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Weak anti-localization in undoped Ge/GeSn heterostructures

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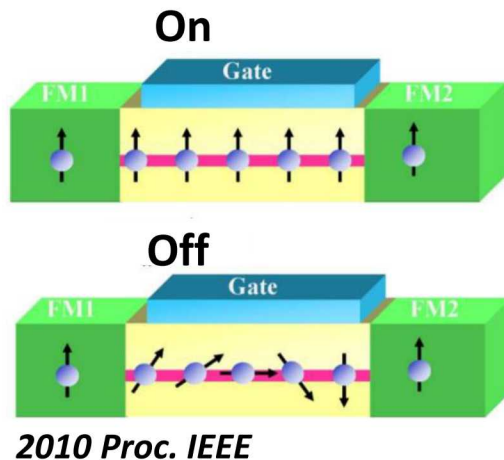
Outline

- ◆ Motivation
- ◆ Epi structures and device structures
- ◆ Magneto-transport at low temperature
- ◆ Weak antilocalization and Spin-orbit Coupling
- ◆ Conclusions

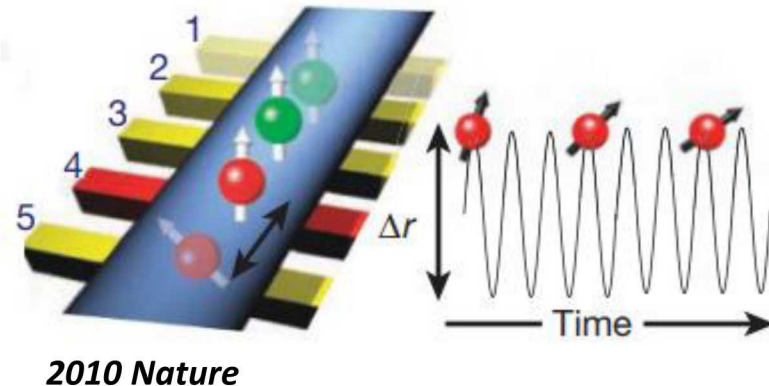
Motivation

Spin computing devices

Spin-FETs



Spin-orbit qubits

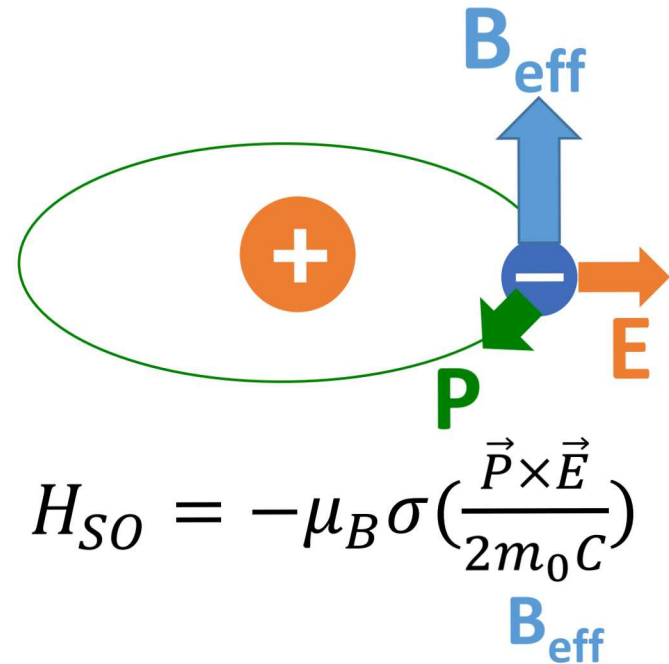
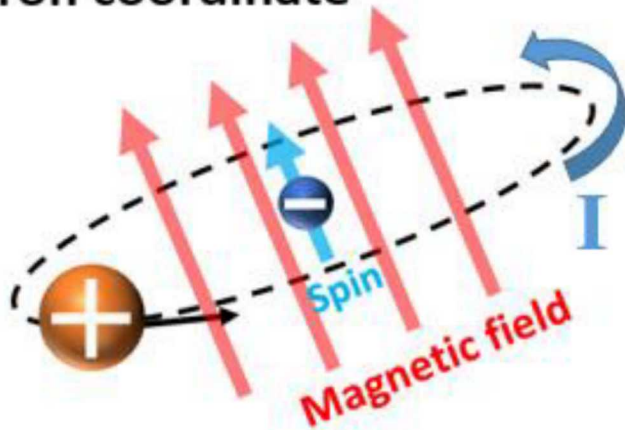


- ◆ Spin FETs → Manipulate spins via SOC to control on/off
- ◆ Spin-based QC → Manipulate spins via SOC to operate qubits

Motivation

Spin manipulation → SOC

Electron coordinate



$$H_{SO} = -\mu_B \sigma \left(\frac{\vec{P} \times \vec{E}}{2m_0 C} \right)$$

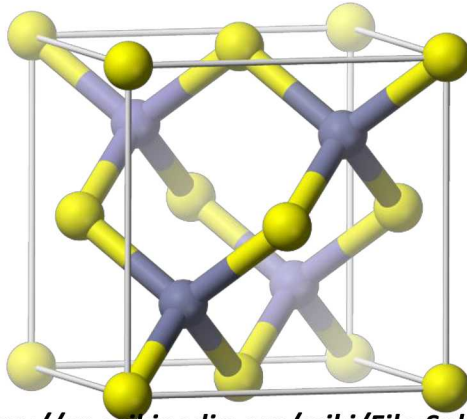
B_{eff}

- ◆ e⁻ will feel an effective magnetic field due to the relative motion between e⁻ and atom
- ◆ Spin orientation changes by the effective magnetic field

