

# The Stabilizer Rank of the T-Gate Magic State

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## Stabilizer states

- given a state  $|\Psi\rangle$ , the stabilizer group is the set of tensor products of Pauli operators that stabilizes  $|\Psi\rangle$   
e.g.  $I|0\rangle = |0\rangle$  and  $Z|0\rangle = |0\rangle$ .
- stabilizer states  $\{|\phi_i\rangle\}_i$  on  $n$   $d$ -dimensional qudits are the set of states stabilized by  $dxn$  commuting stabilizer group elements
- we define the Clifford gateset to be closed for stabilizer states
  - gateset is not universal
- projection onto Paulis also takes stabilizer states to themselves:  
Clifford subtheory = Pauli measurement + Clifford gateset + (convex combos of) stabilizer states

**Gottesman-Knill theorem:** there is a polynomial-time classical algorithm to simulate the Clifford subtheory [1].

- stabilizer states  $\{|\phi_i\rangle\}_i$  form an overcomplete basis.
  - therefore, any state  $|\Psi\rangle$  can be expressed as  $|\Psi\rangle = \sum_i c_i |\phi_i\rangle$

- the **stabilizer rank**  $\chi(\Psi)$  of a pure state  $|\Psi\rangle$  is the minimal number  $\chi$  of states required in a stabilizer state decomposition of  $|\Psi\rangle$ .

**Trivial tensor bound property:** Let  $\chi(n)$  be the stabilizer rank of  $|\Psi\rangle^n$ . Since the tensor product of two stabilizer states is a stabilizer state, it follows that  $\chi(m+n) \leq \chi(m)\chi(n)$ .

## T-Gate Magic State

$T = 2^{-1/2}(|0\rangle + e^{\pi i/4}|1\rangle)$ .

The T-gate magic state extends the Clifford subtheory to universality in the limit of  $t \rightarrow \infty$ , where  $t$  denotes the power of tensoring the  $|T\rangle^t$

- T-gate magic state equivalent to a T-gate by teleportation protocol
- imperfect T-gates can be "distilled" to increase purity by magic state distillation [2]

It is postulated that  $\chi(t)$  grows slowest with increasing number of qubits.

- for  $t = 1$  qubit T gate magic state,  $\chi(1) = 2$ .
- for  $t = 2$  qubit T gate magic states,  $\chi(2) = 2$ .
- for  $t = 3$  qubit T gate magic states,  $\chi(3) = 3$ .
- for  $t = 6$  qubit T gate magic states,  $\chi(6) \leq 7$ .

It is conjectured that this bound is tight [3].

Using the tensor bound property, we can obtain upper bounds on the T-gate stabilizer rank as  $t \rightarrow \infty$  (see Figure 1):

- $\chi(1)^t = 2^t$ .
- $\chi(2)^{t/2} = 2^{0.5t}$  where  $t$  is even.
- $\chi(3)^{t/3} = 3^{t/3} \approx 2^{0.53t}$  where  $t$  is a multiple of 3.
- $\chi(6)^{t/6} = 7^{t/6} \approx 2^{0.47t}$  where  $t$  is a multiple of 6.

The last bound provides the most favorable asymptotic scaling, and so the outcome of this procedure is often reported as a scaling of  $O(2^{0.468t})$  for the qubit T gate magic state.

## Stabilizer Rank as a Computational Cost Metric for Pauli-Based Computation

- the inner product of two stabilizer states  $\langle \phi_i | \phi_j \rangle$  is governed by Gaussian elimination and therefore scales as  $O(n^3)$ .

