

This paper describes objective technical results and analysis. Any subjective views or opinions that might be expressed in the paper do not necessarily represent the views of the U.S. Department of Energy or the United States Government.

SAND2020-9964C

# Cost of Classical Strong Simulation of the T-Gate Magic State

**Lucas Kocia**  
and Mohan Sarovar

<sup>1</sup>Sandia, Livermore

**OVER-QC**  
<https://overqc.sandia.gov>



Sandia  
National  
Laboratories

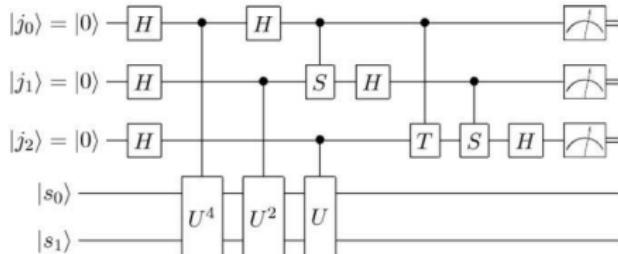


National Nuclear Security Administration



Sandia National Laboratories is a multimission laboratory managed and operated by National Technology & Engineering Solutions of Sandia, LLC, a wholly owned subsidiary of Honeywell International Inc., for the U.S. Department of Energy's National Nuclear Security Administration under contract DE-NA0003525.

# Tensoring Magic State Gives You Universality

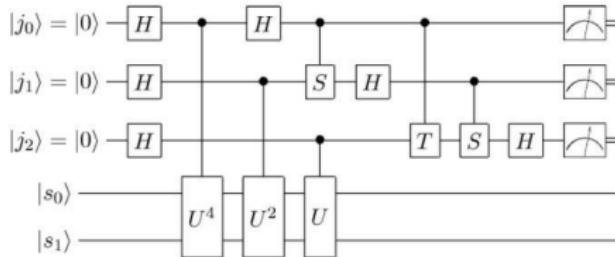


- any universal quantum computation can be written in terms of
  - $k$ -tensored T-gate magic states  $|T\rangle^{\otimes k}$
  - then acted on by Clifford unitaries  $U_C$
  - then partially traced over to obtain a marginal over any qudit

$$P = \text{Tr} \left[ \langle 0 | \hat{U}_C | T \rangle^{\otimes k} | 0 \rangle^{\otimes (n-k)} \right].$$

Depth of simulatable circuit scales with  $k$ .

# Tensoring Magic State Gives You Universality



- any universal quantum computation can be written in terms of
  - $k$ -tensored T-gate magic states  $|T\rangle^{\otimes k}$
  - then acted on by Clifford unitaries  $U_C$
  - then partially traced over to obtain a marginal over any qudit

$$P = \text{Tr} \left[ \langle 0 | \hat{U}_C | T \rangle^{\otimes k} | 0 \rangle^{\otimes (n-k)} \right].$$

Depth of simulatable circuit scales with  $k$ .

Previous methods have relied on Monte Carlo numerics and stop converging at  $k > 7$ .

We show how to get to  $k > 14$  using an algebraic method.

# Stabilizer States as Cost Metric

- Stabilizer states  $\{|\phi_i\rangle\}_i$  form an overcomplete basis.

## Stabilizer States as Cost Metric

- Stabilizer states  $\{|\phi_i\rangle\}_i$  form an overcomplete basis.
- Therefore, any state  $|\Psi\rangle$  can be expressed as  $|\Psi\rangle = \sum_i c_i |\phi_i\rangle_i$ .

## Stabilizer States as Cost Metric

- Stabilizer states  $\{|\phi_i\rangle\}_i$  form an overcomplete basis.
- Therefore, any state  $|\Psi\rangle$  can be expressed as  $|\Psi\rangle = \sum_i c_i |\phi_i\rangle_i$ .

$$|T\rangle = \frac{1}{\sqrt{2}}(|0\rangle + e^{\pi i/4}|1\rangle).$$

The T-gate magic state extends the Clifford subtheory to universality.