Motivation

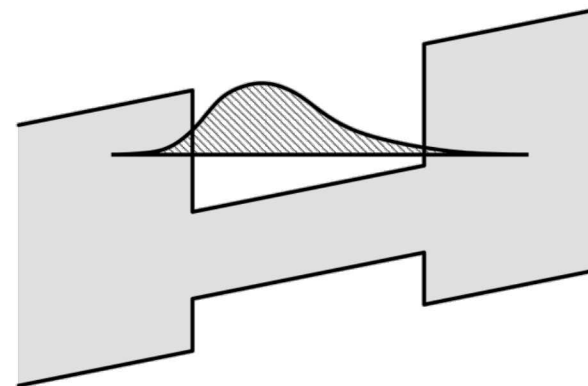
Dresselhaus and Rashba Spin-orbit Coupling

Dresselhaus



<https://en.wikipedia.org/wiki/File:Sphalerite-unit-cell-depth-fade-3D-balls.png>

Rashba



Winkler, Spin-Orbit Coupling Effects in Two-Dimensional Electron and Hole Systems, 2003.

◆ Dresselhaus SOC

- Arising from bulk inversion asymmetry
- Exist in zinc blende structure (III-V material)

◆ Rashba SOC

- Arising from structure inversion asymmetry
- **Tunable by external gating**

Motivation

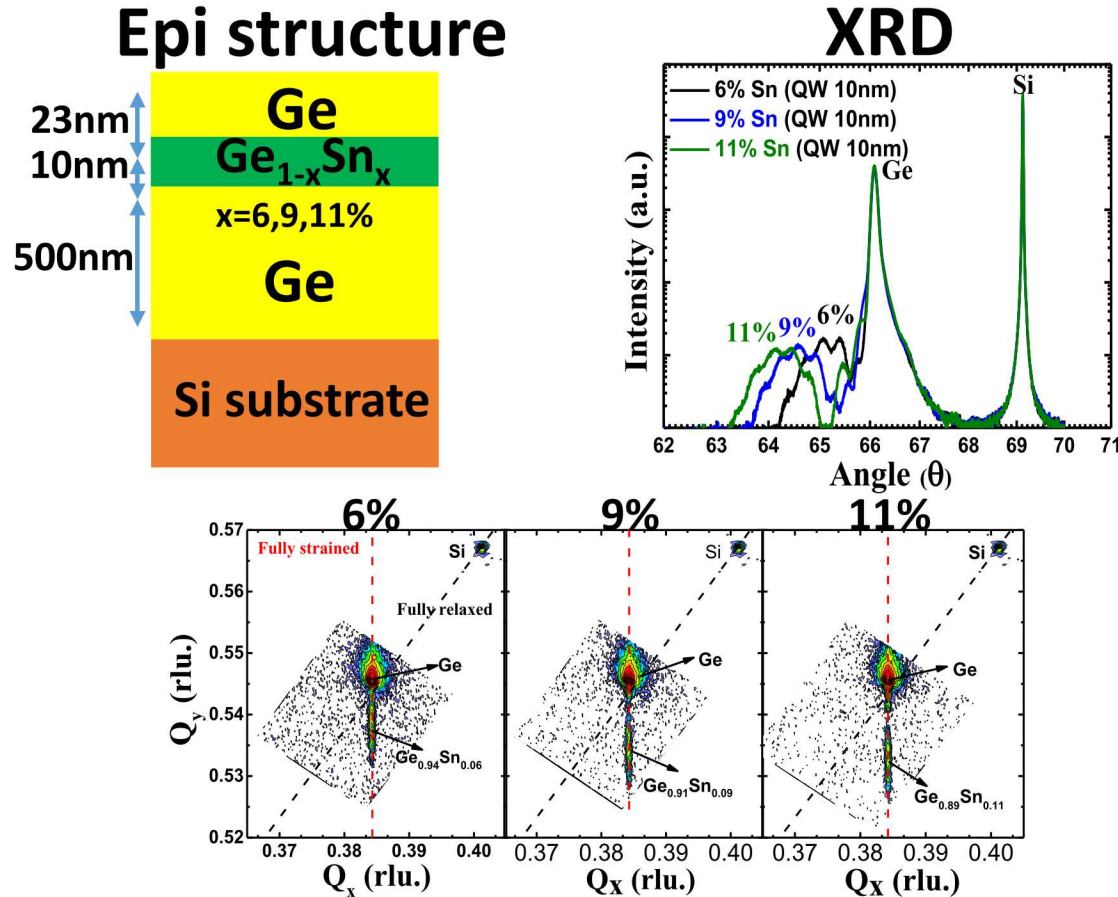
Why focus on GeSn SOC?

