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Hole spins in Ge/SiGe heterostructures

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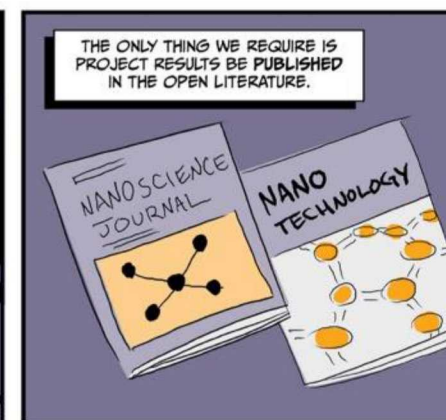
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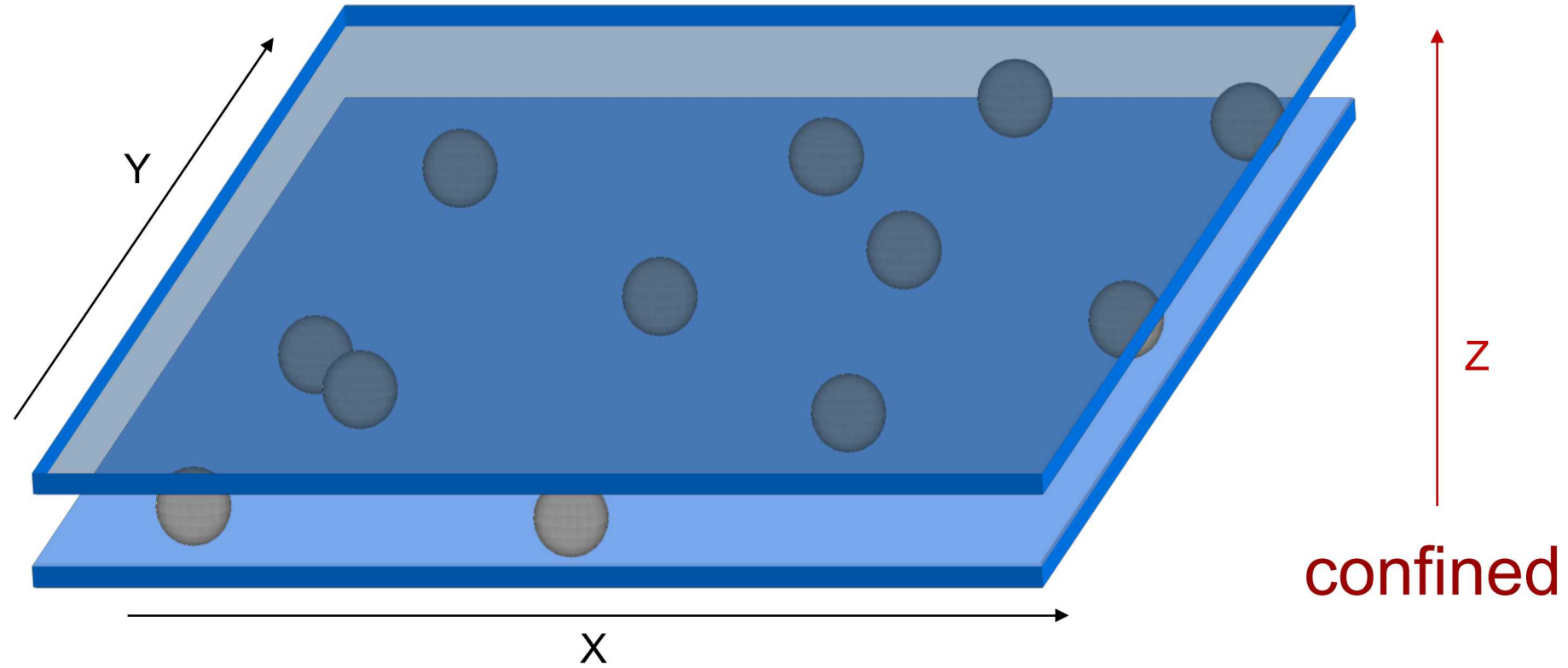
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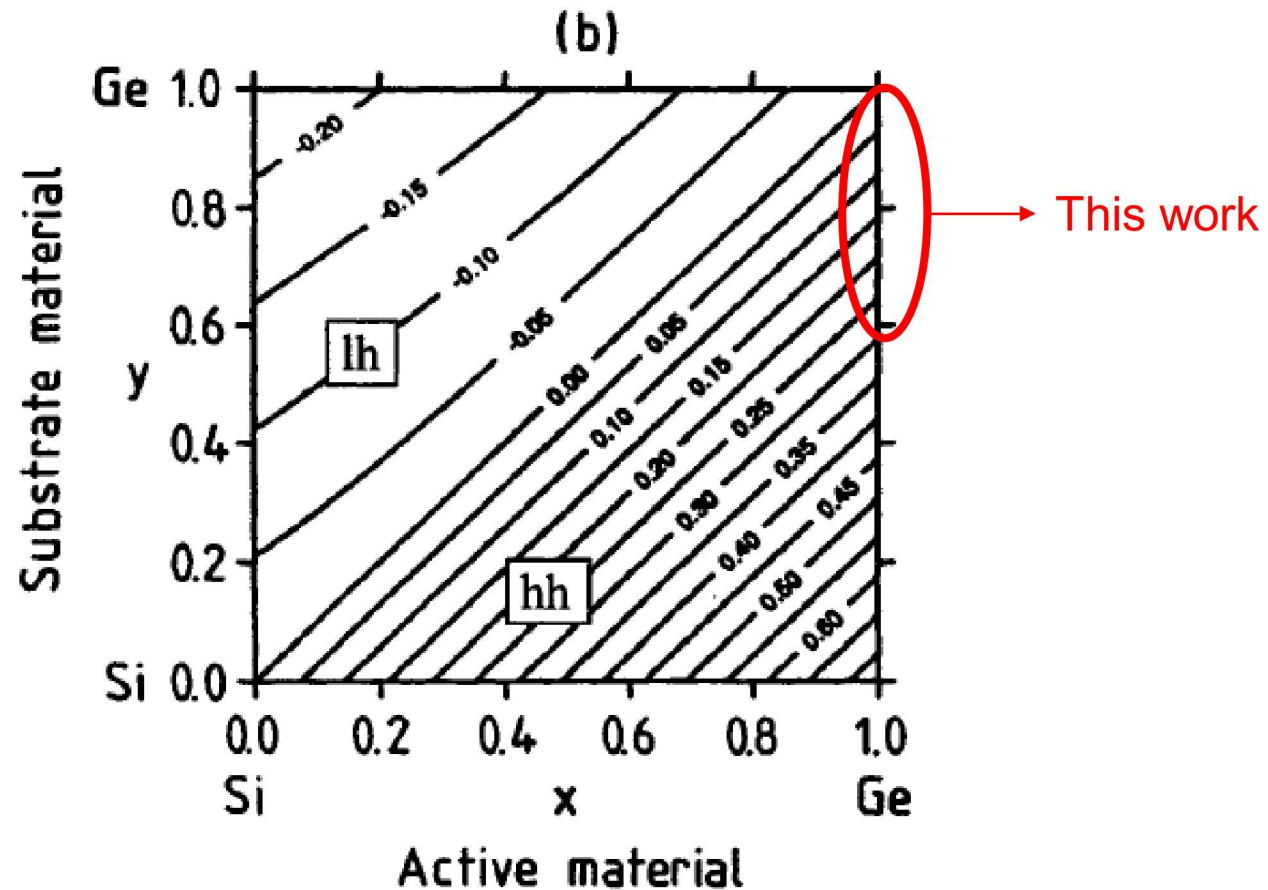
Outline

- Holes in Ge/SiGe heterostructures
 - Physical properties
 - Quantum Hall ferromagnetic transition
 - Spin qubits in Ge/SiGe

2D holes

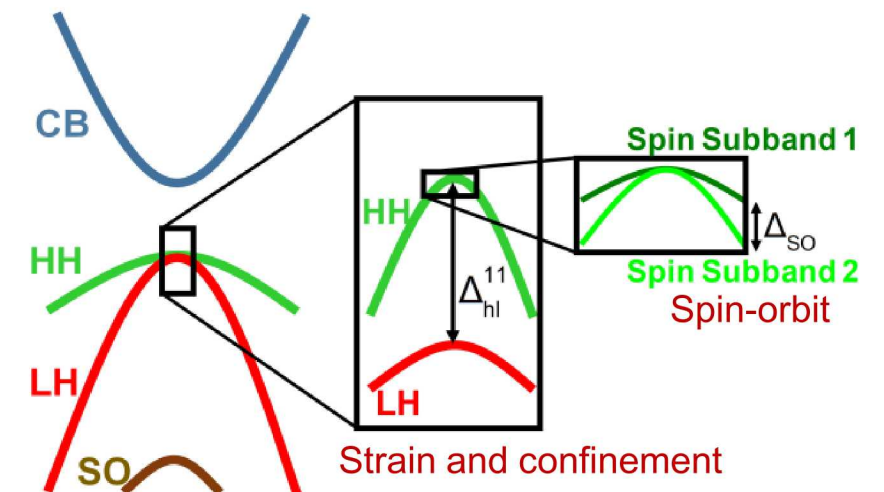
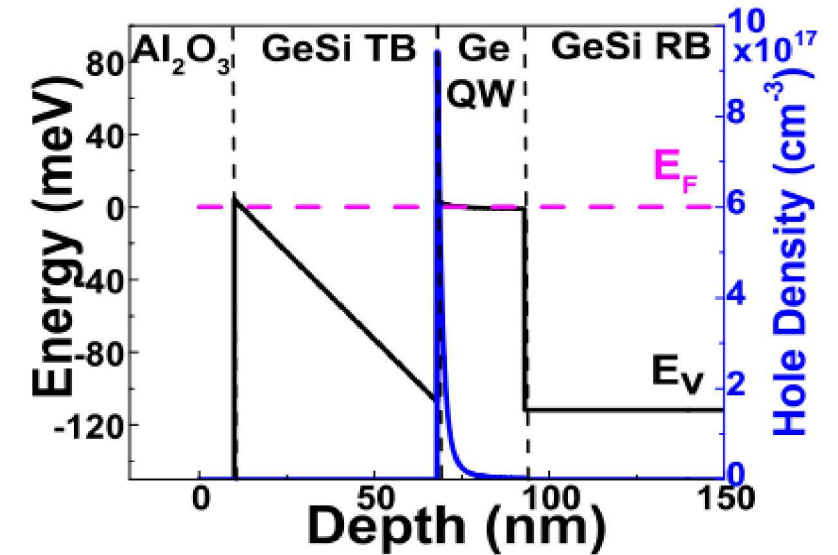
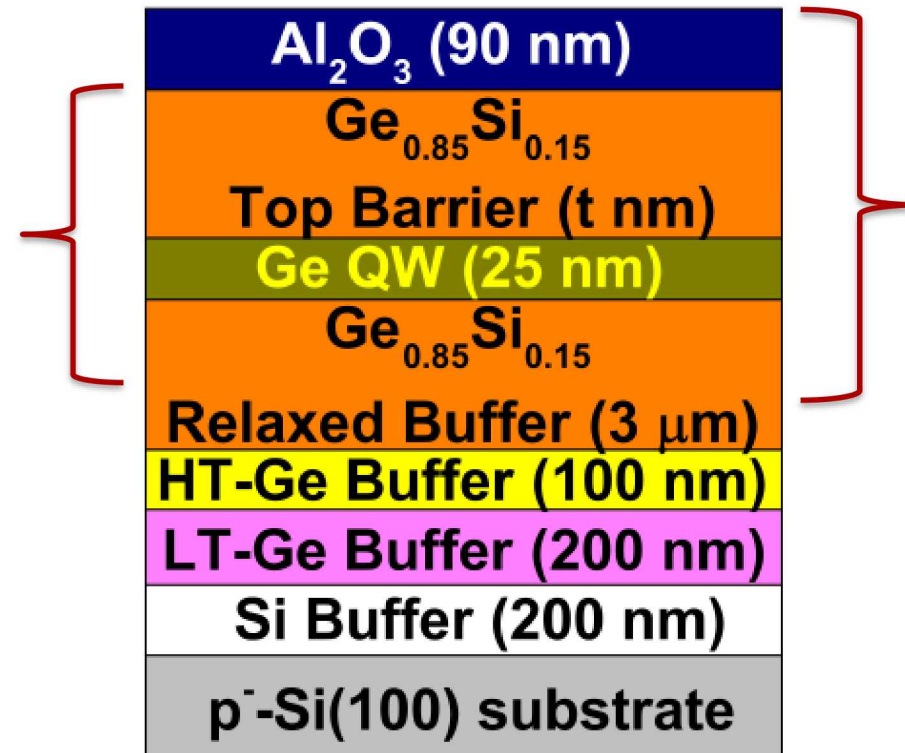
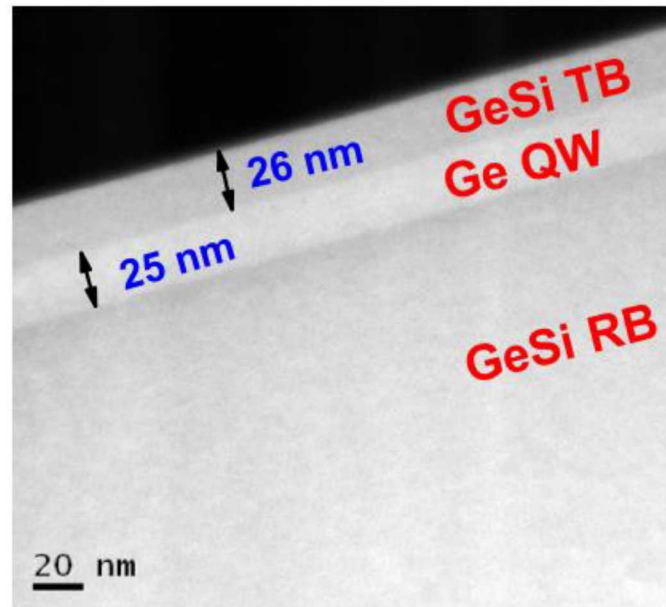