- in a Pauli-based computation, given a projector  $E = \prod_{i=1}^t (I + \sigma_i P_a) / 2$  where  $P_a$  is a Pauli operator, it follows that  $E_{ij}^{\chi(t)} = c_i c_j \langle \phi_i | \prod | \phi_j \rangle$

- since the number of terms is  $\chi(t)^2$ , we want to use the lowest  $\chi(t)$  that we can to simulate this classically

- therefore, stabilizer state decompositions are a good way to measure the cost of strong simulation (i.e. classical computation of the probability outcome of a measurement)
- the stabilizer rank of states (like the T-gate magic state) that extend the Clifford subtheory to a universal set necessarily grow exponentially with tensoring
- the exponential factor of this scaling determines how large of a universal quantum computer can be simulated by today's classical computers

## Current Technique for Finding Stabilizer Rank Not Adequate

The fastest method for finding the stabilizer rank of the T-gate magic state is numerical and based on Monte Carlo (Glauber dynamics) sampling of the full stabilizer state space [3]

- for greater than  $t > 7$ , such an approach ceases to converge

To find a better asymptotic tensor scaling requires reaching larger  $t$ , which therefore requires relying on a new method.

- a non-numerical method would be especially attractive

## Alternative Method: Odd Prime-Dimensional Wigner-Weyl-Moyal (WWM) Formalism

Instead of considering our magic state in terms of vectors in Hilbert space, we can use a kernel (or quasiprobability) representation instead.

Given a set of operators  $R(x)$  indexed by  $x$ , that are Hilbert-Schmidt orthogonal, any operator  $A$  can be represented:

$$A = d^{-1} \sum_x \text{Tr}(R(x)A) R(x) \equiv \sum_x A_x(x) R(x).$$

We choose  $R(x)$  to be Hermitian, self-inverse and unitary, and so  $A_x(x)$  are real.

In particular, we consider the odd-prime- $d$  Wigner "reflection" operators [4].

## Correspondence between Quadratic Gauss Sums in WWM and Orthogonal Clifford-Separable Pointer Stabilizer Rank ( $t < 3$ )

This representation of finite odd-dimensional quantum states is especially simple for the Clifford subtheory:

- stabilizer states  $\rho_x(x)$  are non-negative
- Clifford gates  $U_x(x)$  are symplectic positive maps.

General quantum states  $|\Psi\rangle$  can be expressed in terms of quadratic Gauss sums that correspond to **orthogonal Clifford-separable pointer stabilizer states** [5]:

an equiprobable linear combination of orthogonal stabilizer states, which can be written after some Clifford transformation  $U_C$ , as products of orthogonal single-qudit stabilizer states  $\{|\phi_{ij}\rangle\}_j$  and single-qudit stabilizer states  $\{|\psi_{ik}\rangle\}_k$ :

$$U_C |\Psi\rangle = \sum_i c_i \prod_j |\phi_{ij}\rangle \prod_k |\psi_{ik}\rangle$$

- this is not a numerical approach!
- to find the minimal number of quadratic Gauss sums you search for the number of non-quadratic phase space ( $x$ ) variables after Clifford (unit Jacobian) transformations

- for the T-gate magic state with  $t=1$ , and  $t=2$ , such a decomposition equals the stabilizer rank
  - therefore, the minimal number of quadratic Gauss sums provably equals the stabilizer rank
- however, for  $t > 3$ , the T-gate magic state stabilizer rank is generally non-orthogonal

## Pushing to find Correspondence with (Non-Orthogonal) Stabilizer Rank ( $t > 3$ )

For qutrits ( $d=3$ ), the minimal number of quadratic Gauss sums equals the stabilizer rank for  $t=1$  and  $t=2$  (see Figure 2)

- qualitatively, the qubit ( $d=2$ ) and qutrit ( $d=3$ ) correspond approximately  $2^{\text{at}} \leftrightarrow 3^{\text{at}}$
- however, qutrit stabilizer rank for  $t > 3$  is not known (search space too large for Monte Carlo)

- despite the T-gate magic state's optimal stabilizer decompositions not being orthogonal for  $t > 3$ , the number of quadratic Gauss sums continues to qualitatively correspond to the stabilizer rank for  $t=3$  and  $t=6$  (see Figure 2).

## New Optimal Exponential Coefficient Found for the Asymptotic Stabilizer Rank of the T-Gate Magic State?

Can this correspondence be used to push the search for the T-gate magic state stabilizer rank beyond six qudits and find a better asymptotic scaling?

