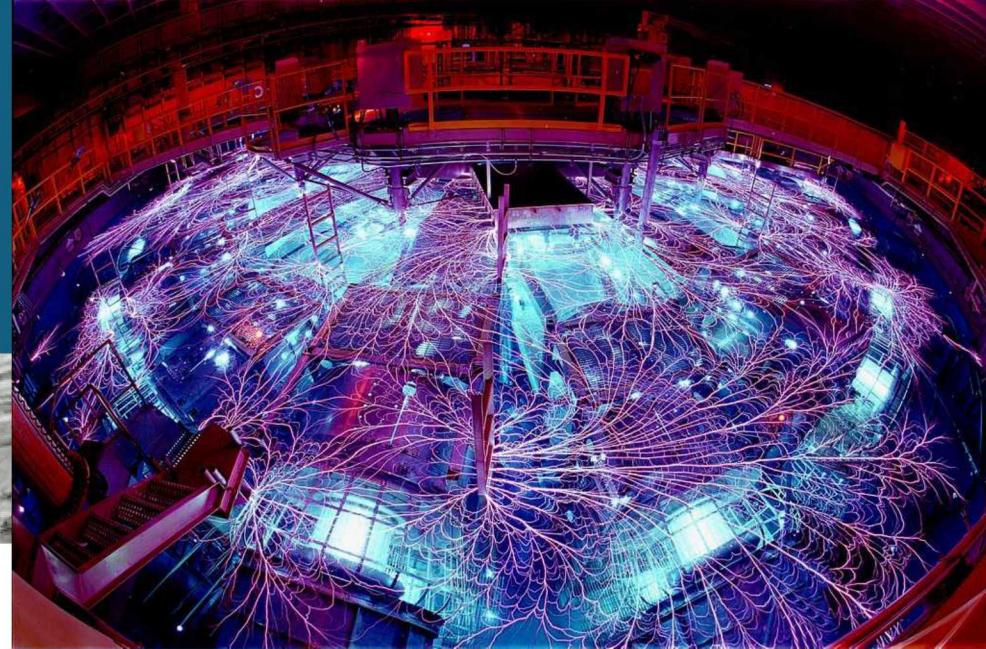
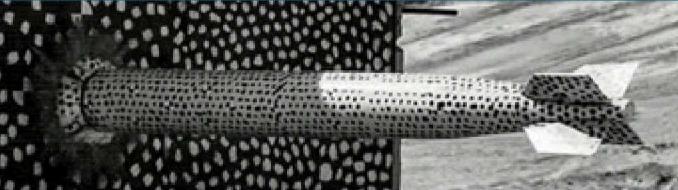


# Optical velocimetry



*Presented by:*

Daniel Dolan

Dynamic Material Properties

Pulsed Power Division



Sandia National Laboratories is a multimission laboratory managed and operated by National Technology & Engineering Solutions of Sandia, LLC, a wholly owned subsidiary of Honeywell International Inc., for the U.S. Department of Energy's National Nuclear Security Administration under contract DE-NA0003525.

# Overview

## Background

- Why we measure velocity
- Transit time measurements

## Optical Doppler shift measurements

- “Displacement” interferometry
- “Velocity” interferometry (aka VISAR)

## Photonic Doppler Velocimetry (PDV)

- Technology overview
- Theory
- Examples

# Why do we measure velocity?

## Why do we measure velocity?

Because we can:

- With fast time resolution (ns and faster)
- With high accuracy (1% or better)
- In harsh conditions (impact, detonation, etc.)
- At many locations simultaneously
- Directly compare with wave codes

Pressure/density linked to jump conditions:

- Shock velocity  $U_s$
- Particle velocity  $u_p$
- Most tabulated data based on this type of measurement

$$\frac{\rho}{\rho_0} = \frac{U_s}{U_s - u_p}$$

$$P = \rho_0 U_s u_p$$

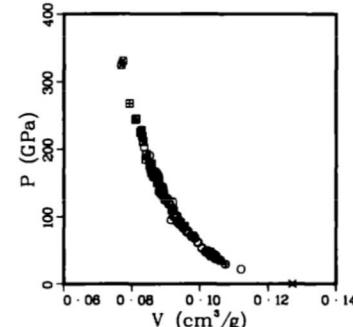
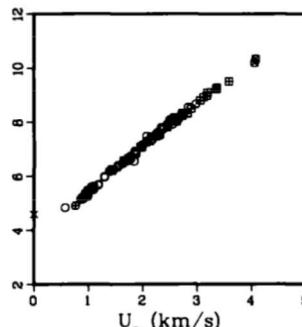
### IRON

Average  $\rho_0 = 7.856 \text{ g/cm}^3$ .

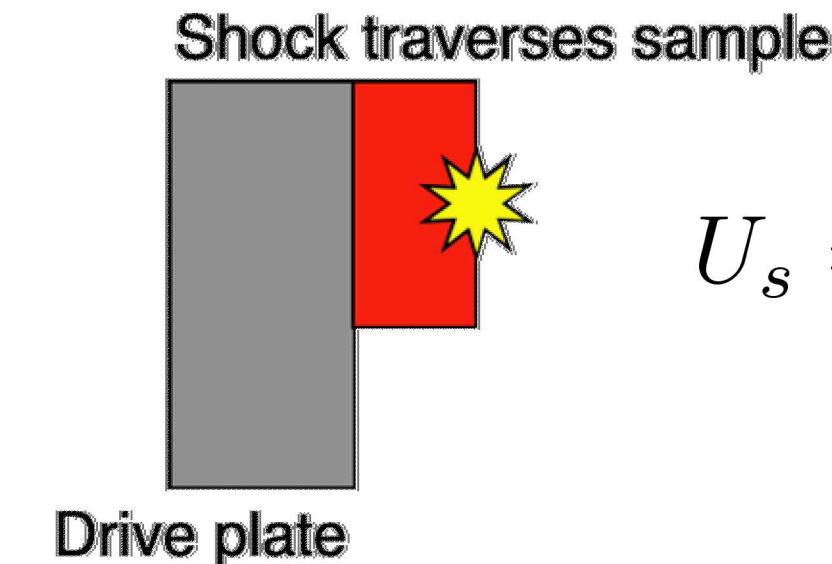
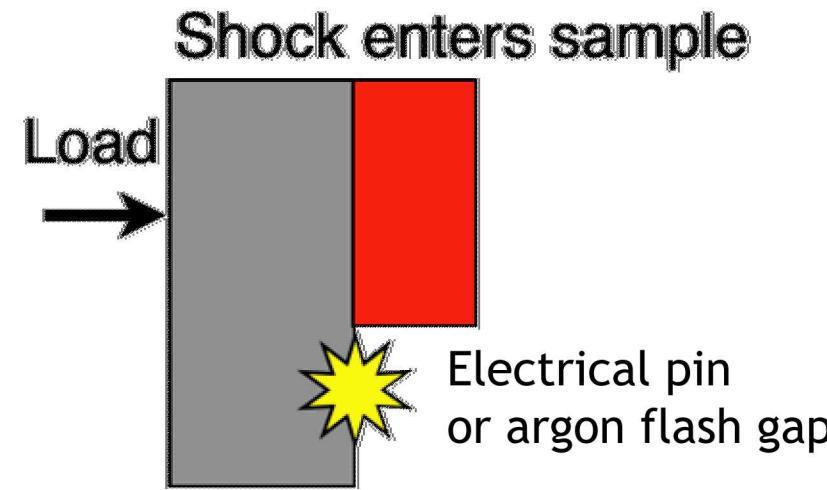
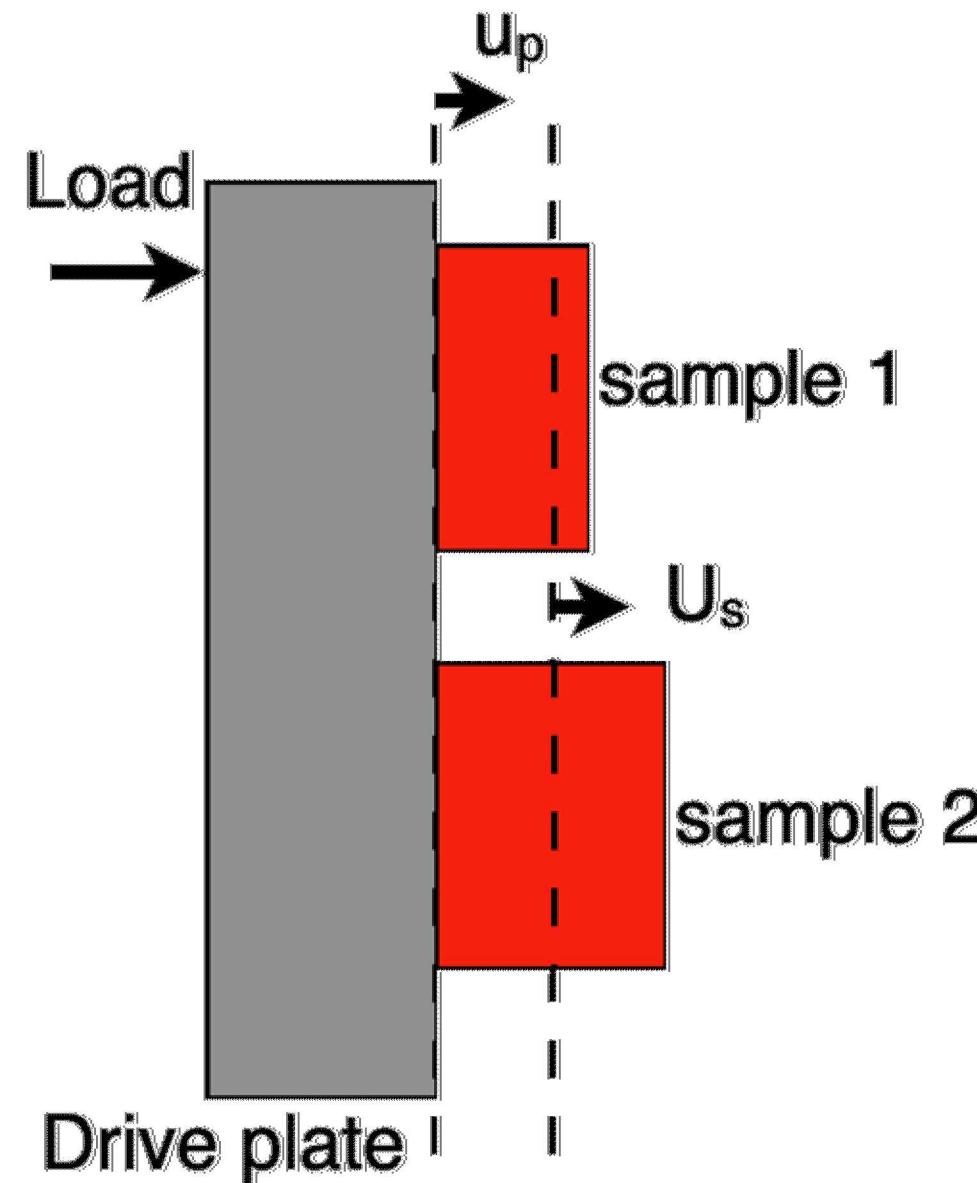
Sound velocities      longitudinal      5.94 km/s.  
 References      shear      3.26 km/s.  
 4, 5, 6, 11, 12, 13, 17, 22

$\rho_0$ ( $\text{g/cm}^3$ )	$U_s$ (km/s)	$U_p$ (km/s)	P (GPa)	V ( $\text{cm}^3/\text{g}$ )	$\rho$ ( $\text{g/cm}^3$ )	$V/V_0$	Exp
7.870	4.595	0.000	0.000	1.271	7.870	1.000	s s p x
7.861	4.838	.573	21.792	1.121	8.917	.882	im1 o
7.850	4.913	.763	29.427	1.076	9.293	.845	s f1 s
7.840	5.144	.867	34.965	1.061	9.429	.831	s f1 s
7.840	5.147	.867	34.986	1.061	9.428	.832	s f1 s
7.840	5.168	.876	35.493	1.059	9.440	.830	s f1 s
7.850	5.190	.881	35.893	1.058	9.455	.830	s f1 s
7.840	5.166	.884	35.803	1.057	9.459	.829	s f1 s
7.882	5.225	.903	37.189	1.049	9.529	.827	im1 o
7.882	5.172	.906	36.934	1.046	9.556	.825	im1 o
7.850	5.328	.948	39.650	1.047	9.549	.822	sp1 s
7.850	5.360	.952	40.056	1.048	9.545	.822	s f1 s
7.840	5.393	.968	40.928	1.047	9.555	.821	im1 o
7.840	5.373	.969	40.818	1.045	9.565	.820	im1 o
7.882	5.339	.984	41.409	1.035	9.663	.816	im1 o
7.840	5.408	.988	41.890	1.042	9.592	.817	im1 o
7.864	5.252	.989	40.847	1.032	9.688	.812	im1 o
7.850	5.443	.995	42.514	1.041	9.606	.817	s f1 s
7.843	5.458	.998	42.721	1.042	9.598	.817	sp1 s

(Continued)



## A traditional shock experiment



$$U_s = \frac{\Delta x}{\Delta t}$$

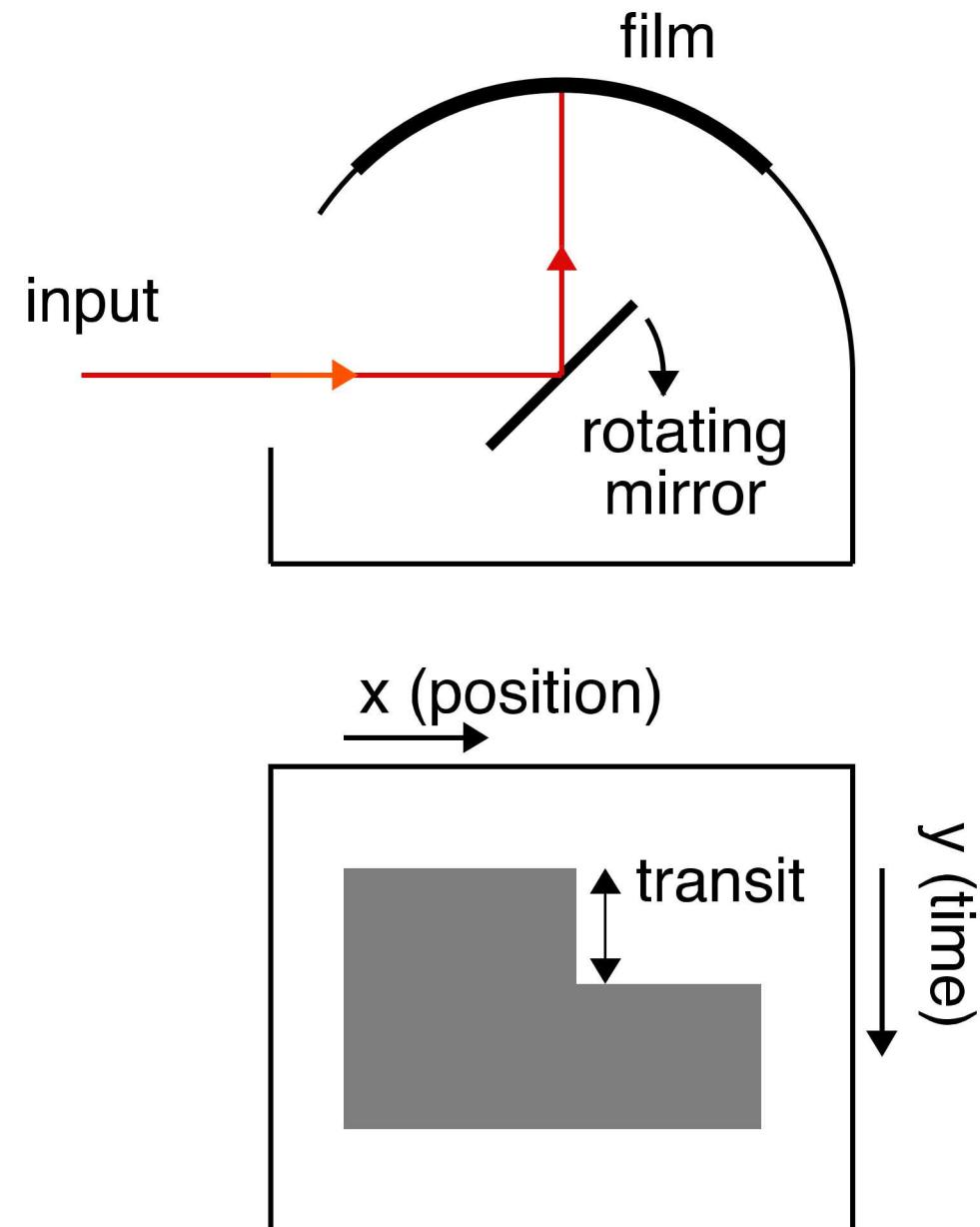
## 6 Recording bandwidth

Rotating mirror streak cameras for Ar flash

- ~100 MHz recording bandwidth
- Rotation must be known precisely
- Film distortion is a problem
- Largely replaced by streak tubes
  - >1 GHz bandwidth

Electrical shorting pins became more common with faster oscilloscopes

- 10-24 MHz from late 1940s to 1960
- ~1000 MHz available in 1960
- Digitizers transition 1960-1995
  - 1995-2005 >1 GHz
  - 2005-2015 >10 GHz



# Clever pin use reveals wave structure

- Bancroft et al., J. Appl Phys 27, 291 (1956).
- Most of us are **not** that clever

## POLYMORPHISM OF

it meets a retreating rarefaction in general to increase the shock velocity. It will be assumed that effects due to the elastic wave are negligible.  $D_{32}'$  can be determined from the coordinate

4. It may easily be shown that

$$t_4 = \frac{d_3 - d_5 + t_3 D_c' + t_5 D_d'}{D_c' + D_d'} \quad \{$$

his same conclusion may be reached in a variety of other ways.

