and the detrended tilt angle (red). Panel e shows the azimuthal phase speed of the dipolarization scaled to the spacecraft location.

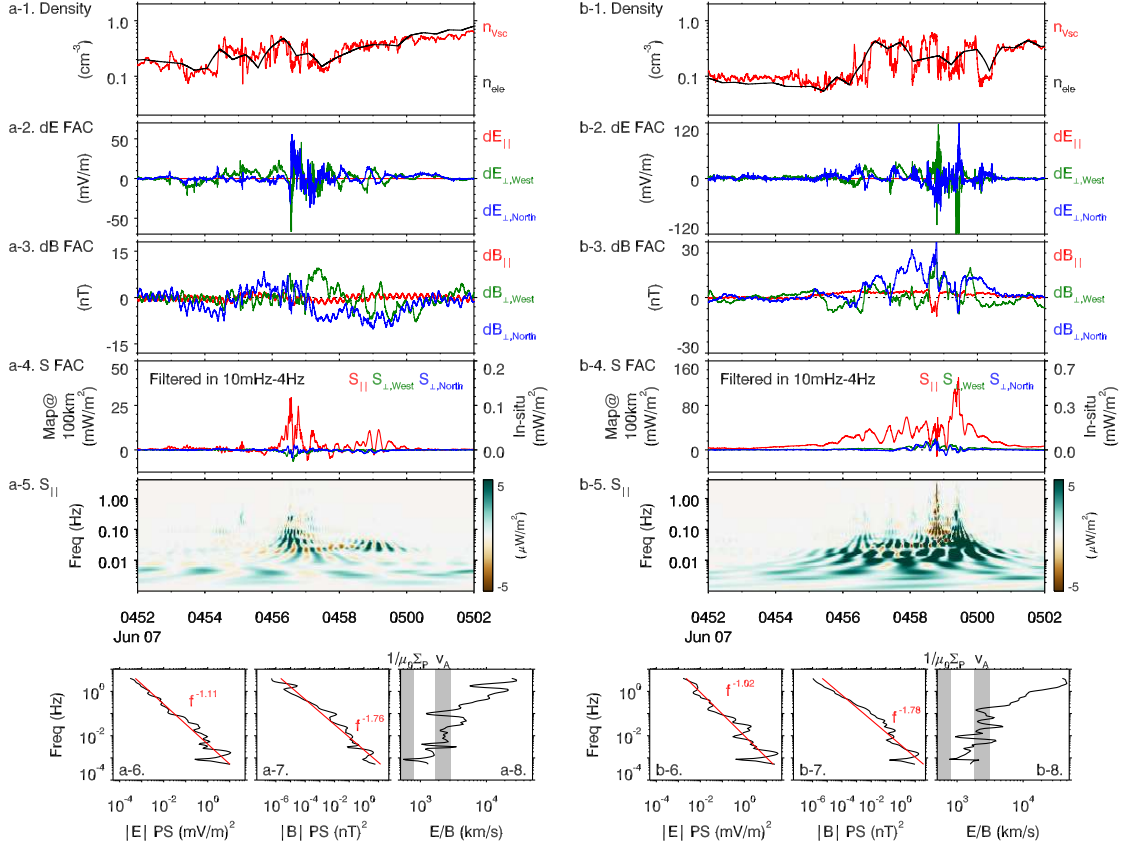


Figure 2. Poynting flux calculation and properties. Data for RBSP-A are presented in the panels on the left and RBSP-B data are on the right in the same format. Panel a-1 shows the same density data in Figure 1c for reference purposes. Panel a-2 and a-3 show the waveforms of the wave electric and magnetic fields in the field-aligned coordinate (FAC), where the components are parallel (red), westward (green) and northward (blue). Panel a-4 shows the Poynting flux filtered in 10 mHz to 4 Hz before and after normalized to 100 km to account for the converging flux tube. Panel a-5 shows the parallel component of the Poynting flux in the time-frequency domain before integration. Panel a-6 and a-7 show the power spectra of the electric and magnetic field. Panel a-8 shows the E/B ratio as a function of frequency. The shaded areas are the E/B ratios expected for the Pedersen conductivity and Alfvén waves, where v_A is calculated from the measured magnetic field, density and ion species. The Poynting fluxes at both spacecraft were predominantly earthward. The main power ranged 10-100 mHz and in this range the E/B ratio is comparable to the local Alfvén speed.

