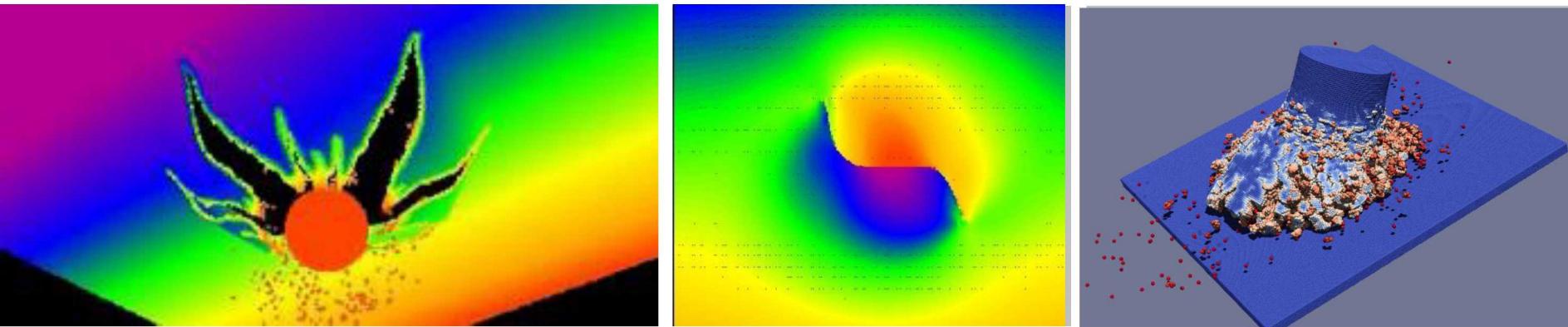


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# Nonlocality in peridynamics

Stewart Silling

Computational Multiscale Department  
Sandia National Laboratories  
Albuquerque, New Mexico

Workshop on Experimental and Computational Fracture Mechanics  
Baton Rouge, LA, February 26, 2020



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# Outline

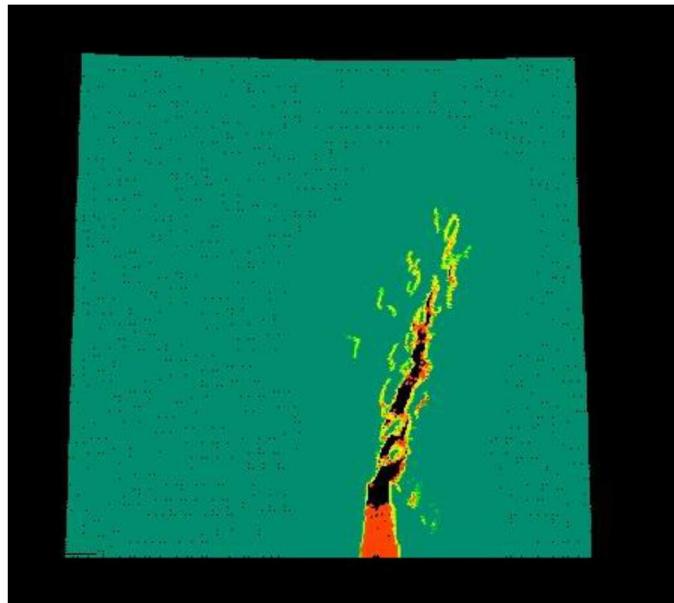
- Nonlocality
  - It's not as weird as everybody thinks
- Peridynamics background
  - All-in on nonlocality
- Can nonlocality be derived or observed?
  - Long-range forces
  - Smoothed degrees of freedom (homogenization)
  - Multiple pathways for flux
  - Wave dispersion



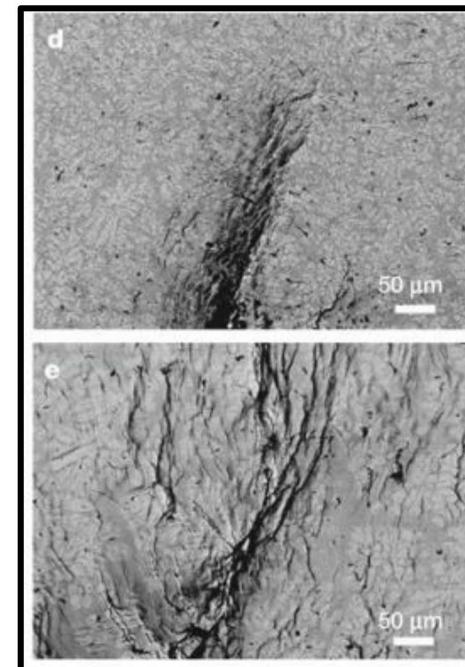
Do we ask too much of the local theory of  
continuum mechanics?

# What peridynamics seeks to accomplish

- Treat material points on or off of evolving discontinuities with the same equations.
- Include long-range forces in the basic equations.
- Fit all this into a thermodynamic framework that's consistent with the mechanics.



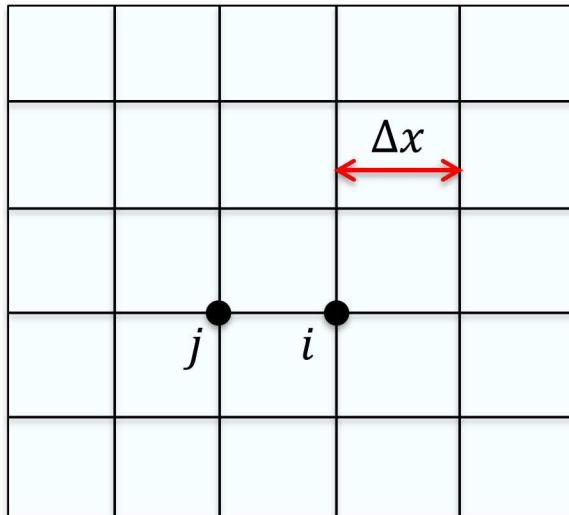
Peridynamic simulation



Metallic glass crack tip\*

\*Hofmann et al, Nature (2008)

# Discretized numerical methods are nonlocal



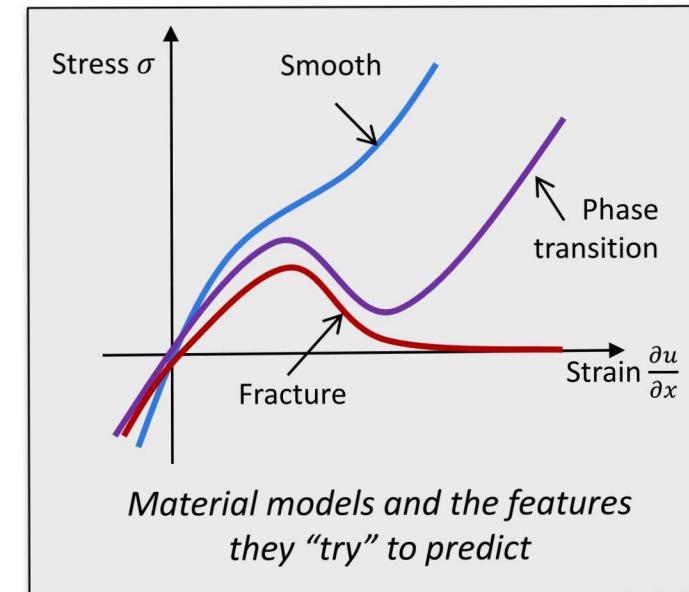
- Node  $i$  interacts directly with node  $j$  through the finite element equations.
- Interaction is across a finite distance  $\Delta x$ .
- This is a form of nonlocality.
  - Notwithstanding that the result converges to the local result as  $\Delta x \rightarrow 0$ .

# Local PDEs get themselves into trouble

- Classical (Cauchy) PDE:

$$\rho \ddot{\mathbf{u}} = \nabla \cdot \boldsymbol{\sigma} \left( \frac{\partial \mathbf{u}}{\partial \mathbf{x}} \right) + \mathbf{b}.$$

- Many material models  $\sigma(\cdot)$  evolve into deformations that are incompatible with the fundamental assumptions.
  - Phase boundaries, shock waves, cracks, ...
- Can't directly treat some important physical effects.
  - Wave dispersion, surface energy, microstructure evolution, long-range forces, ...
- People often take drastic measures if they want to work with this PDE.
  - Element deletion, ...

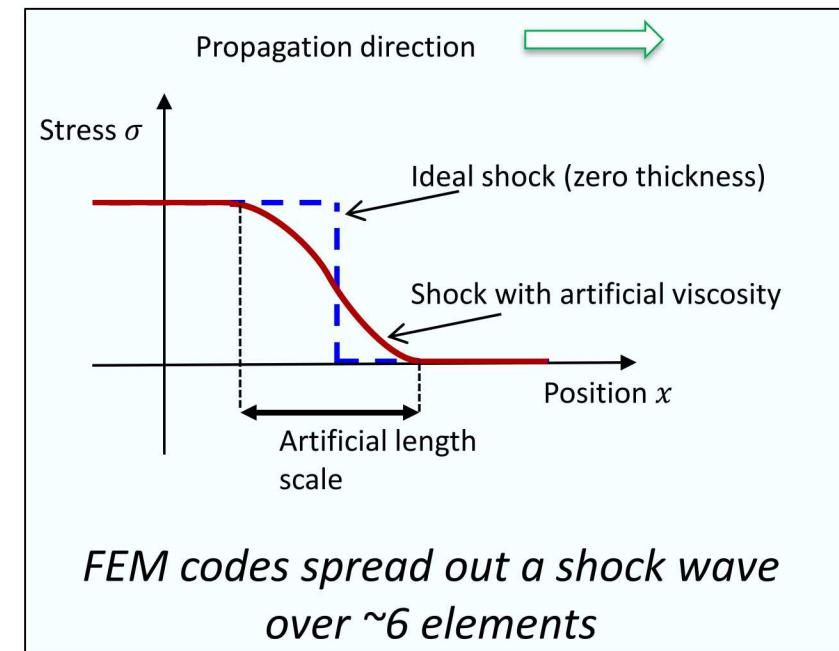


# These drastic measures often involve nonlocality

- Example: Artificial viscosity spreads out a shock wave and dissipates energy.

$$\rho \ddot{\mathbf{u}} = \nabla \cdot \boldsymbol{\sigma} \left( \frac{\partial \mathbf{u}}{\partial \mathbf{x}} \right) + \gamma (\nabla \cdot \dot{\mathbf{u}})^2 + \mathbf{b}.$$

- It avoids the need to apply jump conditions across an ideal shock.
- It allows conventional discretization to be used “within” a shock.
- By spreading out a shock it introduces a length scale.
- This is a type of nonlocality.



- J. Von Neumann & R. D. Richtmyer, *J. Appl. Phys.* 21 (1950). 232

# Peridynamics goes all-in on nonlocality

Classification of some theories with respect to local/nonlocality:

***PDEs with no length scale:***

- Classical continuum mechanics

***PDEs with a length scale:***

- Micropolar
- Mindlin
- Kroner
- Eringen
- Phase field
- Nonlocal damage
- Plate & shell theories
- Gradient theories

***Full nonlocality:***

- Kunin
- Peridynamics

- Every fundamental relation in peridynamics is nonlocal in space:
  - Transport
  - Conservation
  - Material models

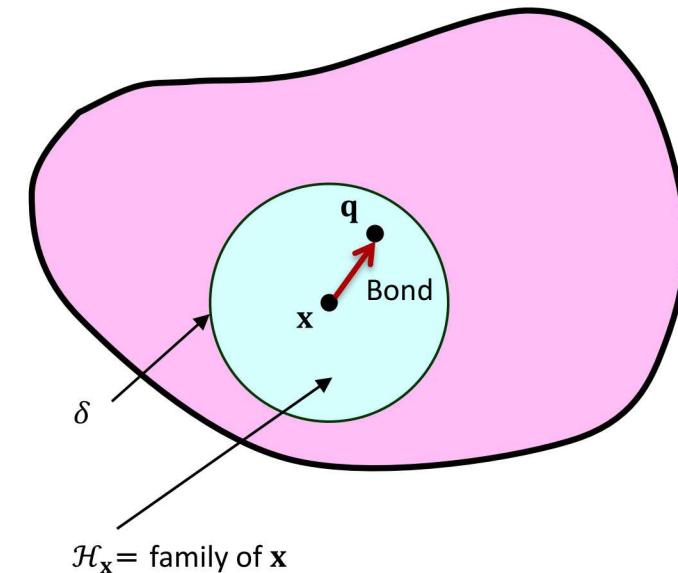
# Peridynamic\* momentum balance

- Any point  $x$  interacts directly with other points within a distance  $\delta$  called the “horizon.”
- The material within a distance  $\delta$  of  $x$  is called the “family” of  $x$ ,  $\mathcal{H}_x$ .

