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¹ **A Variance-Based Decomposition and Global
2 Sensitivity Index Method for Uncertainty
3 Quantification: Application to Retrieved Ice Cloud
4 Properties**

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5 Abstract.

6 This study develops a novel uncertainty quantification (UQ) method for
7 cloud microphysical property retrievals using variance-based decomposition
8 and global sensitivity index. In this UQ framework, empirical orthogonal func-
9 tion (EOF) analysis is applied to the U.S. Department of Energy Atmospheric
10 Radiation Measurement (ARM) ground-based observations, which are the
11 inputs for the cloud retrieval studied here. The principal components (PCs)
12 in the EOF expansion are parameterized as random input variables, and hence
13 the input dimension is greatly reduced (up to a factor of 50), allowing large
14 ensemble of random samplings. The EOF expansion improves the accuracy
15 of the uncertainty estimation by taking into account the cross correlations
16 in the input data profiles. This method enables a probabilistic representa-
17 tion of a retrieval process by adding normally distributed perturbations into
18 PCs of sample-means of input data profiles within a time window. There-
19 fore, it effectively facilitates objective validation of climate models against
20 cloud retrievals under a probabilistic framework for rigorous statistical in-
21 ferences. Moreover, the variance-based global sensitivity index analysis, part
22 of this method, attributes the output uncertainties to each individual source,
23 thus providing directions for improving retrieval algorithms and observation
24 strategies. For demonstration, we apply this method to quantify the uncer-
25 tainties of the ARM program's baseline cloud retrieval algorithm for an ice
26 cloud case observed at the Southern Great Plains site on March 9, 2000.

1. Introduction

27 Cloud microphysical properties such as liquid and ice water contents retrieved from
28 ground-based measurements are important geophysical quantities that are often used in
29 developing and evaluating cloud parameterizations in climate models. However, studies
30 have shown that large differences and uncertainties exist in ground-based cloud retrievals
31 [Comstock *et al.*, 2007; Turner *et al.*, 2007; Zhao *et al.*, 2012; Huang *et al.*, 2012]. The
32 retrieval uncertainties are primarily caused by uncertainties in the retrieval theoretical
33 bases, assumptions, input data, and constraint parameters as indicated in these studies.
34 Quantitative knowledge about the retrieval uncertainties has thus been long desired by
35 the climate modeling community to better constrain model-produced cloud properties
36 [Xie *et al.*, 2005; Xu *et al.*, 2005; Xie *et al.*, 2011].

37 A traditional way to estimate uncertainty is to randomly perturb input data profiles and
38 several key retrieval parameters used in a single cloud retrieval technique. The standard
39 deviation from the ensemble mean of the perturbed retrieval is considered as a proxy of
40 the uncertainty [Zhao *et al.*, 2014]. The other way to estimate uncertainty is to calculate
41 the mean and standard deviation from multiple unperturbed cloud retrievals based on
42 different retrieval techniques and ground-based remote sensors [Comstock *et al.*, 2007]. In
43 recent years, several studies estimate retrieval uncertainties through a radiative transfer
44 model and apply Bayesian calibration to statistically compare the surface and top-of-
45 atmosphere (TOA) radiative fluxes and other properties to observations [Posselt *et al.*,
46 2008; Comstock *et al.*, 2013]. One can further apply multi-model Bayesian model selection,

⁴⁷ and model discrepancy techniques to mitigate the uncertainty estimated by multi-retrieval
⁴⁸ algorithm [*Määttä et al.*, 2014].

⁴⁹ The above methods suffer from but are not restricted to the following limitations: (1)
⁵⁰ the cross correlations in the input data profiles are often not considered; (2) parame-
⁵¹ terizing input data profiles with appropriate cross correlations is not an obvious task;
⁵² (3) sampling random variables amongst vertical layers, which could be on the order of
⁵³ hundreds depending on the vertical resolution, requires an enormous, infeasible sampling
⁵⁴ size; (4) characterizing probability density functions (PDFs) of the random variables of-
⁵⁵ ten require unrealistic statistical hypotheses; (5) attributing the variability in the retrieval
⁵⁶ output to that in each individual uncertainty source (i.e., global sensitivity analysis) is
⁵⁷ not permitted in general; and (6) differences between measurements and the truth (i.e.,
⁵⁸ bias analysis) are usually not considered.

⁵⁹ To address these issues, we propose an uncertainty quantification (UQ) and sensitivity
⁶⁰ analysis methodology based on Karhunen-Loéve expansion (KLE), Central-Limit Theo-
⁶¹ rem (CLT), and Sobol' indices. The KLE [*Kuhunen*, 1947; *Loéve*, 1945] is a principal
⁶² component analysis (PCA) [Wilks, 2011], which transforms a number of possibly corre-
⁶³ lated variables into a smaller number of uncorrelated variables called principal components
⁶⁴ (PCs) through the empirical orthogonal function (EOF) expansion. For the first issue,
⁶⁵ the application of the EOF expansion to ground-based cloud measurements such as those
⁶⁶ from the U.S. Department of Energy (DOE) Atmospheric Radiation Measurement (ARM)
⁶⁷ program allows us to obtain the cross correlations of input data profiles within a given
⁶⁸ time window. In this study, we use a 0.5-hour time window, comparable to the typical
⁶⁹ climate model temporal resolution.

70 For the second issue, a stochastic representation of input data profiles is constructed
71 by using the PCs of the EOF expansion as input random variables, associated with the
72 extracted observational covariance matrix within the 0.5-hour time window. Hence the
73 uncertainty is automatically embedded in the EOF expansion by adding appropriate per-
74 turbations into the PCs. Since random perturbations are added into the PCs instead
75 of into each vertical layer of the input data profiles, the dimensionality of the stochastic
76 space (issue (3)) is significantly reduced so that it reduces the sampling size required to
77 stabilize the statistics of the cloud retrieval.