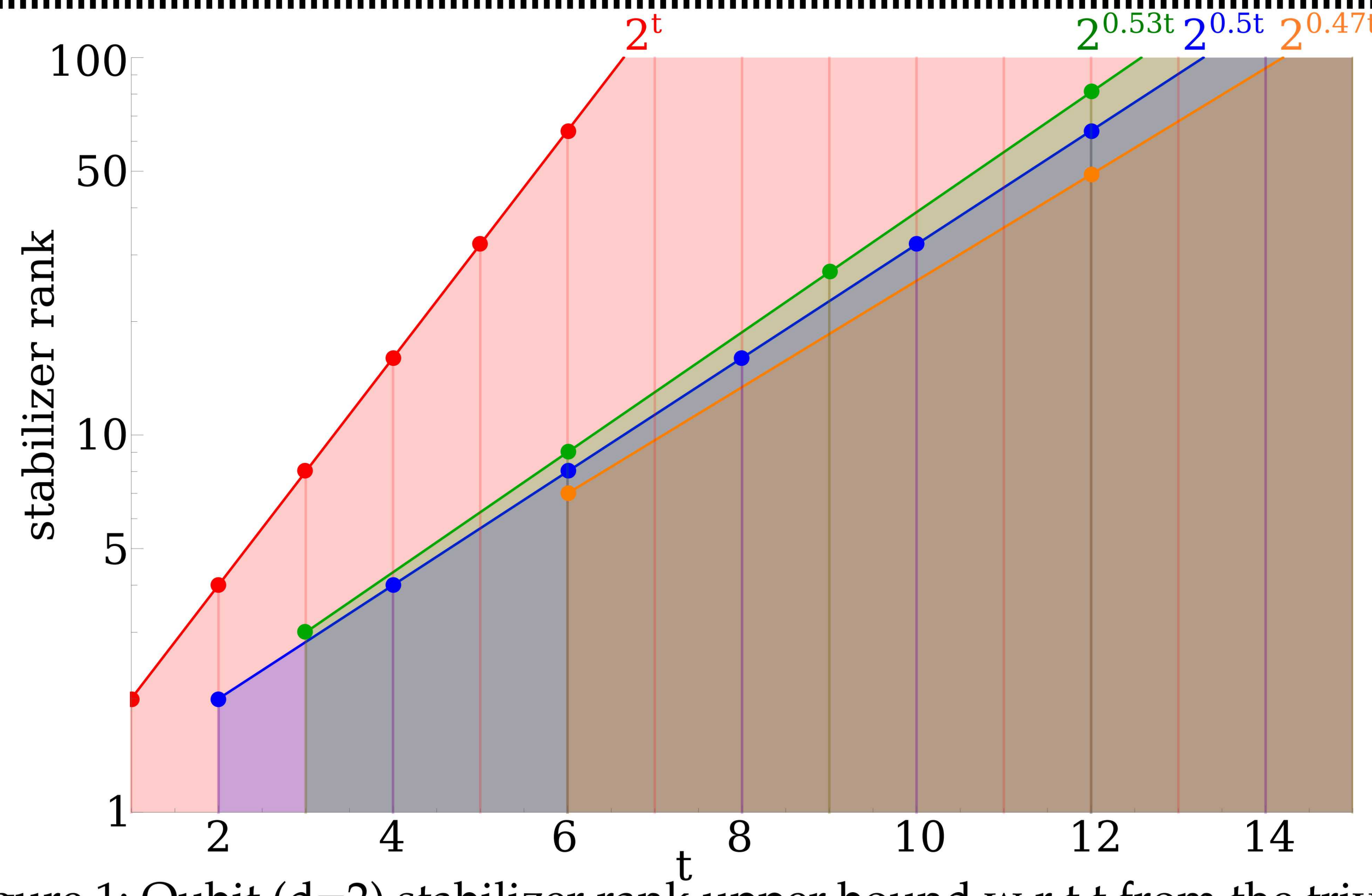


Figure 1: Qubit ( $d=2$ ) stabilizer rank upper bound w.r.t  $t$  from the trivial tensor bound property using the stabilizer rank found for 1, 2, 3, and 6 tensored T-gate magic states