Peridynamic equilibrium equation

$$\int_{\mathcal{H}_x} \mathbf{f}(\mathbf{q}, \mathbf{x}) dV_{\mathbf{q}} + \mathbf{b}(\mathbf{x}) = 0$$

$\mathbf{f}$  = bond force density (from the material model, which includes damage)



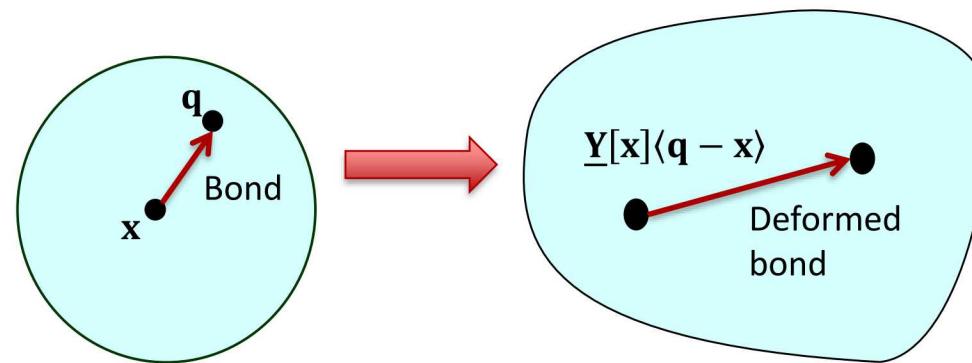
- If  $\mathbf{f}$  satisfies  $\mathbf{f}(\mathbf{x}, \mathbf{q}) = -\mathbf{f}(\mathbf{q}, \mathbf{x})$  for all  $\mathbf{x}, \mathbf{q}$  then linear momentum is conserved.
- SS, JMPS (2000)

\* Peri (near) + dyne (force)

# Formalism for nonlocal interactions: States

- A *state* is a mapping whose domain is all the bonds  $\xi$  in a family.

$$\underline{\mathbf{A}}\langle \xi \rangle = \text{something} \quad \forall \xi \in \mathcal{H}.$$

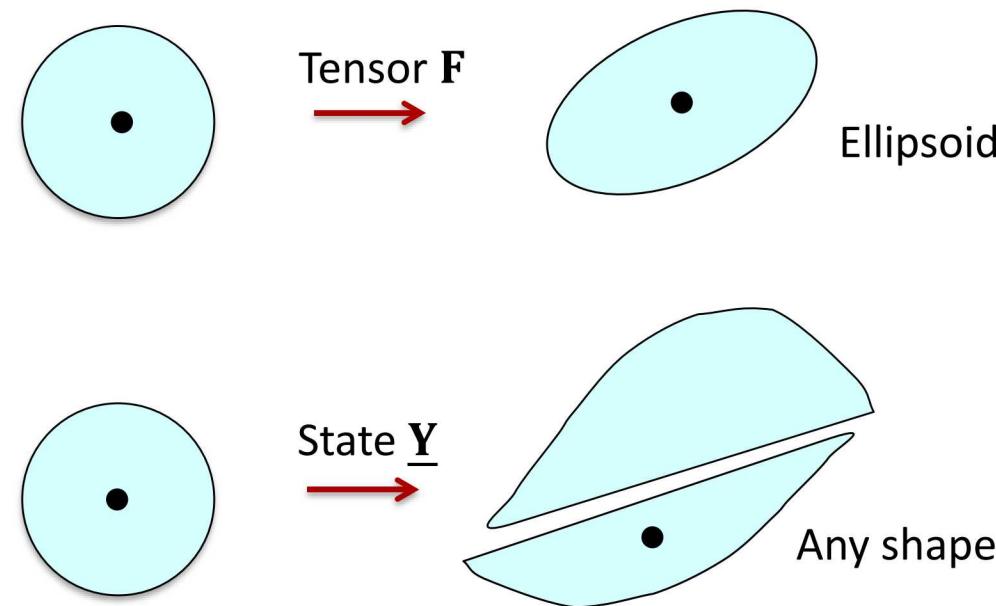


- Deformation state...

$$\underline{\mathbf{Y}}[x]\langle q - x \rangle = \mathbf{y}(q) - \mathbf{y}(x) = \text{deformed image of the bond}$$

# States: Nonlocal analogues of second order tensors

- Classical theory uses tensors (linear mappings from vectors to vectors).
- Peridynamics uses states (nonlinear mappings from vectors to vectors).



# Peridynamic vs. local equations

- Structurally similar but with states instead of local operators.

Relation	<i>Peridynamic theory</i>	<i>Standard theory</i>
Kinematics	$\underline{\mathbf{Y}} \langle \mathbf{q} - \mathbf{x} \rangle = \mathbf{y}(\mathbf{q}) - \mathbf{y}(\mathbf{x})$	$\mathbf{F}(\mathbf{x}) = \frac{\partial \mathbf{y}}{\partial \mathbf{x}}(\mathbf{x})$
Linear momentum balance	$\rho \ddot{\mathbf{y}}(\mathbf{x}) = \int_{\mathcal{H}} \left( \mathbf{t}(\mathbf{q}, \mathbf{x}) - \mathbf{t}(\mathbf{x}, \mathbf{q}) \right) dV_{\mathbf{q}} + \mathbf{b}(\mathbf{x})$	$\rho \ddot{\mathbf{y}}(\mathbf{x}) = \nabla \cdot \boldsymbol{\sigma}(\mathbf{x}) + \mathbf{b}(\mathbf{x})$
Constitutive model	$\mathbf{t}(\mathbf{q}, \mathbf{x}) = \underline{\mathbf{T}} \langle \mathbf{q} - \mathbf{x} \rangle, \quad \underline{\mathbf{T}} = \hat{\underline{\mathbf{T}}}(\underline{\mathbf{Y}})$	$\boldsymbol{\sigma} = \hat{\boldsymbol{\sigma}}(\mathbf{F})$
Angular momentum balance	$\int_{\mathcal{H}} \underline{\mathbf{Y}} \langle \mathbf{q} - \mathbf{x} \rangle \times \underline{\mathbf{T}} \langle \mathbf{q} - \mathbf{x} \rangle dV_{\mathbf{q}} = \mathbf{0}$	$\boldsymbol{\sigma} = \boldsymbol{\sigma}^T$
Elasticity	$\underline{\mathbf{T}} = W_{\underline{\mathbf{Y}}} \text{ (Fréchet derivative)}$	$\boldsymbol{\sigma} = W_{\mathbf{F}} \text{ (tensor gradient)}$
First law	$\dot{\varepsilon} = \underline{\mathbf{T}} \bullet \dot{\underline{\mathbf{Y}}} + q + r$	$\dot{\varepsilon} = \boldsymbol{\sigma} \cdot \dot{\mathbf{F}} + q + r$

$$\underline{\mathbf{T}} \bullet \dot{\underline{\mathbf{Y}}} := \int_{\mathcal{H}} \underline{\mathbf{T}} \langle \boldsymbol{\xi} \rangle \cdot \dot{\underline{\mathbf{Y}}} \langle \boldsymbol{\xi} \rangle dV_{\boldsymbol{\xi}}$$

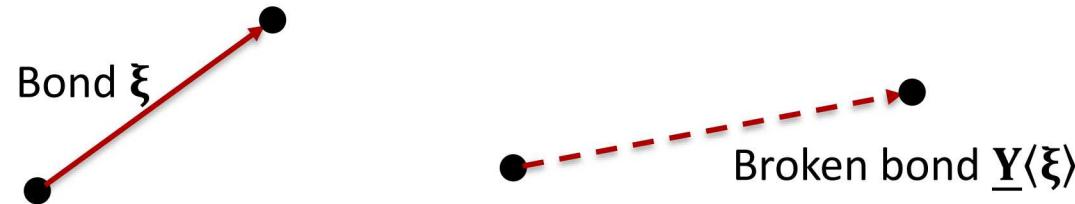
# Damage

- Damage is usually treated through *bond breakage*.
- After a bond  $\xi$  breaks according to some criterion, it no longer carries any force.
- Typical breakage criterion: prescribed *critical bond strain*  $s_0$ :

$$s = \frac{|\mathbf{Y}(\xi)| - |\xi|}{|\xi|} \quad \text{bond strain.}$$

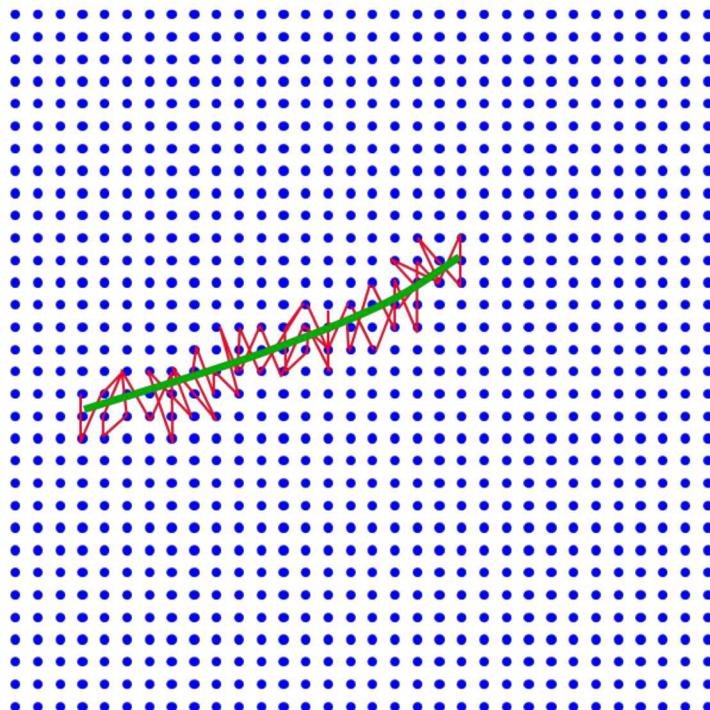
$$s \geq s_0 \text{ at some time } t_0$$

means the bond remains broken for all  $t \geq t_0$ .

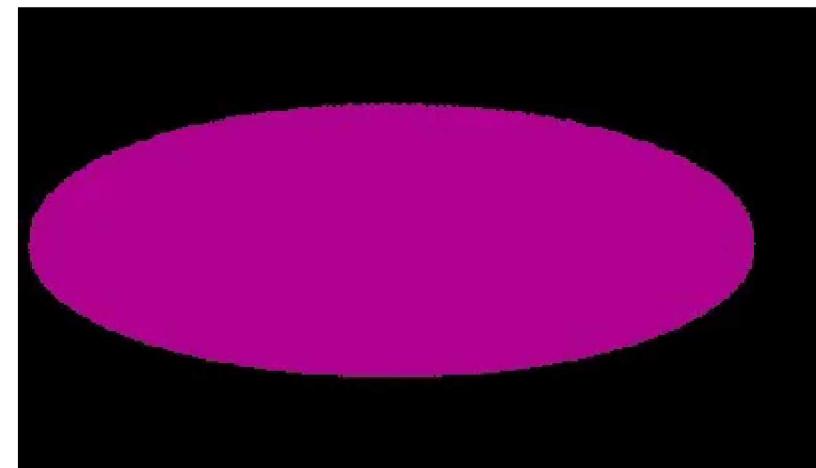


# Autonomous crack growth

- Cracks do what they want (grow, arrest, branch, curve, oscillate, ...)



— Broken bond  
— Crack path



- SS & Askari, *Computers and Structures* (2005)

# Many validation studies have been done

- First issue of the new *Journal of Peridynamics and Nonlocal Modeling* had a review article by Diehl on published validation to date:

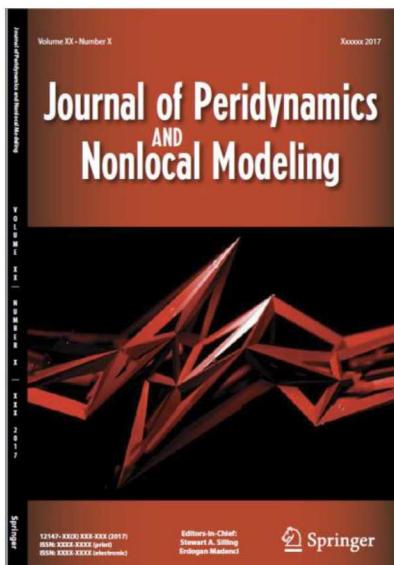
Journal of Peridynamics and Nonlocal Modeling  
<https://doi.org/10.1007/s42102-018-0004-x>

**REVIEWS**

**A Review of Benchmark Experiments for the Validation of Peridynamics Models**

Patrick Diehl<sup>1</sup>  · Serge Prudhomme<sup>2</sup> · Martin Lévesque<sup>1</sup>

Received: 2 November 2018 / Accepted: 25 December 2018  
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**Table 3** Applications of bond-based and state-based peridynamics for the comparison with experimental data

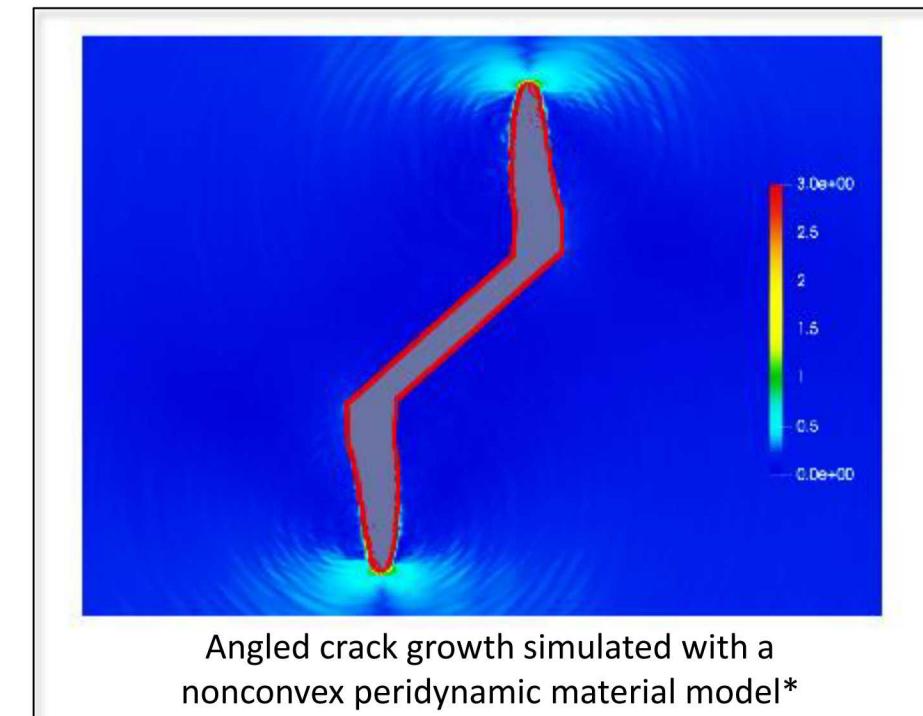
Material	Mechanical test	B	S	Exp	Sim
Composite	Flexural test with an initial crack	✓		[75]	[2]
Composite	Damage growth prediction (six-bolt specimen)	✓		[120]	[96]
Composite	Damage prediction (center-cracked laminates)	✓		[6, 12, 69, 134]	[70]
Composite	Dynamic tension test (prenotched rectangular plate)	✓		[12, 65]	[58]
Steel	Crack growth (Kalthoff-Winkler)	✓	✓	[66–68]	[3, 52, 114, 144]
Aluminum/Steel	Fracture (compact tension test)	✓		[9, 77, 89, 91]	[135, 141, 142]
Aluminum	Taylor impact test	✓		[4, 21]	[3, 43, 45]
Aluminum (6061-T6)	Ballistic impact test	✓		[132]	[127]
Concrete	Lap-splice experiment	✓		[48]	[48]
Concrete	3-point bending beam	✓	✓	[19, 63]	[7, 51]
Concrete	Failure in a Brazilian disk under compression	✓		[51]	[54]
Concrete	Anchor Bolt Pullout	✓		[128]	[83]
Glass	Dynamic crack propagation (prenotched thin rectangular plate)	✓		[15, 36, 100]	[2, 53, 144]
Glass	Impact damage with a thin polycarbonate backing	✓		[8, 20, 40]	[59]
Glass	Single crack paths (quenched glass plate)	✓		[13, 103, 136]	[71]
Glass	Multiple crack paths (quenched glass plate)	✓		[102, 137]	[71]
Glass	Crack tip propagation speed	✓		[15]	[52, 53, 144]
PMMA	Fast cracks in PMMA	✓		[39]	[2]
PMMA	Tensile test	✓		[124]	[32]
Soda-lime glass	Impact on a two-plate system	✓		[16, 130]	[130]

Legend: B refers to bond-based peridynamics, S refers to state-based peridynamics, Exp to experimental data, and Sim to simulation

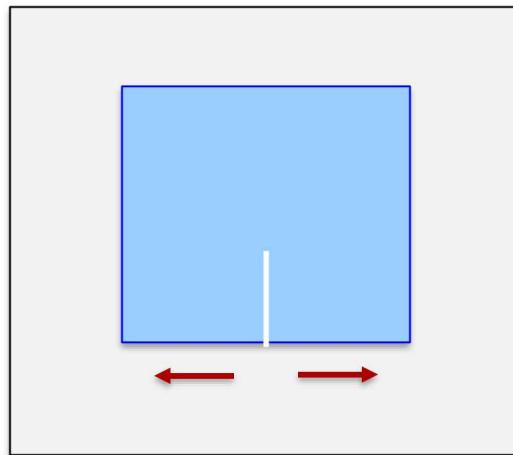
# Peridynamics converges as the horizon $\rightarrow 0$

- Linear peridynamics converges to Navier equations of linear elasticity.
- Linear or nonlinear material models converge to a stress-strain relation.
- Problems with nonconvex elastic peridynamic models can converge to nonlinear elasticity with Griffith cracks.