# Stabilizer States as Cost Metric

- Stabilizer states  $\{|\phi_i\rangle\}_i$  form an overcomplete basis.
- Therefore, any state  $|\Psi\rangle$  can be expressed as  $|\Psi\rangle = \sum_i^\chi c_i |\phi_i\rangle_i$ .

$$|T\rangle = \frac{1}{\sqrt{2}}(|0\rangle + e^{\pi i/4}|1\rangle).$$

The T-gate magic state extends the Clifford subtheory to universality.

It has been postulated that  $\chi(T)$  grows slowest with increasing number of qubits

# Computational Cost

$$P = \text{Tr} \left[ \langle 0 | \hat{U}_C | T \rangle^{\otimes k} | 0 \rangle^{\otimes (n-k)} \right].$$

## Remark: Trivial Tensor Bound

Let  $\chi_k$  be the stabilizer rank of  $|\Psi\rangle^k$ . Since the tensor product of two stabilizer states is a stabilizer state, it follows that  $\chi_t \leq \chi_k^{t/k}$ .

# Computational Cost

$$P = \text{Tr} \left[ \langle 0 | \hat{U}_C | T \rangle^{\otimes k} | 0 \rangle^{\otimes (n-k)} \right].$$

## Remark: Trivial Tensor Bound

Let  $\chi_k$  be the stabilizer rank of  $|\Psi\rangle^k$ . Since the tensor product of two stabilizer states is a stabilizer state, it follows that  $\chi_t \leq \chi_k^{t/k}$ .

[ Bravyi *et al.* PRL 116 (2016): 250501, & Quantum 3 (2019): 181 ]

Numerical Monte Carlo findings:

| qubit<br>$k$   | 1     | 2          | 3           | 4 | 5 | 6           | 7  |
|----------------|-------|------------|-------------|---|---|-------------|----|
| $\chi_k$       | 2     | 2          | 3           | 4 | 6 | 7           | 12 |
| $\chi_k^{t/k}$ | $2^t$ | $2^{0.5t}$ | $2^{0.53t}$ |   |   | $2^{0.47t}$ |    |

# Computational Cost

$$P = \text{Tr} \left[ \langle 0 | \hat{U}_C | T \rangle^{\otimes k} | 0 \rangle^{\otimes (n-k)} \right].$$

## Remark: Trivial Tensor Bound

Let  $\chi_k$  be the stabilizer rank of  $|\Psi\rangle^k$ . Since the tensor product of two stabilizer states is a stabilizer state, it follows that  $\chi_t \leq \chi_k^{t/k}$ .

[ Bravyi *et al.* PRL 116 (2016): 250501, & Quantum 3 (2019): 181 ]

Numerical Monte Carlo findings:

| qubit<br>$k$   | 1     | 2          | 3           | 4 | 5 | 6           | 7  |
|----------------|-------|------------|-------------|---|---|-------------|----|
| $\chi_k$       | 2     | 2          | 3           | 4 | 6 | 7           | 12 |
| $\chi_k^{t/k}$ | $2^t$ | $2^{0.5t}$ | $2^{0.53t}$ |   |   | $2^{0.47t}$ |    |

Can we push this further?! Not with numerical techniques.

Really want to know if can do better than trivial tensor bound for large  $t$ . This is why  $\chi_6$  is a big deal.

# The Wigner Discrete Propagator

Instead of considering our magic state in terms of vectors in Hilbert space, we can use a kernel (or quasi-probability) representation instead;

# The Wigner Discrete Propagator

Instead of considering our magic state in terms of vectors in Hilbert space, we can use a kernel (or quasi-probability) representation instead;

Given a set of operators  $R(\mathbf{x})$ , indexed by  $\mathbf{x} \in (\mathbb{Z}/p\mathbb{Z})^n$ , that are Hilbert-Schmidt orthogonal, any operator can be represented as

$$\hat{A} = d^{-1} \sum_{\mathbf{x}} \text{Tr}(\hat{R}(\mathbf{x}) \hat{A}) \hat{R}(\mathbf{x}) \equiv \sum_{\mathbf{x}} A(\mathbf{x}) \hat{R}(\mathbf{x}).$$