78 Since the target of this study is the 0.5-hour sample-mean of the stochastic data profile,
79 the CLT [Cramér, 1946; Gnedenko *et al.*, 1954; Storch and Zwiers, 2002] can be leveraged
80 to solve the issue (4). Based on the CLT, the normally distributed perturbations are thus
81 added on the extracted input random variables (proof is provided in Appendices A and
82 B). To address the issue (5), we apply the variance-based global sensitivity index analysis
83 [Sobol, 1993] to the retrieval algorithm and attribute the uncertainties of vertically resolved
84 retrieval output to the input random variables as well as retrieval parameters. These
85 sensitivity indices can provide insights for improving retrieval algorithms and observation
86 strategies.

87 In summary, we employ the probabilistic PCA to enable the stochastic cloud retrieval by
88 adding observation-based perturbations to the PCs of the EOF expansion of the sample-
89 mean of input data profiles following normal distributions per CLT. The variance-based
90 sensitivity analysis attributes the vertically resolved retrieval output uncertainties to each
91 individual source. This variance-based UQ method effectively facilitates objective valida-
92 tion of climate models against cloud retrievals under a probabilistic framework for rigorous

₉₃ statistical inferences. It should be noted that the retrieval model structural uncertainty
₉₄ and the bias are not addressed in the current work.

₉₅ The structure of the paper is as follows. In Section 2, details of the probabilistic PCA
₉₆ based uncertainty analysis in cloud microphysical property retrievals are given. In Sec-
₉₇ tion 3, the capability of the UQ method is illustrated with an ice cloud case using the
₉₈ ARM baseline cloud microphysical algorithm (MICROBASE). Results from our uncer-
₉₉ tainty analysis and sensitivity study are shown in Section 4, followed by conclusions and
₁₀₀ discussions in Section 5.

2. Methodology

₁₀₁ The EOF analysis is a variance-based statistical technique designed for decomposition
₁₀₂ of time series in terms of orthogonal basis functions that are determined from the empi-
₁₀₃ rical data. The orthogonal basis functions are chosen to account for as much as variance
₁₀₄ of the empirical data as possible. In this paper, the EOF analysis is applied to the
₁₀₅ ARM measurements required as the MICROBASE inputs with 0.5-hour interval. The
₁₀₆ realizations of the random variables in the EOF expansion are computed by projecting
₁₀₇ empirical ARM measurements on the orthogonal basis functions. In general, the proba-
₁₀₈ bilistic distributions of the random variables cannot be determined by these realizations.
₁₀₉ According to the CLT, however, the sample-means of these random variables are normally
₁₁₀ distributed when the number of measurements is large enough (a sampling size greater
₁₁₁ than 30 is generally considered as large enough) within the time window. Therefore, it en-
₁₁₂ ables the probabilistic representation of a stochastic retrieval process by adding normally
₁₁₃ distributed perturbations to the PCs of the EOF representation of the sample-means of
₁₁₄ the input data profiles (see Appendix B).

115 To demonstrate the probabilistic PCA based method, we start with a temporal-spatial
 116 stochastic process denoted as $Y(\mathbf{x}, t, \theta)$ to represent stochastic input data profiles for the
 117 retrieval algorithm, where \mathbf{x} denotes the height, t the time, and θ a random event. For
 118 example, $Y(\mathbf{x}, t, \theta)$ can be referred as the radar reflectivity profile that will be described
 119 in Section 3. Accordingly, an ensemble of snapshots of the stochastic process $Y(\mathbf{x}, t, \theta)$
 120 observed in the time window $[0, T]$ can be recorded as

$$\{y_1, y_2, \dots, y_n\}, \quad (1)$$

121 where $y_i(\mathbf{x}) = y(\mathbf{x}, t_i)$, $i = 1, \dots, n$, n is the number of snapshots; and the ensemble
 122 average of the snapshots can be defined as $\bar{y}(\mathbf{x}) = \frac{1}{n} \sum_{i=1}^n y_i$.

123 With measurement noises added to the stochastic process $Y(\mathbf{x}, t, \theta)$, we have a per-
 124 turbed input data profile denoted as $Y'(\mathbf{x}, t, \theta)$ and it can be written as

$$Y'(\mathbf{x}, t, \theta) = Y(\mathbf{x}, t, \theta) + \text{noise}. \quad (2)$$

125 The stochastic process $Y(\mathbf{x}, t, \theta)$ can be decomposed into the ensemble average $\bar{y}(\mathbf{x})$,
 126 and an intrinsic unknown random estimation error $\epsilon(\mathbf{x}, t, \theta)$, such that

$$Y(\mathbf{x}, t, \theta) = \bar{y}(\mathbf{x}) + \epsilon(\mathbf{x}, t, \theta). \quad (3)$$

127 Therefore, the perturbed stochastic process $Y'(\mathbf{x}, t, \theta)$ can be decomposed as

$$Y'(\mathbf{x}, t, \theta) = \bar{y}(\mathbf{x}) + \epsilon(\mathbf{x}, t, \theta) + \text{noise}. \quad (4)$$

128 One goal of this paper is to quantify the cloud retrieval uncertainties for climate model
 129 evaluation. The typical climate model temporal resolution is currently $\sim 1 \text{--} 10 \text{--} 100 \text{--}$

130 Thus, the perturbed input data profile $Y'(\mathbf{x}, t, \theta)$ is transformed to a smoother statistic
 131 $\bar{Y}'(\mathbf{x}, t, \theta)$, which is the sample-mean of $Y'(\mathbf{x}, t, \theta)$ within the 0.5-hour time window. Given
 132 a time window, the statistic $\bar{Y}'(\mathbf{x}, t, \theta)$ is also a random variable and it approximately
 133 follows a normal distribution when the sampling size within the time window is large
 134 enough (large number law) [Storch and Zwiers, 2002].