	Ge 2DHG (C)	2014,PRL Ge 2DHG (C) ²⁴	2002,PRB Si 2DEG ²³ (L)
Low Temperature Mobility (cm ² /Vs)	450 000	~5 000	200 000
Rashba spin-orbit parameter α or β	1.0×10^{-28} eVm ³	0.2×10^{-28} eVm ³	5.5×10^{-15} eVm
Rashba spin-orbit energy (meV)	1.4	0.30	~0.001
Lattice mis-match strain	0.8%	2.1%	...
			2014,APL



- ◆ GeSn is expected to have larger SOC, and is compatible to Si VLSI Technology

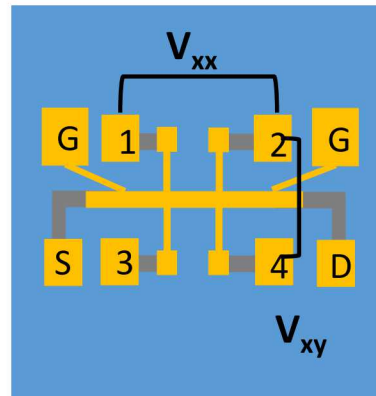
Ge/GeSn Epitaxial Structures



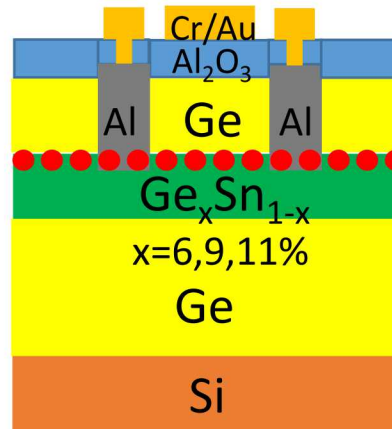
- ◆ [Sn]=6,9,11% and thickness = 10 nm by XRD
- ◆ GeSn is fully strain on Ge layer with $\epsilon = -0.57\%$, -0.96% , -1.24%

Device Structure

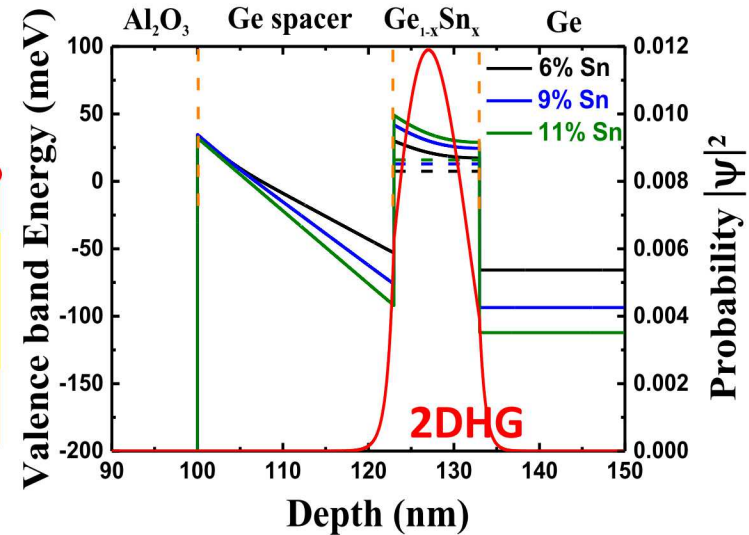
Top view



Side view

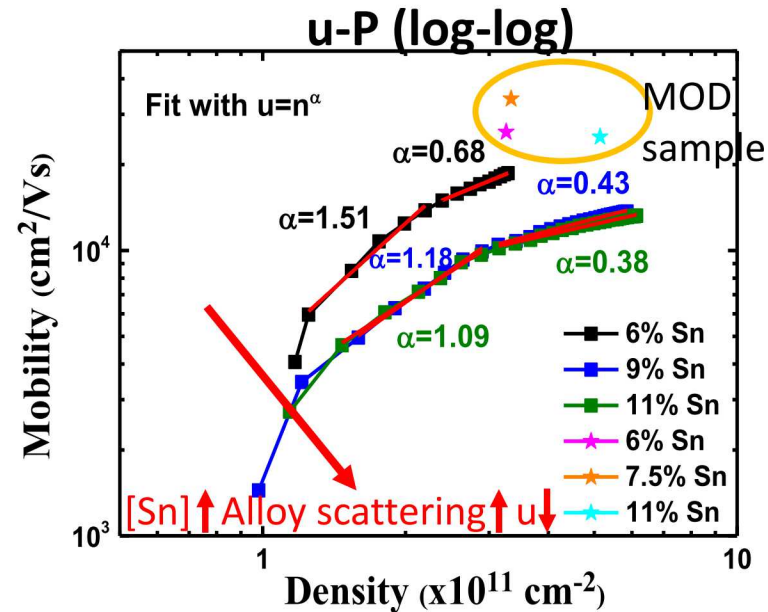
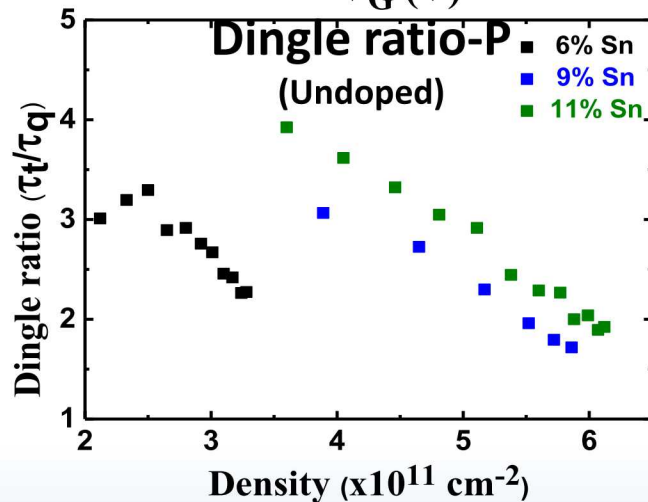
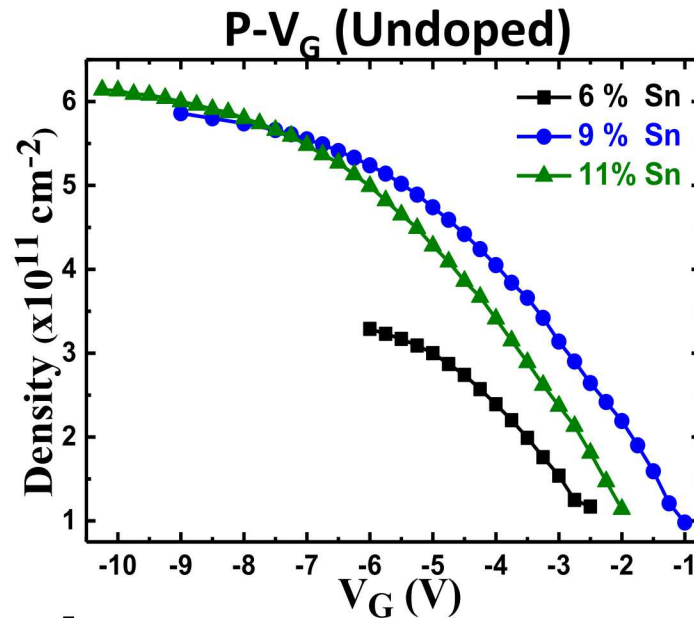


