



LDRD

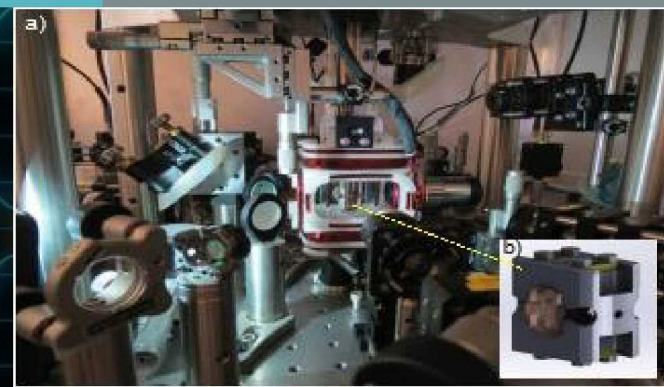
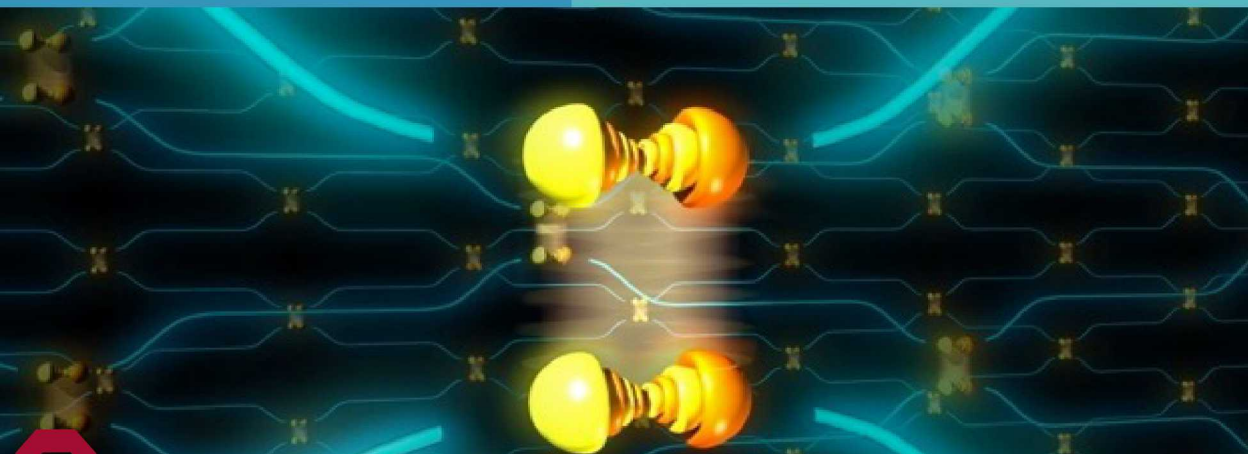
Laboratory Directed Research and Development

Grant Biedermann
University of Oklahoma

Jongmin Lee: Sandia National Laboratories

Team: L.P. Parazzoli, B. Little, M. Chow, J. Lee, P. Schwindt,
C. Brif, B. Ruzic, W. Kendal, A. Orozco

Growing complexity in quantum metrology with neutral atom spins



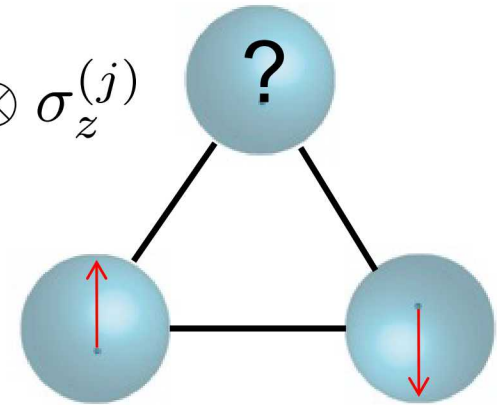
Sandia National Laboratories is a multimission laboratory managed and operated by National Technology & Engineering Solutions of Sandia, LLC, a wholly owned subsidiary of Honeywell International Inc., for the U.S. Department of Energy's National Nuclear Security Administration under contract DE-NA0003525.



Future applications of *Neutral Atom* quantum information systems

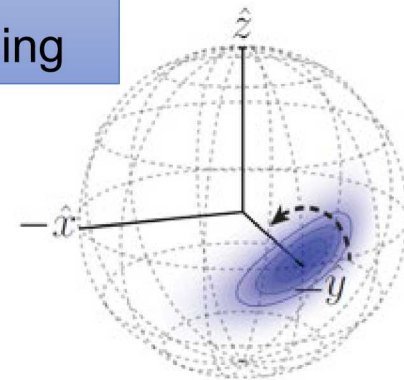
Quantum simulation

$$H = \sum_{i=1}^N h_x^{(i)} \sigma_x^{(i)} + \dots + \sum_{i \neq j}^N J_{i,j} \sigma_z^{(i)} \otimes \sigma_z^{(j)}$$



Frustrated magnetism

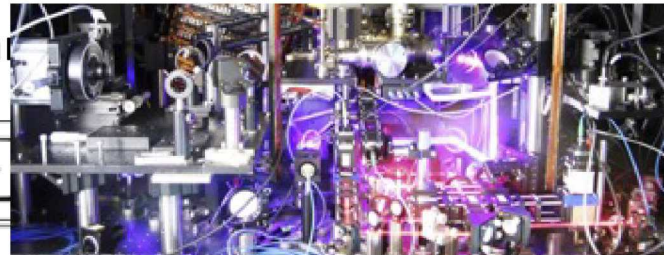
Large-scale/rapid
entanglement for sensing



Sr lattice clock, Ye lab JILA

Quantum Fourier transform

$$QFT : |x\rangle \mapsto \frac{1}{\sqrt{N}} \sum_{k=0}^{N-1} e^{2\pi i x k / N} |k\rangle$$



2015 Guinness book

AI phase shift: a simple perspective

Use atom to stroboscopically measure lateral position (optical phase) at three equal-spaced points in time

$$\varphi_i = \vec{k}_{eff} \cdot \vec{x}_i$$

calculate curvature via finite difference method

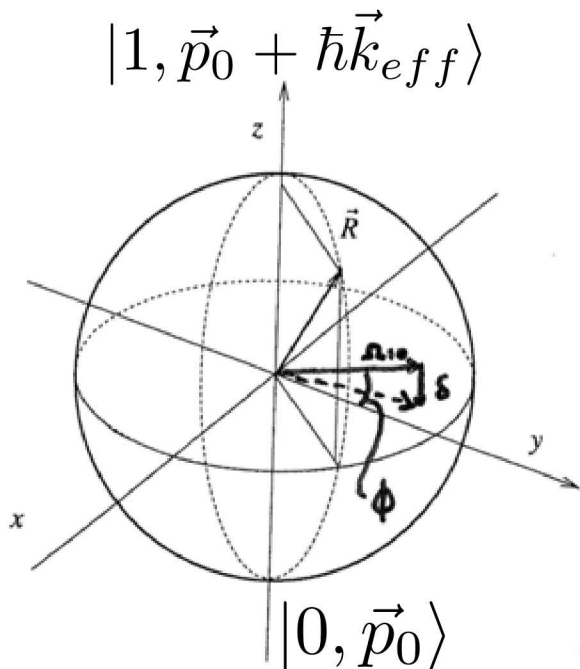
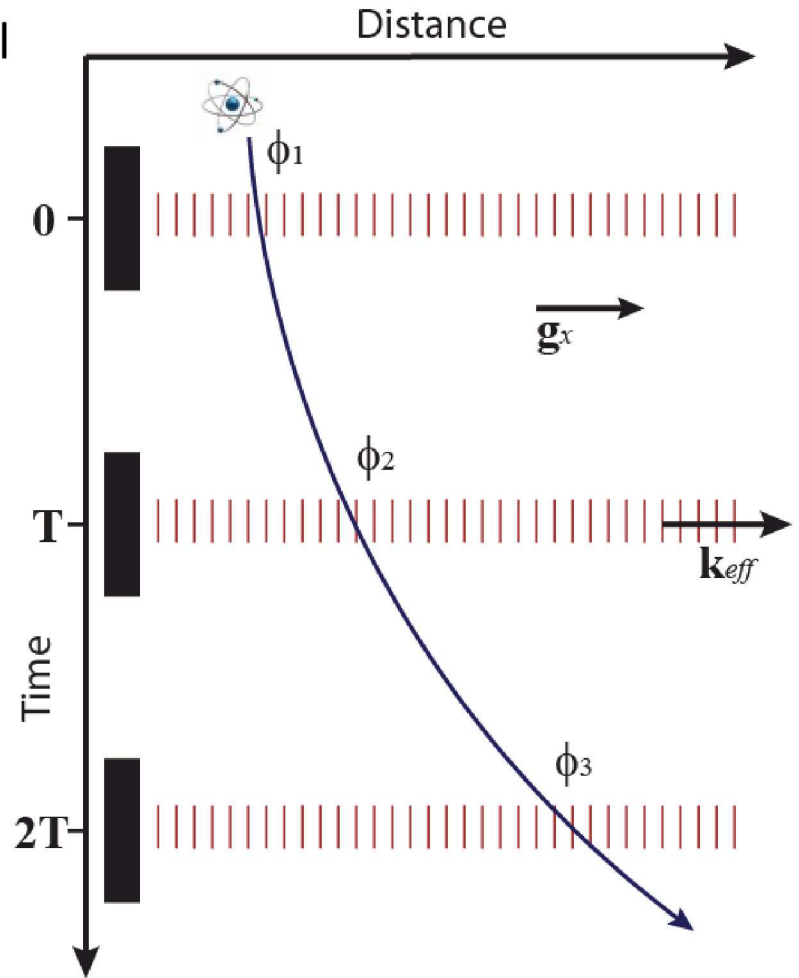
$$\Delta\varphi = \varphi_1 - 2\varphi_2 + \varphi_3$$

...and then

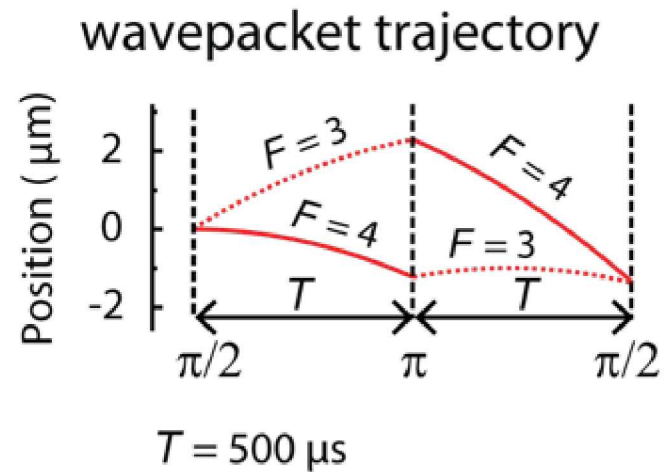
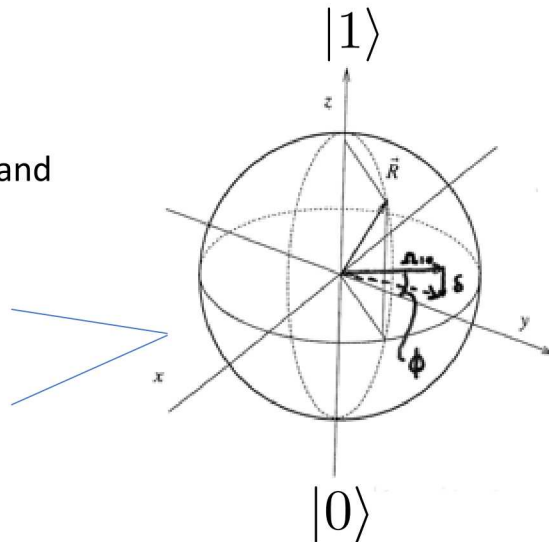
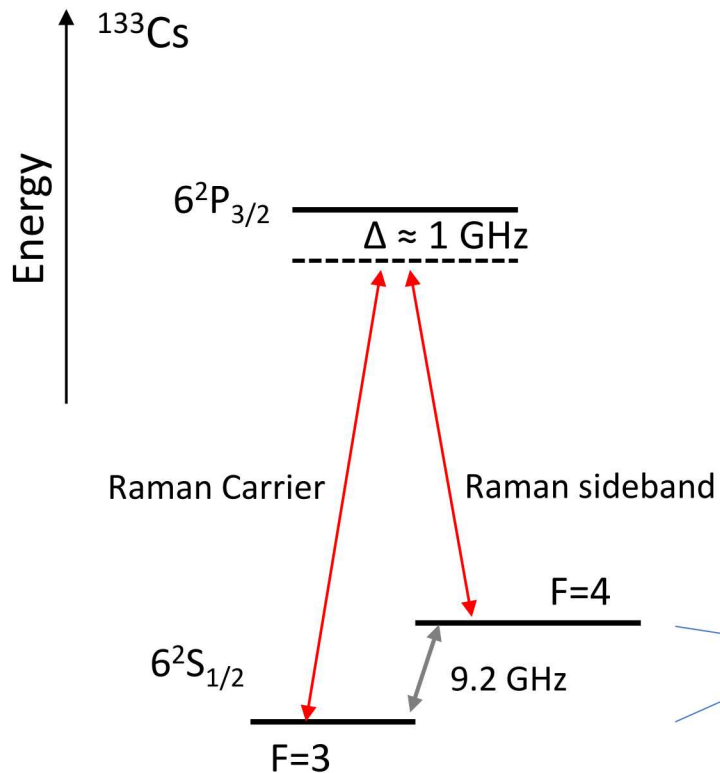
$$g_x = \frac{\Delta\varphi}{k_{eff} T^2}$$