An estimate of  $D_d'$  is also fairly clear. The amplitude of the disturbance at this stage is less than the amplitude to which  $D_{21}$  applies; accordingly one might expect a velocity to be somewhat lower than  $D_{21}$ . However, results from other experiments at shock pressures below 13 megabar include velocities of propagation which do not clearly differ from  $D_{21}$ . Therefore it seems reasonable to assume that  $D_d' = D_{21}$ , so that  $D_d' = D_{21} - 2u_2$ . Thus Eqs. (10) become

$$t_4 = \frac{d_3 - d_5 + D_{21}(t_5 + t_3) + 2u_2(t_5 - t_3)}{2D_{21}} \quad (14)$$

$$d_4 = d_3 - (D_{21} - 2u_2)(t_4 - t_3).$$

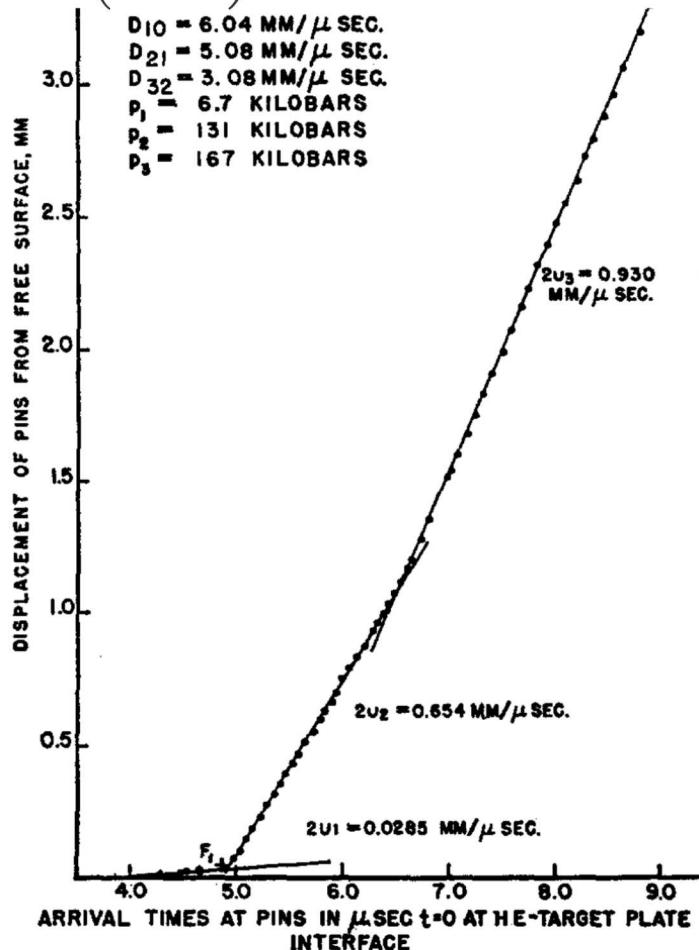


FIG. 8. Displacement-time curve for Exp. 5.  
See text for analysis.

shock system was detected. The pin array was simple

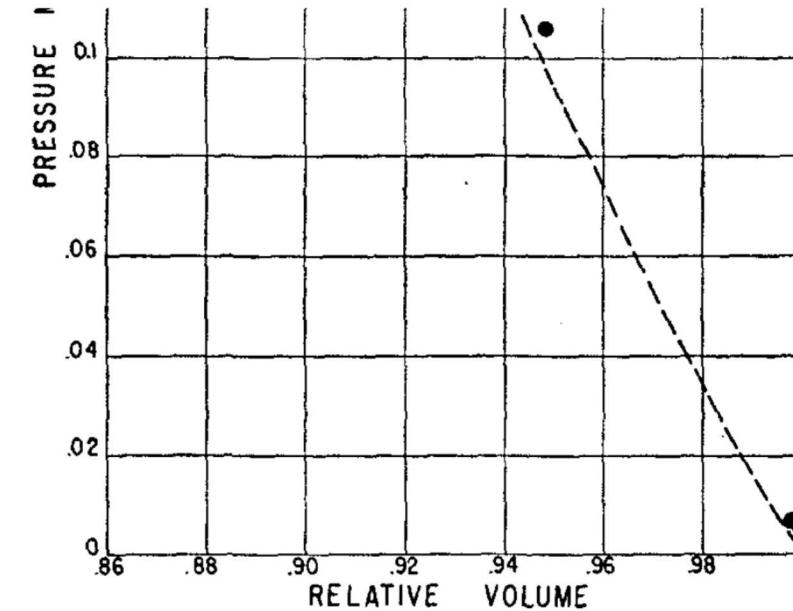


FIG. 10. Relative volume  $v/v_0$  as a function of pressure. Circles represent points computed from measurement of shock wave and surface velocities of Armco iron. The dashed curve is an extrapolation of Bridgman's data [with the help of Eq. (21)].

gabar, which is negligible. Accordingly the average may be used directly for computing a value in the isentropic equation of state

$$p_s = \alpha_s \mu + \beta_s \mu^2, \quad (15)$$

where  $\mu = (v/v_0) - 1$  is the compression. For an

## Non-steady waves

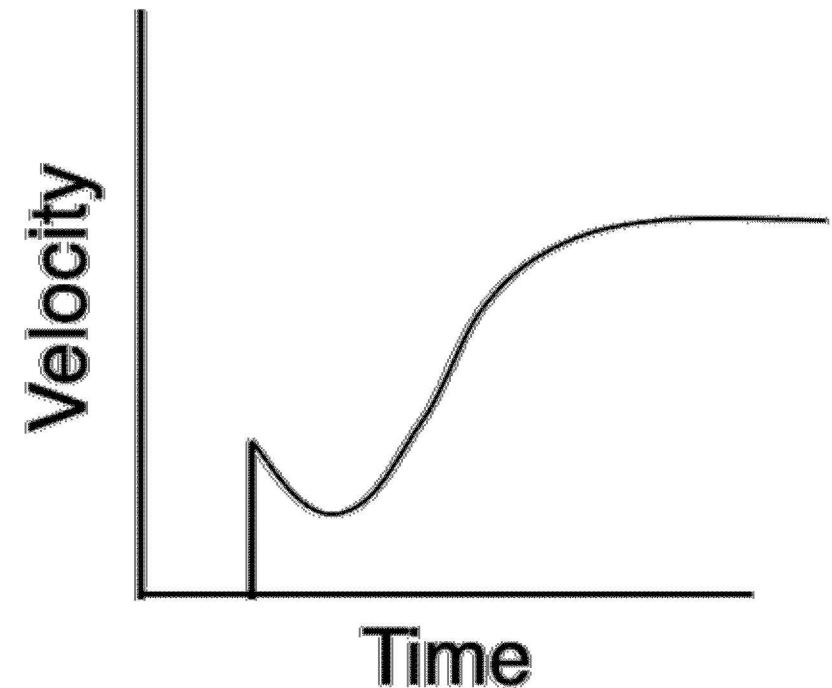
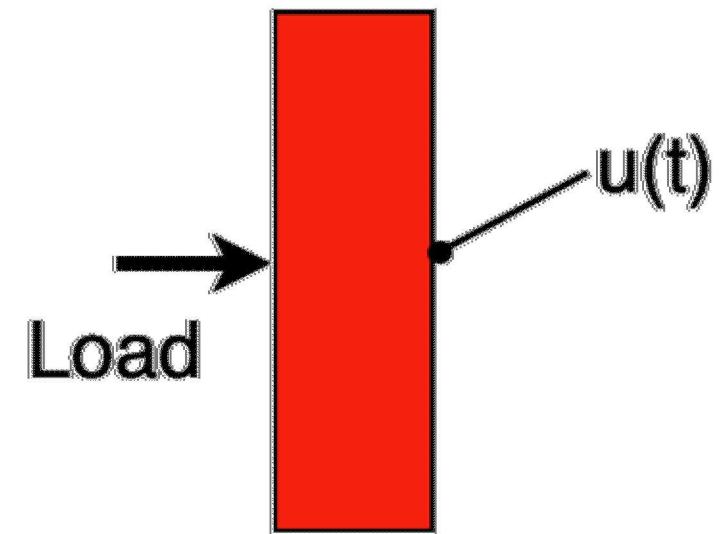


Mechanical waves often contain a lot of structure

- Inelastic compression
- Phase transitions
- Chemical reactions
- Ramp wave (release and/or tailored)

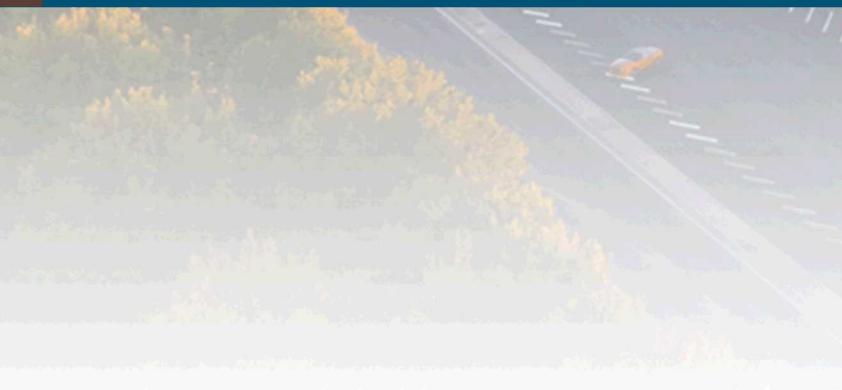
Structure difficult to extract from transit time measurements

- Real-time velocity diagnostics are needed





# Optical velocimetry



## There are many ways to track motion with light



Optical emission: compression creates light

- “Passive” shock breakout

Optical reflection (amplitude or direction)

- “Active” shock breakout
- Beam deflection requires special geometry and specular surfaces

Optical phase techniques are more flexible

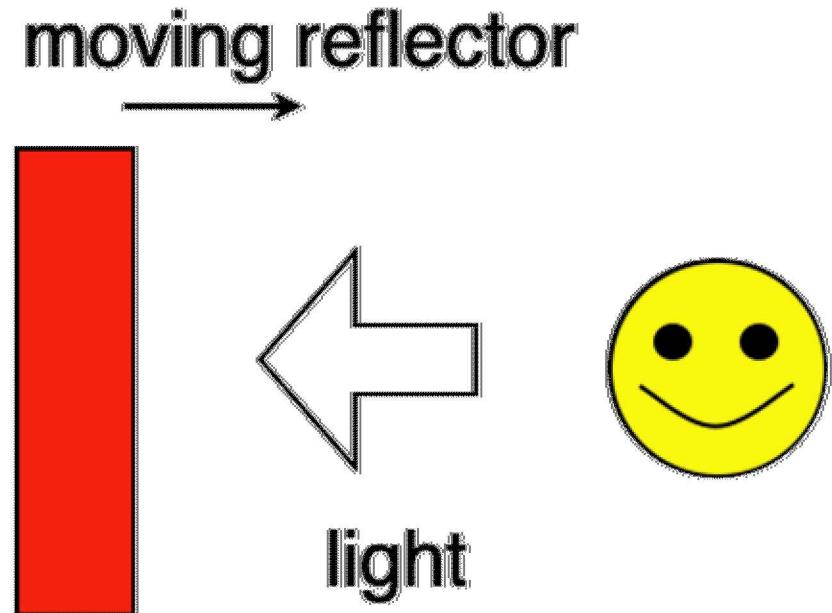
- Optical phase “wags” the electric field
- Frequency is the rate of optical phase change
- Wavelength is the reciprocal of frequency
- 532 nm is 564 THz (1-2 fs)

$$E(t) = A(t) \cos \phi(t)$$

$$\phi(t) \sim 2\pi ft$$

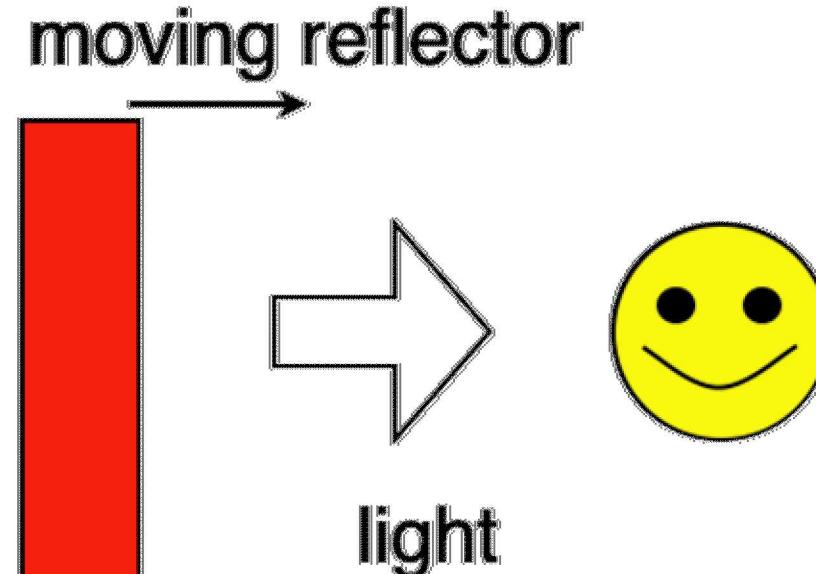
$$f = \frac{c_0}{\lambda}$$

# The optical Doppler effect



Observer sees:  $\lambda_0$

Reflector sees:  $\lambda'$



Reflected light:  $\lambda'$

Observer sees:  $\lambda$

$$\frac{\lambda}{\lambda_0} \approx 1 - \frac{2v}{c_0}$$

Non-relativistic motion  
(100 km/s is 0.03%  $c_0$ )

## Optical velocimetry usually means optical interferometry



Electric field cannot be measured directly

- Direction flips many times over detector response

Wavelength changes cannot be resolved with grating spectrometers

- 6-7 ppm change for 1 km/s (0.004 nm at 532 nm)

Two-beam interferometry is the most common approach

$$I(t) = I_1 + I_2 + 2\sqrt{I_1 I_2} \cos \underbrace{(\phi_1(t) - \phi_2(t))}_{\text{phase difference}}$$

- Fields add coherently
- Intensity (power/area) is time-averaged square of electric field

