

International Conference on Simulation of Semiconductor Processes and Devices (SISPAD 2021)
Dallas, TX, September 27 - 30, 2021

Quantum Transport Simulations for Si: P δ -layer Tunnel Junctions

Juan P. Mendez, Denis Mamaluy, Xujiao (Suzey) Gao and Shashank Misra
jpmende@sandia.gov

Sandia National Laboratories, Albuquerque, New Mexico



Sandia National Laboratories is a multimission laboratory managed and operated by National Technology and Engineering Solutions of Sandia, LLC., a wholly owned subsidiary of Honeywell International, Inc., for the U.S. Department of Energy's National Nuclear Security Administration under contract DE-NA-0003525.

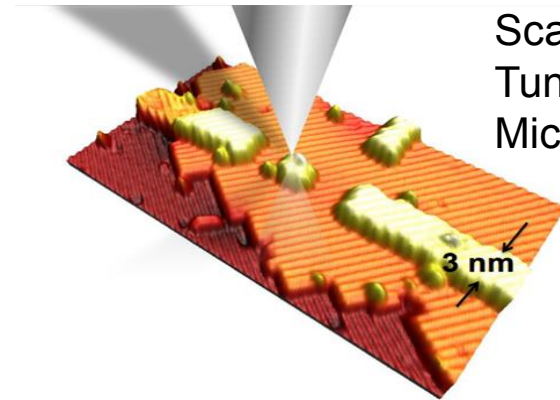
Sandia National Laboratories is a multimission laboratory managed and operated by National Technology and Engineering Solutions of Sandia, LLC., a wholly owned subsidiary of Honeywell International, Inc., for the U.S. Department of Energy's National Nuclear Security Administration under contract DE-NA0003525.



Motivation

Atomic Precision Advanced Manufacturing (APAM)

is a process of area-selective dopant incorporation at the atomic scale



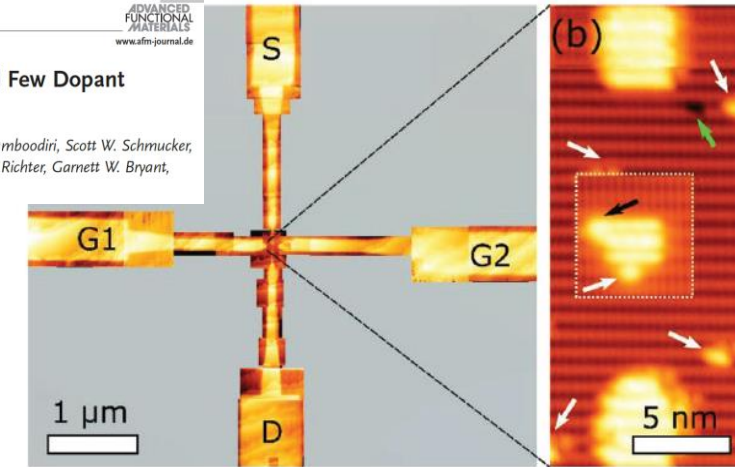
STM =
Scanning
Tunneling
Microscope

FULL PAPER

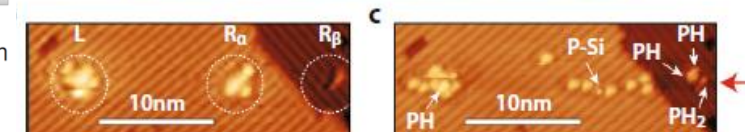
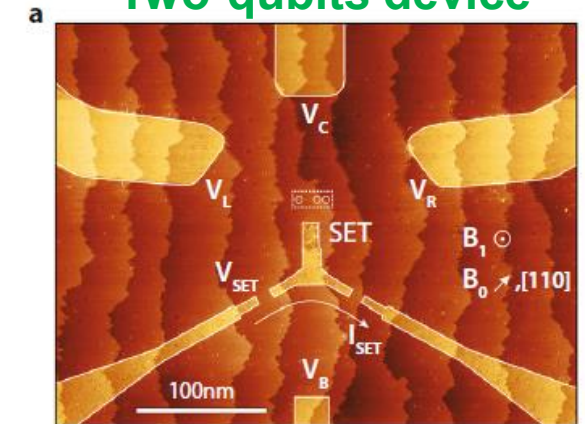
Atom-by-Atom Fabrication of Single and Few Dopant Quantum Devices

Jonathan Wyrick,* Xiqiao Wang, Ranjit V. Kashid, Pradeep Nambodiri, Scott W. Schmucker, Joseph A. Hagmann, Keyi Liu, Michael D. Stewart Jr., Curt A. Richter, Garrett W. Bryant, and Richard M. Silver[†]

Single atom transistor



Two-qubits device



ARTICLE

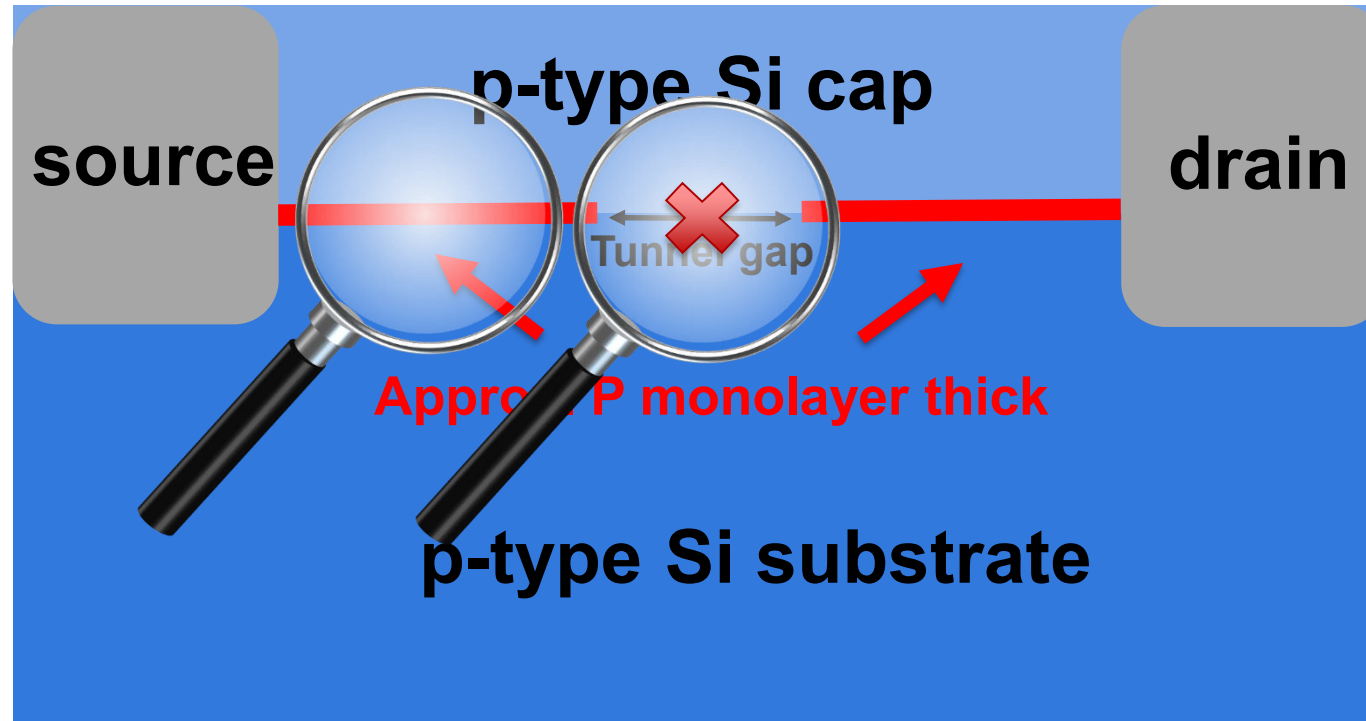
<https://doi.org/10.1038/s41467-021-23662-3> OPEN

Coherent control of a donor-molecule electron spin qubit in silicon

Lukas Fricke^{1,3}, Samuel J. Hille^{1,2,3}, Ludwik Kranz¹, Yousun Chung¹, Yu He¹, Prasanna Pakkiam¹, Matthew G. House¹, Joris G. Keizer¹ & Michelle Y. Simmons¹