135 Due to the high dimensionality of the stochastic space for $Y'(\mathbf{x}, t, \theta)$ (e.g., 512 vertical
 136 layers in the ARM radar reflectivity profiles), it is computationally infeasible to sample all
 137 the vertical layers individually. To reduce the dimensionality, we apply the EOF expansion
 138 to represent the perturbed stochastic process $Y'(\mathbf{x}, t, \theta)$ in terms of eigenfunctions of its
 139 correlation kernel assuming that it is piece-wise constant within the 0.5-hour time window.
 140 The detailed derivations can be found in Appendices A and B.

141 By applying the CLT to the ARM ground-based observations, in Appendix B we show
 142 that random variables appeared in the EOF expansion of $\bar{Y}'(\mathbf{x}, t, \theta)$ approximately follow
 143 normal distribution when the sampling size is large enough. By truncating EOF expansion
 144 of $\bar{Y}'(\mathbf{x}, t, \theta)$ to the order of M , we finally arrive at $\bar{Y}'(\mathbf{x}, t, \theta)$ that is the sample-mean of
 145 the perturbed data profile with white noises added and it can be written explicitly as

$$\bar{Y}'(\mathbf{x}, t, \theta) = \bar{y} + \sum_{i=1}^M \psi_i \sqrt{\frac{\lambda_i}{n}} \sqrt{1 + \left(\frac{\sigma_0}{\sqrt{\frac{\lambda_i}{n}}} \right)^2} \frac{z_i}{\sqrt{n}} + \text{error}(\mathbf{x}, t, \theta), \quad (5)$$

146 where $z = [z_1, z_2, \dots, z_M]^T$ follows a standard multivariate normal distribution, i.e., $z \sim$
 147 $\mathcal{N}(0, \mathbf{I}_M)$ and \mathbf{I}_M is a $M \times M$ identity matrix; (λ_i, ψ_i) are corresponding pairs of eigenvalues
 148 and eigenfunctions; σ_0 denotes the standard deviation of normally distributed random
 149 measurement noises; $\text{error}(\mathbf{x}, t, \theta)$ is the truncation error. The detailed proof is given in

150 the Appendix B. To simplify, other than observation-based input data profiles, uniform
 151 distributions are applied to perturb the retrieval parameters.

152 We apply Sobol's method to derive the global sensitivity analysis of microphysi-
 153 cal properties retrieved by MICROBASE. Sobol's method is a variance-based sensi-
 154 tivity analysis method, which divides the variance $\text{Var}(\bar{Y})$ into fractions attributed
 155 to each input X_i (first-order sensitivity or main effect indices $S_i = \frac{V_i}{\text{Var}(\bar{Y})}$ where
 156 $V_i = \text{Var}_{X_i}(E_{X \sim i}(\bar{Y} | X_i))$ defined as average over variations in other random inputs
 157 or parameters), and their interactions (second-order sensitivity indices S_{ij} or higher-order
 158 indices formed by dividing other terms in the variance decomposition). The fractions
 159 measure the contribution to the output variances of each input variable, including all
 160 interactional variances with any other input variables in all the orders. The sum of all
 161 the Sobol's indices equals to one. Also, Latin Hypercube Sampling (LHS) procedure is
 162 used to draw samples in the designed space for the input random variables and retrieval
 163 parameters. LHS is an effective stratified sampling approach in a high-dimensional space
 164 ensuring that all portions (with equal probability) of a given partition are sampled [McKay
 165 *et al.*, 1979].

3. Case Study

166 For demonstration, we apply the probabilistic PCA to propagate uncertainties from
 167 ARM ground-based measurements as well as empirical parameters used in MICROBASE
 168 into its retrieved products for uncertainty quantification and analysis. MICROBASE is
 169 the ARM base-line cloud microphysical property retrieval algorithm based on the cloud
 170 radar and lidar measurements [Dunn *et al.*, 2011; Zhao *et al.*, 2014]. It derives the cloud
 171 liquid and ice properties using empirical regression equations obtained from in situ aircraft

172 measurements. Liquid water content (LWC) and ice water content (IWC) are derived
 173 from radar reflectivity measured by Millimeter Wavelength Cloud Radar (MMCR) with
 174 a frequency of 35 GHz. In MICROBASE, LWC is retrieved by

$$LWC = LWP \frac{Ze_{Liq}^g}{\sum_{i=1}^n Ze_{Liq}^g \Delta Z}, \quad (6)$$

175 whereas for pure ice clouds, IWC is retrieved by

$$IWC = aZe_{Ice}^b. \quad (7)$$

176 In the equations above, a , b and g are empirical parameters, and ΔZ is the vertical
 177 increment. LWP represents the liquid water path, while Ze_{Liq} and Ze_{Ice} represent the ef-
 178 fective radar reflectivity profile for the liquid and ice, respectively. These Z-IWC empirical
 179 parameters in the retrieval algorithm are determined with certain assumptions about the
 180 ice particle size distribution, ice particle shape, and density [Liu and Illingworth, 2000].

181 Uncertainties in the retrieved quantities come from three sources: input profiles, the re-
 182 trieval algorithms, and parameter assumptions as described in Zhao *et al.* [2012, 2014].

183 The uncertainty of input profiles is present in terms of two components, bias (related to
 184 accuracy measuring difference between measurements and truth), and the unavoidable
 185 random variations in measurements (related to precision). The bias of input data profiles
 186 and retrieval model structural uncertainty are not considered in this research article.