# Computation

$$\begin{aligned} P &= \text{Tr} \left[ \langle 0 | \hat{U}_C | T \rangle^{\otimes k} | 0 \rangle^{\otimes (n-k)} \right] \\ P &\stackrel{\text{WWM}}{=} \sum_{\mathbf{x}' \in D} \left[ \prod_{i=1}^t \rho_T(\mathbf{x}'_i) \prod_{j=t+1}^n \delta(x'_{q_j}) \right]. \end{aligned}$$

for

$$D = \left\{ \mathbf{x}' \left| \left( \mathcal{M}^{-1} \left( \mathbf{x}' - \frac{\alpha}{2} \right) - \frac{\alpha}{2} \right)_{n+1} \bmod 3^2 = 0 \right. \right\}$$

Basic computational primitive is quadratic Gauss sums. Each one costs  $\mathcal{O}(n^3)$ . We want to know how few we need.

# Comparison

| Hilbert space                                       | WWM  |
|---|--|
| Stab state  | classical probability distribution             |
| Clifford  | symplectic transformation                      |
| stabilizer inner product<br>cost $\mathcal{O}(n^3)$ | quadratic Gauss sum<br>cost $\mathcal{O}(n^3)$ |

# Comparison

| Hilbert space                                       | WWM  |
|---|--|
| Stab state  | classical probability distribution             |
| Clifford  | symplectic transformation                      |
| stabilizer inner product<br>cost $\mathcal{O}(n^3)$ | quadratic Gauss sum<br>cost $\mathcal{O}(n^3)$ |
| stabilizer rank $\chi_t$                            | quadratic Gauss sum rank $\xi_t$               |

# Comparison

| Hilbert space                                       | WWM  |
|---|--|
| Stab state  | classical probability distribution             |
| Clifford  | symplectic transformation                      |
| stabilizer inner product<br>cost $\mathcal{O}(n^3)$ | quadratic Gauss sum<br>cost $\mathcal{O}(n^3)$ |
| stabilizer rank $\chi_t$                            | quadratic Gauss sum rank $\xi_t$               |

WWM formalism is simpler for odd-prime-dimensional qudits and so here we focus on qutrits ( $d = 3$ )

# Using WWM to Bound Stab Rank

[ Kocia & Love, arXiv:1810.03622 ] Prior Algebraic Findings:

| $k$           | 1     | 2          |
|---------------|-------|------------|
| qutrit        |       |            |
| $\chi_k$      | 3     | 3          |
| $\xi_k$       | 3     | 3          |
| $\xi_k^{t/k}$ | $3^t$ | $3^{0.5t}$ |

# Using WWM to Bound Stab Rank

[ Kocia & Love, arXiv:1810.03622 ] Prior Algebraic Findings:

| $k$            | 1     | 2          |
|----------------|-------|------------|
| qutrit         |       |            |
| $\chi_k$       | 3     | 3          |
| $\xi_k$        | 3     | 3          |
| $\xi_k^{t/k}$  | $3^t$ | $3^{0.5t}$ |
| qubit          |       |            |
| $\chi_k^{t/k}$ | $2^t$ | $2^{0.5t}$ |

# Pushing Forward

Three Qutrit T-Gate Magic State:

$$\begin{aligned}\rho(\mathcal{M}_{C_{1,3}^2 C_{2,3}} \mathbf{x}) &= 3 \sum_{\substack{y_{q_1}, y_{q_2} \\ \in \mathbb{Z}/3\mathbb{Z}}} \exp \left[ \frac{2\pi i}{9} (7y_{q_1}^3 + 8x_{q_1}^3) \right] \mathcal{A}_3(y_{q_1}, y_{q_2}, \mathbf{x}) \\ &\quad \times [\delta(\neg(y_{q_1} - x_{q_1})) + \delta(y_{q_1} - x_{q_1})\delta(\Delta)],\end{aligned}$$

$y_{q_1}$  and  $y_{q_2}$  index 9 quadratic Gauss sums.