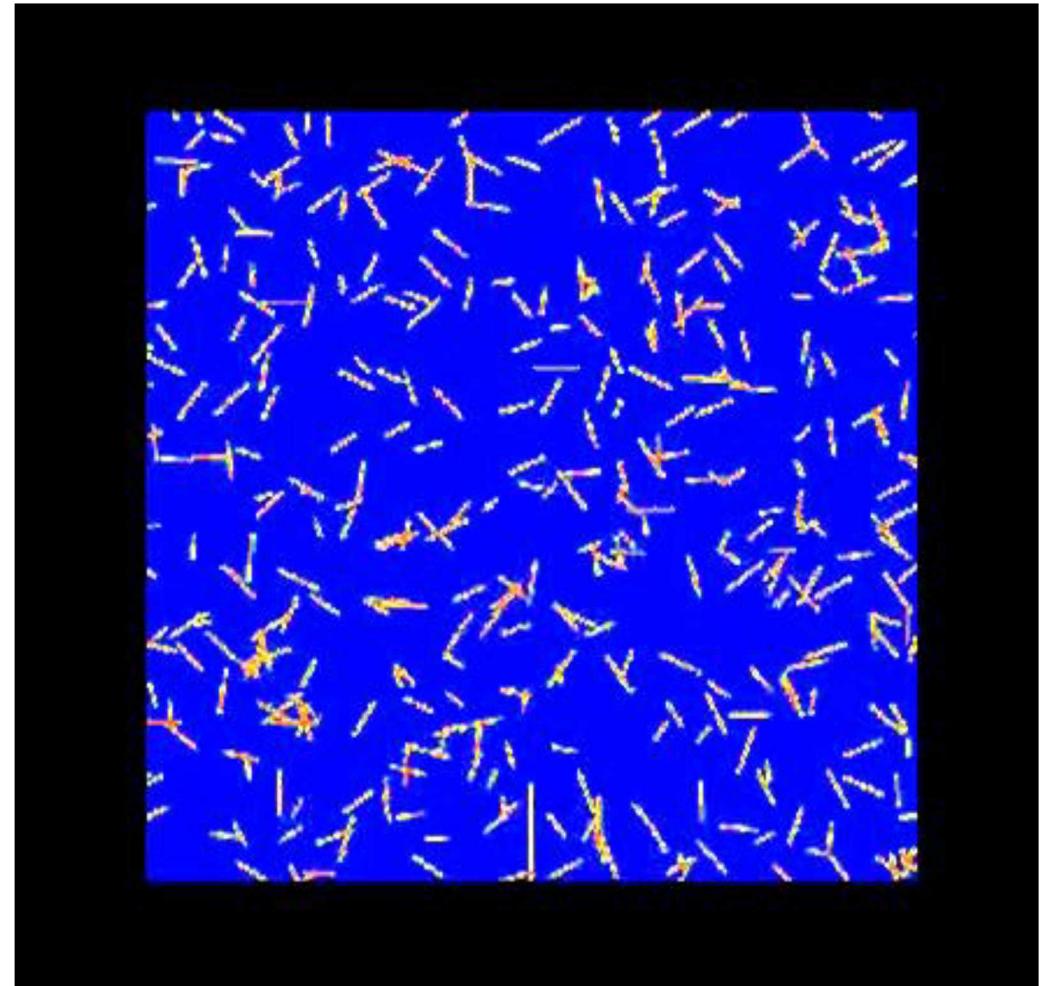
- E. Emmrich & O. Weckner, *Communications in Mathematical Sciences* (2007).
- F. Bobaru et al., *Int. Journal for Numerical Methods in Engineering* (2009).
- T. Mengesha, & Q. Du, *Journal of Elasticity* (2014).
- S.S. & R. B. Lehoucq, *Journal of Elasticity* (2008).
- P. Seleson & D.J. Littlewood, *Computers & Mathematics with Applications* (2016).
- \*R. P. Lipton, R. B. Lehoucq, & P.K. Jha, *Journal of Peridynamics and Nonlocal Modeling* (2019).



# Example: Fracture in a brittle plate with a lot of defects

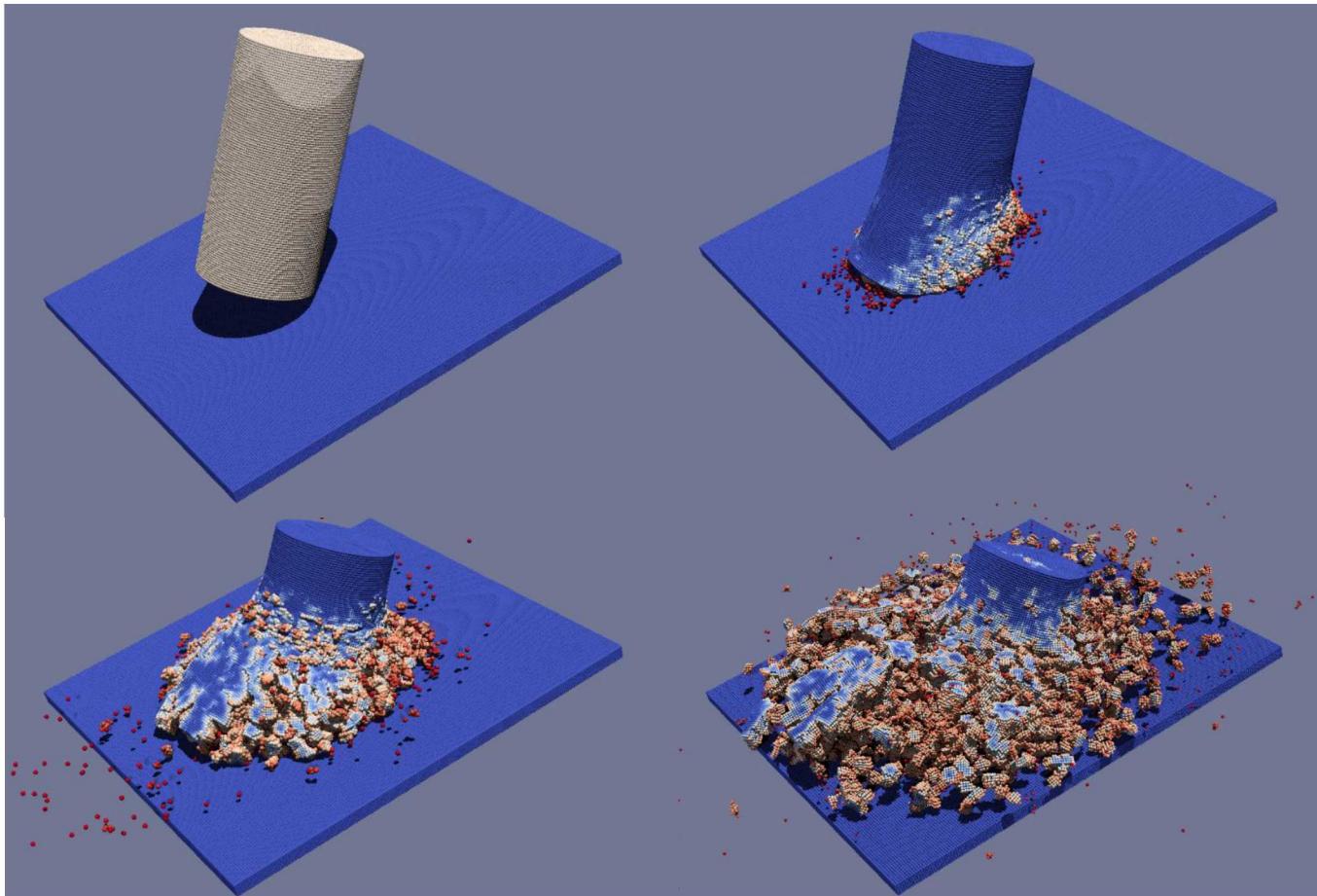


VIDEO



# Example: Fragmentation due to impact

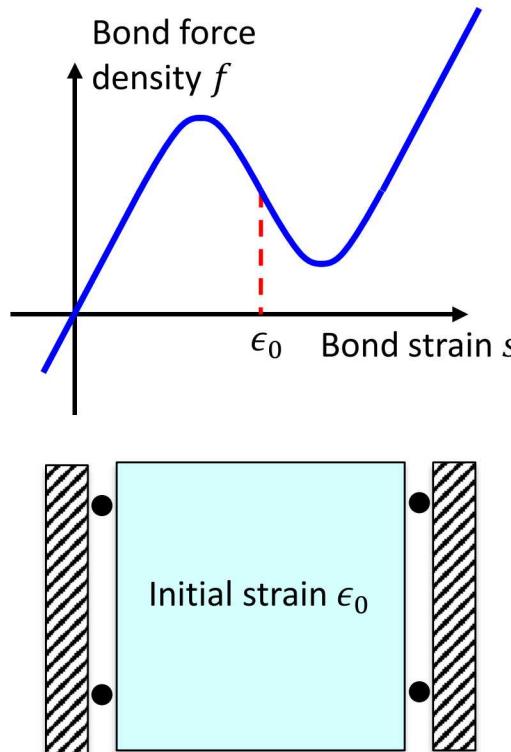
- Brittle cylinder vs. rigid plate at 1km/s.



Colors show damage

# Example: Microstructure evolution

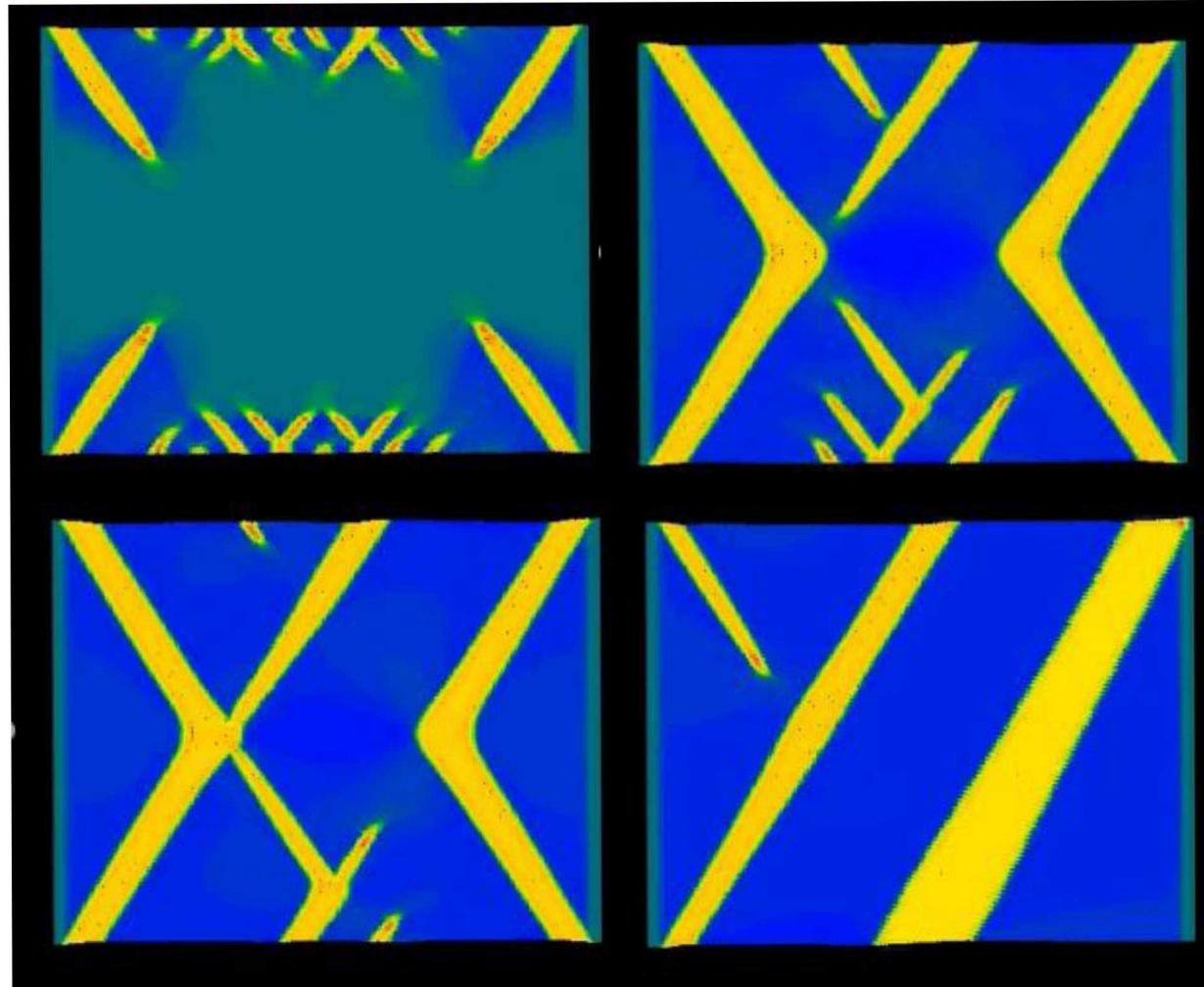
- Plate with ends fixed. Global strain  $\epsilon_0$  is in the unstable part of the material model.
- Complex microstructure appears at first, then simplifies.
- Driving force is the energy stuck in a phase boundary.