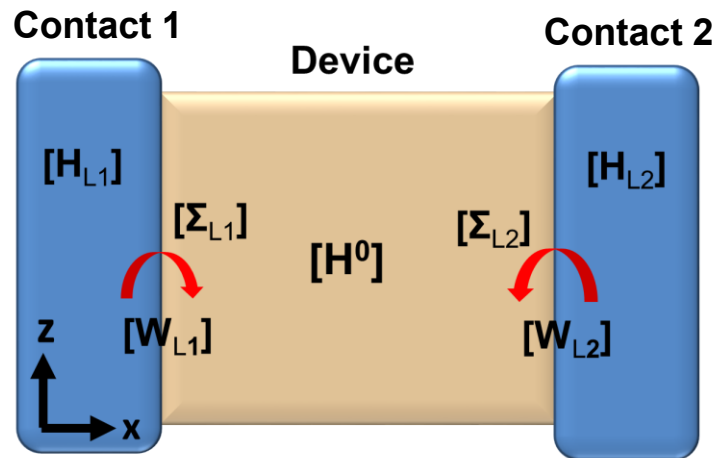
Si: P delta-layer Tunnel Junction

- High potential for **quantum computing** and **advanced microelectronic devices**

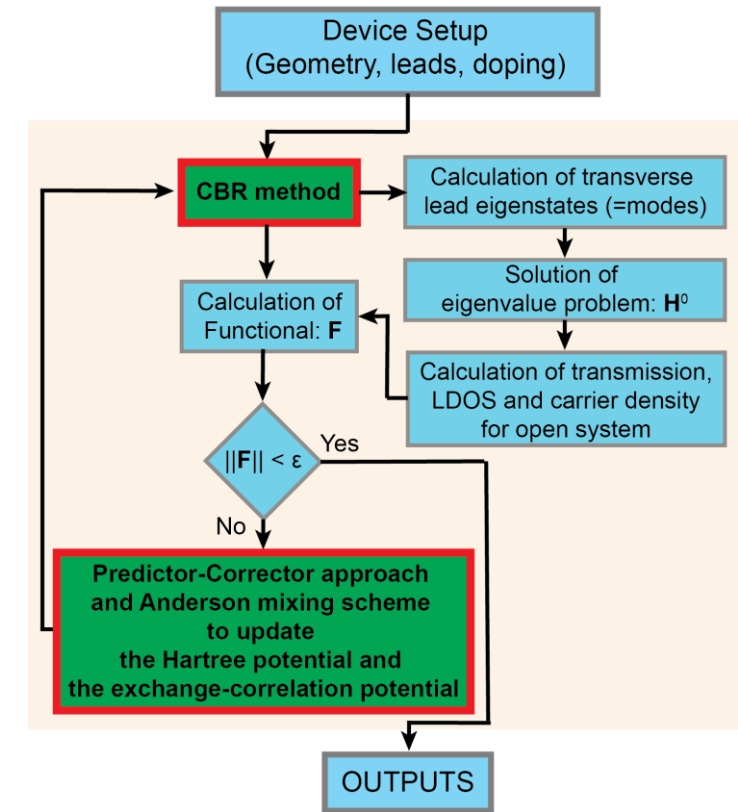


Our quantum transport simulator is based on **Non-Equilibrium Green's function (NEGF)** formalism with

- Fully charge self-consistent solution of Poisson-open system Schrödinger equation
- Single-band (Γ valley) effective mass approximation
- Predictor-corrector approach and Anderson mixing scheme
- Contact Block Reduction (CBR) method for fast numerical efficiency

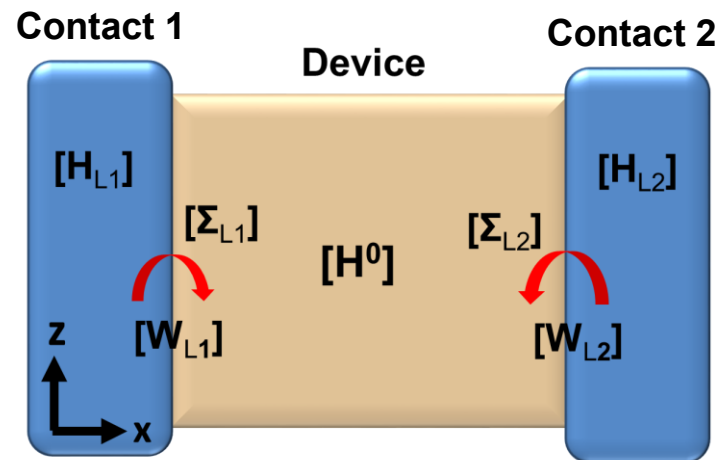


Quantum transport simulator

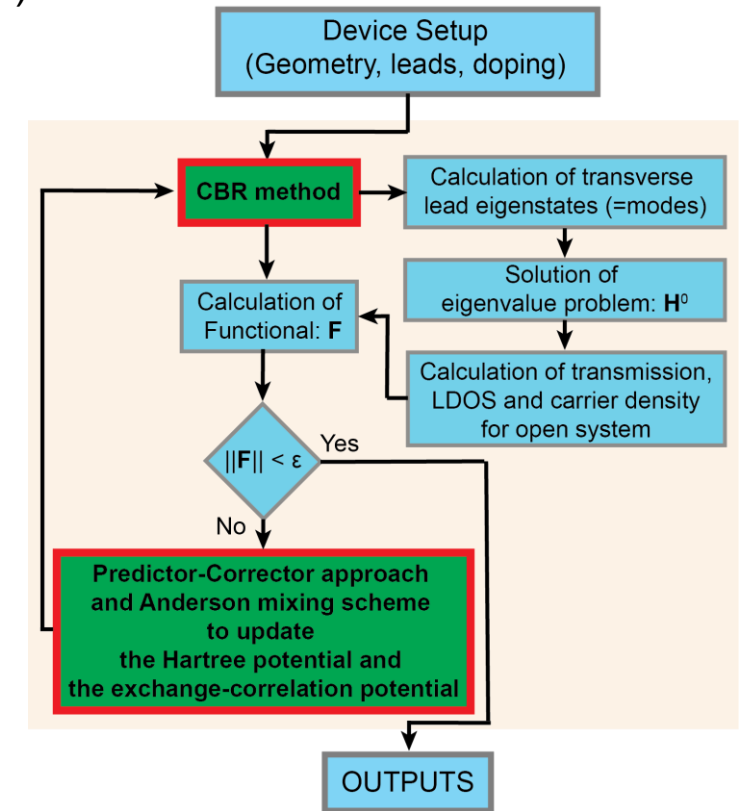


Our quantum transport simulator is based on **Non-Equilibrium Green's function** (NEGF) formalism with

- Fully charge self-consistent solution of Poisson-open system Schrödinger equation
- Single-band (Γ valley) effective mass approximation
- Predictor-corrector approach and Anderson mixing scheme
- **Contact Block Reduction (CBR) method for fast numerical efficiency**



Quantum transport simulator



- D. Mamaluy *et al.*, J. Appl. Phys., vol. 93, no. 8, p. 4628-4633, 2003.
- D. Mamaluy *et al.* Phys. Rev. B, vol. 71, p. 245321, 2005.