$$P_{|1\rangle} = \frac{1}{2} (1 - \cos(\Delta\phi + \phi_0))$$

... $S_{1/2}$ states behave as a pristine oscillator with phase memory



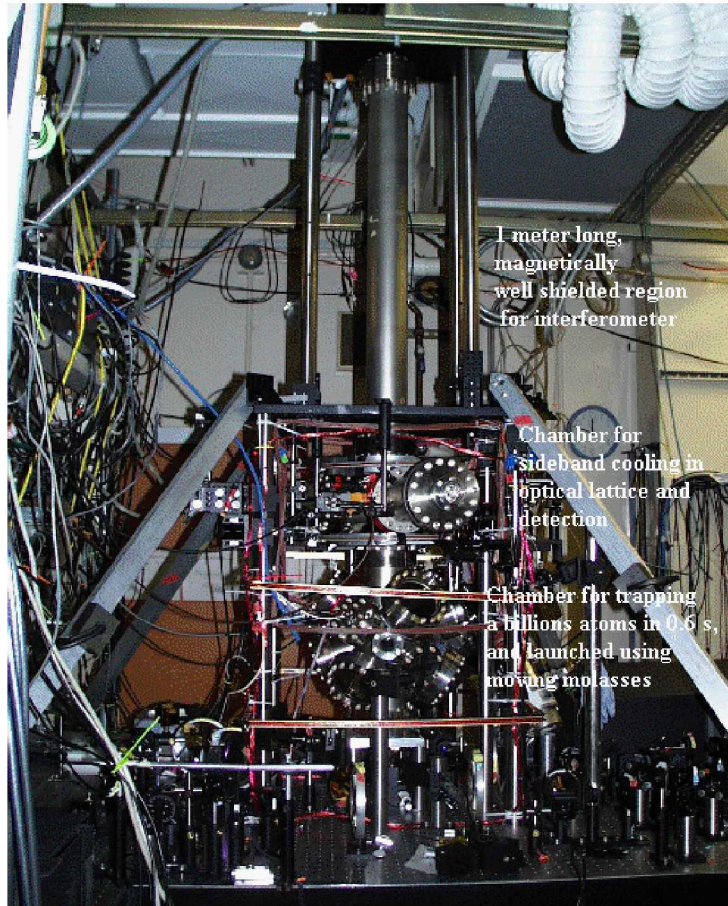
Photon recoil beam splitter



Store "spin" information in ground states



Light-pulse atom interferometers



Stanford atom interferometer gravimeter
Steven Chu, 1990's

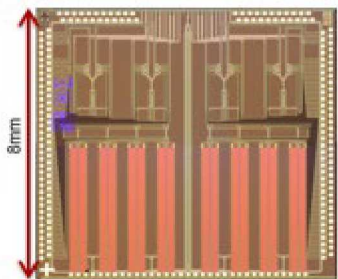


NIST F-1 atomic fountain clock
(credit: NIST website)

Superb gyroscopes and gravity gradiometers
demonstrated as well

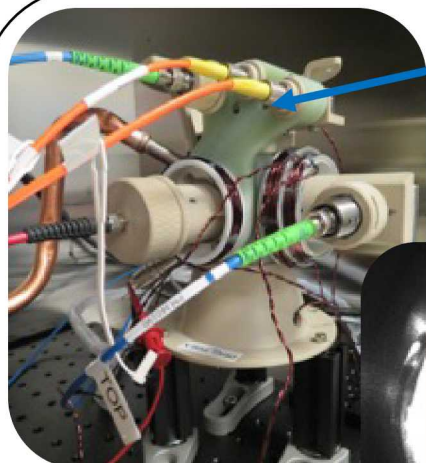
SNL Atom Interferometer Development

Integrated Photonics



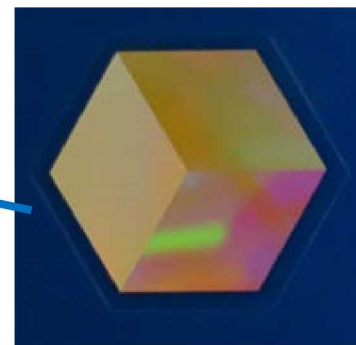
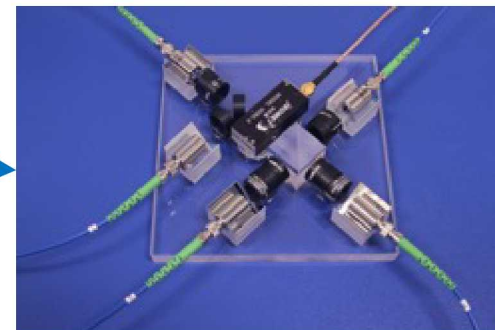
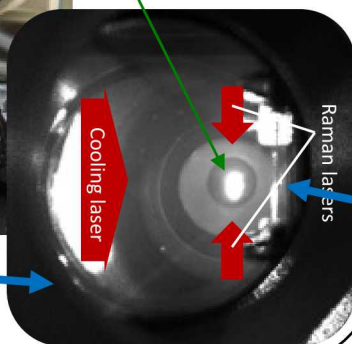
Custom vacuum chamber

Cold Atom Sensor Head with Grating-mirror
MOT and Alignment-Free Optical Package



Alignment-free
optical package

GMOT

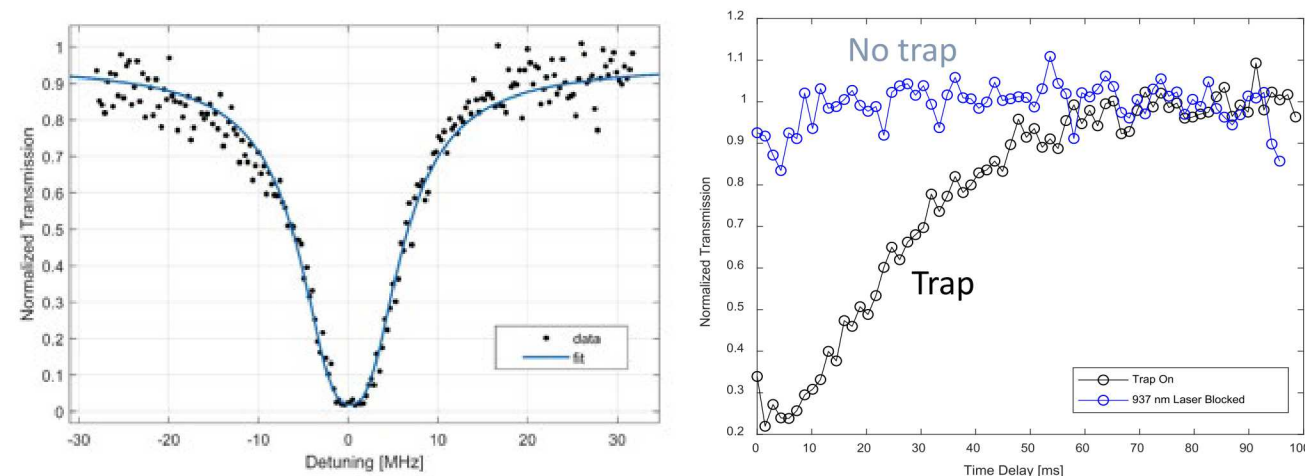
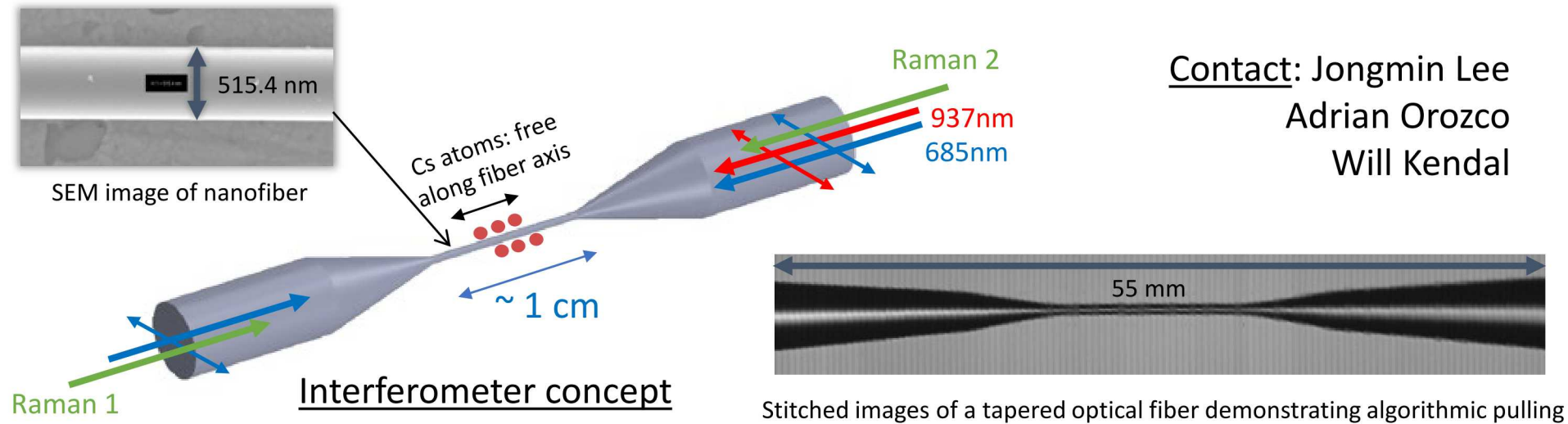


Sandia-fabricated
Grating-mirror chip

Contact: Jongmin Lee
Peter Schwindt

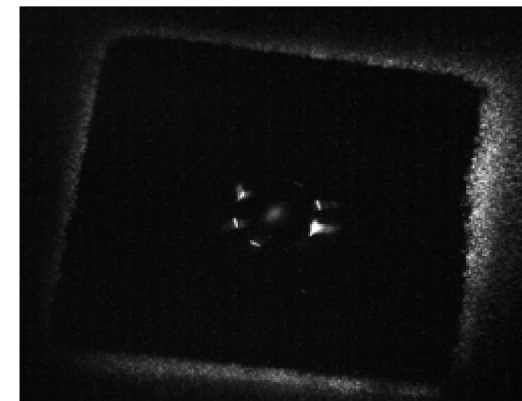
SNL Guided Atom Interferometer

Contact: Jongmin Lee
Adrian Orozco
Will Kendal

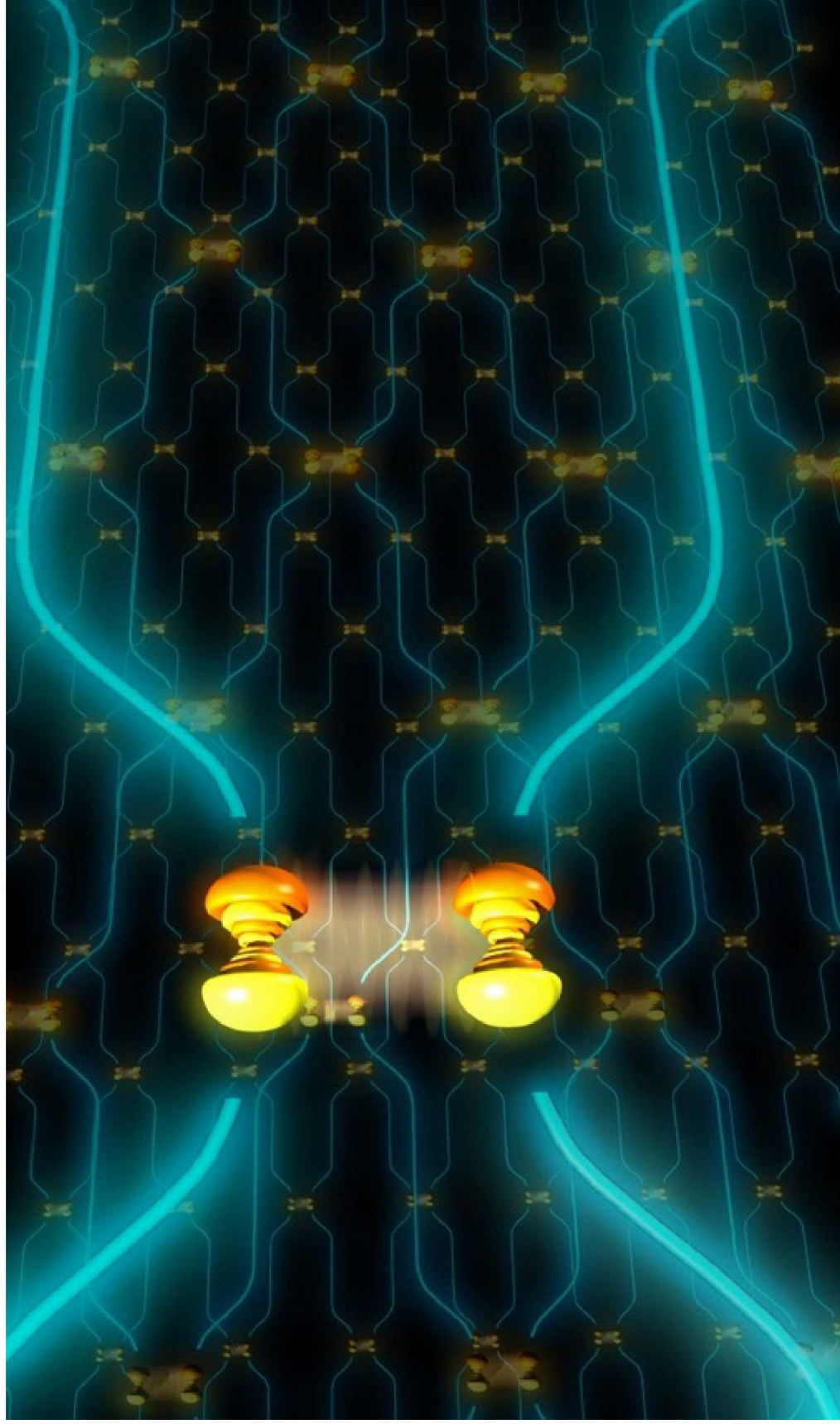


- $OD_N = 3.95$, $OD_1 = 0.078$; absorption measurement with nanofiber probe
- Trapped atom number = 50; Trap lifetime = 30ms