[VIDEO](#)

Colors show bond strain

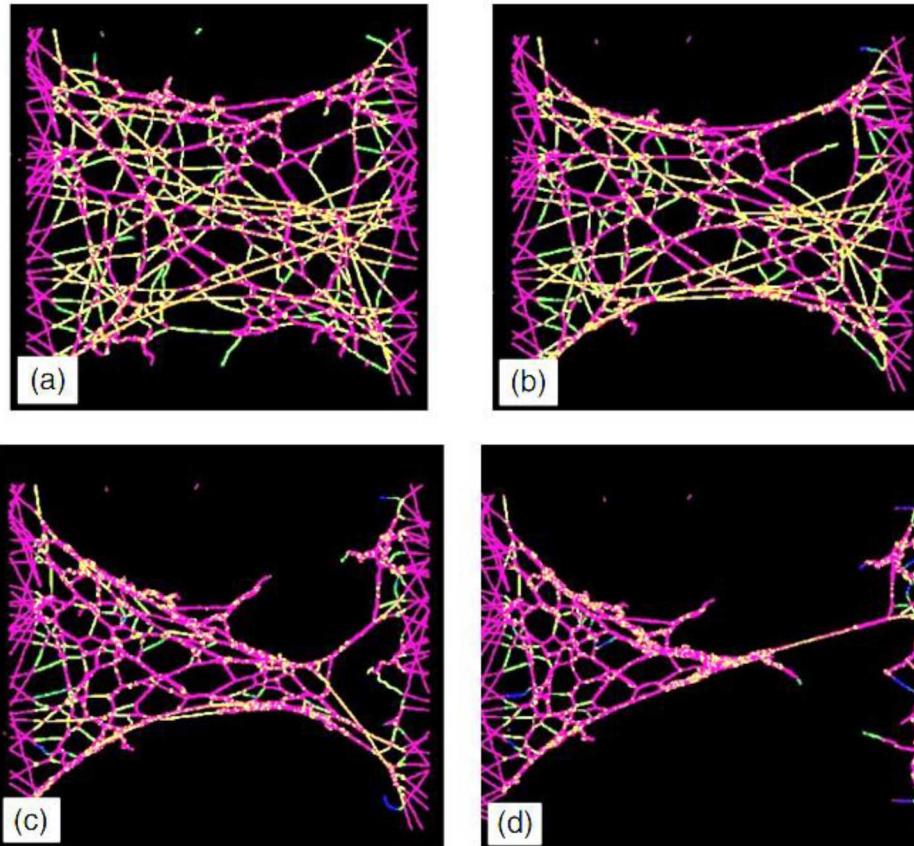
# Example: Microstructure evolution

Colors show bond strain



# Straightforward case for nonlocality: When there really are long-range forces

- Fracture of nanofiber network held together by Van der Waals forces.



F. Bobaru, *Modelling and Simulation in Materials Science and Engineering* 15, no. 5 (2007): 397.

# Smoothing the smallest scale degrees of freedom results in nonlocality

- Try to approximate known, small-scale response (e.g. molecular motion) by a continuous variable, yet retain realistic behavior.
- How to make the connection?
- One approach: Smooth out the small-scale degrees of freedom.
- Example:
  - Heterogeneous infinite bar.



- Small-scale model (local):

$$\rho(x)\ddot{u}(x, t) = \sigma'(x, t) + b(x, t)$$

where  $\rho$ =density,  $u$ =displacement,  $\sigma$ =stress, and  $b$ =body force density.

- Material model:

$$\sigma(x, t) = E(x)u'(x, t)$$

where  $E$ =Young's modulus.

# Define a smoothed displacement field

- Let  $w(z)$  be a smoothing function on  $z \in [-\epsilon, \epsilon]$ ,  $\int w = 1$ ,  $w(-z) = w(z)$ .
- Define the smoothed displacement field  $\bar{u}$  by

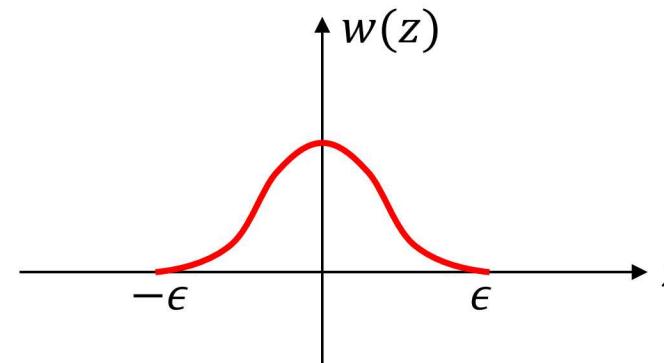
$$\bar{u}(x, t) = \frac{1}{\bar{\rho}(x)} \int_{-\infty}^{\infty} w(p-x) \rho(p) u(p, t) \, dp, \quad \bar{\rho}(x) := \int_{-\infty}^{\infty} w(p-x) \rho(p) \, dp$$

- Recall

$$\rho(x) \ddot{u}(x, t) = \sigma'(x, t) + b(x, t).$$

- Multiply through by  $w$  and integrate, find that

$$\bar{\rho}(x) \ddot{\bar{u}}(x, t) = \int_{-\infty}^{\infty} w(x-p) \sigma'(p, t) \, dp + \bar{b}(x, t), \quad \bar{b}(x, t) := \int_{-\infty}^{\infty} w(x-p) b(p, t) \, dp$$



# Evolution equation for smoothed DOFs

- Recall

$$\bar{\rho}(x)\ddot{u}(x, t) = \int_{-\infty}^{\infty} w(x-p)\sigma'(p, t) \, dp + \bar{b}(x, t).$$

- Integrate by parts (surprise!):

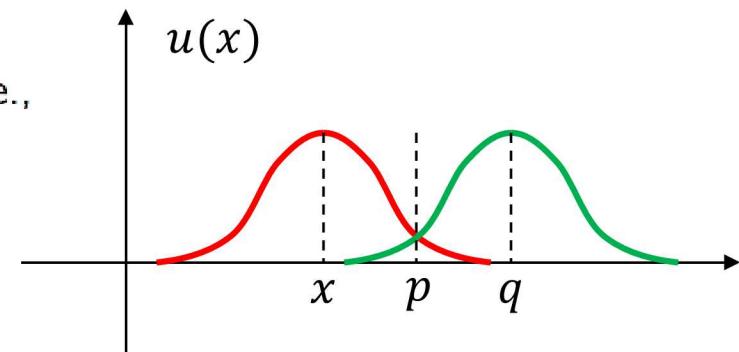
$$\bar{\rho}(x)\ddot{u}(x, t) = - \int_{-\infty}^{\infty} w'(x-p)\sigma(p, t) \, dp + \bar{b}(x, t).$$

- Starting to look nonlocal.
- Let  $q$  be defined so that  $p$  is halfway between  $x$  and  $q$ , i.e.,

$$p = \frac{x+q}{2}.$$

- Then

$$\bar{\rho}(x)\ddot{u}(x, t) = -\frac{1}{2} \int_{-\infty}^{\infty} w' \left( \frac{q-x}{2} \right) \sigma \left( \frac{q+x}{2}, t \right) \, dp + \bar{b}(x, t).$$



# Evolution equation is nonlocal

- Recall

$$\bar{\rho}(x)\ddot{u}(x, t) = -\frac{1}{2} \int_{-\infty}^{\infty} w' \left( \frac{q-x}{2} \right) \sigma \left( \frac{q+x}{2}, t \right) dq + \bar{b}(x, t).$$

- Now define the *pairwise bond force density* by

$$f(q, x) = -\frac{1}{2} w' \left( \frac{q-x}{2} \right) \sigma \left( \frac{q+x}{2}, t \right)$$

and define the *horizon* by

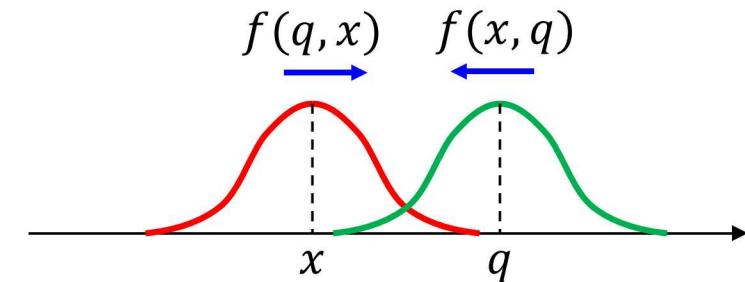
$$\delta = 2\epsilon.$$

- We now have

$$\bar{\rho}(x)\ddot{u}(x, t) = \int_{x-\delta}^{x+\delta} f(q, x) dq + \bar{b}(x, t).$$

- Observe that  $f$  has the required symmetry

$$f(x, q) = -f(q, x).$$



# Need a material model in terms of the smoothed DOFs

- Unfortunately we don't know  $\sigma$ .
- One possibility is to back out  $u'$  from the Fourier transform using the convolution theorem:

$$\mathcal{F}\{\bar{u}\} = \mathcal{F}\{w\}\mathcal{F}\{u\} \quad \Rightarrow \quad u = \mathcal{F}^{-1} \left\{ \frac{\mathcal{F}\{\bar{u}\}}{\mathcal{F}\{w\}} \right\}$$

hence

$$\sigma(x) = E(x) \frac{d}{dx} \mathcal{F}^{-1} \left\{ \frac{\mathcal{F}\{\bar{u}\}}{\mathcal{F}\{w\}} \right\}.$$

- This is too much work!
- Instead come up with a nonlocal material model.

# Bond-based heterogeneous material model

- Observe that in equilibrium with  $b \equiv 0$  and fixed stress  $\sigma_0$ ,

$$u_0(x) = \int_0^x \frac{\sigma_0}{E(z)} dz.$$

- From this compute the smoothed displacements:

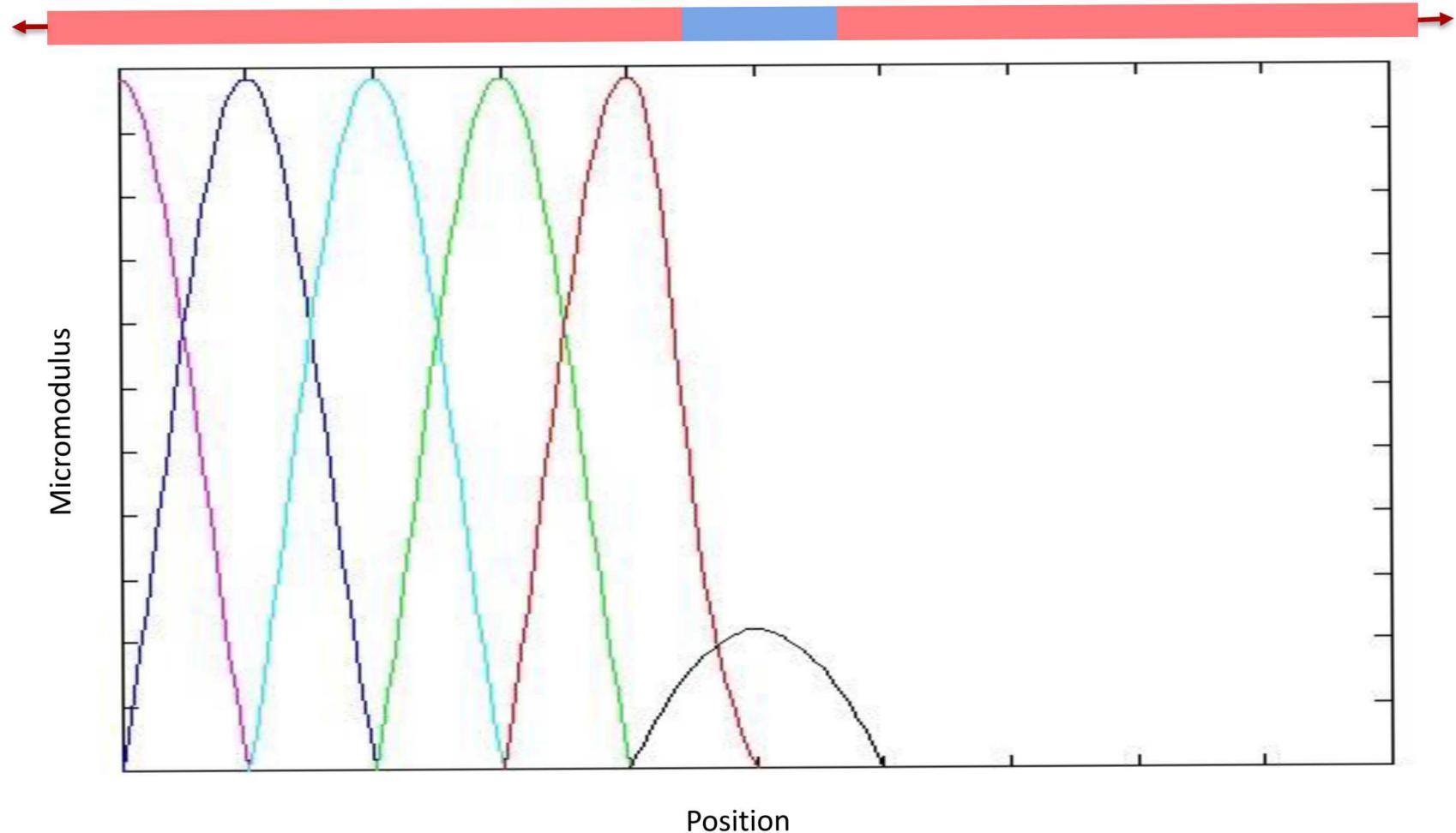
$$\bar{u}(x) = \int_{-\epsilon}^{\epsilon} w(\zeta) u_0(x + \zeta) d\zeta.$$

- Define a nonlocal material model by (omit  $t$ ):