Simulation band diagram



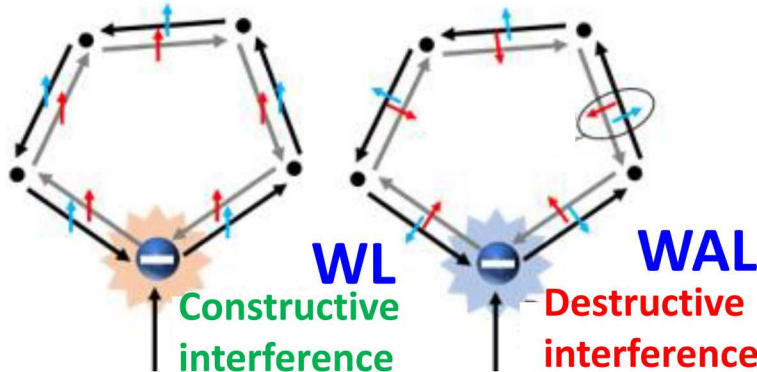
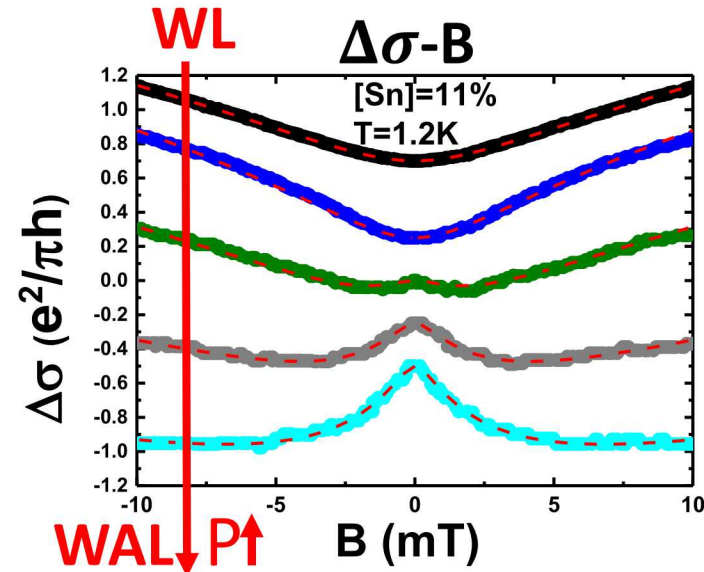
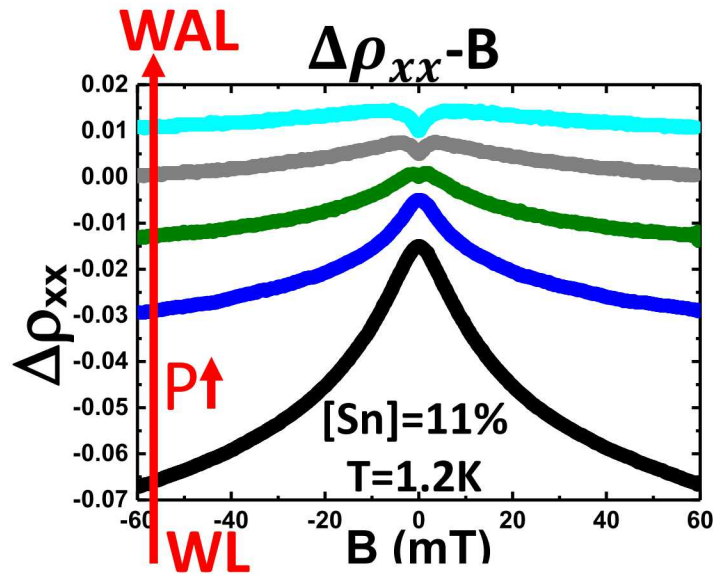
- ◆ Two-dimensional hole gas (2DHG) formed in the GeSn quantum well by gating
- ◆ Band diagram simulated by Schrodinger + Poisson
 - Valence band offset at Ge/GeSn creates a QW
 - [Sn]↑ Valence band offset↑ Quantum confinement↑