187 For the selected case study, we apply our probabilistic PCA based method to quantify
 188 uncertainties in MICROBASE retrieved ice for the high cirrus cloud case observed at the
 189 ARM SGP CF site on March 9, 2000 during the year 2000 cloud intensive observational
 190 period. This cirrus cloud case has been studied comprehensively by Comstock *et al.*

[2007] to examine ice cloud properties from 15 state-of-art cloud retrievals. As described by *Comstock et al.* [2007], the cirrus clouds were associated with the passage of a weak upper-level disturbance over the SGP region and deepened as the disturbance moved northeastward. Accordingly, optically thin ice clouds were observed initially (19:00–19:15 UTC) and between 22:00 and 22:30 UTC. The majority of the observed clouds were optically thick clouds over the 3.5-hour time period as shown in Fig. 1a. *Comstock et al.* [2007] shows large uncertainties in the retrieved ice cloud properties among the tested algorithms for both optically thin (optical depth, $\tau < 0.3$) and thick ($0.3 < \tau < 5.0$) cirrus clouds.

The detailed procedures of applying the probabilistic PCA based method to MI-CROBASE are described as the following. First, in order to efficiently represent uncertainty of input radar reflectivity profiles, we apply the PCA analysis to reduce the dimensions from 512 layers to 10 modes (the first four EOFs or modes of the PCA are shown in Fig. 2) to capture greater after 90% variance in the observed radar reflectivity profiles (Fig. 1a) and extract uncorrelated, independent random variables with orthogonal modes. For this test case, we assume that the correlation kernel of the stochastic radar reflectivity profile is piece-wise constant within each 0.5-hour time window. The details of computing eigenvalues and EOFs based on snapshots taken for an ensemble of relative errors can be found in Appendix A. As a result, we expand the input radar reflectivity profiles in terms of pairs of obtained eigenvalues and EOFs combined with associated PCs. The probabilistic distributions of these PCs in the EOF expansion of the stochastic input radar reflectivity profiles are generally unknown.

213 The stochastic radar reflectivity profile is transformed to a smoother statistic, which
214 is the sample-mean of the observed ones within the 0.5-hour time window. Within each
215 0.5-hour time window, we have 180 times of observations (sampling size > 30). Based on
216 the CLT, the sample-mean of the radar reflectivity profile is thus expanded in terms of
217 independent and normally distributed random variables. The perturbation ranges of the
218 input data and empirical parameters follow *Zhao et al.* [2014] and the cloud retrieval and
219 measurement experts' suggestions (personal communications). The measurement noise
220 of the radar reflectivity profile for each layer, denoted as σ_0 , is chosen to be 1.0 dBZ.
221 The range of the empirical parameter "a" is 0.03–0.22 (g/m³) /dBZ, while the empirical
222 parameter "b" is assumed to be a dimensionless constant 0.59. The value 0.59 is chosen
223 by the original MICROBASE algorithm for the parameter "b". We opt not to perturb
224 the parameter "b", as the main purpose of this paper is to demonstrate the capacity of
225 our UQ method instead of thoroughly exploring the uncertainties in the Z-IWC empirical
226 relationship. Based on Eq. (5), the normally distributed perturbations are added on
227 the independent random variables for the sample-mean of the radar reflectivity profile.
228 Uniform distributions are assumed for the parameter "a".

229 We utilize the Problem Solving environment for Uncertainty Analysis and Design Explo-
230 ration toolkit (PSUADE) [Tong, 2009] to perform the uncertainty and sensitivity analysis.
231 PSUADE is a software toolkit for performing uncertainty analysis, responsive surface anal-
232 ysis, global sensitivity analysis, design optimization, model calibration, with large number
233 of parameters and complex constraint. The samples of the random variables including PCs
234 in the EOF expansion and empirical parameters in MICROBASE are obtained by uniform
235 LH sampling using PSUADE. Based on Eq. (5), the PCs in the EOF expansion follow

236 normal distributions. Therefore, the uniformly generated samples are thus converted to
237 normally distributed ones for the PCs. Based on the uncertainty analysis performed by
238 PSUADE, Figs. 1b and 1c compare the 0.5-hour averages of IWC from the original (un-
239 perturbed) MICROBASE and probabilistic PCA based ensemble means (ensemble size
240 = 1000) of 0.5-hour sample-mean of IWC from the perturbed MICROBASE. Note that
241 the temporal resolution of the original MICROBASE retrievals is 10s. The probabilistic
242 PCA based ensemble means are of the same degree of magnitude as the original 0.5-hour
243 ensemble means, but generally greater for thick clouds. The probabilistic PCA based
244 standard deviation (STD, see Fig. 1d) is about 1/5 of the corresponding mean value in
245 this case.

246 Using our developed UQ methodology, the average (min, max) values of the ice water
247 path (IWP, unit: g/m²) retrieved by MICROBASE are 25.4 (0.8, 119.4), respectively.
248 IWP is derived consistently (i.e., integration over all the layers including the cloudless
249 ones) for different approaches. Accordingly, cloudless layers are included when calculating
250 0.5-hour sample-mean of IWP derived from different approaches. The range is about
251 a factor of 2 greater than the average numbers from 14 different retrievals shown in
252 Table 2 of *Comstock et al. [2007]* (16.4 (0.076, 63.3)). We choose a large perturbation
253 range (0.03–0.22 (g/m³) /dBZ) of parameter “*a*” to cover various ice cloud conditions
254 rather than the one-day case here as the goal is to quantify uncertainties in long-term
255 ARM cloud retrievals. Despite the amplified parametric uncertainty, our IWP range
256 falls into the individual retrieval range in *Comstock et al. [2007]*. It highlights the fact
257 that propagating the uncertainties in the input data as well as the parameters through a
258 single retrieval (i.e., MICROBASE) leads to the uncertainties in the output comparable

259 to the differences amongst different retrievals, many of which are rooted from different
260 theories/hypotheses and even based on different instruments. This implies that it might
261 be possible to partly reconcile different algorithms by understanding the causes of the
262 uncertainty in one of them. Through the variance-based sensitivity analysis performed by
263 PSUADE, it is found that the parameter “*a*” is mainly responsible for the variability of
264 the IWP retrieved by MICROBASE in this one-day case (see the Sobol’s global sensitivity
265 analysis section below for more details). Thus, the retrieval differences may be largely
266 caused by how differently the parameter “*a*” is assumed (or implicit assumptions about
267 the size, shape, and density of the target ice particles) by different algorithms.

