

Quantum transport in Si:P δ -layer wires

Juan P. Mendez, Denis Mamaluy, Xujiao (Suzey) Gao, Evan M. Anderson, DeAnna M. Campbell, Jeffrey A. Ivie, Tzu-Ming Lu, Scott W. Schmucker, and Shashank Misra

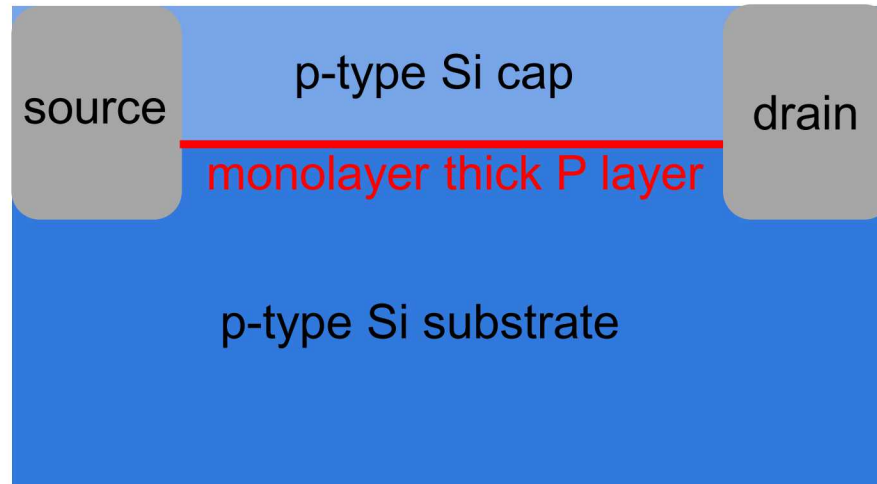
jpmende@sandia.gov

Sandia National Laboratories, Albuquerque, New Mexico

Sandia National Laboratories is a multimission laboratory managed and operated by National Technology and Engineering Solutions of Sandia, LLC., a wholly owned subsidiary of Honeywell International, Inc., for the U.S. Department of Energy's National Nuclear Security Administration under contract DE-NA-0003525



What is a Si: P δ -layer nanodevice?



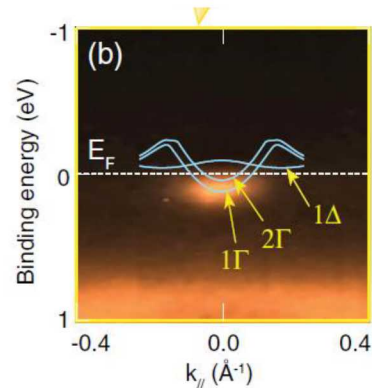
Schematic model of a Si: P δ -layer nanodevice

- High potential for quantum computing and advanced microelectronic devices
- The current advances to place dopants with atomic precision using hydrogen lithography

**Atomic Precision Advanced
Manufacturing (APAM)**

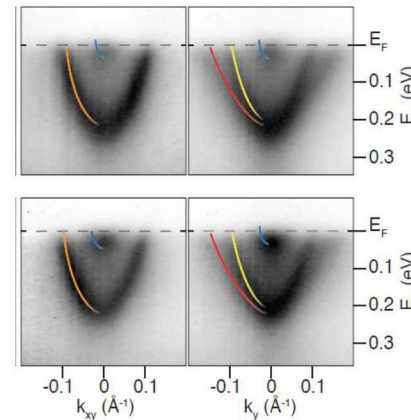
Conductive shallow sub-bands

J. A. Miwa, et al.,
Nano Letters 14, 1515 (2014)



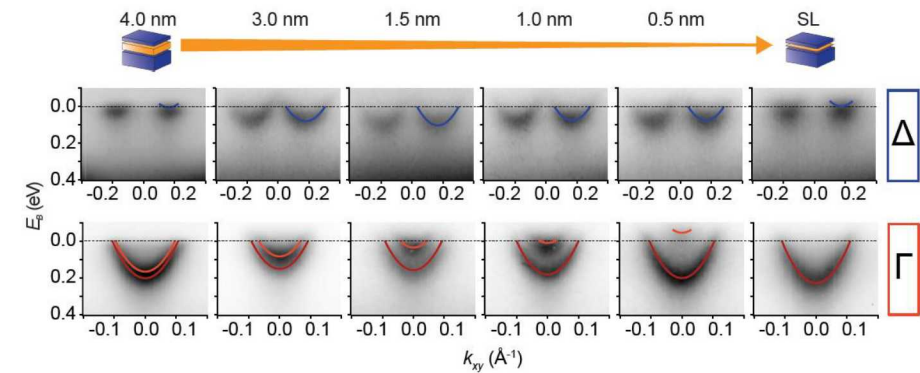
First observation of conductive shallow sub-bands (1Γ and 2Γ)

F. Mazzola, et al.,
npj Quantum Materials 5 (2020)

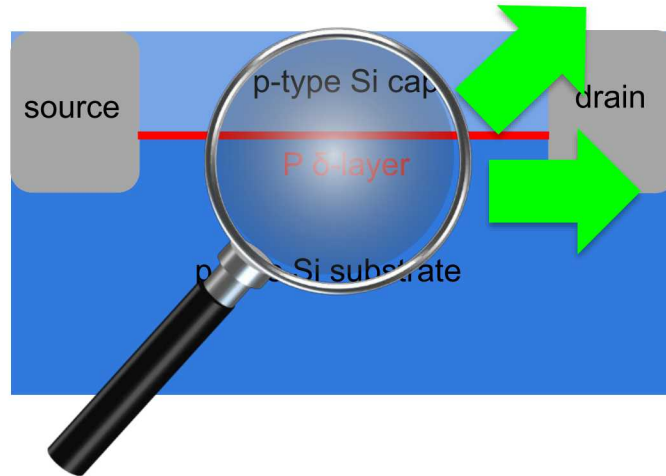


First observation of three sub-bands (1Γ, 2Γ and 3Γ)

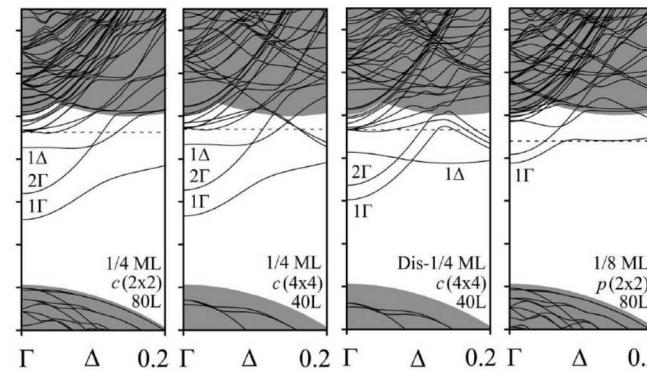
A. J. Holt, et al.,
Phys. Rev. B 101, 121402 (2020)