Hall measurement at T=1.2K



- ◆ Hole density increase with V_G and saturated at high bias
- ◆ $0.5 \cong \alpha \leq 1.5$ and small dingle ratio
→ Background charge scattering (large angle scattering)

Weak Antilocalization Effect

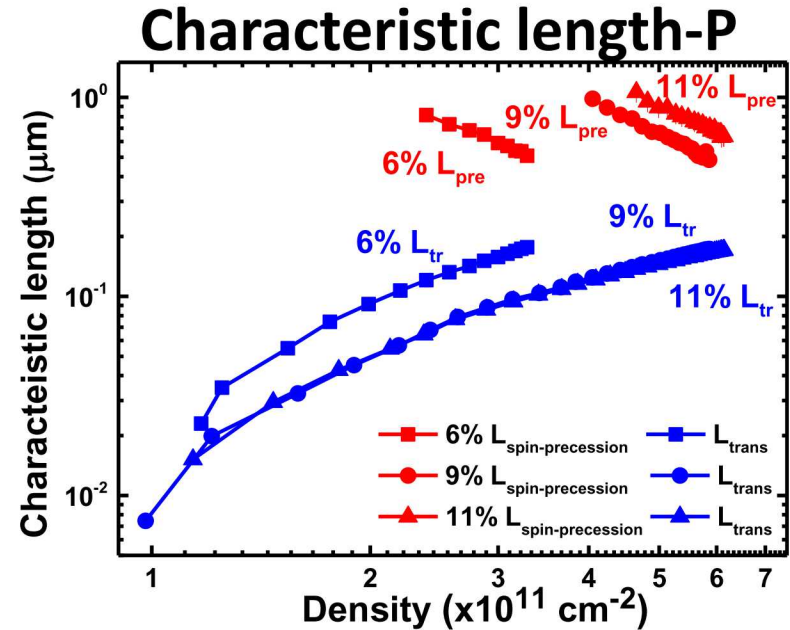
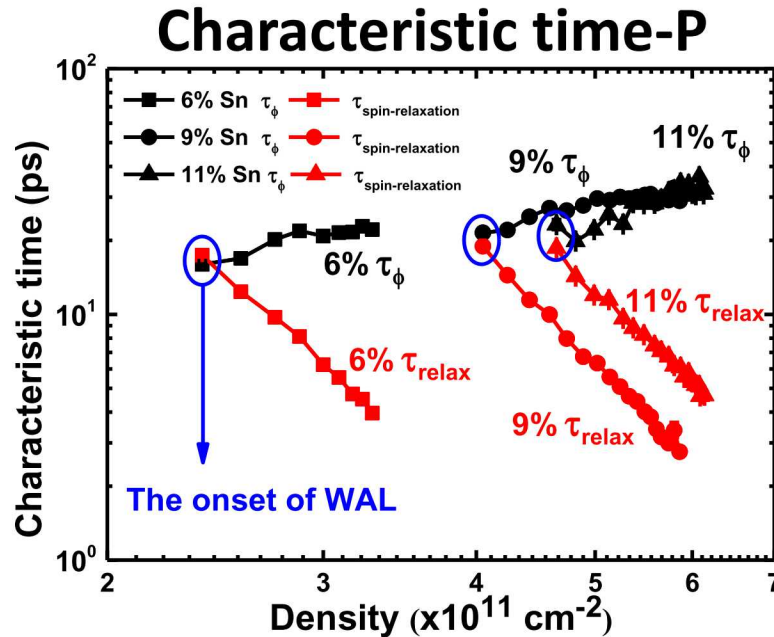


$\Delta\sigma(B) - \Delta\sigma(0)$ HLN formula 2012,PRL

$$= \frac{e^2}{2\pi^2\hbar} \left\{ \psi \left(\frac{1}{2} + \frac{B_\phi}{B} + \frac{B_{SO}}{B} \right) + \frac{1}{2} \psi \left(\frac{1}{2} + \frac{B_\phi}{B} + \frac{2B_{SO}}{B} \right) - \frac{1}{2} \psi \left(\frac{1}{2} + \frac{B_\phi}{B} \right) - \ln \frac{B_\phi + B_{SO}}{B} - \frac{1}{2} \ln \frac{B_\phi + 2B_{SO}}{B} + \frac{1}{2} \ln \frac{B_\phi}{B} \right\}$$

- ◆ The $\Delta\rho_{xx}$ shows the transition from WL $\rightarrow$ WAL when $P \uparrow$
- ◆ The red-dashed line shows the fitting result with HLN formula