*New direction:
microfabricated
waveguides*



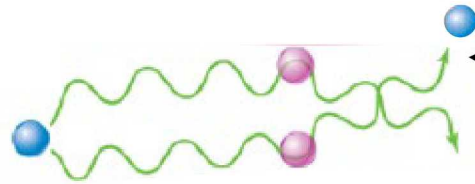
Interferometry with entangled spins





Quantum metrology of inertial forces with neutral atom spins

Matterwave interferometer



E. Rasel, Physics 5, 135 (2012)

$$\langle \Pi \rangle = \cos \theta$$

θ proportional to gravity

GHZ state

$$|\psi\rangle = |0\rangle^{\otimes N} + e^{iN\phi} |1\rangle^{\otimes N}$$

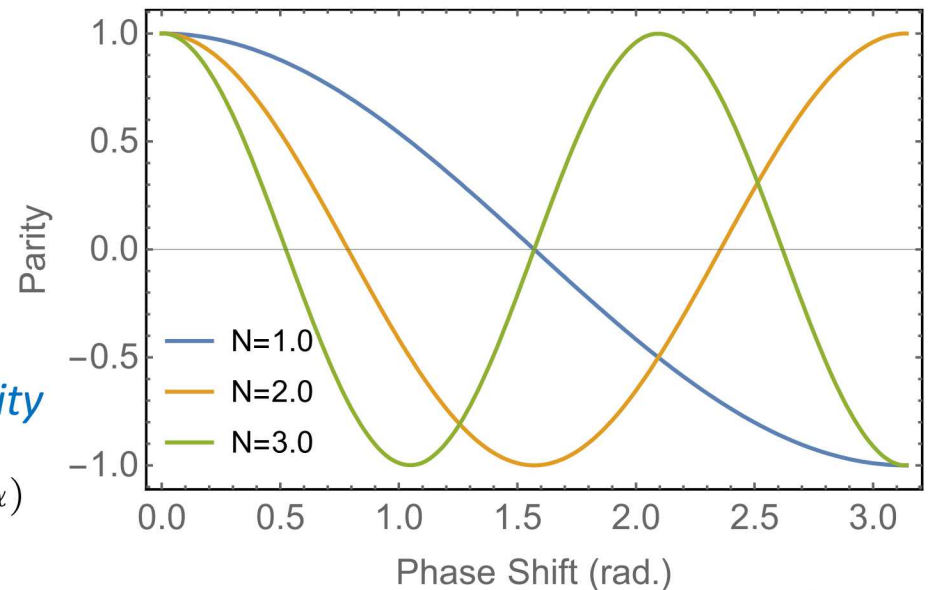
↙

$$\langle \Pi \rangle = \cos N\theta$$

Measure Parity

$$\Pi = \bigotimes_{\alpha=1}^N \sigma_z^{(\alpha)}$$

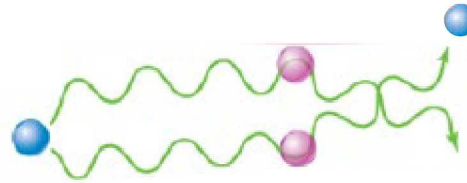
For 2 atoms $\langle \Pi \rangle = P_{ee} + P_{gg} - P_{eg} - P_{ge}$





Quantum metrology of inertial forces with neutral atom spins

Matterwave interferometer



E. Rasel, Physics 5, 135 (2012)

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GHZ state

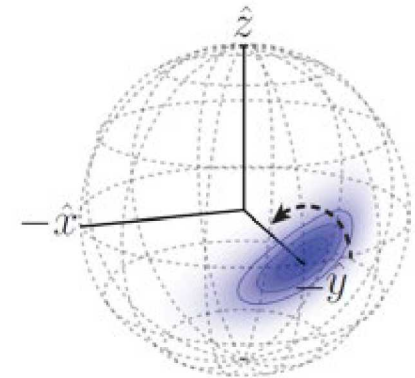
$$|\psi\rangle = |0\rangle^{\otimes N} + e^{iN\phi} |1\rangle^{\otimes N}$$

↙

$$\langle \Pi \rangle = \cos N\theta$$

? optimal states ?

← F Fröwis et al 2014 New J. Phys. 16 083010 →



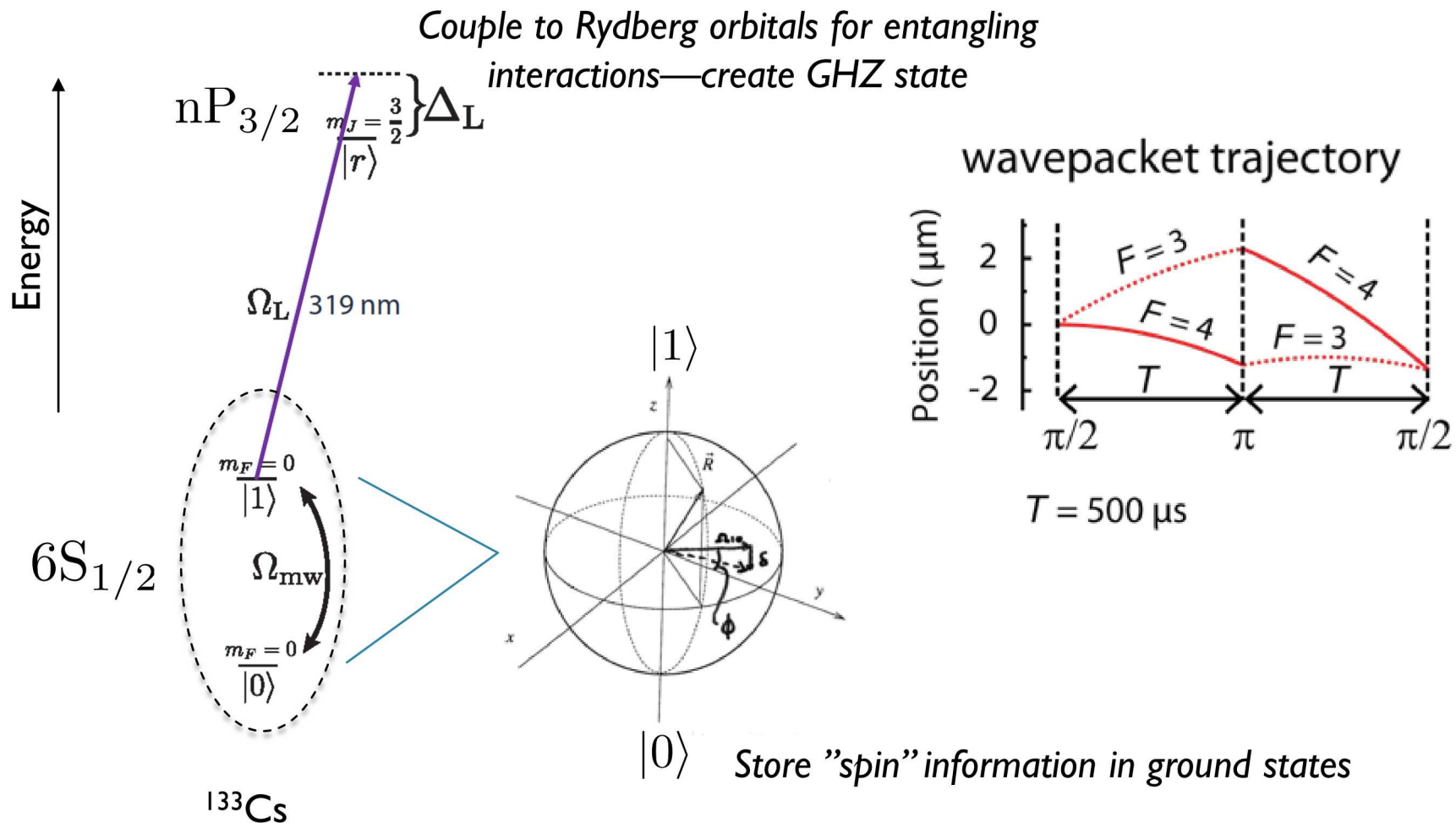
Spin-squeezed state

Wineland, et al., PRA 46 (1992)

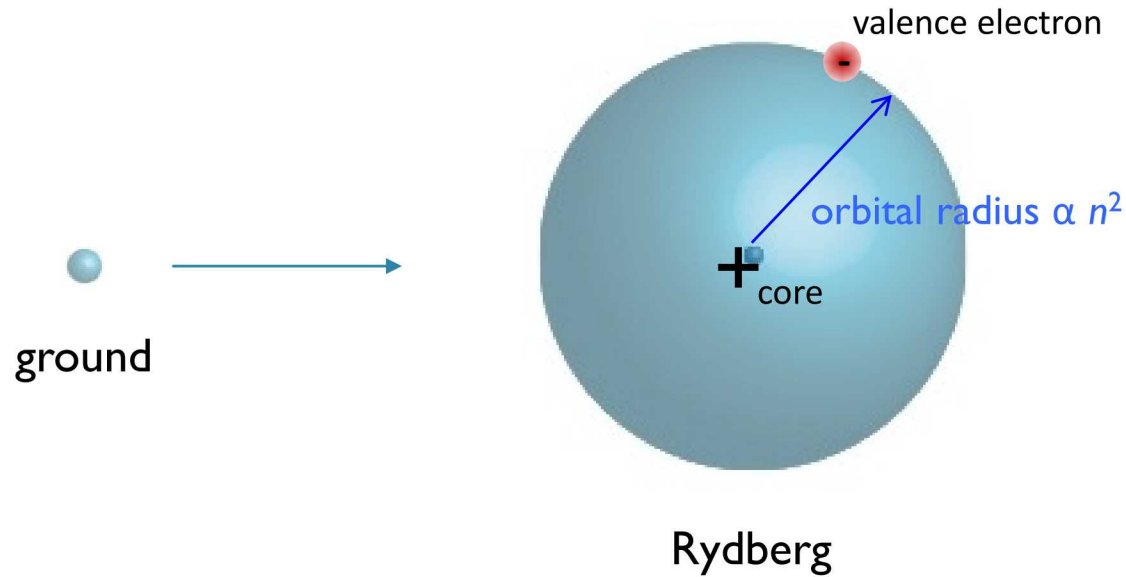
FRAGILE!

ROBUST!