# Pushing Forward

Three Qutrit T-Gate Magic State:

$$\begin{aligned}\rho(\mathcal{M}_{C_{1,3}^2 C_{2,3}} \mathbf{x}) &= 3 \sum_{\substack{y_{q_1}, y_{q_2} \\ \in \mathbb{Z}/3\mathbb{Z}}} \exp \left[ \frac{2\pi i}{9} (7y_{q_1}^3 + 8x_{q_1}^3) \right] \mathcal{A}_3(y_{q_1}, y_{q_2}, \mathbf{x}) \\ &\quad \times [\delta(\neg(y_{q_1} - x_{q_1})) + \delta(y_{q_1} - x_{q_1})\delta(\Delta)],\end{aligned}$$

$y_{q_1}$  and  $y_{q_2}$  index 9 quadratic Gauss sums.

However, due to the additional delta functions, the Wigner function of three tensored qutrit magic states can be expressed in terms of only 8 non-zero quadratic Gauss sums

# Pushing Forward

Three Qutrit T-Gate Magic State:

$$\begin{aligned}\rho(\mathcal{M}_{C_{1,3}^2 C_{2,3}} \mathbf{x}) &= 3 \sum_{\substack{y_{q_1}, y_{q_2} \\ \in \mathbb{Z}/3\mathbb{Z}}} \exp \left[ \frac{2\pi i}{9} (7y_{q_1}^3 + 8x_{q_1}^3) \right] \mathcal{A}_3(y_{q_1}, y_{q_2}, \mathbf{x}) \\ &\quad \times [\delta(\neg(y_{q_1} - x_{q_1})) + \delta(y_{q_1} - x_{q_1})\delta(\Delta)],\end{aligned}$$

$y_{q_1}$  and  $y_{q_2}$  index 9 quadratic Gauss sums.

However, due to the additional delta functions, the Wigner function of three tensored qutrit magic states can be expressed in terms of only 8 non-zero quadratic Gauss sums

Extrapolating to higher  $t$  counts using the trivial tensor bound, this result shows that  $3^{\frac{\log 8}{3 \log 3} t} \approx 3^{0.63t}$  quadratic Gauss sums can represent  $t$  magic states, for  $t$  a multiple of 3.

# Pushing Forward Summary

New algebraic findings:

| $k$    | 1 | 2 | 3  | 4 | 5 | 6 | 7                           |
|--------|---|---|----|---|---|---|-----------------------------|
| qutrit | 3 | 3 | 8? |   |   |   | <i>inaccessible to M.C.</i> |

# Pushing Forward Summary

New algebraic findings:

| $k$      | 1 | 2 | 3  | 4 | 5                           | 6  | 7         |
|----------|---|---|----|---|-----------------------------|----|-----------|
| qutrit   |   |   |    |   |                             |    |           |
| $\chi_k$ | 3 | 3 | 8? |   | <i>inaccessible to M.C.</i> |    |           |
| $\xi_k$  | 3 | 3 | 8  | 9 | 21                          | 23 | $\leq 63$ |

# Pushing Forward Summary

New algebraic findings:

| $k$           | 1     | 2          | 3            | 4 | 5                           | 6            | 7         |
|---------------|-------|------------|--------------|---|-----------------------------|--------------|-----------|
| qutrit        |       |            |              |   |                             |              |           |
| $\chi_k$      | 3     | 3          | 8?           |   | <i>inaccessible to M.C.</i> |              |           |
| $\xi_k$       | 3     | 3          | 8            | 9 | 21                          | 23           | $\leq 63$ |
| $\xi_k^{t/k}$ | $3^t$ | $3^{0.5t}$ | $3^{0.631t}$ |   | $3^{0.554}$                 | $3^{0.476t}$ |           |

# Pushing Forward Summary

New algebraic findings:

| $k$            | 1     | 2          | 3            | 4 | 5                           | 6            | 7         |
|----------------|-------|------------|--------------|---|-----------------------------|--------------|-----------|
| qutrit         |       |            |              |   |                             |              |           |
| $\chi_k$       | 3     | 3          | 8?           |   | <i>inaccessible to M.C.</i> |              |           |
| $\xi_k$        | 3     | 3          | 8            | 9 | 21                          | 23           | $\leq 63$ |
| $\xi_k^{t/k}$  | $3^t$ | $3^{0.5t}$ | $3^{0.631t}$ |   | $3^{0.554}$                 | $3^{0.476t}$ |           |
| qubit          |       |            |              |   |                             |              |           |
| $\chi_k^{t/k}$ | $2^t$ | $2^{0.5t}$ | $2^{0.528t}$ |   |                             | $2^{0.468t}$ |           |