Band alignment of SiGe heterostructures



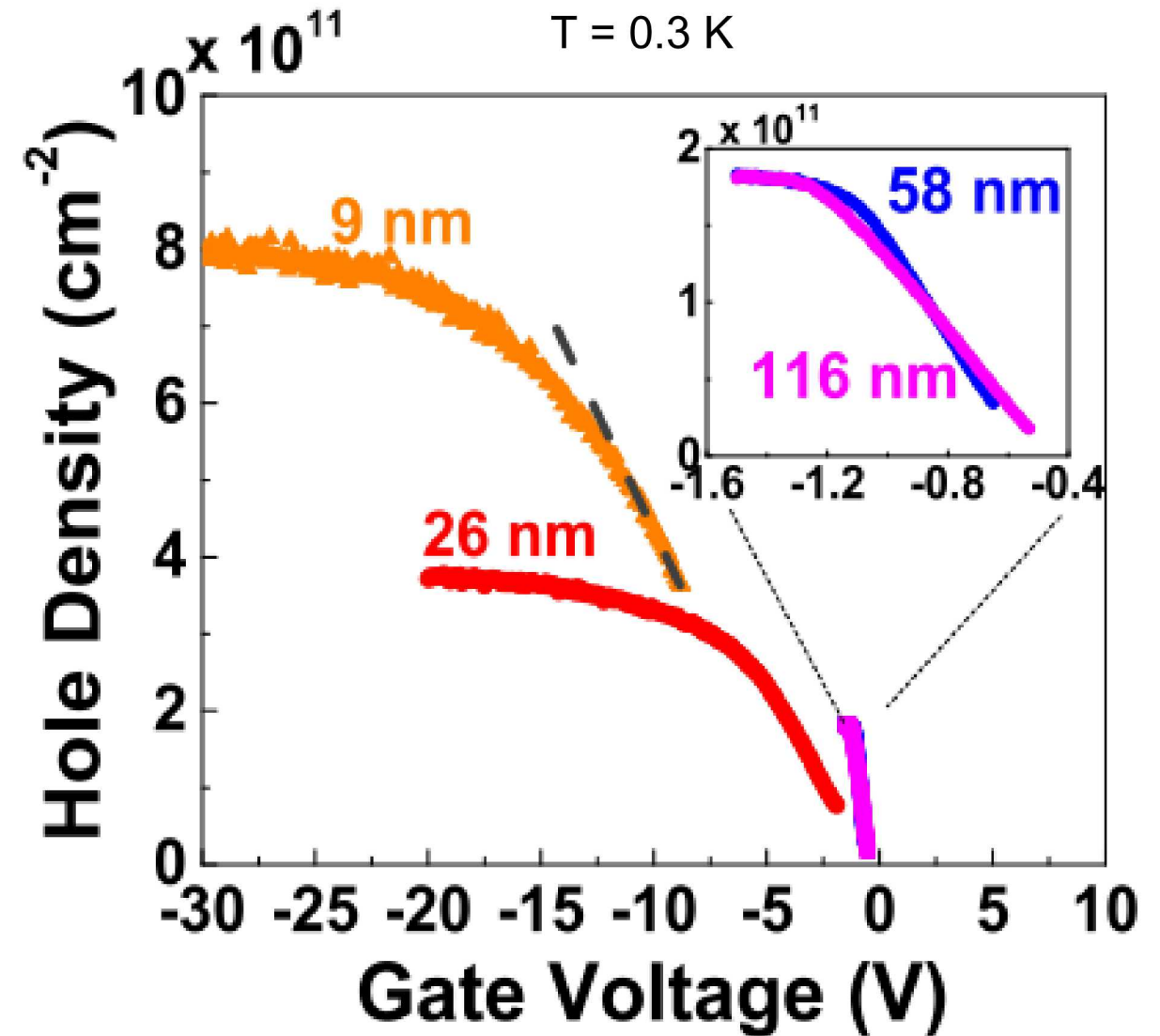
Valance band
2D holes in Ge quantum wells

Undoped Ge/SiGe heterostructure field-effect transistors



Device operation – 2D hole density

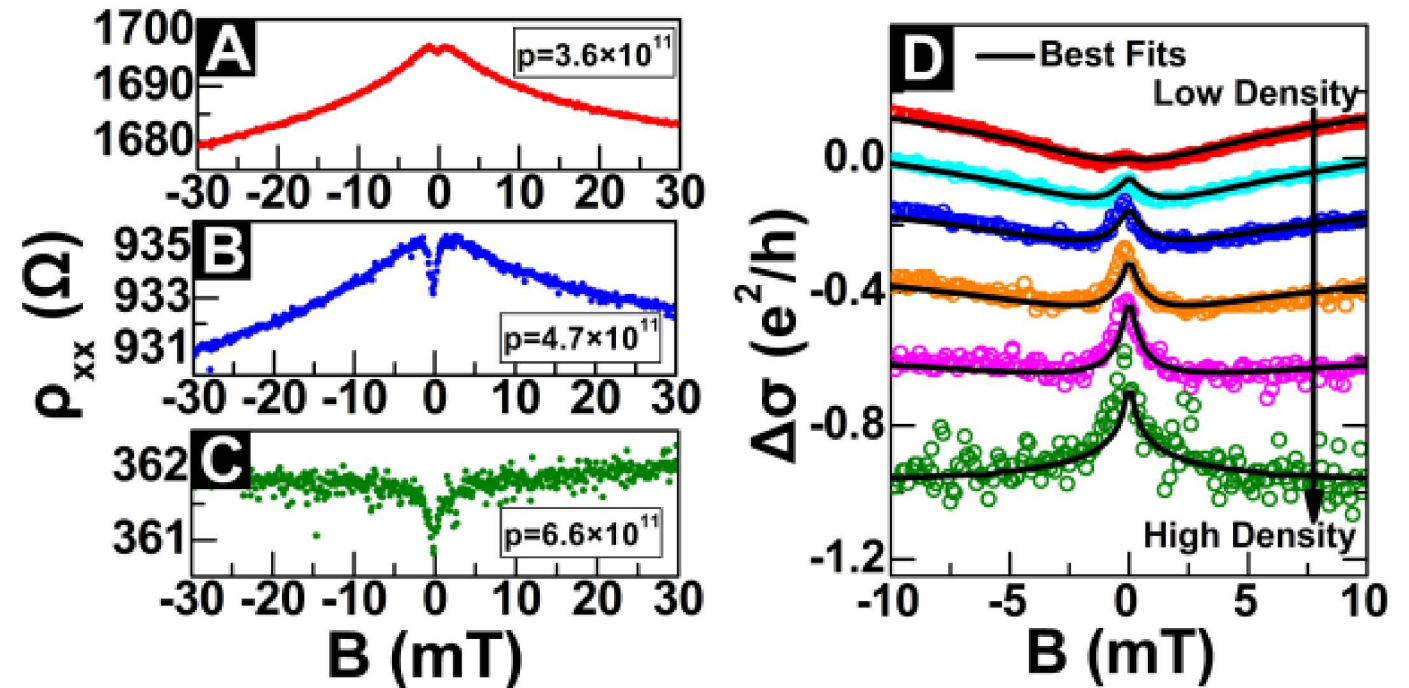
- The 2D hole density saturates.
- Shallow channels
 - High saturation densities, depth dependent.
 - Small slopes (capacitances)
- Deep channels
 - Low saturation densities, depth independent.
 - Large slopes (capacitances)



Physical properties – spin-orbit coupling

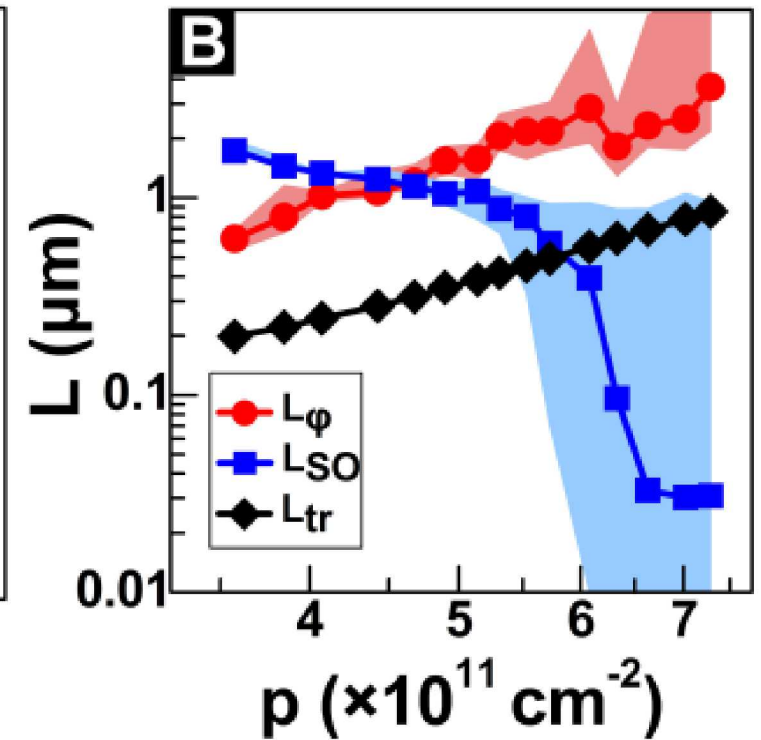
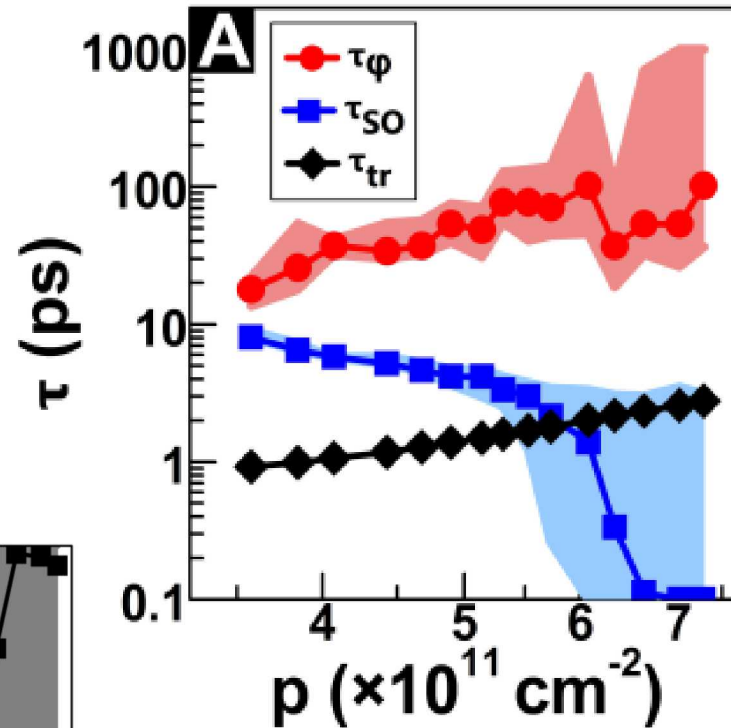
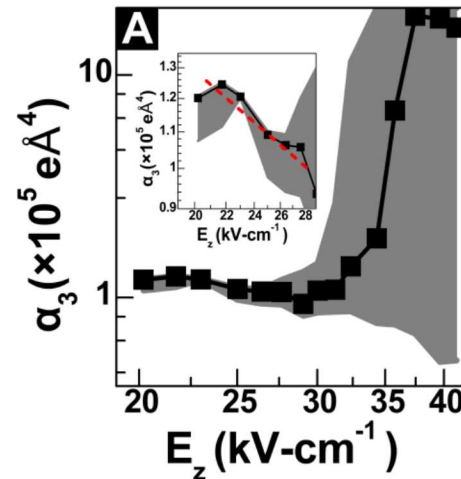
- Low densities
 - Weak localization only
- Intermediate densities
 - Weak anti-localization on top of weak localization only
- High densities
 - Weak anti-localization only