## Some terminology

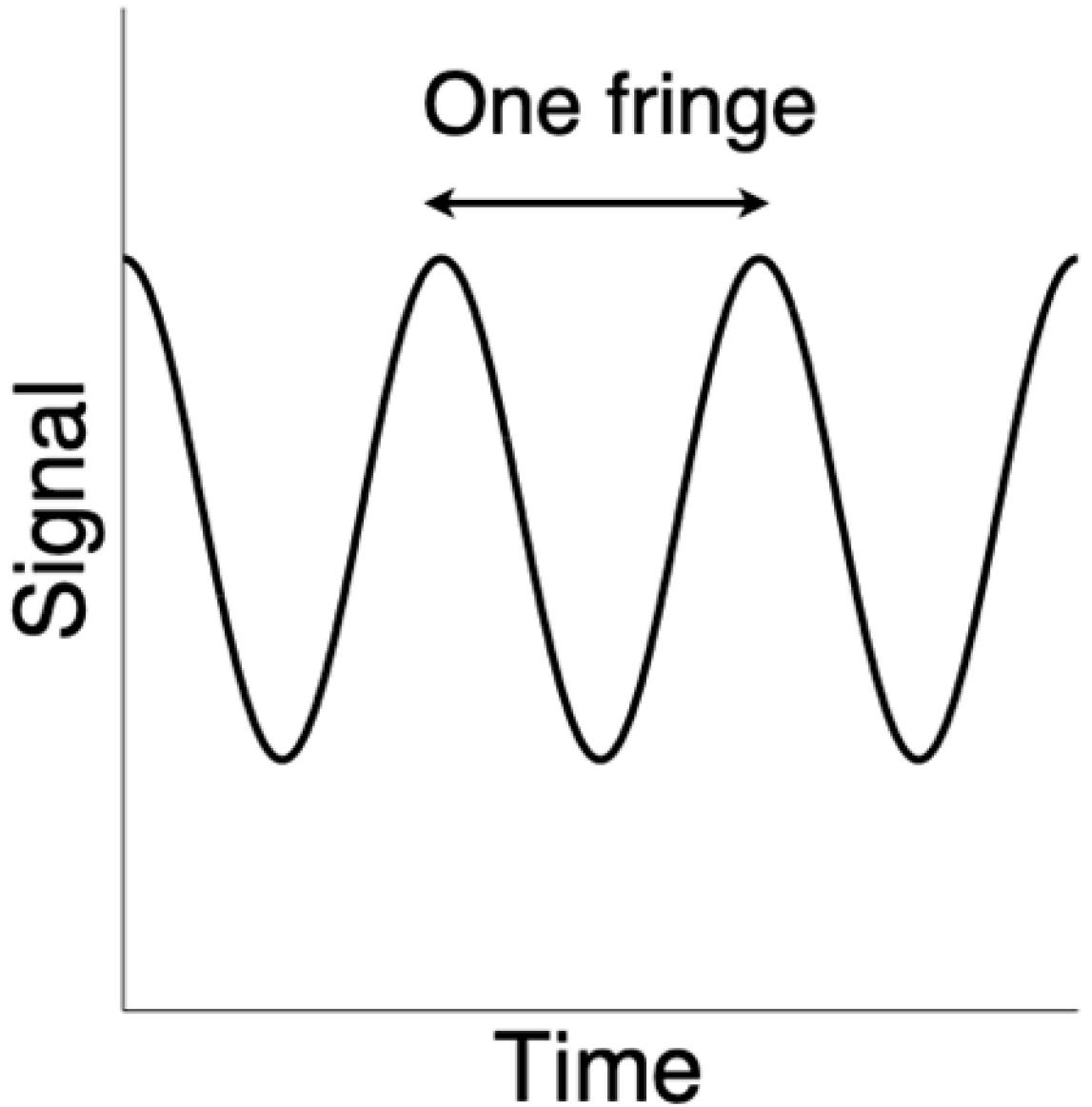


Fringe shift

- Phase difference scaled by  $2\pi$
- Same as signal cycles

Beat/fringe frequency

- Rate of signal cycles
- Not the same as optical frequency



**We can only measure signal cosine functions, not the electric field cosine function**

## A historical detour



L.M. Barker, Experimental Mechanics 12, 209 (1972).

### Sandia displacement interferometer

- Michelson configuration
- 1 fringe = 1/2 wavelength motion

### Problems

- Mirror finish required
  - Surface often changes at shock breakout
- Limited velocity range
  - 1 km/s is 3.2 GHz at 632.8 nm
  - No oscilloscope could follow such frequencies in the 1970s

## ive Research

ibes



# Velocity Interferometer System for Any Reflector (VISAR)



# VISAR solves the bandwidth problem



Barker and Hollenbach, J. Appl. Phys. **43**, 4669 (1972).

Doppler-shifted light mixed with a time-shifted version of itself

- Avoids large steady-state frequencies

Etalon allows diffuse reflectors

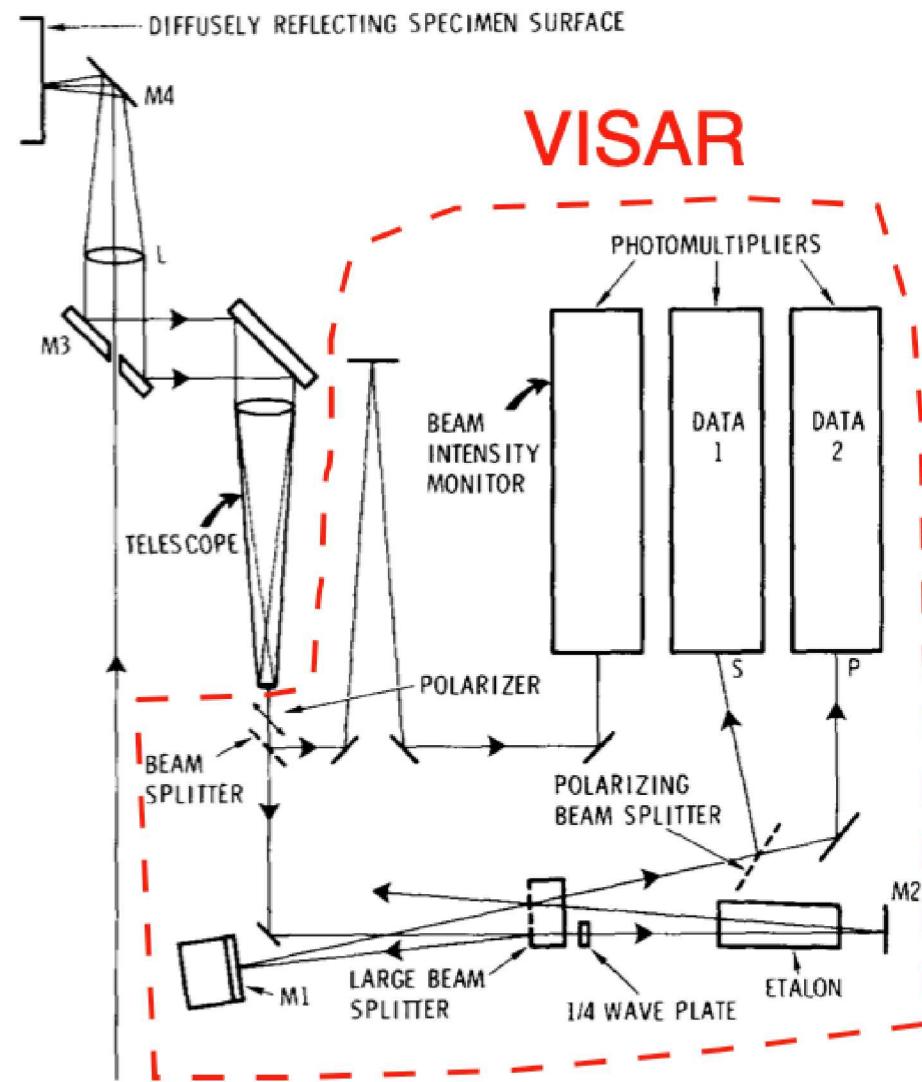
- “Any Reflector”

Velocities in mm/ $\mu$ sec, accelerations in mm/ $\mu$ sec<sup>2</sup>, str

$\frac{u_{fm}}{u_m}$ <sup>a</sup>	$\frac{u_{fs}}{u_m}$ <sup>a</sup> atop P1 wave	Stress atop P1 wave	Peak accel. in P1 wave <sup>b</sup>	Peak accel. in P2 wave <sup>c</sup>
1.982	0.645	13.15	...	1.52
1.976	0.652	13.30	...	4.00
1.971	0.635	12.92	...	0.85
1.968	0.637	12.96	...	4.66
1.972	0.653	13.32	...	7.86
1.978	0.667	13.60	...	26.1
1.964	0.634	12.90	...	5.29
1.967	0.640	13.02	...	14.8
1.976	...	...	...	...
1.971	...	...	...	...
...	...	...	...	...
1.969	...	...	...	...

Continuous wave profile of iron.

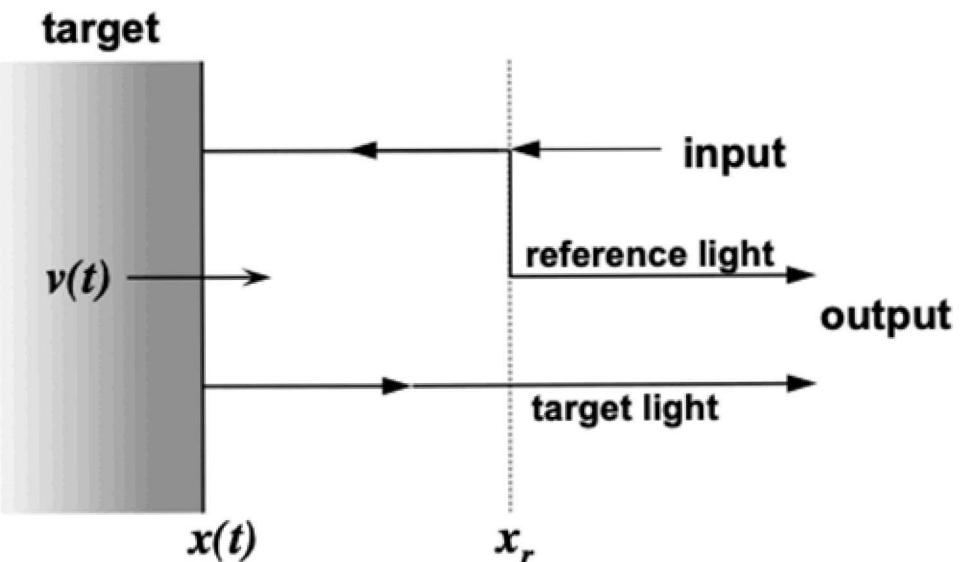
Barker and Hollenbach,  
J. Appl. Phys. **45**, 4872 (1974).



# What is the difference?

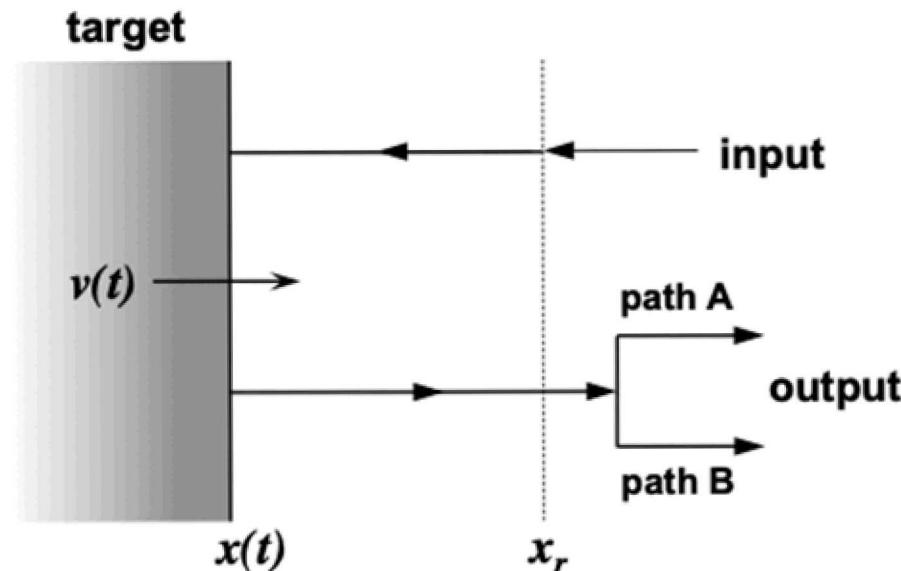
## “Displacement” approach

- One output path contains target (Doppler)
- Other output path does NOT contain the target
- Mixes two different optical frequencies



## “Velocity” approach

- Both output paths contain the target (Doppler)
- Mixes two copies of a single frequency\*



## \*The VISAR approximation

IF velocity changes slowly compared to the delay time:

- Fringe shift scales with velocity (approximation)
- Exact behavior is more complex
  - Explicitly depends on delay time
  - Dolan and Specht, JDBM 3, 407 (2017).

Sensitivity defined by wavelength and delay time

- Fringe constant or Velocity Per Fringe (VPF)
- ~50-5000 m/s per fringe

$$F = \left( \frac{2\tau}{\lambda_0} \right) u$$

$$VPF = \frac{\lambda_0}{2\tau}$$

$$\tau \approx 0.1\text{--}10 \text{ ns}$$

**Smaller hard to characterize**  
**Larger hard to build**

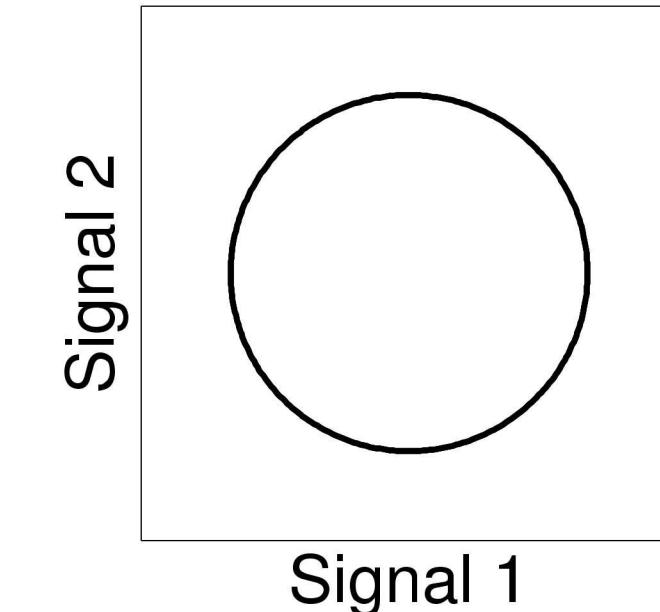
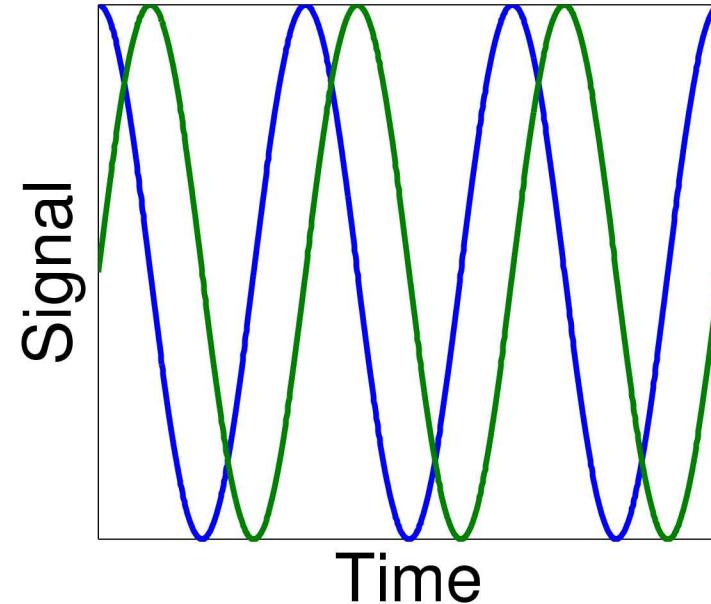
## The need for quadrature