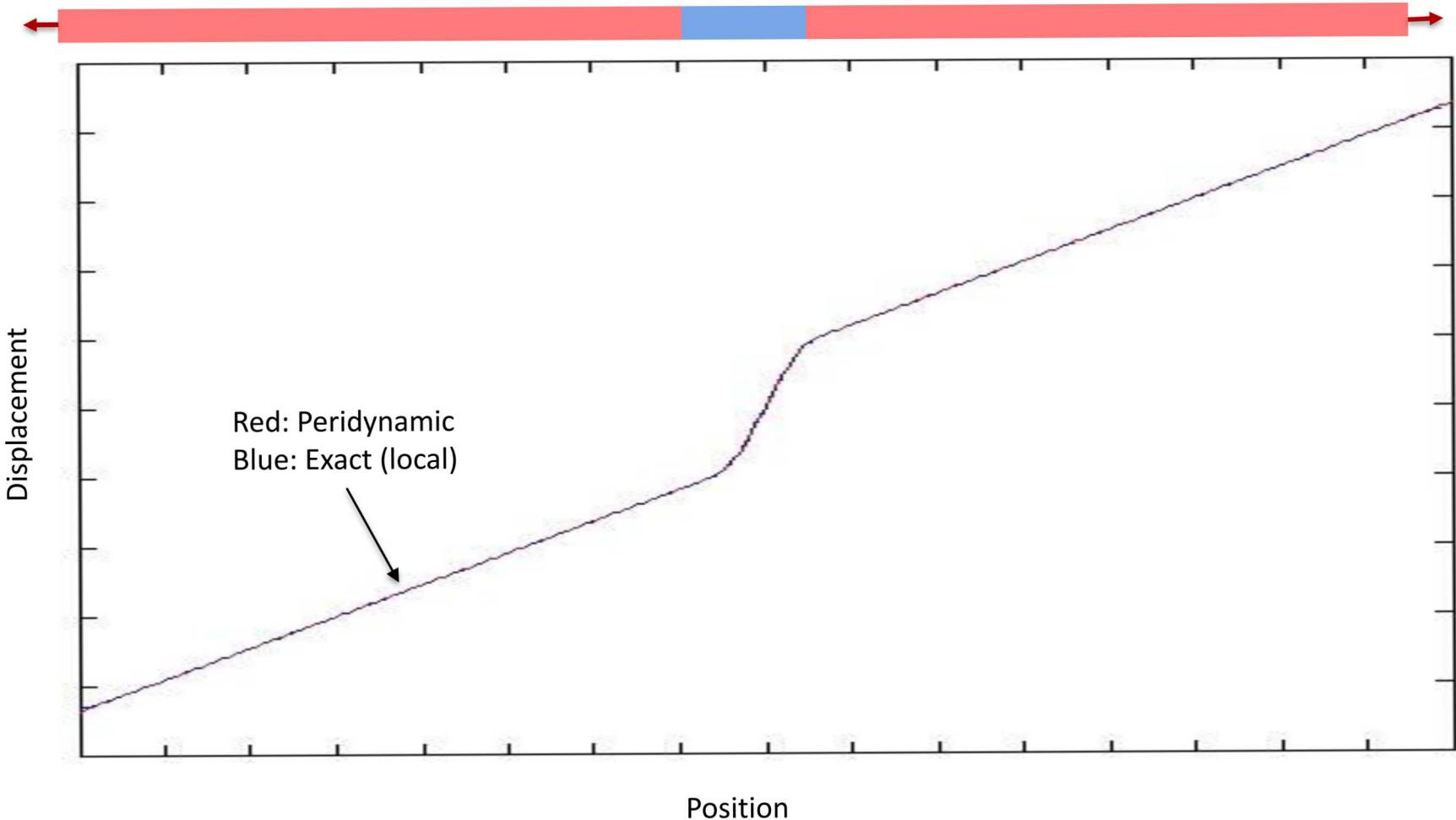
$$f(q, x) = C(q, x)(\bar{u}(q) - \bar{u}(x)), \quad C(q, x) := \frac{\sigma_0 w'((q - x)/2)}{\bar{u}_0(q) - \bar{u}_0(x)}.$$

- This exactly reproduces the local result for equilibrium with  $b \equiv 0$ .
- (But not in general.)

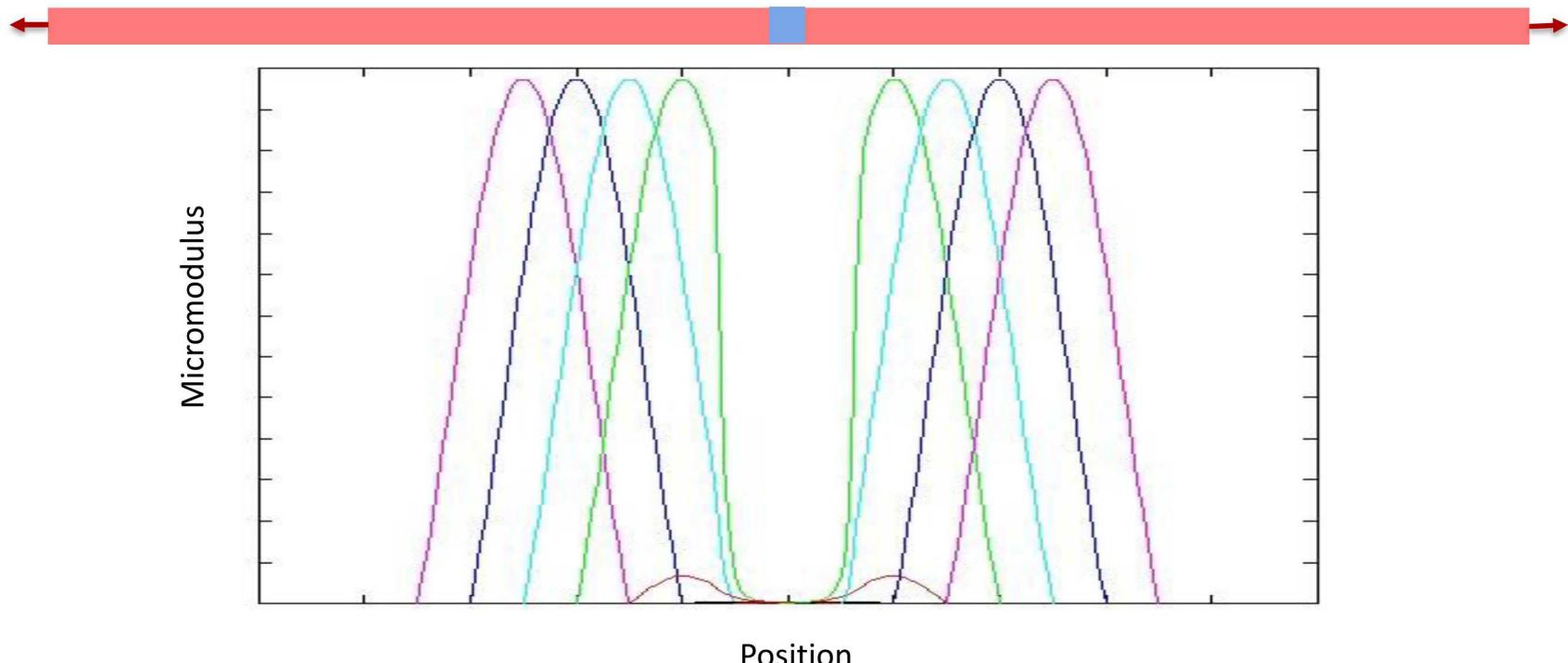
# Bar with a soft spot: Micromodulus



# Bar with a weak spot: Displacement



# Bar with a very weak spot: Micromodulus shows broken bonds



- The heterogeneous peridynamic material model zeroes out the micromodulus for bonds crossing the crack.
- Bond breakage!

# What the preceding analysis shows

- Using smoothed displacements results in a nonlocal evolution law.
- This evolution law is peridynamics provided a material model in terms of  $\bar{u}$  is defined.
- The micromodulus is determined by:
  - The small-scale (local) material model and heterogeneity.
  - The smoothing function  $w$ .
- A nonlocal concept of damage (bond breakage) emerges naturally when the original problem contains a crack.

# A hint of unexpected behavior

- Recall

$$\bar{u}(x) = \int w(x-p)u(p) \, dp.$$

- Fourier transform of any function  $v$ :

$$v^*(k) = \mathcal{F}\{v(x)\} = \int_{-\infty}^{\infty} e^{-ikx}v(x) \, dx.$$

- Convolution theorem

$$\bar{u}^* = w^*u^*$$

so that formally we can derive the small-scale displacements from any given  $\bar{u}$ :

$$u(x) = \mathcal{F}^{-1} \left\{ \frac{\bar{u}^*}{w^*} \right\}.$$

# A hint of unexpected behavior, ctd.

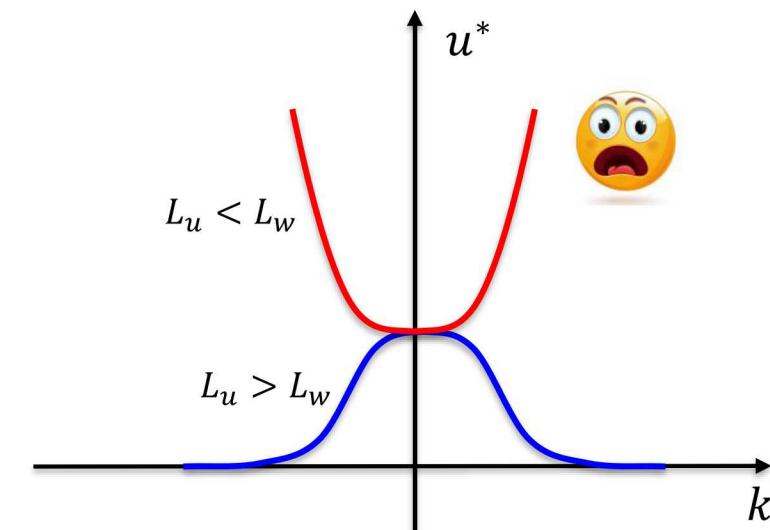
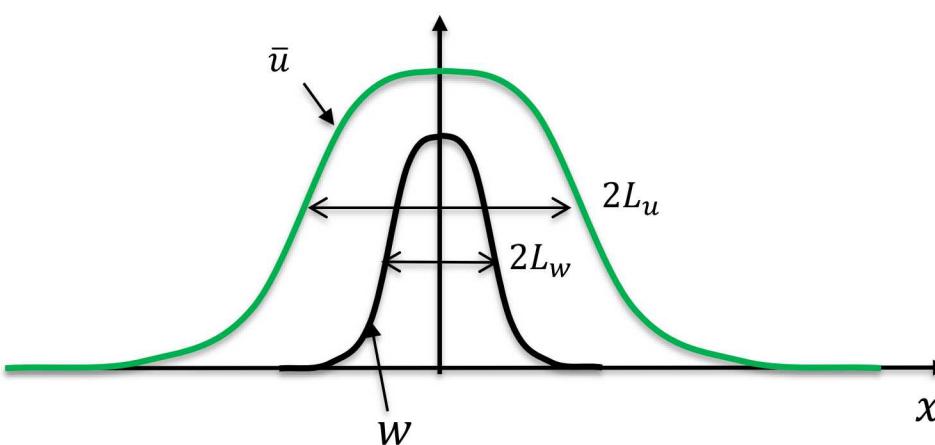
- Can we arbitrarily prescribe  $\bar{u}$ ?
- Suppose  $w$  and  $\bar{u}$  are both Gaussians:

$$\bar{u}(x) = e^{-(x/L_u)^2}, \quad w(x) = e^{-(x/L_w)^2}.$$

- Then

$$u^*(k) = \frac{\bar{u}^2(k)}{w^*(k)} = \sqrt{\frac{L_u}{L_w}} e^{\pi^2(L_w^2 - L_u^2)k^2}$$

- Bad news if  $L_u < L_w$ !

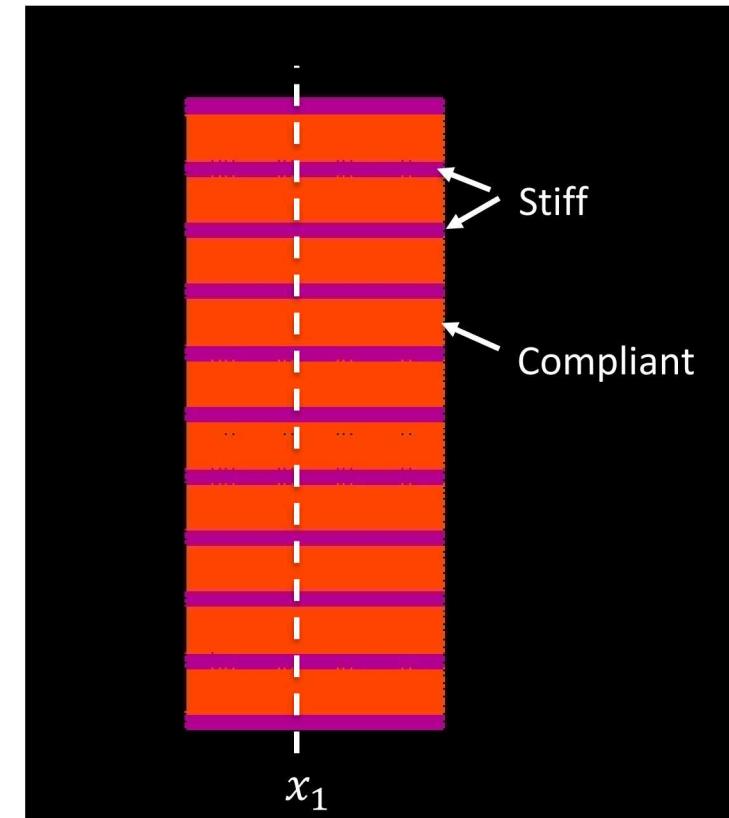


# Can nonlocality be observed experimentally in elastostatics?

- Consider a 2D composite composed of alternating layers of stiff and compliant material.
- Smoothed DOF is the average  $x$  displacement along a vertical line.