Cubic Rashba: $\sim(\sin 3\theta, -\cos 3\theta)$



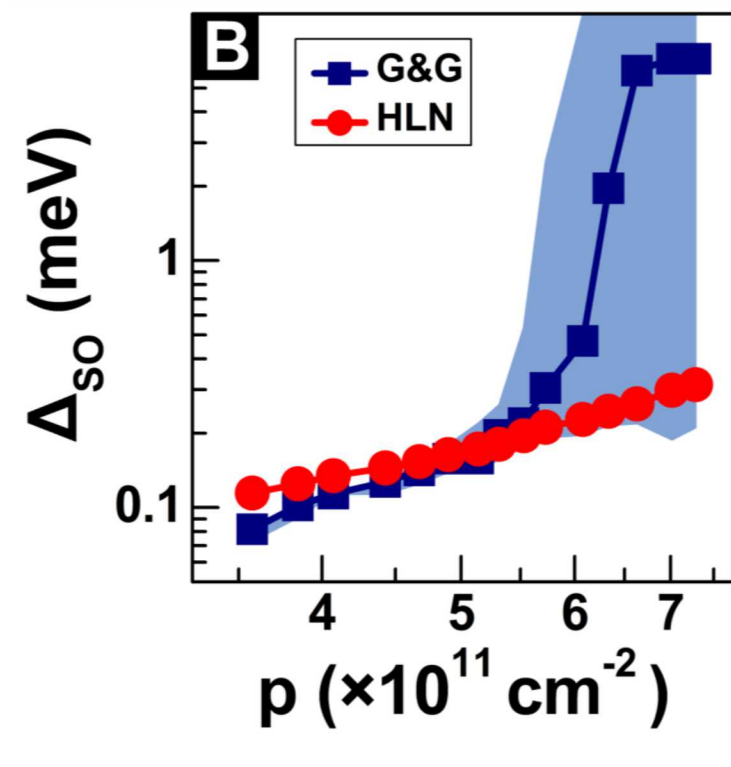
Physical properties – spin-orbit coupling

- Spin-orbit length decreases with density and can be as short as 0.1 μm ($<$ mean free path), while the phase coherence length can be a few microns long ($>$ mean free path)
- This means the hole spin can rotate at a high yet controlled rate, maintain its phase coherence, and suffer no scattering.



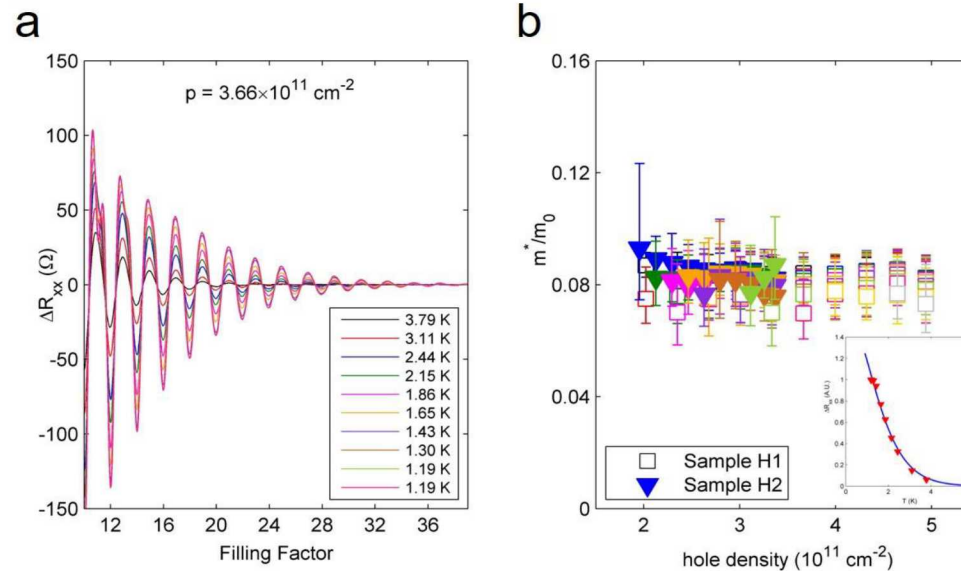
Physical properties – weak antilocalization beyond diffusive regime

- Conventional theories for weak (anti)localization are only valid in the diffusive regime at low magnetic fields.
- Our data lie outside this regime.
- Numerical methods and code for HPC available with paper.
-

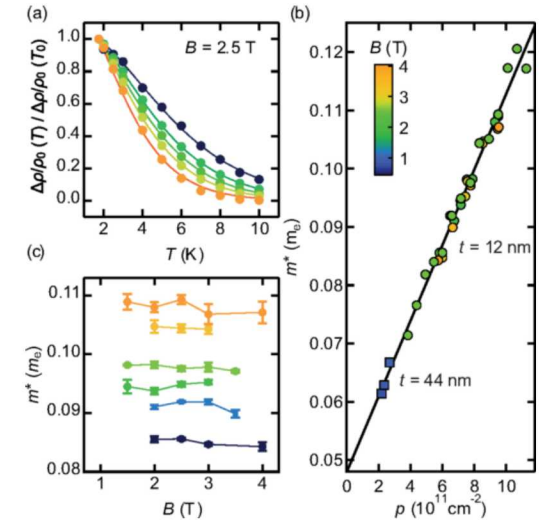


Physical properties – effective mass

- $\sim 0.08 m_0$.
- $\sim$ density independent.
- This mass is small compared to the mass of electrons in Si (0.19), the mass of holes in GaAs (0.2-0.4), and is comparable to the mass of electrons in GaAs (0.07).
- Smaller mass
 - => more extended wave functions
 - => easier gate controls for nanostructures
 - => larger orbital gaps
 - => can use higher T cryostats



Hardy *et al.*, *Nanotechnology*, 30, 215202 (2019)



Lodari *et al.*, *PRB*, 100, 041304(R) (2019)

- Delft results
 - $\sim 0.06-0.12 m_0$.
 - $\sim$ density dependent.

Physical properties – effective mass

Theory

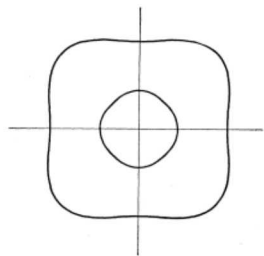
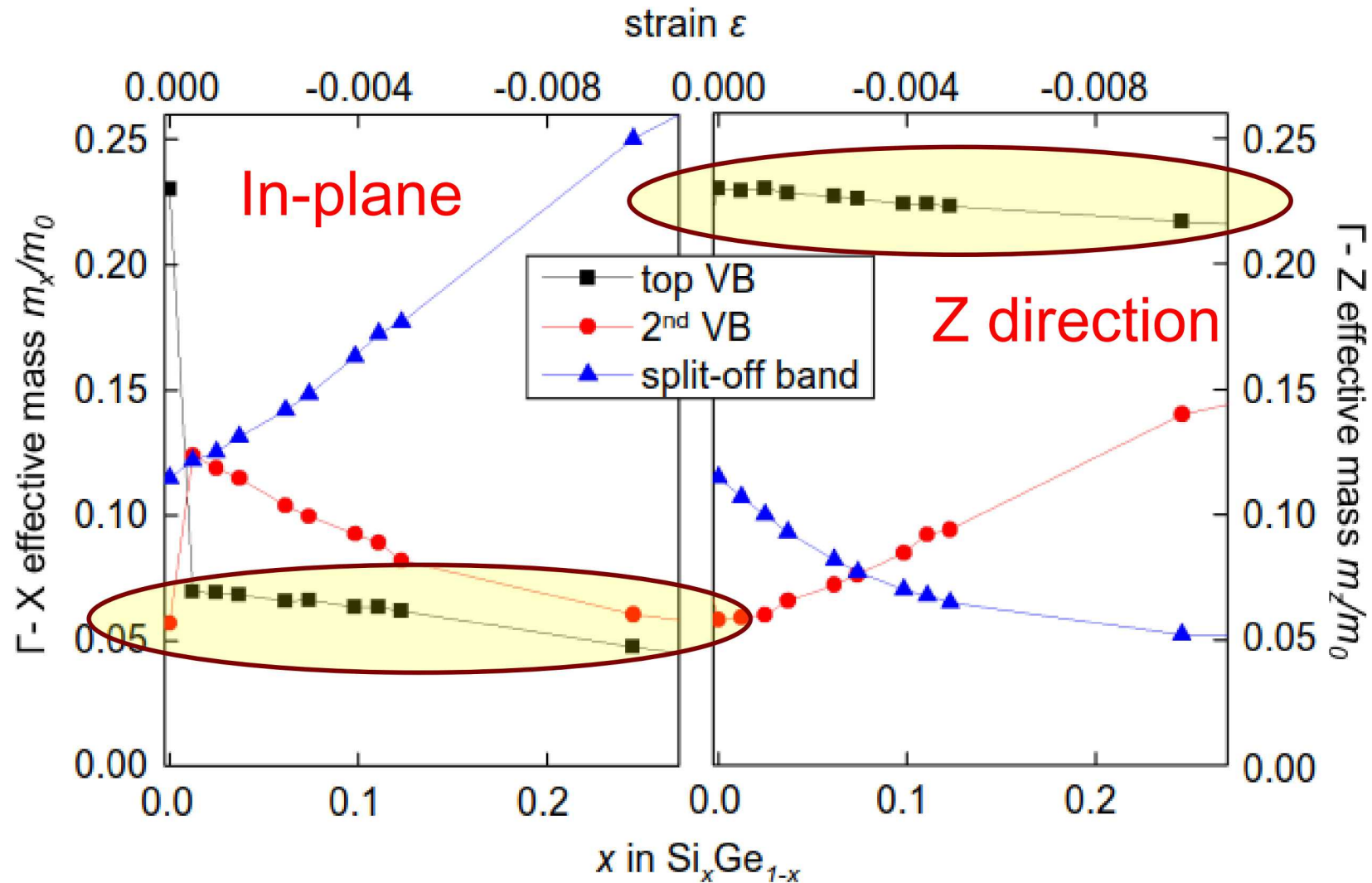
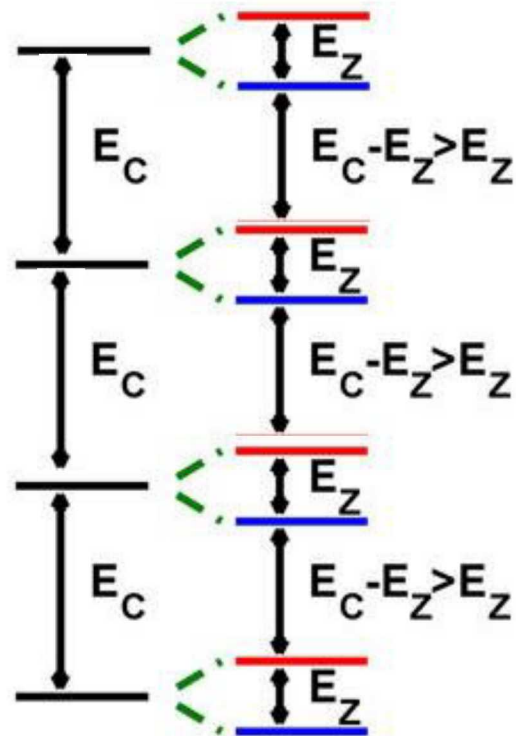


FIG. 8. Figures of constant energy in the (100) plane of k -space for the two fluted energy surfaces which are degenerate at the valence band edge; constants as for germanium.