Inverting a sinusoid is not always easy

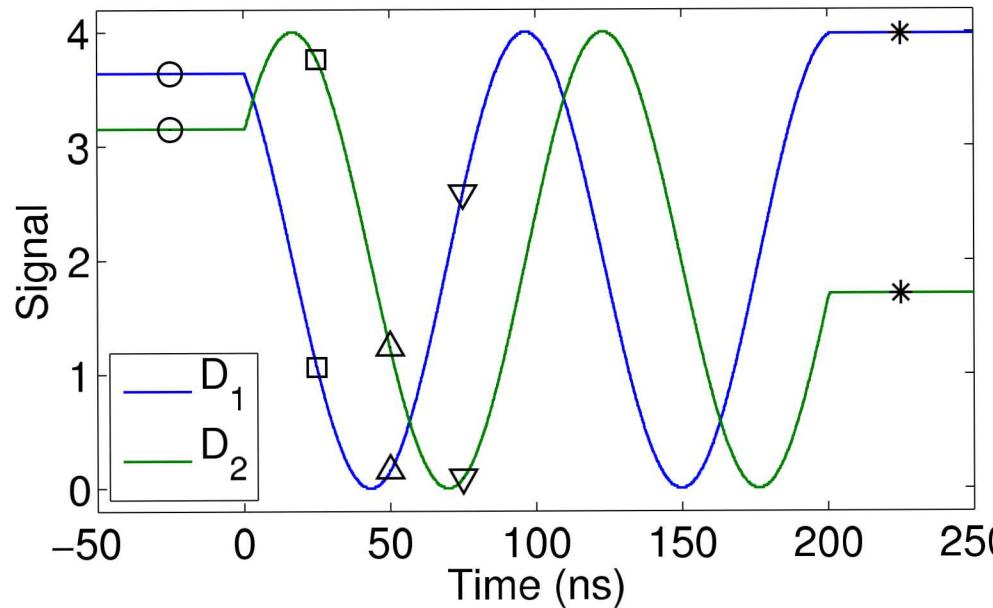
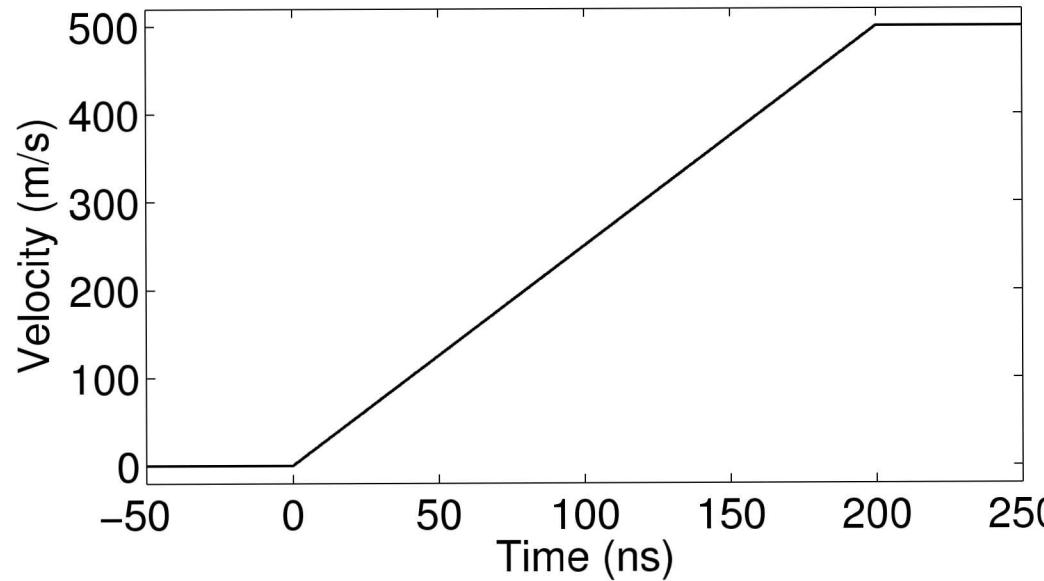
- Arcsine and arccosine defined over 180 degrees
- Steep sections are sensitive
- Peaks/troughs are insensitive

Measuring 2+ quadrature signals provides:

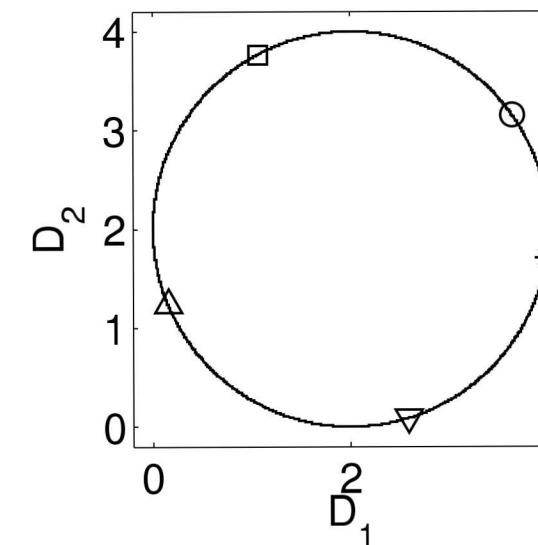
- A robust inversion
- Reduces the effect of amplitude variations
- Arctangent defined over 360 degrees



# Ramp example (1 ns delay)

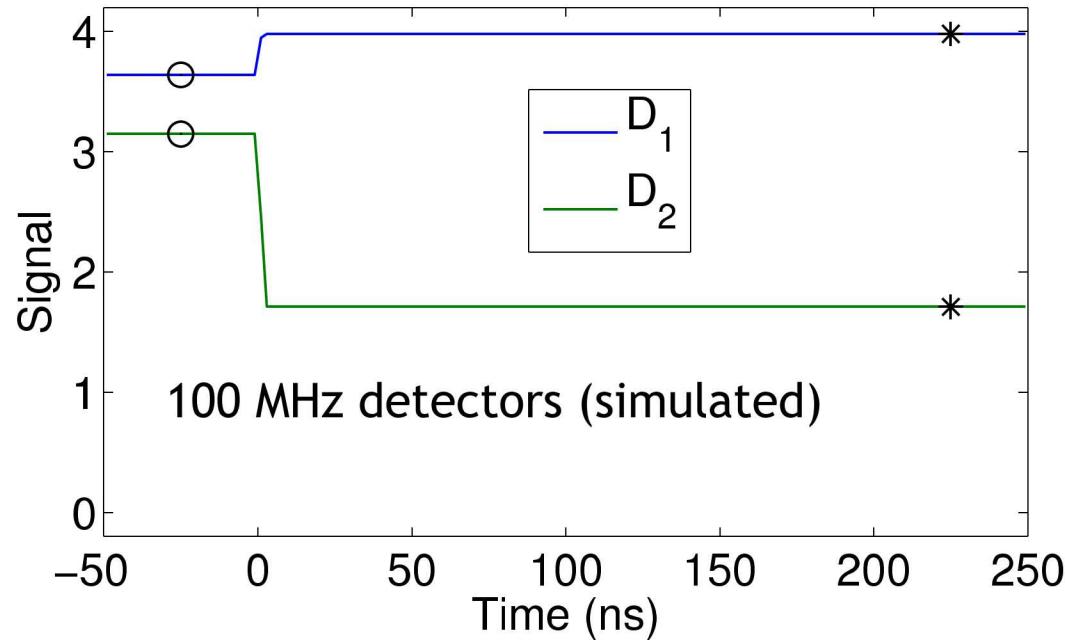
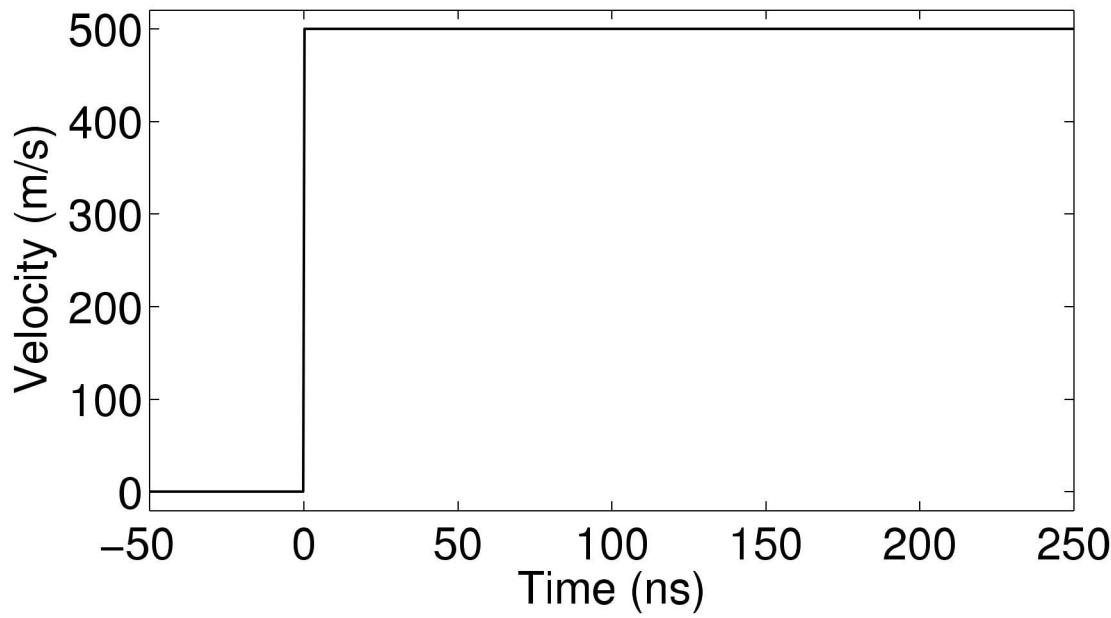


Quadrature resolves  
fringe wrapping

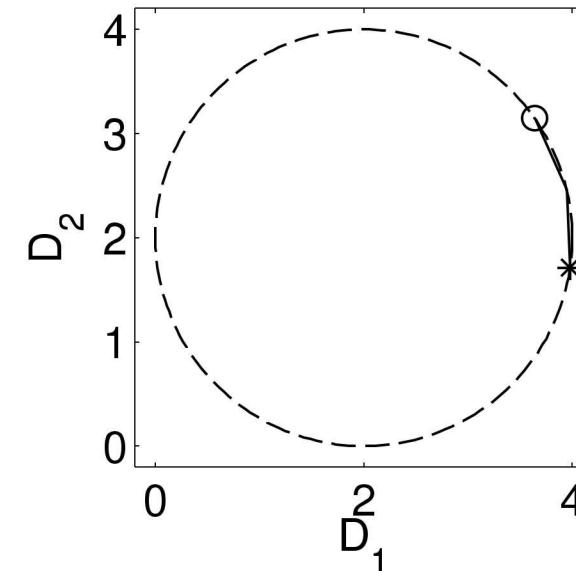


Signals change when  
velocity changes

# Shock example (1 ns delay)



Motion is backwards?



Signals change over etalon delay

## Cold, hard realities of VISAR



Fringe shift only known to an integer offset

- Detectors that cannot keep up with optical signal may “lose/jump fringes”
- Multiple VPFs are generally used to resolve this ambiguity
- VISAR is designed to measure fringe shift from a **single** Doppler shift
  - Multiple velocities cause confusing interference
  - VISAR ellipse collapses to its center

$$\cos(\Phi) = \cos(\Phi + 2\pi) = \dots$$

$$\begin{aligned}u &= K_1(F_1 + N_1) \\&= K_2(F_2 + N_2)\end{aligned}$$

Ellipse position/size changes with light level!



Push-pull quadrature [Hemsing, Rev. Sci. Instrum. **50**, 73 (1979).]

- Four signal phases (0, 90, 180, and 270 degrees)
- Explicitly removes coherent/incoherent light variation

Fiber coupling

- VISAR was originally an open-beam diagnostic
- Open beam now used primarily for line VISAR measurements

Faster detectors

- Optical streak-cameras (ORVIS)
- Improved photodiodes/photodetectors
- Faster digitizers

Incremental changes over the past two decades



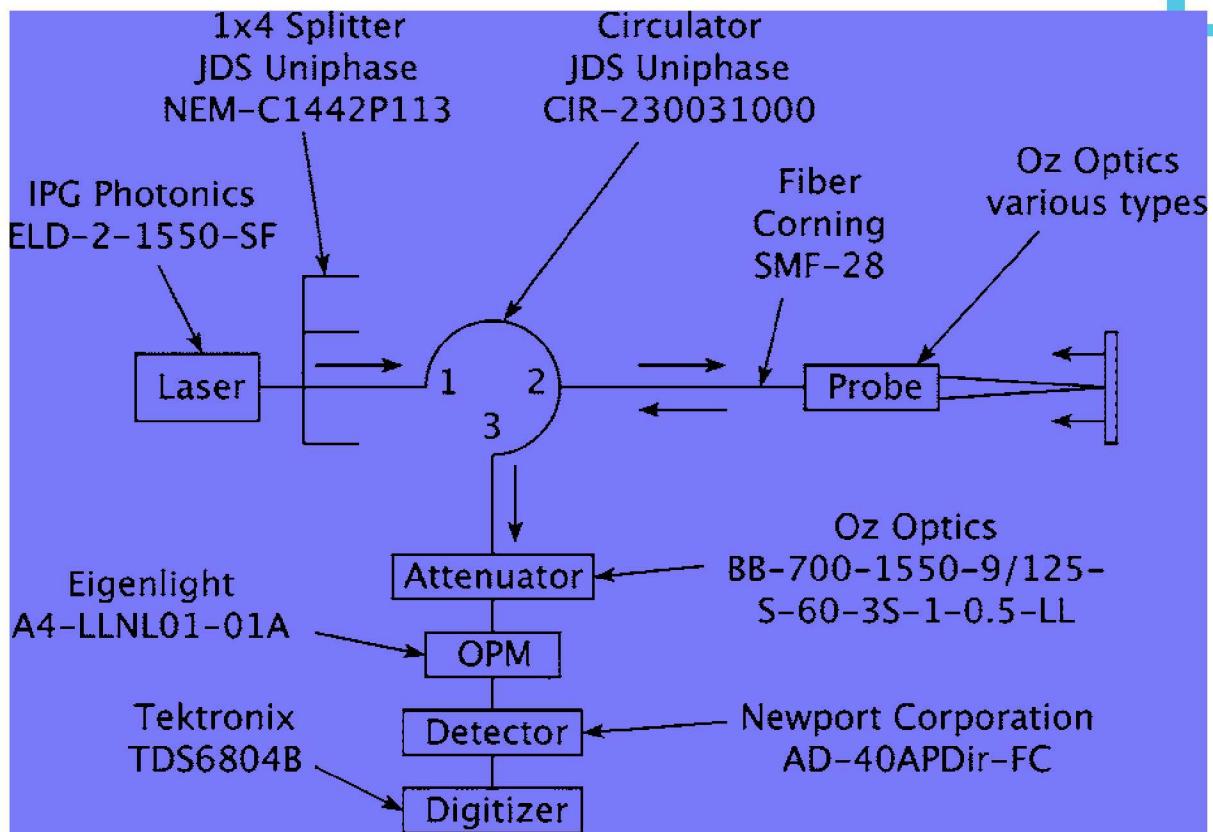
# Photonic Doppler Velocimetry (PDV)



# PDV born at LLNL (2002-2003)

Utilizes advances from the telecommunications industry (1550 nm)

- Compact fiber lasers
- 9  $\mu\text{m}$  core diameter (SMF)
- Narrow line width (<10-100 kHz)
- Three-port circulator (magic!)
- Port 1 input goes to port 2
- Port 2 input goes to port 3
- High speed detectors/digitizers (>10 GHz)



Strand et al, Rev. Sci. Instrum. 77, 83108 (2006).  
Reference light comes from probe's back reflection