Non-Equilibrium Green Function

- Current from leads λ to λ'

$$J_{\lambda\lambda'} = \frac{2e}{h} \int T_{\lambda\lambda'}(E) (f_{\lambda}(E) - f_{\lambda'}(E)) dE$$

$$T_{\lambda\lambda'}(E) = \text{Tr} [\Gamma_{\lambda} G_D \Gamma_{\lambda'} G_D^{\dagger}]$$

Self-energy matrix of lead

$$\Gamma_{\lambda} = i(\Sigma_{\lambda} - \Sigma_{\lambda}^{\dagger})$$

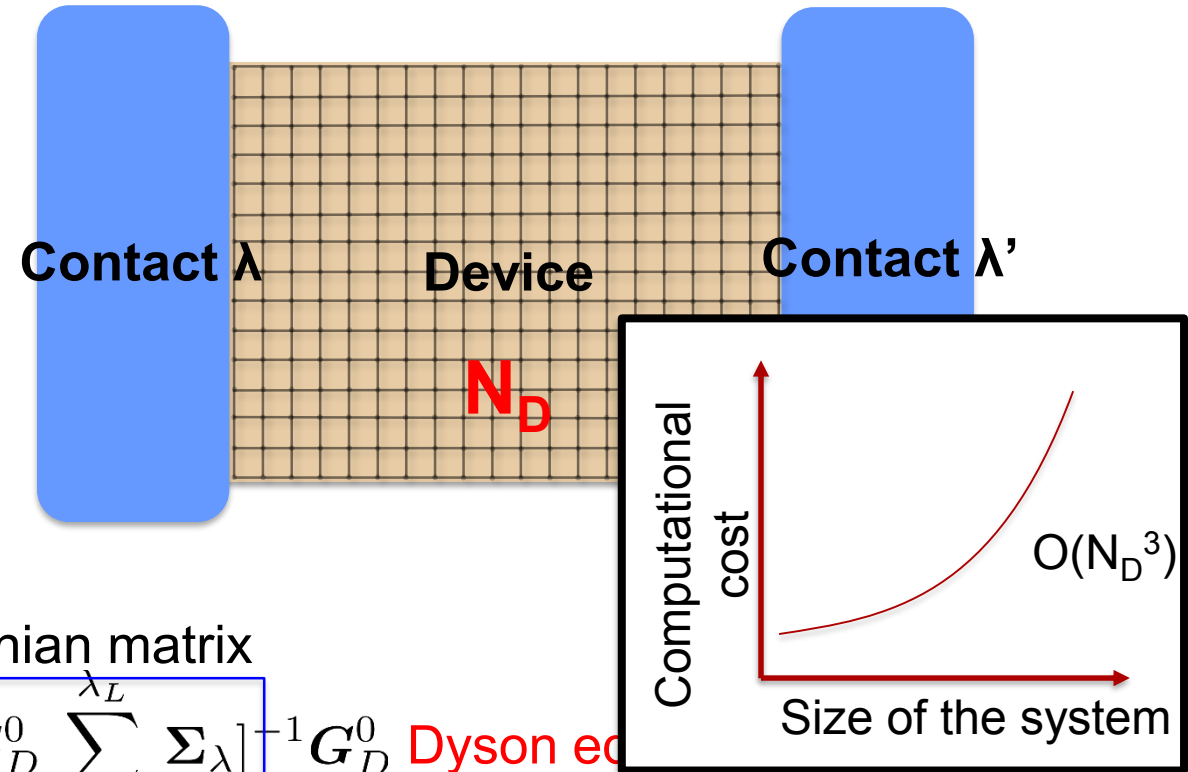
Open-device Hamiltonian matrix

$$G_D = [I - G_D^0 \sum_{\lambda=\lambda_1}^{\lambda_L} \Sigma_{\lambda}]^{-1} G_D^0 \quad \text{Dyson eq}$$

- Electron density

$$\rho(\mathbf{r}_i) = \sum_{\lambda} \int_{-\infty}^{\infty} \rho_{\lambda}(\mathbf{r}_i, E) f_{\lambda}(E) dE$$

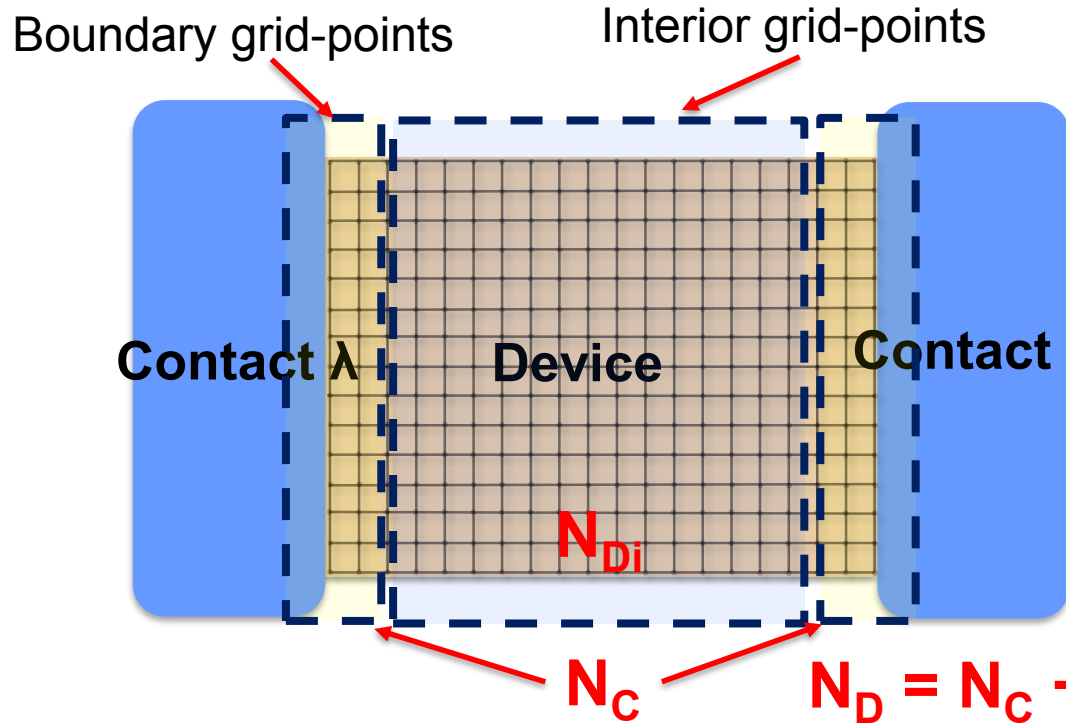
$$\rho_{\lambda}(\mathbf{r}_i, E) = \frac{1}{2\pi} G_D \Gamma_{\lambda} G_D^{\dagger}$$



$$G_D^0 = [E^+ - H_D^0]^{-1} = \sum_{\alpha} \frac{|\phi_{\alpha}\rangle \langle \phi_{\alpha}|}{E^+ - E_{\alpha}}$$

$(N_D \times N_D)$

Contact Block Reduction Method



$$H_D = \begin{pmatrix} \boxed{H_C} & \boxed{H_{CD_i}} \\ \boxed{H_{D_iC}} & \boxed{H_{D_i}} \end{pmatrix} \quad \begin{matrix} (N_C \times N_C) & (N_C \times N_{Di}) \\ (N_{Di} \times N_C) & (N_{Di} \times N_{Di}) \end{matrix}$$

$$G_D = \begin{pmatrix} \boxed{G_C} & \boxed{G_{CD_i}} \\ \boxed{G_{D_iC}} & \boxed{G_{D_i}} \end{pmatrix}$$

$$\Sigma = \sum_{\lambda=\lambda_1}^{\lambda_L} \Sigma_\lambda = \begin{pmatrix} \boxed{\Sigma_C} & \boxed{0} \\ \boxed{0} & \boxed{0} \end{pmatrix}$$

- Electrical current:**

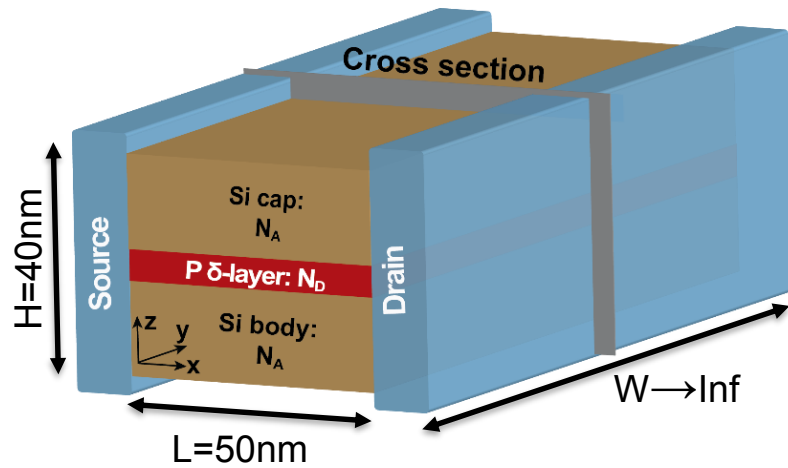
$$J_{\lambda\lambda'} = \frac{2e}{h} \int T_{\lambda\lambda'}(E)(f_\lambda(E) - f_{\lambda'}(E))dE, \quad T_{\lambda\lambda'}(E) = \text{Tr}(\boxed{\Gamma_{C_\lambda} G_C \Gamma_{C_{\lambda'}} G_C^\dagger}), \quad G_C = \boxed{[I - G_C^0 \Sigma_C]^{-1} G_C^0}$$