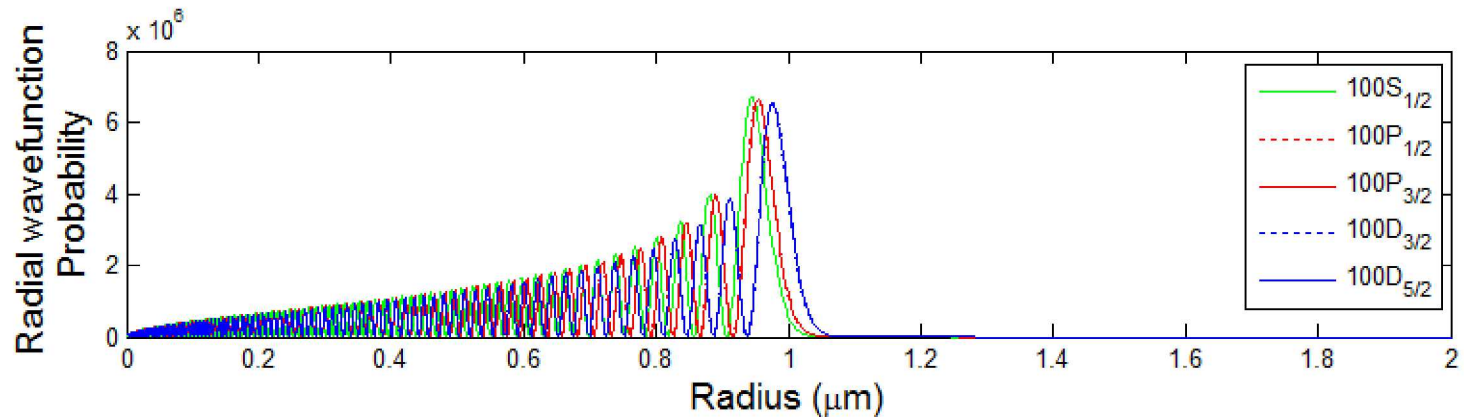
Neutral atom spins



Rydberg orbitals

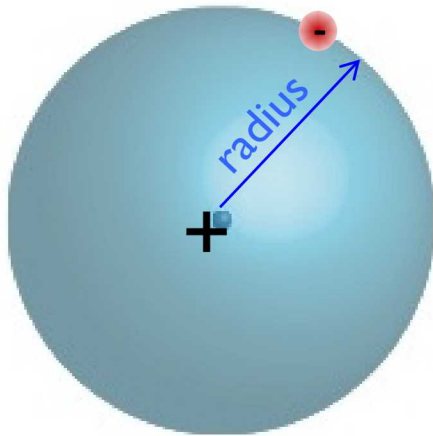


An example of the radial wavefunctions of a Cs atom at $n = 100$:



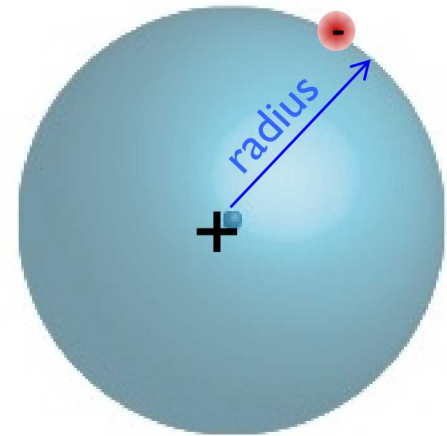
Interaction between *neutral* atoms

Valence electron in
Rydberg state



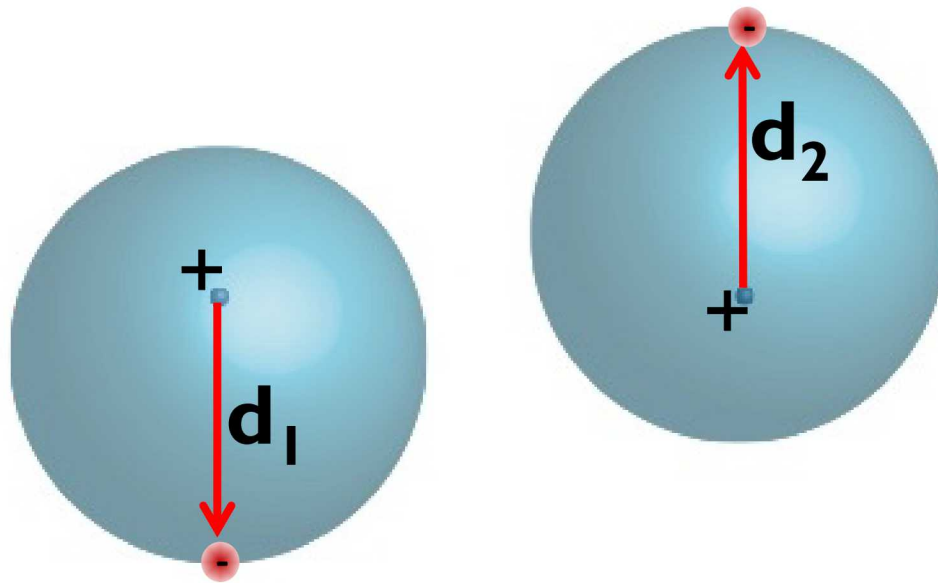
orbital radius $\propto n^2$

Valence electron in
Rydberg state



- Excite valence electron to Rydberg state—nearly ionized
- Atom becomes highly polarizable—strong interactions

Interaction between *neutral* atoms



van der Waals interaction

Parameter scaling

van der Waals

$$U \propto n^{11}$$

Lifetime

$$\tau \propto n^3$$

DC polarizability

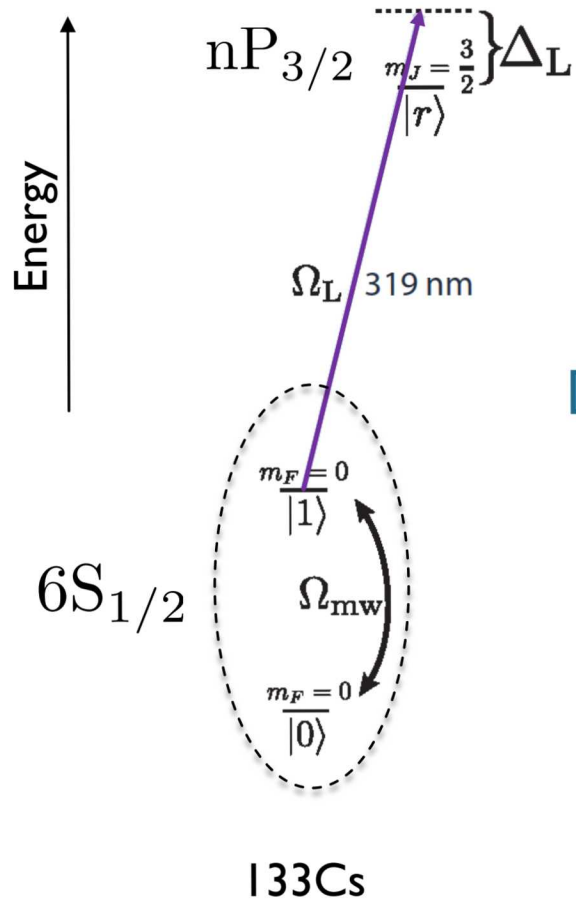
$$\alpha(0) \propto n^7$$

- Even the presence of another atom can cause a massive response $\gg 10$ MHz
- Induced Electric Dipole-Dipole Interaction $\propto 1/r^6$

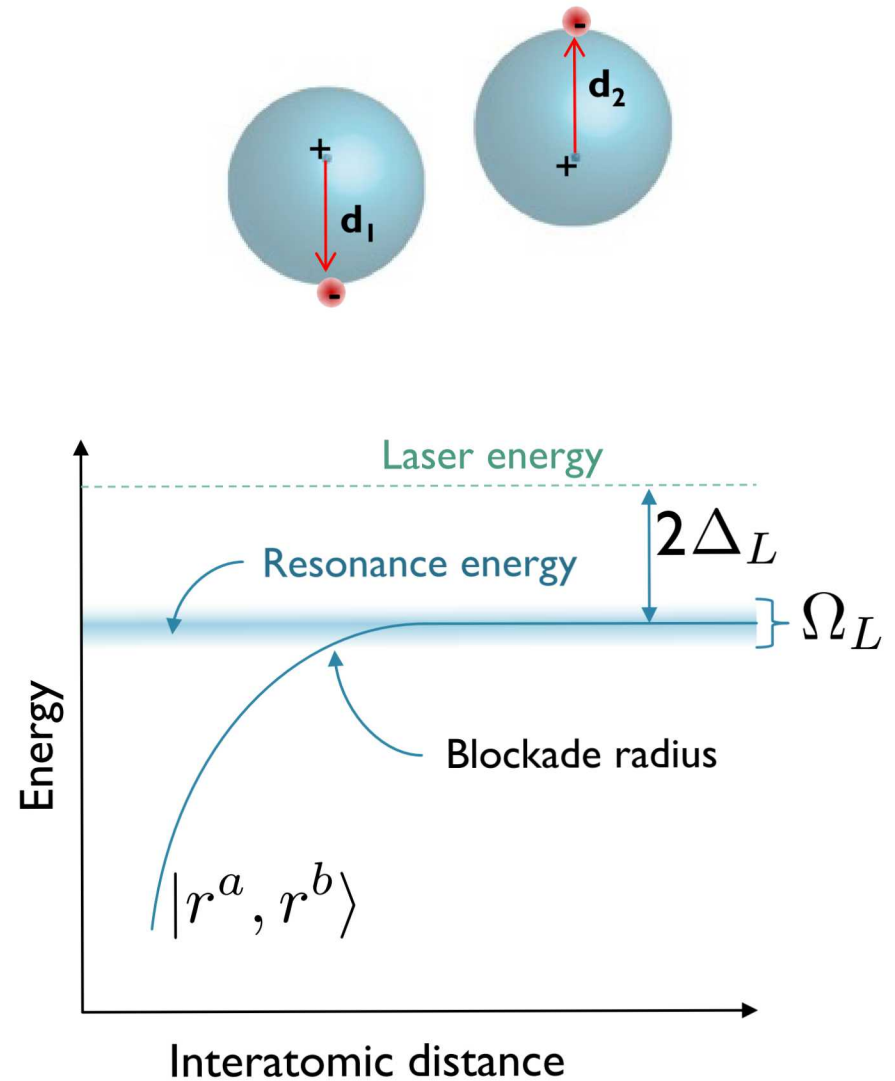
Entanglement demonstrations

- Madison: Phys. Rev. Lett. 104, 010503 (2010)
- Paris: Phys. Rev. Lett. 104, 010502 (2010)