268 Comparisons with independent observations (e.g., aircraft) provide another way to in-
269 terpret our method. Figure 1e compares the IWP from the counterflow virtual impactor
270 (CVI) (black line) [Twohy *et al.*, 1997] observation on the aircraft, original MICROBASE
271 (red line), and our results (blue line). In general, the averages of in situ CVI measure-
272 ments are greater than both retrievals and they agree within a factor of two, which has
273 been revealed from a dozen of state-of-the-art retrievals comparisons (see Comstock *et*
274 *al.*, 2007 Fig. 5a). The differences between observations and retrievals are partly due to
275 different sampling volumes, instrument uncertainties, sensitivities, and limitations [Com-
276 stock *et al.*, 2007]. Our probabilistic PCA based ensemble means of sample-means of
277 retrieval products obtained by sampling perturbed MICROBASE are closer to the CVI
278 probe measurements than the averages obtained from the original MICROBASE, which
279 shows some encouraging signs of improving the retrieval results with our UQ method.
280 This improvement is probably because our methodology parameterizes the input mea-
281 surements based on the facts that (1) PCA extracts uncorrelated, independent random

variables with orthogonal modes and the auto-correlation kernel is relatively more stable than instantaneous measurements within the 0.5-hour window; and (2) sample-mean is a smoother statistical variable and follows a normal distribution when the sampling size is large enough per CLT. In other words, targeting at the 0.5-hour observation window we replace the retrieval input of finite observations with normally distributed random fields. Statistically, it is thus more likely that the probabilistic PCA based ensemble means of the retrieved properties are closer to the reality than the ones using original algorithm.

The vertical bars in Fig. 1e are defined differently, but comparable. The CVI bars (black) represent the STDs of the 2-min IWP observations when the aircraft flew over the SGP site. The raw MICROBASE bars (red) depict the 0.5-hour STDs, while those of our results (blue) represent the STDs of sample-means within 0.5 hour window. Three types of bars overlap, which is consistent with *Comstock et al.* [2007]. Both methods generally show smaller uncertainties than the CVI observations, which likely reflects the large discrepancies in the sample volumes between the in situ observations and radar retrievals. However, to fully evaluate the proposed method and compare it with the original MICROBASE, the analysis needs to be expanded from the 1-day case to a longer time period that covers different seasons and cloud conditions.

It is worth noting that the probabilistic PCA based method includes a uniform perturbation from the parameter “ a ”, whereas MICROBASE uses a constant value for the parameter “ a ”. Nevertheless, targeting at the 0.5-hour time window, the uncertainties (standard deviations) of sample-means quantified by applying probabilistic PCA to the perturbed MICROBASE are generally smaller than those computed by the high-frequency original MICROBASE data. This highlights the fact that our probabilistic PCA based

305 method estimates the uncertainty for the sample-mean of N observations within an in-
306 terval chosen at model temporal resolution. The distribution of such statistic has a mean
307 that equals to the interval population mean of the input data profiles and its variance
308 equal to the variance of each instantaneous observation divided by N . Sample-mean is a
309 good statistical estimator of the population mean of the input data profiles within the time
310 window, where a “good” statistical estimator is defined as being efficient and unbiased in
311 a statistical sense.

312 When keeping the parameter “ a ” as a constant, the probabilistic PCA based error bars
313 are, as expected, much smaller than original MICROBASE (see Fig. 3). Great reduction
314 in the probabilistic PCA based uncertainties when fixing parameter “ a ” (see Fig. 1e and
315 Fig. 3) suggests that the parameter “ a ” is the main source of the uncertainty. In the
316 following section, we will apply Sobol’s sensitivity analysis to quantify the parametric
317 uncertainty from the parameter “ a ”, measurement uncertainty from radar profiles, and
318 their possible interactions.

319 Figure 4 displays the box plot of IWC PDFs at 8 km (panel a) and the IWP PDFs
320 (panel b). The retrievals exhibit larger spread in the probability distribution of both
321 IWC and IWP for the optically thick clouds at 21:00–21:30 UTC. The IWP mean and
322 STD at 21:00–21:30 UTC are 65.4 g/m^2 and 28.7 g/m^2 , respectively; while its mean and
323 STD at 22:00–22:30 UTC are 3.6 g/m^2 and 1.6 g/m^2 , respectively. However, the IWP
324 coefficients of variance defined as fraction of STD over mean are 0.4 for both 0.5-hour
325 windows, whereas IWC at 8 km has slightly larger coefficient of variance at 22:00–22:30
326 UTC (0.5) than at 21:00–21:30 UTC (0.4). These results reinforce the needs of quantifying
327 IWC uncertainties on different vertical layers.

4. Sobol's Sensitivity Analysis

328 Using PSUADE, Sobol's sensitivity analysis with bootstrapping [Tong, 2009] is im-
329 plemented by resampling a response surface. Figures 5a–d show the results of Sobol's
330 first and second order sensitivity analysis for IWP (left column) and IWC at 8 km (right
331 column) at 22:00–22:30 UTC. Results for 21:00–21:30 UTC are similar and not shown.
332 It is found that the parameter “ a ” is the major uncertainty source for both IWP and
333 IWC with close to 1.0 Sobol's index (variance-based first-order sensitivity measure). This
334 means that almost 100% of the output variance is caused by the variance in the parame-
335 ter “ a ”, whereas almost no variance of the output is caused by the variances in the radar
336 reflectivity modes or interactions among them. Since the parameter “ a ” in the Z-IWC
337 relationship is determined by the ice particle size, shape, and density, these characteristics
338 need to be better described with more accurate cloud observations. This should be one
339 emphasis area in future measurements.