- Electron density:**

$$\rho(\mathbf{r}_i) = \int_{-\infty}^{\infty} \Xi(E) f_\lambda(E) dE, \quad \Xi(E) = \frac{1}{2\pi} \frac{\text{Tr}[\boxed{B_C^{-1} \Gamma_C B_C^{-1\dagger}}]}{(E^+ - E_\alpha)(E^- - E_\alpha)}, \quad B_C = \boxed{1_C - \Sigma_C G_C^0} \quad \begin{matrix} \nearrow \\ \searrow \end{matrix} (N_C \times N_C)$$

Last year in SISPAD 2020 ...

Infinite-width δ -layer systems:



$$N_D = 1.2 \times 10^{14} \text{ cm}^{-2}, N_A = 10^{17} \text{ cm}^{-3}$$

communications physics

ARTICLE

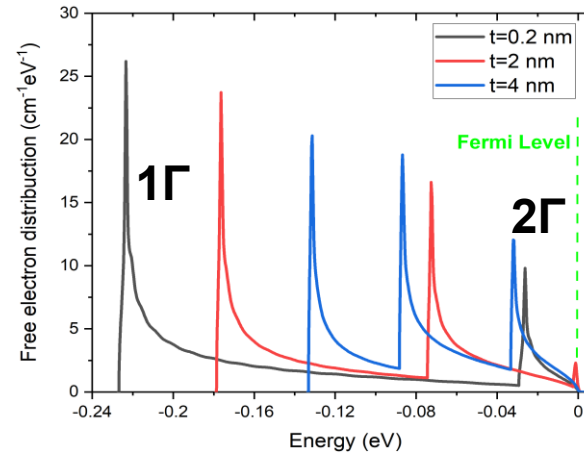
<https://doi.org/10.1038/s42005-021-00705-1> OPEN

Revealing quantum effects in highly conductive δ -layer systems

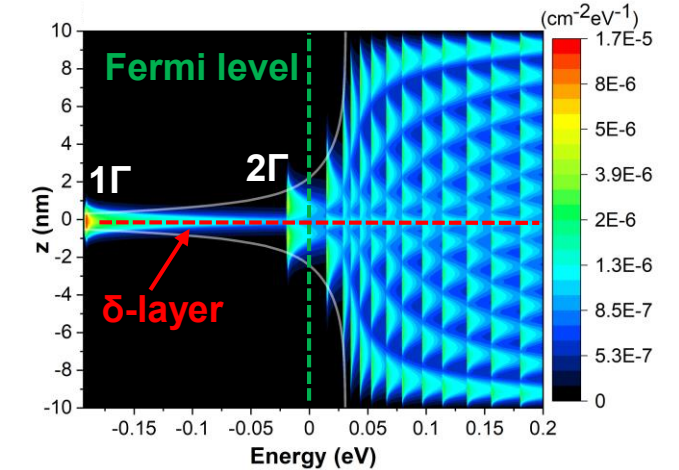
Denis Mamaluy¹, Juan P. Mendez¹, Xujiao Gao¹ & Shashank Misra¹

<https://www.nature.com/articles/s42005-021-00705-1>

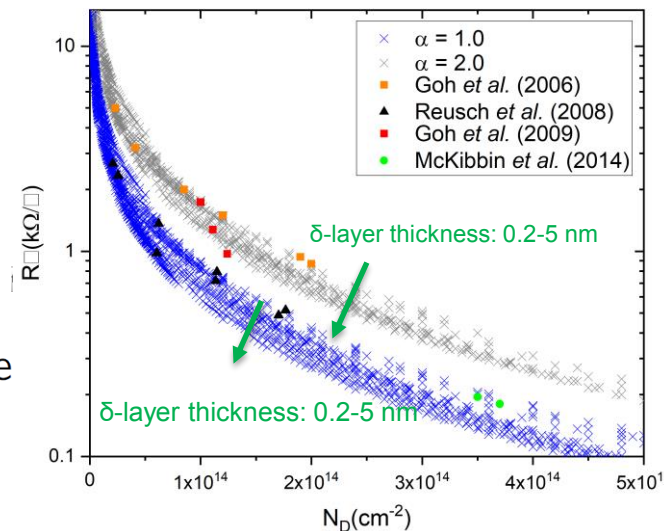
Prediction of shallow Sub-bands



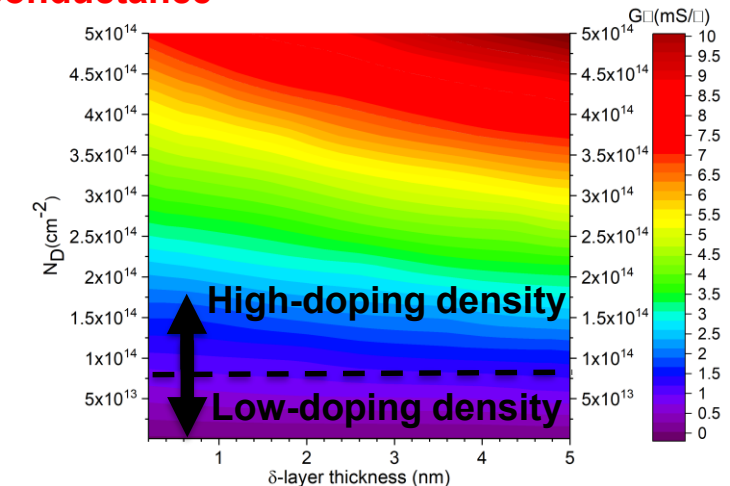
The so-called “quantum menorah”: LDOS(E,z)



Good agreement with experiments at 4K!!