First observation of the Δ sub-band



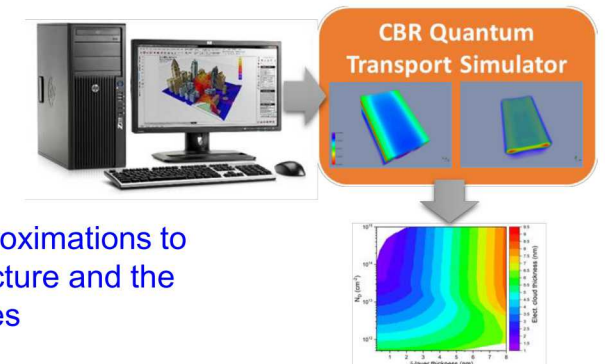
D. J. Carter, et al.,
Phys. Rev. B 79, 033204 (2009)



Study of the electronic structure of Si:P δ-layer systems using DFT

This Work

Open-system quantum transport treatment



- Closed system approximations to study the band structure and the conductive properties

Closed systems vs. Open systems

Closed system BC



- Extract the conductive properties from the additional (semi-)classical approximations (Drude or mobility-based)

$$J = e n_{2D} \mu E \sim n_{2D}$$



$$\sigma = \frac{J}{\mathcal{E}}$$

VS

Open system BC



- Extract the conductive properties from the quantum flux

$$j \sim \Psi \nabla \Psi^* - \Psi^* \nabla \Psi$$

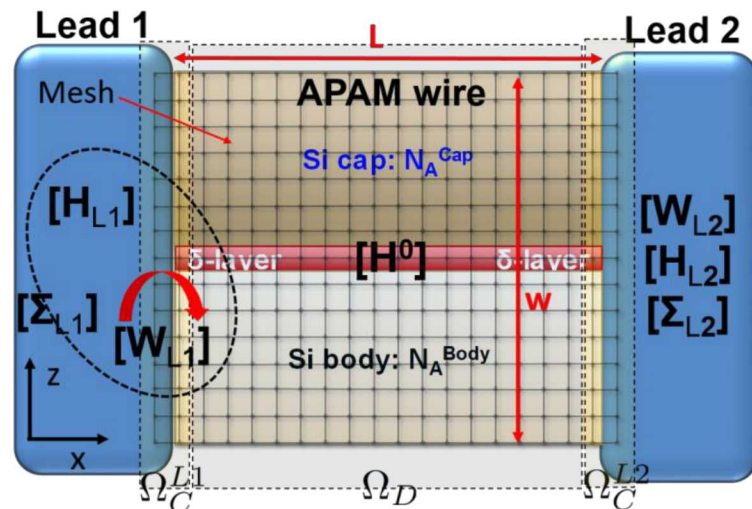


$$G = \sigma \frac{A}{L}$$

Open-System Quantum Transport

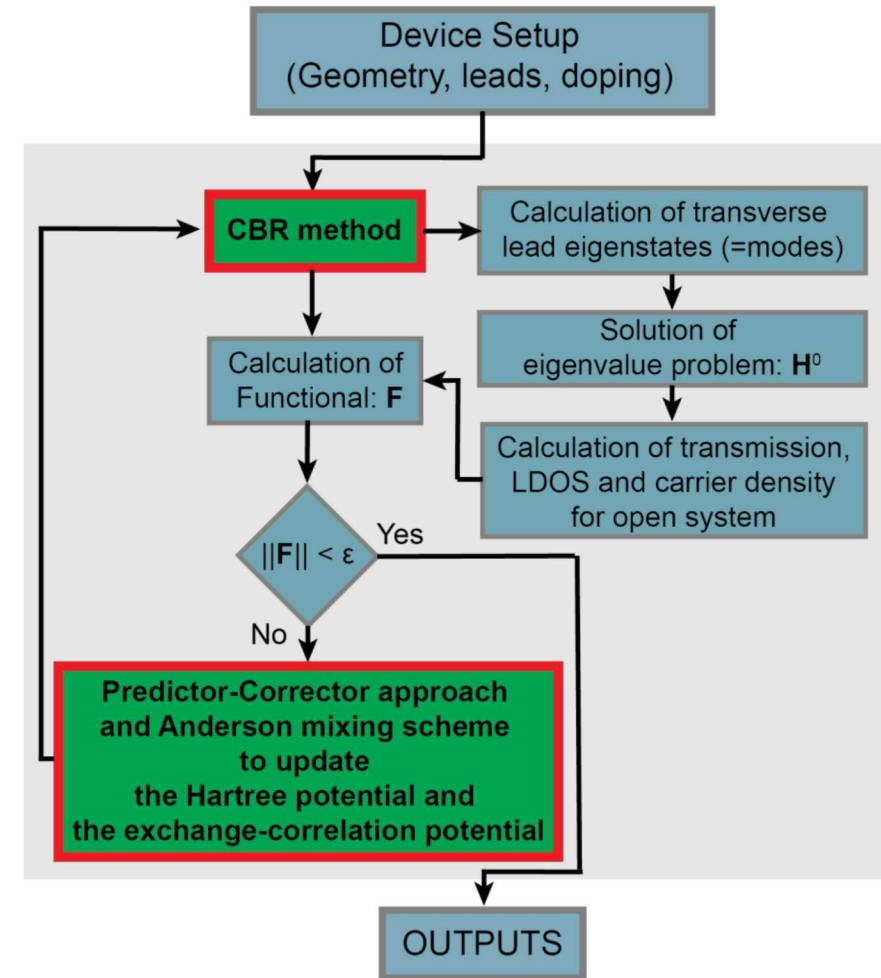
Our quantum transport simulator is based on **Non-Equilibrium Green's function** (NEGF) formalism with

- Fully charge self-consistent solution of Poisson-open system Schrödinger equation
- Single-band (Γ valley) effective mass approximation
- **Contact Block Reduction (CBR) method** ^[1,2] for fast numerical efficiency
- **Predictor-corrector approach** and **Anderson mixing scheme** ^[3]



Simulations were done for the cryogenic temperature of **4 K**.