340 In addition, the first mode of radar reflectivity (Z1) is the second largest uncertainty
341 source (but much smaller than the parameter “ a ”) within the time window 22:00–22:30
342 UTC. There are also small contributions from the interaction between the parameter “ a ”
343 and Z1 (see Fig. 5d). The green points denote the uncertainties in Sobol's index due to
344 statistical errors of resampling a response surface. These errors may cause Sobol's indices
345 larger than one.

346 To separate out the measurement uncertainty (instrument noises) of radar profiles and
347 their possible interactions from the parametric uncertainty, parallel results are shown for
348 holding the parameter “ a ” as a constant in Figs. 6 and 7. Under this scenario, different
349 radar reflectivity modes can be responsible for the IWP and IWC uncertainties at the

350 same time (e.g., Fig. 6b vs. Fig. 7b). It should be noted that radar reflectivity modes
351 with larger eigenvalues does not necessarily mean that they have more contributions on
352 the variance of the IWP or IWC than other modes. For instance, in Fig. 7a (21:00–21:30
353 UTC), the second largest eigenvalue and the value for the corresponding radar reflectivity
354 profile mode of IWC at 8km are 119 and -0.0222, respectively, while the third largest
355 eigenvalue and the value for the corresponding radar reflectivity profile mode of IWC at
356 8km are 105 and -0.1030, respectively. Accordingly, the sensitivity of the second and
357 the third radar reflectivity mode can be computed as $119 \times |-0.0222| = 2.6418$ and
358 $105 \times |-0.1030| = 10.815$, respectively. It is found that the second radar reflectivity
359 mode is less sensitive than the third one in this case. This means that there is no obvious
360 correlation between eigenvalue and sensitivity for the extracted random variables using
361 the EOF expansion. Similar rationales can be applied to study how various uncertainty
362 sources contribute to the variability of IWP as well.

363 Therefore, this probabilistic PCA based sensitivity analysis is determined by both eigen-
364 values and corresponding spatial modes of the observed stochastic input profiles. As far as
365 the sensitivity analysis is concerned, the contribution of variability of IWP due to radar
366 reflectivity mode interactions is larger for the optically thin clouds observed at 22:00–
367 22:30 UTC (Fig. 6d) than other periods such as the one observed at 21:00–21:30 UTC
368 (Fig. 6c). Nevertheless, it is a different kind of variability analysis result for IWC at 8 km
369 (see Figs. 7cd). Such quantitative knowledge with vertically resolved information about
370 the relative contribution of individual error source to the output uncertainties provides
371 valuable insights and clues to improve the both retrieval algorithm and measurements.

5. Conclusions and Discussions

372 Understanding and quantifying uncertainties in cloud retrieval is a subject of many
373 earlier studies [*Comstock et al.*, 2007; *Turner et al.*, 2007; *Posselt et al.*, 2008; *Comstock*
374 *et al.*, 2013; *Zhao et al.*, 2014]. Our contribution here is the development of a general, novel
375 observation-based methodology to quantify the retrieval uncertainties for climate model
376 evaluation. The EOF reduction of dimensions of random inputs enables our probabilistic
377 PCA based approach to quantify vertically resolved uncertainty and conduct global sensi-
378 tivity analysis. The UQ profiles with high vertical resolution are often more desirable for
379 model evaluation as vertical structures of clouds are essential to many important topics
380 such as radiative forcing and climate change [*Schneider*, 1972; *Schneider and Dickinson*,
381 1974; *Zelinka et al.*, 2012].

382 To reduce the dimensionality of random inputs, our method takes into account the
383 correlation between vertical layers in the input data by adopting the EOF expansion.
384 Moreover, by eliminating the unrealistic assumption that different layers are uncorre-
385 lated, the output uncertainty range becomes more accurate and reliable. Besides means
386 and standard deviations, the proposed method also quantifies the full PDFs of retrieved
387 quantities at each vertical layer. This observation-based PDF information can be used as
388 *a priori* for the Bayesian approach [*McFarlane et al.*, 2002; *Posselt et al.*, 2008; *Shen et al.*,
389 2013] to avoid the so-called subjective uncertainty introduced by assuming *a priori* PDF
390 (usually assumed to be uniform), and hence improve the results from Bayesian studies.
391 Besides propagating uncertainties in the input data and the parameters to retrieval
392 outputs, this UQ approach has the capability of attributing the output uncertainties to
393 individual error source, i.e., Sobol's global sensitivity analysis. This capacity is partic-

394 ularly useful when dealing with highly non-linear retrieval algorithms, as various error
395 sources are more likely entangled.

396 Despite of the above advantages, this framework does not cover all the aspects of UQ
397 analysis. For example, it cannot quantify systematic biases and the retrieval model struc-
398 tural uncertainty. The parameters of the retrieval algorithm may not be independent
399 as assumed in this approach. For instance, parameters “ a ” and “ b ” in the Eq. (7) are
400 dependent on each other [Matrosov, 1999]. In addition, some retrievals [McFarlane *et al.*,
401 2002; Turner, 2005; Posselt *et al.*, 2008] have already applied some uncertainty estimation
402 approach, e.g. Bayesian calibration, and thus our approach may not be able to be directly
403 applied to such algorithms.