The quantum Hall effect

In a perpendicular magnetic field, the spectrum of a 2D gas is a series of Landau levels:



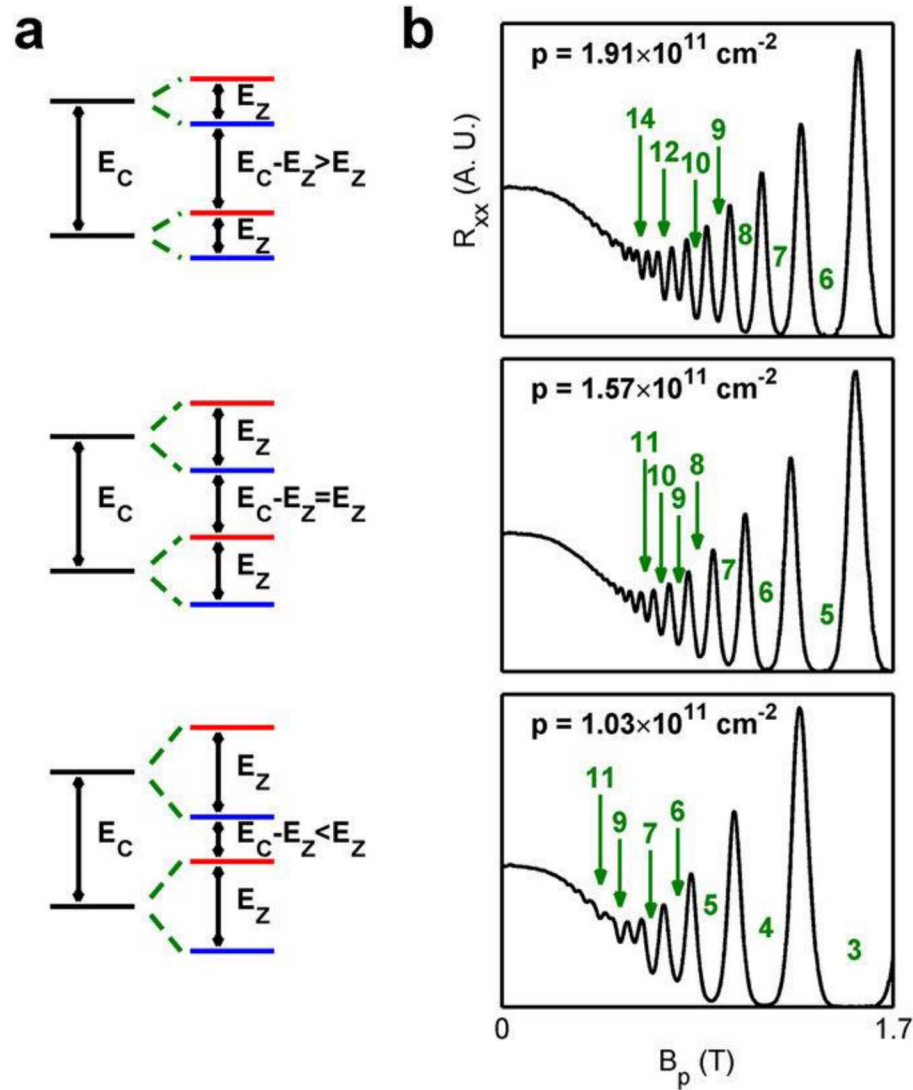
E_C is cyclotron gap: $\hbar e B / (2\pi m^*)$

E_Z is Zeeman gap: $g^* \mu_B$

m^* and g^* are material parameters.

Landau level degeneracy (# electrons / area):
 eB/h

Quantum Hall ferromagnetic transition



In most cases, $E_C \gg E_Z$

$\Rightarrow$ Strong even states, weak odd states

If $E_C \sim 2E_Z$

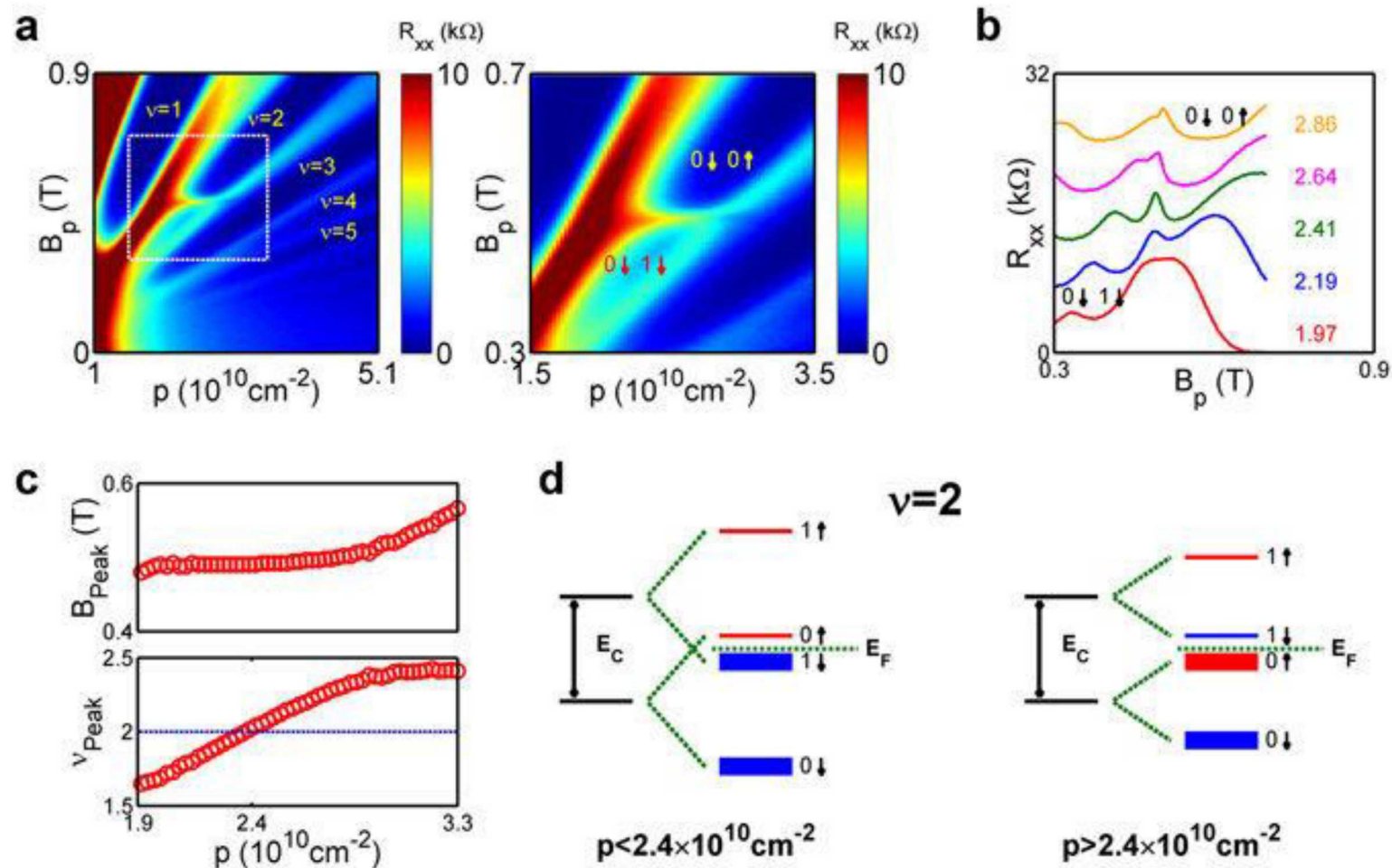
$\Rightarrow$ Strength of even states $\sim$ strength of odd states

If $E_C < 2E_Z$

$\Rightarrow$ Strong odd states, weak even states

E_Z/E_C increases with decreasing density.

Quantum Hall ferromagnetic transition



A spin transition (unpolarized $\leftrightarrow$ polarized) at $\nu=2$ occurs at $p \sim 2.4 \times 10^{10}$ cm $^{-2}$.

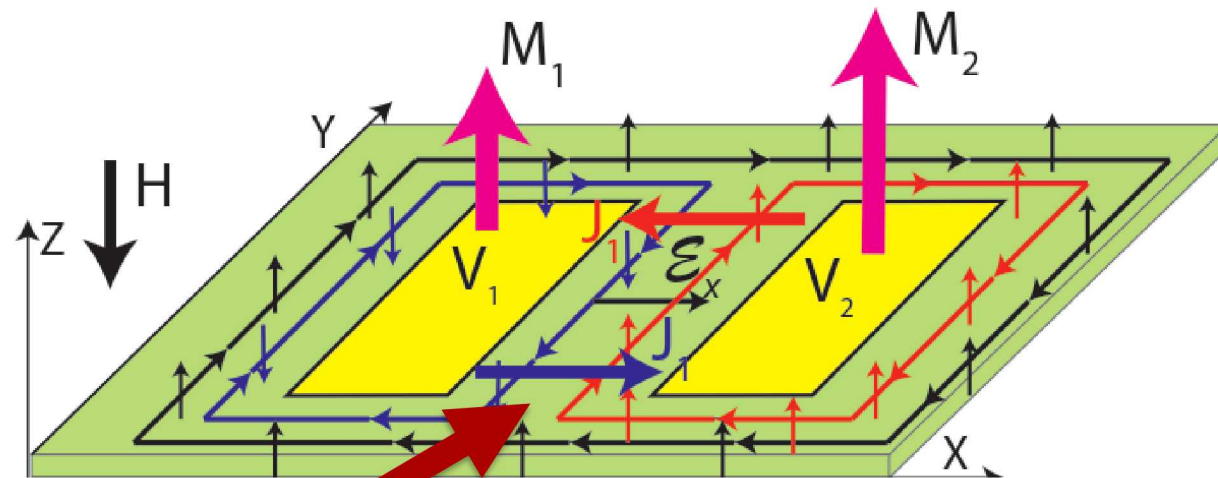
This transition marks the point where $E_c \sim E_z$

$$\frac{\hbar e B}{(2\pi m^*)} = g^* \mu_B$$

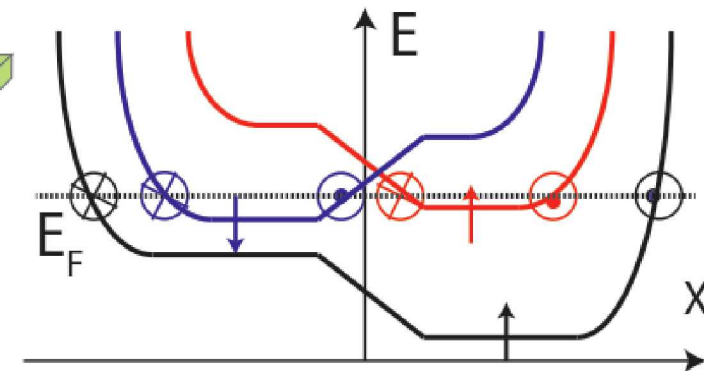
$$\Rightarrow m^* g^* = 2$$

Quantum Hall ferromagnetic transition

- Local gating to create counter-propagating edge states with opposite spins
- Depending on the strength of the effective spin-orbit coupling, topological superconductivity and Majorana modes can emerge.
- $B_c \sim 0.5\text{T}$ -> very friendly for superconductivity.