Quantum transport simulator



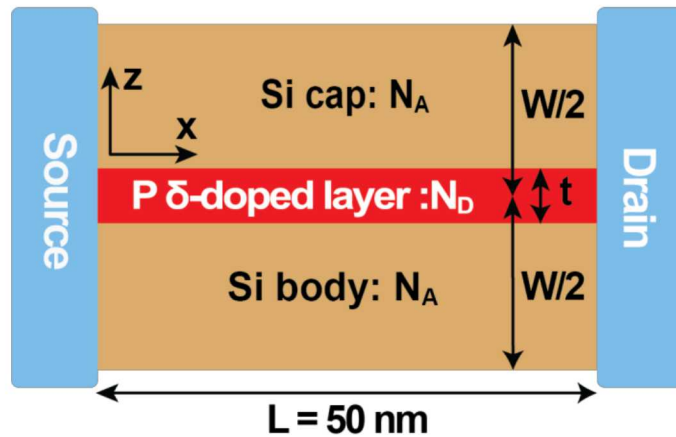
[1] D. Mamaluy *et al.*, J. Appl. Phys., vol. 93, no. 8, pp. 4628–4633, 2003.

[2] D. Mamaluy *et al.* Phys. Rev. B, vol. 71, p. 245321, 2005.

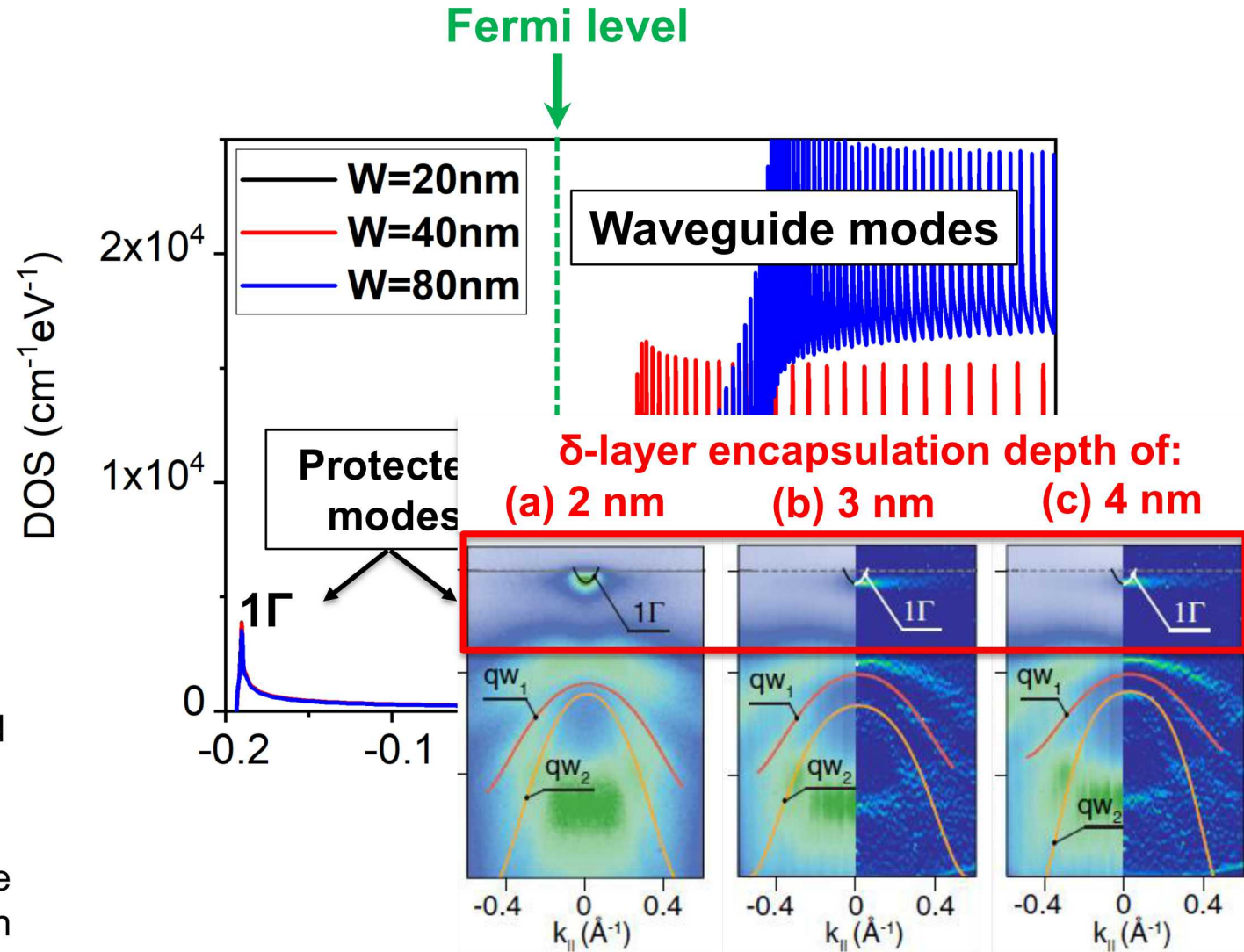
[3] X. Gao *et al.*, J. Appl. Phys., vol. 115, no. 13, p. 133707, 2014.

Prediction of Shallow Sub-bands

$$N_D=10^{14} \text{ cm}^{-2}, N_A=10^{17} \text{ cm}^{-3}, t=0.2 \text{ nm}$$



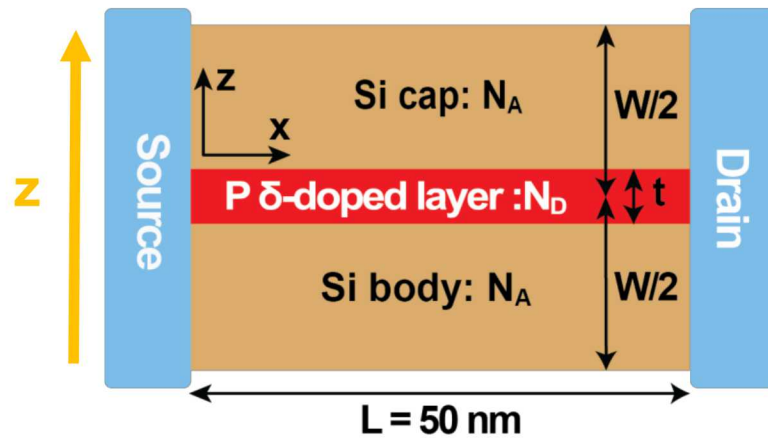
- Our open-system quantum simulations **predict shallow sub-bands**
- 1Γ sub-band is about 190 meV and 2Γ sub-band is about 20 meV below the Fermi level
- Below the Fermi energy level the states are “protected” (independent of encapsulation depth).



