

The Stabilizer Rank of the T-Gate Magic State

Stabilizer states

- given a state $|\Psi\rangle$, the stabilizer group is the set of tensor products of Pauli operators that stabilizes $|\Psi\rangle$
e.g. $I|0\rangle = |0\rangle$ and $Z|0\rangle = |0\rangle$.
- stabilizer states $\{|\phi_i\rangle\}_i$ on n d-dimensional qudits are the set of states stabilized by d^n commuting stabilizer group elements
- we define the Clifford gateset to be closed for stabilizer states
- gateset is not universal
- projection onto Paulis also takes stabilizer states to themselves:
Clifford subtheory = Pauli measurement + Clifford gateset
+ (convex combos of) stabilizer states

Gottesman-Knill theorem: there is a polynomial-time classical algorithm to simulate the Clifford subtheory [1].

- stabilizer states $\{|\phi_i\rangle\}_i$ form an overcomplete basis.
 - therefore, any state $|\Psi\rangle$ can be expressed as $|\Psi\rangle = \sum_i c_i |\phi_i\rangle$
- the **stabilizer rank $\chi(\Psi)$** of a pure state $|\Psi\rangle$ is the minimal number χ of states required in a stabilizer state decomposition of $|\Psi\rangle$.

Trivial tensor bound property: Let $\chi(n)$ be the stabilizer rank of $|\Psi\rangle^n$. Since the tensor product of two stabilizer states is a stabilizer state, it follows that $\chi(m+n) \leq \chi(m)\chi(n)$.

T-Gate Magic State

$$T = 2^{1/2}(|0\rangle + e^{\pi i/4} |1\rangle).$$

The T-gate magic state extends the Clifford subtheory to universality in the limit of $t \rightarrow \infty$, where t denotes the power of tensoring the $|T\rangle^t$

- T-gate magic state equivalent to a T-gate by teleportation protocol
- imperfect T-gates can be "distilled" to increase purity by magic state distillation [2]

It is postulated that $\chi(t)$ grows slowest with increasing number of qubits.

- for $t = 1$ qubit T gate magic state, $\chi(1) = 2$.
- for $t = 2$ qubit T gate magic states, $\chi(2) = 2$.
- for $t = 3$ qubit T gate magic states, $\chi(3) = 3$.
- for $t = 6$ qubit T gate magic states, $\chi(6) \leq 7$.

It is conjectured that this bound is tight [3].

Using the tensor bound property, we can obtain upper bounds on the T-gate stabilizer rank as $t \rightarrow \infty$ (see Figure 1):

- $\chi(1)^t = 2^t$.
- $\chi(2)^{t/2} = 2^{0.5t}$ where t is even.
- $\chi(3)^{t/3} = 3^{t/3} \approx 2^{0.53t}$ where t is a multiple of 3.
- $\chi(6)^{t/6} = 7^{t/6} \approx 2^{0.47t}$ where t is a multiple of 6.

The last bound provides the most favorable asymptotic scaling, and so the outcome of this procedure is often reported as a scaling of $O(2^{0.468t})$ for the qubit T gate magic state.