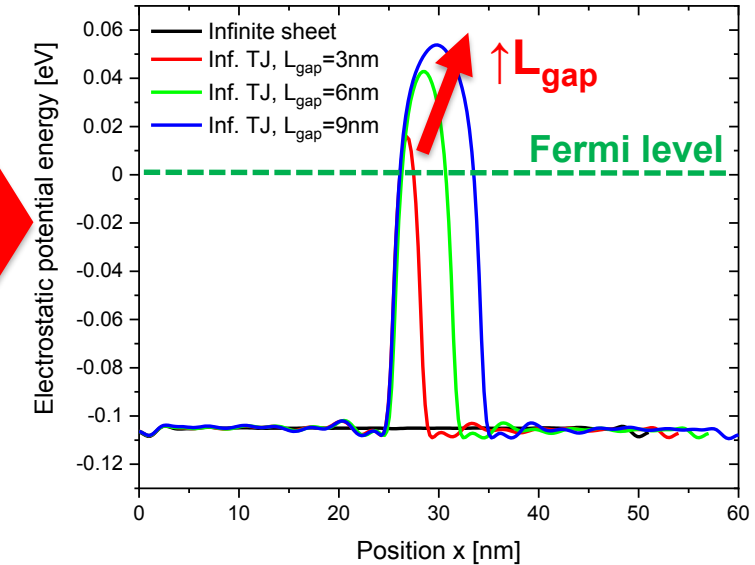
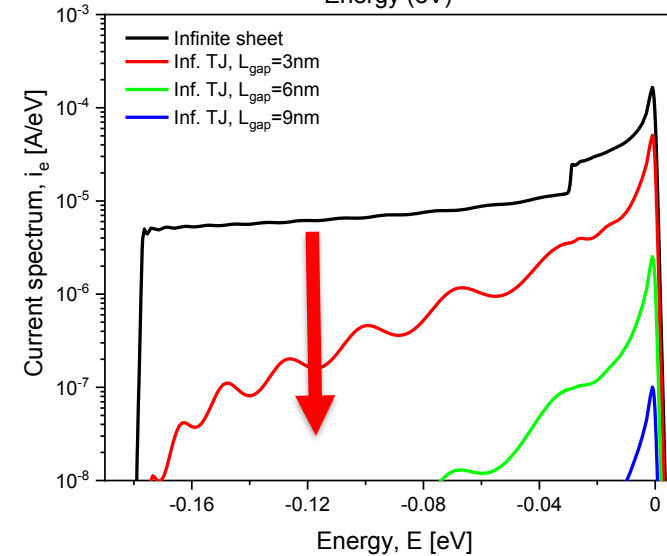
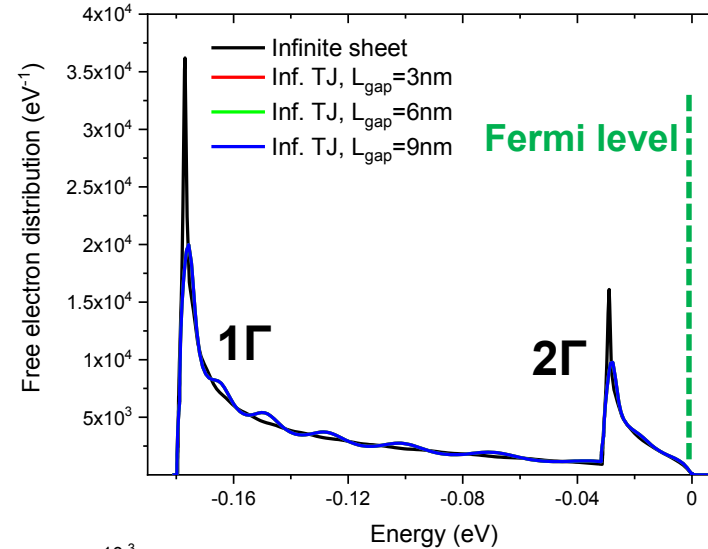
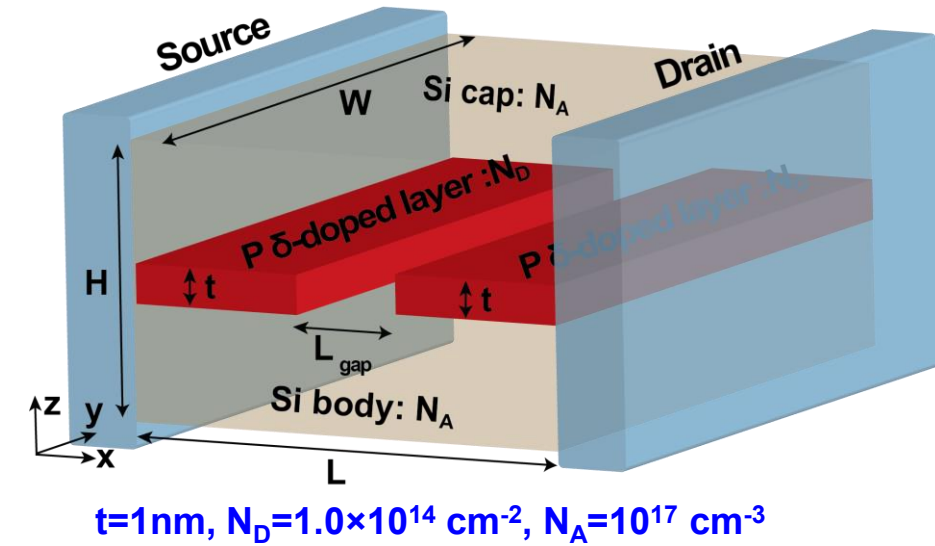
Quantum thickness dependence on the sheet conductance



Conductive Properties for Ideal Tunnel Junctions Sandia National Laboratories

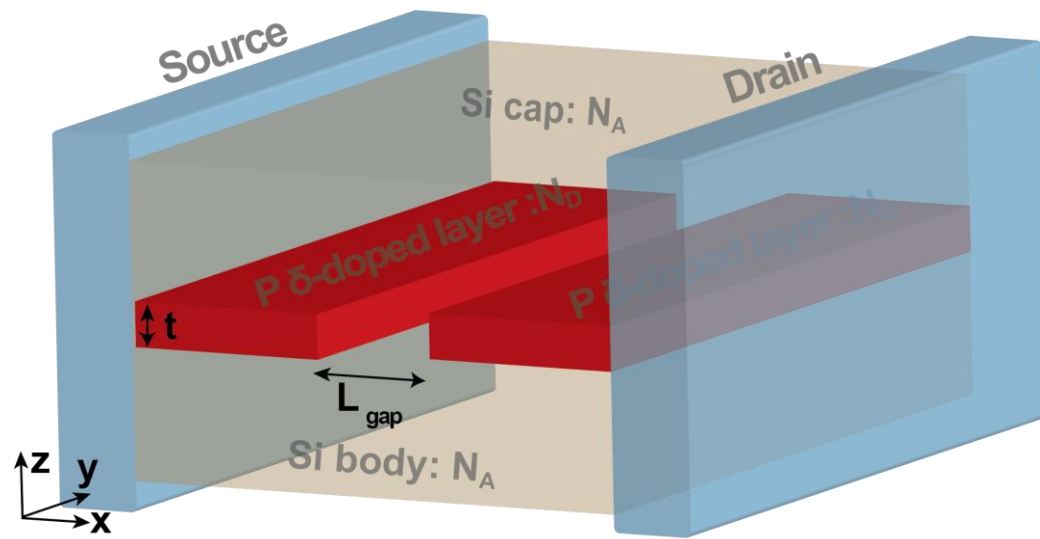
- Low-energy electrons contribution on the current is depressed with the tunnel gap

Infinite-width δ -layer Tunnel Junctions

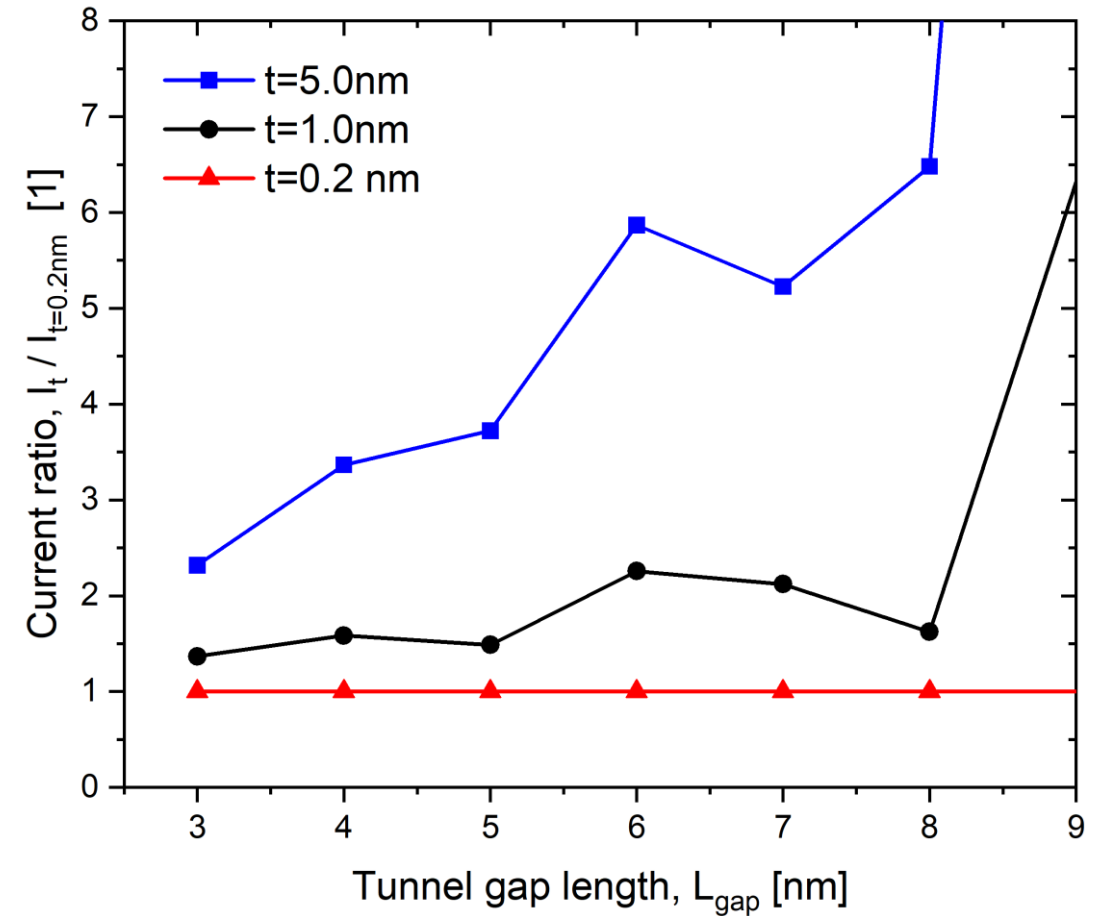


Variation of the δ -layer thickness

- The tunneling rate considerably increases with the δ -layer thickness, specially for larger tunnel gaps