Characteristic Time/Length



$$B_{SO}(\phi) = \frac{\hbar}{4eD\tau_{SO}(\phi)}$$

DP spin-relaxation mechanism

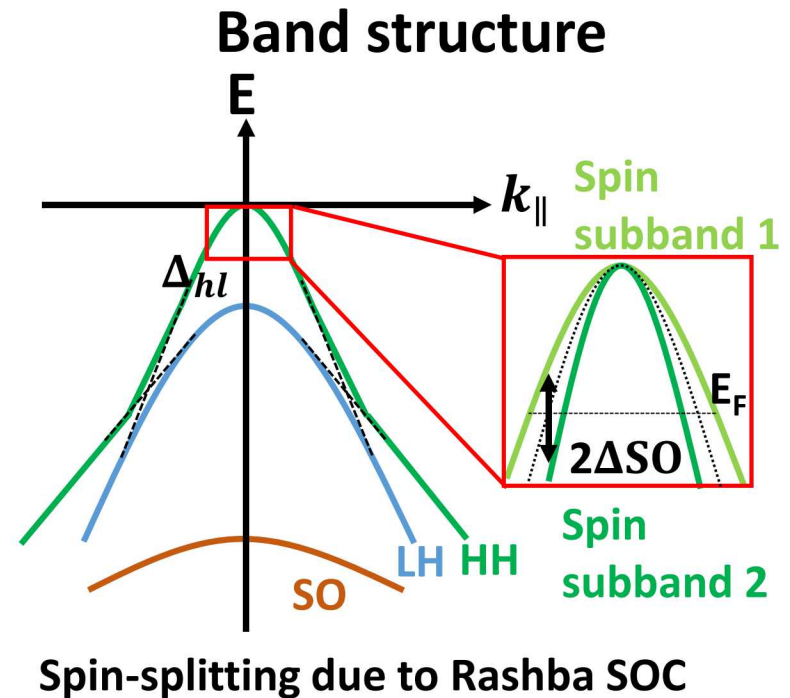
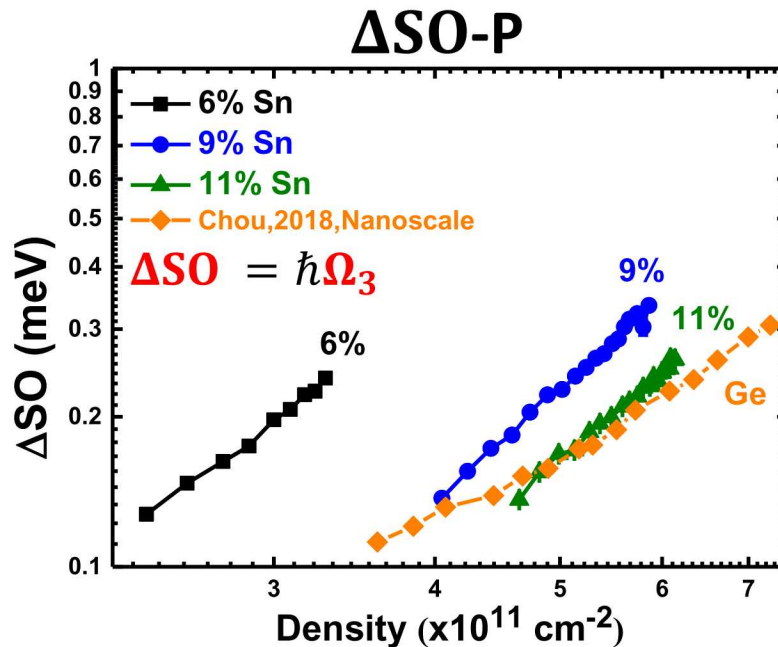
$$\frac{1}{\tau_{SO}} = 2|\Omega_3|^2\tau_{trans}$$

Spin precession frequency

2018, *Nanoscale*.

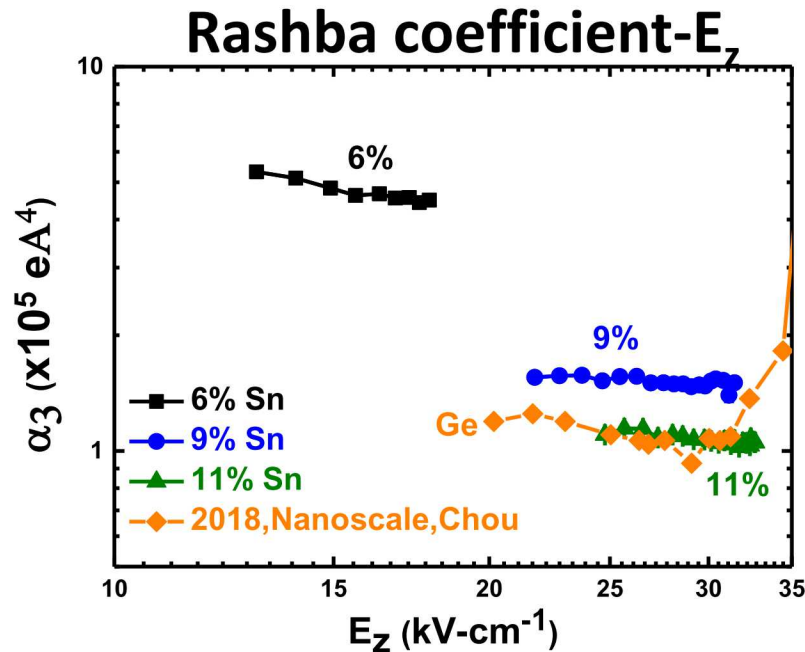
- ◆ [Sn] $\uparrow$, τ_{relax} $\uparrow$, SOC $\downarrow$ $\rightarrow$ Due to strain effect
- ◆ $L_{\text{precess}} < L_{\text{transport}}$ (spin-diffusive regime) $\rightarrow$ HLN is eligible