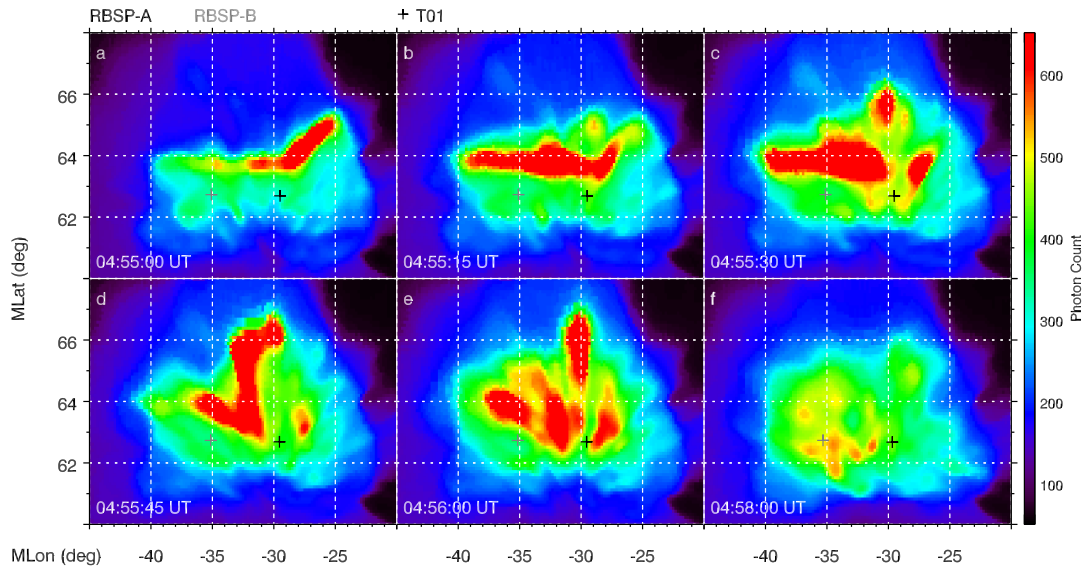


Figure 3. Snapshots of the auroral westward (panels a and b) and poleward (panels b to d) motions and the feinting of the auroral activity (panels e to f) from the Movie S1 of the supporting information. The footprints of RBSP-A (black) and -B (gray) are over-plotted.

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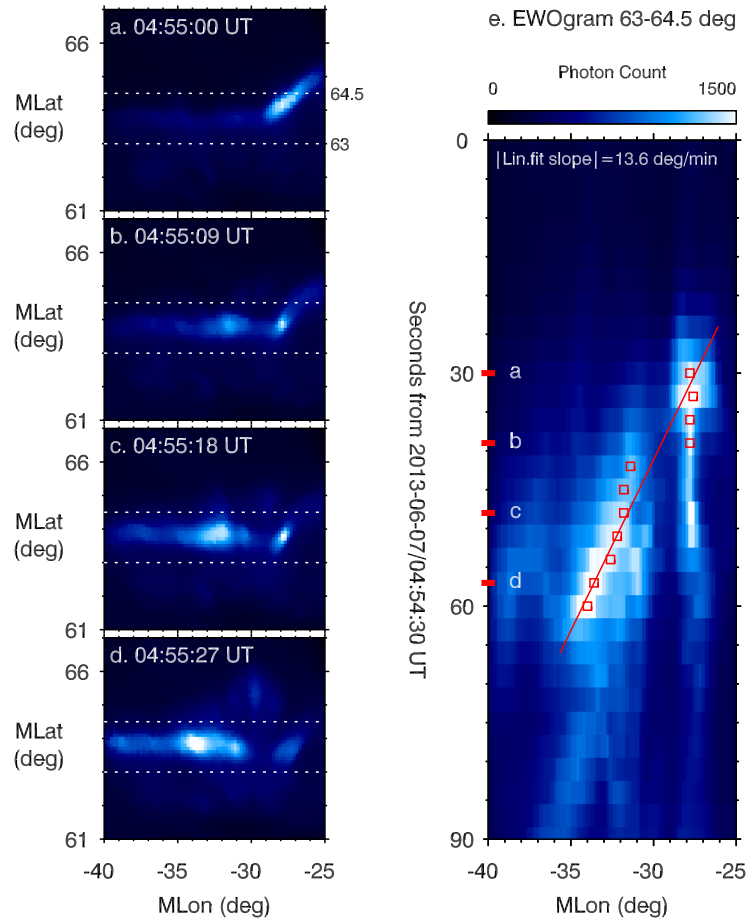


Figure 4. Determining the westward motion through the EWOgram between 63-64.5 deg MLat. The westward motion is 13.6 deg/min, determined through a linear fit of the MLon of the brightest aurora in each time frame. This speed is about 6 times of the westward speed of the dipolarization. The region around -30 deg MLon could be either a ground object blocking the camera or a physical dark region of the aurora. Because of the ambiguity, the four red boxes around the dark region are not used for the linear fit (including them results in a similar value of 14.3 deg/min).