404 The case study in this paper mainly demonstrates the capacities of this newly developed
405 UQ methodology. We will expand the UQ analysis to long-term ARM observations to
406 include different seasons and cloud types. Such comprehensive knowledge about retrieval
407 uncertainties will facilitate the application of retrieval products in model evaluation and
408 can be used to improve instruments, observation strategies as well as retrieval algorithms.
409 We also plan to exploit the uncertainties of other retrieval algorithms. Using multi-
410 retrieval and global model observations, we can further apply multi-model calibration
411 technique to mitigate the uncertainty estimated by each retrieval algorithm.

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⁴²⁵ group research activities. The data for the ice cloud case observed at the ARM SGP CF
⁴²⁶ site can be accessed by linking to <http://www.archive.arm.gov/armlogin/login.jsp>.

Appendix A

⁴²⁷ Subtracting ensemble mean $\bar{y}(\mathbf{x})$ from each snapshot, we obtain a zero-mean $N \times n$
⁴²⁸ snapshot matrix

$$\mathbf{Y} = [y_1 - \bar{y}, y_2 - \bar{y}, \dots, y_n - \bar{y}]. \quad (\text{A1})$$

⁴²⁹ It should be noted that we take snapshots of relative error for radar reflectivity profiles
⁴³⁰ and LWP, i.e., the snapshot matrix above is divided by \bar{y} , for which corresponding formulas
⁴³¹ can be derived similarly.

⁴³² Without loss of generality, the following set of vectors

$$\Psi = \{\psi_1, \psi_2, \dots, \psi_M\} \quad (\text{A2})$$

of order $M \leq n$ provides an optimal representation of the ensemble data in a M -dimensional subspace by minimizing the averaged projection error

$$\begin{aligned} \min_{\{\psi_1, \psi_2, \dots, \psi_M\}} \frac{1}{n} \sum_{i=1}^n \|(y_i - \bar{y}) - \Pi_{\Psi, M}(y_i - \bar{y})\|^2 \\ \text{s.t. } \langle \psi_i, \psi_j \rangle = \delta_{ij} = \begin{cases} 1 & i = j \\ 0 & i \neq j \end{cases}, \end{aligned} \quad (\text{A3})$$

⁴³³ where $\langle \cdot, \cdot \rangle$ represents an inner product, and $\Pi_{\Psi, M} = \sum_{i=1}^M \langle y_i - \bar{y}, \psi_i \rangle \psi_i$ is the projection
⁴³⁴ operator onto the M -dimensional space spanned by Ψ .

⁴³⁵ To compute the EOFs or the modes of PCA $\psi_i \in \mathbb{R}^N$ satisfying Eq. (A3), one solves
⁴³⁶ an N -dimensional eigenvalue problem

$$\mathbf{A}\psi_i = \lambda_i\psi_i, \quad (\text{A4})$$

⁴³⁷ where $\mathbf{A} = \mathbf{Y}\mathbf{Y}^T$ is the spatial correlation matrix.

438 Since in practice the number of snapshots is much less than the the state dimension,
 439 $n \ll N$, an efficient way to compute the reduced basis is to introduce a n -dimensional
 440 matrix $\mathbf{K} = \mathbf{Y}^T \mathbf{Y}$ and compute the eigenvalues $\lambda_1 \geq \lambda_2 \geq \dots \geq \lambda_n \geq 0$ of \mathbf{K} with its
 441 corresponding eigenvectors ϕ_1, \dots, ϕ_n . The corresponding EOF or modes of PCA are thus
 442 obtained by

$$443 \quad \psi_i = \frac{1}{\sqrt{\lambda_i}} \mathbf{Y} \phi_i, \quad i = 1, \dots, M, \quad (A5)$$

443 where $\langle \psi_i, \psi_j \rangle = \delta_{ij}$.

444 One can define a relative information content to choose a low-dimensional basis of size
 445 $M \ll n$ by neglecting modes corresponding to the small eigenvalues. We define

$$446 \quad I(m) = \frac{\sum_{i=1}^{i=m} \lambda_i}{\sum_{i=1}^{i=n} \lambda_i} \quad (A6)$$

446 and choose M such that $M = \arg \min \{I(m) : I(m) > \gamma\}$, where $0 \leq \gamma \leq 1$ is the
 447 percentage of total information retained in the reduced space and the tolerance γ must
 448 be chosen to be close unity in order to capture most of the energy of the snapshots basis.

449 A fast algorithm for eigenvalue calculation using a transposed matrix can be referenced
 450 [Shen et al., 2014].

451 Therefore, for each one observation y_i , it can be expanded in terms of M numbers of
 452 EOFs or modes of PCA written as

$$453 \quad y_i = \bar{y} + \sum_{i=1}^M \psi_i \sqrt{\frac{\lambda_i}{n}} V_i, \quad (A7)$$

453 where modal coefficients V_i computed by

$$V_i = \psi_i^T y_i \sqrt{\frac{n}{\lambda_i}}, \quad (\text{A8})$$

454 such that $\langle V_i, V_j \rangle = \delta_{ij}$.

455 Since mean is subtracted from each snapshots, it can be shown that $\frac{1}{n} \sum_{j=1}^n (V_{ij}) = 0$,

456 where V_{ij} corresponds to the observation $y_j - \bar{y}$ projected onto the mode ψ_i .

457 As a result, $Y(\mathbf{x}, t, \theta)$ can be approximated by EOF expansion to the order of M as

$$Y(\mathbf{x}, t, \theta) = \bar{y} + \sum_{i=1}^M \psi_i \sqrt{\frac{\lambda_i}{n}} \xi_i, \quad (\text{A9})$$

458 such that $E(\xi_i) = 0$ and $E(\xi_i \xi_j) = \delta_{ij}$, $i = 1, \dots, M$, and ξ_i follows some unknown

459 distribution.