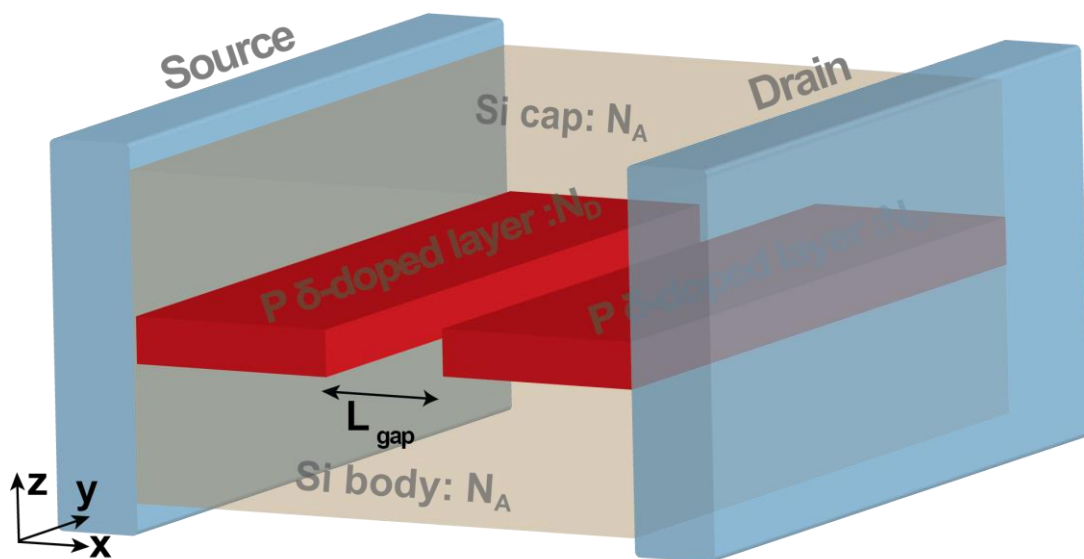


$$N_D = 1.0 \times 10^{14} \text{ cm}^{-2}, N_A = 10^{17} \text{ cm}^{-3}$$

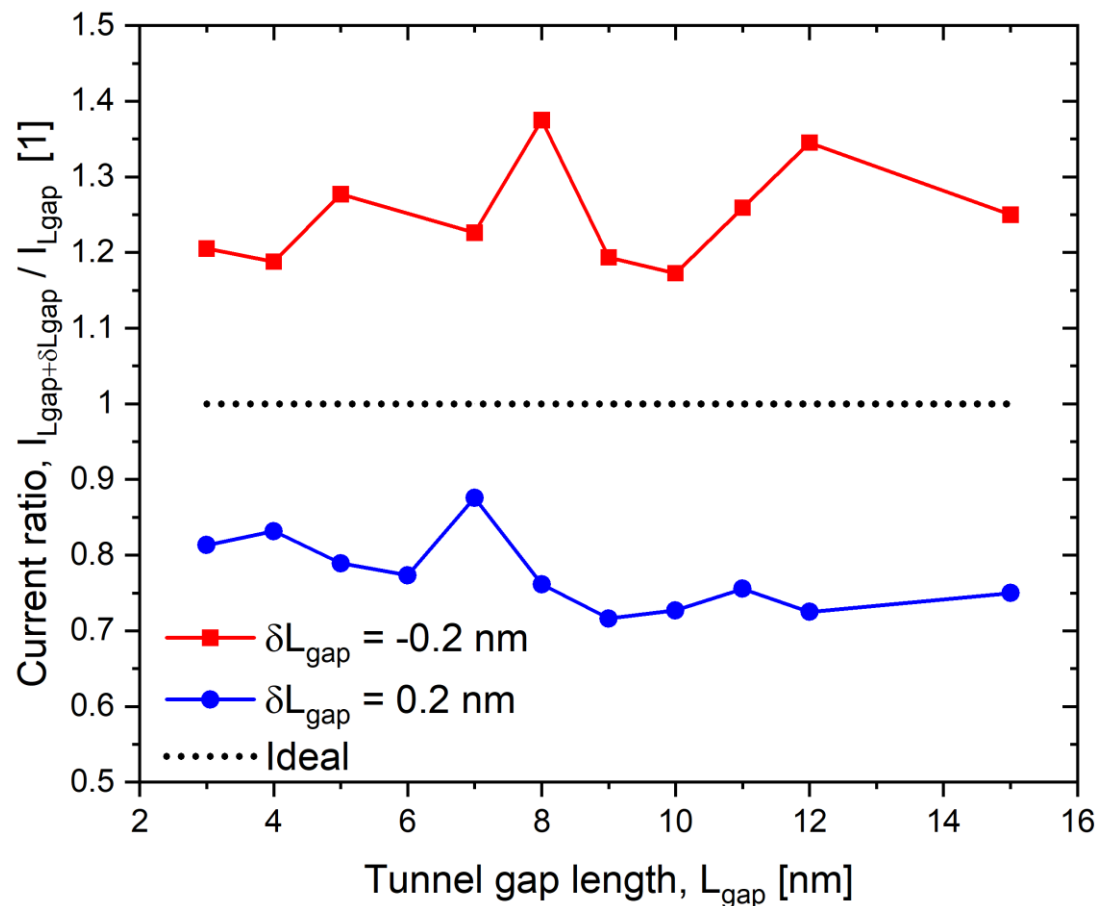


Variation of the tunnel gap length

- Small variations of the gap length (around 0.2 nm) leads to a tunneling rate change of around 20%-30%