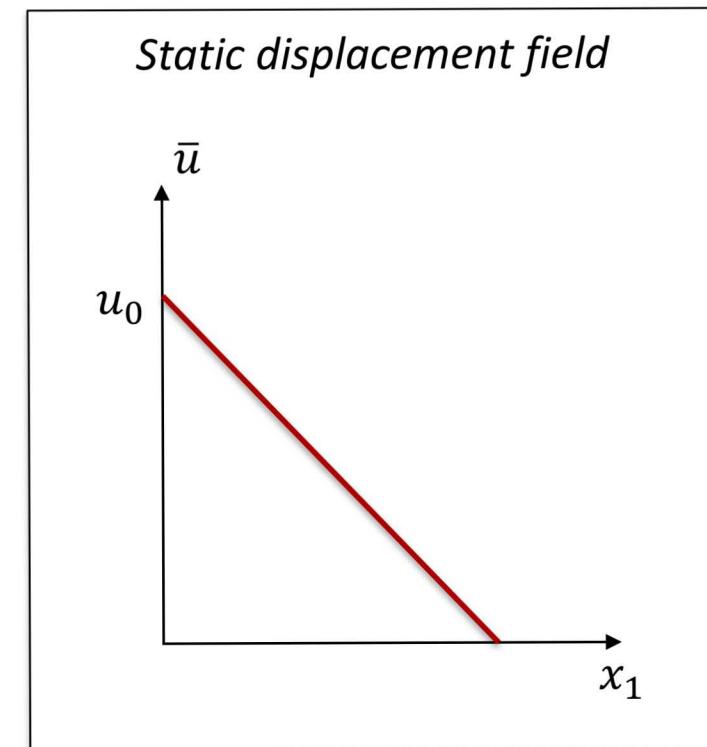
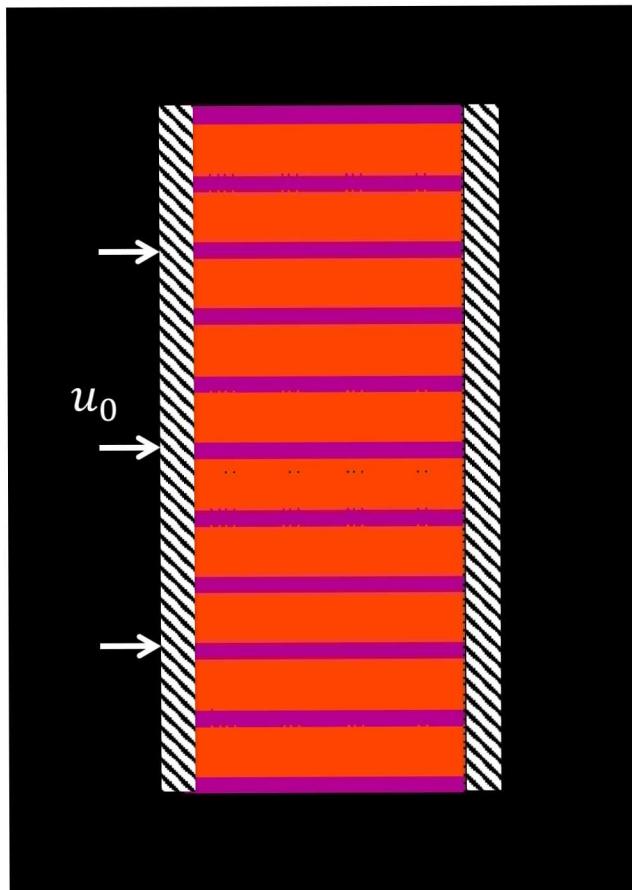
$$\bar{u} = \frac{1}{L} \int_0^L u_1 \, dx_2$$

- We will examine “seemingly” 1D deformations.



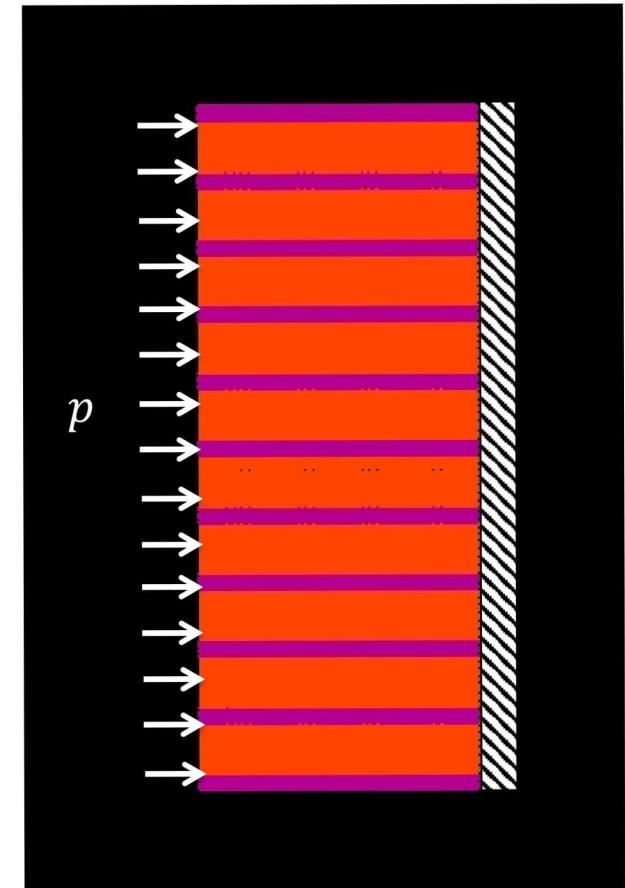
# Static Dirichlet problem for a composite

- Solve for the 2D displacements in the local theory.
- Both phases deform the same way.
- No surprises (yet).



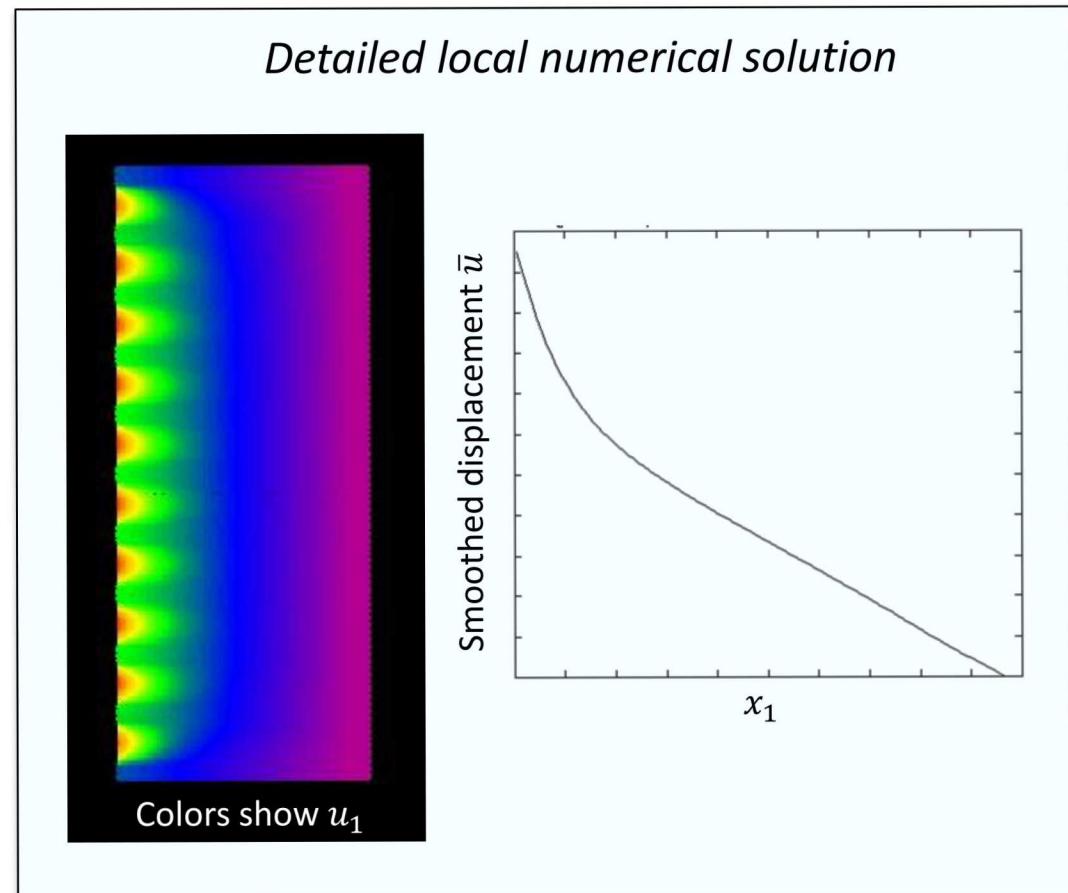
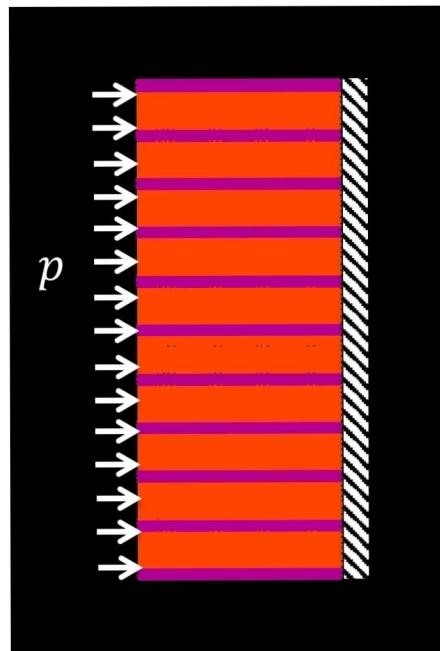
# Now consider a mixed Dirichlet/Neumann static problem

- Apply a constant traction  $p$  along the left surface.
- Still using 2D local theory.
- Should we still expect  $\bar{u}$  to vary linearly with  $x_1$ ?



# Smoothed DOFs show interesting features

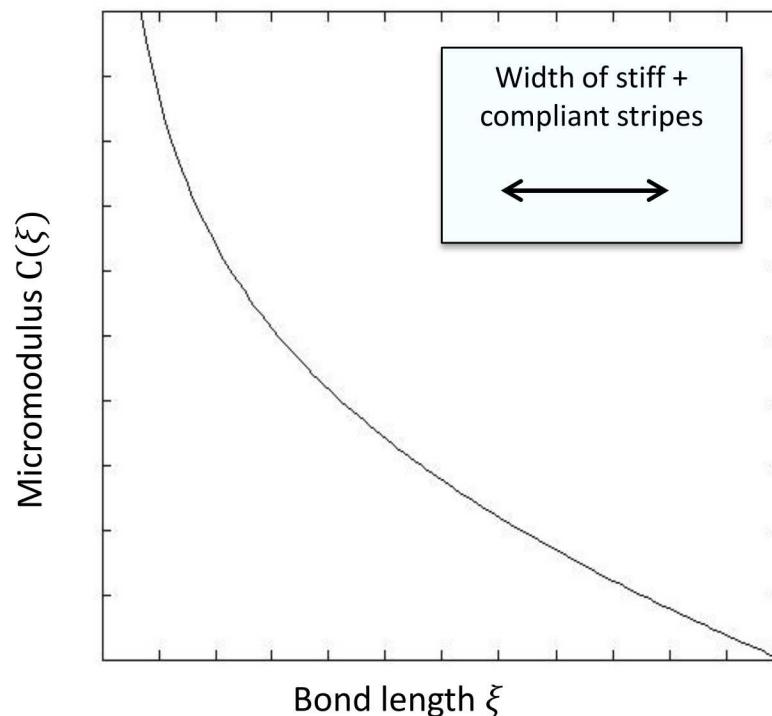
- A detail computational model shows complex behavior near the left edge.
- Smoothing this solution results in nonlinear  $\bar{u}(x_1)$ .



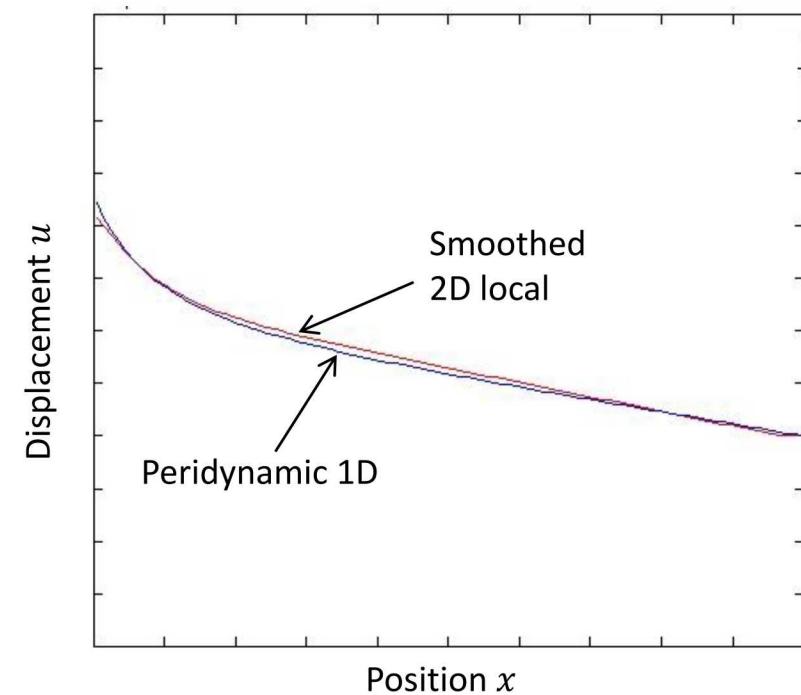
# Nonlocality helps reproduce response near loaded boundary

- Tune a 1D peridynamic microelastic material model.
- Try to reproduce the behavior seen in the detailed 2D local solution..

*Peridynamic material model*

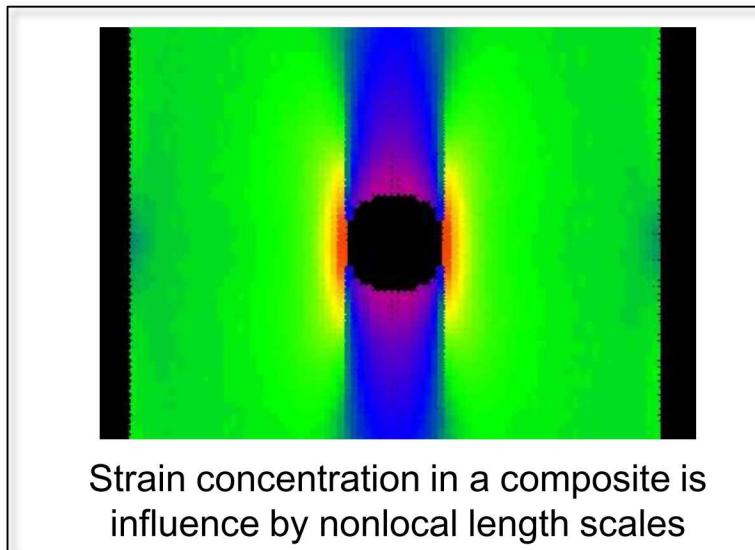


*Predicted displacement*

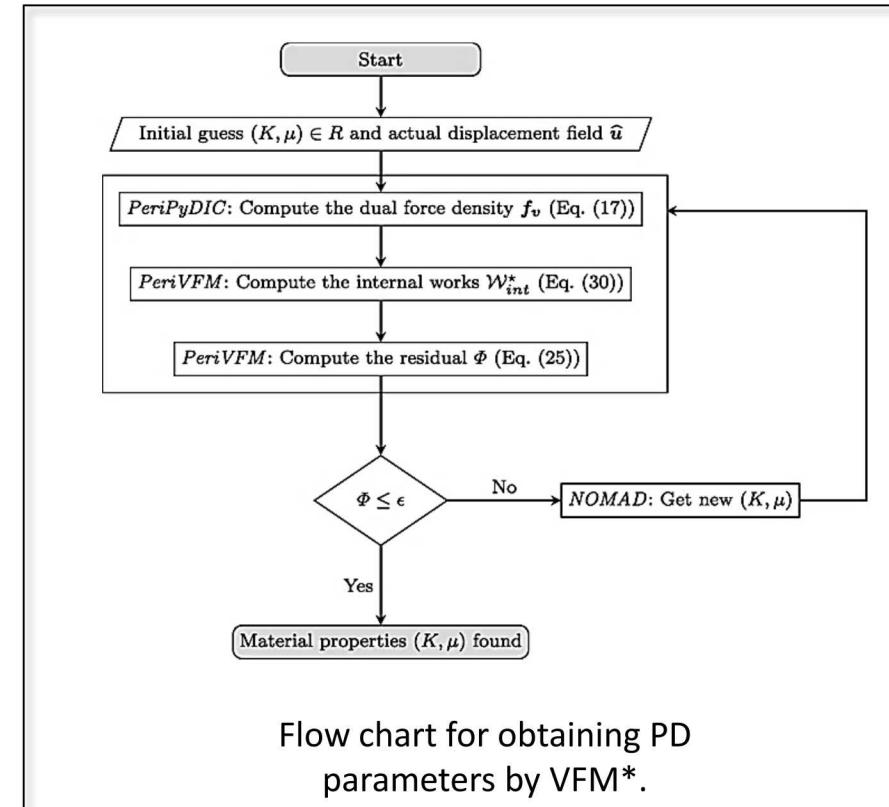


# Nonlocal material parameters can be derived from static full-field data

- Digital image correlation (DIC).
- Virtual field method (VFM).
- Electronic speckle pattern interferometry (ESPI).