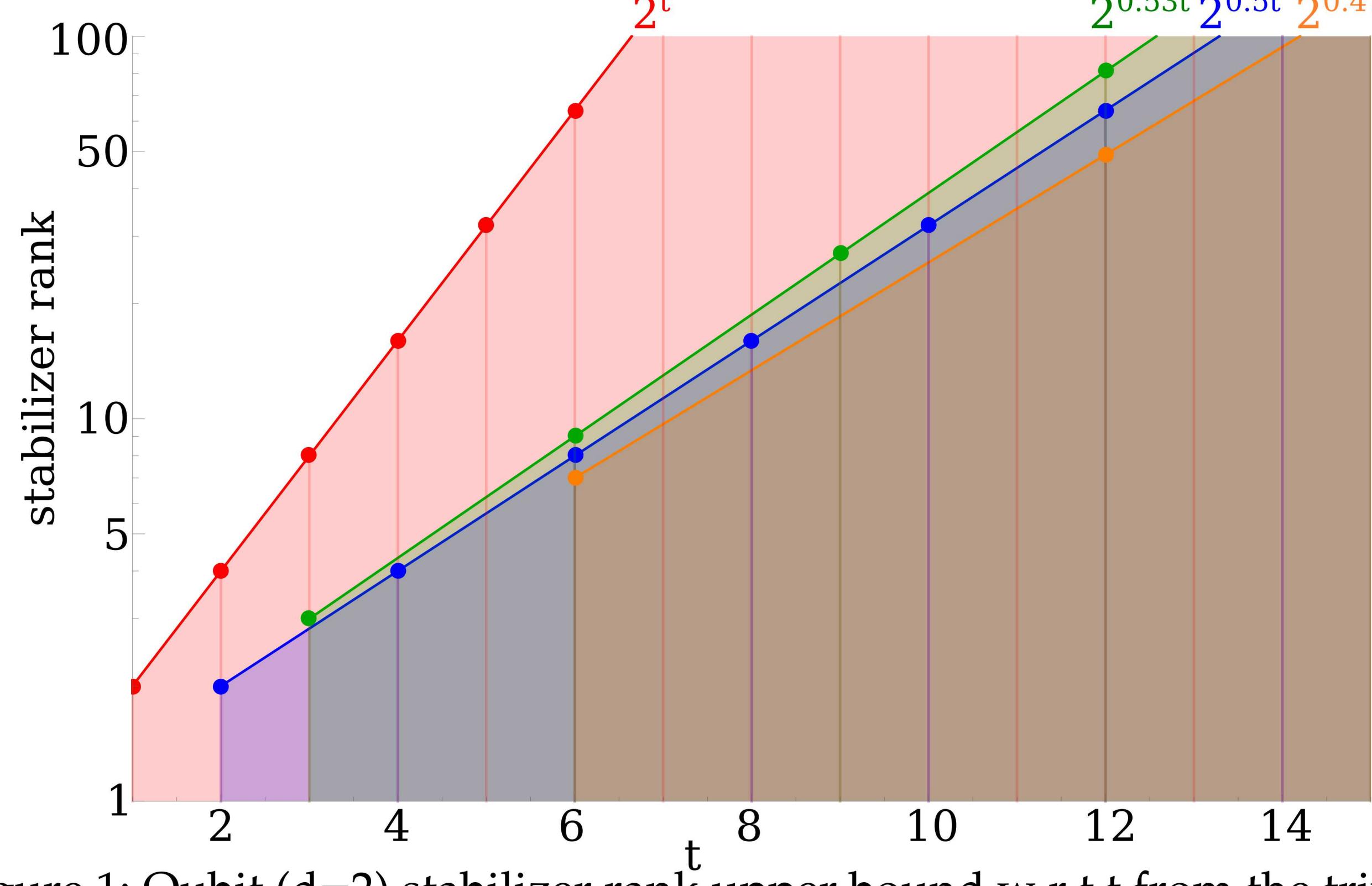


Figure 1: Qubit ($d=2$) stabilizer rank upper bound w.r.t t from the trivial tensor bound property using the stabilizer rank found for 1, 2, 3, and 6 tensored T-gate magic states

Stabilizer Rank as a Computational Cost Metric for Pauli-Based Computation

- the inner product of two stabilizer states $\langle \phi_i | \phi_j \rangle$ is governed by Gaussian elimination and therefore scales as $O(n^3)$.
- in a Pauli-based computation, given a projector $E = \prod_{i=1}^t (I + \sigma_i P_a)/2$ where P_a is a Pauli operator, it follows that $E_{ij}^{(x,t)} = c_i x c_j \langle \phi_i | \prod_{i=1}^t | \phi_j \rangle$
- since the number of terms is $\chi(t)^2$, we want to use the lowest $\chi(t)$ that we can to simulate this classically
- therefore, stabilizer state decompositions are a good way to measure the cost of strong simulation (i.e. classical computation of the probability outcome of a measurement)
 - the stabilizer rank of states (like the T-gate magic state) that extend the Clifford subtheory to a universal set necessarily grow exponentially with tensoring
 - the exponential factor of this scaling determines how large of a universal quantum computer can be simulated by today's classical computers

Current Technique for Finding Stabilizer Rank Not Adequate

The fastest method for finding the stabilizer rank of the T-gate magic state is numerical and based on Monte Carlo (Glauber dynamics) sampling of the full stabilizer state space [3]

- for greater than $t > 7$, such an approach ceases to converge

To find a better asymptotic tensor scaling requires reaching larger t , which therefore requires relying on a new method.