Rydberg blockade

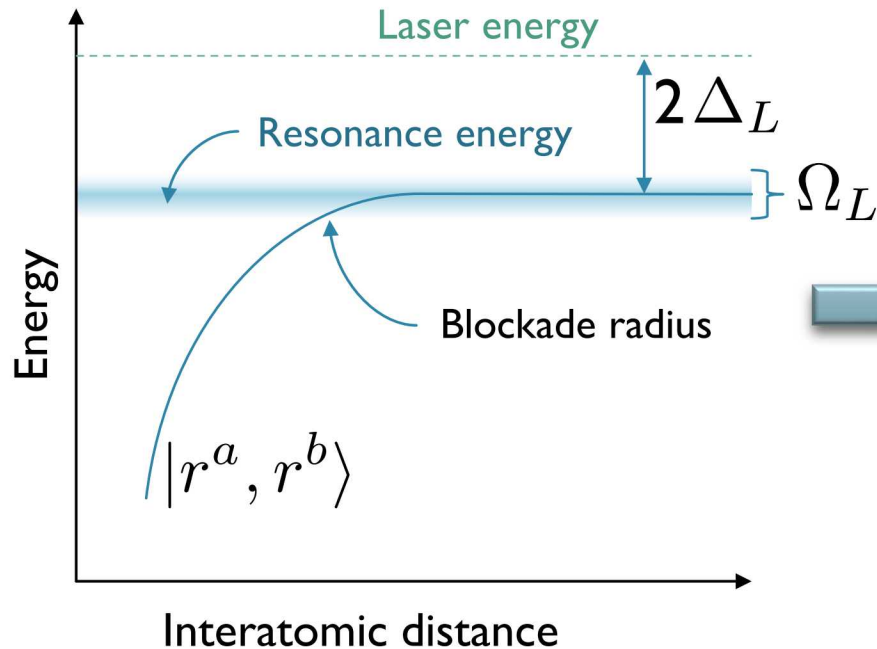


Single atom



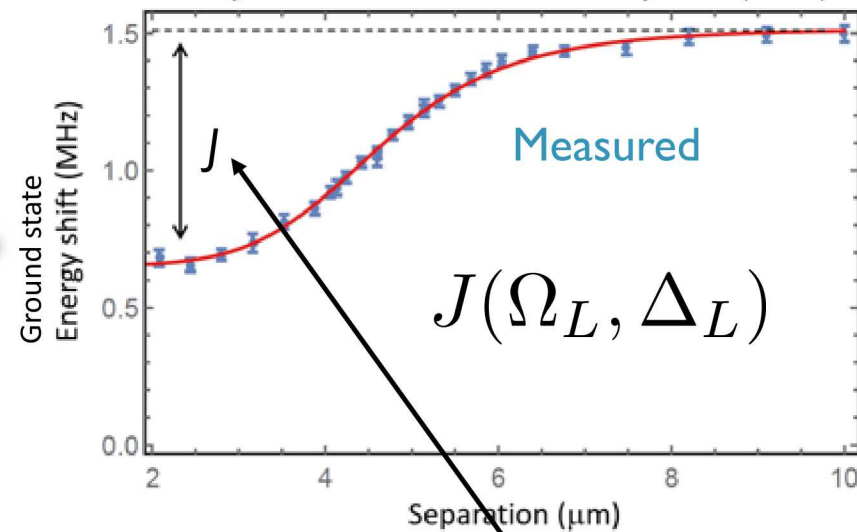
2 atoms

Direct Rydberg → Rydberg-Dressed



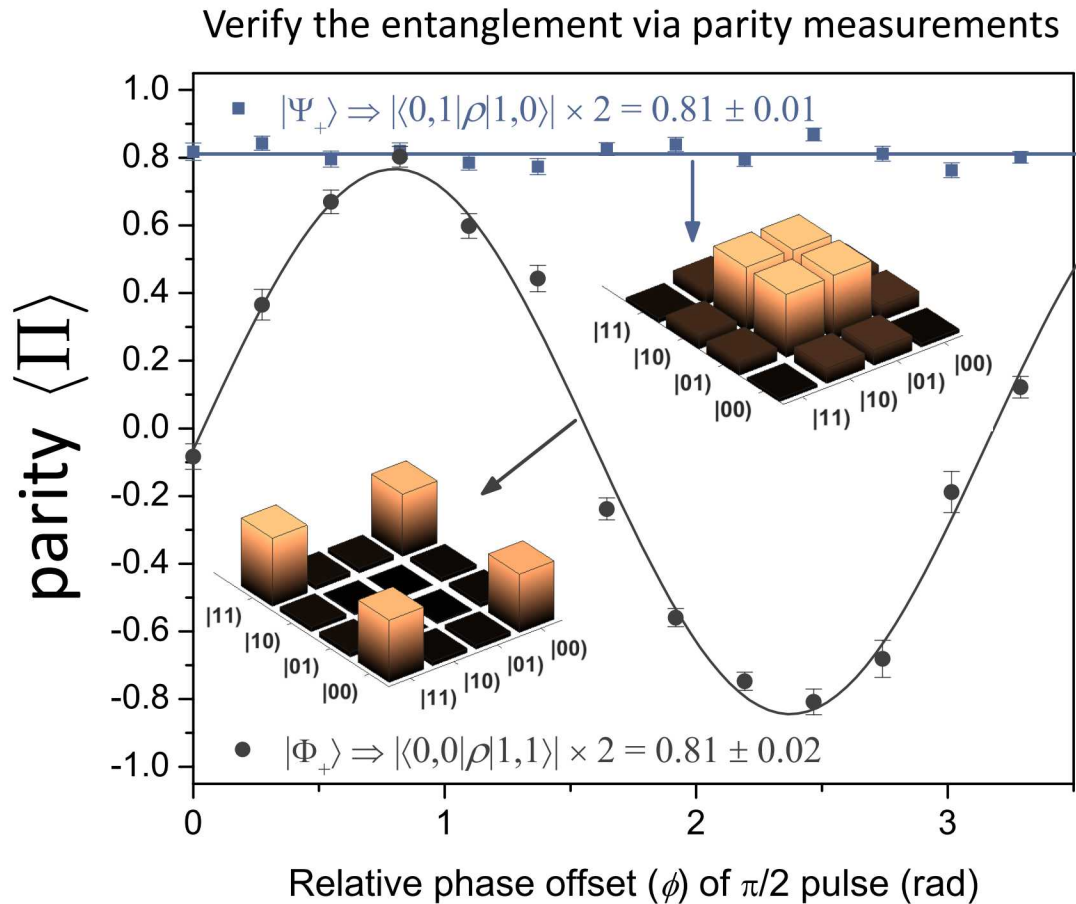
Rydberg-Dressed Interaction

Y.-Y. Jau, ..., I.D., **G.B.**, Nature Phys. **12** (2016)



$$H = \sum_{i=1}^N h_x^{(i)} \sigma_x^{(i)} + \dots + \sum_{i \neq j}^N J_{i,j} \sigma_z^{(i)} \otimes \sigma_z^{(j)}$$

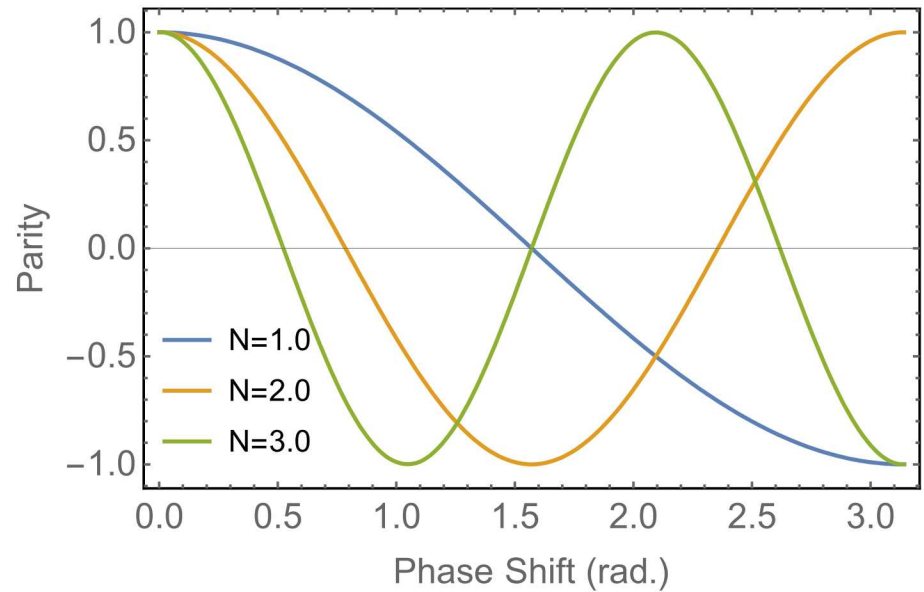
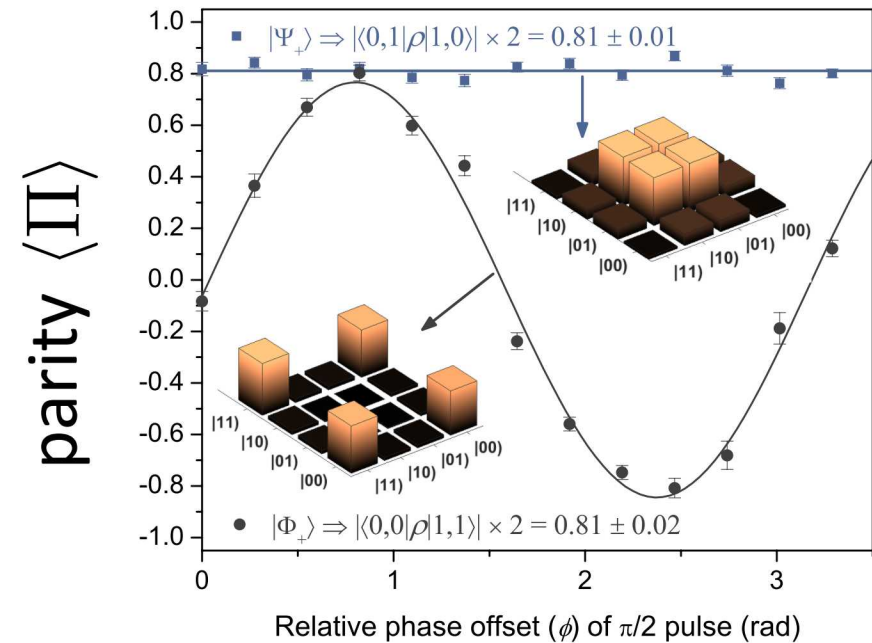
Spin-flip blockade