F. Mazzola, et al. Phys. Rev. Lett. 120, 046403, 2018

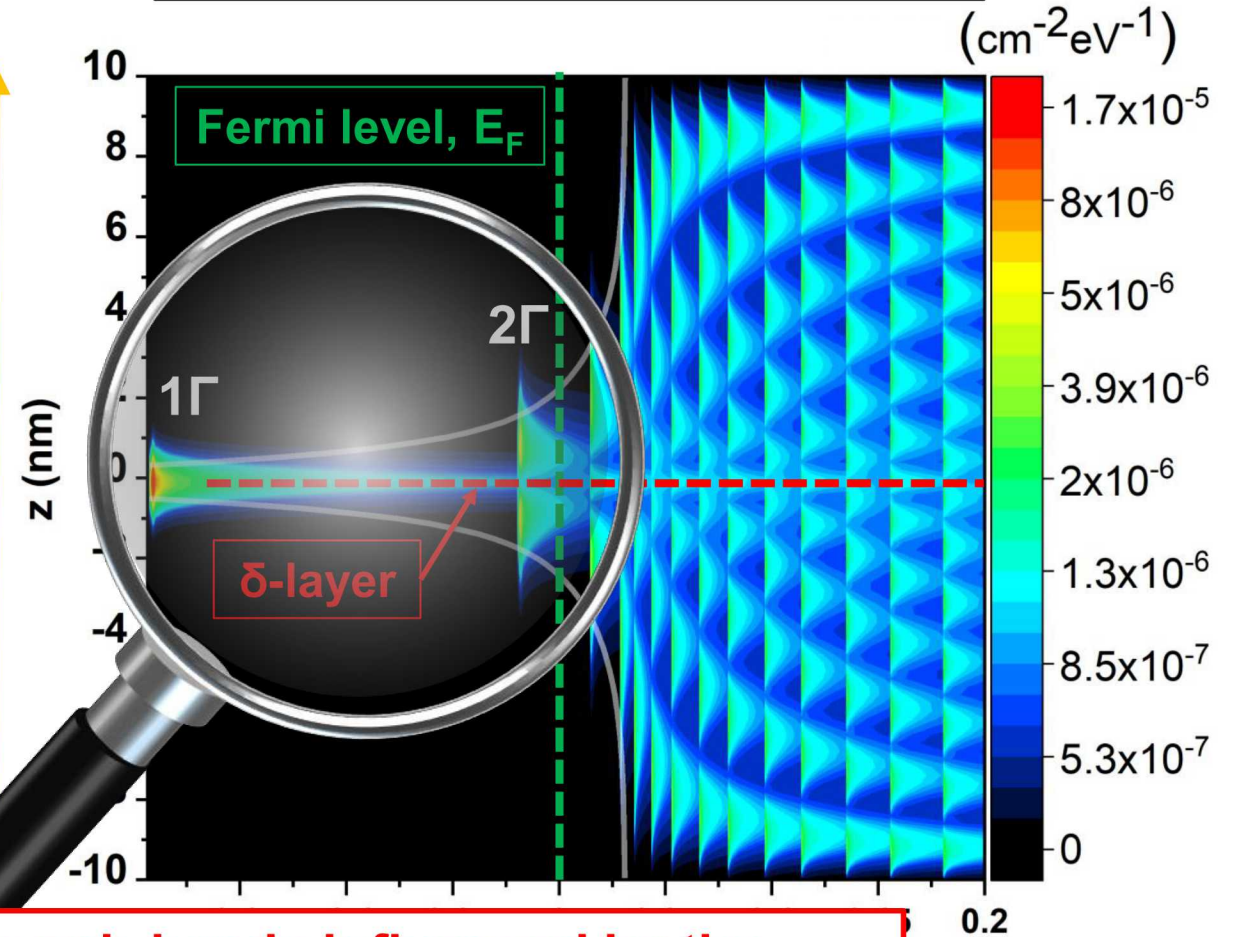
Si: P δ -layer conduction band structure

$N_D=1.2 \times 10^{14} \text{ cm}^{-2}$, $N_A=10^{17} \text{ cm}^{-3}$, $t=0.2 \text{ nm}$, $W=20 \text{ nm}$



- **Space quantization** of the modes: they are separated in space and with distinct structure

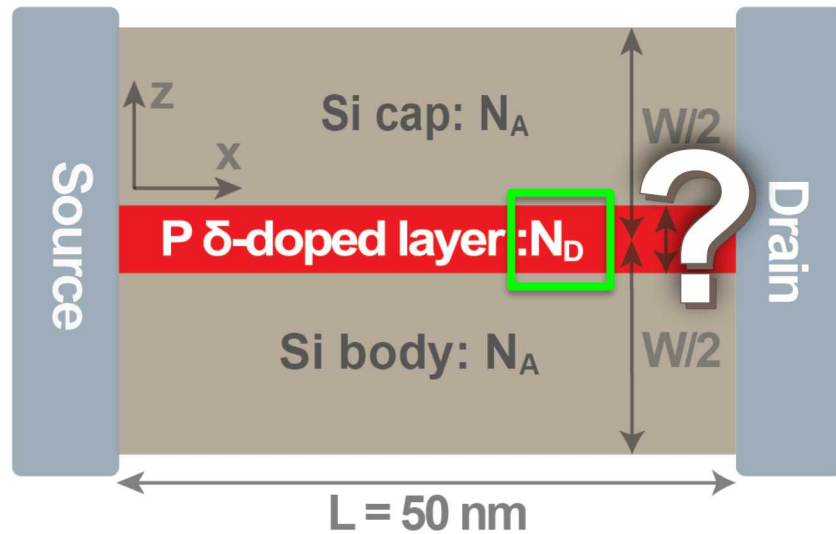
Local Density of States, LDOS(z, E)