Helical domain wall



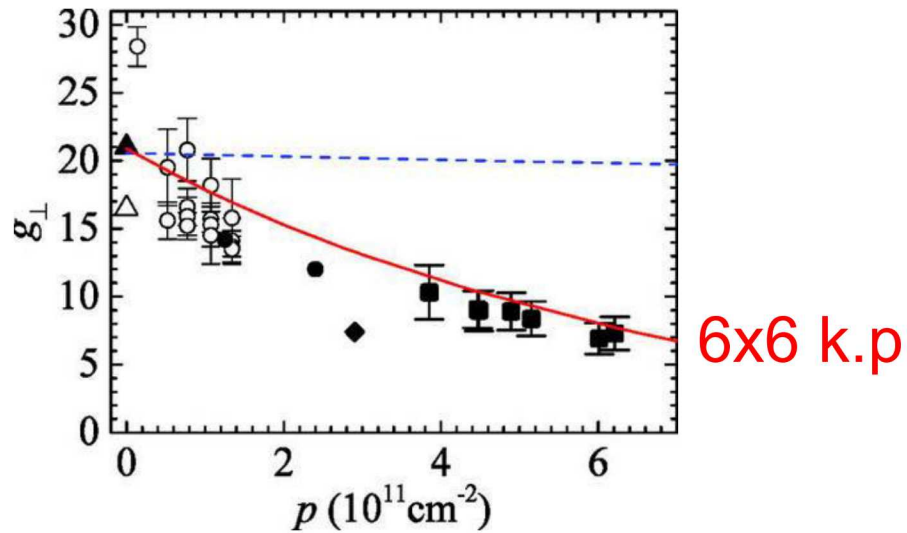
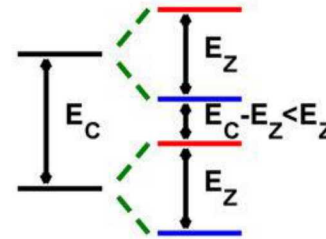
“Impurity-generated non-Abelions”
Simion Phys. Rev. B 97, 245107 (2018)

Physical properties – g factor

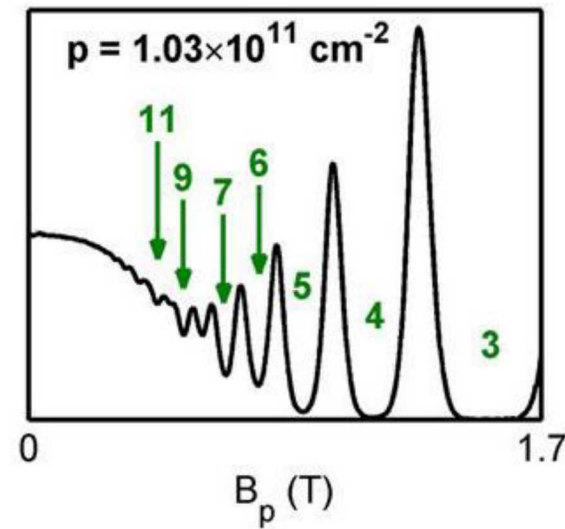
- $\sim 5 - 30$
- $\sim$ density dependent.
- The g factor is large compared to the g factor of electrons in Si (~ 2) and the g factor of electrons in GaAs (~ 0.44).

$$P < 1.6 \times 10^{11} \text{ cm}^{-2}$$

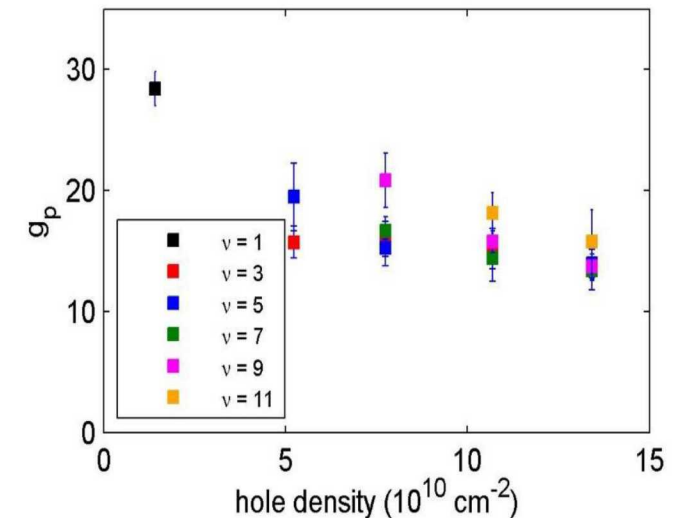
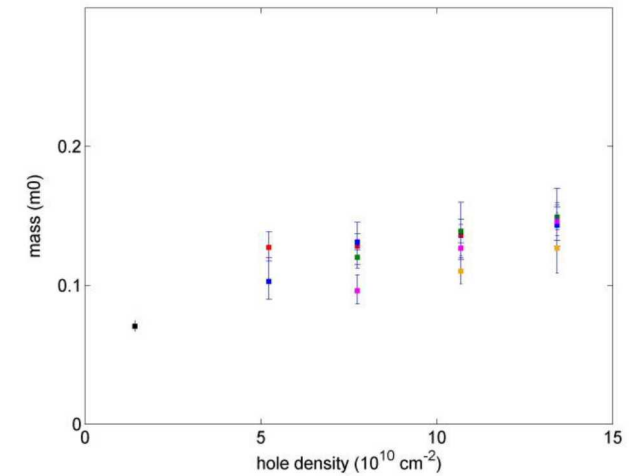
$$0.5 E_c < E_z < E_c$$



Drichko J. Appl. Phys. 123, 165703 (2018)

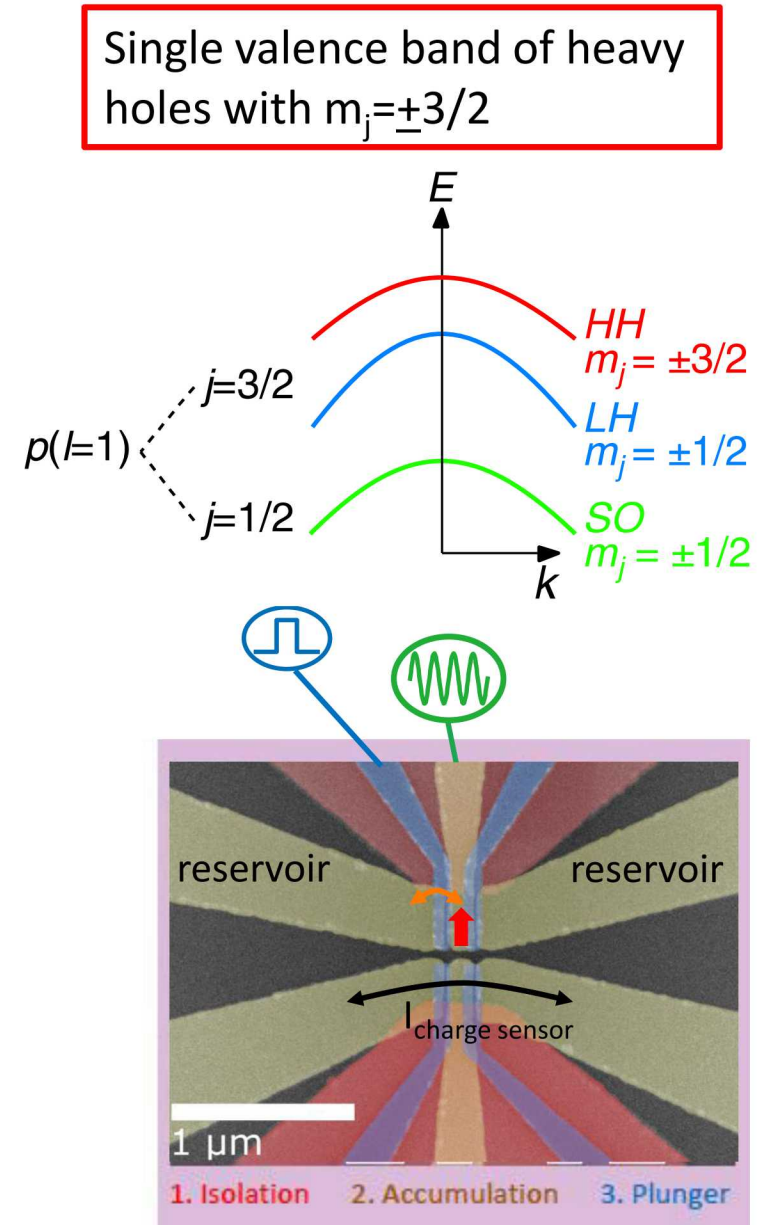


Lu, Appl. Phys. Lett. 111, 102108 (2017)

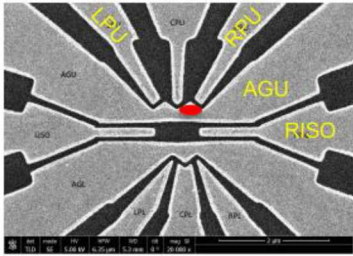


Approach

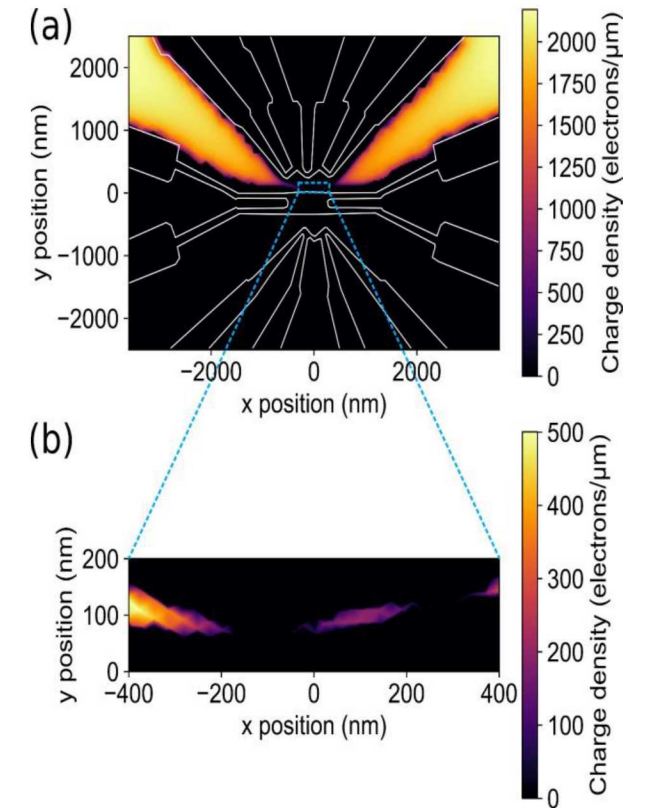
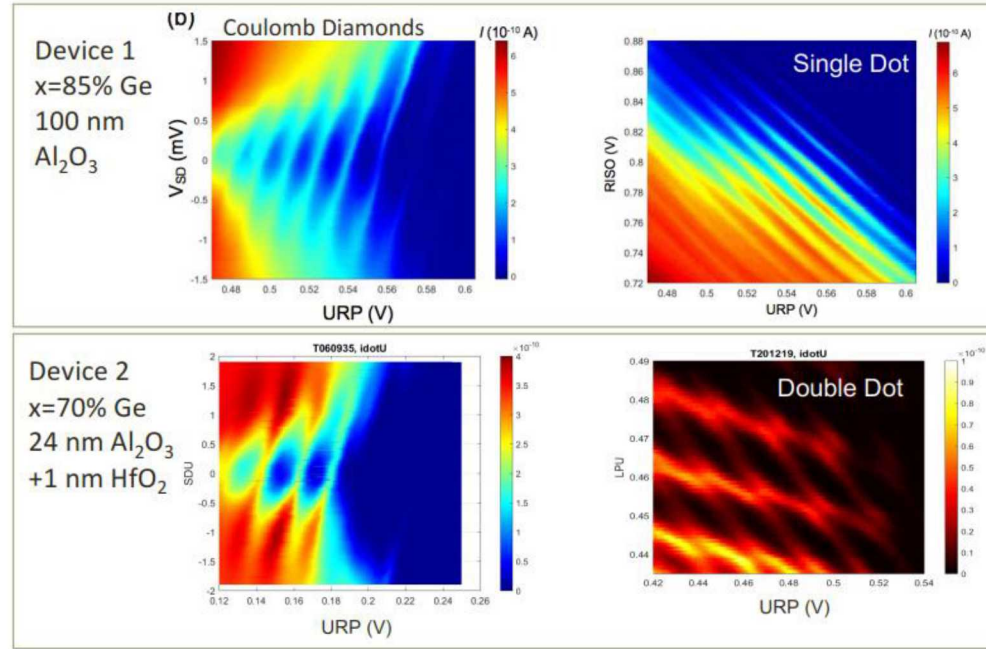
- Single hole confined to lateral quantum dot in Ge heterostructure (small mass)
- Spin qubit states: $m_j = \pm 3/2$ (no valley splitting)
- EDSR qubit control through microwaves applied to gates (strong SOC)
- Qubit readout and initialization through energy selective tunneling to reservoir
- Occupancy detected through nearby charge sensor
- P-orbital nature -> suppressed hyperfine.
- Enrichment is possible.