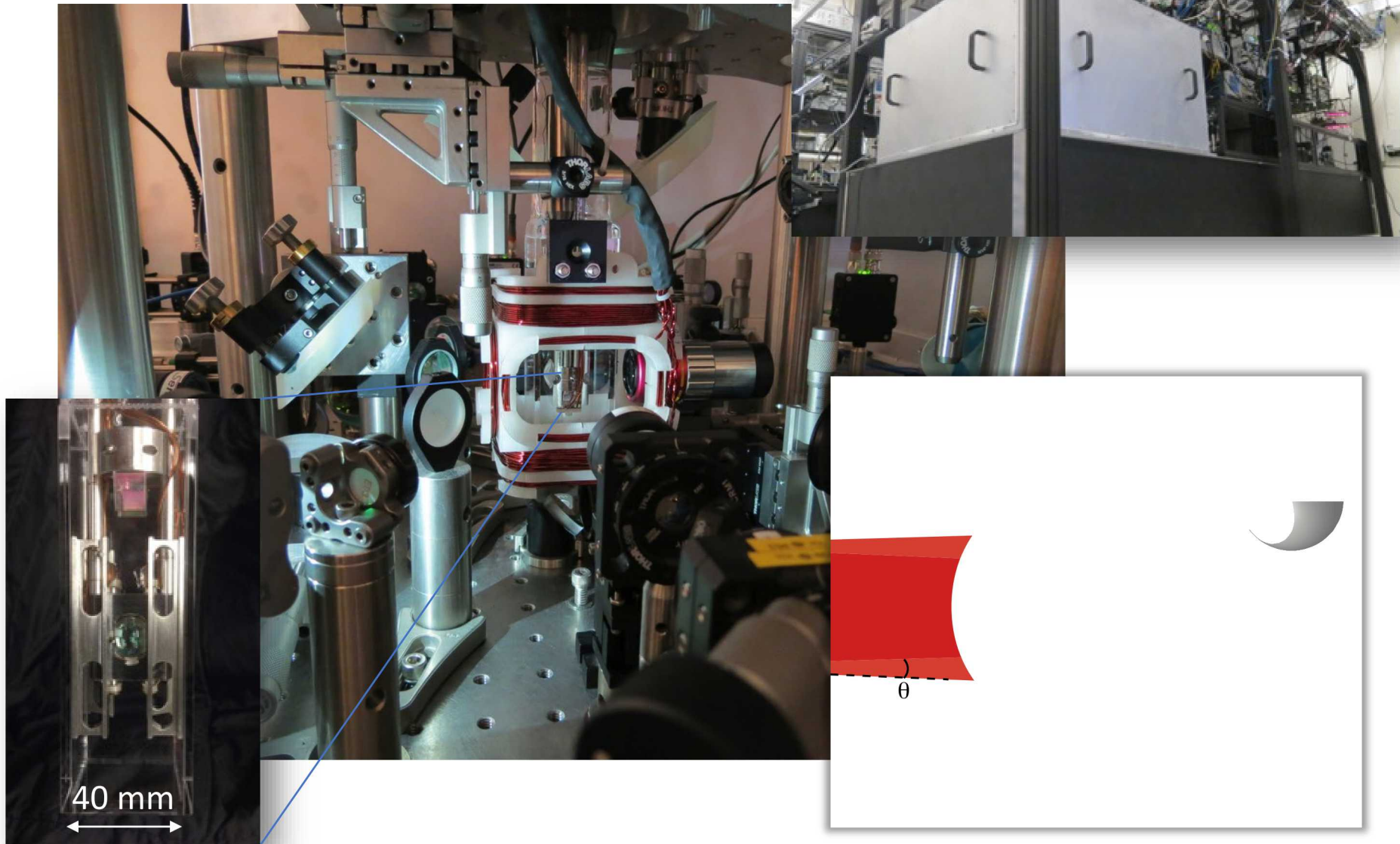
For 2 atoms

$$\langle \Pi \rangle = P_{ee} + P_{gg} - P_{eg} - P_{ge}$$

Spin-flip blockade



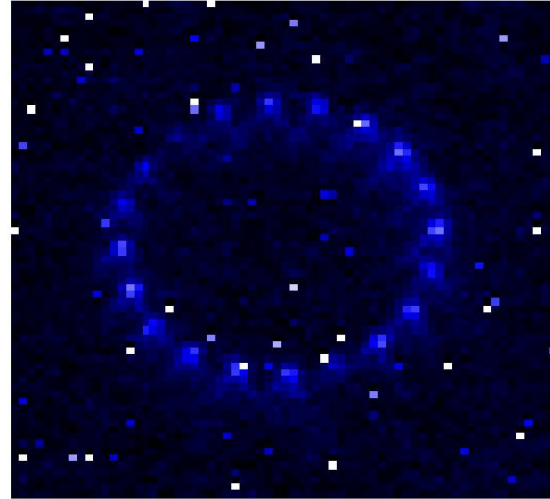
Experiment hardware



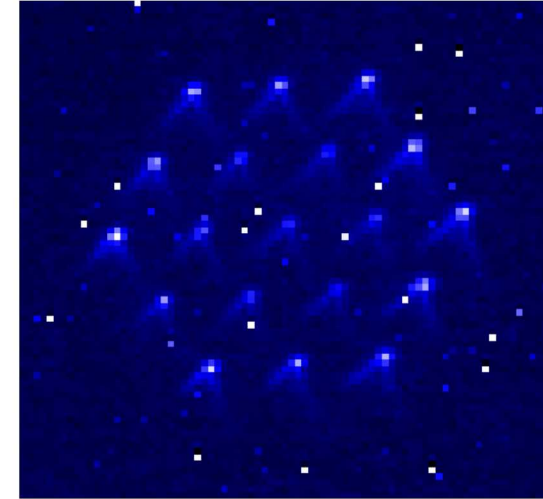
Outlook: large arrays

- Spatial light modulator allows arbitrary patterns
- Large systems > 100

Averaged over many exposures

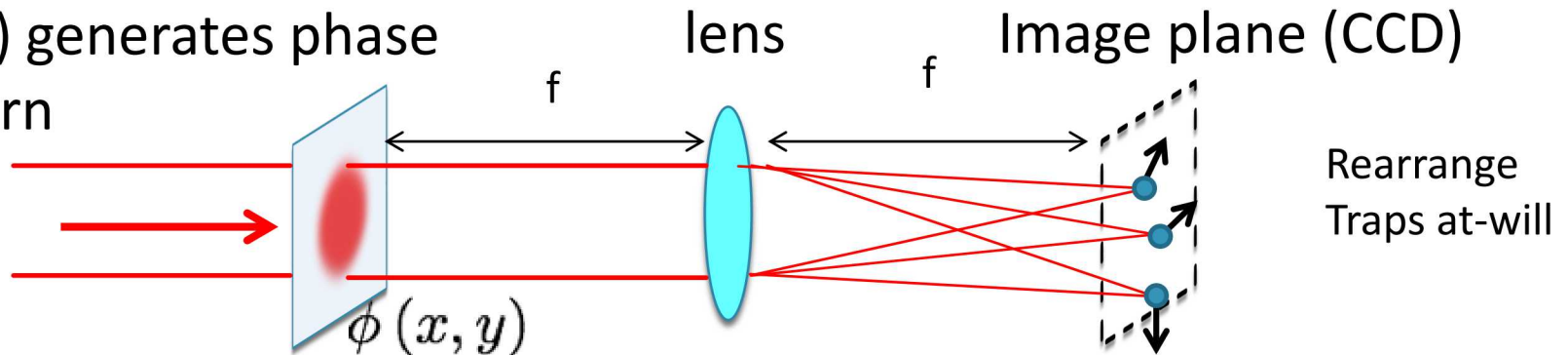


Ring lattice



Triangular lattice

Spatial Light Modulator
(SLM) generates phase
pattern

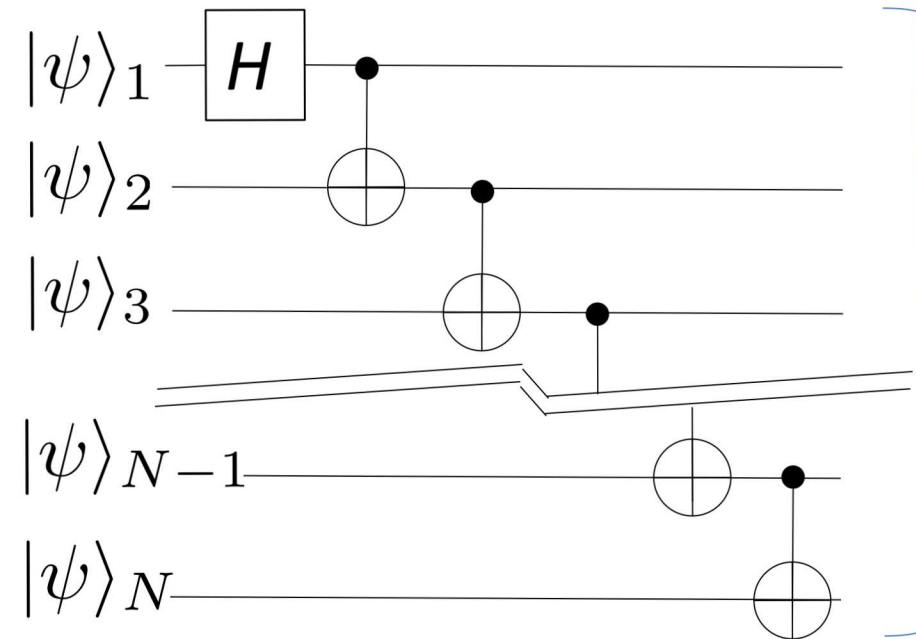


Data and image credit: Michael Martin, Sandia (now at LANL)

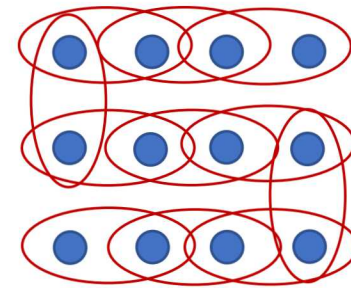
Path to complex quantum systems

Generating an entangled state of N atoms

Sequential gate approach



$$|\psi\rangle = |0\rangle^{\otimes N} + e^{iN\phi} |1\rangle^{\otimes N}$$



Array of atoms

Progress in entangling neutral atoms

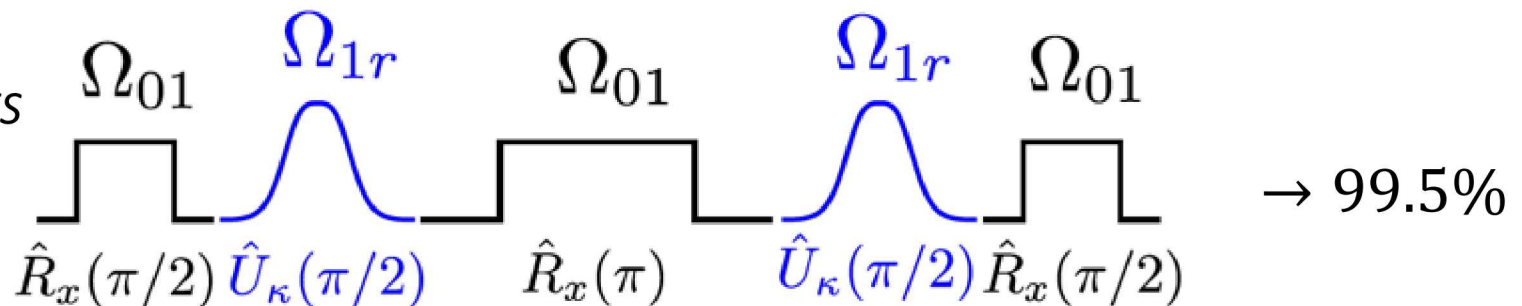
Atom	Method	Bell fidelity	Post selected fidelity	Year and reference
^{87}Rb	Blockade, simultaneous addressing	(0.46)	0.75	2009 [35]
^{87}Rb	Blockade, separate addressing	(0.48)	0.58	2009 [33]
^{87}Rb	Blockade, separate addressing	0.58	0.71	2010 [34]
Cs	Blockade, separate addressing	0.73	0.79	2015 [10]
Cs	Dressing, simultaneous addressing	0.60	0.81	2015 [11]
^{87}Rb	Local spin exchange	(0.44)	0.63	2015 [31]

M. Saffman, J. Phys. B: At. Mol. Opt. Phys. 49 (2016) 202001

Now 89% to 97%

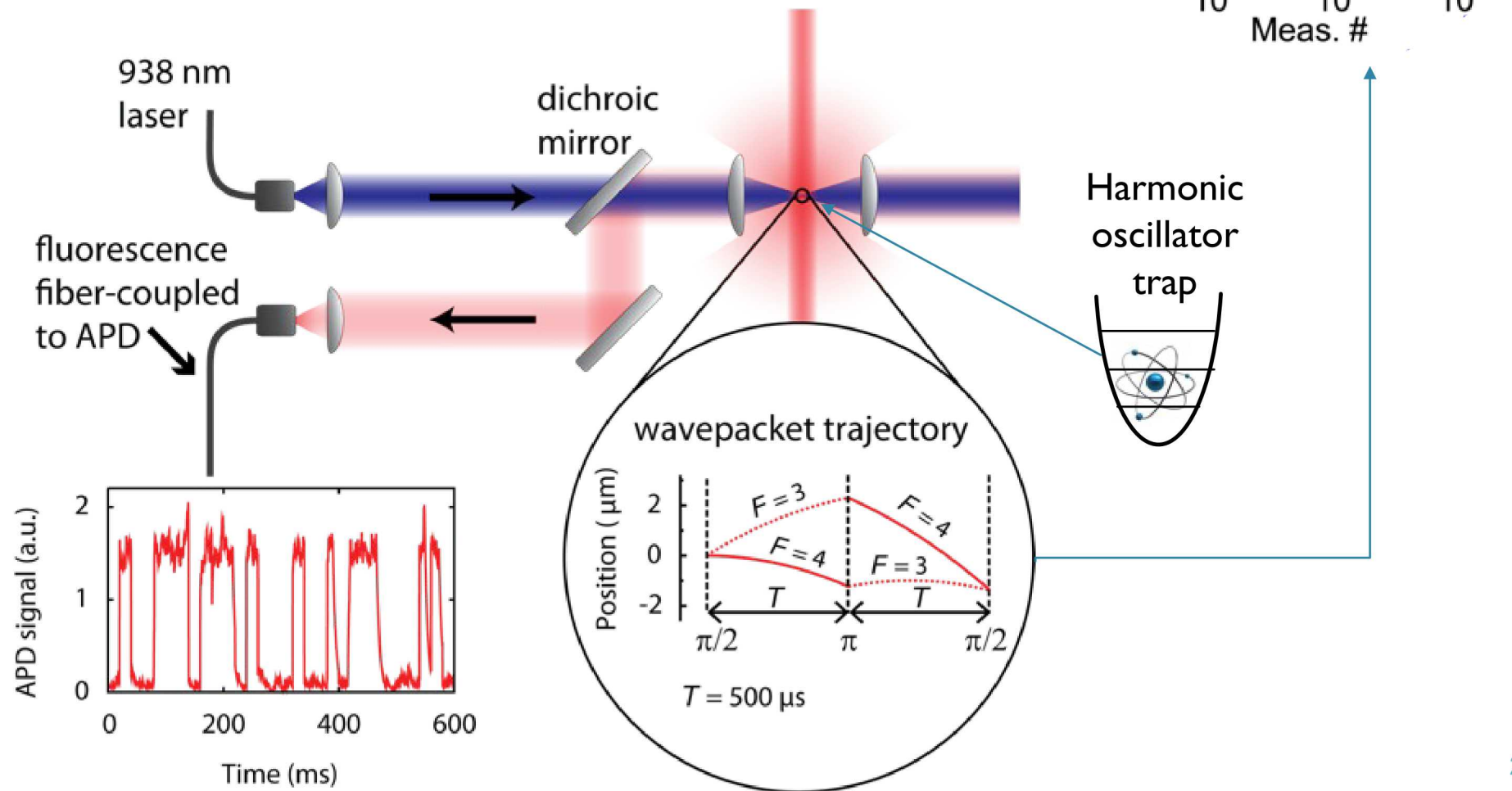
(see new results from Harvard and Wisconsin)

*Dressed-Rydberg
approach ENABLES
this protocol*

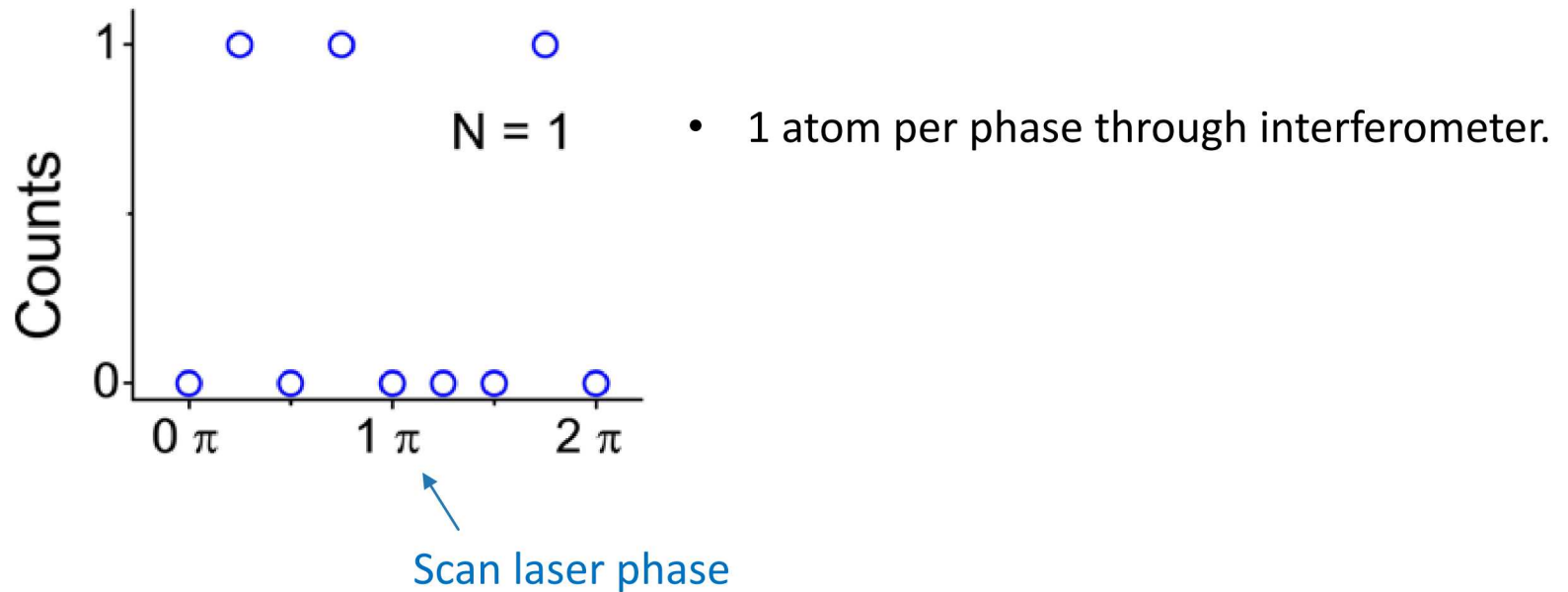


See our new theory work: A. Mitra, ..., G.B., I. Deutsch, *Robust Mølmer-Sørensen gate for neutral atoms using rapid adiabatic Rydberg dressing*, arXiv:1911.04045 (2019)