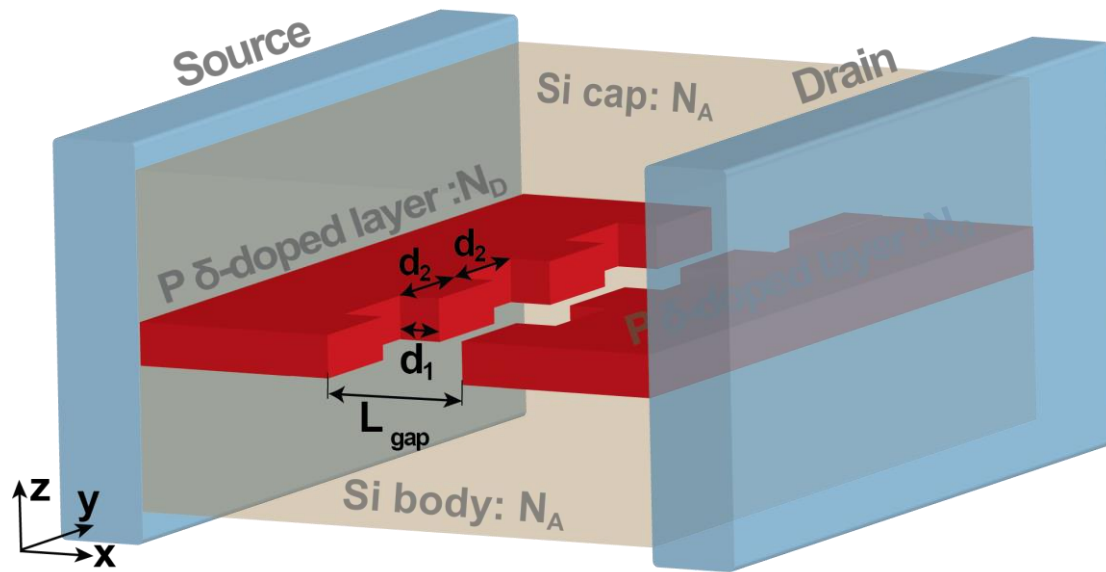


$$N_D = 1.0 \times 10^{14} \text{ cm}^{-2}, N_A = 10^{17} \text{ cm}^{-3}$$



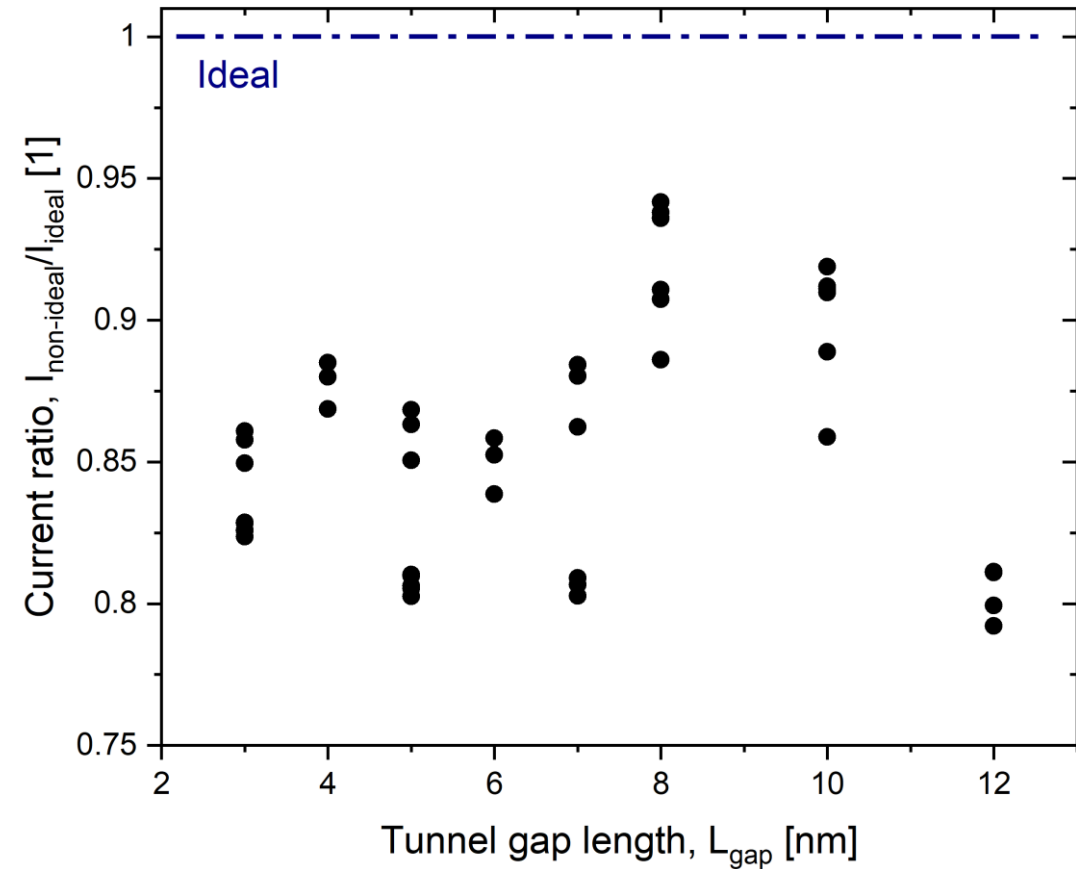
Effects of Edge roughness

- Edge roughness might lead to a decreases of the tunneling rate up to 5%-25%



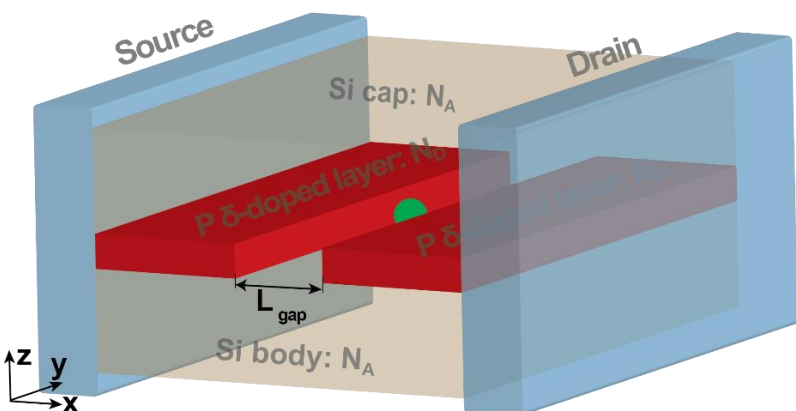
$$N_D = 1.0 \times 10^{14} \text{ cm}^{-2}, N_A = 10^{17} \text{ cm}^{-3}$$