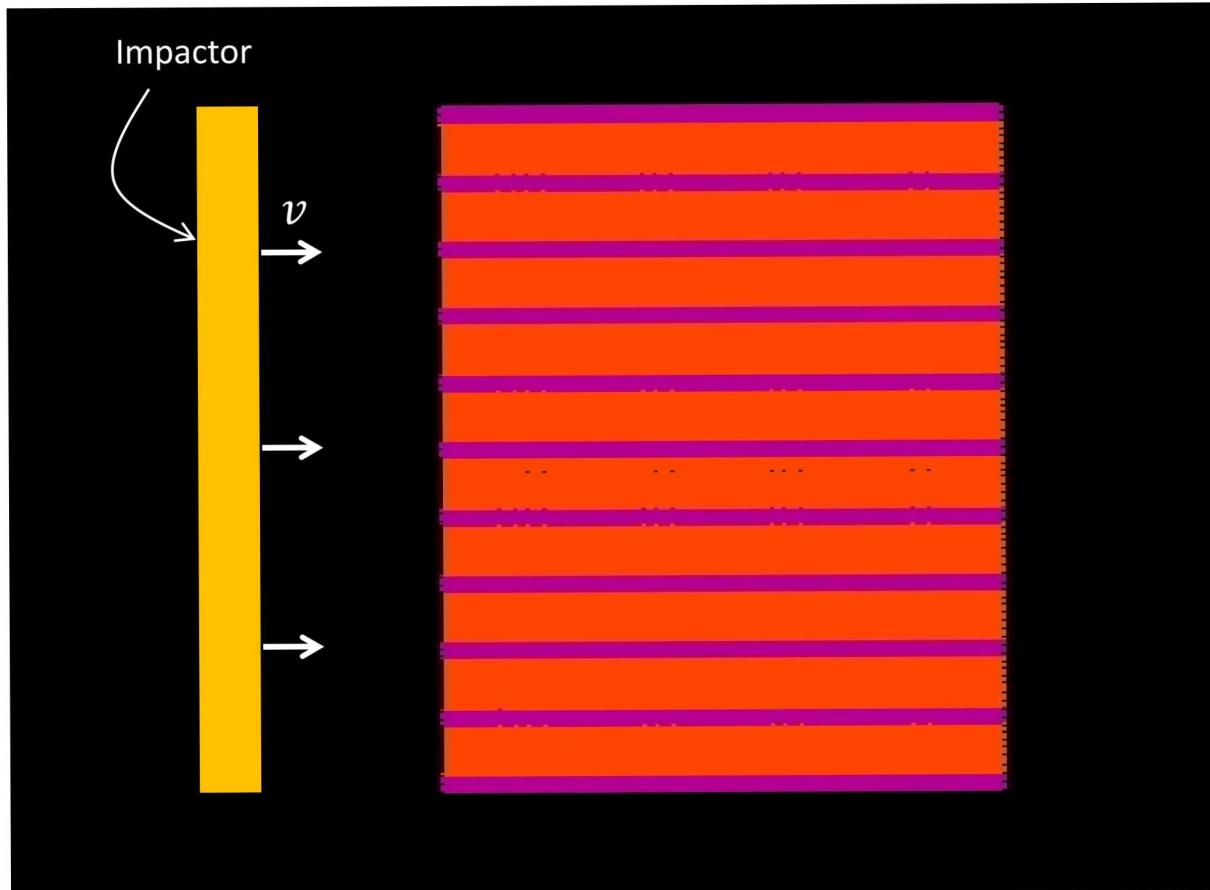


- L. Toubal, M. Karama, & B. Lorrain, *Composite structures*, (2005).
- D. Turner, B. Van Bloemen Waanders, & M. Parks J. *Mechanics of Materials and Structures* (2015).
- D. Turner, J. *Engineering Mechanics* (2015).
- \*Delorme, R., Diehl, P., Tabiai, I. et al., *J Peridyn Nonlocal Model* (2020)



# Dynamics: impact problem

- Impactor strikes the composite edge-on.



# Dynamics: impact problem video

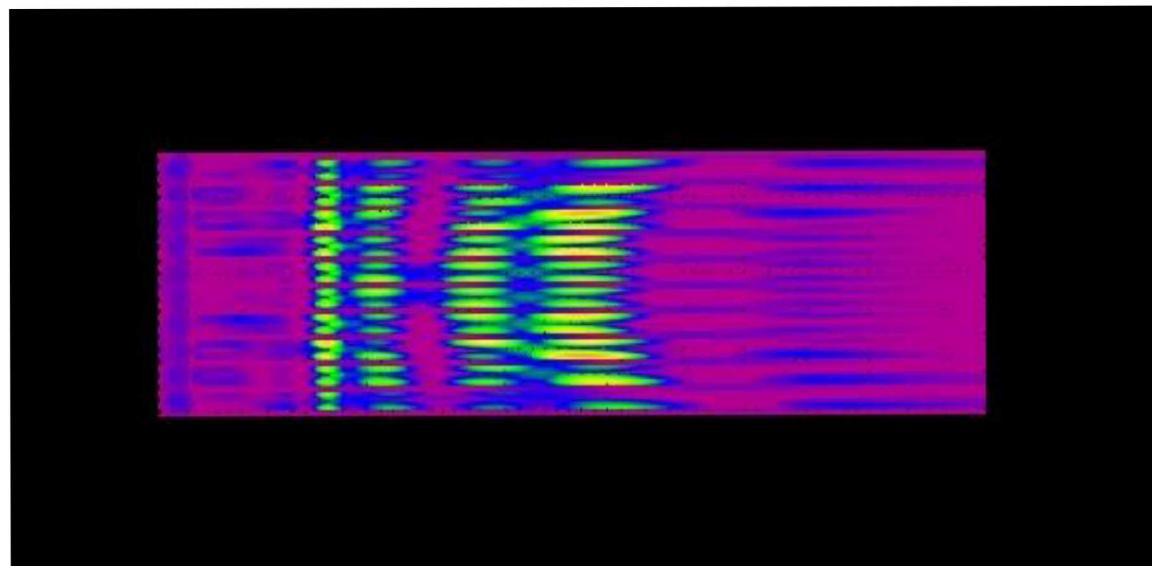
- Detailed 2D local simulation.
- Complex wave structure is created in the composite.



Colors show maximum principal strain

# Dynamics: impact problem

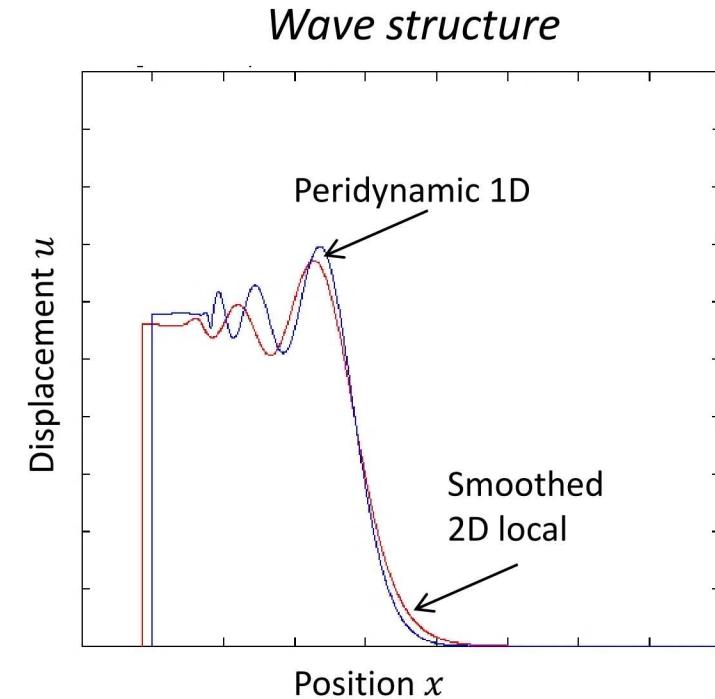
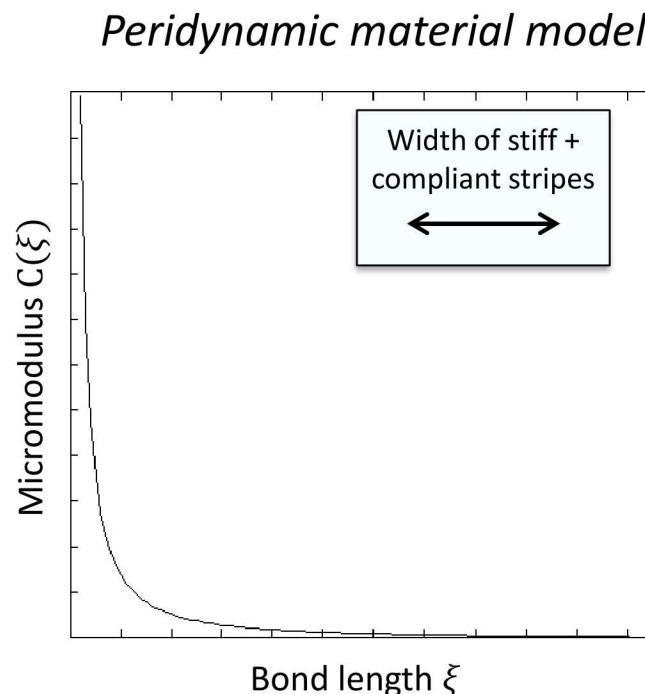
- Detailed 2D local simulation.
- Complex wave structure is created in the composite.



Colors show maximum principal strain

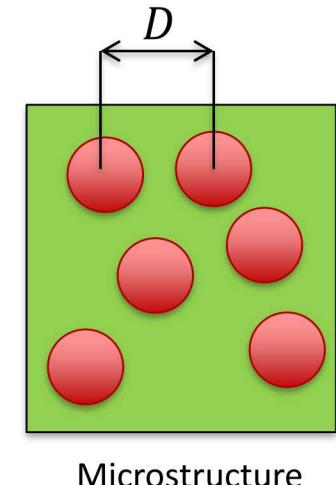
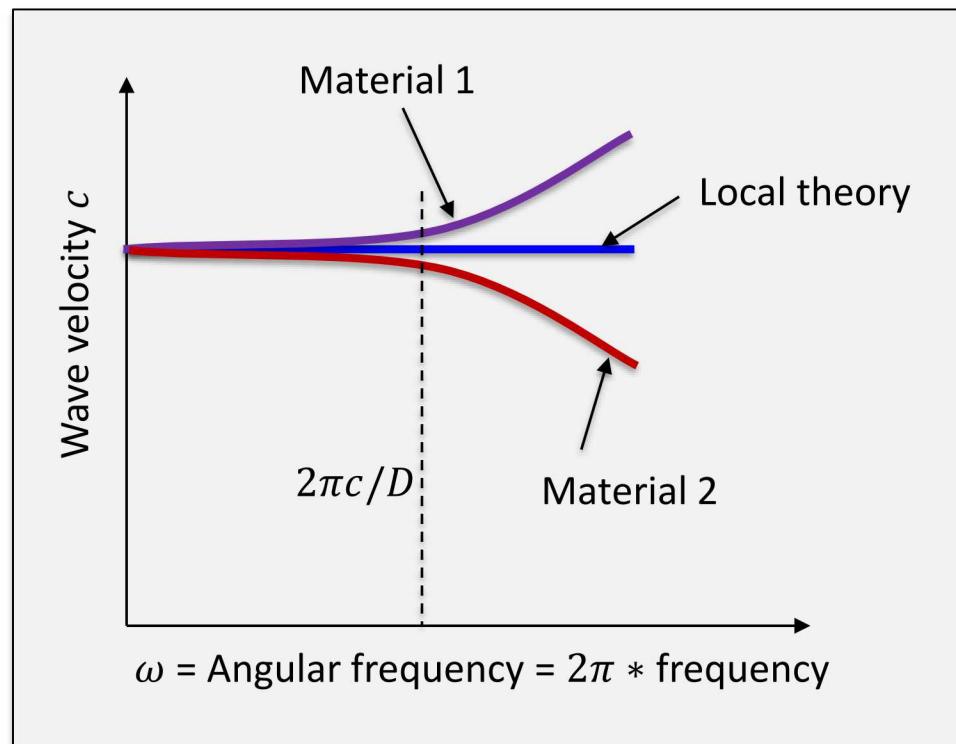
# Nonlocality helps predict the dispersive nature of waves in the composite

- After smoothing the displacement along vertical lines, the complex wave structure is manifested as dispersion.
- A 1D peridynamic model (after tuning of the micromodulus) reproduces some of these features.



# Wave dispersion

- All real solids exhibit dispersion for sufficiently short wavelengths.
- The wavelength depends on the microstructure and composition.
  - Dispersion starts to appear for **wavelengths < microstructure size**.
  - This implies that nonlocality is required to predict dispersion.



Microstructure

# Wave dispersion in linear peridynamics

- Equation of motion with  $b \equiv 0$ :

$$\rho \ddot{u}(x, t) = \int_{-\delta}^{\delta} C(\xi)(u(x + \xi, t) - u(x, t)) d\xi$$

- Look for plane wave solutions of the form

$$u(x, t) = e^{i(kx - \omega t)}$$

where  $k$ =wavenumber and  $\omega$ =angular frequency.