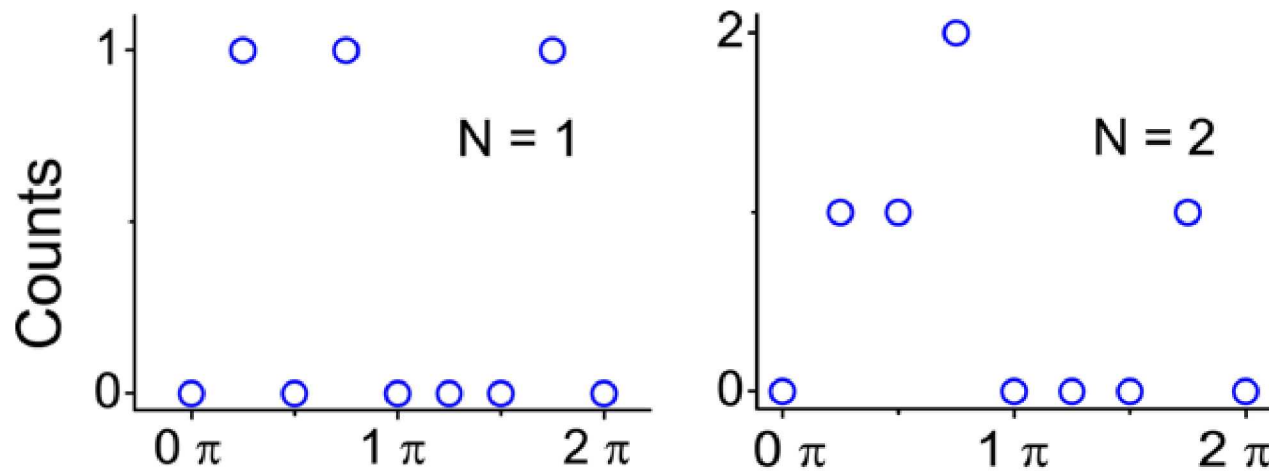
Matter wave control at the single atom level



Building fringe one atom at a time

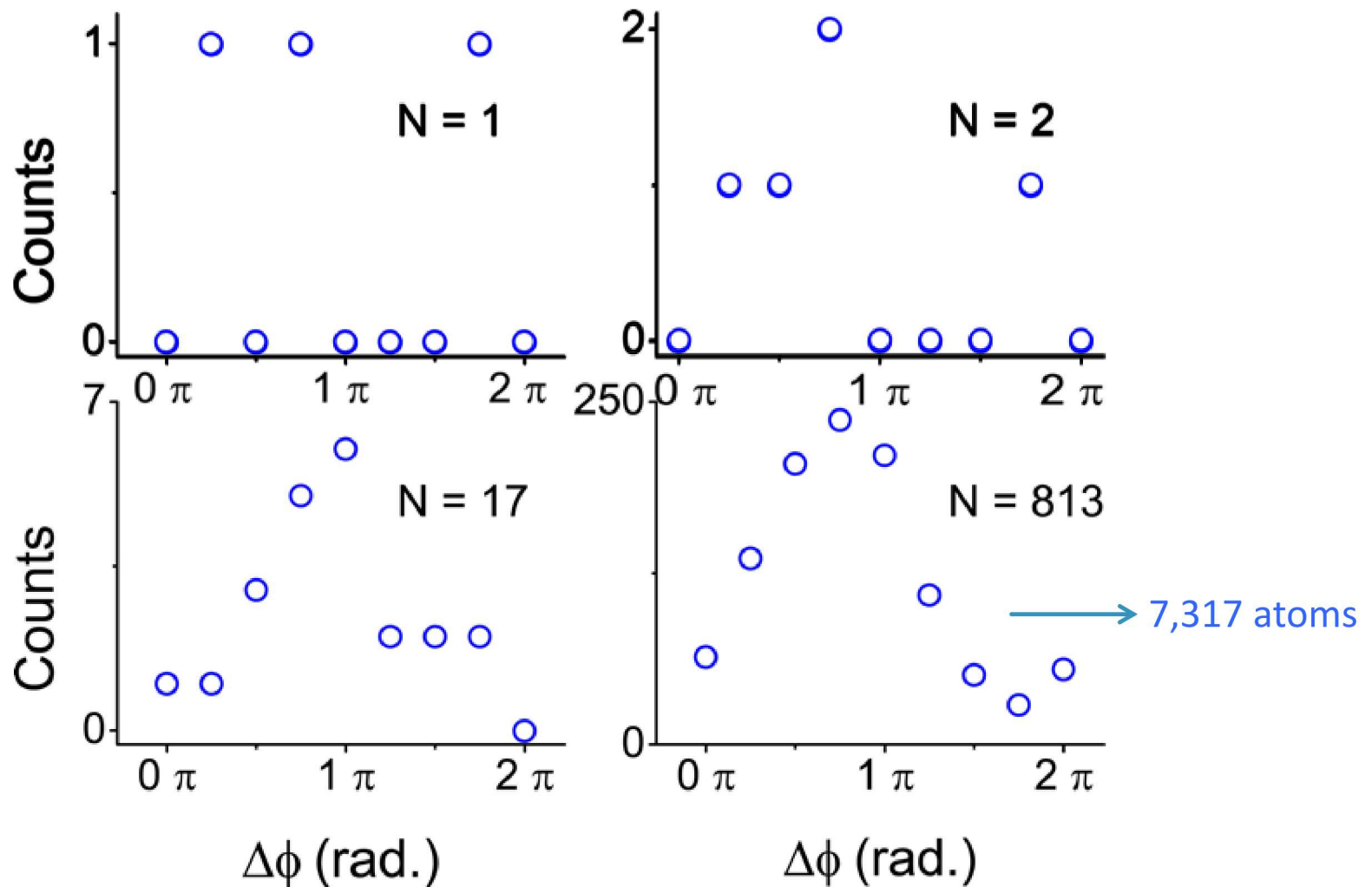


Building fringe one atom at a time

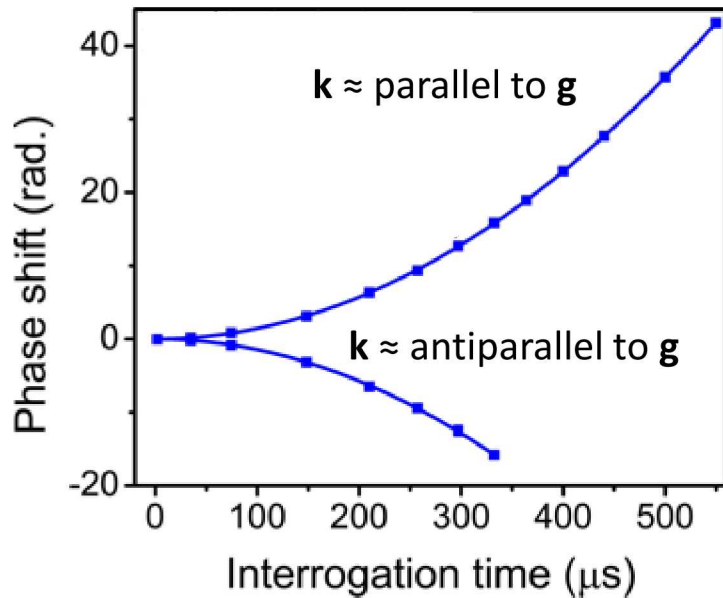
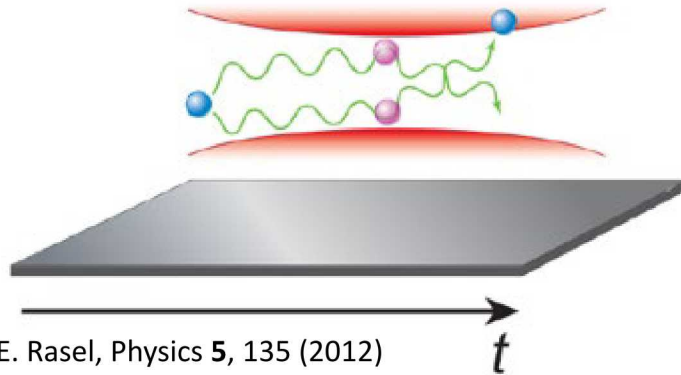


- 2 atoms per phase through interferometer.

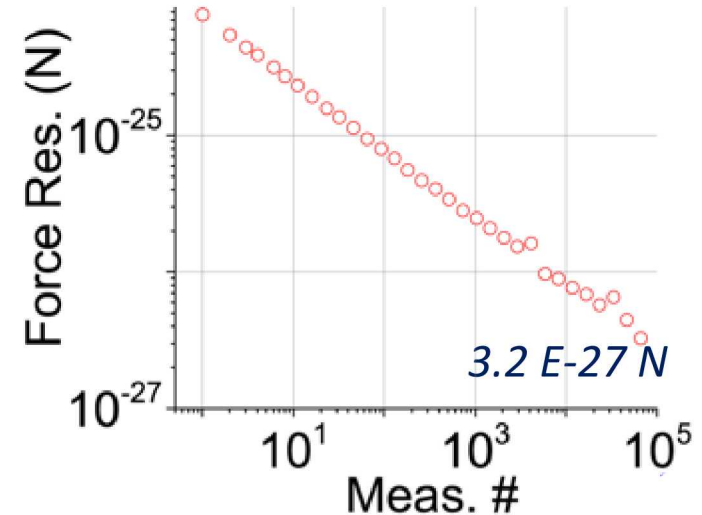
Building fringe one atom at a time



Force resolution of a single atom interferometer

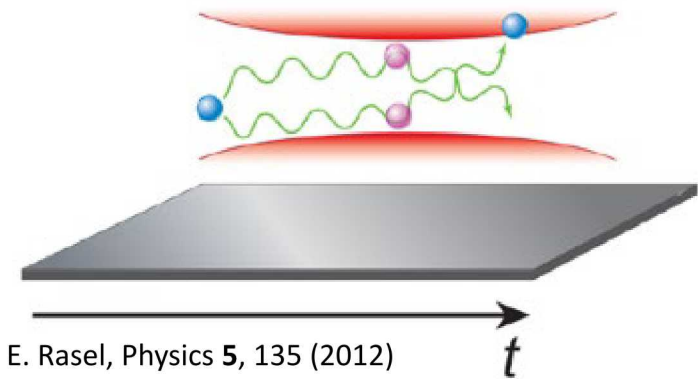


Parazzoli, et al., Phys. Rev. Lett. **109**, 230401 (2012)

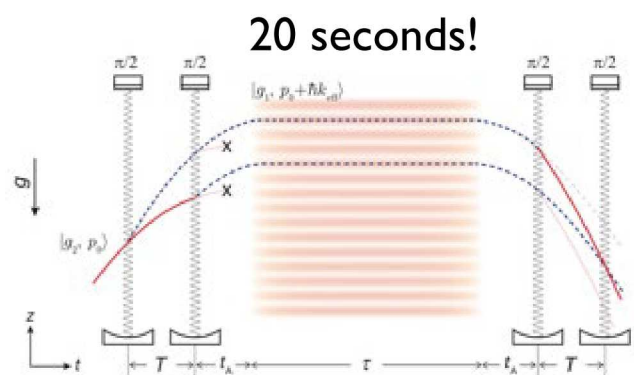
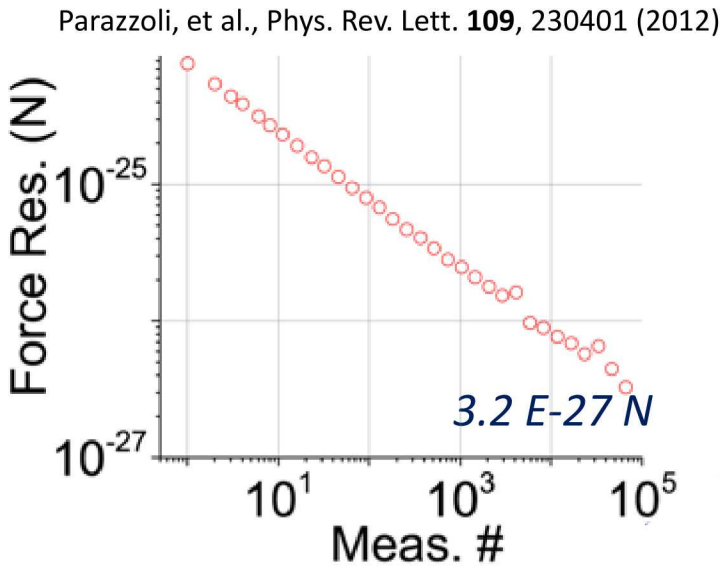
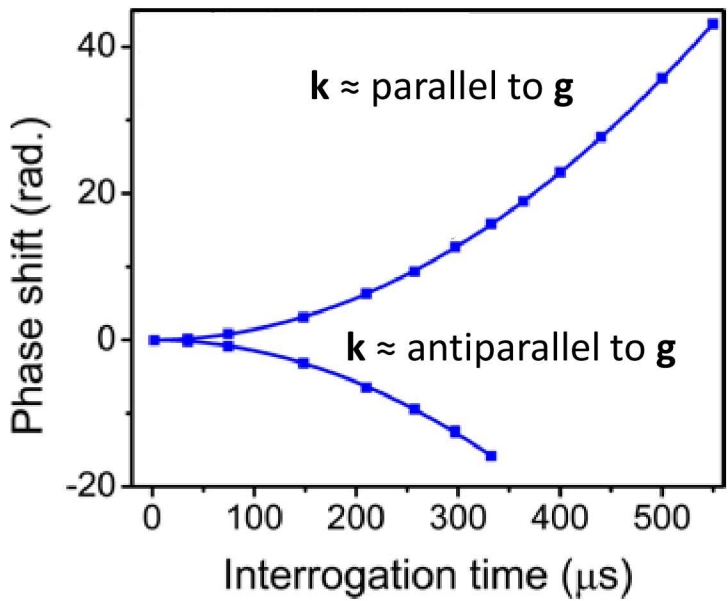