Spin-subband splitting



- ◆ $[Sn] \uparrow, \Delta SO \downarrow, SOC \downarrow \rightarrow$ Due to strain effect
- ◆ Spin degeneracy of HH band broken by SOC

Rashba Coefficient Extraction



$$\Delta SO = \alpha_3 E_z (k_F)^3$$

$$H_{SO} = \alpha_3 E_z i(k_-^3 \sigma_+ - k_+^3 \sigma_-)$$

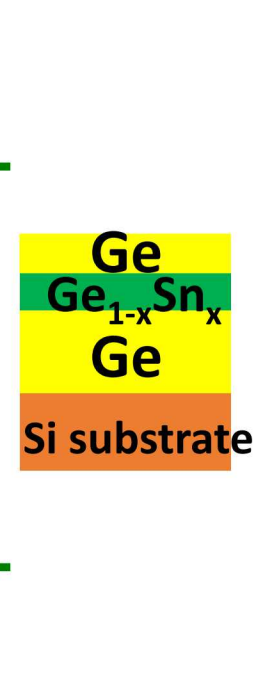
2018, Nanoscale

$$\alpha_3 = a\gamma_3(\gamma_2 + \gamma_3) \frac{e\hbar^4}{m_0^2} \left[\frac{1}{\Delta_{11}^{hl}} \left(\frac{1}{\Delta_{12}^{hl}} - \frac{1}{\Delta_{12}^{hh}} \right) + \frac{1}{\Delta_{12}^{hl} \Delta_{12}^{hh}} \right]$$

2004, APL

Strong
quantization $\parallel Z$
(HH-LH splitting)

Weak
quantization $\perp Z$
(Rashba)



◆ [Sn] $\uparrow$, α_3 $\downarrow$ $\rightarrow$ HH-LH splitting affects Rashba SOC

◆ HH-LH splitting quantized in z direction
 B_{Rashba} quantized in x-y plane } $\Delta_{hl} \uparrow \alpha \downarrow$

Conclusions

- ◆ Demonstration of magneto-transport in Ge/GeSn QW structures
- ◆ Clear transition from WL→WAL, and Rashba coefficient, spin-subband splitting energy has extracted
- ◆ The experiment result shows a negative strain effect on SOC

Acknowledgements



- ◆ National Taiwan University (NTU)
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- ◆ Center for Integrated Nanotechnologies (CINT)
- ◆ Sandia National Laboratories

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Thank you very much!