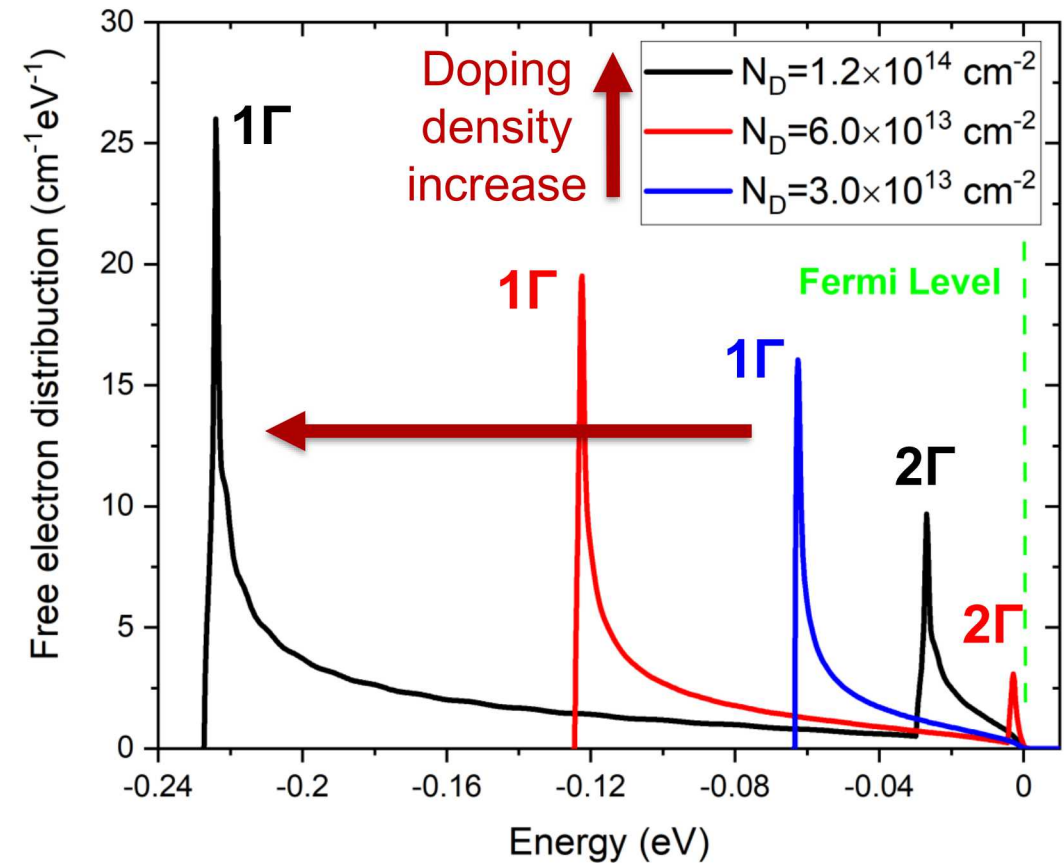
How are these conductive sub-bands influenced by the thickness and doping density of δ -layer?

Influence of δ -layer doping density

$N_A = 10^{17} \text{ cm}^{-3}$, $t = 0.2 \text{ nm}$, $W = 40 \text{ nm}$

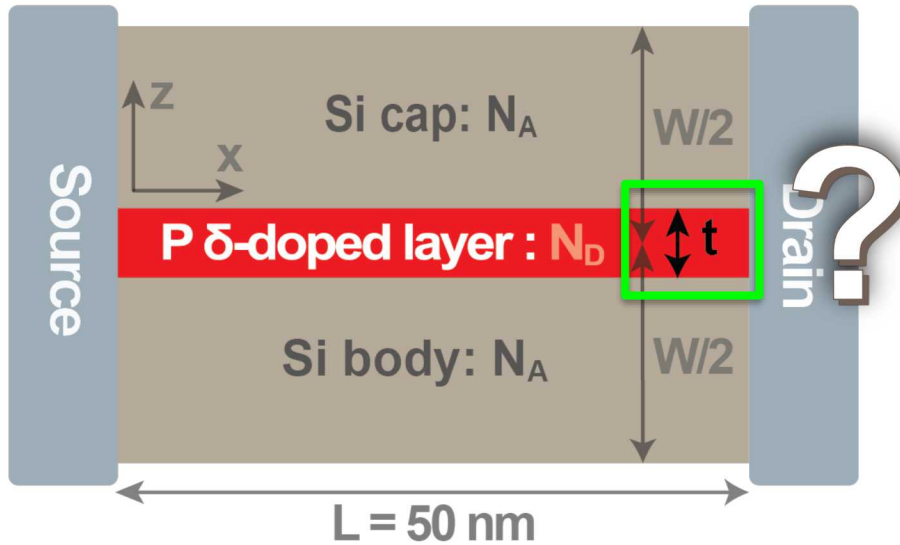


- For a fixed δ -layer thickness, the **increment of the sheet doping N_D increases the number of conducting modes**, as well as the **splitting energy** between them

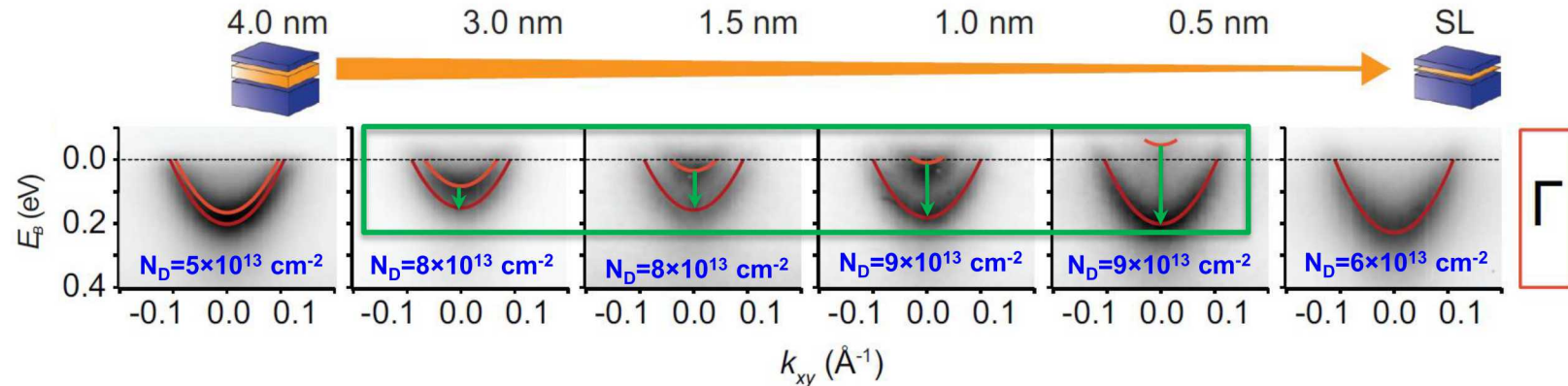
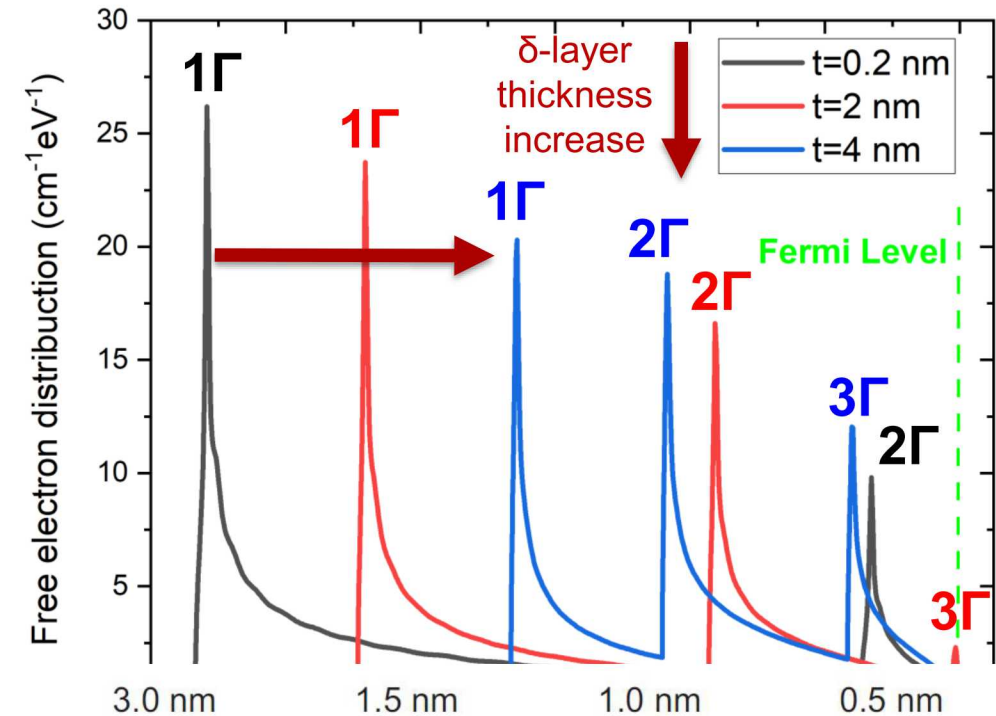