$$\Delta\phi = \vec{k}_{eff} \cdot \vec{g} T^2$$

Force resolution of a single atom interferometer



E. Rasel, Physics 5, 135 (2012)



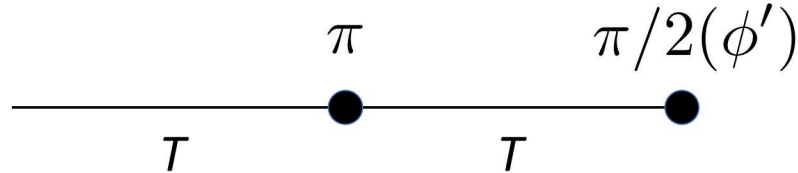
V. Xu, et al., Science 2019

The Ideal Atom Interferometer

N entangled atoms

$$\frac{|g \dots g\rangle + |e \dots e\rangle}{\sqrt{2}}$$

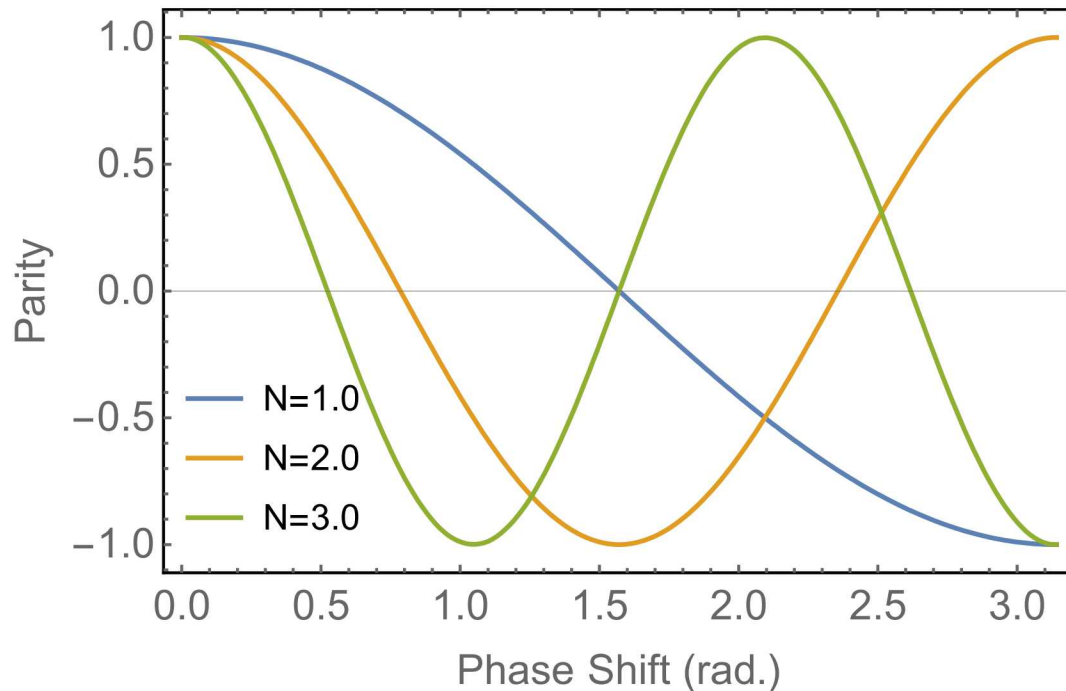
Pulse sequence



Measure Parity

$$\Pi = \bigotimes_{\alpha=1}^N \sigma_z^{(\alpha)}$$

For 2 atoms $\langle \Pi \rangle = P_{ee} + P_{gg} - P_{eg} - P_{ge}$



As *N* grows:

Parity oscillates faster

$$\langle \Pi \rangle = \cos \Theta_N$$

$$\begin{aligned} \Theta_N &= N \mathbf{K} \cdot \mathbf{a} T^2 \\ &= N \phi \end{aligned}$$

Theory by Constantin Brif
&

Brandon Ruzic
Sandia National Laboratories

Imperfect Initial State Preparation

*Random initial
state fraction*

*Uncertainty in
initial phase*

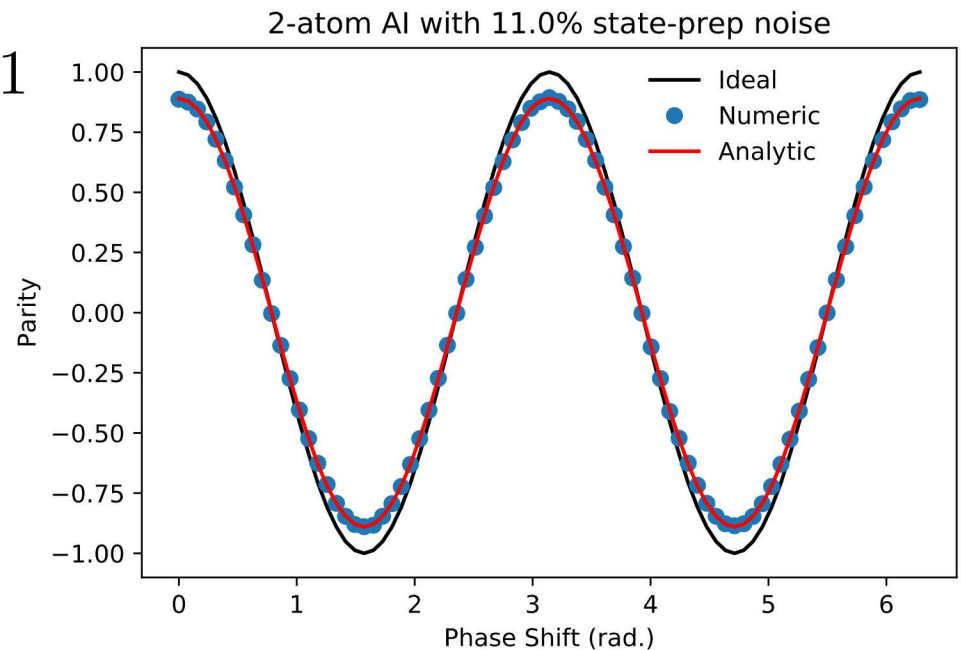
Analytic result when $p + \sigma^2/2 \ll 1$

$$\langle \Pi \rangle \approx (1 - p - \sigma^2/2) \cos \Theta_N$$

$$\Theta_N = N\mathbf{K} \cdot \mathbf{a}T^2 = N\phi$$

Visibility is reduced

$$p + \sigma^2/2 \approx 0.11 \text{ for two atoms}$$

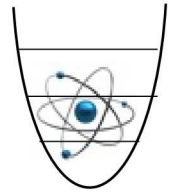


Metrological advantage is only linearly sensitive to entanglement imperfection

Random Doppler Shift

Atoms start in optical tweezers

$$\rho = \rho_{\text{vib}} \otimes |\psi_{\text{spin}}\rangle \langle \psi_{\text{spin}}|$$

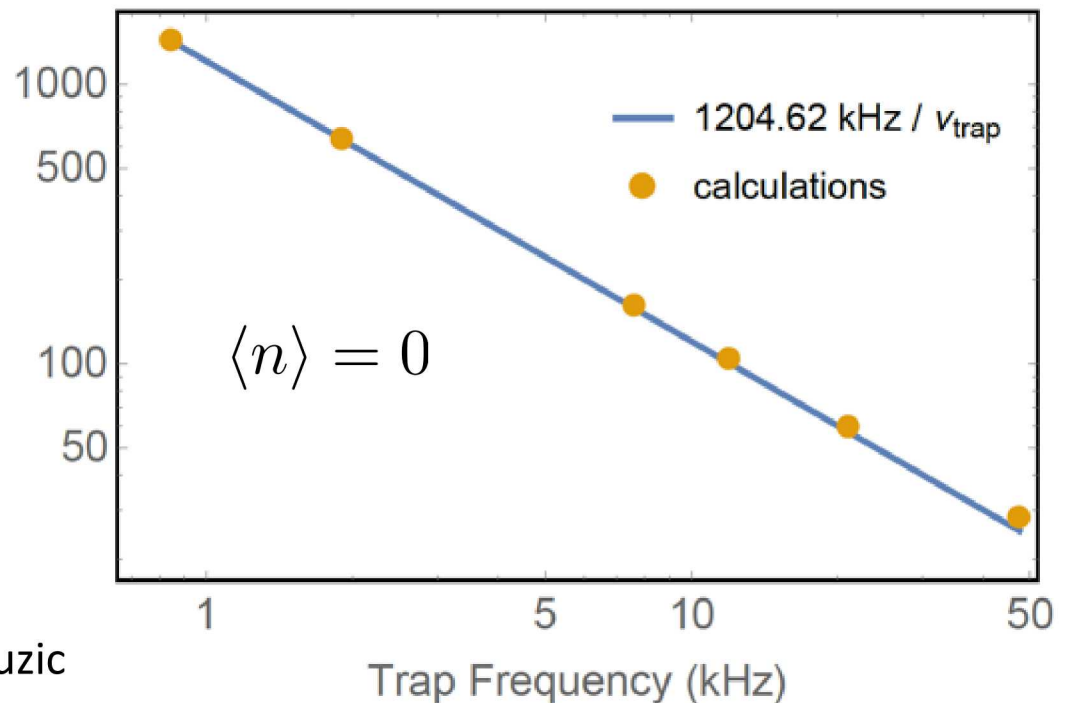


Initial momentum spread

- Each run randomly samples a momentum distribution
- doppler shift
- detuning error in pulses

Max. useful
number of
atoms

$1/\eta$



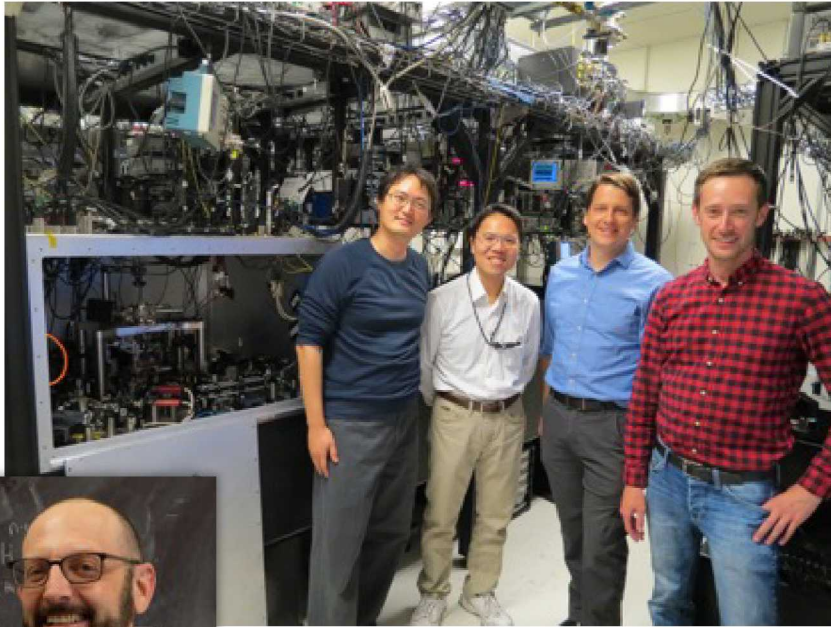
Theory by Constantin Brif & Brandon Ruzic
Sandia National Laboratories

$\propto$ Temperature

Summary

- Sandia is a significant player in the development of atom interferometer technique and instrumentation
- Entangled atom interferometer for > 1000 atoms appears feasible with our current approach
- Doppler broadening is the dominant limitation
- Phase resolution performance of GHZ state with 100 atoms is already comparable to typical LPAI results (10 mrad/shot)

Sandia-OU-UNM Team



J. Lee, Y-Y. Jau, M. Martin, G.B.



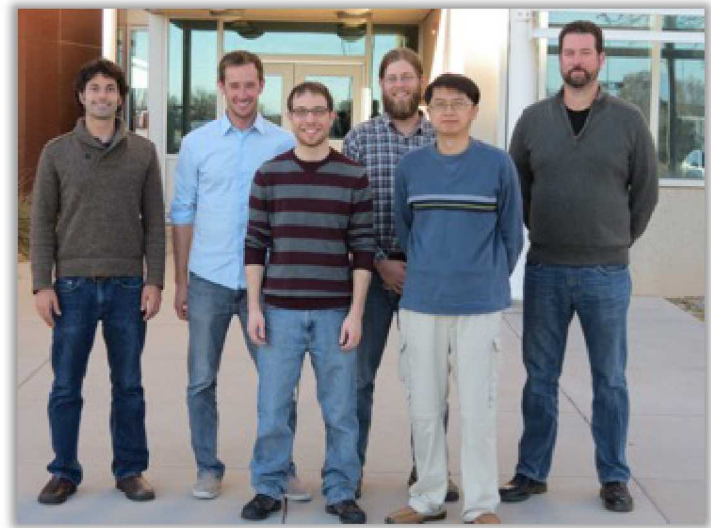
Postdoc & Graduate Research
Positions Available



I. Deutsch (UNM)

Not pictured

- Adrian Orozco
- Will Kendal
- Constantin Brif
- Brandon Ruzic
- Peter Schwindt
- Bethany Little
- Matthew Chow
- Anupam Mitra



L. Parazzoli, G. B., A. Hankin, A. Ferdinand, J. Chou, G. Burns