Back up slide

Fabrication Process flow

Fabrication process flow of Hall-bar Undoped

Undoped

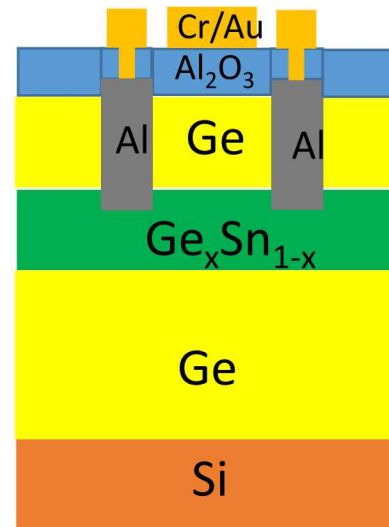
Deposit Al
as contact

Grow Al_2O_3
by ALD

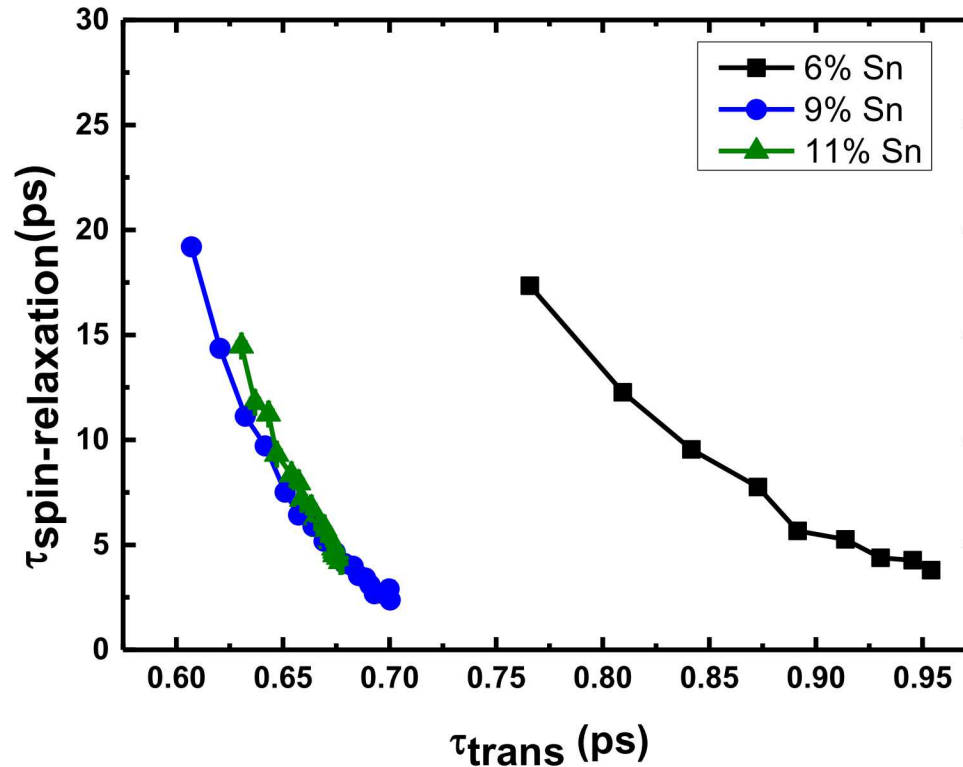
Open via by
BOE

Deposit
Cr/Au

Etch back to
define gate
and contact

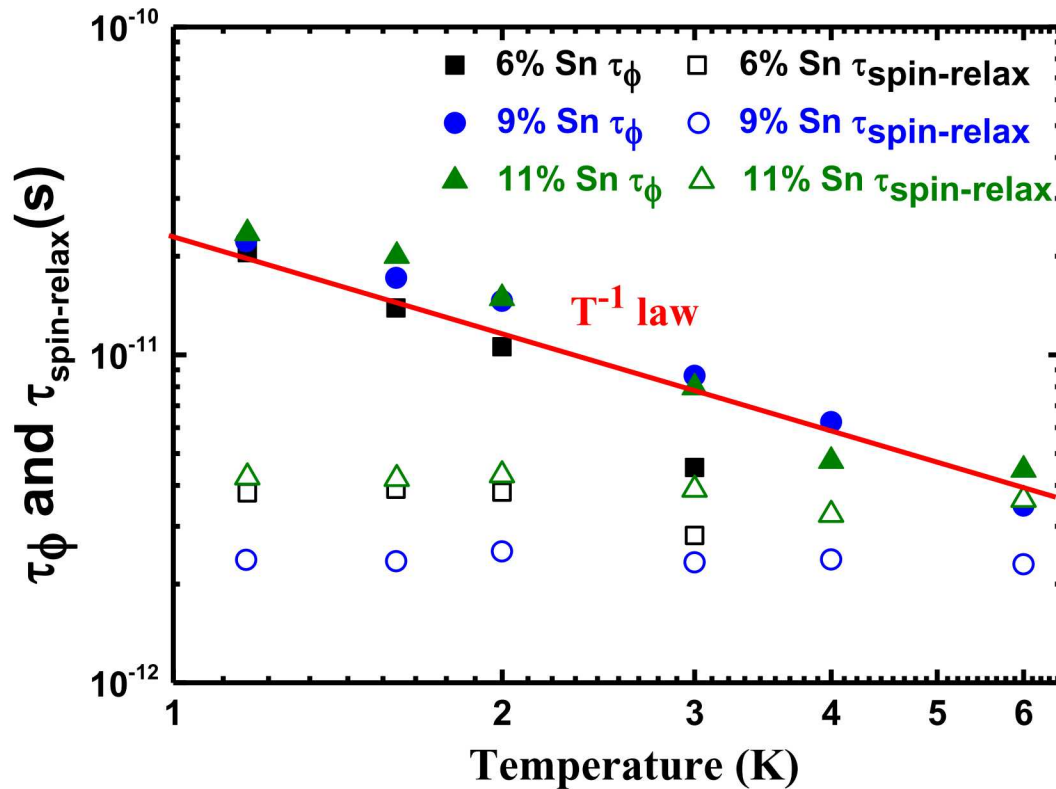


D'yakonov-Perel (DP) spin relaxation



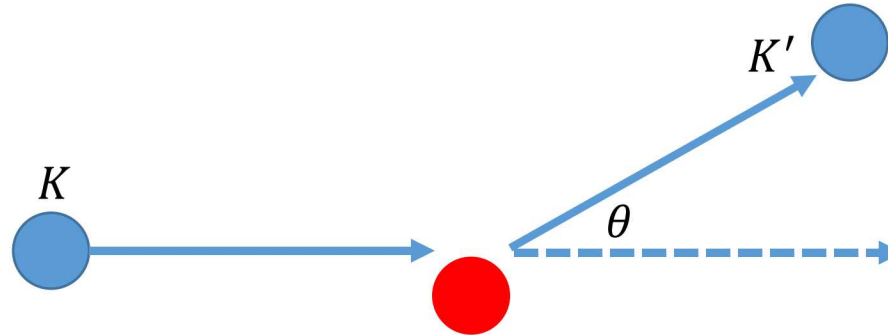
- ◆ $\tau_{spin-relax}$ is inverse proportional to τ_{trans}
→ An indication of DP spin-relaxation mechanism

Temperature Dependence



- ◆ $\tau_\phi \propto T^{-1} \rightarrow e^-e^-$ interaction dephasing is dominant
- ◆ $\tau_{\text{spin-relax}}$ is independent of temperature

Dingle-ratio



$$\alpha = \frac{\tau_{trans}}{\tau_q} = \frac{\frac{1}{\tau_q}}{\frac{1}{\tau_{trans}}} = \frac{\int w(k, k')}{\int w(k, k')(1 - \cos(\theta))}$$

- ◆ Remote charge scattering (small angle scattering)
→ Dingle ratio ↑
- Background charge scattering (large angle scattering)
→ Dingle ratio close to unity