$$d_1 = 0.8\text{-}2.0 \text{ nm and } d_2 = 0.6\text{-}3.4 \text{ nm}$$

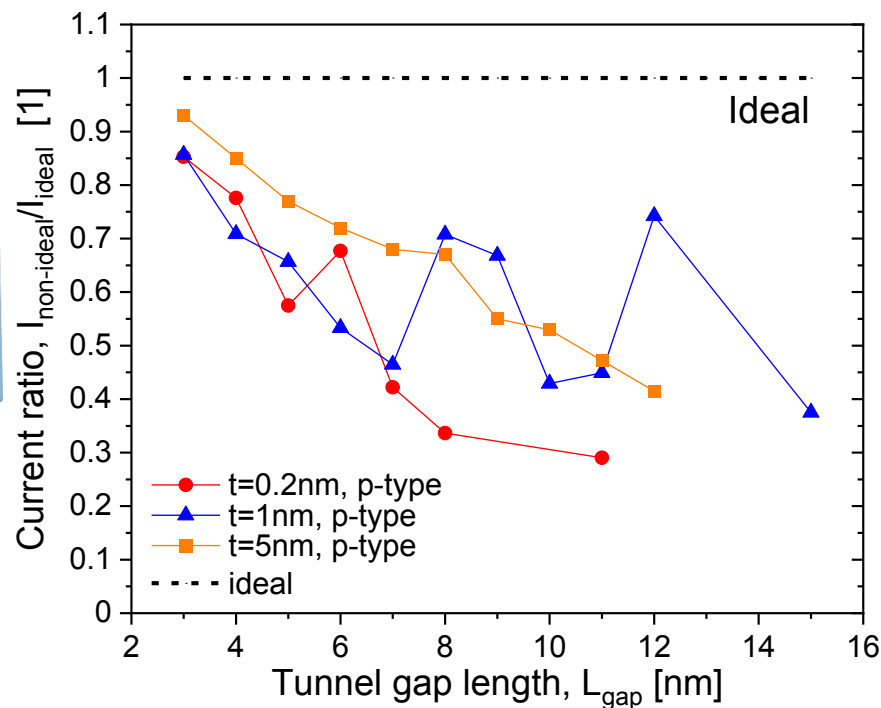


Effects of a Single Impurity in the Tunnel Gap

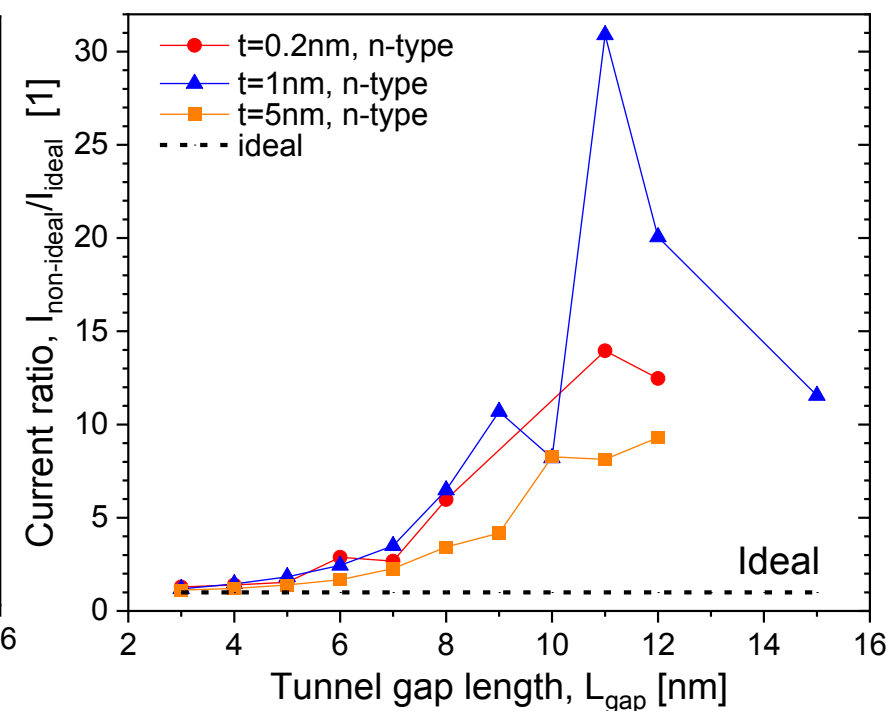
- While p-type impurities might moderately reduce the tunneling rate, n-type impurities might dramatically increase the tunneling rate, specially for tunnel gaps of the order of 10 nm



p-type impurity

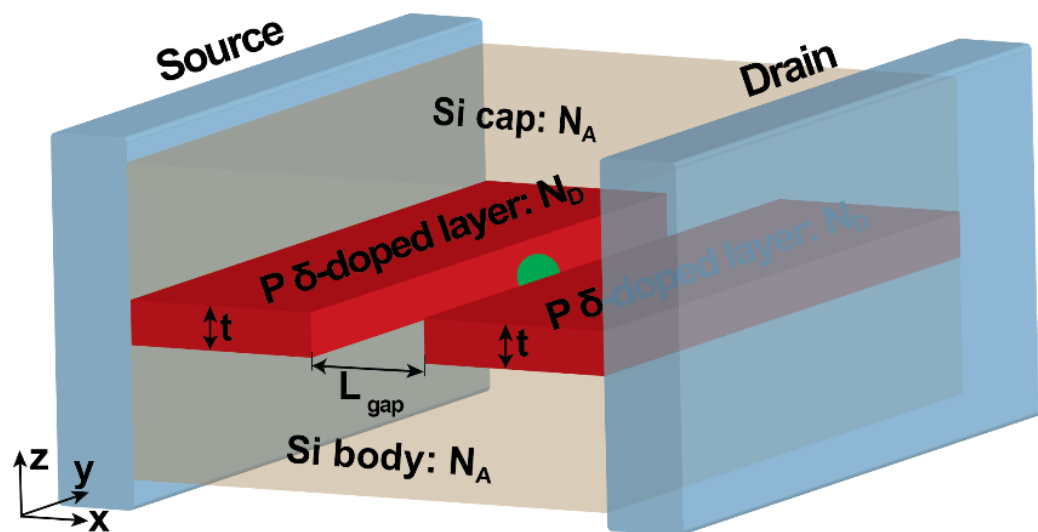


n-type impurity

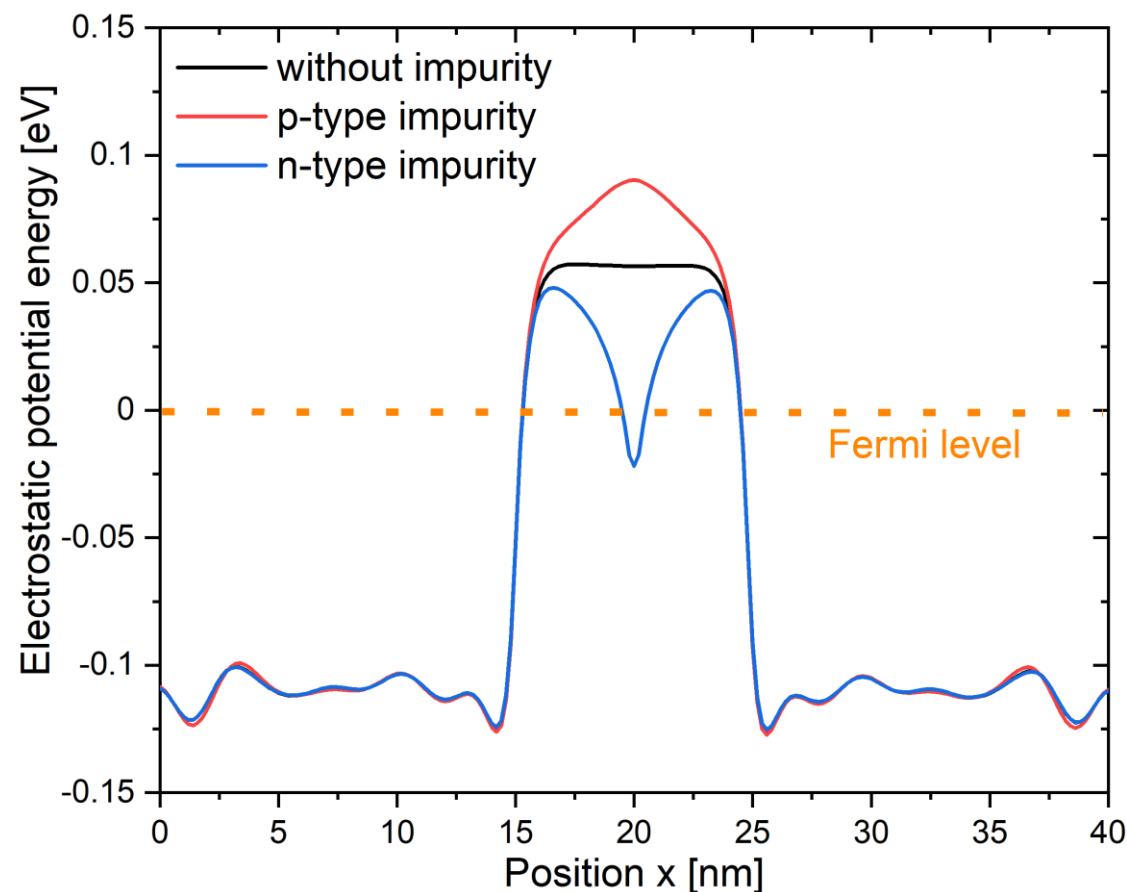


Effects of a Single Impurity in the Tunnel Gap

- While p-type impurities might moderately reduce the tunneling rate, n-type impurities might dramatically increase the tunneling rate, specially for tunnel gaps of the order of 10 nm



$$N_D = 1.0 \times 10^{14} \text{ cm}^{-2}, N_A = 10^{17} \text{ cm}^{-3}$$



Summary

- ❑ Presented an **efficient quantum open-system transport framework** for δ -layer tunnel junctions
- ❑ While most of the non-idealities moderately affect the tunneling rate, a **single charged impurity** in the tunnel gap can alter the tunneling rate by more than an **order of magnitude**, even for relatively large tunnel gaps
- ❑ The electric **sign of impurity** plays an important role in the tunneling rate: the change of current due to an n-type impurity is an order of magnitude stronger than for p-type impurity
- ❑ Overall these simulations suggest that **geometric fidelity of the device fabrication is less important than mitigation of defects inside of the junction**



Quantum Transport Simulations for Si: P δ -layer Tunnel Junctions

Juan P. Mendez, Denis Mamaluy, Xujiao (Suzey) Gao and Shashank Misra
jpmende@sandia.gov

Sandia National Laboratories, Albuquerque, New Mexico



Sandia National Laboratories is a multimission laboratory managed and operated by National Technology and Engineering Solutions of Sandia, LLC., a wholly owned subsidiary of Honeywell International, Inc., for the U.S. Department of Energy's National Nuclear Security Administration under contract DE-NA-0003525