See Dolan RSI 91, 051501 (2020) for comprehensive review

## Example: ramp measurement with PDV



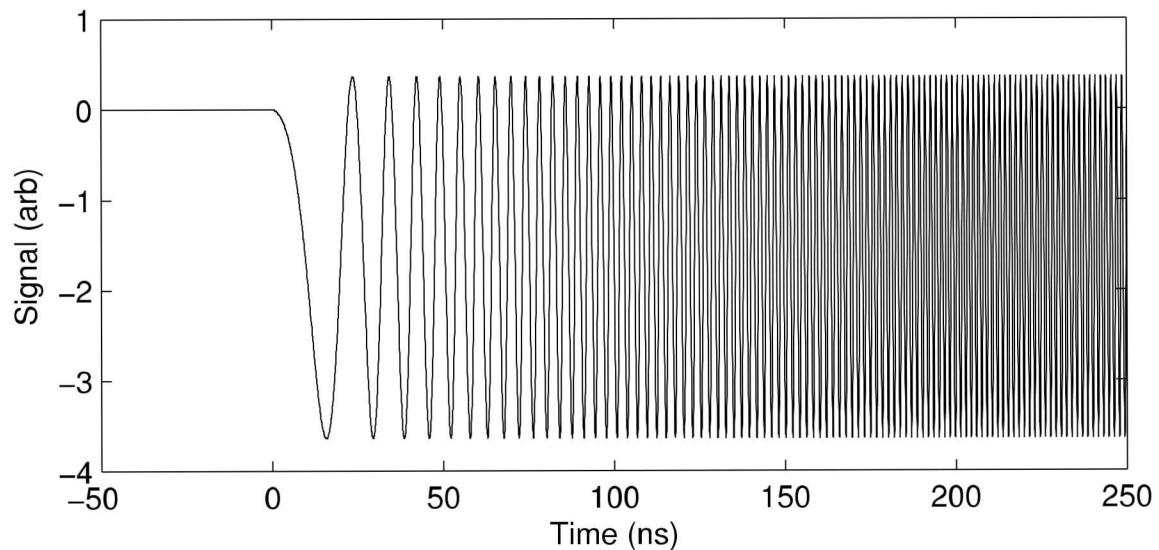
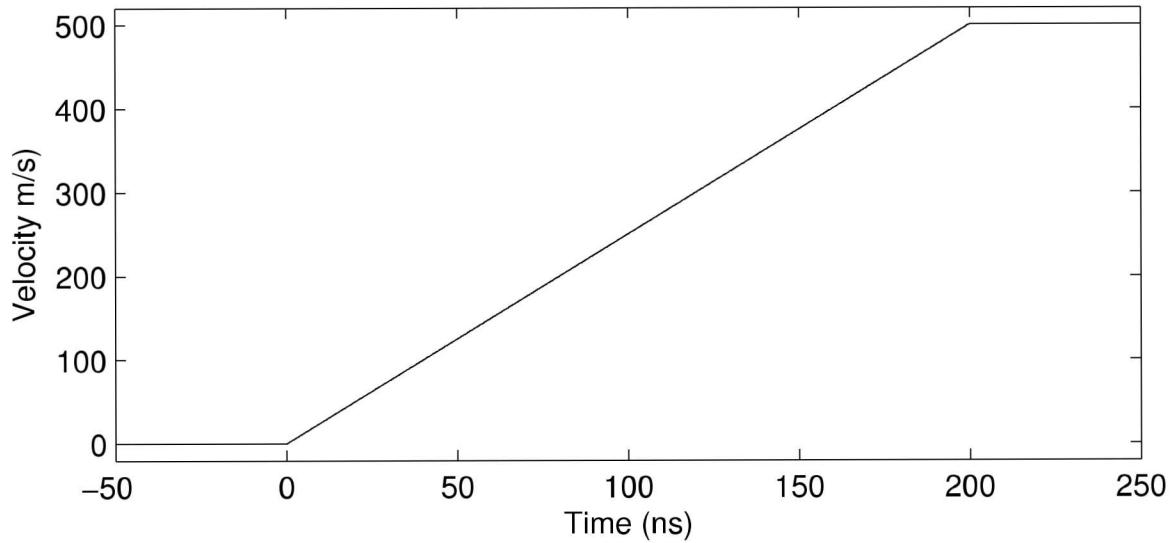
Michelson interferometer

- 775 nm motion = 1 fringe

Signal frequency changes with velocity

- Displacement interferometer lives again (in fiber)!

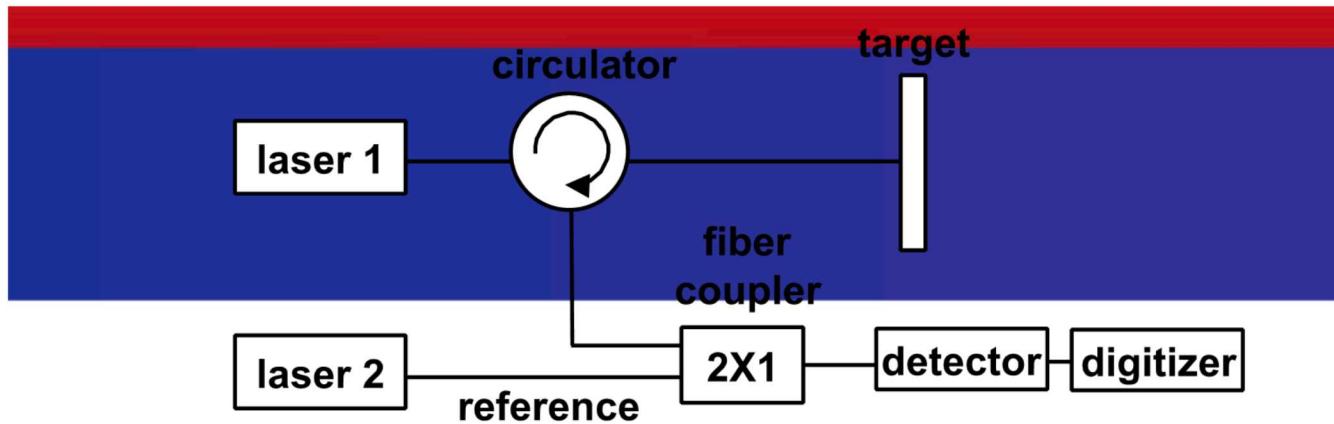
$$\frac{dF}{dt} = B(t) = \frac{2v}{\lambda_0}$$



# There are many names for and configurations of PDV

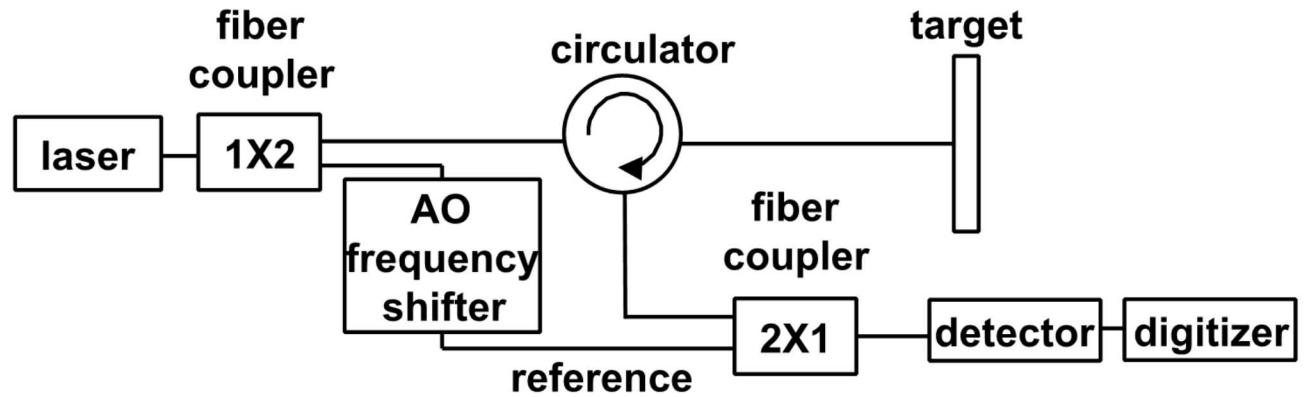
## Conventional PDV

- One laser wavelength
- Confusingly called heterodyne velocimetry (HetV)



## Quadrature PDV

- Phase-based measurement (3 signals)
- Triature, PDI, DISAR, ...



## Frequency-shifted PDV

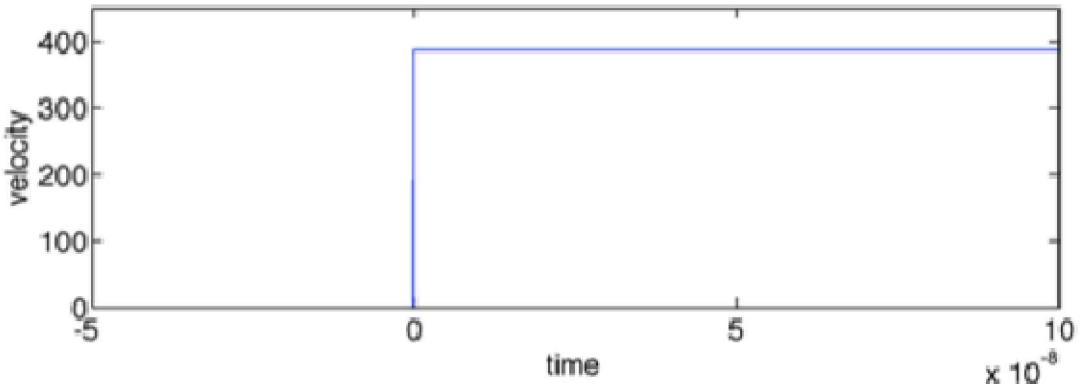
- Two laser wavelengths OR laser + shifter

D. Dolan, “Accuracy and precision in photonic Doppler velocimetry,” Review of Scientific Instruments **81**, 53905 (2010).

## Example: step measurement with PDV

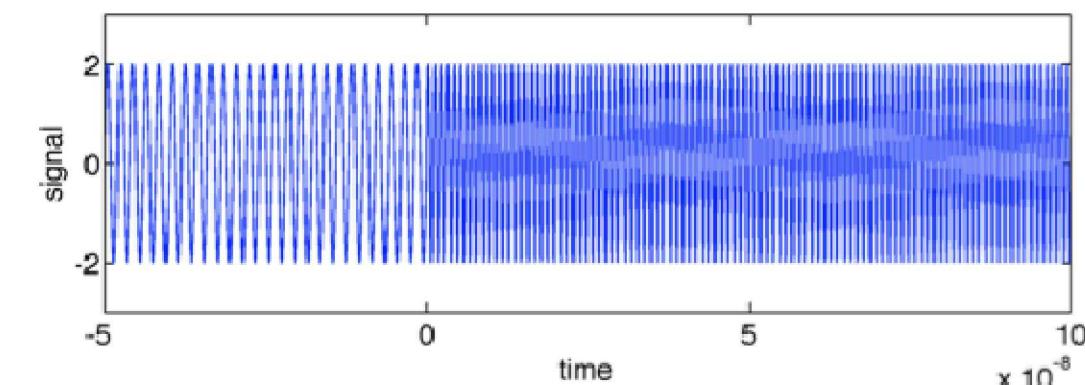
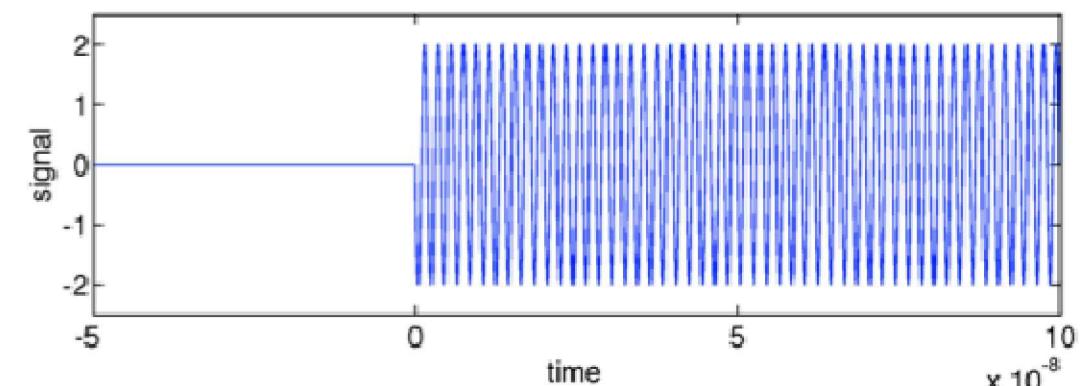
### Conventional PDV

- No motion, no fringes



### Frequency-shifted PDV

- Always fringing, even at rest
- Helpful in digitizer setup



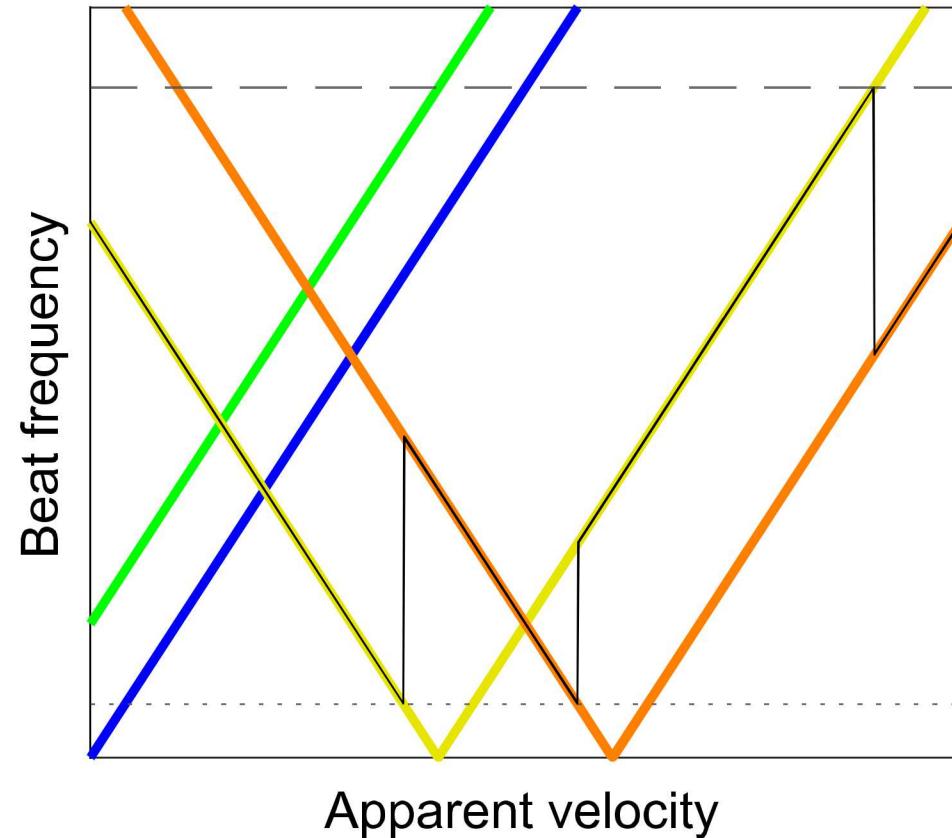
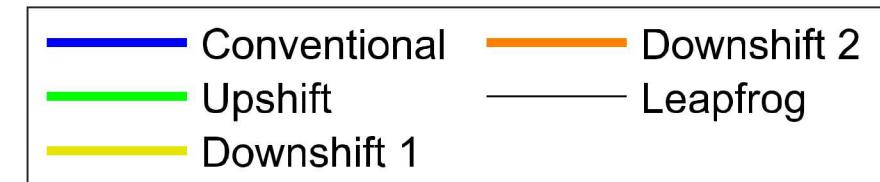
# Mapping velocity to beat frequency

Frequency shifting moves mapping up/down

- This is done by tuning laser wavelengths
- Upshift: reference longer than target
- Downshift: reference shorter than target