# Pushing Forward Summary

New algebraic findings:

| $k$            | 1     | 2          | 3            | 4 | 5                           | 6            | 7         |
|----------------|-------|------------|--------------|---|-----------------------------|--------------|-----------|
| qutrit         |       |            |              |   |                             |              |           |
| $\chi_k$       | 3     | 3          | 8?           |   | <i>inaccessible to M.C.</i> |              |           |
| $\xi_k$        | 3     | 3          | 8            | 9 | 21                          | 23           | $\leq 63$ |
| $\xi_k^{t/k}$  | $3^t$ | $3^{0.5t}$ | $3^{0.631t}$ |   | $3^{0.554}$                 | $3^{0.476t}$ |           |
| qubit          |       |            |              |   |                             |              |           |
| $\chi_k^{t/k}$ | $2^t$ | $2^{0.5t}$ | $2^{0.528t}$ |   |                             | $2^{0.468t}$ |           |

| $k$      | 8 | 9                                  | 10 | 11 | 12 | 13 | 14 |
|----------|---|------------------------------------|----|----|----|----|----|
| qutrit   |   |                                    |    |    |    |    |    |
| $\chi_k$ |   | <i>inaccessible to Monte Carlo</i> |    |    |    |    |    |

# Pushing Forward Summary

New algebraic findings:

| $k$            | 1     | 2          | 3            | 4 | 5                           | 6            | 7         |
|----------------|-------|------------|--------------|---|-----------------------------|--------------|-----------|
| qutrit         |       |            |              |   |                             |              |           |
| $\chi_k$       | 3     | 3          | 8?           |   | <i>inaccessible to M.C.</i> |              |           |
| $\xi_k$        | 3     | 3          | 8            | 9 | 21                          | 23           | $\leq 63$ |
| $\xi_k^{t/k}$  | $3^t$ | $3^{0.5t}$ | $3^{0.631t}$ |   | $3^{0.554}$                 | $3^{0.476t}$ |           |
| qubit          |       |            |              |   |                             |              |           |
| $\chi_k^{t/k}$ | $2^t$ | $2^{0.5t}$ | $2^{0.528t}$ |   |                             | $2^{0.468t}$ |           |

| $k$      | 8  | 9          | 10  | 11                                 | 12  | 13       | 14   |
|----------|----|------------|-----|------------------------------------|-----|----------|------|
| qutrit   |    |            |     |                                    |     |          |      |
| $\chi_k$ |    |            |     | <i>inaccessible to Monte Carlo</i> |     |          |      |
| $\xi_k$  | 69 | $\leq 189$ | 207 | $< 513$                            | 513 | $< 1539$ | 1539 |

# Pushing Forward Summary

New algebraic findings:

| $k$            | 1     | 2          | 3            | 4 | 5                           | 6            | 7         |
|----------------|-------|------------|--------------|---|-----------------------------|--------------|-----------|
| qutrit         |       |            |              |   |                             |              |           |
| $\chi_k$       | 3     | 3          | 8?           |   | <i>inaccessible to M.C.</i> |              |           |
| $\xi_k$        | 3     | 3          | 8            | 9 | 21                          | 23           | $\leq 63$ |
| $\xi_k^{t/k}$  | $3^t$ | $3^{0.5t}$ | $3^{0.631t}$ |   | $3^{0.554}$                 | $3^{0.476t}$ |           |
| qubit          |       |            |              |   |                             |              |           |
| $\chi_k^{t/k}$ | $2^t$ | $2^{0.5t}$ | $2^{0.528t}$ |   |                             | $2^{0.468t}$ |           |

| $k$           | 8  | 9          | 10  | 11                                 | 12             | 13       | 14   |
|---------------|----|------------|-----|------------------------------------|----------------|----------|------|
| qutrit        |    |            |     |                                    |                |          |      |
| $\chi_k$      |    |            |     | <i>inaccessible to Monte Carlo</i> |                |          |      |
| $\xi_k$       | 69 | $\leq 189$ | 207 | $< 513$                            | 513            | $< 1539$ | 1539 |
| $\xi_k^{t/k}$ |    |            |     | $< 3^{0.528t}$                     | $< 3^{0.473t}$ |          |      |