Appendix B

460 Let $w = [w_1, w_2, \dots, w_N]^T$ be a temporally independent Gaussian noise injected into
 461 each one of the observation y_i . Therefore, w follows a multivariate normal distribution
 462 defined as $w \sim \mathcal{N}(0, \sigma_0^2 \mathbf{I}_N)$, where \mathbf{I}_N is a $N \times N$ identity matrix. Since Ψ is an orthogonal
 463 transformation, Ψw follows the same distribution as w , i.e., $\Psi w \sim \mathcal{N}(0, \sigma_0^2 \mathbf{I}_N)$. Therefore,
 464 without loss of generality, adding Ψw to the Equation (A9) and truncating it to the order
 465 of M , we obtain that

$$\begin{aligned} Y'(\mathbf{x}, t, \theta) &= Y(\mathbf{x}, t, \theta) + \Psi w \\ &= \bar{y} + \sum_{i=1}^M \psi_i \sqrt{\frac{\lambda_i}{n}} \xi_i + \sum_{i=1}^M \psi_i w_i \\ &= \bar{y} + \sum_{i=1}^M \psi_i \sqrt{\frac{\lambda_i}{n}} \left(\xi_i + \frac{w_i}{\sqrt{\frac{\lambda_i}{n}}} \right), \end{aligned} \quad (B1)$$

466 where $Y'(\mathbf{x}, t, \theta)$ is a stochastic process representing noisy observations, w_i is the i -th
 467 component of the truncated random vector Ψw .

468 Let ζ_i be $\zeta_i = \xi_i + \frac{w_i}{\sqrt{\frac{\lambda_i}{n}}}$, we obtain

$$Y'(\mathbf{x}, t, \theta) = \bar{y} + \sum_{i=1}^M \psi_i \sqrt{\frac{\lambda_i}{n}} \zeta_i, \quad (B2)$$

469 where $E(\zeta_i) = 0$ and $Var(\zeta_i) = \sqrt{1 + \left(\frac{\sigma_0}{\sqrt{\frac{\lambda_i}{n}}} \right)^2}$.

470 Taking average on both sides of the Eq. (B2) above, it can be rewritten as

$$\bar{Y}'(\mathbf{x}, t, \theta) = \bar{y} + \sum_{i=1}^M \psi_i \sqrt{\frac{\lambda_i}{n}} \bar{\zeta}_i. \quad (B3)$$

⁴⁷¹ Finally, based on CLT [*Cramér*, 1946; *Gnedenko et al.*, 1954] and considering the truncation error, we have

$$\overline{Y'}(\mathbf{x}, t, \theta) = \bar{y} + \sum_{i=1}^M \psi_i \sqrt{\frac{\lambda_i}{n}} \sqrt{1 + \left(\frac{\sigma_0}{\sqrt{\frac{\lambda_i}{n}}} \right)^2} \frac{z_i}{\sqrt{n}} + \text{error}(\mathbf{x}, t, \theta), \quad (\text{B4})$$

where $z_i \sim \mathcal{N}(0, 1)$, and the $\text{error}(\mathbf{x}, t, \theta)$ represents the truncation error incurred in the expansion above, which can be estimated following [*Shen et al.*, 2004, 2014].

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Figure 1. Height-time plots at SGP CF site on March 9, 2000 for (a) MMCR reflectivity (dBZ); (b) 0.5-hour averages of IWC (g/m^3) from raw MICROBASE; (c) 0.5-hour ensemble means (ensemble size=1000), and (d) standard deviations (STDs, σ) of IWC from probabilistic PCA based method; (e) comparison of 0.5-hour IWP (g/m^2) from probabilistic PCA based method (blue), raw MICROBASE (red), and 2-min in-situ CVI measurements (black) as aircraft passed over the SGP CF site. Dashed lines connect the means, and error bars represent $\pm 1\sigma$.

Figure 2. Height-time plots at SGP CF site on March 9, 2000 for the leading four MMCR reflectivity (dBZ) profiles modes, marked with corresponding eigenvalues and weighting percentage to capture the energy of snapshots of radar reflectivity profiles within each time window.

Figure 3. Comparison of 0.5-hour IWP (g/m^2) from probabilistic PCA based method (blue) when keeping a as a constant, raw MICROBASE (red), and 2-min in-situ CVI measurements (black) as aircraft passed over the SGP CF site. Dashed lines connect the means, and error bars represent $\pm 1\sigma$.

Figure 4. Probability density functions derived from probabilistic PCA based method of (a) IWC (g/m^3) at 8 km, 19:00–22:30 shown as box plot (red lines: median; lower/upper blue box lines: lower/upper quartiles; whiskers show the extent of the data); (b) IWP (g/m^2) at 21:00–21:30 UTC (green) and 22:00–22:30 UTC (black) on March 9, 2000.

Figure 5. Sobol's first-order index of (a) IWP and (b) IWC at 8 km, and Sobol's first and second order (i.e., the sum of two different first order index and their joint index, Note that the diagonal and sub-diagonal numbers are not shown by definition.) index of (c) IWP and (d) IWC at 8 km. All the results are for March 9, 2000 22:00–22:30 UTC.

Figure 6. Sobol's first-order index of IWP for (a) 2100–21:30 UTC and (b) 22:00–22:30 UTC, and Sobol's first and second order (i.e., the sum of two different first order index and their joint index, Note that the diagonal and sub-diagonal numbers are not shown by definition.) index of IWP for (c) 2100–21:30 UTC and (d) 22:00–22:30 UTC. All the results are for March 9, 2000.

Figure 7. Sobol's first-order index of IWC at 8 km for (a) 2100–21:30 UTC and (b) 22:00–22:30 UTC, and Sobol's first and second order (i.e., the sum of two different first order index and their joint index, Note that the diagonal and sub-diagonal numbers are not shown by definition.) index of IWP for (c) 2100–21:30 UTC and (d) 22:00–22:30 UTC. All the results are for March 9, 2000.