Mapping is not unique

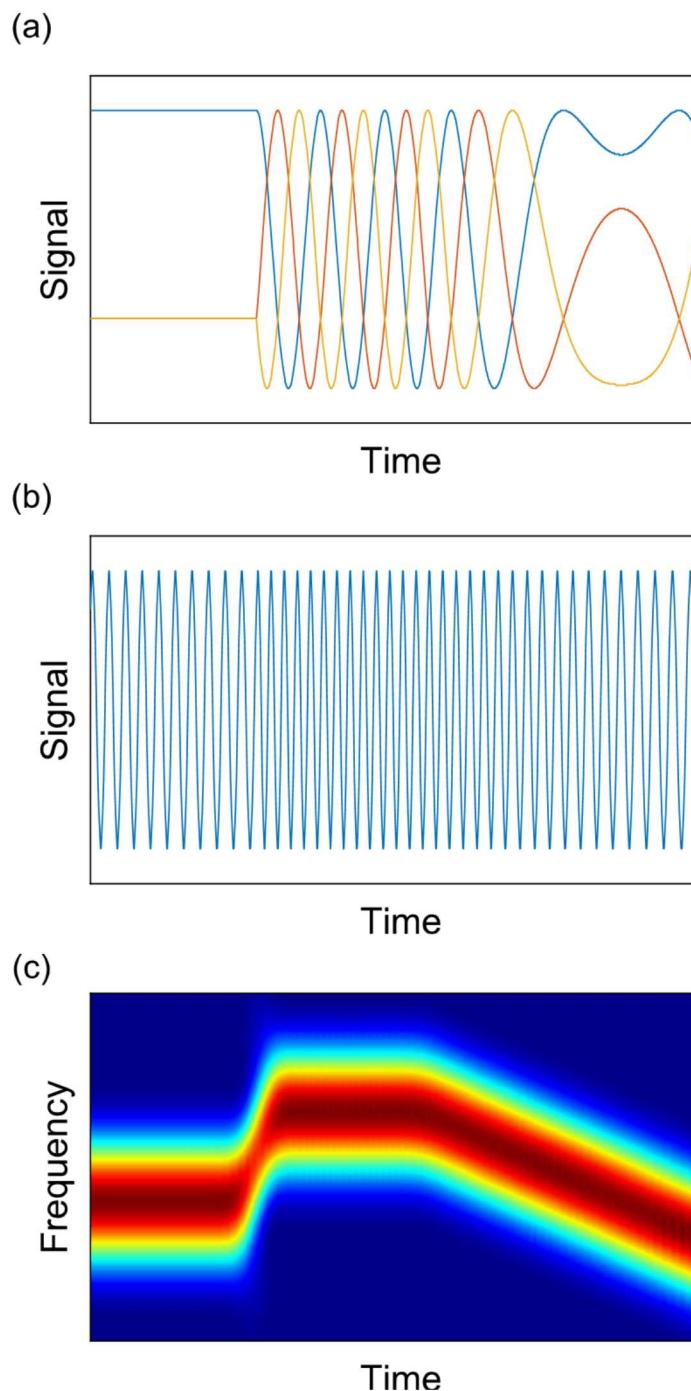
- Each beat frequency is associated with two velocities
- Physical constraints can eliminate one choice
- “Leapfrog” measurements use multiple references to resolve this ambiguity



$$B = \left| v_T - v_R + \frac{2}{\lambda_T} v^* \right|.$$

## A more complicated example

- (a) Quadrature PDV signals
  - 120 degree phase shifts
  - Frequency is time-derivative of phase shift
- (b) Frequency-shifted signal
  - Same content as (a)
  - Difficult to interpret by eye
- (c) Time-frequency representation of (b)
  - History is obvious in the spectrogram



## The PDV approximation

Suppose velocity changes slowly over some small duration.

$$x(t) \approx x(\bar{t}) + \bar{v} \times (t - \bar{t})$$

The optical signal in this duration would be harmonic:

$$I(t) = I_1 + I_2 + 2\sqrt{I_1 I_2} \cos \left[ \bar{\Phi} + 2\pi \left( \frac{2\bar{v}}{\lambda_0} \right) t \right]$$

with a beat frequency proportional to velocity.

This frequency can be determined with a short-time Fourier transform (STFT).

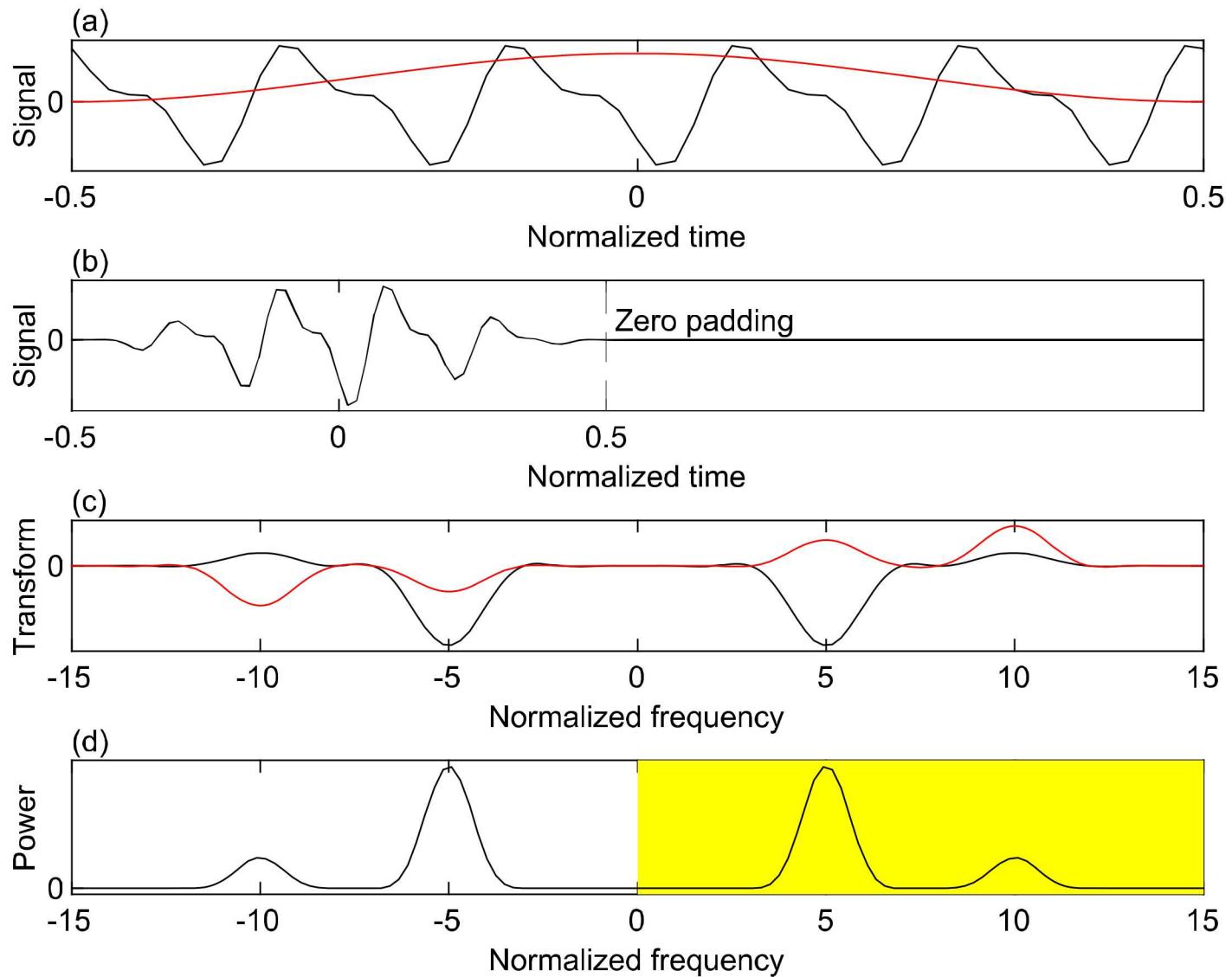
$$S(f, \bar{t}) = \int_{-\infty}^{\infty} s(t) w(t - \bar{t}) e^{-2\pi i f t} dt$$

Window  $w(t)$  selects regions in signal  $s(t)$ .

# Generating a PDV spectrogram

- (a) Extract local region
- (b) Multiple local region by window function and zero pad to a power of two.
- (c) Use FFT to determine complex spectrum
- (d) Calculate the power spectrum for positive frequencies only

Every slice of the spectrogram is built from steps (a)-(d).



## Extracting the velocity history

Region of interest (ROI) selection

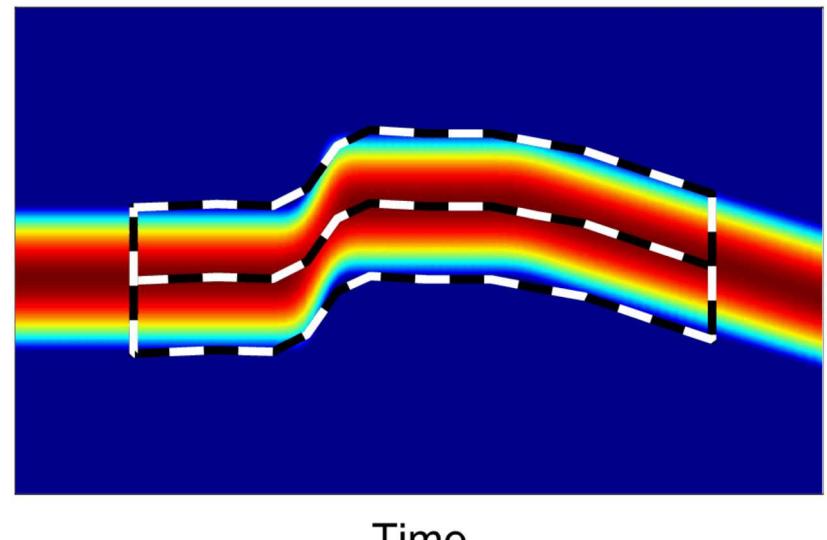
- Usually needs human attention

Determine peak location at each time step,  
subject to ROI

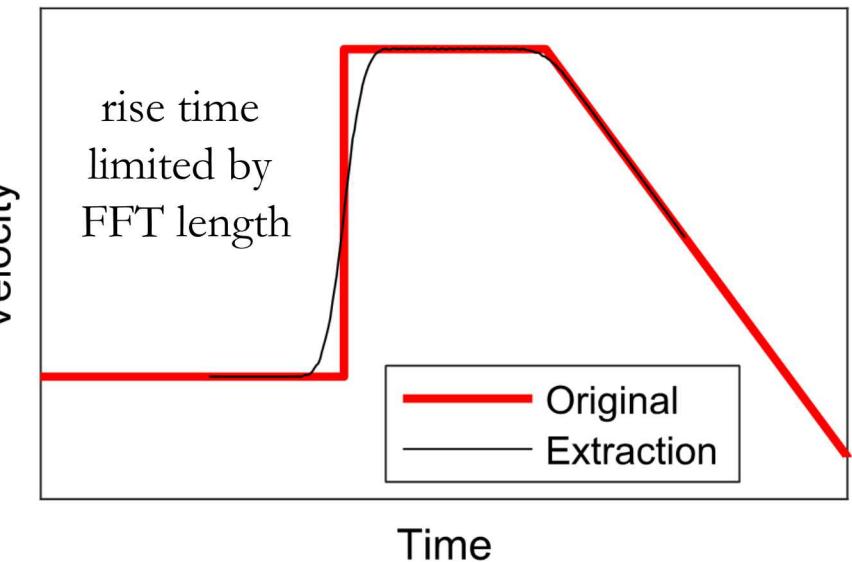
- Centroid, peak fit, etc.
- Subtract reference frequency (as needed)
- Scale by half wavelength

*“Spectrograms don’t lie, but histories sometimes do”*,  
-David Holtkamp

(a)



(b)



# Time resolution versus frequency uncertainty

Standard uncertainty principle applies but is misleading

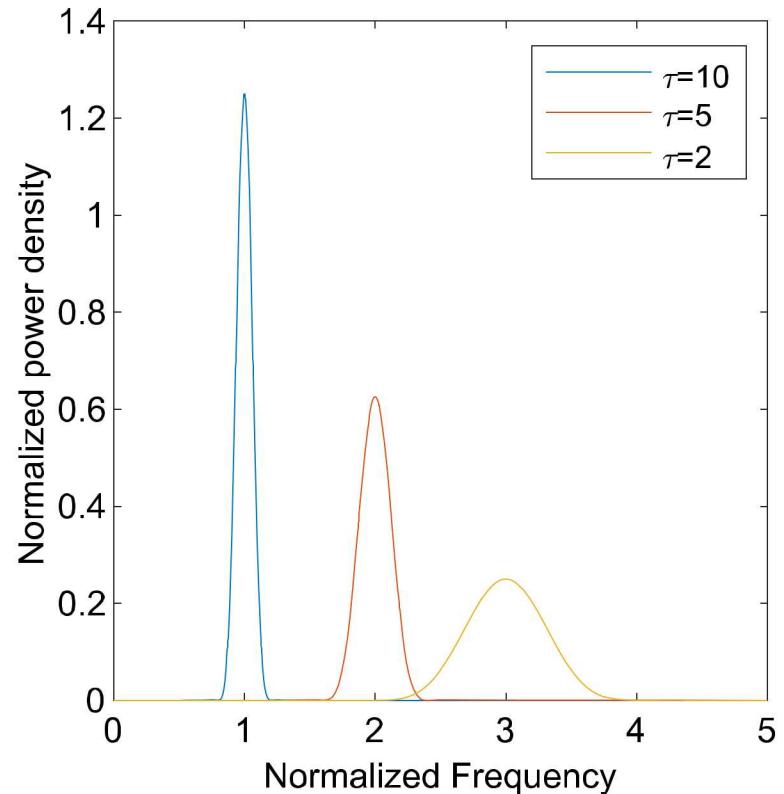
- Time/frequency width product is fixed

Peak locations can be determined more accurately than one width

- Frequency resolution scales with duration to the  $-3/2$  power

Similar limits apply to VISAR

- Etalon delay sets the limiting time scale
- PDV time scale defined in software



$$\sigma_B = \sqrt{\frac{6}{f_s \tau^3}} \frac{\sigma_s}{A} \frac{1}{\pi}.$$

## Limiting resolution in PDV

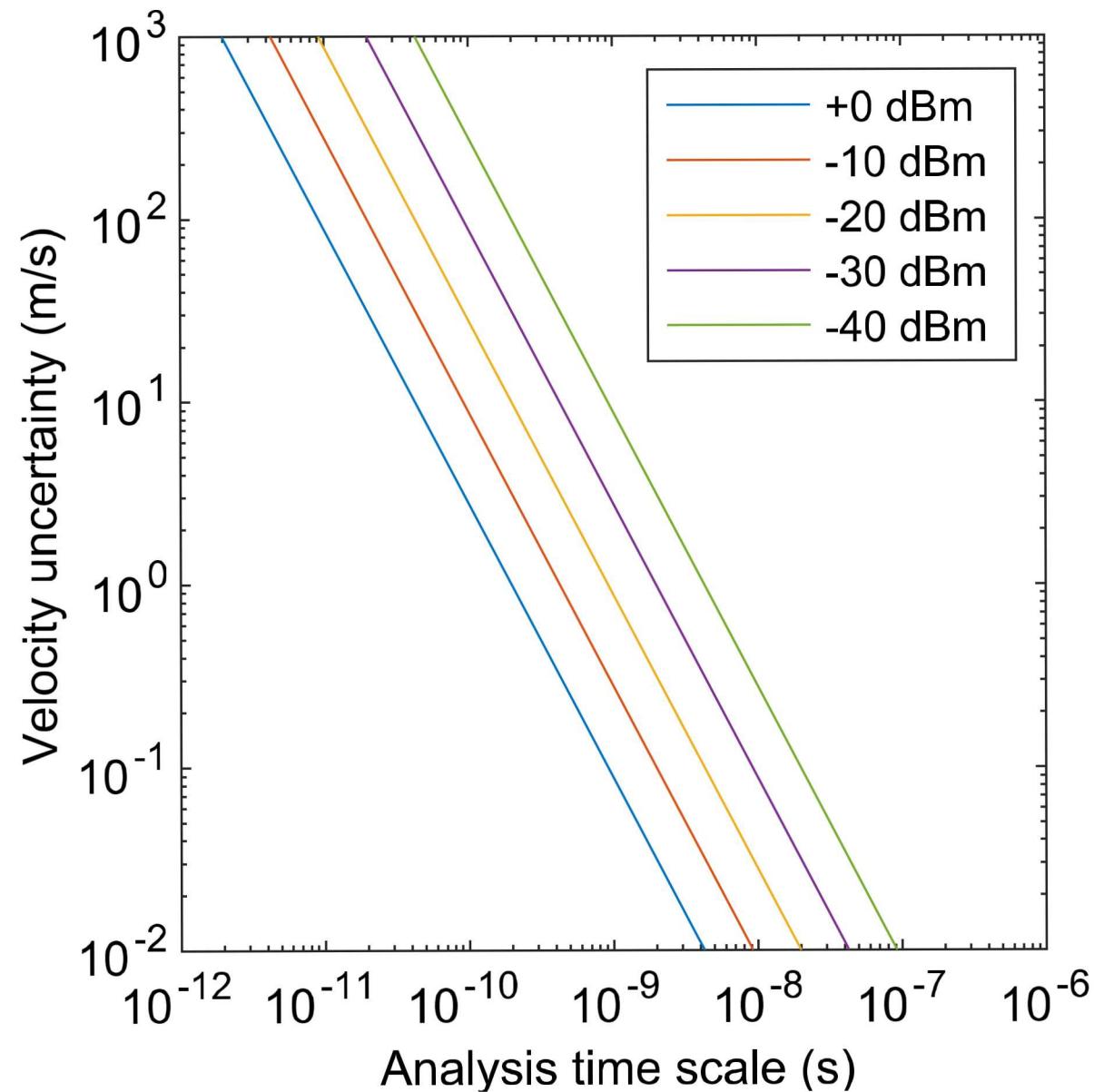
Performance determined by light return

- Usually specified on a log scale
- 0 dBm is 1 mW, -10 dBm is 0.1 mW
- +20 dBm typically sent to probe
- 30 to 60 dB probe return (efficiency)

Time scale plays a significant role

- m/s uncertainty plausible at 1 ns
- Sub m/s uncertainty trivial  $>100$  ns

This assumes well-separated spectra!