# Pushing Forward Summary

New algebraic findings:

| $k$            | 1     | 2          | 3            | 4 | 5                           | 6            | 7         |
|----------------|-------|------------|--------------|---|-----------------------------|--------------|-----------|
| qutrit         |       |            |              |   |                             |              |           |
| $\chi_k$       | 3     | 3          | 8?           |   | <i>inaccessible to M.C.</i> |              |           |
| $\xi_k$        | 3     | 3          | 8            | 9 | 21                          | 23           | $\leq 63$ |
| $\xi_k^{t/k}$  | $3^t$ | $3^{0.5t}$ | $3^{0.631t}$ |   | $3^{0.554}$                 | $3^{0.476t}$ |           |
| qubit          |       |            |              |   |                             |              |           |
| $\chi_k^{t/k}$ | $2^t$ | $2^{0.5t}$ | $2^{0.528t}$ |   |                             | $2^{0.468t}$ |           |

| $k$            | 8  | 9          | 10  | 11                                 | 12             | 13       | 14   |
|----------------|----|------------|-----|------------------------------------|----------------|----------|------|
| qutrit         |    |            |     |                                    |                |          |      |
| $\chi_k$       |    |            |     | <i>inaccessible to Monte Carlo</i> |                |          |      |
| $\xi_k$        | 69 | $\leq 189$ | 207 | $< 513$                            | 513            | $< 1539$ | 1539 |
| $\xi_k^{t/k}$  |    |            |     | $< 3^{0.528t}$                     | $< 3^{0.473t}$ |          |      |
| qubit          |    |            |     |                                    |                |          |      |
| $\chi_k^{t/k}$ |    |            |     | <i>inaccessible to Monte Carlo</i> |                |          |      |

Monte Carlo numeric search for qubit stops converging at  $t = 7$  when the stabilizer state space consists of

$$8.3 \times 10^{130}$$

possible states for a stabilizer rank of 12.

Monte Carlo numeric search for qubit stops converging at  $t = 7$  when the stabilizer state space consists of

$$8.3 \times 10^{130}$$

possible states for a stabilizer rank of 12.

Assuming a similar reduction occurs for qubits at  $t = 12$  as for qutrits, the WWM method is able to algebraically search a space that consists of

$$> 8.3 \times 10^{13000}$$

possible stabilizer states!

## Conclusion

- Converting these results into the discrete setting would likely help validate and simulate NISQ devices.

## Conclusion

- Converting these results into the discrete setting would likely help validate and simulate NISQ devices.
- To show this, we extended of the Wigner-Weyl Moyal (WWM) formalism for discrete odd prime dimensions to higher order  $\hbar$  corrections through uniformization.

## Conclusion

- Converting these results into the discrete setting would likely help validate and simulate NISQ devices.
- To show this, we extended of the Wigner-Weyl Moyal (WWM) formalism for discrete odd prime dimensions to higher order  $\hbar$  corrections through uniformization.
- Produces bounds that follow the stabilizer rank of qubits (after conversion from base  $d$  to base 2).

## Conclusion

- Converting these results into the discrete setting would likely help validate and simulate NISQ devices.
- To show this, we extended of the Wigner-Weyl Moyal (WWM) formalism for discrete odd prime dimensions to higher order  $\hbar$  corrections through uniformization.
- Produces bounds that follow the stabilizer rank of qubits (after conversion from base  $d$  to base 2).

Thanks!

*Sandia National Laboratories is a multimission laboratory managed and operated by National Technology & Engineering Solutions of Sandia, LLC, a wholly owned subsidiary of Honeywell International Inc., for the U.S. Department of Energy's National Nuclear Security Administration under contract DE-NA0003525.*