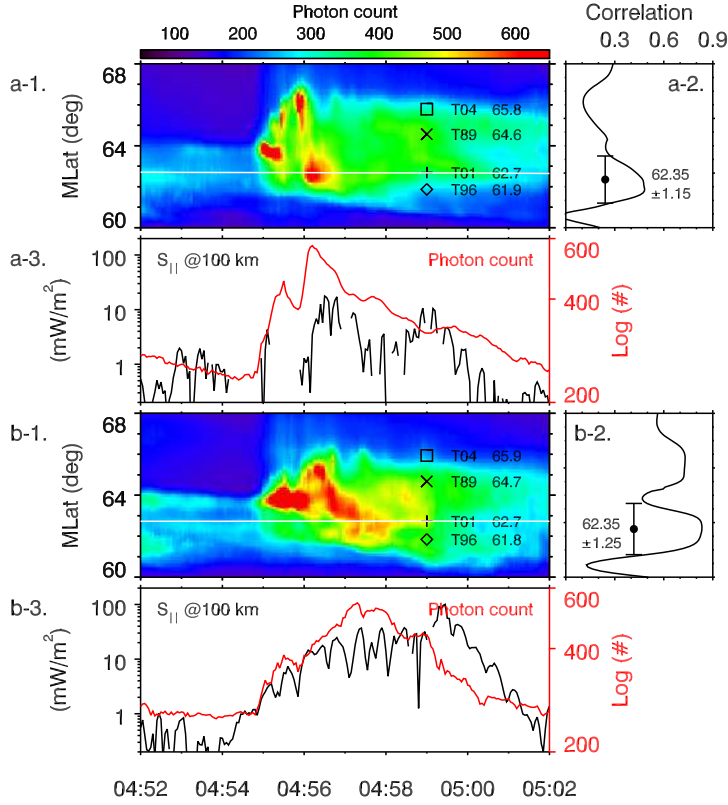


Figure 5. The spatial and temporal correlation between the Poynting flux and the auroral brightness around the footpoint. Panel a-1 shows the auroral KEOgram around the MLon of RBSP-A. The MLat of the footpoint according to 4 Tsyganenko models are over plotted. Although a wide range of MLat is predicted, the T01 model is chosen for two reasons. Firstly, T01 predicted the in-situ magnetic field the best. Secondly, as shown in panel a-2, cross-correlation between the Poynting flux and the auroral brightness at all MLat peaks around the T01 MLat. The T01 MLat is marked by the white line in panel a-1. For the same reason, panel b1- to b-2 show that the T01 model also works for RBSP-B. Panel a-3 and b-3 show the Poynting flux and the auroral brightness at the T01 MLat. After 04:58 UT the correlation weakens somewhat, consistent with the fact that the satellite is moving out of the field of view of the auroral imager.

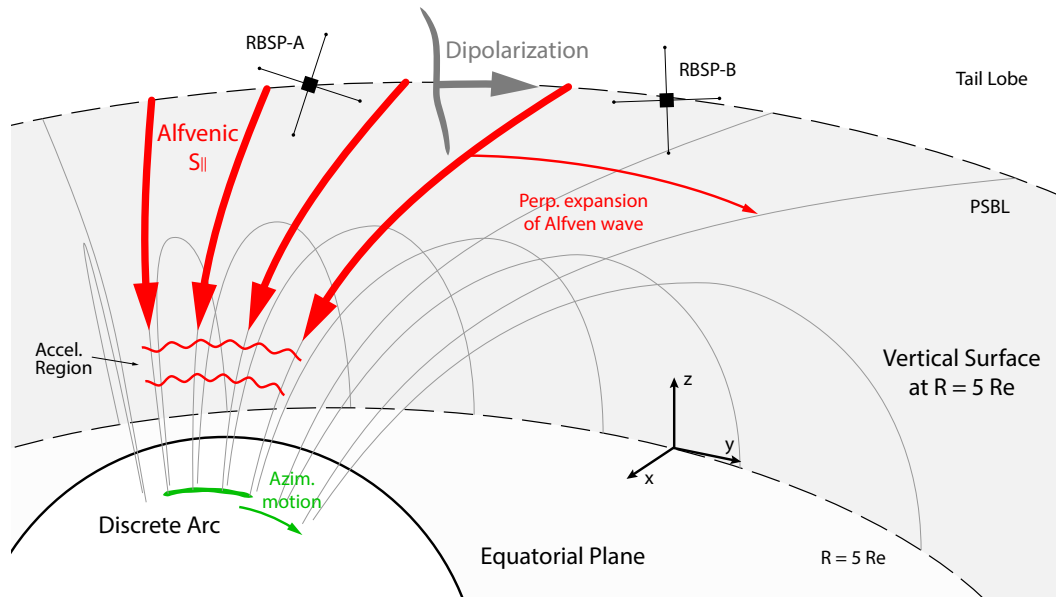


Figure 6. A schematic view to summarize the observations between the auroral ionosphere and the inner-magnetosphere. The curved gray wall is a vertical cylindrical surface at $5 R_E$, containing the two RBSP spacecraft. The earthward Alfvénic Poynting flux transmits electromagnetic energy to the ionosphere, while auroral brightening and a dipolarization are co-observed. A possible chain of energy conversion includes: Poynting flux (shear Alfvén wave) is generated related to the dipolarization; The Alfvén waves propagate earthward to transmit the electromagnetic energy to lower altitude, powering the parallel acceleration of electrons around the auroral acceleration region. The westward motion of the aurora is directly linked to the perpendicular expansion of Alfvén waves.

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