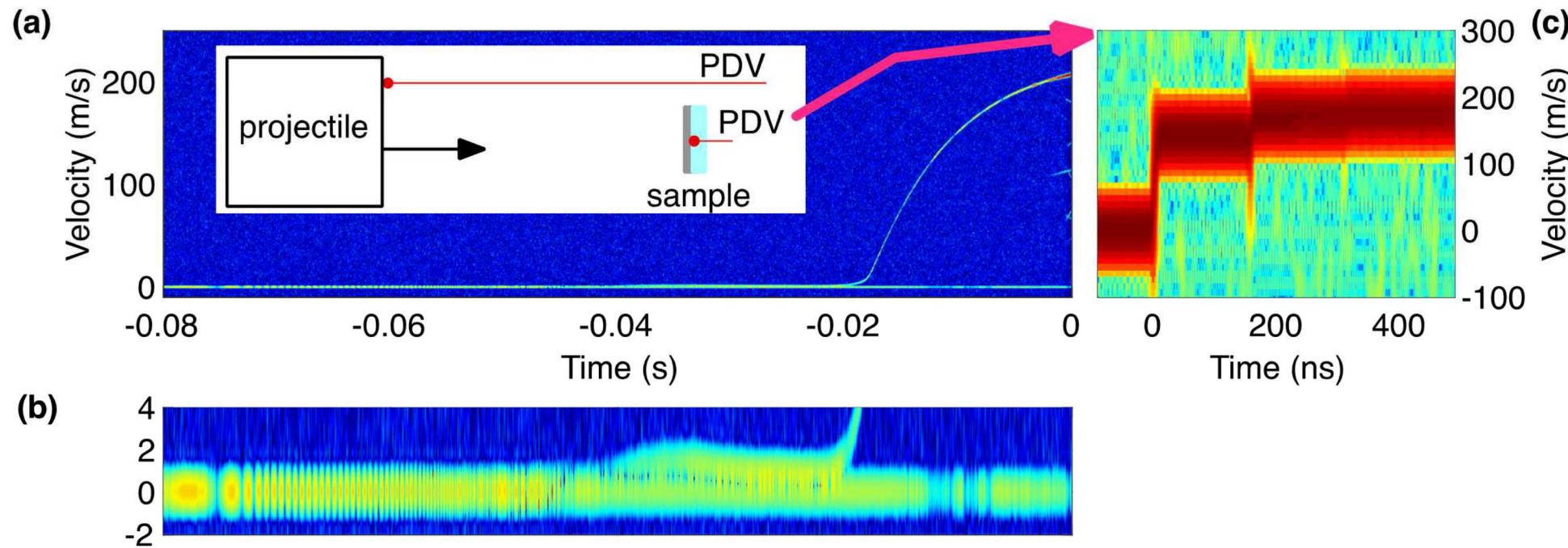




# PDV examples



# Impact experiments

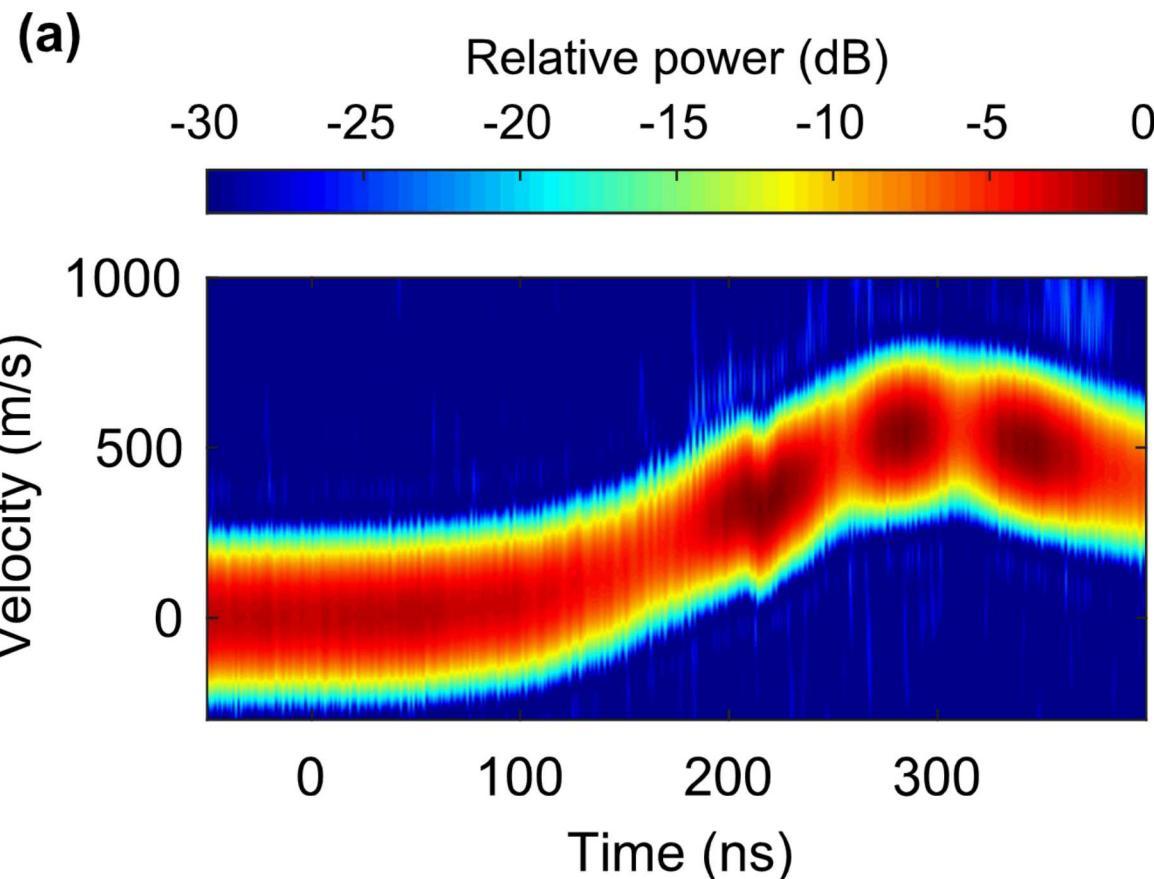


- (a) Spectrogram for projectile launch in a wrap-around gas gun
- (b) Early details show the projectile creeping through the breech before takeoff
- (c) Reverberating shock measurement of the sample during impact

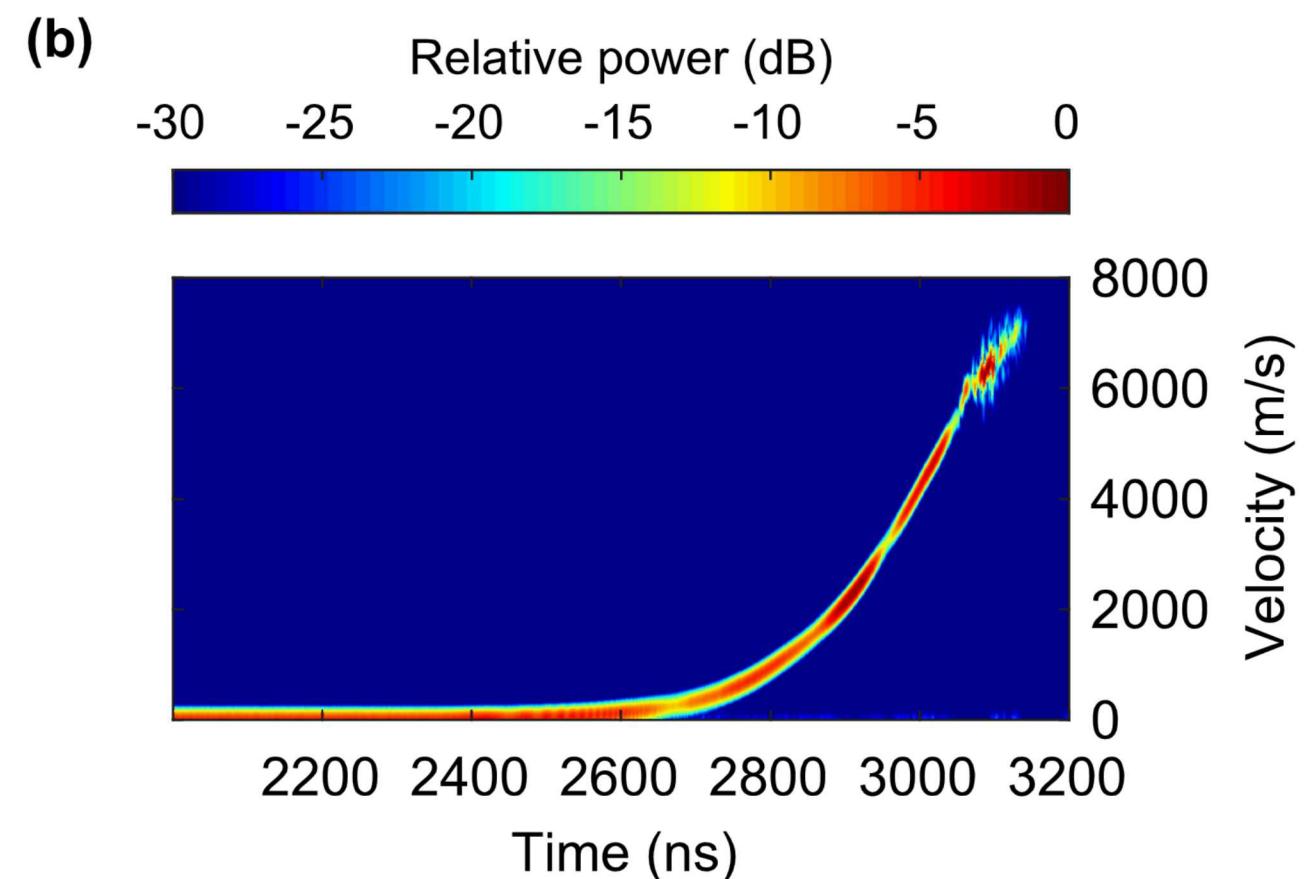
# Ramp wave measurements



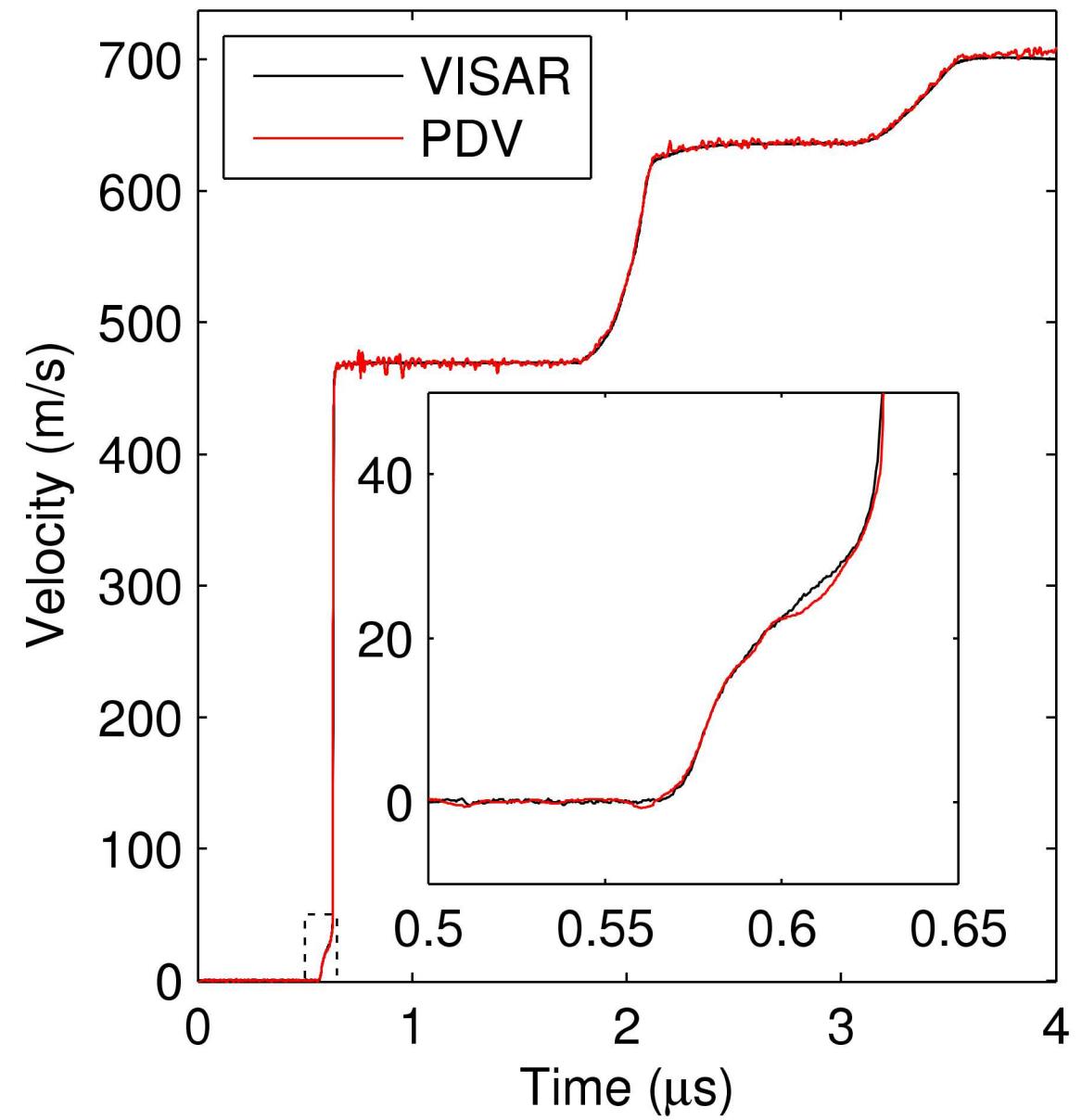
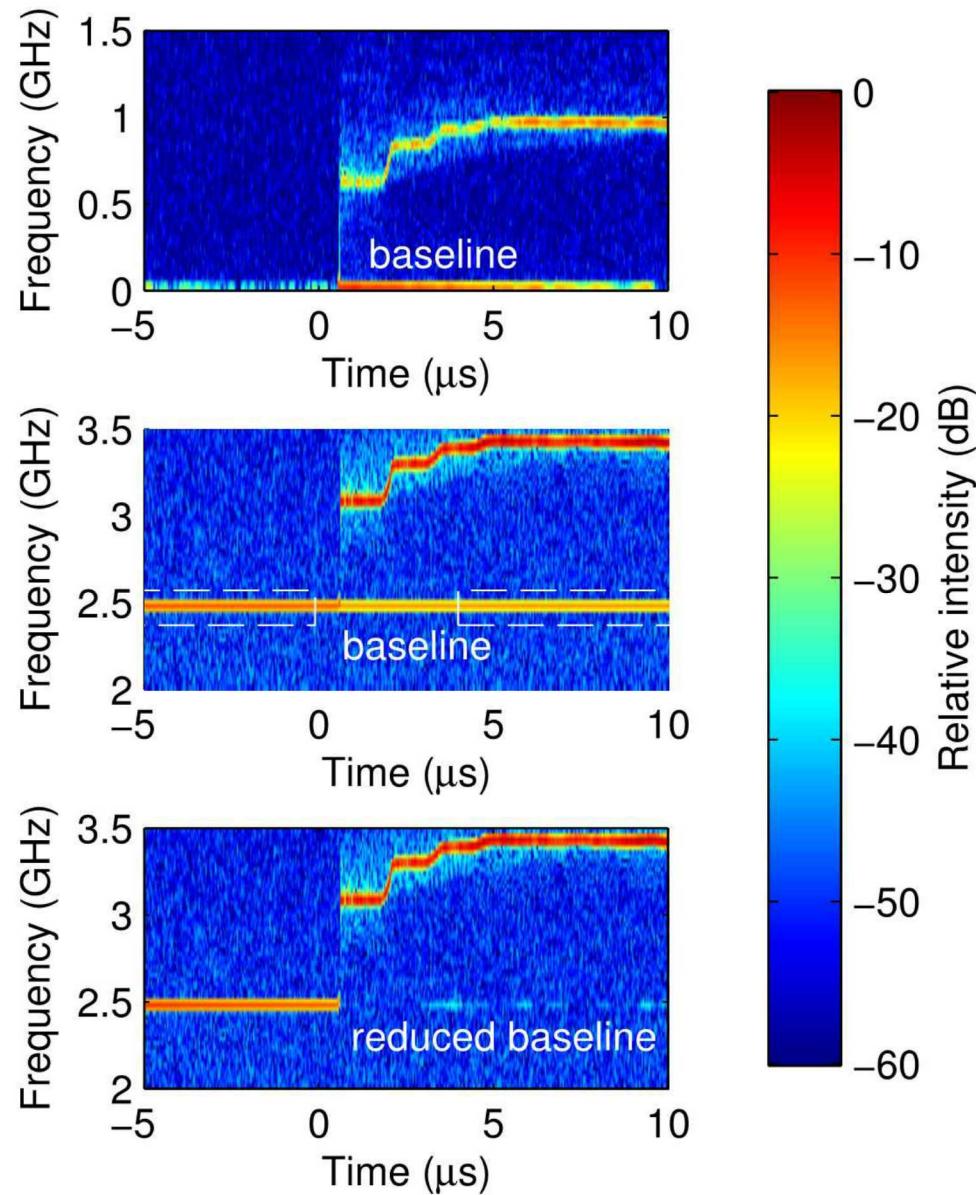
Ramp compression of a water sample