- a non-numerical method would be especially attractive

Alternative Method: Odd Prime-Dimensional Wigner-Weyl-Moyal (WWM) Formalism

Instead of considering our magic state in terms of vectors in Hilbert space, we can use a kernel (or quasiprobability) representation instead.

Given a set of operators $R(x)$ indexed by x , that are Hilbert-Schmidt orthogonal, any operator A can be represented:

$$A = d^{-1} \sum_x \text{Tr}(R(x)A)R(x) \equiv \sum_x A_x(x)R(x).$$

We choose $R(x)$ to be Hermitian, self-inverse and unitary, and so $A_x(x)$ are real.

In particular, we consider the odd-prime- d Wigner "reflection" operators [4].

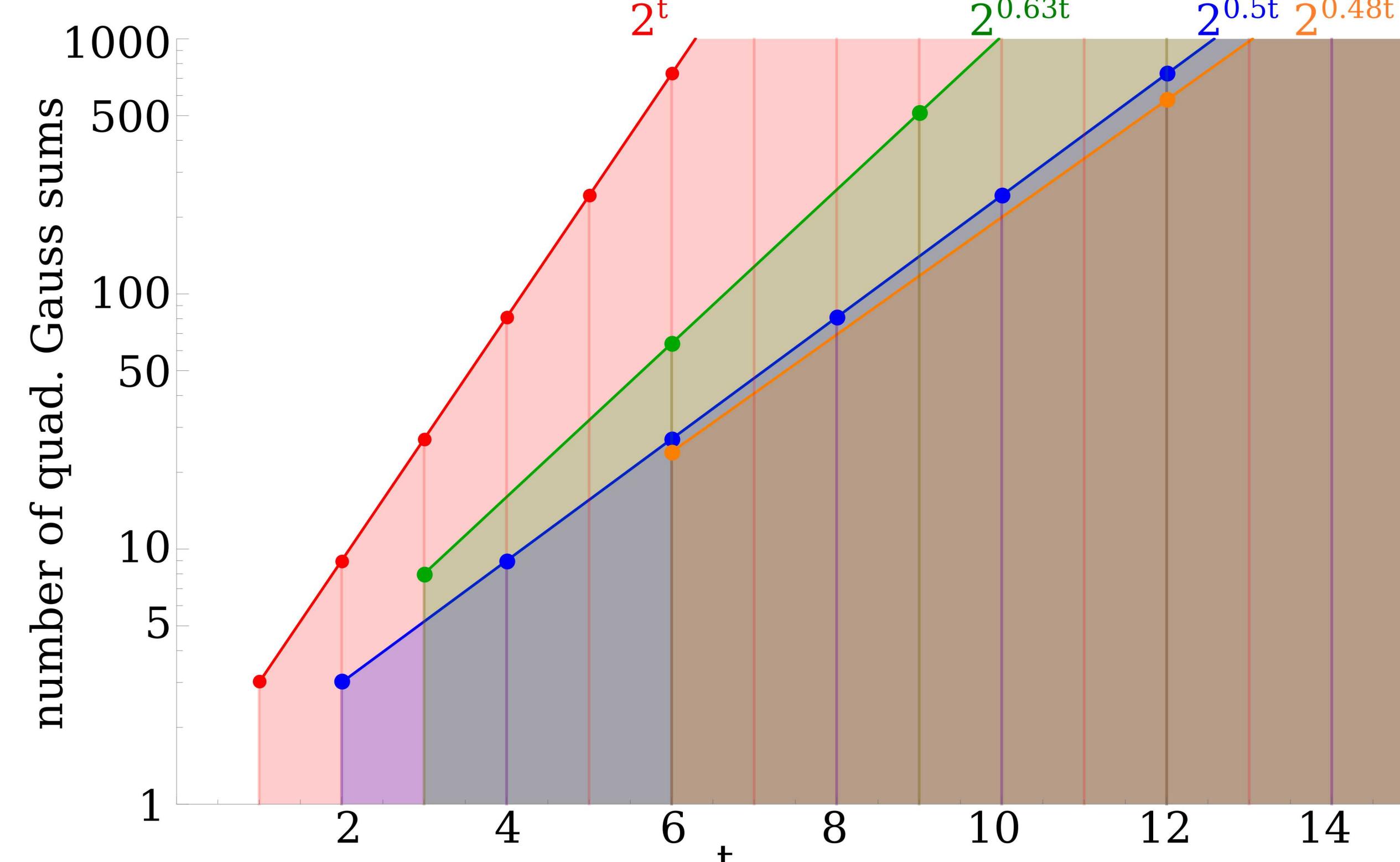


Figure 2: Qutrit ($d=3$) quadratic Gauss sum upper bound w.r.t t from the trivial tensor bound property using the minimal quadratic Gauss sum found for 1, 2, 3, and 6 tensored T-gate magic states

Correspondence between Quadratic Gauss Sums in WWM and Orthogonal Clifford-Separable Pointer Stabilizer Rank ($t < 3$)

This representation of finite odd-dimensional quantum states is especially simple for the Clifford subtheory:

- stabilizer states $\rho_x(x)$ are non-negative
- Clifford gates $U_x(x)$ are symplectic positive maps.

General quantum states $|\Psi\rangle$ can be expressed in terms of quadratic Gauss sums that correspond to **orthogonal Clifford-separable pointer stabilizer states** [5]:

an equiprobable linear combination of orthogonal stabilizer states, which can be written after some Clifford transformation U_C , as products of orthogonal single-qudit stabilizer states $\{|\phi_{ij}\rangle\}_j$ and single-qudit stabilizer states $\{|\psi_{ik}\rangle\}_k$:

$$U_C|\Psi\rangle = \sum_i c_i \prod_j |\phi_{ij}\rangle \prod_k |\psi_{ik}\rangle$$

- this is not a numerical approach!

- to find the minimal number of quadratic Gauss sums you search for the number of non-quadratic phase space (x) variables after Clifford (unit Jacobian) transformations