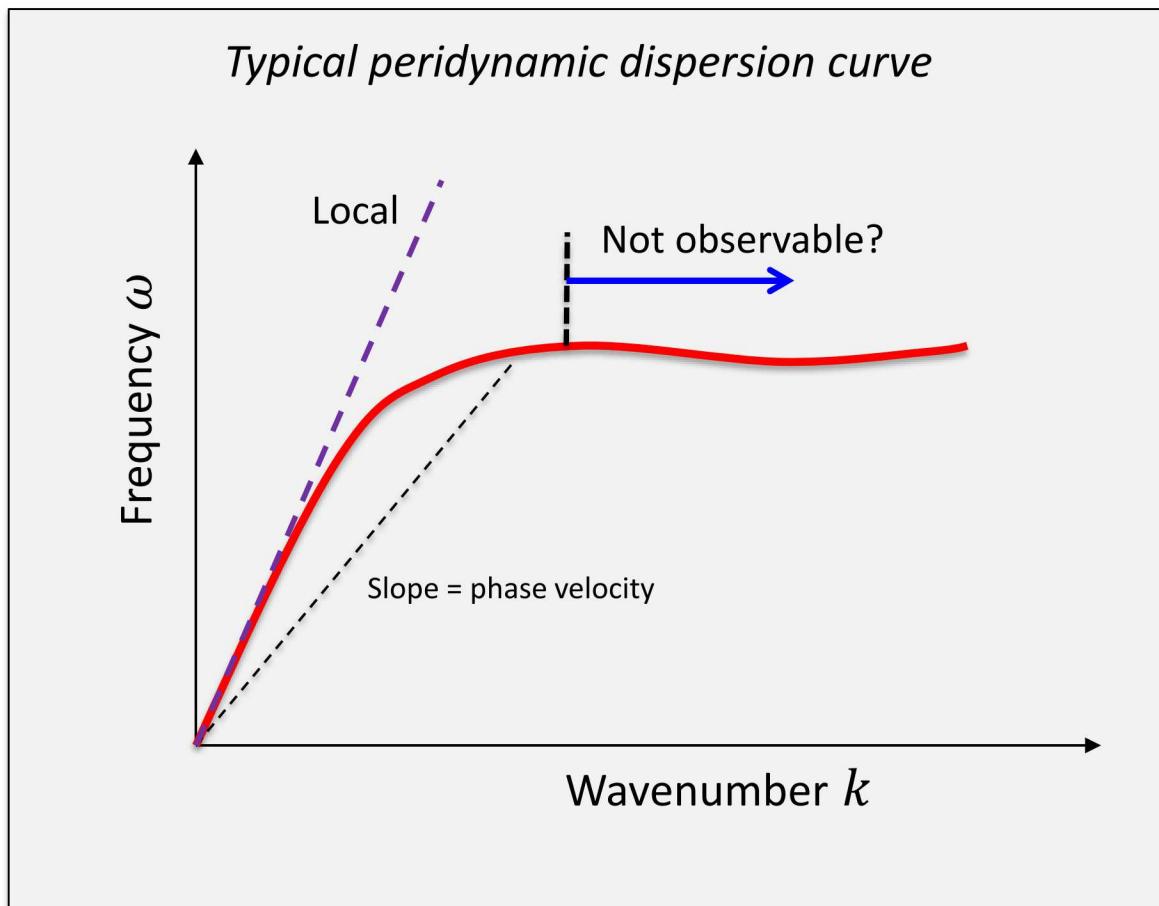
- Condition on  $\omega$  and  $k$ :

$$-\rho\omega^2 = \int_{-\delta}^{\delta} C(\xi) e^{ik\xi} d\xi - P, \quad P := \int_{-\delta}^{\delta} C(\xi) d\xi$$

- or in terms of the Fourier transform  $C^* = \mathcal{F}\{C\}$ ,

$$\rho\omega^2(k) = P - C^*(k)$$

# Wave dispersion in linear peridynamics



- S. N. Butt, J. J. Timothy, & G. Meschke, *Computational Mechanics* (2017).
- V. S. Mutnuri, USNCCM15 presentation (2019).

# Finding peridynamic material properties from measured dispersion data

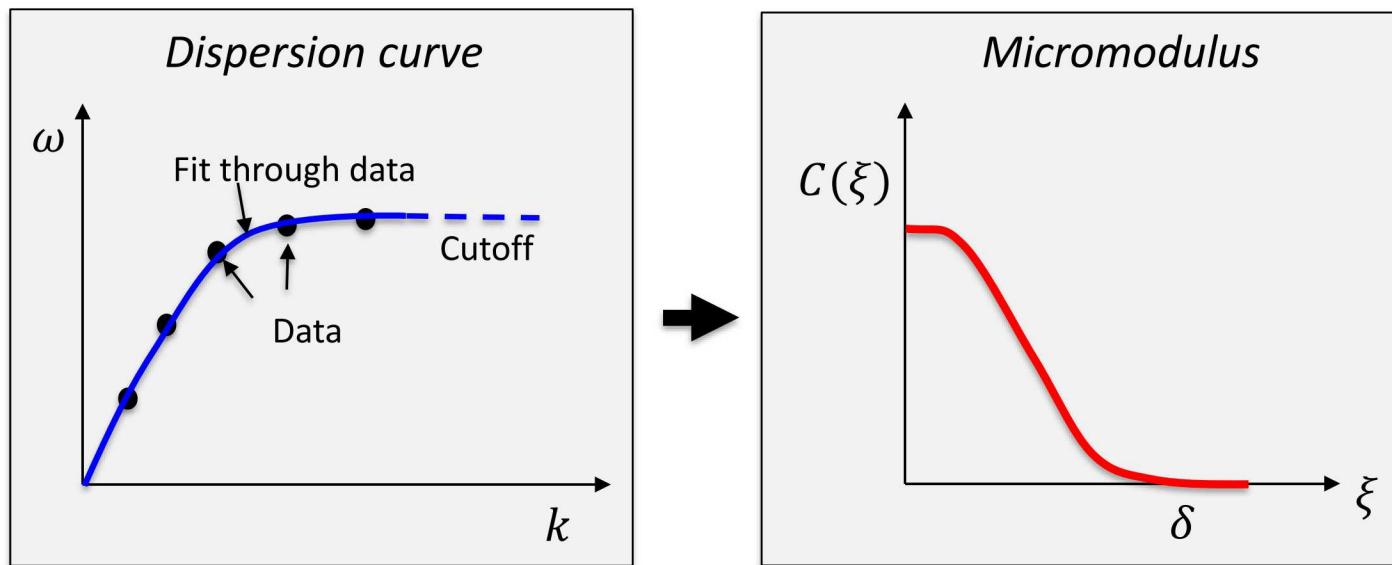
- We found

$$\rho\omega^2(k) = P - C^*(k).$$

- Given measured  $\omega_{exper}(k)$ , formally solve

$$C(\xi) = \mathcal{F}^{-1}\{P - \rho\omega_{exper}^2(k)\}$$

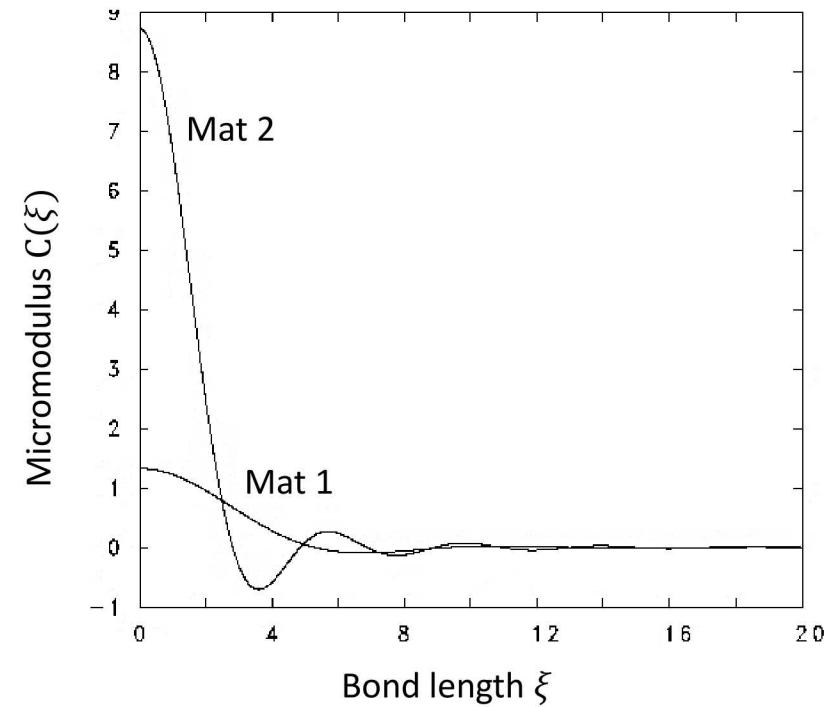
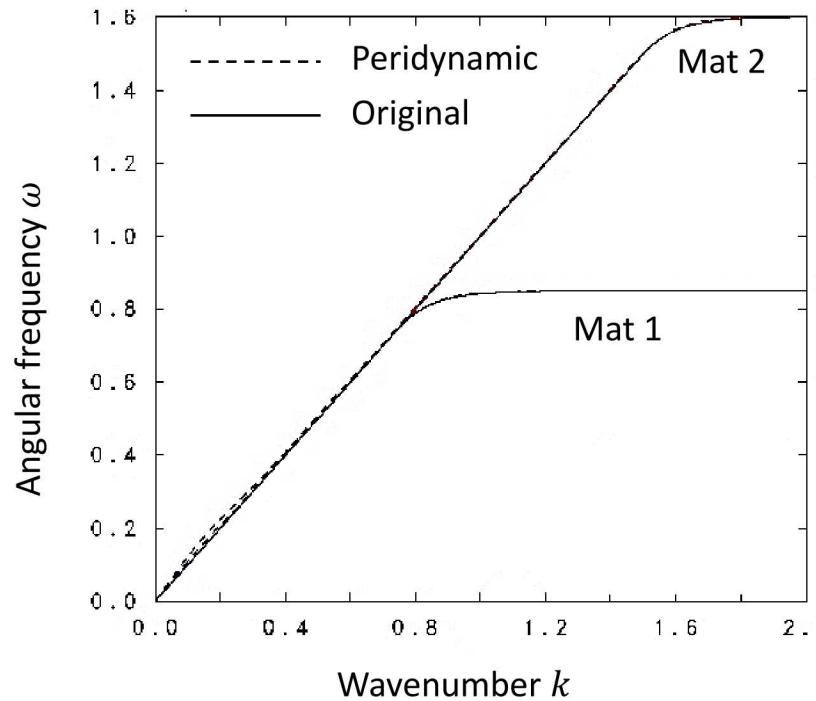
(requires data to be cut off for large  $k$ ).



- O. Weckner & S.S., *Int. J. for Multiscale Computational Engineering* (2011).

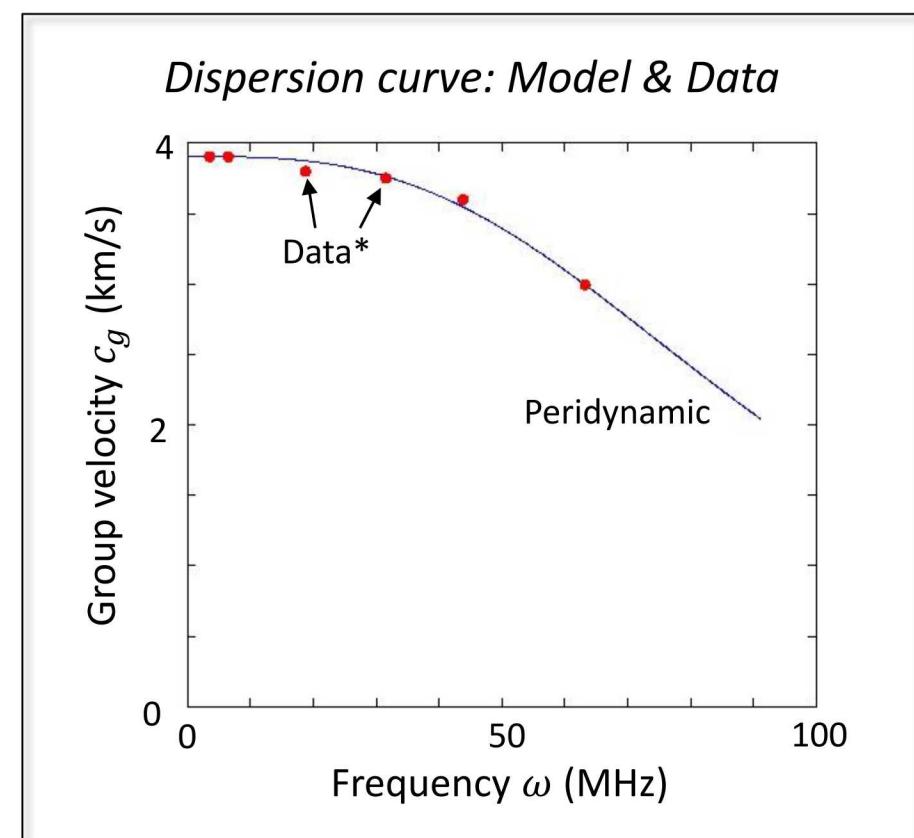
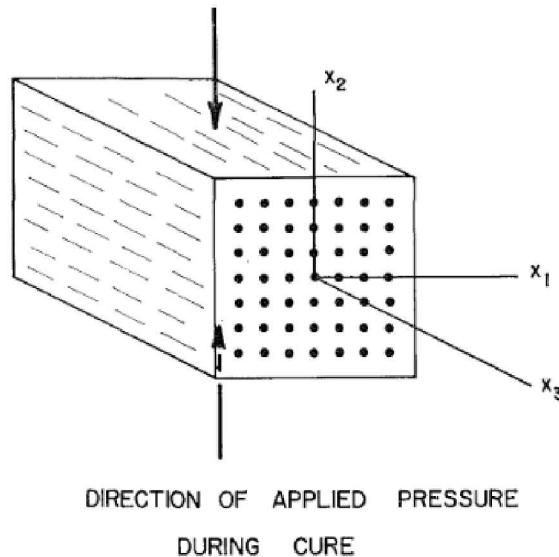
# Higher cutoff frequency leads to narrower micromodulus curve

- The limiting case of micromodulus  $\rightarrow$  delta function corresponds to the local theory.



# Example: PD model calibrated to a composite dispersion curve

- Boron-epoxy composite.
- Longitudinal waves normal to fibers.
- Compare measured ultrasonic group velocity\* with calibrated peridynamic result.



\* T. R. Tauchert & A. N. Guzelsu, *J. Applied Mechanics* (1972).

# Discussion: Nonlocality in peridynamics

- Nonlocality emerges from how we choose to model a problem.
- Origins
  - Long-range forces
  - Smoothed degrees of freedom
  - Multiple pathways for flux (of momentum, heat, mass, ...)
- Consistency
  - Peridynamics uses a consistently nonlocal approach to the evolution of all fields including damage.