Influence of the δ -layer thickness

$$N_D = 1.2 \times 10^{14} \text{ cm}^{-2}, N_A = 10^{17} \text{ cm}^{-3}, W = 40 \text{ nm}$$



- For a fixed sheet doping, the increment of the delta-layer thickness increases the number of modes, but decreases the energy splitting between them



A. J. Holt *et al.*, Phys. Rev. B, vol. 101, p. 121402(R), 2020

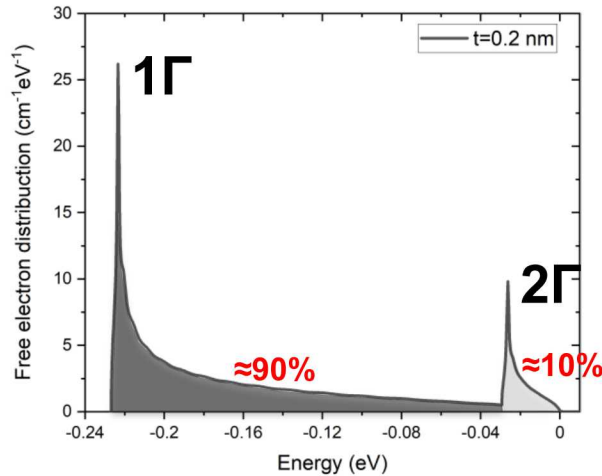
Electron contribution on the current

- Dissimilar contribution of the conductive sub-bands on the electronic current

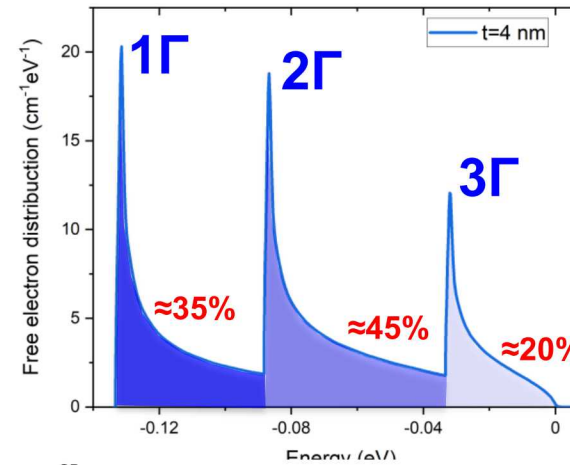
For strong confinement potentials =
Sharp doping profile

For weak confinement potentials =
Smooth doping profile

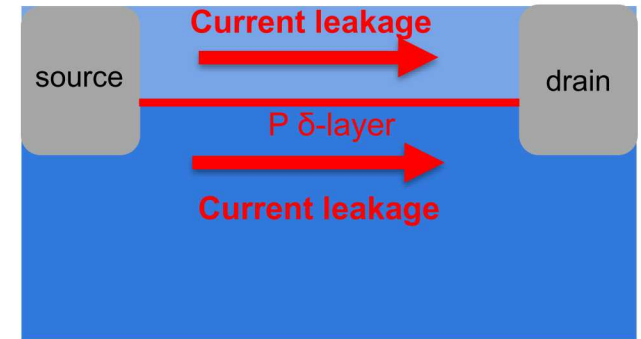
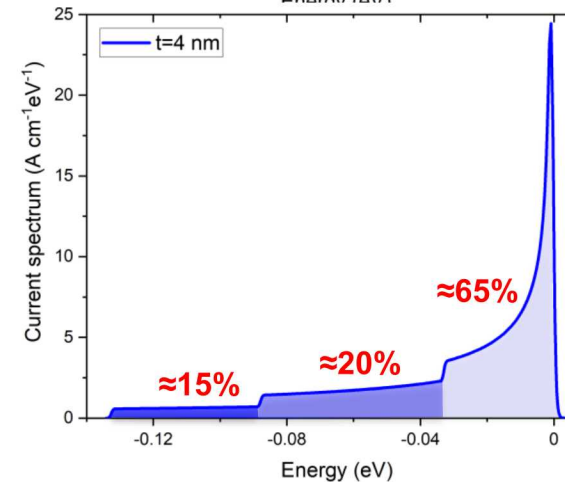
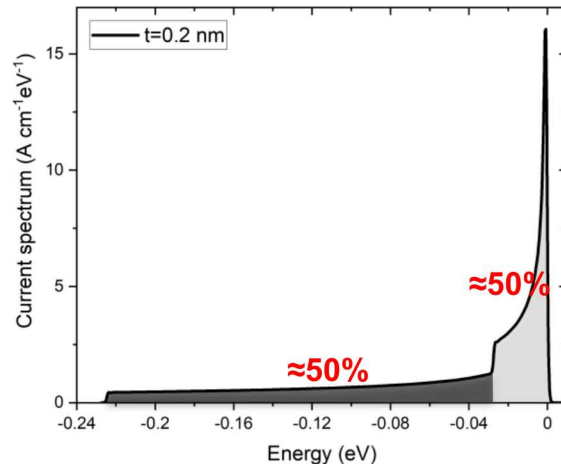
Free
electrons
distribution



VS.