- for the T-gate magic state with $t=1$, and $t=2$, such a decomposition equals the stabilizer rank

- therefore, the minimal number of quadratic Gauss sums provably equals the stabilizer rank
- however, for $t > 3$, the T-gate magic state stabilizer rank is generally non-orthogonal

Pushing to find Correspondence with (Non-Orthogonal) Stabilizer Rank ($t > 3$)

For qutrits ($d=3$), the minimal number of quadratic Gauss sums equals the stabilizer rank for $t=1$ and $t=2$ (see Figure 2)

- qualitatively, the qubit ($d=2$) and qutrit ($d=3$) correspond approximately $2^{at} \leftrightarrow 3^{at}$
- however, qutrit stabilizer rank for $t > 3$ is not known (search space too large for Monte Carlo)

- despite the T-gate magic state's optimal stabilizer decompositions not being orthogonal for $t > 3$, the number of quadratic Gauss sums continues to qualitatively correspond to the stabilizer rank for $t=3$ and $t=6$ (see Figure 2).

New Optimal Exponential Coefficient Found for the Asymptotic Stabilizer Rank of the T-Gate Magic State?

Can this correspondence be used to push the search for the T-gate magic state stabilizer rank beyond six qudits and find a better asymptotic scaling?

Bibliography

[1] Gottesman, Daniel (1998). "The Heisenberg Representation of Quantum Computers." arXiv:quant-ph/9807006v1

[2] Bravyi, Sergey, and Jeongwan Haah. "Magic-state distillation with low overhead." Physical Review A 86.5 (2012): 052329.

[3] Bravyi, Sergey, Graeme Smith, and John A. Smolin. "Trading classical and quantum computational resources." Physical Review X 6.2 (2016): 021043.

[4] Rivas, A. M. F., and AM Ozorio De Almeida. "The Weyl representation on the torus." Annals of Physics 276.2 (1999): 223-256.

[5] Kocia, Lucas, and Peter Love. "Stationary Phase Method in Discrete Wigner Functions and Classical Simulation of Quantum Circuits." arXiv preprint arXiv:1810.03622 (2018)

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