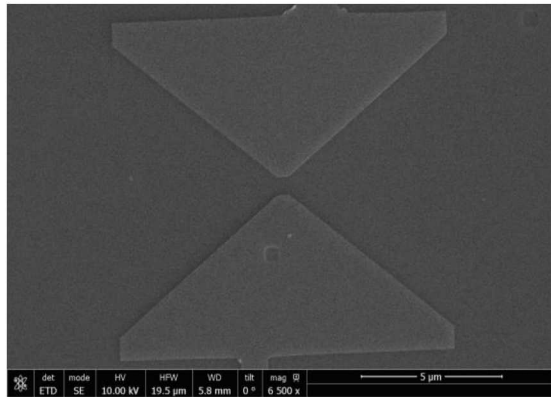
Spin qubits in Ge/SiGe



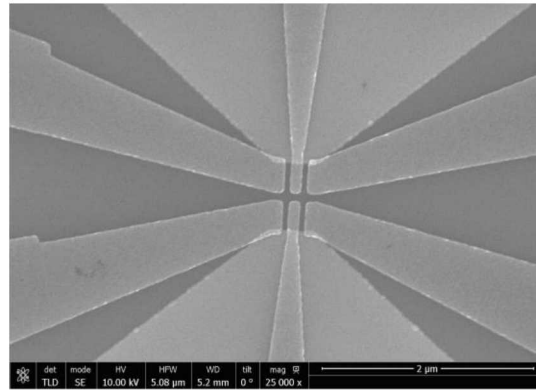
Single Layer Devices
can be tuned to low-
hole regime in
transport



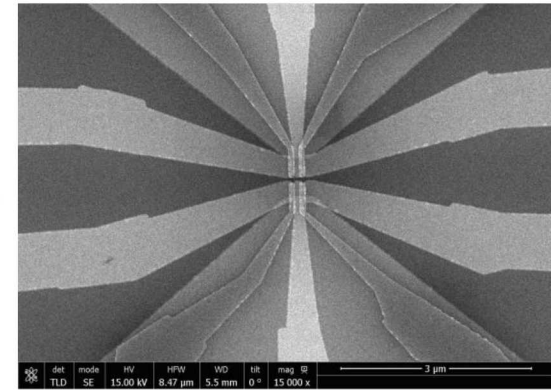
3-layer device fabrication



1: isolation gates

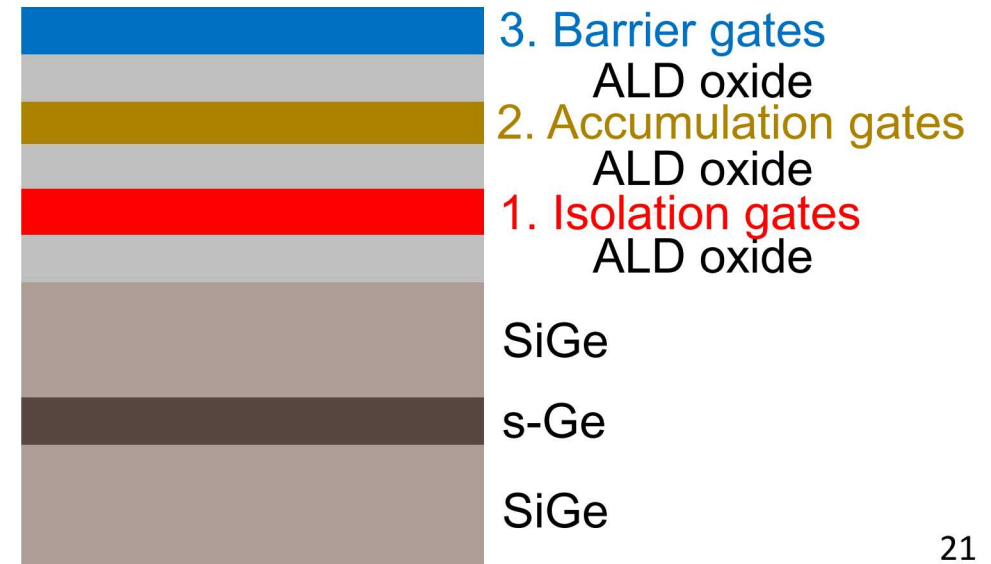
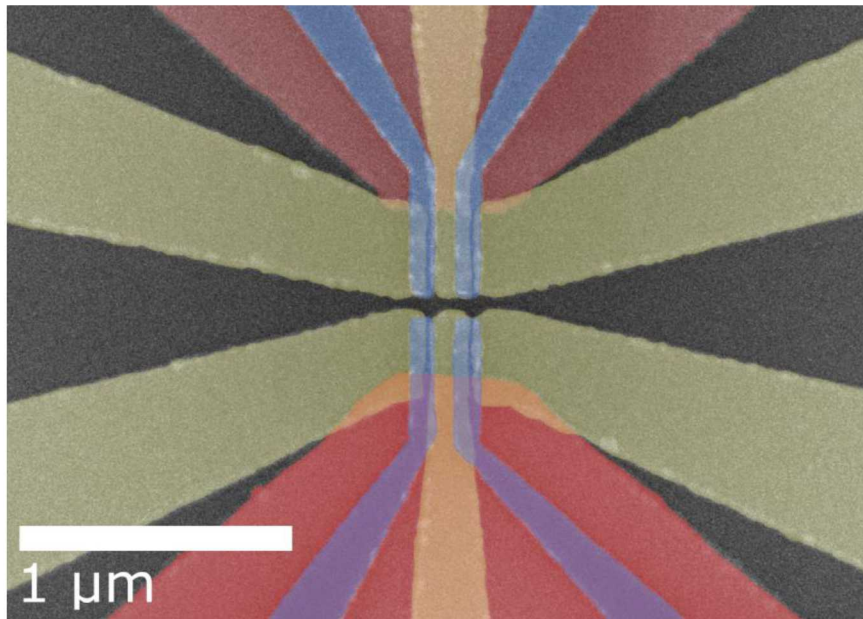


2: accumulation gates



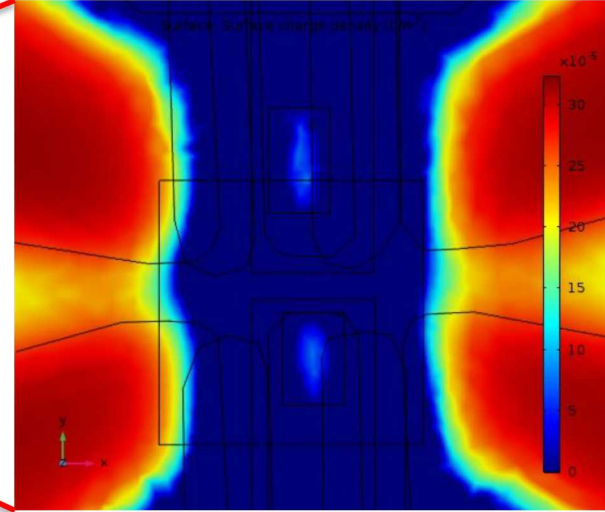
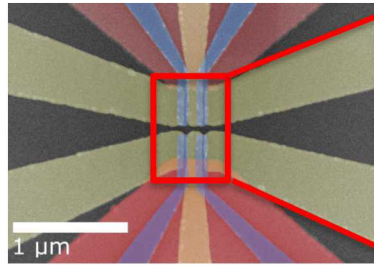
3: barrier gates

E-beam lithography
ALD Al_2O_3 oxide
Ti/Pt gates

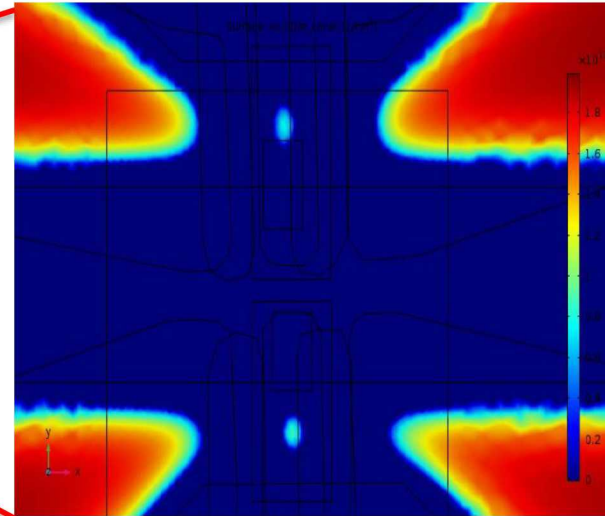
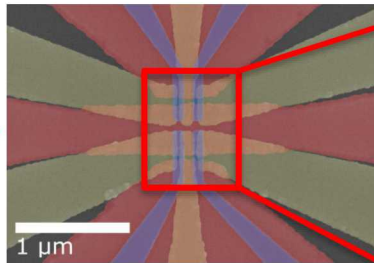


3-layer device modeling: surface charge density

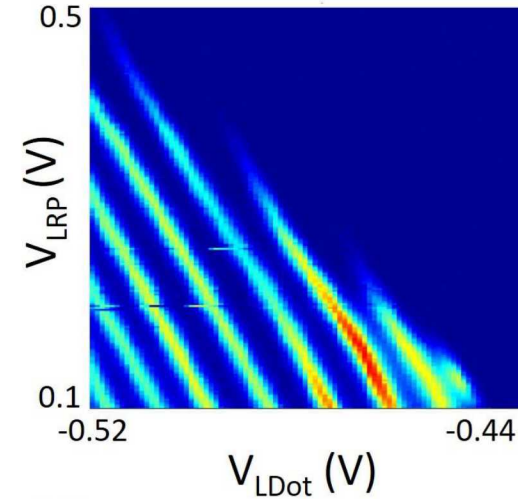
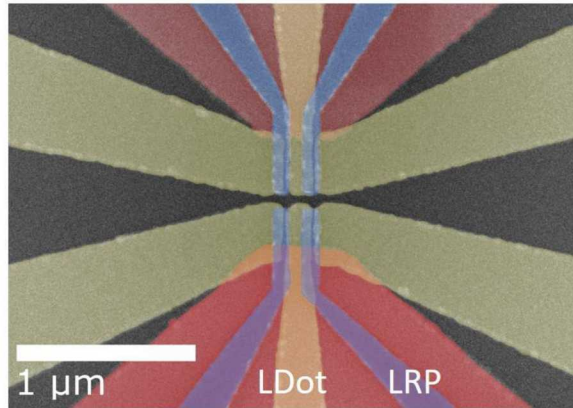
Without horizontal isolation gate:
reservoirs merge