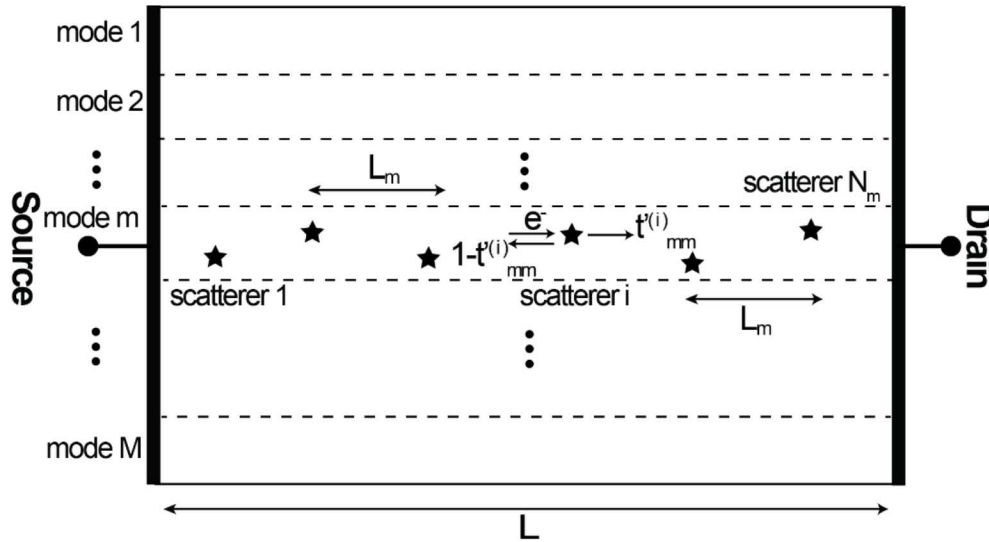
Current
spectrum



- Sharp doping profile reduces the current leakage

$$N_D = 1.2 \times 10^{14} \text{ cm}^{-2}, N_A = 10^{17} \text{ cm}^{-3}, W = 40 \text{ nm}$$

Elastic scattering model



$T_{mm}(E) \equiv$ Electronic transmission
for mode m without scatterers

$t_{mm}^{(i)}(E) \equiv$ Defect transmission probability
due to scatterer i in mode m

$L_m \equiv$ mean free path

- Effective electronic transmission

$$T_{mm}^{\text{eff}}(E) = \frac{1}{\underbrace{1 + \frac{1 - T_{mm}(E)}{T_{mm}(E)}}_{\text{Contact scattering}} + \underbrace{L\nu_m \frac{1 - t_{mm}'(E)}{t_{mm}'(E)}}_{\text{Defect scattering}}}$$

Geometry scattering

$$L\nu_m \frac{1 - t_{mm}'}{t_{mm}'} \equiv \mathcal{O}\left(\frac{L}{L_m}\right) \equiv \alpha \frac{L}{L_m}$$

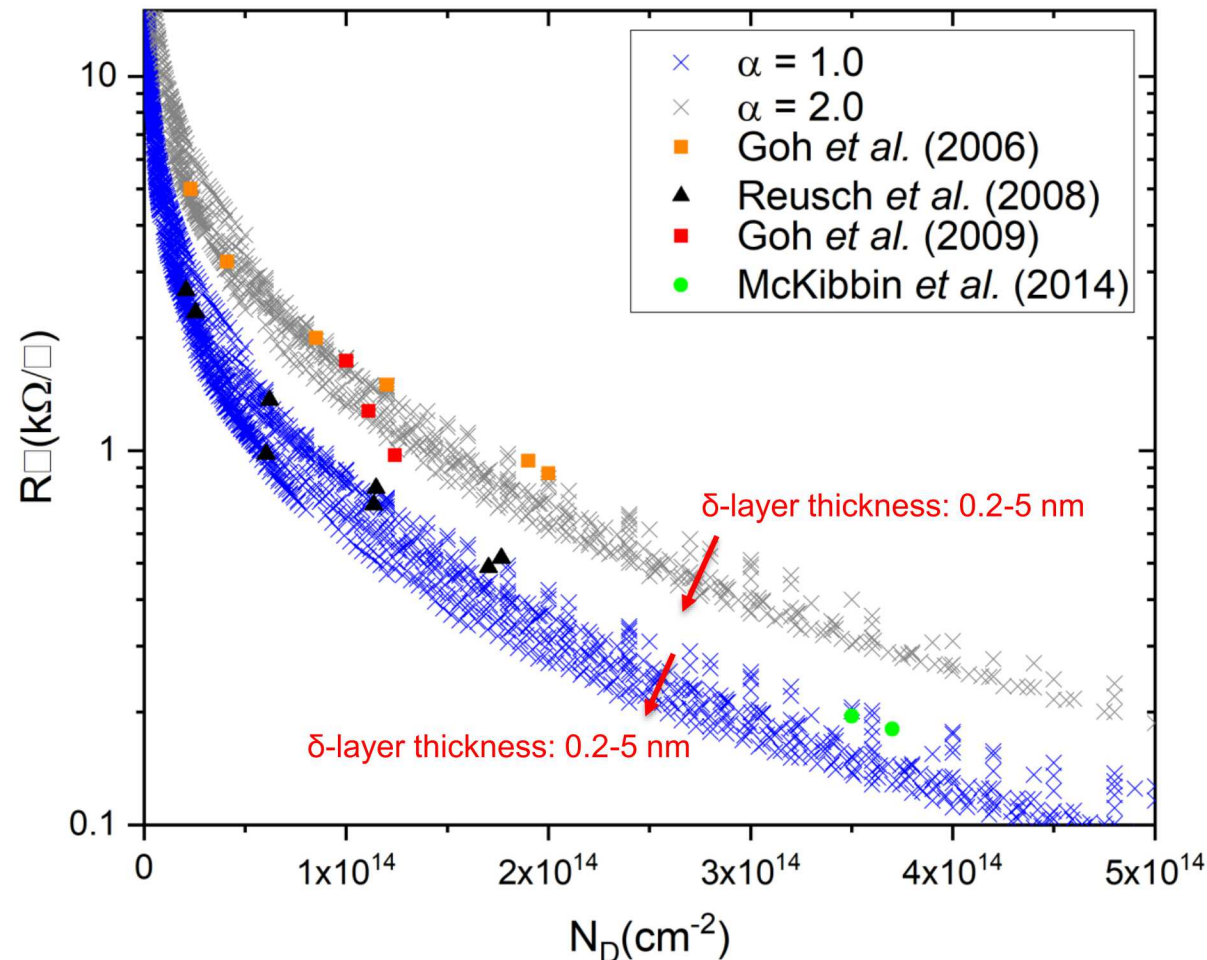
The parameter α is proportional to the linear impurity density

- Current density (Landauer formula)

$$J = \frac{2e}{h} \int \sum_{m=1}^M T_{mm}^{\text{eff}}(E) (f_S(E) - f_D(E)) dE$$

Proof that the Quantum Simulation works!