Implosion of a copper cylinder

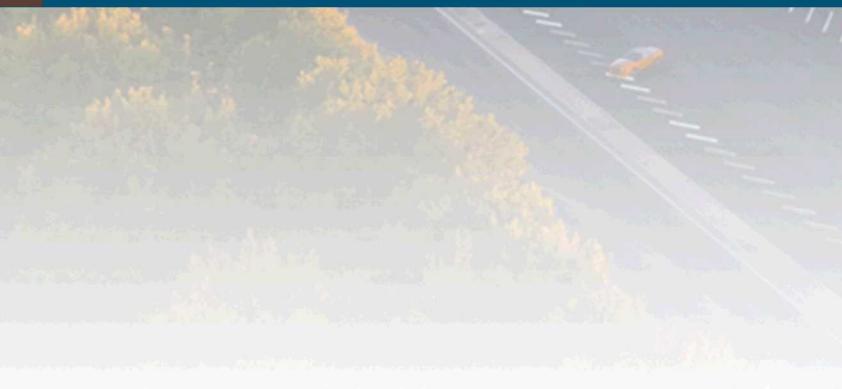


# Velocity history example





# VISAR versus PDV



# Differences between VISAR and PDV

## VISAR

- Open beam or multi-mode fiber
- Usually visible light (532 nm)
- Max. velocity independent of bandwidth
- Time scale defined in hardware
- Usually analyzed in the time domain
- Optimized for single-velocity measurements

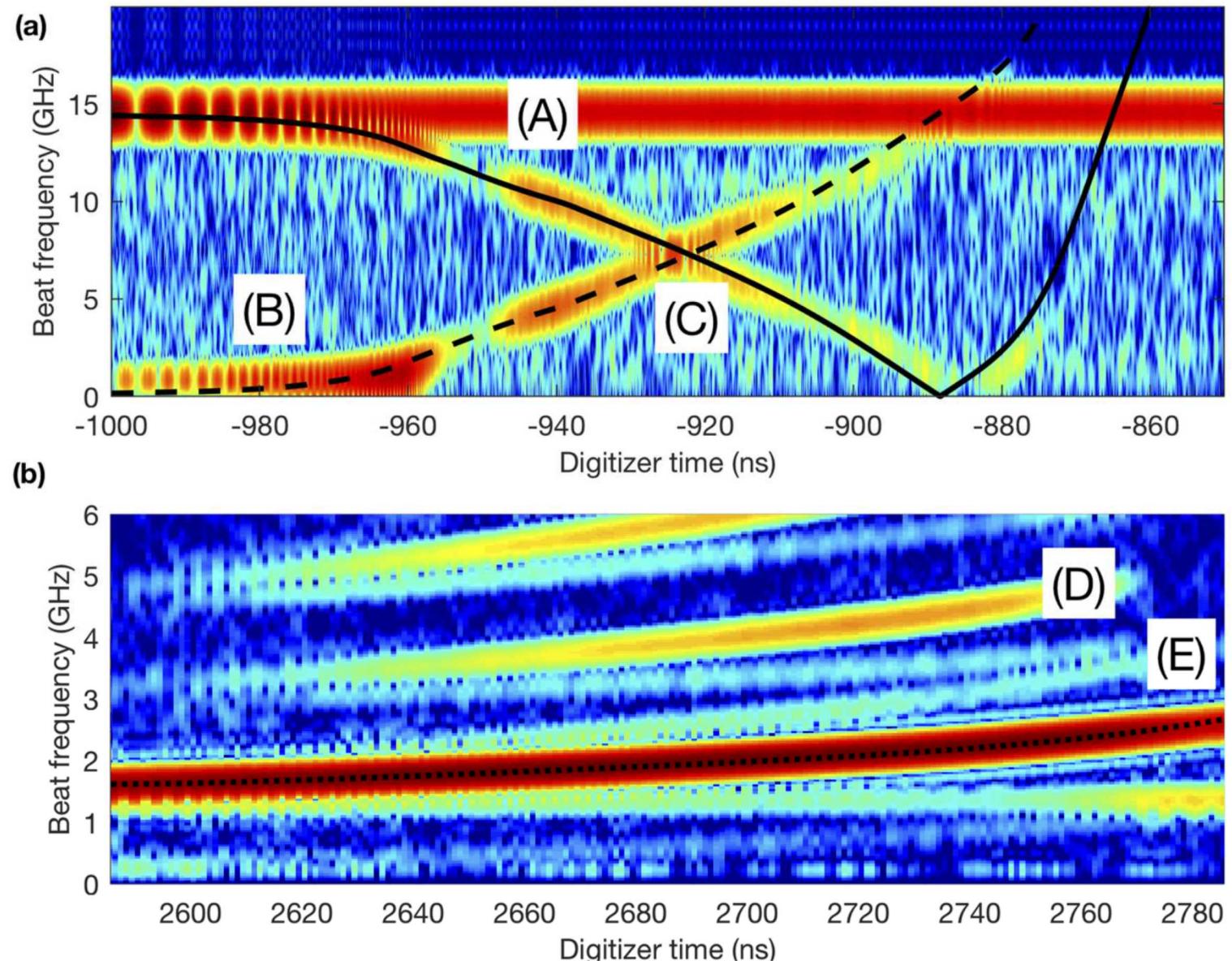
## PDV

- Single-mode fiber
- Usually infrared (1530-1565 nm)
- Max. velocity depends on bandwidth
- Time scale defined in software
- Usually analyzed in the frequency domain
- Can measure multiple simultaneous velocities

# PDV analysis can be tricky

- (a) Measurement with extended baseline (A), conventional artifact (B), and the down-shifted beat of interest (C).
- (b) Measurement with electrical (D) and optical (E) harmonics.

Overlapping features create spectrogram modulations



# Other PDV challenges

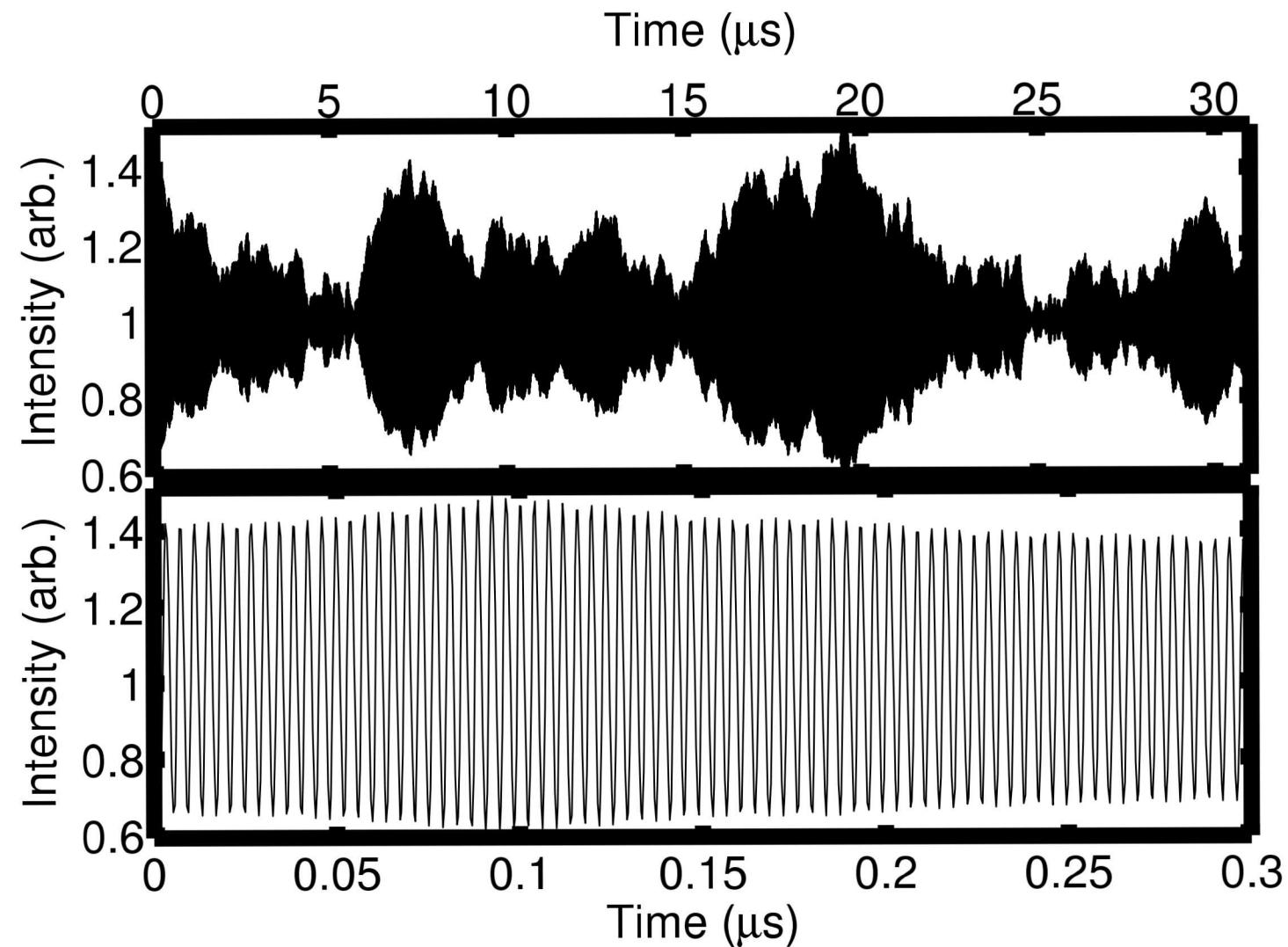


## Dynamic speckle

- Highly coherent lasers + rough reflectors
- Random dropouts
- Can be mitigated with redundant probes

## Fiber limitations

- 9  $\mu\text{m}$  core limits max power
- Stimulated Brillouin scattering limits power over long lengths



## So which is better?

VISAR is technically more sensitive but also more complicated

- Shorter wavelength and multiple signals (reduces noise)
- Open beam alignments and multiple signals (more digitizer channels)

...but both VISAR and PDV have their merits

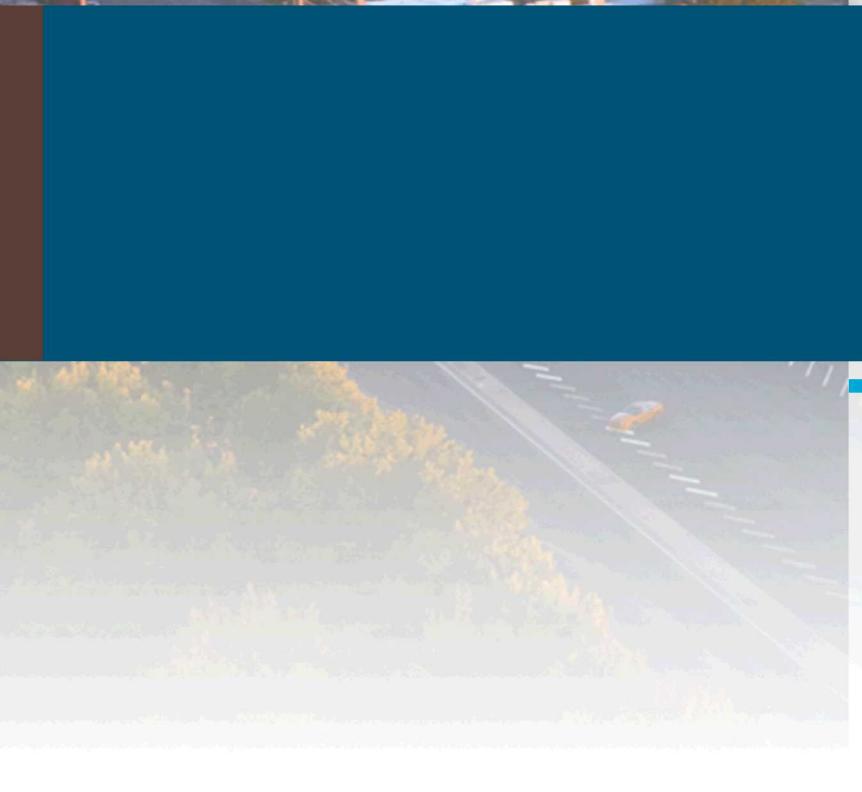
- Use what you have for wave profiles (windowed samples, 1D compression)
- Use PDV for everything else (detonation studies, ejecta measurements,...)

In a new facility, VISAR may not make sense for velocities under 15 km/s

- Digitizer dominates the upfront cost of PDV (can easily exceed \$100,000)
- PDV operation/maintenance are generally much easier than VISAR



# Summary and exercises



## Summary

Velocimetry is a core diagnostic for dynamic compression research

- Virtually every experiment will have VISAR and/or PDV
- Directly tied to the jump conditions and wave simulations
- Historical data based on arrival time measurements

Optical velocimetry now usually based on the Doppler shift

- VISAR encodes that shift as phase on a set of quadrature signals
  - Requires 3-4 measurements, but not high bandwidth
  - Does not tolerate more than one velocity at a time
- PDV encodes that shift as beat frequency change (only 1 signal needed)
  - Only 1 measurement needed, possibly high bandwidth
  - Can handle multiple velocities, though overlapping features are hard to analyze

## Exercise 1



Suppose that you have a 1550 nm PDV and need to measure a projectile moving at 1000 m/s. What is the minimum recording bandwidth needed to track this motion?

For same velocity, what is the etalon delay needed for ten fringes in a 532 nm VISAR?

**Bonus question:** what is the limiting velocity resolution for the above PDV measurement assuming 10% signal noise, 80 GS/s sampling, and 10 ns FFTs? How does this compare to VISAR? Hint: Barker's rule of thumb is 1-2% of the fringe constant.

## Exercise II



Suppose a wedged projectile moves horizontally at velocity  $v$ .

- What velocity would a VISAR at location A measure? What about a PDV at location A?
- Optical measurements at location B see different parts of the projectile as it moves by; the illuminated spot gets closer. What velocity would VISAR/PDV measure at this location?
- Does the distinction that VISAR is a “velocity” interferometer and PDV is a “displacement” interferometer play any role here?