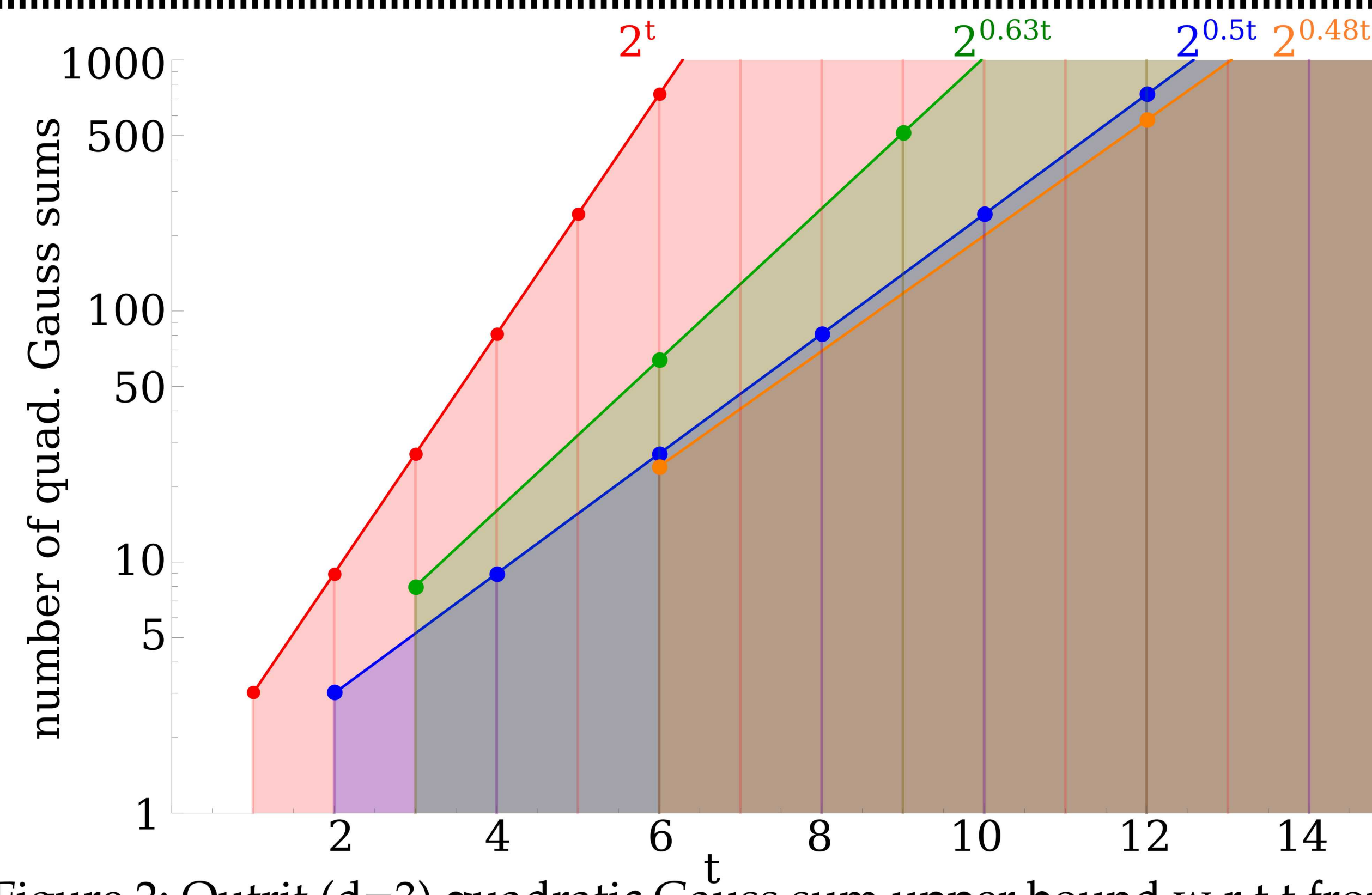


Figure 2: Qutrit ( $d=3$ ) quadratic Gauss sum upper bound w.r.t  $t$  from the trivial tensor bound property using the minimal quadratic Gauss sum found for 1, 2, 3, and 6 tensored T-gate magic states

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