- Our simulations accurately reproduce the experimental sheet resistances at 4 K



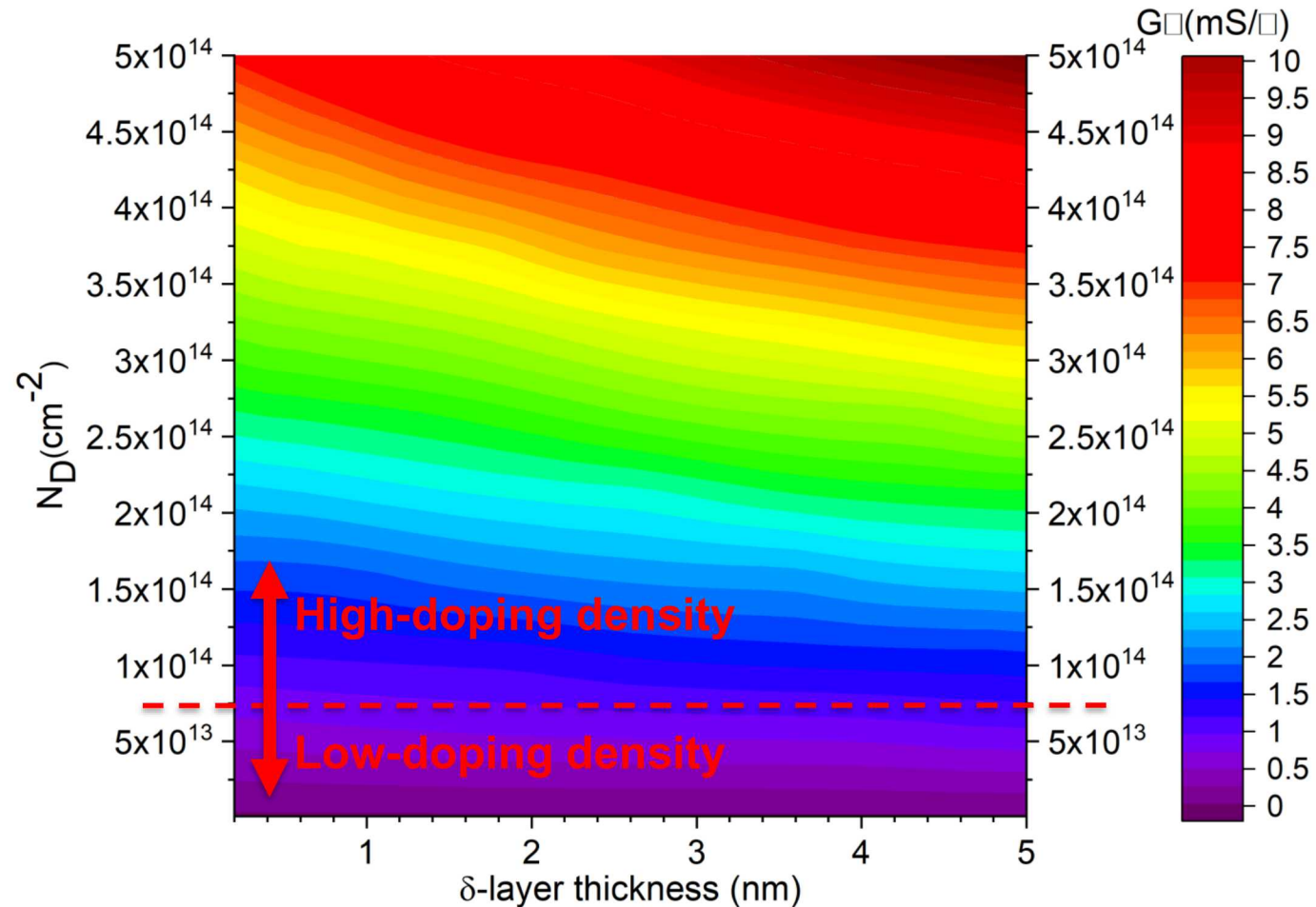
$$T_{mm}^{\text{eff}}(E) = \frac{1}{1 + \frac{1 - T_{mm}(E)}{T_{mm}(E)} + \alpha \frac{L}{L_m(N_D)}}$$

The parameter α is proportional to the linear defect density in the system

K. E. J. Goh *et al.*, Phys. Rev. B, vol. 73, p. 035401, 2006
K. E. J. Goh and M. Y. Simmons, Appl. Phys. Lett., vol. 95, no. 14, p. 142104, 2009
S. R. McKibbin *et al.*, Appl. Phys. Lett., vol. 104, no. 12, p. 123502, 2014
T. C. G. Reusch *et al.*, J. Appl. Phys., vol. 104, no. 6, p. 066104, 2008

Sheet conductance for Si: P δ -layer wires

- Our simulations predicts a quantum thickness dependence on the sheet conductance



Summary

- ❑ Presented a fully quantum **open-system transport simulations** for Si: P δ -layer systems, including a **elastic scattering model**
- ❑ Predicted the **conductive shallow sub-bands** observed experimentally, and evaluated how these conductive sub-bands are influenced by the **donor doping density** and **thickness of the δ -layer**
- ❑ Predicted **quantum thickness dependence** on the conductive properties
- ❑ Excellent agreement with experimental sheet resistance data



Quantum transport in Si:P δ -layer wires

Juan P. Mendez, Denis Mamaluy, Xujiao (Suzey) Gao, Evan M. Anderson,
DeAnna M. Campbell, Jeffrey A. Ivie, Tzu-Ming Lu, Scott W. Schmucker, and
Shashank Misra

ipmende@sandia.gov

Sandia National Laboratories, Albuquerque, New Mexico

Sandia National Laboratories is a multimission laboratory managed and operated by National Technology and Engineering Solutions of Sandia, LLC., a wholly owned subsidiary of Honeywell International, Inc., for the U.S. Department of Energy's National Nuclear Security Administration under contract DE-NA-0003525