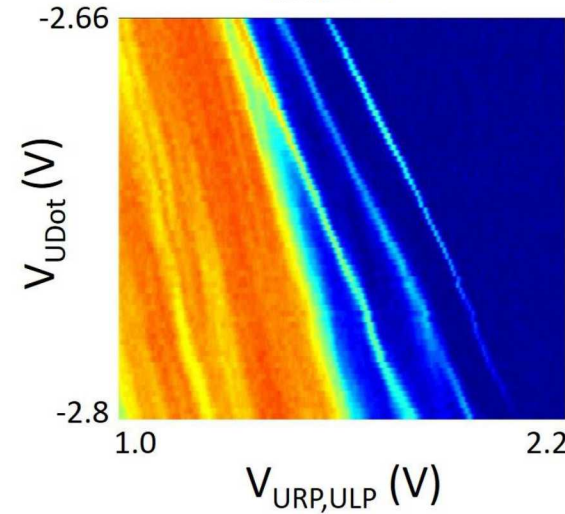
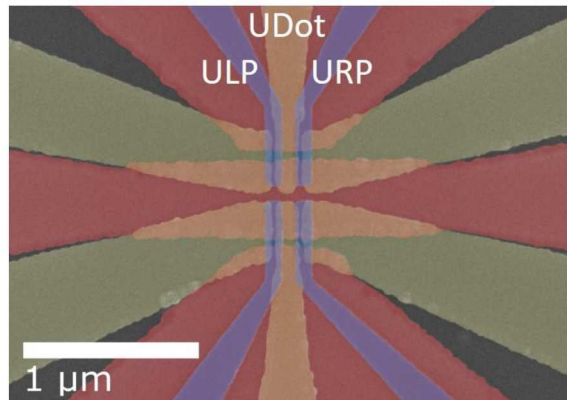
Add horizontal isolation gate:  reservoirs separated



Improved Coulomb blockade by adding center isolation gate



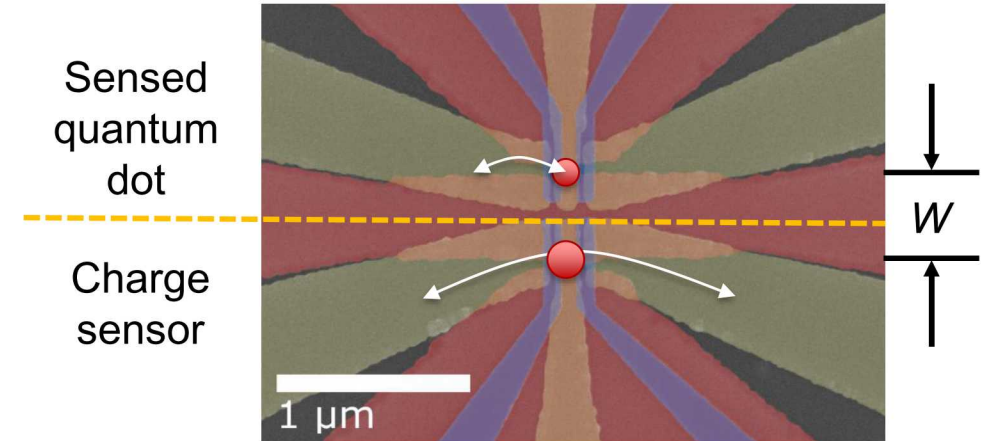
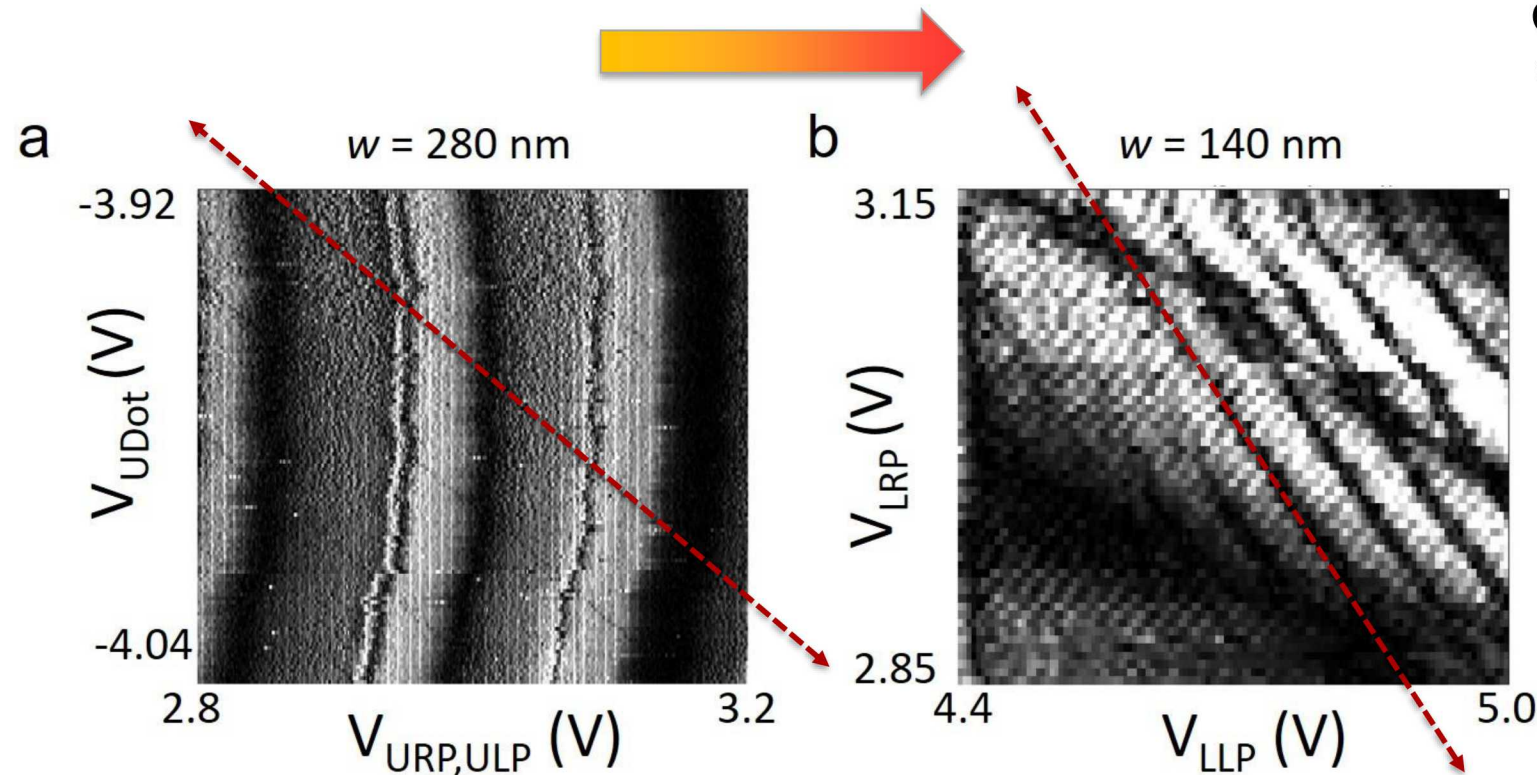
Without horizontal
isolation gate



With horizontal
isolation gate

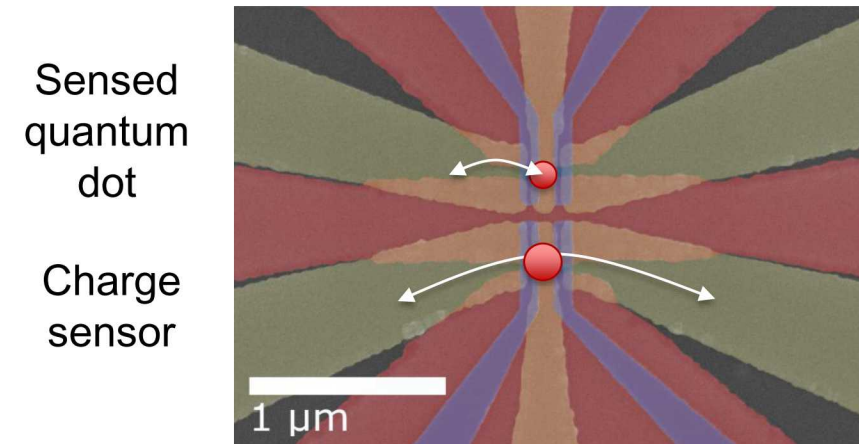
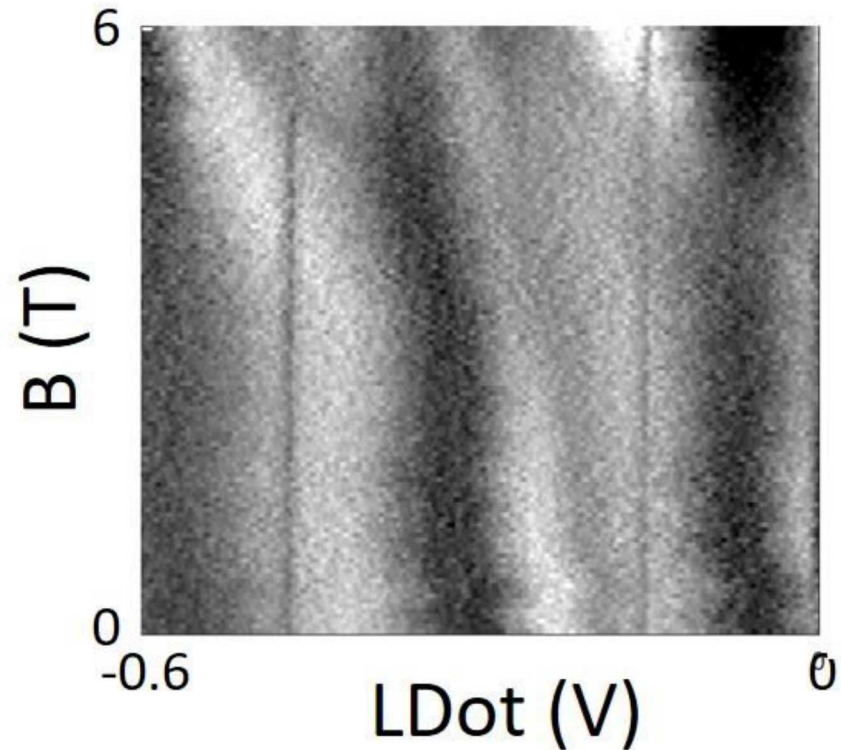
Charge sensing: Comparison of isolation gate widths

150% improvement in charge sensing signal when w is reduced by 2x



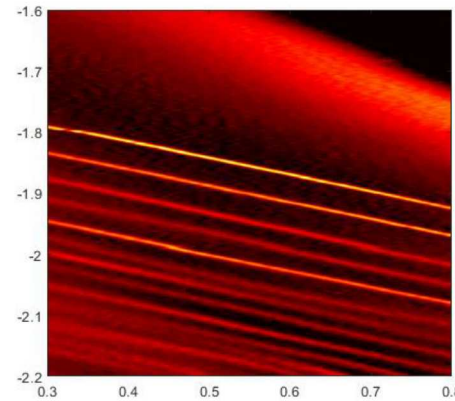
Parallel diagonal lines are sensed charge transitions

In-plane B field



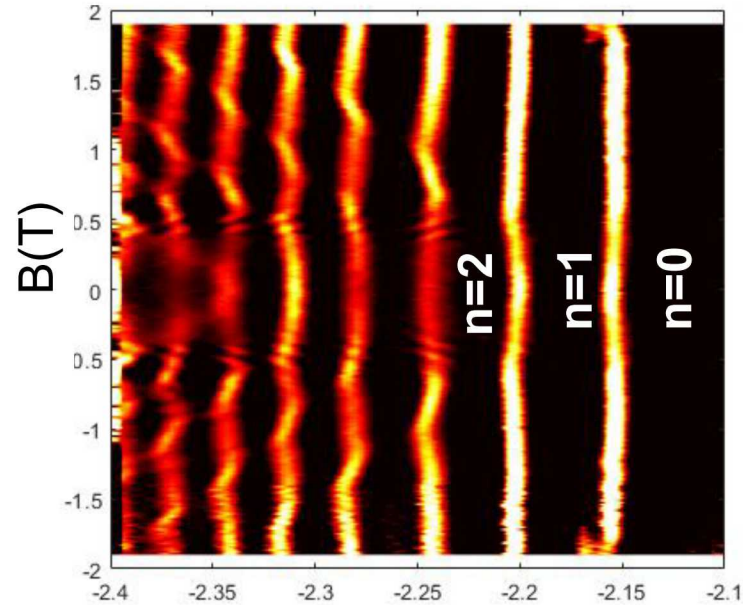
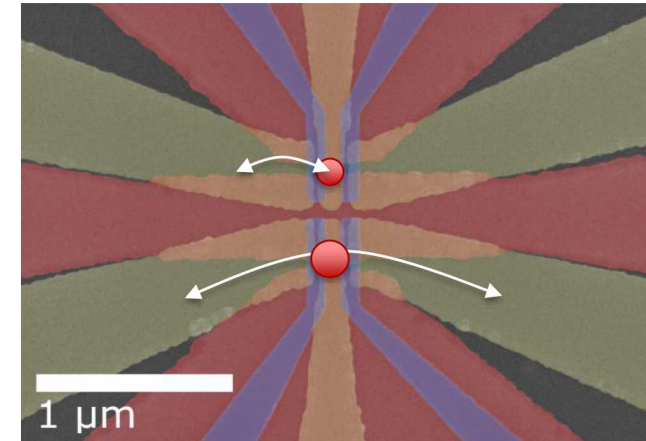
- Charge sensed lines appear insensitive to in-plane B-field
 - $g_{\text{in-plane}}$ is very small (estimate < 0.5)

Perpendicular B field

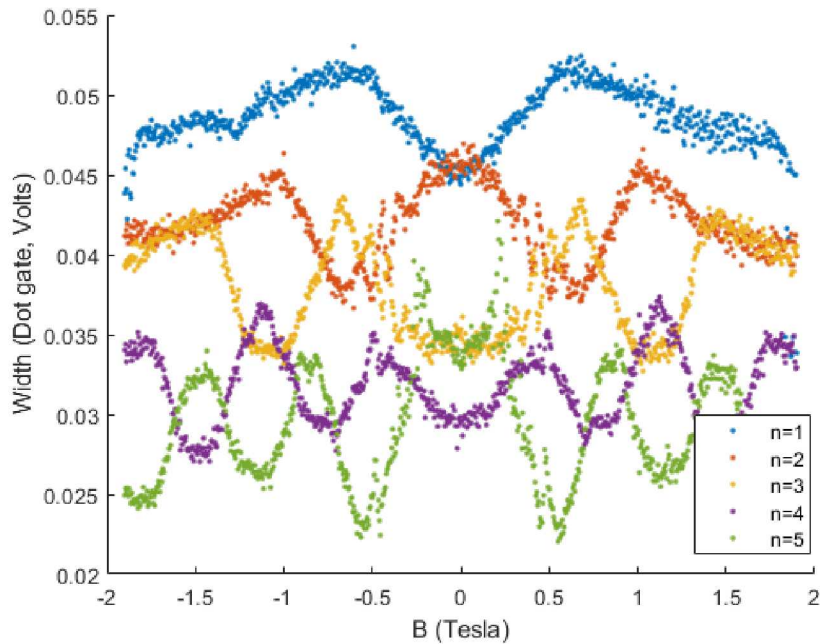


Sensed
quantum
dot

Charge
sensor



Perpendicular B field



- $g_{\text{perpendicular}}$: ~ 15 (based on $N=1$)
 ~ 18 (based on $N=2$)
- Consistent with estimates from transport measurement on 2DHGs.

Modeling SOC in Ge QDs

Goal: Develop device-level models to guide & interpret experiments

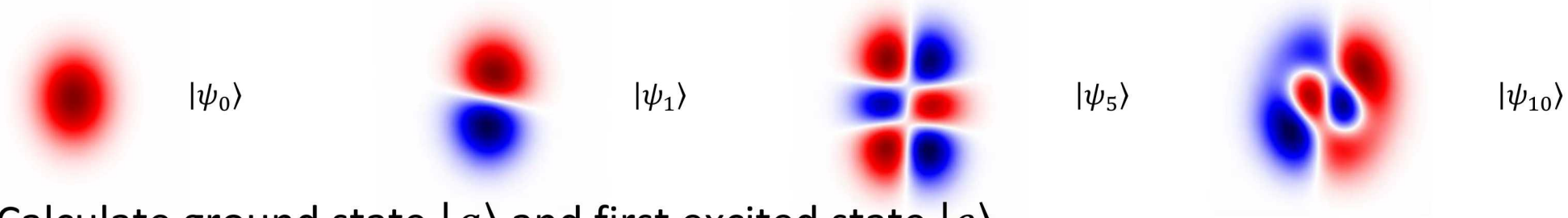
- Challenges
 - The form of the SOC Hamiltonian remains ambiguous in Ge quantum dots
 - No existing software for device-level modeling with SOC
- Our trajectory
 - Use Rabi oscillation frequency as a connection to experiment
 - Calibrate expectations with a simple model
 - Enhance device-level modeling tools to accommodate SOC physics
 - Use tools to explore various SOC models (e.g., single-band vs. multi-band)
 - Integrate more physical details into our theory

Luttinger Hamiltonian for SOC modeling

- Include externally applied potential and (spin-3/2) Zeeman term

$$\hat{H}_Z = \kappa\mu_B \begin{bmatrix} 3B_z & \sqrt{3}B_- & 0 & 0 \\ \sqrt{3}B_+ & B_z & 2B_- & 0 \\ 0 & 2B_+ & -B_z & \sqrt{3}B_- \\ 0 & 0 & \sqrt{3}B_+ & -3B_z \end{bmatrix}; \quad B_{\pm} = B_x \pm iB_y$$

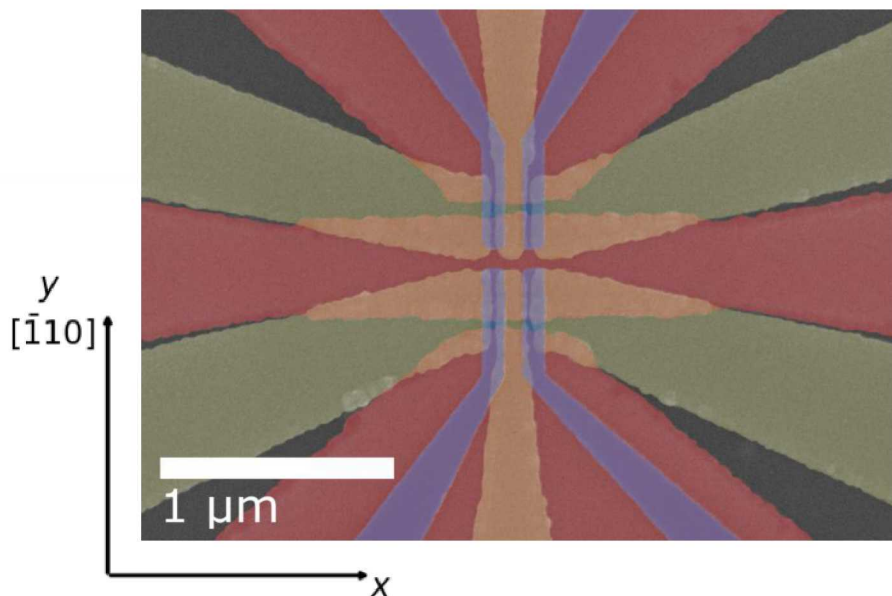
- Select a basis in which to expand $\hat{H}_{LK}$, e.g. eigenstates of the potential



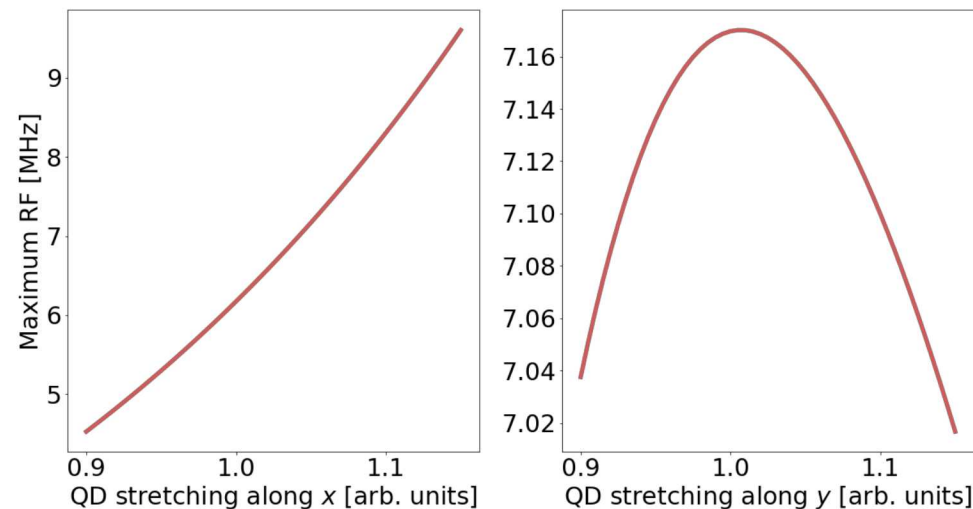
- Calculate ground state $|g\rangle$ and first excited state $|e\rangle$
- Rabi frequency:

$$f_R \propto \langle e | \hat{D} | g \rangle$$

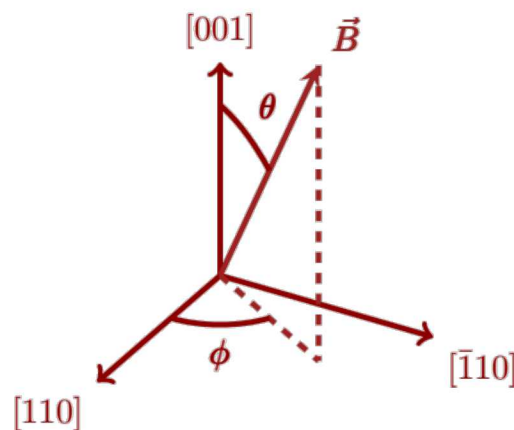
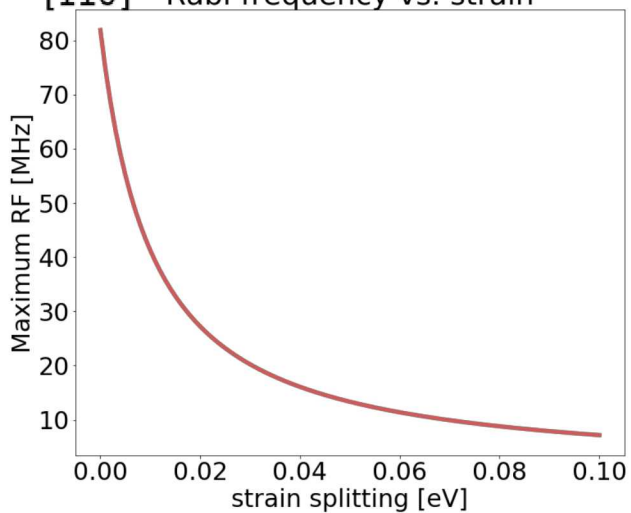
Results!



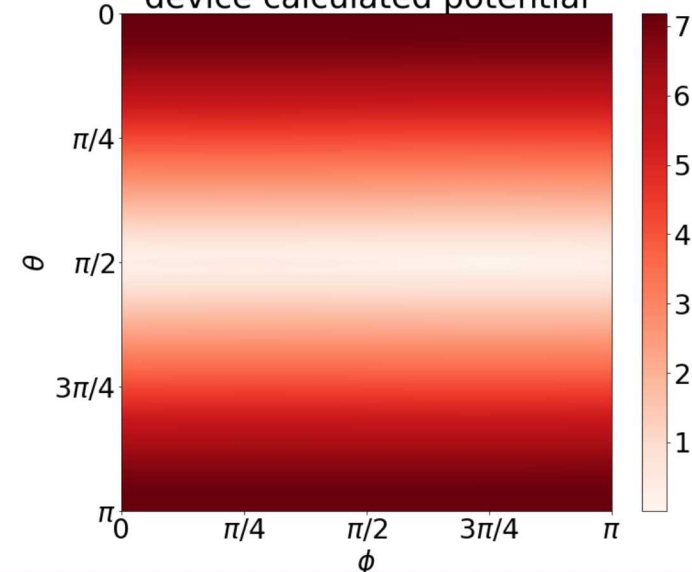
Rabi frequency vs. dot stretching



Rabi frequency vs. strain



Rabi frequencies [MHz] for device calculated potential



Summary

- Induced 2D holes in Ge/SiGe heterostructure field effect transistors demonstrated.
- Physical properties (mass, g factor, spin-orbit coupling strength) characterized.
- Gate controlled quantum Hall ferromagnetic transition observed at low densities. Potential platform for topological superconductivity.
- Development of spin qubits in Ge/SiGe